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Hardness and toughness of sodium borosilicate glasses via Vickers's indentations



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ABSTRACT

This study investigates the mechanical response of sodium borosilicate (SBN) glasses as a function of their chemical composition. Vickers's indentation tests provide an estimate of the material hardness (H_V) and indentation fracture toughness (K_C^{VIF}) plus the amount of densification/shear flow processes. Sodium content significantly impacts the glass behavior under a sharp indenter. Low sodium glasses maintain high connected networks and low Poisson's ratios (ν). This entails significant densification processes during deformation. Conversely, glasses with high sodium content, i.e. large ν , partake in a more depolymerized network favoring deformation by shear flow. As a consequence, indentation patterns differ depending on the processes occurring. Densification processes appear to hinder the formation of half-penny median–radial cracks. Increasing ν favors shear flow and residual stresses enhance the development of half-penny median–radial cracks. Hence, K_C^{VIF} decreases linearly with ν .

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1. Introduction

Portable electronic devices frequently require thin lightweight glass used to protect the internal electronics. As such these glasses need to be resilient to external pressures. A common test to study these protective glasses is micro-indentation from which two important and standard measurements are extracted: (1) hardness (material's resistance to permanent deformation, H_V) and (2) indentation fracture toughness (material's resistance to fracture, K_C). These tests classify the glass's mechanical response into two groups: *anomalous* and *normal* behavior. *Anomalous* glasses predominantly densify under high external pressures. These glasses have a low atomic packing density; thus, the relative movement of the Si–O–Si linkage under pressure leads to the volume shrinkage [1–4]. On the other hand, *normal* behavior implies volume conserving shear flow. This is evidenced by a plastic flow generating pile-up of matter in the vicinity of the indentation without volume change [5–8]. Typically broken bonds and cations favor this phenomenon [9]. The degree at which a glass behaves *normally* and *anomalously* significantly depends on the glass' chemical composition [9–11].

Residual indentation patterns vary significantly with the chemical composition [12]. Typically, *anomalous* glasses exhibit cone crack,

whereas *normal* glasses predominantly exhibit radial–median cracks [9]. Previous, in-situ indentation studies of *normal* and *anomalous* glasses emphasize variations in deformation processes and effects in the residual stress levels [13]. Variations in contribution of densification versus shear flow alter the indentation shape, the crack appearance and the toughness measurements [5,11,13]. Furthermore, Hagan [7,8,14] highlighted that flow lines which appear in the indentation imprints can pile-up to produce seed cracks for median and radial cracks. In order to discriminate between shear flow and densification in glasses, researchers developed a simple test to estimate the amount of permanent densification under an indenter in glasses and the amount of plastic flow [15–18].

This paper investigates the mechanical response due to indentation in eight SBN glasses of modulated chemical composition. The glass's mechanical response depends on the glass structure. Imaging of the indents provides a means to obtain the hardness (H_V), the crack appearance probability (P_C) and the indentation fracture toughness (K_C^{VIF} , VIF implies Vickers's indentation fracture). To understand how matter flows beneath the indenter, AFM imaging before and after annealing discriminates between densification and shear flow mechanisms. For the reader's convenience, Appendix A provides a list of symbols, their meaning, and when appropriate the equation used to calculate them.

The following sections detail experimental techniques: (1) glass fabrication, (2) techniques used in understanding the glass properties (density, elastic moduli, MAS NMR); and (3) measuring and extracting information on the glass's response to Vickers's indenter. The Results

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Table 1

Target and measured ICP-AES molar compositions (where $[\cdot] \equiv \text{mol}\%$) of elaborated glass samples with their R_{SBN} and K_{SBN} values and physical properties: density (ρ); glass transition temperature (T_g); Young's modulus (E); Poisson's ratio (ν); $\langle \text{CN} \rangle$ is the mean coordination number of the boron atoms; the concentration of ^{14}B per volume unit deduced from NMR measurements and ICP-AES results and number of NBO per volume unit (N_{NBO}) deduced from Eq. (4). Target and ICP-AES measured molar composition and densities were previously published by Barlet et al. [19]. It was not possible to fabricate just one batch of glass to produce all samples; thus multiple batches were fabricated. All batches are within the 10% error of the ICP-AES measurement. Each SBN glass composition has an associated symbol in the subsequent figures. Glass's names won't be recalled on the figures below for clarity.

Name	Target values			Measured via ICP-AES			R_{SBN} measured	K_{SBN} measured	ρ ($\frac{\text{g}}{\text{cm}^3}$)	T_g ($^\circ\text{C}$)	E (GPa)	ν	$\langle \text{CN} \rangle$	$^{14}\text{B} \cdot 10^{21}$ (cm^{-3})	$N_{\text{NBO}} \cdot 10^{21}$ (cm^{-3})	Symbols
	$[\text{SiO}_2]$	$[\text{B}_2\text{O}_3]$	$[\text{Na}_2\text{O}]$	$[\text{SiO}_2]$	$[\text{B}_2\text{O}_3]$	$[\text{Na}_2\text{O}]$										
SBN 12	59.6	28.2	12.2	59.6	23.9	16.5	0.69	2.5	2.463	543	80.1	0.209 \pm 0.004	3.60	6.78	1.022	\triangle
SBN 25	50.7	23.9	25.4	52.6	20.6	26.8	1.30	2.5	2.545	535	80.3	0.238 \pm 0.001	3.70	7.22	5.91	\triangle
SBN 30	47.3	22.3	30.4	51	20.1	28.6	1.44	2.5	2.541	494	74.7	0.255 \pm 0.002	3.68	6.77	7.36	\triangle
SBN 35	44	20.6	35.4	46.9	18.6	34.5	1.85	2.5	2.537	467	76.7	0.264 \pm 0.0014	3.62	5.65	11.21	\triangle
SBN 14	67.8	18	14.2	70	15.8	14.2	0.89	4.4	2.474	588	81.8	0.212 \pm 0.004	3.72	5.49	1.34	*
SBN 63	63.2	16.8	20.0	66.7	14.1	19.2	1.35	4.7	2.524	573	81.9	0.226 \pm 0.001	//	//	//	\diamond
SBN 59	59.2	15.8	25	61.1	13.3	25.5	1.91	4.5	2.534	539	77.2	0.230 \pm 0.01	3.79	5.22	7.39	\bullet
SBN 55	55.3	14.7	30	58.0	12.9	29.1	2.25	4.5	2.538	505	72.8	0.251 \pm 0.006	3.76	4.86	7.49	\star

section presents structural properties, H_V , P_C , K_C^{VF} , and variations in contribution of densification and shear flow processes in the permanent deformation of sodium borosilicate (SBN) glasses. The Discussion section expounds the glass structure with their mechanical response to a Vickers's indenter. This part also compares and contrasts results presented in Sellappan et al. paper [18]. Furthermore, this section estimates the residual stresses induced during loading and after total unloading.

2. Experimental procedure

This section contains three subsections. It first describes the elaboration process of the glasses studied herein. Then, it details tests used to analyze the glasses' structural/material properties. Finally, it presents tests to understand the material response to microindentation.

2.1. Glass elaboration

Studies herein employ eight sodium borosilicate (SBN) glasses elaborated in-house [19]. During the elaboration process, manual homogenization of the silica (SiO_2), orthoboric acid (H_3BO_3), and sodium carbonate (Na_2CO_3) powders occurs. Platinum/gold (Pt/Au) crucibles retain the homogenized powder during the glass melting process. The formation of the glass melt undergoes three principle steps. Initially, the dehydration of the orthoboric acid takes place at 200 $^\circ\text{C}$ for 2 h. Next, the decarbonation of sodium carbonate Na_2CO_3 occurs at 800 $^\circ\text{C}$ for 3 h to avoid bubble formation. The final stage produces the glass melt. Depending on the glass composition, this stage occurs between 1100 $^\circ\text{C}$ and 1300 $^\circ\text{C}$ for 3 h. To avoid residual stress during the cooling process, the glass melt is transferred into a

preheated carbon crucible whose temperature is approximately T_g (glass transition temperature). Subsequently, the glass melt enters in a second furnace and cools at a slower rate (10 $^\circ\text{C}/\text{h}$) to release the residual stresses.

ICP-AES measurements (conducted by a third party, Prime Verre) verify the chemical compositions of the SBN glasses. Table 1 summarizes the target and measured values. ICP-AES measurements give approximately 10% error for each oxide. Several batches were fabricated to produce all samples. All of them are within the 10% error of the ICP-AES measurement. The glasses are classified depending on their $R_{\text{SBN}} = \frac{[\text{Na}_2\text{O}]}{[\text{B}_2\text{O}_3]}$ and $K_{\text{SBN}} = \frac{[\text{SiO}_2]}{[\text{B}_2\text{O}_3]}$ ratios.

2.2. Structural investigation

A glass's mechanical response is linked intrinsically to its structure. Thus, it is important to understand and quantify several glass parameters including density, elastic moduli, and the environment around the boron atoms.

2.2.1. Density, ρ

The densities of the glasses are estimated by Archimedes' principle. The geometry is a cylinder of thickness 10 mm and diameter 30 mm. Tests are conducted at ambient conditions using a hydrostatic balance. Initially, the glass samples are weighted in air (m_d) and water (m_w). Then by multiplying by the density of water (ρ_w) one can arrive at the density of the sample (ρ):

$$\rho = \frac{m_d}{m_d - m_w} \times \rho_w \quad (1)$$

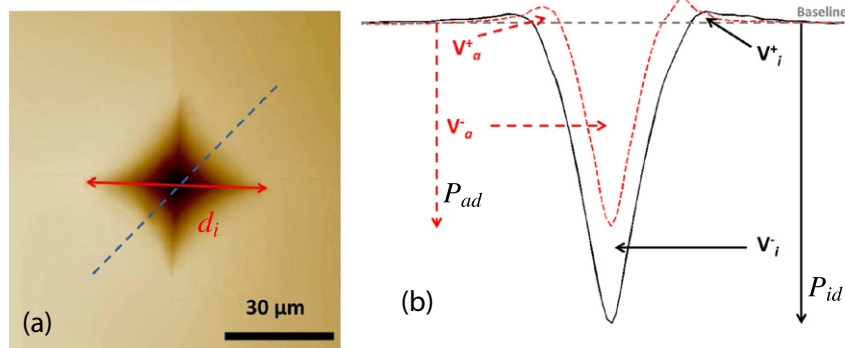


Fig. 1. (a) A typical indentation imprint used to determine the indentation diagonal length, (d_i and marked by a continuous line) and to estimate the pile-up profile (dotted line). (b) Sketch showing the evolution of indentation prints before (black solid line) and after annealing (red dotted line). V^+ and V^- represent the volume above and below the baseline (gray dotted line), respectively.

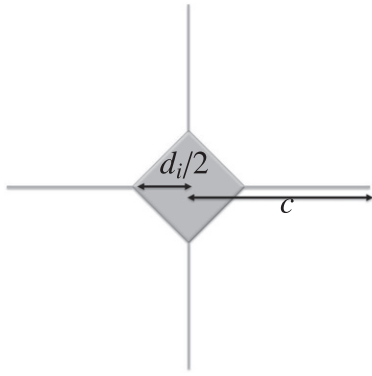


Fig. 2. Sketch of a Vickers's indent and notations used to define K_{IC}^{VIF} : d_i is the indentation diagonal length from one indent and c_i is the crack length measured from the indentation center.

2.2.2. Young modulus and Poisson's ratio

Ultrasonic echography techniques provide a means to estimate the elastic moduli and Poisson's ratio using a 5 MHz piezoelectric transducer. Tests invoke the same specimens as for the density tests. The specimen thickness is precisely determined using a digital micrometer with an accuracy of $\pm 1 \mu\text{m}$. The ultrasonic velocities of longitudinal (V_L) and transverse (V_T) acoustic waves are calculated from the thickness and transit time values. The Young modulus (E) and Poisson's ratio (ν) are then related to V_L and V_T via:

$$E = \rho * \frac{(3V_L^2 - 4V_T^2)}{(V_L/V_T)^2 - 1} \quad (2)$$

$$\nu = \frac{V_L^2 - 2V_T^2}{2(V_L^2 - V_T^2)} \quad (3)$$

Table 1 presents E and ν for the eight different SBN glasses.

