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RESEARCH ARTICLE

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Quantization effects and convergence properties of rigid formation control systems with quantized distance measurements

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Summary

In this paper, we discuss quantization effects in rigid formation control systems when target formations are described by interagent distances. Because of practical sensing and measurement constraints, we consider in this paper distance measurements in their quantized forms. We show that under gradient-based formation control, in the case of uniform quantization, the distance errors converge locally to a bounded set whose size depends on the quantization error, while in the case of logarithmic quantization, all distance errors converge locally to zero. A special quantizer involving the signum function is then considered with which all agents can only measure coarse distances in terms of binary information. In this case, the formation converges locally to a target formation within a finite time. Lastly, we discuss the effect of asymmetric uniform quantization on rigid formation control.

KEYWORDS

binary measurement, quantization effect, rigid formation control

1 | INTRODUCTION

Quantized control has been an active research topic in the recent decade, motivated by the fact that digital sensors and numerous industrial controllers can only generate quantized measurements or feedback signals.^{1,2} Recent years have also witnessed extensive discussions on quantized control for networked control systems. This is because data exchange and transmission over networks often occurs in a digitally quantized manner, thus giving rise to coarse and imperfect information; see, eg, other works.³⁻⁸

In this paper, we aim to discuss the quantization effect on rigid formation control. Formation control based on graph rigidity is a typical networked control problem involving interagent measurements and cooperations. There have been

many papers in the literature focusing on control performance and convergence analysis for rigid formation control systems (see, eg, other works⁹⁻¹⁴), with virtually all assuming that all agents can acquire the relative position measurements to their neighbors perfectly. We remark that there are some recent works on linear consensus–based formation control with quantized measurements. An exemplary paper along this line of research is the work of Jafarian and Persis,¹⁵ which showed that, by using very coarse measurements (ie, measurements in terms of binary information), the formation stabilization task can still be achieved. The case of coarse measurements can be seen as a special (or extreme) quantizer, which generates quantized feedbacks in the form of binary signals. However, in the work of Jafarian and Persis¹⁵ and similar works on linear consensus–based formation control, a common knowledge of the global coordinate frame orientation is required for all the agents to implement the control law. This is a strict assumption and is not always desirable in practical formation control systems. Actually, it has been shown in the work of Meng et al¹⁶ that coordinate orientation mismatch may also cause undesired formation motions in linear consensus–based formation systems. All these restrictions and disadvantages are known to be avoidable in rigid formation control systems, in which any common knowledge of the global coordinate system is not required.

In the framework of quantized formation control, we also consider in the latter part of this paper a special quantizer described by the *signum* function. This part is motivated by the previous work of Liu et al,¹⁷ which discussed *triangular* formation control with coarse distance measurements involving the signum function. In this paper, we will consider a more general setting, which extends the discussions from the triangular case in the work of Liu et al¹⁷ to more general rigid formations.

The aim of this paper is to explore whether the introduction of quantized measurement and feedback can still guarantee the success of formation control, and to what extent the controller performance limits exist. Our broad conclusion is that quantization is not fatal, but may reduce performance in achieving a target formation. Generally speaking, quantized measurements are a type of approximation of actual measurements, and such approximations bring about bounded measurement errors that depend on quantizer functions. We remark that formation shape control with distance measurement errors or biases was discussed in other works. Measurement errors due to quantizations are different to measurement noises, in the sense that measurement errors induced by quantizations are deterministic, and some quantizers (especially logarithmic quantizers and binary quantizers) can also distinguish whether the quantity under quantization (distance error in the context of formation control) is zero. In this respect, measurements might be coarse from quantization, but the most important information (distance error being zero or not) is known without any noise. Furthermore, if quantization errors are not large and do not exceed the convergence region for a formation control system, then the local stability and convergence of formation systems are still guaranteed. For different types of quantizers applied in a distance measurement, we will present a detailed analysis to characterize the actual convergence and the boundedness of formation errors.

In this paper, we focus on local convergence of target formations with general shapes by assuming that initial formations are minimally and infinitesimally rigid, and are close to a target formation, which is a common assumption that has been widely used in many papers on rigid formation control; see, eg, other works. 9,12,14,19,22 We note that local stability is also a practical problem, arising when wind disturbs a formation away from its desired shape, and the original shape has to be recovered. Local convergence of formation shape also has practical significance. For example, agents can firstly assemble an approximate formation close to the desired shape, and then apply the control law to achieve the target formation guaranteed by the local convergence. Global analysis of formation convergence and stability is however only available for some particular and simple formation shapes (see, eg, the work of Dörfler and Francis¹⁰ for 2D triangular shape, the work of Park et al²³ for 3D tetrahedral shape, and the work of Chen et al²⁴ for 2D triangulated formations), while global analysis for rigid formation systems with general shapes is a challenging and open problem. Global analysis of formation convergence is therefore beyond the scope of this paper and will not be discussed here.

A preliminary version of this paper was presented in the work of Sun et al.²⁵ The extensions of this paper compared with the work of the aforementioned authors²⁵ include detailed proofs for all the key results, which were omitted in the work of Sun et al,²⁵ examples on several quantizer functions, and a new section to discuss the formation convergence when using an asymmetric uniform quantizer. For formation control with binary distance measurement, we also present several new properties for such formation control systems, including the independence of a global coordinate frame, boundedness of the control input, and a subspace-preserving principle, which will help analyze the transient behavior and implementation issue of such control laws. Furthermore, several simulation results, which support the theoretical analysis, are provided in this extended paper.

The remaining parts are organized as follows. Section 2 briefly reviews some background on graph rigidity and two commonly used quantizer functions. Section 3 discusses the convergence of the formation systems under two quantized formation controllers. In Section 4, we show a special quantized formation controller with binary distance information. Section 5 focuses on the case of an asymmetric uniform quantizer and its performance. Some illustrative examples are provided in Section 6. Section 7 concludes this paper.

2 | BACKGROUND AND PRELIMINARIES

2.1 | Notations

Most notations used in this paper are fairly standard, and here we introduce some special notations that will find use in later analysis. The operator $\operatorname{col}(\cdot)$ defines the stacked column vector. For a given matrix $A \in \mathbb{R}^{n \times m}$, define $\overline{A} := A \otimes I_d \in \mathbb{R}^{nd \times md}$, where the symbol \otimes denotes the Kronecker product and I_d is the d-dimensional identity matrix with $d = \{2, 3\}$. We denote by ||x|| the Euclidean norm of a vector x by $\hat{x} := \frac{x}{||x||}$ the unit vector of $x \neq 0$, and by $\tilde{x} := \frac{1}{||x||}$ the reciprocal of the norm of $x \neq 0$. For a stacked vector $x := [x_1^\top, x_2^\top, \dots, x_k^\top]^\top$ with $x_i \in \mathbb{R}^l, i \in \{1, \dots, k\}$, we define the block diagonal matrix $D_x := \operatorname{diag}\{x_i\}_{i \in \{1, \dots, k\}} \in \mathbb{R}^{kl \times k}$.

2.2 | Graph rigidity

Now, we review some basic graph theoretic tools for modeling formation control systems. For more background on graph theory in multiagent systems and formation control, we refer the reader to the work of Mesbahi and Egerstedt. Consider an undirected graph with m edges and n vertices, denoted by $G = (\mathcal{V}, \mathcal{E})$ with vertex set $\mathcal{V} = \{1, 2, ..., n\}$ and edge set $\mathcal{E} \subseteq \mathcal{V} \times \mathcal{V}$. The neighbor set \mathcal{N}_i of node i is defined as $\mathcal{N}_i := \{j \in \mathcal{V} : (i, j) \in \mathcal{E}\}$. We define an oriented incidence matrix $B \in \mathbb{R}^{n \times |\mathcal{E}|}$ for the undirected graph G by assigning an arbitrary orientation for each edge. Note that, for a rigid formation modeled by an undirected graph considered in this paper, the orientation of each edge for writing the incidence matrix can be defined arbitrarily and the stability analysis in the next sections remains unchanged. By doing this, we define the entries of B as $B_{ik} = -1$, if $B = \mathcal{E}_k^{\text{tail}}$, or $B_{ik} = +1$, if $B = \mathcal{E}_k^{\text{head}}$, or $B_{ik} = 0$ otherwise, where B_k^{tail} and B_k^{head} denote the tail and head nodes, respectively, of the edge B_k^{tail} , if B_k^{tail} , we have B_k^{tail} and B_k^{tail} are B_k^{tail} and B_k^{tail} and B_k^{tail} and B

We denote by $p = [p_1^T, p_2^T, \dots, p_n^T]^T \in \mathbb{R}^{dn}$ the stacked vector of all the agents' positions $p_i \in \mathbb{R}^d$. We also define *noncollocated* positions for all agents as those positions for which $p_i \neq p_j$ for all $(i, j) \in \mathcal{E}$. The pair (\mathcal{G}, p) is said to be a framework of \mathcal{G} in \mathbb{R}^d . The incidence matrix \mathcal{B} defines the sensing topology of the formation, ie, it encodes the set of available relative positions that can be measured by the agents. One can construct the stacked vector of available relative positions by

$$z = \overline{B}^{\mathsf{T}} p,\tag{1}$$

where each element $z_k \in \mathbb{R}^d$ in z is the relative position vector for the vertex pair defined by the edge \mathcal{E}_k .

