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# Optical/NIR stellar absorption and emission-line indices from luminous infrared galaxies 

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#### Abstract

We analyse a set of optical-to-near-infrared long-slit nuclear spectra of 16 infrared-luminous spiral galaxies. All of the studied sources present $\mathrm{H}_{2}$ emission, which reflects the star-forming nature of our sample, and they clearly display Hi emission lines in the optical. Their continua contain many strong stellar absorption lines, with the most common features due to $\mathrm{CaI}, \mathrm{CaII}$, FeI, $\mathrm{NaI}, \mathrm{Mg} \mathrm{I}$, in addition to prominent absorption bands of $\mathrm{TiO}, \mathrm{VO}, \mathrm{ZrO}, \mathrm{CN}$, and CO . We report a homogeneous set of equivalent width (EW) measurements for 45 indices, from optical to NIR species for the 16 star-forming galaxies as well as for 19 early-type galaxies where we collected the data from the literature. This selected set of emission and absorptionfeature measurements can be used to test predictions of the forthcoming generations of stellar population models. We find correlations among the different absorption features and propose here correlations between optical and NIR indices, as well as among different NIR indices, and compare them with model predictions. Although for the optical absorption features the models consistently agree with the observations, the NIR indices are much harder to interpret. For early-type spirals the measurements agree roughly with the models, while for star-forming objects they fail to predict the strengths of these indices.


Key words: stars: AGB and post-AGB - galaxies: bulges - galaxies: evolution-galaxies: stellar content.

## 1 INTRODUCTION

One challenge in modern astrophysics is to understand galaxy formation and evolution. Both processes are strongly related to the

[^0]star-formation history (SFH) of galaxies. Thus, the detailed study of the different stellar populations found in galaxies is one of the most promising ways to shed some light on their evolutionary histories. So far, stellar population studies have been concentrated mainly in the optical spectral range (e.g. Bica 1988; Worthey et al. 1994; Trager et al. 2000; Sánchez-Blázquez et al. 2006; González Delgado et al. 2015; Goddard et al. 2017; Martín-Navarro et al. 2018). In the near-infrared, ( $0.8-2.4 \mu \mathrm{~m}, \mathrm{NIR}$ ) even with some work dating back to the 1980s (e.g. Rieke et al. 1980), stellar population studies have just started to become more common in the last two decades

Table 1. Near-infrared observation $\log$ and basic sample properties.

| Source | $\alpha$ | $\delta$ | z | Obs. date | Exp. time <br> (s) | Airmass | $\begin{gathered} \text { PA } \\ (\mathrm{deg}) \end{gathered}$ | $\begin{gathered} \text { Size } \\ (\mathrm{pc} \times \mathrm{pc}) \end{gathered}$ | Activity | $\begin{gathered} \log \left(\frac{L_{\mathrm{R}}}{L_{\odot}}\right)^{\star} \\ m \end{gathered}$ | Morphology |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| NGC 23 | 00h09m53.4s | +25d55m26s | 0.0157202 | 2010-10-07 | $29 \times 120$ | 1.04 | 330 | $1348 \times 270$ | $\mathrm{SFG}^{a}$ | 11.05 | SBa |
| NGC 520 | 01 h 24 m 35.1 s | $+03 \mathrm{~d} 47 \mathrm{~m} 33 \mathrm{~s}$ | 0.0080367 | 2010-10-04 | $16 \times 120$ | 1.04 | 300 | $689 \times 138$ | $\mathrm{SFG}^{b}$ | 10.91 | S0 |
| NGC 660 | 01 h 43 m 02.4 s | $+13 \mathrm{~d} 38 \mathrm{~m} 42 \mathrm{~s}$ | 0.0029152 | 2010-10-06 | $24 \times 120$ | 1.01 | 33 | $237 \times 50$ | Sy2/H $\mathrm{II}^{\text {b, }, ~}{ }^{\text {d }}$ d | 10.49 | SBa pec |
| NGC 1055 | 02 h 41 m 45.2 s | $+00 \mathrm{~d} 26 \mathrm{~m} 35 \mathrm{~s}$ | 0.0036267 | 2010-10-04 | $16 \times 120$ | 1.07 | 285 | $466 \times 62$ | LINER/H $\mathrm{II}^{\text {b, }, ~}{ }^{\text {d }}$ | 10.09 | Sbc |
| NGC 1134 | 02h53m41.3s | +13d00m51s | 0.0129803 | 2010-10-04 | $16 \times 120$ | 1.11 | 0 | $1113 \times 223$ | $\mathrm{SFG}^{e}$ | 10.83 | S? |
| NGC 1204 | 03h04m39.9s | $-12 \mathrm{~d} 20 \mathrm{~m} 29 \mathrm{~s}$ | 0.0154058 | 2010-10-07 | $16 \times 120$ | 1.23 | 66 | $1321 \times 264$ | $\mathrm{LINER}^{f}$ | 10.88 | S0/a |
| NGC 1222 | 03h08m56.7s | $-02 \mathrm{~d} 57 \mathrm{~m} 19 \mathrm{~s}$ | 0.0082097 | 2010-10-06 | $24 \times 120$ | 1.13 | 315 | $598 \times 141$ | $\mathrm{SFG}^{g}$ | 10.60 | S0 pec |
| NGC 1266 | 03h16m00.7s | -02d25m38s | 0.0077032 | 2010-10-07 | $18 \times 120$ | 1.09 | 0 | $661 \times 132$ | LINER ${ }^{\text {g }}$ | 10.46 | SB0 pec |
| UGC 2982 | 04h12m22.4s | +05d32m51s | 0.0177955 | 2010-10-04 | $9 \times 120$ | 1.11 | 295 | $1526 \times 305$ | $\mathrm{SFG}^{h}$ | 11.30 | SB |
| NGC 1797 | 05h 07 m 44.9 s | -08d01m09s | 0.0154111 | 2010-10-07 | $16 \times 120$ | 1.23 | 66 | $1321 \times 264$ | $\mathrm{SFG}^{a}$ | 11.00 | SBa |
| NGC 6814 | 19h42m40.6s | -10d19m25s | 0.0056730 | 2010-10-07 | $16 \times 120$ | 1.17 | 0 | $486 \times 97$ | Sy $1^{g}$ | 10.25 | SBbc |
| NGC 6835 | 19 h 54 m 32.9 s | $-12 \mathrm{~d} 34 \mathrm{~m} 03 \mathrm{~s}$ | 0.0057248 | 2010-10-06 | $22 \times 120$ | 1.21 | 70 | $368 \times 98$ | $\mathrm{SFG}^{i}$ | 10.32 | SBa |
| UGC 12150 | 22 h 41 m 12.2 s | +34d14m57s | 0.0214590 | 2010-10-04 | $15 \times 120$ | 1.08 | 37 | $1656 \times 368$ | LINER/H $\mathrm{II}^{j}$ | 11.29 | SB0/a |
| NGC 7465 | 23h02m01.0s | +15d57m53s | 0.0066328 | 2010-10-06 | $12 \times 120$ | 1.03 | 340 | $569 \times 114$ | LINER/Sy $2^{k}$ | 10.10 | SB0 |
| NGC 7591 | 23 h 18 m 16.3 s | +06d35m09s | 0.0165841 | 2010-10-07 | $16 \times 120$ | 1.03 | 0 | $1422 \times 284$ | LINER ${ }^{\text {g }}$ | 11.05 | SBbc |
| NGC 7678 | 23h28m27.9s | +22d25m16s | 0.0120136 | 2010-10-04 | $16 \times 120$ | 1.01 | 90 | $927 \times 206$ | $\mathrm{SFG}^{l}$ | 10.77 | SBc |

Note. SFG: Star-forming galaxies (Starburst or H iI galaxies). LINER/H II were assumed to be pure LINERs in the text. The galaxies are listed in order of right ascension, and the number of exposures refers to on-source integrations. The slit width is 0.8 arcsec .
${ }^{a}$ Balzano (1983); ${ }^{b}$ Ho, Filippenko \& Sargent (1997a); ${ }^{c}$ Ho et al. (1997b); ${ }^{d}$ Filho et al. (2004); ${ }^{e}$ Condon, Cotton \& Broderick (2002); ${ }^{f}$ Sturm et al. (2006); ${ }^{g}$ Pereira-Santaella et al. (2010); ${ }^{h}$ Schmitt et al. (2006); ${ }^{i}$ Coziol $+98 ;{ }^{j}$ Veilleux et al. (1995); ${ }^{k}$ Ferruit, Wilson \& Mulchaey (2000); ${ }^{l}$ Gonçalves, Veron \& Veron-Cetty (1998); ${ }^{m}$ Sanders et al. (2003).
(Origlia, Moorwood \& Oliva 1993; Origlia et al. 1997; Riffel et al. 2007, 2008, 2009; Cesetti et al. 2009; Lyubenova et al. 2010; ChiesSantos et al. 2011a,b; Riffel et al. 2011c; Kotilainen et al. 2012; La Barbera et al. 2013; Martins et al. 2013b; Noël et al. 2013; Zibetti et al. 2013; Dametto et al. 2014; Riffel et al. 2015; Baldwin et al. 2017; Alton, Smith \& Lucey 2018; Dahmer-Hahn et al. 2018, 2019; Francois et al. 2019; Dametto et al. 2019, for example). Models have shown that the NIR spectral features provide very important insights, particularly into the stellar populations dominated by cold stars (e.g. Maraston 2005; Riffel et al. 2007; Conroy \& van Dokkum 2012; van Dokkum \& Conroy 2012; Zibetti et al. 2013; Riffel et al. 2015; Röck 2015; Röck et al. 2016). For example, the stars in the thermally pulsing asymptotic giant branch (TP-AGB) phase may be responsible for nearly half of the luminosity in the $K$ band for stellar populations with an age of $\sim 1 \mathrm{Gyr}$ (Maraston 1998, 2005; Salaris et al. 2014).

One common technique to study the unresolved stellar content of galaxies is the fitting of a combination of simple stellar populations (SSPs) to obtain the SFH. However, due to difficulties in theoretical treatment (Maraston 2005; Marigo et al. 2008; Noël et al. 2013) and the lack of complete empirical stellar libraries in the NIR (Lançon et al. 2001; Chen et al. 2014; Riffel et al. 2015; Villaume et al. 2017) the available SSP models produce discrepant results (e.g. Baldwin et al. 2017), thus making it very difficult to reliably analyse the SFH in the NIR.

On the other hand, the stellar content and chemical composition of the unresolved stellar populations of galaxies can also be obtained by the study of the observed absorption features present in their integrated spectra. So far, we still lack a comprehensive NIR data set to compare with model predictions, required to make improvements to the models and to lead to a better understanding of the role played by the cooler stellar populations in the integrated spectra of galaxies.

Among the best natural laboratories to study these kinds of stellar content are infrared galaxies, sources that emit more energy in the infrared $(\sim 5-500 \mu \mathrm{~m})$ than at all the other wavelengths combined (Sanders \& Mirabel 1996; Sanders et al. 2003). The relevance of
studying these galaxies lies particularly in the fact that they are implicated in a variety of interesting astrophysical phenomena, including the formation of quasars and elliptical galaxies (e.g. Genzel et al. 2001; Veilleux 2006; Wang et al. 2006). When studying luminous infrared galaxies in the Local Universe, it is possible to obtain high-angular-resolution observations of these objects, thus allowing the investigation of their very central regions. Comparison of such objects with those at higher redshifts may help to understand the SFH over cosmic times.

With the above in mind, we obtained optical and NIR spectra of a subsample of galaxies selected from the IRAS Revised Bright Galaxy Sample present in the Local Universe. These galaxies are believed to be experiencing massive star formation, making them suitable for studying their most important spectral features that can be used as proxies to test and constrain stellar-population models. As part of a series of papers aimed at studying the stellar population and gas emission features, here we provide measurements for the most conspicuous emission and absorption features, and present new correlations between absorption features. The outline of the paper is as follows: in Section 2 we describe the observations and data reduction. The results are presented and discussed in Section 3. Final remarks are made in Section 5.

## 2 OBSERVATIONS AND DATA REDUCTION

Our sample is composed of 16 Local Universe ( $\mathrm{v}_{\mathrm{r}} \lesssim 6400 \mathrm{~km}$ $\mathrm{s}^{-1}$ ) galaxies that are very bright in the infrared (see Table 1). They were selected from the IRAS Revised Bright Galaxy Sample, which is regarded as a statistically complete sample of 629 galaxies, with $60 \mu \mathrm{~m}$ flux density $\gtrsim 5.24 \mathrm{Jy}$. Galaxies chosen for this study were those with $\log \left(L_{\text {fir }} / L_{\odot}\right) \gtrsim 10.10$, accessible from the Infrared Telescope Facility (IRTF) and the Wyoming Infrared Observatory (WIRO, see below), and bright enough to reach an $\mathrm{S} / \mathrm{N} \sim 100$ in the $K$ band within a reasonable on-source integration time.





NGC1222


NGC1266


Figure 1. Final reduced and redshift-corrected spectra for NGC 1134, NGC 1204, NGC 1222, and NGC 1266. For each galaxy we show from top to bottom the optical, $z+J, H$, and $K$ bands, respectively. The flux is in units of $10^{-15} \mathrm{erg} \mathrm{cm}^{-2} \mathrm{~s}^{-1}$. The shaded grey area represents the uncertainties and the brown area indicates the poor transmission regions between different bands. The remaining spectra are shown in online material.

