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Grid-size reduction in flow calculations on infinite domains by higher-order far-field asymptotics in numerical boundary conditions *

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Summary

An error analysis is presented of the numerical calculation of the steady flow on an infinite domain around a given airfoil by a domain-splitting (= zonal) method. This method combines a fully-conservative finite-difference approximation on a finite domain around the airfoil with an approximate asymptotic solution outside this finite domain. The errors are analyzed as a function of the accuracy of the approximate asymptotic expansion, of the distance to the airfoil of the far-field boundary of the finite domain and of the mesh size. The numerical experiments show that, for a given desired accuracy level, large reduction in grid sizes are possible, if the usual far-field asymptotic approximation (uniform flow plus first-order perturbation by a circulation vortex at infinity) is augmented by only a few extra terms in the approximate asymptotic far-field solution. In this way, considerable numerical efficiency improvements can be realized. It is expected that this conclusion can be generalized to many other applications where computations on infinite domains are performed.

1. Introduction

In computational fluid dynamics problems on infinite domains, where the far-field flow is a small perturbation of an inviscid uniform parallel flow, the infinite size of the flow domain can be handled by what nowadays is called domain-splitting [2] or zonal method [3]. The numerical calculations are then restricted to a finite inner domain, while outside this inner domain a simple model of the flow is used. Examples are aerodynamic calculations of steady compressible subsonic or transonic flows around aircraft or airfoils by finite-difference methods. The finite-difference calculations are restricted to a large but finite "inner" computational domain, while on the corresponding "outer" domain the flow is assumed to be a steady uniform parallel flow, or a small perturbation of such a flow in which nonlinear compressibility effects and aircraft or airfoil thickness effects are neglected. In order to obtain a well-posed finite-difference problem on the finite inner domain, a suitable numerical boundary condition on the far-field boundary can be formulated from this assumption.

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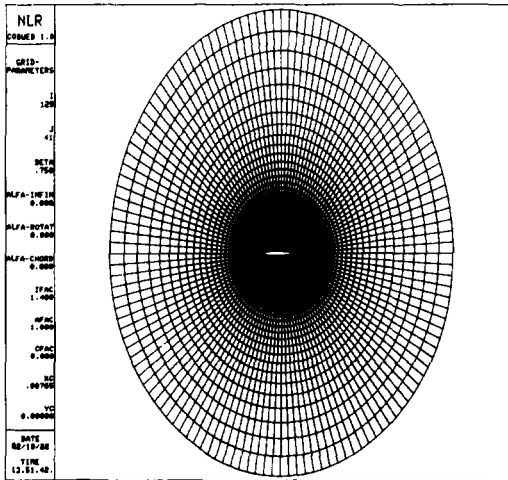


Figure 1a. Far-field part of 130x42 grid around NACA 0012 airfoil.

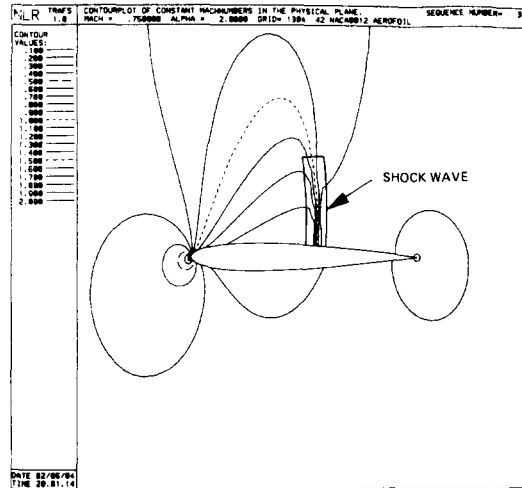


Figure 1b. Lines of constant Mach numbers in computed transonic flow around NACA 0012 airfoil ($M_\infty = 0.75$, $\alpha = 2^\circ$).

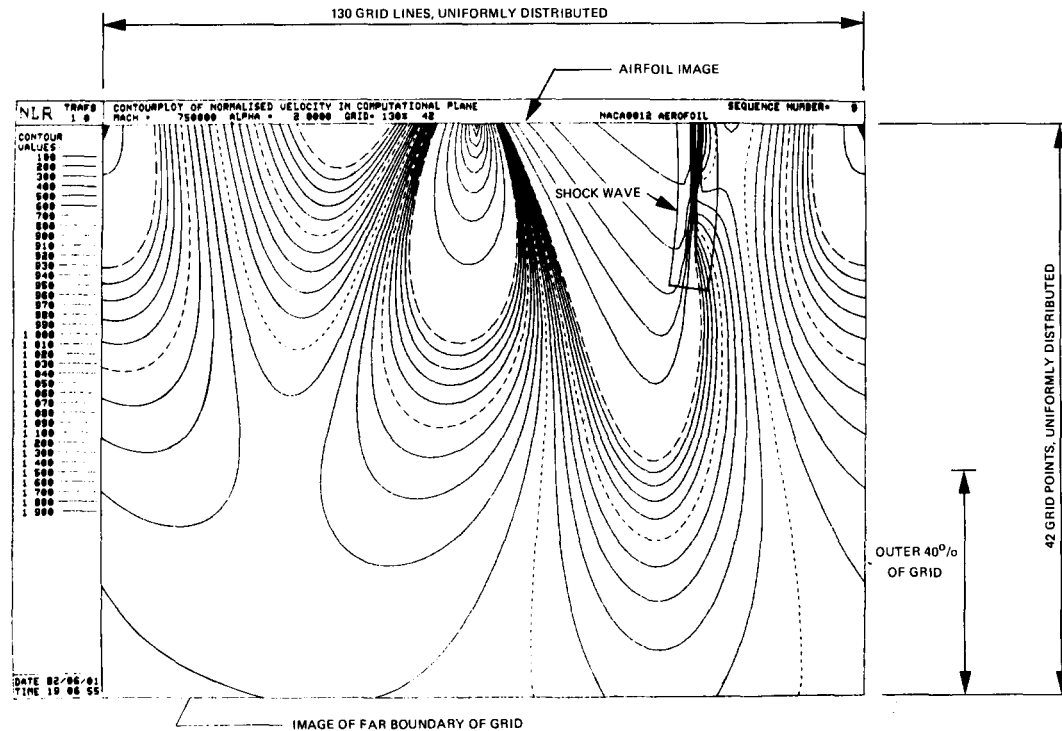


Figure 1c. Lines of constant speed on uniform computational grid in computed transonic flow around NACA 0012 airfoil ($M_\infty = 0.75$, $\alpha = 2^\circ$).

There are various reasons to investigate whether this technique can be improved by using more accurate solutions of the flow equations on the infinite domain.

- Recently it has become clear that, on large parts of grids in calculations of transonic flows around airfoils, the velocity is a smooth function of the grid coordinates. A typical example from current engineering practice is presented in Fig. 1 [4]. On about 40% of the grid near the far-field boundary, the lines of constant speed show a very smooth and regular pattern around the free-stream speed value 1.0, see Fig. 1.c. It may be expected that such velocity fields can be approximated more efficiently by a small number of appropriately chosen asymptotic solutions of the flow equations at infinity, instead of by a grid function.

- In time-dependent flow calculations, recently so-called non-reflecting or weakly reflecting numerical boundary conditions have been proposed and applied with great success to prevent nonphysical numerical reflections of outgoing waves at the far-field boundary of inner domains (see e.g. [5–7]). The boundary conditions, which are generalizations of well-known radiation conditions, have the form of differential expressions, allowing only physically-relevant outgoing waves. Similarly, in steady-flow problems differential expressions can be used, which allow a more accurate asymptotic expansion as its solution.

- A third reason concerns the calculation of drag. Detailed analysis of the calculation of transonic flows around airfoils [4] has shown that reliable calculation of drag is difficult because of a too crude approximation of the flow on the infinite outer domain, causing effects which are known to aerodynamicists as wind-tunnel blockage effects. This situation may be remedied by more accurate solutions on the outer domain.

