Relativistic calculations of the isotope shifts in highly charged Li-like ions

N. A. Zubova^{1,2}, Y. S. Kozhedub^{1,2}, V. M. Shabaev¹, I. I. Tupitsyn¹, A.V. Volotka^{1,3}, G. Plunien³, C.Brandau^{4,5,6}, and Th. Stöhlker^{4,7,8}

¹St. Petersburg State University, Russia; ²SSC RF ITEP of NRC "Kurchatov Institute", Russia; ³TU Dresden,

Germany; ⁴GSI Darmstadt, Germany; ⁵ExtreMe Matter Institute EMMI and Research Division GSI Darmstadt,

Germany; ⁶Justus-Liebig-University Giessen, Germany; ⁷Helmholtz-Institut Jena, Germany;

⁸Friedrich-Schiller-Universität Jena, Germany

In the past years a great progress was achieved in experimental study of the isotope shifts in highly charged ions. This was done in the dielectronic recombination (DR) experiments at ESR (GSI) [1] as well as applying a highresolution grating spectrometer at EBIT [2]. The joint analysis of the results obtained in the high-precision theoretical and experimental investigations of isotope shifts in highly charged ions provides determination of nuclear charge radii differences and enables tests of quantum electrodynamics (QED) at strong fields. With the FAIR facilities the isotope shift measurements in heavy ions will be further improved in accuracy and extended to radioactive isotopes with a lifetime longer than about 10 s.

From the theoretical side, to obtain the total value of the isotope shift one needs to evaluate the nuclear size (field shift) and nuclear recoil (mass shift) contributions, including the relativistic and QED effects. Since there exists some discrepancy in the results obtained in Refs. [3] and [4] for heavy ions, we develop an alternative method which merges the perturbative and CI-DFS (configurationinteraction Dirac-Fock-Sturm) calculations [5]. Namely, we calculate the nuclear recoil contributions within the Breit approximation to zeroth and first orders in 1/Z and add the related contributions of second and higher orders in 1/Z, obtained using the CI-DFS method. To derive the nuclear recoil contributions to the binding energies of Li-like ions by perturbation theory, we use the two-time Green's function method [6] with the $(1s)^2$ shell regarded as belonging to a redefined vacuum. The obtained non-QED results are combined with the corresponding QED contributions of the zeroth order in 1/Z to get the most accurate theoretical data for the mass shifts in highly charged Li-like ions. In addition, the field shifts are calculated in the framework of the Dirac-Coulomb-Breit Hamiltonian. These calculations, being performed by the CI-DFS method, are compared with the corresponding MCDF (multiconfiguration Dirac-Fock) calculations of Ref. [3]. The QED corrections to the field shifts are also evaluated. In addition, we consider the nuclear deformation and nuclear polarization corrections. As the result, the most precise theoretical values of the isotope shifts for the $2p_{1/2} - 2s$ and $2p_{3/2} - 2s$ transitions in Li-like ions are presented in our paper [5]. The theoretical contributions to the isotope shifts in Li-like thorium and uranium are given in Table 1.

Table 1: Individual contributions to the isotope shifts for the $2p_{1/2} - 2s$ and $2p_{3/2} - 2s$ transitions in Li-like 232,230 Th⁸⁸⁺ and 238,234 U⁸⁹⁺(in meV) with given values of $\delta \langle r^2 \rangle$.

