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Automatic Multi-Objective Optimization of Coarse-Grained Lipid Force Fields Using SwarmCG

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ABSTRACT The development of coarse-grained (CG) molecular models typically requires a timeconsuming iterative tuning of parameters in order to have the approximated CG models behaving correctly and consistently with, e.g., available higher-resolution simulation data and/or experimental observables. Automatic data-driven approaches are increasingly used to develop accurate models for molecular dynamics simulations. But the parameters obtained via such automatic methods often make use of specificallydesigned interaction potentials, and are typically poorly transferable to molecular systems or conditions other than those used for training them. Using a multi-objective approach in combination with an automatic optimization engine (SwarmCG), here we show that it is possible to optimize CG models that are also transferable, obtaining optimized CG force fields (FFs). As a proof of concept, here we use lipids, for which we can avail of reference experimental data (area per lipid, bilayer thickness) and reliable atomistic simulations to guide the optimization. Once the resolution of the CG models (mapping) is set as an input, SwarmCG optimizes the parameters of the CG lipid models iteratively and simultaneously against higherresolution simulations (bottom-up) and experimental data (top-down references). Including different types of lipid bilayers in the training set in a parallel optimization guarantees the transferability of the optimized lipid FF parameters. We demonstrate that SwarmCG can reach satisfactory agreement with experimental data for different resolution CG FFs. We also obtain stimulating insights on the precision-resolution balance of the FFs. The approach is general and can be effectively used to develop new FFs, as well as to improve



existing ones.

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I. INTRODUCTION

Molecular dynamics (MD) has become a fundamental tool in the study of complex molecular systems, providing high-resolution insights often inaccessible via experimental techniques. One of the main limitations of all-atom (AA) MD simulations is the space and time scales accessible with current computational capabilities. Coarse-grained (CG) molecular models, in which groups of atoms are represented as larger CG particles (or beads), may alleviate these issues and are increasingly employed to study systems of interest in structural biology^{1–3}, drug discovery^{4,5}, biophysics^{6–8} and nanomaterials design^{9–} ¹¹. Lipid bilayers, key components of the cell membranes, are a notable example of supramolecular systems exhibiting properties which, to a large extent, require CG models to be studied effectively¹².

Different approaches are typically used for the development of CG force fields (FFs)^{13–26}. Top-down strategies essentially aim at reproducing molecular properties observed experimentally with the CG models. Conversely, *bottom-up* approaches rely on calibrating CG model parameters using equilibrium simulations of higher-resolution molecular models (e.g., AA). Widely used for the simulation of lipids, the popular Martini²⁴ FF, for example, presents aspects of both. The possibility to parametrize molecular models for a variety of molecules using transferable CG beads makes the Martini²⁴ FF versatile. However, general CG FFs remain intrinsically approximated in modeling specific molecular systems compared to CG models that, albeit less general and transferable, are optimized ad hoc to this end^{18,27}.

Methods such as, e.g., Inverse Monte Carlo (IMC)¹³, Iterative Boltzmann Inversion (IBI)¹⁵, Multi-State IBI (MS-IBI)¹⁹, Force Matching (FM)^{16–18,28}, ForceBalance^{29,30}, Relative Entropy Minimization (REM)³¹, the generalized Yvon-Born-Green (g-YBG)³² equation, and different flavors of Particle Swarm Optimization (PSO)^{21,22,33,34} have been used as basis to build *bottom-up* and/or *top-down* FF parametrization approaches to calibrate AA FFs relying, e.g., on quantum mechanical data^{30,35,36}, or to calibrate CG FFs based on AA MD trajectories^{37–42,34,43,44}. For what pertains to the *bottom-up* route, their parameters extraction schemes are based on, e.g., reproducing pair distribution functions 13,15,19, matching forces 16-18,28-30, minimizing the information loss in terms of relative entropy³¹, or on the liquid state theory³² (see, e.g., Kmiecik et al. 45 for an exhaustive review of these approaches).

More recently, the evolution of machine learning approaches is considerably accelerating the development of accurate CG molecular models. Deshmukh et al. developed a CG FF for different solvents^{21,22}, hydrocarbons⁴⁶, small peptides³³ and several polymers⁴⁷, optimizing the interaction parameters to reproduce experimental observables (exclusively top-down) using particle swarm optimization⁴⁸ (PSO) and artificial neural networks (ANN)-assisted PSO²¹. Force matching^{16–18,28} has been reformulated as a supervised learning problem in CGNet²⁵, using ANN and point forces as features to learn the potential of mean force of a polypeptide in water. Automatic learning of both the CG FF and its "functional form" (abstract featurization) was then introduced in CGSchNet²⁶ using graph convolutional neural networks (GCNN).

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Such approaches, relying exclusively on equilibrium AA trajectories for CG FF calibration (exclusively bottom-up), are exposed to MD sampling problems and potential inaccuracies in the reference AA FF (which are then transferred into the CG FF). Automatic approaches capable of optimizing the accuracy of the CG models, while at the same time guaranteeing their transferability to different conditions or system variants than those used during FF calibration, would be fundamental for the development of nextgeneration transferable CG FFs.

Using phosphatidylcholine (PC) lipids as a test case, here we describe an automatic multi-objective optimization approach that allows to develop accurate and transferable CG lipid FFs. We build on SwarmCG, 23 a CG FF optimization algorithm based on fuzzy self-tuning PSO⁴⁹ (FST-PSO), recently developed to optimize bonded parameters in CG molecular models. We designed a comprehensive general strategy that now allows SwarmCG to optimize also the non-bonded parameters of a FF in order to improve the accuracy of the CG models. We chose lipids as a test case, as an example of molecular systems for which we can avail of experimental data (e.g., area per-lipid, bilayer thickness, etc.) and of reliable all atom FFs. Here, in a new multi-objective version of SwarmCG (https://github.com/GMPavanLab/SwarmCGM), we provide the opportunity to combine bottom-up and top-down reference information for calibrating CG lipid FFs using simultaneously high-resolution (AA) MD simulations and the experimental data. Using new metrics based on optimal transport (OT)⁵⁰ for deriving the bonded and non-bonded interaction terms of the CG FF, we show that SwarmCG can simultaneously and iteratively optimize CG models of different types of PC lipids in parallel, improving the transferability of the optimized CG parameters among different lipid types, and also to those which are not included in the training set. Several FF calibration experiments demonstrate that this multi-objective approach can be successfully applied for generating new and custom lipid FFs across different resolutions. Furthermore, a stress-test of SwarmCG against a state-of-the-art CG lipid FF (Martini 3.0)²⁴ proves the robustness of the software.

II. METHODOLOGY

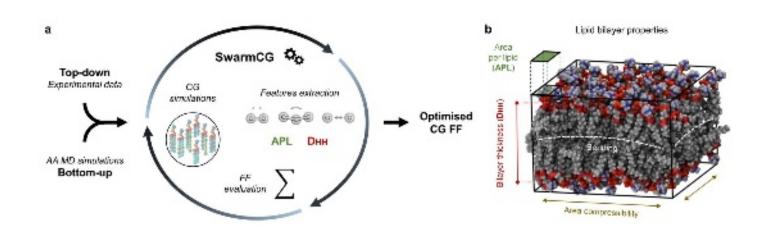
A. Optimization strategy

Our multi-objective CG FF optimization strategy relies on the complementary use of structure-based information from high-resolution molecular simulations (bottom-up: AA MD), providing knowledge on the submolecular structure and dynamics of the systems, and of experimental data (top-down: e.g., area per lipid, bilayer thickness), used to guide the calibration of the models on a larger scale (Fig. 1a). Similar multi-objective strategies, based on the simultaneous combination of simulation and experimental data used as the references for FFs fitting, have been employed also by others for the development, e.g., of different types of AA FFs^{29,36,51,52}. In our approach, the discrepancies observed between the data provided by the CG models and the *bottom-up* and *top-down* reference data are measured *via* a global scoring function (Eq. 1) and minimized through an iterative optimization procedure. Executed in parallel using in the training set

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multiple bilayers, each composed of a different type of lipid, our procedure can output optimal versions of CG lipid FFs, offering optimal consistency with the experimental data set as target. Importantly, the quality and completeness of the information embedded in the training set (e.g., number of different types of lipids, number of different temperatures used, resolution and topology of the CG representations, accuracy of the AA MD simulations and of the experimental data set used as the targets) directly conditions the accuracy of the CG FFs optimized via SwarmCG, as well as their capacity to transfer to other types of lipids (not included in the training set).