2.2.3. ^{11}B MAS NMR

^{11}B magic angle spinning (MAS) nuclear magnetic resonance (NMR) spectra were collected on a Bruker AVANCE II 500WB spectrometer operating at a Larmor frequency of 160.14 MHz (magnetic field 11.72 T) using a 4 mm (outer diameter of ZrO_2 rotor) Bruker (boron-free) CPMAS probe. For each SBN composition, cylinder samples ($0.8 \times 2.96 \text{ mm}^2$) were spun at a frequency of 14 kHz, and a recycle delay of 2 s and a pulse length of $1 \mu\text{s}$ (tip angle of about 20°) were used to ensure quantitative data of the spectra. Data processing occurred via an in-house code (for details see reference [20]). ^{11}B magic angle spinning nuclear magnetic resonance (MAS NMR) experiments provide complementary information on the glass structure via revealing the boron environment. The high symmetry around ^{14}B

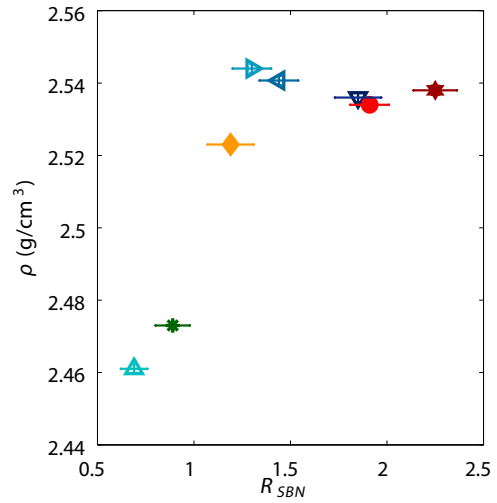


Fig. 3. Density, ρ , as a function of R_{SBN} .

leads to a SMALL quadrupolar coupling (0.2–0.4 MHz) in contrast to the trigonal ^{13}B unit (2.5–2.7 MHz). Thus, the boron environment reveals the proportions of ^{13}B and ^{14}B .

Two assumptions enable scientist to estimate the number of non-bridging oxygen atoms (note a NBO can be either Si or B atoms) in a SBN glass. First, each ^{14}B structural unit has one Na^+ ion attaching to it, and the Na^+ ion acts as a network compensator. Second, all the other Na^+ ions act as network modifiers. This gives one NBO per network modifying Na^+ ion. The equation to calculate this is:

$$N_{\text{NBO}} = N_{\text{Na}} - N_{^{14}\text{B}} \quad (4)$$

N_{Na} , N_{NBO} , and $N_{^{14}\text{B}}$ are the total number of Na^+ ions, the number of NBO, and the number of Na^+ acting as network compensators (equivalently the number of ^{14}B) in the system, respectively. **Table 1** presents these results.

2.3. Mechanical testing

This subsection details the experiments used to obtain the glass's mechanical response due to a Vickers's micro-indenter. The first subsection presents the geometry of the samples. It first presents the instrument and techniques used for the Vickers's indentations. Then, it details techniques for imaging imprints. The following sections detail how analyzing the images reveals the hardness, crack resistance, and indentation fracture toughness. Finally, the last subsection details the annealing procedures and how to extract the information concerning the densification and shear flow processes.

Table 2

Experimental values (P_{id} , V_i^+ (μm^3), V_i^- (μm^3), V_i^+/V_i^- (%), V_a^- , V_a^+ and P_{ad}) obtained from the AFM images of the indentation prints in the various SBN glasses. V_R and V_P present the quantification of the densification and the shear flow process from Eqs. (8) and (9), respectively. d_1 and d_2 record the asymmetry of the indenter; henceforth, these values will be replaced by d_i , the average of the two.

Name	Diagonals d_i (μm)		c (μm)	H_V (MPa)	C_R (g)	K_{IC}^{VIF} (MPa·m ^{1/2})	P_{id} (μm)	V_i^+ (μm^3)	V_i^- (μm^3)	V_i^+/V_i^- (%)	V_a^- (μm^3)	V_a^+ (μm^3)	P_{ad} (μm)	V_R	V_P
	d_1	d_2													
SBN 12	10.4	10.6	$37_{\pm 2}^{300 \text{ g}}$	$8200_{\pm 300}$	270	$0.98_{\pm 0.05}^{300 \text{ g}}$	1.012	5	36.7	13.6	4	14.4	0.665	52.8	16.3
SBN 25	12.0	12.2	$14.1_{\pm 0.6}$	$6340_{\pm 100}$	26	$0.69_{\pm 0.07}$	1.037	13	54.8	23.7	12	31.1	0.805	30.8	25.6
SBN 30	12.7	13.3	$15_{\pm 1}$	$5400_{\pm 200}$	32	$0.61_{\pm 0.06}$	1.19	17	62	27.4	16.3	47.2	0.912	12.6	28.6
SBN 35	13.5	13.7	$16.7_{\pm 0.8}$	$4900_{\pm 300}$	29	$0.56_{\pm 0.06}$	1.2	19.1	70.1	27.3	15.5	60.2	1.08	9	32.4
SBN 14	11.2	11.3	$11.9_{\pm 3}$	$7200_{\pm 300}$	45	$0.81_{\pm 0.08}$	0.88	7	35.2	19.8	10	26.1	0.822	34.3	11.3
SBN 63	12	11.9	$13_{\pm 3}$	$6400_{\pm 200}$	38	$0.74_{\pm 0.05}$	0.99	10.8	49.3	21.9	9	31.7	0.844	32	25.6
SBN 59	13.1	13.1	$13.9_{\pm 0.9}$	$5370_{\pm 90}$	37	$0.71_{\pm 0.07}$	1.03	12	55.3	21.7	15	41.7	0.846	30	16.3
SBN 55	14	13.25	$15.1_{\pm 4}$	$4900_{\pm 100}$	40	$0.65_{\pm 0.05}$	1.1	12	56.8	21.1	15	43.9	0.857	28	15.8

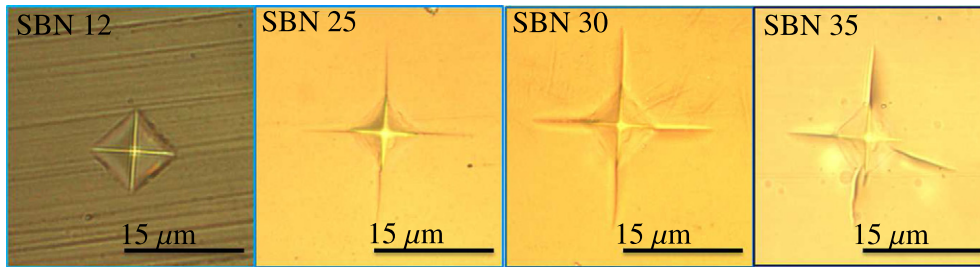


Fig. 4. Evolution of indentation prints for the $K_{SBN} \sim 2.5$ series (from left to right: SBN 12, SBN 25, SBN 30, and SBN 35) at 50 g in ambient conditions. An optical camera coupled with the Vickers's indenter captures the imprints within 5 min after indentation.

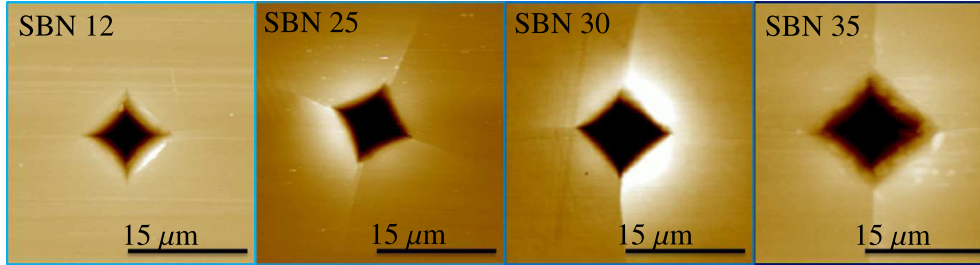


Fig. 5. Evolution of indentation prints for the $K_{SBN} \sim 2.5$ series (from left to right: SBN 12, SBN 25, SBN 30, and SBN 35) at 50 g in ambient conditions. An AFM scans the indents after indentation.

2.3.1. Sample preparation

Indentation experiments require samples with two parallel surfaces. Samples herein are rectangular parallelepipeds of size $5 \times 5 \times 25 \text{ mm}^3$. Indentations take place on one of the $5 \times 25 \text{ mm}^2$ faces. Surface quality affects the crack behavior of glasses; therefore diamond grinding disks optically polish the samples.

2.3.2. Mechanical testing by Vickers's indentation

Vickers's microindentation tests customarily give sample hardness (H_V) and indentation fracture toughness (K_{IC}^{VIF}). Indentations occur at $22 \pm 1 \text{ }^\circ\text{C}$ in air (ambient conditions for humidity) with an Aton Paar MHT-10 hardness tester. Before each set of indents, an Aton Paar MHT-10 hardness tester underwent a series of tests on standard samples to ensure proper regulations. The indentation load (P), varies from 25 g to 300 g. The dwelling time at maximum load is 15 s for all tests. Experimental conditions remain the same during the tests.

2.3.3. Image acquisition of indentation imprints

An atomic force microscopy (Dimension Icon Nanoscope V Bruker) images the indentation imprints in PeakForce Tapping mode. The AFM cantilever (Bruker's RTESPA) is a silicon tip with a nominal radius of 8 nm and nominal force constant of 40 N/m. The scan size is sufficiently large with respect to the imprint size to permit correcting the AFM image planarity. Typically, the image sizes are greater than twice the indentation diagonal. The estimated errors for measuring the length by AFM are less than 2% in the X and Y directions and about 2% in the Z direction. Post-image analysis reveals 4 parameters per indent (Fig. 1):

1. Residual indentation depth: P_{id}
2. Indentation diagonal length: d_i
3. Indentation volume: V_i^- , volume measured below the free surface
4. Volume of pile-up: V_i^+ , volume measured above the free surface.

2.3.4. Hardness determination

Vickers's hardness measures the material ability to resist permanent deformation induced by a harder material. Analyzing the indentation

imprints reveals a permanent deformation due to the Vickers's indenter at specific loads. Knowing the diagonal length of the imprint indentation (d_i , Fig. 1, a) and P , the Vickers's hardness is:

$$H_V = \frac{1.8544P}{d_i^2}. \quad (5)$$

At least 10 indentations occur at 50 g (0.49 N) maximum load. Indentations which do not form well-developed cracks are rejected.

2.3.5. Crack resistance, C_R

Vickers's indentation can cause cracks to initiate off the corners of the indenter. The probability of crack appearance, P_C , is the average number of radial cracks per corner for a given load. This is conducted

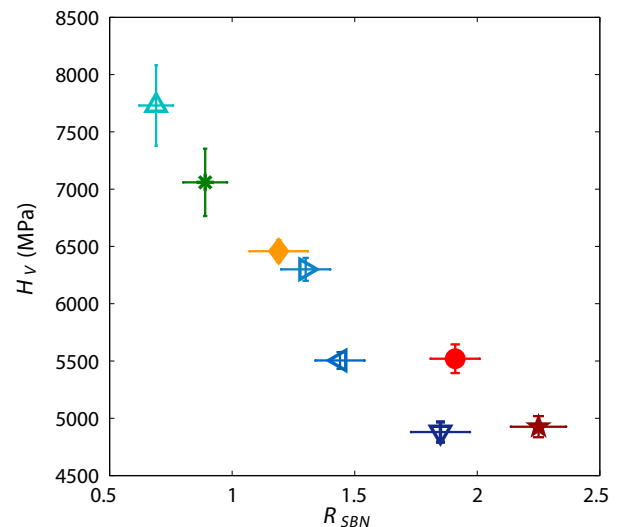


Fig. 6. Hardness values, H_V , versus R_{SBN} for SBN glasses (at 50 g during 15 s in ambient conditions).

