

On the classification of Darboux integrable chains

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We study a differential-difference equation of the form $t_x(n+1)=f(t(n),t(n+1),t_x(n))$ with unknown $t=t(n,x)$ depending on x and n . The equation is called a Darboux integrable if there exist functions F (called an x -integral) and I (called an n -integral), both of a finite number of variables $x,t(n),t(n\pm 1),t(n\pm 2),\dots,t_x(n),t_{xx}(n),\dots$, such that $D_x F=0$ and $DI=I$, where D_x is the operator of total differentiation with respect to x and D is the shift operator: $Dp(n)=p(n+1)$. The Darboux integrability property is reformulated in terms of characteristic Lie algebras that give an effective tool for classification of integrable equations. The complete list of equations of the form above admitting nontrivial x -integrals is given in the case when the function f is of the special form $f(x,y,z)=z+d(x,y)$. © 2008 American Institute of Physics. [DOI: 10.1063/1.2992950]

I. INTRODUCTION

In this paper we study integrable semidiscrete chains of the following form:

$$t_x(n+1) = f(t(n), t(n+1), t_x(n)), \quad (1)$$

where the unknown $t=t(n,x)$ is a function of two independent variables: discrete n and continuous x . Chain (1) can also be interpreted as an infinite system of ordinary differential equations for the sequence of the variables $\{t(n)\}_{n=-\infty}^{\infty}$. Here $f=f(t,t_1,t_x)$ is assumed to be a locally analytical function of three variables satisfying at least locally the condition

$$\frac{\partial f}{\partial t_x} \neq 0. \quad (2)$$

For the sake of convenience we introduce subindex denoting shifts $t_k=t(n+k,x)$ (keep $t_0=t$) and derivatives $t_x=(\partial/\partial x)t(n,x)$, $t_{xx}=(\partial^2/\partial x^2)t(n,x)$, and so on. We denote through D and D_x the shift operator and, correspondingly, the operator of total derivative with respect to x . For instance, $Dh(n,x)=h(n+1,x)$ and $D_x h(n,x)=(\partial/\partial x)h(n,x)$. Set of all the variables $\{t_k\}_{k=-\infty}^{\infty}$; $\{D_x^m t\}_{m=1}^{\infty}$ constitutes the set of dynamical variables. Below we consider the dynamical variables as independent ones. Since in the literature the term “integrable” has various meanings let us specify the meaning used in the article. Introduce first notions of n - and x -integrals.¹

Functions I and F , both depending on x and a finite number of dynamical variables, are called, respectively, n - and x -integrals of (1) if $DI=I$ and $D_x F=0$.

Definition: Chain (1) is called integrable (Darboux integrable) if it admits a nontrivial n -integral and a nontrivial x -integral.

Darboux integrability implies the so-called C -integrability. Knowing both integrals F and I a Cole–Hopf-type differential substitution $w=F+I$ reduces Eq. (1) to the discrete version of the D’Alembert wave equation, $w_{1x}-w_x=0$. Indeed, $(D-1)D_x(w)=(D-1)D_x F+D_x(D-1)I=0$.

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It is remarkable that an integrable chain is reduced to a pair consisting of an ordinary differential equation and an ordinary difference equation. To illustrate it, note first that any n -integral might depend only on x - and x -derivatives of the variable t , $I=I(x,t,t_x,t_{xx},\dots)$, and similarly any x -integral depends only on x and the shifts, $F=F(x,t,t_{\pm 1},t_{\pm 2},\dots)$. Therefore each solution of integrable chain (1) satisfies two equations:

$$I(x,t,t_x,t_{xx},\dots)=p(x), \quad F(x,t,t_{\pm 1},t_{\pm 2},\dots)=q(n),$$

with properly chosen functions $p(x)$ and $q(n)$.

Nowadays the discrete phenomena are studied intensively due to their various applications in physics. For the discussions and references we refer to the articles in Refs. 1–5.

Chain (1) is very close to a well studied object—the partial differential equation of the hyperbolic type

$$u_{xy}=f(x,y,u,u_x,u_y). \quad (3)$$

The definition of integrability for Eq. (3) was introduced by Darboux. The famous Liouville equation $u_{xy}=e^u$ provides an illustrative example of the Darboux integrable equation. An effective criterion of integrability of (3) was discovered by Darboux himself: Eq. (3) is integrable if and only if the Laplace sequence of the linearized equation terminates at both ends (see Refs. 6–8). This criterion of integrability was used in Ref. 8 where the complete list of all Darboux integrable equations of form (3) is given.

An alternative approach to the classification problem based on the notion of the characteristic Lie algebra of hyperbolic-type systems was introduced years ago in Refs. 9 and 10. In these articles an algebraic criterion of Darboux integrability property has been formulated. An important classification result was obtained in Ref. 9 for the exponential system

$$u_{xy}^i = \exp(a_{i1}u^1 + a_{i2}u^2 + \dots + a_{in}u^n), \quad i = 1, 2, \dots, n. \quad (4)$$

It was proved that system (4) is a Darboux integrable if and only if the matrix $A=(a_{ij})$ is the Cartan matrix of a semisimple Lie algebra. Properties of the characteristic Lie algebras of the hyperbolic systems

$$u_{xy}^i = c_{jk}^i u^j u^k, \quad i, j, k = 1, 2, \dots, n, \quad (5)$$

have been studied in Refs. 11 and 12. Hyperbolic systems of general form admitting integrals are studied in Ref. 13. A promising idea of adopting the characteristic Lie algebras to the problem of classification of the hyperbolic systems which are integrated by means of the inverse scattering transforms method is discussed in Ref. 14.

The method of characteristic Lie algebras is closely connected with the symmetry approach¹⁵ which is proved to be a very effective tool to classify integrable nonlinear equations of evolutionary type^{16–20} (see also the survey in Ref. 3 and references therein). However, the symmetry approach meets very serious difficulties when applied to hyperbolic-type models. After the papers in Refs. 21 and 22 it became clear that this case needs alternative methods.

In this article an algorithm of classification of integrable discrete chains of form (1) is suggested based on the notion of the characteristic Lie algebra (see also Refs. 23–25).

To introduce the characteristic Lie algebra L_n of (1) in the direction of n , note that

$$D^{-j} \frac{\partial}{\partial t_1} D^j I = 0 \quad (6)$$

for any n -integral I and $j \geq 1$. Indeed, the equation $DI=I$ can be rewritten in an enlarged form as

$$I(x, n+1, t_1, f, f_x, f_{xx}, \dots) = I(x, n, t, t_x, t_{xx}, \dots). \quad (7)$$

The left hand side DI of equality (7) contains the variable t_1 , while the right hand side does not. Hence, $(\partial/\partial t_1)(DI)=0$, which implies

$$D^{-1} \frac{\partial}{\partial t_1} D I = 0.$$

Proceeding this way one can easily prove (6) from the equality $D^j I = I$, $j \geq 1$.

Define vector fields

$$Y_j = D^{-j} \frac{\partial}{\partial t_1} D^j, \quad j \geq 1, \quad (8)$$

and

$$X_j = \frac{\partial}{\partial t_{-j}}, \quad j \geq 1. \quad (9)$$

We have $Y_j I = 0$ and $X_j I = 0$ for any n -integral I of (1) and $j \geq 1$. The following theorem (see Ref. 24) defines the characteristic Lie algebra L_n of (1).

Theorem 1: Equation (1) admits a nontrivial n -integral if and only if the following two conditions hold:

- (1) Linear space spanned by the operators $\{Y_j\}_{j=1}^{\infty}$ is of finite dimension. We denote this dimension by N .
- (2) Lie algebra L_n generated by the operators $Y_1, Y_2, \dots, Y_N, X_1, X_2, \dots, X_N$ is of finite dimension. We call L_n the characteristic Lie algebra of (1) in the direction of n .

To introduce the characteristic Lie algebra L_x of (1) in the direction of x , note that Eq. (1) due to (2) can be rewritten as $t_x(n-1) = g(t(n), t(n-1), t_x(n))$. An x -integral $F(x, t, t_{\pm 1}, t_{\pm 2}, \dots)$ solves the equation $D_x F = 0$, i.e., $K_0 F = 0$, where

$$K_0 = \frac{\partial}{\partial x} + t_x \frac{\partial}{\partial t} + f \frac{\partial}{\partial t_1} + g \frac{\partial}{\partial t_{-1}} + f_1 \frac{\partial}{\partial t_2} + g_{-1} \frac{\partial}{\partial t_{-2}} + \dots \quad (10)$$

Since F does not depend on the variable t_x one gets $X F = 0$, where

$$X = \frac{\partial}{\partial t_x}. \quad (11)$$

Therefore, any vector field from the Lie algebra generated by K_0 and X annihilates F . This algebra is called the characteristic Lie algebra L_x of chain (1) in the x -direction.

The following result is essential. Its proof is a simple consequence of the famous Jacobi theorem (the Jacobi theorem is discussed, for instance, in Ref. 10).

Theorem 2: Equation (1) admits a nontrivial x -integral if and only if its Lie algebra L_x is of finite dimension.

In the present paper we restrict ourselves to consideration of the existence of x -integrals for a particular kind of chain (1), namely, we study chains of the form

$$t_{1x} = t_x + d(t, t_1) \quad (12)$$

admitting nontrivial x -integrals. The main result of the paper, Theorem 3 below, is the complete list of chains (12) admitting nontrivial x -integrals.

Theorem 3: Chain (12) admits a nontrivial x -integral if and only if $d(t, t_1)$ is one of the following kinds:

- (1) $d(t, t_1) = A(t - t_1)$,
- (2) $d(t, t_1) = c_0 t(t - t_1) + c_2(t - t_1)^2 + c_3 t - c_3 t_1$,
- (3) $d(t, t_1) = A(t - t_1) e^{at}$,

- (4) $d(t, t_1) = c_4(e^{\alpha t_1} - e^{\alpha t}) + c_5(e^{-\alpha t_1} - e^{-\alpha t})$,
 where $A = A(t - t_1)$ is a function of $\tau = t - t_1$ and $c_0 - c_5$ are some constants with $c_0 \neq 0$, $c_4 \neq 0$, and $c_5 \neq 0$, and α is a nonzero constant. Moreover, x -integrals in each of the cases are
- (i) $F = x + \int^\tau du / A(u)$ if $A(u) \neq 0$ and $F = t_1 - t$ if $A(u) \equiv 0$,
 (ii) $F = (1 / (-c_2 - c_0)) \ln |(-c_2 - c_0) \tau_1 / \tau_2 + c_2| + 1 / c_2 \ln |c_2 \tau_1 / \tau - c_2 - c_0|$ for $c_2(c_2 + c_0) \neq 0$, $F = \ln \tau_1 - \ln \tau_2 + \tau_1 / \tau$ for $c_2 = 0$, and $F = \tau_1 / \tau_2 - \ln \tau + \ln \tau_1$ for $c_2 = -c_0$,
 (iii) $F = \int^\tau e^{-\alpha u} du / A(u) - \int^{\tau_1} du / A(u)$, and
 (iv) $F = [(e^{\alpha t} - e^{\alpha t_2})(e^{\alpha t_1} - e^{\alpha t_3})] / [(e^{\alpha t} - e^{\alpha t_3})(e^{\alpha t_1} - e^{\alpha t_2})]$.

The n -integrals of chain (12) can be studied in a similar way by using Theorem 1, but this problem is out of the frame of the present article.

The article is organized as follows. In Sec. II, by using the properly chosen sequence of multiple commutators, a very rough classification result is obtained: function $d(t, t_1)$ for chain (12) admitting x -integrals is a quasipolynomial on t with coefficients depending on $\tau = t - t_1$. Then it is observed that the exponents $\alpha_0 = 0, \alpha_1, \dots, \alpha_s$ in expansion (26) cannot be arbitrary. For example, if the coefficient before $e^{\alpha_0 t} = 1$ is not identically zero, then the quasipolynomial $d(t, t_1)$ is really a polynomial on t with coefficients depending on τ . In Sec. III we prove that the degree of this polynomial is at most 1. If d contains a term of the form $\mu(\tau) t^j e^{\alpha_k t}$ with $\alpha_k \neq 0$, then $j = 0$ (Sec. IV). In Sec. V it is proved that if d contains terms with $e^{\alpha_k t}$ and $e^{\alpha_j t}$ having nonzero exponents, then $\alpha_k = -\alpha_j$. This last case contains chains having infinite dimensional characteristic Lie algebras for which the sequence of multiple commutators grows very slowly. They are studied in Secs. VI and VII. One can find the well known semidiscrete version of the sine-Gordon (SG) model among them. It is worth mentioning that in Sec. VII the characteristic Lie algebra L_x for semidiscrete SG is completely described. The last section, Sec. VIII, contains the proof of the main theorem, Theorem 3, and here the method of constructing x -integrals is also briefly discussed.

II. A NECESSARY INTEGRABILITY CONDITION

Define a class \mathbf{F} of locally analytical functions each of which depends only on a finite number of dynamical variables. In particular, we assume that $f(t, t_1, t_x) \in \mathbf{F}$. We will consider vector fields given as an infinite formal series of the form

$$Y = \sum_{k=-\infty}^{\infty} y_k \frac{\partial}{\partial t_k}, \quad (13)$$

with coefficients $y_k \in \mathbf{F}$. Introduce notions of linearly dependent and independent sets of vector fields (13). Denote through P_N the projection operator acting according to the rule

$$P_N(Y) = \sum_{k=-N}^N y_k \frac{\partial}{\partial t_k}. \quad (14)$$

First we consider finite vector fields as

$$Z = \sum_{k=-N}^N z_k \frac{\partial}{\partial t_k}. \quad (15)$$

We say that a set of finite vector fields Z_1, Z_2, \dots, Z_m is linearly dependent in some open region \mathbf{U} if there is a set of functions $\lambda_1, \lambda_2, \dots, \lambda_m$ defined on \mathbf{U} such that the function $|\lambda_1|^2 + |\lambda_2|^2 + \dots + |\lambda_m|^2$ does not vanish identically and the condition

$$\lambda_1 Z_1 + \lambda_2 Z_2 + \dots + \lambda_m Z_m = 0 \quad (16)$$

holds for each point of region \mathbf{U} .

We call a set of the vector fields Y_1, Y_2, \dots, Y_m of form (13) linearly dependent in region \mathbf{U} if for each natural N the set of finite vector fields $P_N(Y_1), P_N(Y_2), \dots, P_N(Y_m)$ is linearly dependent in this region. Otherwise we call the set Y_1, Y_2, \dots, Y_m linearly independent in \mathbf{U} .

The following proposition is very useful. Its proof is almost evident.

Proposition: *If a vector field Y is expressed as a linear combination,*

$$Y = \lambda_1 Y_1 + \lambda_2 Y_2 + \dots + \lambda_m Y_m, \quad (17)$$

where the set of vector fields Y_1, Y_2, \dots, Y_m is linearly independent in \mathbf{U} and the coefficients of all the vector fields Y, Y_1, Y_2, \dots, Y_m belonging to \mathbf{F} are defined in \mathbf{U} , then the coefficients $\lambda_1, \lambda_2, \dots, \lambda_m$ are in \mathbf{F} .

Below we concentrate on the class of chains of form (12). For this case the Lie algebra L_x splits down into a direct sum of two subalgebras. Indeed, since $f = t_x + d$ and $g = t_x - d_{-1}$ one gets $f_k = t_x + d + \sum_{j=1}^k d_j$ and $g_{-k} = t_x - \sum_{j=1}^{k+1} d_{-j}$ for $k \geq 1$, where $d = d(t, t_1)$ and $d_j = d(t_j, t_{j+1})$. Due to this observation the vector field K_0 can be rewritten as $K_0 = t_x \tilde{X} + Y$, with

$$\tilde{X} = \frac{\partial}{\partial t} + \frac{\partial}{\partial t_1} + \frac{\partial}{\partial t_{-1}} + \frac{\partial}{\partial t_2} + \frac{\partial}{\partial t_{-2}} + \dots \quad (18)$$

and

$$Y = \frac{\partial}{\partial x} + d \frac{\partial}{\partial t_1} - d_{-1} \frac{\partial}{\partial t_{-1}} + (d + d_1) \frac{\partial}{\partial t_2} - (d_{-1} + d_{-2}) \frac{\partial}{\partial t_{-2}} + \dots$$

Due to the relations $[X, \tilde{X}] = 0$ and $[X, Y] = 0$ we have $\tilde{X} = [X, K_0] \in L_x$; hence $Y \in L_x$. Therefore $L_x = \{X\} \oplus L_{x1}$, where L_{x1} is the Lie algebra generated by the operators \tilde{X} and Y .

Lemma 1: *If Eq. (12) admits a nontrivial x -integral, then it admits a nontrivial x -integral F such that $\partial F / \partial x = 0$.*

Proof: Assume that a nontrivial x -integral of (12) exists. Then the Lie algebra L_{x1} is of finite dimension. One can choose a basis of L_{x1} in the form

$$T_1 = \frac{\partial}{\partial x} + \sum_{k=-\infty}^{\infty} a_{1,k} \frac{\partial}{\partial t_k},$$

$$T_j = \sum_{k=-\infty}^{\infty} a_{j,k} \frac{\partial}{\partial t_k}, \quad 2 \leq j \leq N.$$

Thus, there exists an x -integral F depending on $x, t, t_1, \dots, t_{N-1}$, satisfying the system of equations

$$\frac{\partial F}{\partial x} + \sum_{k=0}^{N-1} a_{1,k} \frac{\partial F}{\partial t_k} = 0,$$

$$\sum_{k=0}^{N-1} a_{j,k} \frac{\partial F}{\partial t_k} = 0, \quad 2 \leq j \leq N.$$

Due to the famous Jacobi theorem¹⁰ there is a change of variables $\theta_j = \theta_j(t, t_1, \dots, t_{N-1})$ that reduces the system to the form

$$\frac{\partial F}{\partial x} + \sum_{k=0}^{N-1} \tilde{a}_{1,k} \frac{\partial F}{\partial \theta_k} = 0,$$

$$\frac{\partial F}{\partial \theta_k} = 0, \quad 2 \leq j \leq N - 2,$$

which is equivalent to

$$\frac{\partial F}{\partial x} + \tilde{a}_{1,N-1} \frac{\partial F}{\partial \theta_{N-1}} = 0$$

for $F = F(x, \theta_{N-1})$.

There are two possibilities: (1) $\tilde{a}_{1,N-1} = 0$ and (2) $\tilde{a}_{1,N-1} \neq 0$. In case (1), we at once have $\partial F / \partial x = 0$. In case (2), $F = x + H(\theta_{N-1}) = x + H(t, t_1, \dots, t_{N-1})$ for some function H . Evidently, $F_1 = DF = x + H(t_1, t_2, \dots, t_N)$ is also an x -integral, and $F_1 - F$ is a nontrivial x -integral not depending on x . □

Below we look for x -integrals F depending on dynamical variables $t, t_{\pm 1}, t_{\pm 2}, \dots$ only (not depending on x). In other words, we study the Lie algebra generated by vector fields \tilde{X} and \tilde{Y} , where

$$\tilde{Y} = d \frac{\partial}{\partial t_1} - d_{-1} \frac{\partial}{\partial t_{-1}} + (d + d_1) \frac{\partial}{\partial t_2} - (d_{-1} + d_{-2}) \frac{\partial}{\partial t_{-2}} + \dots \tag{19}$$

One can prove that the linear operator $Z \rightarrow DZD^{-1}$ defines an automorphism of the characteristic Lie algebra L_x . This automorphism plays the crucial role in all of our further considerations. Further we refer to it as the shift automorphism. For instance, direct calculations show that

$$D\tilde{X}D^{-1} = \tilde{X}, \quad D\tilde{Y}D^{-1} = -d\tilde{X} + \tilde{Y}. \tag{20}$$

Lemma 2: Suppose that a vector field of the form $Z = \sum a(j)(\partial / \partial t_j)$ with the coefficients $a(j) = a(j, t, t_{\pm 1}, t_{\pm 2}, \dots)$ depending on a finite number of the dynamical variables solves an equation of the form $DZD^{-1} = \lambda Z$. If for some $j = j_0$ we have $a(j_0) \equiv 0$, then $Z = 0$.