Summarizing, application of higher-order asymptotic theory seems to promise improved numerical approximation efficiency. It should allow the replacement of the approximation of the flow on the outer 40% of a grid by a combination of a small number (5 to 10) of asymptotic solutions on the other domain, leading to grid-size reduction. This grid-size reduction should then give substantial savings in computer runtime, because the total number of unknowns in the computational problem is also reduced by say 40%. Such savings help to bring down the calculation costs of three-dimensional flows around aircraft down to a level where repeated calculations for design purposes become feasible. Moreover, the accuracy of the calculation of drag should improve considerably.

The present study has been undertaken to explore the first of these speculations: the possible grid-size reduction, as a function of the accuracy of far-field solutions, for a fixed error level in the overall numerical simulation. Incompressible flow around lifting airfoils was taken as a test problem. The nonlinearity in computational aerodynamics (due to compressibility) is only a minor additional complication, from the point of view of approximation accuracy, it only complicates the algorithm, and has thus been removed from the model problem.

The study was thus primarily meant to resolve accuracy and approximation questions concerning converged steady flows, and not to design fast and robust solution algorithms. However, in the concluding discussion it will be argued that efficient solution procedure for steady-flow problems may be expected in combination with multigrid techniques.

The material of the study is presented as follows.

- Section 2: Test-problem formulation for incompressible flow around lifting airfoil.
- Section 3: Discretization of test problem by domain splitting, fully-conservative finite differencing on the inner domain, and construction of asymptotic expansion on the infinite outer domain. Formulation of numerical boundary conditions, and of techniques to determine or eliminate the unknown coefficients in the asymptotic expansion. Outline of solution process.

- Section 4: Numerical experiments, test strategy, dependence of error level on domain size, number of terms in asymptotic expansion, and mesh size.
- Section 5: Conclusions, and discussion of generalizations.

2. Description of test problem

The test problem is an inviscid, irrotational, and incompressible flow around a lifting airfoil. Introducing a velocity potential, φ , the mass conservation in the flow is governed by Laplace's equation:

$$\Delta\varphi = 0. \quad (1)$$

At the airfoil, the mass-flux is zero:

$$\frac{\partial\varphi}{\partial n} = 0. \quad (2)$$

At infinity, a uniform parallel flow is assumed, and from the general solution of (1) we then may derive the following expansion for φ at infinity:

$$\varphi = x + \frac{\Gamma\theta}{2\pi} + O(r^{-1}). \quad (3)$$

where Γ , the circulation, is an unknown scalar, $\theta = \arctan(y/x)$, and $r^2 = x^2 + y^2$. The angle θ is not single-valued, and therefore a slit has been made along the positive x -axis, where φ makes a jump of magnitude Γ . The value of Γ should be chosen such that the velocity remains bounded, especially at the trailing edge of the airfoil (Kutta condition). It can be shown that the above problem is well-posed and has a unique solution for φ and Γ .

3. Numerical discretization and solution process

The problem described in Section 2 has been solved using a domain-splitting (zonal) method. For the discretization, a finite-difference approach has been used. The resulting set of equations has been solved iteratively by successive-line relaxation, accelerated by Richardson extrapolation.

3.1. Domain-splitting

In domain-splitting methods, the physical domain is split in several coupled subdomains, in which different physical models and/or numerical solution methods are used. In the test problem, a simple split into two domains is introduced (fig. 2); a bounded inner domain A, where a fully-conservative finite-difference method will be applied, and an unbounded outer domain B, where an analytic approximation of the flow will be used. In both domains, Laplace's equation must be satisfied, and the solutions on the two domains must be appropriately coupled at the boundary C – which is assumed to be regular – so that a unique solution is obtained on the entire domain. It can be shown that the continuity of φ and φ_n over C is sufficient for continuity of all higher derivatives of φ over

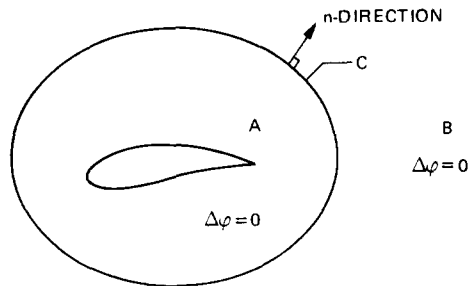


Figure 2. Splitting into two domains A and B.

C. Three discrete methods of coupling have been investigated; these will be identified as the φ , φ_n -method, the φ , φ_n -method, and the DE-(differential equation) method.

3.1.1. φ , φ_n - and the φ , φ_n -method

On the outer domain B, the solution of (1) will be approximated by a part of the general solution,

$$\varphi = x + \frac{\Gamma\theta}{2\pi} + \sum_{n=1}^N (a_n \cos n\theta + b_n \sin n\theta)r^{-n}. \tag{4}$$

Observe that for $N < \infty$ this is a finite linear combination of particular solutions of Laplace's equation, with a_n, b_n free constants to be suitably determined. On the inner domain a finite-difference approximation will be constructed (Section 3.3).

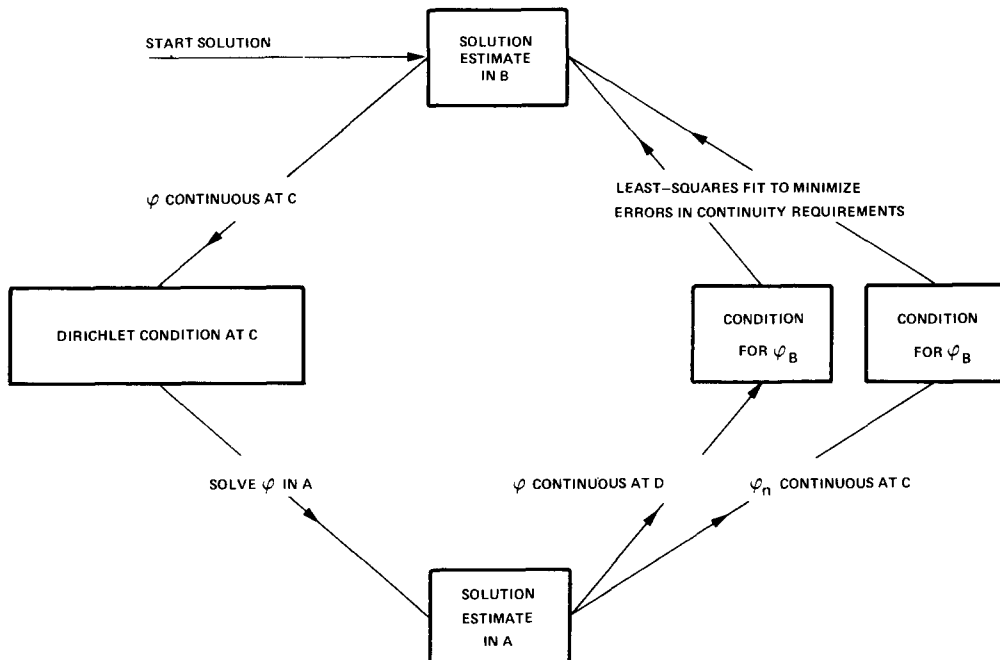


Figure 3. Summary of φ , φ_n - and φ , φ_n -method.

Continuity of φ and φ_n over C can numerically be assured by taking N large enough in (4), i.e. equal to the number of mesh points of the finite-difference approximation on the boundary C . In practice, a certain discretization error in the continuity requirements is acceptable. This freedom will be used to limit the number of terms N in (4) as much as possible. The continuity requirement for φ or φ_n was chosen to be imposed in a least-squares sense, giving the algorithm of Fig. 3.

The algorithm works as follows: starting with an approximate solution on the outer domain B , the solution is restricted to the boundary C , giving a Dirichlet condition for the inner domain A , where the problem is solved by a finite-difference method. In the φ , φ_n -method, a normal derivative at curve C is now computed from the infinite-difference solution; in the φ , φ -method a potential φ is computed on a curve D , slightly different from C . This finite-difference representation of φ or φ_n is fitted to the φ or φ_n on the outer domain B by a least-squares fit, giving a new approximate solution in B . The φ or φ_n will be continuous over C only in a least-squares sense.