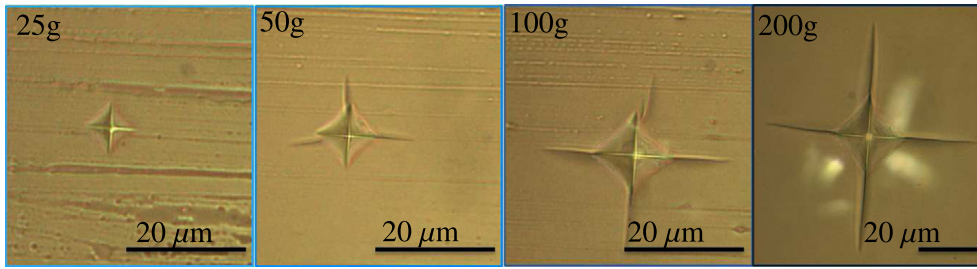


Fig. 7. Set of optical images in SBN 14 acquired at increasing loads. As the load increases, the number of cracks initiating off the corners of the indenter increases.

at multiple loads. The load where P_C exceeds 50% defines the crack resistance, C_R [21].

2.3.6. K_C measurements from indentation

Fracture toughness, K_C , defines the material's resistance to failure under vacuum. Herein, estimates for K_C arise by studying the length of the cracks generated off the corners of the indenter (see Fig. 2). (Note: Henceforth, K_C arising from indentation studies will be referred to as K_C^{IND}) Many models exist in literature to estimate K_C^{IND} by indentation [12,22–31]. There are two predominate models: half-penny median–radial cracks ($c > 1.25d$; c being the average distance ($\langle c_i \rangle$) from the indentation center to the crack tip and d the mean of $\langle d_i \rangle$ values obtained for one indent) and Palmqvist cracks ($c < 1.25d$).

Lawn's et al. model assumes well-developed half-penny median–radial cracks [12,23]. These cracks emerge when the load is sufficiently high insuring $c > 1.25d$. Palmqvist cracks, on the other hand, exists when $c < 1.25d$. These cracks extend from opposite indentation corners; however, they do not connect to each other [32]. In general, half-penny median–radial cracks models follow the following relationship [12,23,32,33]:

$$K_C^{IND} = \alpha \left(\frac{E}{H_V} \right)^n \frac{P}{(c)^{3/2}} \quad (6)$$

where α is a calibration constant. n depends on the model. Standard values of n are 0 [27], 0.4 [26,32], 0.5 [28], and 2/3 [12]. Weber et al. uses Lawn's ASTM report to arrive at the following equation:

$$K_C^{VIF} = 0.057 \times H_V \times \sqrt{d/2} \left(\frac{E}{H_V} \right)^{0.4} \times \left(\frac{c}{(d/2)} \right)^{-3/2} \quad (7)$$

This is predominantly valid for well-defined indentation patterns with crack lengths fulfilling the condition $c > 1.25d$. However, Weber suggests that it is also acceptable for Palmqvist cracks as it resembles

Lankford et al. equation [29,30,32]. This is favorable as SBN 12 does not fulfill the condition of $c > 1.25d$ at 50 g (0.49 N); yet, all other samples meet this requirement. If the indents do not have 4 cracks, the indentation pattern is not included in the averages. (It is note worthy that error bars in the figures correspond to one standard deviation of the mean values [34].) A force of 50 g will be invoked for all K_C^{VIF} measurements (with an exception for SBN 12 where higher loads are applied). Due to stress corrosion cracking (ambient conditions for humidity), c can vary with time after indentation [28]. To avoid substantial errors in the K_C^{VIF} evaluation, an optical microscope images the surface crack pattern within 5 min after indentation.

2.4. Complementary measurements of Vickers' indentations: densification and shear flow processes

The high pressures induced by the sharp indenter yields both reversible (elastic) response and irreversible (plastic) deformations. The irreversible component is mainly due to densification or volume conserving shear flow [5,16,35–39]. Yoshida et al. [16] propose a 3D technique to estimate the fraction of densification from the comparison of the indentation imprints before and after the annealing: Upon annealing, the indent shape will change due to the recovery of densified areas [16,37,40,41]. This procedure has recently been extended by Sellappan et al. [18] to quantify the amount of volume conserving shear flow (V_p). The following subsections present the methods applied to quantify first the densification and subsequently the isochoric shear flow process.

2.4.1. Densification estimation

Studying the volume ratio before and after annealing provides an easy way to estimate the amount of densified volume. As revealed above, post-processing of AFM scans of the indent imprints reveals 4 parameters (Fig. 1): P_{id} , d_i , V_i^- , and V_i^+ . Subsequently, annealing the samples at $0.9 T_g$ for 2 h in air relieves a significant proportion of the densified regions [16,37,40,41]. After annealing, imprints are imaged.

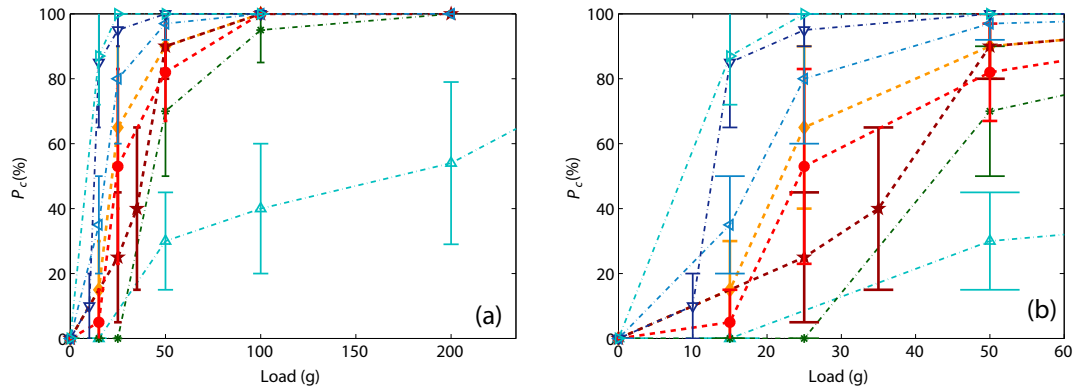


Fig. 8. (a) Crack appearance probability for all the SBN glasses $K_{SBN} = 2.1$ series (blue) and $K_{SBN} = 3.7$ series (red) at different loads. (b) A zoom of the crack appearance probability for all the SBN glasses $K_{SBN} = 2.1$ series (blue) and $K_{SBN} = 3.7$ series (red) for loads between 0 g and 50 g.

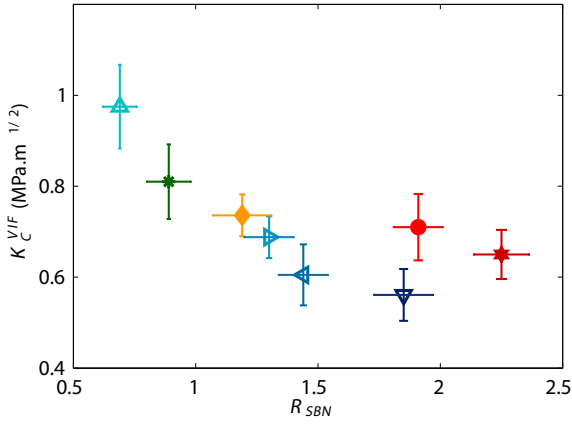


Fig. 9. Toughness values K_C^{VIF} of SBN glasses depending on R_{SBN} for $K_{SBN} \sim 2.5$ series (blue) and $K_{SBN} \sim 4.5$ series (red).

Post-treatment of the images reveals d_a , P_{ad} , V_a^- and V_a^+ (Fig. 1). Table 2 presents these parameters. The densified volume is then estimated through the recovered volume ratio (V_R) of the indent below the surface:

$$V_R = \frac{(V_i^- - V_a^-) + (V_a^+ - V_i^+)}{V_i^-} \quad (8)$$

2.4.2. Shear flow

Shear flow represents a plastic flow generating a displacement of matter without volume change. The pile-up around the indents exposes the amount of shear flow. Profiles extracted from Fig. 1 follow the dotted blue line. Sellappan et al. [18] propose V_p to study the volume-conserving flow ratio:

$$V_p = \frac{(2V_i^+ - V_a^+)}{V_i^-} \quad (9)$$

$(V_a^+ - V_i^+)$ represents the densified volume contribution in the pile-up area. The denominator in Eq. (9) normalizes the shear flow.

3. Results

This section details the structural and mechanical property investigations. First, this section presents results concerning the structural

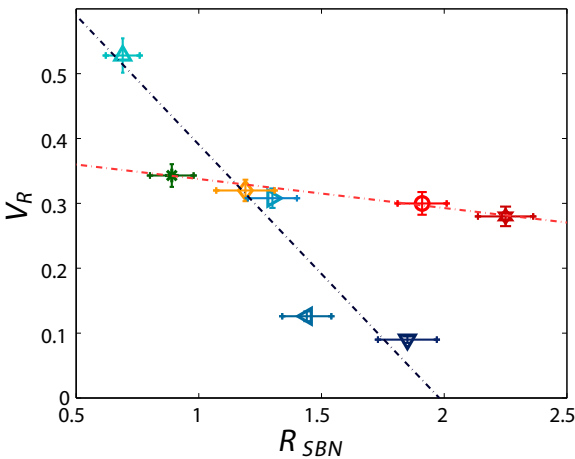


Fig. 10. The recovery volume decreases with R_{SBN} . Yet, the slopes ($m_{K_{SBN} \sim 2.5} = -0.4$ and $m_{K_{SBN} \sim 4.5} = -0.04$) depend significantly on K_{SBN} .

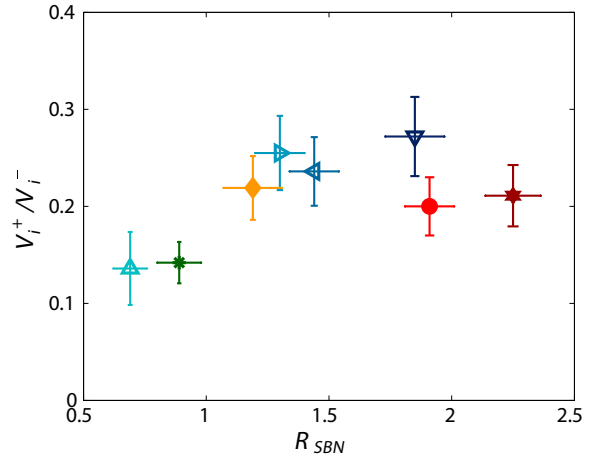


Fig. 11. The normalized ratio (V_i^+ / V_i^-) of the glasses pile-up with R_{SBN} displays two trends depending on K_{SBN} . For $K_{SBN} \sim 2.5$, the ratio increases initially then appears to stabilize. For $K_{SBN} \sim 4.5$, the ratio decreases slightly.

properties of SBN glasses. Following this, the section provides results concerning the indentation. The end of the section presents the variations of V_R and V_p .

3.1. Structural investigation

3.1.1. Density, ρ

Archimedes' principle gives the density, ρ , for the different glasses. Fig. 3 shows how ρ varies with $R_{SBN} = \frac{[Na_2O]}{[B_2O_3]}$ (where $[\cdot] \equiv \text{mol\%}$). Previous literature details a complex 3D analysis concerning density's dependence on $K_{SBN} = \frac{[SiO_2]}{[B_2O_3]}$ and R_{SBN} [19,42–44]. When possible the same batches of glasses were used herein as in the Barlet et al. [19] study.

3.1.2. Elastic moduli, E and ν

As mentioned in Subsection 2.2.2, E and ν result from both measurements of the longitudinal wave (V_L) and of the transverse wave (V_T) speeds. Appendix B contains these raw measurements. Eqs. (2) and (3) give way to the estimates of E and ν . Table 1 presents the results of E and ν . Clearly E and ν depend on the chemical composition of the SBN glasses. Conversely, E and ν do not vary monotonically with R_{SBN} .

3.1.3. ^{11}B MAS NMR

Table 1 presents results on the amount of ^{14}B in each SBN sample. The chemical composition of silicate glasses significantly affects the ^{14}B count [45,46].