Proof: By applying the shift automorphism to the vector field Z one gets $DZD^{-1} = \sum D(a(j))(\partial / \partial t_{j+1})$. Now, to complete the proof, we compare the coefficients of $\partial / \partial t_j$ in the equation $\sum D(a(j))(\partial / \partial t_{j+1}) = \lambda \sum a(j)(\partial / \partial t_j)$. □

Construct an infinite sequence of multiple commutators of the vector fields \tilde{X} and \tilde{Y} ,

$$\tilde{Y}_1 = [\tilde{X}, \tilde{Y}], \quad \tilde{Y}_k = [\tilde{X}, \tilde{Y}_{k-1}] \quad \text{for } k \geq 2. \tag{21}$$

Lemma 3: We have

$$D\tilde{Y}_kD^{-1} = -\tilde{X}^k(d)\tilde{X} + \tilde{Y}_k, \quad k \geq 1. \tag{22}$$

Proof: We prove the statement by induction on k . The base of induction holds. Indeed, by (20) and (21), we have

$$D\tilde{Y}_1D^{-1} = D[\tilde{X}, \tilde{Y}]D^{-1} = [D\tilde{X}D^{-1}, D\tilde{Y}D^{-1}] = [\tilde{X}, -d\tilde{X} + \tilde{Y}] = -\tilde{X}(d)\tilde{X} + \tilde{Y}_1.$$

Assuming Eq. (22) holds for $k = n - 1$, we have

$$D\tilde{Y}_nD^{-1} = [D\tilde{X}D^{-1}, D\tilde{Y}_{n-1}D^{-1}] = [\tilde{X}, -\tilde{X}^{n-1}(d)\tilde{X} + \tilde{Y}_{n-1}] = -\tilde{X}^n(d)\tilde{X} + \tilde{Y}_n,$$

which finishes the proof of the lemma. □

Since vector fields X, \tilde{X} , and \tilde{Y} are linearly independent, then the dimension of Lie algebra L_x is at least 3. By (22), case $\tilde{Y}_1 = 0$ corresponds to $\tilde{X}(d) = 0$, or $d_t + d_{t_1} = 0$, which implies $d = A(t - t_1)$, where $A(\tau)$ is an arbitrary differentiable function of one variable.

Assume Eq. (12) admits a nontrivial x -integral and $\tilde{Y}_1 \neq 0$. Consider the sequence of the vector fields $\{\tilde{Y}_1, \tilde{Y}_2, \tilde{Y}_3, \dots\}$. Since L_x is of finite dimension, then there exists a natural number N such that

$$\tilde{Y}_{N+1} = \gamma_1 \tilde{Y}_1 + \gamma_2 \tilde{Y}_2 + \dots + \gamma_N \tilde{Y}_N, \quad N \geq 1, \tag{23}$$

and $\tilde{Y}_1, \tilde{Y}_2, \dots, \tilde{Y}_N$ are linearly independent. Therefore,

$$D\tilde{Y}_{N+1}D^{-1} = D(\gamma_1)D\tilde{Y}_1D^{-1} + D(\gamma_2)D\tilde{Y}_2D^{-1} + \dots + D(\gamma_N)D\tilde{Y}_ND^{-1}, \quad N \geq 1.$$

Due to Lemma 3 and (23) the last equation can be rewritten as

$$\begin{aligned} -\tilde{X}^{N+1}(d)\tilde{X} + \gamma_1 \tilde{Y}_1 + \gamma_2 \tilde{Y}_2 + \dots + \gamma_N \tilde{Y}_N &= D(\gamma_1)(-\tilde{X}(d)\tilde{X} + \tilde{Y}_1) + D(\gamma_2)(-\tilde{X}^2(d)\tilde{X} + \tilde{Y}_2) \\ &+ \dots + D(\gamma_N)(-\tilde{X}^N(d)\tilde{X} + \tilde{Y}_N). \end{aligned}$$

Comparing coefficients before linearly independent vector fields $\tilde{X}, \tilde{Y}_1, \tilde{Y}_2, \dots, \tilde{Y}_N$, we obtain the following system of equations:

$$\tilde{X}^{N+1}(d) = D(\gamma_1)\tilde{X}(d) + D(\gamma_2)\tilde{X}^2(d) + \dots + D(\gamma_N)\tilde{X}^N(d),$$

$$\gamma_1 = D(\gamma_1), \quad \gamma_2 = D(\gamma_2), \quad \dots, \quad \gamma_N = D(\gamma_N).$$

Since the coefficients of the vector fields \tilde{Y}_j depend only on the variables $t, t_{\pm 1}, t_{\pm 2}, \dots$ the factors γ_j might depend only on these variables (see the proposition above). Hence the system of equations implies that all coefficients $\gamma_k, 1 \leq k \leq N$, are constants, and $d = d(t, t_1)$ is a function that satisfies the following differential equation:

$$\tilde{X}^{N+1}(d) = \gamma_1 \tilde{X}(d) + \gamma_2 \tilde{X}^2(d) + \dots + \gamma_N \tilde{X}^N(d), \tag{24}$$

where $\tilde{X}(d) = d_t + d_{t_1}$. Using the substitution $s = t$ and $\tau = t - t_1$, Eq. (24) can be rewritten as

$$\frac{\partial^{N+1}d}{\partial s^{N+1}} = \gamma_1 \frac{\partial d}{\partial s} + \gamma_2 \frac{\partial^2 d}{\partial s^2} + \dots + \gamma_N \frac{\partial^N d}{\partial s^N}, \tag{25}$$

which implies that

$$d(t, t_1) = \sum_k \left(\sum_{j=0}^{m_k-1} \lambda_{k,j}(t-t_1)^j \right) e^{\alpha_k t} \tag{26}$$

for some functions $\lambda_{k,j}(t-t_1)$, where α_k are roots of multiplicity m_k for a characteristic equation of (25).

Let $\alpha_0 = 0, \alpha_1, \dots, \alpha_s$ be the distinct roots of characteristic equation (24). Equation (24) can be rewritten as

$$\Lambda(\tilde{X})d := \tilde{X}^{m_0}(\tilde{X} - \alpha_1)^{m_1}(\tilde{X} - \alpha_2)^{m_2} \dots (\tilde{X} - \alpha_s)^{m_s}d = 0 \tag{27}$$

and $m_0 + m_1 + \dots + m_s = N + 1, m_0 \geq 1$.

Initiated by formula (19) define a map $h \rightarrow Y_h$ which assigns to any function $h = h(t, t_{\pm 1}, t_{\pm 2}, \dots)$ a vector field

$$Y_h = h \frac{\partial}{\partial t_1} - h_{-1} \frac{\partial}{\partial t_{-1}} + (h + h_1) \frac{\partial}{\partial t_2} - (h_{-1} + h_{-2}) \frac{\partial}{\partial t_{-2}} + \dots$$

For any polynomial with constant coefficients $P(\lambda) = c_0 + c_1 \lambda + \dots + c_m \lambda^m$ we have the formula

$$P(ad_{\tilde{X}})\tilde{Y} = Y_{P(\tilde{X})h}, \quad \text{where } ad_{\tilde{X}}Y = [X, Y], \tag{28}$$

which establishes an isomorphism between the linear space V of all solutions of Eq. (25) and the linear space $\tilde{V} = \text{span}\{\tilde{Y}, \tilde{Y}_1, \dots, \tilde{Y}_N\}$ of the corresponding vector fields.

Represent function (26) as the sum $d(t, t_1) = P(t, t_1) + Q(t, t_1)$ of the polynomial part $P(t, t_1) = \sum_{j=0}^{m_0-1} \lambda_{0,j}(t-t_1)^j$ and the ‘‘exponential’’ part $Q(t, t_1) = \sum_{k=1}^s (\sum_{j=0}^{m_k-1} \lambda_{k,j}(t-t_1)^j) e^{\alpha_k t}$.

Lemma 4: Assume that Eq. (12) admits a nontrivial x -integral. Then one of the functions $P(t, t_1)$ and $Q(t, t_1)$ vanishes.

Proof: Assume in contrary that neither of the functions vanish. First we show that in this case algebra L_x contains vector fields $T_0 = Y_{A(\tau)e^{\alpha_k t}}$ and $T_1 = Y_{B(\tau)}$ for some functions $A(\tau)$ and $B(\tau)$. Indeed, take $T_0 := \Lambda_0(ad_{\tilde{X}})\tilde{Y} = Y_{\Lambda_0(\tilde{X})d} \in L_x$, where $\Lambda_0(\lambda) = \Lambda(\lambda)/(\lambda - \alpha_k)$. Evidently the function $\tilde{A}(t, t_1) = \Lambda_0(\tilde{X})d$ solves the equation $(\tilde{X} - \alpha_k)\tilde{A}(t, t_1) = \Lambda(\tilde{X})d = 0$ which implies immediately that $\tilde{A}(t, t_1) = A(\tau)e^{\alpha_k t}$. In a similar way one shows that $T_1 \in L_x$. Note that due to our assumption the functions $A(\tau)$ and $B(\tau)$ cannot vanish identically.

Consider an infinite sequence of the vector fields defined as follows:

$$T_2 = [T_0, T_1], \quad T_3 = [T_0, T_2], \quad \dots, \quad T_n = [T_0, T_{n-1}], \quad n \geq 3.$$

One can show that

$$[\tilde{X}, T_0] = \alpha_k T_0, \quad [\tilde{X}, T_1] = 0, \quad [\tilde{X}, T_n] = \alpha_k(n-1)T_n, \quad n \geq 2,$$

$$DT_0 D^{-1} = -Ae^{\alpha_k t} \tilde{X} + T_0, \quad DT_1 D^{-1} = -B\tilde{X} + T_1,$$

$$DT_n D^{-1} = T_n - \frac{(n-1)(n-2)}{2} \alpha_k A e^{\alpha_k t} T_{n-1} + b_n \tilde{X} + \sum_{k=0}^{n-2} a_k^{(n)} T_k, \quad n \geq 2.$$

Since algebra L_x is of finite dimension, then there exists number N such that

$$T_{N+1} = \lambda \tilde{X} + \mu_0 T_0 + \mu_1 T_1 + \dots + \mu_N T_N, \tag{29}$$

and vector fields $\tilde{X}, T_0, T_1, \dots, T_N$ are linearly independent. We have

$$DT_{N+1} D^{-1} = D(\lambda)\tilde{X} + D(\mu_0)\{-Ae^{\alpha_k t} \tilde{X} + T_0\} + \dots + D(\mu_N)\left\{T_N - \frac{(N-1)(N-2)}{2} \alpha_k A e^{\alpha_k t} T_{N-1} + \dots\right\}.$$

By comparing the coefficients before T_N in the last equation one gets

$$\mu_N - \frac{N(N-1)}{2} \alpha_k A(\tau) e^{\alpha_k t} = D(\mu_N).$$

It follows that μ_N is a function of variable t only. Also, by applying $ad_{\tilde{X}}$ to both sides of Eq. (29), one gets

$$N\alpha_k T_{N+1} = [\tilde{X}, T_{N+1}] = \tilde{X}(\lambda)\tilde{X} + (\tilde{X}(\mu_0) + \mu_0\alpha_k)T_0 + \dots + (\tilde{X}(\mu_N) + \mu_N(N-1)\alpha_k)T_N.$$

Again, by comparing coefficients before T_N , we have

$$N\alpha_k \mu_N = \tilde{X}(\mu_N) + (N-1)\alpha_k \mu_N, \quad \text{i. e. ,} \quad \tilde{X}(\mu_N) = \alpha_k \mu_N.$$

Therefore, $\mu_N = A_1 e^{\alpha_k t}$, where A_1 is some nonzero constant, and thus $A(\tau) e^{\alpha_k t} = A_2 e^{\alpha_k t} - A_2 e^{\alpha_k t}$.

Here A_2 is some constant. We have $T_0 = A_2 e^{\alpha_k t} \tilde{X} - A_2 S_0$, where

$$S_0 = \sum_{j=-\infty}^{\infty} e^{\alpha_k t_j} \frac{\partial}{\partial t_j}.$$

Also,

$$[\tilde{X}, S_0] = \alpha_k S_0, \quad DS_0 D^{-1} = S_0.$$

Consider a new sequence of vector fields,

$$P_1 = S_0, \quad P_2 = [T_1, S_0], \quad P_3 = [T_1, P_2], \quad P_n = [T_1, P_{n-1}], \quad n \geq 3.$$

One can show that

$$[\tilde{X}, P_n] = \alpha_k P_n, \quad DP_n D^{-1} = P_n - \alpha_k(n-1)BP_{n-1} + b_n \tilde{X} + a_n S_0 + \sum_{j=2}^{n-2} a_j^{(n)} P_j, \quad n \geq 2.$$

Since algebra L_x is of finite dimension, then there exists number M such that

$$P_{M+1} = \lambda^* \tilde{X} + \mu_2^* P_2 + \cdots + \mu_M^* P_M, \quad (30)$$

and fields $\tilde{X}, P_2, \dots, P_M$ are linearly independent. Thus,

$$DP_{M+1} D^{-1} = D(\lambda^*) \tilde{X} + D(\mu_2^*) \{P_2 + \cdots\} + \cdots + D(\mu_M^*) \{P_M - \alpha_k(M-1)BP_{M-1} + \cdots\}.$$

We compare the coefficients before P_M in the last equation and get

$$\mu_M^* - M\alpha_k B(\tau) = D(\mu_M^*), \quad (31)$$

which implies that μ_M^* is a function of variable t only. Also, by applying $ad_{\tilde{X}}$ to both sides of (30), one gets

$$\alpha_k P_{M+1} = [\tilde{X}, P_{M+1}] = \tilde{X}(\lambda^*) \tilde{X} + (\tilde{X}(\mu_2^*) + \alpha_k \mu_2^*) P_2 + \cdots + (\tilde{X}(\mu_M^*) + \alpha_k \mu_M^*) P_M.$$

Again, we compare the coefficients before P_M and have $\alpha_k \mu_M^*(t) = \tilde{X}(\mu_M^*(t)) + \alpha_k \mu_M^*(t)$, which implies that μ_M^* is a constant. It follows then from (31) that $B(\tau) = 0$. This contradiction shows that our assumption that both functions are not identically zero was wrong. \square

III. MULTIPLE ZERO ROOT

In this section we assume that Eq. (12) admits a nontrivial x -integral and that $\alpha_0 = 0$ is a root of the characteristic polynomial $\Lambda(\lambda)$. Then, due to Lemma 4, zero is the only root and therefore $\Lambda(\lambda) = \lambda^{m+1}$. It follows from formula (26) with $m_0 = m+1$ that

$$d(t, t_1) = a(\tau)t^m + b(\tau)t^{m-1} + \cdots, \quad m = m_0 - 1 \geq 0.$$

The case $m=0$ corresponds to a very simple equation, $t_{1x} = t_x + A(t-t_1)$, which is easily solved in quadratures, so we concentrate on the case $m \geq 1$. For this case the characteristic algebra L_x contains a vector field $T = Y_{\tilde{\kappa}}$ with

$$\tilde{\kappa} = a(\tau)t + \frac{1}{m}b(\tau).$$

Indeed,

$$T = \frac{1}{m!} ad_{\tilde{X}}^{m-1} \tilde{Y} = Y_{\tilde{\kappa}}. \quad (32)$$

Introduce a sequence of multiple commutators defined as follows:

$$T_0 = \tilde{X}, \quad T_1 = [T, T_0] = Y_{-a(\tau)}, \quad T_{k+1} = [T, T_k], \quad k \geq 0, \quad T_{k,0} = [T_0, T_k].$$

Note that $T_{1,0}=0$. We will see below that the linear space spanned by this sequence is not invariant under the action of the shift automorphism $Z \rightarrow DZD^{-1}$ introduced above. We extend the sequence to provide the invariance property. We define T_α with the multi-index α . For any sequence $\alpha = k, 0, i_1, i_2, \dots, i_{n-1}, i_n$, where k is any natural number, $i_j \in \{0; 1\}$, denote

$$T_\alpha = \begin{cases} [T_0, T_{k,0,i_1,\dots,i_{n-1}}] & \text{if } i_n = 0 \\ [T, T_{k,0,i_1,\dots,i_{n-1}}] & \text{if } i_n = 1, \end{cases}$$

$$m(\alpha) = \begin{cases} k & \text{if } \alpha = k \\ k & \text{if } \alpha = k, 0 \\ k + i_1 + \dots + i_n & \text{if } \alpha = k, 0, i_1, \dots, i_n, \end{cases}$$

$$l(\alpha) = k + n + 1 - m(\alpha).$$

The multi-index α is characterized by two quantities, $m(\alpha)$ and $l(\alpha)$, which allow to order partially the sequence $\{T_\alpha\}$. We have

$$DT_0D^{-1} = T_0, \quad DTD^{-1} = T - \tilde{\kappa}T_0, \quad DT_1D^{-1} = T_1 + aT_0.$$

One can prove by induction on k that

$$DT_kD^{-1} = T_k + aT_{k-1} - \tilde{\kappa} \sum_{m(\beta)=k-1} T_\beta + \sum_{m(\beta)\leq k-2} \eta(k, \beta)T_\beta. \tag{33}$$

In general, for any α ,

$$DT_\alpha D^{-1} = T_\alpha + \sum_{m(\beta)\leq m(\alpha)-1} \eta(\alpha, \beta)T_\beta. \tag{34}$$

We can choose a system P of linearly independent vector fields in the following way:

- (1) T and T_0 are linearly independent. We take them into P .
- (2) We check whether T, T_0 , and T_1 are linearly independent or not. If they are dependent, then $P = \{T, T_0\}$ and $T_1 = \mu T + \lambda T_0$ for some functions μ and λ .
- (3) If T, T_0 , and T_1 are linearly independent, then we check whether T, T_0 , and T_1, T_2 are linearly independent or not. If they are dependent, then $P = \{T, T_0, T_1\}$.
- (4) If T, T_0, T_1 , and T_2 are linearly independent, we add vector fields T_β , $m(\beta)=2$, $\beta \in I_2$ (actually, by definition I_2 is the collection of such β), in such a way that $J_2 := \{T, T_0, T_1, T_2, \cup_{\beta \in I_2} T_\beta\}$ is a system of linearly independent vector fields and for any T_γ with $m(\gamma) \leq 2$ we have $T_\gamma = \sum_{\beta \in J_2} \mu(\gamma, \beta)T_\beta$.
- (5) We check whether $T_3 \cup J_2$ is a linearly independent system. If it is not, then P consists of all elements from J_2 , and $T_3 = \sum_{\beta \in J_2} \mu(\gamma, \beta)T_\beta$. If it is, then to the system $T_3 \cup J_2$ we add vector fields T_β , $m(\beta)=3$, $\beta \in I_3$, in such a way that $J_3 := \{T_3, J_2, \cup_{\beta \in I_3} T_\beta\}$ is a system of linearly independent vector fields and for any T_γ with $m(\gamma) \leq 3$ we have $T_\gamma = \sum_{\beta \in J_3} \mu(\gamma, \beta)T_\beta$.

