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## Distributional inference

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## Appendices

## Appendix A

## Proof of Theorem 4.1

The bivariate normality follows from general $U$-statistic theory, and the central limit theorem. Theorem 4.1 is essentially equal to the following three lemmata.

Lemma A. 1 Under $H_{0}: f=\psi$ we have $\operatorname{Var} \varepsilon=\frac{1}{n} \sigma^{2}=\frac{1}{12 n}$.
Lemma A. 2 Under $H_{0}: f=\psi$ we have

$$
\operatorname{Var} \delta=\frac{1}{45} \frac{n+3}{n(n-1)} \rightarrow \tau^{2}=\frac{1}{45}
$$

Lemma A. 3 Under $H_{0}: f=\psi$ we have $\operatorname{Cov}(\varepsilon, \delta)=0$.
Proof of Lemma A. 1
If $Y_{i} \sim \mathrm{U}(0,1)$ i.i.d. then $\operatorname{Var}\left(Y_{i}-\frac{1}{2}\right)=\frac{1}{12}$ and $\operatorname{Var}\left(\frac{1}{n} \sum\left(Y_{i}-\frac{1}{2}\right)\right)=\frac{1}{12 n}$.
Proof of Lemma A. 2
This result can be found in Nair (1936) ${ }^{1}$ Alternatively, it can be derived along the lines of Hoeffding (1948, p. 308) and Fraser (1959, p.230), which will be repeated here. The variance of Gini's mean difference can be written as

$$
\text { Var } g=\frac{2}{n(n-1)}\left[2 \zeta_{1}(n-2)+\zeta_{2}\right]
$$

with $^{2}$

$$
\begin{aligned}
\Delta & =\int_{0}^{1} \int_{0}^{1}\left|y_{1}-y_{2}\right| \mathrm{d} y_{1} \mathrm{~d} y_{2}=\frac{1}{3} \\
\zeta_{1} & =\int_{0}^{1}\left[\int_{0}^{1}\left|y_{1}-y_{2}\right| \mathrm{d} y_{2}\right]^{2} \mathrm{~d} y_{1}-\Delta^{2} \\
& =\int_{0}^{1}\left(y_{1}^{2}-y_{1}+\frac{1}{2}\right)^{2} \mathrm{~d} y_{1}-\Delta^{2}=\frac{7}{60}-\frac{1}{9}=\frac{1}{180} \\
\zeta_{2} & =\int_{0}^{1} \int_{0}^{1}\left(y_{1}-y_{2}\right)^{2} \mathrm{~d} y_{1} \mathrm{~d} y_{2}-\Delta^{2} \\
& =2 \operatorname{Var} y_{i}-\left(\frac{1}{3}\right)^{2}=\frac{1}{18} .
\end{aligned}
$$

[^0]So,

$$
\begin{aligned}
\operatorname{Var} g & =\frac{2}{n(n-1)}\left(\frac{2(n-2)}{180}+\frac{1}{18}\right)=\frac{1}{45} \frac{n+3}{n(n-1)} \\
\tau^{2} & =\lim _{n \rightarrow \infty} n \operatorname{Var} g=\frac{1}{45}
\end{aligned}
$$

Proof of Lemma A. 3
Approach 1
Of course, $\operatorname{Cov}(\varepsilon, \delta)=\operatorname{Cov}(\bar{u}, g)$. That $\operatorname{Cov}(\bar{u}, g)=0$ is easily seen with this symmetry-argument: define $y_{i}^{*}=1-y_{i}$. Then, of course, $\overline{u^{*}}=1-\bar{u}$ is also transformed, yet $g^{*}=g$ remains the same. By definition

$$
\operatorname{Cov}(\bar{u}, g)=\operatorname{Cov}\left(1-\bar{u}^{*}, g^{*}\right)=-\operatorname{Cov}\left(\bar{u}^{*}, g^{*}\right) .
$$

Therefore, the covariance is equal to zero.
Approach 2
The covariance between $\bar{u}$ and $g$ can be written as a function of the covariances of the order statistics,

$$
\begin{aligned}
\operatorname{Cov}(\bar{u}, g) & =\operatorname{Cov}\left(\sum \frac{u_{[i]}}{n}, \sum \frac{2(2 i-n-1)}{n(n-1)} u_{[i]}\right) \\
& =\sum_{i=1}^{n} \sum_{j=1}^{n} \frac{2(2 j-n-1)}{n^{2}(n-1)} \operatorname{Cov}\left(u_{[i]}, u_{[j]}\right) \\
& =\frac{2(2 j-n-1)}{n^{2}(n-1)} \sum_{i=1}^{n} \sum_{j=1}^{n}(2 j-n-1) \operatorname{Cov}\left(u_{[i]}, u_{[j]}\right)
\end{aligned}
$$

