

# Automatic Self-consistent Gain-Matching of DSSSD Detector Channels\*

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**Double-sided silicon strip detectors (DSSSD) are a widely used type of detector in nuclear and particle physics experiments for position and energy measurements. This report describes an automatic method that allows to gain-match all strips of DSSSD detectors without the need of dedicated in-beam calibration.**

## Introduction

DSSSD detectors are constructed as large area silicon diodes with electrically segmented p and n-side contacts. Signals are read out on both, p and n-side simultaneously. The segmentation is usually such that there are unique intersection points of opposite side's contacts, the measurement of one p and on n-side signal allows to reconstruct the two dimensional spatial position of an event inside the detector.

## Method

In the following, it will be assumed that the deposited energy  $E = s \cdot A$ , in the detector is proportional to the amplitude, i.e. there are no offsets or nonlinearities in the electronics.  $A$  is the measured amplitude of the detected signal and  $s$  the slope factor for calibration.

Given a DSSSD detector with  $N$  strips on each side, each event that is registered in a given pixel will create a signal with amplitude  $A_p$  in the strip number  $p$  on the p-side and a signal with amplitude  $A_n$  in strip number  $n$  on the n-side ( $n, p = 1 \dots N$ ). Assuming that both strips measure the same energy  $E$  deposited in the active area of the detector, one can write

$$E_p = s_p A_p, \quad E_n = s_n A_n \quad \text{and} \quad E_p = E_n = E, \quad (1)$$

with  $s_p$  and  $s_n$  being the calibration coefficients (slopes) for the  $p$ -th p-side strip and  $n$ -th n-side strip, respectively.

For each pixel of the detector, the values of  $A_p$  and  $A_n$  for various values of energy depositions  $\Delta E$  are assumed to be linearly related. For each pixel, the relation between the amplitudes is

$$A_p = S_{pn} A_n, \quad (2)$$

with the slope parameter  $S_{pn}$  which can be determined from the data (see below). The set of  $N^2$  slope parameters

can be used to deduce a set of  $2N$  calibration coefficients  $\{s_p, s_n\}$  by minimizing the expression

$$\chi^2 = \sum_{p,n} \left( \frac{S_{pn} - \frac{s_n}{s_p}}{\Delta S_{pn}} \right)^2, \quad (3)$$

that follows directly from Eqs. (1,2).

## Implementation

There are two independent steps: First, the determination of the coefficients  $S_{pn}$ . Second, finding a set of calibration coefficients  $\{s_p, s_n\}$  based on  $S_{pn}$ . The former is done using a Bayesian approach, updating the posterior probability distribution of the quantity of interest for each detected event. The latter is done by a nonlinear least squares fit algorithm.

*Determination of  $S_{pn}$ :* The posterior probability distribution of the quantity of interest is in this case  $p(S_{pn} | \{d_{pn}\})$ , with  $\{d_{pn}\}$  being the set of all measured data points, i.e. the ratio of amplitudes  $d_{pn} = A_p/A_n$ . Knowing this distribution allows to get the most likely value for each slope parameter  $S_{pn}$  and its variance  $\Delta S_{pn}$ . Starting with an initial guess for this distribution, i.e. uniform within reasonable limits, one can refine it by iterating over the measured data points, each time applying Bayes' theorem [4, 5]

$$p(S_{pn} | d_{pn}) = \frac{p(S_{pn}) p(d_{pn} | S_{pn})}{p(d_{pn})}, \quad (4)$$

with the commonly used terminology [6]:  $p(S_{pn} | d_{pn})$  is called posterior distribution,  $p(S_{pn})$  is the prior distribution,  $L(d_{pn} | S_{pn})$  the likelihood function and  $p(d_{pn})$  the evidence of the measured data. After each treated data point the normalized posterior distribution becomes the prior for the next data point.

The likelihood function is chosen to be a Lorentzian-shaped distribution with width  $w$

$$L(d_{pn} | S_{pn}) \propto \frac{1}{w^2 + (\log d_{pn} - \log S_{pn})^2}. \quad (5)$$

Such a distribution makes the result less sensitive to abundant background events as it would be the case with a Gaussian-shaped likelihood function. It depends on the logarithm of the slope parameter, because the slope is a scale parameter with a possible range inside  $[0, \infty[$ .

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While iterating over the data, the posterior distribution is represented as a discrete number  $N_p$  of points between the limits  $S_{\min}$  and  $S_{\max}$ . The density of points depends on the desired accuracy of the final result and can be in the order of a few thousands. After all data is processed, the best estimates for all slopes  $S_{pn}$  and their uncertainties  $\Delta S_{pn}$  are determined from the mean and variance of the final posterior probability distribution. For the example data shown here, the set of values was:  $N_p = 1000$ ,  $S_{\min} = 0.3$ ,  $S_{\max} = 3$  and  $w = 0.1$ .