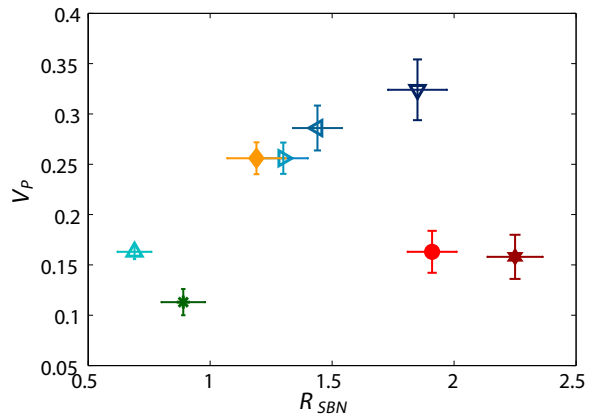


Fig. 12. The pile-up ratio (V_p) which is related to the volume-conserving shear flow as a function of R_{SBN} . The tendencies differ for approximately constant K_{SBN} .

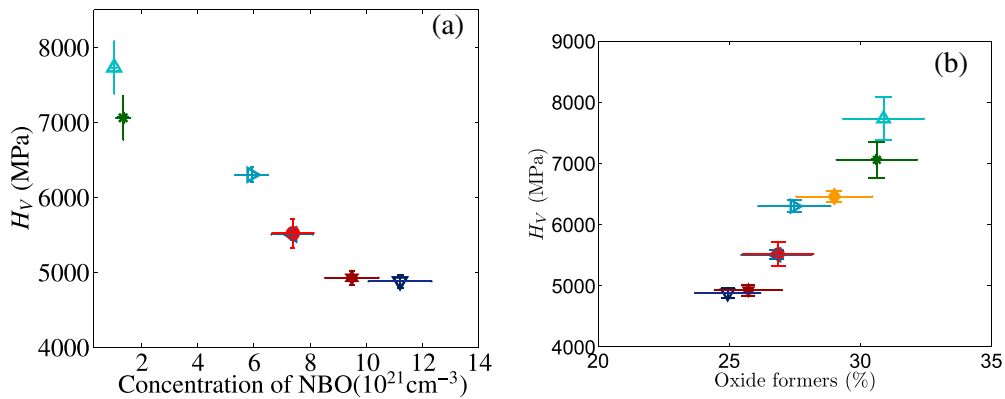


Fig. 13. (a) Evolution of hardness, H_V (at 50 g in ambient conditions), versus the concentration of non-bridging oxygen (NBO) deduced from NMR measurements. (b) Evolution of hardness, H_V (at 50 g in ambient conditions), versus percentage of oxide formers determined from ICP measurements.

3.2. Mechanical testing results

Immediately after unloading the Vickers' indenter, an optical microscope coupled with the Vickers's indenter visually captures the features of the indents. Fig. 4 displays indentation imprints for the $K_{\text{SBN}} \sim 2.5$ series (from left to right: SBN 12, SBN 25, SBN 30, and SBN 35). A clear discrepancy appears according to the crack behavior. For low amounts of sodium (i.e. SBN 12) at 50 g, no cracking occurs. On the other hand, cracks emerge at the surface in glasses with higher sodium content. In addition, the light reflection seen in SBN 35 glass can be a sign of pile-up.

Fig. 5 displays indentation prints for the $K_{\text{SBN}} \sim 2.5$ series (from left to right: SBN 12, SBN 25, SBN 30, and SBN 35) at 50 g in ambient conditions via an AFM. Subsequently, these AFM images aid in enumerating P_{id} , d_i , V_i^- and V_i^+ . Table 2 presents the average results of at least 10 different indents.

3.2.1. Hardness, H_V

Knowing the maximum load applied ($P = 50 \text{ g}$) during the indent along with d_i , Eq. (5) gives H_V for each sample. H_V results herein represent the average of at least 10 tests. Table 2 presents the obtained values. Fig. 6 presents the evolution of H_V (MPa) as a function of R_{SBN} . The chemical composition of silicate glasses significantly affects the hardness values. For constant $K_{\text{SBN}} = \frac{[\text{SiO}_2]}{[\text{B}_2\text{O}_3]}$ and increasing R_{SBN} , the hardness values decrease. The decrease is less pronounced for

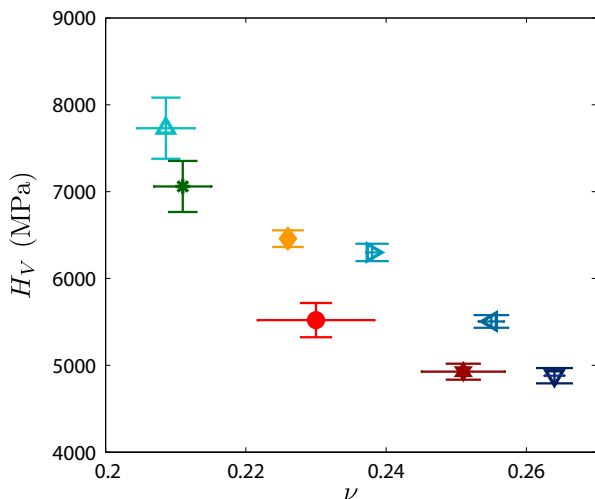


Fig. 14. Evolution of hardness values as a function of the Poisson ratio (ν) (indentation at 50 g in ambient conditions).

$K_{\text{SBN}} \sim 4.5$ (i.e. samples with more SiO_2). This decrease in H_V corresponds to an increasing residual penetration depth (P_{id}) as determined from AFM imaging (see Table 2). As the sodium content increases, the indenter penetrates deeper into the glass.

3.2.2. Crack resistance

Fig. 7 displays optical images of SBN 14 at different loads. As seen in the figure, at low loads ($P \leq 25 \text{ g}$) cracks do not propagate off the corners of the Vickers' indenter. The average number of cracks initiating off the corners of the Vickers' indenter increases between 25 g and 200 g. At 200 g, all four corners have, on average, one crack.

The probability of crack appearance, P_C , is the average number of cracks per corner out of at least 10 indents at a given load. Fig. 8 exemplifies the probability of crack appearance in SBN 14 (dark green stars) along with the other SBN samples. All SBN glasses, except SBN 12, rapidly increase their crack appearance probability at low loads. 100% P_C never occurred for SBN 12 due to limitations of the experimental setup (i.e. higher loads were not feasible with the hardness tester).

As stated in the previous section the crack resistance, C_R , is the load at which the probability of crack appearance equals 50% (estimated by looking at when the data in Fig. 8 crosses the 50% line). Table 2 presents C_R for the different SBN samples. SBN 12 is significantly different from the other glasses.

3.2.3. Indentation fracture toughness, K_C^{VIF}

Section 2.3.6 details the requirements for radial/median cracks (i.e. $c > 1.25d$). The formation of radial/median cracks permits the use of Evans and Marshall's model. On the other hand, to compare the K_C^{VIF} values in the different glasses, the maximum load needs to be the same in all glasses. For almost all the samples, reasonable indentation imprints plus sufficient energy to propagate medial/radial cracks ($C_R > 50\%$) occur at 50 g. The only exception herein is SBN 12, which neither forms well-developed cracks nor satisfies the criterion $c > 1.25d$ at 50 g. Hence, Evans and Marshall's model applies well to all samples but SBN 12. On the other hand, comparing and contrasting K_C^{VIF} results estimated through different models also lead to discrepancies. Thus, results herein all use the Evans and Marshall's model, but care should be taken when examining SBN 12 results.

Fig. 9 shows that K_C^{VIF} decreases as the sodium content increases for constant K_{SBN} . It should be noted that K_C^{VIF} for SBN 12 probably overestimates the correct value due to the lack of well-developed cracks.

3.3. Complementary measurements of Vickers' indentations

The high pressure under the indenter can cause material compaction [2–4,7,14]. The original structural configuration can be recovered if sufficient activation energy is supplied to the material (for instance by heating) as demonstrated by Bridgman and Simon [38] and later by

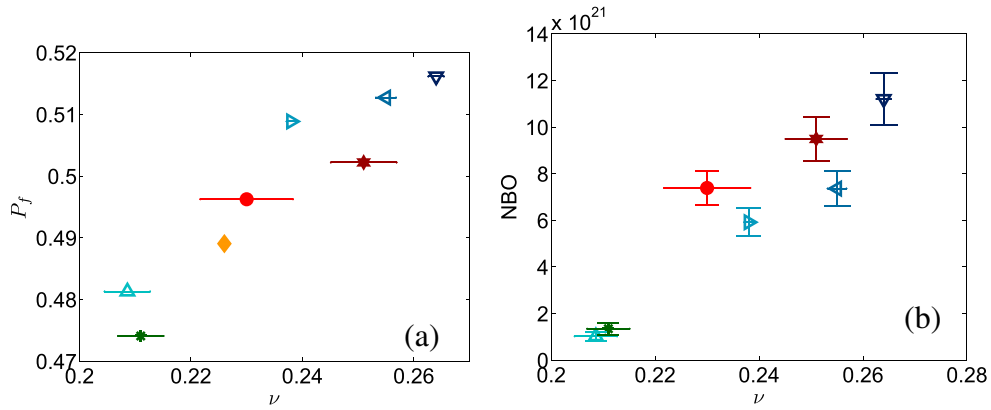


Fig. 15. (a) Evolution of glass packing density (P_f) versus the Poisson's ratio (ν); (b) evolution of the concentration of NBO from NMR measurements as a function of ν .

Mackenzie [15]. Table 2 presents parameters acquired from AFM images before (d_i, P_{id}, V_i^- and V_i^+) and after (d_a, P_{ad}, V_a^- and V_a^+) annealing. These values are the average values of at least ten indents. V_R from Eq. (8) provides a means to calculate the amount of volume recovered due to annealing. This indicates the amount of densification due to an indenter. Table 2 reveals the average V_R for the various SBN samples. It is well noted that annealing may not relieve all regions of densified material due to geometrical constraints. In $K_{SBN} \sim 4.5$ the residual stress field can lead to the development of cracks during unloading [18]. Confirming this hypothesis requires further studies.

Fig. 10 unveils two scenarios depending on the value of K_{SBN} . For samples with low K_{SBN} values, the recovery volume decreases roughly linearly with R_{SBN} . ($V_R = m_{K_{SBN} \sim 2.5} R_{SBN} + b_{K_{SBN} \sim 2.5}$ where $m_{K_{SBN} \sim 2.5} = -0.409 \pm 0.005$ and $b_{K_{SBN} \sim 2.5} = 0.80 \pm 0.002$). On the other hand, the $K_{SBN} \sim 4.6$ series exhibits a significantly less pronounced decrease but approximately linearly ($V_R = m_{K_{SBN} \sim 4.6} R_{SBN} + b_{K_{SBN} \sim 4.6}$ where $m_{K_{SBN} \sim 4.6} = -0.0431 \pm 0.003$ and $b_{K_{SBN} \sim 4.6} = 0.38 \pm 0.005$).

In addition to densification processes, pile-up occurs. The glasses' pile-up (V_i^+) response depends on the chemistry of the glass. Fig. 11 presents the normalized ratio (V_i^+/V_i^-) of the pile-up as a function of R_{SBN} . Variations in (V_i^+/V_i^-) convey the glass's tendency to pile-up around the indenter under pressure [47]. Fig. 11 unveils two different scenarios depending on the value of K_{SBN} . For $K_{SBN} \sim 2.5$, the ratio

initially increases and subsequently stabilizes. For $K_{SBN} \sim 4.5$, the ratio decreases, and then remains stable within the error bars. There is a chance that it too increases; however, determining this requires smaller R_{SBN} values.

Within the pile-up region, there is a chance that densification occurs. Eq. (9) permits the removal of the densified portion of the pile-up by imaging before and after annealing. Annealing relieves densified areas this highlighting variations in isochoric shear flow. Fig. 12 depicts how V_p varies with R_{SBN} for the different series of glasses. For $K_{SBN} \sim 2.5$, the ratio increases; yet, for $K_{SBN} \sim 4.6$, the ratio decreases.

A possible explanation as to why SBN 55 and SBN 59 decrease is that annealing may not fully relieve densified regions, and it may also induced some plastic flow. Confirming this hypothesis requires more research.

