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Effect of Heat Treatments under High Isostatic Pressure on the Transport Critical Current Density at 4.2 K and 20 K in Doped and Undoped MgB₂ Wires

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Abstract: Annealing undoped MgB₂ wires under high isostatic pressure (HIP) increases transport critical current density (J_{tc}) by 10% at 4.2 K in range magnetic fields from 4 T to 12 T and significantly increases J_{tc} by 25% in range magnetic fields from 2 T to 4 T and does not increase J_{tc} above 4 T at 20 K. Further research shows that a large amount of 10% SiC admixture and thermal treatment under a high isostatic pressure of 1 GPa significantly increases the J_{tc} by 40% at 4.2 K in magnetic fields above 6 T and reduces J_{tc} by one order at 20 K in MgB₂ wires. Additionally, our research showed that heat treatment under high isostatic pressure is more evident in wires with smaller diameters, as it greatly increases the density of MgB₂ material and the number of connections between grains compared to MgB₂ wires that smaller wire diameters and high isostatic pressure do not lead to a higher density of MgB₂ material and more connections between grains during the liquid-state Mg reaction.

Keywords: MgB₂ superconducting wires; high isostatic pressure; critical current density

1. Introduction

MgB₂ superconductors have many advantages in that they are cheap components [1] and have a low specific weight [2], low anisotropy [3], high critical temperature [1], low resistivity in the normal state [4], and high B_{c2} [5]. However, they have one large disadvantage: shrinkage during the synthesis reaction. The analysis of the MgB₂ structure shows that the volume decreases by 25% after a reaction [6]. This in turn causes a reduction in the connections between grains and creates surface pinning centers that significantly decrease the critical current density (J_c) in middle and high magnetic fields [7,8]. Doping, e.g., SiC and C, promotes middle and high pinning centers, which increase J_c in middle and high magnetic fields; however, it decreases the critical temperature (T_c) [8,13,14] and increases the irreversibility magnetic fields (B_{irr}) and upper magnetic fields (B_{c2}) [8,10,15–17]. A 2002, an investigation of SiC showed that the addition of nanometer-scale SiC can effectively increase B_{c2} and transport the critical current density (J_c) at high temperatures and fields



Citation: Gajda, D.; Zaleski, A.J.; Morawski, A.J.; Małecka, M.; Nenkov, K.; Rindfleisch, M.; Hossain, M.S.A.; Czujko, T. Effect of Heat Treatments under High Isostatic Pressure on the Transport Critical Current Density at 4.2 K and 20 K in Doped and Undoped MgB₂ Wires. *Materials* **2021**, 14, 5152. https://doi.org/10.3390/ ma14185152

Academic Editors: Gaia Grimaldi, Antonio Leo and Armando Galluzzi

Received: 27 July 2021 Accepted: 6 September 2021 Published: 8 September 2021

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Copyright: © 2021 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). and decrease the anisotropy [11,13,18,19]. In contrast, large SiC grains lead to higher J_{ct} at low temperatures [19]. Additionally, the results showed that SiC nanoparticles with excess Mg improve J_c in low and high magnetic fields [14,18,20–22]. Furthermore, the examination showed that SiC doping increases the magnetic critical current density (J_{cm}) in low and middle magnetic fields between 5 K and 25 K [9]. However, measurements indicated that SiC doping slightly increased J_{cm} in high magnetic fields [9]. Serrano et al. suggested that J_c is increased by strong pinning centers [23]. The measurements indicate that nano SiC increased J_c at 20 K more than other types of doping (e.g., carbon nanotubes) [23] and yielded a high irreversibility magnetic field (B_{irr}) of 7.3 T [16]. An investigation conducted by Li et al. showed that B_{irr} is responsible for improved J_c performance in high magnetic fields [20]. Qu et al. suggested that impurities can increase or decrease J_c in MgB₂ materials [16]. Moreover, Liang et al. indicated that the negative effects on J_c could be attributed to the absence of significant effective pinning centers (Mg₂Si) due to the high chemical stability of crystalline SiC particles [24]. Moreover, more crystalline SiC nanoparticles located at the grain boundaries means more degradation of J_c [24]. The results suggest that the reaction of SiC with Mg might be a necessary condition for the enhancement of J_c in SiC-doped MgB₂ wires [24]. Asthana et al. showed that ethyltoluene and SiC doping significantly increases I_c and B_{c2} in MgB₂ tape more than ethyltoluene doping alone [25]. Additionally, measurements indicate that SiC-doped samples sintered at 650 °C have a better I_c than those sintered at 1000 °C [26]. The results for the SiC-doped MgB₂ material showed that a lower heating time improves J_c while slightly degrading T_c [27].

An investigation performed by Shcherbakova et al. indicated that a long annealing time enhanced the connections between grains [27]. Further results showed that a higher density of MgB₂ material led to better connections between grains [22]. Zhang et al. stated that leftover B constituted an impurity in MgB₂ and degraded the connectivity [22]. The number of connections between grains is strongly related to the amount of impurities [20,22]. Moreover, the results indicated that Mg addition with 10% SiC increased the number of connections between grains [20]. However, Asthana et al. suggested that a smaller amount of MgO led to better connectivity of the MgB₂ grains with SiC doping [25]. Serrano et al. showed that SiC doping weakly increased the number of connections between grains [23]. In addition, the results showed that SiC doping did not change the density of MgB₂ material [13].

Jung et al. indicated that SiC nanoparticles react with Mg above 600 °C and form Mg₂Si compounds [14]. Shcherbakova et al. showed that sintering a SiC-doped sample near the melting point of Mg (640–650 °C) resulted in large numbers of nanoprecipitates and grain boundaries [27]. Furthermore, the investigation showed that a longer cooling time created more Mg₂Si impurities, producing more Mg vacancies [27]. Li et al. stated that liquid Mg first reacts with SiC to form Mg₂Si and is then released free with B to form MgB₂ and released free with Si at high sintering temperatures, leading to strong flux-pinning centers due to small particles being created [8]. The results showed that SiC doping created intragrain defects and a high density of nanoinclusions and impurities (Mg_2Si , Mg_2C_3 , and MgO unreacted SiC), which are effective pinning centers [9,13,21,28]. The results indicate that Mg₂Si inclusions inside MgB₂ grains significantly enhance the pinning strength in high magnetic fields [27]. Furthermore, the results showed that large impure particles such as unreacted SiC would not be effective pinning centers [21]. Dou et al. suggest that nanoscale precipitates with sizes below 10 nm (e.g., Mg₂Si and BC) create strong pinning centers [21]. Shimada et al. showed that nano-SiC easily formed Mg₂Si [27]. Moreover, analysis of the structure of MgB_2 material made by using the internal Mg diffusion (IMD) method showed that large Mg₂Si regions are distributed parallel to the drawing direction [29]. This may suggest that the SiC admixture will behave similarly in MgB2 wires made by the PIT method.