For our demonstrations, here we use in the training sets up to 5 PC lipids that span a range of different tail characteristics (length, unsaturation, and combinations of those), for which accurate experimental measurements for the area per lipid (APL) and phosphate-to-phosphate bilayer thickness (DHH) are available from lamellar bilayer isolates in the liquid phase^{53,54} (Fig. 1b, Table S2). As proof of concept of this method for optimizing CG FFs for lipid models with different resolutions, we first focus on parametrizing custom CG models where the lipid molecules are represented at low resolution using either 5 beads per lipid (lowest resolution), or 6-8 beads per lipid (slightly higher resolution). We then test the same approach on high resolution PC lipid models (3-5 heavy atoms per bead). In such a case, as a stress-case for the method, we start from the state-of-the-art version of the widely used Martini FF (i.e., Martini 3.0) 24 . Given the notable level of usage and testing of this CG FF (especially for the simulation of lipids), we use this as a control case to check that SwarmCG does not produce CG parameters deviating much from a FF that is already quite evolute in terms of accuracy and reliability. All the models in these demonstrations make use of the same simple and computationally-efficient FF functional form (i.e., the standard one used in Martini²⁴ lipid models, see also Sec. S1), with *bonded* interactions described by harmonic terms for bonds and angles, and non-bonded interactions described by Lennard-Jones (LJ) and Coulomb potentials. The parameters of the CG FFs are iteratively optimized using FST-PSO⁴⁹ (one of the most efficient PSO variant to date⁵⁵) and by running at each iteration 200 ns of CG MD simulation of lipid bilayer patches composed of 128 lipids (which preliminary tests indicated being enough to reach successfully the MD equilibrium and sufficient sampling in the simulated CG bilayer systems), from which the scores of the CG models are measured and used for improving the FFs accuracy according to a loss function.



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Figure 1. Multi-objective optimization process implemented in SwarmCG. (a) SwarmCG simultaneously relies on bottom-up and top-down references to iteratively optimize CG lipid FFs using higher-resolution AA MD simulations and experimental data. (b) Properties of lipid bilayers (APL and D_{HH}) used for calculating the top-down component of the loss function.

In classic PSO⁴⁸, a swarm of individuals (referred to as "particles", each representing a set of parameters to be optimized) moves iteratively inside a bounded multidimensional search space and cooperates to identify the best solution for a problem according to a loss function, without using analytical gradients. Settings referred to as "social" and "cognitive" attraction respectively favor the collaboration among particles and their tendency to rely on individual experience. In FST-PSO⁴⁹, faster convergence is achieved with the introduction of fuzzy logic for adjusting attraction settings independently for each particle and dynamically during optimization. This approach is particularly competitive in computationally expensive black-box optimization problems and effectively handles noisy data. Applied to our CG FF calibration problem, this enables using a variational principle for high-dimensional FF parameterization with limited concern over the impact of the noise originating from MD sampling. In our demonstrations, the loss function evaluates the CG FF parameters of up to 5 different lipids (Fig. 2a,b) simultaneously (i.e., each particle of the swarm is tasked with running up to 5 CG MD simulations). The optimization problem is thus formulated for maximizing the thermodynamic consistency of the optimized CG FF for all tested PC lipids (potentially simulated using different temperatures), increasing the sampling of the data while at the same time limiting the number of local minima on the loss surface. Therefore, the computational cost of running multiple simulations for calculating the relevance of a single set of FF parameters (i.e., a particle of the swarm) is compensated by an even faster convergence, enabled by an information-rich loss function. Because swarm optimization already constitutes an embarrassingly parallel workload, and each particle is tasked with running independent MD simulations (i.e., another layer of parallelization), this approach also efficiently leverages high-performance computing (HPC) resources (see Sec. S2).

B. Loss function and optimal-transport-based metrics

1. Top-down components

Our process is based on the minimization of a loss function encompassing the distances from the bottom-up and top-down target objectives. Our construction of the loss function aims at: (i) reducing a many-objective optimization problem to a single-objective one (global FF accuracy score); and (ii) using as priority reference the available experimental data (APL and DHH), while the features calculated from AA MD simulations (which may suffer from sampling issues or FF inaccuracies) are used only as guidance during optimization, and for restricting the number of possible solutions to a given optimization problem. The loss function takes the form

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where ΔAPL_{global} and ΔD HH $_{global}$ are the aggregated APL and DHH deviations with respect to experimental values, calculated across all CG lipid bilayers used for an optimization (loss components 1 and 2) as

$$\Delta APL_{global} = \sqrt{\frac{\sum (w_1 + \min(max(0, \Delta APL_{lt} - E_{tol}), \Delta APL_{cap}))^2}{L_q}},$$
(2)

and

$$\Delta D_{\text{HH}_{global}} = \sqrt{\frac{\sum (w_1 + \min (max(0, \Delta D_{\text{HH}_{lt}} - E_{tol}), \Delta D_{\text{HH}_{cap}}))^2}{L_g}},$$
(3)

where ΔAPL_{lt} and ΔD_{HH}_{lt} are the APL and DhH absolute % deviations with respect to the experimental values for a bilayer composed of lipid type l simulated at temperature t, E_{tol} is set to 1.5 and represents the tolerated error in ΔAPL_{lt} and ΔD_{HH}_{lt} accounting for experimental error in measurements, w_1 is a weight (empirically set to 10) that prioritizes using the target experimental data over the AA reference data during the FF optimization, L_g is the number of different lipid bilayer simulations used in the training set (a bilayer of lipid type l can be simulated at different temperatures), and ΔAPL_{cap} and ΔDHH_{cap} are set to 30 and used to cap ΔAPL_{lt} and ΔD_{HH}_{lt} values for limiting noise during the first steps of an optimization procedure, by allowing to disregard uninformative loss values that can be produced when putative CG FF parameters would induce a disassembly or explosion of the bilayer during MD simulations. By capping ΔAPL_{global} and ΔD нн $_{global}$ values, then OT- B_{global} and OT- NB_{global} (Eq. 8, 10) are able to guide the optimization even in otherwise potentially uninformative conditions. We define the convergence criterion as 10 swarm iterations without improving loss.

2. Optimal-transport-based metrics

Regarding the bottom-up component, as preliminary steps we obtain well-sampled equilibrium AA MD trajectories of lipid bilayers to be used as references for each lipid to be used in the training sets (see Sec. S3), and map the AA lipid models at the desired CG resolution (Fig. 2a). The AA-to-CG mapping determines the chemical identity/correspondence of each CG bead, bond and angle, defining also the number of parameters to be optimized in the CG FF (Figs. 2c, additional details provided in Sec. S4). Reference bond and angle distributions, as well as the distance distributions between each type of particle (within a 25 Å cut-off), are computed from each AA-mapped MD trajectory, and compared to those calculated using the corresponding CG models at each iteration during optimization (inexpensive via MDAnalysis^{56,57}). Altogether, the discrepancies between such average AA and CG quantities measure how closely a putative CG FF matches the AA description of the molecular systems¹⁸.

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For evaluating the mismatch between corresponding AA vs. CG bond and angle distributions, we employ the OT-based Wasserstein distance^{58,59} (a.k.a., Earth Movers' Distance, EMD) with an underlying symmetric and positive-definite distance matrix (hereafter referred to as "OT-B metrics", Fig. 2d). This metrics has been already proven well-suited for parametrizing the bonded terms of CG models of complex and flexible molecules in a previous version of SwarmCG²³. Noteworthy, OT-based metrics offer several interesting features: (i) multimodal distributions are properly handled; (ii) distances are robust to noise; (iii) distances are quantified in interpretable units (e.g., Å, degrees); and (iv) their computations are inexpensive. In particular, here we introduce a new metrics that relies on OT for comparing the spatial distribution of particles in equilibrium MD trajectories (hereafter referred to as "OT-NB metrics", Fig. 2e). The OT-NB metrics employs the Wasserstein distance^{58,59} on the distance distribution between particles, with an underlying distance matrix accounting for the differences in between radial shell volumes. This metrics can be considered as an OT-based adaptation of the Kirkwood-Buff integrals⁶⁰, which is particularly well-suited for quantifying differences in the spatial organization of particles in molecular systems described at different resolutions (e.g., AA vs. CG).