Table 2. Optical observation log. The slit was oriented north-south.

| Source | Obs. date | $\begin{aligned} & \text { Exp. } \\ & \text { time (s) } \end{aligned}$ | Airmass | $\begin{gathered} \text { Size } \\ (\mathrm{pc} \times \mathrm{pc}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: |
| NGC 23 | 2010-10-04 | 600 | 1.20 | $2359 \times 1348$ |
| NGC 520 | 2010-10-03 | 600 | 1.28 | $4307 \times 689$ |
| NGC 660 | 2010-10-04 | 600 | 1.18 | $437 \times 250$ |
| NGC 1055 | 2010-10-04 | 600 | 1.36 | $544 \times 311$ |
| NGC 1134 | 2010-10-04 | 600 | 1.14 | $12243 \times 1113$ |
| NGC 1204 | 2010-10-04 | 600 | 1.68 | $2312 \times 1321$ |
| NGC 1222 | 2010-10-03 | 600 | 1.42 | $3344 \times 704$ |
| NGC 1266 | 2010-10-03 | 600 | 1.39 | $6440 \times 661$ |
| UGC 2982 | 2010-10-04 | 600 | 1.27 | $2670 \times 1526$ |
| NGC 1797 | 2010-10-02 | 600 | 1.56 | $11893 \times 1321$ |
| UGC 12150 | 2010-10-03 | 600 | 1.03 | $18401 \times 1840$ |
| NGC 7465 | 2010-10-03 | 600 | 1.12 | $3270 \times 569$ |
| NGC 7591 | 2010-10-03 | 600 | 1.21 | $13154 \times 1422$ |
| NGC 7678 | 2010-10-04 | 600 | 1.18 | $1803 \times 1030$ |

Note. The slit width is 4 arcsec.

### 2.1 Near-infrared data

Cross-dispersed NIR spectra in the range $0.8-2.4 \mu \mathrm{~m}$ were obtained on 2010 October 4, 6, and 7 with the SpeX spectrograph (Rayner et al. 2003) attached to the NASA 3 m IRTF telescope at the Mauna Kea observing site. The detector is a $1024 \times 1024$ ALADDIN 3 InSb array with a spatial scale of 0.15 arcsec pixel $^{-1}$. A $0.8 \mathrm{arcsec} \times$ 15 arcsec slit was used during the observations, giving a spectral resolution of $\mathrm{R} \sim 1000$ (or $\sigma=127 \mathrm{~km} \mathrm{~s}^{-1}$ ). Both the arc lamp spectra and the night-sky spectra are consistent with this value (Riffel et al. 2013a). The observations were done by nodding in an Object-Sky-Object pattern with typical individual integration times of 120 s and total on-source integration times between 18 and 58 min . During the observations, A0 V stars were observed near each target to provide telluric standards at similar air masses. These stars were also used to flux calibrate the galaxy spectra using blackbody functions to calibrate the observed spectra of the standard stars. The seeing varied between 0.4 and 0.7 arcsec over the different nights and there were no obvious clouds.

We reduced the NIR observations following the standard data reduction procedures given by Riffel, Rodríguez-Ardila \& Pastoriza (2006) and Riffel et al. (2013b). In short, spectral extraction and wavelength calibration were performed using SPEXTOOL, software developed and provided by the SpeX team for the IRTF community (Cushing, Vacca \& Rayner 2004). The area of the integrated region is listed in Table 1. Each extraction was centred at the peak of the continuum-light distribution for every object of the sample. No effort was made to extract spectra at positions different from the nuclear region, even though some objects show evidence of extended emission, as this goes beyond the scope of this analysis. Telluric absorption correction and flux calibration were applied to the individual 1D spectra by means of the IDL routine xtellcor (Vacca, Cushing \& Rayner 2003).

### 2.2 Optical data

For completeness, the same sample was also observed in the optical range on nearly the same dates as the NIR data were collected with the WIRO long-slit spectrograph. The instrument is attached to the University of Wyoming's 2.3-m telescope, located on Jelm Mountain at WIRO. The Cassegrain-mounted instrument uses a Marconi $2 k \times 2 k$ CCD detector. During our observations we used a
$9001 / \mathrm{mm}$ grating in first order to obtain spectra from approximately 4000-7000 Å calibrated with a CuAr comparison lamp. Given our 4 -arcsec slit oriented north-south, the resolution was R $\sim 1200$. Due to the relatively large spatial extent of these low-redshift objects, we offset the telescope pointing by 2 arcmin to obtain sky spectra uncontaminated by galaxy light. The seeing varied between 1 and 2 arcsec during the nights of observation. We reduced the spectra using standard techniques in IRAF. ${ }^{1}$ Table 1 shows the observation log along with extraction apertures. The 1D wavelength and fluxcalibrated spectra were then corrected for redshift, determined from the average $z$ measured from the position of [S III] $0.953 \mu \mathrm{~m}, \mathrm{~Pa} \delta$, He I $1.083 \mu \mathrm{~m}, \mathrm{~Pa} \beta$, and $\mathrm{Br} \gamma$.

Examples of the final reduced spectra, from optical to NIR ( $\sim 0.4-$ $2.4 \mu \mathrm{~m}$ ) are presented in Fig. 1, for the remaining galaxies see Appendix A. For each galaxy we show the optical, $z+J, H$, and $K$ bands, from top to bottom, respectively. It is worth mentioning that the optical and NIR data do not share the same apertures, and the slit was not generally oriented at the same position angles. However, since we are interested in the nuclear region, the different slit orientations should not introduce large discrepancies in the measurements. The ordinate axis represents the monochromatic flux in units of $10^{-15} \mathrm{erg} \mathrm{cm}^{-2} \mathrm{~s}^{-1} \AA^{-1}$. The position of the most common and expected emission and absorption lines are indicated as dotted (red) and dashed (blue) lines, respectively.

## 3 RESULTS

### 3.1 Emission-line spectra

A visual inspection of the data reveals a wide diversity of emissionline strengths and species. The most common emission features detected are: $\mathrm{H} \beta$, [O III] 4959, $5007 \AA$, [N II] 6548, $6583 \AA$, $\mathrm{H} \alpha$, [S II] $6716,6730 \AA$ A, [S III] $9531 \AA, \mathrm{~Pa} \delta,\left[\right.$ C I $^{2} 9824,9850 \AA, \mathrm{~Pa} \beta, \mathrm{He} \mathrm{I}$ $10830 \AA \AA^{\circ}$, [P II] $11886 \AA$ A, [Fe II] 12570, $16436 \AA$, $\mathrm{Pa} \alpha \mathrm{H}_{2} 19570 \AA$, $\mathrm{H}_{2} 21218 \AA$, and $\mathrm{Br} \gamma$.

Emission-line fluxes for each object of the sample were measured by fitting a Gaussian function to the observed profile and then integrating the flux under the curve. The LINER software (Pogge \& Owen 1993) was used for this purpose. No attempt to correct for stellar absorption was made before measuring the emission lines. This was done because NIR models with adequate spectral resolution (to allow the measurements of the weaker emission lines) are not available for the younger ages. Martins et al. (2013a) have shown that the underlying stellar population has only a strong effect on the hydrogen recombination emission lines, with the largest differences in fluxes being about 25 per cent. This value is within the largest uncertainties on the fluxes values too. For completeness, we have not subtracted the stellar features from the optical range too.

The results, including $3 \sigma$ uncertainties, are listed in Tables 3 and 4. For most of our targets, these measurements are made public for the first time. In addition, we computed the extinction coefficient, $\mathrm{C}_{\mathrm{ext}}$, for the NIR using the intrinsic value of 5.88 for the flux ratio of $\mathrm{Pa} \beta / \mathrm{Br} \gamma$ (Hummer \& Storey 1987, using case B). The Cardelli, Clayton \& Mathis (1989) extinction law was used, and the values obtained for the coefficients are listed in Tables 3 and 4.

[^1]Table 3. Emission-line fluxes in units of $1 \times 10^{-15} \mathrm{erg} \mathrm{cm}^{-2} \mathrm{~s}^{-1}$.

| Line | Ion | NGC 23 | NGC 520 | NGC 660 | NGC 1055 | NGC 1134 | NGC 1204 | NGC 1222 | NGC 1266 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $C_{\text {ext }}$ | - | - | $3.42 \pm 0.13$ | - | - | $2.95 \pm 0.14$ | $1.06 \pm 0.05$ | $6.74 \pm 0.73$ |
| 4861 | $\mathrm{H} \beta$ | $20.40 \pm 1.16$ | - | $0.86 \pm 0.26$ | - | $6.06 \pm 0.37$ | $0.30 \pm 0.06$ | $49.30 \pm 0.65$ | - |
| 4959 | [ O III] | $30.70 \pm 3.13$ | - | $3.37 \pm 0.41$ | - | $8.25 \pm 0.72$ | $1.69 \pm 0.48$ | $47.30 \pm 0.75$ | - |
| 5007 | [ OIII ] $^{\text {I }}$ | $10.50 \pm 3.13$ | $2.16 \pm 0.30$ | $1.15 \pm 0.41$ | - | $2.81 \pm 0.72$ | $0.58 \pm 0.48$ | $140.00 \pm 0.75$ | $4.27 \pm 0.64$ |
| 6548 | [ NiI ] | $108.00 \pm 3.26$ | $1.60 \pm 0.31$ | $6.25 \pm 0.53$ | $1.41 \pm 0.32$ | $12.90 \pm 1.41$ | $5.75 \pm 0.43$ | $14.20 \pm 0.40$ | $14.30 \pm 1.07$ |
| 6563 | $\mathrm{H} \alpha$ | $83.80 \pm 3.89$ | $8.68 \pm 0.44$ | $17.20 \pm 0.46$ | $4.97 \pm 0.34$ | $40.50 \pm 0.95$ | $20.50 \pm 0.42$ | $222.00 \pm 0.43$ | $11.40 \pm 0.76$ |
| 6583 | [ $\mathrm{NiI}^{\text {] }}$ | $27.90 \pm 3.89$ | $6.96 \pm 0.49$ | $18.60 \pm 0.48$ | $4.56 \pm 0.39$ | $38.30 \pm 1.21$ | $18.60 \pm 0.42$ | $49.70 \pm 0.45$ | $30.60 \pm 0.87$ |
| 6716 | [S II] | $46.50 \pm 1.92$ | $3.29 \pm 0.46$ | $5.15 \pm 0.48$ | $1.25 \pm 0.29$ | $14.60 \pm 1.23$ | $5.76 \pm 1.23$ | $20.70 \pm 0.48$ | $18.70 \pm 1.16$ |
| 6730 | [S II] | $34.70 \pm 1.92$ | $3.07 \pm 0.55$ | $4.88 \pm 0.61$ | $1.31 \pm 0.44$ | $12.90 \pm 1.50$ | $5.69 \pm 1.23$ | $19.90 \pm 0.57$ | $21.10 \pm 1.16$ |
| 9069 | [S III] | $14.80 \pm 3.96$ | - | - | - | - | $10.30 \pm 1.40$ | $32.80 \pm 1.59$ | - |
| 9531 | [S III] | $15.70 \pm 3.96$ | - | $23.90 \pm 0.94$ | - | - | $11.90 \pm 0.54$ | $72.60 \pm 1.49$ | - |
| 9824 | [CI] | $1.57 \pm 0.60$ | - | $2.10 \pm 0.27$ | - | - | $1.15 \pm 0.30$ | $1.55 \pm 0.49$ | - |
| 9850 | [CI] | $5.15 \pm 0.60$ | - | $2.53 \pm 0.27$ | - | - | $1.88 \pm 0.30$ | $0.97 \pm 0.49$ | $9.21 \pm 1.06$ |
| 10049 | $\mathrm{Pa} \delta$ | - | - | $2.55 \pm 0.22$ | - | - | $1.03 \pm 0.08$ | $4.05 \pm 0.31$ | - |
| 10122 | He II | - | - | $2.78 \pm 0.22$ | - | - | $1.34 \pm 0.08$ | - | - |
| 10830 | He I | $22.30 \pm 2.60$ | - | $13.30 \pm 0.41$ | - | - | $9.24 \pm 0.56$ | $53.80 \pm 0.77$ | $6.26 \pm 0.83$ |
| 10938 | $\mathrm{Pa} \gamma$ | $5.65 \pm 1.58$ | - | $8.57 \pm 0.31$ | - | - | $2.78 \pm 0.27$ | $12.90 \pm 0.76$ | - |
| 11470 | [P II] | - | - | $1.82 \pm 1.09$ | - | - | $1.17 \pm 0.27$ | - | - |
| 11886 | [ $\mathrm{PII}^{\text {I }}$ ] | - | - | $3.68 \pm 1.09$ | - | - | $1.70 \pm 0.19$ | - | - |
| 12567 | [ Fe II] | $10.50 \pm 0.87$ | - | $13.90 \pm 0.65$ | - | - | $5.16 \pm 0.20$ | $6.32 \pm 0.43$ | $3.19 \pm 0.51$ |
| 12820 | $\mathrm{Pa} \beta$ | - | - | $29.00 \pm 0.60$ | - | - | $12.10 \pm 0.20$ | $28.20 \pm 0.37$ | $0.75 \pm 0.08$ |
| 12950 | [ Fe II] | - | - | $1.21 \pm 0.15$ | - | - | $1.35 \pm 0.25$ | - | - |
| 13209 | [ Fe II] | - | - | $6.60 \pm 0.24$ | - | - | $4.49 \pm 0.39$ | - | - |
| 15342 | [ Fe II] | - | - | - | - | - | $1.52 \pm 0.36$ | - | - |
| 16436 | [Fe II] | $14.20 \pm 3.46$ | $3.87 \pm 0.14$ | $15.70 \pm 0.90$ | - | - | $6.51 \pm 0.36$ | $4.65 \pm 0.18$ | $2.90 \pm 0.38$ |
| 16773 | $\left[\mathrm{Fe}_{\mathrm{II}}\right]+\mathrm{Br} 11$ | $31.50 \pm 1.34$ | - | - | - | - | $2.65 \pm 0.50$ | - | - |
| 17360 | Br10 | - | - | - | - | - | $1.60 \pm 0.13$ | - | - |
| 18750 | Pa $\alpha$ | - | $35.40 \pm 0.30$ | $60.20 \pm 1.64$ | - | - | $61.10 \pm 0.33$ | $69.30 \pm 0.57$ | $8.03 \pm 0.34$ |
| 19446 | Br $\delta$ | - | $2.07 \pm 0.28$ | $5.67 \pm 0.53$ | - | - | $1.80 \pm 0.34$ | $3.78 \pm 0.20$ | - |
| 19570 | $\mathrm{H}_{2}$ | $20.20 \pm 4.60$ | $2.41 \pm 0.39$ | $8.87 \pm 0.80$ | $1.24 \pm 0.3$ | - | $4.17 \pm 0.51$ | $1.95 \pm 0.34$ | $15.10 \pm 0.47$ |
| 20332 | $\mathrm{H}_{2}$ | $4.94 \pm 0.57$ | $1.13 \pm 0.10$ | $3.67 \pm 0.56$ | - | $1.14 \pm 0.3$ | $1.81 \pm 0.34$ | $0.95 \pm 0.14$ | $5.10 \pm 0.44$ |
| 20580 | $\mathrm{H}_{2}$ | - | $2.47 \pm 0.09$ | $5.32 \pm 0.64$ | - | - | $1.47 \pm 0.27$ | $4.68 \pm 0.14$ | - |
| 21218 | $\mathrm{H}_{2}$ | $10.00 \pm 1.26$ | $2.40 \pm 0.18$ | $6.91 \pm 0.72$ | $0.6 \pm 0.08$ | $2.37 \pm 0.4$ | $3.61 \pm 0.25$ | $0.84 \pm 0.12$ | $13.70 \pm 0.25$ |
| 21654 | $\mathrm{Br} \gamma$ | - | $6.67 \pm 0.19$ | $16.60 \pm 0.71$ | - | - | $5.88 \pm 0.27$ | $6.98 \pm 0.07$ | $1.40 \pm 0.33$ |
| 22230 | $\mathrm{H}_{2}$ | $4.53 \pm 2.97$ | $0.67 \pm 0.19$ | $2.01 \pm 1.20$ | - | - | $0.87 \pm 0.09$ | $0.49 \pm 0.19$ | $3.26 \pm 0.12$ |
| 22470 | $\mathrm{H}_{2}$ | $1.47 \pm 0.37$ | $0.83 \pm 0.12$ | $1.18 \pm 0.12$ | - | - | $0.72 \pm 0.10$ | $0.29 \pm 0.04$ | $1.51 \pm 0.14$ |