To compute the Cov $\left(u_{[i]}, u_{[j]}\right)$, we need the joint distribution function of $U_{[i]}$ and $U_{[j]}$, which is, according to $\operatorname{WILKS}$ (1962, eq. 7.7.6), for $i<j$

$$
\begin{aligned}
f_{U_{[i]}, U_{[j]}}(u, v) & =c_{1} u^{i-1}(v-u)^{j-i-1}(1-v)^{n-j} \\
\text { with } c_{1} & =\frac{\Gamma(n+1)}{\Gamma(i) \Gamma(j-i) \Gamma(n+1-j)}
\end{aligned}
$$

and $0<u<v<1,1 \leq i<j \leq n$. This joint probability density function is that of the ordered bivariate Dirichlet $(i, j-i, n+1-j$ ) distribution. (This result can also be found using Kendall and Stuart, 1958, p. 325, who provides the formula for the joint distribution for general $F$.) Hence

$$
\begin{aligned}
\operatorname{Cov}\left(u_{[i]}, u_{[j]}\right) & =\mathbf{E} u_{[i]} u_{[j]}-\mathbf{E} u_{[i]} \mathbf{E} u_{[j]} \\
& =\mathbf{E} u_{[i]} u_{[j]}-\frac{i j}{(n+1)^{2}}
\end{aligned}
$$

The expectation of the product being

$$
\begin{aligned}
\mathbf{E} U_{[i]} U_{[j]} & =\int_{0}^{1} \int_{0}^{v} u v f_{U_{[i]}, U_{[j]}}(u, v) \mathrm{d} u \mathrm{~d} v \\
& =c_{1} \int_{0}^{1} v(1-v)^{n-j}\left[\int_{0}^{v} u^{i}(v-u)^{j-i-1} \mathrm{~d} u\right] \mathrm{d} v \\
& =c_{1} \int_{0}^{1} v(1-v)^{n-j}\left[v^{s} \operatorname{Beta}(r+1, s-r)\right] \mathrm{d} v \\
& =c_{1} c_{2} \int_{0}^{1} v^{j+1}(1-v)^{n-j} \mathrm{~d} v \\
& =\frac{\Gamma(n+1)}{\Gamma(i) \Gamma(j-i) \Gamma(n-j+1)} \frac{\Gamma(i+1) \Gamma(j-i)}{\Gamma(j+1)} \frac{\Gamma(j+2) \Gamma(n-j+1)}{\Gamma(n+3)} \\
& =\frac{i(j+1)}{(n+1)(n+2)}
\end{aligned}
$$

That $c_{2}=\operatorname{Beta}(r+1, s-r)$ can be seen through

$$
\int_{0}^{v} u^{i}(v-u)^{j-i-1} \mathrm{~d} u=v^{j} \int_{0}^{1} w^{i}(1-w)^{j-i-1} \mathrm{~d} w=\mathrm{B}(i+1, j-i) v^{j}
$$

Using $\mathbf{E} U_{[i]} U_{[j]}$, we can calculate the covariance

$$
\operatorname{Cov}\left(U_{[i]} U_{[j]}\right)= \begin{cases}\frac{i(j+1)}{(n+1)(n+2)}-\frac{i j}{(n+1)^{2}}=\frac{i(n+1-j)}{(n+1)^{2}(n+2)} & \text { if } i<j \\ \frac{i(n+1-i)}{(n+1)^{2}(n+2)} & \text { if } i=j \\ \frac{j(n+1)}{(n+1)^{2}(n+2)} & \text { if } i>j\end{cases}
$$

Note that the formula for $i>j$ immediately follows from that for $i<j$ by swapping the $i$ and $j$. The formula for $i=j$ is easily derived by having that $\operatorname{Cov}\left(U_{[i]} U_{[j]}\right)=$ $\operatorname{Var}\left(U_{[i]}\right)$ and $U_{[i]} \sim \operatorname{Beta}(i, n+1-i)$. Note that this formula coincides with the formula for $i<j$ as well as the formula for $i>j$. The final step is to use all these covariances to derive that $\operatorname{Cov}(\bar{u}, g)=0$, and, hence, $\rho=0 . \operatorname{Cov}(\bar{u}, g)=$

$$
\begin{aligned}
= & \frac{2}{n^{2}(n-1)} \sum_{j=1}^{n} \sum_{i=1}^{n}(2 j-n-1) \operatorname{Cov}\left(u_{[i]}, u_{[j]}\right) \\
= & \frac{2}{(n-1) n^{2}(n+1)^{2}(n+2)}\left(\sum_{j=1}^{n} \sum_{i=1}^{j-1}(2 j-n-1) i(n+1-j)+\right. \\
& \left.+\sum_{j=1}^{n}(2 j-n-1) j(n+1-j)+\sum_{j=1}^{n} \sum_{i=j+1}^{n}(2 j-n-1) j(n+1-i)\right) \\
= & \frac{2}{(n-1) n^{2}(n+1)^{2}(n+2)}\left(\frac{(n-1) n(n+1)(n+2)(n+3)}{120}+0-\frac{(n-1) n(n+1)(n+2)(n+3)}{120}\right) \\
= & 0
\end{aligned}
$$