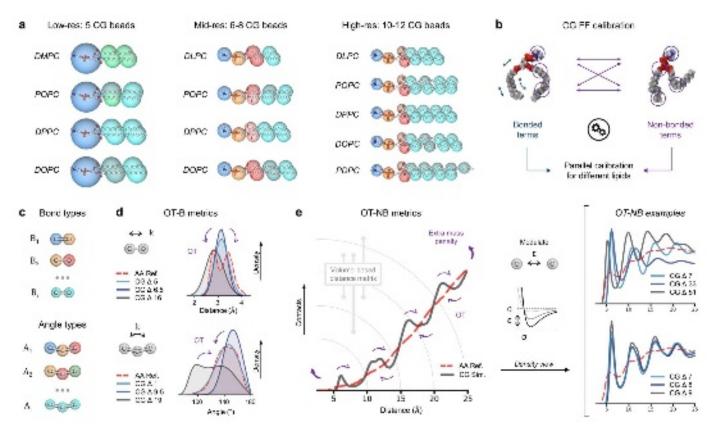


Figure 2. Description of the CG representations applied and the metrics used in this modified version of *SwarmCG*. (a) AA vs. CG mapping for the benchmark lipids used in the training set for optimization at different resolutions: 1,2dilaurovl-sn-glycero-3-phosphocholine (DLPC), 1.2-dimyristovl-sn-glycero-3-phosphocholine (DMPC), 1.2dipalmitoyl-sn-glycero-3-phosphocholine (DPPC), 1-palmitoyl-2-oleoyl-glycero-3-phosphocholine (POPC), 1,2dioleoyl-sn-glycero-3-phosphocholine Δ9-Cis (DOPC) and 1-palmitoyl-2-docosahexaenoyl-sn-glycero-3phosphocholine $\Delta 4,7,10,13,16,19$ -Cis (PDPC, 16:0-22:6). (b) Bottom-up components of the scoring function: SwarmCG optimizes in concert the bonded and non-bonded terms of a CG FF, iterating CG MD simulations of bilayers composed of different types of lipids. (c) CG bonds and angles are classified according to the CG beads involved, and are attributed specific parameters. (d) Principle of the OT-B metrics used for structure-based information related to bonded FF terms. (e) Principle of the OT-NB metrics used for structure-based information related to non-bonded FF terms.

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Given 2 histograms P, Q the EMD as initially defined by Rubner $et\ al.^{59}$ is:

$$EMD(P,Q) = \frac{\min\limits_{\{f_{ij}\}} \sum_{i,j} f_{ij} d_{ij}}{\sum_{i,j} f_{ij}},$$
(4)

with $f_{ij} \ge 0$, $\sum_j f_{ij} \le P_i$, $\sum_i f_{ij} \le Q_j$ and $\sum_{i,j} f_{ij} = \min(\sum_i P_i, \sum_j Q_j)$, where $\{f_{ij}\}$ represents the optimal transport plan, each f_{ij} represents the amount transported from the i supply bin to the j demand bin, and d_{ij} is the distance matrix between bin i and bin j in the histograms. In this study, the domains of the bond and angle distributions D_b and D_a are set to [0, 50] Å and [0, 180] degrees, respectively, for the distributions of all AA-mapped and CG bond and angle, with histogram bandwidths set to 0.1 Å and 2 degrees. The domain of all the AA-mapped and CG distance distributions between pairs of bead types D_p is set to [0, 15] Å and the bandwidth used is 0.25 Å. All EMD calculations used in this study rely on the implementation of PyEMD^{58,61}.

To calculate the EMD between 2 corresponding AA-mapped vs. CG bonds or angles, noted OT- $B_{bond(i,l)}$ and OT – $B_{angle(i,l)}$ in the next section, we use normalized histograms ($\sum_i {P'}_i = 1$ and $\sum_j {Q'}_j = 1$) and a symmetric positive-definite distance matrix. For obtaining the EMD between 2 corresponding AA-mapped vs. CG distance distributions between pairs of bead types, noted $OT-NB_{bead\ pair(l)}$ in the next section, we first normalize the bins in each of the 2 histograms so that $\sum_i P'_i + \sum_j Q'_j = 2$, with $P'_i = \frac{2P_i}{\sum_i P_i + \sum_j Q_j}$ and $Q'_j = \frac{2 Q_j}{\sum_i P_i + \sum_i Q_i}$. We then define a distance matrix that accounts for the differences in between radial shell volumes as:

$$d_{ij} = \begin{cases} i < j, & \frac{V_j}{V_i} \\ i = j, & 0 \\ i > j, & \frac{V_i}{V_j} \end{cases}, \tag{5}$$

with V_i and V_j the volumes of the radial shells for bin i and j. The optimal transport plan $\{f_{ij}\}$ calculated here consists in a partial matching of the 2 compared histograms^{58,61} because of the normalization we apply. We finally account for the extra or missing mass that is left out of $\{f_{ij}\}$ using:

$$OT - NB_{bead\ pair(l)} = \left(EMD(P', Q') + \left(\frac{\max(\sum_{i} P_{i}, \sum_{j} Q_{j})}{\min(\sum_{i} P_{i}, \sum_{j} Q_{j})} - 1\right)\right) * 100.$$
 (5)

Because the distance matrices d_{ij} are symmetric, the OT-B and OT-NB distances inherits the properties of metrics^{58,61}.



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Finally, for the definition of our global loss function (Eq. 1), we aggregate the OT-B and OT-NB distances obtained from the comparison between a set of CG MD simulations (i.e., a particle of the swarm) and their corresponding reference AA MD trajectories (bottom-up), together with the discrepancies observed between CG vs. experimental APL and Dнн measurements (top-down). The OT-B distances from the reference AA MD trajectories are calculated as:

$$OT - B_{bond\ type} = \sqrt{\frac{\sum (w_2 \times OT - B_{bond(i,l,t)})^2}{B_{i,l,t}}},$$
(6)

where OT- $B_{bond\ type}$ quantifies the deviation of the CG models from the reference AA trajectories in terms of bond distributions for a given bond type, $B_{i,l,t}$ is the number of instances of this bond type in CG model topologies across all simulations used in an optimization, OT- $B_{bond(i,l,t)}$ is the OT-B distance for each instance of this bond type across all simulations used in an optimization (a bond of a given type can be present multiple times in the topology of a single lipid model), w_2 is a weight that prioritizes minimizing the OT-B distances of the bonds over those of the angles (see below); and

$$OT - B_{angle\ type} = \sqrt{\frac{\sum (OT - B_{angle(i,l,t)})^2}{A_{i,l,t}}},\tag{7}$$

where OT- $B_{angle\ type}$ quantifies the deviation of the CG models from the reference AA trajectories in terms of angle distributions for a given angle type, $A_{i,l,t}$ is the number of instances of this angle type in CG model topologies across all simulations used in an optimization, OT – $B_{angle(i,l,t)}$ is the OT-B distance for each instance of this angle type across all simulations used in an optimization (an angle of a given type can be present multiple times in the topology of a single lipid model); and

$$OT - B_{global} = \sqrt{\frac{\sum OT - B_{bond \ type}^2 + \sum OT - B_{angle \ type}^2}{B_g + A_g}},$$
 (8)

where OT- B_{global} is the global OT-B deviation score for the CG FF being optimized (loss component 3), B_g is the number of different bond types in this FF and A_g the number of different angle types in this FF. The OT-NB distances from the reference AA MD trajectories are calculated as:

$$OT - NB_{bead\ pair} = \sqrt{\frac{\sum (OT - NB_{bead\ pair(l)})^2}{P_l}},$$
(9)

where OT-NB_{bead pair} quantifies the deviation of the CG models from the reference AA trajectories in terms of distance distributions between pairs of bead types, $OT-NB_{bead\ pair(l)}$ is the OT-NB distance for this pair of bead types calculated from each CG MD simulation in which this interaction is sampled, P_l is the number of instances of this pair of bead types across all lipids used in an optimization; and

$$OT - NB_{global} = \sqrt{\frac{\sum OT - NB_{bead\ pair}^2}{P_g}},\tag{10}$$



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where OT-NB_{global} is the global OT-NB deviation score for the CG FF being optimized (loss component 4), and P_a is the number of different pair types in this FF.

Therefore, the components of the loss function are empirically weighted according to only 2 parameters: (i) W_1 slightly prioritizes minimizing the APL and DHH discrepancies over the OT-B and OT-NB distances, which regulates the extent to which structure-based information is discarded for better fitting experimental measurements, and (ii) w₂ allows to obtain comparable OT-B metrics for the bond and angle deviations and is set to 50, meaning that an OT-B of 0.4 Å between corresponding CG vs. AA-mapped bond distributions (noted OT – $B_{bond(i,l)}$) is considered equivalent to an OT-B of 20 degrees between corresponding CG vs. AA-mapped angle distributions (noted OT – $B_{anale(i,l)}$).