3.1.2. DE-method

In the DE (differential equation) method, the series expansion (4) is replaced by a differential equation, applied along the curve C , and having a number of terms of the expansion as its solution. The expansion (4) may be slightly generalized to

$$\varphi = x + \frac{\Gamma\theta}{2\pi} + \sum_{n=1}^{\infty} f_n(\theta)r^{-n}. \quad (5)$$

It can be shown [1] that the DE

$$B\varphi = \left(\prod_{n=1}^N B_n \right) \varphi = 0, \quad B_n = \partial/\partial r + (2n-1)r^{-1}, \quad (6)$$

has the first N terms of the sum of (5) as its solution. Therefore, if the DE

$$B\varphi = B\left(x + \frac{\Gamma\theta}{2\pi}\right) \quad (7)$$

is applied to the finite-difference approximation along the curve C , the DE effectively imposes a boundary condition, which has (5) up to term N as its solution.

The DE-method is, in contrast to φ , φ - or φ , φ_n -methods, implicit, i.e. it does not need explicit evaluations of coefficients a_n , b_n in (4), and specification of the form of the functions $f_n(\theta)$ in (5). However, it is impractical to handle many terms of (5), as the DE will then be of a high order, and is hard to discretize. The φ , φ - and φ , φ_n -methods have more flexibility in this respect.

3.2. Application of φ , φ - and φ , φ_n -method to the flow around a circle

To obtain some insight in the convergence of the described iteration algorithm of Fig. 3, the application of this algorithm to the calculation of the steady flow around a circle has been analyzed. The analytic solution is known:

$$\varphi = r \cos \theta + \frac{r_1^2}{r} \cos \theta, \quad (8)$$

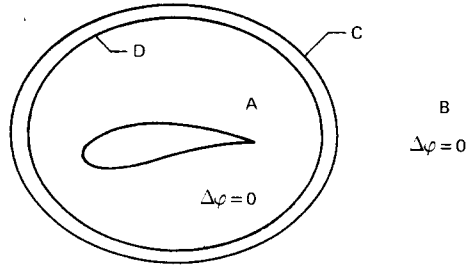


Figure 4. The curves C and D, used to couple the solutions in A and B.

where r_1 is the radius of the circle. This circle-airfoil and the curves C and D (Fig. 4) are assumed to be concentric circles with radius r_1, r_2, r_3 , respectively ($r_1 \ll r_3 < r_2$).

The solution of Laplace's equation in the inner domain A may be given the general form:

$$\varphi = r \cos \theta + \frac{r_1^2}{r} \cos \theta + \sum_{n=-\infty}^{\infty} (a_n \cos n\theta + b_n \sin n\theta) r^n. \quad (9)$$

This approximation has the form of the exact solution plus a perturbation which should be shown to tend to zero. Using the boundary condition $\partial\varphi/\partial r = 0$ at the circle, this approximate solution becomes

$$\varphi = r \cos \theta + \frac{r_1^2}{r} \cos \theta + \sum_{n=0}^{\infty} (a_n \cos n\theta + b_n \sin n\theta) \left(r^n + \frac{r_1^{2n}}{r^n} \right). \quad (10)$$

This solution must be coupled with an approximate solution in domain B, which may be written as

$$\varphi = r \cos \theta + \frac{r_1^2}{r} \cos \theta + \sum_{n=1}^N (\tilde{a}_n \cos n\theta + \tilde{b}_n \sin n\theta) r^{-n}. \quad (11)$$

Here, we have omitted the particular solutions with $n < 0$, because $\nabla\varphi$ must be bounded for $r \rightarrow \infty$ (free-stream boundary condition).

When the continuity requirement in φ is applied on the curve C, the following relations between coefficients on the inner domain and on the outer domain are obtained:

$$a_0^k = 0, \quad b_0^k = 0, \quad (12.a)$$

$$a_n^k \left(r_2^n + \frac{r_1^{2n}}{r_2^n} \right) = \tilde{a}_n^k r_2^{-n}, \quad n = 1, \dots, N, \quad (12.b)$$

$$a_n = b_n = 0 \quad \text{for } n > N; \quad (12.c)$$

(12.b) holds also between b_n^k and \tilde{b}_n^k . The superscript k has been appended to indicate that the coefficients are those of the k -th iteration of the iteration algorithm of Fig. 3.

From Fig. 3 it is seen that in the φ, φ_n -method the outer solution is further coupled to the inner one by the demand that φ_n is continuous over C. This gives the equation

$$\begin{aligned} \cos \theta - \frac{r_1^2}{r_2^2} \cos \theta + \sum_{n=1}^N (a_n^k \cos n\theta + b_n^k \sin n\theta) \left(nr_2^{n-1} - n \frac{r_1^{2n}}{r_2^{n+1}} \right) \\ = \cos \theta - \frac{r_1^2}{r_2^2} \cos \theta + \sum_{n=1}^N (\tilde{a}_n^{k+1} \cos n\theta + \tilde{b}_n^{k+1} \sin n\theta) (-nr_2^{n-1}), \end{aligned} \quad (13)$$

where a_n^k and b_n^k are defined by (12.b). This gives new relations between the coefficients of inner and outer domain:

$$a_n^k \left(r_2^{n-1} - \frac{r_1^{2n}}{r_2^{n+1}} \right) = -\tilde{a}_n^{k+1} r_1^{n-1}, \quad n = 1, \dots, N, \quad (14)$$

with the same formula for b_n^k . We are interested in reduction factors of the coefficients after one iteration cycle of the algorithm of Fig. 3. In the case of the φ, φ_n -method, substitution of (12.b) in (14) gives the reduction factors

$$\frac{\tilde{a}_n^{k+1}}{\tilde{a}_n^k} = - \left(\frac{r_2^{2n} - r_1^{2n}}{r_2^{2n} + r_1^{2n}} \right), \quad n = 1, \dots, N. \quad (15)$$

This formula shows, that the coefficients tend to zero (as desired) in an oscillatory manner, slowly for $r_1 \ll r_2$. This process can be accelerated by underrelaxation ($a_n^{k+1} = a_n^k + \sigma(a_n^{k+1} - a_n^k)$). In a straight-forward manner, the relaxation factors

$$\sigma_n = \frac{1}{2} + \frac{r_1^{2n}}{2r_2^{2n}}, \quad n = 1, \dots, N \quad (16)$$

are found. So $\sigma = \frac{1}{2} + r_1^2/(2r_2^2)$ will provide a good relaxation factor for all coefficients, when the φ, φ_n -method is applied.

The φ, φ -method differs from the φ, φ_n -method in the computation of the outer solution from the inner solution. This time, φ should be made continuous on the circle D. This means that the additional continuity relation is resembling equation (12.b), but it should now hold on circle D with radius r_3 , i.e.

$$a_n^k \left(r_3^n + \frac{r_1^{2n}}{r_3^{2n}} \right) = \tilde{a}_n^{k+1} r_3^{-n}, \quad n = 1, \dots, N. \quad (17)$$

When this is combined with (12), the reduction factors

$$\frac{\tilde{a}_n^{k+1}}{\tilde{a}_n^k} = \frac{r_3^{2n} + r_1^{2n}}{r_2^{2n} + r_1^{2n}}, \quad n = 1, \dots, N \quad (18)$$

are found. This equation shows that the lower-order terms converge slowly. Acceleration is again possible by relaxation, here overrelaxation. The relaxation factor is

$$\sigma_n = \frac{r_2^{2n} + r_1^{2n}}{r_2^{2n} - r_3^{2n}}, \quad n = 1, \dots, N. \quad (19)$$

An overrelaxation factor for all n may be chosen to be

$$\sigma = \frac{r_2^2 + r_1^2}{r_2^2 - r_3^2}. \quad (20)$$

However, this factor should not be so large that it causes higher-order terms to diverge. This is assured if

$$\sigma < 2. \quad (21)$$

This restriction results in a slower convergence rate for the φ , φ_n -method compared with the φ , φ_n -method.

