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#### RESONANCE TWO PHOTON IONIZATION STUDY OF BINARY CLUSTERS OF

#### STYRENE WITH POLAR MOLECULES

A Dissertation submitted in partial fulfillment of the requirements for the degree of doctor of philosophy at Virginia Commonwealth University.

by

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#### Abstract

# RESONANCE TWO PHOTON IONIZATION STUDY OF BINARY CLUSTERS OF

#### STYRENE WITH POLAR MOLECULES

By Hatem Ahmed Mahmoud, Ph. D.

A Dissertation submitted in partial fulfillment of the requirements for the degree of Doctor of Philosophy at Virginia Commonwealth University.

Virginia Commonwealth University, 2004

Major Director: M. Samy El-Shall Professor, Department of Chemistry

One-color resonance two-photo ionization (R2PI) spectra of mixed clusters of styrene molecule (S) with polar molecules [water (W), methanol (M), ethanol (E), and trifuoroethanol (T)] were measured through the  $S_1 \leftarrow S_0$  transition of the styrene molecule. The spectra reveal a rapid increase in complexity with the number of polar molecules in the cluster, associated with van der Waal modes and isomeric forms. The spectral shifts of the cluster origins from the  $S_1$ - $S_0$  transition of the bare styrene molecule reflect the nature of the intermolecular interactions within the binary clusters. The obtained R2PI spectra

were compared with the spectra of the analogous benzene containing clusters. The styrenewater and the styrene-methanol complexes exhibited very different spectral shifts and structures as compared to the benzene-water and benzene-methanol complexes, respectively. The favorable interactions of polar molecules with the olefin group of styrene may explain the strong inhibition effects of exerted by small concentrations of water and alcohols on the cationic polymerization of styrene. Size-specified intracluster proton transfer reactions were observed for mixed clusters of styrene dimer with water, methanol and ethanol molecules. It was proposed that the polar molecules tend to aggregate around the olefin center, which promotes the transfer of the charge from the propagating chain to the hydrogen-bonded polar molecules subcluster. The minimum number of polar molecules required for a proton transfer to take place exothermically depends on the proton affinity of the polar molecules subcluster. The estimated upper limit value for the proton affinity of styrene dimer radical was evaluated based on the energetic of the proton transfer reaction to be  $\leq$  220.4 kcal/mol. No intracluster reaction was observed for styrenetrifluoroethanol (ST<sub>n</sub>) system. In order to provide a comparison between the styrene and benzene systems, the benzene-ethanol (BE<sub>n</sub>) and benzene-trifluoroethanol (BT<sub>n</sub>) clusters were studied by using the R2PI technique via the  $6_0^1$  transition of the benzene molecule. Both dissociative electron transfer and dissociative proton transfer reactions were observed within the  $BE_n$  clusters, where n = 2 and 3, respectively. Proton transfer reactions were observed following dissociative electron transfer reactions within the (BT<sub>n</sub>) clusters, where n = 4, to generate the protonated clusters ( $H^+T_n$ ).

#### **Chapter 1 Introduction**

A molecular level understanding of the solvent effects in chemical reactions requires detailed knowledge of solute-solvent interactions including the size dependence of the spectral shifts and the evolution of intermolecular mode frequencies.[1-4] Gas phase molecular clusters provide models for such interactions since the size effect and the degree of complexity of the interaction potentials can be systematically varied in well-defined systems.[5-10] The influence of stepwise solvation on both the ground and the excited electronic states of a solute molecule can be probed by studying molecular clusters composed of a single "solute" molecule and selected numbers of "solvent" molecules. The solvent molecules can be chosen in such a way to present a wide range of intermolecular interactions which may result in small perturbations of the electronic states of the solute molecule or in large changes in the structure and electronic properties such as those associated with charge-transfer and excimer interactions. Such studies have been advanced by the development of powerful techniques, which enables one to isolate ultracold clusters containing a "solute" chromophore and selected numbers of "solvent" molecules. Among these techniques, Resonant Two- Photon Ionization (R2PI) has been extensively used as a selective spectroscopic tool for the study of molecular clusters of various sizes and compositions.[6-26]

The study of molecular clusters of varying sizes and compositions constitutes an active and fascinating area of research. These clusters provide microscopic models for a molecular level understanding of a variety of important condensed phase phenomena such as solvation, phase transitions, nucleation, microphase separation, polymerization reactions, protein folding, and molecular recognition.[2, 3, 5-10, 27]

Binary clusters composed of a non-polar aromatic center and polar or hydrogenbonded solvent molecules represent an important class of molecular clusters. Several examples of these clusters have been investigated in recent years using combinations of spectroscopic and theoretical techniques.[6-10] The majority of the work has been focused on benzene in the presence of various polar and hydrogen bonding molecules.[6-10, 20-23, 25, 28-32] For example, extensive studies have been conducted on the benzene-water  $(BW_n)$  and benzene-methanol  $(BM_n)$  clusters using spectroscopic techniques such as Resonant Two- Photon Ionization (R2PI), hole burning, and resonant ion-dip infrared spectroscopy. [10, 20-23, 25, 28, 29, 33, 34] These studies have established several characteristic features of the intermolecular interactions between aromatic  $\pi$ -systems and hydrogen bonded solvent molecules. Among these features are (a) the blue spectral shifts in small clusters, which seem to increase with increasing hydrogen bonding ability of the solvent molecules, (b) the gradual switch to red shifts in larger clusters, (c) the tendency of the solvent molecules to form a hydrogen bonding network in a one-sided structural type with the benzene ring, and (d) the very efficient fragmentation of the photoionized clusters, which has been interpreted as a direct consequence of the hydrogen bonding to the benzene ring.[10] For a complete understanding of these systems, it is desirable to study other aromatic systems, including in which an additional  $\pi$ -interaction site is available.

Styrene is the simplest molecule that contains an unsaturated group covalently linked to an aromatic ring, and therefore represents a model system for studying the interactions between different  $\pi$ -systems with H-bonded molecules such as water and methanol. The unsaturated side chain provides an additional site for  $\pi$ -interaction with polar molecules that competes with the aromatic ring. These interactions are relevant to a wide variety of chemical, photochemical and biological processes in several fields such as photoconductivity, light induced charge and electron transfer reactions, and polymer chemistry.[1-3, 27, 35, 36] Styrene polymerization is among the most extensively studied polymerization process and continues to attract interest. The molecule is known to undergo polymerization in bulk monomer or in solution by free radical, cationic and anionic mechanisms.[37, 38] The thermally-initiated gas phase polymerization has been demonstrated and efforts to achieve a molecular level understanding of the various processes taking place during the polymerization continue to increase the potential of designing novel polymeric materials with unusual advanced properties.

In condensed phases, the crucial role played by trace amounts of water and other protic impurities on the mechanism and rate of polymerization in styrene is widely recognized. Under very dry conditions, radiation-induced polymerization of the bulk monomer proceeds rapidly [38, 39], but under "wet" conditions ( $[H_2O] > 10^{-3}$  M) polymerization proceeds more slowly.[39] Small amounts of methanol are also known to suppress polymerization and enhance dimerization in styrene.[40]

The particular systems of interest in the present study consist of a single styrene molecule (S) and several solvent molecules: water (W), methanol (M), ethanol (E), and trifuoroethanol (T). These systems have never been studied before. The study of styrenewater  $(SW_n)$ , styrene-methanol  $(SM_n)$ , styrene-ethanol  $(SE_n)$ , and styrene-trifluoroethanol  $(ST_n)$  clusters can provide information on the interactions involved in these systems and a greater a molecular level understanding of the retardation and inhibition effects of water and alcohols on the polymerization of styrene. The addition of the first few solvent molecules to the styrene molecule provides information on the preferred solvent sites and their relative strengths. The addition of further molecules yields information on solvent cooperative effects, coordination, and formation of solvent shells around the reactive monomer species. Like water, alcohol can participate in relatively weak hydrogen bonding to the  $\pi$ -system of styrene. However, the presences of the alkyl and fluoroalkyl groups provide favorable sites for hydrophobic interactions with the  $\pi$ -system of styrene. Detailed knowledge of these interactions is a prerequisite for a molecular level understanding of the retardation, inhibition and solvation effects of water and alcohols on the polymerization of styrene.

In order to provide a comparison between the styrene and benzene systems, we also study the benzene-ethanol,  $BE_n$ , and benzene-trifluoroethanol,  $BT_n$ , clusters. Again, these cluster systems have not been studied before.

The work described in this thesis is presented in five chapters. In chapter 2, a review of the spectroscopy of benzene and benzene-containing clusters is provided. In addition, the spectroscopy of the styrene molecule is presented and discussed. Chapter 3

deals with the experimental methods used in this research, which include cluster beam, time-of-flight mass spectrometry (TOF), resonant enhanced multiphoton ionization (REMPI) and non-resonance multiphoton ionization (MPI). Chapter 4 includes the results of the REMPI studies of the SW<sub>n</sub>, SM<sub>n</sub>, SE<sub>n</sub>, ST<sub>n</sub>, BE<sub>n</sub> and BT<sub>n</sub> clusters. These results include both the spectroscopy of the neutral clusters and the reactions observed following the ionization of the clusters. In Chapter 5, we present the overall conclusions from these studies.

#### **Chapter 2 Literature Survey**

#### 2.1 The Spectroscopy of Styrene and Styrene Containing Clusters

The 288 nm  $\widetilde{A} - \widetilde{X}$  absorption spectra of styrene and styrene- $\beta$ -D<sub>2</sub> were recorded by Hollas et al.[41] The spectra resembled that of  $\widetilde{A} - \widetilde{X}$  systems of aniline, phenol, chlorobenzene and ethynylbenezene and indicated that the  $\widetilde{A} - \widetilde{X}$  excitation involved primarily the benzene ring with some participation from the ethylene group. They reported that no vibrational evidence was found to support the proposal by Hui and Rice[42] that the  $CH_2$ -group is perpendicular to the plane of the rest of the molecule (see Figure 1). The C(1)- $C(\alpha)$  torsional vibration (v<sub>42</sub>) was found to be strongly mixed in the Å state with v<sub>41</sub>, out-of-plane vibration. The  $0_0^0$  band of styrene- $\beta$ -D<sub>2</sub> exhibited a small blue shift of +19 cm<sup>-1</sup> from that of styrene, which is incompatible with such configuration. The cross-sequence  $41_0^1 42_1^0$ ,  $41_0^2 42_2^0$ , ... in the X state was stronger than the sequence  $42_1^1$ ,  $42_2^2$ , ... The values for the separation of the first three levels of  $v_{42}$  in X state was deduced and were in disagreement with those obtained from Raman spectrum in the liquid phase.[43] The problem of torsional potential function for  $v_{42}$  remained unsolved until the single vibronic level (SVL) spectra of styrene following laser excitation into  $0_0^0$  band were recorded by Hollas and Ridley[44]. Also, they were recorded SVL spectra following laser excitation into more bands of the styrene  $\widetilde{A} - \widetilde{X}$  system.[45] Interpretation of those spectra led to many new assignments and reassignment of  $41_0^1 42_0^1$  which was previously assigned to  $40_0^2$ . The normal coordinates of  $v_{41}$  and  $v_{42}$  were found to be heavily mixed in the excited state A. Again no spectroscopic evidence was found to support the configuration of Hui and Rice.[42]

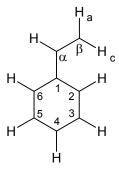


Figure 1: The atom numbering convention for styrene.

Zewail et al[46] were studied the excitation and dispersed fluorescence spectra for the  $\tilde{A}^1A' \leftarrow \tilde{X}^1A'$  transition of jet-cooled styrene. Although the obtained frequencies were in good agreement with Hollas and co-workers[41], their studies revealed new features of the mode structure and coupling for styrene in the absence of spectral congestion and sequence bands present in the previous bulb work. The excess energy and intramolecular vibrational redistribution were discussed in relevance to the dynamics of photoisomerization. The results of these studies were in support of planner configuration with significant rate of isomerization at low excess energies. The barrier to isomerization in the excited state  $\tilde{A}$  was estimated to be larger than that of t-stilbene ( $\approx 1200$  cm<sup>-1</sup>).

The binding energy of jet-cooled styrene-Ar cluster was determined by using resonant two-photon ionization.[47] The spectral shift of STY-Ar relative to the  $0_0^0$  band of  ${}^{1}B_{2} \leftarrow {}^{1}A_{1}$  of bare styrene was found to be -30 cm<sup>-1</sup>. The main feature in the R2PI spectra of STY-Ar was the disappearance of electronic vibrational transitions at excess vibrational energy  $E_v > 410 \text{ cm}^{-1}$  with respect to the  $0_0^0$  origin of styrene. This result was attributed to the vibrational predissociation on the  ${}^{1}B_{2}$  potential surface. The lower and upper limits for the dissociation energy in the excited state were determined from the R2PI spectra of STY-Ar. The dissociation energy  $D(^{1}B_{2})$  of STY-Ar was estimated to be 426±16 cm<sup>-1</sup>. Since the spectral shift  $\delta v$  corresponded to the extra stabilization of STY-Ar, the value of the dissociation energy in the ground state was thus;  $D({}^{1}A_{2}) = D({}^{1}B_{2}) - \delta v = 396 \pm 16 \text{ cm}^{-1}$ . This value was in good agreement with the calculated value of benzene-Ar complex [48] (D = 366± cm<sup>-1</sup>). The first adiabatic ionization energies of styrene (STY), phenylacetylene (PA), styrene-argon (STY-Ar) and phenylacetylene-argon (PA-Ar) were measured using (1+1') resonance enhanced multiphoton ionization (REMPI) threshold photoelectron spectroscopy (TES) technique as  $68.267\pm 5$ ,  $71.175\pm 5$ ,  $68.151\pm 5$ , and  $71.027\pm 5$  cm<sup>-1</sup>, respectively[49]. By this technique a specific molecule can be selected by tuning the excitation laser (frequency  $\omega_1$ ) to a resonant excited state and then ionizing from that state with a second laser (frequency  $\omega_2$ ). The threshold electrons were collected by applying a delayed pulsed field extraction as  $\omega_2$  was scanned. So, the photoelectron spectrum was obtained for the molecule in a chosen vibrational level of its excited state and the sum of the laser frequencies  $(\omega_1 + \omega_2)$  represents the energy axis of the spectrum. Ionization from a

particular  $S_1$  mode enhanced the intensity of the photoelectron spectrum and vibrational assignment of this structure could be achieved if the assignment of the vibronic level was already known. This was done for three modes in both styrene and phenylacetylene and the obtained ionic vibrational frequencies were compared with those known for  $S_0$  and  $S_1$ states.

RE2PI spectra were recorded for styrene-Ar<sub>n</sub> clusters (n = 2,3) and for 4fluorostyrene-Ar<sub>n</sub> clusters (n = 2-5).[50] The  $0_0^0$  electronic origins of STY-Ar, STY-Ar<sub>2</sub> and STY-Ar<sub>3</sub> were red shifted by 30, 63 and 33 cm<sup>-1</sup>, respectively, with respect to the  $0_0^0$ band of styrene. It was important to notice that the additive rule governs for STY-Ar<sub>n</sub> clusters. Therefore, the structure of STY-Ar<sub>2</sub> was believed to be (1+1) adduct. On the other hand, two possible structures could explain the red shifted origin for STY-Ar<sub>3</sub> by 33 cm<sup>-1</sup> which was nearly the same for STY-Ar; the first was a (3+0) adduct and the second was a (2+1) adduct with the argon atom that caused the blue shift attached to the other side of the styrene molecule. The trend in the R2PI spectra for 4-FSTY-Ar<sub>n</sub> clusters was very similar to that for STY-Ar<sub>2</sub>. These observations were attributed to different isomeric structures, with argon atom tied to one or both sides of the aromatic ring. The quantum calculations of the van der Waals vibrational frequencies for mono-Ar complexes matched with the observed spectrum for 4-FSTY-Ar. In a recent study[51], the one-color REMPI of styrene-Ar<sub>n</sub>, n = 4-12, were reported. The spectra of these clusters exhibited a single band, red shifted from the styrene origin transition by 24-37 cm<sup>-1</sup>.

#### 2.2 The spectroscopy of Benzene and Benzene Containing Clusters

The vibronic spectroscopy of benzene has received comprehensive studies. Moreover, the electronic structure of benzene is one of the most important polyatomic structures in the modern structural chemistry. The recent developed in the experimental techniques has resolved many of outstanding problems in benzene spectroscopy such as; the mechanism of vibronic coupling in  $\pi$  E<sub>2g</sub> singlet state, the absence of  $\sigma \rightarrow \pi^*$  or  $\pi \rightarrow \sigma^*$  transitions, the absence of *n*s Rydberg transitions and, the nature of rapid radiationless process above the origin of the first transition "channel three".