4. Modulation of the loss function for mixed bottom-up and top-down FF calibration

The simple form used for the aggregation of the different objectives in the loss function (Eq. 1) enables modularity. In this modified version of SwarmCG, if no AA trajectories are provided for a given optimization procedure, the bottom-up components of the loss function described in Eq. 1 will be discarded and the loss will automatically become:

$$loss' = \sqrt{\frac{\Delta APL_{global}^2 + \Delta DHH_{global}^2}{2}}.$$
 (11)

User-defined configuration files allow complete modularity of the parameters being optimized and of the reference data being used, notably for performing CG lipid FF calibrations in a mixed bottom-up and topdown fashion. For example, this allows to make use of bottom-up reference AA trajectories only for lipids for which an accurate AA FF is available (see Sec. S3 and Table S1), while parameters that are specific to CG lipids for which the AA FFs are still inaccurate to date (e.g., containing highly unsaturated tails, such as SDPC and PDPC⁵³) can be calibrated exclusively with respect to top-down experimental data. In this case, these parameters can be evaluated in parallel and in context with other FF parameters subjected to both bottom-up and top-down reference calibration.

III. RESULTS

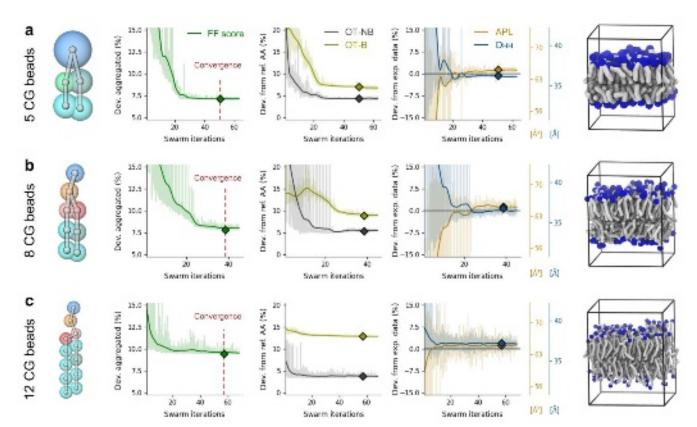
A. Multi-objective automatic optimization of CG models of POPC at different resolutions

As a first demonstration of this approach, we use SwarmCG to perform individual multi-objective (bottomup plus top-down) optimizations of 3 different-resolution CG models of POPC. In particular, for this first test we use: (I) a low-resolution (5 beads) implicit-solvent CG model, (II) a mid-resolution (8 beads) implicit-solvent CG model, and (III) a high-resolution (12 beads) explicit-solvent CG model of a POPC

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bilayer composed of 128 lipids (64 per leaflet). These three CG models are then optimized individually, with the prime purpose to check how SwarmCG copes with the different resolution of the models.

In the case of (I) and (II), the parametrization of the coarser CG models is performed completely ab-initio (Fig. 2a,b). In such cases we initially set all parameters randomly, while in the first particle in the first swarm iteration all LJ ε are set to 4 kJ/mol (it can be a random number, but too low or too high ε values should be avoided in order to prevent all CG MD simulations to crash in the first swarm iteration). All parameters are then iteratively changed by SwarmCG with the goal of minimizing the loss function. In particular, the following terms of the CG FFs are calibrated: (i) equilibrium values for bonds and angles; (ii) force constants for bonds and angles (bonded terms); and (iii) LJ σ and ε parameters defining all interactions between pairs of CG bead types (non-bonded terms). For obtaining the LJ σ parameters, we optimize the values of the radii attributed to each bead type and apply the Lorentz-Berthelot rule⁶² (see also Sec. S5 and S6 in the SI). As initial parameters, we use: (i) the average equilibrium values of the bonds and angles computed from AA-mapped MD trajectories; (ii) arbitrary values for all force constants; and (iii) arbitrary values for the bead radii and LJ ε parameters. The green curves in Fig. 3a,b (loss) indicate that the optimizations converged after ~50 and ~40 swarm iterations, respectively. Independently of the resolution applied, at convergence the models correctly reproduce the APL and DHH experimental data used as target (Fig. 3a,b: yellow and blue curves converging to reference black line, set to 0). The OT-B and OT-NB distances are effectively minimized (Fig. 3a,b: black and olive curves), indicating that both CG models globally reproduce the structural features present in the reference AA MD trajectories (Figs. S3-4: lowresolution).



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Figure 3. Multi-objective optimization of CG FF parameters in models of POPC described at different resolutions, at 303K in the liquid phase. For each modelling resolution: (a) 5 CG beads, (b) 8 CG beads, (c) 12 CG beads – (Left panels) green curve: loss during optimization; (Middle panels) black and olive curves: OT-B and OT-NB metrics for the specific lipid; (Right panels) yellow and blue lines: APL and DhH for the specific lipid model during optimization, displayed with window-averaging (solid) and without (shaded); (Right images) CG MD snapshots of the lipid bilayers for the FF obtained at each resolution. The horizontal black lines set at 0 identify the target experimental APL and Dhh values. Diamonds represent values at convergence, obtained with the selected optimized CG FF parameters.

In the case of (III), we test our approach in explicit solvent starting from the current high-resolution Martini 3.0²⁴ model of POPC (Fig. 2c). The following terms of the CG FF are iteratively optimized: (i) the equilibrium values for bonds and for angles not initially set to 180; (ii) force constants for bonds and angles; and (iii) LJ ε parameters defining all solute-solute interactions between pairs of CG bead types. LJ σ parameters, solute-solvent and solvent-solvent interactions are unchanged. As initial parameters, we use the existing ones in Martini 3.0^{24} for this model. The optimization procedure converged after ~60 swarm iterations (Fig. 3c). At convergence the model of POPC correctly reproduces the APL and DHH experimental data used as target (Fig. 3a,b: yellow and blue curves), whereas the original Martini 3.0²⁴ model produces a small offset (~6 %) on DhH with respect to experimental data. The OT-B distances reached a higher plateau (Fig. 3c: olive curve) than in the 2 previous cases, due to the CG topologies used in Martini 3.0²⁴ lipid models, which creates small offsets in the CG vs. AA-mapped fits of the bond distributions (see Sec. S7). This has no effect on OT-NB distances, which were minimized to a plateau similar to the 2 previous cases (Fig. 3c: black curve).

This first experiment demonstrates the ability of SwarmCG to balance bonded and non-bonded interaction terms to optimize the CG lipid models, independently of the resolution used. In all three cases, the obtained CG models for POPC self-assemble into bilayers and allow to simulate vesicle fusion (Fig. 4: e.g., 5 and 8 CG beads models) consistently with previous studies^{63,64}. In terms of computational time, the *ab-initio* calibration of the low-resolution POPC model (Fig. 3a, 5 CG beads and 26 FF parameters) required 36 hours (wall-clock time) to reach 60 swarm iterations using 20 particles in the swarm and using 20 CPUs (each CG simulation running on a single CPU, allowing for a complete parallelization of the swarm of particles). The ab-initio calibration of the mid-resolution POPC model (Fig. 3b, 8 CG beads and 55 FF parameters) required 2 days to reach 50 swarm iterations using 25 particles in the swarm and requesting 25 CPUs. The optimization starting from the existing Martini 3.0²⁴ POPC model (Fig. 3c, 12 CG beads and 40 FF parameters) required 8 days to reach 70 swarm iterations using 23 particles in the swarm and requesting 23 CPUs. We underline that while the time and computational cost required for the optimization may seem non-negligible, the benefit makes the process balance favorable, especially considering that the result of SwarmCG in this case, as it will be better described in the next section, is an optimized CG FF.

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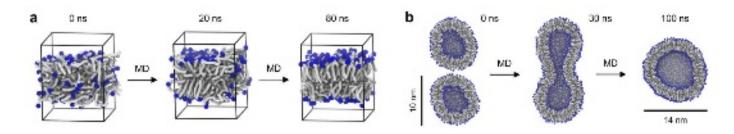


Figure 4. Behavior of the POPC models calibrated *ab-initio* in different configurations of unbiased MD simulations in implicit solvent. (a) CG MD snapshots of the low-resolution 8 beads POPC models undergoing self-assembly into a bilayer starting from a dispersed configuration (128 lipids, 310K). (b) CG MD snapshots of the mid-resolution 5 beads POPC models undergoing vesicular fusion starting from 2 pre-assembled vesicles (1196 lipids each, 310K).