Let us now briefly recall the notions of infinitesimally rigid framework and minimally rigid framework from the works of Hendrickson²⁷ and Anderson et al.²⁸ Let us define the edge function $f_G(p) := \mathop{\rm col}(\|z_k\|^2)$. We denote the Jacobian of $\frac{1}{2}f_G(p)$ by R(z), which is called the *rigidity matrix*. An easy calculation shows that $R(z) = D_z^\top \overline{B}^\top$. A framework (G,p) is *infinitesimally rigid* if $\mathop{\rm rank}(R(z)) = 2n - 3$ when it is embedded in \mathbb{R}^2 or if $\mathop{\rm rank}(R(z)) = 3n - 6$ when it is embedded in \mathbb{R}^3 . If additionally $|\mathcal{E}| = 2n - 3$ in the 2D case or $|\mathcal{E}| = 3n - 6$ in the 3D case, then the framework is called *minimally rigid*. In this paper, we assume that the target formation is infinitesimally and minimally rigid, while the convergence results obtained in this paper can be extended to nonminimally but still infinitesimally rigid target formations by following the analysis in the works of Sun et al¹⁴ or Mou et al.¹⁹

2.3 | Quantizer functions

In this paper, we mainly consider two types of quantizers, ie, the uniform quantizer and the logarithmic quantizer.⁴⁻⁸ In later sections, we will also consider two special quantizers, namely, a quantizer involving the signum function and the asymmetric uniform quantizer, derived from these quantizer functions.

2.3.1 | Definition of the quantizers

The *symmetric* uniform quantizer is a map $q_u: \mathbb{R} \to \mathbb{R}$ such that

$$q_u(x) = \delta_u \left(\left\lfloor \frac{x}{\delta_u} \right\rfloor \right),\tag{2}$$

where δ_u is a positive number and $\lfloor a \rceil$, $a \in \mathbb{R}$ denotes the nearest integer to a. We also define $\left\lfloor \frac{1}{2} + h \right\rfloor = h$ for any $h \in \mathbb{Z}$.

The logarithmic quantizer is an odd map $q_l: \mathbb{R} \to \mathbb{R}$ such that

$$q_{l}(x) = \begin{cases} \exp(q_{u}(\ln x)), & \text{when } x > 0; \\ 0, & \text{when } x = 0; \\ -\exp(q_{u}(\ln(-x))), & \text{when } x < 0, \end{cases}$$
 (3)

where $\exp(\cdot)$ denotes the exponential function.

2.3.2 | Properties of the quantizers

For the uniform quantizer, the quantization error is always bounded by $\delta_u/2$, namely, $|q_u(x)-x| \leq \frac{\delta_u}{2}$, for all $x \in \mathbb{R}$. For the logarithmic quantizer, it holds that $q_l(x)x \geq 0$, for all $x \in \mathbb{R}$, and the equality holds if and only if x = 0. The quantization error for the logarithmic quantizer is bounded by $|q_l(x)-x| \leq \delta_l|x|$, where the parameter δ_l is determined by $\delta_l = \exp\left(\frac{\delta_u}{2}\right) - 1$ (note that $\delta_l > 0$ because $\delta_u > 0$). The aforementioned definitions for scalar-valued uniform and logarithmic quantizers can be generalized to

The aforementioned definitions for scalar-valued uniform and logarithmic quantizers can be generalized to vector-valued quantizers for a vector in a component-wise manner. For an illustration of a logarithmic quantizer function, see Figure 1. Note that, in Section 5, we will further consider an *asymmetric* uniform quantizer and will provide some comparisons between a *symmetric* uniform quantizer and an *asymmetric* uniform quantizer.

2.4 | Nonsmooth analysis

Consider a differential equation

$$\dot{x}(t) = X(x(t)),\tag{4}$$

where $X: \mathbb{R}^d \to \mathbb{R}^d$ is a vector field, which is measurable but discontinuous. The existence of a continuously differentiable solution to (4) cannot be guaranteed due to the discontinuity of X(x(t)). In addition, as shown in the work of Ceragioli et al,⁵ the Caratheodory solutions (for definitions, see the work of Cortés²⁹) may not exist from a set of initial conditions of measure zero in quantized control systems. Therefore, we understand the solutions to the quantized rigid formation system in the sense of Filippov.³⁰ We first introduce the Filippov set-valued map.

Definition 1. Let $\mathcal{D}(\mathbb{R}^d)$ denote the collection of all subsets of \mathbb{R}^d . The Filippov set-valued map $F[X]: \mathbb{R}^d \to \mathcal{D}(\mathbb{R}^d)$ is defined by

$$\mathcal{F}[X](x) \triangleq \bigcap_{\delta > 0} \bigcap_{\mu(S) = 0} \overline{\operatorname{co}} \left\{ X \left(\mathbb{B}(x, \delta) \setminus \mathcal{S} \right) \right\}, \quad x \in \mathbb{R}^d,$$
 (5)

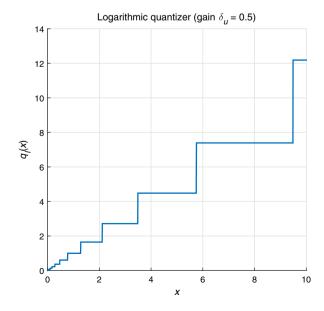
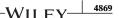


FIGURE 1 Logarithmic quantizer function with the gain $\delta_u = 0.5$, defined in (3) [Colour figure can be viewed at wileyonlinelibrary.com]



where $\overline{\text{co}}$ denotes convex closure, S is the set of x at which X(x) is discontinuous, $\mathbb{B}(x, \delta)$ is the open ball of radius δ centered at x, and $\bigcap_{u(S)=0}$ denotes the intersection over all sets S of Lebesgue measure zero.

Because of the way the Filippov set-valued map is defined, the value of $\mathcal{F}[X]$ at a discontinuous point x is independent of the value of the vector field X at x. Filippov solutions are absolutely continuous curves that satisfy almost everywhere the differential inclusion $\dot{x}(t) \in \mathcal{F}[X](x)$ defined. To keep this paper concise, we do not review further background knowledge on nonsmooth analysis and Filippov solutions. More properties of the Filippov solution and examples of how to compute a Filippov set-valued map can be found in the tutorial paper.²⁹

3 | FORMATION CONTROL WITH QUANTIZED MEASUREMENTS

3.1 | Quantized formation controllers

In rigid formation control, the target formation shape is described by a certain set of distances and all agents aim to stabilize cooperatively their distances to the desired ones. If one assumes perfect measurements, a commonly used formation controller can be written as (see, eg, the works of Garcia de Marina et al²² and Anderson et al³¹)

$$\dot{p}_i = -\sum_{k=1}^{|\mathcal{E}|} b_{ik} (\|z_k\| - d_k) \, \hat{z}_k, \tag{6}$$

where d_k is the desired distance for edge k, which is adjacent with agent i; $z_k := p_j - p_i$ is the relative position vector for edge k that associates agents i and j; $\hat{z}_k := \frac{z_k}{\|z_k\|}$ is the unit vector denoting the bearing information for edge k; and other symbols appeared in (6) have been introduced in Sections 2.1 and 2.2. It is clear from (6) that each agent needs to obtain relative position information to its neighbors. More precisely, agents in many, if not most, practical situations do not measure relative positions z_k directly but the distance $\|z_k\|$ and the direction \hat{z}_k as independent terms for the control action (6). For example, in ground and aerial robotic applications, interagent distances and bearings are measured by different kinds of sensors, eg, measuring the round trip time of a radio signal for observing distances, and vision cameras or directional antennas for observing bearings. In particular, as we will show in Lemma 2, the relative direction \hat{z}_k can be given with respect to an agent's local coordinate frame.

Since the interagent distances and bearings are so often observed by different sensors, it is reasonable to quantize independently the variables $||z_k||$ and \hat{z}_k in (6). Consequently, in the presence of quantized sensing, the right-hand side of the control action (6) needs modification. Here, we assume that the distance measurement is quantized, and the bearing measurement is unquantized. This assumption is not arbitrary but based on recent results in distance-based formation control. The work of Bishop et al³² shows that, in distance-based formation control, accuracy in bearing measurements are not critical. In particular, deviations of up to but less than ± 90 degrees with respect to the optimal direction, ie, the one linking two (or more) neighboring agents, do not compromise the stability and convergence of the formation. On the other hand, a distance-based formation is very sensitive to measurement errors, biases or inconsistencies of an interagent distance by its two adjacent agents as has been rigorously shown in the works of Mou et al. and Garcia de Marina et al. a Therefore, we focus our analysis on studying the quantization in sensors that measure the distances and not the bearings, ie, in this paper, we consider the control term with quantized distance measurements in the form as $q(||z_k|| - d_k)\hat{z}_k$. In practice, this choice of nonquantization of the bearing is also reasonable because the bearing measurement is always bounded (described by a unit vector, or by an angle in $[-\pi,\pi]$ in the 2D case). A normal digital sensor, say a 10-bit uniform quantizer, applying to bearing measurements gives rather accurate measurement with very small error to the true bearing, and as discussed in the work of Bishop et al,³² this is not an important issue regarding the stability of the formation. This is not the case for distance measurements, which may have larger magnitudes.

