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Laser written waveguides in glasses and crystals

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Waveguides for integrated optical applications are usually made by indiffusion of ions into crystal or glass materials. However, the generation of three-dimensional structures is hard to achieve This problem can be overcome by laser generated

subclines is hard to achieve this problem can be overcome by laser generated waveguides. Ultrashort laser pulses focused inside optical transparent bulk materials create local structural changes and in consequence of this a refractive index gradient. The result of moving the sample through the Laserbeam is a refractive index profile like in a buried waveguide

This direct writing method offers the potential for three-dimensional waveguide

This order writing method others the potential for thereafterins intal waveguide structures and allows integrated optics in different materials. The experimental setup is shown in Fig 1 The laserpulses that we used to induce refractive index changes were obtained from a regeneratively amplified $Ti^3 A_{\rm ICO}$ laser The duration of the 780 nm pulse was 12 kTz. The laserbearn with a pulse energy of a few microjoules was focused into the material. The sample was fixed on a computercontrolled XYZ stage



Generation of waveguides in fused and crystal silica is presented. Properties of the waveguides are discussed

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Photon Echo Spectroscopy in the Single Optical-Cycle Regime

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A very high temporal resolution and a broad bandwidth are but two advantages provided by the use of extremely short sub-5-fs pulses [1] in a nonlinear spectroscopic experiment. However, the applicability of the standard theoretical description becomes questionable for the pulses that consist merely of a couple of optical oscillations. For instance, the conventionally employed slowly varying envelope approximation, implying that the change of the pulse amplitude inself, can no longer be maintained. Furthermore, the phase-matching bandwidth that is limited due to dispersion in the nonlinear medium rapidly gains importance with the increase of the spectral width of the pulse. Another point of serious concern is the frequency-dependent variation in the sensitivity of signal photodetectors. In combination, the above listed features result in what is known as a spectral-filter effect. Finally, artificial lengthening of the experimental transients is a direct consequence of the noncollinear geometry employed in spectroscopic experiments

known as a spectral-filter effect. Finally, artificial lengthening of the experimental transients is a direct consequence of the noncollinear geometry employed in spectroscopic experiments. In this Contribution we present a theoretical analysis which throughly reexamines the formalism of ultrafast photon echo spectroscopy. We obtain a general expression for the echo signal, which is valid even for single-eycle-pulse applications. The derived formalism is applied to photon-echo spectroscopy on the hydrated electron with 5-fs pulses. According to our calculations, a careful choice of the beam geometry and selection of a photodetector with the suitable spectral-filter effect as well as that of geometrical smearing. Importantly for the weak-signal applications such as photon echo spectroscopy, the absence of spectral filtering eliminates the otherwise unavoidable requirement to frequency-resolve the signals. signals

Figures 1 presents the two-pulse photon echo signals obtained Figures 1 presents the two-pulse photon echo signals obtained from the near twater and electrons, solvated in water. A minute difference in the widths of these two traces suggests that the electronic dephasing of the hydrated electrons is extremely fast. The finite population lifetime of the electrons in the excited state causes the delay of the echo trace in Fig.2b. The best fit to the causes the used of the card order in Figure 1 both T_{t} because the experimental data yields a pure dephasing time of $T_{2} = 16$ fs. With this value, the absorption spectrum of the hydrated electron can be successfully modeled provided the conventionally used rotating wave approximation is abandoned [2].

References 1. A. Battuska et al., Opt. Lett 22, 102 (1997); M Nisoli et al, Opt. Lett 22, 522 (1997); A Shirakawa et al., Appl. Phys. Lett. 74, 2268 (1999). 2. A. Baltuska et al., J. Phys. Chem. A103, 10065 (1999).



Fig 2. Results of two-pulse Fig.2. Results of two-pulse photon echo experiments on water alone (a) and hydrated electrons (b). Circles represent experimental data points and solid curves show results of simulations.

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