Such individual optimizations demonstrate that, from a technical point of view, *SwarmCG* can produce CG models fitting with the experimental data. However, the following comments on the robustness of the results are necessary. As said above, there is no guarantee that optimized CG parameters obtained in such a way are transferable to simulate other lipid systems (*vice versa*, typically, they do not). Furthermore, the individual optimization of multiple parameters (*e.g.*, all *non-bonded* LJ terms) to fit, *e.g.*, the APL of POPC, could lead to potential artifacts. In fact, automatic optimization approaches, such as this one, cope badly with resolving a high-dimensional problem using low-dimensional criteria. Typically, the method can find multiple combinations of LJ parameters that allow reaching the APL target. While the use of multiple references (*bottom-up* and *top-down*) may to some extent alleviate such limitation, here we designed a different approach that imparts not only robustness, but also transferability to the results of *SwarmCG* in this sense (see next sections).

B. Parallel ab-initio calibration of low-resolution implicit-solvent CG lipid force fields

In order to provide more information for guiding the optimization, here we introduce an additional transferability constraint. This consist in the parallel optimization of multiple lipid bilayer systems of different types, which are optimized as in the previous section in an iterative way against *bottom-up* and *top-down* references. We conducted this test using, *e.g.*, POPC, DMPC, DOPC and DPPC lipid bilayers (Fig. 5a). The difference from the previous experiment is that in this case, at every iteration, the parameters change affects all CG beads of the same type in all 4 lipid systems in the same way. Furthermore, the discrepancies between all 4 systems and their *bottom-up* and *top-down* references become all part of the scoring function which is iteratively minimized by *SwarmCG*. This introduces an automatic transferability constraint, as the number of parameters combinations that can minimize the distance from all objectives in all cases are strongly reduced. In a sense, this is an endogenous constraint, in that the automatic optimization self-regulates under the condition that this has to satisfy and optimize all systems (and not only one, POPC, as in the previous tests of Fig. 3).

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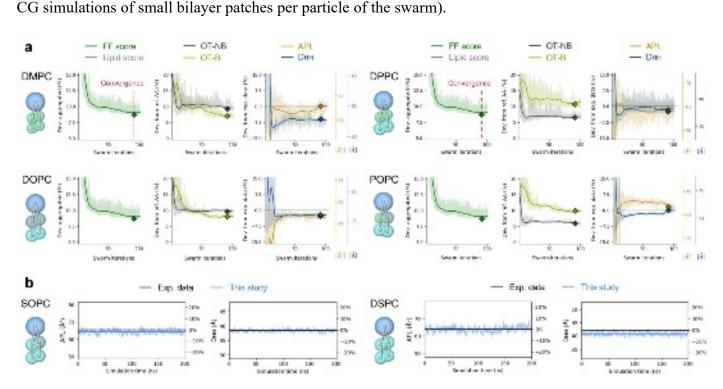


Figure 5. Multi-objective ab-initio calibration of CG FF parameters using 4 low-resolution models of PC lipids in the training set, in the liquid phase and in implicit solvent (DMPC: 303K, DPPC: 323K, DOPC: 303K, POPC: 303K). (a) For each lipid – (Left panels) green curve: loss during optimization, grey curve: loss calculated per lipid; (Middle panels) black and olive curves: OT-B and OT-NB metrics for the specific lipid; (Right panels) yellow and blue lines:

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APL and DHH for the specific lipid during optimization, displayed with window-averaging (solid) and without (shaded). The horizontal black lines set at 0 identify the target experimental APL and DHH values. Diamonds represent values at convergence, obtained with the selected optimized CG FF parameters. (b) APL and Dhh measured across 200 ns for SOPC (303K) and DSPC (333K) using transferred FF parameters.

Considering that the obtained CG FF parameters are transferable between 4 different PC lipids, a relevant question is to what extent these are transferable also to other PC lipids which are not in the training set, and which structure can be described by combinations of the same CG beads (and by their bonded parameters). To tackle such a question, starting from the optimized CG FF parameters obtained at convergence with SwarmCG in Fig. 5a, we apply these parameters a posteriori to parametrize 2 additional low-resolution models of SOPC and DSPC lipids, and simulate the corresponding bilayers composed of 128 lipids in the liquid phase, at 303K and 333K respectively. It is worth noting that these models can be assembled using the previously calibrated CG beads as building blocks (non-bonded parameters), but most of the bonded parameters are still unknown. According to our topology definition, only those of the unsaturated tail of POPC can be recycled (see Sec. S4). This test therefore required another preliminary optimization procedure for calibrating exclusively the 12 unknown bonded parameters (3 bonds and 3 angles: equilibrium values and force constants), which we perform this time in a fully top-down fashion (no AA trajectories are provided for reference, as described in Sec. II B 4). The calibration of the additional 12 bonded parameters required 10 hours of computation to reach 20 swarm iterations using 17 particles and requesting 34 CPUs (2 CG simulations of small bilayer patches per particle of the swarm).

The APL and DHH measured for the obtained CG models of SOPC and DSPC across 200 ns of MD simulation of patches of bilayers of 128 lipids (parametrized using the transferred CG parameters obtained from the optimization of Fig. 5a) are found in very good agreement with experimental data (Fig. 5b: blue data from the CG models vs. black experimental target APL and DhH data), with a maximum deviation of <3% in the calculation of the Dнн of DSPC. This is a remarkable result, as this second test using the lowresolution models of Fig. 5 demonstrates the ability of our procedure to generate ab-initio CG lipid FFs even using incomplete training sets. Namely, the CG parameters obtained by optimizing the FF including some lipids in the training set can be then transferred also to lipid molecules that are outside the training set. It is worth noting that here we used as training set 4 CG simulations of small bilayer patches composed of 4 different types of lipids. In this way, we could further calibrate FF parameters for modelling different types of lipids (e.g., SOPC and DSPC) with a reduced computational effort, by transferring some of the previously calibrated FF parameters, and calculating only the missing parameters. In comparison, obtaining the single POPC model in Fig. 3a required the calibration of 26 FF parameters, while obtaining 4 and then 6 different lipid models required the calibration of only 42 and 54 FF parameters in total, respectively. While this is not the central purpose of this work, this demonstrates how the present approach could be scaled for the calibration of general FFs for many different types of lipids. Importantly, it is not forcefully required to

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C. Parallel ab-initio calibration of mid-resolution implicit-solvent CG lipid force fields

In this section we also demonstrate the *ab-initio* calibration of an implicit-solvent PC lipid CG FF using 4 different lipid types in the training set, while in this case the lipids are represented by mid-resolution CG models (6-8 CG beads per lipid). The combined bottom-up and top-down multi-objective approach is always the same, and the conditions of the optimization are identical to these of the mid-resolution experiment in Sec. III A, for the types of parameters being optimized in the FF and the initialization of the swarm of particles. Also in this case, SwarmCG outputs transferable optimized CG parameters for POPC, DMPC, DOPC and DPPC lipid bilayers matching the experimental APL and DHH data for all 4 lipids within the experimental error (Fig. 6a: yellow and blue curves). The OT-B and OT-NB distances are effectively minimized (Fig. 6a: black and olive curves), with the structural features of the AA trajectories being reproduced in the CG models (Figs. S5-6). This ab-initio FF calibration required 40 hours of computation to reach 40 swarm iterations using 27 particles and requesting 54 CPUs (79 FF parameters, 4 CG simulations of small bilayer patches per particle of the swarm).