Remark 1. One may wonder why there is not use of the quantization feedback in the form of $q(||z_k||)$, ie, the direct quantized distance measurement, in the control (7) as follows. We note three reasons for choosing $q(||z_k|| - d_k)$ instead of $q(||z_k||)$, from the viewpoint of using quantization as a necessity (arising from limited measurement capabilities), and the advantages that are brought about by adopting such a quantization strategy.

• In rigid formation control, the control objective is to stabilize the actual distances between neighboring agents to prescribed values. If one chooses the quantization strategy in the form of $q(||z_k||)$, then this control objective may not be achieved. To this end, the quantization strategy $q(||z_k|| - d_k)$ used in (7) can be interpreted as arising from

a digital distance sensor with an embedded or prescribed offset (where the offset is the desired distance d_k), which is practical in real-world applications.

- In the case of nonuniform quantizers (eg, logarithmic quantizer), the quantization accuracy (or resolution) usually increases when the quantizer input approaches closer to the desired state (which is the origin in this case). Thus, when the formation approaches closer to the target formation, a higher quantization accuracy (if possible) is required, and this cannot be achieved if one uses the quantization function (eg, logarithmic quantization) on the actual distance in the form of $q(||z_k||)$. Furthermore, such a quantization is appealing as a design choice as one can have finite-time convergence (as proved in the later section), which is another advantage from the formation convergence viewpoint.
- We will further show in Section 4 that the chosen quantization strategy $q(||z_k|| d_k)$ will specialize to a simple and effective quantizer with *coarse binary* distance measurement. This also brings about the possibility of finite time formation convergence, as will be discussed in Section 4.

With the aforementioned considerations, we rewrite the control action (6) as

$$\dot{p}_i = -\sum_{k=1}^{|\mathcal{E}|} b_{ik} \, q \, (\|z_k\| - d_k) \, \hat{z}_k, \tag{7}$$

where q is a quantization function, which can be the uniform quantizer or the logarithmic quantizer defined in Section 2.3. We also assume that all the agents use the same quantizer $q(\cdot)$, and their initial positions start with noncollocated positions (which ensures $z_k(0) \neq 0$ for all k).

In the presence of quantized measurement and feedback, the right-hand side of (7) is discontinuous and we will consider the following differential inclusion:

$$\dot{p}_i \in \mathcal{F} \left[\sum_{k=1}^{|\mathcal{E}|} b_{ik} q \left(\| z_k \| - d_k \right) \hat{z}_k \right]. \tag{8}$$

In the following, we define the distance error for edge k as $e_k = ||z_k|| - d_k$. We then calculate the differential inclusion $\mathcal{F}(q(e_k))$, which will be used in later analysis. In the case of a symmetric uniform quantizer, the differential inclusion $\mathcal{F}(q_u(e_k))$ can be calculated as

$$\mathcal{F}(q_{u}(e_{k})) = \begin{cases} h\delta_{u}, & e_{k} \in \left(\left(h - \frac{1}{2}\right)\delta_{u}, \left(h + \frac{1}{2}\right)\delta_{u}\right), h \in \mathbb{Z}; \\ \left[h\delta_{u}, (h+1)\delta_{u}\right], & e_{k} = \left(h + \frac{1}{2}\right)\delta_{u}, h \in \mathbb{Z}. \end{cases}$$

Note that $e_k \mathcal{F}(q_u(e_k)) \ge 0$ for all e_k , and $e_k \mathcal{F}(q_u(e_k)) = 0$ if and only if $e_k \in \left[-\frac{\delta_u}{2}, \frac{\delta_u}{2}\right]$. We refer the reader to Section 2.3 for the definition and meaning of notations such as δ_u and h.

In the case of a logarithmic quantizer, the differential inclusion $\mathcal{F}(q_l(e_k))$ can be calculated as

$$\mathcal{F}(q_l(e_k)) = \begin{cases} \operatorname{sign}(e_k) \exp(q_u (\ln |e_k|)), & e_k \neq e^{\left(h + \frac{1}{2}\right)\delta_u}, h \in \mathbb{Z}; \\ \left[\exp(h\delta_u), \exp((h+1)\delta_u)\right], & e_k = e^{\left(h + \frac{1}{2}\right)\delta_u}, h \in \mathbb{Z}. \end{cases}$$

In addition, note that $e_k \mathcal{F}(q_l(e_k)) > 0$ for all $e_k \neq 0$, and $e_k \mathcal{F}(q_l(e_k)) = 0$ if and only if $e_k = 0$.

We define the distance error vector as $e = [e_1, e_2, \dots, e_m]^{\mathsf{T}}$. Then, in a compact form, one can rewrite the dynamics of (8) as

$$\dot{p} \in \mathcal{F}\left[-\overline{B}D_{\dot{z}}q\left(e\left(\operatorname{col}\left(\|z_{k}\|\right)\right)\right)\right]. \tag{9}$$

In order not to overload the notation, here, by \hat{z} , we exclusively mean the vector-wise normalization of z; therefore, $D_{\hat{z}}$ in the aforementioned equation and in the sequel is defined as $D_{\hat{z}} = \text{diag}\{\hat{z}_1, \dots, \hat{z}_m\}$. This notation rule will also be applied to \tilde{z} in the sequel. Note that the differential inclusion $\mathcal{F}(q(e))$ with the vector e is defined according to the product rule of Filippov's calculus properties (see the work of Paden and Sastry³³).

Example 1. We show an example to illustrate the formation control system with quantized measurements. Consider a formation system aiming to achieve a double tetrahedron shape in 3D space (see Figure 2), which consists of five agents labeled by 1, 2, 3, 4, 5, and nine edges. This formation is minimally rigid. For the purpose of writing an oriented

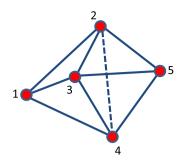


FIGURE 2 A 3D double tetrahedron formation shape, with five agents and nine distances [Colour figure can be viewed at wileyonlinelibrary.com]

incidence matrix, suppose that the nine edges are oriented from i to j just when i < j. Then, we can number the edges in the following order: 12, 13, 14, 23, 34, 24, 25, 35, and 45. Thus, the following oriented incidence matrix for the undirected graph in Figure 2 can be obtained:

$$B = \begin{bmatrix} -1 & -1 & -1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & -1 & 0 & -1 & -1 & 0 & 0 \\ 0 & 1 & 0 & 1 & -1 & 0 & 0 & -1 & 0 \\ 0 & 0 & 1 & 0 & 1 & 1 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 1 & 1 \end{bmatrix}.$$
 (10)

The composite relative position vector z is then defined according to the orientation for each edge, as in the incidence matrix B, in the form of $z = \overline{B}^T p$. As an example, one has $z_1 = p_2 - p_1$, ie, the vector z_1 at edge 1 is defined by the relative position between agent 2 and agent 1. The formation dynamics for agent 1 can be written as

$$\dot{p}_1 \in \mathcal{F} \left[q(\|z_1\| - d_1) \, \hat{z}_1 + q(\|z_2\| - d_2) \, \hat{z}_2 + q(\|z_3\| - d_3) \, \hat{z}_3 \right], \tag{11}$$

and similarly, one can obtain the system equation for other agents.

By defining the matrix \overline{B} and $D_{\hat{z}}$, one can obtain compact equations of system dynamics in the compact form of (9).

3.2 | Properties of quantized formation control systems

We first discuss the solution issue of the formation control system (8). However, it is more convenient to focus on the dynamics of the relative position vector z, which can be derived from (8) as follows:

$$\dot{z} = \overline{B}^{\mathsf{T}} \dot{p} \in \mathcal{F} \left[-\overline{B}^{\mathsf{T}} \overline{B} D_{\hat{z}} q \left(e \left(\operatorname{col}_{k} (\|z_{k}\|) \right) \right) \right]. \tag{12}$$

First, note that, at any noncollocated finite initial point p(0), the right-hand side of (8) and of (12) is measurable and locally essentially bounded. Thus, the existence of a local Filippov solution of (8) and of (12) starting at such initial points is guaranteed.