We continue the construction of system P . Since L_x is of finite dimension, then there exists such a natural number N such that we have the following:

- (i) $T_k \in P, k \leq N$.
- (ii) $m(\beta) \leq N$ for any $T_\beta \in P$.
- (iii) For any T_γ with $m(\gamma) \leq N$ we have $T_\gamma = \sum_{\beta \in P, m(\beta) \leq m(\gamma)} \mu(\gamma, \beta)T_\beta$ and also $T_{N+1} = \mu(N+1, N)T_N + \sum_{\beta \in P, m(\beta) \leq N} \mu(N+1, \beta)T_\beta$.

We then have the following:

- (iv) For any vector field T_α with $m(\alpha)=N$ that does not belong to P , the coefficient $\mu(\alpha, N)$ before T_N in the expansion

$$T_\alpha = \mu(\alpha, N)T_N + \sum_{T_\beta \in P} \mu(\alpha, \beta)T_\beta \tag{35}$$

is constant. Indeed, by (34),

$$DT_\alpha D^{-1} = T_\alpha + \sum_{m(\beta) \leq N-1} \eta(\alpha, \beta)T_\beta = \mu(\alpha, N)T_N + \sum_{T_\beta \in P} \mu(\alpha, \beta)T_\beta + \sum_{m(\beta) \leq N-1} \eta(\alpha, \beta)T_\beta.$$

From (35) we also have

$$DT_\alpha D^{-1} = D(\mu(\alpha, N))DT_N D^{-1} + \sum_{T_\beta \in P} D(\mu(\alpha, \beta))DT_\beta D^{-1} = D(\mu(\alpha, N))\{T_N + \dots\} + \sum_{T_\beta \in P} D(\mu(\alpha, \beta))\{T_\beta + \dots\}.$$

By comparing the coefficients before T_N in these two expressions for $DT_\alpha D^{-1}$, we have

$$\mu(\alpha, N) = D(\mu(\alpha, N)),$$

which implies that $\mu(\alpha, N)$ is a constant indeed.

Lemma 5: We have $a(\tau) = c_0\tau + c_1$, where c_0 and c_1 are some constants.

Proof: Since

$$T_{N+1} = \mu(N+1, N)T_N + \sum_{T_\beta \in P} \mu(N+1, \beta)T_\beta,$$

then

$$DT_{N+1} D^{-1} = D(\mu(N+1, N))\{T_N + \dots\} + \sum_{T_\beta \in P} D(\mu(N+1, \beta))\{T_\beta + \dots\}.$$

On the other hand,

$$DT_{N+1} D^{-1} = T_{N+1} + aT_N - \tilde{\kappa} \sum_{m(\beta)=N} T_\beta + \sum_{m(\beta) \leq N-1} \eta(N+1, \beta)T_\beta.$$

We compare the coefficients before T_N in the last two expressions. For $N \geq 0$ the equation is

$$\mu(N+1, N) + a - \tilde{\kappa} \sum_{T_\beta \in P, m(\beta)=N} \mu(\beta, N) = D(\mu(N+1, N)). \tag{36}$$

Denote by $c = -\sum_{T_\beta \in P, m(\beta)=N} \mu(\beta, N)$ and by $\mu_N = \mu(N+1, N)$. By property (iv), c is a constant. It follows from (36) that μ_N is a function of variables t and n only. Therefore,

$$a(\tau) + c \left(a(\tau)t + \frac{1}{m}b(\tau) \right) = \mu_N(t_1, n+1) - \mu_N(t, n).$$

By differentiating both sides of the equation with respect to t and then t_1 , we have

$$-a''(\tau) - c \left(a''(\tau)t + a'(\tau) + \frac{1}{m}b''(\tau) \right) = 0,$$

which implies that $a''(\tau) = 0$, or the same, $a(\tau) = c_0\tau + c_1$ for some constants c_0 and c_1 . □

Vector fields T_1 and T in new variables are rewritten as

$$T_1 = \sum_{j=-\infty}^{\infty} a(\tau_j) \frac{\partial}{\partial \tau_j}, \quad (37)$$

$$\begin{aligned} T &= - \sum_{j=-\infty}^{\infty} \left\{ a(\tau_j) t_j + \frac{1}{m} b(\tau_j) \right\} \frac{\partial}{\partial \tau_j} = - \sum_{j=-\infty}^{\infty} \left\{ a(\tau_j) (t + \rho_j) + \frac{1}{m} b(\tau_j) \right\} \frac{\partial}{\partial \tau_j} \\ &= -tT_1 - \sum_{j=-\infty}^{\infty} \left\{ a(\tau_j) \rho_j + \frac{1}{m} b(\tau_j) \right\} \frac{\partial}{\partial \tau_j}, \end{aligned} \quad (38)$$

where

$$\rho_j = \begin{cases} -\tau - \tau_1 - \cdots - \tau_{j-1} & \text{if } j \geq 1 \\ 0 & \text{if } j = 0 \\ \tau_{-1} + \tau_{-2} + \cdots + \tau_j & \text{if } j \leq -1. \end{cases}$$

The following two lemmas are to be useful.

Lemma 6: *If the Lie algebra generated by the vector fields $S_0 = \sum_{j=-\infty}^{\infty} \partial / \partial w_j$ and $P = \sum_{j=-\infty}^{\infty} c(w_j) (\partial / \partial w_j)$ is of finite dimension, then $c(w)$ is one of the following forms:*

- (1) $c(w) = c_2 + c_3 e^{\lambda w} + c_4 e^{-\lambda w}$, $\lambda \neq 0$, and
- (2) $c(w) = c_2 + c_3 w + c_4 w^2$, where $c_2 - c_4$ are some constants.

Proof: Introduce vector fields

$$S_1 = [S_0, P], \quad S_2 = [S_0, S_1], \quad \dots, \quad S_n = [S_0, S_{n-1}], \quad n \geq 3.$$

Clearly, we have

$$S_n = \sum_{j=-\infty}^{\infty} c^{(n)}(w_j) \frac{\partial}{\partial w_j}, \quad n \geq 1. \quad (39)$$

Since all vector fields S_n are elements of L_x and L_x is of finite dimension, then there exists a natural number N such that

$$S_{N+1} = \mu_N S_N + \mu_{N-1} S_{N-1} + \cdots + \mu_1 S_1 + \mu_0 P + \mu S_0, \quad (40)$$

and S_0, P, S_1, \dots, S_N are linearly independent. (Note that we may assume that S_0 and P are linearly independent). Since $DS_0 D^{-1} = S_0$, $DPD^{-1} = P$, and $DS_n D^{-1} = S_n$ for any $n \geq 1$, then it follows from (40) that

$$S_{N+1} = D(\mu_N) S_N + D(\mu_{N-1}) S_{N-1} + \cdots + D(\mu_1) S_1 + D(\mu_0) P + D(\mu) S_0$$

and together with (40), it implies that $\mu, \mu_0, \mu_1, \dots, \mu_N$ are all constants.

By comparing the coefficients before $\partial / \partial w$ in (40) one gets, with the help of (39), the following equality:

$$c^{(N+1)}(w) = \mu_N c^{(N)}(w) + \cdots + \mu_1 c'(w) + \mu_0 c(w) + \mu.$$

Thus, $c(w)$ is a solution of the nonhomogeneous linear differential equation with constant coefficient whose characteristic polynomial is

$$\Lambda(\lambda) = \lambda^{N+1} - \mu_N \lambda^N - \cdots - \mu_1 \lambda - \mu_0.$$

Denote by $\beta_1, \beta_2, \dots, \beta_t$ the characteristic roots and by m_1, m_2, \dots, m_t their multiplicities. Following are the possibilities:

- (i) There exists a nonzero characteristic root, say, β_1 , and its multiplicity $m_1 \geq 2$.
- (ii) There exists zero characteristic root, say, β_1 , and $m_1 \geq 3$, $\mu = 0$ or $m_1 \geq 2$, $\mu \neq 0$.
- (iii) There are two distinct characteristic roots, say, β_1 and β_2 with $\beta_1 \neq 0$ and $\beta_2 = 0$.
- (iv) There are two nonzero distinct characteristic roots, say, β_1 and β_2 .

In case (i), consider

$$\Lambda_1(\lambda) = \frac{\Lambda(\lambda)}{\lambda - \beta_1} \quad \text{and} \quad \Lambda_1^{(2)}(\lambda) = \frac{\Lambda(\lambda)}{(\lambda - \beta_1)^2}.$$

Then $\Lambda_1(S_0)c(w) = \alpha_1 e^{\beta_1 w} + \alpha_2$ and $\Lambda_1^{(2)}(S_0)c(w) = (\alpha_3 w + \alpha_4)e^{\beta_1 w} + \alpha_5$, where α_j , $1 \leq j \leq 5$, are some constants with $\alpha_1 \neq 0$ and $\alpha_3 \neq 0$. We have

$$\begin{aligned} \Lambda_1(ad_{S_0})P &= \sum_{j=-\infty}^{\infty} (\alpha_1 e^{\beta_1 w_j} + \alpha_2) \frac{\partial}{\partial w_j} = \alpha_1 \left(\sum_{j=-\infty}^{\infty} e^{\beta_1 w_j} \frac{\partial}{\partial w_j} \right) + \alpha_2 S_0 = \alpha_1 P_1 + \alpha_2 S_0, \\ \Lambda_1^{(2)}(ad_{S_0})P &= \sum_{j=-\infty}^{\infty} ((\alpha_3 w_j + \alpha_4)e^{\beta_1 w_j} + \alpha_5) \frac{\partial}{\partial w_j} = \alpha_3 \left(\sum_{j=-\infty}^{\infty} w_j e^{\beta_1 w_j} \frac{\partial}{\partial w_j} \right) + \alpha_4 P_1 + \alpha_5 S_0 \\ &= \alpha_3 P_2 + \alpha_4 P_1 + \alpha_5 S_0 \end{aligned}$$

are elements from L_x and therefore vector fields $P_1 = \sum_{j=-\infty}^{\infty} e^{\beta_1 w_j} (\partial / \partial w_j)$ and $P_2 = \sum_{j=-\infty}^{\infty} w_j e^{\beta_1 w_j} (\partial / \partial w_j)$ belong to L_x . Since P_1 and P_2 generate an infinite dimensional Lie algebra L_x , then case (i) fails to be true.

In case (ii), consider

$$\Lambda_1^{(3)}(\lambda) = \frac{\Lambda(\lambda)}{\lambda^3} \quad \text{and} \quad \Lambda_1^{(2)}(\lambda) = \frac{\Lambda(\lambda)}{\lambda^2} \quad \text{if } \mu = 0$$

or

$$\Lambda_1^{(3)}(\lambda) = \frac{\Lambda(\lambda)}{\lambda^2} \quad \text{and} \quad \Lambda_1^{(2)}(\lambda) = \frac{\Lambda(\lambda)}{\lambda} \quad \text{if } \mu \neq 0.$$

We have

$$\Lambda_1^{(3)}(S_0)c(w) = \alpha_1 w^3 + \alpha_2 w^2 + \alpha_3 w + \alpha_4 \quad \text{and} \quad \Lambda_1^{(2)}(S_0)c(w) = \alpha_5 w^2 + \alpha_6 w + \alpha_7,$$

where α_j , $1 \leq j \leq 7$, are some constants with $\alpha_1 \neq 0$ and $\alpha_5 \neq 0$. Straightforward calculations show that vector fields

$$\begin{aligned} \Lambda_1^{(3)}(ad_{S_0})P &= \sum_{j=-\infty}^{\infty} (\alpha_1 w_j^3 + \alpha_2 w_j^2 + \alpha_3 w_j + \alpha_4) \frac{\partial}{\partial w_j} \quad \text{and} \quad \Lambda_1^{(2)}(ad_{S_0})P \\ &= \sum_{j=-\infty}^{\infty} (\alpha_5 w_j^2 + \alpha_6 w_j + \alpha_7) \frac{\partial}{\partial w_j} \end{aligned}$$

generate an infinite dimensional Lie algebra. It proves that case (ii) fails to be true.

In case (iii), consider

$$\Lambda_1(\lambda) = \frac{\Lambda(\lambda)}{\lambda - \beta_1} \quad \text{and} \quad \Lambda_2(\lambda) = \frac{\Lambda(\lambda)}{\lambda}.$$

We have

$$\Lambda_1 c(w) = \alpha_1 e^{\beta_1 w} + \alpha_2 \quad \text{and} \quad \Lambda_2 c(w) = \alpha_3 w + \alpha_4 \quad \text{if } \mu = 0$$

or

$$\Lambda_1(S_0)c(w) = \alpha_1 e^{\beta_1 w} + \alpha_2 \quad \text{and} \quad \Lambda_2(S_0)c(w) = \alpha_5 w^2 + \alpha_6 w + \alpha_7 \quad \text{if } \mu \neq 0,$$

where $\alpha_j, 1 \leq j \leq 7$, are constants with $\alpha_1 \neq 0, \alpha_3 \neq 0$, and $\alpha_5 \neq 0$. Since vector fields $\Lambda_1(ad_{S_0})P$ and $\Lambda_2(ad_{S_0})P$ generate an infinite dimensional Lie algebra, then case (iii) also fails to exist.

In case (iv), consider

$$\Lambda_1(\lambda) = \frac{\Lambda(\lambda)}{\lambda - \beta_1} \quad \text{and} \quad \Lambda_2(\lambda) = \frac{\Lambda(\lambda)}{\lambda - \beta_2}.$$

We have $\Lambda_1(S_0)c(w) = \alpha_1 e^{\beta_1 w} + \alpha_2$ and $\Lambda_2(S_0)c(w) = \alpha_3 e^{\beta_2 w} + \alpha_4$, where $\alpha_1 \neq 0, \alpha_2, \alpha_3 \neq 0$, and α_4 are some constants. Note that

$$\Lambda_1(ad_{S_0})P = \alpha_1 \left(\sum_{j=-\infty}^{\infty} e^{\beta_1 w_j} \frac{\partial}{\partial w_j} \right) + \alpha_2 S_0 \quad \text{and} \quad \Lambda_2(ad_{S_0})P = \alpha_3 \left(\sum_{j=-\infty}^{\infty} e^{\beta_2 w_j} \frac{\partial}{\partial w_j} \right) + \alpha_4 S_0,$$

and vector fields $\sum_{j=-\infty}^{\infty} e^{\beta_1 w_j} (\partial / \partial w_j)$ and $\sum_{j=-\infty}^{\infty} e^{\beta_2 w_j} (\partial / \partial w_j)$ generate an infinite dimensional Lie algebra if $\beta_1 + \beta_2 \neq 0$.

It follows from (i)–(iv) that $c(w)$ is one of the following forms:

- (1) $c(w) = c_2 + c_3 e^{\lambda w} + c_4 e^{-\lambda w}, \lambda \neq 0.$
- (2) $c(w) = c_2 + c_3 w + c_4 w^2$, where $c_2 - c_4$ are some constants.

□

Lemma 7: *If the Lie algebra generated by the vector fields $S_0 = \sum_{j=-\infty}^{\infty} \partial / \partial w_j, Q = \sum_{j=-\infty}^{\infty} q(w_j) (\partial / \partial w_j)$, and $S_1 = \sum_{j=-\infty}^{\infty} \{ \tilde{p}_j + \tilde{b}(w_j) \} (\partial / \partial w_j)$ is of finite dimension, then $q(w)$ is a constant function.*

Proof: It follows from Lemma 6 that

- (1) $q(w) = c_2 + c_3 w + c_4 w^2$ or
- (2) $q(w) = c_2 + c_3 e^{\lambda w} + c_4 e^{-\lambda w}, \lambda \neq 0,$
where $c_2 - c_4$ are some constants.

Consider case (1). We have

$$[S_0, Q] = c_3 \sum_{j=-\infty}^{\infty} \frac{\partial}{\partial w_j} + 2c_4 \sum_{j=-\infty}^{\infty} w_j \frac{\partial}{\partial w_j} = c_3 S_0 + 2c_4 \sum_{j=-\infty}^{\infty} w_j \frac{\partial}{\partial w_j}.$$

If $c_4 \neq 0$, then $\sum_{j=-\infty}^{\infty} w_j (\partial / \partial w_j) \in L_x$ and $\sum_{j=-\infty}^{\infty} w_j^2 (\partial / \partial w_j) \in L_x$.

If $c_4 = 0$ and $c_3 \neq 0$, then $\sum_{j=-\infty}^{\infty} w_j (\partial / \partial w_j) = (1/c_3)(Q - c_2 S_0) \in L_x$.

If $c_3 = c_4 = 0$, then $q(w) = c_2$ and there is nothing to prove.

Assume $c_4^2 + c_3^2 \neq 0$. Denote by $P = \sum_{j=-\infty}^{\infty} w_j (\partial / \partial w_j)$. Construct the vector fields

$$P_1 = [P, S_1], \quad P_n = [P, P_{n-1}], \quad n \geq 2.$$

We have

$$DS_0 D^{-1} = S_0,$$

$$DS_1 D^{-1} = S_1 - (e^w - \tilde{c}) S_0,$$

$$DPD^{-1} = P,$$

$$DP_1D^{-1} = P_1 + (-we^w + e^w - \bar{c})S_0,$$

$$DP_2D^{-1} = P_2 + (-w^2e^w + we^w - e^w + \bar{c})S_0.$$

In general,

$$DP_nD^{-1} = P_n + (-w^n e^w + R_{n-1}(w)e^w + c_n)S_0, \quad n \geq 3,$$

where R_{n-1} is a polynomial of degree $n-1$ and c_n is a constant. Since L_x is of finite dimension, then there exists a natural number N such that

$$P_{N+1} = \mu_N P_N + \cdots + \mu_1 P_1 + \mu_0 S_0,$$

and S_0, P_1, \dots, P_N are linearly independent. Thus

$$DP_{N+1}D^{-1} = D(\mu_N)DP_ND^{-1} + \cdots + D(\mu_1)DP_1D^{-1} + D(\mu_0)S_0,$$

or the same,

$$\begin{aligned} & \mu_N P_N + \cdots + \mu_1 P_1 + \mu_0 S_0 + (-w^{N+1}e^w + R_N(w)e^w + c_{N+1})S_0 \\ &= D(\mu_N)\{P_N + (-w^N e^w + R_{N-1}(w)e^w + c_N)S_0\} + \cdots + D(\mu_1)\{P_1 + (-we^w + e^w - \bar{c})S_0\} \\ & \quad + D(\mu_0)S_0. \end{aligned}$$

By comparing the coefficients before P_N, \dots, P_1 we have

$$\mu_N = D(\mu_N), \dots, \mu_1 = D(\mu_1),$$

which implies that μ_N, \dots, μ_1 are all constants. By comparing the coefficients before S_0 we have

$$\begin{aligned} \mu_0 - w^{N+1}e^w + R_N(w)e^w + c_{N+1} &= \mu_N(-w^N e^w + R_{N-1}(w)e^w + c_N) + \cdots + \mu_1(-we^w + e^w - \bar{c}) \\ & \quad + D(\mu_0). \end{aligned}$$

The last equality shows that $D(\mu_0) - \mu_0$ is a function of w only. Thus $D(\mu_0) - \mu_0$ is a constant; we denote it by d_0 . The last equality becomes a contradictory one:

$$w^{N+1}e^w = R_N(w)e^w + c_{N+1} - \mu_N(-w^N e^w + R_{N-1}(w)e^w + c_N) - \cdots - \mu_1(-we^w + e^w - \bar{c}) - d_0.$$

This contradiction proves that $c_3^2 + c_4^2 = 0$, i.e., $c_3 = c_4 = 0$ in case (1). Therefore, $q(w) = c_2$.