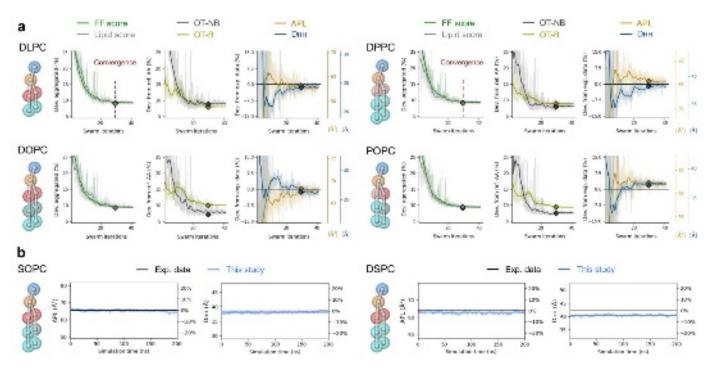


Figure 6. Multi-objective ab-initio calibration of CG FF parameters using 4 mid-resolution models of PC lipids in the training set, in the liquid phase and in implicit solvent (DLPC: 303K, DPPC: 323K, DOPC: 303K, POPC: 303K). (a) For each lipid – (Left panels) green curve: loss during optimization, grey curve: loss calculated per lipid; (Middle panels) black and olive curves: OT-B and OT-NB metrics for the specific lipid; (Right panels) yellow and blue lines: APL and DHH for the specific lipid during optimization, displayed with window-averaging (solid) and without (shaded). The horizontal black lines set at 0 identify target experimental APL and Dнн values. Diamonds represent values at convergence, obtained with the selected optimized CG FF parameters. (b) APL and DHH measured across 200 ns for SOPC (303K) and DSPC (333K) using transferred FF parameters.



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Also in this test, we evaluate the transferability of the optimized FF parameters obtained at convergence by building mid-resolution models of SOPC and DSPC and simulating bilayers composed of 128 lipids in the

building mid-resolution models of SOPC and DSPC and simulating bilayers composed of 128 lipids in the

liquid phase, at 303K and 333K respectively. In this case, this requires calibrating just one additional angle

(equilibrium value and force constant), which we did again in a fully top-down fashion. The APL and Dнн

measured for the obtained models of SOPC and DSPC using transferred FF parameters are found again in

very good agreement with experimental data (Fig. 6b), and results are overall comparable to those of the 5

beads per-lipid model. The calibration of the 2 additional bonded parameters required 3 hours of

computation to reach 15 swarm iterations using 3 particles and requesting 16 CPUs (2 CG simulations of

small bilayer patches per particle of the swarm).

Altogether, the results presented in this and in the previous section demonstrates the effectiveness of SwarmCG, used with the parallel multi-objective paradigm presented herein, for producing optimized CG lipid FF of variable resolution. The input required for the process are: (i) a pre-defined CG mapping, which defines the resolution in the description of the lipids in the system (in this sense, a too low resolution -e.g., below 5 beads per lipid, was found to produce poor results and inefficiency in reaching the objectives), (ii) available experimental data and (iii) reliable higher-resolution for the systems that one wants to model. As anticipated, the computational cost for the process, which may seem non-negligible, is fully compensated by the fact that, in a few days, it is possible to obtain fairly accurate and transferable CG FFs. In this sense, we point out that the information present in the training set is important. For example, in the tests described above we obtain optimized CG FF parameters that can be transferred among PC lipids (and in particular, for those lipids which share the lipid types optimized herein). More complete FFs can be for sure obtained by including additional systems in the training set (e.g., PA, PG lipids, etc.). In such a case, the cost/benefit balance of the process is of prima importance. Noteworthy, the fact that SwarmCG is demonstrated to produce good results already with a minimal training set is very promising, as this could open the possibility to enrich the obtained FFs adding a limited number of systems to the training set. While future employment and testing of SwarmCG to develop new FFs for a growing number of classes of molecular systems will provide insights onto the limits and opportunities in terms of obtaining reliable FFs even using incomplete information, the results reported herein aim at demonstrating the potential of the approach, and to prove the robustness of the parameters that can be obtained with SwarmCG. An important point to guarantee is that such powerful data-driven approaches do better, comparably, or at least not worse than humans. With this aim, we designed the last control in silico experiment described in the following section.

D. Control test: using SwarmCG starting from the Martini lipid force field parameters

As a last test case, to control the robustness of the approach, we applied the same approach using the current version of the (explicit solvent) lipid models available in Martini 3.0²⁴. This time, we used 5 different high-resolution PC lipids to be included in the training set, which optimization is again iteratively conducted in



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The conditions of the optimization are identical to these of the high-resolution experiment in Sec. III A, for the types of parameters being optimized in the FF and the initialization of the swarm of particles. After ~30 swarm iterations the procedure reaches convergence (Fig. 7a: green curves), identifying a set of FF parameters that moderately improves the overall matching with experimental APL and Dhh values set as target for the 5 lipids in the dataset (Fig. 7a: yellow and blue curves). This FF optimization required 7 days of computation to reach 30 swarm iterations using 23 particles and requesting 60 CPUs (42 FF parameters, 5 CG simulations of small bilayer patches per particle of the swarm). We point out that here we are not aiming to provide an updated version of Martini 3.0 for lipids. A richer training set would be needed to this end. Furthermore, one should guarantee that the changes in the optimized LJ ε parameters obtained for the beads do not produce an impairment in the partition free-energies of the Martini beads (which have been developed based on such concept). However, the results of such test encourage us on the robustness of SwarmCG. For example, in all 5 cases we obtain results that variate just by some percentage point compared to the results obtained from the same 5 lipids with Martini 3.0. The result of Fig. 7a demonstrates that while in some cases we obtain a slight improvement (e.g., PDPC, POPC and DPPC), the optimized CG parameters perform slightly worse than Martini 3.0 parameters for other ones (e.g., DOPC and DMPC). This shows that in this case *SwarmCG* struggles in finding a way to improve these parameters, and this is as expected, considered the evolute version of this FF. This demonstrates that SwarmCG is found robust in this control experiment. The fact that the software cannot change much FFs, such as Martini 3.0, which are already quite optimized, provides even more value to the results obtained in the tests discussed in the two previous sections.



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Figure 7. Multi-objective optimization of CG FF parameters using 5 high-resolution models of PC lipids in explicit solvent in the training set (DMPC: 303K, DPPC: 323K, DOPC: 303K, POPC: 303K, PDPC: 303K). (a) For each lipid – (Left panels) green curve: loss during optimization, grey curve: CG FF score for the specific lipid; (Middle panels) black and olive curves: OT-B and OT-NB metrics for the specific lipid; (Right panels) yellow and blue lines: APL and Dhh for the specific lipid during optimization, displayed with window averaging (solid) and without (shaded), and for Martini models (dotted). The horizontal black lines set at 0 identify target experimental APL and Dhh values. Diamonds represent values at convergence, obtained with the selected optimized CG FF parameters. (b) APL and Dhh measured across 200 ns for SDPC (303K), DLPC (303K), SOPC (303K) and DSPC (333K) using transferred FF parameters (blue) and for Martini models (yellow).

Also after such a control test, we tried to transfer the obtained FF parameters for modelling DLPC, DSPC, SOPC and SDPC and simulate bilayers in the liquid phase (Fig. 7b), staying within the framework of Martini 3.0²⁴ lipid models. A known limit with the current CG representations employed in Martini is that they do not allow differentiating between some types of lipids^{27,65}. Here, the same models are used for representing, enumerated by pairs: DLPC and DMPC, SDPC and PDPC, DPPC and DSPC, and POPC and SOPC (Table S2). Therefore, while at convergence the experimental APL and Dhh values were better fitted for the models of DPPC and POPC (Fig. 7a), the deviation from experimental data is increased when considering the same models for the representation of DSPC and SOPC, respectively (Fig. 7b). Similarly, the model obtained at convergence for DMPC slightly decreased the agreement with experimental Dhh value (Fig. 7a), with respect to the original Martini 3.0 model, while the optimized model applied for

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modeling DLPC yielded better agreement for this observable (Fig. 7b). Such additional tests demonstrate again that the SwarmCG optimized parameters perform comparably to the Martini 3.0 ones even when transferred to model lipids not included in the training set.

IV. CONCLUSIONS

We introduce an automated multi-objective FF optimization strategy and demonstrate its efficiency for developing optimized CG lipid FFs of variable resolution. The approach is general and can be used for optimizing explicit- or implicit-solvent FFs, with a variable resolution in the CG representation of the lipids. Newly devised OT-based metrics allow to quantify differences in the spatial organization of particles in molecular systems described at different resolutions. The parallel optimization of multiple lipid systems maximizes the transferability of the CG FF, and mitigates potential inaccuracies in the structure-based information (suboptimal AA FFs or limited MD sampling).