	$^{232,230}\mathrm{Th}^{87+}$	
	$^{232,230}\delta\langle r^{2}\rangle {=} 0.2050{\rm fm}^{2}$	
	$2p_{1/2}-2s$	$2p_{3/2}-2s$
Main contributions		
Field shift	-116.1	-128.7
Mass shift	0.1	0.3
FS plus MS	-116.0	-128.4
QED		
Field shift	0.6	0.9
Mass shift	0.4	0.4
Others		
Nuclear polarization	1.6	1.7
Nuclear deformation	1.5	1.5
Total IS theory	-111.9(22)	-123.9(22)
	$^{238,234}\mathrm{U}^{89+}$	
	238,234	$^{4}\mathrm{U}^{89+}$
	$^{238,234}_{238,234}\delta\langle r^2 angle$	$^{4}\mathrm{U}^{89+}$
	238,234 $238,234 \delta \langle r^2 \rangle$ $2p_{1/2} - 2s$	$^{4}\mathrm{U}^{89+}$ =0.334 fm ² $2\mathrm{p}_{3/2}-2\mathrm{s}$
Main contributions	238,234 $^{238,234}\delta\langle r^{2} angle$ $^{2}\mathrm{p}_{1/2}-^{2}\mathrm{s}$	${}^{4}\mathrm{U}^{89+}$ =0.334 fm ² $2\mathrm{p}_{3/2} - 2\mathrm{s}$
Main contributions Field shift	238,234 $238,234 \delta \langle r^2 \rangle$ $2p_{1/2} - 2s$ -227.8	$^{4}U^{89+}$ =0.334 fm ² $2p_{3/2} - 2s$ -254.5
Main contributions Field shift Mass shift	$238,234$ $238,234 \delta \langle r^2 \rangle$ $2p_{1/2} - 2s$ -227.8 0.2	$^{4}U^{89+}$ =0.334 fm ² $2p_{3/2} - 2s$ -254.5 0.6
Main contributions Field shift Mass shift FS plus MS	$\begin{array}{c} 238,234\\ 238,234 \delta \langle r^2 \rangle\\ 2p_{1/2}-2s\\ -227.8\\ 0.2\\ -227.6\end{array}$	$^{4}U^{89+}$ =0.334 fm ² $2p_{3/2} - 2s$ -254.5 0.6 -253.9
Main contributions Field shift Mass shift FS plus MS QED	$238,234 \\ 238,234 \\ \delta \langle r^2 \rangle \\ 2p_{1/2} - 2s \\ -227.8 \\ 0.2 \\ -227.6 \\ \end{cases}$	$^{4}U^{89+}$ =0.334 fm ² $2p_{3/2} - 2s$ -254.5 0.6 -253.9
Main contributions Field shift Mass shift FS plus MS QED Field shift	$\begin{array}{c} 238,234\\ 238,234 \delta \langle r^2 \rangle\\ 2p_{1/2}-2s\\ -227.8\\ 0.2\\ -227.6\\ 1.2\end{array}$	${}^{4}U^{89+}$ =0.334 fm ² 2p _{3/2} - 2s -254.5 0.6 -253.9 1.8
Main contributions Field shift Mass shift FS plus MS QED Field shift Mass shift	$\begin{array}{c} 238,234\\ 238,234 \delta \langle r^2 \rangle\\ 2p_{1/2}-2s\\ -227.8\\ 0.2\\ -227.6\\ 1.2\\ 0.9\end{array}$	${}^{4}U^{89+}$ =0.334 fm ² 2p _{3/2} - 2s -254.5 0.6 -253.9 1.8 0.8
Main contributions Field shift Mass shift FS plus MS QED Field shift Mass shift Others	$\begin{array}{c} 238,234\\ 238,234 \delta \langle r^2 \rangle\\ 2p_{1/2}-2s\\ -227.8\\ 0.2\\ -227.6\\ 1.2\\ 0.9\end{array}$	${}^{4}U^{89+}$ =0.334 fm ² 2p _{3/2} - 2s -254.5 0.6 -253.9 1.8 0.8
Main contributions Field shift Mass shift FS plus MS QED Field shift Mass shift Others Nuclear polarization	$\begin{array}{c} 238,234\\ 238,234 \delta \langle r^2 \rangle\\ 2p_{1/2}-2s\\ -227.8\\ 0.2\\ -227.6\\ 1.2\\ 0.9\\ 2.3\end{array}$	${}^{1}U^{89+}$ =0.334 fm ² 2p _{3/2} - 2s -254.5 0.6 -253.9 1.8 0.8 2.6
Main contributions Field shift Mass shift FS plus MS QED Field shift Mass shift Others Nuclear polarization Nuclear deformation	$\begin{array}{c} 238,234\\ 238,234 \delta \langle r^2 \rangle\\ 2p_{1/2}-2s\\ -227.8\\ 0.2\\ -227.6\\ 1.2\\ 0.9\\ 2.3\\ -2.4\end{array}$	${}^{4}U^{89+}$ =0.334 fm ² 2p _{3/2} - 2s -254.5 0.6 -253.9 1.8 0.8 2.6 -2.7

References

- [1] C. Brandau et al., Phys. Rev. Lett. 100, 073201 (2008).
- [2] R. Soria Orts et al., Phys. Rev. Lett. 97, 103002 (2006).
- [3] J. Li et al., Phys. Rev. A 86, 022518 (2012).
- [4] Y. S. Kozhedub et al., Phys. Rev. A 81, 042513 (2010).
- [5] N. A. Zubova et al., Phys. Rev. A 90, 062512 (2014).
- [6] V. M. Shabaev, Phys. Rep. 356, 119 (2002).