### 3.2 The continuum spectra

The main goal of this section is to characterize the continuum emission observed in our sample and compare it to other data in the literature. To help in the visual inspection ${ }^{2}$ of the individual spectra, we normalized the continuum emission to unity in two regions free from emission/absorption features taken from Riffel et al. (2011b). The NIR spectra were normalized at $20925 \AA$ and then sorted according to their continuum shapes. For a proper comparison with the optical portion of the spectrum, we normalized the optical spectra at $5300 \AA$ and plotted them in the same order as the NIR spectra (Figs 4 and 5).

A first-order inspection of Fig. 4 allows us to infer that, contrary to what happens in Seyfert galaxies (Riffel et al. 2006), there seems to be no correlation between activity type (LINERs or SFGs) and continuum shape. In fact, these very bright infrared galaxies present a continuum shape very similar to what is found in fainter H II sources and normal galaxies, as reported by Martins et al. (2013a), which may indicate that the LINER spectrum of these galaxies is powered by starburst instead of a low-luminosity AGN. In addition, the continua of all the optical spectra look very similar.

[^2]A large diversity of atomic absorption lines and molecular bands is also apparent in the spectra. These features are seen from the very blue optical end to the red end of the observed NIR spectral region. The most common atomic absorption features are due to CaI, Ca II, Fe I, Si I, Na I, and Mg I, besides the prominent absorption bands of $\mathrm{CH}, \mathrm{MgH}, \mathrm{TiO}, \mathrm{VO}, \mathrm{ZrO}$, and CO . These features are identified in Fig. 1. It is clear in these figures that some of the most important features predicted for intermediate-age stellar populations, which are expected to be enhanced in the RGB and TP-AGB stellar phases (Maraston 2005; Riffel et al. 2007, 2015), are detected in the spectra. Among these features are the $\mathrm{ZrO} / \mathrm{CN} / \mathrm{VO}$ at $9350 \AA$, the $10560 \AA$ VO, $1.1 \mu \mathrm{~m} \mathrm{CN}$, and 1.6 and $2.3 \mu \mathrm{~m}$ CO bands.

### 3.2.1 Towards a homogeneous NIR index definition

The equivalent widths (EWs) of these features offer coarse but robust information about the stellar content of a galaxy spectrum, and therefore they can be used as powerful diagnostics of the stellar content of galaxies. Contrary to the optical range, where there exist indices defined in a homogeneous way by the Lick group (see Worthey et al. 1994, and references), in the NIR there is no such homogeneous set of definitions covering the full NIR wavelength range, and authors tend to use their own definitions (e.g. Riffel

Table 4. Continuation of Table 3.

| Line | Ion | UGC2982 | NGC 1797 | NGC 6814* | NGC 6835 | UGC12150 | NGC 7465 | NGC 7591 | NGC 7678 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $C_{\text {ext }}$ | - | $2.68 \pm 0.08$ | 0.00 | $4.87 \pm 0.19$ | $2.16 \pm 0.17$ | $2.35 \pm 0.93$ | $2.53 \pm 0.07$ | $1.39 \pm 0.23$ |
| 4861 | $\mathrm{H} \beta$ | $0.49 \pm 0.15$ | $13.10 \pm 0.71$ | - | - | - | $21.80 \pm 1.05$ | - | $9.45 \pm 0.42$ |
| 4959 | [O III] | $0.35 \pm 0.16$ | - | - | - | - | $31.40 \pm 2.27$ | - | - |
| 5007 | [O III] | $0.76 \pm 0.29$ | $3.11 \pm 0.44$ | - | - | - | $56.00 \pm 1.85$ | - | $2.83 \pm 0.54$ |
| 6548 | [ NiI ] | $1.64 \pm 0.39$ | $11.20 \pm 0.59$ | - | - | - | $29.50 \pm 1.03$ | $8.15 \pm 1.36$ | $7.88 \pm 0.82$ |
| 6563 | $\mathrm{H} \alpha$ | $13.90 \pm 0.40$ | $88.70 \pm 0.74$ | - | - | - | $122.00 \pm 0.89$ | $21.10 \pm 1.00$ | $53.40 \pm 1.00$ |
| 6583 | [ $\mathrm{NiI}^{\text {I }}$ | $5.62 \pm 0.39$ | $46.40 \pm 0.76$ | - | - | - | $71.10 \pm 0.98$ | $18.60 \pm 1.04$ | $27.60 \pm 1.00$ |
| 6716 | [S II] | $2.77 \pm 0.45$ | $10.80 \pm 0.41$ | - | - | - | $39.70 \pm 1.17$ | $4.87 \pm 1.94$ | $9.23 \pm 0.97$ |
| 6730 | [S II] | $2.66 \pm 0.75$ | $10.30 \pm 0.46$ | - | - | - | $34.10 \pm 1.22$ | $3.46 \pm 1.94$ | $10.30 \pm 1.37$ |
| 9069 | [S III] | - | $5.80 \pm 0.59$ | $23.70 \pm 0.49$ | - | $5.16 \pm 0.27$ | $18.00 \pm 1.43$ | - | $6.59 \pm 0.64$ |
| 9531 | [S III] | - | $12.10 \pm 0.59$ | $55.50 \pm 0.58$ | - | $5.22 \pm 0.20$ | $38.20 \pm 0.78$ | $10.00 \pm 0.32$ | $15.60 \pm 0.64$ |
| 9824 | [CI] | - | $1.71 \pm 0.16$ | - | - | - | $2.41 \pm 0.76$ | $0.98 \pm 0.15$ | $0.82 \pm 0.09$ |
| 9850 | [CI] | - | $1.67 \pm 0.16$ | - | - | $2.49 \pm 0.17$ | $3.32 \pm 0.76$ | $2.58 \pm 0.15$ | $1.74 \pm 0.09$ |
| 10049 | $\mathrm{Pa} \delta$ | - | - | - | - | - | - | $3.85 \pm 0.44$ | - |
| 10122 | He II | - | - | - | - | - | - | $3.28 \pm 0.23$ | - |
| 10830 | He I | - | $7.72 \pm 1.08$ | - | - | $8.40 \pm 1.23$ | $25.60 \pm 1.52$ | $8.41 \pm 0.88$ | $9.47 \pm 0.77$ |
| 10938 | $\mathrm{Pa} \gamma$ | - | $3.98 \pm 1.08$ | - | - | $2.57 \pm 0.57$ | $8.07 \pm 1.20$ | $3.02 \pm 0.47$ | $4.08 \pm 0.55$ |
| 11470 | [P II] | - | $1.30 \pm 0.37$ | - | - | $1.17 \pm 0.14$ | - | $2.50 \pm 0.87$ | - |
| 11886 | [P II] | - | $1.80 \pm 0.37$ | $3.83 \pm 1.08$ | - | $1.46 \pm 0.14$ | - | $4.98 \pm 0.87$ | - |
| 12567 | [ Fe II] | - | $5.04 \pm 0.25$ | $4.64 \pm 0.52$ | $2.06 \pm 0.23$ | $4.73 \pm 0.33$ | $11.80 \pm 0.66$ | $6.80 \pm 0.21$ | $3.28 \pm 0.49$ |
| 12820 | $\mathrm{Pa} \beta$ | - | $12.10 \pm 0.26$ | $3.53 \pm 0.57$ | $4.86 \pm 0.19$ | $7.86 \pm 0.30$ | $9.57 \pm 2.61$ | $9.74 \pm 0.21$ | $9.20 \pm 0.51$ |
| 12950 | [ $\mathrm{Fe}_{\text {II] }}$ ] | - | $1.18 \pm 0.42$ | - | - | - | - | - | - |
| 13209 | [ Fe II] | - | $2.94 \pm 0.42$ | - | - | - | $5.60 \pm 0.34$ | - | - |
| 15342 | [Fe II] | - | - | - | - | - | - | - | - |
| 16436 | [ Fe II] | - | $4.65 \pm 0.47$ | $5.53 \pm 0.48$ | $3.24 \pm 0.08$ | $3.56 \pm 0.18$ | $9.20 \pm 0.28$ | $6.05 \pm 0.60$ | $3.09 \pm 0.21$ |
| 16773 | $[\mathrm{Fe} \mathrm{II}]+\mathrm{Br} 11$ | - | - | - | - | - | - | - | - |
| 17360 | Br10 | - | - | - | - | - | - | - | - |
| 18750 | Pa $\alpha$ | $4.70 \pm 0.18$ | $60.40 \pm 0.46$ | $82.50 \pm 3.49$ | $32.60 \pm 0.29$ | $34.30 \pm 0.17$ | $31.60 \pm 3.36$ | $36.70 \pm 0.67$ | $23.30 \pm 0.32$ |
| 19446 | Br $\delta$ | - | $2.12 \pm 0.39$ | - | $2.36 \pm 0.22$ | - | - | - | $1.05 \pm 0.10$ |
| 19570 | $\mathrm{H}_{2}$ | $1.79 \pm 0.49$ | $4.10 \pm 0.52$ | $3.36 \pm 0.54$ | $3.33 \pm 0.44$ | $5.25 \pm 0.05$ | $4.33 \pm 0.16$ | $8.17 \pm 0.26$ | $1.45 \pm 0.10$ |
| 20332 | $\mathrm{H}_{2}$ | $0.61 \pm 0.08$ | $1.48 \pm 0.30$ | $1.25 \pm 0.20$ | $0.86 \pm 0.05$ | $1.65 \pm 0.07$ | $2.36 \pm 0.39$ | $2.46 \pm 0.20$ | $2.03 \pm 0.20$ |
| 20580 | $\mathrm{H}_{2}$ | - | $1.40 \pm 0.26$ | - | $1.98 \pm 0.05$ | - | $1.33 \pm 0.34$ | $1.53 \pm 0.18$ | $1.24 \pm 0.20$ |
| 21218 | $\mathrm{H}_{2}$ | $0.60 \pm 0.02$ | $3.21 \pm 0.28$ | $2.14 \pm 0.20$ | $1.34 \pm 0.21$ | $4.00 \pm 0.13$ | $3.92 \pm 0.27$ | $4.80 \pm 0.36$ | $0.88 \pm 0.14$ |
| 21654 | $\mathrm{Br} \gamma$ | $0.80 \pm 0.04$ | $5.33 \pm 0.09$ | $0.52 \pm 0.32$ | $4.66 \pm 0.25$ | $2.88 \pm 0.13$ | $3.75 \pm 0.69$ | $4.07 \pm 0.05$ | $2.56 \pm 0.15$ |
| 22230 | $\mathrm{H}_{2}$ | - | $1.20 \pm 0.17$ | $1.05 \pm 0.23$ | $0.65 \pm 0.18$ | $1.70 \pm 0.85$ | $1.71 \pm 0.10$ | $2.25 \pm 0.25$ | $0.71 \pm 0.03$ |
| 22470 | $\mathrm{H}_{2}$ | - | $1.18 \pm 0.16$ | $0.61 \pm 0.13$ | - | $0.55 \pm 0.12$ | $0.49 \pm 0.06$ | $1.30 \pm 0.26$ | $0.32 \pm 0.14$ |

et al. 2007, 2008; Silva, Kuntschner \& Lyubenova 2008; Cesetti et al. 2009; Mármol-Queraltó et al. 2009; Riffel et al. 2011a, 2015; Kotilainen et al. 2012; Röck et al. 2017), and therefore it is very difficult to compare results from different investigations.

