Memory cost of quantum contextuality

Matthias Kleinmann, Otfried Guehne, Jose R. Portillo, Jan-Åke Larsson and Adan Cabello

Linköping University Post Print

N.B.: When citing this work, cite the original article.

Original Publication:

Matthias Kleinmann, Otfried Guehne, Jose R. Portillo, Jan-Åke Larsson and Adan Cabello, Memory cost of quantum contextuality, 2011, New Journal of Physics, (13), 113011. <u>http://dx.doi.org/10.1088/1367-2630/13/11/113011</u> Licensee: Institute of Physics <u>http://www.iop.org/</u>

Postprint available at: Linköping University Electronic Press http://urn.kb.se/resolve?urn=urn:nbn:se:liu:diva-73336

New Journal of Physics

The open-access journal for physics

Memory cost of quantum contextuality

Matthias Kleinmann^{1,2,8}, Otfried Gühne^{1,2,3}, José R Portillo⁴, Jan-Åke Larsson⁵ and Adán Cabello^{6,7}

¹ Universität Siegen, Department Physik, Walter-Flex-Straße 3, D-57068 Siegen, Germany
² Institut für Quantenoptik und Quanteninformation, Österreichische Akademie der Wissenschaften, Technikerstraße 21A, A-6020 Innsbruck, Austria
³ Institut für Theoretische Physik, Universität Innsbruck, Technikerstraße 25, A-6020 Innsbruck, Austria
⁴ Departamento de Mathemática Aplicada I, Universidad de Sevilla, E-41012 Sevilla, Spain
⁵ Institutionen för Systemteknik och Matematiska Institutionen, Linköpings Universitet, SE-581 83 Linköping, Sweden
⁶ Departamento de Física Aplicada II, Universidad de Sevilla, E-41012 Sevilla, Spain
⁷ Department of Physics, Stockholm University, S-10691 Stockholm, Sweden E-mail: matthias.kleinmann@uni-siegen.de

New Journal of Physics **13** (2011) 113011 (20 Received 8 July 2011 Published 9 November 2011 Online at http://www.njp.org/ doi:10.1088/1367-2630/13/11/113011

Abstract. The simulation of quantum effects requires certain classical resources, and quantifying them is an important step to characterize the difference between quantum and classical physics. For a simulation of the phenomenon of state-independent quantum contextuality, we show that the minimum amount of memory used by the simulation is the critical resource. We derive optimal simulation strategies for important cases and prove that reproducing the results of sequential measurements on a two-qubit system requires more memory than the information-carrying capacity of the system.

⁸ Author to whom any correspondence should be addressed.

IOP Institute of Physics **O**DEUTSCHE PHYSIKALISCHE GESELLSCHAFT

Contents

1. Introduction		2
2. Scenario		3
3. A first model		3
4. Memory cost of classical models		4
5. Contextuality conditions		5
6. Compatibility conditions		5
7. The extended Peres–Mermin (PM) s	quare	6
8. Conclusions		7
Acknowledgments		8
Appendix A. A_3 is optimal		8
Appendix B. A_4 obeys $\mathscr{Q}'_{\text{context}}$ and $\mathscr{Q}'_{\text{cont}}$	nnat	9
Appendix C. Definitions and basic rule		9
Appendix D. A_4 is memory-optimal		10
Appendix E. Proof that the classical si	mulation of the extended PM square re	quires
more than two bits of memo		11
Appendix F. A ten-state automaton ob	eying all sequences	19
References		19

1. Introduction

According to quantum mechanics (QM), the result of a measurement may depend on which other compatible observables are measured simultaneously [1-3]. This property is called contextuality and is in contrast to classical physics, where the answer to a single question does not depend on which other compatible questions are asked at the same time.

Contextuality can be seen as complementary to the well-known nonlocality of distributed quantum systems [4]. Both phenomena can be used for information-processing tasks, albeit the applications of contextuality are far less explored [5-12]. Although contextuality and nonlocality can be considered as signatures of nonclassicality, they can be simulated by classical models [3, 13–15]. However, while nonlocal classical models violate a fundamental physical principle (the bounded speed of information), it is not clear whether contextual classical models violate any fundamental principle. Moreover, while the resources needed to imitate quantum nonlocality have been extensively investigated [16–18], there is no such knowledge about the resources needed to simulate quantum contextuality.

In any model that exhibits contextuality in sequential measurements, the system will eventually attain different internal states during certain measurement sequences. These states can be considered as *memory*—a model attaining the minimum number of states is then memory-optimal and defines the memory cost. In this paper, we investigate the memory cost as the critical resource in a classical simulation of quantum contextuality and we construct memory-optimal models for relevant cases. The amount of memory required increases as we consider more and more contextuality constraints. This can be used to quantify contextuality in a given quantum setting. We show that certain scenarios breach the amount of two-bits needed for

2

3

the simulation of two qubits. This demonstrates that the memory needed to simulate only a small set of measurements on a quantum system may exceed the information that can be transmitted using this system (given by the Holevo bound [19])—a similar effect has been observed so far only for the classical simulation of a *unitary* evolution [20].

2. Scenario

We focus on the following set of two-qubit observables, also known as the Peres–Mermin (PM) square [4, 21],

$$\begin{bmatrix} A & B & C \\ a & b & c \\ \alpha & \beta & \gamma \end{bmatrix} = \begin{bmatrix} \sigma_z \otimes \mathbb{1} & \mathbb{1} \otimes \sigma_z & \sigma_z \otimes \sigma_z \\ \mathbb{1} \otimes \sigma_x & \sigma_x \otimes \mathbb{1} & \sigma_x \otimes \sigma_x \\ \sigma_z \otimes \sigma_x & \sigma_x \otimes \sigma_z & \sigma_y \otimes \sigma_y \end{bmatrix},$$
(1)

where σ_x , σ_y and σ_z denote the Pauli operators. The square is constructed such that the observables within each row and column commute and are hence compatible, and the product of the operators in a row or column yields 1, except for the last column where it yields -1. Thus, the product of the measurement results for each row and column will be +1 except in the third column, where it will be -1. In contrast, for a noncontextual model the measurement result for each observable must not depend on whether the observable is measured in the column or row context. Hence, the number of rows and columns yielding a product of -1 is always even, as any observable appears twice.

Similar to the Bell inequalities for local models, any noncontextual model satisfies the inequality

$$\langle \chi \rangle \equiv \langle ABC \rangle + \langle abc \rangle + \langle \alpha\beta\gamma \rangle + \langle Aa\alpha \rangle + \langle Bb\beta \rangle - \langle Cc\gamma \rangle \leqslant 4, \tag{2}$$

while for perfect observables QM predicts $\langle \chi \rangle = 6$ [22]. Here, the term $\langle ABC \rangle$ denotes the average value of the product of the outcomes of *A*, *B* and *C* if these observables are measured simultaneously or in sequence on the same quantum system. The violation is independent of the quantum state, which emphasizes that the phenomenon is a property of the set of observables rather than of a particular quantum state.

Recently, this inequality was experimentally tested using trapped ions [23], photons [24] and nuclear magnetic resonance systems [25]. The results show good agreement with the quantum predictions. In these experiments, the observables are measured in a sequential manner. Since the observed results cannot be explained by a model using only preassigned values, the system necessarily attains different states during some particular sequences, i.e. the system memorizes previous events. (Note that in QM the system also attains different states during the measurement sequences.) This leads to our central question: *how much memory is required to simulate quantum contextuality*?

3. A first model

Before we formulate the previous question more precisely, let us provide an example of a model that simulates the contextuality in the PM square. We assume that the system can only attain three different physical states S_1 , S_2 and S_3 (e.g. discrete points in phase space). Let us associate a table with each state via

$$S_{1}:\begin{bmatrix} + & + & (+, 2) \\ + & + & (+, 3) \\ + & + & + \end{bmatrix}, \quad S_{2}:\begin{bmatrix} + & (+, 1) & + \\ - & + & - \\ - & (-, 3) & + \end{bmatrix}, \quad S_{3}:\begin{bmatrix} + & - & - \\ (+, 1) & + & + \\ (-, 2) & - & + \end{bmatrix}.$$
(3)

Those tables define the model's behavior in the following way: if, e.g., the system is in state S_1 and we measure the observable γ , consider the first table at the position of γ (i.e. the last entry in the third row). The + sign at this position indicates that the measurement result will be +1, while the system stays in state S_1 . If we continue and measure *C*, we encounter the entry (+, 2), which indicates the measurement result +1 and a subsequent change to the state S_2 . Being in state S_2 , the second table defines the behavior for the next measurement: for instance, a measurement of *c* yields now the result -1 and the system stays in state S_2 .

Thus, starting from state S_1 , the measurement results for the sequence γCc are +1, +1, -1, so that the product is -1 in accordance with the quantum prediction. It is straightforward to verify that this model yields $\langle \chi \rangle = 6$. In addition, the observables within each context are compatible in the sense that in sequences of the form AA, ABA or A αaA , the first and last measurement of A yields the same output. In fact this particular model is *memory-optimal* (cf theorem 1) and we assign the symbol A_3 to it.

4. Memory cost of classical models

Any model that reproduces contextuality eventually predicts that the system attains different states during some measurement sequences. As an omniscient observer one would know the state prior to each measurement and could include it in the measurement record. Thus, knowing the state of the system, one can predict the measurement outcome as well as the state of the system that will occur prior to the next measurement. Thus, we can write any model that explains the outcomes of sequential measurements in the same fashion as we did for A_3 .