Research shows that both SiC and C doping reduces the a-axis parameter but does not change the c-axis parameter [10,15,30]. The reduction of the a-axis might cause Mg deficiency, oxygen occupancy on B sites, strain (pressure effect—0.5 GPa), and C doping

on B sites [10]. However, Wang et al. stated that Mg vacancies cannot be the origin of a lattice reduction [10]. Furthermore, Dou et al. stated that, because the atomic radius of C is 0.077 nm, that of Si is 0.11 nm and that of B is 0.097 nm [13,28], carbon substitution for boron in the MgB₂ lattice can lead to a large lattice distortion because of the shorter C-B bonds (1.71 A) [17,20,22,26]. Additionally, Kazakov et al. stated that the disorder should lead to a shortening of the Mg-C distance (in the range of 2.25–2.52 Å) [31]. Moreover, carbon substitution for boron leads to increased band scattering [17,20] and significantly enhances B_{c2} and B_{irr} [20]. Further research showed that the proportion of C-added SiC to substitute for B in the lattice is small compared to the pure C substitution case [21,27]. In addition, a small proportion of C substituted for B forms nano-domains in the MgB_2 material [21]. These nanodomains are rectangular and approximately 2–4 nm [21]. Dou et al. stated that C substitution for B caused a reduction in the grain size [26]. Moreover, Serrano et al. suggested that C improved the connections between grains [23]. Others showed that C substitution for B in SiC-doped MgB_2 led to an increase in the normal state resistivity [13,18,20,21]. This high resistivity might suggest that SiC doping creates poor connections between grains [22]. Kazakov et al. showed that C-substituted crystals decrease the coherence length [31]. C substitution for B occurs at temperature as low as $650 \,^{\circ}\text{C}$ [26]. Dou et al. showed that a higher annealing temperature led to a higher level of C substitution than a low annealing temperature [26]. Moreover, Dou et al. and Kazakov et al. indicated that an increase in the C content in MgB₂ material led to an increasingly larger reduction in T_c [26,31]. SiC and organic co-additions may increase the amount of C substitution for B [32]. Mg deficiency will increase as the carbon content increases [31].

Wang et al. suggested that strains are created by a chemical addition and are not influenced by the synthesis parameters because the reduction of the lattice is not observed in pure MgB₂ material [10]. A longer annealing time decreases the lattice strain from 0.45 to 0.4% for pure MgB₂ material [27]. In addition, the quenching of the samples creates more strains in the MgB₂ lattice in both pure and SiC-doped MgB₂ [27]. In contrast, long annealing durations in SiC-doped MgB₂ create more strains (which are formed by Mg₂Si) [27]. The strains lead to a number of crystal defects, such dislocations, which are strong pinning centers [27]. Transmission electron microscopy images showed that nano-SiC doping created a high density of dislocations, stacking faults, a large number of 10 nm-sized inclusions inside the grains [13,21,26,28], and inclusions and nanoscale impurities between grains, e.g., MgB₄ and MgO [21]. Li et al. indicated that thermal strains originating from the interface of SiC and MgB₂ are one of most effective sources of flux pinning centers to improve the supercurrent critical density [33]. Kazakov et al. indicated that local disorder introduced by carbon substitution increased the pinning force [31].

Previous results suggest that Si may also substitute into the crystalline lattice [13]. However, even a high pressure (30 kbar) and high annealing temperature (above 1900 °C) do not lead to a substitution of Si into the crystalline lattice [21]. Wang et al. and Ghorbani et al. also showed that Si cannot be incorporated into the crystal lattice [10,18].

Nanocrystalline SiC does not react with Mg [15,22,24] and is mostly located on the surface boundaries of MgB₂ grains [24]. Unreactive crystalline SiC creates a high density of defects (dislocations and lattice distortion) in the structure of MgB₂ material [15]. However, amorphous SiC reacts with Mg and creates nano-Mg₂Si [15,24]. Amorphous SiC in MgB₂ thick films creates strong pinning centers and improves intergrain connectivity [14]. Additionally, TEM images show strong bonding of MgB₂ with crystalline SiC [15]. Zeng et al. showed that cooling MgB₂ with unreactive crystalline SiC creates thermal strains, which are effective pinning centers [15]. Moreover, Zeng et al. indicated that strain effectively increases J_c [15]. Li et al. suggested that the combination of high connectivity and strong disorder led to high critical parameters [8,20]. Liang et al. suggested that, unlike crystalline SiC nanoparticles located inside MgB₂ grains, crystalline SiC nanoparticles located at grain boundaries may not act as effective pinning centers [24].

Hot pressing increases the reaction rate between Mg and B [16]. Qu et al. showed that the density of the sample increases as the hot pressing temperature increases [16].

Moreover, the hot pressing process can reduce the size of pores, produce small Mg₂Si particles (from 35 nm to 230 nm) with a homogeneous distribution [14], and increase J_c from 5 K to 30 K [16]. Moreover, Qu et al. showed that Mg₂Si is formed at a temperature of 500 °C (hot pressing process) [16]. However, Mg₂Si and MgB₂ are created only at 550 °C [16]. Furthermore, there is no reaction between B and Mg SiC during a hot pressing process at 450 °C [16].

The results showed that annealing under high isostatic pressure (HIP) significantly improved the structure of the MgB₂ material in wires made by using the powder in tube (PIT) method (e.g., small grains, small size, fewer voids, increased connection between grains, increased MgB₂ density, enhanced homogeneity of MgB₂ material, and increased density of structural defects) [34–39]. Additionally, the HIP process increases the density of pinning centers [35] and mainly increases the density of high-field pinning centers [36], leading to a significant increase in J_c under a high magnetic field and B_{irr} and a decrease in T_c [35]. Moreover, HIP increases the rate of reaction and accelerates the rate of carbon substitution for boron [36–38]. Furthermore, annealing under high isostatic pressure increases the melting point of pure Mg, e.g., 0.1 MPa at 650 °C and 1 GPa at 720 °C [40]. This is very important because it permits heat treatment at higher temperatures for pure Mg in the solid state. Thermal treatments in the solid state will increase the density of MgB₂ material [35–39].