4. Discussion

The experimental observations reported in the previous section can be qualitatively understood by invoking Dell and Bray's scenario [43,45]. SBN glasses with a low concentration of sodium ions ($R_{SBN} < R_{max}^{SBN} = 0.5 + K_{SBN}/16$) consist predominantly of oxide formers (Si and B atoms). The small amount of Na^+ available in the glass transforms fully coordinated ^{13}B into fully coordinated ^{14}B [43,45]. When $R_{SBN} > R_{max}$, the extra Na^+ begins to form NBOs initially in the Si network and then in the B network. A side effect to the formation of the NBOs is the reversal of fully coordinated ^{14}B units to ^{13}B units with at least one NBO. Table 1 confirms the decrease in ^{14}B units as $[Na_2O]$ increases (or as R_{SBN} increases for constant K_{SBN}).

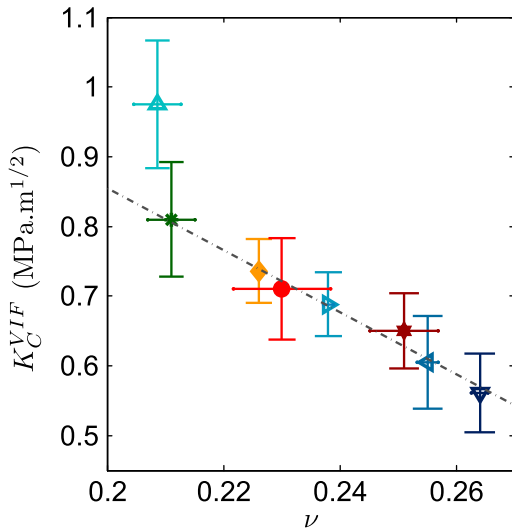


Fig. 16. Toughness values (K_C^{VIF}) of all SBN glasses as a function of ν . The line ($K_C^{VIF} = m\nu + b$ where $m = -4.46 \pm 0.27$ and $b = 1.75 \pm 0.06$) depicts a linear fit through all of the data excluding SBN 12.

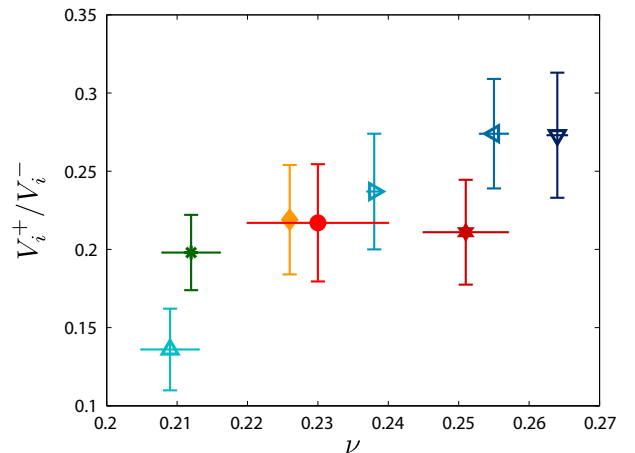


Fig. 17. V_i^+/V_i^- versus ν before annealing the sample.

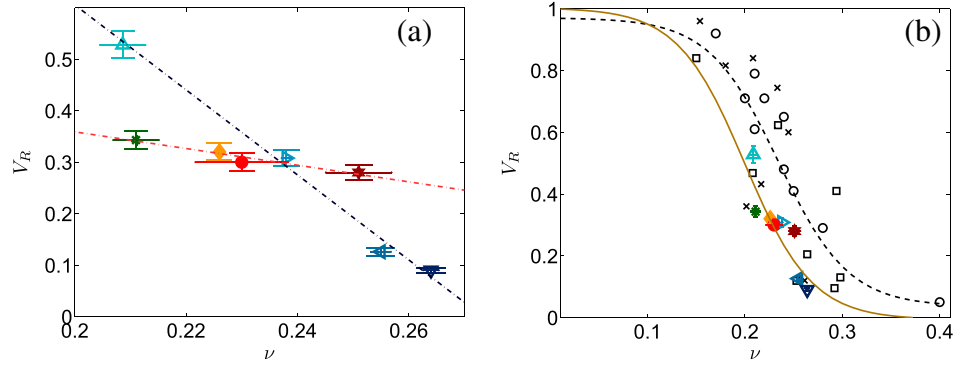


Fig. 18. (a) The figure depicts the evolution the volume recovery ratio (V_R) versus the Poisson's ratio (ν) for SBN data. The blue color line depicts a linear fit to the data for $K_{SBN} \sim 2.5$ ($V_R = m_{K_{SBN} \sim 2.5} \nu + b_{K_{SBN} \sim 2.5}$ where $m_{K_{SBN} \sim 2.5} = -8.3 \pm 0.6$ and $b_{K_{SBN} \sim 2.5} = 2.26 \pm 0.13$). The red color line depicts a linear fit to the data for $K_{SBN} \sim 4.5$ ($V_R = m_{K_{SBN} \sim 4.5} \nu + b_{K_{SBN} \sim 4.5}$ where $m_{K_{SBN} \sim 4.5} = -1.6 \pm 0.3$ and $b_{K_{SBN} \sim 4.5} = 0.68 \pm 0.07$). (b) The figure depicts a comparison of V_R versus ν relationships found herein along with published data, black data points [16,18,55]. The black dashed line depicts the sigmoid fit proposed by Sellappan et al. The brown solid fit depicts the best sigmoid fit, $(\alpha_R, \beta_R, \chi_R, \hat{\nu}_R) = (.99, .003, 29, -.007)$, for the SBN data assuming 2 sduo points: (0, 1) and (.5, 0).

4.1. Linking hardness (H_V) to the glass structure

The results in Section 3.2.1 reveal a decrease in the hardness of the SBN glass as R_{SBN} increases (Fig. 6). Linking this variation to the structure is the key in understanding the glasses' behavior. High levels of $[\text{Na}_2\text{O}]$ reduce the SBN glasses resistance to Vickers's indentation and forms NBO on the network formers when $R_{SBN} > R_{\text{max}}$. This displays an inverse correlation between H_V and NBO in the glassy network (Fig. 13).

NBO reduces the connectivity of the glass, which in turn alters the short- and medium-range order. Poisson's ratio (ν) provides an interesting insight into the glass's short- and medium-range order. The Poisson's ratio is the negative of the ratio between the tensile and lateral strain. The SBN glasses presented herein exhibit a range of values from 0.21 to 0.27 (Table 1). H_V decreases almost linearly with ν (Fig. 14). When ν is low (i.e. the glass network is reticulate), H_V is high.

As stated above, additional $[\text{Na}_2\text{O}]$ alters both the number of NBOs and ν . Moreover, the NBOs concentration is linearly linked to ν . It also affects the packing fraction (P_f) which is a measure of how densely the system is packed [48,49]. The ratio between the minimal theoretical density (ρ_t) and the actual density of the glass (ρ) gives P_f :

$$P_f = \frac{\rho}{\rho_t} \quad (10)$$

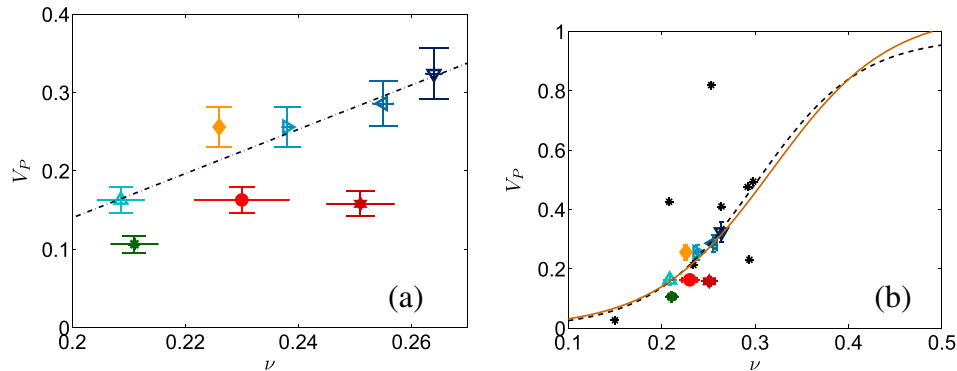


Fig. 19. (a) The figure depicts the evolution of V_P versus the ν for SBN data. The purple color line depicts a linear fit to the data for $K_{SBN} \sim 2.5$ ($V_P = m_{K_{SBN} \sim 2.5} \nu + b_{K_{SBN} \sim 2.5}$ where $m_{K_{SBN} \sim 2.5} = 0.30 \pm 0.022$ and $b_{K_{SBN} \sim 2.5} = -0.42 \pm 0.05$). For $K_{SBN} \sim 4.5$, the data does not fit a linear extrapolation. (b) The figure depicts a comparison of V_P versus ν relationships found here in along with published data [16,18,55]. The black dashed line depicts the sigmoid fit proposed by Sellappan et al. The beige color fit depicts the best sigmoid fit, $(\alpha_R, \beta_R, \chi_R) = (.94, 152, -16)$, for the SBN data herein assuming 2 sduo points: (0, 0) and (.5, 1).

where

$$\rho_t = \frac{M}{V} = \frac{\sum_{i=1}^n (f_i M_i)}{\sum_{i=1}^n (f_i V_i)} \quad (11)$$

Here, the sums occur over the different oxides in the systems. f_i is the molar fraction of each oxide. For oxide $[\text{A}_x\text{O}_y]$, the ideal volume is [50]:

$$V_i = \frac{4\pi N_A}{3} (x r_A^3 + y r_O^3) \quad (12)$$

N_A is Avogadro number. r_A and r_O are the ionic radii of the cation and anion oxides respectively. Fig. 15 reveals P_f increases as ν increases. For low $[\text{Na}_2\text{O}]$ concentrations and ν , the glass structure has a large fraction of free volume (i.e. low glass packing density). On the other hand, increasing the $[\text{Na}_2\text{O}]$ concentrations gives higher ν and high P_f ; consequently the free volume decreases in these glass structures. Moreover, H_V decreases.

Thus, it appears that the atomic bond strength and the reticulation of the glassy network control the behavior upon indentation. These results corroborate Kilymis et al. recent studies on similar glasses. MD simulations on similar glasses reveal that H_V decreases as the number of ^{13}B and NBOs increases, as observed here [51,52].

H_V sheds light on a material's resistance to permanent deformation. Yet, two glasses with the same H_V may exhibit significantly different overall behaviors upon indentation: resistance to cracking, pile-up, etc. The subsequent parts of this section address these differences.

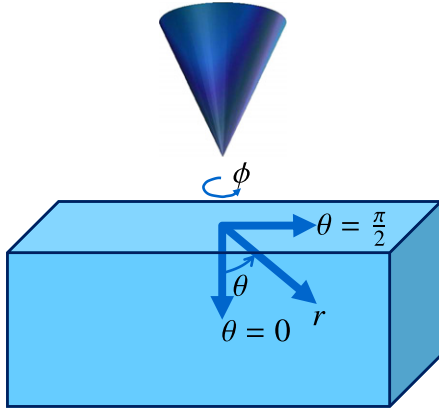


Fig. 20. Depicts spherical polar coordinates (r, θ, ϕ) used in Yoffe's stress field equations for a conic indenter [56].

4.2. Linking the resistance to cracking C_R to the glass structure

As observed in Fig. 8, the C_R value for SBN 12 is significantly higher than C_R for other SBN glasses. Table 2 provides C_R values for the SBN series. Crack appearance depends on how the glass deforms under the indenter. Processes occurring during indentation do not favor growth of median/radial cracks in SBN 12 as observed for other SBN glasses. Thus, complex processes such as densification and shear flow are non-negligible for C_R estimates. Section 4.4 examines densification and shear flow processes under the indenter.

4.3. Linking the indentation fracture toughness, K_C^{VIF} , to the glass structure

Over the past several decades, Evans and Marshall's model (Eq. (7)) and other models provided a method to estimate the indentation fracture toughness of brittle materials (especially ceramics). Yet, these techniques are essentially linked to the occurrence of a plastic zone beneath the contact area and isochoric deformation. Figs. 10, 11, and 12 all reveal varying deformation mechanisms in the SBN ternary systems as the composition changes. This calls for caution in interpreting indentation cracking patterns. Nevertheless, in the absence of a better indentation tool to probe the cracking resistance, Evans and Marshall's model (Eq. (7)) can be used for comparison purposes.