*Computing a set of calibration coefficients:* Minimizing (3) is done using the implementation of a nonlinear least squares fit provided by [2]. The set of fit parameters is  $\{s_p, s_n\}$ , and the data is the complete set of measured slope parameters  $\{S_{pn}\}$ . The algorithm performs the minimization of (3). After convergence is reached, the parameter set describes the best calibration coefficients for the data set on a common arbitrary scale.

## Results

The described procedure was applied to data from the PreSPEC-AGATA setup [7]. The DSSSD detector was mounted close to the position of the secondary target and had  $2 \times 32$  strips, read out with two 32-channel ADCs. It is part of the Lund-York-Cologne CALorimeter (LYCCA) [1] that is tracking and identifying heavy ions.

The result of the procedure is best summarized in a two-dimensional sum histogram of all p-side vs. all n-side amplitudes, without and with calibration coefficients as determined by the described method. In a correctly calibrated detector, each event should have equal calibrated amplitudes for both sides of the detector. That is confirmed by the Fig. 1.

## References

- [1] P. Golubev, et al., Nucl. Instr. and Meth. A 723, (2013), 55-66.
- [2] M. Galassi et al, GNU Scientific Library Reference Manual (3rd Ed.), ISBN 0954612078.
- [3] O.B. Tarasov, et al., Nucl. Instr. and Meth. B 266 (2008) 4657-4664.
- [4] T. Bayes, Philosophical Transactions of the Royal Society, pp. 370-286.
- [5] P.S. Laplace Mémoires de l'Académie royale des sciences, 6, 621-656 (1774).
- [6] D.S. Sivia and J. Skilling, Data Analysis, (2006).
- [7] P. Boutachkov et al., GSI Report 2012 (2013) 148.

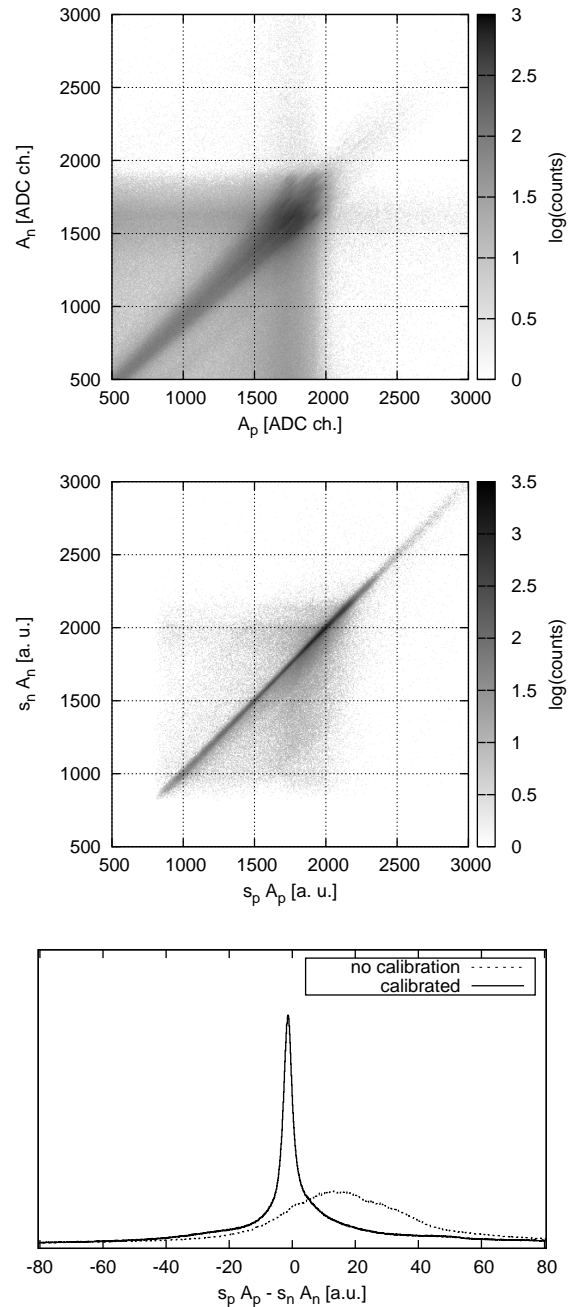


Figure 1: Sum histogram of the p and n-side amplitude distributions before (top) and after (middle) the calibration procedure. The picture on the bottom shows the difference of p and n-side amplitude before and after calibration.