With this in mind, here we create a set of definitions for absorption features found in the NIR. We used two SSPs from the IRTFbased EMILES models (Röck 2015; Röck et al. 2016; Vazdekis et al. 2016), with 1.0 and 10 Gyr , solar metallicity, calculated with the PADOVA evolutionary tracks and with $\sigma=228 \mathrm{~km} \mathrm{~s}^{-1}$. We added up their light fractions (normalized to unity at $\lambda=12230 \AA$ ) as follows:
$F_{\text {comb }}=0.5 \frac{F_{\lambda}^{1 \mathrm{Gyr}}}{F_{\lambda=12230}^{1-\mathrm{Gyr}}}+0.5 \frac{F_{\lambda}^{10 \mathrm{Gyr}}}{F_{\lambda=12230}^{10 \mathrm{Gyr}}}$.
To this resulting spectrum we added Gaussians to model emissionlines profiles. These lines are located at the wavelengths of the most common emission lines detected in galaxies in this spectral region (see Section 3.1) with full width at half-maximums (FWHMs) characteristic of galaxies observed with SpeX with the configuration used here ( $25 \AA \lesssim$ FWHM $\lesssim 40 \AA$ ) with arbitrary flux values. We employed the ELPROFILE routine of the IFSCUBE package ${ }^{3}$ (Ruschel-

Dutra, in preparation). Using this simulated spectrum we defined the line limits and continuum band passes as illustrated in Fig. 3 and listed in Table 5.

We have measured the EWs for the most prominent absorption features using an updated PYTHON version of the PACCE code (Riffel \& Borges Vale 2011). In this code version, the EW uncertainties are assumed to be the standard deviation of 1000 EWs measurements of simulated spectra created by perturbing each flux point by its uncertainty through a Monte Carlo approach. The line definitions used are listed in Table 5, and the measured values are in Tables 6 and 7. In order to have a sample of early-type galaxies (ETGs) to compare our results with, we have collected NIR spectra from the literature and measured the EW of the absorption features with the same definitions used for our sample. Tables B1 and B2 present the measurements for the sample of galaxies presented in Baldwin et al. (2017). For four of the galaxies we were able to find Sloan Digital Sky Survey data used to measure the optical EW, while for the remaining objects we collected the values of $\mathrm{Fe} 5015, \mathrm{Mg}_{b}$ and Fe 5270 from McDermid et al. (2015). We also measured the values from the spectra presented by Dahmer-Hahn et al. (2018), which values are listed in Table B3.

[^3]Table 5. Line limits and continuum bandpasses.

| Centre (A) | Main absorber | Index name | Line limits <br> (A) | Blue continuum (A) | Red continuum (A) | Reference |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 4228.5 | CaI | Ca4227 | 4222.250-4234.750 | 4211.000-4219.750 | 4241.000-4251.000 | Worthey et al. (1994) |
| 4298.875 | CH ( $G$ band) | G4300 | 4281.375-4316.375 | 4266.375-4282.625 | 4318.875-4335.125 | Worthey et al. (1994) |
| 4394.75 | Fel | Fe4383 | 4369.125-4420.375 | 4359.125-4370.375 | 4442.875-4455.375 | Worthey et al. (1994) |
| 4463.375 | CaI | Ca4455 | 4452.125-4474.625 | 4445.875-4454.625 | 4477.125-4492.125 | Worthey et al. (1994) |
| 4536.75 | Fe I | Fe4531 | 4514.250-4559.250 | 4504.250-4514.250 | 4560.500-4579.250 | Worthey et al. (1994) |
| 4677.125 | $\mathrm{C}_{2}$ | Fe4668 | 4634.000-4720.250 | 4611.500-4630.250 | 4742.750-4756.500 | Worthey et al. (1994) |
| 5015.875 | FeI | Fe5015 | 4977.750-5054.000 | 4946.500-4977.750 | 5054.000-5065.250 | Worthey et al. (1994) |
| 5101.625 | MgH | Mg1 | 5069.125-5134.125 | 4895.125-4957.625 | 5301.125-5366.125 | Worthey et al. (1994) |
| 5175.375 | MgH | Mg2 | 5154.125-5196.625 | 4895.125-4957.625 | 5301.125-5366.125 | Worthey et al. (1994) |
| 5176.375 | Mg b | Mgb | 5160.125-5192.625 | 5142.625-5161.375 | 5191.375-5206.375 | Worthey et al. (1994) |
| 5265.65 | FeI | Fe5270 | 5245.650-5285.650 | 5233.150-5248.150 | 5285.650-5318.150 | Worthey et al. (1994) |
| 5332.125 | Fe I | Fe5335 | 5312.125-5352.125 | 5304.625-5315.875 | 5353.375-5363.375 | Worthey et al. (1994) |
| 5401.25 | Fel | Fe5406 | 5387.500-5415.000 | 5376.250-5387.500 | 5415.000-5425.000 | Worthey et al. (1994) |
| 5708.5 | Fer | Fe5709 | 5696.625-5720.375 | 5672.875-5696.625 | 5722.875-5736.625 | Worthey et al. (1994) |
| 5786.625 | FeI | Fe5782 | 5776.625-5796.625 | 5765.375-5775.375 | 5797.875-5811.625 | Worthey et al. (1994) |
| 5893.125 | Na I | NaD | 5876.875-5909.375 | 5860.625-5875.625 | 5922.125-5948.125 | Worthey et al. (1994) |
| 5965.375 | TiO | TiO1 | 5936.625-5994.125 | 5816.625-5849.125 | 6038.625-6103.625 | Worthey et al. (1994) |
| 6230.875 | TiO | TiO2 | 6189.625-6272.125 | 6066.625-6141.625 | 6372.625-6415.125 | Worthey et al. (1994) |
| 8498.0 | Ca II | CaT1 | 8476.000-8520.000 | 8110.000-8165.000 | 8786.000-8844.000 | Bica \& Alloin (1987) ( $\dagger$ ) |
| 8542.0 | Ca II | CaT2 | 8520.000-8564.000 | 8110.000-8165.000 | 8786.000-8844.000 | Bica \& Alloin (1987) ( $\dagger$ ) |
| 8670.0 | Ca II | CaT3 | 8640.000-8700.000 | 8110.000-8165.000 | 8786.000-8844.000 | Bica \& Alloin (1987) ( $\dagger$ ) |
| 9320.0 | ZrO/TiO/CN | ZrO | 9170.000-9470.000 | 8900.000-8960.000 | 9585.000-9615.000 | New Definition ( $\alpha$ ) |
| 10560.0 | VO | VO | 10470.000-10 650.000 | 10430.000-10 465.000 | 10660.000-10 700.000 | New Definition ( $\alpha$ ) |
| 11000.0 | CN | CN11 | 10910.000-11090.000 | 10705.000-10730.000 | 11310.000-11345.000 | New Definition ( $\beta$ ) |
| 11390.0 | NaI | NaI1.14 | 11350.000-11430.000 | 11310.000-11345.000 | 11450.000-11515.000 | New Definition ( $\beta$ ) |
| 11605.0 | FeI | FeI1.16 | $11580.000-11630.000$ | $11450.000-11515.000$ | $11650.000-11690.000$ | Roeck (2015) |
| 12430.0 | MgI | MgI1.24 | 12405.000-12455.000 | 12335.000-12365.000 | 12 465.000-12 490.000 | Roeck (2015) |
| 12944.0 | MnI | MnI1.29 | 12893.000-12995.000 | 12858.000-12878.000 | 13026.000-13 068.000 | New Definition |
| 13132.5 | Al I | AlI1.31 | 13 095.000-13170.000 | $13000.000-13070.000$ | $13175.000-13215.000$ | Roeck (2015) |
| 14875.0 | Mg I | MgI1.48 | 14850.000-14900.000 | 14750.000-14800.000 | 14910.000-14 950.000 | New Definition |
| 15032.5 | Mg I | MgI1.50 | 14995.000-15070.000 | 14910.000-14950.000 | $15150.000-15200.000$ | New Definition |
| 15587.5 | $\mathrm{CO}+\mathrm{Mg}$ I | CO1.5a | 15555.000-15620.000 | 15470.000-15 500.000 | 15700.000-15730.000 | New Definition ( $\epsilon$ ) |
| 15780.0 | $\mathrm{CO}+\mathrm{Mg}$ I | CO1.5b | 15750.000-15810.000 | 15700.000-15730.000 | 16095.000-16145.000 | New Definition ( $\epsilon$ ) |
| 15830.0 | Fel | FeI1.58 | 15 810.000-15 850.000 | $15700.000-15730.000$ | 16090.000-16140.000 | New Definition |
| 15890.0 | $\mathrm{Si} \mathrm{I}+\mathrm{Mg} \mathrm{I}$ | SiI1.58 | 15850.000-15930.000 | 15700.000-15730.000 | 16090.000-16140.000 | New Definition |
| 15985.0 | $\mathrm{CO}+\mathrm{SiI}$ | CO1.5c | 15950.000-16020.000 | $15700.000-15730.000$ | 16090.000-16 140.000 | New Definition ( $\epsilon$ ) |
| 16215.0 | $\mathrm{CO}+\mathrm{SiI}+\mathrm{CaI}$ | CO1.6a | 16145.000-16285.000 | 16090.000-16140.000 | 16290.000-16340.000 | New Definition ( $\epsilon$ ) |
| 17064.0 | $\mathrm{CO}+\mathrm{Fe} \mathrm{I}$ | CO1.6b | 17035.000-17093.000 | 16970.000-17 025.000 | $17140.000-17200.000$ | New Definition ( $\epsilon$ ) |
| 17111.5 | Mg I | MgI1.7 | 17093.000-17 130.000 | 16970.000-17 025.000 | 17 140.000-17 200.000 | Roeck (2015) |
| 22073.5 | NaI | NaI2.20 | 22 040.000-22 107.000 | $21910.000-21966.000$ | 22 125.000-22 160.000 | Frogel et al. (2001) |
| 22634.5 | CaI | CaI2.26 | $22577.000-22692.000$ | $22530.000-22560.000$ | $22700.000-22720.000$ | Frogel et al. (2001) ( $\gamma$ ) |
| 22820.0 | Mg I | MgI2.28 | $22795.000-22845.000$ | $22700.000-22720.000$ | $22850.000-22865.000$ | New Definition ( $\delta$ ) |
| 23015.0 | CO | CO2.2 | 22 870.000-23 160.000 | $22700.000-22790.000$ | $23655.000-23680.000$ | New Definition ( $\epsilon$ ) |
| 23290.0 | CO | CO2.3a | 23160.000-23 420.000 | $22700.000-22790.000$ | $23655.000-23680.000$ | New Definition ( $\epsilon$ ) |
| 23535.0 | CO | CO 2.3 b | 23 420.000-23 650.000 | $22700.000-22790.000$ | $23655.000-23680.000$ | New Definition ( $\epsilon$ ) |

Note. The optical indices are those of the LICK observatory (Worthey et al. 1994, and references). The CaT indices are those of Bica \& Alloin (1987) with a change in the blue continuum band passes in order to fit in our spectral region; $\alpha$ Based on Riffel et al. (2015) with small changes on the line limits; $\beta$ New continuum limits with central bandpasses from Roeck (2015); $\epsilon$ adapted from Riffel et al. (2007) with fixed continuum band passes, with better identifications of the main absorbers as well as better constraints of the line limits; $\gamma$ We made a small change on the blue continuum band pass to remove possible $\mathrm{H}_{2}$ emission lines.; $\delta$ Adapted from Silva et al. (2008) in order to better accomodate the continuum regions for the CO lines.