The vibrational pattern of an electronic excitation consists of two components; the vibronically induced component and the Frank-Condon component. Frank-Condon component forms progressions added to the vibronic inducing modes and possibly to the vibrationless origin. The analyses of vibrational pattern gave more information about the structure of the molecule in the excited state.[52]

#### 2.2.1 The Spectroscopy of Benzene

Spears and Rice[53] were studied the lifetime of 22 vibronic states of benzene and the relative quantum yield of SVL fluorescence for transitions from these states. The data showed the dependence of radiative and nonradiative lifetimes on excess vibrational energy and on vibrational symmetry type. The nonradiative lifetime  $\tau_{nr}$  decreased as the number of vibrational quanta increases in the prepared stats. The rate of change of  $\tau_{nr}$  per quantum of vibrational energy was almost the same for all vibrational modes. They found that the nonradiative rate constants increased linearly with increasing vibrational energy for the  $6^{n}1^{m}$  and  $1^{n}$  levels. They found a pressure dependence of the observed lifetime to a lower pressure limit of 0.006 torr, which was due to hard-core collision of benzene.

Schuyler et al[54, 55] analyzed the SVL fluorescence spectra of benzene vapor in the <sup>1</sup>B<sub>2u</sub> state. The fluorescence spectra were found to be in good agreement with Herzberg-Teller description of the <sup>1</sup>B<sub>2u</sub>-<sup>1</sup>A<sub>1g</sub> vibronically induced transition. A single first-order Herzberg-Teller selection rule described the dominant feature in every obtained SVL spectrum. Higher order Herzberg-Teller selection rules provided a consistent account for minor fluorescence structure. Qualitative agreement was found between observed and Herzberg-Teller predicted fluorescence band intensities. The relative intensities in  $v_1$ progressions of the types  $1_n^o$ ,  $1_n^1$  and  $1_n^2$  were qualitatively different, and each matches the approximate Frank-Condon calculations. Four new vibrational frequencies in the <sup>1</sup>B<sub>2u</sub> state were located:  $v_{4'} = 365 \text{ cm}^{-1}$ ,  $v_{5'} = 749 \text{ cm}^{-1}$ ,  $v_{9'} = 1148 \text{ cm}^{-1}$  and  $v_{17'} = 712 \text{ cm}^{-1}$ . Six bands were found among the prominent structure in  ${}^{1}B_{2u}$ - ${}^{1}A_{1g}$  absorption spectrum.[55] The absorption spectrum and SVL fluorescence spectra showed that a set of 12 vibrations was responsible for most the structure in  ${}^{1}B_{2u}$ - ${}^{1}A_{1g}$  radiative transitions. The "active" vibrations encompassed six of the ten vibrational symmetry species. Each vibration in these six species was "active".

The MPI spectrum of benzene was first reported by Johnson.[56, 57] A low-lying state of benzene was observed at 391 nm, which had not been previously seen, in the three-photon ionization spectrum. This new state was an allowed two-photon state of probable symmetry  ${}^{1}E_{1g}$  or  ${}^{1}E_{2g}$ . The three-photon Rydberg structure was found to be similar to the one-photon absorption spectrum. This similarity indicated a strong  ${}^{1}A_{1g}$  component in the two-photon virtual state in agreement with the two-photon vibronic structure of the  ${}^{1}B_{2u}$  state, which was also strongly  ${}^{1}A_{1g}$ .

Schlag et al[58] reported the two-photon excitation spectra of benzene and benzene- $d_6$  in the gas phase with a high power tunable laser. The two-photon inducing modes  $v_{14}$ ,  $v_{15}$  ( $b_{2u}$  symmetry) and  $v_{18}$  ( $e_{1u}$  symmetry) were determined experimentally by new technique based on the zero point level of S<sub>1</sub> state of benzene. The hot band spectrum was assigned from these transitions. The two strongest inducing modes in the excited state were found to be the  $v_{14}$  mode ( $b_{2u}$  symmetry) and the  $v_{18}$  mode ( $e_{1u}$  symmetry). The frequency of the  $v_{14}$  showed a strong blue shift by 257 cm<sup>-1</sup> in the excited state. The  $v_{19}$  was found to be extremely weak or absent in disagreement with the theoretical predictions.

The absorption spectra of benzene- $h_6$  and benzene- $d_6$  were measured photoelectrically.[59] The data were in consistent with Frank-Condon factors predicted on the basis of Herzberg-Teller theory for a C-C change in the range 0.032 to 0.035 Å. The relative intensities for a large number of progressions showed small deviations from prediction and not all of which could be assigned to Fermi resonance.

The resonance enhanced multiphoton ionization and fragmentation of benzene beam in mass spectrometer were reported by Bernstein et al.[60, 61] The fragmentation pattern of benzene ion, corresponding to m-photon ionization and two-photon excitation, were measured near 391 nm. Increasing the laser pulse peak energy increased the average number of photon absorbed per molecule and thus the extent of fragmentation. However, the ionization and the fragmentation were found to be wavelength dependence even at the highest laser peak power intensities. Under strong focusing conditions,  $C^+$  ion was the most abundant ion indicating the absorption of at least 9 photons per molecule during the 6 ns laser pulse.

Atkinson and Parmenter studied the 260 nm absorption of benzene. They were remeasured the positions and intensities of most band maxima in the  ${}^{1}B_{2u}$ - ${}^{1}A_{1g}$  transition of benzene vapor[62] at 300 K. They examined the role of vibrational angular momentum, vibrational degeneracies and the selection rules in the 260 nm absorption spectrum of benzene[63]. They developed a new parameter,  $\chi$ , which can be used to obtain the rotational band contour of any vibronic component transition. Maxima were sorted into v<sub>1</sub> progressions and then most of these progressions assigned to first-, second-, and third order Herzberg-Teller transitions[64]. Third-order transitions were found to be few in number with intensities that comparable to those of second-order transitions. New nine fundamentals in the  ${}^{1}B_{2u}$  state and two in the  ${}^{1}A_{1g}$  state were specified.

The FL spectrum of free jet benzene at a rotational temperature of 0.3 K was studied.[65] The band contour of rotationally cooled benzene displayed well-resolved single rovibronic transitions in contrast with the band contour obtained for a room temperature sample.[66] The derived rotational and Coriolis constants were in agreement

with those derived from room temperature band contour analyses. It was found that more than 95% of all benzene molecules were in the lowest possible rovibronic level.

Two time delayed laser pulses of different frequencies were used to study the mechanism of multiphoton ionization and fragmentation of jet cooled benzene.[67] Data showed that the mechanism of multiphoton ionization of benzene includes two different energy states. First, the neutral molecules absorbed photons until the ionization potential was reached. Then ions were formed with little excess energy. Photoexcitation of these cold ions to states with higher excess energy produced fragmentation.

Johnson et al[68] studied R2PI of benzene-h<sub>6</sub> and benzene-d<sub>6</sub> The effect of channel three[69], which was the increase in the radiationless rates in the S<sub>1</sub> state of benzene when the vibrational excess energy exceeded about 3000 cm<sup>-1</sup>, was followed 3000 cm<sup>-1</sup> past the fluorescence region. Under rotationally cooled conditions, a new progression in C<sub>6</sub>H<sub>6</sub> spectrum was revealed next to the  $14_0^11_0^n$  progression (n = 1-5). Also, new and wellresolved structure was obtained for C<sub>6</sub>D<sub>6</sub> as the  $14_0^11_0^n$  progression was followed.

Benzene, aniline, N,N-dimethyl and 2,4-dimethyl aniline were studied by resonance enhanced laser ionization mass spectrometry under collision free conditions.[70] At lower laser intensities, these molecules showed efficient R2PI producing entirely the parent ion. At higher laser intensities, extensive fragmentation was produced by R3PI or R4PI. For benzene, the fragmentation pattern showed  $C_3H_3^+$ ,  $C_4H_2^+$ ,  $C_5H_3^+$  and  $C_6H_6^+$ peaks. All of these ionic fragments were observed in conventional electron impact spectra. The cross section for ionization of  ${}^{1}B_{2}$  excited state in benzene and aniline were found to be equal and independent of the wavelength in the range 250-300 nm.

The  ${}^{1}B_{2u}$  transition of benzene at room temperature was very congested. The rotational envelopes were broad and it was hard to resolve the rotational band because of Doppler broadening. Using Doppler free two-photon FL spectroscopy Schlag and co-workers[71] resolved a part of one rotational band. In this study, it was difficult to determine the band origin. By using supersonic molecular beam, most of the rotational envelope was eliminated providing accurate assignment of band origins. Johnson et al[72] obtained the  ${}^{1}B_{2u}$ - ${}^{1}A_{1g}$  two-photon spectra of several isotopes of benzene using supersonic beam-multiphoton ionization spectroscopy. New assignments and reassignments were introduced especially in transitions concerning the v<sub>12</sub> vibration. Also, they measured the isotope shifts due to naturally occurring C<sup>13</sup> in benzene and perdeuterated benzene. New feature due to benzene dimmers was revealed in all the deuterium isotopes.

An extensive investigation of the single vibronic level fluorescence spectroscopy of jet-cooled benzene was carried out by Rice co-workers.[73] The intramolecular energy transfer in the  ${}^{1}B_{2u}$  of benzene was investigated. The most important result was the assignment of the  $8_{0}^{1}$  transition. The assignment for v<sub>8</sub> in the  ${}^{1}B_{2u}$  state led to new assignments involving the activity in v<sub>18</sub> and v<sub>4</sub>; one-photon transitions involving the e<sub>1u</sub> mode v<sub>18</sub> were observed and, two new progressions involving activity of the b<sub>2g</sub> mode v<sub>4</sub> were recognized. Several spectroscopic controversies of the  ${}^{1}B_{2u}$ - ${}^{1}A_{1g}$  spectra of benzene were resolved. The spectral simplification accomplished by using supersonic expansion

allowed measurements of the lifetimes of several vibrational levels in  ${}^{1}B_{2u}$  benzene.[74] Non-exponential fluorescence decays were observed from vibrational levels near the "channel three" region.

Smalley et al[16] obtained the  $S_1 \leftarrow S_0$  spectrum of (benzene)<sub>n</sub>, where n=2-4, in a supersonic beam using 2C-R2PI technique. Weakly induced  $0_0^0$  transitions red shifted from the forbidden  $0_0^0$  of benzene were observed for all clusters. The red shifts were -40, -115, and -149 cm<sup>-1</sup> from the monomer origin for the dimer, trimer, and the tetramer, respectively. The splitting of the  $0_0^0$  band was found to be less than 6 cm<sup>-1</sup> for all the clusters. The trimer and teramer showed strong vdW progressions. The fluorescence measurement data showed that the dimer rearrange from the T-shaped structure into sandwich excimer with net D<sub>6h</sub> symmetry within several picosecond after laser excitation.

Levy et al[75] reported the  ${}^{1}B_{2u}$ - ${}^{1}A_{1g}$  ultraviolet spectra of benzene clusters. Sharp features were observed in the fluorescence excitation spectrum for dimers, trimers and tetramers. The dispersed fluorescence spectra of the trimer showed identical structures from excitation of  $6^{1}$  and  $0^{0}$  levels.

#### 2.2.2 Spectroscopy of Benzene Containing Clusters:

The spectral shift of the cluster's origin with respect to the rare benzene gives a direct measure of the relative binding energies of the cluster in the ground and excited states,  $S_0$  and  $S_1$  respectively. This shift is known as solvent shift. The electronic frequency shifts for most benzene-(solvent)<sub>n</sub> clusters are on the order of a few hundred wavenumber. The solvent can induce transitions that are forbidden in the bare benzene molecule. For

example, the  $0_0^0$  transition is electronic dipole forbidden in the absence of external perturbations. If the solvent bounds to benzene in such a way to reduce the sixfold symmetry of benzene to lower than threefold symmetric, the  $0_0^0$  transition becomes allowed. Therefore, studying of the vibrations that gain intensity upon clustering can be used as a probe for the symmetry and structure of benzene-X complexes.

#### 2.2.2.1 Benzene-Rare Gas Clusters:

The benzene-rare gas clusters have been the subject of numerous experiments and several theoretical studies. Shown in Table 1 are the spectral shifts of benzene-rare gas clusters relative to the  $6_0^1$  transition of bare benzene molecule. Illustrated in Figure 2 is the dependence of the spectral shift of benzene-rare gas clusters on the polarizability of the rare gas. One-sided benzene-Ar<sub>n</sub> clusters (n=1-7) showed ideal size dependence on the spectral shift as shown in Figure 3.

Bz-R	Shift, $cm^{-1}$ , $6_0^1$	Ionization potential, ev	References
He	2.30	21.56454	[65]
Ar	-21.0	15.759	[76]
Kr	-33.397	13.99961	[77]
Xe	-53.666	12.12987	[77]

Table 1: Collective data for benzene-rare gas binary clusters.

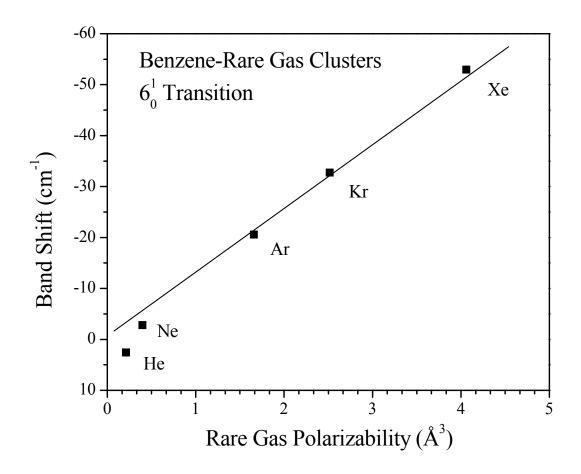


Figure 2: Spectral shift of the band as a function of the rare gas polarizability of the measured benzene-rare gas complexes. (Adapted from reference [77])

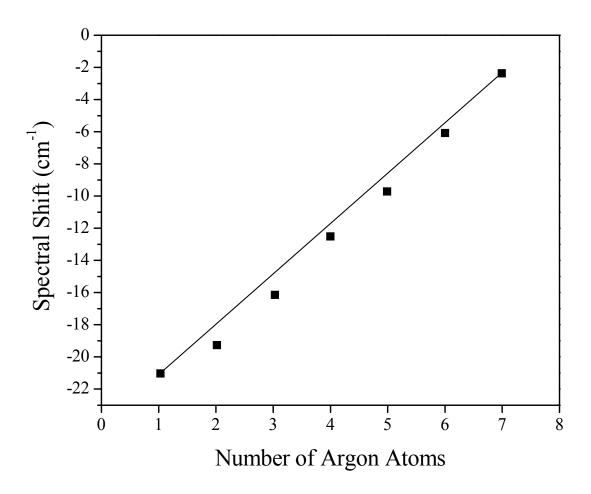


Figure 3: Size dependence of the spectral shift of the one-sided Bz-Ar<sub>n</sub> clusters (where n=1-7). (Adapted from reference [76])

#### 2.2.2.2 Benzene-Nonpolar Molecule and Benzene-Polar Molecules Clusters

The collective data for benzene-nonpolar molecule and benzene-polar molecule binary clusters are given in Table 2 and Table 3, respectively.

Bz-X	Shift,	Binding	Dipole	Ionization	Polarizability	References
	$cm^{-1}$	energy,	moment[78],	potential,	$10^{-25} \text{ cm}^3$	
	$6^{1}_{0}$	cm <sup>-1</sup>	Debye	ev		
$C_2H_2$	-127		0	11.40		[79]
$C_3H_8$	-72.0	1044	0.084	10.94		[80]
$CCl_4$	-68	839-	0	8.11	105	[21],
		1119				[81],[82]
$C_2H_6$	-57(I)	778	0	11.52		
	-					[80]
	51(II)					
Cyclohexane	-47		0	9.88		[83]
$C_6H_6$	-43/-		0	9.24	103	[31]
	39					[31]
$CD_4$	-41.2	540	0	10.53		[26]
$CH_4$	-41	599	0	12.61	26.0	[84]
$SF_6$	-11.5		0	15.32		[85], [86]
$N_2$	-6.8	-501	0	15.581	17.6	[87], [88]
$CO_2$	4	-868	0	13.78	26.5	[87], [88]
CF <sub>4</sub>	6.1		0	10.14		[26]

Table 2: Collective data for benzene-nonpolar molecule binary clusters.

Table 3: Collective data for benzene-polar molecule binary clusters.