In the case of airfoil calculations, it was found that under- or overrelaxation improved the convergence rate of the iteration processes significantly. The convergence analysis of the section thus seems to be indicative for what has to be done in more general situations.

3.3. Finite-difference approximation

The equation (1) has been discretized on an 0-type grid (Fig. 1a). In this grid, potentials are given at midpoints of cells, and velocities at midpoints of cell boundaries. Using tensor-calculus, a fully-conservative finite-difference method has been formulated [1]. At the airfoil, virtual points are introduced to achieve second-order accurate approximations of the boundary conditions.

The transformation of an 0-type grid to an equidistant computational grid in a computational rectangle was based on an approximately conformal mapping of an airfoil to a circle producing near-orthogonal grids. The mapping is singular at the trailing edge. It may be shown that this implies that the transformed velocity in the computational rectangle then must be zero at the trailing edge to guarantee finite velocities at the trailing edge. Thus, the Kutta condition has been transformed from a qualitative to a quantitative condition. In the computations, the zero transformed velocity at the trailing edge is accomplished by means of a Newton iteration on the circulation.

3.4. Solution method

The solution method of the finite-difference equations was based on successive line relaxation (SLOR), because for this method the theory is well-developed. The SLOR iteration matrix was derived in the usual way directly from the system matrix of the problem. This gave iteration matrices that were diagonally dominant, except possibly at the boundaries. Here, diagonal dominance is often lost due to central differencing in the discrete boundary conditions. As convergence of SLOR is usually realized by using diagonally-dominant iteration matrices, the iteration matrices were made diagonally

dominant at the boundaries, too, by making $O(h)$ modifications to the finite-difference boundary conditions for SLOR solution-corrections. These modifications are more drastic for DE-methods than for φ , φ_n - and φ , φ -methods – where always a Dirichlet condition at boundary C is used – and have slowed down the convergence of the DE-methods. For details about this non-standard issue, see [1].

The line relaxation has been accelerated by Richardson extrapolation, which will briefly explained here. After a sufficiently large number of SLOR iterations, the following relation approximately holds in practical calculations:

$$\Delta\varphi^{n+1} = \lambda \Delta\varphi^n, \quad \Delta\varphi^{n+1} = \varphi^{n+1} - \varphi^n, \quad (22)$$

where λ is the largest eigenvalue of the SLOR iteration matrix (usually slightly less than unity). If at certain stages of the iteration process, the above relation holds to a sufficient degree, an estimate of the ultimate $\Delta\varphi = \varphi^\infty - \varphi^n$ can be based on knowledge of the value of this eigenvalue. We assume the SLOR process to converge with fixed λ , so that

$$\Delta\varphi = \sum_{k=n+1}^{\infty} \Delta\varphi^k. \quad (23)$$

This relation can, together with (22), be used to define

$$\Delta\varphi = \frac{1}{1-\lambda} \Delta\varphi^{n+1}. \quad (24)$$

In our numerical calculation, this extrapolation on the largest eigenvalue has been applied when the variance of the last three eigenvalue estimates was below a certain given value. A sample convergence history is shown in Fig. 5. The vertical lines indicate the iterations where the extrapolation is applied. It is clear that the well-known slowdown of the convergence due to the largest eigenvalue of the iteration matrix is remedied by the extrapolation (see e.g. iterations 20, 52). If properly applied, the number of iterations can be reduced by more than a factor of two [8].

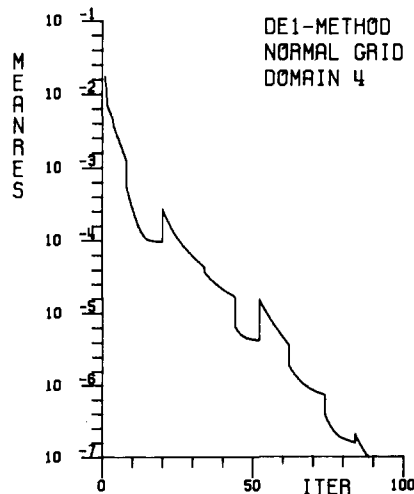


Figure 5. Convergence history of SLOR combined with Richardson extrapolation.

4. Numerical experiments

The objective of the present investigation is the relation between the number of N of terms in the asymptotic far-field solutions (4) or (5) used in the numerical boundary conditions, the size of the computational domain, and the mesh size in the finite-difference approximation. In this respect, we are interested in:

1. The error, when the size of the domain is varied (using a fixed number N of terms in the boundary condition).
2. The error, when the number of N terms of the expansion is varied (on a fixed domain).
3. The comparison of the errors for the DE-, φ , φ^* - and φ^* , φ_n^* -methods, both qualitatively and quantitatively.
4. The error as a function of the mesh size in the finite-difference calculation.

4.1. Test case

Computations have been done for the von Karman-Trefftz airfoil, for which the solution is analytically known. The results will be presented by means of plots of C_p -errors, which is the difference between the analytical and the numerical pressure coefficient on the airfoil surfaces. A comparison in terms of C_p -errors has been chosen as the pressure distribution on the airfoil is of great practical importance in aerodynamics, and errors tend to be largest on the airfoil surface. The angle of attack of the airfoil is 4 degrees. The C_p -distribution of this case (Fig. 6) has a form representative for engineering applications. The circulation Γ is a direct measure of the amount of lift on the airfoil.

Four differently-sized grid domains have been studied, and on each domain, seven different boundary conditions have been imposed at the outer artificial boundary. These conditions were formed from the basic three methods (φ , φ^- , φ , φ_n^- , and DE-method) by varying the number of terms N , giving $\varphi\varphi 0^-$, $\varphi\varphi 2^-$, $\varphi\varphi_n 2^-$, $\varphi\varphi 5^-$, $\varphi\varphi_n 5^-$, DE1-, DE2-methods, where the sequence number indicates how many terms of the expansion have been used.

The sizes of the grids, corresponding to the domains 1 through 4, are respectively 64×18 , 64×14 , 64×10 and 64×8 points. Some computations have been performed on domain 3 and 4 on fine grids, containing 128×18 and 128×14 points. to assess the influence of the mesh size. From preliminary experiments, it was concluded that these

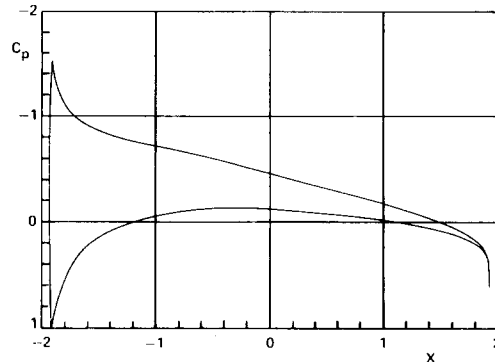


Figure 6. Analytical C_p -distribution. Von Karman-Trefftz airfoil, $\alpha = 4^\circ$, $\Gamma = 0.92918$.

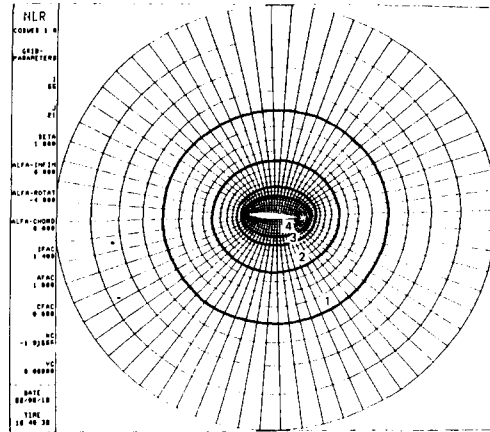


Figure 7a. Normal grid.