We then derive a dynamical system from (9) to describe the evolution of the distance error vector e. According to the definition of the distance error e_k , e_k is a smooth function of z_k . Thus, by using the calculus property (see the work of Cortés²⁹) and the set-valued Lie derivative computation theorem (see the work of Bacciotti and Ceragioli³⁴), one can show \dot{e}_k exists and $\dot{e}_k = \frac{1}{\|z_k\|} z_k^{\mathsf{T}} \dot{z}_k$ holds almost everywhere. The dynamics for the distance error vector e can be obtained in a compact form as

$$\dot{e} = D_{\tilde{z}} D_{z}^{\mathsf{T}} \dot{z} = D_{\tilde{z}} D_{z}^{\mathsf{T}} \overline{B}^{\mathsf{T}} \dot{p} = D_{\tilde{z}} R(z) \dot{p}, \quad \text{a. e.}$$

$$\in -\mathcal{F} \left[D_{\tilde{z}} R(z) R^{\mathsf{T}}(z) D_{\tilde{z}} q(e) \right], \quad \text{a. e.}$$
(13)

A more general compact form of the system equation \dot{e} can be found in section III in the work of Garcia de Marina et al.²² Again, the existence of a local Filippov solution of (13) starting with a noncollocated finite initial point p(0) is guaranteed. In the next section, we will also show that the solutions to (13) (and the solutions to (8) and (12) are *bounded* and can be extended to $t \to \infty$ when agents' initial positions are chosen noncollocated and close to a target formation shape. In addition, as shown in the work of Mou et al,¹⁹ when the formation shape is close to the desired one, the entries of the matrix $R(z)R^{T}(z)$ are continuously differentiable functions of e. Since the nonzero entries of the diagonal matrix $D_{\bar{z}}$ are of

the form $\frac{1}{\|z_k\|}$, which are also continuously differentiable functions of e, we conclude that the system described in (13) is a self-contained system, and we will call it *the distance error* system in the sequel.

Example 2. (Continued) We again use the formation shape shown in Figure 2 to illustrate the derivation of the aforementioned system equations.

According to the definition of the relative position vector z, one can derive the compact form of the system dynamic equation for z, as shown in (12). From the construction of the incidence matrix B and the relative position vector z, the rigidity matrix can be obtained as

$$R = \begin{bmatrix} -z_{1}^{\mathsf{T}} & z_{1}^{\mathsf{T}} & 0 & 0 & 0 \\ -z_{2}^{\mathsf{T}} & 0 & z_{2}^{\mathsf{T}} & 0 & 0 \\ -z_{3}^{\mathsf{T}} & 0 & 0 & z_{3}^{\mathsf{T}} & 0 \\ 0 & -z_{4}^{\mathsf{T}} & z_{4}^{\mathsf{T}} & 0 & 0 \\ 0 & 0 & -z_{5}^{\mathsf{T}} & z_{5}^{\mathsf{T}} & 0 \\ 0 & -z_{6}^{\mathsf{T}} & 0 & z_{6}^{\mathsf{T}} & 0 \\ 0 & -z_{7}^{\mathsf{T}} & 0 & 0 & z_{7}^{\mathsf{T}} \\ 0 & 0 & -z_{8}^{\mathsf{T}} & 0 & z_{8}^{\mathsf{T}} \\ 0 & 0 & 0 & -z_{9}^{\mathsf{T}} & z_{9}^{\mathsf{T}} \end{bmatrix}.$$

$$(14)$$

From the expression of the matrix R(z) in (14), it is obvious that the entries of the matrix product $R(z)R^{T}(z)$ are either zero, or inner products of relative position vectors in the form of $z_i^T z_j$, which are functions of the distance error vector e (for detailed analysis, see, eg, the works of Sun et al¹⁴ and Mou et al¹⁹). Since the diagonal matrix $D_{\bar{z}}$ is defined as $D_{\bar{z}} = \text{diag}\{\bar{z}_1, \ldots, \bar{z}_9\}$ with $\bar{z}_k = \frac{1}{\|z_k\|}$, it is clear that the entries of the matrix $D_{\bar{z}}$ are also functions of e, and therefore, the entries in the matrix product $D_{\bar{z}}R(z)R^T(z)D_{\bar{z}}$ are functions of e. Hence, the distance error system in (13) is a self-contained system, for which we can apply the Lyapunov argument to show its stability. Note the compact form of the error system (13) can be derived by the definition of the distance error e.

Finally, we show some additional properties of the formation control system with quantized information. Note that, through this paper, we assume that the underlying graph modeling interagent interactions is undirected.

Lemma 1. In the presence of the uniform/logarithmic quantizer, the formation centroid remains stationary.

Proof. Denote by $p_c \in \mathbb{R}^d$ the center of the mass of the formation, ie, $p_c = \frac{1}{n} \sum_{i=1}^n p_i = \frac{1}{n} (\mathbf{1}_n \otimes I_{d \times d})^{\mathsf{T}} p$. By applying the calculus property for the set-valued Lie derivative (see the works of Cortés²⁹ or Bacciotti and Ceragioli³⁴), one has

$$\dot{p}_{c}(t) = \frac{1}{n} (\mathbf{1}_{n} \otimes I_{d \times d})^{\mathsf{T}} \dot{p}$$

$$\in -\frac{1}{n} (\mathbf{1}_{n} \otimes I_{d \times d})^{\mathsf{T}} R^{\mathsf{T}}(z) D_{\tilde{z}} \mathcal{F} \left[q \left(e(z) \right) \right] \quad \text{for a.e. } t. \tag{15}$$

Note that $(\mathbf{1}_n \otimes I_{d \times d})^{\mathsf{T}} R^{\mathsf{T}}(z) = \mathbf{0}$. Therefore,

$$\dot{p}_c(t) \in -\frac{1}{n} (\mathbf{1}_n \otimes I_{d \times d})^\top R^\top (z) D_{\bar{z}} \mathcal{F} \left[q \left(e(z) \right) \right] = \{0\} \text{ for a.e. } t.$$
 (16)

Thus, $\dot{p}_c = 0$ for a.e. t, which indicates that the position of the formation centroid remains constant.

Lemma 2. To implement the control, each agent can use its own local coordinate system to measure the relative position (quantized distance and unquantized bearing) of its neighbors, and a global coordinate system is not involved.

Proof. The key part in the proof of local coordinate requirement is to show that the control function for all the agents is an SE(N)-invariant function.* The control function for agent i is $f_i = \sum_{k=1}^{|\mathcal{E}|} b_{ik} q(\|z_k\| - d_k) \hat{z}_k$. Given an arbitrary coordination rotation $R \in SO(N)$ and displacement of origin $\omega \in \mathbb{R}^N$, there holds $\|Rp_i + \omega - (Rp_j + \omega)\| = \|Rp_i - Rp_i\| = \|p_i - p_i\|$, and therefore, $q(\|z_k\| - d_k)$ is invariant under any action of rotation R and translation ω .

^{*}A function f is said to be SE(N)-invariant, if for all $R \in SO(N)$ and all $x_1, \ldots, x_n, \omega \in \mathbb{R}^N$, there holds $Rf(x_1, \ldots, x_n) = f(Rx_1 + \omega, \ldots, Rx_n + \omega)$.

WILFY 4873

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Thus, one has

$$f_{i}(Rp_{1} + \omega, \dots, Rp_{n} + \omega) = \sum_{k=1}^{|\mathcal{E}|} b_{ik} q(\|z_{k}\| - d_{k}) R\hat{z}_{k}$$

$$= R \sum_{k=1}^{|\mathcal{E}|} b_{ik} q(\|z_{k}\| - d_{k}) \hat{z}_{k}$$

$$= R f_{i}(p_{1}, \dots, p_{n})$$
(17)

and the statement is proved.

Note that the aforementioned lemma implies the SE(N) invariance (ie, translational and rotational invariance)³⁵ of the proposed formation controller, which enables a convenient implementation of the quantized formation control law without coordinate frame alignment for all of the agents. We refer the readers to the work of Vasile et al35 for a general treatment on coordinate frame issues in networked control systems.

3.3 | Convergence analysis

In this section, we aim to prove the following convergence result.

Theorem 1. Suppose the target formation is infinitesimally and minimally rigid and the formation controller with quantized measurement is applied.

- In the case of a uniform quantizer, the formation converges locally to an approximately correct and static shape defined by the set $F_{approx} = \left\{ e | e_k \in \left[-\frac{\delta_u}{2}, \frac{\delta_u}{2} \right], k \in \{1, \dots, |\mathcal{E}|\} \right\};$ • In the case of a logarithmic quantizer, the formation converges locally to a correct and static formation shape.

In the proof, we will use the Lyapunov theory of nonsmooth analysis, for which we construct a Lyapunov function candidate as

$$V(e) = \sum_{k=1}^{|\mathcal{E}|} V_k(e_k), \text{ with } V_k(e_k) = \int_0^{e_k} q(s) ds.$$
 (18)

Before giving the proof of Theorem 1, we first show some properties of the function V defined in (18). For the definition of function regularity in nonsmooth analysis, see, eg, chapter 2 in the work of Clarke³⁶ or the work of Cortés.^{29, p57}

Lemma 3. The function V constructed in (18) is positive semidefinite, and is regular everywhere.

Proof. The positive semidefiniteness of V is obvious from the property of the quantization functions q_u and q_l . Note that V(e)=0 if and only if $e\in\left\{e|e_k\in\left[-\frac{\delta_u}{2},\frac{\delta_u}{2}\right],k\in\{1,\ldots,|\mathcal{E}|\}\right\}$ for a uniform quantizer q_u , or when e=0 for a logarithmic quantizer q_l . The proof for the regularity is omitted here but follows similarly to the proof of lemma 6 in the previous paper of Liu et al. We note a key fact that supports the regularity statement of V, ie, V is continuously differentiable almost everywhere, while, at the nondifferentiable points, V has corners of convex type. From the sufficient condition of regular functions stated in the work of Clarke. 37 , 200 This key fact implies that V is regular everywhere.