Bz-X	Shift, cm <sup>-1</sup>	Binding	Dipole	Ionization	Polarizability 10 <sup>-25</sup> cm <sup>3</sup>	References
		energy,	moment,	potential,	10 cm	
	$6^{1}_{0}$		Debye	ev		
NH <sub>3</sub>	-25	-677	1.47	10.07	23.0	[89], [15]
CO	-16	-612	0.112	14.01	19.5	[87], [88]
CH <sub>3</sub> Cl	8	524-	1.87	11.37	45.3	[21],[81],[82]
		1189				[21],[01],[02]
CH <sub>3</sub> CN	37.6	559.4	3.9	12.20		[90], [31]
CH <sub>3</sub> OH	44		1.70	10.84	32.0	[25], [91]
$H_2O$	52	559-	1.85	12.62	14.8	[21],
		1468				[81],[82]
HC1	125.1	629-	1.08	12.74	26.0	[21],
		1329				[81],[82]
DCl	132.5		1.084	12.76		[20]
CH <sub>3</sub> COOH	152		1.74	10.65		[30]
HCCl <sub>3</sub>	179	699-	1.01	11.37	85.0	[21],
		1608				[81],[82]

#### **Chapter 3 Experimental**

#### 3.1 Cluster Beam Apparatus

The generation of clusters involved the expansion of a gas from a high-pressure source (stagnant pressure) through a narrow hole into an expansion chamber of lower pressure. The thermal energy in the high-pressure source is converted into a directed velocity u as the gas expands. If Cp is the heat capacity of the expanding gas at source temperature  $T_s$ , the energy balance of an ideal gas of molecular weight m is given by:

$$C_{\rm P}T_{\rm S} = C_{\rm P}T_{\rm B} + {\rm mu}^2/2$$
 Eq. 3-1

Where  $T_B$  is the beam temperature. Since mu<sup>2</sup>/2 is positive,  $T_B$  is lower than the source temperature  $T_S$ . So the adiabatic gas expansion provides cooling of gas molecules, which promotes clustering. The expansion of the gas into a collision free region (a few nozzle diameters) creates a directed motion of the gas. This directed motion provides cooling in the vibrational, rotational and translational temperatures. The quenched temperatures are in the order  $T_{vib}>T_{rot}>T_{trans}$ . The liberation of latent heat of condensation during the clustering increases the internal temperature of the clusters. However, collisions with an inert gas provide cooling of the internal temperature of the clusters.[92]

The cluster beam apparatus consists of two cubic chambers connected by a tube (Figure 5). The chambers are isolated through a conical skimmer orifice (1 mm or 3 mm). The first chamber (expansion or jet chamber) is used mainly for the generation of clusters

and the second chamber (beam chamber) houses the TOF mass spectrometer. The expansion chamber is pumped by a diffusion pump (Varian VHS-6, pumping speed = 3000 l/s in He), achieving a background pressure of  $10^{-7}$ - $10^{-8}$  Torr, and a working pressure of  $\approx 5 \times 10^{-5}$  Torr when the nozzle is operated with a backing pressure of  $10^2$ - $5 \times 10^3$  Torr. The beam chamber is pumping by a diffusion pump (Varian VHS-6) equipped with nitrogen trap to condense the background gas, hence lowering the ultimate vacuum that can achieved (low  $10^{-8}$  Torr). The background pressure of the beam chamber rises to  $\approx 10^{-7}$  Torr when the nozzle is operated.

### **3.2 TOF Mass Spectrometry**

The common configuration in TOFMS is the two-stage Wiley-McLaren configuration.[93] The principle of mass separation in this design relies on the ions being created at the same time and location within an electrostatic field. Then the ions are accelerated in an electrostatic field. The two-stage accelerating field consists of three parallel plates (Figure 5). The electrostatic field accelerates all ions into a field-free drift region with a kinetic energy of zV, where z is the ion charge and V is the applied voltage. Since the ion kinetic energy is  $\frac{1}{2}$ mv<sup>2</sup>, lighter ions reach the detector at the end of drift region sooner than the heavier ions.

Focused laser beams are ideal for TOF ionization sources due to their short pulse width and small beam waist. The placement and overlap of the cluster and laser beams within the repeller and accelerator fields define the ionization region. However, since both laser and cluster beams have finite temporal and spatial widths, broadening occurs in the measured ion flight time distribution. The laser plus width is less than 10 ns, and therefore the spatial effects are more significant than the temporal effects. The spatial distribution of ions results in a spread of ion kinetic energies since the ions are created with different potential energies (determined by the ion distance from the field). The space focus design reduces the broadening by creating ions within a differential positive field, and hence, focusing the ions. The focusing region is defined by the differential voltage between the repeller and accelerator fields ( $\Delta V \approx 500 \text{ V}$ ). The focused ions enter the second stage and are accelerated ( $\Delta V \approx 4000 \text{ V}$ ) and proceed to travel a field free region ( $\approx 110 \text{ cm}$ ) until detected by two-stage microchannel-plate detector. Two deflection plates (0-350 V) are placed in the free field region to correct the trajectories of ions.

The current output of the detector is amplified and then digitized by 500 MHz (LeCroy 9350A). In principle, a complete TOF mass spectrum can be obtained from a single ionization event, but in practice an average mass spectrum is from 50-1000 ionization events.

Pulsed adiabatic expansion of a sample seeded into an inert carrier gas was employed to generate the neutral clusters. A vapor of mixture sample was expanded through a conical nozzle (200 or 500  $\mu$ m diameter) in pulses of 300-550  $\mu$ s duration at a repetition rate of 8-20 Hz into an evacuated chamber at  $\approx 10^{-5}$  Torr. The jet was skimmed and passed into a high vacuum chamber, which is maintained at  $2x10^{-7}$ - $8x10^{-8}$  Torr. The collimated cluster beam passed into the ionization region of the TOF mass spectrometer where it intersects the laser beam. A XeCl excimer pumped dye laser combination (Lambda Physik LPX-101 and FL3002) was used to provide tunable light for photoionization of clusters. Coumarin 540A (Exciton;  $\lambda_{range}$ = 513-612 nm) and coumarin 307 (Lambda Physik;  $\lambda_{range}$ = 485-546 nm) were used with  $\beta$ -BaB<sub>2</sub>O<sub>4</sub> crystal (CSK Co.) to generate a tunable frequency-doubled output of 10<sup>-8</sup> s pulses for styrene and benzene clusters systems, respectively. The spatially ultraviolet radiation (using a set of four quartz Pellin-Broca prisms) of average pulse energies 0.1-0.3 mj per pulse was shaped with a 3 cm lens to provide 1-2 mm<sup>2</sup> beam in the ionization region.

# 3.3 Resonant Enhanced Multiphoton Ionization REMPI

In the REMPI process, the neutral molecule or cluster in the ground electronic state is excited to an intermediate state by absorbing a resonant *m*-photon. The presence of a resonant intermediate state enhances the probability of ionization. Absorbing of one or more *n*-photon ionizes the molecule. This resonant ionization process is labeled (m+n)R(m+n)PI or 1C-R(m+n)PI for one-color photoionization scheme and (m+n')R(m+n')PI or 2C-R(m+n')PI for two-color photoionization scheme (Figure 4). The combination of REMPI, supersonic beam and mass-selective TOF spectrometry provides a powerful spectroscopic technique.

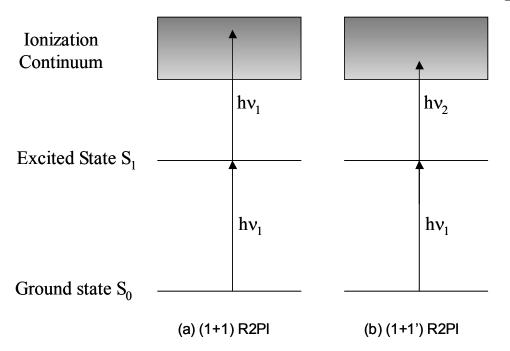


Figure 4: (a) One-color photoinoization scheme or (1+1)R2PI or 1C-R2PI. (b) two-color photoionization sheeme or (1+1')R2PI or 2C-R2PI.

## 3.4 Multiphoton Ionization MPI

The process at which two or more photons are absorbed in concerted interaction with individual atoms or molecules is called "multiphoton". The cluster beam was prepared as previously described above. The ionization was accomplished by absorbing photons of monochromatic light. The MPI was achieved by one of the following laser: ArF (193 nm) provided by excimer laser (Lambda Physik COMPex 102).

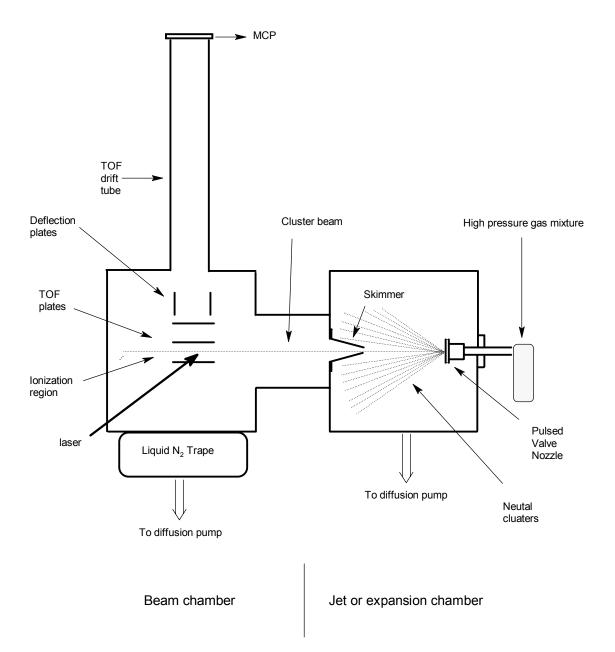


Figure 5: Schematic representation of the cluster beam apparatus combined with the TOF mass spectrometer.

#### **Chapter 4 Results and Discussion**

### 4.1 Styrene Spectroscopy of Clusters

Displayed in Figure 6 are the R2PI spectra obtained by monitoring the mass channel corresponding to styrene (S) in the  $0_0^0$  region between 34,730 and 35,820 cm<sup>-1</sup>. The spectrum in the bottom panel shows the  $0_0^0$  origin, the overtone band  $41_0^2$  and the vibronic band  $29_0^1$  as well as the cross sequence bands  $41_0^142_0^1$  between the  $v_{42}$  and  $v_{41}$  modes (ethylene group torsional rotation and out-of-plane bend, respectively) and  $40_0^141_0^1$ . The middle panel shows strong activities of three vibronic bands  $(28_0^1, 27_0^1 \text{ and } 26_0^1)$ , two overtone bands  $(42_0^2 \text{ and } 40_0^2)$  and three cross-sequence bands  $(29_0^141_0^142_0^1, 41_0^242_0^2 \text{ and } 28_0^129_0^1)$ . The spectrum region in the top panel shows seven fundamental vibrations ( $v_{19}$ - $v_{25}$ ). The vibration modes  $v_{19}$ ,  $v_{20}$ ,  $v_{21}$  and  $v_{23}$  were attributed to in plane C-H bending vibrations, while the vibration mode  $v_{19}$  corresponds to  $-CH_2$  rocking vibration. The spectrum is in good agreement ( $\pm 0.7 \text{ cm}^{-1}$ ) with previous fluorescence excitation study by Zewail et al.[46] The relative frequencies and intensities obtained from this work together with data obtained by Zewail et al [46] are tabulated in Table 4.

The styrene  $0_0^0$  band obtained by monitoring the mass channels corresponding to  $C_8H_8$ ,  ${}^{13}CC_7H_8$  and  ${}^{13}C_2C_6H_8$  [zero, one and two  ${}^{13}C$  isotopes in  $C_8H_8$ , m/z = 104, 105 and

106; respectively] (Figure 7). The replacement of one and two <sup>12</sup>C with <sup>13</sup>C atoms results in small blue shifts of the  $0_0^0$  band by 2.5 and 5 cm<sup>-1</sup> in the spectra, respectively. Similarly, the shift of the  $0_0^0$  can also be induced by the clustering of atoms or molecules around the styrene molecule.

Shown in Figure 8 is the R2PI spectra of (styrene)<sub>n</sub> clusters, where n = 1-4, in the  $S_1 \leftarrow S_0 \ 0_0^0$  band region of styrene monomer. The R2PI excitation spectrum of styrene dimer,  $S_2$ , exhibits a broad red-shifted band with several spectral features on the base line. These features are reproducible under various expansion conditions and tabulated in Table 5. The red-most peak at  $-252.4 \text{ cm}^{-1}$  may be the origin of  $S_2$ . The fragmentations from higher styrene clusters greatly contribute to the broadening in the R2PI spectrum of styrene dimer. The spectra of trimer,  $S_3$ , and tetramer,  $S_4$ , show unresolvable and broad bands with maximum red shifted by  $-107 \text{ cm}^{-1}$  ahead from the origin of styrene monomer. The obtained mass spectrum at this maximum is displayed in Figure 9. The red spectral shifts by the styrene clusters are due the intramolecular dispersion forces within these clusters.

A typical mass spectrum obtained following the multiphoton ionization of styrene clusters produced by 193 nm is displayed in Figure 10. The MPI mass spectrum of styrene clusters illustrates an enhanced intensity for  $S_2$  with a distinguished even/odd alternation in the ion intensities, up to n = 10. Similar findings have been reported by our group using different laser wavelength (248 nm).[94] The styrene dimer cation ( $C_{16}H_{16}^{+.}$ ) is believed to possess covalently bonded structure.[37, 38, 94-96]

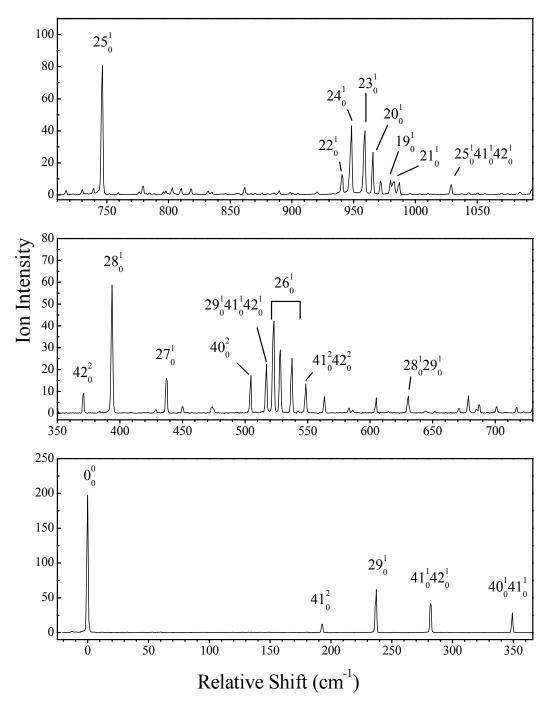


Figure 6: Three successive panels of the one-color R2PI spectrum of the jet-cooled styrene molecule in the supersonic expansion at He carrier gas (P = 100 psi). The relative shift is scaled with respect to the  $0_0^0$  transition at 34,758.79 cm<sup>-1</sup>.

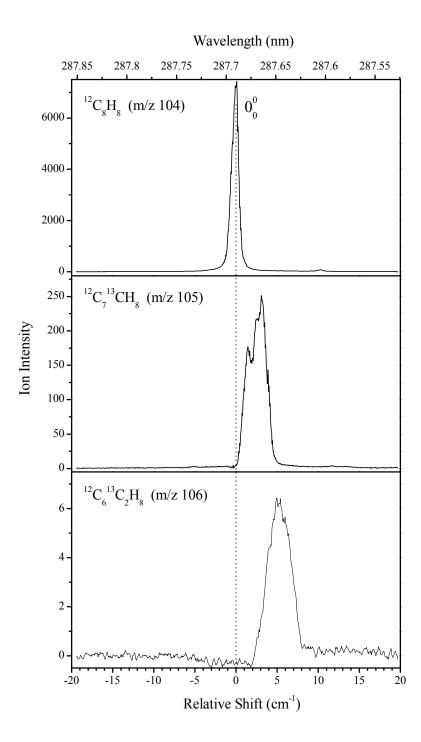


Figure 7: The effect of  ${}^{13}C$  isotope substitution on the  $0_0^0$  band of styrene molecule.

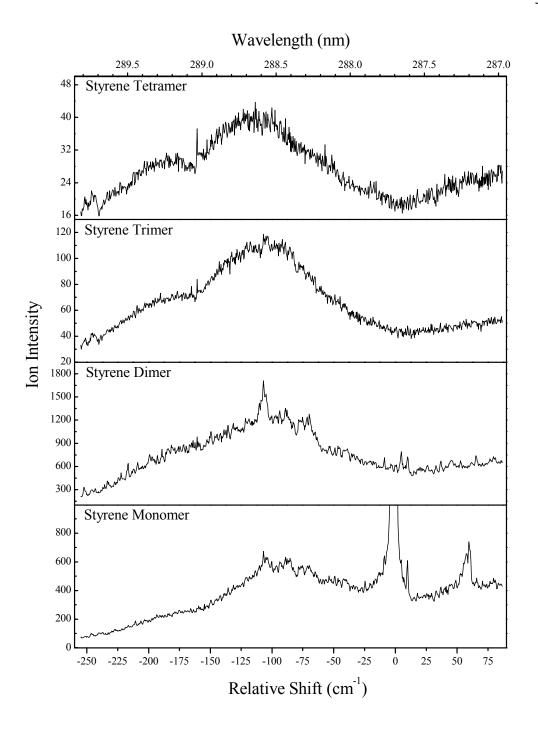


Figure 8: One-color R2PI measured in the (styrene)<sub>n</sub> mass channels, with (n = 1-4), [S, S<sub>2</sub>, S<sub>3</sub> and S<sub>4</sub>] relative to the electronic origin band of the S<sub>1</sub>  $\leftarrow$  S<sub>0</sub> transition of the styrene molecule at 34,758.79 cm<sup>-1</sup>.