An intrinsic drawback of multi-parametric FF optimization is that multiple sets of parameters may produce the same results (i.e., a high-dimensional problem is here formulated as a low-dimensional one). The transferability constraint, here induced by the fact that the building blocks of the FF (bonded and nonbonded interactions between CG particle types) are conserved in different lipid systems and are optimized in parallel, efficiently guides the optimization and discards spurious solutions. As demonstrated, this allows to build a solid framework of shared and sampled interactions, eventually allowing to transfer the optimized CG parameters to model other lipids. While an optimized CG FF becomes more complete as the diversity of the training set increases (more interactions between different beads become better sampled in the context of different molecular systems), our demonstrations indicate that the multiple interactions present in such complex bilayer systems guarantee a satisfactory transferability even using a reduced number of lipids in the training set. For the lipids within the training set, SwarmCG generates optimized FF parameters that systematically reproduce target experimental observables. As a proof-of-concept, here we optimize CG FFs for PC lipids described at different resolutions. At the same time, the obtained parameters can be transferred to other types of lipids, suggesting that substantial improvements could be achieved also in high-resolution models and even using limited data with this procedure. Increasing the diversity of the training set, including also other types of lipids, should allow to obtain even more general and accurate CG lipid FFs.

This automated multi-objective CG FF optimization strategy is general, and essentially requires that: (i) reference AA MD trajectories can be obtained for part of the lipids used in the training set (bottom-up) requirement); (ii) reliable experimental data are available (top-down requirement); and (iii) the CG MD simulations used for testing the FF parameters are computationally accessible and sufficiently informative of the quality of the FF being optimized. Lipids thus constitute an appropriate use case, but the process can also be extended to parametrize FFs also for other classes of molecules for which reference experimental/simulation data are available. Because the form of the loss function used in this study is

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simple, it would be straightforward to input additional top-down observables to be evaluated in simulations, and to be quantified as percentages of deviation from target values (i.e. similar to modifying the loss function from Eq. 11 to Eq. 1). Moreover, this approach makes an efficient use of HPC resources and scales efficiently. In fact, while enlarging the training set may seem to increase the computational burden, this makes the transferability constraints more informative and the convergence of the optimization faster. The cost of the optimization process is fully compensated by the benefit, considered that optimized and transferable force fields can be obtained in a few days of calculation. Given the revolutions that transferable FFs brought to the molecular modeling community, we envisage that multi-objective optimization approaches, such as that presented here based on SwarmCG, will have great impact in the evolution towards next generation FFs.

SUPPLEMENTARY MATERIAL

Supplementary material includes details on the functional form of the CG FFs for which the parameters are optimized, the various-resolution molecular models used in these experiments, their topologies, their optimized FF parameters obtained with SwarmCG, as well as the implementation for usage with HPC resources. Additional details are also provided concerning the submolecular features observed in the CG models obtained at the end of different FF optimization experiments, underlining the relevance of the OT-B and OT-NB metrics.

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AUTHOR DECLARATIONS

Conflicts of Interest

The authors have no conflicts of interest to disclose.

Author Contributions

C.E.M. and R.C. devised the OT-based metrics, the algorithm and built the molecular models. R.C., M.P., C.C. and G.D. performed the experiments. C.E.M. implemented the algorithm. C.E.M. and G.M.P. wrote the paper and designed the research. G.M.P. supervised the work. All authors agreed on the final form of the paper.

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DATA AVAILABILITY

The code used in this study for the optimization of the CG lipid models is available at:

https://github.com/GMPavanLab/SwarmCGM, together with the models obtained at convergence for each experiment, as well as the configuration files used in this study, and all material necessary for running the software and for reproducibility testing.

REFERENCES

- 1. Abellón-Ruiz, J. et al. Structural basis for maintenance of bacterial outer membrane lipid asymmetry. Nat. Microbiol. 2, 1616-1623 (2017).
- 2. Van Eerden, F. J., Melo, M. N., Frederix, P. W. J. M., Periole, X. & Marrink, S. J. Exchange pathways of plastoquinone and plastoquinol in the photosystem II complex. Nat. Commun. 8, 15214 (2017).
- 3. Yen, H.-Y. et al. PtdIns(4,5)P 2 stabilizes active states of GPCRs and enhances selectivity of G-protein coupling. Nature 559, 423-427 (2018).
- Hoffmann, C., Centi, A., Menichetti, R. & Bereau, T. Molecular dynamics trajectories for 630 coarse-grained drug-membrane permeations. Sci. Data 7, 51 (2020).
- 5. Souza, P. C. T. et al. Protein-ligand binding with the coarse-grained Martini model. Nat. Commun. 11, 3714 (2020).
- 6. Heidenreich, M. et al. Designer protein assemblies with tunable phase diagrams in living cells. Nat. Chem. Biol. **16**, 939–945 (2020).
- 7. Vögele, M., Köfinger, J. & Hummer, G. Hydrodynamics of Diffusion in Lipid Membrane Simulations. Phys. Rev. Lett. 120, 268104 (2018).
- 8. D'Agostino, M., Risselada, H. J., Lürick, A., Ungermann, C. & Mayer, A. A tethering complex drives the terminal stage of SNARE-dependent membrane fusion. *Nature* **551**, 634–638 (2017).
- Bochicchio, D., Salvalaglio, M. & Pavan, G. M. Into the Dynamics of a Supramolecular Polymer at Submolecular Resolution. *Nat. Commun.* **8**, 1–11 (2017).
- 10. Sarkar, A. et al. Self-Sorted, Random, and Block Supramolecular Copolymers via Sequence Controlled, Multicomponent Self-Assembly. J. Am. Chem. Soc. 142, 7606–7617 (2020).
- 11. Datta, S. et al. Self-assembled poly-catenanes from supramolecular toroidal building blocks. Nature 583, 400-405 (2020).

PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0079044

- 12. Bottaro, S. & Lindorff-Larsen, K. Biophysical experiments and biomolecular simulations: A perfect match? Science **361**, 355–360 (2018).
- 13. Lyubartsey, A. P. & Laaksonen, A. Calculation of effective interaction potentials from radial distribution functions: A reverse Monte Carlo approach. Phys. Rev. E Stat. Phys. Plasmas Fluids Relat. Interdiscip. Top. 52, 3730-3737 (1995).
- 14. Tschöp, W., Kremer, K., Batoulis, J., Bürger, T. & Hahn, O. Simulation of polymer melts. I. Coarse-graining procedure for polycarbonates. Acta Polym. 49, 61-74 (1998).
- 15. Reith, D., Pütz, M. & Müller-Plathe, F. Deriving effective mesoscale potentials from atomistic simulations. J. Comput. Chem. 24, 1624–1636 (2003).
- 16. Izvekov, S., Parrinello, M., Burnham, C. J. & Voth, G. A. Effective force fields for condensed phase systems from ab initio molecular dynamics simulation: A new method for force-matching. J. Chem. Phys. 120, 10896-10913 (2004).
- 17. Izvekov, S. & Voth, G. A. A Multiscale Coarse-Graining Method for Biomolecular Systems. J. Phys. Chem. B 109, 2469-2473 (2005).
- 18. Noid, W. G. et al. The multiscale coarse-graining method. I. A rigorous bridge between atomistic and coarsegrained models. J. Chem. Phys. 128, 244114 (2008).
- 19. Moore, T. C., lacovella, C. R. & McCabe, C. Derivation of coarse-grained potentials via multistate iterative Boltzmann inversion. J. Chem. Phys. 140, (2014).
- 20. Rudzinski, J. F. & Noid, W. G. A generalized-Yvon-Born-Green method for coarse-grained modeling. Eur. Phys. J. Spec. Top. **224**, 2193–2216 (2015).
- 21. Bejagam, K. K., Singh, S., An, Y. & Deshmukh, S. A. Machine-Learned Coarse-Grained Models. J. Phys. Chem. Lett. 9, 4667-4672 (2018).
- 22. Bejagam, K. K., Singh, S., An, Y., Berry, C. & Deshmukh, S. A. PSO-Assisted Development of New Transferable Coarse-Grained Water Models. J. Phys. Chem. B 122, 1958–1971 (2018).
- 23. Empereur-Mot, C. et al. Swarm-CG: Automatic Parametrization of Bonded Terms in MARTINI-Based Coarse-Grained Models of Simple to Complex Molecules via Fuzzy Self-Tuning Particle Swarm Optimization. ACS Omega **5**, 32823–32843 (2020).
- 24. Souza, P. C. T. et al. Martini 3: a general purpose force field for coarse-grained molecular dynamics. Nat. Methods 18, 382-388 (2021).