Taking a different point of view, any such model can be considered to be an automaton with finitely many internal states, taking inputs (measurement settings) and yielding outputs (measurement results). In our notation, the output depends not only on the internal state, but also on the input. Such automatons are known as Mealy machines [26, 27].

The quantum predictions add restrictions to such an automaton and thus increase the number of internal states needed. As a simple example we could require that an automaton reproduces the quantum predictions from the rows and columns in the PM square. That is, for all sequences in the set

$$\mathcal{Q}_{\rm rc} = \{ABC, abc, \alpha\beta\gamma, Aa\alpha, Bb\gamma, Cc\gamma \text{ and permutations}\},\tag{4}$$

we require that the automaton must yield an output that matches the quantum prediction. For example, QM predicts for the sequence *ABC* that the output is either +1, +1, +1 or one of the permutations of +1, -1, -1.

More generally, if \mathscr{Q} denotes a set of measurement sequences, we say that an automaton \mathcal{A} obeys the set \mathscr{Q} if the output for any sequence in \mathscr{Q} matches the quantum prediction—i.e., if for any sequence in \mathscr{Q} , the output of \mathcal{A} could have occurred with a nonvanishing probability according to the quantum scenario. We say that a sequence *yields a contradiction* if the output of this sequence cannot occur according to QM. Hereby we consider all quantum predictions from any initial state (it would also suffice to only consider the completely mixed state $\varrho = 1/\text{tr}[1]$). Furthermore, we assume that prior to the measurement of a sequence, the automaton is always re-initialized. This ensures that the output of the automaton is independent of any action prior to the selected measurement sequence. Note that we only consider the *certain* quantum predictions that occur with a probability of 1, whereas, e.g., in the PM square we do not require that for the sequence $A\gamma$ the probability of obtaining +1, +1 is equal to the probability of obtaining +1, -1.

Finally, if an automaton with k states S_1, \ldots, S_k obeys \mathcal{Q} and there exists no automaton with fewer states obeying \mathcal{Q} , we define the *memory cost* of \mathcal{Q} to be $M(\mathcal{Q}) = \log_2(k)$.

5. Contextuality conditions

Our definition of memory cost so far applies to arbitrary situations, even those in which contextuality does not directly play a role. In contrast, contextuality of sequential measurements corresponds to the particular feature that certain sequences of mutually compatible observables cannot be explained by a model with preassigned values (cf [28] for a detailed discussion). The contextuality conditions for observables X_1, X_2, \ldots thus arise from the set of all sequences of mutually compatible observables,

$$\mathscr{Q}_{\text{context}} = \{X_1 X_2 \dots \mid X_\ell \text{ mutually compatible}\}.$$
(5)

If the choice of observables $X_1, X_2, ...$ exhibits contextuality, then $M(\mathscr{Q}_{context}) > 0$. In the case of the PM square, $\mathscr{Q}_{context}$ surely contains all the row and column sequences that we included in \mathscr{Q}_{rc} . In addition, however, $\mathscr{Q}_{context}$ contains, e.g., the sequences AA, ABA and A αaA , for which QM predicts with certainty a repetition of the value of A in the first and last instance. Note that the set $\mathscr{Q}_{context}$ also contains more complicated sequences like ACABCA for which QM predicts with certainty that the values of A (C) in the first, third and sixth (second and fifth) measurements coincide and that the product of the outcome for ABC yields +1.

A particular feature of contextuality is that one can find observables that exhibit contextuality independently of the actual preparation (the initial state) of the quantum system. Consequently, one may consider an extended preparation procedure of the automaton, where the experimenter carries out additional measurements between the initialization of the automaton and the actual sequence. The experimenter would, e.g., measure the sequence bABC but consider the measurement of the observable b to be actually part of the preparation procedure. We write $\lfloor b \rfloor$ for a sequence where we are not interested in the result of b. If \mathcal{Q}_{all} denotes the set of all sequences with observables X_1, X_2, \ldots , we write

$$\mathscr{Q}_{\text{context}}' = \{ \lfloor T \rfloor S \mid S \in \mathscr{Q}_{\text{context}}, T \in \mathscr{Q}_{\text{all}} \}$$
(6)

for the set of all sequences in $\mathscr{Q}_{\text{context}}$, including arbitrary preparation procedures.

For the contextuality in the PM square, the automaton A_3 obeys $\mathcal{Q}_{context}$, while no automaton with fewer than three states can obey $\mathcal{Q}_{context}$; cf appendix A for details. We did not specify an initial state for A_3 and indeed the contextuality conditions are obeyed for any initial state. We summarize:

Theorem 1. The memory cost for the contextuality correlations $\mathscr{Q}'_{\text{context}}$ in the PM square is $\log_2(3) \approx 1.58$ bits; $M(\mathscr{Q}'_{\text{context}}) = M(\mathscr{Q}_{\text{context}}) = \log_2(3)$. Consequently, the automaton \mathcal{A}_3 is memory-optimal.

6. Compatibility conditions

The set $\mathscr{Q}_{context}$ contains all sequences of mutually compatible observables, but does not contain sequences like *ABaA*, for which QM also predicts that both occurring values of *A* are the same. Sequences of this form enforce that all observables compatible with an observable *Y*

must not change the measurement result of Y. This can be covered by the set of all compatibility conditions

$$\mathcal{Q}_{\text{compat}} = \{ Y \lfloor X_1 X_2 \dots \rfloor Y \mid X_\ell \text{ compatible to } Y \},$$
(7)

and a convincing test of contextuality must also test the correlations due to this set of sequences. Again we define $\mathscr{Q}'_{\text{compat}}$ to include arbitrary preparation procedures.

The automaton \mathcal{A}_3 does not obey \mathcal{Q}_{compat} , since e.g. starting with state S_1 , the sequence $B \lfloor C\beta \rfloor B$ yields the record +1, $\lfloor +1, -1, \rfloor -1$ and hence violates the assumption of compatibility; similar sequences can be found for any initial state. We show in appendix D that no automaton with three states can obey simultaneously \mathcal{Q}'_{compat} and $\mathcal{Q}'_{context}$ and hence $M(\mathcal{Q}'_{compat} \text{ and } \mathcal{Q}'_{context}) \ge 2$. However, automata with four internal states exist that obey \mathcal{Q}'_{compat} and \mathcal{Q}'_{compat} . As an example of such an automaton, we define \mathcal{A}_4 via

$$S_{1}:\begin{bmatrix} + & + & (+, 2) \\ + & + & (+, 3) \\ + & + & + \end{bmatrix}, S_{2}:\begin{bmatrix} + & + & + \\ - & + & - \\ (-, 4) & (+, 1) & + \end{bmatrix},$$

$$S_{3}:\begin{bmatrix} + & - & - \\ + & + & + \\ (+, 1) & (-, 4) & + \end{bmatrix}, S_{4}:\begin{bmatrix} + & - & (-, 3) \\ - & + & (-, 2) \\ - & - & + \end{bmatrix}.$$
(8)

Similar to the situation for A_3 , the initial state for the automaton A_4 can be chosen freely; for details see appendix B. So we have:

Theorem 2. The memory cost for the contextuality and compatibility correlations in the PM square is two bits; $M(\mathscr{Q}'_{compat} \text{ and } \mathscr{Q}'_{context}) = 2$. Consequently, the automaton \mathcal{A}_4 is memory-optimal.

7. The extended Peres–Mermin (PM) square

There are, however, further contextuality effects for two qubits, which then require more than two bits for a simulation. Namely, in [29] an extension of the PM square has been introduced, involving 15 different observables in 15 different contexts. The argument goes as follows: consider the 15 observables of the type $\sigma_{\mu} \otimes \sigma_{\nu}$ where $\mu, \nu \in \{0, x, y, z\}$ and $\sigma_0 = 1$ and the case $\mu = \nu = 0$ is excluded. In this set, there are 12 trios of mutually compatible observables such that the product of their results is always +1, such as $[\sigma_x \otimes 1, 1 \otimes \sigma_y, \sigma_x \otimes \sigma_y]$ and $[\sigma_x \otimes \sigma_y, \sigma_y \otimes \sigma_x, \sigma_z \otimes \sigma_z]$, and three trios of mutually compatible observables such that the product of their results is always -1, like $[\sigma_x \otimes \sigma_y, \sigma_y \otimes \sigma_z, \sigma_z \otimes \sigma_x]$. This leads to 15 contexts in total. Similarly to the usual PM square, one can derive a state-independent inequality. For this inequality, QM predicts a value of 15 for the total sum, whereas noncontextual models have a maximal value of 9; cf [29] and appendix E for details.

One may argue that this new contextuality argument is stronger than the usual PM square [29]. Does a simulation of it require more memory than the original PM square? Indeed, this is the case:

Theorem 3. The memory cost for the contextuality and compatibility correlations in the extended version of the PM square is strictly larger than two bits.

More precisely, according to equations (5) and (7) we define the contextuality sequences $\mathscr{Q}'_{compat,15}$ and compatibility sequences $\mathscr{Q}'_{context,15}$ for the 15 observables in the extension of the PM square. Then, the theorem states that $M(\mathscr{Q}'_{compat,15} \text{ and } \mathscr{Q}'_{context,15}) > 2$. The proof is based on the following idea: if one considers the 15 contexts in the extended square, then they can be arranged in a collection of ten distinct squares, each similar to the usual PM square. The contextuality in this arrangement is strong enough that for each fixed assignment of the output, one must have *three* contradictions for one of the ten usual PM squares. One can show, however, that any four-state solution obeying $\mathscr{Q}'_{context}$ and \mathscr{Q}'_{compat} is similar to \mathcal{A}_4 , in which for no fixed state one has three contradictions. The full proof is given in appendix E.

