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Spectroscopy around $^{36}\text{Ca}^*$

A. BÜRGER¹, F. AZAIEZ², M. STANOIU³, ZS. DOMBRÁDI⁴,
A. ALGORA⁴, A. AL-KHATIB¹, B. BASTIN⁵, G. BENZONI⁶,
R. BORCEA⁷, CH. BOURGEOIS², P. BRINGEL¹, E. CLÉMENT⁸,
J.-C. DALOUZY⁹, Z. DLOUHÝ¹⁰, A. DROUART⁸, C. ENGELHARDT¹,
S. FRANCHOO², ZS. FÜLÖP⁴, A. GÖRGEN⁸, S. GRÉVY⁹, H. HÜBEL¹,
F. IBRAHIM², W. KORTEN⁸, J. MRÁZEK¹⁰, A. NAVIN⁹, F. ROTARU⁷,
P. ROUSSEL-CHOMAZ⁹, M.-G. SAINT-LAURENT⁹, G. SLETTEN¹²,
D. SOHLER⁴, O. SORLIN⁹, C. TIMIS¹¹, CH. THEISEN⁸, D. VERNEY² AND
S. WILLIAMS¹¹.

¹Helmholtz-Institut für Strahlen- und Kernphysik, Univ. Bonn, Germany
²Institut de Physique Nucléaire, IN2P3-CNRS, Orsay, France
³GSI, Darmstadt, Germany
⁴Institute of Nuclear Research, Debrecen, Hungary
⁵Laboratoire de Physique Corpusculaire, Caen, France
⁶Università degli studi e INFN sezione di Milano, Italy
⁷IFIN-HH, Bucharest-Magurele, Romania
⁸DAPNIA/SPhN, CEA Saclay, France
⁹GANIL, Caen, France
¹⁰Nuclear Physics Institute of ASCR, Řež, Czech Republic
¹¹Department of Physics, University of Surrey, UK
¹²Niels Bohr Institute, University of Copenhagen, Denmark

23 An experiment was performed to study excited states in neutron-deficient
24 nuclei around Ca. A one-neutron knockout reaction was used to produce
25 ^{36}Ca ions from a ^{37}Ca secondary beam, and in-beam γ -rays were measured.
26 The 2^+ energy in ^{36}Ca is compared to the mirror nucleus ^{36}S to deduce
27 information on the isospin dependence of the nuclear force near the proton
28 drip line. The energy of the first excited 2^+ state in ^{36}Ca and the cross
29 section for the 1-neutron knock-out reaction from ^{37}Ca at $\approx 45 \cdot A$ MeV
30 were obtained. Furthermore, for two other $T_z = -2$ nuclei, ^{28}S and ^{32}Ar ,
31 the de-excitation of the first 2^+ state has been observed.

32 PACS numbers: 21.10.-k, 23.20.Lv, 27.30+t

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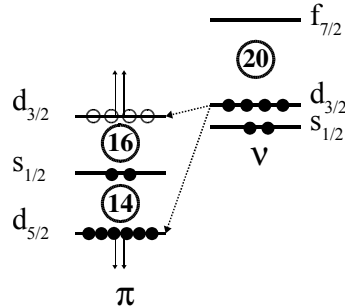


Fig. 1. Illustration of the effect of the $\nu d_{3/2}$ orbital filling in ^{36}S . Due to the tensor interaction, the $\pi d_{3/2}$ level is shifted up and the $\pi d_{5/2}$ level is shifted down in energy.