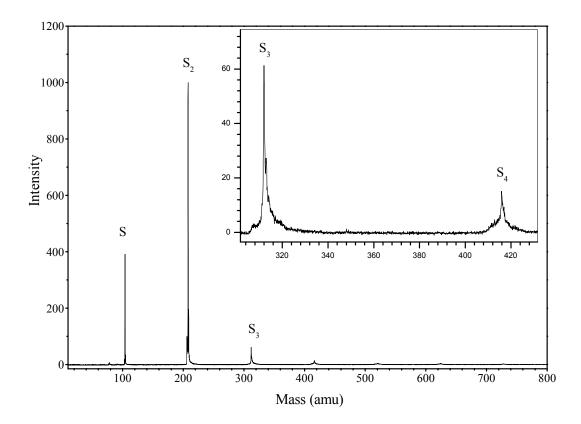


Figure 9: R2PI mass spectrum of the styrene cluster beam obtained at 288.577 nm.

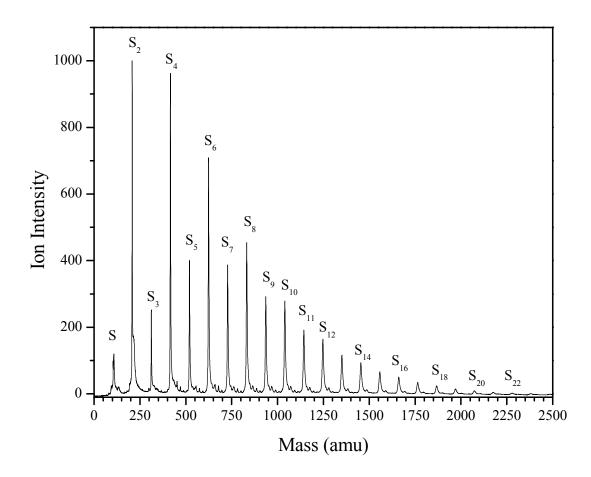


Figure 10: Multiphoton ionization mass spectrum of styrene cluster beam (S<sub>n</sub>), n up to 22, obtained by using 193 nm excitation wavelength.

Relative Shift (this work)	Relative Intensity (this work)	Relative Shift (Ref [46])	Relative Intensity (Ref [46])	Assignment (Ref [46])
0.0	100.0	0.0	100	$0_{0}^{0}$
60.0	0.6	59.5	0.6	$41_0^142_1^0$
192.8	6.3	193.2	0.9	$41_0^2$
237.3	31.4	237.2	8.9	$29^{1}_{0}$
282.0	21.2	282.4	9.3	$41_0^142_0^1$
349.3	14.4	349.1	4.0	$40^1_041^1_0$
371.0	4.7	371.0	0.8	$42_{0}^{2}$
393.7	29.8	394.5	9.6	$28^{1}_{0}$
437.3	8.1	437.1	2.0	$27^{1}_{0}$
504.8	8.8	505.5	0.6	$40^{2}_{0}$
517.2	11.4	517.5	0.6	$29^1_0 41^1_0 42^1_0$
523.2	21.4	523.2	4.6	
528.1	14.7	528.2	2.9	$26^{1}_{0}$
537.6	12.8	537.3	2.0	
563.5	3.9	563.6	0.2	$41_0^2 42_0^2$
630.5	4.0	630.0	0.3	$28^1_0 29^1_0$
678.6	4.0	678.7	0.3	$28_0^1 41_0^1 42_0^1, 27_0^1 29_0^1 (?)$
746.5	40.9	745.8	23.6	$25^{1}_{0}$
861.8	2.4	862.2	0.2	$27_0^2, \\ 28_0^1 29_0^2(?)$
940.6	6.5	940.4	3.4	$22^{1}_{0}$
948.3	21.9	947.8	22.6	$24^{1}_{0}$
959.0	20.4	959.1	22.8	$23_{0}^{1}$
965.6	13.5	965.1	8.3	$20_{0}^{1}$
979.9	4.7	980.7	3.1	$19^{1}_{0}$
982.0	4.2	983.0	3.0	$21_0^1(?)$

Table 4: Assignment of the R2PI spectrum of the jet-cooled styrene molecule.

987.1	4.0	987.3	1.9	$25_0^1 29_0^1(?)$
1029.0	3.3	1029.2	1.7	$25^{1}_{0}41^{1}_{0}42^{1}_{0}$

Table 5: Spectral features observed in  $S_2$  mass channel (S = styrene)

Relative Shift	<b>Relative Intensity</b>	<b>Relative Shift</b>	<b>Relative Intensity</b>
65.1	738	-141.8	1031
52.7	674	-144.3	1018
37.0	670	-149.4	1050
9.8	725	160.7	983
4.7	792	-164.0	935
-8.8	716	-169.5	863
-35.5	766	-171.6	882
-69.8	1280	-174.5	858
-71.5	1241	-181.7	878
-75.2	1240	-189.2	825
-82.8	1178	-193.0	743
-89.2	1354	-197.2	727
-93.8	1322	-199.3	755
-106.9	1709	-208.9	651
-110.3	1421	-216.5	639
-118.7	1217	-222.3	519
-121.6	1224	-224.4	456
-123.7	1104	-233.6	410
-131.3	1158	-252.4	330
-135.5	1136		

### 4.2 R2PI of Styrene-Water Clusters

### 4.2.1 R2PI spectra of styrene (water)<sub>n</sub> clusters, n=1-2

The R2PI spectra obtained by monitoring the mass channels corresponding to S, SW and SW<sub>2</sub> in the  $0_0^0$  region of the styrene monomer between 34,700 cm<sup>-1</sup> and 34,900 cm<sup>-1</sup> are presented in Figure 11. The electronic origin band of the S<sub>1</sub>  $\leftarrow$  S<sub>0</sub> transition of styrene is located at 34,758.79 cm<sup>-1</sup>.[46, 97] The excitation spectrum obtained in the S mass channel shows the  $0_0^0$  origin and the vibronic band  $42_0^2$  as well as the cross–sequence band  $41_0^142_0^1$ between the v<sub>42</sub> and v<sub>41</sub> modes (ethylene group torsional rotation and out-of-plane bend, respectively). The other small features shown in the spectrum of Figure 11 are due to hot bands ( $29_1^1$  and  $42_1^1$ ). The spectrum is in very good agreement with previous fluorescence excitation and REMPI studies.[46, 98]

The SW mass channel has several strong features identical to the features that appear in the SW<sub>2</sub> channel. These features were assigned to the SW<sub>2</sub> cluster with a sharp origin peak blue shifted relative to the  $0_0^0$  of the styrene by 45.7 cm<sup>-1</sup>. In addition, a week peak at 21.5 cm<sup>-1</sup> was assigned to the origin of the SW complex. Two other small peaks at 25 and 35 were assigned to low frequency van der Waals (vdW) intermolecular vibrations of the SW complex. The SW<sup>+</sup> cluster ion fragments efficiently into the styrene channel as indicated by the significant intensity of the 21.5 cm<sup>-1</sup> peak found in the S channel shown in Figure 11.

In contrast to the simplicity of the spectral features assigned to the SW complex, a very rich spectrum is observed in the SW<sub>2</sub> mass channel with all the features appearing in the SW channel as a result of efficient loss of a water molecule from the SW<sub>2</sub> cluster upon photoionization. The origin of the SW<sub>2</sub> cluster possesses two long progressions of bands assigned to the excitation of the intermolecular vibrational modes [the first progression appears at 73, 100, 125 and 148 cm<sup>-1</sup> and the second progression corresponds to peaks at 80, 107, 133 and 156 cm<sup>-1</sup>. The very rich intermolecular Franck-Condon activity of the SW<sub>2</sub> cluster indicates a large change in the geometry of the cluster following the electronic excitation of the cluster.

The tentative assignments of the SW and SW<sub>2</sub> origins are complicated by the efficient fragmentation of the ionized clusters as a direct result of the structural change between the neutral and the ionized cluster. As the cluster ion relaxed to a more stable structure, the relaxation energy released promotes fast evaporation of water molecules from the cluster. The effect of fast ion fragmentation can be clearly seen in Figure 12 and Figure 13, which display the mass spectra of the styrene/water clusters obtained at the resonance features assigned the SW and SW<sub>2</sub> cluster origins, respectively. It is clear that the ion intensity of the SW complex obtained at the origin of the SW<sub>2</sub> cluster is much higher than that obtained at the SW origin. Fragmentation probabilities for the SW<sup>+</sup> and SW<sup>+</sup><sub>2</sub> cluster ions were measured by comparing the absolute integrated intensities I<sub>n</sub> at various mass channels (parent and daughter) at the same wavelength of an identified spectral resonance and normalizing to the sum over I<sub>n</sub>. The results expressed as a percent of the parent ion intensity are shown in Table 7. Within the two photon's energy limit,

93% of the SW complex fragments into the styrene mass channel (S<sup>+</sup>) upon photoionization. The fragmentation efficiency of the SW<sub>2</sub><sup>+</sup>  $\rightarrow$  SW<sup>+</sup> process is estimated as 63% from the origin peak and between 67%-73% from the vdW progressions. The SW<sub>2</sub><sup>+</sup> cluster ion also fragments into the S<sup>+</sup> mass channel by the loss of two water molecules (29% at the SW<sub>2</sub> origin at 45.7 cm<sup>-1</sup>).

In order to support the assignments of the cluster origins, the dependence of the intensities of the spectral features assigned to the origins of the SW and SW<sub>2</sub> clusters on the water vapor concentration in the pre-expansion mixture was investigated. The results, shown in Figure 14, indicate that the peak at 45.7 cm<sup>-1</sup> (assigned to the SW<sub>2</sub> origin) decreases much more rapidly with decreasing the water concentration as compared to the peak at 21.5 cm<sup>-1</sup> (assigned to the SW origin). This behavior indicates that the two peaks must be due to different water containing clusters. The very weak dependence of the peak at 21.5 cm<sup>-1</sup> and the strong dependence of the peak at 45.7 cm<sup>-1</sup> on the water concentration support the assignments of these peaks to the SW and SW<sub>2</sub> clusters, respectively.

## 4.2.2 R2PI spectra of styrene (water)<sub>n</sub> clusters, n=3-5

The spectra detected at the mass channels of SW<sub>3</sub>, SW<sub>4</sub> and SW<sub>5</sub> clusters are displayed in Figure 15. The origin of SW<sub>3</sub> is blue shifted by 64.2 cm<sup>-1</sup> with respect to the  $0_0^0$  of styrene at 34,758.79 cm<sup>-1</sup>. Two peaks at 84.7 and 106.1 cm<sup>-1</sup> from the  $0_0^0$  of styrene were assigned the vdW progressions developed on the SW<sub>3</sub> origin. There are several other features present in the R2PI scan of the SW<sub>3</sub> mass channel as well as the main peaks described. The peaks at 31.8 and 77.3 cm<sup>-1</sup> appear to be coming from the features of SW<sub>5</sub>

mass channels. The other features in the  $SW_3$  mass channels are not assigned. The R2PI mass spectrum obtained at the resonance wavelength of the  $SW_3$  origin at 287.165 nm is shown in Figure 16.

The spectral shift reverses its direction for the origin band of SW<sub>4</sub> at 20.7 cm<sup>-1</sup> as shown in the SW<sub>4</sub> mass channel in Figure 15. This observation suggests that the structure of SW<sub>4</sub> incorporates dispersion interaction between styrene molecule and water tetramer subcluster. Furthermore, this assumption is supported by the low fragmentation probability of SW<sub>4</sub><sup>+</sup> cluster ion into SW<sub>3</sub> mass channel (Table 7). Two van der Waals progressions were assigned to SW<sub>4</sub> at 39.7 and 58.1 cm<sup>-1</sup> relative to the  $0_0^0$  band of styrene molecule (Table 6). The other spectral features at 31.8, 62.3 and 77.3 cm<sup>-1</sup> in SW<sub>4</sub> mass channel have the same fingerprints of SW<sub>5</sub> cluster (Figure 15).

In SW<sub>5</sub> mass channel, three major spectral features were assigned to SW<sub>5</sub> neutral cluster. The  $0_0^0$  origin band of SW<sub>5</sub> was blue shifted at 31.8 cm<sup>-1</sup> with respect to the S<sub>1</sub> $\leftarrow$ S<sub>0</sub> transition of styrene molecule as shown in Figure 15. The built-in progressions of van der Waals modes are located at 62.3 and 77.3 cm<sup>-1</sup>. The SW<sub>5</sub><sup>+</sup> cluster ion experiences considerable fragmentation efficiencies into SW<sub>4</sub> mass channel as shown in Table 7. The tentative assignments of the R2PI spectra of SW<sub>n</sub> (n=1-5) clusters are listed in Table 6

The overall R2PI scan of styrene-(water)<sub>n</sub> clusters, where n = 1-5, is displayed on Figure 17. Figure 18 presents the observed spectral shifts of the styrene(water)<sub>n</sub> cluster origins, with respect to the  $0_0^0$  band of styrene molecule, as a function of of water molecules. The first three clusters of SW<sub>n</sub> (n=1-3) showed linear increases in the direction

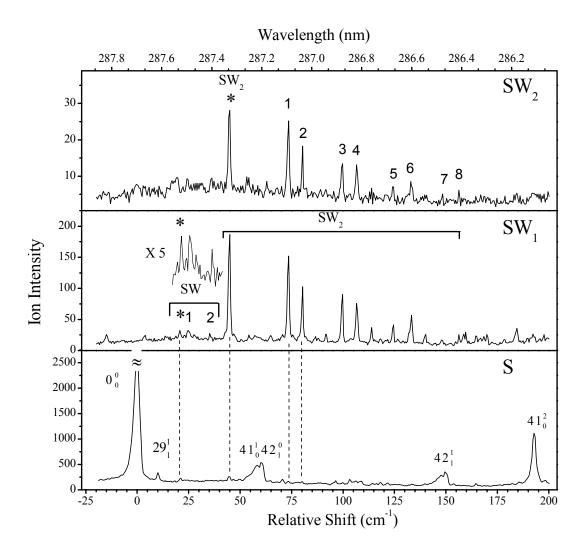


Figure 11: One-color R2PI spectra measured in the styrene mass channel [S] and styrene(water)<sub>n</sub> mass channels, with n = 1 and 2, [SW<sub>1</sub> and SW<sub>2</sub>] relative to the electronic origin band of the S<sub>1</sub>  $\leftarrow$  S<sub>0</sub> transition of styrene at 34,758.79 cm<sup>-1</sup>. In SW<sub>1</sub> and SW<sub>2</sub> mass channels, the origin of each cluster isomer is marked with (\*) and the peaks labeled with numbers following each cluster's origin represent vdW bands associated with the cluster origin (see Table 6).

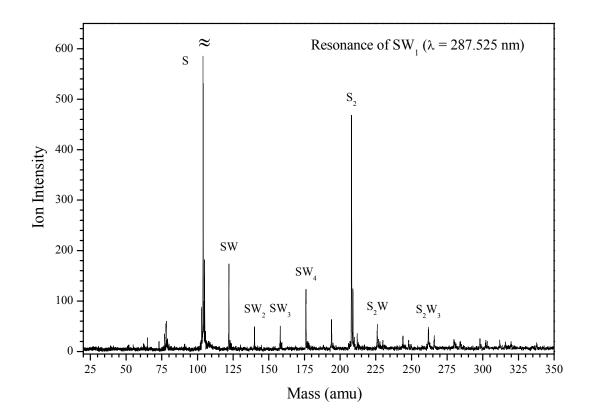


Figure 12: R2PI mass spectra of the styrene-water  $(SW_n)$  cluster beam obtained at the resonance ionization assigned to the origin of  $SW_1$ .