Furthermore, according to the definition of generalized derivative (see, eg, chapter 2 in the work of Clarke³⁶), one can calculate the generalized derivative of V_k (for the case of a uniform quantizer) as

$$\partial V_k = \left\{ \begin{array}{ll} \left[h \delta_u, (h+1) \delta_u \right], & e_k = \left(h + \frac{1}{2} \right) \delta_u, h \in \mathbb{Z} \\ q(e_k), & \text{elsewise.} \end{array} \right.$$

Similarly, one can also calculate the generalized derivative of $V_k(e_k)$ for the case of a logarithmic quantizer (which is omitted here). The generalized derivative of V(e) can be obtained by the product rule (see the work of Cortés^{29, p50}). Now, we are ready to prove Theorem 1.

[†] "Roughly speaking, we can think of regular functions as those that, at each point, are either smooth, or else have a corner of convex type." 37, p200

Proof. We choose the Lyapunov function constructed in (18) for the distance error system (13) with discontinuous right-hand side. Note that $R(z)R^{\top}(z)$ and $D_{\bar{z}}$ are positive definite matrices at the desired formation shape. Similar to the analysis in the works of Sun et al¹⁴ and Garcia de Marina et al,²² we define a sublevel set $\mathcal{B}(\rho) = \{e : V(e) \leq \rho\}$ for some suitably small ρ such that, when $e \in \mathcal{B}(\rho)$, the formation is infinitesimally minimally rigid and the initial formation shape is close enough to the prescribed shape (which implies that interagent collisions cannot be possible). Note that all these imply that $R(z)R^{\top}(z)$ and $D_{\bar{z}}$ are *positive definite* when $e \in \mathcal{B}(\rho)$. Note also that the defined sublevel set $\mathcal{B}(\rho)$ is compact, and the matrix $Q(e) := D_{\bar{z}}R(z)R^{\top}(z)D_{\bar{z}}$ is also *positive definite* when $e \in \mathcal{B}(\rho)$. As a consequence, in the following, we rewrite the distance error system as $\dot{e} \in \mathcal{F}[-Q(e)q(e)]$.

The regularity of V shown in Lemma 3 allows us to employ the nonsmooth Lyapunov theorem (see section 2 in the work of Bacciotti and Ceragioli³⁴) to develop the stability analysis. We calculate the set-valued derivative of V along the trajectory of the distance error system (13). By applying the calculation rule for the set-valued derivative (see the work of Cortés²⁹, P^{62-63}), one can obtain

$$\dot{V}(e)_{(13)} \in \tilde{\mathcal{L}}_{(13)} V(e) = \left\{ a \in \mathbb{R} | \exists v \in \dot{e}_{(13)}, \text{ such that } \zeta^{\top} v = a, \forall \zeta \in \partial V(e) \right\}. \tag{19}$$

Note that the set $\tilde{\mathcal{L}}_{(13)}V(e)$ could be empty, and in this case, we adopt the convention that $\max(\emptyset) = -\infty$. When it is not empty, there exists $v \in -Q(e)q(e)$ such that $\zeta^\top v = a$ for all $\zeta \in \partial V(e)$. A natural choice of v is to set $v \in -Q(e)\zeta$, with which one can obtain $a = -q^\top(e)Q(e)q(e)$. Let $\bar{\lambda}_{\min}$ denote the smallest eigenvalue of Q(e) when e(p) is in the compact set B (ie, $\bar{\lambda}_{\min} = \min_{e \in B} \lambda(Q(e)) > 0$). Note that $\bar{\lambda}_{\min}$ exists because the set B is a compact set and the eigenvalues of a matrix are continuous functions of the matrix elements, and $\bar{\lambda}_{\min} > 0$ because Q(e) is positive definite with $e \in B(\rho)$ as aforementioned. Then, if the set $\tilde{\mathcal{L}}_{(13)}V(e)$ is not empty, one can show

$$\max\left(\tilde{\mathcal{L}}_{(13)}V(e)\right) \le -\bar{\lambda}_{\min}q(e)^{\mathsf{T}}q(e),\tag{20}$$

and if the set $\tilde{\mathcal{L}}_{(13)}V(e)$ is empty, one has $\max(\tilde{\mathcal{L}}_{(13)}V(e)) = -\infty$. Note that both cases imply that V is nonincreasing, and consequently, the Filippov solution e(t) of (13) is bounded. Thus, all solutions to (13) (and the solutions to (8) are bounded and can be extended to $t = \infty$ (ie, there is no finite escape time).

We now divide the rest of the proof into two parts according to different quantizers.

- The case of uniform quantizers: It can be seen that $\max(\tilde{\mathcal{L}}_{(13)}V(e)) \leq 0$ for all $e \in \mathcal{B}(\rho)$ and $0 \in \max(\tilde{\mathcal{L}}_{(13)}V(e))$ if and only if $e \in F_{\text{approx}}$. Moreover, note that F_{approx} is compact, and is positively invariant for the distance error system (13) (ie, if the initial formation is such that $e(0) \in F_{\text{approx}}$, then all agents are static and $e(t) \in F_{\text{approx}}$ for all t). According to the nonsmooth invariance principle (see theorem 3 in the work of Bacciotti and Ceragioli³⁴), the first part of the convergence result is proved. Since this is a convergence to a *closed and bounded* set F_{approx} (ie, a compact set), and outside this set the set-valued derivative of V along the trajectory of the distance error system is always negative (ie, $\max(\tilde{\mathcal{L}}_{(13)}V(e)) < 0$ for $e \in \mathcal{B}(\rho) \setminus F_{\text{approx}}$) while $\mathcal{B}(\rho)$ is also a compact set, the convergence to F_{approx} is achieved within a *finite time*. Note also from (7) the final formation is stationary because $\dot{p}(t) = 0$ for $e(t) \in F_{\text{approx}}$.
- The case of logarithmic quantizers: It can be seen that $\max(\tilde{\mathcal{L}}_{(13)}V(e)) \leq 0$ for all $e \in \mathcal{B}(\rho)$ and $0 \in \max(\tilde{\mathcal{L}}_{(13)}V(e))$ if and only if e = 0. According to the nonsmooth invariance principle (see theorem 3 in the work of Bacciotti and Ceragioli³⁴), the second part of the convergence result is proved. In addition, note from (7) that the final formation is stationary.

The proof is thus completed.

Remark 2. (**Finite-time convergence to a compact set**) In the aforementioned statement, we have shown the trajectories of distance errors in the formation system under uniform quantization converge to a bounded and closed set F_{approx} within a finite time, the size of which also depends on the uniform quantizer errors. The key recipes to guarantee the finite-time convergence are the following: (i) the set F_{approx} and the sublevel sets V(e) are compact sets, (ii) outside the set F_{approx} the set-valued derivative of V along the trajectories of the distance error system is always negative, and (iii) the function V(e) is a strictly increasing function of e. An intuitive illustration of the finite-time convergence of distance error trajectories to the set F_{approx} is shown in Figure 3.

Remark 3. We now show a stronger convergence result (ie, convergence to a point in the set) in addition to the finite-time convergence in the case of uniform quantizers (2). We observe that a sufficient condition for the position p_i of agent i to converge to a fixed point is that $\int_0^\infty \dot{p}_i(t) dt < \infty$, which is true since (i) initially all agents are at finite

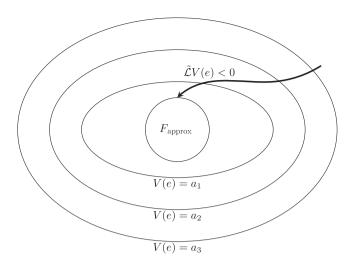


FIGURE 3 Illustration of finite-time convergence to a compact set $F_{\text{approx}(e)}$ centered at e = 0. Outside this set F_{approx} , the set-valued derivative of V along the trajectory of the distance error system is always negative. In the figure, three level sets of V(e) with $A_3 > A_2 > A_1$ are shown, which are compact sets with respect to e. Note, as also shown in the figure, V(e) is a positive definite and strictly increasing function of e

positions (ie, $p_i(0) < \infty$), and (ii) all $\dot{p}_i(t)$ (associated with the control input) are upper bounded and converge to the origin in *finite time*. By the integration law, this implies $p_i(t)_{t>T}$ is constant at a fixed position when $e(t)_{t>T} \in F_{\text{approx}}$, where T is the finite settling time of convergence, which further implies that the distance error e(t) converges to a fixed point in the set F_{approx} .