Although this paper is mainly concerned with the memory cost of contextuality, we mention that simulation of all certain quantum predictions of the PM square already requires more than two bits of memory. In fact, any four-state automaton that obeys \mathscr{Q}'_{compat} and $\mathscr{Q}'_{context}$ is of the form \mathcal{A}_4 , up to some symmetries (cf appendix E, proposition 5), but for \mathcal{A}_4 the sequence $\lfloor \beta \rfloor ab \lfloor C \rfloor c$ yields a contradiction. This proves that $M(\mathscr{Q}_{all}) > 2$.

On the other hand, QM itself suggests an automaton for simulating contextuality. If, e.g., we choose the pure state $|\phi\rangle\langle\phi|$ defined by $A|\phi\rangle = B|\phi\rangle = |\phi\rangle$ as the initial state, then this state and all the states occurring during measurement sequences define a (nondeterministic) automaton. By a straightforward calculation one finds that this automaton attains 24 different states if we consider the set of all sequences \mathcal{Q}_{all} . By a suitable elimination of the nondeterminism, we can readily reduce the number of states to ten (cf appendix F), yielding an upper bound on the required memory and hence $2 < M(\mathcal{Q}_{all}) \leq \log_2(10) \approx 3.32$.

8. Conclusions

We have investigated the amount of memory needed in order to simulate quantum contextuality in sequential measurements. We determined the memory-optimal automata for important cases and have proven that the simulation of contextuality phenomena for two qubits requires more than two classical bits of memory. However, the maximum amount of classical information that can be stored and retrieved in two qubits is well known to be limited to two bits [19]. This implies that any classical model of such a system would either allow storage and retrieval of more than two bits, or would have inaccessible degrees of freedom. (An example of the latter is A_3 , since one cannot perfectly infer the initial state from the results of any measurement sequence.)

It should be emphasized that our analysis is about the memory that is needed to classically simulate the *certain* predictions from measurement *sequences* on a quantum system. In contrast, one may ask: how many different states are needed to merely explain the observed expectation values [30–35]? However, the number of states needed in this scenario measures the number of different initial configurations of the system, while we have shown that even for a fixed initial configuration, the system must eventually attain a certain number of states during measurement sequences. Similarly, it has been demonstrated that a hybrid system of one qubit and one classical bit of memory is *on average* superior to a classical system having access only to a single bit of memory [36], while we show in theorem 1 that for a two-qubit system even the certain predictions cannot be simulated with one classical bit of memory.

Our work provides a link between information theoretical concepts, on the one hand, and quantum contextuality and the Kochen–Specker theorem, on the other. Whereas for Bell's theorem such connections are well explored and have given as deep insights into

QM [18, 37, 38], for contextuality many questions remain open: if an experiment violates some noncontextuality inequality up to a certain degree, but not maximally, what memory is required to simulate this behavior? Can nondeterministic machines help us to simulate contextuality? What amount of memory and randomness is required to simulate all quantum effects in the PM square, especially in the distributed setting [12]? Finally, for quantum nonlocality it has been extensively investigated why QM does not exhibit the maximal nonlocality [37, 38]. A similar situation occurs for quantum contextuality—can concepts from information theory also help us to understand the nonmaximal violation in this situation?

Acknowledgments

The authors thank E Amselem, P Badziąg, J Barrett, I Bengtsson, M Bourennane, Č Brukner, P Horodecki, A R Plastino, M Rådmark and V Scholz for discussions. This work was supported by the Austrian Science Fund (FWF) through Y376 N16 (START Prize) and SFB FOQUS, by the MICINN projects MTM2008-05866 and FIS2008-05596, by the Wenner-Gren Foundation and by the EU (QICS, NAMEQUAM, Marie Curie CIG 293993/ENFOQI).

Appendix A. A_3 is optimal

We have already defined the set of row and column sequences \mathcal{Q}_{rc} in equation (4). Another natural constraint is given by the set of repeated measurements

$$\mathcal{Q}_{\text{repeat}} = \{AA, BB, CC, aa, bb, \text{c.c.}, \alpha\alpha, \beta\beta, \gamma\gamma\},\tag{A.1}$$

where we expect for any of these pairs that the results in the first and the second measurement coincide. Both sets \mathcal{Q}_{rc} and \mathcal{Q}_{repeat} are obviously subsets of the set of contextuality sequences $\mathcal{Q}_{context}$ of the PM square. Nevertheless, an automaton that simultaneously obeys \mathcal{Q}_{rc} and \mathcal{Q}_{repeat} already possesses more than two internal states, i.e. $M(\mathcal{Q}_{rc} \text{ and } \mathcal{Q}_{repeat}) > 1$. In order to see this, assume that the automaton has only two internal states and without loss of generality that it starts in state S_1 . We consider the case when in the last column there must be a prescribed state change in order to avoid a contradiction, i.e. in S_1 the product of the assignments of $Cc\gamma$ is +1, contrary to the quantum prediction. Note that there always exists at least one row or column with such a contradiction and that the proof for any row or column follows the same lines. If there is only one state change (say, after a measurement of γ), then while measuring the sequence $Cc\gamma$, the automaton would remain in S_1 until after the last column (say, c and γ), both must go to S_2 . Then, the constraints from \mathcal{Q}_{repeat} require that γ has the same values in S_1 and S_2 (this is also true for c). But then the sequence $Cc\gamma$ in \mathcal{Q}_{rc} will yield a contradiction. Thus a two-state automaton cannot obey both \mathcal{Q}_{rc} and \mathcal{Q}_{repeat} .

On the other hand, \mathcal{A}_3 is an example of a three-state automaton, which obeys \mathcal{Q}_{rc} and \mathcal{Q}_{repeat} . In fact, \mathcal{A}_3 obeys $\mathcal{Q}'_{context}$. In order to see this, it is enough to show that for any choice of the initial state, the automaton will obey $\mathcal{Q}_{context}$. So, we assume that S_1 is the initial state; the reasoning for S_2 and S_3 is similar. If we now measure a sequence with observables from the first row only, we may jump between the states S_1 and S_2 , but the output for all observables in the first row are the same for either state. A similar argument holds for all rows and the first and second columns. For a sequence with measurements from the third column, assume that the first observable in the sequence that is not γ , is the observable c. Then the state changes to S_3 , in

which the last column does not yield a contradiction. Since only the output C was changed, but C was not measured so far, we cannot get any contradiction. A similar argument can be used for the case when the first observable in the sequence that is not γ , is the observable C.

In summary, since any automaton that obeys $\mathscr{Q}_{\text{context}}$ has at least three states and \mathcal{A}_3 is a three-state automaton obeying the larger set $\mathscr{Q}'_{\text{context}}$, we have shown that \mathcal{A}_3 is memory-optimal for either set.

Appendix B. A_4 obeys $\mathscr{Q}'_{\text{context}}$ and $\mathscr{Q}'_{\text{compat}}$

In this appendix, we demonstrate that the automaton \mathcal{A}_4 indeed obeys $\mathscr{Q}'_{\text{context}}$ and $\mathscr{Q}'_{\text{compat}}$. The proof for $\mathscr{Q}'_{\text{context}}$ is completely analogous to the one in appendix A.

For \mathscr{Q}'_{compat} , we consider a fixed observable, e.g. *B*. Then S_1 and S_2 yield +1, whereas S_3 and S_4 give -1. However, using arbitrary measurements compatible with *B* (i.e. *A*, *B*, *C*, *b* and β), we can never reach S_3 or S_4 if we start from S_1 or S_2 and vice versa. Hence no contradiction occurs for any sequence of the type $\lfloor T \rfloor B \lfloor X_1 X_2 \ldots \rfloor B$. A similar argument holds for all observables if we note, in addition, that e.g. after a measurement of *C* the automaton can only be in S_2 or S_3 .

Appendix C. Definitions and basic rules used in the optimality proofs

As we have already done in the main text, we denote the observables from the PM square by

$$\begin{bmatrix} A & B & C \\ a & b & c \\ \alpha & \beta & \gamma \end{bmatrix}.$$
 (C.1)

Furthermore, we denote the rows of the square by \mathbf{R}_i and the columns by \mathbf{C}_i . The value table of each memory state *i* is denoted by T_i and the update table by U_i . We write an entry of zero in U_i if the state does not change for that observable. Furthermore, we write measurement sequences as $A_1^+B_2^-C_2^-a_3^+$ meaning that when the sequence ABCa was measured, the results were +, -, -, +, and the memory was initially in state S_1 and changed like $S_1 \mapsto S_2 \mapsto S_2 \mapsto S_3$.

It will be useful for our later discussion to note some rules about the structure of the value and update tables.

Sign flips. Let us assume that we have an automaton obeying ℒ_{context} and ℒ_{compat} (or some subset of those sets) and pick a 2 × 2 square of observables (e.g. the set {A, B, a, b} or {A, B, α, β} or {A, C, α, γ}). Then, if we flip in each T_i the signs corresponding to these observables, we will obtain another valid automaton.

This holds true, because the mentioned sign flips do not change any of the certain quantum predictions from $\mathscr{Q}'_{\text{context}}$ or $\mathscr{Q}'_{\text{compat}}$. This rule will allow us to later fix one or two entries in a given value table T_i .

