

# Observation of Unbound States in $^{16}\text{Ne}$ / $^{15}\text{Ne}$ via 1n- / 2n-Knockout on $^{17}\text{Ne}$ \*

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In recent years, experiments investigating the driplines have unearthed rich evidence on the peculiarities of the nuclear force, in particular those connected to weak binding and large proton-neutron asymmetry. While for very neutron-rich systems, e.g., various manifestations for 1n and 2n halo systems are found, the existence of the dripline is known only up to  $Z = 8$ . Quite reversed due to the additional presence of the Coulomb barrier, the proton dripline is sharp and known up to  $Z = 91$ , whereas halo formation is suppressed and just a few cases exist. In the recent past we have investigated proton-knockout reactions studying the borromean 2p-halo nucleus  $^{17}\text{Ne}$  [1, 2], and here we have used the dataset from the same experiment as a stepping stone to reach beyond the proton dripline and explore yet unknown regions of the nuclear landscape.

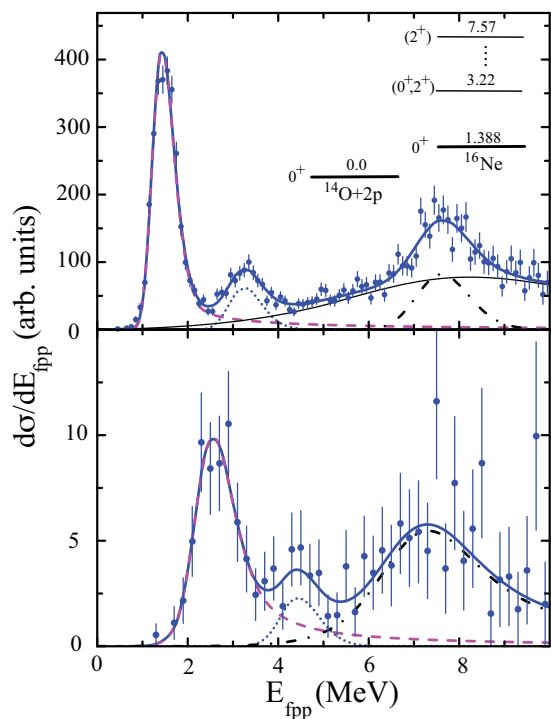


Fig. 1: Relative-energy ( $E_{fpp}$ ) of  $^{16}\text{Ne}$  (top) and  $^{15}\text{Ne}$  (bottom), after 1n-/2n-knockout from  $^{17}\text{Ne}$  projectiles.

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This report presents data of 1n- and 2n- knockout reactions on  $^{17}\text{Ne}$  projectiles in light targets (C,  $\text{CH}_2$ ), populating states in the unbound nuclei  $^{16}\text{Ne}$  and the yet unobserved  $^{15}\text{Ne}$ . In a simple picture, the respective neutrons were removed from the  $^{17}\text{Ne}$  core,  $^{15}\text{O}$ , thus creating  $^{14,13}\text{O}$  fragments coupled to the two remaining  $s^2/d^2$  protons, all travelling under forward angles and being detected in coincidence. The data analysis procedure, via 4-momentum reconstruction and invariant-mass technique, is equivalent to the description in [1]. The excitation spectra, in terms of f-2p relative-energy spectra, of the  $^{16}\text{Ne}$  and  $^{15}\text{Ne}$  systems are shown in Fig. 1. The shown data (full dots with errorbars) have been corrected for experimental acceptance, and the peaks have been fitted by Coulomb-Breit-Wigner functions (dashed, dotted lines) folded with the experimental resolution, and in the case of  $^{16}\text{Ne}$  in addition by a non-resonant background (full line). The experimental acceptance and the  $E_{rel}$  calibration and resolution have been obtained from R3BROOT-based simulations [2] in combination with the width and position of the known  $5/2^-$  state in the  $^{15}\text{O}+2p$  continuum of  $^{17}\text{Ne}$  as a reference.

For  $^{16}\text{Ne}$  ( $^{14}\text{O}+2p$ ) we have extracted the positions and widths of the ground state and the first two excited states as  $E_r(\text{g.s.}) = 1.388(15)$  MeV,  $\Gamma_r(\text{g.s.}) = 0.082(15)$  MeV;  $E_r(1.x.) = 3.22(5)$  MeV,  $\Gamma_r(1.x.) \leq 0.05$  MeV;  $E_r(2.x.) = 7.57(6)$  MeV,  $\Gamma_r(2.x.) \leq 0.1$  MeV. These values are in good agreement with previous publications on  $^{16}\text{Ne}$ , e.g. [4], and confirm the validity of our technique and calibrations.

For the first time, the unbound isotope  $^{15}\text{Ne}$  has been observed, as coincidences between  $^{13}\text{O}$  and two beam-like protons (see Fig. 1(bottom)). The same type of analysis as for the  $^{16}\text{Ne}$  case has been performed for the  $^{15}\text{Ne}$  ( $E_{fpp}$ ) spectrum. We identified the ground state and the first excited state with parameters of  $E_r(\text{g.s.}) = 2.522(66)$  MeV,  $\Gamma_r(\text{g.s.}) = 0.59(23)$  MeV;  $E_r(1.x.) = 4.42(4)$  MeV,  $\Gamma_r(1.x.) \leq 0.1$  MeV. The observed position of the  $^{15}\text{Ne}$  ground state at  $S_{2p} = -2.522(66)$  MeV is in good agreement to a recent model prediction [5].

## References

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