4 | A SPECIAL QUANTIZER: FORMATION CONTROL WITH BINARY DISTANCE INFORMATION

4.1 | Rigid formation control with coarse distance measurements

In this section, we consider the special case in which each agent uses very coarse distance measurements, in the sense that it only needs to detect whether the current distance to each of its neighbors is greater or smaller than the desired distance. This gives rise to a special quantizer defined by the following *signum* function:

$$\operatorname{sign}(x) = \begin{cases} 1, & \text{when } x > 0; \\ 0, & \text{when } x = 0; \\ -1, & \text{when } x < 0. \end{cases}$$

Accordingly, we obtain the following rigid formation control system with binary distance measurements:

$$\dot{p}_i = -\sum_{k=1}^{|\mathcal{E}|} b_{ik} \text{sign}(||z_k|| - d_k) \,\hat{z}_k. \tag{21}$$

Remark 4. Formation control using the signum function has been discussed in several previous papers. In the work of Zhao et al,³⁸ a finite-time convergence was established for stabilization of cyclic formations using binary bearing only measurements. A *linear consensus–based* formation control with coarsely quantized measurements was discussed in the work of Jafarian and Persis,¹⁵ while the implementation of the controller requires each agent to have knowledge of the global coordinate frame orientation. To be specific, the formation control law proposed in the work of the aforementioned authors¹⁵ is described by

$$u_i = \sum_{k=1}^{|\mathcal{E}|} b_{ik} \operatorname{sign} \left(z_k - z_k^* \right), \tag{22}$$

where $z_k^* \in \mathbb{R}^d$ is the desired relative position for edge k and the $sign(\cdot)$ function operates on each element of the d-dimensional vector $z_k - z_k^*$. In order to interpret the desired relative position vector z_k^* , all agents need to agree on a

common orientation in their coordinate frames. In contrast, the proposed control law with coarse distance measurements in the form of (21) does not require an orientation alignment of the agents' local coordinate frames. This is summarized in the following lemma.

Lemma 4. To implement the control of (21), each agent can use its own local coordinate system to measure the binary distances and bearings with respect to its neighbors, and a global coordinate system is not involved.

The proof of the aforementioned lemma is similar to the proof of Lemma 2 and is therefore omitted here.

Remark 5. The work closest to the controller setting in this section is the paper of Liu et al,¹⁷ which studied the stabilization control of a *cyclic triangular* formation modeled in a directed graph with the controller (21). Here, we extend such controllers to stabilize a general undirected formation, which is minimally and infinitesimally rigid. The aforementioned controller (21) can also be seen as a high-dimensional extension of the 1D formation controller studied in the work of De Persis et al.³⁹

Remark 6. (**Boundedness of the control input** (21)) Note that (since $b_{ij} \in \{-1,0,1\}$) the right-hand side of (21) is composed of the sum of a number of unit vectors multiplied by a signum function. This implies that the formation controller (21) is of special interest in practice since the control action is explicitly upper bounded by the cardinality of the set of neighbors for each agent i, which prevents potential implementation problems due to saturation. To be precise, the bound of the magnitude of the control input for agent i is derived by

$$||u_{i}|| = \left\| \sum_{k=1}^{|\mathcal{E}|} b_{ik} \operatorname{sign}(||z_{k}|| - d_{k}) \, \hat{z}_{k} \right\| \leq \sum_{k=1}^{|\mathcal{E}|} ||b_{ik} \operatorname{sign}(||z_{k}|| - d_{k}) \, \hat{z}_{k}||$$

$$\leq \sum_{k=1}^{|\mathcal{E}|} |b_{ik}| \, ||\operatorname{sign}(||z_{k}|| - d_{k})|| \, ||\hat{z}_{k}||$$

$$\leq \sum_{k=1}^{|\mathcal{E}|} |b_{ik}|$$

$$\leq |\mathcal{N}_{i}|. \tag{23}$$

Again, we consider the Filippov solution to the formation control system (21). The differential inclusion $\mathcal{F}(\text{sign}(e_k))$ can be calculated as

$$\mathcal{F}\left(\text{sign}(e_k)\right) = \begin{cases} 1 & ||z_k|| > d_k, \\ [-1, 1], & ||z_k|| = d_k, \\ -1 & ||z_k|| < d_k. \end{cases}$$

In a compact form, the rigid formation system (21) can be rewritten as

$$\dot{p} \in \mathcal{F}\left[-R^{\mathsf{T}}(z)D_{\bar{z}}\mathrm{sign}(e)\right],$$
 (24)

where sign(e) is defined in a component-wise way.

Note that the right-hand side of (24) is measurable and essentially bounded at any noncollocated and finite point p, and the existence of a local Filippov solution to (24) is guaranteed from such an initial point p(0). In the following analysis, we will also show that the solutions are bounded and complete.

We now show a subspace-preserving property for the formation control system with the control law (24).

Lemma 5. For the formation control system described by (24), the (affine) subspace spanned by agents' solutions is invariant over time, ie, the same as the (affine) subspace spanned by their initial positions. To be precise, there holds

$$\operatorname{span}([p_1(t), p_2(t), \dots, p_n(t)]) = \operatorname{span}([p_1(0), p_2(0), \dots, p_n(0)]), \quad \forall t \ge 0,$$
(25)

and

$$\operatorname{span}\left(\left[z_{1}(t), z_{2}(t), \dots, z_{|\mathcal{E}|}(t)\right]\right) = \operatorname{span}\left(\left[z_{1}(0), z_{2}(0), \dots, z_{|\mathcal{E}|}(0)\right]\right), \quad \forall t \ge 0.$$
(26)

Proof. In the work of Sun and Yu,⁴⁰ it has been proved that any networked dynamical system with scalar couplings in the following form:

$$\dot{x}_i(t) = \sum_{i=1}^n w_{ij}(t) x_j(t), \tag{27}$$

where $x_i \in \mathbb{R}^d$ and w_{ij} is a scalar (constant or time varying) coupling weight between agents j and i, possesses a subspace-preserving property, in the sense that $\text{span}([x_1(t), x_2(t), \dots, x_n(t)]) = \text{span}([x_1(0), x_2(0), \dots, x_n(0)]), \forall t \geq 0$ (see theorem 1 and corollary 1 in the work of Sun and Yu⁴⁰). Note that the formation system with the control law (24) can be equivalently written as

$$\dot{p}_i = \sum_{j \in \mathcal{N}_i} \frac{\text{sign}(\|p_j - p_i\| - d_{ji})}{\|p_j - p_i\|} (p_j - p_i), \tag{28}$$

where \mathcal{N}_i denotes agent i's neighbor set and d_{ji} denotes the desired distance between agents i and j. Denote $\mu_{ij} := \frac{\text{sign}(\|p_j - p_i\| - d_{ji})}{\|p_j - p_i\|}$. One can observe that

$$w_{ij} = \mu_{ij}, \text{ if } (i,j) \in \mathcal{E}; \text{ or } w_{ij} = 0, \text{ if } (i,j) \notin \mathcal{E}$$

$$w_{ii} = -\sum_{j \in \mathcal{N}_i} \mu_{ij}. \tag{29}$$

Therefore, the formation control system (24) can be rewritten in the form as $\dot{p}_i = \sum_{j=1}^n w_{ij}(t)p_j(t)$, which can be considered as a special case of the general coupled system described by (27). The subspace-preserving property then follows from the work of Sun and Yu.⁴⁰

Now, we define a weighted Laplacian matrix as $L_{\omega} = B\Omega B^{\mathsf{T}}$, with the state-dependent diagonal weight matrix Ω defined as $\Omega = \mathrm{diag}(\omega_1, \omega_2, \ldots, \omega_{|\mathcal{E}|})$. Then, a compact form of the formation position system can be obtained as

$$\dot{p} = -(L_{\omega} \otimes I_d)p. \tag{30}$$

Furthermore, from (1), there holds

$$\dot{z} = \overline{B}^{\mathsf{T}} \dot{p} = -\left((B^{\mathsf{T}} B \Omega B^{\mathsf{T}}) \otimes I_d \right) p$$

$$= -\left((B^{\mathsf{T}} B \Omega) \otimes I_d \right) z. \tag{31}$$

Then again, the affine-subspace-preserving property of the z system in the sense of (26) follows from theorem 1 in the work of Sun and Yu.⁴⁰

Remark 7. An intuitive interpretation of the (affine) subspace-preserving property shown in Lemma 5 is the invariance of collinear or coplanar positions for formation control systems in 2D/3D spaces. That is, for 2D formations, if all agents start with collinear positions, they will always be in that collinear subspace spanned by their initial positions under the general control law described by (24). Similarly, for 3D formations, if all the agents start with coplanar (respectively, collinear) positions, then their positions will always be coplanar (respectively, collinear) under the control (24). The paper of Liu et al¹⁷ presented a detailed analysis on such a collinear invariance principle for a triangular formation system under the control (24). In this sense, Lemma 5 presents a more general invariance result, which applies for any 2D/3D formation control systems with the control law (24).

Similar to the analysis in deriving the distance error system shown in Section 3.2, the distance error system with binary distance information can be obtained as

$$\dot{e} \in \mathcal{F}\left[-D_{\bar{z}}R(z)R^{\mathsf{T}}(z)D_{\bar{z}}\mathrm{sign}(e)\right], \quad \text{a. e.}$$
(32)

Again, similar to the analysis for (13), one can also show that (32) is a self-contained system when e takes values locally around the origin.

4.2 □ Convergence analysis

The main result in this section is stated in the following convergence theorem for the formation controller (24) with binary distance information.