Consider case (2). Since

$$[S_0, Q] = \lambda c_3 \sum_{j=-\infty}^{\infty} e^{\lambda w_j} \frac{\partial}{\partial w_j} - \lambda c_4 \sum_{j=-\infty}^{\infty} e^{-\lambda w_j} \frac{\partial}{\partial w_j},$$

$$[S_0, [S_0, Q]] = \lambda^2 c_3 \sum_{j=-\infty}^{\infty} e^{\lambda w_j} \frac{\partial}{\partial w_j} + \lambda^2 c_4 \sum_{j=-\infty}^{\infty} e^{-\lambda w_j} \frac{\partial}{\partial w_j},$$

then vector fields $Q_\lambda = c_3 \sum_{j=-\infty}^{\infty} e^{\lambda w_j} (\partial / \partial w_j)$ and $Q_{-\lambda} = c_4 \sum_{j=-\infty}^{\infty} e^{-\lambda w_j} (\partial / \partial w_j)$ both belong to L_x . We have $DQ_\lambda D^{-1} = Q_\lambda$ and $DQ_{-\lambda} D^{-1} = Q_{-\lambda}$.

Assume $c_3 \neq 0$. Construct vector fields

$$Q_1 = [Q_\lambda, S_1], \quad Q_n = [Q_\lambda, Q_{n-1}], \quad n \geq 2.$$

Direct calculations show that

$$DQ_1 D^{-1} = Q_1 - c_3 e^{(1+\lambda)w} S_0 + (e^w - \bar{c}) \lambda Q_\lambda,$$

$$DQ_2D^{-1} = Q_2 - c_3^2(1 + \lambda)e^{(1+2\lambda)w}S_0 + 2\lambda c_3e^{(1+\lambda)w}Q_\lambda.$$

It can be proved by induction on n that

$$DQ_nQ^{-1} = Q_n - p_nS_0 + q_nQ_\lambda, \quad n \geq 2,$$

where

$$p_n = c_3^n(1 + \lambda)(1 + 2\lambda) \cdots (1 + (n - 1)\lambda)e^{(1+n\lambda)w},$$

$$q_n = nc_3^{n-1}\lambda(1 + \lambda) \cdots (1 + (n - 2)\lambda)e^{(1+(n-1)\lambda)w}.$$

Since L_x is of finite dimension, there exists such a natural number N that

$$Q_{N+1} = \mu_NQ_N + \cdots + \mu_1Q_1 + \mu_\lambda Q_\lambda + \mu_0S_0,$$

and $S_0, Q_\lambda, Q_1, \dots, Q_N$ are linearly independent. Then

$$DQ_{N+1}D^{-1} = D(\mu_N)DQ_ND^{-1} + \cdots + D(\mu_0)DS_0D^{-1}$$

or

$$\begin{aligned} \mu_NQ_N + \cdots + \mu_1Q_1 + \mu_\lambda Q_\lambda + \mu_0S_0 - p_{N+1}S_0 + q_{N+1}Q_\lambda &= D(\mu_N)\{Q_N - p_NS_0 + q_NQ_\lambda\} \\ &+ \cdots + D(\mu_1)\{Q_1 - p_1S_0 + q_1Q_\lambda\} + D(\mu_\lambda)Q_\lambda + D(\mu_0)S_0. \end{aligned}$$

By comparing the coefficients before Q_N, \dots, Q_1 , we have that $\mu_k, 1 \leq k \leq N$, are all constants. Comparing coefficients before S_0 gives

$$\mu_0 - p_{N+1} = -\mu_Np_N - \cdots - \mu_2p_2 - \mu_1p_1 + D(\mu_0). \tag{41}$$

Since $p_k, 1 \leq k \leq N+1$, depend on w only, then $D(\mu_0) - \mu_0$ is a function of w , and therefore $D(\mu_0) - \mu_0$ is a constant; we denote it by d_0 .

If $\lambda \neq -1/r$ for all $r \in \mathbb{N}$, then $p_k \neq 0$ for all $k \in \mathbb{N}$, and Eq. (41) fails to be true.

Consider the case when $\lambda = -1/r$ for some $r \in \mathbb{N}$. Substitution $u_j = e^{-\lambda w_j}$ transforms vector fields $(-1/\lambda c_3)Q_\lambda, (-1/\lambda)S_1$, and $(-1/\lambda)S_0$ into vector fields

$$Q_\lambda^* = \sum_{j=-\infty}^{\infty} \frac{\partial}{\partial u_j},$$

$$S_1^* = \sum_{j=-\infty}^{\infty} \{\tilde{\rho}_j^* + \tilde{b}^*(u_j)\}u_j \frac{\partial}{\partial u_j},$$

$$S_0^* = \sum_{j=-\infty}^{\infty} u_j \frac{\partial}{\partial u_j},$$

where

$$\tilde{\rho}_j^* = \begin{cases} \sum_{k=0}^{j-1} (u_k^r - \tilde{c}) & \text{if } j \geq 1 \\ 0 & \text{if } j = 0, \\ -\sum_{k=j}^{-1} (u_k^r - \tilde{c}) & \text{if } j \leq -1. \end{cases} \quad \tilde{b}^*(u_j) = \tilde{b}(r \ln u_j)$$

First consider the case $r=1$. We have

$$T := [Q_\lambda^*, S_1^*] = \sum_{j=-\infty}^{\infty} \{ju_j + \tilde{\rho}_j^* + \tilde{b}^*(u_j) + u_j \tilde{b}^{*'}(u_j)\} \frac{\partial}{\partial u_j},$$

$$K := \frac{1}{2}[Q_\lambda^*, T] = \sum_{j=-\infty}^{\infty} \{j + c(u_j)\} \frac{\partial}{\partial u_j},$$

where $c(u_j) = \tilde{b}^{*'}(u_j) + \frac{1}{2}u_j \tilde{b}^{*''}(u_j)$,

$$T_1 = [T, K] = \gamma_1 \sum_{j=-\infty}^{\infty} \{j^2 + jg_{1,1}^{(j)}(u_j) + g_{1,0}^{(j)}(u, u_1, \dots, u_j)\} \frac{\partial}{\partial u_j},$$

$$T_2 = [T, T_1] = \gamma_2 \sum_{j=-\infty}^{\infty} \{j^3 + j^2g_{2,2}^{(j)}(u_j) + jg_{2,1}^{(j)}(u, u_1, \dots, u_j) + g_{2,0}^{(j)}(u, u_1, \dots, u_j)\} \frac{\partial}{\partial u_j},$$

where $\gamma_1 = -\frac{3}{2}$ and $\gamma_2 \neq 0$.

Construct vector fields, $T_n = [T, T_{n-1}]$, $n \geq 3$. Direct calculations show that

$$T_n = \gamma_n \sum_{j=0}^{\infty} \left\{ j^{n+1} + j^n g_{n,n}(u_j) + \sum_{k=0}^{n-1} j^k g_{n,k}(u, u_1, \dots, u_j) \right\} \frac{\partial}{\partial u_j} + \sum_{j=-\infty}^{-1} a_j \frac{\partial}{\partial u_j}, \quad n \geq 1.$$

Since $\{T_n\}_{n=1}^{\infty}$ is an infinite sequence of linearly independent vector fields from L_x , then case $r = 1$ fails to exist.

Consider case $r \geq 2$. We have

$$ad_{Q_\lambda^*} S_1^* = [Q_\lambda^*, S_1^*] = \sum_{j=-\infty}^{\infty} \left\{ \text{sgn}(j)r \left(\sum_{k=0}^{j-1} u_k^{r-1} \right) u_j + \tilde{\rho}_j^* + \tilde{b}^*(u_j) + u_j \tilde{b}^{*'}(u_j) \right\} \frac{\partial}{\partial u_j}$$

and

$$ad_{Q_\lambda^*}^r S_1^* = \sum_{j=-\infty}^{\infty} \left\{ r! ju_j + \text{sgn}(j)r! \sum_{k=0}^{j-1} u_k + d(u_j) \right\}$$

for some function d ,

$$ad_{Q_\lambda^*}^{r+1} S_1^* = \sum_{j=-\infty}^{\infty} \{2r! j + d'(u_j)\} \frac{\partial}{\partial u_j}.$$

Note that vector fields $ad_{Q_\lambda^*}^r S_1^*$ and $ad_{Q_\lambda^*}^{r+1} S_1^*$ have coefficients of the same kind as vector fields T and K (from case $r=1$) have. It means that $ad_{Q_\lambda^*}^r S_1^*$ and $ad_{Q_\lambda^*}^{r+1} S_1^*$ generate an infinite dimensional Lie algebra. This contradiction implies that case $r \geq 2$ also fails to exist.

Thus, $c_3=0$. By interchanging λ with $-\lambda$, we obtain that $c_4=0$ also. Hence $c_3=c_4=0$ and $q(w)=c_2$. □

We already know that $a(\tau)=c_0\tau+c_1$. The next lemma shows that $c_0 \neq 0$.

Lemma 8: c_0 is a nonzero constant.

Proof: Assume the contrary. Then $a(\tau)=c_1$ and $c_1 \neq 0$. Vector fields (37) and (38) become

$$T_1 = c_1 \sum_{j=-\infty}^{\infty} \frac{\partial}{\partial \tau_j} = c_1 \tilde{T}_1$$

and

$$T = -tT_1 - c_1 \sum_{j=-\infty}^{\infty} \left\{ \rho_j + \frac{1}{mc_1} b(\tau_j) \right\} \frac{\partial}{\partial \tau_j} = -c_1 t \tilde{T}_1 - c_1 \tilde{T},$$

where

$$\tilde{T}_1 = \sum_{j=-\infty}^{\infty} \frac{\partial}{\partial \tau_j}, \quad \tilde{T} = \sum_{j=-\infty}^{\infty} \left\{ \rho_j + \frac{1}{mc_1} b(\tau_j) \right\} \frac{\partial}{\partial \tau_j}.$$

Since

$$[\tilde{T}_1, [\tilde{T}_1, \tilde{T}]] = \frac{1}{mc_1} \sum_{j=-\infty}^{\infty} b''(\tau_j) \frac{\partial}{\partial \tau_j}$$

and \tilde{T}_1 both belong to a finite dimensional L_x , then, by Lemma 6, (1) $b''(\tau) = \tilde{C}_1 + \tilde{C}_2 e^{\lambda\tau} + \tilde{C}_3 e^{-\lambda\tau}$ or (2) $b''(\tau) = \tilde{C}_1 + \tilde{C}_2 \tau + \tilde{C}_3 \tau^2$ for some constants $\tilde{C}_1 - \tilde{C}_3$.

In case (1), $b(\tau) = C_1 + C_2 e^{\lambda\tau} + C_3 e^{-\lambda\tau} + C_4 \tau^2 + C_5 \tau$ and

$$[\tilde{T}_1, [\tilde{T}_1, \tilde{T}]] - \lambda^2 \tilde{T} - \frac{2C_4 - \lambda^2 C_1}{mc_1} \tilde{T}_1 = -\lambda^2 \sum_{j=-\infty}^{\infty} \left\{ \rho_j + \frac{C_4 \tau_j^2 + C_5 \tau_j}{mc_1} \right\} \frac{\partial}{\partial \tau_j}$$

is an element in L_x .

In case (2), $b(\tau) = C_1 + C_2 \tau + C_3 \tau^2 + C_4 \tau^3 + C_5 \tau^4$ and

$$\tilde{T} - \frac{C_1}{mc_1} \tilde{T}_1 = \sum_{j=-\infty}^{\infty} \left\{ \rho_j + \frac{C_2 \tau_j + C_3 \tau_j^2 + C_4 \tau_j^3 + C_5 \tau_j^4}{mc_1} \right\} \frac{\partial}{\partial \tau_j}$$

belongs to L_x .

To finish the proof of the lemma it is enough to show that vector fields

$$\tilde{T}_2 := \sum_{j=-\infty}^{\infty} \{ \rho_j + C_2 \tau_j + C_3 \tau_j^2 + C_4 \tau_j^3 + C_5 \tau_j^4 \} \frac{\partial}{\partial \tau_j}$$

and

$$\tilde{T}_1 = \sum_{j=-\infty}^{\infty} \frac{\partial}{\partial \tau_j}$$

produce an infinite dimensional Lie algebra L_x for any fixed constants $C_2 - C_5$. One can prove it by showing that L_x contains vector fields $\sum_{j=-\infty}^{\infty} j^k (\partial / \partial \tau_j)$ for all $k = 1, 2, \dots$. Note that

$$[\tilde{T}_1, \tilde{T}_2] = \sum_{j=-\infty}^{\infty} (-j + C_2 + 2C_3 \tau_j + 3C_4 \tau_j^2 + 4C_5 \tau_j^3) \frac{\partial}{\partial \tau_j}.$$

There are four cases: (a) $C_5 \neq 0$, (b) $C_5 = 0, C_4 \neq 0$, (c) $C_5 = C_4 = 0, C_3 \neq 0$, and (d) $C_5 = C_4 = C_3 = 0$.

In case (a),

$$[\tilde{T}_1, [\tilde{T}_1, [\tilde{T}_1, \tilde{T}_2]]] - 6C_4 \tilde{T}_1 = \sum_{j=-\infty}^{\infty} 24C_5 \tau_j \frac{\partial}{\partial \tau_j} = 24C_5 P_1 \in L_x, \quad P_1 = \sum_{j=-\infty}^{\infty} \tau_j \frac{\partial}{\partial \tau_j},$$

$$[\tilde{T}_1, [\tilde{T}_1, \tilde{T}_2]] = \sum_{j=-\infty}^{\infty} \{2C_3 + 6C_4\tau_j + 12C_5\tau_j^2\} \frac{\partial}{\partial \tau_j} \in L_x,$$

and therefore,

$$P_2 := \sum_{j=-\infty}^{\infty} \tau_j^2 \frac{\partial}{\partial \tau_j} \in L_x$$

and

$$\tilde{T}_3 := [\tilde{T}_1, \tilde{T}_2] - C_2\tilde{T}_1 - 2C_3P_1 - 3C_4P_2 = \sum_{j=-\infty}^{\infty} (-j + 4C_5\tau_j^3) \frac{\partial}{\partial \tau_j} \in L_x.$$

We have

$$J_1 := -\frac{1}{3}([\tilde{T}_3, P_1] + 2\tilde{T}_3) = \sum_{j=-\infty}^{\infty} j \frac{\partial}{\partial \tau_j} \in L_x.$$

Now,

$$[J_1, [J_1, P_2]] = \frac{1}{2} \sum_{j=-\infty}^{\infty} j^2 \frac{\partial}{\partial \tau_j} \in L_x.$$

Assuming $J_k = \sum_{j=-\infty}^{\infty} j^k (\partial / \partial \tau_j) \in L_x$ we have that

$$J_{k+1} := \frac{1}{2}[J_1, [J_k, P_2]] = \sum_{j=-\infty}^{\infty} j^{k+1} \frac{\partial}{\partial \tau_j} \in L_x.$$

In case (b) we have

$$P_1 := \frac{1}{6C_4} \{[\tilde{T}_1, [\tilde{T}_1, \tilde{T}_2]] - 2C_3T_1\} = \sum_{j=-\infty}^{\infty} \tau_j \frac{\partial}{\partial \tau_j} \in L_x$$

and

$$\tilde{T}_3 = [\tilde{T}_1, \tilde{T}_2] - C_2\tilde{T}_1 - 2C_3P_1 = \sum_{j=-\infty}^{\infty} (-j + 3C_4\tau_j^2) \frac{\partial}{\partial \tau_j} \in L_x.$$

We have

$$J_1 := -\frac{1}{2}([\tilde{T}_3, P_1] + \tilde{T}_3) = \sum_{j=-\infty}^{\infty} j \frac{\partial}{\partial \tau_j} \in L_x$$

and

$$P_2 = \frac{1}{6C_4} (\tilde{T}_3 - [\tilde{T}_3, P_1]) = \sum_{j=-\infty}^{\infty} \tau_j^2 \frac{\partial}{\partial \tau_j} \in L_x.$$

As it was shown in the proof of case (a), J_1 and P_2 produce an infinite dimensional Lie algebra.

In case (c),

$$\tilde{T}_3 = [\tilde{T}_1, \tilde{T}_2] - C_2\tilde{T}_1 = \sum_{j=-\infty}^{\infty} (-j + 2C_3\tau_j) \frac{\partial}{\partial \tau_j} \in L_x,$$

$$\tilde{T}_4 = [\tilde{T}_3, \tilde{T}_2] = \sum_{j=-\infty}^{\infty} \left(\frac{j(j-1)}{2} - jC_2 - 2C_3j\tau_j + 2C_3^2\tau_j^2 \right) \frac{\partial}{\partial \tau_j} \in L_x.$$

Also,

$$\tilde{T}_5 = [\tilde{T}_3, \tilde{T}_4] = 2C_3 \sum_{j=-\infty}^{\infty} \left(\frac{j(j+1)}{2} + C_2j - 2C_3j\tau_j + 2C_3^2\tau_j^2 \right) \frac{\partial}{\partial \tau_j} \in L_x.$$

Since \tilde{T}_4 and \tilde{T}_5 both belong to L_x , then either

(i)

$$J_1 = \sum_{j=-\infty}^{\infty} j \frac{\partial}{\partial \tau_j} \in L_x, \quad \tilde{T}_6 = \sum_{j=-\infty}^{\infty} \left(\frac{j^2}{2} - 2C_3j\tau_j + 2C_3^2\tau_j^2 \right) \frac{\partial}{\partial \tau_j} \in L_x$$

or

(ii)

$$C_2 = -\frac{1}{2}, \quad \tilde{T}_6 = \sum_{j=-\infty}^{\infty} \left(\frac{j^2}{2} - 2C_3j\tau_j + 2C_3^2\tau_j^2 \right) \frac{\partial}{\partial \tau_j} \in L_x.$$

In case (c) (i),

$$P_1 = \frac{1}{4C_3^2} \{ [\tilde{T}_1, \tilde{T}_6] + 2C_3J_1 \} = \sum_{j=-\infty}^{\infty} \tau_j \frac{\partial}{\partial \tau_j} \in L_x.$$

Since

$$[P_1, \tilde{T}_6] = \sum_{j=-\infty}^{\infty} \left(-\frac{j^2}{2} + 2C_3^2\tau_j^2 \right) \frac{\partial}{\partial \tau_j}$$

and

$$[P_1, [P_1, \tilde{T}_6]] = \sum_{j=-\infty}^{\infty} \left(\frac{j^2}{2} + 2C_3^2\tau_j^2 \right) \frac{\partial}{\partial \tau_j}$$

both belong to L_x , then

$$J_2 = \sum_{j=-\infty}^{\infty} j^2 \frac{\partial}{\partial \tau_j} \in L_x, \quad P_2 = \sum_{j=-\infty}^{\infty} \tau_j^2 \frac{\partial}{\partial \tau_j} \in L_x.$$

P_2 and J_1 generate an infinite dimensional Lie algebra.