The aim of this research was to show the impact of large amounts of 10% SiC doping and high isostatic pressure on J_c at 4.2 K and 20 K and pinning centers. The research shows that the large amount of admixture (10 wt.% SiC) and annealing under high isostatic pressure leads to a significant reduction in the transport critical current density (2 T-100 A/mm²) at 20 K. In addition, the research shows that the HIP process allows one to increase the J_c at 20K in undoped MgB₂ wires. Moreover, the measurements indicated that large amounts of SiC doping and high isostatic pressure lead to a high J_c at 4.2 K.

2. Materials and Methods

2.1. Materials

Multifilament MgB₂ wires with a Nb barrier, Cu matrix, and Monel sheath were produced at Hyper Tech Research, Inc. in Columbus, OH, USA [41]. These wires were made by using the continuous tube forming and filling (CTFF) method. The in situ MgB_2 precursor materials were manufactured with Mg and B powders with 99% purity (with a nominal atomic ratio of 1.1:2). The amorphous boron and magnesium grain sizes were \sim 50 nm and \sim 40 μ m, respectively. The 10 wt.% nano SiC doping MgB₂ wires had 6 filaments and a diameter of 0.83 mm and 0.63 mm. On the other hand, the undoped MgB_2 wire had 18 filaments and a diameter of 0.83 mm. The fill factor was approximately 15%, and the wires were sized to either 0.63 mm or 0.83 mm in diameter. The 0.63 mm samples were obtained from further processing of the 0.83 mm diameter wire. Cold drawing of MgB₂ wire reduced the cross-sectional area of MgB₂ material by approximately 42%, from 0.81 mm² to 0.46 mm². Both 0.63 mm and 0.83 mm diameter samples were heated together at the same temperature (from 680 °C to 725 °C), pressure (1 GPa), and duration (from 15 min to 25 min)—see Tables 1 and 2 [42,43]. For the HIP process, the isostatic pressure is first increased, and then the annealing temperature is increased in the second step. After the required time of the HIP process, the temperature was reduced before the isostatic pressure was reduced. All samples had a length of 100 mm. The HIP process was performed by the Institute of High Pressure Physics, PAS, in Warsaw, Poland.

No.	Annealing Temperature (°C)	Annealing Time (min)	Isostatic Pressure	Wire Diameter (mm)	B _{irr} (4.2 K) T	B _{irr} (10 K) T	<i>B</i> _{irr} (20 K) T	B _{irr} (25 K) T
А	700	15	0.1 MPa	0.83	14.5	10	5.8	3.8
В	700	15	1 GPa	0.83	15	11	6.2	4
С	725	15	1 GPa	0.83	15	11	6.2	4
D	700	25	1 GPa	0.83	15	11	6.2	4

Table 1. Heat treatment parameters of the undoped MgB₂ superconductor wires.

Table 2. Heat treatment parameters of the 10% SiC doped MgB₂ superconductor wires.

No.	Annealing Temperature (°C)	Annealing Time (min)	Isostatic Pressure	Wire Diameter (mm)	B _{irr} (4.2 K)T	B _{irr} (10 K) T	<i>B</i> _{irr} (20 K) T	B _{irr} (25 K) T
Е	700	15	0.1 MPa	0.83	15.2	11	8	
F	700	15	1 GPa	0.83	16	12	7.8	
G	680	15	1 GPa	0.83	16.5	12	7.8	
Η	725	15	1 GPa	0.83	16	12	8	
Ι	700	25	1 GPa	0.83	16	12	8	
J	700	15	0.1 MPa	0.63	15	11	8	3.8
Κ	700	15	1 GPa	0.63	16	12	8	4
L	725	15	1 GPa	0.63	16	12	8	4
М	700	25	1 GPa	0.63	16	12	8	4

2.2. Methods

Critical current (I_c) measurements were made using the four-probe resistive method for a sample of 20 mm length. All samples were measured in a perpendicular magnetic field (Bitter type magnet—14 T) at liquid helium temperature and in accordance with the 1 μ V/cm criterion [44,45]. Measurements were made using a DC source in the range from 0 A to 150 A. All critical current measurements carried out using current sweep type—constant magnetic field and increasing current (from 0 A to 150 A by 90 s). The $I_{\rm c}$ measurements were conducted by the Institute of Low Temperature and Structure Research, PAS, in Wroclaw, Poland. Furthermore, the I_c measurements at 10 K, 20 K, and 25 K were made at the Leibniz Institute for Solid State and Materials Research, Dresden, Germany [46,47]. The integrity of the Nb barriers was checked using the field sweep method and temperature sweep method by the Institute of Low Temperature and Structure Research, PAS, in Wroclaw, Poland [44,45]. The B_{irr} at 4.2 K was determined from the Kramer model ($B^{0.25} \times J_c^{0.5}$). On the other hand, B_{irr} at 10 K, 20 K, and 25 K were specified from the measurements of the critical current ($I_c = 0$ A). Microstructure and composition analyses were performed using an FEI Nova Nano SEM 230 SEM (Hillsboro, OR, USA) by the Institute of Low Temperature and Structure Research, PAS, in Wroclaw, Poland. The structure of the MgB_2 wires studies were investigated using the secondary electron (SE) method. All samples were embedded in carbon resin and were then polished using standard sandpaper and cleaned in isopropanol.

3. Results—Structure

3.1. Microstructure of Undoped MgB₂ Wires

The results in Figure 1a,b show that thermal treatment at 700 °C under a low isostatic pressure of 0.1 MPa leads to the presence of a large number of large voids (up to 5 μ m) with a non-uniform distribution and to a decrease in the MgB₂ material density in superconducting wires. The longitudinal section for sample A in Figure 1b is very important because it helps to better investigate the connection between grains. In MgB₂, the connection between grains significantly influences the transport critical current density properties. Additionally, the SEM images in Figure 1b show that low isostatic pressure annealing creates a thick layer structure (from 2 μ m to 3 μ m). The studies in Figure 1c,d show that heat treatment under high isostatic pressure increases the density of the MgB₂ material, reduces the void number by 90%, improves the homogeneity of the superconducting material, and creates a structure with a thinner layer thickness (below 1 μ m).

between the grains and more connections between the layers. Further studies show that annealing at 725 °C under an isostatic pressure of 1 GPa leads to the creation of more voids (up to 1 μ m) and a structure with thicker layers (up to 2 μ m). This reduces the number of connections between the grains and layers and reduces the density and uniformity of the MgB₂ material in the wire (Figure 1e,f). The tests conducted for sample D show that the longer annealing times at 700 °C and 1 GPa create a superconducting material structure similar to the structure of sample B.