Through Evans and Marshall's model, Vickers's indentation provides a method to estimate K_C^{VIF} . Fig. 9 reveals a decreasing trend of K_C^{VIF} with R_{SBN} . Yet, the trend depends on the K_{SBN} series, which implies that the underline structure of the glass should play an important role. Conversely, the variations of K_C^{VIF} with ν , plotted in Fig. 16 reveal a fairly good collapse. A notable exception is SBN 12, which as already mentioned in Section 3, does not develop well defined cracks. Densification under the indenter is a possible source of the discrepancy (to be studied in Section 4.4). Otherwise, Fig. 16 reveals K_C^{VIF} decreases linearly with ν ($K_C^{VIF} = m\nu + b$ where $m = -4.46 \pm 0.27$ and $b = 1.75 \pm 0.06$).

Table 3

Sellappan et al. [18] proposed these parameters for Yoffe's stress field. B is the strength of the Blister field; and β is Sellappan's blister field strength. V_i^m is an estimate of the indent volume disregarding elastic recover from assumptions made by [18]. a is an estimate of the half diagonal of the projected surface area of an indent, Ψ is the angle of the indenter.

Variable	symbol	function
Strength of blister field	B	$\frac{3E}{4\pi(1-\nu^2)(1+\nu)}\beta V_i^m$
Sellappan's blister field strength	β	$(1 - \nu_R - \nu_P)$
Volume of indent disregarding elastic recovery	V_i^m	$\frac{2a^3}{3 \tan(\Psi)}$
half the diagonal of the projected surface area	a	$\left(\frac{P}{2H}\right)^{0.5}$
apical angle of the indenter	Ψ	70.3

Table 4

Table of predominate opening mode cracks due to a conic indenter and their corresponding stresses [56,18]. Stresses are normalized by the H of a conical indenter.

Type of crack	Driving stress field	θ	Boussinesq term	Blister term
Ring	σ_{rr}/H	$\pi/2$	$0.5 - \nu$	$\frac{2(\nu-2)\beta}{\pi(1+\nu)(1-2\nu)\tan(\Psi)} \frac{E}{H}$
Radial	$\sigma_{\phi\phi}/H$	$\pi/2$	$\nu - 0.5$	$\frac{2\beta}{\pi(1+\nu)\tan(\Psi)} \frac{E}{H}$
Median	$\sigma_{\theta\theta}/H$	0	$0.25 - 0.5\nu$	$-\frac{\beta}{\pi(1+\nu)\tan(\Psi)} \frac{E}{H}$
Lateral	σ_{rr}/H	0	-1.5	$\frac{6\beta}{\pi(1+\nu)(1-2\nu)\tan(\Psi)} \frac{E}{H}$

4.4. Contribution of densification and shear flow to the permanent deformation

Beyond H_V and K_C^{VIF} , information on the processes occurring during indentations can be extracted from analyzing post-indentation AFM images. For example a good correlation is found between V_i^+/V_i^- and ν (Fig. 17). This shows that pile-up processes favor glasses with a lower degree of polymerization.

Current literature suggests that the chemical composition impacts more V_R than H_V [16,17,53,54]. Sellappan et al. [18] evidenced a sigmoid relation between ν and V_R . As ν increases V_R decreases, thus the contribution of densification occurring in the system decreases.

Fig. 18 unveils two different scenarios when looking at approximately constant K_{SBN} . For samples with $K_{SBN} \sim 2.5$, the recovery volume decreases with ν ($V_R = m_{K_{SBN} \sim 2.5} \nu + b_{K_{SBN} \sim 2.5}$ where $m_{K_{SBN} \sim 2.5} = -8.3 \pm 0.6$ and $b_{K_{SBN} \sim 2.5} = 2.26 \pm 0.13$). As the sodium content increases for low $K_{SBN} \sim 2.5$, V_R decreases quickly. On the other hand, the $K_{SBN} \sim 4.5$ series exhibits a significantly less pronounced decrease ($V_R = m_{K_{SBN} \sim 4.5} \nu + b_{K_{SBN} \sim 4.5}$ where $m_{K_{SBN} \sim 4.5} = -1.6 \pm 0.3$ and $b_{K_{SBN} \sim 4.6} = 0.68 \pm 0.07$).

Nevertheless, Sellappan et al. attempted to develop a universal behavior for V_R with ν [18]. Fig. 18 presents the data collected herein, their data, and their sigmoid fit (dash black line):

$$V_R = \frac{1}{\alpha_R + \beta_R \exp(\chi_R \nu)} + \delta_R \quad (13)$$

where $(\alpha_R, \beta_R, \chi_R, \delta_R)$ are fitting parameters. The data for the SBN glasses does fit the spread of data presented in Sellappan et al. [18]. Yet, it has a tendency to be shifted left (i.e. smaller values of ν) as compared to the glasses of Sellappan et al. A better fit (solid brown line in Fig. 18) for the SBN glasses would occur with $(\alpha_R, \beta_R, \chi_R, \delta_R) = (.99, .003, 29, -.007)$.

Understanding shear flow requires V_P . Sellappan et al. [18] evidenced a sigmoid increasing trend between ν and V_P . Thus the amount of volume-conserving shear flow increases with ν . Fig. 19 (left) unveils two different scenarios when looking at approximately constant K_{SBN} . For samples with $K_{SBN} \sim 2.5$, V_P increases with ν ($V_P = m_{K_{SBN} \sim 2.5} \nu + b_{K_{SBN} \sim 2.5}$ where $m_{K_{SBN} \sim 2.5} = 0.30 \pm 0.022$ and $b_{K_{SBN} \sim 2.5} = -0.42 \pm 0.05$). On the other hand, the $K_{SBN} \sim 4.5$ series exhibits a drastic decrease in V_P with ν . Sellappan et al. attempted to universalize the behavior of V_P as they did with V_R [18]. In doing this, they concentrated on the pile-up. Fig. 19 presents the data collected herein, their data (black point), and their sigmoid fit (dash black line):

$$V_P = \frac{1}{\alpha_P + \beta_P \exp(\chi_P \nu)} \quad (14)$$

where $(\alpha_P, \beta_P, \chi_P)$ are fitting parameters. The data for the SBN glasses does fit the spread of data presented in Sellappan et al. [18]. A better fit for the SBN glasses would occur with $(\alpha_P, \beta_P, \chi_P) = (.94, 152, -16)$.

SBN glasses do fit the general spread of data when considering previously published V_P verse ν data. On the other hand, a sigmoid fit misses the subtleties of the SBN glasses for constant K_{SBN} . Moreover, it is difficult to understand why the trends of V_P are so drastically different: $K_{SBN} \sim 2.5$ has a tendency to increase and $K_{SBN} \sim 4.5$ has a tendency to decrease. Understanding this requires more research.

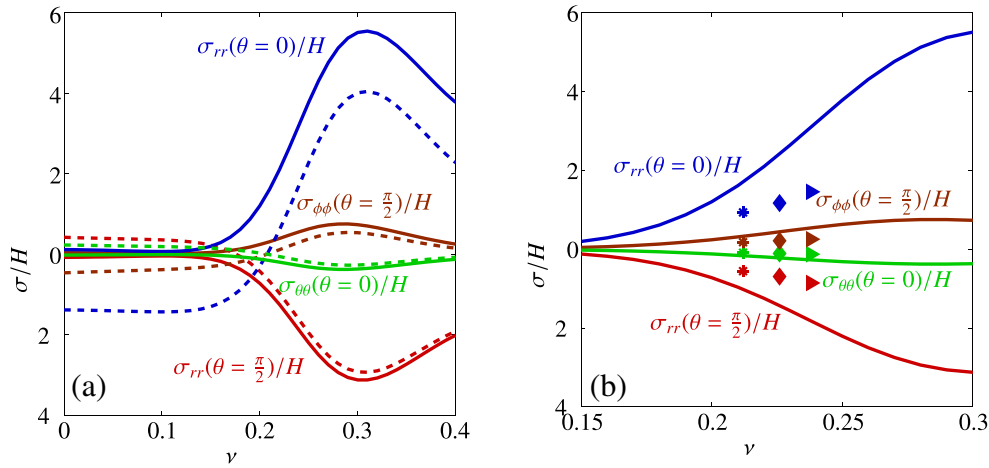


Fig. 21. (a) Estimates for the stress field calculated from equations in Table 3 for $E/H = 11.6$. Hashed line corresponds to stress during loading. Solid lines correspond to stresses during unloading. (b) A zoom of x axis of the figure on the left side with only the fully unloaded curves. The figure highlights data points herein that have E/H ratios between 10.3 and 12.8 (symbols remain the same as in Table 1, but the colors correspond to stresses). Three data points herein fit this requirement (SBN 14, SBN 25 and SBN 63). Colors correspond as follows: blue lines $\sigma_{rr}(r = a, \theta = 0)/H$; red lines $\sigma_{rr}(r = a, \theta = \pi/2)/H$; green lines $\sigma_{\theta\theta}(r = a, \theta = 0)/H$; and brown lines $\sigma_{\phi\phi}(r = a, \theta = \pi/2)/H$.

4.5. Residual stresses

Upon indentation, the indenter subjects a small region beneath itself to elastic/plastic processes [12]. Yet, sufficiently far from the indenter, the material obeys linear elastic dynamics. Assuming a conical indenter for symmetry purposes and to reduce the problems arising due to edges, Yoffe estimates the stresses induced during loading and unloading via Boussinesq and blister fields in spherical polar coordinates (r, θ, ϕ) , see Fig. 20 [56]. The following equations (with Table 3 for parameters) present Yoffe's stress fields:

$$\sigma_{rr} = \frac{P}{2\pi r^2} [1 - 2\nu - 2(2 - \nu) \cos\theta] + \frac{4B}{r^3} [(5 - \nu) \cos^2\theta - 2 + \nu] \quad (15)$$

$$\sigma_{\theta\theta} = \frac{P(1 - 2\nu) \cos^2\theta}{2\pi r^2(1 + \cos\theta)} - \frac{2B}{r^3} (1 - 2\nu) \cos^2\theta \quad (16)$$

$$\sigma_{\phi\phi} = \frac{P(1 - 2\nu)}{2\pi r^2} \left[\cos\theta - \frac{1}{1 + \cos\theta} \right] - \frac{2B}{r^3} (1 - 2\nu) (2 - 3 \cos^2\theta) \quad (17)$$

$$\sigma_{r\theta} = \frac{P(1 - 2\nu) \sin\theta \cos\theta}{2\pi r^2(1 + \cos\theta)} + \frac{4B}{r^3} (1 + \nu) \sin\theta \cos\theta \quad (18)$$

$$\sigma_{r\phi} = \sigma_{\theta\phi} = 0 \quad (19)$$

B is the strength of the Blister field. P is the applied load of the indenter. Sellappan et al. proposed a set of parameters for Yoffe's stresses [18]. Table 4 presents these parameters. Several input parameters in Sellappan et al. model include: (1) β governs the

Blister fields strength; (2) V_i^m is an estimate of the indent volume disregarding elastic recover; (3) a is an estimate of the half diagonal of the projected surface area of an indent; and (4) ψ is the angle of the indenter.

Fig. 21 (a) reveals theoretical stress calculations as presented by Sellappan et al. [18] for loading (hashed line) and unloading (solid line) stresses when $E/H = 11.6$. Sellappan et al. proposed this fit for E/H_V values ranging between 10.3 and 12.8. Three data points herein fit this requirement (SBN 14, SBN 25 and SBN 63). Fig. 21 (b) depicts the loading stress for these three data points using experimental results for V_i^- , a , V_R , V_P , ν , E and H_V (Tables 5, 1 and 2). This permits a comparison with Sellappan's model. Stresses for $\sigma_{\theta\theta}(r = a, \theta = 0)/H$ and $\sigma_{\phi\phi}(r = a, \theta = \pi/2)/H$ correspond well. However, discrepancies arise between theoretical values and experimental data for $\sigma_{rr}(r = a, \theta = \pi/2)/H$ and $\sigma_{rr}(r = a, \theta = 0)/H$. The signs of the stresses are consistent, but the magnitude of calculated experimental values is less than theoretical values.

Several factors could cause these discrepancies. Table 5 highlights differences in input parameters of stress equations, which are calculated from experiments and Sellappan's model. One difference of particular importance is the leftwards shift in the V_R versus ν as seen in Fig. 14. This implies that the volume recovery in SBN glasses is less than complex glasses used by Sellappan et al. On the other hand, V_P versus ν fits the spread of the data, yet there are still large differences between the model and experimental values.