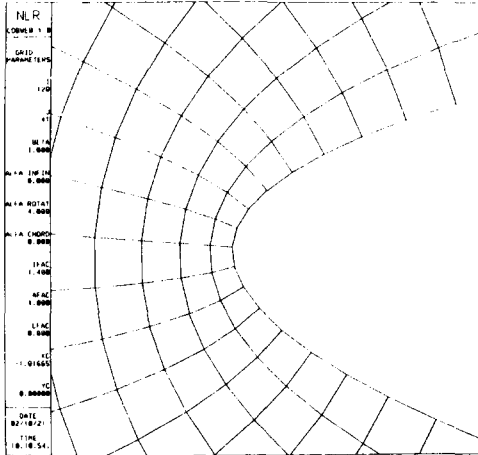
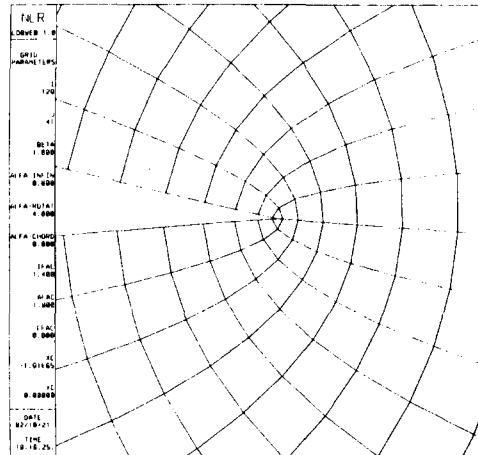


Figure 7b. Enlargement of fine grid at leading edge.



the upper side and ends at the trailing edge on the lower side. In all graphs, the values $CPE - CPC = \pm 0.02$ are indicated by two horizontal lines. Ideally, $CPE - CPC$ curves should be at least between these two horizontal lines, because computational results not

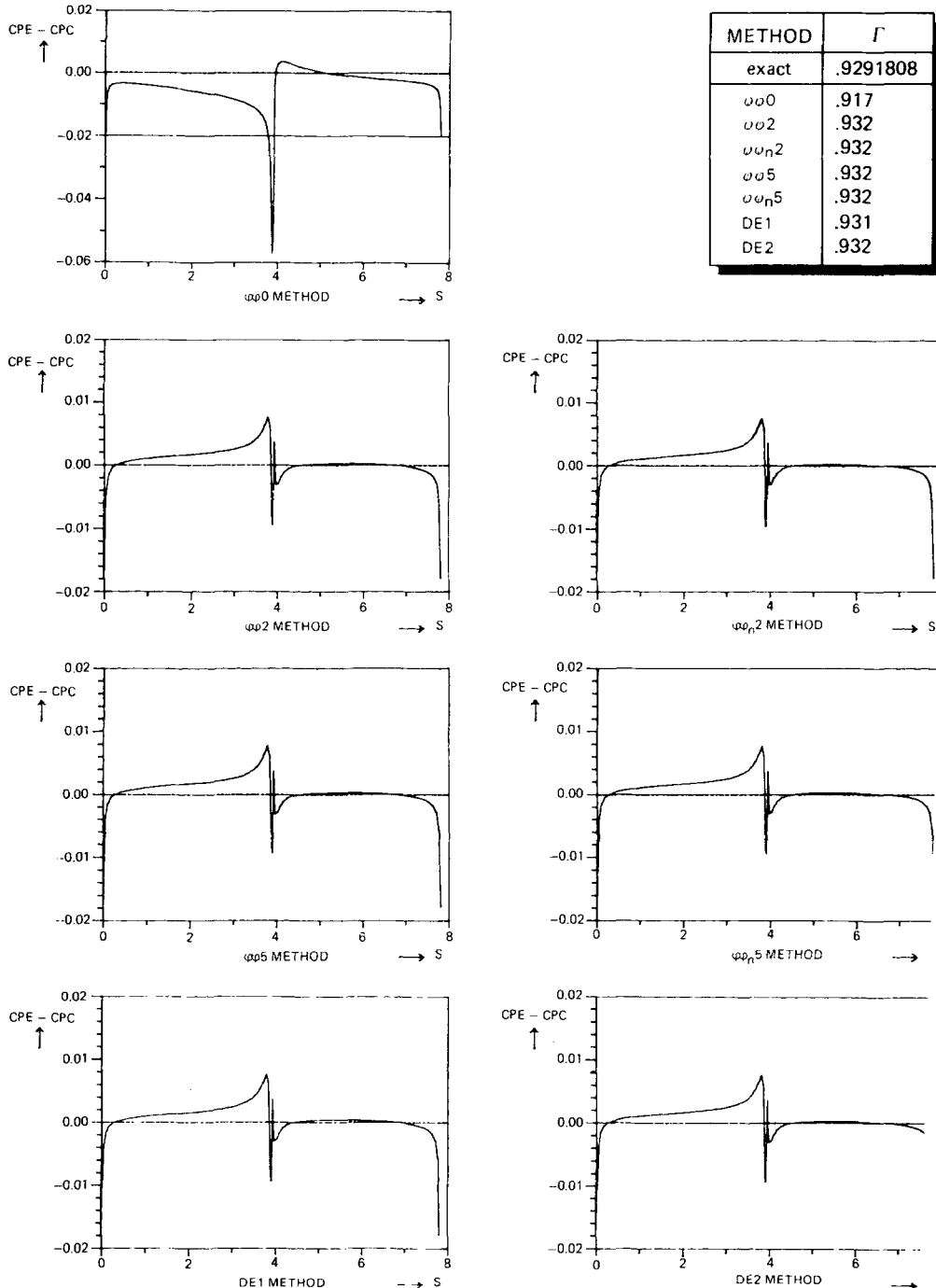
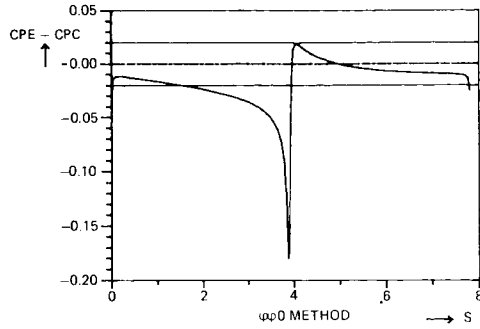


Figure 8a. C_p error results for: $\alpha = 4^\circ$, domain 1, $M = 65$, $N = 17$.



METHOD	Γ
exact	.92918
ϕ_0	.883
ϕ_2	.931
ϕ_{n2}	.930
ϕ_5	.932
ϕ_{n5}	.932
DE1	.928
DE2	.932

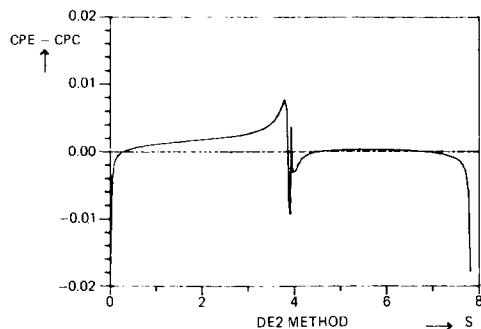
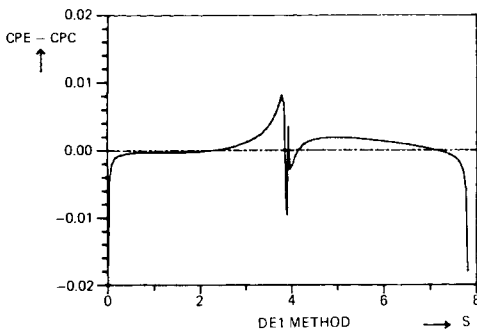
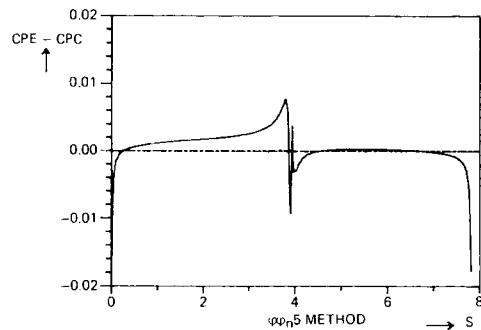
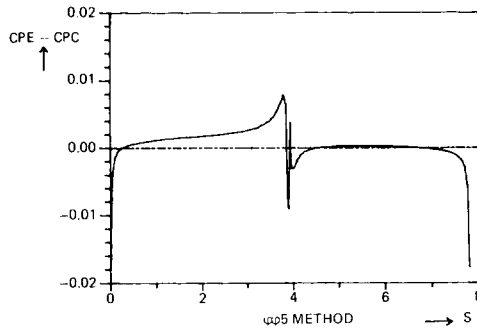
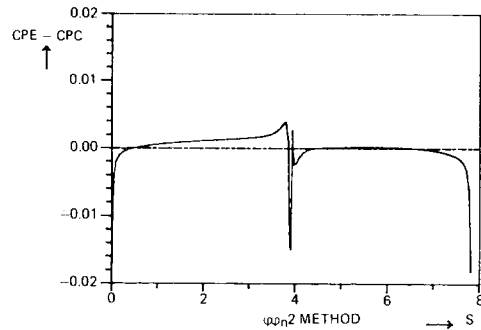
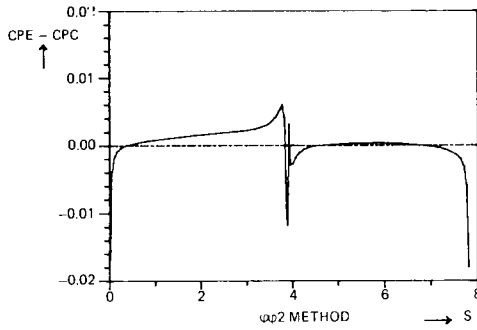