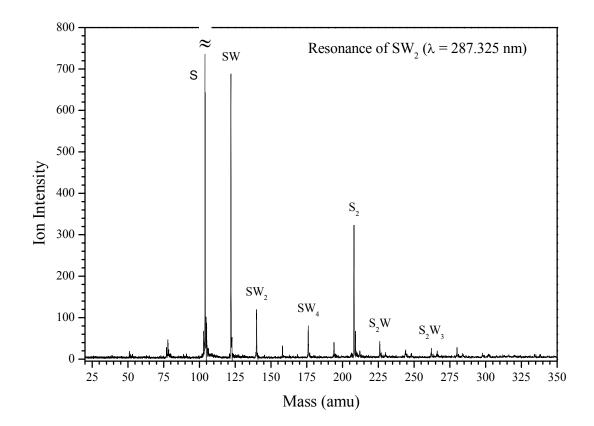


Figure 13: R2PI mass spectra of the styrene-water  $(SW_n)$  cluster beam obtained at the resonance ionization assigned to the origin of  $SW_2$ .

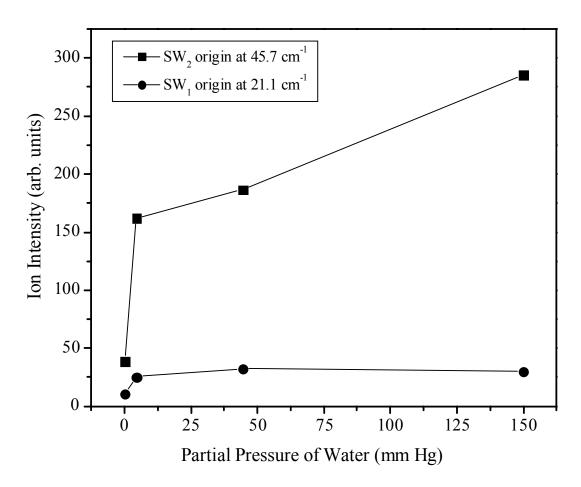


Figure 14: The dependence of the ion intensity assigned to the SW and SW<sub>2</sub> clusters on the water vapor pressure in the pre-expansion mixture. The plot points represent water vapor pressures of 0.005, 4.43, 44.56 and 149.97 mm Hg.

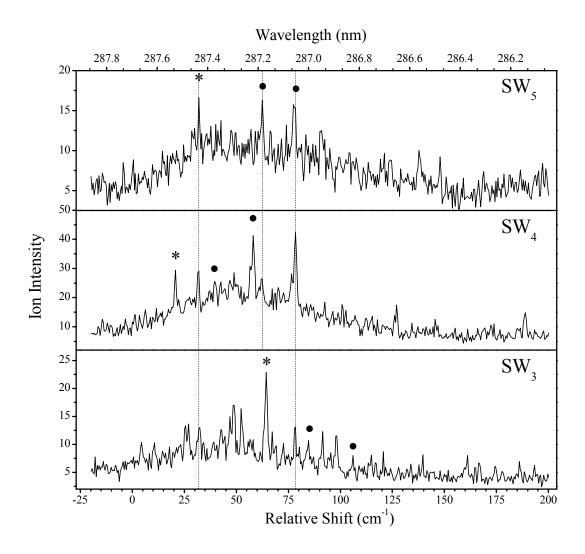


Figure 15: One-color R2PI spectra measured in the styrene(water)<sub>n</sub> mass channels, with n = 3-5, [SW<sub>3</sub>, SW<sub>4</sub> and SW<sub>5</sub>] relative to the electronic origin band of the S<sub>1</sub>  $\leftarrow$  S<sub>0</sub> transition of styrene at 34,758.79 cm<sup>-1</sup>. The origin of each cluster isomer is marked with (\*). The peaks labeled with symbols following each cluster's origin represent vdW bands associated with the cluster origin (see Table 6).

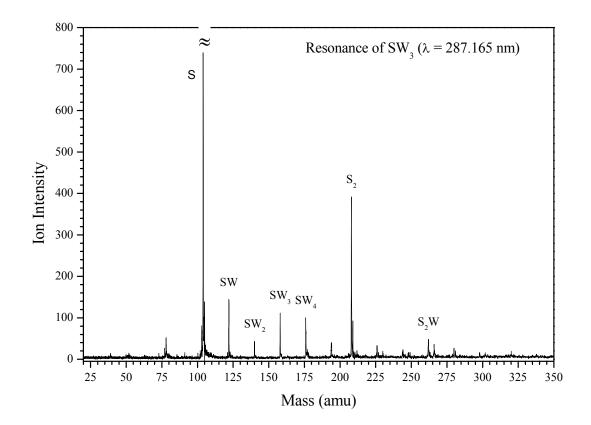


Figure 16: R2PI mass spectra of the styrene-water  $(SW_n)$  cluster beam obtained at the resonance ionization assigned to the origin of  $SW_3$ .

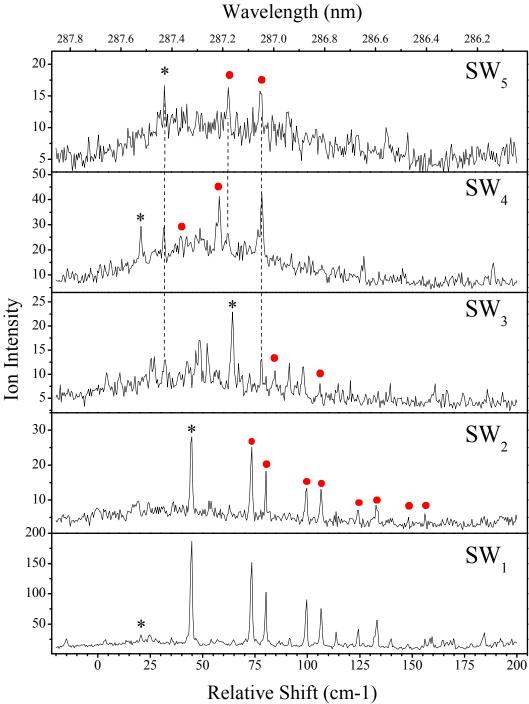


Figure 17: Overall R2PI spectra measured in the styrene(water)<sub>n</sub> mass channels, with n = 1-5, relative to the electronic origin band of the  $S_1 \leftarrow S_0$  transition of styrene at 34,758.79 cm<sup>-1</sup>. The origin of each cluster isomer is marked with (\*). The peaks labeled with symbols following each cluster's origin represent vdW bands associated with the cluster origin (see Table 6).

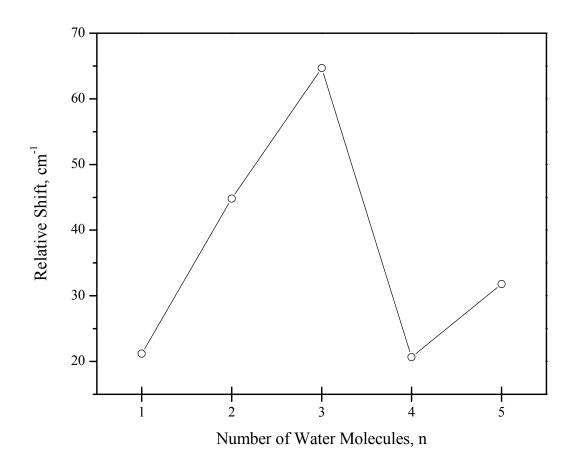


Figure 18: Spectral shifts of the origins of the  $SW_n$  clusters, relative to the  $0_0^0$  band of styrene, as a function of number of water molecules, n.

	vdW = van der	Waals Mode)			
Cluster	Shift ( $cm^{-1}$ )	Observed	Relative	Assignments	Shift $(cm^{-1})$
	from the $0_0^0$ of	mass channel	intensity		from cluster's
	styrene				origin
SW	21.2	SW	30.4	SW origin	0
		<b></b>			<u>^</u>
$SW_2$	44.8	SW	285.7	SW <sub>2</sub> origin	0
	44.8	$SW_2$	40.3	SW <sub>2</sub> origin	0
	73.4	SW	220.8	vdW	28.6
	73.4	$SW_2$	30.7	vdW	28.6
	99.6	SW	171.3	vdW	54.8
	99.6	$SW_2$	23.7	vdW	54.8
	106.7	SW	149.8	vdW	61.9
	106.7	$SW_2$	18.8	vdW	61.9
	124.3	SW	103.5	vdW	79.5
	124.3	$SW_2$	15.4	vdW	79.5
	133.1	SW	155.5	vdW	88.3
	133.1	$SW_2$	18.0	vdW	88.3
	148.1	SW	43.3	vdW	103.3
	147.6	$SW_2$	7.8	vdW	102.8
	156.2	SW	49.8	vdW	111.4
	156.0	$SW_2$	7.6	vdW	111.2
$SW_3$	64.2	$SW_3$	22.6	SW <sub>3</sub> origin	0
5113	84.7	SW3	107	vdW	20.5
	106.1	SW <sub>3</sub>	8.0	vdW	41.9
$SW_4$	20.7	$SW_4$	29.4	SW4 origin	0
<b>S W</b> 4			29.4	vdW	
	39.7	$SW_4$			19.0
	58.1	$SW_4$	41.3	vdW	37.4
$SW_5$	31.8	$SW_3$	12.8	SW5 origin	0
	31.8	$\mathrm{SW}_4$	29.1	SW5 origin	0
	31.8	$SW_5$	16.6	SW5 origin	0
	62.3	$SW_4$	26.4	vdW	305
	62.3	$SW_5$	16.3	vdW	30.5
	77.3	$SW_3$	13.1	vdW	45.5
	77.3	$SW_4$	42.2	vdW	45.5
	77.3	$SW_5$	15.8	vdW	45.5

Table 6: Spectral features observed in the  $SW_n$  mass channels (S = Styrene, W = Water, vdW = van der Waals Mode)

Fragmentation channel	Parent, shift (cm <sup>-1</sup> )	%
$SW^+ \rightarrow S^+$	SW origin (21.2)	93.2
$SW_2^+ \rightarrow SW^+$	SW <sub>2</sub> origin (44.8)	62.6
	vdW 1 (73.4)	67.5
	vdW 2 (80.2)	67.7
	vdW 3 (99.6)	66.5
	vdW 4 (106.6)	63.0
$SW_2^+ \rightarrow S^+$	$SW_2$ origin (44.8)	29.3
2	vdW 1 (73.4)	21.6
	vdW 2 (80.2)	21.6
	vdW 3 (99.6)	24.6
	vdW 4 (106.6)	24.6
$SW_3^+ \rightarrow SW_2^+$	SW <sub>3</sub> origin (64.7)	39.3
$SW_4^+ \rightarrow SW_3^+$	SW <sub>4</sub> origin (20.7)	19.6
$SW_5^+ \rightarrow SW_4^+$	SW <sub>5</sub> origin (31.8)	43.9

Table 7: Fragmentations probabilities of the  $SW_n$  cluster ions

### 4.3 R2PI of Styrene-Methanol Clusters

Shown in Figure 19 are the R2PI spectra obtained by monitoring the mass channel corresponding to styrene (S) in the  $0_0^0$  region between 34,700 cm<sup>-1</sup> and 35,000 cm<sup>-1</sup>. The spectrum has the  $0_0^0$  origin and the vibronic bands  $41_0^2$  and  $29_0^1$  as well as the cross – sequence bands  $41_0^142_0^1$  between the  $v_{42}$  and  $v_{41}$  modes (ethylene group torsional rotation and out-of-plane bend, respectively),  $41_0^142_1^0$  and  $41_0^342_1^0$ . The spectrum is in good agreement with previous fluorescence excitation and REMPI studies.[46, 97, 98] The other small features observable in Figure 19 are hot bands ( $29_1^1$  and  $42_1^1$ ) and some minor features (marked with asterisks) correspond to SM<sub>n</sub> clusters that fragment to the S mass channel. These features will be discussed in the next section.

#### 4.3.1 Spectra of Small SM<sub>n</sub> Clusters, n = 1-3

Displayed in Figure 20 are the spectra obtained by monitoring the mass channels corresponding to SM, SM<sub>2</sub> and SM<sub>3</sub> clusters. The SM channel shows several strong features, which appear in the SM<sub>2</sub> channel with very small intensities. The peak at 63.5 cm<sup>-1</sup> relative to the  $0_0^0$  of styrene (not present in the SM<sub>2</sub> mass channel) was assigned to the  $0_0^0$  of the SM complex with a small peak at 78 cm<sup>-1</sup> assigned to a vdW vibration of the complex. Several other peaks at 95, 106, 116 and 126.5 cm<sup>-1</sup> from the  $0_0^0$  of styrene were assigned to a vdW progression built on the SM origin.

The stronger peak at 83.8 cm<sup>-1</sup> accompanied by two rich vdW progressions was assigned to one isomer of the SM<sub>2</sub> cluster (SM<sub>2</sub>-I). The SM<sub>2</sub>-I cluster possesses two long

progressions of bands assigned to the excitation of intermolecular vibrational modes. The first progression appears at 102, 121, 140.6, 157 and 176.6 cm<sup>-1</sup> (corresponding to peaks 1-5, respectively in the SM channel in Figure 20). The second progression corresponds to peaks a, b, c, d and e shown in the SM channel at frequencies 96, 106, 116, 126.5 and 136 cm<sup>-1</sup>, respectively. It should be noted that the peaks a, b, c and d are relatively broad and they overlap with the vdW progression assigned to the SM cluster. Therefore, it appears that the series of peaks a-e contain contributions from both the vdW progressions of the SM and SM<sub>2</sub> clusters. The appearance of strong vibronic intensities away from the origin of the SM and SM<sub>2</sub>-I clusters indicates that Franck-Condon factors do not favor the cluster origin intensity, thus the ground state and excited state geometries of the SM and SM<sub>2</sub>-I isomer are quite different.

The SM<sub>2</sub> mass channel shows three sets of peaks. The small red-shifted peaks, relative to the  $0_0^0$  of the styrene, are due to larger clusters containing the styrene dimer S<sub>2</sub>. These features were found in the mass channels corresponding to S<sub>2</sub>M<sub>2</sub> and S<sub>2</sub>M<sub>3</sub> clusters at the same frequencies as they appear in the SM<sub>2</sub> channel. The second group consists of a strong peak that blue shifted by 3.4 cm<sup>-1</sup> from the  $0_0^0$  of styrene. This peak was assigned to a second isomer of the SM<sub>2</sub> cluster (SM<sub>2</sub>-II) with a vdW progression at 32, 49, 60 and 72 cm<sup>-1</sup> (corresponding to peaks 1-4 following the SM<sub>2</sub>-II origin in the SM<sub>2</sub> channel in Figure 20). It is interesting that the second isomer SM<sub>2</sub>-II does not show significant fragmentation upon photoionization. This suggests that the SM<sub>2</sub>-II isomer possesses a structure distinct

from that of the  $SM_2$ -I isomer, and that this structure does not appear to undergo major changes upon photoionization.

The third group of peaks observed in the  $SM_2$  mass channel starts with a sharp peak at 182 cm<sup>-1</sup> followed by a vdW progression of very low frequency modes at 189, 199, 204, 213, 232 and 341 cm<sup>-1</sup> (corresponding to peaks 1-6 following the  $SM_3$ -I origin in the  $SM_2$ channel in Figure 20). Some of these peaks appear with very small intensities in the  $SM_3$ mass channel, and they may be tentatively assigned to one isomer of the  $SM_3$  cluster ( $SM_3$ -I).

The most prominent features in the SM<sub>3</sub> mass channel in Figure 20 could be assigned to a second isomer of the SM<sub>3</sub> cluster with a strong origin peak blue shifted from the styrene origin by 16.5 cm<sup>-1</sup>. This origin is accompanied by a strong vibronic intensity with a well-defined vdW progression of doublets at 24, 33, 43, 52, 62, 70 and 80 cm<sup>-1</sup>. In addition, two other weaker peaks appear at 28 and 36 cm<sup>-1</sup>. Several weaker features to the red of the SM<sub>3</sub>-II isomer are due to fragmentation from the SM<sub>4</sub> cluster.

The blue shift observed for the SM cluster (63 cm<sup>-1</sup>) is consistent with a weak hydrogen bonding interaction between the OH group of methanol and the styrene  $\pi$ system. For example, benzene complexes with H<sub>2</sub>O, CH<sub>3</sub>OH, CH<sub>3</sub>COOH, HCl and CHCl<sub>3</sub> exhibit blue shifts of 52, 44, 152, 125 and 179 cm<sup>-1</sup>, respectively.[20-23, 25, 30] It is interesting to note that the addition of a second methanol molecule to styrene resulted in a greater blue shift in the SM<sub>2</sub>-I cluster (84 cm<sup>-1</sup>). This is very similar to the benzene (methanol)<sub>n</sub> clusters, BM<sub>n</sub>, where the spectral shifts of BM and BM<sub>2</sub> clusters relative to the  $6_0^1$  origin of the isolated benzene molecule are 44 and 80 cm<sup>-1</sup>, respectively.[25] Also, the spectrum of the SM<sub>2</sub> cluster exhibits several vdW transitions indicating that the 1:2 clusters undergo major changes in geometries upon the electronic excitation of styrene. This is exactly the same trend found in the BM<sub>2</sub> clusters.[25] Furthermore, in both the benzene and styrene systems extensive fragmentation was observed following photoionization of the neutral clusters. The high fragmentation probability of the ionized clusters is consistent with the hydrogen bonding interaction with the  $\pi$ -system in benzene or styrene.