Figure 1. SEM images of undoped MgB₂ wires with a diameter of 0.83 mm: (**a**) cross-section and (**b**) longitudinal section for sample A (0.1 MPa and 700 °C), (**c**) cross-section and (**d**) longitudinal section for sample B (1 GPa and 700 °C), and (**e**) cross-section and (**f**) longitudinal section for sample C (1 GPa and 725 °C). The metal sheath is the Nb barrier.

Jung et al. [6] showed that the MgB₂ material shrinks by 25% during the formation of the superconducting phase. This leads to a significant reduction in the density and uniformity of the superconducting material and intergrain connections. This reduces the transport critical current density in MgB₂ wires. Currently, voids are the biggest problem in MgB₂ materials and limits the application possibilities. Our results show that this problem can be solved by using annealing under high isostatic pressure (1 GPa) in the solid state of Mg. This process leads to a significant reduction in the number of voids. Additionally, our research indicates that thermal treatment under high isostatic pressure and the liquid state of Mg slightly reduce the number of voids in the MgB₂ superconducting material.

3.2. Microstructure of SiC Doped MgB₂ Wires

The results in Figure 2a show that annealing at low pressure creates a larger and greater number of voids (up to $10 \ \mu$ m), leading to a reduction in the connection between grains. Moreover, these voids are heterogeneously distributed in the structure of the MgB₂ material. Previous research has shown that low pressure forms an inhomogeneous distribution of Si particles [37].



Figure 2. SEM images of 10% SiC-doped MgB₂ wires with a 0.83 mm diameter: (**a**) cross-section and (**b**) longitudinal section for sample E (0.1 MPa and 700 $^{\circ}$ C) and (**c**) cross-section and (**d**) longitudinal section for sample F (1 GPa and 700 $^{\circ}$ C). The metal sheath is the Nb barrier.

The SEM image in Figure 2b indicates that low pressure creates large and long voids in a lamellar structure with thick layers (up to 5 μ m). The results in [37] showed that Si particles are located outside and between the lamellar structure. This may be due to the non-substitution of Si into the crystal lattice [10,18,21]. This also leads to reduced connections between grains. Li et al. [20] and Zhang et al. [22] indicated that impurities weaken intergrain connections. In addition, Serrano et al. suggested that a SiC admixture weakly increases the number of connections between grains. This may suggest that the connections between grains are created mainly by Mg grains. On this basis, we can deduce that Mg grains mainly influence the efficiency of connections between grains. Furthermore, the investigation indicated that annealing under a high isostatic pressure of 1 GPa decreased the size by less than 1 μ m and the number (by 90%) of voids, increased the number of connections between grains, created more connections between grains, created a lamellar structure with thin layers (Figure 2c,d), produced smaller Si particles [37], and produced a more homogeneous distribution of Si particles [37]. The SEM images in Figure 3a,b show that a higher annealing temperature (725 °C) and high isostatic pressure of 1 GPa form larger (up to 5 μ m), longer, and an increased number of connections between grains. Furthermore, the SEM images in Figure 3c,d indicate that a longer annealing time at 700 °C and a high isostatic pressure create longer, larger (up to 3 μ m), and an increased number of voids; decrease the density of the MgB₂ material; and decrease the connections between grains to sample F.



Figure 3. SEM images of 10% SiC MgB₂ wires with a 0.83 mm diameter: (**a**) cross-section and (**b**) longitudinal section for sample H (1 GPa, 725 °C, 15 min) and (**c**) cross-section and (**d**) longitudinal section for sample I (1 GPa, 700 °C, 25 min). The metal sheath is the Nb barrier.

However, the longer annealing time produces a higher density structure compared to sample H, which was heated at a higher temperature. The investigation showed that pure Mg at 1 GPa transitions from the solid state to the liquid state at approximately 720 °C [40]. These results indicate that high isostatic pressure increases the effective density of the MgB₂ material during the reaction in solid-state Mg. Moreover, the results indicate that annealing

under high isostatic pressure in the liquid state of Mg produces a similar structure to sample E (annealing under low pressure).

The results in Figure 4a,b show that annealing under low isostatic pressure forms long, large (up to $5 \mu m$) and more numerous voids; reduces the connection between grains; and creates a lamellar structure with thick layers (up to $5 \mu m$). This in turn reduces the density and homogeneity of the MgB₂ material. The SEM images in Figure 4c,d indicate that annealing under high isostatic pressure significantly reduces the number, size and length of voids; increases the connections between grains; and decreases the thickness of layers, increasing the density and homogeneity of the MgB₂ material. A higher annealing temperature with high isostatic pressure forms long, large (up to $5 \,\mu$ m), and more numerous voids; decreases the connections between grains; and creates a lamellar structure with thick layers (Figure 5a,b), decreasing the density and homogeneity of the MgB₂ material. Moreover, the SEM images in Figure 5c,d show that a long annealing time and high isostatic pressure create a small number of large voids (from 1µm to 3µm) and voids with a short length, form a lamellar structure with thin layers (up to 2µm), and slightly reduce the connections between grains compared to sample K. These temperature and pressure parameters only slightly decreased the density and homogeneity of the MgB₂ material for sample M.



Figure 4. SEM images of 10% SiC-doped MgB₂ wires with a 0.63 mm diameter: (**a**) cross-section and (**b**) longitudinal section for sample J (0.1 MPa and 700 °C) and (**c**) cross-section and (**d**) longitudinal section for sample K (1 GPa and 700 °C). The metal sheath is the Nb barrier.



Figure 5. SEM images of 10% SiC-doped MgB₂ wires of 0.63 mm diameter: (**a**) cross-section and (**b**) longitudinal section for sample L (1 GPa, 725 °C, 15 min) and (**c**) cross-section and (**d**) longitudinal section for sample M (1 GPa, 700 °C, 25 min). The metal sheath is the Nb barrier.