Figure 8b. C_p error results for: $\alpha = 4^\circ$, domain 2, $M = 65$, $N = 13$.

satisfying this criterion may be considered to be too inaccurate for engineering applications.

Figure 8 shows 28 plots of the pressure-coefficient error $CPE-CPC$ versus s , as a

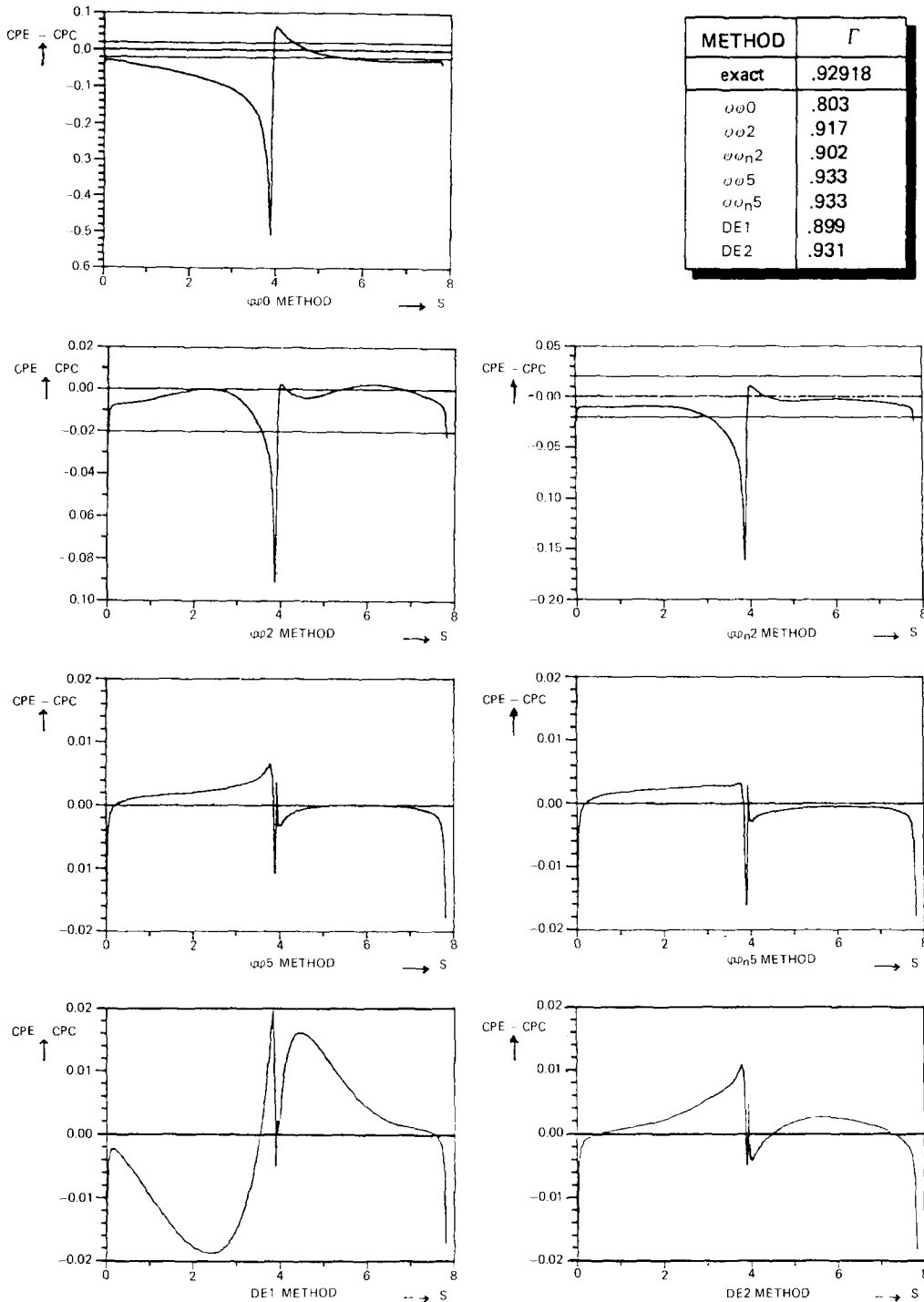


Figure 8c. C_p error results for: $\alpha = 4^\circ$, domain 3, $M = 65$, $N = 9$.

function of the four domain sizes and the seven far-field boundary-condition discretizations. Every figure also shows a table of calculated circulation values Γ . From all graphs it may be seen that the greatest errors occur at the trailing edge and at the leading edge of