1 In recent years, an intensive research activity was devoted to the study
 2 of nuclear structure of extremely neutron- or proton-rich nuclei, both the-
 3 oretically and experimentally. In this context, we aimed in the present
 4 experiment to measure the excitation energy of the first 2^+ state in ^{36}Ca
 5 and compare it to its mirror nucleus ^{36}S . In the ground state of ^{36}S , the
 6 $\pi d_{5/2}$ and $s_{1/2}$ as well as the $\nu d_{3/2}$ orbitals are completely filled. In ^{36}Ca ,
 7 the same orbitals are occupied with neutron and proton shells exchanged.
 8 Due to the tensor interaction between the proton spin-orbit partners $d_{5/2}$
 9 and $d_{3/2}$ and the neutron $d_{3/2}$ orbital, the proton $d_{5/2}$ orbital becomes more
 10 bound whereas the $\pi d_{3/2}$ orbital becomes less bound than for nuclei where
 11 the $\nu d_{3/2}$ shell is not completely filled (1). Assuming that the effect of the
 12 filling of the $\nu d_{3/2}$ on the $\pi s_{1/2}$ is weak, this enlarges the gaps between the
 13 $\pi s_{1/2}$ and $\pi d_{3/2}$ levels and between the $\pi s_{1/2}$ and $\pi d_{5/2}$ levels, as illustrated
 14 in fig. 1. These shifts lead to high excitation energies for the first 2^+ states in
 15 both ^{36}S and ^{34}Si , which from this point of view reflects a spherical rigidity
 16 comparable to the doubly magic nucleus ^{40}Ca . For ^{36}Ca , the mirror nucleus
 17 of ^{36}S , the same picture should apply with protons and neutrons exchanged,
 18 so that also in this case a high excitation energy can be expected for the 2^+
 19 state.

20 The experiment was performed at the GANIL in Caen, France. The
 21 two-step fragmentation technique was used (2) to populate excited states in
 22 ^{36}Ca . A primary beam of ^{40}Ca with an energy of 95.4 MeV was fragmented
 23 on a carbon foil in the SISSI target device (3). The Alpha spectrometer,
 24 optimised for ^{37}Ca or, in a different setting, ^{36}Ca , was used to purify the
 25 resulting beam cocktail with the help of a degrader. Event-by-event iden-
 26 tification of the beam particles was achieved using a time measurement
 27 between the high frequency of the accelerator and the time signal from a
 28 CATS detector (4), that was placed just in front of the secondary target. In

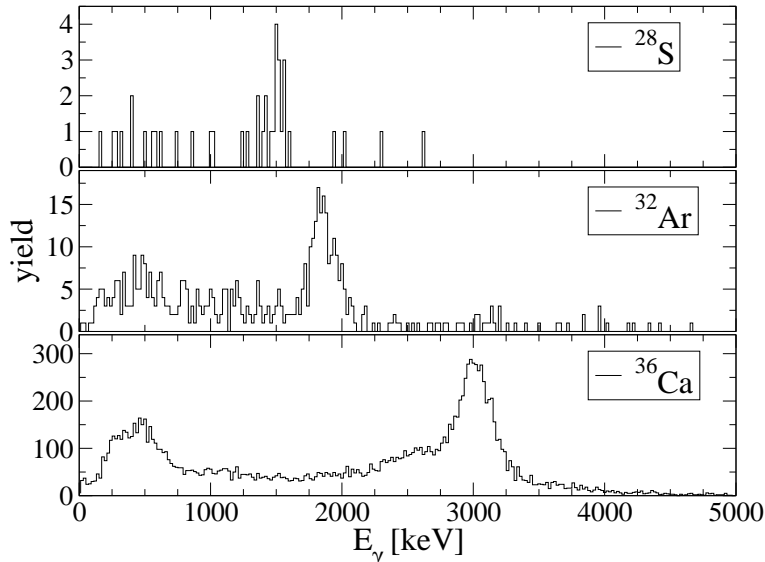


Fig. 2. Gamma-ray spectra for the nuclei ^{36}Ca , ^{32}Ar and ^{28}S . The energies of the 2^+ states have been determined to be 3036(11) keV, 1873(20) keV and 1525(30) keV, respectively.

1 the secondary target, a ^9Be foil of 200 mg/cm² thickness, further nucleons
 2 were removed at energies between $60 \cdot A$ MeV before and $35 \cdot A$ MeV after the
 3 target. Behind the secondary target, the produced fragments were identi-
 4 fied through time-of-flight, $B\rho$ and energy-loss measurements in the SPEG
 5 spectrometer (5). For some settings, suppression of the secondary beam in
 6 the focal plane necessitated the placement of an additional slit in SPEG.