Currently, the model assumes that densification and shear flow processes are independent of loading. Thus, this could be one source of uncertainty in the model. Whatever the values of V_R and V_P , the trends are coherent with [18].

Table 5

Comparison of measured and values obtained in [18] (Table 3 and Eqs. (13) and (14)). Negative difference indicates that the model overestimates measured values.

Name	Experimental					Model from [18]				Percent difference			
	d_i (μm)	V_i^- (μm^3)	V_R	V_P	β	V_i^- (μm^3)	V_R	V_P	β	V_i^-	V_R	V_P	β
SBN 12	10.5	36.7	0.528	0.163	0.31	34.54	0.704	0.157	0.14	6%	-33%	4%	55%
SBN 25	12.025	54.8	0.308	0.256	0.44	51.88	0.508	0.239	0.25	5%	-65%	7%	42%
SBN 30	13	62	0.126	0.286	0.59	65.55	0.388	0.298	0.31	-6%	-208%	-4%	47%
SBN 35	13.6	70.1	0.090	0.324	0.59	75.05	0.329	0.333	0.34	-7%	-265%	-3%	42%
SBN 14	11.25	35.2	0.343	0.113	0.54	42.48	0.686	0.164	0.15	-21%	-100%	-45%	72%
SBN 63	11.95	49.3	0.320	0.256	0.42	50.91	0.593	0.202	0.20	-3%	-85%	21%	52%
SBN 59	13.1	55.3	0.300	0.163	0.54	67.07	0.565	0.214	0.22	-21%	-88%	-31%	59%
SBN 55	13.625	56.8	0.280	0.158	0.56	75.46	0.415	0.283	0.30	-33%	-48%	-79%	46%

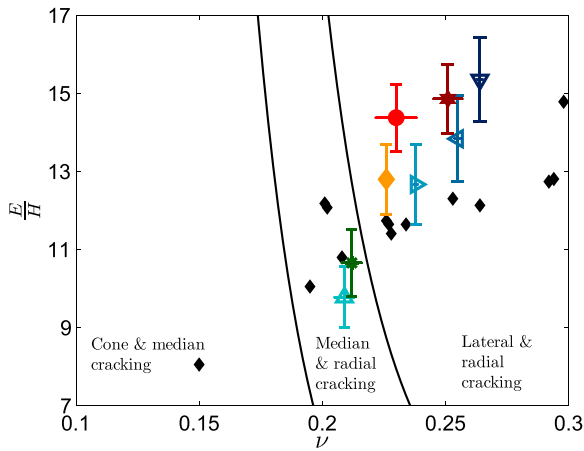


Fig. 22. Dependence of ν on the transitions from cone/median to median cracks and from median to lateral radial cracking. Black points represent values from Sellappan's paper [18].

Last but not the least, it is not unexpected that calculated stresses overestimate experimental values. The current expression for the Blister field (presented in Table 3) does not account for elastic recovery. Hence, the indentation volume (as estimated from the geometrical volume with contact at maximum load) is larger than the recovered volume. This means that both V_i^- and β are overestimated.

Variations in the chemical composition change the morphology of indentation print, specially the nature and the probability of crack appearance, as presented in Fig. 4. For example the nature and the intensity of the residual stresses impact directly median cracks. This is because they originate from the subsurface and are directly impacted by the intrinsic deformation [57]. Table 5 reveals variations in densification and shear flow processes. These variations lead to variations in the stress field (Fig. 21) which alters the crack appearance probability for SBN glasses (Fig. 8). This can explain the drastic variations in the crack appearance behavior observed from the SBN 12 to SBN 25. Thus, small variations in the chemical composition lead to a significant change in the cracking behavior [54].

Fig. 22 presents the data along with the types of cracking predicted by Sellappan et al. [18]. SBN 12 and SBN 14 require higher forces to initiate cracks off the corners of the indenter with SBN 12 never reaching 100% P_C . Thus a transition in cracking behavior from SBN 12 to the other glasses (which obtain a $P_C = 100\%$ at low forces) is feasible.

The SBN 12 presents a singular behavior regarding the SBN series herein. No cracks are visible after Vickers's indentations at 50 g. Furthermore, increasing the load did not facilitate the development of cracks popping up from the indentation corners. Complementary measurements of V_R and V_P highlight the important contribution of the densification process during the indentation. This changes the mechanism yielding to the crack formation: (1) change in the residual stresses resulting from these processes, and (2) appearance of other types of cracks under indentation which would limit the propagation of visible cracks.

5. Conclusion

A glass's chemical composition impacts its behavior under a sharp indenter. Glasses with a higher concentration of network formers exhibit higher hardness values. The addition of sodium in the composition (for glasses studied herein) induces NBOs in the silicate/borate network and changes the coordination of boron. As a consequence, the indenter penetrates deeper at a given load.

A good macroscopic parameter to link the glass structural variations to its continuum scale properties is ν . Densification processes favor smaller values of ν , i.e. strongly connected network, whereas shear flow process becomes predominant as the alkali content increases.

The two processes impact the indentation pattern and the cracking behavior. Sellappan et al. [18] model provides an estimation of stresses during loading and unloading. Because shear flow process increases with sodium concentration, stresses during loading and unloading increasingly favor lateral–radial cracks. As a consequence, K_C^{VIF} is indirectly proportional to ν .

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Appendix A. Table of symbols and their meanings

Table 6

Table of symbols.

Symbol	Meaning	Equation
SBN	Sodium BoroSilicate ($[\text{SiO}_2]-[\text{B}_2\text{O}_3]-[\text{Na}_2\text{O}]$)	
NBO	Non-bridging oxygen atoms	
H_V	Material hardness	5
K_C	Fracture toughness as measured in vacuum	
K_C^{IND}	K_C arises from indentation studies.	6
K_C^{VIF}	Vickers's indentation fracture toughness	7
ν	Poisson's ratio	3
[·]	mol%	
R_{SBN}	$\frac{[\text{Na}_2\text{O}]}{[\text{B}_2\text{O}_3]}$, ratio of mol% of $[\text{Na}_2\text{O}]$ to mol% of $[\text{B}_2\text{O}_3]$	
R_{max}^{SBN}	For $R_{SBN} < R_{max}^{SBN}$, Na^+ ions predominantly cause fully coordinated $^{[3]}\text{B}$ to transform into fully coordinated $^{[4]}\text{B}$. For $R_{SBN} > R_{max}^{SBN}$, extra Na^+ ions cause fully coordinated $^{[4]}\text{B}$ to transform into $^{[3]}\text{B}$ with NBO and it causes NBO in the silica network.	$R_{max}^{SBN} = 0.5 + \frac{K_{SBN}}{16}$
K_{SBN}	$\frac{[\text{SiO}_2]}{[\text{B}_2\text{O}_3]}$, ratio of mol % of $[\text{SiO}_2]$ to mol % of $[\text{B}_2\text{O}_3]$	
T_g	Glass transition temperature	
E	Young's modulus	2
$\langle \text{CN} \rangle$	Mean coordination number of the boron atoms	
$^{[4]}\text{B}$	The concentration of $^{[4]}\text{B}$ per volume unit deduced from NMR measurements and ICP-AES results	
N_{NBO}	Number of NBO per volume unit	4
N_{Na}	Number of Na^+ ions acting as network modifiers	
$N_{^{[4]}\text{B}}$	Each $^{[4]}\text{B}$ structural unit has one Na^+ ion attaching to it. Thus, this equates to the number of Na^+ ion acting as a network compensator.	
ρ	Density of sample	1
ρ_w	Density of water	
m_d	Mass of samples in air at 24 °C	
m_w	Mass of samples in water at 24 °C	
V_L	longitudinal Velocities	
V_T	transverse Velocities	
P_C	Crack appearance probability	
P	Indentation load	
P_{id}	Residual indentation depth	
P_{ad}	Residual indentation depth after annealing	
d_i	An indentation diagonal length on one indent (note there are 2 per indent)	
d_a	An indentation diagonal length on one indent (note there are 2 per indent) after annealing	
d	The mean of $\langle d_i \rangle$ values obtained for one indent.	
V_i^-	Indentation volume below the free surface	
V_i^+	Indentation volume above the free surface, i.e. pile-up	
V_a^-	Indentation volume below the free surface after annealing	
V_a^+	Indentation volume above the free surface after annealing, i.e. pile-up	
c_i	The crack length measured from the indentation center.	
c	The average distance ($\langle c_i \rangle$) from the indentation center to the crack tip	

(continued on next page)

Table 6 (continued)

Symbol	Meaning	Equation
C_R	Crack resistance: the load where P_C exceeds 50%	
V_R	Recovered volume ratio	8
V_P	Volume-conserving shear flow ratio as determined in the pile-up	9
P_f	Packing fraction	10
ρ_t	Minimal theoretical density	11
V_i	Ideal volume for an oxide	12
N_A	Avogadro number	
r_A	Ionic radii of the cation in the oxide $[A_xO_y]$	
r_O	The ionic radii of anion in the oxides $[A_xO_y]$	
$(\alpha_R, \beta_R, \chi_R, \delta_R)$	Fitting parameters for the sigmoid fit of V_R for Eq. (13)	
$(\alpha_P, \beta_P, \chi_P)$	Fitting parameters for the sigmoid fit of V_P for Eq. (14)	
$\sigma_{rr}, \sigma_{\theta\theta}, \sigma_{\phi\phi}, \sigma_{r\theta}, \sigma_{r\phi}, \sigma_{\theta\phi}$	Yoffe's stress field in spherical polar coordinates	15, 16, 17, 18, 19
B	Strength of blister field	Table 3
β	Sellappan's blister field strength	Table 3
V_i^m	Volume of indent disregarding elastic recovery	Table 3
a	Half the diagonal of the projected surface area	Table 3
ψ	Apical angle of the indenter	Table 3

Appendix B. Data used to calculate E and ν

Table 7

V_L , V_T , E , and ν calculated for each SBN glass. Average values result from 3 measurements of V_L and V_T during different times of the year. Standard deviations result from these measurements. Error bars in the figures depict one standard deviation.

Name	V_L (m/s)	V_T (m/s)	E (GPa)	ν
SBN 12	3074 \pm 17	1835 \pm 6	80.1 \pm 0.03	0.208 \pm 0.06
SBN 25	3053 \pm 10	1789 \pm 8	80.3 \pm 0.04	0.238 \pm 0.008
SBN 30	2938 \pm 6	1713 \pm 8	74.7 \pm 0.05	0.255 \pm 0.001
SBN 35	3057 \pm 15	1727 \pm 6	76.7 \pm 0.04	0.264 \pm 0.002
SBN 14	3049 \pm 21	1851 \pm 32	81.8 \pm 1.02	0.210 \pm 0.015
SBN 63	3046 \pm 11	1809 \pm 8	81.9 \pm 0.06	0.226 \pm 0.009
SBN 59	2968 \pm 24	1767 \pm 15	77.2 \pm 0.03	0.230 \pm 0.007
SBN 55	2931 \pm 20	1689 \pm 4	72.8 \pm 0.06	0.251 \pm 0.005