In case (c) (ii),

$$\tilde{T}_1 = \sum_{j=-\infty}^{\infty} \frac{\partial}{\partial \tau_j}, \quad \tilde{T}_2 = \sum_{j=-\infty}^{\infty} \left(C_3\tau_j^2 - \frac{1}{2}\tau_j + \rho_j \right) \frac{\partial}{\partial \tau_j}.$$

Note that the Lie algebra generated by the vector fields

$$\tilde{T}_2^* = \tilde{T}_2 - \left(C_3\tau^2 - \frac{1}{2}\tau \right) \tilde{T}_1 = d(\tau, \tau_1) \frac{\partial}{\partial \tau_1} - d(\tau_{-1}, \tau) \frac{\partial}{\partial \tau_{-1}} + (d(\tau, \tau_1) + d(\tau_1, \tau_2)) \frac{\partial}{\partial \tau_2} + \dots$$

and

$$\tilde{T}_1 = \sum_{j=-\infty}^{\infty} \frac{\partial}{\partial \tau_j}$$

is infinite dimensional. It can be proved by comparing this algebra with the infinite dimensional characteristic Lie algebra of the chain

$$t_{1x} = t_x + C_3(t_1^2 - t^2) - \frac{1}{2}(t_1 + t). \quad (42)$$

Indeed, the Lie algebra L_{x1} for (42) is generated by operators (18) and (19) with $d(t, t_1) = C_3(t_1^2 - t^2) - \frac{1}{2}(t_1 + t)$. To keep standard notations we set $a(\tau) = -2C_3\tau - 1$ and $b(\tau) = C_3\tau^2 + \frac{1}{2}\tau$. Note that since $C_3 \neq 0$, function $a(\tau)$ is not a constant. It follows from Theorem 3 proved below that the characteristic Lie algebras L_x (and therefore algebra L_{x1}) for Eq. (42) is of infinite dimension. Thus, in case (c) (ii) we also have an infinite dimensional Lie algebra L_x .

In case (d),

$$\tilde{T}_2 = \sum_{j=-\infty}^{\infty} (-\tau - \tau_1 - \dots - \tau_{j-1} + C_2\tau_j) \frac{\partial}{\partial \tau_j} \in L_x.$$

Then

$$J_1 = c_2\tilde{T}_1 - [\tilde{T}_1, \tilde{T}_2] = \sum_{j=-\infty}^{\infty} j \frac{\partial}{\partial \tau_j} \in L_x$$

and

$$J_2 = -2 \left([J_1, \tilde{T}_2] - \left(\frac{1}{2} + C_2 \right) J_1 \right) = \sum_{j=-\infty}^{\infty} j^2 \frac{\partial}{\partial \tau_j} \in L_x.$$

Assuming that J_k , $1 \leq k \leq n$, belong to L_x , by considering $[J_n, \tilde{T}_2]$ one may show that $J_{n+1} = \sum_{j=-\infty}^{\infty} j^{k+1} (\partial / \partial \tau_j) \in L_x$. It implies that L_x is of infinite dimension. \square

Let us introduce new variables,

$$w_j = \ln \left(\tau_j + \frac{c_1}{c_0} \right).$$

Vector fields T_1 and T in variables w_j can be rewritten as

$$T_1 = c_0 \sum_{j=-\infty}^{\infty} \frac{\partial}{\partial w_j} = c_0 S_0,$$

$$T = -tc_0 S_0 + c_0 \sum_{j=-\infty}^{\infty} \{ \tilde{\rho}_j + \tilde{b}(w_j) \} \frac{\partial}{\partial w_j} = -c_0 t S_0 + c_0 S_1,$$

where

$$S_0 = \sum_{j=-\infty}^{\infty} \frac{\partial}{\partial w_j}, \quad S_1 = \sum_{j=-\infty}^{\infty} \{ \tilde{\rho}_j + \tilde{b}(w_j) \} \frac{\partial}{\partial w_j},$$

$$\tilde{\rho}_j = \begin{cases} \sum_{k=0}^{j-1} (e^{w_k} - \tilde{c}) & \text{if } j \geq 1 \\ 0 & \text{if } j = 0, \\ -\sum_{k=j}^{-1} (e^{w_k} - \tilde{c}) & \text{if } j \leq -1. \end{cases} \quad \tilde{c} = \frac{c_1}{c_0}, \quad \tilde{b}(w_j) = -\frac{1}{m} \left(\frac{b(\tau_j)}{c_0 \tau_j + c_1} \right)$$

We have

$$DS_0 D^{-1} = S_0, \quad DS_1 D^{-1} = S_1 - (e^w - \tilde{c})S_0.$$

These lemmas allow one to prove the following theorem.

Theorem 4: *If the equation*

$$t_{1x} = t_x + a(\tau)t^m + b(\tau)t^{m-1} + \dots, \quad m \geq 1,$$

admits a nontrivial x -integral, then

- (1) $a(\tau) = c_0 \tau$ and $b(\tau) = c_2 \tau^2 + c_3 \tau$, where c_0, c_2 , and c_3 are some constants.
- (2) $m = 1$.

Proof: Consider case (1). Define vector field

$$Q = [S_0, [S_0, S_1]] - [S_0, S_1] = \sum_{j=-\infty}^{\infty} (\tilde{b}''(w_j) - \tilde{b}'(w_j)) \frac{\partial}{\partial w_j}.$$

By Lemma 7, $\tilde{b}''(w) - \tilde{b}'(w) = C$ for some constant C . Thus, $\tilde{b}(w) = C_0 + C_1 e^w + C_2 w$ for some constants C_1, C_2 , and C_0 . Consider vector fields

$$P = (C_2 - C_0)S_0 + S_1 - [S_0, S_1] = \sum_{j=-\infty}^{\infty} (C_2 w_j + \tilde{c} j) \frac{\partial}{\partial w_j},$$

$$R = [S_0, [S_0, S_1]] = \sum_{j=1}^{\infty} \left\{ \left(\sum_{k=1}^j e^{w_k} \right) + C_1 e^{w_j} \right\} \frac{\partial}{\partial w_j} + C_1 e^w \frac{\partial}{\partial w} - \sum_{j=-\infty}^{-1} \left\{ \left(\sum_{k=j}^{-1} e^{w_k} \right) + C_1 e^{w_j} \right\} \frac{\partial}{\partial w_j},$$

$$R_1 = [P, R], \quad R_{n+1} = [P, R_n], \quad n \geq 1.$$

Then

$$R_n = \sum_{j \geq 0} \{ e^{w_j} (C_1 C_2^n w_j^n + P_{n,j}) + r_{n,j}(w, w_1, \dots, w_{j-1}) \} \frac{\partial}{\partial w_j} + \sum_{j \leq -1} \{ e^{w_j} ((C_1 - 1) C_2^n w_j^n + P_{n,j}) + r_{n,j}(w_{-1}, w_{-2}, \dots, w_{j+1}) \} \frac{\partial}{\partial w_j},$$

where $P_{n,j} = P_{n,j}(w_j, j)$ is a polynomial of degree $n - 1$ whose coefficients depend on j ; $r_{n,j}$ are the functions that do not depend on w_j . Since all vector fields R_n belong to a finite dimensional Lie algebra L_x , then $C_1 C_2 = (C_1 - 1) C_2 = 0$, or the same, $C_2 = 0$. Therefore,

$$\tilde{b}(w) = C_0 + C_1 e^w.$$

Since $C_2 = 0$, then

$$P = \tilde{c} \sum_{j=-\infty}^{\infty} j \frac{\partial}{\partial w_j},$$

$$R = \sum_{j=1}^{\infty} \left\{ \binom{j}{k=1} e^{w_k} + C_1 e^{w_j} \right\} \frac{\partial}{\partial w_j} + C_1 e^w \frac{\partial}{\partial w} - \sum_{j=-\infty}^{-1} \left\{ \binom{-1}{k=j} e^{w_k} + C_1 e^{w_j} \right\} \frac{\partial}{\partial w_j}$$

and

$$R_n = \tilde{c}^n \sum_{j=1}^{\infty} \{ e^{w_1} + 2^n e^{w_2} + (j-1)^n e^{w_{j-1}} + j^n C_1 e^{w_j} \} \frac{\partial}{\partial w_j} \\ - \tilde{c}^n \sum_{j=-\infty}^{-1} \{ (-1)^n e^{w_{-1}} + (-2)^n e^{w_{-2}} + (j)^n e^{w_j} + j^n C_1 e^{w_j} \} \frac{\partial}{\partial w_j}.$$

Again, vector fields R_n belong to a finite dimensional Lie algebra only if $\tilde{c}=0$, or the same, $c_1=0$. It implies that

$$a(\tau) = c_0 \tau, \quad b(\tau) = c_2 \tau^2 + c_3 \tau.$$

Consider case (2). Assume the contrary, that is, $m \geq 2$. Then the following vector field:

$$\frac{1}{m!} ad_X^{m-2}(\tilde{Y}) = Y_{(1/2)a(\tau)t^2 + (1/m)b(\tau)t + \{1/[m(m-1)]\}c(\tau)} = - \sum_{j=-\infty}^{\infty} \left(\frac{1}{2} a(\tau_j) t_j^2 + \frac{1}{m} b(\tau_j) t_j + \frac{1}{m(m-1)} c(\tau_j) \right) \frac{\partial}{\partial \tau_j} \\ = - \sum_{j=-\infty}^{\infty} \left(\frac{1}{2} a(\tau_j) (t + \rho_j)^2 + \frac{1}{m} b(\tau_j) (t + \rho_j) + \frac{1}{m(m-1)} c(\tau_j) \right) \frac{\partial}{\partial \tau_j} - \frac{t^2}{2} \sum_{j=-\infty}^{\infty} a(\tau_j) \frac{\partial}{\partial \tau_j} \\ - t \sum_{j=-\infty}^{\infty} \left\{ a(\tau_j) \rho_j + \frac{1}{m} b(\tau_j) \right\} \frac{\partial}{\partial \tau_j} - \sum_{j=-\infty}^{\infty} \left\{ \frac{1}{2} a(\tau_j) \rho_j^2 + \frac{1}{m} b(\tau_j) + \frac{1}{m(m-1)} c(\tau_j) \right\} \frac{\partial}{\partial \tau_j},$$

is in L_X . In variables $w_j = \ln \tau_j$,

$$\frac{1}{m!} ad_X^{m-2}(\tilde{Y}) = - \frac{t^2}{2} c_0 S_0 + t c_0 S_1 - c_0 S_2,$$

where

$$S_2 = \sum_{j=-\infty}^{\infty} \left\{ \frac{1}{2} \tilde{\rho}_j^2 - \tilde{b}(w_j) \tilde{\rho}_j + \tilde{c}(w_j) \right\} \frac{\partial}{\partial w_j}, \quad \tilde{c}(w_j) = \frac{c(\tau_j)}{m(m-1)\tau_j}.$$

The vector fields S_0 and S_1 are as in Lemma 7. We have

$$[S_0, S_2] = 2S_2 + C_0 S_1 + P, \quad P = \sum_{j=-\infty}^{\infty} r(w_j) \frac{\partial}{\partial w_j}, \quad r(w) = \tilde{c}'(w) - 2\tilde{c}(w) - C_0 \tilde{b}(w).$$

Construct the sequence

$$S_3 = [S_1, S_2], \quad S_{n+1} = [S_1, S_n], \quad n \geq 2.$$

One can prove by induction on n that

$$[S_0, S_n] = nS_n + \sum_{k=0}^{n-1} \nu_{n,k} S_k$$

and

$$DS_n D^{-1} = S_n + \left\{ \frac{n(n-1)}{2} - 1 \right\} e^w S_{n-1} + \sum_{k=0}^{n-2} \eta(n,k) S_k, \quad n \geq 3.$$

Since L_x is of finite dimension, then there exists a natural number N such that

$$S_{N+1} = \mu_N S_N + \mu_{N-1} S_{N-1} + \dots + \mu_0 S_0.$$

Then

$$DS_{N+1} D^{-1} = D(\mu_N) DS_N D^{-1} + D(\mu_{N-1}) DS_{N-1} D^{-1} + \dots + D(\mu_0) DS_0 D^{-1}.$$

On the other hand,

$$DS_{N+1} D^{-1} = S_{N+1} + \left\{ \frac{(N+1)N}{2} - 1 \right\} e^w S_N + \dots.$$

We compare the coefficients before S_N and have two equations:

$$D(\mu_N) = \mu_N + \left\{ \frac{(N+1)N}{2} - 1 \right\} e^w, \quad N \geq 2,$$

and

$$D(\mu_1) = \mu_1 + e^w, \quad N = 1.$$

Both equations are contradictory. Therefore, our assumption that $m \geq 2$ was wrong. □

IV. NONZERO ROOT

Lemma 9: Assume that Eq. (12) admits a nontrivial x -integral. Then the characteristic polynomial of Eq. (25) can only have simple nonzero roots.

Proof: Assume that $m_1 \geq 2$. Introduce polynomials

$$\Lambda_{\alpha_1}^{(2)}(\lambda) = \frac{\Lambda(\lambda)}{(\lambda - \alpha_1)^2}, \quad \Lambda_{\alpha_1}(\lambda) = \frac{\Lambda(\lambda)}{(\lambda - \alpha_1)}.$$

Consider vector fields

$$S_0^* = \Lambda_{\alpha_1}^{(2)}(ad_{\tilde{X}}) Y_d = Y_{A(\tau)e^{\alpha_1 t}},$$

$$S_1^* = \Lambda_{\alpha_1}(ad_{\tilde{X}}) Y_d = Y_{(A(\tau)t + B(\tau))e^{\alpha_1 t}}$$

from the Lie algebra L_x .

In variables $\tau_j = t_j - t_{j+1}$, vector fields S_0^* and S_1^* become

$$S_0^* = -e^{\alpha_1 t} \sum_{j=-\infty}^{\infty} A(\tau_j) e^{\alpha_1 \rho_j} \frac{\partial}{\partial \tau_j} = -e^{\alpha_1 t} S_0,$$

$$S_1^* = -te^{\alpha_1 t} S_0 - e^{\alpha_1 t} \sum_{j=-\infty}^{\infty} \{A(\tau_j)\rho_j + B(\tau_j)\} e^{\alpha_1 \rho_j} \frac{\partial}{\partial \tau_j} = -te^{\alpha_1 t} S_0 - e^{\alpha_1 t} S_1,$$

with $S_0 = \sum_{j=-\infty}^{\infty} A(\tau_j) e^{\alpha_1 \rho_j} (\partial / \partial \tau_j)$ and $S_1 = \sum_{j=-\infty}^{\infty} \{A(\tau_j)\rho_j + B(\tau_j)\} e^{\alpha_1 \rho_j} (\partial / \partial \tau_j)$.

Direct calculations show that

$$DS_0 D^{-1} = e^{\alpha_1 \tau} S_0, \quad DS_1 D^{-1} = e^{\alpha_1 \tau} S_1 + \tau e^{\alpha_1 \tau} S_0.$$

Define the sequence

$$S_2 = [S_0, S_1], \quad S_{n+1} = [S_0, S_n], \quad n \geq 2.$$

One can easily show that

$$DS_2 D^{-1} = e^{2\alpha_1 \tau} S_2 + \alpha_1 e^{2\alpha_1 \tau} A(\tau) S_1 + e^{2\alpha_1 \tau} (A(\tau) - \alpha_1 B(\tau)) S_0.$$

It can be proved by induction on n that

$$DS_n D^{-1} = e^{n\alpha_1 \tau} S_n + \alpha_1 \frac{n(n-1)}{2} e^{n\alpha_1 \tau} A(\tau) S_{n-1} + \sum_{k=0}^{n-2} \gamma(n, k) S_k.$$

Since the dimension of L_x is finite and S_0, S_1, \dots are elements of L_x , then there exists a natural number N such that

$$S_{N+1} = \mu_N S_N + \mu_{N-1} S_{N-1} + \dots + \mu_0 S_0,$$

and S_0, S_1, \dots, S_N are linearly independent. Therefore,

$$DS_{N+1} D^{-1} = D(\mu_N) DS_N D^{-1} + D(\mu_{N-1}) DS_{N-1} D^{-1} + \dots + D(\mu_0) DS_0 D^{-1}.$$

On the other hand,

$$DS_{N+1} D^{-1} = e^{(N+1)\alpha_1 \tau} S_{N+1} + \alpha_1 \frac{(N+1)N}{2} e^{(N+1)\alpha_1 \tau} A(\tau) S_N + \sum_{k=0}^{N-1} \gamma(N+1, k) S_k.$$

By comparing the coefficients before S_N in the last two equations we have

$$e^{(N+1)\alpha_1 \tau} \mu_N + \frac{\alpha_1 (N+1)N}{2} e^{(N+1)\alpha_1 \tau} A(\tau) = D(\mu_N) e^{N\alpha_1 \tau}.$$

It follows at once that μ_N is a constant and then

$$A(\tau) = C(e^{-\alpha_1 \tau} - 1), \quad C = \frac{2\mu_N}{\alpha_1 N(N+1)}.$$

Let us construct a new infinite sequence of vector fields belonging to L_x , enumerated by a multi-index,

$$T_0 := S_1, \quad T_1 := S_0, \quad T_2 = [S_1, T_1], \quad T_{n+1} = [S_1, T_n], \quad n \geq 2, \quad T_{n,0} = [S_0, T_n],$$

$$T_{n,0,i_1,\dots,i_{n-1},i_n} = [S_{i_n}, T_{n,0,i_1,\dots,i_{n-1}}], \quad i_j \in \{0; 1\}.$$

Direct calculations show that

$$DT_2 D^{-1} = e^{2\alpha_1 \tau} T_2 + e^{2\alpha_1 \tau} (\alpha_1 B - A) T_1 - \alpha_1 e^{2\alpha_1 \tau} A T_0,$$

$$DT_3D^{-1} = e^{3\alpha_1\tau}T_3 + e^{3\alpha_1\tau}(3\alpha_1B - A + 3\alpha_1\tau A)T_2 + \tau e^{3\alpha_1\tau}T_{2,0} + \sum_{m(\beta)<2} \nu(3,\beta)T_\beta.$$

Here and below we use functions $m=m(\beta)$ and $l=l(\beta)$ defined in Sec. III. It can be proved by induction on n that

$$DT_nD^{-1} = e^{n\alpha_1\tau}T_n + e^{n\alpha_1\tau}\{c_nB - A + c_n\tau A\}T_{n-1} + \tau e^{n\alpha_1\tau} \sum_{m(\beta)=n-1, l(\beta)=1} \nu^*(n,\beta)T_\beta + \sum_{m(\beta)\leq n-2} \nu(n,\beta)T_\beta,$$

where

$$c_n = \frac{\alpha_1 n(n-1)}{2},$$

and $\nu^*(n,\beta)$ are constants for any β with $m(\beta)=n-1$ and $l(\beta)=1$.