A noticeable difference between the benzene and styrene systems is the assignment of two structural isomers to the SM<sub>2</sub> cluster with quite different spectral shifts (84 and 3.4 cm<sup>-1</sup> relative to the 0<sup>0</sup><sub>0</sub> transition of styrene). This is in contrast to the BM<sub>2</sub> cluster, where only one isomer was assigned.[25] The interactions between the ethylene group of styrene and the methyl group for methanol are considered to be purely dispersive. This suggests that the SM<sub>2</sub>-II isomer may represent a structure that maximizes induced dipole-induced dipole forces between the ethylene group of styrene and the methyl groups of the two methanol molecules. This interaction is known to lead to red spectral shifts, and, therefore, one expects the SM<sub>2</sub>-II isomer to exhibit a significant red shift relative to the SM complex ( $\Delta v = -60$  cm<sup>-1</sup>). The dominance of the dispersion interaction and the small  $\pi$ -hydrogen bonding interaction in the SM<sub>2</sub>-II isomer following photoionization.

The blue-shifted isomer assigned to the SM<sub>3</sub> cluster (SM<sub>3</sub>-I) may have a hydrogenbonded methanol chain structure H-bonded to the styrene  $\pi$ -system. This structure is compatible with the strong blue shift observed (182 cm<sup>-1</sup> with respect to the styrene origin). The second isomer (SM<sub>3</sub>-II) may reflect a cyclic methanol structure with no hydrogen bonding to styrene since all the OH groups would be used in the formation of the methanol trimer ring. This isomer has a smaller blue shift of 16 cm<sup>-1</sup>, which is consistent with the absence of hydrogen bonding interaction to the styrene  $\pi$ -system. The calculated structures for the SM<sub>n</sub> clusters are consistent with the structural predictions presented here based on the observed spectral shifts and the degree of cluster's fragmentation following ionization.[99]

## 4.3.2 Spectra of $SM_n$ Clusters, n = 4-6

The R2PI spectra obtained by monitoring the mass channels corresponding to the  $SM_4$ ,  $SM_5$  and  $SM_6$  clusters in the region of the styrene  $0_0^0$  transition are displayed in Figure 21. A remarkable shift in the spectra from blue to red, with respect to the styrene origin, is observable at the start of the  $SM_4$  cluster. In additions, ion fragmentations from the S<sup>+</sup>M\_5 and S<sup>+</sup>M\_6 clusters to the S<sup>+</sup>M\_4 channel appear to be pronounced since several features belonging to the  $SM_5$  and  $SM_6$  clusters are present in the S<sup>+</sup>M\_4 mass channel. The set of strong peaks in the -98 to -60 cm<sup>-1</sup> region were assigned to the  $SM_4$  cluster with the origin assigned to the singlet at -94 cm<sup>-1</sup>. A vdW progression associated with the cluster origin also appears at -88, -79, -72, -66 and -62 cm<sup>-1</sup> (corresponding to peaks 1-5 following the  $SM_4$  origin in the  $SM_2$  channel in Figure 21). The cluster origin and the vdW modes also appear in the  $SM_3$  channel as a result of the ion fragmentation process S<sup>+</sup>M\_4  $\rightarrow$  S<sup>+</sup>M\_3.

Other features present in the  $S^+M_4$  mass channel, as a result of dissociation of the ionized clusters, include two isomers of each of the  $SM_5$  and  $SM_6$  clusters.

The red shifted spectrum of the  $SM_4$  cluster, unlike the SM,  $SM_2$  and  $SM_3$  clusters, may indicate a major departure from the structure that allows hydrogen bonding interaction between the OH group of methanol and the styrene  $\pi$ -system. A possible structure may involve a cyclic methanol tetramer placed above the styrene molecular plane. This structure is expected to enhance the dispersion and induction interactions, which typically result in red spectral shifts. The calculated lowest energy structures of the SM<sub>4</sub> cluster[99] are consistent with the proposed cyclic methanol tetramer above the styrene plane.

Two geometrical isomers are assigned to the  $SM_5$  cluster, which appear in both the  $SM_5$  and  $SM_4$  mass channels. The first isomer ( $SM_5$ -I) has a strong origin at -59 cm<sup>-1</sup> and possesses two vdW peaks at -41 and -25 cm<sup>-1</sup> (corresponding to peaks 1 and 2 following the  $SM_5$ -I origin in the  $SM_5$  channel in Figure 21). The second isomer ( $SM_5$ -II) displays a weak low energy origin at -119 cm<sup>-1</sup> followed by a stronger vdW mutiplet at -115, -110, - 105 and -103 cm<sup>-1</sup> (corresponding to peaks 1-4 following the  $SM_5$ -II origin in the  $SM_5$  channel in Figure 21). Other weak vdW peaks appear at -108 and -98 cm<sup>-1</sup> and possess significant intensities in the  $SM_4$  mass channel. The origins of both isomers and their vdW features also appear in the  $S^+M_4$  mass channel as a result of ion fragmentation.

The  $SM_6$  mass channel contains features belonging to both the  $SM_6$  and  $SM_7$  clusters. Two isomers  $SM_6$ -I and  $SM_6$ -II can be assigned to the  $SM_6$  cluster with origins at -46 and -113 cm<sup>-1</sup>, respectively. Several weaker features with a long progression are

assigned to vdW modes associated with the  $SM_6$ -I isomer and they appear at -31, -14, -2, 10 and 21 cm<sup>-1</sup> from the styrene origin (corresponding to peaks 1-5 following the  $SM_6$ -I origin in the  $SM_6$  channel in Figure 21). The  $SM_6$ -I isomer undergoes extensive fragmentation to both the  $SM_5$  and the  $SM_4$  mass channels following photoionization of the cluster. The fragmentation also proceeds through the associated vdW modes as evident in the  $SM_4$  and  $SM_5$  mass channels. The  $SM_6$ -II origin (-113 cm<sup>-1</sup>) has two vdW peaks at -100 and -96 cm<sup>-1</sup> and both fragment to the  $SM_5$  mass channel.

# 4.3.3 Spectra of $SM_n$ Clusters, n = 7-9

Displayed in Figure 22 are the R2PI spectra obtained by monitoring the mass channels corresponding to the  $SM_7$ ,  $SM_8$  and  $SM_9$  clusters. The  $SM_7$  cluster shows two distinct isomers ( $SM_7$ -I and  $SM_7$ -II) with origins at +54 and -89 cm<sup>-1</sup>, respectively. Both clusters fragment efficiently upon photoionization into the  $SM_6$  mass channel. The  $SM_7$ -I isomer exhibits a small vdW multiplet at 64, 66, 70, 73 and 79 cm<sup>-1</sup> with greater intensity in the  $SM_6$  mass channel due to ion fragmentation (corresponding to peaks 2-6 following the  $SM_7$ -I origin in the  $SM_6$  channel in Figure 21). The second isomer  $SM_7$ -II shows small vdW peaks at -80/-79, -75 and -73 cm<sup>-1</sup>. The remaining features in the  $SM_7$  mass channel belong to the  $SM_8$  and  $SM_9$  clusters. The  $SM_8$  cluster origin appears in the  $SM_8$  mass channel at -36/-34 cm<sup>-1</sup>. A vdW progression associated with the  $SM_8$  cluster appears at -28, -23, -14, -8 and -3 cm<sup>-1</sup>, corresponding to peaks 1-5, respectively following the  $SM_8$  origin in the  $SM_8$  channel in Figure 22. The  $SM_8$  cluster has a doublet origin and shows strong

fragmentation into the SM<sub>7</sub> mass channel through the vdW modes which appear at -34, -28, -23, -14, -8 and -5 cm<sup>-1</sup>. The other features in the  $SM_{g}$  mass channel can be assigned to the SM<sub>9</sub> cluster as shown in Figure 22. Based on the comparison of the spectral features in the  $SM_{o}$  and  $SM_{o}$  mass channels, we assign the peak at -120 cm<sup>-1</sup> to one isomer of the  $SM_{o}$ cluster (SM<sub>o</sub>-II). This isomer shows several vdW peaks at -111, -105, -100, -94, -87, -78, -74, -68 and -63 cm<sup>-1</sup> corresponding to peaks 1-9, respectively following the SM<sub>9</sub>-II origin in the SM<sub>9</sub> channel in Figure 22. Another origin in the SM<sub>9</sub> mass channel at +24  $\rm cm^{-1}$ followed by a doublet at +29/30 cm<sup>-1</sup> and a vdW peak at +46 cm<sup>-1</sup> is assigned to a second isomer of the SM<sub>o</sub> cluster (SM<sub>o</sub>-I). It is interesting to note that all the assigned isomers of the  $SM_n$  clusters with n = 4-9 exhibit red spectral shifts, relative to the styrene origin, except for the SM7-I and SM9-I isomers which display blue shifts. Figure 23 exhibits the overall R2PI spectra recorded in the SM<sub>n</sub> mass channels with n=1-9, displayed relative to the  $0_0^0$  origin of styrene. The relative frequencies and tentative assignments of the SM<sub>n</sub> cluster isomers are listed in Table 8.

### 4.3.4 Spectral Shifts

The spectral shifts relative to the  $0_0^0$  origin of the isolated styrene molecule imposed by methanol clusters provide information on the nature of the intermolecular interactions within the binary clusters. Spectral shifts are due to different cluster binding energies in the ground and excited states. A red shift implies that the excited state is more tightly bound than the ground state, and a blue shift is associated with stronger interaction in the ground state. A red shift is usually observed in clusters where the dispersion energy is dominant due to the increase in the molecular polarizability in the excited state relative to the ground state.[100-102] For example, ring-ring interactions result in red shifts since dispersive forces are stronger for the more delocalized excited states, as in benzene and other aromatic clusters.[17, 103] On the other hand, hydrogen-bonding to the  $\pi$ -system results in a blue shift, which tends to increase with increasing the H-bond donating capacity of the solvent molecules.[10]

In contrast to benzene, few studies of clusters containing styrene have been reported.[49-51, 104-111] Clusters of styrene with Ar, N<sub>2</sub>, CO<sub>2</sub>, NH<sub>3</sub> and trimethylamine were studied by laser-induced fluorescence (LIF) and REMPI.[47, 49-51, 105, 106, 108, 110, 112] The origin bands of these clusters are shifted by -31 cm<sup>-1</sup>, -25.5 cm<sup>-1</sup>, +51 cm<sup>-1</sup>, +52 cm<sup>-1</sup>, and -24 cm<sup>-1</sup> relative to the  $0_0^0$  band of the bare styrene molecule. The spectra of the styrene-TMA 1:1 cluster were assigned to two isomeric forms, both capable of forming a charge transfer complex upon excitation to the locally excited state.[105, 106] We recently reported the REMPI spectra of styrene (water)<sub>n</sub> clusters, SW<sub>n</sub>, with n = 1-2, which indicate that the origin of the SW complex is blue shifted by -22 cm<sup>-1</sup> relative to the  $0_0^0$  band of the bare styrene molecule.[104]

In the case of methanol clusters, the binding energy is mainly due to hydrogen bonding. Monte Carlo, molecular dynamics, density functional and *ab initio* calculations have shown that the intermolecular interactions in small methanol clusters (n = 4-9) are dominated by configurations in which hydrogen-bonded chains are formed with cyclic structures starting at n = 3.[113-117] If it is assumed that the styrene molecule does not significantly affect hydrogen bonding in methanol clusters, then shifts of the electronic origin of  $SM_n$  clusters can be related to the styrene  $\leftrightarrow$  (methanol subcluster) interaction.

Displayed in Figure 24 are the observed spectral shifts of the SM<sub>n</sub> cluster origins, relative to the  $0_0^0$  transition of styrene, as a function of n. The observed pattern reveals some interesting correlation with the proposed structures of the clusters. Figure 23 and Figure 24 show that the first isomers of SM<sub>n</sub> (n=1-3) exhibit successively greater increases in blue shift, which reaches a maximum of 182 cm<sup>-1</sup> for the SM<sub>3</sub>-I isomer. The shifts of these isomers are consistent with H-bonding to the  $\pi$ -system of styrene. The SM<sub>2</sub>-II and SM<sub>3</sub>-II isomers exhibit smaller blue shifts, which may reflect more dispersive interactions suggesting structures with no directional hydrogen bonding to the  $\pi$ -system of styrene.[99]

The spectral shift reverses direction in the single isomer observed for SM<sub>4</sub>. The successive addition of methanol molecules to the SM complex leads to a stable cyclic methanol tetramer in SM<sub>4</sub> that is unable to H-bond to the styrene  $\pi$ -system, consistent with the observed red shift.[99]

The two isomers assigned to each of the  $SM_5$  and  $SM_6$  clusters may contain cyclic methanol subclusters. Similarly, the two  $SM_7$  cluster isomers could have cyclic five- or sixmember rings with the other molecule(s) as ring branches. The single isomer assigned to the  $SM_8$  cluster could involve a bicyclic methanol ring structure. Also, the two  $SM_9$  isomers could represent a larger bicyclic ring containing nine methanol molecules, and a bicyclic eight-member ring on one side of styrene and a free methanol molecule on the other side of the styrene plane. Interestingly, the addition of the red shift of -36 cm<sup>-1</sup> imposed by the  $SM_8$  cluster and the blue shift of  $+63 \text{ cm}^{-1}$  of the SM cluster results in a  $+27 \text{ cm}^{-1}$  shift, which is similar to the observed shift of the SM<sub>9</sub>-I isomer. This may suggest that the SM<sub>9</sub>-I isomer consists of a M<sub>8</sub> subcluster on one side of styrene and a single methanol molecule on the other. The arrangement of the methanol molecules with respect to the styrene plane has been investigated using pair potentials.[99]

# 4.3.5 Fragmentation of the Ionized Clusters

The ionization potential of the styrene molecule is 68,267 cm<sup>-1</sup>.[49] Since the  $0_0^0$  transition is located at 34,758.79 cm<sup>-1</sup>, two photons resonant with the origin transition suffice to ionize styrene in a one-color REMPI experiments. The ionic state of the styrene molecule in the SM<sub>n</sub> cluster is prepared according to the Franck-Condon distribution, while the rest of the cluster relaxes structurally in a time-dependent manner about the styrene ion. This relaxation releases energy and thus lowers the cluster ionization threshold relative to that of the isolated styrene molecule, and the additional energy imparted to the S<sup>+</sup>M<sub>n</sub> cluster ion promotes fragmentation via solvent evaporation from the cluster. However, since a large fraction of the energy usually goes to the electron, some cluster ions are formed with little internal energy and may not dissociate while in the acceleration fragmentation processes [S<sup>+</sup>M<sub>n</sub>  $\rightarrow$  S<sup>+</sup>M<sub>n-m</sub>] are faster than the time of acceleration in the mass spectrometer, then the resonant intensity of cluster n will contribute to the observed intensity of cluster n-m. These time scales will in general depend on the cluster's internal energy prior to excitation, the final

energy state produced by ionization, and the timescale of energy relaxation into the intermolecular modes.

Fragmentation probabilities are measured by comparing the absolute integrated intensities  $I_n$  at various mass channels (parent  $I_n$  and daughters  $I_{n-1}$ ,  $I_{n-2}$ , etc.) at the laser wavelength of a specific spectral resonance and normalizing to the sum over  $I_n + I_{n-1} + ...$  etc. The results expressed as a percent of the parent ion intensity are shown in Table 9.

The relatively efficient fragmentation observed for the S<sup>+</sup>M complex following photoionization is a direct consequence of the  $\pi$ -hydrogen-bonded geometry of the neutral complex. The structural change in going from the  $\pi$ -hydrogen-bonding geometry in the neutral species to the predominantly ion-dipole interaction in the ionized species involves a great strain, which leads to efficient fragmentation. However, the fragmentation of the S<sup>+</sup>M complex is significantly less than that of the B<sup>+</sup>M complex (where almost 100% fragmentation was observed). This is consistent with the methanol interacting with both the ethylene double bond and the aromatic ring.