References

- [1] T. Rouxel, T. Hammouda, A. Moréac, Poisson's ratio and the densification of glass under high pressure, *Phys. Rev. Lett.* 100 (2008) 225501.
- [2] D. Uhlmann, Densification of alkali silicate glasses at high pressure, *J. Non-Cryst. Solids* 13 (1973) 89–99.
- [3] F. Ernsberg, Role of densification in deformation of glasses under point loading, *J. Am. Ceram. Soc.* 51 (1968) 545.
- [4] J. Mackenzie, High-pressure effects on oxide glasses. 1. densification in rigid state, *J. Am. Ceram. Soc.* 46 (1963) 461–470.
- [5] K. Peter, Densification and flow phenomena of glass in indentation experiments, *J. Non-Cryst. Solids* 5 (1970) 103–115.
- [6] J. Hagan, S. Van der Zwaag, Plastic processes in a range of soda–lime–silica glasses, *J. Non-Cryst. Solids* 64 (1984) 249–268.
- [7] J. Hagan, Shear deformation under pyramidal indentations in soda–lime glass, *J. Mater. Sci.* 15 (1980) 1417–1424.
- [8] J.T. Hagan, M.V. Swain, The origin of median and lateral cracks around plastic indents in brittle materials, *J. Phys. D: Appl. Phys.* 11 (15) (1978) 2091. <http://dx.doi.org/10.1088/0022-3727/11/15/007>.
- [9] M. Bertoldi, M. Sglavo, Soda–borosilicate glass: normal or anomalous behavior under Vickers indentation, *J. Non-Cryst. Solids* 344 (2004) 51–59.
- [10] G. Buscarino, S. Agnello, F.M. Gelardi, Structural modifications induced by electron irradiation in SiO₂ glass: local densification measurements, *Europhys. Lett.* 87 (2009) 26007.
- [11] T. Rouxel, J.-C. Sangleboeuf, C. Moysan, B. Truffin, Indentation topometry in glasses by atomic force microscopy, *J. Non-Cryst. Solids* 344 (2004) 26–36.
- [12] B. Lawn, A.G. Evan, D.B. Marshall, Elastic–plastic indentation damage in ceramics – the median–radial crack system, *J. Am. Ceram. Soc.* 63 (9–10) (1980) 574–581. <http://dx.doi.org/10.1111/j.1151-2916.1980.tb10768.x>.
- [13] A. Arora, D. Marshall, B. Lawn, Indentation deformation/fracture of normal and anomalous glasses, *J. Non-Cryst. Solids* 31 (1979) 415–428.
- [14] J. Hagan, Micromechanics of crack nucleation during indentations, *J. Mater. Sci.* 14 (1979) 2975–2980.
- [15] J.D. Mackenzie, High-pressure effects on oxide glasses, *J. Am. Ceram. Soc.* 46 (10) (1963) 461–470.
- [16] S. Yoshida, J.-C. Sangleboeuf, T. Rouxel, Quantitative evaluation of indentation-induced densification in glass, *J. Mater. Res.* 20 (2005) 3404–3412.
- [17] S. Yoshida, J.-C. Sangleboeuf, T. Rouxel, Indentation-induced densification of soda–lime silicate glass, *Int. J. Mater. Res.* 98 (2007) 360–364.
- [18] P. Sellappan, T. Rouxel, F. Celarie, E. Becker, P. Houizot, R. Conradt, Composition dependence of indentation deformation and indentation cracking in glass, *Acta Mater.* 61 (2013) 5949–5965.
- [19] M. Barlet, A. Kerrache, J.-M. Delaye, C. Rountree, SiO₂–Na₂O–B₂O₃ density: a comparison of experiments, simulations, and theory, *J. Non-Cryst. Solids* 382 (0) (2013) 32–44. <http://dx.doi.org/10.1016/j.jnoncrysol.2013.09.022> (URL <http://www.sciencedirect.com/science/article/pii/S0022309313004869>).
- [20] F. Angelini, T. Charpentier, D. de Ligny, C. Cailleteau, Boron speciation in soda–lime borosilicate glasses containing zirconium, *J. Am. Ceram. Soc.* 93 (9) (2010) 2693–2704. <http://dx.doi.org/10.1111/j.1551-2916.2010.03771.x>.
- [21] C. Hermansen, J. Matsuoka, S. Yoshida, H. Yamazaki, Y. Kato, Y. Yue, Densification and plastic deformation under microindentation in silicate glasses and the relation to hardness and crack resistance, *J. Non-Cryst. Solids* 364 (0) (2013) 40–43. <http://dx.doi.org/10.1016/j.jnoncrysol.2012.12.047> (URL <http://www.sciencedirect.com/science/article/pii/S0022309313000094>).
- [22] P. Chantikul, G. Anstis, B. Lawn, D. Marshall, A critical-evaluation of indentation techniques for measuring fracture-toughness. 2. Strength method, *J. Am. Ceram. Soc.* 64 (9) (1981) 539–543. <http://dx.doi.org/10.1111/j.1151-2916.1981.tb10321.x>.
- [23] D.B. Marshall, A.G. Evans, Reply to comment on elastic/plastic indentation damage in ceramics: the median/radial crack system, *J. Am. Ceram. Soc.* 64 (12) (1981) C-182–C-183. <http://dx.doi.org/10.1111/j.1151-2916.1981.tb15909.x>.
- [24] D. Marshall, B. Lawn, Residual stress effects in sharp contact cracking part 1 indentation fracture mechanics, *J. Mater. Sci.* 14 (8) (1979) 2001–2012. <http://dx.doi.org/10.1007/BF00551043>.
- [25] K. Niihara, A fracture-mechanics analysis of indentation-induced cs palmqvist crack in ceramics, *J. Mater. Sci. Lett.* 2 (5) (1983) 221–223. <http://dx.doi.org/10.1007/BF00725625>.
- [26] K. Niihara, R. Morena, D. Hasselman, Evaluation of klc of brittle solids by the indentation method with low crack-to-indent ratios, *J. Mater. Sci. Lett.* 1 (1) (1982) 13–16. <http://dx.doi.org/10.1007/BF00724706>.
- [27] K. Niihara, A. Nakahira, T. Hirai, The effect of stoichiometry on mechanical-properties of boron–carbide, *J. Am. Ceram. Soc.* 67 (1) (1984) C13–C14.
- [28] G. Anstis, P. Chantikul, B. Lawn, D. Marshall, A critical evaluation of indentation technique for measuring fracture toughness: I, direct crack measurement, *J. Am. Ceram. Soc.* 64 (1981) 533–538. <http://dx.doi.org/10.1111/j.1151-2916.1981.tb10320.x>.
- [29] J. Lankford, Indentation microfracture in the Palmqvist crack regime – implications for fracture-toughness evaluation by the indentation method, *J. Mater. Sci. Lett.* 1 (11) (1982) 493–495. <http://dx.doi.org/10.1007/BF00721938>.
- [30] M. Laugier, The elastic plastic indentation of ceramics, *J. Mater. Sci. Lett.* 4 (12) (1985) 1539–1541. <http://dx.doi.org/10.1007/BF00721390>.
- [31] M. Laugier, New formula for indentation toughness in ceramics, *J. Mater. Sci. Lett.* 6 (3) (1987) 355–356. <http://dx.doi.org/10.1007/BF01729352>.
- [32] W.J. Weber, H. Matzke, J.L. Roubert, Indentation testing of nuclear-waste glasses, *J. Mater. Sci.* 19 (8) (1984) 2533–2545. <http://dx.doi.org/10.1007/BF00550807>.
- [33] R.D. Dukino, M.V. Swain, Comparative measurement of indentation fracture toughness with Berkovich and Vickers indenters, *J. Am. Ceram. Soc.* 75 (12) (1992) 3299–3304. <http://dx.doi.org/10.1111/j.1151-2916.1992.tb04425.x>.
- [34] J.R. Taylor, *An Introduction to Error Analysis*, University Science Books, 1982.
- [35] A. Perriot, E. Barthel, G. Kermouche, G. Querel, D. Vandembroucq, On the plastic deformation of soda–lime glass – a Cr³⁺ luminescence study of densification, *Philos. Mag.* 91 (2011) 1245–1255.
- [36] A. Perriot, D. Vandembroucq, E. Barthel, V. Martinez, L. Grosvalet, C. Martinet, B. Champagnon, Raman microspectroscopic characterization of amorphous silica plastic behavior, *J. Am. Ceram. Soc.* 89 (2006) 596–601.
- [37] J. Mackenzie, High-pressure effects on oxide glasses. 2. Subsequent heat treatment, *J. Am. Ceram. Soc.* 46 (1963) 470.
- [38] P.W. Bridgman, I. Simon, Effect of very high pressure on glass, *J. Appl. Phys.* 24 (1953) 405.
- [39] F.M. Ernsberger, Role of densification in deformation of glasses under point loading, *J. Am. Ceram. Soc.* 51 (1968) 545–547.
- [40] S. Yoshida, S. Isono, J. Matsuoka, N. Soga, Shrinkage behavior of Knoop indentations in silica and soda–lime–silica glasses, *J. Am. Ceram. Soc.* 84 (9) (2001) 2141–2143.
- [41] G. Hetherington, The viscosity of vitreous silica, *Phys. Chem. Glasses* 5 (1964) 130.
- [42] H. Inoue, A. Masuno, Y. Watanabe, K. Suzuki, T. Iseda, Direct calculation of the physical properties of sodium borosilicate glass from its chemical composition using the concept of structural units, *J. Am. Ceram. Soc.* 95 (1) (2012) 211–216. <http://dx.doi.org/10.1111/j.1551-2916.2011.04964.x>.
- [43] D. Feil, S. Feller, The density of sodium borosilicate glasses related to atomic arrangements, *J. Non-Cryst. Solids* 119 (1) (1990) 103–111. [http://dx.doi.org/10.1016/0022-3093\(90\)90246-l](http://dx.doi.org/10.1016/0022-3093(90)90246-l).
- [44] K. BUDHWANI, S. Feller, Density model for the lithium, sodium and potassium borosilicate glass systems, *Phys. Chem. Glasses* 36 (4) (1995) 183–190.
- [45] W.J. Dell, P.J. Bray, S.Z. Xiao, B-11 NMR-studies and structural modeling of Na₂O–B₂O₃–SiO₂ glasses of high soda content, *J. Non-Cryst. Solids* 58 (1) (1983) 1–16. [http://dx.doi.org/10.1016/0022-3093\(83\)90097-2](http://dx.doi.org/10.1016/0022-3093(83)90097-2).
- [46] B.C. Bunker, D.R. Tallant, R.J. Kirkpatrick, G. Turner, Multinuclear nuclear–magnetic-resonance and Raman investigation of sodium borosilicate glass structures, *Phys. Chem. Glasses* 31 (1) (1990) 30–41.

- [47] H. Ji, E. Robin, T. Rouxel, Compressive creep and indentation behavior of plasticine between 103 and 353 K, *Mech. Mater.* 41 (2009) 199–209.
- [48] R.L. Henderson, N.W. Ashcroft, Perturbation theory of structure in mixtures near phase separation, *Phys. Rev. A* 13 (1976) 859–871.
- [49] T. Rouxel, Elastic properties and short- to medium-range order in glasses, *J. Am. Ceram. Soc.* 10 (2007) 3019–3039. <http://dx.doi.org/10.1111/j.1551-2916.2007.01945.x>.
- [50] A. Makishima, J.D. Mackenzie, Direct calculation of Young's modulus of glass, *J. Non-Cryst. Solids* 12 (1) (1973) 35–45. [http://dx.doi.org/10.1016/0022-3093\(73\)90053-7](http://dx.doi.org/10.1016/0022-3093(73)90053-7).
- [51] D. Kilymis, J.-M. Delaye, Deformation mechanisms during nanoindentation of sodium borosilicate glasses of nuclear interest, *J. Chem. Phys.* 141 (1) (2014) 014504.
- [52] D. Kilymis, J.-M. Delaye, Nanoindentation studies of simplified nuclear glasses using molecular dynamics, *J. Non-Cryst. Solids* 401 (2014) 147–153.
- [53] T. Rouxel, H. Ji, V. Keryvin, T. Hammouda, S. Yoshida, Poisson's ratio and the glass network topology — relevance to high pressure densification and indentation behavior, *Adv. Mater. Res.* 39–40 (2008) 137–146.
- [54] P. Sellappan, Fissuration par indentation des verres: influence de la composition et renforcement par des particules de diamants (Ph.D. thesis) Université de Rennes 1, 2011.
- [55] Y. Kato, H. Yamazaki, S. Yoshida, J. Matsuoka, Effect of densification on crack initiation under Vickers indentation test, *J. Non-Cryst. Solids* 356 (35–36) (2010) 1768–1773. <http://dx.doi.org/10.1016/j.jnoncrysol.2010.07.015>.
- [56] E. Yoffe, Elastic stress fields caused by indentation brittle materials, *Philos. Mag. A* 46 (1982) 617–627.
- [57] H. Ishikawa, N. Shinkai, Critical load for median crack initiation in Vickers indentation of glasses, *J. Am. Ceram. Soc.* 65 (8) (1982) C124–C127. <http://dx.doi.org/10.1111/j.1151-2916.1982.tb10496.x>.