2. *Number of contradictions*. Any table T_i contains either one, three or five contradictions to the row and column constraints. This follows directly from the fact that any fixed assignment fulfills $\prod_k \mathbf{R}_k \mathbf{C}_k = +1$, while

the row and column constraints require $\prod_k \mathbf{R}_k \mathbf{C}_k = -1$.

3. Condition for fixing the memory. Let us assume that we have an automaton obeying $\mathscr{Q}'_{\text{context}}$ and let there be a table T_i which assigns to an observable (say A) a value different from all

other tables. Then, the update table U_i must contain only zeros in the corresponding row and column (here, \mathbf{R}_1 and \mathbf{C}_1).

The observables in the row and column correspond to compatible observables, which are not allowed to change the value of the first observable. However, any change of the memory state would change the value, as T_i is the only table with the initial assignment.

- 4. Contradictions and transformations. Let us assume that we have an automaton obeying $\mathscr{Q}'_{\text{context}}$ and let there be in T_i some contradiction in the column \mathbf{C}_j (or the row \mathbf{R}_j). Then, in the update table U_i there cannot be two zeros in the column \mathbf{C}_j (or the row \mathbf{R}_j). If there were two zeros, it could happen that one measures two entries of \mathbf{C}_j without changing the memory state. But then measuring the third one will reveal the contradiction in T_i . (Note that the automaton first provides the result and then updates its state.)
- 5. Contradictions and other tables. Let us assume that we have an automaton obeying $\mathscr{Q}'_{\text{context}}$ and let there be in T_i a contradiction in the column \mathbf{C}_j (or the row \mathbf{R}_j). Then, there must be two different tables T_k and T_l where in both the column \mathbf{C}_j has no contradictions anymore, but the assignments of T_k and T_l differ in two observables of \mathbf{C}_j . Furthermore, in the column \mathbf{C}_i of the update table U_i there must be two entries leading to two different states.

First, note that there must be at least one other table T_k where the contradiction does not exist anymore. This follows from the fact that we may measure C_j starting from the memory state *i*. After having made these measurements, we arrive at some state *k*, and from the contextuality correlations $\mathscr{Q}'_{\text{context}}$ it follows that C_j in T_k has no contradiction.

The table T_k differs from T_i in at least one observable X in C_j . On the other hand, starting from T_i one might measure X as a first observable. Then, making further measurements on C_j one must arrive at a table T_l without a contradiction. Since T_k and T_l have both no contradiction, they must differ in at least two places, one of them being X. Finally, if the column C_j in U_i would only have entries of zero and k, then C_j in T_k could not differ from T_i . This eventually leads to a contradiction and hence proves the last assertion.

Appendix D. A_4 is memory-optimal

Here, we prove the optimality of the four-state automaton A_4 , in the sense of obeying $\mathscr{Q}'_{\text{context}}$ and $\mathscr{Q}'_{\text{compat}}$ with a minimum number of states. We use the definitions and rules as introduced in appendix C.

Let us assume that we would have a three-state automaton obeying $\mathscr{Q}'_{context}$. T_1 has a contradiction, and we can assume, without loss of generality, that it is C_3 . Then, according to rule 1 we can, without loss of generality, assume that all entries in C_3 are '+ '. Together with rule 5 this leads to the conclusion that the three states T_i are, without loss of generality, of the form

$$T_1: \begin{bmatrix} & + \\ & + \\ & + \end{bmatrix}, \quad T_2: \begin{bmatrix} & + \\ & + \\ & - \end{bmatrix}, \quad T_3: \begin{bmatrix} & + \\ & - \\ & + \end{bmatrix}, \quad (D.1a)$$

$$U_1:\begin{bmatrix} & 2\\ & 3 \end{bmatrix}, \quad U_2:\begin{bmatrix} & 0\\ & 0\\ 0 & 0 & 0 \end{bmatrix}, \quad U_3:\begin{bmatrix} & 0\\ 0 & 0 & 0\\ & 0 \end{bmatrix}.$$
 (D.1*b*)

Here, empty places in the tables mean that the corresponding entries are not yet fixed. The table U_1 follows from rule 5, and U_2 and U_3 follow from rule 3.

Which can be the entries corresponding to the observables a and b in U_2 ? Since T_3 assigns a different value to c than T_2 , there cannot be a '3' at these entries; otherwise, a sequence like $c_2^+a_2^2c_3^-$ would lead to a contradiction to the conditions of $\mathscr{Q}_{\text{context}}$.

But there can also not be a '1' at these entries, because then the sequence $c_2^+ a_2^2 \gamma_1^+ c_3^-$ yields a contradiction to \mathcal{Q}_{compat} , since *a* and γ are compatible with *c*. So the entries of \mathbf{R}_2 in U_2 must be zero, as there are only three states in the memory. A similar argument can be applied to U_3 , showing that here \mathbf{R}_3 must be zero.

So the tables have to be of the form

$$T_1:\begin{bmatrix} & +\\ & +\\ & +\end{bmatrix}, \quad T_2:\begin{bmatrix} & +\\ & +\\ & -\end{bmatrix}, \quad T_3:\begin{bmatrix} & +\\ & -\\ & +\end{bmatrix}, \quad (D.2a)$$

$$U_1:\begin{bmatrix} & 2\\ & 3 \end{bmatrix}, \quad U_2:\begin{bmatrix} & 0\\ 0 & 0 & 0\\ 0 & 0 & 0 \end{bmatrix}, \quad U_3:\begin{bmatrix} & 0\\ 0 & 0 & 0\\ 0 & 0 & 0 \end{bmatrix}.$$
(D.2b)

Now, according to rule 4, the contradictions in T_2 as well as in T_3 can only be in \mathbf{R}_1 . But, according to rule 5, if there is a contradiction in \mathbf{R}_1 of T_2 , there must be two different T_i and T_j where there is no contradiction in \mathbf{R}_1 . But there is only one table left, namely T_1 , and we arrive at a contradiction.

Appendix E. Proof that the classical simulation of the extended PM square requires more than two bits of memory

Let us now discuss the extended PM square from [29]. Again, we refer to appendix C for basic definitions and rules. As already mentioned, one considers for that the array of observables

$$\begin{bmatrix} \chi_{01} & \chi_{02} & \chi_{03} \\ \chi_{10} & \chi_{11} & \chi_{12} & \chi_{13} \\ \chi_{20} & \chi_{21} & \chi_{22} & \chi_{23} \\ \chi_{30} & \chi_{31} & \chi_{32} & \chi_{33} \end{bmatrix} = \begin{bmatrix} \mathbb{1} \otimes \sigma_x & \mathbb{1} \otimes \sigma_y & \mathbb{1} \otimes \sigma_z \\ \sigma_x \otimes \mathbb{1} & \sigma_x \otimes \sigma_x & \sigma_x \otimes \sigma_y & \sigma_x \otimes \sigma_z \\ \sigma_y \otimes \mathbb{1} & \sigma_y \otimes \sigma_x & \sigma_y \otimes \sigma_y & \sigma_y \otimes \sigma_z \\ \sigma_z \otimes \mathbb{1} & \sigma_z \otimes \sigma_x & \sigma_z \otimes \sigma_y & \sigma_z \otimes \sigma_z \end{bmatrix}.$$
(E.1)

These observables can be grouped into trios, in which the observables commute and their product equals $\pm \mathbb{1}$. Nine trios are of the form { χ_{k0} , χ_{kl} , χ_{0l} }; three trios where the product equals + $\mathbb{1}$ are { χ_{11} , χ_{23} , χ_{32} }, { χ_{12} , χ_{21} , χ_{33} } and { χ_{13} , χ_{22} , χ_{31} }. Three trios where the product equals - $\mathbb{1}$ are { χ_{11} , χ_{22} , χ_{33} }, { χ_{12} , χ_{23} , χ_{31} } and { χ_{13} , χ_{21} , χ_{32} }. From this, one can derive the inequality

$$\sum_{k,l} \langle \chi_{k0} \chi_{kl} \chi_{0l} \rangle + \langle \chi_{11} \chi_{23} \chi_{32} \rangle + \langle \chi_{12} \chi_{21} \chi_{33} \rangle + \langle \chi_{13} \chi_{22} \chi_{31} \rangle - \langle \chi_{11} \chi_{22} \chi_{33} \rangle - \langle \chi_{12} \chi_{23} \chi_{31} \rangle - \langle \chi_{13} \chi_{21} \chi_{32} \rangle \leqslant 9$$
(E.2)

for noncontextual models, while QM predicts a value of 15, independently of the state.

First note that in this new inequality 15 terms (or contexts) occur but any noncontextual model can fulfill the quantum prediction for only 12 of them at most, so three contradictions cannot be avoided. One can directly check that in the whole construction of the inequality, ten different PM squares occur. Nine of them are a simple rewriting of the usual PM square, while the 10th comes from the observables χ_{kl} with $k, l \neq 0$. Any of the 15 terms in the inequality contributes to four of these PM squares.

Any value table for the 15 observables leads to assignments to the 15 contexts, but it has at least three contradictions. As any context contributes to four PM squares, this would lead to 12 contradictions in the 60 contexts of the ten PM squares, if we consider them separately. Since in a PM square the number of contradictions cannot be two (rule 2), this means that one of the PM squares has to have three contradictions.

Let us now assume that we have a valid automaton for this extended PM square with four memory states. Of course, this would immediately give a valid four-state automaton of any of the ten PM squares. For one of these PM squares, at least one table has to have *three* contradictions. So it suffices to prove the following lemma:

Lemma 4. There is no four-state automaton obeying \mathscr{Q}'_{compat} and $\mathscr{Q}'_{context}$, where one table T_i has three contradictions.