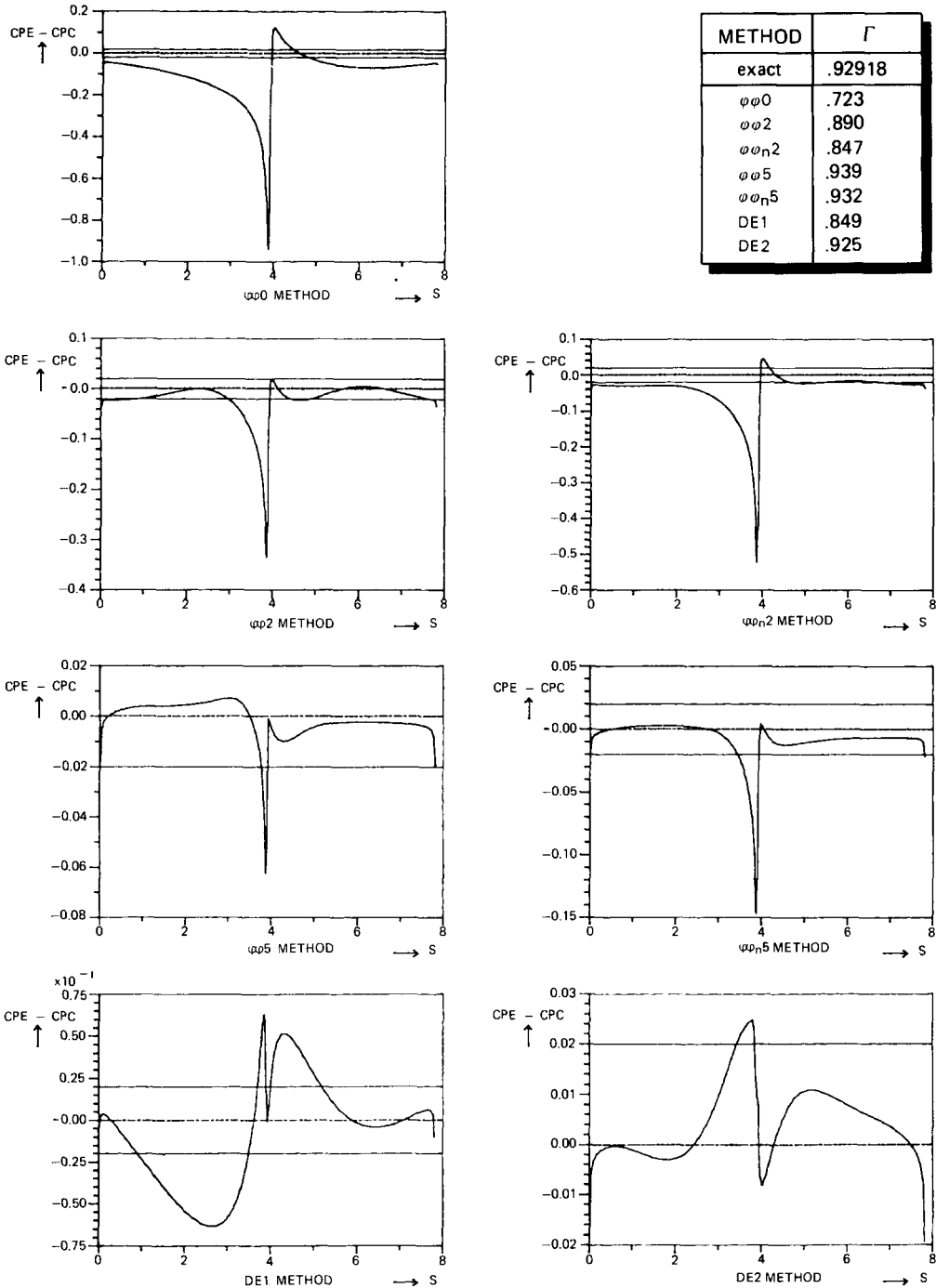


Figure 8d. C_p error results for: $\alpha = 4^\circ$, domain, 4, $M = 65$, $N = 7$.

the airfoil. The first error is a result of the singularity of the transformation at the trailing edge (see Section 3.2). This error is not of $O(h^2)$. The error at the leading edge is due to strong gradients. This error reduces well with grid refinement.

Figure 8a (artificial boundary at about $2\frac{1}{2}$ chords (= airfoil-lengths) distance, Fig. 7a) shows that the C_p -error can be reduced to an acceptable engineering level by only one term (DE1-method). Because all graphs, except $\varphi\varphi_0$, are identical we can conclude that the error in these results of Fig. 8a is mainly caused by the mesh width, and that second- and higher-order terms in the asymptotic expansion are not significant, when the boundary is this far away. Notice also that the calculated circulation values are about an order of magnitude more accurate than those of the standard $\varphi\varphi_0$ -method. Figure 8b (artificial boundary at about $1\frac{1}{2}$ chords distance) shows that the $\varphi\varphi_0$ -method results have become considerably worse (compare with $\varphi\varphi_0$ in Fig. 8a.) The other results are still excellent, and the error is an order of magnitude smaller than those of the $\varphi\varphi_0$ -method. Figure 8c (artificial boundary at about 1 chord distance) shows that also the accuracy of the $\varphi\varphi_2$ -, $\varphi\varphi_n$ 2- and DE1-method becomes considerably worse. Apparently, on a domain of this size, the higher-order terms in the asymptotic far-field expansion are going to play a role. Figure 8d (artificial boundary at about $3/4$ chord distance) shows, that only the results of the $\varphi\varphi_5$ -, and DE2-method are still acceptable.

It is remarkable, that the DE1- and DE2-method give better results compared with the φ -, φ - and φ -, φ_n -method, for the same number of terms used in the expansion. This is probably due to the fact that the DE n -operator allows a more general solution, because it does not prescribe the precise form of the functions of θ (compare (4) and (5)), so the φ -, φ - and φ -, φ_n -method use a smaller class of functions and will in general require more terms to obtain a given accuracy.

4.3. Some results on the fine grid

From the results of Fig. 8a it follows that, for all methods, except the $\varphi\varphi_0$ -method, the error level is dominated by the mesh size, and not by the number of terms in the far-field expansion of φ . Hence, it seems appropriate to check this by investigating the effect of mesh-size halving. Not all experiments need to be repeated on a fine grid, because sufficient information may be obtained from a selection. Here, the calculations with the $\varphi\varphi_2$ -, $\varphi\varphi_n$ 2-, $\varphi\varphi_5$ -, $\varphi\varphi_n$ 5-, DE1- and DE2-method are reported on the fine grid on domain 3 (128×18 points), and the $\varphi\varphi_5$ - and $\varphi\varphi_n$ 5-method also on the fine grid on domain 4 (128×14 points).

The error graphs are presented in Fig. 9a–9b. From these figures, it may be seen that there are two dominant errors. The first one is caused by the truncation of the expansion on a fixed domain. The second is the numerical error caused by the mesh size. The first error can be seen clearly by comparing, on domain 3, $\varphi\varphi_2$ -, $\varphi\varphi_n$ 2-, DE1 in Fig. 8c to Fig. 9a. There is no significant altering of the C_p -error so two terms are too few on this domain. On domain 4, the same is true for $\varphi\varphi_5$ and $\varphi\varphi_n$ 5, see Fig. 8d and Fig. 9b. The second error, caused by the mesh width, is dominant in the results for the $\varphi\varphi_5$ - and $\varphi\varphi_n$ 5-method on domain 3 (Fig. 8c and Fig. 9a). It depends thus on the importance of the latter error whether better results can be achieved on a final grid. The DE2-method in Fig. 8c and Fig. 9a shows a case, where both error have about the same importance.

Note that the error in the circulation value is not directly determining the overall accuracy of the method. For instance, in Fig. 8d we find for the $\varphi\varphi_n$ 5-method a good value for Γ , but the C_p -error graph of $\varphi\varphi_n$ 5 is worse than the error graph of $\varphi\varphi_5$.

4.4. Remarks concerning convergence rates

Concerning numerical convergence aspects, the following conclusions were arrived at. For details, see [1].

- A grid-size reduction was usually accompanied by a proportional decrease in computation time, which means that the more complicated far-field approximation methods do not give a substantial overhead.

METHOD	Γ
exact	.92918
$\psi\psi_2$.915
$\psi\psi_{n2}$.898
$\psi\psi_5$.931
$\psi\psi_{n5}$.932
DE1	.901
DE2	.930