3.3. Comparison of the Structure of Undoped and SiC-Doped MgB₂ Wires

Cold drawing of in situ MgB₂ wires causes the Mg grains to elongate [48]. This means that the 0.63 mm wires contain smaller Mg grains than the 0.83 mm wires. The SEM images in Figures 2a,b and 4a,b show that sample J (0.63 mm, 0.1 MPa) has longer, larger, and an increased number of voids than sample E (0.83 mm, 0.1 MPa), indicating that smaller Mg grains react faster than large Mg grains at 700 $^{\circ}$ C (Mg was in the liquid state). The SEM images in Figures 2c,d and 4c,d indicate that the MgB₂ material density is higher in sample K (0.63 mm, 1 GPa) than in sample F, indicating that high isostatic pressure increases the MgB₂ density with smaller Mg grains and a smaller wire diameter (reaction in the solid-state Mg). Moreover, these results suggest that the size of Mg grains does not influence the rate of reaction for the solid-state Mg. Furthermore, the results show that sample L (0.63 mm, 1 GPa, 725 °C) has larger, longer, and an increased number of voids than sample H (0.83 mm, 1 GPa, 725 °C). This indicates that smaller Mg increases the rate of reaction for liquid-state Mg (Figures 2a,b and 4a,b). Furthermore, these results indicate that high isostatic pressure does not increase the density of MgB2 material in liquid-state Mg (smaller diameter and smaller Mg grains). Additionally, the results in Figures 3c,d and 5c,d show that a long annealing time, high isostatic pressure, reaction in the solid-state Mg, and smaller Mg grains increase the density of MgB₂ material in wires. Shcherbakova et al. [27] indicated that MgB2 grains grow faster in the c direction than in the ab direction. This also leads to the formation of a lamellar structure in MgB₂ wires. Studies have shown that the lamellar structure significantly increases J_c and B_{irr} and improves the connections between grains [39]. Our results showed that thinner and longer Mg grains obtain a lamellar structure with thinner layers and improve the longitudinal connection

between grains. The research performed for undoped MgB₂ wires shows that the HIP process yields a higher MgB₂ material density and creates a greater number of connections between grains in smaller diameter wires (d = 0.63 mm [49]) than for wires with a larger diameter of 0.83 mm (Figure 1). Additionally, the results show that the diameter of the undoped MgB₂ wire barely influences the structure of the superconducting material for the synthesis reaction in the Mg liquid state (Figure 1 and [49]).

4. Results—Transport Measurements

4.1. Analysis of the Critical Current Density and Pinning Centers for Undoped MgB₂ Wires with the Diameter of 0.83 mm

The critical current density in MgB₂ wires depends on two factors: the density of pinning centers and the connection between grains. Different types of pinning centers can effectively trap lattice vortices in various ranges of magnetic fields [7,49,50]. Better understanding and explanations of the factors that influence the critical current density [7,50] will be useful for improving superconducting wires for commercial applications. The results in Figure 6a at 4.2 K and 10 K show that thermal treatment under a high isostatic pressure of 1 GPa will increase J_{tc} by 35% in all magnetic fields. These results indicate that the 1 GPa isostatic pressure at 4.2 K and 10 K increase the density of low, middle, and high-field pinning centers (low-field pinning centers to anchor vortex lattice in low magnetic fields, middle-field pinning centers to trap vortex lattice in middle magnetic fields, and high-field pinning centers to anchor vortex lattice in high magnetic fields). Further studies at 20 K show that the HIP process improves J_{tc} by 30% in magnetic fields from 0 T to 4 T. Currently, the Dew-Hughes model is the most used for the study of pinning centers in superconducting wires [51]. This model demonstrates the dominant pinning mechanism. The dominant pinning mechanism indicates the type of pinning centers that are the most abundant in the superconducting material. This allows one to identify structural defects, which in turn allows one to achieve the highest transport critical current density in superconducting wires and tape. The results in Figure 6b indicate that heat treatment under a high isostatic pressure of 1 GPa does not change the dominant pinning mechanism at 20 K. This indicates that the isostatic pressure of 1 GPa at 20 K increases the density of the low- and middle-field pinning centers and slightly increases the density of the high-field pinning centers. Additionally, the results (Figure 6a,b) show that high isostatic pressure annealing poorly increases J_{tc} at 25 K and does not change the dominant pinning mechanism. These results indicate that the HIP process slightly increase the density and improves the pinning centers at 25 K. The results in Figure 6c,d show that higher annealing temperatures and a high isostatic pressure of 1 GPa slightly increase J_{tc} in the range of 4.2 K to 25 K and does not change the dominant pinning mechanism at 20 K and 25 K. The measurement results in Figure 6e indicate that a longer annealing time and an isostatic pressure of 1 GPa slightly increase J_{tc} in the temperature range from 4.2 K to 25 K. In addition, the Dew–Hughes model [51] shows that a longer annealing time under an isostatic pressure of 1 GPa does not lead to changes of the dominant pinning mechanism at 20 K and 25 K (Figure 6f). These results show that increasing the annealing temperature by 25 °C with an annealing time of 10 min under an isostatic pressure of 1 GPa slightly increases the pinning center density in all regimes (low, middle, and high magnetic fields).



Figure 6. (**a**,**c**,**e**) The transport critical current density (J_{tc}) as a function of the perpendicular magnetic field (B) at 4.2 K, 10 K, and 20 K for undoped MgB₂ wires with a 0.83 mm diameter. (**b**,**d**,**f**) The reduced pinning force depending on the reduced magnetic fields for undoped MgB₂ wires.

The results presented in [51] (undoped MgB₂ wires with a diameter of 0.63 mm) and Figure 6a (undoped MgB₂ wires with a diameter of 0.83 mm) show that thermal treatment under high isostatic pressure significantly increases J_{tc} at 20 K in undoped MgB₂ wires with a diameter of 0.63 mm. The measurements shown in [51] (undoped MgB₂ wires with a diameter of 0.63 mm) and Figure 6a (undoped MgB₂ wires with a diameter of 0.83 mm) and Figure 6a (undoped MgB₂ wires with a diameter of 0.63 mm) and Figure 6a (undoped MgB₂ wires with a diameter of 0.83 mm)

indicate that thermal treatment under high and low isostatic pressure increases J_{tc} at 4.2 K in undoped MgB₂ wires with a diameter of 0.83 mm. The SEM images in [52,53] show that 0.63 mm wires after the HIP process have a higher density of undoped MgB₂ material than 0.83 mm wires. The SEM images in [52,53] indicate that 0.63 mm wires after annealing at low isostatic pressure have a slightly lower density for undoped MgB₂ wires with a diameter of 0.83 mm. These results indicate that Mg grains with a longer length and smaller diameter create better connections between grains and more effective pinning centers at 20 K. In addition, the reaction of solid-state Mg (pressure 1 GPa) produces a higher density of low- and middle-field pinning centers. The large Mg grains with a smaller length create better connections and a higher density of pinning centers at 4.2 K. The Dew–Hughes analysis [51] shows that the high isostatic pressure, wire diameter, annealing time, and annealing temperature did not have an influence on the dominant pinning mechanism in undoped MgB₂ wires.