Theorem 2. Suppose the target formation is infinitesimally and minimally rigid, the initial formation shape is close to the target formation shape, and the formation controller (21) with binary distance information is applied.

- The formation converges locally to a static target formation shape.

$$\bar{\lambda}_{\min} = \min_{e \in \mathcal{B}} \lambda\left(Q(e)\right) > 0,\tag{33}$$

where $Q(e) := D_{\bar{e}}R(z)R^{T}(z)D_{\bar{e}}$ and $B(\rho) = \{e : V(e) \le \rho\}$ is a sublevel set of some suitably small constant $\rho > 0$ such that, when $e \in \mathcal{B}(\rho)$, the formation is infinitesimally minimally rigid.

Proof. Choose the Lyapunov function defined as $V = \sum_{k=1}^{m} V_k(e_k)$ with $V_k(e_k) = |e_k|$ for the distance error system (32). Note that V is a convex and regular function of e. In addition, V is locally Lipschitz at e = 0 and is continuously differentiable at all other points. The generalized derivative of $V_k(e_k)$ can be calculated as

$$\partial V_k = \begin{cases} 1, & e_k > 0; \\ [-1, 1], & e_k = 0; \\ -1, & e_k < 0, \end{cases}$$

and the generalized derivative of V can be calculated similarly via the product rule (see the work of Cortés²⁹). We define a sublevel set $\mathcal{B}(\rho) = \{e : V(e) \le \rho\}$ for some suitably small ρ such that, when $e \in \mathcal{B}(\rho)$, the formation is infinitesimally minimally rigid and $R(z)R^{T}(z)$ and $D_{\bar{z}}$ are positive definite. Now, the matrix $Q(e) := D_{\bar{z}}R(z)R^{T}(z)D_{\bar{z}}$ is also positive definite when $e \in B(\rho)$. Let $\bar{\lambda}_{\min}$ denote the smallest eigenvalue of Q(e) when e(p) is in the compact set $\mathcal{B}(\text{ie}, \bar{\lambda}_{\min} = \min_{e \in \mathcal{B}} \lambda(Q(e)) > 0).$

In the following, we calculate the set-valued derivative of V along the trajectory described by the differential inclusion (32). The argument follows similarly to the analysis in the proof of Theorem 1. By applying the calculation rule for the set-valued derivative (see the work of Cortés^{29, pp62-63}), one can obtain

$$\dot{V}(e)_{(32)} \in \tilde{\mathcal{L}}_{(32)} V(e) = \left\{ a \in \mathbb{R} | \exists \nu \in \dot{e}_{(32)}, \text{ such that } \zeta^{\mathsf{T}} \nu = a, \forall \zeta \in \partial V(e) \right\}. \tag{34}$$

If the set $\tilde{\mathcal{L}}_{(32)}V(e)$ is not empty, there exists $v \in -Q(e)\operatorname{sign}(e)$ such that $\zeta^{\mathsf{T}}v = a$ for all $\zeta \in \partial V(e)$. A natural choice of v is to set $v \in -Q(e)\zeta$, with which one can obtain $a = -\operatorname{sign}^{\mathsf{T}}(e)Q(e)\operatorname{sign}(e)$. Then, one can further show

$$\max\left(\tilde{\mathcal{L}}_{(32)}V(e)\right) \le -\bar{\lambda}_{\min}\operatorname{sign}(e)^{\mathsf{T}}\operatorname{sign}(e),\tag{35}$$

if the set is not empty, while if it is empty, we adopt the convention $\max(\tilde{\mathcal{L}}_{(32)}V(e)) = -\infty$. Note that this implies that V is nonincreasing, and consequently, the Filippov solution e(t) is bounded. Thus, all solutions to (32) (and the solutions to (24) are complete and can be extended to $t = \infty$ (ie, there is no finite escape time). It can be seen that $\max(\tilde{\mathcal{L}}_{(32)}V(e)) \leq 0$ for all $e \in \mathcal{B}(\rho)$ and $0 \in \max(\tilde{\mathcal{L}}_{(32)}V(e))$ if and only if e = 0. According to the nonsmooth invariance principle (see theorem 3 in the work of Bacciotti and Ceragioli³⁴), the asymptotic convergence is proved.

We then prove the stronger convergence result, ie, the finite-time convergence. From the definition of the sign function in (21), there holds $sign(e)^{T}sign(e) > 1$ for any $e \neq 0$, which implies

$$\max\left(\tilde{\mathcal{L}}_{(32)}V(e)\right) \le -\bar{\lambda}_{\min} \tag{36}$$

for any $e \neq 0$. Thus, by applying the finite-time Lyapunov theorem, ⁴¹ any solution starting at $e(0) \in \mathcal{B}(\rho)$ reaches the origin in finite time, and the convergence time is upper bounded by $T^* = V(e(0))/\bar{\lambda}_{\min} = ||e(0)||_1/\bar{\lambda}_{\min}$.

Remark 8. (Finite-time formation convergence) Different to the finite-time convergence to an approximate formation shape under uniform quantizers as shown in Theorem 1, in Theorem 2, it is shown the formation system converges locally to a correct formation shape under binary distance measurements, which is a more desirable convergence result. Moreover, compared with the finite time formation controller discussed in the paper of Sun et al⁴² in which a sig function is used, the finite time formation controller in (21) requires less information in the distance measurements, in which very coarse measurements in terms of binary signals are sufficient.

Remark 9. (Effects on finite settling time) According to the formula of the convergence time associated with (33), the upper bound on the finite setting time is affected by the initial shape (in the form of 1-norm of the distance error vectors) and the least singular value of the matrix $R^{T}(z)D_{\bar{z}}$ when the formation shape is evolved in the set B.

Roughly speaking, when a shape is close enough to the target shape, the least singular value $\bar{\lambda}_{\min}$ could be approximated by $\lambda(Q(0))$, ie, the least singular value when the formation is at the desired shape. In this sense, a formation control system with a large $\lambda(Q(0))$ at the desired shape will generally have a shorter convergence time under the control law (21) when the initial formation shape is close to the target shape.

Remark 10. (**Dealing with chattering**) In the controller (21), the sign function is used, which may cause chattering of the solutions to the formation system when the formation is very close to the desired one (ie, when e is very close to the origin). This is because, in practice, imperfections (eg, perturbations in measurements or delays) could cause agents' state trajectories to "chatter" across the discontinuity surface (see, eg, chapter 3.5 in the work of Sastry⁴³). Possible solutions to eliminate the chattering include the following.

- Add deadzone (approximated by smooth functions) to the sign function around the origin (similar to the case of uniform quantizers; see Part 1 of Theorem 1). This will give rise to a trade-off in the convergence, ie, the distance error does not converge to the origin but to a bounded set, the size of which depends on (for a fixed number of agents) how large the deadzone parameter is chosen (see, eg, the works of Lee and Utkin⁴⁴ and Gupta and Ghosh⁴⁵).
- Use the hysteresis principle in the quantization function design.⁵

The adoption of the aforementioned techniques to avoid chattering will be discussed in future research.

5 | ASYMMETRIC UNIFORM QUANTIZER

In the work of Liu et al,⁶ it has been shown that when an *asymmetric* uniform quantizer (defined as follows) is applied to double-integrator consensus dynamics some undesirable motions may occur. In this section, we investigate whether there are undesired motions for rigid formation control in the presence of an asymmetric uniform quantizer.

We consider the following asymmetric uniform quantizer (the same as in the work of Liu et al⁶), defined by

$$q_u^*(x) = \delta_u \left(\left| \frac{x}{\delta_u} \right| \right), \tag{37}$$

where δ_u is a positive number and $\lfloor a \rfloor$, $a \in \mathbb{R}$ denotes the greatest integer that is less than or equal to a. For a comparison of the uniform quantizers defined in (2) and in (37), see Figure 4.

5.1 | Motivating example: two-agent formation case

We first consider a two-agent formation case. Suppose two agents are controlled to achieve an interagent distance of d_{12} with the quantization function (37). The system dynamics for agents 1 and 2 can be described, respectively, as

$$\dot{p}_1 = q_u^* (\|z_1\| - d_{12}) \, \hat{z}_1 \tag{38}$$

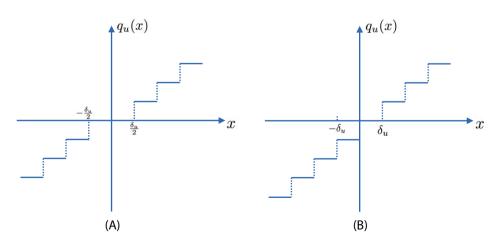


FIGURE 4 A, Symmetric uniform quantizer function, defined in (2); B, Asymmetric uniform quantizer function, defined in (37) [Colour figure can be viewed at wileyonlinelibrary.com]

and

$$\dot{p}_2 = -q_u^* (\|z_1\| - d_{12}) \, \hat{z}_1, \tag{39}$$

where $z_1 = p_2 - p_1$, and $q_u^*(\cdot)$ denotes the asymmetric uniform quantizer in (37).

Lemma 6. Consider the two-agent formation control system (38) and (39) with the asymmetric quantization function (37).