Displayed in Figure 25 are the examples of the mass spectra obtained at some selected resonance features assigned to the  $SM_n$  cluster origins. In the two examples displayed in Figure 25, the S<sup>+</sup>M<sub>4</sub> and S<sup>+</sup>M<sub>5</sub>-I clusters show little fragmentation following photoionization, and the mass peaks corresponding to the parent resonance cluster show enhanced intensity. However, when fast ion fragmentation is predominant, significant intensity of the SM<sub>n-1</sub> ions are observed following the resonant ionization of the SM<sub>n</sub> clusters. This was clearly observed in the SM<sub>2</sub>-I, SM<sub>3</sub>-I and SM<sub>7</sub>-I cluster isomers. The small fragmentation probability of the

SM<sub>4</sub> cluster ion supports the assumption of a ring structure of the methanol tetramer placed above the plane of the styrene molecule.[99]

### 4.3.6 Intracluster Reactions following Photoionization

Mass spectra of the  $SM_n$  clusters obtained at selected resonance ionization energies reveal the generation of protonated methanol clusters  $H^+M_n$  starting at  $n \ge 2$ . Since pure methanol clusters (in the absence of styrene) could not be ionized with the same laser power used in the R2PI experiments, the origin of the protonated clusters must be attributed to the  $SM_n$  clusters. The  $H^+M_n$  clusters could be produced by either dissociative electron transfer (DET) or dissociative proton transfer (DPT), or combination of both mechanisms. The small intensity of the  $H^+M_n$  ions (due to the necessity of low laser power in order to stay within the two photon process) prevented correlation of their spectral features with those of the parent  $SM_n$  clusters and the unambiguous determination the origin of the protonated methanol clusters.

Shown in Figure 26 are the mass spectra obtained at wavelengths corresponding to the assigned SM<sub>3</sub>-I and SM<sub>3</sub>-II cluster origins. It was noted that the protonated methanol clusters are formed exclusively at the resonance for the SM<sub>3</sub>-I isomer (286.202 nm) and not at the SM<sub>3</sub>-II resonance. However, it should be noted that at this wavelength the mass spectrum contains significant contributions from the styrene dimer ( $S_2^{+}$ ) and the  $S_2M_n^{+}$  series. The styrene dimer cation has been a subject of several studies that have established it to be covalently bonded ( $C_{16}H_{16}^{+}$ ) and probably cyclic in structure.[37, 38, 94-96] The correlation between the generation of the  $S_2M_n^{+}$  series and the protonated methanol clusters suggests that

DPT takes place from the styrene dimer cation  $C_{16}H_{16}^{+}$  to the methanol sub-cluster  $M_n$  within the  $S_2M_n^{+}$  cluster to generate the  $H^+M_n$  species, as shown in Eq. 4-1. At lower laser power, the protonated methanol clusters are generated starting at n = 3. This suggests that the protonated methanol dimer is produced by evaporation from larger  $H^+M_n$  clusters with  $n \ge 3$ as shown in Eq. 4-2.

$$C_{16}H_{16}^{+}(CH_{3}OH)_{n} \rightarrow H^{+}(CH_{3}OH)_{n} + C_{16}H_{15} \cdots n \ge 3 \qquad \text{Eq. 4-1}$$
$$H^{+}(CH_{3}OH)_{n} \rightarrow H^{+}(CH_{3}OH)_{n-1} + CH_{3}OH \qquad \text{Eq. 4-2}$$

The proposed mechanism involves structural changes following the ionization of the styrene dimer resulting in the formation of a covalent bonded styrene dimer cation  $(C_{16}H_{16}^+)$  solvated by methanol molecules in the cluster. DPT between  $C_{16}H_{16}^+$  and methanol subcluster then, results in the generation of protonated methanol clusters  $[H^+(CH_3OH)_n]$  and the  $C_{16}H_{15}$  radical. It should be noted the DPT channel becomes exothermic in the  $S_2M_3^+$  cluster if the proton affinity (PA) of  $C_{16}H_{15}$  [ $C_{16}H_{16}^+ \rightarrow C_{16}H_{15} + H^+$ ] is less than that of the methanol trimer [ $H^+M_3 \rightarrow M_3 + H^+$ ]. Based on the estimation of the PA of M<sub>3</sub> as 224 kcal/mol, [82] the upper limit of the PA of the styrene dimer radical ( $C_{16}H_{15}$ ) can be estimated as  $\leq 224$  kcal/mol.

The DPT mechanism suggested for the generation of the protonated methanol clusters is consistent with the study of the styrene/water binary clusters, where isotopic experiments involving D<sub>2</sub>O have revealed that the protonated water clusters  $H^+(H_2O)_n$  were produced via a DPT mechanism involving the styrene dimer series  $[C_{16}H_{16}^+(D_2O)_n]$  with n≥3. The observed intracluster DPT reactions in  $S_2M_n^+$  series are consistent with the detrimental effect of methanol on the bulk cationic polymerization of methanol. For example, it is well established that radiation-induced polymerization of styrene proceeds via a cationic mechanism in pure, dry bulk styrene.[38-40, 118, 119] The presence of a small amount of methanol suppresses the cationic polymerization and enhances both the dimerization and the radical polymerization. It has been shown that methanol reacts with the dimer radical cation and converts it to a neutral radical that initiates the radical polymerization.[39] This is exactly the result obtained in the present work via the DPT mechanism within the S<sub>2</sub>M<sub>n</sub><sup>+</sup> clusters. The similarity between the cluster and the bulk phase results provides strong evidence for the validity of the cluster approach in elucidating the mechanisms of initiation, propagation and inhibition of cationic polymerization.

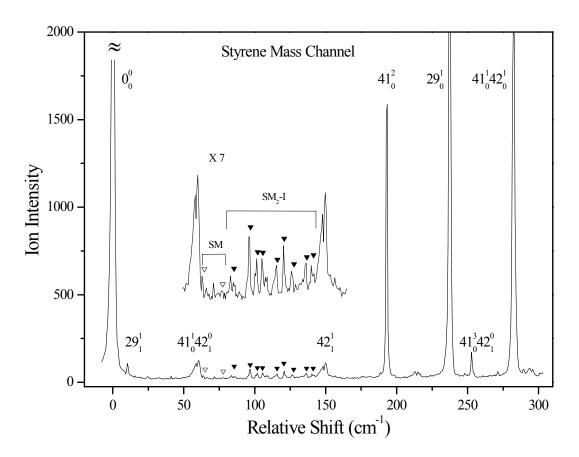


Figure 19: One-color R2PI spectra of styrene (S) mass in the region of  $S_1 \leftarrow S_0$  transition of the styrene molecule at 34,758.79 cm<sup>-1</sup>.

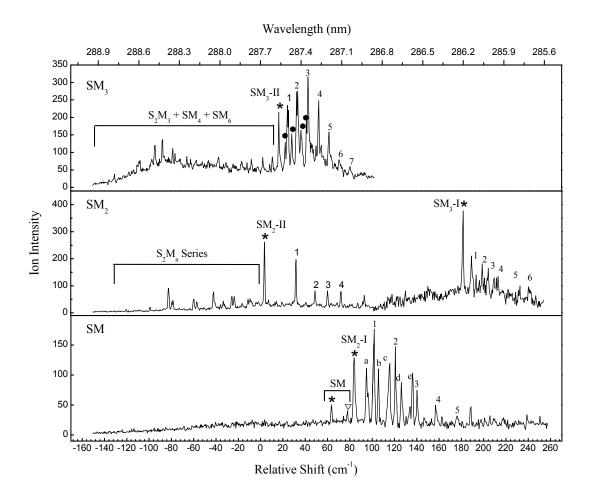


Figure 20: One-color R2PI spectra measured in the styrene (methanol)<sub>n</sub> mass channels with n = 1-3, [SM, SM<sub>3</sub> and SM<sub>3</sub>] relative to the electronic origin band of the S<sub>1</sub> $\leftarrow$ S<sub>0</sub> transition of the styrene molecule at 34,758.79 cm<sup>-1</sup>. The origin of each cluster isomer is marked with a (\*). The peaks labeled with numbers, letters or symbols following each cluster's origin represent vdW bands associated with the cluster origin (see Table 8).

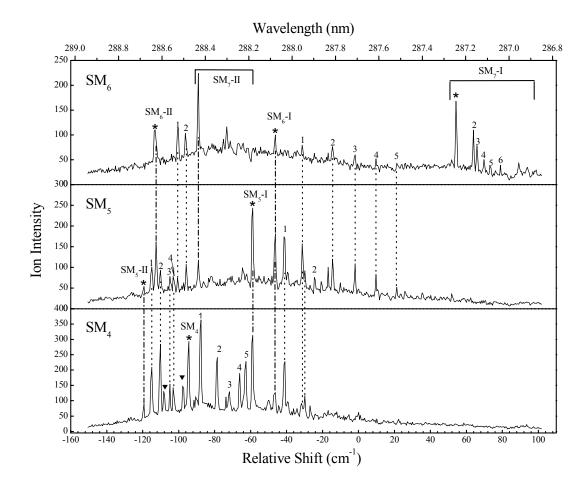


Figure 21: One-color R2PI spectra measured in the styrene (methanol)<sub>n</sub> mass channels with n = 4-6, [SM<sub>4</sub>, SM<sub>5</sub> and SM<sub>6</sub>] relative to the electronic origin band of the S<sub>1</sub>←S<sub>0</sub> transition of the styrene molecule at 34,758.79 cm<sup>-1</sup>. The origin of each cluster isomer is marked with a (\*). The peaks labeled with numbers following each cluster's origin represent vdW bands associated with the cluster origin (see Table 1). The two peaks labeled with V at -108 cm<sup>-1</sup> and -98 cm<sup>-1</sup> from the styrene origin in the SM<sub>4</sub> mass channel are unassigned.

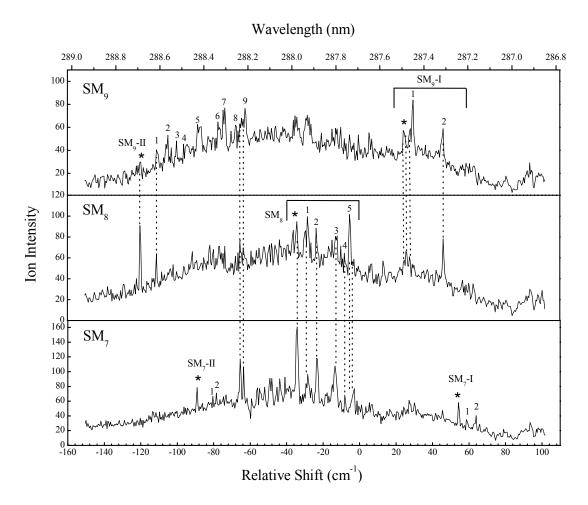


Figure 22: One-color R2PI spectra measured in the styrene (methanol)<sub>n</sub> mass channels with n = 7-9, [SM<sub>7</sub>, SM<sub>8</sub> and SM<sub>9</sub>] relative to the electronic origin band of the S<sub>1</sub>  $\leftarrow$  S<sub>0</sub> transition of the styrene molecule at 34,758.79 cm-1. The origin of each cluster isomer is marked with a (\*). The peaks labeled with numbers following each cluster's origin represent vdW bands associated with the cluster origin (see Table 8).

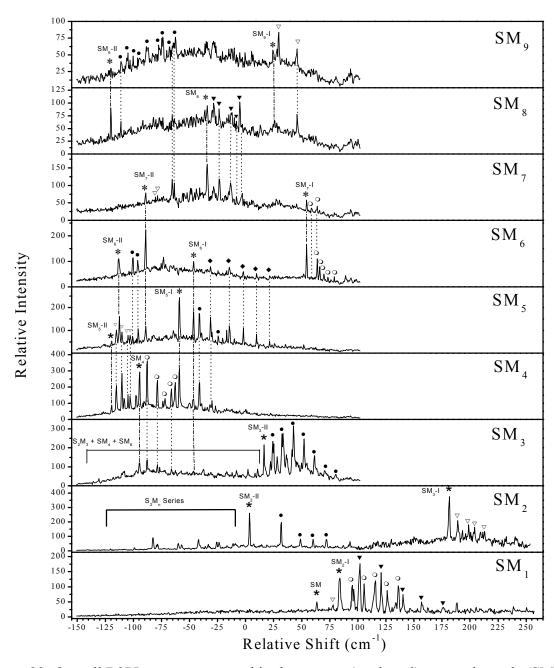
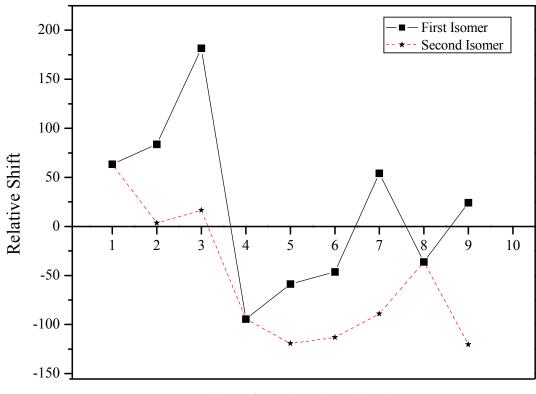


Figure 23: Overall R2PI spectra measured in the styrene (methanol)<sub>n</sub> mass channels (SM<sub>n</sub>), n = 1-9, relative to the electronic origin band of the S<sub>1</sub> $\leftarrow$ S<sub>0</sub> transition of the styrene molecule at 34,758.79cm<sup>-1</sup>. The origin of each cluster isomer is marked with a (\*). The peaks labeled with symbols following each cluster's origin represent vdW bands associated with the cluster origin (see Table 8).



Number of Methanol Molecules

Figure 24: Spectral shifts of the origins of the SMn clusters, relative to the  $0_0^0$  band of styrene, as a function of n.

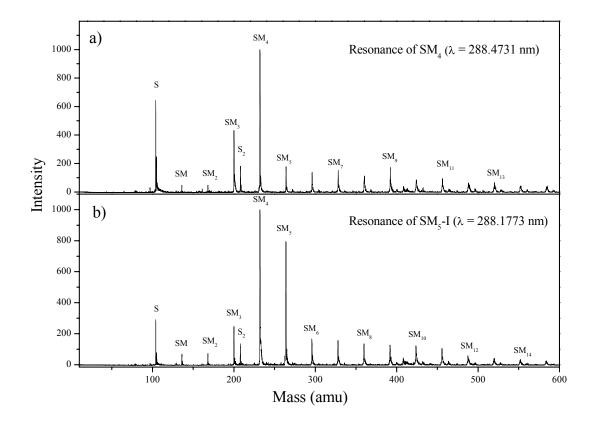


Figure 25: R2PI mass spectra of the styrene-methanol  $(SM_n)$  cluster beam obtained at the resonance ionizations assigned to the origins of (a)  $SM_4$  and (b)  $SM_5$ –I clusters.

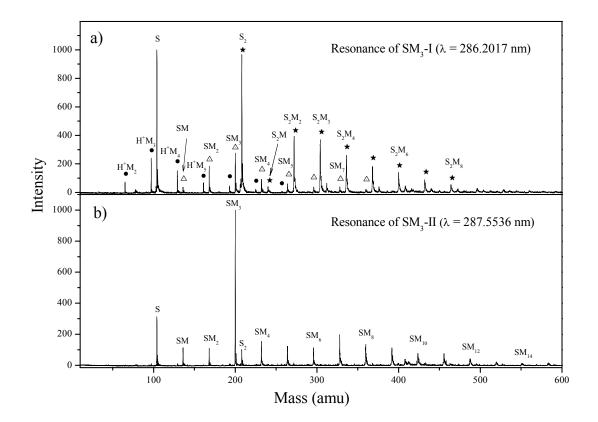


Figure 26: R2PI mass spectra of the styrene-methanol (SM<sub>n</sub>) cluster beam obtained at the resonance ionizations assigned to the origins of (a) SM<sub>3</sub>–I and (b) SM<sub>3</sub>–II clusters. The SM<sub>n</sub> series is labeled as  $\triangle$ . Note the correlation between the generations of the protonated methanol clusters H<sup>+</sup>M<sub>n</sub> (labeled as  $\bullet$ ) and the styrene dimer series S<sub>2</sub>M<sub>n</sub> (labeled as  $\bigstar$ ) as shown in (a).

vď	W = van der Waa	ıls mode).			
Cluster	Shift (cm <sup>-1</sup> ) from the $0_0^0$ of Styrene	Observed Mass Channel	Relative Intensity	Assignment	Shift (cm <sup>-1</sup> ) from Cluster's Origin
SM	63.4	SM	50	SM origin	0
	78.0	"	41	vdW	14.6
SM <sub>2</sub> -I	83.8	SM	128	SM <sub>2</sub> –I origin	0
	95.2/96.4	"	104/92	vdW-a	11.4/12.6
	102.2	>>	176	vdW-1	18.4
	105.8	"	110	vdW-b	22.0
	116.2	"	120	vdW-c	32.4
	121.1	"	147	vdW-2	37.3
	126.5	"	88	vdW-d	42.7
	136.3	"	103	vdW-e	52.6
	140.6	>>	75	vdW-3	56.8
	157.0	>>	50	vdW-4	73.2
	176.6	>>	32	vdW-5	92.6
SM <sub>2</sub> -II	3.4	$SM_2$	262	SM <sub>2</sub> –II origin	0
	32.0	"	197	vdW-1	28.6
	48.9	"	81	vdW-2	45.5
	60.1	"	91	vdW-3	56.7
	72.2	>>	79	vdW-4	68.8
SM <sub>3</sub> -I	181.6	$SM_2$	377	SM <sub>3</sub> –I origin	0
5	188.9		210	vdW-1	7.3
	198.7	>>	181	vdW-2	17.1
	204.1	>>	165	vdW-3	22.5
	212.9	>>	135	vdW-4	31.3
	232.4	>>	100	VdW-5	50.8
	241.2	>>	86	VdW-6	59.6
					_
SM <sub>3</sub> –II	16.5	$SM_3$	216	SM <sub>3</sub> –II origin	0
	24.2	>>	235	vdW-1	7.7
	28.1	>>	153	vdW	11.6
	33.4	>>	276	vdW-2	16.9
	36.3	"	164	vdW	19.8
	42.6	>>	315	vdW-3	26.1

Table 8: Spectral features observed in the  $SM_n$  mass channels (S = styrene, M = methanol, vdW = van der Waals mode).