## 4 DISCUSSION

### 4.1 Emission lines

In order to compare the frequency of occurrence of the emission lines in our sample with what is seen in Seyfert galaxies, we show a histogram in Fig. 2 where the lines found here are compared to those of Riffel et al. (2006). What clearly emerges from this figure is that [S III], He I, and $\mathrm{Pa} \beta$ lines are less frequent in our sample (occurring in $\sim 60$ per cent of the sources) than in Seyferts (present in almost
all of the objects). On the other hand, we find a higher frequency of occurrence of lines of [C I] ( $\sim 65$ per cent), [P II] ( $\sim 40$ per cent), and [ Fe II] ( $\sim 65$ per cent) than in Sy 1 objects, and a similar rate as in Sy 2 s . The remaining emission lines occur with similar frequencies in the present sample and in Seyferts (see also Lamperti et al. 2017). Lines that are less frequent in the present sample compared to AGNs are located in regions with strong stellar features. Thus, it is possible that the absence of these features is because they are intrinsically weaker than in AGNs and/or diluted by the broad

Table 6. Absorption feature EWs (in $\AA$ ).

| Line | NGC 23 | NGC 520 | NGC 660 | NGC 1055 | NGC 1134 | NGC 1204 | NGC 1222 | NGC 1266 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Ca4227 | $0.36 \pm 0.06$ | - | $0.53 \pm 0.61$ | - | $1.14 \pm 0.16$ | - | - | - |
| G4300 | $1.56 \pm 0.19$ | - | $5.18 \pm 1.62$ | - | $2.01 \pm 0.56$ | - | - | - |
| Fe4383 | $1.67 \pm 0.22$ | - | $7.23 \pm 0.68$ | - | $3.84 \pm 1.05$ | - | - | $7.53 \pm 2.22$ |
| Ca4455 | $0.49 \pm 0.1$ | $3.22 \pm 0.47$ | $2.01 \pm 0.3$ | - | $0.28 \pm 0.22$ | - | - | - |
| Fe4531 | $2.26 \pm 0.16$ | - | - | - | $1.2 \pm 0.46$ | - | - | - |
| $C_{2} 4668$ | $3.88 \pm 0.2$ | - | - | - | $4.94 \pm 0.55$ | - | - | - |
| Fe5015 | - | - | - | - | - | - | - | - |
| $\mathrm{Mg}_{1}$ | $3.84 \pm 0.19$ | - | $5.26 \pm 0.3$ | $5.2 \pm 0.33$ | $4.06 \pm 0.31$ | $4.39 \pm 0.27$ | - | $3.27 \pm 0.74$ |
| $\mathrm{Mg}_{2}$ | $4.98 \pm 0.13$ | - | $6.46 \pm 0.21$ | $6.73 \pm 0.2$ | $5.36 \pm 0.18$ | $5.24 \pm 0.18$ | - | $4.14 \pm 0.37$ |
| $\mathrm{Mg}_{b}$ | $2.52 \pm 0.15$ | - | $2.71 \pm 0.25$ | $3.08 \pm 0.26$ | $2.9 \pm 0.27$ | $3.03 \pm 0.31$ | $0.85 \pm 0.19$ | $3.81 \pm 0.5$ |
| Fe5270 | $1.91 \pm 0.12$ | - | $2.53 \pm 0.35$ | $1.92 \pm 0.32$ | $2.44 \pm 0.2$ | $2.09 \pm 0.3$ | - | $2.43 \pm 0.45$ |
| Fe5335 | $1.73 \pm 0.12$ | - | $2.02 \pm 0.24$ | $1.64 \pm 0.34$ | $2.17 \pm 0.2$ | $1.48 \pm 0.33$ | $0.66 \pm 0.28$ | $2.38 \pm 0.67$ |
| Fe5406 | $0.95 \pm 0.04$ | - | $1.0 \pm 0.15$ | $0.57 \pm 0.27$ | $1.25 \pm 0.14$ | $0.68 \pm 0.3$ | $0.09 \pm 0.1$ | $2.33 \pm 0.32$ |
| Fe5709 | $0.64 \pm 0.04$ | - | $0.77 \pm 0.1$ | $0.72 \pm 0.13$ | $0.79 \pm 0.07$ | $0.75 \pm 0.14$ | $0.45 \pm 0.04$ | $0.74 \pm 0.4$ |
| Fe5782 | $0.41 \pm 0.04$ | - | $1.11 \pm 0.08$ | - | $0.68 \pm 0.06$ | $0.05 \pm 0.16$ | $0.68 \pm 0.18$ | $0.98 \pm 0.22$ |
| NaD | $4.25 \pm 0.08$ | $1.91 \pm 0.25$ | $5.07 \pm 0.19$ | $4.02 \pm 0.27$ | $4.71 \pm 0.12$ | $3.56 \pm 0.21$ | - | $6.37 \pm 0.23$ |
| $\mathrm{TiO}_{1}$ | $0.57 \pm 0.09$ | - | - | - | - | $0.51 \pm 0.24$ | - | - |
| $\mathrm{TiO}_{2}$ | $3.82 \pm 0.11$ | - | $5.79 \pm 0.2$ | $7.96 \pm 0.28$ | $6.13 \pm 0.2$ | $6.07 \pm 0.27$ | $3.87 \pm 0.19$ | $8.34 \pm 0.51$ |
| CaT1 | $4.13 \pm 0.11$ | - | - | - | $3.62 \pm 0.20$ | $1.16 \pm 0.36$ | $3.93 \pm 0.13$ | $5.74 \pm 1.14$ |
| CaT2 | $5.48 \pm 0.09$ | - | - | - | $7.51 \pm 0.17$ | $3.11 \pm 0.31$ | $5.46 \pm 0.13$ | $6.16 \pm 1.13$ |
| CaT3 | $3.22 \pm 0.16$ | - | - | - | $3.07 \pm 0.50$ | - | $2.97 \pm 0.17$ | - |
| ZrO | $16.76 \pm 0.27$ | - | - | - | $13.09 \pm 1.29$ | $15.60 \pm 1.78$ | $14.06 \pm 0.56$ | $6.70 \pm 2.53$ |
| VO | $0.05 \pm 0.31$ | - | $1.41 \pm 0.63$ | - | - | - | - | - |
| CN11 | $12.32 \pm 0.14$ | - | $3.76 \pm 0.28$ | - | $11.15 \pm 0.17$ | $6.48 \pm 0.31$ | $6.37 \pm 0.27$ | $12.41 \pm 0.92$ |
| NaI1.14 | $1.74 \pm 0.08$ | - | $2.43 \pm 0.18$ | - | $1.48 \pm 0.11$ | $1.34 \pm 0.18$ | $2.01 \pm 0.06$ | $3.46 \pm 0.44$ |
| FeI1.16 | $0.71 \pm 0.05$ | - | $0.44 \pm 0.07$ | - | $0.21 \pm 0.06$ | $0.71 \pm 0.07$ | $0.65 \pm 0.06$ | - |
| MgI1.24 | $0.99 \pm 0.06$ | - | $0.74 \pm 0.05$ | - | $1.68 \pm 0.10$ | $1.15 \pm 0.07$ | - | $0.57 \pm 0.15$ |
| MnI1.29 | $0.03 \pm 0.15$ | - | $0.28 \pm 0.14$ | - | $0.06 \pm 0.25$ | $0.80 \pm 0.35$ | $0.55 \pm 0.10$ | $3.32 \pm 0.14$ |
| AlI1.31 | $1.54 \pm 0.07$ | - | $1.93 \pm 0.57$ | - | $2.17 \pm 0.10$ | $2.16 \pm 0.34$ | $1.90 \pm 0.07$ | $3.04 \pm 0.15$ |
| MgI1.48 | $1.80 \pm 0.03$ | $2.94 \pm 0.17$ | $1.96 \pm 0.03$ | $1.14 \pm 0.25$ | $1.67 \pm 0.04$ | $1.16 \pm 0.06$ | $1.07 \pm 0.07$ | $1.15 \pm 0.11$ |
| MgI1.50 | $3.77 \pm 0.07$ | $3.28 \pm 0.30$ | $2.43 \pm 0.10$ | - | $4.35 \pm 0.09$ | $3.46 \pm 0.08$ | $2.44 \pm 0.08$ | $2.81 \pm 0.13$ |
| CO1.5a | $3.52 \pm 0.09$ | $6.51 \pm 0.23$ | $4.33 \pm 0.10$ | $5.12 \pm 0.36$ | $2.61 \pm 0.23$ | $3.18 \pm 0.13$ | $2.66 \pm 0.05$ | $5.24 \pm 0.15$ |
| CO1.5b | $4.26 \pm 0.11$ | $6.76 \pm 0.23$ | $4.94 \pm 0.10$ | $4.71 \pm 0.19$ | $4.44 \pm 0.19$ | $4.35 \pm 0.08$ | $2.50 \pm 0.05$ | $5.28 \pm 0.07$ |
| FeI1.58 | $1.50 \pm 0.06$ | $3.64 \pm 0.13$ | $2.11 \pm 0.06$ | $0.45 \pm 0.10$ | $0.89 \pm 0.11$ | $1.85 \pm 0.06$ | $1.01 \pm 0.04$ | $1.66 \pm 0.05$ |
| SiI1.58 | $3.66 \pm 0.10$ | $3.65 \pm 0.24$ | $4.03 \pm 0.13$ | $4.40 \pm 0.20$ | $3.77 \pm 0.18$ | $4.25 \pm 0.11$ | $3.00 \pm 0.08$ | $4.63 \pm 0.10$ |
| CO1.5c | $3.50 \pm 0.06$ | $3.39 \pm 0.20$ | $3.83 \pm 0.11$ | $3.16 \pm 0.17$ | $2.72 \pm 0.10$ | $4.07 \pm 0.11$ | $2.22 \pm 0.12$ | $4.49 \pm 0.12$ |
| CO1.6a | $4.97 \pm 0.11$ | $7.60 \pm 0.39$ | $6.70 \pm 0.15$ | $7.21 \pm 0.44$ | $5.43 \pm 0.16$ | $7.73 \pm 0.25$ | $4.29 \pm 0.29$ | $4.73 \pm 0.30$ |
| CO1.6b | $2.32 \pm 0.05$ | $3.21 \pm 0.19$ | $0.79 \pm 0.07$ | $1.11 \pm 0.62$ | $1.57 \pm 0.08$ | $1.34 \pm 0.13$ | $1.75 \pm 0.07$ | $2.92 \pm 0.22$ |
| MgI1.7 | $1.68 \pm 0.05$ | $0.76 \pm 0.21$ | $1.37 \pm 0.04$ | $0.94 \pm 0.45$ | $0.91 \pm 0.04$ | $1.35 \pm 0.08$ | $1.16 \pm 0.04$ | $0.23 \pm 0.21$ |
| NaI2.20 | $3.45 \pm 0.08$ | $2.51 \pm 0.10$ | $2.88 \pm 0.07$ | $4.26 \pm 0.26$ | $2.82 \pm 0.08$ | $3.24 \pm 0.04$ | $1.29 \pm 0.08$ | $3.31 \pm 0.12$ |
| CaI2.26 | $3.07 \pm 0.07$ | $4.07 \pm 0.12$ | $2.40 \pm 0.18$ | $3.87 \pm 0.39$ | $1.78 \pm 0.07$ | $2.21 \pm 0.14$ | $2.25 \pm 0.10$ | $2.15 \pm 0.08$ |
| MgI2.28 | $1.14 \pm 0.02$ | $1.20 \pm 0.09$ | $0.60 \pm 0.04$ | $5.57 \pm 0.10$ | $0.62 \pm 0.04$ | $0.37 \pm 0.04$ | $0.31 \pm 0.06$ | $0.03 \pm 0.01$ |
| CO2.2 | $20.26 \pm 0.56$ | $23.16 \pm 0.36$ | $12.23 \pm 0.39$ | $23.11 \pm 1.22$ | $19.33 \pm 0.41$ | $21.80 \pm 0.58$ | $18.03 \pm 0.61$ | $24.34 \pm 0.61$ |
| CO2.3a | $19.07 \pm 0.38$ | $26.71 \pm 0.35$ | $12.45 \pm 0.66$ | $22.69 \pm 0.80$ | $21.93 \pm 0.26$ | $22.02 \pm 0.43$ | $18.71 \pm 0.64$ | $24.94 \pm 0.40$ |
| CO2.3b | $21.53 \pm 0.37$ | $27.00 \pm 0.42$ | $12.50 \pm 0.78$ | $18.28 \pm 0.56$ | $24.03 \pm 0.13$ | $22.57 \pm 0.44$ | $13.99 \pm 0.85$ | $24.52 \pm 0.26$ |

absorption features that dominate the $z+J$ band. Note though that for three objects (NGC 1055, NGC 6835, and NGC 520, see Fig. 4), our spectral range excludes the [S III], He I, and [C I] emission lines. If present in these spectra, they would show up in $\sim 80$ per cent of our sample.

It is worth mentioning that the kinematics of the [ S III], $[\mathrm{Fe}$ II], and $\mathrm{H}_{2}$ lines as well as the excitation mechanisms of the [Fe II] and $\mathrm{H}_{2}$ lines of the galaxies of this sample were explored in Riffel et al. (2013b). However, the low-ionization forbidden lines of [C I] (i.e. $\lambda 9850 \AA$ ) and [P II] (i.e. $\lambda 11886 \AA$ ), also detected in our sample, were not yet analysed. Although the [PII] line is stronger compared to [Fe II] $\lambda 12570 \AA$ in Sy 2s (Riffel et al. 2006) than in the other types of galaxies, the detection of [PII] lines is surprising here. This is because at Solar metallicity, Phosphorus is about 1000 times less
abundant than Carbon (Ferguson et al. 1997) and 100 times less abundant than Iron (Oliva et al. 2001). Hence, if the P/C abundance is near to solar, the [PII] lines should not be present, unless other strong abundant elements are much more optically thick than they appear. A similar problem is found in some quasars for which broad absorption lines of PV $\lambda \lambda 1118,1128 \AA$ are detected and extreme abundances ratios for P/C are found (Hamann 1998; Hamann et al. 2001; Borguet et al. 2012). According to Oliva et al. (2001) for a solar $\mathrm{Fe} / \mathrm{P} \sim 100$ abundance ratio, one expects that $\frac{\left[\mathrm{Fe} \mathrm{II}^{2}\right]}{[\mathrm{P} I]}=50$, similar to what is expected for supernova remnants. The NIR [P II] emission lines may probably help to set some constraints on the abundance of Phosphorus in galaxies.