## Appendix B

## Critical values for the $\left\|f_{n}^{(m)}-\psi\right\|_{1}$-test statistic

The Tables B.1, B.2, and B. 3 are each constructed on the basis of the following simulation experiment. For $m=2,3$, and 4 , and for each value of $n$, a sample of size $n$ was drawn from the standard uniform distribution providing an outcome $t_{n}^{(m)}$ of the test statistic $T_{n}^{(m)}$. This process is repeated 100000 times. (Simulation studies showed that this number of replications suffices to obtain reliable critical values.) The percentiles are taken from the empirical distribution of $T_{n}^{(m)}$. (Note: these values are computed using the true $t_{n}^{(m)}$, not the approximation mentioned in Section 4.4 for $m=2$.)

| $n$ | .90 | .75 | .50 | .25 | .10 | .05 | .025 | .01 |
| ---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 5 | .0499 | .0838 | .1422 | .2311 | .3241 | .3801 | .4302 | .4865 |
| 6 | .0439 | .0746 | .1283 | .2093 | .2949 | .3482 | .3954 | .4466 |
| 7 | .0394 | .0675 | .1177 | .1934 | .2717 | .3215 | .3641 | .4133 |
| 8 | .0365 | .0627 | .1102 | .1807 | .2539 | .3007 | .3436 | .3906 |
| 9 | .0336 | .0582 | .1029 | .1700 | .2401 | .2842 | .3224 | .3664 |
| 10 | .0320 | .0552 | .0979 | .1613 | .2276 | .2684 | .3055 | .3473 |
| 12 | .0286 | .0494 | .0889 | .1468 | .2075 | .2463 | .2802 | .3207 |
| 14 | .0259 | .0455 | .0818 | .1357 | .1914 | .2268 | .2574 | .2950 |
| 16 | .0241 | .0424 | .0765 | .1274 | .1798 | .2125 | .2415 | .2769 |
| 18 | .0225 | .0395 | .0718 | .1198 | .1700 | .2013 | .2286 | .2620 |
| 20 | .0212 | .0374 | .0685 | .1132 | .1604 | .1899 | .2164 | .2486 |
| 22 | .0201 | .0354 | .0651 | .1085 | .1531 | .1816 | .2065 | .2379 |
| 24 | .0191 | .0339 | .0623 | .1037 | .1471 | .1741 | .1990 | .2283 |
| 26 | .0185 | .0326 | .0596 | .0995 | .1406 | .1667 | .1908 | .2178 |
| 28 | .0176 | .0312 | .0576 | .0962 | .1361 | .1613 | .1838 | .2118 |
| 30 | .0171 | .0303 | .0555 | .0922 | .1310 | .1562 | .1788 | .2050 |
| 35 | .0155 | .0278 | .0512 | .0849 | .1209 | .1434 | .1633 | .1890 |
| 40 | .0147 | .0262 | .0482 | .0798 | .1133 | .1345 | .1538 | .1759 |
| 45 | .0137 | .0244 | .0454 | .0754 | .1066 | .1266 | .1452 | .1677 |
| 50 | .0130 | .0232 | .0428 | .0715 | .1014 | .1211 | .1379 | .1580 |
| 60 | .0118 | .0211 | .0391 | .0651 | .0925 | .1101 | .1259 | .1437 |
| 70 | .0109 | .0195 | .0362 | .0607 | .0856 | .1015 | .1160 | .1331 |
| 80 | .0103 | .0183 | .0339 | .0564 | .0802 | .0954 | .1084 | .1248 |
| 90 | .0096 | .0172 | .0319 | .0532 | .0756 | .0898 | .1024 | .1175 |
| 100 | .0090 | .0162 | .0302 | .0506 | .0718 | .0852 | .0971 | .1115 |
| 120 | .0082 | .0148 | .0276 | .0459 | .0653 | .0777 | .0889 | .1024 |
| 140 | .0077 | .0138 | .0257 | .0425 | .0606 | .0722 | .0823 | .0940 |
| 160 | .0072 | .0129 | .0238 | .0399 | .0567 | .0675 | .0770 | .0881 |
| 180 | .0067 | .0121 | .0225 | .0378 | .0535 | .0637 | .0729 | .0843 |
| 200 | .0064 | .0115 | .0215 | .0355 | .0504 | .0600 | .0688 | .0788 |
| 250 | .0057 | .0103 | .0190 | .0319 | .0454 | .0538 | .0614 | .0708 |
| 300 | .0052 | .0093 | .0174 | .0291 | .0414 | .0493 | .0562 | .0646 |
| 400 | .0045 | .0081 | .0151 | .0254 | .0360 | .0429 | .0490 | .0560 |
| 500 | .0040 | .0072 | .0135 | .0226 | .0321 | .0382 | .0436 | .0498 |
|  |  |  |  |  |  |  |  |  |