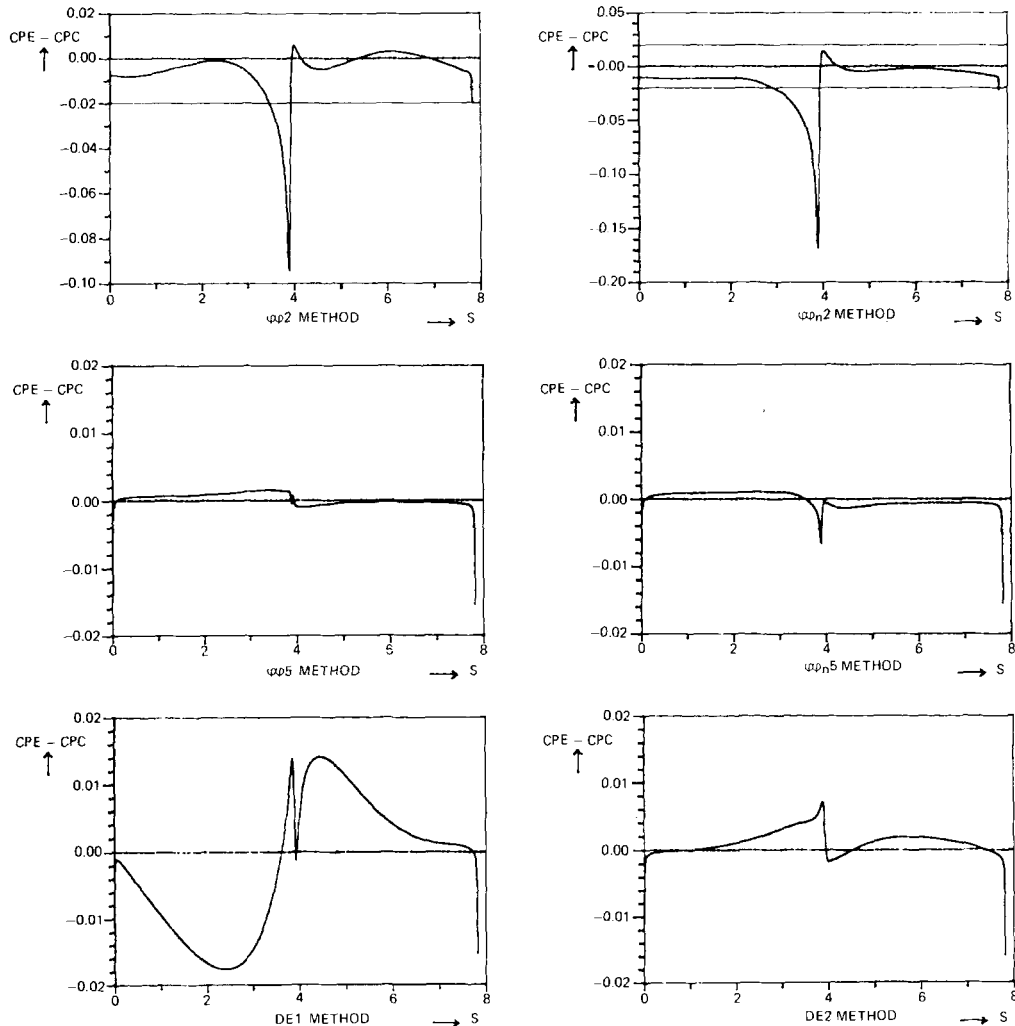


Figure 9a. C_p error results on fine grid, $\alpha = 4^\circ$, domain 3, $M = 129$, $N = 17$.

METHOD	Γ
exact	.92918
$\varphi\varphi 5$.939
$\varphi\varphi_n 5$.933

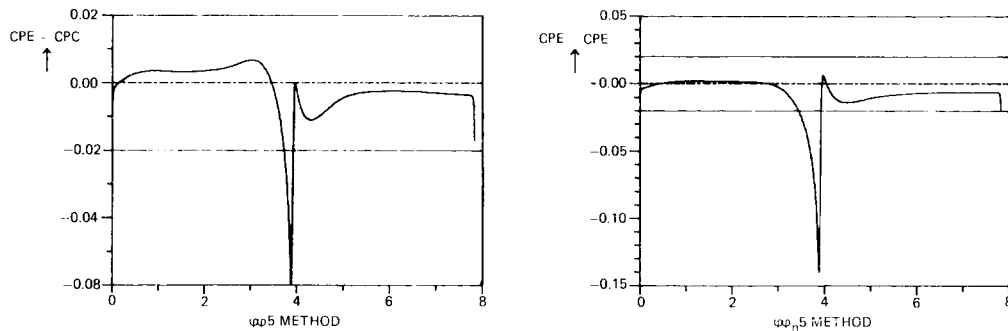


Figure 9b. C_p error results on fine grid, $\alpha = 4^\circ$, domain 4, $M = 129$, $N = 13$.

- Application of over- (under-) relaxation for the φ , φ - and φ , φ_n -methods resulted in a factor 2 to 3 decrease of computation times.
- φ , φ_n -methods converged faster than φ , φ -methods, due to the better estimation of the relaxation factor (see subsection 3.2).
- Richardson-extrapolation appeared to be successful, and the number of pure SLOR iterations could be reduced by a factor of 3.

4.5. Error analysis

An analytic analysis of the errors, as a function of the physical size of the computational domain and the number of terms used in the asymptotic far-field expansion, is given by Bayliss, Turkel and Gunzburger [7]. They have derived error approximations for the DE-methods which majorize errors due to a certain domain size and the number of terms used. This analysis has been extended by us [1] to errors in the potential and the C_p , also for the φ , φ - and the φ , φ_n -methods. Generally holds:

$$|CPE - CPC|_{max} < \frac{C}{r^{N+2}},$$

where N is the number of terms, used in the expansion. This estimate has been verified with the computed results, and was found to hold when the error is not dominated by errors due to the mesh size.

4.6. Summary of results

An overall appreciation of the accuracy of the results of the various methods will now be presented in Table 1. This judgement will be based on inspection of C_p -error plots and

Table 1. Appreciation table of computational results

Domain	Method						
	$\varphi\varphi 0$	$\varphi\varphi 2$	$\varphi\varphi_n 2$	$\varphi\varphi 5$	$\varphi\varphi_n 5$	DE1	DE2
1	0	+	+	+	+	+	+
2	-	+	+	+	+	+	+
3	-	-	-	+	+	-	-
4	-	-	-	0	-	-	-

circulation values. The engineering accuracy of each result is indicated by a +, 0, or - sign. The three signs have the following meaning:

- + : the result is clearly better than that of the $\varphi\varphi 0$ -method on domain 1,
- 0 : the result is of about the same quality as that of the $\varphi\varphi 0$ -method on domain 1,
- : the result is clearly worse than that of the $\varphi\varphi 0$ -method on domain 1.

From this table it is clear that the $\varphi\varphi 5$ -, $\varphi\varphi_n 5$ - and DE2-method are very interesting: they give good results on domain 3, which means that a grid size reduction of 50% is easily achieved for this test problem using these methods.

5. Concluding discussion

For the test problem of incompressible flow around fitting airfoils, the main conclusions may be summarized as follows.

- A substantial reduction in total number of grid points is possible when numerical boundary conditions are applied, which allow (DE-methods) or contain (φ , φ_n - and φ , φ -methods) a numerical far-field expansion with more terms than usual. Good results may still be found when the total number of grid points is reduced by 50%, which corresponds to a reduction of the physical domain by 85% (Section 4).
- The φ , φ_n - and φ , φ -methods are more general and easier to implement than the DE-methods, because they lead to simple Dirichlet or Neumann boundary conditions at the numerical boundary (Section 3). The key to the success is evidently the approximation efficiency of the potential on the other infinite domain by a finite number of terms, accounting for the uniform parallel flow, a vortex term subtracting out the singularity of the potential, and further consisting of the first few terms of a series approximating an arbitrary analytic potential on the outer domains. Subtraction of the singularity is obviously required to achieve good approximation efficiency.
- The different kinds of approximations of the flow (grid function on the inner domain, asymptotic expansion on the finite outer domain) may be efficiently matched by least-squares approximations of continuity conditions over the common boundary.

Although optimization of the numerical solution algorithm was not the primary goal of the present study, the following conclusions can be drawn:

- Underrelaxation (φ , φ_n -method) or overrelaxation (φ , φ -method) accelerates the convergence of the iterative algorithm updating in turn the finite-difference approximation on the finite inner domain and the coefficients in the finite combination of asymptotic solutions on the outer domain.

- When designing such iterative algorithms, it should be kept in mind that these coefficients should only be updated, when the large wave-length part in the finite-difference approximation has been sufficiently improved. It is expected that efficient alternating update techniques may be realized in conjunction with FAS multigrid relaxation algorithms (which allow a fast update of the long wave-length part on the coarser grids of the algorithm).

What can be learned from the study in view of other applications? It is expected that the φ , φ_n - and the φ , φ -methods may be generalized to handle inviscid steady compressible potential flows around airfoils. In that case, it is necessary to expand also the nonlinear potential equation in a systematic way in terms of the well-known distance variable of linearized Prandtl-Glauert theory [9], and the condition of continuity of φ , φ_n has to be replaced by a condition of continuous normal mass flux.

A more ambitious generalization is inviscid compressible steady potential flow around aircraft configurations. The vortex sheet downstream of the aircraft configuration generates a singularity that has to be subtracted out in a sufficiently accurate way.

Other generalizations coming into mind are unsteady potential-flow and Euler-flow calculations around airfoils and aircraft configurations. The authors are aware of the application of higher-order asymptotic methods in other branches of mathematical physics where infinite domains occur, e.g. ship dynamics, sound radiation, and electro-magnetic field theory. It is hoped that the experiences reported here are also useful to researchers in these other branches.

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