In general, for any γ ,

$$DT_\gamma D^{-1} = e^{(m(\gamma)+l(\gamma))\alpha_1\tau}T_\gamma + \sum_{m(\beta)\leq m(\gamma)-1} \nu(\gamma,\beta)T_\beta.$$

Among the vector fields T_β we choose a system P of linearly independent vector fields in such a way that for some natural number N , we have the following:

- (i) $T_k \in P, k \leq N$.
- (ii) $m(\beta) \leq N$ for any $T_\beta \in P$.
- (iii) For any T_γ with $m(\gamma) \leq N$ we have $T_\gamma = \sum_{T_\beta \in P, m(\beta) \leq m(\gamma)} \mu(\gamma,\beta)T_\beta$. Also

$$T_{N+1} = \mu(N+1,N)T_N + \sum_{T_\beta \in P} \mu(N+1,\beta)T_\beta.$$

- (iv) For any $T_\gamma \notin P$ with $m(\gamma)=N$ and $l(\gamma)=1$, we have $\mu(\gamma,N)=0$.

Indeed,

$$DT_\gamma D^{-1} = D(\mu(\gamma,N))DT_N D^{-1} + \sum_{T_\beta \in P, \beta \neq N} D(\mu(\gamma,\beta))DT_\beta D^{-1}.$$

On the other hand,

$$\begin{aligned} DT_\gamma D^{-1} &= e^{(m(\gamma)+l(\gamma))\alpha_1\tau}T_\gamma + \sum_{m(\beta)\leq N-1} \nu(\gamma,\beta)T_\beta \\ &= e^{(N+1)\alpha_1\tau} \left\{ \mu(\gamma,N)T_N + \sum_{T_\beta \in P, m(\beta)\leq N, \beta \neq N} \mu(\gamma,\beta)T_\beta \right\} + \sum_{m(\beta)\leq N-1} \nu(\gamma,\beta)T_\beta. \end{aligned}$$

By comparing the coefficients before T_N we have

$$e^{(N+1)\alpha_1\tau}\mu(\gamma,N) = D(\mu(\gamma,N))e^{N\alpha_1\tau},$$

which proves $\mu(\gamma,N)=0$ for any γ with $m(\gamma)=N$ and $l(\gamma)=1$. We have

$$T_{N+1} = \mu_N T_N + \sum_{T_\beta \in P} \mu(N+1,\beta)T_\beta.$$

Here $\mu_N = \mu(N+1,N)$. Then

$$DT_{N+1}D^{-1} = D(\mu_N)DT_N D^{-1} + \sum_{T_\beta \in P} D(\mu(N+1,\beta))DT_\beta D^{-1}.$$

We continue and have

$$\begin{aligned}
 & e^{(N+1)\alpha_1\tau} \left\{ \mu_N T_N + \sum_{T_\beta \in P} \mu(N+1, \beta) T_\beta \right\} + e^{(N+1)\alpha_1\tau} \{ c_{N+1} B - A + c_{N+1} \tau A \} T_N \\
 & + \tau e^{(N+1)\alpha_1\tau} \sum_{m(\beta)=N, l(\beta)=1} v^*(N+1, \beta) T_\beta + \sum_{m(\beta) \leq N-1} v(N+1, \beta) T_\beta \\
 & = D(\mu_N) \left\{ e^{N\alpha_1\tau} T_N + \sum_{m(\beta) \leq N-1} v(N, \beta) T_\beta \right\} \\
 & + \sum_{T_\beta \in P} D(\mu(N+1, \beta)) \left\{ e^{(m(\beta)+l(\beta))\alpha_1\tau} T_\beta + \sum_{m(r) \leq N-1} v(\beta, r) T_r \right\}.
 \end{aligned}$$

We compare the coefficients before T_N and get

$$e^{(N+1)\alpha_1\tau} \mu_N + e^{(N+1)\alpha_1\tau} \{ c_{N+1} B - A + c_{N+1} \tau A \} = e^{N\alpha_1\tau} D(\mu_N).$$

Note that, by property (iv), we do not have term $\tau e^{(N+1)\alpha_1\tau}$ in the left side of the last equality. Thus, using the expression for $A(\tau) = C(e^{-\alpha_1\tau} - 1)$ and the fact that μ_N is a constant, we have

$$B(\tau) = C_1 A + C_2 \tau A = C_1 (e^{-\alpha_1\tau} - 1) + C_2 \tau (e^{-\alpha_1\tau} - 1),$$

where

$$C_1 = \frac{\mu_N}{C c_{N+1}} + \frac{1}{c_{N+1}}, \quad C_2 = -1.$$

We introduce new vector fields

$$\tilde{S}_0 = \frac{1}{C} S_0 = (e^{-\alpha_1\tau} - 1) \frac{\partial}{\partial \tau} + \dots, \quad \tilde{S}_1 = \frac{1}{C} S_1 + \frac{C_1}{C} S_0 = \tau (e^{-\alpha_1\tau} - 1) \frac{\partial}{\partial \tau} + \dots.$$

$$\tilde{S}_2 = [\tilde{S}_0, \tilde{S}_1], \quad \tilde{S}_{n+1} = [\tilde{S}_0, \tilde{S}_n], \quad n \geq 2.$$

We have

$$D\tilde{S}_0 D^{-1} = e^{\alpha_1\tau} \tilde{S}_0, \quad D\tilde{S}_1 D^{-1} = e^{\alpha_1\tau} \tilde{S}_1 - \tau e^{\alpha_1\tau} \tilde{S}_0,$$

$$D\tilde{S}_n D^{-1} = \sum_{k=0}^n \tilde{\gamma}(n, k) \tilde{S}_k, \quad \tilde{\gamma}(n, n) = e^{n\alpha_1\tau},$$

where $\tilde{\gamma}(n, k)$ are functions of τ only. Since all vector fields \tilde{S}_k belong to a finite dimensional Lie algebra L_x , then there exists such a natural number M that

$$\tilde{S}_{M+1} = \tilde{\mu}_M \tilde{S}_M + \dots + \tilde{\mu}_0 \tilde{S}_0, \tag{43}$$

and $\tilde{S}_0, \dots, \tilde{S}_M$ are linearly independent. Then

$$D\tilde{S}_{M+1} D^{-1} = D(\tilde{\mu}_M) D\tilde{S}_M D^{-1} + \dots + D(\tilde{\mu}_0) D\tilde{S}_0 D^{-1}$$

and

$$\tilde{\gamma}(M+1, M+1) \{ \tilde{\mu}_M \tilde{S}_M + \dots + \tilde{\mu}_0 \tilde{S}_0 \} + \sum_{k=0}^M \tilde{\gamma}(M+1, k) \tilde{S}_k = D(\tilde{\mu}_N) \{ \tilde{\gamma}(M, M) \tilde{S}_M + \dots \} + \dots.$$

By comparing the coefficients before \tilde{S}_M , we have

$$e^{(M+1)\alpha_1\tau}\tilde{\mu}_M + \tilde{\gamma}(M+1, M) = D(\tilde{\mu}_M)e^{M\alpha_1\tau},$$

which implies that $\tilde{\mu}_M$ is a constant. In the same way, by comparing the coefficients before \tilde{S}_{M-1} and then before \tilde{S}_{M-2} and so on, one can show that all coefficients $\tilde{\mu}_k$ are constants.

One can show by induction on n that for $n \geq 2$,

$$\tilde{S}_n = \left\{ \alpha_1^{n-2}(-1)^{n-2}(n-2)! e^{-n\alpha_1\tau} + \sum_{k=0}^{n-1} r(n, k)e^{-\alpha_1 k\tau} \right\} \frac{\partial}{\partial \tau} + \dots,$$

where $r(n, k)$ are some constants. Return to equality (43) with constant coefficients $\tilde{\mu}_k$ and compare the coefficients before $\partial/\partial\tau$:

$$\begin{aligned} &\alpha_1^{M-1}(-1)^{M-1}(M-1)! e^{-(M+1)\alpha_1\tau} + \sum_{k=0}^M r(M+1, k)e^{-\alpha_1 k\tau} = \tilde{\mu}_M(\alpha_1^{M-2}(-1)^{M-2}(M-2)! e^{-M\alpha_1\tau} \\ &+ \sum_{k=0}^{M-1} r(M, k)e^{-\alpha_1 k\tau}) + \dots + \tilde{\mu}_0(e^{-\alpha_1\tau} - 1). \end{aligned}$$

The last equality fails to be true. It shows that our assumption that multiplicity m_1 of a nonzero root α_1 can be 2 or more was wrong. \square

If the characteristic polynomial of (25) has only one nonzero root α , then $d(t, t_1) = A(t - t_1)e^{\alpha t}$. In this case Eq. (12) admits a nontrivial x -integral (see Sec. I, Theorem 3). In Sec. V we consider a case when the characteristic polynomial of (25) has at least two nonzero roots.

V. TWO NONZERO ROOTS

Let α and β be two nonzero roots. Consider the vector fields

$$S_0 = \sum_{j=-\infty}^{\infty} A(\tau_j)e^{\alpha\rho_j} \frac{\partial}{\partial \tau_j}, \quad S_1 = \sum_{j=-\infty}^{\infty} B(\tau_j)e^{\beta\rho_j} \frac{\partial}{\partial \tau_j}$$

from the Lie algebra L_x and construct a new sequence of vector fields

$$S_2 = [S_0, S_1], \quad S_{n+1} = [S_0, S_n], \quad n \geq 1.$$

We have

$$DS_0D^{-1} = e^{\alpha\tau}S_0, \quad DS_1D^{-1} = e^{\beta\tau}S_1,$$

$$DS_2D^{-1} = e^{(\alpha+\beta)\tau}S_2 + \beta Ae^{(\alpha+\beta)\tau}S_1 - \alpha Be^{(\alpha+\beta)\tau}S_0.$$

In general, for any $n \geq 3$,

$$DS_nD^{-1} = e^{((n-1)\alpha+\beta)\tau} \left\{ S_n + (c_n\alpha + d_n\beta)AS_{n-1} + (p_nA' + q_nA)AS_{n-2} + \sum_{k=0}^{n-2} \nu(n, k)S_k \right\},$$

where

$$c_n = \frac{(n-1)(n-2)}{2}, \quad d_n = n-1, \quad p_{n+1} = \frac{n(n-1)}{2} \left\{ \frac{n-2}{3}\alpha + \beta \right\}, \quad n \geq 2,$$

$$q_{n+1} = \frac{n(n-2)(n-1)(3n-1)}{24}\alpha^2 + \frac{(n-1)^2n}{2}\alpha\beta + \frac{n(n-1)}{2}\beta^2, \quad n \geq 2.$$

Let us consider a particular case when

$$S_2 = \mu_0 S_0 + \mu_1 S_1. \quad (44)$$

We have

$$\begin{aligned} DS_2 D^{-1} &= D(\mu_0)e^{\alpha\tau}S_0 + D(\mu_1)e^{\beta\tau}S_1 = e^{(\alpha+\beta)\tau}S_2 + \beta A e^{(\alpha+\beta)\tau}S_1 - \alpha B e^{(\alpha+\beta)\tau}S_0 \\ &= e^{(\alpha+\beta)\tau}\{\mu_0 S_0 + \mu_1 S_1\} + \beta A e^{(\alpha+\beta)\tau}S_1 - \alpha B e^{(\alpha+\beta)\tau}S_0. \end{aligned}$$

Comparing coefficients before S_0 and S_1 produces the following two equations:

$$e^{(\alpha+\beta)\tau}\mu_0 - \alpha B e^{(\alpha+\beta)\tau} = D(\mu_0)e^{\alpha\tau}, \quad e^{(\alpha+\beta)\tau}\mu_1 + \beta A e^{(\alpha+\beta)\tau} = D(\mu_1)e^{\beta\tau}.$$

It follows that μ_0 and μ_1 are constants and

$$B(\tau) = -\frac{\mu_0}{\alpha}(e^{-\beta\tau} - 1), \quad A(\tau) = \frac{\mu_1}{\beta}(e^{-\alpha\tau} - 1).$$

Finally, comparing coefficients before $\partial/\partial\tau$ in Eq. (44) implies that $\alpha = -\beta$.

Let us return to the general case. Since L_x is of finite dimension, then there exists such number N that S_0, S_1, \dots, S_N are linearly independent and

$$S_{N+1} = \mu_N S_N + \mu_{N-1} S_{N-1} + \dots + \mu_0 S_0.$$

Then

$$DS_{N+1} D^{-1} = D(\mu_N)DS_N D^{-1} + D(\mu_{N-1})DS_{N-1} D^{-1} + \dots + D(\mu_0)DS_0 D^{-1}$$

and therefore,

$$\begin{aligned} &e^{(N\alpha+\beta)\tau}\{(\mu_N S_N + \mu_{N-1} S_{N-1} + \dots) + A(c_{N+1}\alpha + d_{N+1}\beta)S_N + A(p_{N+1}A' + q_{N+1}A)S_{N+1} + \dots\} \\ &= D(\mu_N)\{e^{(N-1)\alpha+\beta}\tau(S_N + A(c_N\alpha + d_N\beta)S_{N-1} + \dots)\} + D(\mu_{N-1})\{e^{(N-2)\alpha+\beta}\tau S_{N-1} + \dots\} + \dots \end{aligned}$$

By comparing the coefficients before S_N we have

$$e^{(N\alpha+\beta)\tau}\{\mu_N + A(c_{N+1}\alpha + d_{N+1}\beta)\} = D(\mu_N)e^{(N-1)\alpha+\beta}\tau.$$

It follows that μ_N is a constant and then

$$A(c_{N+1}\alpha + d_{N+1}\beta) = \mu_N(e^{-\alpha\tau} - 1).$$

If $c_{N+1}\alpha + d_{N+1}\beta = N\{[(N-1)/2]\alpha + \beta\} \neq 0$, then

$$A(\tau) = C_1(e^{-\alpha\tau} - 1)$$

for some constant C_1 .

If $c_{N+1}\alpha + d_{N+1}\beta = N\{[(N-1)/2]\alpha + \beta\} = 0$ (in this case $\mu_N = 0$) we compare coefficients before S_{N-1} and have

$$e^{(N\alpha+\beta)\tau}\{\mu_{N-1} + A(p_{N+1}A' + q_{N+1}A)\} = D(\mu_{N-1})e^{(N-2)\alpha+\beta}\tau.$$

It follows that μ_{N-1} is a constant and

$$p_{N+1}AA' + q_{N+1}A^2 = \mu_{N-1}(e^{-2\alpha\tau} - 1).$$

Note that if $c_{N+1}\alpha + d_{N+1}\beta = N\{[(N-1)/2]\alpha + \beta\} = 0$, then $p_{N+1} = -[N(N-1)(N+1)/12]\alpha \neq 0$ and $q_{N+1} = -[(N-1)N(N+1)/24]\alpha^2 \neq 0$ for $N \geq 2$. Therefore, $(2/q_{N+1})p_{N+1} = \alpha$. Case $N=1$ should be studied separately ($S_2 = \mu_1 S_1 + \mu_0 S_0$) and it was already. Let us solve the equation

$$p_{N+1}AA' + q_{N+1}A^2 = \mu_{N-1}(e^{-2\alpha\tau} - 1).$$

Denote by $y = A^2$. We have

$$y' + \alpha y = k_1 e^{-2\alpha\tau} - k_1$$

for some constant k_1 . It follows that

$$A^2(\tau) = K_1(e^{-2\alpha\tau} + K_2 e^{-\alpha\tau} + 1)$$

for some constants K_1 and K_2 .

Construct a new sequence of vector fields,

$$S_2^* = [S_1, S_0], \quad S_{n+1}^* = [S_1, S_n^*], \quad n \geq 2.$$

Note that $S_2^* = -S_2$. Since L_x is of finite dimension, then there exists number M such that S_0, S_1, \dots, S_M^* are linearly independent and

$$S_{M+1}^* = \mu_M^* S_M^* + \mu_{M-1}^* S_{M-1}^* + \dots + \mu_0^* S_0.$$

Following are the possibilities:

(1)

$$A(\tau) = K_1(e^{-\alpha\tau} - 1),$$

$$B(\tau) = K_3(e^{-\beta\tau} - 1),$$

(2)

$$A(\tau) = K_1(e^{-\alpha\tau} - 1),$$

$$B^2(\tau) = K_3^2(e^{-2\beta\tau} + K_4 e^{-\beta\tau} + 1), \quad S_{M+1}^* = \mu_M^* S_M^* + \mu_{M-1}^* S_{M-1}^* + \dots + \mu_0^* S_0, \quad \frac{M-1}{2}\beta + \alpha = 0,$$

(3)

$$B(\tau) = K_3(e^{-\beta\tau} - 1),$$

$$A^2(\tau) = K_1^2(e^{-2\alpha\tau} + K_2 e^{-\alpha\tau} + 1), \quad S_{N+1} = \mu_N S_N + \mu_{N-1} S_{N-1} + \dots + \mu_0 S_0, \quad \frac{N-1}{2}\alpha + \beta = 0,$$

and

(4)

$$A^2(\tau) = K_1^2(e^{-2\alpha\tau} + K_2 e^{-\alpha\tau} + 1), \quad S_{N+1} = \mu_N S_N + \mu_{N-1} S_{N-1} + \dots + \mu_0 S_0, \quad \frac{N-1}{2}\alpha + \beta = 0,$$

$$B^2(\tau) = K_3^2(e^{-2\beta\tau} + K_4 e^{-\beta\tau} + 1), \quad S_{M+1}^* = \mu_M^* S_M^* + \mu_{M-1}^* S_{M-1}^* + \dots + \mu_0^* S_0, \quad \frac{M-1}{2}\beta + \alpha = 0,$$

where $K_1, K_2 \neq -2, K_3$, and $K_4 \neq -2$ are some constants, $M, N \geq 2$.

In case (1), vector fields S_0 and S_1 generate an infinite dimensional Lie algebra L_x unless $\alpha + \beta = 0$.

In case (2), we make a substitution $1 - e^{\alpha\tau} = e^{-\alpha w}$. Vector fields S_0 and S_1 become

$$S_0 = K_1 \frac{\partial}{\partial w} + \dots,$$

$$S_1 = \{K_3^2((1 - e^{-\alpha w})^{-2\beta/\alpha} + K_4(1 - e^{-\alpha w})^{-\beta/\alpha} + 1)\}^{1/2} \frac{\partial}{\partial w} + \dots = g(w) \frac{\partial}{\partial w} + \dots.$$

Note that if

$$S_{M+1}^* = \mu_M^* S_M^* + \mu_{M-1}^* S_{M-1}^* + \dots + \mu_0^* S_0,$$

then all coefficients μ_k^* are constants. By comparing coefficients before $\partial/\partial w$ in both sides of the last equation we obtain that $g(w)$ is a solution of linear differential equation with constant coefficients, that is,

$$g(w) = \{K_3^2((1 - e^{-\alpha w})^{-2\beta/\alpha} + K_4(1 - e^{-\alpha w})^{-\beta/\alpha} + 1)\}^{1/2} = \sum_k R_k(w) e^{\nu_k w}, \tag{45}$$

where $R_k(w)$ are some polynomials. One can show that equality (45) holds only if $B(\tau) = K_3(e^{\alpha\tau} + 1)$. It can be shown that in case (3) $A(\tau) = K_1(e^{\beta\tau} + 1)$. In case (4) we make substitution $e^{\alpha\tau} + K_1/2 + \sqrt{e^{2\alpha\tau + K_1} e^{\alpha\tau} + 1} = e^{\alpha w}$. Then

$$S_0 = K_1 \frac{\partial}{\partial w} + \dots,$$

$$S_1 = \left\{ K_3^2 \left(\frac{1}{2} e^{\alpha w} - \frac{K_1}{2} + \left(\frac{K_1^2}{8} - \frac{1}{2} \right) e^{-\alpha w} \right)^{-2\beta/\alpha} + K_4 \left(\frac{1}{2} e^{\alpha w} - \frac{K_1}{2} + \left(\frac{K_1^2}{8} - \frac{1}{2} \right) e^{-\alpha w} + 1 \right) \right\}^{1/2} \frac{\partial}{\partial w} + \dots = g(w) \frac{\partial}{\partial w} + \dots.$$

For function $g(w)$ to be of form $\sum_k R_k(w) e^{\nu_k w}$, where $R_k(w)$ are polynomials, function $B(\tau)$ has to be of form $B(\tau) = K_3(e^{\alpha\tau} + 1)$. Then, by case (3), $A(\tau) = K_1(e^{-\alpha\tau} + 1)$.