In the course of proving this lemma we will also prove the following:

Proposition 5. The four-state automaton A_4 is unique, up to some permutation or sign changes.

To prove the lemma, we proceed in the following way. Without loss of generality, we can assume that the first three tables T_i look like the T_i in equation (D.1a). Then, we can add a fourth table T_4 . For the last column of this table, there are $2^3 = 8$ possible values. We will investigate all eight possibilities and show either that we arrive directly at a contradiction or that only an automaton similar to A_4 is possible, in which any table has only one contradiction. This proves the lemma.

We will first deal with the four cases where T_4 has also a contradiction in C_3 . This will lead to observation 6, which will be useful in the following four cases.

Case 1: *For* T_4 *one has* $[C, c, \gamma] = [+++]$.

In this case, a simple application of the previous rules implies that several entries are fixed:

$$T:\begin{bmatrix} & +\\ & +\\ & +\end{bmatrix}, \begin{bmatrix} & +\\ & +\\ & -\end{bmatrix}, \begin{bmatrix} & +\\ & -\\ & +\end{bmatrix}, \begin{bmatrix} & +\\ & +\\ & +\end{bmatrix}, \quad (E.3)$$
$$U:\begin{bmatrix} & 2\\ & 2\end{bmatrix}, \begin{bmatrix} 0 & 0\\ 0 & 0\end{bmatrix}, \begin{bmatrix} 0 & 0\\ 0 & 0\end{bmatrix}, \begin{bmatrix} 0 & 0\\ 0 & 0\end{bmatrix}, \begin{bmatrix} 2\\ & 2\end{bmatrix}, \quad (E4)$$

$$U: \begin{bmatrix} 2\\ 3 \end{bmatrix}, \begin{bmatrix} 0 & 0 & 0\\ 0 & 0 & 0 \end{bmatrix}, \begin{bmatrix} 0 & 0 & 0\\ 0 & 0 & 0 \end{bmatrix}, \begin{bmatrix} 2\\ 3 \end{bmatrix}.$$
(E.4)

Here and in the following, we write the T_i and U_i just as a row for notational simplicity, starting from T_1 to T_4 . The entries in U_1 and U_4 are fixed from the following reasoning: let us assume that one measures c in T_1 , then, since the values $C(T_i)$ are the same in all T_i , one has to change immediately to a table with no contradiction in C_3 , and where the value of c is still the same. The only possibility is T_2 . Furthermore, \mathbf{R}_2 in U_2 and \mathbf{R}_3 in U_3 must be zero due to the same argument that led to equation (D.2b).

It follows (rule 4) that T_2 and T_3 have both exactly one contradiction, which must be in \mathbf{R}_1 . So, in $\mathbf{R}_1(U_2)$ there must be the entries '1' and '4' (an entry '3' would not solve the problem, because in $\mathbf{R}_1(T_3)$ has also a contradiction). As we can still permute the first and the second column, we can without loss of generality assume that the first row in U_2 is [1 4 0]. Due to

rule 1, we can also assume, without loss of generality, that $A(T_2) = +$. Similarly, in $\mathbf{R}_1(U_3)$ there must be the entries '1' and '4', resulting in two different cases:

If $\mathbf{R}_1(U_3) = [140]$, we must have the following tables,

$$T:\begin{bmatrix} + & + \\ + & + \\ + & + \end{bmatrix}, \begin{bmatrix} + & - & + \\ - & + \\ - & - \end{bmatrix}, \begin{bmatrix} + & - & + \\ - & - \\ + & + \end{bmatrix}, \begin{bmatrix} - & + \\ + & + \\ + & + \end{bmatrix}, \quad (E.5)$$
$$U:\begin{bmatrix} & 2\\ & 2\\ & 3\end{bmatrix}, \begin{bmatrix} 1 & 4 & 0\\ 0 & 0 & 0\\ 0 & 0 & 0 \end{bmatrix}, \begin{bmatrix} 1 & 4 & 0\\ 0 & 0 & 0\\ 0 & 0 & 0 \end{bmatrix}, \begin{bmatrix} 1 & 4 & 0\\ 0 & 0 & 0\\ 0 & 0 & 0 \end{bmatrix}, \begin{bmatrix} 2\\ & 2\\ & 3 \end{bmatrix}, \quad (E.6)$$

where the added values in \mathbf{R}_1 of the T_i follow from $\mathbf{R}_1(U_2)$ and $\mathbf{R}_1(U_3)$.

Now, if we start from T_2 and measure the sequence $a_2A_2a_1$, we see that we must have $a(T_1) = a(T_2)$. Similarly, from T_3 we can measure $a_3A_3a_1$, implying that $a(T_1) = a(T_2) = a(T_3)$. Similarly, we find that $b(T_2) = b(T_3) = b(T_4)$. But this gives a contradiction: in $\mathbf{R}_2(T_2)$ and $\mathbf{R}_2(T_3)$ there is no contradiction and $c(T_2) \neq c(T_3)$. Therefore, it cannot be that $a(T_2) = a(T_3)$ and at the same time $b(T_2) = b(T_3)$.

As the second case, we have to consider the possibility that $\mathbf{R}_1(U_3) = [4\ 1\ 0]$. Then, also the values of $\mathbf{R}_1(T_3)$ must be interchanged, $\mathbf{R}_1(T_3) = [-++]$. Then, starting from T_2 , the sequence $\alpha_2 A_2 \gamma_1 \alpha_3$ shows directly that $\alpha(T_2) = \alpha(T_3)$. Similarly, starting from T_3 , the sequence $a_3 A_3 c_4 a_2$ shows that $a(T_2) = a(T_3)$. But since $A(T_2) \neq A(T_3)$, this is a contradiction.

Case 2: *For* T_4 *one has* $[C, c, \gamma] = [+ - -].$

As in case 1, one can directly see that several entries are fixed:

$$T:\begin{bmatrix} & +\\ & +\\ & +\end{bmatrix}, \begin{bmatrix} & +\\ & +\\ & -\end{bmatrix}, \begin{bmatrix} & +\\ & -\\ & +\end{bmatrix}, \begin{bmatrix} & +\\ & -\\ & +\end{bmatrix}, \begin{bmatrix} & +\\ & -\\ & -\end{bmatrix}, \quad (E.7)$$

$$U: \begin{bmatrix} 2\\ 3\\ 3 \end{bmatrix}, \begin{bmatrix} 0\\ 0& 0\\ 0& 0\\ 0& 0 \end{bmatrix}, \begin{bmatrix} 0\\ 0& 0\\ 0& 0\\ 0& 0 \end{bmatrix}, \begin{bmatrix} 3\\ 2\\ 2 \end{bmatrix}.$$
(E.8)

The zeros in U_2 and U_3 come from the following argumentation: starting from T_1 , the measurement sequence $c_1^+X_2c_2$ with X compatible with c shows that in $\mathbf{R}_2(U_2)$ and $\mathbf{C}_3(U_2)$ there can be no '3' or '4'. But there can also be no '1', because then the sequence $c_1^+X_2\gamma_1c_3^-$ would lead to a contradiction. Therefore, $\mathbf{R}_2(U_2)$ and $\mathbf{C}_3(U_2)$ have to be zero. Starting from T_4 and measuring γ one can similarly prove that the entries for $\mathbf{R}_3(U_2)$ have to be zero and analogous arguments also prove the zeros in U_3 .

It is now clear (rule 4) that the contradictions in T_2 and T_3 have to be in \mathbf{R}_1 and the missing entries in U_2 and U_3 can only be '4' and '1'. As we still can permute the first and the second column, there are only two possibilities.

Case 2A: Firstly, we consider the case when $\mathbf{R}_1(U_2) = \mathbf{R}_1(U_3) = [1 4 0]$.

As in case 1, we can directly see that $a(T_2) = a(T_1) = a(T_3)$ and $b(T_2) = b(T_4) = b(T_3)$. Hence, $\mathbf{R}_2(T_2)$ and $\mathbf{R}_2(T_3)$ differ exactly in the value of *c*, but in both cases there is no contradiction in \mathbf{R}_2 . This is not possible.

Case 2B: Secondly, we consider the case when the first rows of U_2 and U_3 differ, and we take $\mathbf{R}_1(U_2) = [140]$ and $\mathbf{R}_1(U_3) = [410]$. Then, we apply rule 1 to fix for $A(T_3) = a(T_3) = +$. Then, the tables have to be

$$T:\begin{bmatrix} -& -& +\\ &-& +\\ &+& + \end{bmatrix}, \begin{bmatrix} -& +& +\\ &+& +\\ &-& -\end{bmatrix}, \begin{bmatrix} +& -& +\\ +& -& -\\ +& +& + \end{bmatrix}, \begin{bmatrix} +& +& +\\ +& -& -\\ +& -& -\end{bmatrix}, \quad (E.9)$$
$$U:\begin{bmatrix} && 2\\ && 3\end{bmatrix}, \begin{bmatrix} 1& 4& 0\\ 0& 0& 0\\ 0& 0& 0\end{bmatrix}, \begin{bmatrix} 4& 1& 0\\ 0& 0& 0\\ 0& 0& 0\end{bmatrix}, \begin{bmatrix} && 3\\ && 2\end{bmatrix}. \quad (E.10)$$

Here, $C_2(T_1)$ and $C_1(T_4)$ come from measurement sequences like $a_3^+A_3^+a_4^+$, starting from T_3 .