As discussed in Oliva et al. (2001) bright [ Fe II] lines can only be formed in regions where hydrogen is partially ionized. Such regions

Table 7. Absorption feature EWs (in $\AA$ ).

| Line | UGC2982 | NGC 1797 | NGC 6814 | NGC 6835 | UGC12150 | NGC 7465 | NGC 7591 | NGC 7678 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Ca4227 | - | $0.91 \pm 0.12$ | - | - | - | $0.47 \pm 0.1$ | - | - |
| G4300 | - | $1.65 \pm 0.44$ | - | - | - | $2.84 \pm 0.52$ | - | $0.78 \pm 0.34$ |
| Fe4383 | - | - | - | - | - | $2.77 \pm 0.27$ | - | $2.24 \pm 0.4$ |
| Ca4455 | - | - | - | - | - | $0.53 \pm 0.16$ | - | $1.11 \pm 0.12$ |
| Fe4531 | - | $1.93 \pm 0.37$ | - | - | - | $2.67 \pm 0.16$ | - | $1.48 \pm 0.32$ |
| $C_{2} 4668$ | - | - | - | - | - | - | - | $1.1 \pm 0.42$ |
| Fe5015 | - | - | - | - | - | - | - | - |
| $\mathrm{Mg}_{1}$ | - | $2.35 \pm 0.21$ | - | - | - | - | - | $3.61 \pm 0.33$ |
| $\mathrm{Mg}_{2}$ | - | $3.77 \pm 0.13$ | - | - | $4.28 \pm 0.26$ | $4.07 \pm 0.24$ | $5.14 \pm 0.36$ | $4.23 \pm 0.16$ |
| $\mathrm{Mg}_{b}$ | - | $2.12 \pm 0.24$ | - | - | $2.04 \pm 0.27$ | $2.56 \pm 0.23$ | $3.09 \pm 0.28$ | $1.96 \pm 0.23$ |
| Fe5270 | - | $1.32 \pm 0.19$ | - | - | $1.34 \pm 0.28$ | $1.78 \pm 0.17$ | $1.82 \pm 0.29$ | $1.31 \pm 0.16$ |
| Fe5335 | - | $1.22 \pm 0.13$ | - | - | $1.11 \pm 0.37$ | $1.85 \pm 0.15$ | $2.32 \pm 0.3$ | $1.5 \pm 0.29$ |
| Fe5406 | - | $0.87 \pm 0.11$ | - | - | $0.44 \pm 0.2$ | $0.78 \pm 0.15$ | $1.31 \pm 0.21$ | $1.04 \pm 0.12$ |
| Fe5709 | - | $0.35 \pm 0.06$ | - | - | $0.64 \pm 0.1$ | $0.7 \pm 0.06$ | $1.14 \pm 0.17$ | $0.78 \pm 0.11$ |
| Fe5782 | - | $0.31 \pm 0.06$ | - | - | - | $0.57 \pm 0.03$ | $0.61 \pm 0.14$ | $0.86 \pm 0.09$ |
| NaD | $4.78 \pm 0.37$ | $6.27 \pm 0.19$ | - | - | $7.47 \pm 0.52$ | $1.67 \pm 0.27$ | $3.86 \pm 0.19$ | $4.17 \pm 0.25$ |
| $\mathrm{TiO}_{1}$ | - | - | - | - | - | $1.4 \pm 0.16$ | - | - |
| $\mathrm{TiO}_{2}$ | - | $2.08 \pm 0.23$ | - | - | $4.88 \pm 0.52$ | $4.63 \pm 0.15$ | $3.4 \pm 0.25$ | $4.6 \pm 0.21$ |
| CaT1 | - | $1.13 \pm 0.39$ | - | - | $2.22 \pm 0.42$ | $5.17 \pm 0.17$ | $0.80 \pm 0.33$ | $2.73 \pm 0.26$ |
| CaT2 | - | $3.16 \pm 0.34$ | $0.36 \pm 0.19$ | - | $1.55 \pm 0.43$ | $6.28 \pm 0.14$ | $5.37 \pm 0.26$ | $4.29 \pm 0.24$ |
| CaT3 | - | - | $3.35 \pm 0.06$ | - | - | $5.93 \pm 0.33$ | $3.73 \pm 0.39$ | $5.90 \pm 0.26$ |
| ZrO | - | $15.11 \pm 2.04$ | $1.20 \pm 1.74$ | - | $12.38 \pm 0.98$ | $9.16 \pm 0.75$ | $14.76 \pm 1.61$ | $6.18 \pm 1.10$ |
| VO | $7.10 \pm 1.35$ | - | $3.89 \pm 0.43$ | - | $1.13 \pm 0.43$ | - | - | - |
| CN11 | $20.92 \pm 0.74$ | $6.49 \pm 0.33$ | - | - | $7.95 \pm 0.42$ | $6.91 \pm 0.37$ | $11.93 \pm 0.34$ | $5.08 \pm 0.72$ |
| NaI1.14 | - | $1.38 \pm 0.19$ | $3.99 \pm 0.13$ | $7.92 \pm 0.83$ | $1.75 \pm 0.15$ | $1.66 \pm 0.18$ | $1.17 \pm 0.22$ | $0.78 \pm 0.11$ |
| FeI1.16 | $0.17 \pm 0.20$ | $0.70 \pm 0.07$ | $0.53 \pm 0.06$ | $2.78 \pm 0.11$ | $0.68 \pm 0.08$ | $0.76 \pm 0.06$ | $1.25 \pm 0.06$ | - |
| MgI1. 24 | $1.40 \pm 0.06$ | $1.16 \pm 0.07$ | $0.76 \pm 0.05$ | $1.49 \pm 0.06$ | $0.71 \pm 0.15$ | $0.63 \pm 0.04$ | $0.67 \pm 0.02$ | $0.41 \pm 0.06$ |
| MnI1.29 | $3.11 \pm 0.62$ | $0.77 \pm 0.33$ | $10.00 \pm 0.43$ | $1.53 \pm 0.11$ | $1.74 \pm 0.14$ | $1.12 \pm 0.10$ | - | $2.18 \pm 0.20$ |
| AlI1.31 | $0.46 \pm 0.26$ | $2.15 \pm 0.35$ | $1.33 \pm 0.11$ | - | $3.16 \pm 0.34$ | $2.29 \pm 0.15$ | $2.89 \pm 0.16$ | $2.84 \pm 0.17$ |
| MgI1.48 | $0.57 \pm 0.10$ | $1.16 \pm 0.07$ | $0.94 \pm 0.03$ | $1.83 \pm 0.14$ | $0.80 \pm 0.10$ | $1.83 \pm 0.09$ | $1.38 \pm 0.05$ | $1.77 \pm 0.08$ |
| MgI1.50 | $2.41 \pm 0.10$ | $3.46 \pm 0.08$ | $2.07 \pm 0.07$ | $1.82 \pm 0.25$ | $2.88 \pm 0.18$ | $3.12 \pm 0.17$ | $3.11 \pm 0.14$ | $3.53 \pm 0.14$ |
| CO1.5a | $2.83 \pm 0.05$ | $3.21 \pm 0.10$ | $2.53 \pm 0.07$ | $3.06 \pm 0.07$ | $4.03 \pm 0.11$ | $3.89 \pm 0.07$ | $4.39 \pm 0.13$ | $2.98 \pm 0.09$ |
| CO1.5b | $3.46 \pm 0.04$ | $4.34 \pm 0.07$ | $3.10 \pm 0.05$ | $4.63 \pm 0.19$ | $5.34 \pm 0.13$ | $3.48 \pm 0.09$ | $4.55 \pm 0.10$ | $3.52 \pm 0.06$ |
| FeI1.58 | $0.73 \pm 0.03$ | $1.85 \pm 0.05$ | $0.77 \pm 0.03$ | $2.46 \pm 0.11$ | $1.80 \pm 0.09$ | $1.35 \pm 0.06$ | $1.81 \pm 0.07$ | $1.18 \pm 0.05$ |
| SiI1.58 | $2.81 \pm 0.07$ | $4.25 \pm 0.12$ | $2.57 \pm 0.07$ | $4.45 \pm 0.19$ | $4.58 \pm 0.17$ | $3.43 \pm 0.11$ | $3.71 \pm 0.15$ | $2.01 \pm 0.14$ |
| CO1.5c | $2.99 \pm 0.07$ | $4.09 \pm 0.10$ | $2.51 \pm 0.06$ | $3.46 \pm 0.18$ | $3.39 \pm 0.13$ | $2.93 \pm 0.13$ | $4.03 \pm 0.19$ | $2.70 \pm 0.16$ |
| CO1.6a | $6.02 \pm 0.19$ | $7.73 \pm 0.27$ | $4.26 \pm 0.20$ | $9.42 \pm 0.45$ | $7.24 \pm 0.18$ | $6.17 \pm 0.20$ | $7.73 \pm 0.26$ | $6.50 \pm 0.28$ |
| CO1.6b | $2.64 \pm 0.13$ | $1.36 \pm 0.12$ | $0.92 \pm 0.05$ | $2.07 \pm 0.28$ | - | $1.13 \pm 0.09$ | $1.66 \pm 0.10$ | - |
| MgI1.7 | $0.85 \pm 0.11$ | $1.35 \pm 0.09$ | $1.11 \pm 0.03$ | $2.21 \pm 0.29$ | $1.34 \pm 0.14$ | $1.82 \pm 0.05$ | $1.76 \pm 0.07$ | $1.73 \pm 0.04$ |
| NaI2.20 | $4.15 \pm 0.15$ | $3.24 \pm 0.05$ | $1.54 \pm 0.03$ | $3.41 \pm 0.04$ | $4.10 \pm 0.07$ | $2.70 \pm 0.04$ | $3.63 \pm 0.08$ | $2.88 \pm 0.07$ |
| CaI2.26 | $6.74 \pm 0.13$ | $2.21 \pm 0.16$ | $1.04 \pm 0.04$ | $2.43 \pm 0.04$ | $4.37 \pm 0.25$ | $1.87 \pm 0.09$ | $4.35 \pm 0.11$ | $2.49 \pm 0.13$ |
| MgI2.28 | $0.71 \pm 0.18$ | $0.36 \pm 0.04$ | $0.05 \pm 0.04$ | $1.05 \pm 0.04$ | - | $0.33 \pm 0.05$ | $1.58 \pm 0.05$ | $0.86 \pm 0.06$ |
| CO2.2 | $17.96 \pm 0.81$ | $21.54 \pm 0.56$ | $6.24 \pm 0.12$ | $22.26 \pm 0.54$ | $22.08 \pm 1.02$ | $16.14 \pm 0.67$ | $23.89 \pm 0.47$ | $17.53 \pm 0.83$ |
| CO2.3a | $22.66 \pm 1.19$ | $22.13 \pm 0.39$ | $2.96 \pm 0.16$ | $23.54 \pm 0.41$ | $22.84 \pm 0.95$ | $14.90 \pm 0.82$ | $26.34 \pm 0.40$ | $24.72 \pm 0.77$ |
| CO2.3b | $27.53 \pm 1.39$ | $22.58 \pm 0.44$ | $4.36 \pm 0.20$ | $24.65 \pm 0.41$ | $20.18 \pm 1.07$ | $14.97 \pm 1.47$ | $27.79 \pm 0.50$ | $20.27 \pm 0.89$ |

of hot, partially ionized gas can only be produced in an efficient way by shocks and/or photoionization by soft X-rays. According to these authors, $\left[\mathrm{Fe}_{\text {II }}\right] /\left[\mathrm{P}_{\text {II }}\right]$ can be used to distinguish between shocks (ratio $\gtrsim 20$ ) and photoionization (ratio $\lesssim 2$ ). In order to test this hypothesis, we plotted in Fig. $6[\mathrm{Fe}$ II $] /[\mathrm{PII}] \times[\mathrm{CI}] /[\mathrm{PII}]$ for our sample as well as the Seyfert galaxies of Riffel et al. (2006). As can be seen in this figure, there is a good correlation and no clear separation between the SFGs and the Seyferts, suggesting that the dominant excitation mechanism is the same for the three ions. Furthermore, due to the low values derived for the $\left[\mathrm{Fe}_{\text {II }}\right] /\left[\mathrm{P}_{\text {II }}\right]$ ratio, that excitation mechanism might be expected to be photoionization based on the arguments of Oliva et al. (2001). To test this, we have
computed photoionization models using CLOUDY/C17.01 ${ }^{4}$ (Ferland et al. 2017) updated with the next release of collisional strengths for [PII] (taken from Tayal 2004) as well as with new transition probabilities ${ }^{5}$ (private communication), however these models are not able to reproduce the observed line ratios, underestimating both (the models values for both ratios are nearly zero). This may be due to the fact that these lines are not in fact excited by photoionization, but mostly driven by shocks.

[^4]

Figure 2. Histogram showing statistics of the most common NIR emission lines.

### 4.2 Absorption features

It is crucial to be able to derive ages and chemical composition in order to understand the dominant underlying unresolved stellar content of galaxies (Röck et al. 2017). So far, the NIR is lacking a clear procedure based on absorption-line strengths. The obvious choice to do this kind of study is using stellar clusters as probes, instead of the use of more complex star-forming objects. However, while observations of the integrated spectra of stellar clusters in the optical region have been available for almost 30 yr (e.g. Bica 1988), in the NIR such observations are very difficult since the light emitted by the stars of the clusters in the NIR bands is dominated by a few very bright stellar phases making it difficult to get reliable integrated spectra of such objects in the NIR (e.g. Lyubenova et al. 2010; Riffel et al. 2011c).