It has been proved that in cases (1)–(4) one has
(1*)

$$A(\tau) = K_1(e^{-\alpha\tau} - 1),$$

$$B(\tau) = K_3(e^{\alpha\tau} - 1),$$

(2*)

$$A(\tau) = K_1(e^{-\alpha\tau} - 1),$$

$$B(\tau) = K_3(e^{\alpha\tau} + 1),$$

(3*)

$$\begin{cases} A(\tau) = K_1(e^{-\alpha\tau} + 1), \\ B(\tau) = K_3(e^{\alpha\tau} - 1), \end{cases}$$

and

(4*)

$$A(\tau) = K_1(e^{-\alpha\tau} + 1),$$

$$B(\tau) = K_3(e^{\alpha\tau} + 1).$$

In case (1*) function $d(t, t_1)$ in (12) has the form $d(t, t_1) = c_4(e^{\alpha t_1} - e^{\alpha t}) + c_5(e^{-\alpha t_1} - e^{-\alpha t})$, where c_4 and c_5 are some constants. Equation (12) with such function $d(t, t_1)$ admits a nontrivial x -integral (see Sec. I, Theorem 3 and Sec. VIII).

In Secs. VI and VII we show that cases (3*) and (4*) both correspond to infinite dimensional Lie algebra L_x . Case (2*) also produces an infinite dimensional Lie algebra L_x . It can be proved in the same way as it is proved for case (3*).

VI. CHARACTERISTIC LIE ALGEBRA L_x OF THE CHAIN $t_{1x} = t_x + A_1(e^{\alpha t_1} + e^{\alpha t}) - A_2(e^{-\alpha t} - e^{-\alpha t_1})$

Since $A(\tau) = A_1(e^{-\alpha\tau} + 1)$ and $B(\tau) = A_2(e^{\alpha\tau} - 1)$, then

$$A(\tau)e^{\alpha t} + \sum_{j=1}^k A(\tau_j)e^{\alpha t_j} = A_1 \left(e^{\alpha t} + \left(2 \sum_{j=1}^{k-1} e^{\alpha t_j} \right) + e^{\alpha t_k} \right)$$

and

$$B(\tau)e^{-\alpha t} + \sum_{j=1}^k B(\tau_j)e^{-\alpha t_j} = A_2(e^{-\alpha t} - e^{-\alpha t_k}).$$

We have

$$\frac{1}{A_1}S_0 = (e^{\alpha t} + e^{\alpha t_1}) \frac{\partial}{\partial t_1} + \sum_{k=1}^{\infty} \left(e^{\alpha t} + \left(2 \sum_{j=1}^{k-1} e^{\alpha t_j} \right) + e^{\alpha t_k} \right) \frac{\partial}{\partial t_k} + \sum_{k=1}^{\infty} \left(e^{\alpha t} + \left(2 \sum_{j=1}^{k-1} e^{\alpha t_j} \right) + e^{\alpha t_k} \right) \frac{\partial}{\partial t_{-k}}$$

and

$$\frac{1}{A_2}S_1 = e^{-\alpha t} \tilde{X} - \sum_{k=-\infty}^{\infty} e^{-\alpha t_k} \frac{\partial}{\partial t_k} = e^{-\alpha t} \tilde{X} - \tilde{S}_1,$$

where

$$\tilde{S}_1 = \sum_{k=-\infty}^{\infty} e^{-\alpha t_k} \frac{\partial}{\partial t_k}.$$

In variables $w_j = (1/\alpha)e^{\alpha t_j}$ vector fields \tilde{S}_1 and $(1/A_1)S_0$ can be rewritten as

$$\tilde{S}_1 = \sum_{k=-\infty}^{\infty} \frac{\partial}{\partial w_j},$$

$$\frac{1}{A_1}S_0 = \alpha^2 \sum_{k=1}^{\infty} \left\{ w_k \left(w + 2 \sum_{j=1}^{k-1} w_j \right) + w_k^2 \right\} \frac{\partial}{\partial w_k} + \alpha^2 \sum_{k=1}^{\infty} \left\{ w_{-k} \left(w + 2 \sum_{j=1}^{k-1} w_{-j} \right) + w_{-k}^2 \right\} \frac{\partial}{\partial w_{-k}}.$$

We have

$$T_1 = \left[\tilde{S}_1, \left[\tilde{S}_1, \frac{1}{\alpha^2 A_1} S_0 \right] \right] = 4 \sum_{k=-\infty}^{\infty} k \frac{\partial}{\partial w_k} = 4\tilde{T}_1, \quad \tilde{T}_1 = \sum_{k=-\infty}^{\infty} k \frac{\partial}{\partial w_k},$$

$$T_2 = \left[\tilde{S}_1, \left[\tilde{T}_1, \frac{1}{\alpha^2 A_1} S_0 \right] \right] = 3 \sum_{k=1}^{\infty} \{k^2 - k + 1\} \left(\frac{\partial}{\partial w_k} + \frac{\partial}{\partial w_{-k}} \right) = 3\tilde{T}_2 - 3\tilde{T}_1 + 3\tilde{S}_1, \quad \tilde{T}_2 = \sum_{k=-\infty}^{\infty} k^2 \frac{\partial}{\partial w_k}.$$

Assume that $\tilde{T}_m = \sum_{k=-\infty}^{\infty} k^m (\partial / \partial w_k)$, $m = 1, 2, \dots, n$, are vector fields from L_x . Then

$$\begin{aligned}
T_{m+1} &= \left[\tilde{S}_1, \left[\tilde{T}_m, \frac{1}{\alpha^2 A_1} S_0 \right] \right] = \sum_{k=1}^{\infty} \{2(1 + 2^m + 3^m + \cdots + k^m) + 2k^{m+1} - k^m\} \left(\frac{\partial}{\partial w_k} + \frac{\partial}{\partial w_{-k}} \right) \\
&= \sum_{k=1}^{\infty} \left\{ 2 \left(\frac{k^{m+1}}{m+1} + d_{m,m+1} k^m + \cdots + d_{1,m+1} k + d_{0,m+1} \right) + 2k^{m+1} - k^m \right\} \left(\frac{\partial}{\partial w_k} + \frac{\partial}{\partial w_{-k}} \right)
\end{aligned}$$

and therefore, $\tilde{T}_{m+1} = \sum_{k=-\infty}^{\infty} k^{m+1} \frac{\partial}{\partial w_k} \in L_x$. It shows that $\tilde{T}_n = \sum_{k=-\infty}^{\infty} k^n \frac{\partial}{\partial w_k} \in L_x$ for all $n = 1, 2, 3, \dots$ and L_x is of infinite dimension.

VII. CHARACTERISTIC LIE ALGEBRA L_x OF THE CHAIN $t_{1x} = t_x + A_1(e^{\alpha t_1} + e^{\alpha t}) + A_2(e^{-\alpha t} + e^{-\alpha t_1})$

It was observed in previous studies (see, for instance, Ref. 10) that S -integrable models have the characteristic Lie algebra of finite growth. The chain studied in this section can easily be reduced to the semidiscrete SG model $t_{1x} = t_x + \sin t + \sin t_1$ which belongs to the S -integrable class. It is remarkable that its algebra L_x is of finite growth, or, more exactly, the dimension of the linear space of multiple commutators grows linearly with the multiplicity. Below we prove that the linear space V_n of all commutators of multiplicity $\leq n$ has a basis $\{P_1, P_2, P_3, \dots, P_{2k}; Q_2, Q_4, \dots, Q_{2k}\}$ for $n = 2k$ and a basis $\{P_1, P_2, P_3, \dots, P_{2k+1}; Q_2, Q_4, \dots, Q_{2k}\}$ for $n = 2k + 1$, where the operators P_j and Q_j are defined consecutively,

$$P_1 = [S_0, S_1] + \alpha S_0 + \alpha S_1, \quad Q_1 = P_1,$$

$$P_2 = [S_1, P_1], \quad Q_2 = [S_0, Q_1],$$

$$P_3 = [S_0, P_2] + \alpha P_2, \quad Q_3 = [S_1, Q_2] - \alpha Q_2,$$

$$P_{2n} = [S_1, P_{2n-1}], \quad Q_{2n} = [S_0, Q_{2n-1}],$$

$$P_{2n+1} = [S_0, P_{2n}] + \alpha P_{2n}, \quad Q_{2n+1} = [S_1, Q_{2n}] - \alpha Q_{2n}$$

for $n \geq 1$. Direct calculations show that

$$DP_1 D^{-1} = P_1 - 2\alpha(S_0 + S_1),$$

$$DP_2 D^{-1} = e^{-\alpha\tau}(P_2 + 2\alpha P_1 - 2\alpha^2(S_0 + S_1)),$$

$$DP_3 D^{-1} = P_3 + 2\alpha Q_2 - 2\alpha P_2 - 4\alpha^2 P_1 + 4\alpha^3(S_0 + S_1),$$

$$DP_4 D^{-1} = e^{-\alpha\tau}(P_4 + 2\alpha Q_3 - 4\alpha^2 P_2 + 4\alpha^2 Q_2 - 4\alpha^3 P_1 + 4\alpha^4(S_0 + S_1)),$$

$$DQ_2 D^{-1} = e^{\alpha\tau}(Q_2 - 2\alpha P_1 + 2\alpha^2(S_0 + S_1)),$$

$$DQ_3 D^{-1} = Q_3 + 2\alpha Q_2 - 2\alpha P_2 - 4\alpha^2 P_1 + 4\alpha^3(S_0 + S_1),$$

$$DQ_4 D^{-1} = e^{\alpha\tau}(Q_4 - 2\alpha P_3 + 2\alpha^2(P_2 - Q_2) + 4\alpha^3 P_1 - 4\alpha^4(S_0 + S_1)),$$

$$P_3 = Q_3, \quad [S_1, P_2] = -\alpha P_2, [S_0, Q_2] = \alpha Q_2, [S_1, P_4] = -\alpha P_4, [S_0, Q_4] = \alpha Q_4. \quad (46)$$

The coefficient before $\partial/\partial\tau$ in all vector fields $DP_i D^{-1}$ and $DQ_i D^{-1}$, $1 \leq i \leq 4$, is zero.

Lemma 10: For $n \geq 1$ we have

- (1) $DP_{2n+1}D^{-1} + 2\alpha e^{\alpha\tau}DP_{2n}D^{-1} = P_{2n+1} + 2\alpha Q_{2n}$,
- (2) $e^{\alpha\tau}DP_{2n+2}D^{-1} - \alpha DP_{2n+1}D^{-1} = P_{2n+2} + \alpha Q_{2n+1}$,
- (3) $DQ_{2n+1}D^{-1} - 2\alpha e^{-\alpha\tau}DQ_{2n}D^{-1} = Q_{2n+1} - 2\alpha P_{2n}$,
- (4) $e^{-\alpha\tau}DQ_{2n+2}D^{-1} + \alpha DQ_{2n+1}D^{-1} = Q_{2n+2} - \alpha P_{2n+1}$,
- (5) $P_{2n+1} = Q_{2n+1}$,
- (6) $[S_1, P_{2n+2}] = -\alpha P_{2n+2}$, and
- (7) $[S_0, Q_{2n+2}] = \alpha Q_{2n+2}$.

Moreover, the coefficient before $\partial/\partial\tau$ in all vector fields DP_kD^{-1} and DQ_kD^{-1} is zero.

Proof: We prove the lemma by induction on n . It follows from (46) that the base of induction holds for $n=1$. Assume that (1)–(7) are true for all n , $1 \leq n \leq k$. Let us prove that (1) is true for $n=k+1$.

$$\begin{aligned}
 DP_{2n+3}D^{-1} &= D([S_0, P_{2n+2}] + \alpha P_{2n+2})D^{-1} = [e^{\alpha\tau}S_0, DP_{2n+2}D^{-1}] + \alpha DP_{2n+2}D^{-1} \\
 &= [e^{\alpha\tau}S_0, \alpha e^{-\alpha\tau}DP_{2n+1}D^{-1} + e^{-\alpha\tau}P_{2n+2} + \alpha e^{-\alpha\tau}Q_{2n+1}] + \alpha DP_{2n+2}D^{-1} \\
 &= -\alpha^2(1 + e^{-\alpha\tau})DP_{2n+1}D^{-1} + \alpha e^{-\alpha\tau}[e^{\alpha\tau}S_0, DP_{2n+1}D^{-1}] \\
 &\quad - \alpha(1 + e^{-\alpha\tau})P_{2n+2} - \alpha^2(1 + e^{-\alpha\tau})Q_{2n+1} + P_{2n+3} - \alpha P_{2n+2} + \alpha Q_{2n+2} + \alpha DP_{2n+2}D^{-1} \\
 &= -\alpha^2(1 + e^{-\alpha\tau})DP_{2n+1}D^{-1} + \alpha e^{-\alpha\tau}D[S_0, Q_{2n+1}]D^{-1} - \alpha(2 + e^{-\alpha\tau})P_{2n+2} \\
 &\quad - \alpha^2(1 + e^{-\alpha\tau})Q_{2n+1} + P_{2n+3} + \alpha Q_{2n+2} + \alpha DP_{2n+2}D^{-1} \\
 &= -\alpha^2(1 + e^{-\alpha\tau})DP_{2n+1}D^{-1} + \alpha Q_{2n+2} - \alpha^2 P_{2n+1} - \alpha^2 DQ_{2n+1}D^{-1} \\
 &\quad - \alpha(2 + e^{-\alpha\tau})P_{2n+2} - \alpha^2(1 + e^{-\alpha\tau})Q_{2n+1} - 2\alpha^2 Q_{2n+1} - 2\alpha P_{2n+2} + P_{2n+3} \\
 &= -2\alpha^2 DP_{2n+1}D^{-1} + 2\alpha Q_{2n+2} - 2\alpha^2 Q_{2n+1} - 2\alpha P_{2n+2} + P_{2n+3} \\
 &= 2\alpha P_{2n+2} + 2\alpha^2 Q_{2n+1} - 2\alpha e^{\alpha\tau}DP_{2n+2}D^{-1} + 2\alpha Q_{2n+2} - 2\alpha^2 Q_{2n+1} - 2\alpha P_{2n+2} + P_{2n+3} \\
 &= -2\alpha e^{\alpha\tau}DP_{2n+2}D^{-1} + 2\alpha Q_{2n+2} + P_{2n+3}.
 \end{aligned}$$

The proof of (3) is the same as the proof of (1). Let us show that (5) is true for $n=k+1$. We have

$$\begin{aligned}
 DP_{2n+3}D^{-1} &= -2\alpha e^{\alpha\tau}DP_{2n+2}D^{-1} + 2\alpha Q_{2n+2} + P_{2n+3} = -2\alpha(\alpha DP_{2n+1}D^{-1} + P_{2n+2} + \alpha Q_{2n+1}) \\
 &\quad + 2\alpha Q_{2n+2} + P_{2n+3}
 \end{aligned}$$

and

$$\begin{aligned}
 DQ_{2n+3}D^{-1} &= 2\alpha e^{-\alpha\tau}DQ_{2n+2}D^{-1} - 2\alpha P_{2n+2} + Q_{2n+3} = 2\alpha(-\alpha DQ_{2n+1}D^{-1} + Q_{2n+2} - \alpha P_{2n+1}) \\
 &\quad - 2\alpha P_{2n+2} + Q_{2n+3}.
 \end{aligned}$$

By (5), $P_{2n+1} = Q_{2n+1}$ and therefore

$$D(P_{2n+3} - Q_{2n+3})D^{-1} = -2\alpha P_{2n+2} - 2\alpha Q_{2n+2} + 2\alpha Q_{2n+2} + 2\alpha P_{2n+2} = 0.$$

Hence, $P_{2n+3} = Q_{2n+3}$.

Let us prove that (2) is true for $n=k+1$. We have

$$\begin{aligned}
 e^{\alpha\tau}DP_{2n+1}D^{-1} &= e^{\alpha\tau}D[S_1, P_{2n+3}]D^{-1} = e^{\alpha\tau}[e^{-\alpha\tau}S_1, DP_{2n+3}D^{-1}] \\
 &= e^{\alpha\tau}[e^{-\alpha\tau}S_1, -2\alpha e^{\alpha\tau}DP_{2n+2}D^{-1} + 2\alpha Q_{2n+2} + P_{2n+3}] \\
 &= e^{\alpha\tau}(-2\alpha^2(1 + e^{\alpha\tau})DP_{2n+2}D^{-1}) - 2\alpha e^{2\alpha\tau}[e^{-\alpha\tau}S_1, DP_{2n+2}D^{-1}] + P_{2n+4} + 2\alpha Q_{2n+3} \\
 &\quad + 2\alpha^2 Q_{2n+2} \\
 &= -2\alpha^2(e^{\alpha\tau} + e^{2\alpha\tau})DP_{2n+2}D^{-1} + 2\alpha^2 e^{2\alpha\tau}DP_{2n+2}D^{-1} + P_{2n+4} + 2\alpha Q_{2n+3} + 2\alpha^2 Q_{2n+2} \\
 &= -2\alpha^2 e^{\alpha\tau}DP_{2n+2}D^{-1} + P_{2n+4} + 2\alpha Q_{2n+3} + 2\alpha^2 Q_{2n+2} \\
 &= \alpha DP_{2n+3}D^{-1} - \alpha P_{2n+3} - 2\alpha^2 Q_{2n+2} + P_{2n+4} + 2\alpha Q_{2n+3} + 2\alpha^2 Q_{2n+2} \\
 &= \alpha DP_{2n+3}D^{-1} + \alpha Q_{2n+3} + P_{2n+4}.
 \end{aligned}$$

The proof of (4) is similar to the proof of (2).