Again, we have two possibilities for the value of b in T_2 . If we set $b(T_2) = -$, then all values in all T_i are fixed and each table has exactly one contradiction. This is, up to some relabeling, the four-state automaton A_4 from the main text (indeed, this is the way how this solution was found). If we set $b(T_2) = +$, then also all T_i can be filled, and we must have

$$T:\begin{bmatrix} -& -& +\\ +& -& +\\ -& +& + \end{bmatrix}, \begin{bmatrix} -& +& +\\ +& +& +\\ -& +& -\end{bmatrix}, \begin{bmatrix} +& -& +\\ +& -& -\\ +& +& + \end{bmatrix}, \begin{bmatrix} +& +& +\\ +& +& -\\ +& +& -\end{bmatrix}, (E.11)$$
$$U:\begin{bmatrix} 3 & 2\\ 2 & 3 \end{bmatrix}, \begin{bmatrix} 1 & 4 & 0\\ 0 & 0 & 0\\ 0 & 0 & 0 \end{bmatrix}, \begin{bmatrix} 4 & 1 & 0\\ 0 & 0 & 0\\ 0 & 0 & 0 \end{bmatrix}, \begin{bmatrix} 2 & 3\\ 3 & 2 \end{bmatrix}. (E.12)$$

Here, the tables T_1 and T_4 have three contradictions (two new ones in \mathbf{R}_2 and \mathbf{R}_3) and the new entries in U_1 and U_4 must be introduced according to rule 5 (note that $a(T_i)$ and $\beta(T_i)$ are for all tables the same). Then, however, starting from T_1 , the sequence $\alpha_1^- A_2^- \gamma_1^+ \alpha_3^+$ shows that this is not a valid solution.

Case 3: *For* T_4 *one has* $[C, c, \gamma] = [-+-]$.

In this case, a simple reasoning according to the usual rules fixes the entries:

$$T: \begin{bmatrix} & +\\ & +\\ & +\end{bmatrix}, \begin{bmatrix} & +\\ & +\\ & -\end{bmatrix}, \begin{bmatrix} & +\\ & -\\ & +\end{bmatrix}, \begin{bmatrix} & -\\ & +\\ & -\end{bmatrix}, \quad (E.13)$$

$$U:\begin{bmatrix} & & \\ & & 3 \end{bmatrix}, \begin{bmatrix} & & \\ & & \end{bmatrix}, \begin{bmatrix} & & 0 \\ 0 & 0 & 0 \\ & & 0 \end{bmatrix}, \begin{bmatrix} 0 & 0 & 0 \\ & & 0 \\ & & 0 \end{bmatrix}.$$
(E.14)

Here we have an obvious contradiction in T_4/U_4 : $C_3(T_4)$ contains a contradiction, but (due to rule 3) one is not allowed to change the memory state when measuring it. Therefore, the memory can never be in state 4. But then, one would have effectively a three-state solution, which is not possible, as we already know.

Case 4: *For* T_4 *one has* $[C, c, \gamma] = [--+]$ *.*

This is the same as case 3, where \mathbf{R}_2 and \mathbf{R}_3 have been interchanged.

Now we have dealt with all the cases where T_4 contains a contradiction in \mathbb{C}_3 , just as T_1 . We have seen that in these cases there can only be a solution if each table contains exactly one contradiction, and this solution is unique, up to some permutations or sign flips. Moreover, we could have made the same discussion with rows instead of columns. Therefore from the first four cases, we can state an observation that will be useful in the remaining four cases:

Observation 6. If in any four-state solution two tables T_i and T_j have both a contradiction in the same column C_k (or row R_k), then there has to be exactly one contradiction in each value table of the automaton.

So, if there is a four-state solution where one table has three contradictions, then it cannot be that two tables have both a contradiction in the same column or row.

Then we can proceed with the remaining cases.

Case 5: *For* T_4 *one has* $[C, c, \gamma] = [++-]$.

This is the critical case, as it is difficult to distinguish the tables T_2 and T_4 here. First, the following entries are directly fixed:

$$T: \begin{bmatrix} & +\\ & +\\ & +\end{bmatrix}, \begin{bmatrix} & +\\ & +\\ & -\end{bmatrix}, \begin{bmatrix} & +\\ & -\\ & +\end{bmatrix}, \begin{bmatrix} & +\\ & +\\ & -\end{bmatrix}, \quad (E.15)$$

$$U: \begin{bmatrix} 2\\ 3 \end{bmatrix}, \begin{bmatrix} 0|4 & 0|4\\ 0|4 & 0|4 \end{bmatrix}, \begin{bmatrix} 0\\ 0 & 0 & 0\\ 0 & 0 & 0 \end{bmatrix}, \begin{bmatrix} 0|2\\ 0|2 \end{bmatrix}.$$
(E.16)

Here, $c(U_1) = 2$ has been chosen without loss of generality. It is clear that $c(U_1) = 2$ or $c(U_1) = 4$; as T_2 and T_4 are equivalent at the beginning, we can choose T_2 here. The entries of the type i|j in U_2 and U_4 mean that the numbers can be *i* or *j*, but nothing else. The values of $c(U_2)$ (and $c(U_4)$) cannot be 1, because then the sequence $c_2^+\gamma_1^+c_3^-$ directly reveals a contradiction. Furthermore, the zeros in $\mathbf{R}_3(U_3)$ and $\mathbf{R}_2(U_3)$ follow similarly as equation (D.2*b*) or from rule 3. In addition, $C(U_2) \neq 1$, because otherwise the sequence $c_1^+C_2^+\gamma_1^+$ reveals a contradiction to the PM conditions. Also, $C(U_2) \neq 3$, because of $c_1^+C_2^+c_3^-$. Similarly, 1 and 3 are excluded as values for $a(U_2)$ and $b(U_2)$, due to the sequences $c_1^+a_2\gamma_1^+c_3^-$ and $c_1^+a_2^+c_3^-$.

Furthermore, we can use our observation 6: if in a four-state solution one column has a contradiction in two of the T_i , then there can be only one contradiction in any T_i . Here we can use it as follows: it is clear that T_3 has its contradiction in \mathbf{R}_1 . Since we aim to rule out a four-state solution where one table has three contradictions, we can assume that there is no contradiction in \mathbf{R}_1 in all the other T_i (especially in T_2 and T_3). Otherwise, we would already know that no solution exists with three contradictions in a table. We can distinguish two cases.

Case 5A: Let us assume that $\gamma(U_2) = 0$. Then, the tables must read:

$$T:\begin{bmatrix} & +\\ & +\\ & +\end{bmatrix}, \begin{bmatrix} & +\\ & +\\ & -\end{bmatrix}, \begin{bmatrix} & +\\ & -\\ & +\end{bmatrix}, \begin{bmatrix} & +\\ & +\\ & -\end{bmatrix}, \quad (E.17)$$

IOP Institute of Physics **D**EUTSCHE PHYSIKALISCHE GESELLSCHAFT

$$U:\begin{bmatrix} 2\\ 3\end{bmatrix}, \begin{bmatrix} 0|4\\ 0|4 & 0|4\\ 0|4 & 0|4\\ 0|4 & 0|4 & 0\end{bmatrix}, \begin{bmatrix} 0\\ 0 & 0 & 0\\ 0 & 0 & 0\end{bmatrix}, \begin{bmatrix} 0|2\\ 0|2\\ 0|2\end{bmatrix}.$$
(E.18)

The new entries in U_2 follow from $\gamma(U_2) = 0$ in combination with $\gamma(T_1) = \gamma(T_3) \neq \gamma(T_2)$.

Due to rule 5, the table T_2 must have a contradiction in C_1 , C_2 or R_1 . From observation 6, we can assume that it is not in R_1 . Due to possible permutations of C_1 and C_2 we further assume without loss of generality that the contradiction is in C_1 . Then we have

$$U:\begin{bmatrix} 2\\ 3\end{bmatrix}, \begin{bmatrix} 1 & 0|4\\ 4 & 0|4 & 0|4\\ 0|4 & 0|4 & 0\end{bmatrix}, \begin{bmatrix} 0\\ 0 & 0 & 0\\ 0 & 0 & 0\end{bmatrix}, \begin{bmatrix} 0|2\\ 0|2\end{bmatrix}.$$
 (E.19)

We cannot have $A(U_2) = 3$, since there is a contradiction in $\mathbf{R}_1(T_3)$ and $C(T_i) = +$ for all tables. In addition, due to rule 5, it is not possible that $A(U_2) = 4$. Finally, we choose $a(U_2) = 4$; the other option would be $\alpha(U_2) = 4$; this will be discussed below.

From observation 6 we can conclude that $C_1(T_1)$ and $C_1(T_4)$ do not contain contradictions, since $C_1(T_2)$ contains already a contradiction. So $C_1(T_1)$ and $C_1(T_4)$ must differ in two places (rule 5). One of these places must be $A(T_1) \neq A(T_4)$. Let us assume that the second one is $a(T_1) \neq a(T_4)$; the other case $[\alpha(T_1) \neq \alpha(T_4)]$ will be discussed below. Then, we can conclude that in $\mathbf{R}_1(U_1)$ and $\mathbf{C}_1(U_1)$ we cannot have the entries '2' and '4', and in $\mathbf{R}_2(U_4)$ and $\mathbf{C}_1(U_4)$ we cannot have the entries '2' and '1'. To see this, note that we must have $A(T_2) = A(T_1) \neq A(T_4)$ and, if $B(U_1) = 2$, we can consider the measurement sequence $A_2B_1a_2A_4$ or, if $B(U_1) = 4$, the sequence $A_2B_1A_4$, etc. Hence, we have

$$U:\begin{bmatrix} 0 & 0 & 0 \\ 0|3 & 2 \\ 0|3 & 3 \end{bmatrix}, \begin{bmatrix} 1 & 0|4 \\ 4 & 0|4 & 0|4 \\ 0|4 & 0|4 & 0 \end{bmatrix}, \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}, \begin{bmatrix} 0|3 & 0 \\ 0|3 & 0|3 & 0 \\ 0|3 & 0|2 \end{bmatrix}.$$

Here, we used in $\mathbf{R}_1(U_1)$ that $\mathbf{R}_1(T_3)$ has a contradiction and $C(T_i) = +$ for all tables, so it is not possible to go there.