4.2. Analysis of the Critical Current Density and Pinning Centers for SiC-Doped MgB₂ Wires with a Diameter of 0.83 mm

The results in Figure 7a show that annealing under high isostatic conditions increases J_{tc} at 4.2 K in high magnetic fields (above 8 T) but decreases J_{tc} at 4.2 K in low magnetic fields. High isostatic pressure increases the density of high-field pinning centers due to the increasing dislocations, rate of substitution of C for B, strain, and precipitation inside grains [7]. Moreover, annealing with high isostatic pressure decreases the density of lowfield pinning centers (surface pinning centers, e.g., voids [49]). The reduction in voids is caused by the increase in MgB₂ material density (Figure 2). Figure 7a shows that annealing under high isostatic pressure slightly decreases J_{tc} by 8% at 10 K. Further measurements indicate that thermal treatment under high isostatic pressure significantly decreases J_{tc} by a factor of three in low, middle, and high magnetic fields at 20 K (Figure 7a). The transport measurements showed that the HIP process increased J_{tc} at 20 K in undoped MgB₂ wires (Figure 6a). Further studies have shown that SiC doping and thermal treatment under low isostatic pressure (sample E—Figure 7a) increases J_{tc} in the middle and high magnetic fields at 20 K compared to an undoped MgB₂ wire (sample A). This indicates that a large reduction in J_{tc} at 20 K in sample F is the result of the HIP process and a large amount of SiC additive. However, the SEM images show that sample F contains more connections between grains than sample E. The analysis of pinning mechanisms [7,50,53] indicates that samples E and F have the same dominant pinning mechanism at 10 K and 20 K (Figure 7b). In addition, increasing the temperature from 4.2 K to 20 K does not remove the structural defects in the MgB₂ material, which creates pinning centers. Increasing the temperature can only lead to the clustering of several pinning centers into one pinning center. Moreover, increasing the temperature from 4.2 K to 20 K only slightly increases the coherence length [49,54]. This indicates that increasing the temperature does not affect the parameters and efficiency of the pinning centers (e.g., precipitation and impurities). Si particles accumulate at the grain boundaries [5] because Si particles do not substitute in the critical lattice [10,18,21]. Moreover, the Si particles located on the grain boundaries might create dislocations and strains in the MgB_2 grains and connections [19]. Serquis et al. [32] stated that the HIP process creates dislocations in the MgB₂ grains and connections, indicating that SiC doping and the HIP process mainly form pinning centers in the connections between grains and on the grain boundaries. Furthermore, Kazakov et al. [31] and Dou et al. [26] showed that strains (e.g., substitution of C for B) led to a local reduction in T_c . These results indicate that the significant reduction in J_{tc} as a result of increasing the temperature from 4.2 K to 20 K is the result of weakening the connections between MgB₂ grains. The connections worsened due to the clustering of structural defects. Zhang et al. [22] and Li et al. [20] indicated that precipitation led to the weakening of connections between grains. The second reason for the reduction of J_{tc} at 20 K in sample F may be the state of Mg during the synthesis reaction. Research shows that the synthesis reaction in sample E was liquid-state Mg [40]. However, the synthesis reaction in sample F was solid-state Mg [40]. This indicates that liquid-state



Mg creates more effective connections between grains than solid-state Mg at 20 K for MgB₂ wires with a large amount of SiC impurities.

Figure 7. (**a**,**c**,**e**) The transport critical current density (*J*_{tc}) as a function of the perpendicular magnetic field (*B*) at 4.2 K, 10 K, and 20 K for SiC doped MgB₂ wires with a 0.83 mm diameter. (**b**,**d**,**f**) The reduced pinning force depending on the reduced magnetic fields for SiC doped MgB₂ wires.

Figure 7b shows that annealing at 680 °C under a high isostatic pressure of 1 GPa significantly increases J_{tc} in middle and high magnetic fields but decreases J_{tc} in low magnetic fields at 4.2 K compared to samples F and H. Moreover, these results indicate that

annealing at higher temperature (725 °C) leads to increasing J_{tc} in low magnetic fields at 4.2 K. Furthermore, the investigation showed that annealing at 680 °C under an isostatic pressure of 1 GPa leads to a significant decrease in J_{tc} at 10 K and 20 K compared with J_{tc} in samples F and H. However, for sample H (725 °C and 1 GPa), the higher annealing temperature slightly increases J_{tc} at 10 K and significantly increases J_{tc} at 20 K compared with J_{tc} in sample F.

The calculation of the dominant pinning mechanism by using the Dew–Hughes model [51] shows that sample G (1 GPa and 680 $^{\circ}$ C) has a dominant point pinning mechanism between 4.2 K to 10 K (from 0 to 0.35—Figure 7d) and a dominant surface pinning mechanism (from 0 to 0.35—Figure 7d) at 20 K. Figure 7d shows that the longer annealing time (sample H) and the high isostatic pressure of 1 GPa does not change the dominant pinning mechanism at 20 K. The results for sample H suggest that J_{tc} in SiC-doped MgB₂ wires mainly depends on the density and type of pinning centers and does not depend on the connections between grains because sample H has fewer connections between grains compared to samples F and G. Based on [37], sample G is an MgB₂ material with high homogeneity and high density, leading to a large number of connections between the grains. Increasing the temperature from 4.2 K to 20 K does not break the connections between the grains. This indicates that a significant reduction in J_{tc} , similar to sample F, is not related to the destruction of intergrain connections but only to the decrease in the density of pinning centers (the clustering of pinning centers) or a change in the dominant pinning mechanism (e.g., dominant point pinning mechanism at 4.2 K to the dominant surface pinning mechanism (20 K)—Figure 7d). Additionally, the research shows that Si particles do not substitute into the crystal lattice [10,18,21]. This indicates that Si particles are on grain boundaries and connections between grains. Additionally, annealing under an isostatic pressure of 1 GPa creates structural defects [34]. This indicates that the HIP process and 10% SiC doping create a large number of structural defects at grain boundaries and connections between grains (similar to sample F). The clustering of pinning centers, the high density of pinning centers, and the change in the dominant pinning mechanism may lead to a weakening of the connections between grains at 20 K. Furthermore, a low annealing temperature (680 °C), 10% SiC doping, and reaction in solid-state Mg might create a weak connection between grains at 20 K.

The results in Figure 7e show that a longer annealing time and high isostatic pressure significantly increase J_{tc} at 4.2 K in low and middle magnetic fields. Additionally, a longer annealing time and high isostatic pressure appear to not change J_{tc} at 10 K and 20 K. Moreover, the Dew–Hughes model [51] indicated that a longer annealing time and the isostatic pressure of 1 GPa does not significantly change the dominant pinning mechanism at 20 K (Figure 7f). However, a longer annealing time and high isostatic pressure only slightly increase the density of high-field pinning centers. The SEM images showed that the density of MgB₂ and the size and number of voids in samples F and I are similar, indicating that the substitution of C for B during reaction in solid-state Mg requires a longer annealing time (more than 25 min).