7 Gamma-ray energies were measured with the *Château de Cristal*, an
 8 array of 74 BaF₂ detectors (6), that was placed around the Be target. The
 9 γ -ray detectors were calibrated using a ^{22}Na source and well separated and
 10 sufficiently intense known transitions in the nuclei ^{28}Si , ^{32}S , ^{34}Ar , ^{29}Si and
 11 ^{33}Cl , which were also produced in the secondary target from different beam
 12 components. The Doppler-correction for γ -ray energies from in-flight decays
 13 used the momentum measured in SPEG, assuming that the decays took
 14 place in the middle of the target. An add-back procedure was applied to
 15 reconstruct Compton-scattered γ -ray energies. Gamma-ray spectra for the
 16 three nuclei ^{36}Ca , ^{32}Ar and ^{28}S are shown in fig. 2. The energy of the
 17 2^+ state in ^{36}Ca has been determined to be $E(2^+) = 3036(11)$ keV, in
 18 agreement with the value measured at GSI in a similar experiment (7). The
 19 estimated $E(2^+)$ for ^{28}S is $\approx 1525(30)$ keV, and $\approx 1873(20)$ keV for ^{32}Ar ,
 20 which is 50 keV above the value reported by Cottle *et al.* (8).

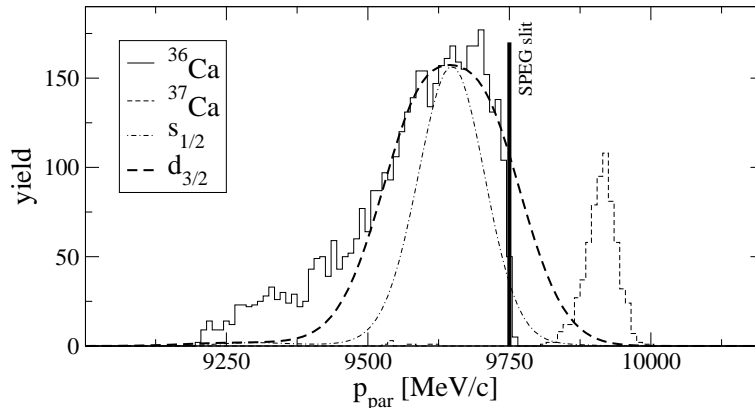


Fig. 3. Inclusive momentum distributions of ^{36}Ca and ^{37}Ca as measured in SPEG. The distribution for ^{36}Ca is cut by a slit that was installed in SPEG to suppress the secondary beam; in the dedicated run for ^{37}Ca this slit was not installed. Included are calculated momentum distributions for one-neutron removal from $d_{3/2}$ or $s_{1/2}$ states, folded with the distribution of the secondary beam.

1 The measured value for the energy of the first 2^+ state in ^{36}Ca is 266 keV
 2 lower than that in the mirror nucleus, ^{36}S . This is, besides ^{14}C - ^{14}O where
 3 the difference is 422(11) keV, one of the largest mirror energy differences
 4 observed so far for a first excited 2^+ state. Qualitatively, this might be
 5 explained as the combined effect of: (i) an almost pure ν nature of the 2^+
 6 state in ^{36}Ca due to the $Z = 20$ gap, (ii) an almost pure π nature of the 2^+
 7 state in ^{36}S due to the $N = 20$ gap, (iii) the almost pure 1-particle 1-hole
 8 configurations of the 2^+ states in ^{36}Ca and ^{36}S due to the large $Z, N = 16$
 9 gaps, and (iv) the Coulomb energy difference between typical s and d states.

10 Figure 3 shows the momentum distribution for ^{36}Ca and a comparison
 11 with calculated momentum distributions (9; 10; 11) as expected for neutron
 12 knock-out from the valence orbits $d_{3/2}$ and $s_{1/2}$. The width of the inclu-
 13 sive experimental momentum distribution fits well to the neutron knock-out
 14 from a $d_{3/2}$ state. From the integral of the extrapolated distribution, the
 15 number of ^{36}Ca ions was determined. Using the number of incident ^{37}Ca
 16 ions and the target thickness, a preliminary experimental cross section for
 17 the one-neutron removal $^{37}\text{Ca} \rightarrow ^{36}\text{Ca}$ of 5.3 (20) mb was obtained, while the
 18 calculated cross section is 18.6 mb assuming a knock-out from $\nu d_{3/2}$. This
 19 represents a quenching of $\approx 30\%$ similar to what has been found in the case
 20 of one-neutron knockout from ^{32}Ar , a nucleus which has a similarly large
 21 neutron separation energy (12).

22 We wish to thank the accelerator crew of GANIL for their smooth run-
 23 ning of the secondary beam production and the SPEG technical staff for

1 their support during the experiment.

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