PLEASE CITE THIS ARTICLE AS DOI:10.1063/5.0079044

- 25. Wang, J. et al. Machine Learning of Coarse-Grained Molecular Dynamics Force Fields. ACS Cent. Sci. 5, 755–767 (2019).
- 26. Husic, B. E. et al. Coarse graining molecular dynamics with graph neural networks. J. Chem. Phys. 153, 194101 (2020).
- 27. Jarin, Z., Newhouse, J. & Voth, G. A. Coarse-Grained Force Fields from the Perspective of Statistical Mechanics: Better Understanding of the Origins of a MARTINI Hangover. J. Chem. Theory Comput. 17, 1170–1180 (2021).
- 28. Noid, W. G. et al. The multiscale coarse-graining method. II. Numerical implementation for coarse-grained molecular models. J. Chem. Phys. 128, 244115 (2008).
- 29. Wang, L.-P., Martinez, T. J. & Pande, V. S. Building Force Fields: An Automatic, Systematic, and Reproducible Approach. J. Phys. Chem. Lett. 5, 1885–1891 (2014).
- 30. Wang, L.-P., Chen, J. & Van Voorhis, T. Systematic Parametrization of Polarizable Force Fields from Quantum Chemistry Data. J. Chem. Theory Comput. 9, 452–460 (2013).
- 31. Shell, M. S. The relative entropy is fundamental to multiscale and inverse thermodynamic problems. J. Chem. Phys. 129, 144108 (2008).
- 32. Mullinax, J. W. & Noid, W. G. A Generalized-Yvon-Born-Green Theory for Determining Coarse-Grained Interaction Potentials. J. Phys. Chem. C 114, 5661–5674 (2010).
- 33. Conway, O., An, Y., Bejagam, K. K. & Deshmukh, S. A. Development of transferable coarse-grained models of amino acids. Mol. Syst. Des. Eng. 5, 675–685 (2020).
- 34. Yang, S., Cui, Z. & Qu, J. A Coarse-Grained Model for Epoxy Molding Compound. J. Phys. Chem. B 118, 1660-1669 (2014).
- 35. Hofmann, F. D., Devereux, M., Pfaltz, A. & Meuwly, M. Toward force fields for atomistic simulations of iridiumcontaining complexes. J. Comput. Chem. 35, 18–29 (2014).
- 36. Huang, L. & Roux, B. Automated Force Field Parameterization for Nonpolarizable and Polarizable Atomic Models Based on Ab Initio Target Data. J. Chem. Theory Comput. 9, 3543–3556 (2013).
- 37. Rudzinski, J. F. & Noid, W. G. Bottom-Up Coarse-Graining of Peptide Ensembles and Helix-Coil Transitions. J. Chem. Theory Comput. **11**, 1278–1291 (2015).
- 38. Bayramoglu, B. & Faller, R. Coarse-Grained Modeling of Polystyrene in Various Environments by Iterative Boltzmann Inversion. *Macromolecules* **45**, 9205–9219 (2012).

PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0079044

- 39. Bayramoglu, B. & Faller, R. Modeling of Polystyrene under Confinement: Exploring the Limits of Iterative Boltzmann Inversion. *Macromolecules* **46**, 7957–7976 (2013).
- 40. Ganguly, P. & van der Vegt, N. F. A. Representability and Transferability of Kirkwood-Buff Iterative Boltzmann Inversion Models for Multicomponent Aqueous Systems. J. Chem. Theory Comput. 9, 5247–5256 (2013).
- 41. Banerjee, P., Roy, S. & Nair, N. Coarse-Grained Molecular Dynamics Force-Field for Polyacrylamide in Infinite Dilution Derived from Iterative Boltzmann Inversion and MARTINI Force-Field. J. Phys. Chem. B 122, 1516–1524 (2018).
- 42. Prasitnok, K. & Wilson, M. R. A coarse-grained model for polyethylene glycol in bulk water and at a water/air interface. Phys. Chem. Chem. Phys. 15, 17093 (2013).
- 43. Caceres-Delpiano, J., Wang, L.-P. & Essex, J. W. The automated optimisation of a coarse-grained force field using free energy data. Phys. Chem. Chem. Phys. 23, 24842-24851 (2021).
- 44. Izvekov, S. & Voth, G. A. Multiscale Coarse-Graining of Mixed Phospholipid/Cholesterol Bilayers. J. Chem. Theory Comput. 2, 637-648 (2006).
- 45. Kmiecik, S. et al. Coarse-Grained Protein Models and Their Applications. Chem. Rev. 116, 7898–7936 (2016).
- 46. An, Y., Bejagam, K. K. & Deshmukh, S. A. Development of New Transferable Coarse-Grained Models of Hydrocarbons. J. Phys. Chem. B 122, 7143–7153 (2018).
- 47. An, Y., Singh, S., Bejagam, K. K. & Deshmukh, S. A. Development of an Accurate Coarse-Grained Model of Poly(acrylic acid) in Explicit Solvents. *Macromolecules* **52**, 4875–4887 (2019).
- 48. Kennedy, J. & Eberhart, R. Particle swarm optimization. in *Proceedings of ICNN'95 International Conference on* Neural Networks vol. 4 1942–1948 (1995).
- 49. Nobile, M. S. et al. Fuzzy Self-Tuning PSO: A settings-free algorithm for global optimization. Swarm Evol. Comput. 39, 70-85 (2018).
- 50. Villani, C. Optimal Transport: Old and New. (Springer-Verlag, 2009).
- 51. Lu, C. et al. OPLS4: Improving Force Field Accuracy on Challenging Regimes of Chemical Space. J. Chem. Theory Comput. 17, 4291–4300 (2021).
- 52. Law, M. M. & Hutson, J. M. I-NoLLS: A program for interactive nonlinear least-squares fitting of the parameters of physical models. Comput. Phys. Commun. 102, 252–268 (1997).
- 53. Marguardt, D. et al. The structures of polyunsaturated lipid bilayers by joint refinement of neutron and X-ray scattering data. Chem. Phys. Lipids 229, 104892 (2020).

PLEASE CITE THIS ARTICLE AS DOI: 10.1063/5.0079044

- 54. Kučerka, N., Nieh, M.-P. & Katsaras, J. Fluid phase lipid areas and bilayer thicknesses of commonly used phosphatidylcholines as a function of temperature. Biochim. Biophys. Acta BBA - Biomembr. 1808, 2761–2771 (2011).
- 55. Sengupta, S., Basak, S. & Peters, R. A. Particle Swarm Optimization: A Survey of Historical and Recent Developments with Hybridization Perspectives. Mach. Learn. Knowl. Extr. 1, 157-191 (2019).
- 56. Michaud-Agrawal, N., Denning, E. J., Woolf, T. B. & Beckstein, O. MDAnalysis: A toolkit for the analysis of molecular dynamics simulations. J. Comput. Chem. 32, 2319-2327 (2011).
- 57. Gowers, R. J. et al. MDAnalysis: A Python Package for the Rapid Analysis of Molecular Dynamics Simulations. *Proc.* 15th Python Sci. Conf. 98–105 (2016).
- 58. Pele, O. & Werman, M. Fast and robust Earth Mover's Distances. in 2009 IEEE 12th International Conference on Computer Vision 460-467 (2009).
- 59. Rubner, Y., Tomasi, C. & Guibas, L. J. The Earth Mover's Distance as a Metric for Image Retrieval. Int. J. Comput. Vis. 40, 99-121 (2000).
- 60. Kirkwood, J. G. & Buff, F. P. The Statistical Mechanical Theory of Solutions. I. J. Chem. Phys. 19, 774–777 (1951).
- 61. Pele, O. & Werman, M. A Linear Time Histogram Metric for Improved SIFT Matching. in Computer Vision ECCV 2008 (eds. Forsyth, D., Torr, P. & Zisserman, A.) 495–508 (Springer, 2008).
- 62. Lorentz, H. A. Ueber die Anwendung des Satzes vom Virial in der kinetischen Theorie der Gase. Ann. Phys. 248, 127-136 (1881).
- 63. Marrink, S. J. & Mark, A. E. The Mechanism of Vesicle Fusion as Revealed by Molecular Dynamics Simulations. J. Am. Chem. Soc. 125, 11144-11145 (2003).
- 64. Smeijers, A. F., Markvoort, A. J., Pieterse, K. & Hilbers, P. A. J. A Detailed Look at Vesicle Fusion. J. Phys. Chem. B **110**, 13212–13219 (2006).
- 65. Carpenter, T. S. et al. Capturing Phase Behavior of Ternary Lipid Mixtures with a Refined Martini Coarse-Grained Force Field. J. Chem. Theory Comput. 14, 6050-6062 (2018).