Let us prove that (6) is true for $n=k+1$,

$$\begin{aligned}
 D[S_1, P_{2n+4}]D^{-1} &= [e^{-\alpha\tau}S_1, \alpha e^{-\alpha\tau}DP_{2n+3}D^{-1} + e^{-\alpha\tau}P_{2n+4} + \alpha e^{-\alpha\tau}Q_{2n+3}] \\
 &= [e^{-\alpha\tau}S_1, \alpha e^{-\alpha\tau}(-2\alpha e^{\alpha\tau}DP_{2n+2}D^{-1} + P_{2n+3} + 2\alpha Q_{2n+2}) + e^{-\alpha\tau}P_{2n+4} + \alpha e^{-\alpha\tau}Q_{2n+3}] \\
 &= [e^{-\alpha\tau}S_1, -2\alpha^2 DP_{2n+2}D^{-1} + 2\alpha e^{-\alpha\tau}P_{2n+3} + 2\alpha^2 e^{-\alpha\tau}Q_{2n+2} + e^{-\alpha\tau}P_{2n+4}] \\
 &= -2\alpha^2 D[S_1, P_{2n+2}]D^{-1} - 2\alpha^2 e^{-2\alpha\tau}(1 + e^{\alpha\tau})P_{2n+3} - 2\alpha^3 e^{-2\alpha\tau}(1 + e^{\alpha\tau})Q_{2n+2} \\
 &\quad + 2\alpha e^{-2\alpha\tau}P_{2n+4} + 2\alpha^2 e^{-2\alpha\tau}Q_{2n+3} + 2\alpha^3 e^{-2\alpha\tau}Q_{2n+2} - \alpha e^{-2\alpha\tau}(1 + e^{\alpha\tau})P_{2n+4} \\
 &\quad + e^{-2\alpha\tau}[S_1, P_{2n+4}] \\
 &= 2\alpha^3 DP_{2n+2}D^{-1} - 2\alpha^2 e^{-\alpha\tau}P_{2n+3} + \alpha(e^{-2\alpha\tau} - e^{-\alpha\tau})P_{2n+4} - 2\alpha^3 e^{-\alpha\tau}Q_{2n+2} \\
 &\quad + e^{-2\alpha\tau}[S_1, P_{2n+4}] \\
 &= \alpha^2 e^{-\alpha\tau}P_{2n+3} + 2\alpha^3 e^{-\alpha\tau}Q_{2n+2} - \alpha^2 e^{-\alpha\tau}DP_{2n+3}D^{-1} - 2\alpha^2 e^{-\alpha\tau}P_{2n+3} \\
 &\quad + \alpha(e^{-2\alpha\tau} - e^{-\alpha\tau})P_{2n+4} - 2\alpha^3 e^{-\alpha\tau}Q_{2n+2} + e^{-2\alpha\tau}[S_1, P_{2n+4}] \\
 &= -\alpha^2 e^{-\alpha\tau}P_{2n+3} + \alpha(e^{-2\alpha\tau} - e^{-\alpha\tau})P_{2n+4} - \alpha DP_{2n+4}D^{-1} + \alpha e^{-\alpha\tau}P_{2n+4} + \alpha^2 e^{-\alpha\tau}Q_{2n+3} \\
 &\quad + e^{-2\alpha\tau}[S_1, P_{2n+4}].
 \end{aligned}$$

Thus,

$$D[S_1, P_{2n+4}]D^{-1} = e^{-2\alpha\tau}[S_1, P_{2n+4}] + \alpha e^{-2\alpha\tau}P_{2n+4} - \alpha DP_{2n+4}D^{-1},$$

$$D([S_1, P_{2n+4}] + \alpha P_{2n+4})D^{-1} = e^{-2\alpha\tau}([S_1, P_{2n+4}] + \alpha P_{2n+4}).$$

Hence, $[S_1, P_{2n+4}] = -\alpha P_{2n+4}$. □

The proof of (7) is similar to the proof of (6).

Corollary 1: We have

$$e^{-\alpha\tau}DQ_{2n}D^{-1} + e^{\alpha\tau}DP_{2n}D^{-1} = Q_{2n} + P_{2n},$$

$$DP_{2n+1}D^{-1} = P_{2n+1} + \sum_{k=1}^n (\mu_{2k}^{(2n+1)} P_{2k} + \nu_{2k}^{(2n+1)} Q_{2k}) + \sum_{k=0}^{n-1} \mu_{2k+1}^{(2n+1)} P_{2k+1} + \mu_0^{(2n+1)} S_0 + \nu_0^{(2n+1)} S_1,$$

$$DP_{2n}D^{-1} = e^{-\alpha\tau} \left(P_{2n} + \sum_{k=1}^{n-1} (\mu_{2k}^{(2n)} P_{2k} + \nu_{2k}^{(2n)} Q_{2k}) + \sum_{k=0}^{n-1} \mu_{2k+1}^{(2n)} P_{2k+1} + \mu_0^{(2n)} S_0 + \nu_0^{(2n)} S_1 \right),$$

$$DQ_{2n}D^{-1} = e^{\alpha\tau} \left(Q_{2n} - \sum_{k=1}^{n-1} (\mu_{2k}^{(2n)} P_{2k} + \nu_{2k}^{(2n)} Q_{2k}) - \sum_{k=0}^{n-1} \mu_{2k+1}^{(2n)} P_{2k+1} - \mu_0^{(2n)} S_0 - \nu_0^{(2n)} S_1 \right).$$

Moreover, $\mu_{2n}^{(2n+1)} = -2\alpha$, $\nu_{2n}^{(2n+1)} = 2\alpha$, and $\mu_{2n-1}^{(2n)} = 2\alpha$.

Assume that L_x is of finite dimension. There are three possibilities:

- (1) $S_0, S_1, P_1, P_2, Q_2, P_3, P_4, Q_4, \dots, P_{2n-1}$ are linearly independent and $S_0, S_1, P_1, P_2, Q_2, P_3, P_4, Q_4, \dots, P_{2n-1}, P_{2n}$ are linearly dependent,
- (2) $S_0, S_1, P_1, P_2, Q_2, P_3, P_4, Q_4, \dots, P_{2n-1}, P_{2n}$ are linearly independent and $S_0, S_1, P_1, P_2, Q_2, P_3, P_4, Q_4, \dots, P_{2n-1}, P_{2n}, Q_{2n}$ are linearly dependent, and
- (3) $S_0, S_1, P_1, P_2, Q_2, P_3, P_4, Q_4, \dots, P_{2n}, Q_{2n}$ are linearly independent and $S_0, S_1, P_1, P_2, Q_2, P_3, P_4, Q_4, \dots, P_{2n}, Q_{2n}, P_{2n+1}$ are linearly dependent.

In case (1),

$$P_{2n} = \gamma_{2n-1} P_{2n-1} + \gamma_{2n-2} P_{2n-2} + \eta_{2n-2} Q_{2n-2} + \dots$$

and

$$DP_{2n}D^{-1} = D(\gamma_{2n-1})DP_{2n-1}D^{-1} + D(\gamma_{2n-2})DP_{2n-2}D^{-1} + D(\eta_{2n-2})DQ_{2n-2}D^{-1} + \dots \quad (47)$$

We use Corollary 1 to compare the coefficients before P_{2n-1} in (47) and have the contradictory equality

$$e^{-\alpha\tau}(\gamma_{2n-1} + 2\alpha) = D(\gamma_{2n-1}).$$

It shows that case (1) is impossible to have.

In case (2),

$$Q_{2n} = \gamma_{2n} P_{2n} + \gamma_{2n-1} P_{2n-1} + \eta_{2n-2} Q_{2n-2} + \dots$$

and

$$DQ_{2n}D^{-1} = D(\gamma_{2n})DP_{2n}D^{-1} + D(\gamma_{2n-1})DP_{2n-1}D^{-1} + D(\eta_{2n-2})DQ_{2n-2}D^{-1} + \dots \quad (48)$$

We use Corollary 1 to compare the coefficients before P_{2n-1} in (48) and have the contradictory equation

$$e^{\alpha\tau}(\gamma_{2n-1} - 2\alpha) = D(\gamma_{2n-1}).$$

It shows that case (2) is impossible to have.

In case (3),

$$P_{2n+1} = \eta_{2n} Q_{2n} + \gamma_{2n} P_{2n} + \dots$$

and

$$DP_{2n+1}D^{-1} = D(\eta_{2n})DQ_{2n}D^{-1} + D(\gamma_{2n})DP_{2n}D^{-1} + \dots \quad (49)$$

We use Corollary 1 to compare the coefficients before P_{2n} in (49) and have the contradictory equation

$$(\gamma_{2n} - 2\alpha) = D(\gamma_{2n})e^{-\alpha\tau}.$$

It shows that case (3) also fails to be true. Therefore, characteristic Lie algebra L_x is of infinite dimension.

VIII. FINDING x -INTEGRALS

Now we are ready to prove the main theorem, Theorem 3, formulated in Sec. I. Really, in the previous sections we proved that if chain (12) admits a nontrivial x -integral, then it is one of forms (1)–(4). List (i)–(iv) allows one to prove the inverse statement: each of the equations from the list admits indeed a nontrivial x -integral. \square

Let us explain briefly how we found list (i)–(iv). Since for each equation (1)–(4) we have constructed the related characteristic Lie algebra to find x -integral F one has to solve the corresponding system of the first order partial differential equations. Below we illustrate the method with case (2), for which the basis of the characteristic algebra L_x is given by the vector fields

$$\tilde{Y} = \partial_x + Y_{a(\tau)t+b(\tau)}, \quad T_1 = Y_{-a(\tau)}, \quad \tilde{X} = \frac{\partial}{\partial t} + \frac{\partial}{\partial t_1} + \frac{\partial}{\partial t_{-1}} + \frac{\partial}{\partial t_2} + \frac{\partial}{\partial t_{-2}} + \dots,$$

where $a(\tau) = c_0\tau$ and $b(\tau) = c_2\tau^2 + c_3\tau$. Note that x -integral F of (2) should satisfy the equations $\tilde{Y}F=0$, $T_1F=0$, and $\tilde{X}F=0$. Introduce new variables $t, w, w_{\pm 1}, \dots$ where $w_j = \ln(\tau_j)$ and $\tau_j = t_j - t_{j+1}$. Vector fields \tilde{X} , T_1 , and \tilde{Y} in new variables are rewritten as

$$\tilde{X} = \frac{\partial}{\partial t}, \quad T_1 = \sum_{j=-\infty}^{\infty} c_0 \frac{\partial}{\partial w_j},$$

$$\tilde{Y} = \frac{\partial}{\partial x} - t \sum_{j=-\infty}^{\infty} c_0 \frac{\partial}{\partial w_j} + c_0 \sum_{j=-\infty}^{\infty} \{\tilde{\rho}_j + \tilde{b}(w_j)\} \frac{\partial}{\partial w_j} = \frac{\partial}{\partial x} - tT_1 + c_0 \sum_{j=-\infty}^{\infty} \{\tilde{\rho}_j + \tilde{b}(w_j)\} \frac{\partial}{\partial w_j},$$

where

$$\tilde{\rho}_j = \begin{cases} \sum_{k=0}^{j-1} e^{w_k} & \text{if } j \geq 1 \\ 0 & \text{if } j = 0, \\ -\sum_{k=j}^{-1} e^{w_k} & \text{if } j \leq -1. \end{cases} \quad \tilde{b}(w_j) = -\frac{1}{c_0}(c_2 e^{w_j} + c_3),$$

Note that since we have $\tilde{X}F=0$, F does not depend on t . Now let us consider the vector field

$$\tilde{Y} + tT_1 = A = \frac{\partial}{\partial x} + c_0 \sum_{j=-\infty}^{\infty} \{\tilde{\rho}_j + \tilde{b}(w_j)\} \frac{\partial}{\partial w_j}.$$

We can write the vector field A explicitly as

$$A = \frac{\partial}{\partial x} + \sum_{j=-\infty}^{\infty} \left\{ \left(c_0 \sum_{k=0}^{j-1} e^{w_k} \right) - c_2 e^{w_j} - c_3 \right\} \frac{\partial}{\partial w_j} = \frac{\partial}{\partial x} - \frac{c_3}{c_0} T_1 + \sum_{j=-\infty}^{\infty} \left\{ \left(c_0 \sum_{k=0}^{j-1} e^{w_k} \right) - c_2 e^{w_j} \right\} \frac{\partial}{\partial w_j}.$$

The commutator $[T_1, A]$ gives

$$[T_1, A] = c_0 A - c_0 \frac{\partial}{\partial x} + c_3 T_1.$$

Thus we have three vector fields,

$$A - \frac{\partial}{\partial x} + \frac{c_3}{c_0} T_1 := \tilde{A} = \sum_{j=-\infty}^{\infty} \left\{ \left(c_0 \sum_{k=0}^{j-1} e^{w_k} \right) - c_2 e^{w_j} \right\} \frac{\partial}{\partial w_j},$$

$$\frac{T_1}{c_0} := T_1 = \sum_{j=-\infty}^{\infty} \frac{\partial}{\partial w_j}, \quad \tilde{X}_1 = \frac{\partial}{\partial x},$$

which solve $\tilde{A}F=0$, $\tilde{T}_1F=0$, and $\tilde{X}_1F=0$. Note that $[\tilde{T}_1, \tilde{A}]=\tilde{A}$. Since $\tilde{X}_1F=0$, F does not depend on x . Hence we end up with two equations. By the Jacobi theorem the system of equations has a nontrivial solution $F(w, w_1, w_2)$ depending on three variables. Therefore we need first three terms of \tilde{A} and \tilde{T}_1 ,

$$\tilde{A} = -c_2 w \frac{\partial}{\partial w} + (c_0 e^w - c_2 e^{w_1}) \frac{\partial}{\partial w_1} + (c_0 e^w + c_0 e^{w_1} - c_2 e^{w_2}) \frac{\partial}{\partial w_2},$$

$$\tilde{T}_1 = \frac{\partial}{\partial w} + \frac{\partial}{\partial w_1} + \frac{\partial}{\partial w_2}.$$

Now we again introduce new variables $w = \epsilon$, $w - w_1 = \epsilon_1$, and $w_1 - w_2 = \epsilon_2$. Vector fields \tilde{A} and \tilde{T}_1 in new variables are rewritten as

$$\tilde{A} = e^\epsilon \left\{ -c_2 \frac{\partial}{\partial \epsilon} + ((-c_2 - c_0) + c_2 e^{-\epsilon_1}) \frac{\partial}{\partial \epsilon_1} + ((-c_2 - c_0) e^{-\epsilon_1} + c_2 e^{-\epsilon_1 - \epsilon_2}) \frac{\partial}{\partial \epsilon_2} \right\}, \quad \tilde{T}_1 = \frac{\partial}{\partial \epsilon}.$$

To find the x -integral (ii) in Theorem 3 one has to solve the equation

$$\left\{ ((-c_2 - c_0) + c_2 e^{-\epsilon_1}) \frac{\partial}{\partial \epsilon_1} + e^{-\epsilon_1} ((-c_2 - c_0) + c_2 e^{-\epsilon_2}) \frac{\partial}{\partial \epsilon_2} \right\} F = 0.$$

IX. CONCLUSION

In this article the problem of classification of Darboux integrable nonlinear semidiscrete chains of hyperbolic type was studied. An approach based on the notion of characteristic Lie algebra was properly modified and successfully used. We gave a complete list of hyperbolic-type chains $t_{1x} = t_x + d(t, t_1)$ admitting nontrivial x -integrals. We demonstrated that the method of characteristic Lie algebras provides an effective tool to classify integrable discrete chains as well. The method did not get much attention in the literature. To our knowledge there are only two studies (see Refs. 9 and 14) where the characteristic Lie algebras are applied for solving the classification problem for the partial differential equations and systems. Surprisingly the first of them was published in 1981 and the second one only 25 years later.

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¹V. E. Adler and S. Y. Startsev, *Teor. Mat. Fiz.* **121**, 271 (1999); [*Theor. Math. Phys.* **121**, 1484 (1999)].

²A. V. Zabrodin, *Teor. Mat. Fiz.* **113**, 179 (1997); [*Theor. Math. Phys.* **113**, 1347 (1997)].

³R. I. Yamilov, *J. Phys. A* **39**, R541 (2006).

⁴F. W. Nijhoff and H. W. Capel, *Acta Appl. Math.* **39**, 133 (1995).

⁵B. Grammaticos, G. Karra, V. Papageorgiou, and A. Ramani, *Integrability of Discrete-Time Systems, Chaotic Dynamics*, NATO Advanced Institute, Series B: Physics Vol. 298 (Plenum, New York, 1992), pp. 75–90.

⁶G. Darboux, *Leçons sur la théorie générale des surfaces et les applications géométriques du calcul infinitesimal*, T.2 (Gautier-Villars, Paris, 1915).

⁷I. M. Anderson and N. Kamran, *Duke Math. J.* **87**, 265 (1997).

⁸A. V. Zhiber and V. V. Sokolov, *Usp. Mat. Nauk* **56**, 63 (2001); [*Russ. Math. Surveys* **56**, 61 (2001)].

⁹A. B. Shabat and R. I. Yamilov, *Exponential Systems of Type I and the Cartan Matrices*, preprint ed. (Bashkirian Branch of Academy of Science of the USSR, Ufa, 1981) (in Russian).

¹⁰A. N. Leznov, V. G. Smirnov, and A. B. Shabat, *Teor. Mat. Fiz.* **51**, 10 (1982).

- ¹¹ A. V. Zhiber and F. Kh. Mukminov, in *Problems of Mathematical Physics and Asymptotics of their Solutions*, edited by L. A. Kalyakin (Institute of Mathematics, RAN, Ufa, 1991) pp. 13–33.
- ¹² A. A. Bormisov and F. Kh. Mukminov, *Teor. Mat. Fiz.* **127**, 47 (2001); [[Theor. Math. Phys.](#) **127**, 446 (2001)].
- ¹³ V. V. Sokolov and S. Ya. Startsev, [Theor. Math. Phys.](#) **155**, 802 (2008).
- ¹⁴ A. V. Zhiber and R. D. Murtazina, *Fundam. Prikl. Mat.* **12**, 65 (2006) (in Russian).
- ¹⁵ N. Kh. Ibragimov and A. B. Shabat, *Funkc. Anal. Priloz.* **14**, 25 (1980).
- ¹⁶ A. V. Mikhailov, A. B. Shabat, and R. I. Yamilov, *Usp. Mat. Nauk* **42**, 3 (1987).
- ¹⁷ R. I. Yamilov and D. Levi, *J. Nonlinear Math. Phys.* **11**, 75 (2004).
- ¹⁸ M. Gürses and A. Karasu, [Phys. Lett. A](#) **251**, 247 (1999); e-print arXiv:solv-int/9811013.
- ¹⁹ M. Gürses, A. Karasu, and R. Turhan, [J. Phys. A](#) **34**, 5705 (2001); e-print arXiv:nlin.SI/0101031.
- ²⁰ S. I. Svinolupov, [Phys. Lett. A](#) **135**, 32 (1989).
- ²¹ A. V. Zhiber and A. B. Shabat, *Dokl. Akad. Nauk SSSR* **247**, 1103 (1979); [*Sov. Phys. Dokl.* **24**, 607 (1979)].
- ²² A. V. Zhiber and A. B. Shabat, *Dokl. Akad. Nauk SSSR* **277**, 29 (1984); [*Sov. Math. Dokl.* **30**, 23 (1984)].
- ²³ I. T. Habibullin, *Symmetry, Integr. Geom.: Methods Appl.* **1**, 23 (2005); e-print arXiv:nlin.SI/0506027.
- ²⁴ I. Habibullin and A. Pekcan, [Theor. Math. Phys.](#) **151**, 781 (2007); e-print arXiv:nlin/0610074.
- ²⁵ I. Habibullin, N. Zheltukhina, and A. Pekcan, *Turk. J. Math.* **32**, 1 (2008); e-print arXiv:nlin/0703065.