Based on the Dew–Hughes model [51], it can be indicated that a significant reduction in J_{tc} at 20K is caused by the change of the dominant pinning mechanism from point to surface and a decrease of the density of pinning centers.

4.3. Analysis of the Critical Current Density and Pinning Centers for SiC-Doped MgB₂ Wires with a Diameter of 0.63 mm

The results in Figure 8a (0.63 mm diameter) show that annealing under a high isostatic pressure of 1 GPa increases J_{tc} in high magnetic fields and decreases J_{tc} in low and middle magnetic fields by between 4.2 K and 25 K. The analysis of the pinning centers (Figure 8b) by using the Dew–Hughes model [51] shows that thermal treatment under high isostatic pressure of 1 GPa does not change the dominant pinning mechanism in a range from 4.2 K to 25 K. This indicates that the reduction of J_{tc} (Figure 8a) in low and middle magnetic fields is mainly due to the reduction of the pinning center density. In addition, the Dew–Hughes model [51] indicated the appearance of the dominant point pinning mechanism in a range

from 0.75 to 1. This leads to the increase of J_{tc} in the high magnetic fields. The SEM images in Figure 4 show that sample K has a much higher density of superconducting material than sample J. These results indicate that a higher density of MgB₂ material might decrease the density of surface and point pinning centers.



Figure 8. (**a**,**c**,**e**) The transport critical current density (J_{tc}) as a function of the perpendicular magnetic field (B) at 4.2 K, 10 K, and 20 K for SiC doped MgB₂ wires with a 0.63 mm diameter. (**b**,**d**,**f**) The reduced pinning force depending on the reduced magnetic fields for SiC doped MgB₂ wires.

The results in Figure 8c (0.63 mm diameter) show that a higher annealing temperature and higher isostatic pressure slightly increase J_{tc} at 4.2 K and significantly decrease J_{tc} by between 10 K and 25 K. The pinning center analysis indicates that a higher annealing temperature and higher isostatic pressure slightly increase the density of pinning centers and does not change the dominant pinning mechanism at 4.2 K (Figure 8d). From 10 K to 25 K, the higher annealing temperature and high isostatic pressure decrease the density of each type of pinning center. The SEM images in Figure 4a,b and Figure 5a,b show that sample J (0.1 MPa, 700 °C) and sample L (1 GPa, 725 °C) have a similar number of connections between grains. The Dew–Hughes analysis [51] indicates that the reduction of J_{tc} at 20 K and 25 K is mainly caused by the reduction of the pinning center density (Figure 8d), because the dominant pinning mechanism does not change significantly at 20 K and 25 K.

The results in Figure 8e (0.63 mm diameter) show that a longer annealing time and high isostatic pressure significantly increase J_{tc} in low and middle magnetic fields between 4.2 K and 25 K and decrease J_{tc} in high magnetic fields in the same temperature range. The Dew–Hughes model shows that the longer annealing time and 1 GPa isostatic pressure does not lead to changes of the dominant pinning mechanism at 4.2 K, 20 K, and 25 K (Figure 8f). This indicates that the longer annealing time and HIP process allows one to increase the pinning center density. The SEM images in Figures 4c,d and 5c,d show that sample K has more connections between grains than sample M. These results indicate that J_{tc} in sample M is dependent on pinning centers rather than the connection between grains. Furthermore, the SEM images (Figure 5c,d) indicate that voids, Mg₂Si particles, and Si particles can create low and middle pinning centers [49].

4.4. Comparison of the Results of J_{tc} and Pinning Mechanisms for Undoped and 10 wt.% MgB_2 Wires

Earlier studies have shown that MgB₂ wires with a large amount of SiC admixture should be fabricated with excess Mg or by maximizing the amount of MgB₂ superconducting phase [14,18,20–22]. Studies conducted by Li et al. showed that the Mg₂Si phase is formed first, followed by the MgB₂ phase and free Si in SiC-doped MgB₂ superconductors [8]. In addition, studies have shown that hot pressing (HP) accelerates the formation of the Mg₂Si phase (500 °C [16]). This indicates that the heat treatment under isostatic pressure accelerates the formation of the MgB₂ phase. Images from a scanning electron microscope with a backscattering electron (BSE) function do not show pure Mg.

The results presented by Susner et al. [48] showed that cold drawing of unreacted MgB_2 wires reduced the thickness and increased the length of Mg grains. The results for sample E (0.1 MPa, 0.83 mm diameter) and sample J (0.1 MPa, 0.63 mm diameter) indicate that J_{tc} in the temperature range from 4.2 K to 20 K is similar. J_{tc} is similar in undoped MgB₂ wires with diameters of 0.63 mm and 0.83 mm after heat treatment at low isostatic pressure. This indicates that the size and length of Mg grains during reaction in the liquid state of Mg does not influence J_{tc} or the different types of pinning centers in SiC-doped and undoped MgB2 wires. Furthermore, for annealing sample F (0.83 mm) and sample K (0.63 mm) under high isostatic pressure, the smaller Mg grains produced in the reaction in solid-state Mg does not increase J_{tc} at 4.2 K but significantly increases J_{tc} at 10 K and 20 K. On the basis of the SEM images of samples F and K, one might deduce that the higher density of the MgB₂ material and more connections between grains increase the density of low- and middle-field pinning centers and significantly increases the density of high-field pinning centers. The measurements for undoped wires (sample B—0.83 mm and [52]—d = 0.63 mm) show that the smaller wire diameter and HIP process significantly increases J_{tc} in low and medium magnetic fields at 20 K but does not improve J_{tc} in high magnetic fields at 20 K. In addition, the investigation for sample H (0.83 mm, 725 °C) and sample L (0.63 mm, 725 $^{\circ}$ C) showed that a higher annealing temperature, high isostatic pressure, smaller Mg grains, and a reaction in the liquid state of Mg had no influence on J_{tc} at 4.2 K but significantly decreased J_{tc} at 10 K and 20 K for both samples because the structures of samples H and L are similar. These results suggest that J_{tc} is mainly dependent on the density of pinning centers. The low-field pinning centers do not act in high magnetic fields [49] and high temperatures. The measurements indicate that a higher annealing temperature (725 °C), HIP process, and smaller wire diameter do not increase Jtc at 4.2 K and 20 K in undoped MgB2 wires. However, a very high annealing temperature ([52]—740 °C) leads to a significant reduction in J_{tc} at 20 K in undoped MgB₂ wires. The results for sample I (0.83 mm) and sample M (0.63 mm) showed that a longer annealing time, high isostatic pressure, smaller Mg grains, and reaction in solid-state Mg slightly increased J_{tc} at 4.2 K and significantly increase J_{tc} at 10 K and 20 K. The SEM images showed that sample M has a higher density of MgB₂ material than sample I. These results indicate that a higher density of MgB₂ material and more connections between grains produces a high density of pinning centers.