Now, by rule 1, we may fix $A(T_2) = a(T_2) = +$. Then we arrive at

$$T:\begin{bmatrix} + & + & + \\ & & + \\ & & + \end{bmatrix}, \begin{bmatrix} + & + & + \\ + & + & + \\ - & + & - \end{bmatrix}, \begin{bmatrix} & & + \\ & - \\ & + \end{bmatrix}, \begin{bmatrix} - & - & + \\ + & & + \\ - & - & - \end{bmatrix}, \quad (E.20)$$

$$U:\begin{bmatrix} 0 & 0 & 0 \\ 0|3 & 2 \\ 0|3 & 3 \end{bmatrix}, \begin{bmatrix} 1 & 0|4 \\ 4 & 0|4 & 0|4 \\ 0|4 & 0|4 & 0 \end{bmatrix}, \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}, \begin{bmatrix} 0|3 & 0|2 \\ 0 & 0 & 0 \\ 0|3 & 0|2 \end{bmatrix}.$$
 (E.21)

Here, we must have $A(T_4) = \alpha(T_4)$ since $C_1(T_4)$ has no contradiction. Furthermore, $\mathbf{R}_1(T_4)$ has no contradiction due to observation 6. The values of $\mathbf{R}_2(U_4)$ are determined by considering sequences like $c_4a_4c_7$; and $C(U_4) \neq 3$, because of $c_4^+C_4c_3^-$, and $C(U_4) \neq 1$, because of $c_4^+C_4\gamma_1c_3^-$.

In addition, we can conclude that $A(U_4) = 0$ and $B(U_4) = 0$, since $\mathbf{R}_1(T_3)$ has a contradiction and $C(T_i) = +$ for all tables, so it is not possible to go there. Then we can fill T_4 completely. Then, also $C(U_4) = 0$; otherwise the sequence $B_4^- C_4^+ B_2^+$ gives a contradiction. If we had $\alpha(U_4) = 3$, then we must have $A(T_4) = A(T_3) = -$ and, consequently (rule 5) $B(U_3) = 1$

16

or 2, but then the sequence $A_4^- \alpha_4^- B_3^+ A_{1|2}^+$ leads to a contradiction, so $\alpha(U_4) = 0$. In summary, we have

$$T:\begin{bmatrix} + & + & + \\ & & + \\ & & + \end{bmatrix}, \begin{bmatrix} + & + & + \\ + & + & + \\ - & + & - \end{bmatrix}, \begin{bmatrix} & & + \\ - & + \\ + & + \end{bmatrix}, \begin{bmatrix} - & - & + \\ + & + & + \\ - & - & - \end{bmatrix}, \quad (E.22)$$

$$U:\begin{bmatrix} 0 & 0 & 0\\ 0|3 & 2\\ 0|3 & 3 \end{bmatrix}, \begin{bmatrix} 1 & 0|4\\ 4 & 0|4 & 0|4\\ 0|4 & 0|4 & 0 \end{bmatrix}, \begin{bmatrix} 0 & 0 & 0\\ 0 & 0 & 0\\ 0 & 0 & 0 \end{bmatrix}, \begin{bmatrix} 0 & 0 & 0\\ 0 & 0 & 0\\ 0 & 0|2 \end{bmatrix}.$$
 (E.23)

Now T_1 is the only candidate for a table with three contradictions. In order to obey observation 6, the only possibilities for contradictions are C_2 , C_3 and R_2 , since T_4 has its contradiction in R_3 . In particular, there must be a contradiction in $C_2(T_1)$. Then, in order to obey rule 5, we must have

However, if $b(U_1) = 3$, then the sequence $B_1^+ b_1 A_3 B_4^-$ leads to a contradiction, while if $\beta(U_1) = 3$, then the sequence $B_1^+ \beta_1 A_3 B_4^-$ leads to a problem.

Finally, if we had taken $\alpha(U_2) = 4$ or $\alpha(T_1) \neq \alpha(T_4)$ the proof would proceed along the same lines, but this time the contradiction in T_4 would be in the second row.

Case 5B: Let us assume that $\gamma(U_2) = 4$. Then, many entries on U_4 are fixed and we have

$$T:\begin{bmatrix} & +\\ & +\\ & +\\ & +\end{bmatrix}, \begin{bmatrix} & +\\ & +\\ & -\end{bmatrix}, \begin{bmatrix} & +\\ & -\\ & +\end{bmatrix}, \begin{bmatrix} & +\\ & +\\ & -\end{bmatrix}, (E.26)$$
$$U:\begin{bmatrix} & 2\\ & 2\\ & 3\end{bmatrix}, \begin{bmatrix} 0|4 & 0|4 & 0|4\\ & 4\end{bmatrix}, \begin{bmatrix} 0 & 0\\ 0 & 0 & 0\\ 0 & 0 & 0\end{bmatrix}, \begin{bmatrix} 0|2\\ 0|2 & 0|2 & 0|2\\ 0|2 & 0|2 & 0|2\end{bmatrix}. (E.27)$$

Here we cannot have $a(U_4) = 1$, due the sequences $c_2\gamma_2 a_4c_1$ (if $c(U_2) = 0$) or $c_2 a_4c_1$ (if $c(U_2) = 4$), and also not $a(U_4) = 3$, due to similar sequences. The same arguments apply to $b(U_4)$. The entries in $\mathbf{R}_3(U_4)$ and $\mathbf{C}_3(4)$ come from possible sequences like $\gamma_4 \alpha_4 \gamma_2$ if $\gamma(U_4) = 0$ or $\gamma_4 \gamma_2 \alpha_4 \gamma_2$ if $\gamma(U_4) = 2$.

But then the proof can proceed exactly as in case 5A, with T_2 and T_4 interchanged: the only significant difference comes from $c(U_1) = 2 \neq 4$, but this was never used in the proof.

Case 6: For T_4 one has $[C, c, \gamma] = [+-+]$: this is the same as case 5 with a permutation of \mathbf{R}_2 and \mathbf{R}_3 .

Case 7: *For* T_4 *one has* $[C, c, \gamma] = [-++]$ *.*

In this case, the tables read

$$T:\begin{bmatrix} & +\\ & +\\ & +\\ & +\end{bmatrix}, \begin{bmatrix} & +\\ & +\\ & -\end{bmatrix}, \begin{bmatrix} & +\\ & -\\ & +\\ & +\end{bmatrix}, \begin{bmatrix} & -\\ & +\\ & +\\ & +\end{bmatrix}, \quad (E.28)$$
$$U:\begin{bmatrix} & 2\\ & 2\\ & 3\end{bmatrix}, \begin{bmatrix} & 0\\ & 0|4 & 0|4 & 0\\ & 0 & 0\end{bmatrix}, \begin{bmatrix} & 0\\ & 0 & 0\\ & 0|4 & 0|4 & 0\\ & 0|4 & 0|4 & 0\end{bmatrix}, \begin{bmatrix} & 0\\ & 0 & 0\\ & 0|2 & 0|2 & 0\\ & 0|3 & 0|3 & 0\\ \end{bmatrix}. \quad (E.29)$$

Here, the entries in U_1 have been chosen without loss of generality: From rules 4 and 5 it follows that one can restrict the attention to the cases where $C_3(U_1) = [, 2, 3]$, $C_3(U_1) = [, 2, 4]$ or $C_3(U_1) = [, 4, 3]$. We only consider the first possibility; in the other cases the proof is analogous and is left to the gentle reader as an exercise. The zeros in U_2 , U_3 and U_4 come from rule 3. The entries 0|2 in U_4 come from possible measurement sequences such as $c_4a_4c_3$ or $c_4a_4\gamma_1c_3$ which prove that there cannot be the entries '3' or '1'. The other entries can be derived accordingly.

From rule 5, it follows that in T_4 the contradiction cannot be in the rows, so it has to be in the first or second column. Let us assume, without loss of generality, that it is in $C_1(T_4)$. Further, we can assume without loss of generality that the values A and a in T_4 are both '+'. Then, the tables can be more specified as

$$T:\begin{bmatrix} & +\\ & +\\ & +\\ & +\end{bmatrix}, \begin{bmatrix} + & +\\ + & +\\ + & -\\ & -\end{bmatrix}, \begin{bmatrix} + & +\\ - & +\\ - & -\\ - & +\end{bmatrix}, \begin{bmatrix} + & -\\ - & +\\ - & -\\ + & +\\ - & -\\ + \end{bmatrix}, \quad (E.30)$$

$$U:\begin{bmatrix} 2\\ 3\end{bmatrix}, \begin{bmatrix} 0|1 & 0\\ 0 & 0|4 & 0\\ 0 & 0 & 0\end{bmatrix}, \begin{bmatrix} 0|1 & 0\\ 0 & 0 & 0\\ 0 & 0|4 & 0\end{bmatrix}, \begin{bmatrix} 0&0 & 0\\ 2 & 0|2 & 0\\ 3 & 0|3 & 0\end{bmatrix}.$$
 (E.31)

To see this, one first fills T_4 ; then, together with the entries of $C_1(U_4)$, many values of T_2 and T_3 are fixed. The entries out are justified similar to the reasoning above.