Table B.1: Critical values for $m=2$

| $n$ | .90 | .75 | .50 | .25 | .10 | .05 | .025 | .01 |
| ---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 5 | .0784 | .1267 | .1957 | .2899 | .3976 | .4656 | .5253 | .5930 |
| 6 | .0684 | .1119 | .1747 | .2611 | .3595 | .4228 | .4776 | .5393 |
| 7 | .0618 | .1007 | .1590 | .2411 | .3331 | .3915 | .4422 | .5001 |
| 8 | .0565 | .0924 | .1469 | .2246 | .3101 | .3641 | .4131 | .4730 |
| 9 | .0517 | .0860 | .1376 | .2114 | .2932 | .3456 | .3918 | .4444 |
| 10 | .0491 | .0811 | .1302 | .1992 | .2767 | .3265 | .3712 | .4234 |
| 12 | .0440 | .0728 | .1176 | .1817 | .2526 | .2980 | .3398 | .3879 |
| 14 | .0397 | .0665 | .1081 | .1679 | .2336 | .2760 | .3137 | .3592 |
| 16 | .0366 | .0616 | .1003 | .1567 | .2188 | .2584 | .2938 | .3360 |
| 18 | .0344 | .0576 | .0947 | .1477 | .2061 | .2440 | .2779 | .3180 |
| 20 | .0327 | .0546 | .0892 | .1397 | .1952 | .2310 | .2625 | .3015 |
| 22 | .0310 | .0519 | .0851 | .1327 | .1863 | .2204 | .2512 | .2869 |
| 24 | .0293 | .0494 | .0812 | .1272 | .1777 | .2103 | .2403 | .2735 |
| 26 | .0284 | .0474 | .0778 | .1223 | .1714 | .2027 | .2302 | .2645 |
| 28 | .0270 | .0454 | .0747 | .1182 | .1642 | .1949 | .2228 | .2554 |
| 30 | .0259 | .0437 | .0723 | .1137 | .1590 | .1880 | .2147 | .2465 |
| 35 | .0240 | .0403 | .0668 | .1051 | .1475 | .1748 | .1992 | .2287 |
| 40 | .0223 | .0375 | .0623 | .0985 | .1380 | .1640 | .1870 | .2144 |
| 45 | .0210 | .0353 | .0585 | .0923 | .1287 | .1532 | .1749 | .2009 |
| 50 | .0198 | .0336 | .0557 | .0877 | .1230 | .1457 | .1657 | .1889 |
| 60 | .0181 | .0306 | .0508 | .0799 | .1119 | .1329 | .1516 | .1733 |
| 70 | .0165 | .0282 | .0471 | .0743 | .1043 | .1237 | .1405 | .1609 |
| 80 | .0155 | .0263 | .0440 | .0694 | .0972 | .1155 | .1318 | .1519 |
| 90 | .0147 | .0248 | .0414 | .0655 | .0917 | .1085 | .1232 | .1411 |
| 100 | .0138 | .0232 | .0390 | .0619 | .0868 | .1032 | .1172 | .1343 |
| 120 | .0126 | .0214 | .0357 | .0565 | .0794 | .0940 | .1073 | .1234 |
| 140 | .0117 | .0197 | .0330 | .0523 | .0735 | .0870 | .0992 | .1142 |
| 160 | .0108 | .0185 | .0310 | .0488 | .0687 | .0814 | .0924 | .1058 |
| 180 | .0102 | .0173 | .0291 | .0461 | .0649 | .0767 | .0871 | .0996 |
| 200 | .0097 | .0165 | .0276 | .0439 | .0616 | .0730 | .0832 | .0954 |
| 250 | .0087 | .0148 | .0248 | .0394 | .0552 | .0652 | .0744 | .0854 |
| 300 | .0079 | .0133 | .0225 | .0357 | .0502 | .0594 | .0677 | .0774 |
| 400 | .0068 | .0116 | .0195 | .0310 | .0435 | .0516 | .0587 | .0673 |
| 500 | .0061 | .0104 | .0174 | .0278 | .0389 | .0461 | .0525 | .0600 |

Table B.2: Critical values for $m=3$


[^0]:    ${ }^{1}$ Lomnicki (1952) showed that there was an error in Nair's formulations, though in the uniform case the results were true (cf. Kendall and Stuart (1958, p. 241-242)).
    ${ }^{2}$ Note that the formula for $\zeta_{1}$ in Fraser (1959, p.230) is wrong.