In order to have a more homogeneous data set, in addition to the data set we present here representing complex SFHs of SFGs (Section 2), we collected spectra of nearby ETGs (which tend to have less complex SFHs than our sample) observed similarly as those in this work. Our final data set representing the older stellar population is composed of 12 ETGs selected in order to span a wide range of ages ( $1-15 \mathrm{Gyr}$ ) at approximately solar metallicity and observed by Baldwin et al. (2017) using Gemini/GNIRS in the cross-dispersed mode $(\sim 0.8-2.5 \mu \mathrm{~m} ; R \sim 1700 ; \sigma \sim 75$ $\mathrm{km} \mathrm{s}^{-1}$ ) plus six ETGs selected from the Calar Alto Legacy Integral Field Area Survey (CALIFA Sánchez et al. 2016) and observed by Dahmer-Hahn et al. (2018) using the TripleSpec spectrograph attached to the Astrophysical Research Consortium (ARC) $3.5-\mathrm{m}$ telescope ( $\sim 0.95-2.45 \mu \mathrm{~m} ; R \sim 2000 ; \sigma \sim 64 \mathrm{~km} \mathrm{~s}^{-1}$ ). In addition to these NIR spectra, we also collected, when available, the optical spectra of the sources. In the case of Baldwin et al. (2017) galaxies, the optical spectra were taken from the Sloan Digital Sky Survey (Ahn et al. 2014), while for the sample of Dahmer-Hahn et al. (2018) we took the data from the CALIFA Sánchez et al. (2016). The optical and NIR indices were measured by us using the definitions of Table 5 and are listed as online material in Tables B1, B2, and B3.

### 4.2.1 Previous NIR index-index correlations

Due to the lack of adequate data sets to test predictions of NIR data, compared to the optical (see Thomas, Maraston \& Bender 2003, for example), there are only a few studies trying to understand the behaviour of NIR $\times$ NIR indices. For instance, Mármol-Queraltó et al. (2009) studied a sample of ETGs and found a strong correlation between $\mathrm{C}_{2} 4668$ and NaI2.20 indices. In Fig. 7a we show the Mármol-Queraltó et al. (2009) measurements (open diamonds) and the literature compilation presented by Röck et al. (plus symbols, 2017) together with our data (squares). Even though we only measured both indices for four sources, this correlation seems to still hold for SFGs, which populate the lower left end of the correlation (Fig. 7a).

Using a similar approach, Cesetti et al. (2009) reported a trend of correlation of the optical $\mathrm{Mg}_{2}$ band with NIR indexes, such as NaI2.20, CaI2.26, and CO2.2 for ETGs. In Figs 7(b)-(d) we plotted our sample (filled squares), together with those of Cesetti et al. (2009, open diamonds) and Kotilainen et al. (2012, open triangles) for early-type sources. Additionally we also added the inactive spirals (octagons, LTG-K12) of Kotilainen et al. (2012). From Fig. 7d we have excluded the two Seyfert galaxies (NGC 660 and NGC 6814) since the CO band can be very diluted in these kind of sources (Riffel et al. 2009; Burtscher et al. 2015).

From Fig. 7 it is clear that the trend seems to hold for NaI2.20 $\times$ $\mathrm{Mg}_{2}$, while for $\mathrm{CaI} 2.26 \times \mathrm{Mg}_{2}$ there is no clear correlation, and in the case of $\mathrm{CO} 2.2 \times \mathrm{Mg}_{2}$ instead of a positive correlation there seems to be an inverse correlation. Additionally there seems to be a segregation between early- and late-type galaxies in this plot (panel d). This indicates that CO is enhanced in younger stellar populations, in agreement with the predictions of the Maraston (2005) models as shown in Riffel et al. (2007).

To help in the interpretation of these results, on these index-index diagrams we have overplotted the new optical-to-NIR IRTF-based stellar population synthesis models of the E-MILES team (Vazdekis et al. 2012, 2016; Röck et al. 2016). The models employed are those computed using the PADOVA isochrones (Girardi et al. 2000), with ages in the range $0.3 \mathrm{Gyr}<t<15.0 \mathrm{Gyr}$ and metallicities within $[\mathrm{Fe} / \mathrm{H}]=-0.40,[\mathrm{Fe} / \mathrm{H}]=0.00$, and $[\mathrm{Fe} / \mathrm{H}]=0.22$ with two different spectral resolutions ( $\sigma=60 \mathrm{~km} \mathrm{~s}^{-1}$ and $\sigma=228$ $\mathrm{km} \mathrm{s}^{-1}$, the shaded area represents the differences caused by $\sigma$ ). We also plotted TP-AGB heavy (see Zibetti et al. 2013, for a comparison between TP-AGB heavy and light models), Picklesbased models of Maraston \& Strömbäck (2011, M11 hereafter), which do have the same prescription than Maraston (2005) models but with a higher spectral resolution $(\mathrm{R}=500)$ than the 2005 models, therefore making them more suitable for our comparisons. However, it is important to have in mind that M11 models do have a poorer spectral resolution than our data, the effects on the indices strengths by degrading the resolution to M11 models is within the uncertainties of our measurements. These models are shown as open brown stars and are only available for solar metalicity. What emerges from this exercise is that the models in general are not able to predict the NIR indices and that there is a segregation between early- (open diamonds and plus markers) and late-type (filled squares and octagons) galaxies in these diagrams. The upper panels show significantly larger NaI2.20 index values than predicted by the models with standard initial mass function (IMF). Both the optical Mg - and C -dominated indices are stronger than the models for the most massive galaxies (i.e. the ones with the largest index values). In the case of the NaI2.20 index, Röck et al. (2017) concluded that for early-type sources the large values


Figure 3. Simulated spectrum showing the NIR indices definitions. The blue and red continuum band passes are in grey and line limits in red. Regions of strong (transmission $<20$ per cent) telluric absorption are shaded with an ' X ' pattern, while regions of moderate (transmission $<80$ per cent) telluric absorption are shaded with a line pattern. Emission lines and absorption features are labelled. See the text for details.


Figure 4. Near-infrared normalized spectra ordered according to their shapes from steeper (top) to flatter (bottom). The data were normalized at $20925 \AA$. Activity types are listed ( $\mathrm{S}=\mathrm{SFG}$ and $\mathrm{L}=\mathrm{LINER}$ ).


Figure 5. Same order as Fig. 4, but for the optical spectral range. The data were normalized at $5300 \AA$.


Figure 6. The correlation between the emission-line ratios of [C I] $\lambda 9850 \AA$ and $[\mathrm{Fe}$ II] $\lambda 12570 \AA$ relative to [P II] $\lambda 11886 \AA$.
obtained for this index are due to a combination of a bottom-heavy IMF and the $[\mathrm{Na} / \mathrm{Fe}]$ abundances. On the other hand, Alton et al. (2018) found that their sample of massive ETGs is consistent with having a Milky Way-like IMF, or at most a modestly bottom-heavy IMF, and suggested that their extreme abundance values for Na , in the cores of massive ETGs, may be explained by the metallicitydependent nucleosynthetic yield of Na .

The lower panels of Fig. 7 show that the ETGs are in better agreement with the predicted values. However, about half of our SFGs sample show stronger CaI2.26 and CO 2.2 values than predicted by the models. From these plots, we can also infer that the TP-AGB phase does not change substantially the CO index, once the solar metalicity M11 models are in agreement with the E-MILES
ones for the younger ages ( $t \lesssim 1 \mathrm{Gyr}$ ), with a large discrepancy for the older ages. Besides age, metallicity appears as an additional discriminator for the measured strengths of CO bands, with low metalicity $([\mathrm{Fe} / \mathrm{H}]=-0.40)$ and intermediate ages $(\sim 350 \mathrm{Myr})$ showing the largest values for the CO 2.2 index. This is in agreement with the previous findings of Kotilainen et al. (2012), who found that the evolved red stars completely dominate the NIR spectra, and that in this age range, the hot, young stars contribution to the EWs is virtually non-existent. So far, to fully access these younger stellar content of the galaxies it is necessary to fit the full spectrum, taking the continuum into account (see Baldwin et al. 2017; Dahmer-Hahn et al. 2018, for example). However, this is beyond the scope of this paper and will be the subject of a future investigation (Riffel et al., in preparation). On the other hand, the lower values of the CO index presented by the ETGs are also not explained by the models, with M11 models underestimating and E-MILES models overestimating the main locus occupied by these sources.

### 4.2.2 New index-index correlations

Because we measured a large set of lines for our sample, we have tried to find new correlations among the different absorption features by plotting all the EWs listed in Tables 6 and 7, as well as literature data (Tables B1 to B3) against each other. From these, we removed the correlations already discussed above (Fig. 7) as well as the optical $\times$ optical indices correlations since these are well studied. ${ }^{6}$ Since the CaT lines are correlated (e.g. Cenarro et al. 2001), we only used CaT2 in our search for correlations. The final set of optical

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Figure 7. Index-index diagrams. The filled boxes are from this work. The plus markers indicate early-type objects indices presented in Röck et al. (2017, panel a, R17). The open diamonds are indices measured in ETGs taken from Mármol-Queraltó et al. (panel a, ETG-MQ09 2009) and Cesetti et al. (panels b, c, and d, ETG-C09 2009). The triangles and octagons represent, respectively, the late and ETGs studied by Kotilainen et al. (2012, LTG-K12 and ETG-K12). The filled green stars represent our new measurements for the ETGs of Dahmer-Hahn et al. (2018, ETG-DH18) in all panels. Note that the literature data may have different definitions among themselves and with the measurements we present here. The open brown stars represent Maraston \& Strömbäck (2011) solar metalicity, Pickles-based models, with the size of the points scaling with ages (smallest points for 300 Myr and largest for 15 Gyr ). The shaded areas represent IRTF-based Emiles models (Röck 2015; Röck et al. 2016; Vazdekis et al. 2016) with red, grey, and blue indicating $[\mathrm{Fe} / \mathrm{H}]=-0.40,[\mathrm{Fe} / \mathrm{H}]=0.00$, and $[\mathrm{Fe} / \mathrm{H}]$ $=0.22$, respectively. The shaded area represents models with a spectral resolution of $\sigma=60 \mathrm{~km} \mathrm{~s}^{-1}$ (the lowest available) to $\sigma=228 \mathrm{~km} \mathrm{~s}^{-1}$. The age range used is between 0.3 and 15.0 Gyr , with arrows, triangles, diamonds, and pentagons representing $0.3,1,5$, and 10 Gyr , respectively. The E-MILES models with ages smaller than 1 Gyr should be taken with caution. For more details see the text.
versus NIR and NIR versus NIR indices correlations are shown in Figs 8 and 9, together with a linear regression using the orthogonal distance regression (ODR) method that takes errors both in the $x$ and $y$ variables into account (Boggs \& Rogers 1990). We note that when it was not possible to measure one of the indices used in the correlations, we have removed the galaxy from the plots and regression. In addition, we only considered the cases where both indices were measured at least for six sources. To help understand these plots we have overplotted the same model set as discussed above.

What emerges from Fig. 8 is that both model sets are able to predict well all the measured values for the optical indices. In the NIR, however, the models fail in their predictions, except for CO2.2 and ZrO , with E-MILES making better predictions of strengths
than M11, especially in the case of atomic absorption features. In addition, there is a clear separation of the ETGs and SFGs on the $\mathrm{G} 4300 \times \mathrm{MgI1.7}, \mathrm{G} 4300 \times \mathrm{NaI} 2.20, \mathrm{Fe} 4531 \times \mathrm{MgI} 1.7$, and $\mathrm{Mg}_{1} \times$ NaI 2.20 diagrams with ETGs in general showing higher values for both optical and NIR indices. A less evident separation of ETGs and SFGs is observed on the $\mathrm{G} 4300 \times \mathrm{CO} 2.2$ and $\mathrm{Mg}_{b} \times \mathrm{NaI} 2.20$ diagrams, while no separation is observed for the $\mathrm{Fe} 5782 \times \mathrm{ZrO}$ and $\mathrm{TiO}_{1} \times \mathrm{CO} 2.2$ diagrams.

The optical indices ( $\mathrm{G} 4300, \mathrm{Fe} 4531$, and $\mathrm{Mg}_{1}$ ) are not very sensitive to the $\alpha / \mathrm{Fe}$ ratio while G4300 is mainly sensitive to the C and O abundances (Thomas et al. 2003). This may indicate that the MgI1.7 and NaI2.20 indices are also sensitive to C and/or O abundances. This is also in agreement with the findings of Röck et al. (2017) who suggested that [C/Fe] enhancement might contribute to the values observed for NaI2.20 in ETGs. However, the good correlation of NaI with $\mathrm{Mg}_{b}$ may also indicate that this index is $\alpha / F e$ dependent, since $\mathrm{Mg}_{b}$ is sensitive to changes in the $\alpha / F e$ ratio (Thomas et al. 2003). The CO2.2 index values are well described by the model predictions for both SFGs and ETGs, with an agemetallicity dependence for the SFGs and no evidence of strong changes on their strengths caused by the amount of TP-AGB stars (see above). This is additionally supported by the $\mathrm{CO} 2.2 \times \mathrm{TiO}_{1}$ diagram, where M11 models, independent of age, do populate the locus filled by the ETGs, while E-MILES models do not reproduce the larger TiO and smallest CO strengths. The CO and $\mathrm{TiO}_{1}$ correlation is not unexpected since these absorptions depend on O being available. The models do show that ZrO is more metallicity dependent while $\mathrm{TiO}_{1}$ seems to be age dependent. In addition, some ETG show $\mathrm{TiO}_{1}$ values larger than the models (specially E-MILES models), which can be interpreted as an IMF effect (see La Barbera et al. 2013). In the case of the Mg -dominated indices (in the NIR and optical) the large values for these indices can be associated with the most massive ETGs, and can be explained by an $[\mathrm{Mg} / \mathrm{Fe}]$ enhancement (e.g. Worthey, Faber \& Gonzalez 1992; Martín-Navarro et al. 2018).