In T_2 as well as in T_3 the contradiction has to be in either \mathbf{R}_1 or \mathbf{C}_2 . However, there cannot be a contradiction in \mathbf{R}_1 . To see this, assume that there was a contradiction in $\mathbf{R}_1(T_2)$. Then, starting from T_2 we may measure the sequence C_2A_2B or C_2B_2A . According to rule 5, we must end in two different T_i . But the memory state can never change to T_4 (because $C(T_4) = -$). So we must have $B(U_2) = 3$, but this will not escape the contradiction, since the values for A and C coincide in T_2 and T_3 . So there is only T_1 left, and we arrive at a contradiction.

Consequently, the contradictions have to be in both $C_2(T_2)$ and $C_2(T_3)$. In principle, our observation 6 implies already that we cannot find a solution with three contradictions in one table. But one can also directly prove that there is no solution at all. We have

$$T:\begin{bmatrix} + & + & + \\ & + & + \\ & - & + \end{bmatrix}, \begin{bmatrix} + & + & + \\ + & + & + \\ + & - & - \end{bmatrix}, \begin{bmatrix} + & + & + \\ - & + & - \\ - & - & + \end{bmatrix}, \begin{bmatrix} + & - & - \\ + & + & + \\ - & - & + \end{bmatrix}, (E.32)$$
$$U:\begin{bmatrix} & & \\ & 2 \\ & & 3 \end{bmatrix}, \begin{bmatrix} 0 & 1 & 1 & 0 \\ 0 & 4 & 0 \\ 0 & 0 & 0 \end{bmatrix}, \begin{bmatrix} 0 & 1 & 1 & 0 \\ 0 & 0 & 0 \\ 0 & 4 & 0 \end{bmatrix}, \begin{bmatrix} 0 & 0 & 0 \\ 2 & 0 & 2 \\ 3 & 0 & 3 & 0 \end{bmatrix}. (E.33)$$

IOP Institute of Physics DEUTSCHE PHYSIKALISCHE GESELLSCHAFT

Here, we must have $B(T_1) = B(T_2) = B(T_3) = +$ due to measurement sequences such as $B_2^+ B_1^+$ or $B_3^+ B_1^+$ and $\beta(T_1) = \beta(T_2)$ due to $\beta_2^- B_2^+ \beta_1^-$ and $b(T_1) = b(T_3)$ due to $b_3^+ B_3^+ b_1^+$. But then, starting from T_2 , the sequence $\beta_2^- B_2^+ b_1^+$ reveals a contradiction to the PM conditions.

Case 8: *For* T_4 *one has* $[C, c, \gamma] = [---].$

In this case, we directly have

$$T:\begin{bmatrix} & +\\ & +\\ & +\end{bmatrix}, \begin{bmatrix} & +\\ & +\\ & -\end{bmatrix}, \begin{bmatrix} & +\\ & -\\ & +\end{bmatrix}, \begin{bmatrix} & -\\ & -\\ & -\end{bmatrix}, \quad (E.34)$$

$$U:\begin{bmatrix} & 2\\ & 3\end{bmatrix}, \begin{bmatrix} & 0\\ 0 & 0 & 0\\ 0 & 0 & 0\end{bmatrix}, \begin{bmatrix} & 0\\ 0 & 0 & 0\\ 0 & 0 & 0\end{bmatrix}, \begin{bmatrix} 0 & 0 & 0\\ & 0\\ & 0\end{bmatrix}.$$
(E.35)

Starting from T_2 we may measure the sequence C_2A_2B or C_2B_2A . According to rule 5, we must end in two different T_i . But the memory state can change neither to T_4 (because $C(T_4) = -$) nor to T_3 (as $\mathbf{R}_1(T_3)$ contains a contradiction). So there is only T_1 left, and we arrive at a contradiction.

In summary, by considering all eight different cases we have shown that no four-state solution exists in which one table has three contradictions. This proves the claim.

Appendix F. A ten-state automaton obeying all sequences

In this appendix, we show an example of a ten-state automaton that obeys the set of all sequences \mathcal{Q}_{all} . For that, we define ten eigenstates of two compatible observables. We let $|A^-B^+\rangle$ be a quantum state with $A|A^-B^+\rangle = -|A^-B^+\rangle$ and $B|A^-B^+\rangle = +|A^-B^+\rangle$. In this fashion we define the ten states $|A^+B^+\rangle$, $|A^-B^+\rangle$, $|C^+c^+\rangle$, $|C^-c^+\rangle$, $|\gamma^+\beta^+\rangle$, $|\gamma^-\beta^+\rangle$, $|\alpha^+a^+\rangle$, $|\alpha^-a^+\rangle$, $|a^+b^+\rangle$ and $|B^+b^+\rangle$. Any measurement of an observable from the PM square projects with finite probability any state of the set onto another state of the set. If, e.g., the automaton is in state $|A^-B^+\rangle$ and we measure *c*, QM predicts a chance of 50% to get the outcome +1 yielding the state $|C^-c^+\rangle$, and a 50% chance to obtain -1 and the state $|C^-c^-\rangle$. The former state is in the set of the ten states and hence our automaton would return +1 and change to the state $|C^-c^+\rangle$. We furthermore define that, if both states predicted by QM are in the set of the ten states, then we prefer the state corresponding to the output of +1. Together with an arbitrary choice of the initial state, this completes the definition of the automaton. By construction, this automaton is deterministic and observables the definition of the automaton. By construction, this automaton is deterministic and observables the definition of the automaton. By construction, this automaton is deterministic and observables the definition of the automaton. By construction, this automaton is deterministic and observables the definition of the automaton. By construction, this automaton is deterministic and observables the definition of the automaton.

References

- [1] Specker E 1960 Dialectica 14 239–46
- [2] Kochen S and Specker E P 1967 J. Math. Mech. 17 59-87
- [3] Bell J S 1966 *Rev. Mod. Phys.* **38** 447–52
- [4] Mermin N D 1990 Phys. Rev. Lett. 65 3373-76
- [5] Heywood P and Redhead M L G 1983 Found. Phys. 13 481–99
- [6] Bechmann-Pasquinucci H and Peres A 2000 Phys. Rev. Lett. 85 3313-6
- [7] Spekkens R W 2008 Phys. Rev. Lett. 101 020401

- [8] Cabello A 2010 Phys. Rev. Lett. 104 220401
- [9] Aharon N and Vaidman L 2008 Phys. Rev. A 77 052310
- [10] Svozil K 2009 Phys. Rev. A 79 054306
- [11] Svozil K 2009 CDMTCS Research Report Series 353 March 2009
- [12] Horodecki K, Horodecki M, Horodecki P, Horodecki R, Pawłowski M and Bourennane M 2010 arXiv:1006.0468
- [13] Bell J S 1982 Found. Phys. 12 989–99
- [14] La Cour B R 2009 Phys. Rev. A 79 012102
- [15] Khrennikov A Y 2009 Contextual Approach to Quantum Formalism (Berlin: Springer)
- [16] Toner B F and Bacon D 2003 Phys. Rev. Lett. 91 187904
- [17] Pironio S 2003 Phys. Rev. A 68 062102
- [18] Buhrman H, Cleve R, Massar S and de Wolf R 2010 Rev. Mod. Phys. 82 665-98
- [19] Holevo A S 1973 Probl. Inf. Trans. 9 177-83
- [20] Galvão E F and Hardy L 2003 Phys. Rev. Lett. 90 087902
- [21] Peres A 1990 Phys. Lett. A 151 107-8
- [22] Cabello A 2008 Phys. Rev. Lett. 101 210401
- [23] Kirchmair G et al 2009 Nature 460 494–7
- [24] Amselem E, Rådmark M, Bourennane M and Cabello A 2009 Phys. Rev. Lett. 103 160405
- [25] Moussa O, Ryan C A, Cory D G and Laflamme R 2010 Phys. Rev. Lett. 104 160501
- [26] Mealy G H 1955 Bell Syst. Tech. J. 34 1045-79
- [27] Roth Jr C H 2009 Fundamentals of Logic Design (Stanford, CT: Thomson)
- [28] Gühne O et al 2010 Phys. Rev. A 81 022121
- [29] Cabello A 2010 *Phys. Rev.* A **82** 032110
- [30] Harrigan N, Rudolph T and Aaronson S 2007 arXiv:0709.1149
- [31] Dakić B, Šuvakov M, Paterek T and Brukner Č 2008 Phys. Rev. Lett. 101 190402
- [32] Vértesi T and Pál K F 2009 Phys. Rev. A 79 042106
- [33] Brukner Č and Zeilinger A 2009 Found. Phys. 39 677-89
- [34] Galvão E F 2009 Phys. Rev. A 80 022106
- [35] Gallego R, Brunner N, Hadley C and Acín A 2010 Phys. Rev. Lett. 105 230501
- [36] Brukner Č, Taylor S, Cheung S and Vedral V 2004 arXiv:quant-ph/0402127
- [37] Brassard G et al 2006 Phys. Rev. Lett. 96 250401
- [38] Pawłowski M et al 2009 Nature 461 1101-4

20