- If the initial distance between agents 1 and 2 is greater than $d_{12} + \delta_u$, then the interagent distance ||z|| will converge to $d_{12} + \delta_u$ and the final formation will be stationary.
- If the initial distance between agents 1 and 2 is smaller than the desired distance d_{12} , then the interagent distance ||z|| will converge to the desired distance d_{12} and the final formation will be stationary.
- If the initial distance between agents 1 and 2 is between d_{12} and $d_{12} + \delta_u$, then both agents 1 and 2 remain stationary and the interagent distance ||z|| remains unchanged.

The proof is obvious and is omitted here as it can be inferred from previous proofs.

Remark 11. In the aforementioned example, it can be seen that, in the case of an asymmetric uniform quantizer, there exist no undesired motions, which is different to the result observed in the work of Liu et al,⁶ which showed unbounded velocities. Apart from the difference in system dynamics under discussions, the key difference that leads to the distinct behaviors is that, when the asymmetric quantizer is applied to the consensus dynamics (which is to quantize a vector), there holds $\mathcal{F}[q_u^*(r_i - r_j)] + \mathcal{F}[q_u^*(r_j - r_i)] = -\delta_u$ when $r_j - r_i \neq k\delta_u$, where $r_j - r_i$ denotes the relative position vector (see section 5 in the work of Liu et al⁶). Note that, in the aforementioned formation controller, the quantization applies only to the distance error term (ie, $q_u^*(||p_2 - p_1|| - d_{12}))$), which is a scalar, and the asymmetric property of the quantizer only affects the convergence of the distance term.

5.2 | General formation case

We consider the general formation case with more than two agents, in which each agent employs asymmetric uniform quantizers in individual controllers.

Theorem 3. Suppose each individual agent takes the asymmetric uniform quantizer (37) in the quantized formation controller (7). Then, the interagent distances converge within a finite time to the set

$$F_{\text{avm}} = \{e | e_k \in [0, \delta_u], k \in \{1, \dots, |\mathcal{E}|\}\}.$$

The proof is omitted here as it can be directly inferred from the previous proof of Theorem 1.

6 | ILLUSTRATIVE EXAMPLES AND SIMULATIONS

In this section, we show several numerical examples to illustrate the theoretical results obtained in previous sections. In the following illustrative examples, we consider the stabilization control of a five-agent minimally rigid formation in the 3D space as a continuation of Example 1. The underlying graph describes a double tetrahedron shape of nine edges (see Figure 2 for an illustration), and the desired distances for all edges are set as six.[‡] The initial positions are chosen such that the initial formation is infinitesimally rigid and is close to a target formation shape. For all simulations, we set the quantization gain as $\delta_u = 0.5$.

Agents trajectories, the final formation shape, and the evolutions of nine distance errors under symmetric uniform quantization and under logarithmic quantization are shown in Figures 5 and 6, respectively. It is obvious from these two figures that, with symmetric uniform quantizer, the formation errors converge to the bounded set $F_{\text{approx}} = \{e | e_k \in [-0.25, 0.25], k \in \{1, ..., m\}\}$ in a *finite time*, and with the logarithmic quantizer, the formation converges to the target shape asymptotically, which are consistent with the theoretical results in Theorem 1.

The formation convergence behavior with binary distance measurements under the quantization strategy (21) is depicted in Figure 7. It can be seen from Figure 7 that, with very coarsely quantized distance measurement via a simple

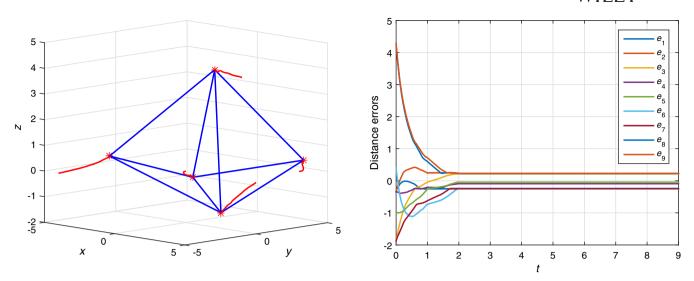


FIGURE 5 Stabilization control of a rigid formation: symmetric uniform quantization case. Left: the trajectories of five agents and the final formation shape. Right: time evolutions of the distance errors. It is obvious from the right figure that the formation shape converges to an approximately correct one in a finite time [Colour figure can be viewed at wileyonlinelibrary.com]

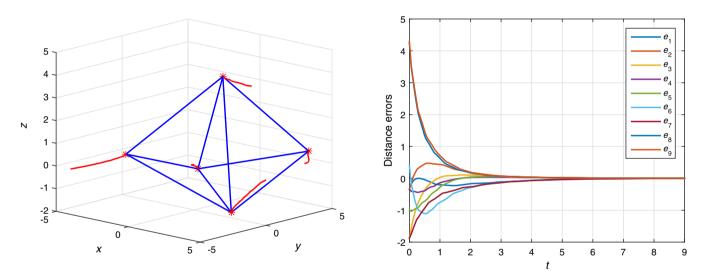


FIGURE 6 Stabilization control of a rigid formation: logarithmic quantization case. Left: the trajectories of five agents and the final formation shape. Right: time evolutions of the distance errors [Colour figure can be viewed at wileyonlinelibrary.com]

signum function as in (21), the formation converges to the target shape within a finite time, but the price to be paid is the occurrence of chattering (as shown in the right part of Figure 7).

Finally, when the asymmetric uniform quantizer (37) is used in the formation control system (7), the formation converges to an approximate one with all distance errors converging to the bounded set $F_{\text{aym}} = \{e | e_k \in [0, 0.5], k \in \{1, \dots, |\mathcal{E}|\}\}$ within a finite time, as shown in Figure 8. This supports the conclusion of Theorem 3.

For a comparison of formation convergences with different quantization functions, we plot the trajectories of the Lyapunov functions (which are functions of distance errors) for the formation system under the aforementioned four different quantization strategies, as shown in Figure 9. In the simulations, the initial conditions were chosen to be the same for the four cases. It can be observed from Figure 9 that, under symmetric uniform quantization, binary quantization, and asymmetric uniform quantization, the convergences of Lyapunov functions (as functions of distance errors) are achieved within a finite time (though the final values depend on different quantization functions), and that under logarithmic quantization, the distance errors converge to zero asymptotically. All these results are consistent with the aforementioned simulations and support the theoretical developments in the previous sections.

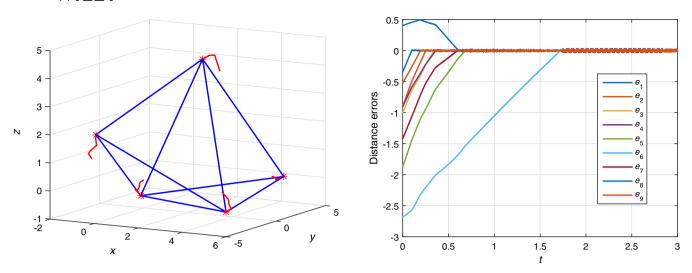


FIGURE 7 Stabilization control of a rigid formation: binary measurement case. Left: the trajectories of five agents and the final formation shape. Right: time evolutions of the distance errors [Colour figure can be viewed at wileyonlinelibrary.com]

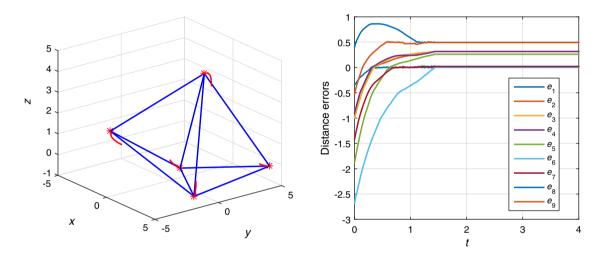


FIGURE 8 Stabilization control of a rigid formation: asymmetric uniform quantization case. Left: the trajectories of five agents and the final formation shape. Right: time evolutions of the distance errors [Colour figure can be viewed at wileyonlinelibrary.com]

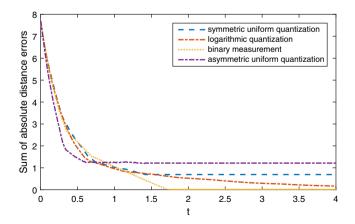


FIGURE 9 Comparisons of convergences of Lyapunov functions (as functions of formation distance errors) under different quantization functions [Colour figure can be viewed at wileyonlinelibrary.com]

7 | CONCLUDING REMARKS

In this paper, we have considered the rigid formation control problem with quantized distance measurements. We have discussed in detail the quantization effect on the convergence of rigid formation shapes under two commonly used quantizers. In the case of the symmetric uniform quantizer, all distances will converge locally to a bounded set, the size of which depends on the quantization error. In the case of the logarithmic quantizer, all distances converge locally to the desired values. We also consider a special quantizer with a signum function, which allows each agent to use very coarse distance measurements (ie, binary information on whether it is close or far away to neighboring agents with respect to the desired distances). We show in this case the formation shape can still be achieved within a finite time. We further discuss the case of an asymmetric quantizer applied in rigid formation control system and analyze the convergence property of distance errors.

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