The correlations found from this exercise for the NIR indices are shown in Fig. 9. One particularly relevant correlation is $\mathrm{CO} 1.6 \mathrm{~b} \times$ CN11, as the CN11 index is believed to be heavily dominated by the AGB evolutionary phase and particularly by C stars (Maraston 2005). Almost 50 per cent of our SFG do show CN11 $\gtrsim 10 \AA$, with a mean value $\sim 20$ per cent larger than in ETG (see Fig. 13) and are consistent with the intermediate-age ( $0.3-2 \mathrm{Gyr}$ ) models. M11 models do cover better the space of values of the measurements, but all the older ages M11 models ( $\mathrm{t} \gtrsim 3 \mathrm{Gyr}$ ) do predict more or less constant values for CN11 (the same hapens for CO1.6b). The ETGs are more or less matched by SSP models with old ages and no indication of an intermediate-age population is required to explain the absorption features of these sources, once, their strengths in some cases are smaller than those of the older E-MILES SSPs. The CO 2.2 and $\mathrm{CO} 2.3 \mathrm{a}, \mathrm{b}$ (also CO 1.5 a and CO 1.5 b ) indices are to some extent described by the models, with larger values predicted for intermediate-age SSPs.The remaining strengths are not predicted by the models and no clear separation is found for SFGs and ETGs.

With the aim of understanding the behaviour of the NIR indices, we plotted them against the $[\mathrm{MgFe}]^{\prime}$ index of Thomas et al. (2003) defined as:
$[\mathrm{MgFe}]^{\prime} \equiv \sqrt{\mathrm{Mg} b(0.72 \times \mathrm{Fe} 5270+028 \times \mathrm{Fe} 5335)}$
which, for this sample with a small range in metallicity, is basically an age-indicator and is completely independent of the $\alpha / F e$ ratio. Assuming that the ETGs used here are objects with relatively normal old stellar populations, which is a valid assumption since their full


Figure 8. Index-index correlations. The black squares are the data points of this work. The green diamonds are from Baldwin et al. (2017) and filled green stars are from Dahmer-Hahn et al. (2018). The open brown stars represent Maraston \& Strömbäck (2011) solar metalicity, Pickles-based models, with the size of the points scaling with ages (smallest points for 300 Myr and largest for 15 Gyr ). The shaded areas represent IRTF-based EMILES models (Röck 2015; Röck et al. 2016; Vazdekis et al. 2016) with red, grey, and blue representing $[\mathrm{Fe} / \mathrm{H}]=-0.40,[\mathrm{Fe} / \mathrm{H}]=0.00$, and $[\mathrm{Fe} / \mathrm{H}]=0.22$, respectively. The shaded area represents models with a spectral resolution of $\sigma=60 \mathrm{~km} \mathrm{~s}^{-1}$ (the lowest available) to $\sigma=228 \mathrm{~km} \mathrm{~s}^{-1}$. The age range used is between 0.3 and 15.0 Gyr , with arrows, triangles, diamonds, and pentagons representing $0.3,1,5$, and 10 Gyr , respectively. The models with ages smaller than 1 Gyr should be taken with caution. For more details see the text.
spectra can be well fitted with SSP models (see Baldwin et al. 2017; Dahmer-Hahn et al. 2018, for details). This can also be seen in Figs 10 to 12, where the ETGs do in general show less scatter in the $[\mathrm{MgFe}]^{\prime}$ index than the SFGs. This indicates a more complex SFH for the latter, most likely with a strong contribution from intermediate ( $\sim 1 \mathrm{Gyr}$ ) stellar populations. In order to test the effect of a more complex SFH on the NIR strengths, we show in Fig. 13 histograms comparing the strength distributions between

SFG and ETG. Except for a few indices (ZrO, MgI1.48, MgI1.50, CO1.5a, FeI1.58, CO1.5c, MgI1.7, NaI2.20), the mean value for SFG is larger than that for ETG. This more complex SFH can also explain why the CN and CO bands are in general stronger for the SFGs than the ETGs. These bands are enhanced by the short-lived younger red giant branch (RGB) and TP-AGB stars (Maraston 2005; Riffel et al. 2007, 2015). According to Maraston (1998), these stars can be responsible for up to 70 per cent of the


Figure 9. Same as Fig. 8 but for $\mathrm{NIR} \times$ NIR indices.


Figure 10. Comparison of NIR indices with $[\mathrm{MgFe}]^{\prime}$. The models are the same as Fig. 8.


Figure 11. Comparison of NIR indices with $[\mathrm{MgFe}]^{\prime}$. The models are the same as Fig. 8.


Figure 12. Comparison of NIR indices with $[\mathrm{MgFe}]^{\prime}$. The models are the same as Fig. 8.
total flux in the NIR. However, for the case of the NaI2.20 index, Röck et al. (2017) constructed models using enhanced contribution from AGB stars and found that these stars have only a very limited effect on the model predictions and do not improve significantly the fit of the model NaI2.20 indices. They also show that small fractions ( 3 per cent) do have a similar impact on NaI2.20 than those with larger amounts of these stars. This result is consistent
with our findings that NaI2.20 index has a mean value $\sim 20$ per cent larger in ETG than in SFG.

In general, the NIR line strengths are not well reproduced by any set of models, suggesting that the SFH of the galaxies cannot be recovered when only using NIR indices. Our results are in agreement with the finding of Baldwin et al. (2017) who have studied the SFH of a sample of ETG by fitting different SSP models and found


Figure 13. Comparison of NIR indices strengths between SFG and ETG.
that the SFH vary dramatically among the different EPS models when fitting NIR data, with higher spectral resolution models producing more consistent results. They also found variations in ages in the NIR tend to be small, and largely encoded in the shape of the continuum. This was also noticed in Riffel et al. (2015) who suggested that TP-AGB stars contribute noticeably to a mean stacked NIR spectrum made up with mostly late-type galaxies hosting a low-luminosity AGN, from the Palomar survey (Mason et al. 2015). This result was obtained by fitting a mix of individual IRTF stars to the mean galaxy spectrum. Nevertheless, in this same work we have shown that other evolved stars (red giants, C-R, and EAGB stars) can reproduce most of the absorption features detected, without having to resort to stars in the TP-AGB phase.

## 5 FINAL REMARKS

We analysed long-slit spectra spanning optical to NIR wavelengths of 16 infrared-luminous star-forming galaxies with the aim of offering the community a set of emission and absorption feature measurements that can be used to test the predictions of the forthcoming generations of stellar population models. The optical and NIR spectra were obtained at WIRO and at SpeX/IRTF, respectively. In addition to these, we collected literature spectra of ETGs and performed the EW measurements using a new homogeneous set of continuum and band pass definitions. The main findings can be summarized as follows:
(i) All our sources display $\mathrm{H}_{2}$ emission, characteristic of the starforming nature of our sample. In the optical they clearly display $\mathrm{H}_{\mathrm{I}}$ emission lines. However, NGC 1055 and NGC 1134 show an NIR spectrum free of HI emission lines. We interpret this latter result as the result of the low sensitivity of the NIR detector in this wavelength interval, thus the expected $\mathrm{Br} \gamma$ fluxes are below the detection limit.
(ii) The continua are dominated by stellar absorption features. The most common features are due to CaI , $\mathrm{CaII}, \mathrm{Fe} \mathrm{I}, \mathrm{NaI}, \mathrm{Mg}$ I, plus prominent absorption bands of: $\mathrm{TiO}, \mathrm{VO}, \mathrm{ZrO}$, and CO . In most cases ( 70 per cent) the stellar continua also show evidence of dust extinction.
(iii) We present new definitions of continuum and line band passes for the NIR absorption lines. These definitions were made taking into account the position of the most common emission lines detected in this wavelength range.
(iv) We report EW measurements for 45 indices, including both optical and NIR features. We also present measurements for most of these indices in spectra of ETGs taken from literature. To the best of our knowledge, they represent the most complete set of EW measurements reported in the literature to date, and can be used to test the predictions of stellar population models from the optical to the NIR.
(v) We looked for correlations among the different absorption features, presenting as the most robust ones those with a Pearson correlation coefficient $\mathrm{r}>0.6$. In addition to the already known correlations in the optical region, we propose here correlations between optical and NIR indices, as well as correlations be-
tween different NIR indices, and compare them with model predictions.
(vi) While for the optical absorption features the new generation of models, with scaled-solar abundance ratios and standard IMF, share the same locus as the observed data points, they fail to predict the strengths of most of the NIR indices for the SFGs, while in the case of the early-type sources they roughly reproduce the observations. This may indicate more complex SFHs for the SFGs, which we interpreted as a strong contribution from the younger stellar populations, thus explaining the fact that the CN and CO bands are in general larger for the SFGs than the ETGs. These bands are enhanced in stars in the TP-AGB phase, however, they seems to have a limited impact on the indices of ETGs.

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## SUPPORTING INFORMATION

Supplementary data are available at $M N R A S$ online.
Appendix A. Final Reduced Spectra.
Appendix B. Literature Data.
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## APPENDIX A: FINAL REDUCED SPECTRA

Final reduced and redshift-corrected spectra for the remaining sample. Available as online material.

## APPENDIX B: LITERATURE DATA

Here we present the measurements using the index definitions listed in Table 5 for the literature data. The data used here are those of Dahmer-Hahn et al. (2018) and Baldwin et al. (2017). For the latter we found optical Sloan Digital Sky Survey data (Ahn et al. 2014) for four sources. For the remaining objects we collected the values of $\mathrm{Fe} 5015, \mathrm{Mg}_{b}$, and Fe 5270 from McDermid et al. (2015), while for the sources of Dahmer-Hahn et al. (2018) the optical data were taken from the Calar Alto Legacy Integral Field Area Survey (CALIFA Sánchez et al. 2016) and we measured the EW of the optical lines.

Table B1. Absorption feature EWs (in $\AA$ ) from the sample of Baldwin et al. (2017). The full table is avaliable as online material.

| Line | IC0719 | NGC 3032 | NGC 3098 | NGC 3156 | NGC 3182 | NGC 3301 |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: |
| Ca4227 | $1.68 \pm 0.01$ | - | - | $0.56 \pm 0.00$ | $1.41 \pm 0.00$ | - |
| G4300 | $6.23 \pm 0.04$ | - | - | $1.94 \pm 0.00$ | $5.57 \pm 0.01$ | - |
| Fe4383 | $5.99 \pm 0.05$ | - | - | $1.89 \pm 0.00$ | $5.33 \pm 0.01$ | - |
| Ca4455 | $2.07 \pm 0.02$ | - | - | $1.08 \pm 0.00$ | $1.68 \pm 0.00$ | - |
| Fe4531 | $3.67 \pm 0.04$ | - | - | $2.99 \pm 0.00$ | $3.37 \pm 0.01$ | - |

Note. The values of $\mathrm{Fe} 5015, \mathrm{Mg}_{b}$, and Fe5270 for NGC 3032, NGC 3098, and NGC 3301 were taken from McDermid et al. (2015) for Re/8.

Table B2. Absorption feature EWs (in $\AA$ ) from the sample of Baldwin et al. (2017). The full table is avaliable as online material.

| Line | NGC 3489 | NGC 4379 | NGC 4578 | NGC 4608 | NGC 4710 | NGC 5475 |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: |
| Ca4227 | - | - | - | - | - | $1.32 \pm 0.00$ |
| G4300 | - | - | - | - | - | $5.76 \pm 0.01$ |
| Fe4383 | - | - | - | - | $5.43 \pm 0.01$ |  |
| Ca4455 | - | - | - | - | $1.87 \pm 0.00$ |  |
| Fe4531 | - | - | - | - | $3.79 \pm 0.01$ |  |

Table B3. Absorption feature EWs (in $\AA$ ) from the sample of Dahmer-Hahn et al. (2018). The full table is avaliable as online material.

| Line | N4636 | N5905 | N5966 | N6081 | N6146 | N6338 | UGC08234 |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Ca4227 | - | - | $1.23 \pm 0.10$ | $1.21 \pm 0.07$ | $1.00 \pm 0.03$ | $1.16 \pm 0.06$ | $0.76 \pm 0.05$ |
| G4300 | - | - | $5.54 \pm 0.28$ | $5.59 \pm 0.30$ | $5.30 \pm 0.16$ | $5.82 \pm 0.23$ | $3.01 \pm 0.21$ |
| Fe4383 | - | - | $4.46 \pm 0.35$ | $4.56 \pm 0.29$ | $4.47 \pm 0.24$ | $4.33 \pm 0.29$ | $2.63 \pm 0.25$ |
| Ca4455 | - | - | $1.36 \pm 0.16$ | $1.14 \pm 0.15$ | $1.04 \pm 0.13$ | $1.30 \pm 0.17$ | $0.92 \pm 0.15$ |
| Fe4531 | - | - | $3.32 \pm 0.14$ | $3.04 \pm 0.21$ | $3.15 \pm 0.14$ | $3.24 \pm 0.12$ | $2.73 \pm 0.08$ |

Note. The optical data were taken from CALIFA survey (Sánchez et al. 2016). N means NGC.

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[^2]:    ${ }^{2}$ Emission lines and equivalent widths of the absorption features were measured on the spectra previous to normalization.

[^3]:    ${ }^{3}$ Available at: https://bitbucket.org/danielrd6/ifscube.git.

[^4]:    ${ }^{4}$ Available at: https://www.nublado.org.
    ${ }^{5}$ They are a combination of data taken from the MCHF/MCDHF data base at http://nlte.nist.gov/MCHF/ and data from the NIST Atomic Spectra Database at https://www.nist.gov/pml/atomic-spectra-database.

[^5]:    ${ }^{6}$ The NaI2.20 and CO2.2 are well studied, however, we decided to keep them here for diagrams distinct from those presented in Fig. 7 because correlations with other lines may help to shed some light in the understanding of the mechanisms driving these lines.