Comparing the results of SiC-doped and undoped [52] MgB₂ 0.63 mm wires after heat treatment (700 °C for 15 min) at low isostatic pressure shows that doping increases J_{tc} in high magnetic fields (e.g., SiC doped at 4.2 K, 100 A/mm² in 10 T and undoped at 4.2 K, 100 A/mm² at 8 T [51]) and slightly decreases J_{tc} in low magnetic fields between 4.2 K and 20 K. SEM images show that the structures of doped and undoped samples are similar (grain size and number of connections between grains), indicating that the SiC dopant creates high-field pinning centers and reduces the density of low-field pinning centers. Similar results were obtained for SiC-doped and undoped MgB₂ wires sized 0.83 mm (Figures 6a and 7a).

Comparing the results for doped and undoped [52] MgB₂ 0.63 mm wires after heat treatment at 725 °C under high isostatic pressure shows that there is a significant reduction in J_{tc} between 10 K and 20 K in MgB₂-doped wire. These samples have a similar structure according to the SEM images, suggesting that, in both doped and undoped [51] MgB₂ wires, the reaction in liquid-state Mg is mainly dependent on the density of pinning centers and much less dependent on intragrain connections. Figures 6e and 7e indicate that a longer annealing time under high isostatic pressure slightly increases J_{tc} in doped and undoped MgB₂ 0.83 mm wires.

On the basis of the Dew–Hughes model [51], it can be indicated that the large amount of SiC doping creates the greater number of low-field pinning centers [49] at 20 K than undoped MgB₂ material. Additionally, the Dew–Hughes analysis [51] shows that wires with a 0.63 mm diameter and a large amount of dopant have more low-field pinning centers than doped MgB₂ wires with 0.83 mm diameter at 20 K.

Kim et al. [55] indicated that amorphous boron and excess Mg yield higher J_{tc} than crystalline boron. Comparing MgB₂ wires with a 10% SiC admixture (amorphous boron, Mg_{1.1}B₂) and MgB₂ wires with a 2% C admixture (crystalline boron, MgB₂ [36]), the J_{tc} of the SiC-doped sample is higher than the J_{tc} in 2% C-doped MgB₂ wires between 10 K and 25 K after heat treatment at low isostatic pressure. Further results indicate that heat treatment under high isostatic pressure significantly increases J_{tc} in 2% C-doped MgB₂ wires (crystalline boron) between 10 K and 25 K compared to SiC-doped and undoped MgB₂ wires [51]. These results show that too much dopant and too many structural defects (HIP process, cooling, without additives) lead to a reduction in J_{tc} at high temperatures.

5. Conclusions

The combination of 10 wt.% SiC doping and low isostatic pressure increases the density of high-field pinning centers and reduces the density of low- and middle-pinning centers between 4.2 K and 25 K. The SiC doping and high isostatic pressure significantly reduce the density of low- and middle-field pinning centers and significantly increase the density of high-field pinning centers in wires with a diameter of 0.63 mm between 4.2 K and 25 K. The efficiency of high isostatic pressure applied to 0.83 mm wires is much lower than that applied to 0.63 mm wires.

Mg grains with a smaller diameter and longer length correspond to higher J_{tc} at 20 K compared to SiC-doped and undoped wires with Mg grains with a larger diameter and shorter length. This indicates that longer Mg grains with smaller thicknesses create better connections between the grains and more efficient pinning centers at 20 K. Large grains with a shorter length create better connections between grains and more pinning centers at 4.2 K. This may be related to the lower density of MgB₂ material with a 0.83 mm diameter (more voids that form both low- and middle-fields pinning centers). A larger number of voids at 20 K may lead to the weakening of the connections between grains because several voids and weak areas of superconductivity are combined into one large non-superconducting area.

The large amount of 10 wt.% SiC admixture decreases J_{tc} at 20 K in both 0.63 mm and 0.83 mm wires after the HIP process. This may be associated with too many structural

defects at grain boundaries and connections between grains. These defects at 4.2 K create a high density of different types of pinning centers. Increasing the temperature from 4.2 K to 20 K causes the clustering of several structural defects into one structural defect (transition of pinning centers, e.g., from dominant point pinning mechanism to dominant surface pinning mechanism), leads to a decrease in the density of the pinning centers, and weakens the connections between grains at 20 K.

The J_{tc} in SiC-doped and undoped MgB₂ wires after reaction in the liquid state of Mg is mainly dependent on the density of the pinning centers. High isostatic pressure during the liquid-state reaction does not increase the density of MgB₂ material in 0.83 mm and 0.63 mm wires. The J_{tc} in SiC-doped and undoped wires after the Mg solid-state reaction is mainly dependent on the connections between the grains and the pinning centers. This is especially evident at 20 K as SEM studies show that high isostatic pressure increases the density of MgB₂ material more strongly in SiC-doped and undoped wires with a smaller diameter (0.63 mm) than in wires with a larger diameter (0.83 mm).

Our research shows that MgB₂ wires used in devices operating at the temperature of liquid helium can be designed with a large amount of admixture and can be thermally treated under high isostatic pressure for optimal superconductivity properties. However, MgB₂ wires used in devices operating at 20 K must have a small amount of admixture when heated under high isostatic pressure.

Author Contributions: D.G.: conceptualization, investigation, formal analysis, and writing—review and editing; A.J.Z.: conceptualization and formal analysis; A.J.M.: investigation methodology and visualization; M.M.: investigation and formal analysis; K.N.: investigation and resources; M.R.: formal analysis and resources; M.S.A.H.: formal analysis and visualization; and T.C.: conceptualization, visualization, and formal analysis. All authors have read and agreed to the published version of the manuscript.

Funding: This research was funded by the Institute of Low Temperature and Structure Research, PAS; the Institute of High Pressure Physics, PAS; and the Ministry of National Defense Republic of Poland Program—Research Grant MUT Project 13-995.

Institutional Review Board Statement: Not applicable.

Informed Consent Statement: Not applicable.

Data Availability Statement: The data presented in this study are available on reasonable request from the corresponding authors.

Acknowledgments: The authors hereby thank Wolfgang Häßler with the Institute for Solid State and Materials Research Dresden, Germany for his help in performing the critical current measurements in 10 K–25 K.

Conflicts of Interest: The authors declare no conflict of interest.

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