			• 10		
	52.3	>>	248	vdW-4	35.8
	61.5	"	158	vdW-5	45
	70.2	"	79	vdW-6	53.7
	80.4	"	60	vdW-7	63.9
$SM_4$	-94.4	$SM_3$	121	SM4 origin	0
	"	$SM_4$	293	"	"
	-87.6	$SM_3$	137	vdW-1	6.8
	>>	$SM_4$	363	"	,,
	-78.5	$SM_3$	112	vdW-2	15.9
	>>	$SM_4$	241	"	"
	-71.8	$SM_3$	81	vdŴ-3	22.6
	>>	$SM_4$	130	"	"
	-66.0	$SM_3$	88	vdW-4	28.4
	"	$SM_4$	189	"	"
	-62.6	$SM_3$	81	vdW-5	31.8
	22	$SM_4$	228	"	"
SM5–I	-58.8	$SM_4$	314	SM <sub>5</sub> –I origin	0
51015-1	-50.0	SM <sub>4</sub> SM <sub>5</sub>	244	Sivis-i oligili	0
	-41.4	SM3	66	vdŴ-1	17.4
		SM <sub>3</sub> SM <sub>4</sub>	228		
	>>	SM <sub>4</sub> SM <sub>5</sub>	174	"	"
	-24.6	SM <sub>3</sub>	67	VdW-2	34.2
		SM <sub>5</sub>	76		
	>>	51415	70	>>	>>
SM <sub>5</sub> –II	-119.3	$SM_4$	89	SM5-II origin	0
	"	$SM_5$	54	"	"
	-115.5	$SM_3$	47	vdW-1	4.0
	"	$SM_4$	210	>>	"
	"	$SM_5$	102	"	"
	-110.2	$SM_3$	64	vdW-2	9.1
	"	$SM_4$	284	"	"
	"	$SM_5$	93	"	"
	-108.3	$SM_4$	131	vdW	11.0
	-104.9	"	155	vdW-3	14.4
	" 102.1	SM <sub>5</sub>	78	)) 1777 - 4	" 1 ~ <b>0</b>
	-103.1	$SM_4$	143	vdW-4	16.2
	" ~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~	$SM_5$	78	<b>))</b>	" • 1
	-97.7	$SM_4$	148	vdW	21.6
SM <sub>6</sub> -I	-46.2	$SM_3$	77	(SM <sub>6</sub> –I) origin	0
0 -		$SM_4$	127	· · · -	
	>>	- т		"	"

	"	$SM_5$	181	22	"
	"	$SM_6$	100	"	"
	-30.8	$SM_3$	71	vdW-1	15.4
	"	$SM_4$	92	"	"
	"	$SM_5$	157	"	"
	"	$SM_6$	80	"	"
	-14.4	$SM_5$	120	vdW-2	31.8
	22	$SM_6$	77	"	"
	-1.9	$SM_5$	109	vdW-3	44.3
	"	$SM_6$	60	"	"
	9.7	$SM_5$	83	vdW-4	55.9
	"	$SM_6$	52	"	"
	21.3	$SM_5$	52	vdW-5	67.5
	"	$SM_6$	44	"	"
	~~			~~	
SM <sub>6</sub> -II	-112.6	$SM_5$	162	SM <sub>6</sub> –II origin	0
	"	$SM_6$	110	"	"
	-100.6	$SM_5$	74	vdW-1	12.0
	"	$SM_6$	114	"	,,
	-95.8	$SM_5$	107	vdW-2	16.8
	"	$SM_6$	104	"	"
SM <sub>7</sub> -I	54.2	$SM_6$	168	SM7–I origin	0
2111/1		$SM_7$	58	-	
	58.6	$SM_6$	41	vdŴ-1	," 4.4
		$SM_{0}$	35		
	63.9	$SM_6$	110	vdŴ-2	9.7
		$SM_{0}$	41		
	65.9	$SM_6$	79	vdŴ-3	11.7
	69.8		50	vdW-4	15.6
	73.1	"	38	vdW-5	18.9
	79.0	"	39	vdW-6	24.8
	19.0	>>	57		21.0
SM7–II	-89.1	$SM_5$	118	SM7–II origin	0
	"	$SM_6$	224	"	"
	>>	$SM_7$	78	"	>>
	-80.4	$SM_6$	81	vdW	8.7
	"	$SM_7$	67	"	"
	-78.5	>>	71	vdW-2	10.6
	-73.2	$SM_6$	117	vdW	15.9
	>>	$SM_7$	69	>>	>>
a) -			0.6/0.7		
$SM_8$	36.1/34.2	$SM_8$	86/95	SM8 origin	-

	-34.2	$SM_7$	161	22	0
	-28.4	, , , , , , , , , , , , , , , , , , , ,	96	VdW-1	7.7
	"	$SM_8$	100		"
	-23.1	$\widetilde{SM_7}$	117	VdŴ-2	13
	>>	$SM_8$	89	>>	"
	-13.5	$SM_7$	108	VdW-3	22.6
	"	$SM_8$	81	"	"
	-8.2	$SM_7$	68	VdW-4	27.9
	>>	$SM_8$	65	>>	"
	-2.84	$SM_7$	77	VdW-5	33.26
SM <sub>9</sub> -I	24.2/25.7	$SM_8$	66	SM9-I origin	-
	24.2	$SM_9$	58	>>	0
	29.5	"	84	VdW-1	5.3
	46.0	$SM_8$	78	VdW-2	20.3
	"	$SM_9$	59	"	"
SM <sub>9</sub> -II	-120.3	$SM_8$	91	SM9–II origin	0
	"	$SM_9$	30	"	"
	-111.2	$SM_8$	64	VdW-1	9.1
	"	$SM_9$	40	"	"
	-104.9	"	53	vdW-2	15.4
	-100.1	"	48	vdW-3	20.2
	-94.4	"	45	vdW-4	25.9
	-86.7	"	61	vdW-5	33.6
	-77.5	"	65	vdW-6	42.8
	-73.7	"	77	vdW-7	46.6
	-67.9	"	61	vdW-8	52.4
	-62.6	>>	77	vdW-9	57.7

Fragmentation channel	Parent, shift (cm <sup>-1</sup> )	%
SM→S	63.4	65
SM₂-I→SM	83.8	96
SM₂-II→SM	3.4	4
$SM_3$ -I $\rightarrow$ $SM_2$	181.6	82
SM <sub>3</sub> -II→SM <sub>2</sub>	16.5	2
SM₄→SM <sub>3</sub>	-94.4	20
SM₅-I→SM₄	-58.8	55
SM₅-II→SM₄	-119.3	67
SM <sub>6</sub> -I→SM <sub>5</sub>	-46.2	48
SM <sub>6</sub> -II→SM <sub>5</sub>	-113.1	63
$SM_7-I \rightarrow SM_6$	54.2	81
$SM_7-II \rightarrow SM_6$	-89.1	61
SM <sub>7</sub> -II→SM <sub>5</sub>	-89.1	26
$SM_8 \rightarrow SM_7$	-36.1	11
SM <sub>9</sub> -I→SM <sub>8</sub>	24.2	27
$SM_9$ -II $\rightarrow SM_8$	-120.3	74

Table 9: Fragmentations probabilities of the SM<sub>n</sub> cluster ions.

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# 4.4 R2PI of Styrene-Ethanol Clusters

# 4.4.1 Spectra of $SE_n$ Clusters, n = 1-2

The R2PI absorption spectra obtained by monitoring the styrene, S, and styrene- $(C_2H_5OH)_n$ , SE<sub>n</sub>, (n = 1 and 2) mass channels are represented in Figure 27. In each mass channel, the assignments have been carefully designated to the observed spectral features taking into account the possible fragmentations from higher mass channels. The R2PI spectrum obtained for SE mass channel exhibits very rich spectral features starting form 37 cm<sup>-1</sup> to 104 cm<sup>-1</sup>. Several peaks with small intensities on the red side of styrene's origin (-10 to -80 cm<sup>-1</sup>) were assigned to the S<sub>2</sub>E cluster. The two isomeric clusters of SE are distinguishable in the R2PI spectrum of SE mass channel. The first isomer (SE-I), with an origin at 50 cm<sup>-1</sup>, has very loaded vdW progression. This van der Waals progression of SE-I cluster is located at 58.1, 64.8, 72.1, 83.0, 88.5 and 93.0 cm<sup>-1</sup> (corresponding to the peaks marked with solid circles in SE mass channel in Figure 27). The R2PI mass spectrum that was obtained at a resonance wavelength corresponding to the assigned  $0_0^0$  origin of SE-I is displayed in Figure 28. The intensity gained by the  $SE_2$  mass peak is in the order of the background intensity that was measured in SE<sub>2</sub> mass channel. The second isomer (SE-II) possesses a low frequency origin at 11.2 cm<sup>-1</sup> with one vdW peak at 16.6 cm<sup>-1</sup>.

The assigned absorption features of SE<sub>2</sub>-I isomer are in close proximity and interfere with those assigned to SE-I isomer as shown in SE mass channel in Figure 27. The origin transition of SE<sub>2</sub>-I was blue shifted by 36.9 cm<sup>-1</sup> relative to the  $0_0^0$  origin of styrene molecule. The van der Waals structure built on the origin transition of the SE<sub>2</sub>-I is

assigned to peaks 1-9 and extends over 55 wavenumbers as illustrated in Figure 27. The presence of rich van der Waals progression indicates that the geometry of the  $SE_2$ -I neutral cluster is changed significantly upon the electronic excitation.

In SE<sub>2</sub> mass channel, very congested spectral features are observed over the scan range (-100 to 180 cm<sup>-1</sup>). Most of the major features could be assigned either by searching for fragmented clusters in the higher mass channels or by observing the dissociative proton transfer reaction within S<sub>2</sub>E<sub>n</sub> clusters as will be discussed later. A second isomer for SE<sub>2</sub> cluster (SE<sub>2</sub>-II) is detected with an origin at 16.9 cm<sup>-1</sup> and two vdW modes at 20.3 and 23.9 cm<sup>-1</sup>. In contrast to the first isomer, SE<sub>2</sub>-II isomer does not show fragmentation into SE mass channel. This means that the structure of SE<sub>2</sub>-II isomer does not seem to be affected upon photoionization.

Two sets of interfered spectral features belonging to two neutral clusters are recognized to the blue side in SE<sub>2</sub> mass channel. One set is assigned to S<sub>2</sub>E<sub>2</sub>-I cluster with peaks at 111.0, 138.0, 145.4, 151.5, 154.9 and 156.7 cm<sup>-1</sup> (marked with open circles in SE<sub>2</sub> mass channel in Figure 27). The other set was assigned to SE<sub>3</sub>-I cluster and composed of an origin at 108.9 cm<sup>-1</sup> and four-associated vdW modes at 114.0, 120.4, 128.6 and 133.5 cm<sup>-1</sup> (corresponding to the peaks marked with  $\blacktriangle$  in SE<sub>2</sub> mass channel in Figure 27). The remaining absorption features to the red side in the SE<sub>2</sub> mass channel was assigned to S<sub>2</sub>E<sub>2</sub> and S<sub>2</sub>E<sub>n</sub>, where n > 2.

# 4.4.2 Spectra of $SE_n$ Clusters, n = 3-5

The R2PI spectra obtained by monitoring the mass channels corresponding to BE<sub>3</sub>, BE<sub>4</sub> and BE<sub>5</sub> clusters in the region of the styrene  $S_1 \leftarrow S_0$  transition are illustrated in Figure 29. The R2PI spectrum of SE<sub>3</sub> mass channel exhibits envelope shape with several spectral features appearing in the base line. The envelope shape could be due to the fragmentation processes from higher styrene/ethanol clusters into SE<sub>3</sub> mass channel. However, several features could be identified as fragmentations from S<sub>2</sub>E<sub>3</sub> and SE<sub>4</sub>-I clusters as shown in SE<sub>3</sub> mass channel in Figure 29. A set of four peaks in the red side to the 0<sup>0</sup><sub>0</sub> of styrene molecule was assigned to a second isomer of SE<sub>3</sub> (SE<sub>3</sub>-II). The 0<sup>0</sup><sub>0</sub> origin of SE<sub>3</sub>-II cluster is located at -46.2 cm<sup>-1</sup> with the associated vdW modes at -39.6, -32.9 and -26.6 cm<sup>-1</sup>. The remaining features could not be assigned because of the interference problem arising from the fragmentation processes.

In SE<sub>4</sub> mass channel, three isomers of SE<sub>4</sub> cluster have been identified. The first isomer (SE<sub>4</sub>-I) has a weak blue-shifted origin at 107.4 cm<sup>-1</sup> followed by two vdW peaks at 120.1 and 123.8 cm<sup>-1</sup>. The origin peak together with one vdW peak appears in SE<sub>3</sub> mass channel and they are obscured by the background intensity measured in that channel. The second isomer (SE<sub>4</sub>-II) possesses a red-shifted origin at -65.2 cm<sup>-1</sup> and three vdW peaks at -62.5, -51.9 and -43.2cm<sup>-1</sup>. The third isomer (SE<sub>4</sub>-III) is the most red-shifted cluster with strong origin at -118.4 cm<sup>-1</sup> only one vdW mode at -99.8 cm<sup>-1</sup>. Figure 30 displays the R2PI mass spectrum that obtained at the resonance wavelength corresponding to the origin of SE<sub>4</sub>-III isomer.

The most predominant feature at  $-104.3 \text{ cm}^{-1}$  in the SE<sub>5</sub> mass channel is assigned to the origin of SE<sub>5</sub> cluster. This origin is followed by a vdW triplet at -100.7, -88.9 and -80.8 cm<sup>-1</sup>. The R2PI mass spectrum at the origin of SE<sub>5</sub> is displayed in Figure 31.

The overall R2PI spectra of  $SE_n$  clusters are shown in Figure 32. The relative frequencies and tentative assignments of the  $SE_n$ , where n = 1-5, cluster isomers are recorded in Table 10.

### 4.4.3 Intracluster Reactions following Photoionization

No ion-molecule chemistry is observed in  $SE_n$  cluster series. However, the dissociative proton transfer (DPT) reactions are noticed in  $S_2E_n$  clusters ( $n \ge 2$ ). Figure 33 shows that the spectral features of  $E_2H$  are in resonance with those of  $S_2E$ . This means that  $E_2H$  and  $S_2E$  are coming from the same reaction precursor, which is  $S_2E_2$ . These features also appear in  $S_2$  and  $SE_2$  mass channels. This suggests that  $S_2E_n$  clusters undergo DPT reactions (Eq. 4-3) in competition with the fragmentation reactions by either evaporating a styrene molecule (Eq. 4-4) or evaporating one or two ethanol molecules (Eq. 4-8 and Eq. 4-9, respectively). Figure 34 illustrates the observation of DPT reaction initiating from  $S_2E_3$  cluster.

$$C_{16}H_{16}^{+}(C_{2}H_{5}OH)_{n} \rightarrow H^{+}(C_{2}H_{5}OH)_{n} + C_{16}H_{15} \cdots n \ge 2$$
 Eq. 4-3

$$C_{16}H_{16}^+(C_2H_5OH)_n \rightarrow C_8H_8^+(C_2H_5OH)_n + C_8H_8$$
 Eq. 4-4

$$C_{16}H_{16}^{+}(C_{2}H_{5}OH)_{n} \rightarrow C_{16}H_{16}^{+}(C_{2}H_{5}OH)_{n-1} + C_{2}H_{5}OH$$
 Eq. 4-5

$$C_{16}H_{16}^+(C_2H_5OH)_n \rightarrow C_{16}H_{16}^+(C_2H_5OH)_{n-2} + 2C_2H_5OH$$
 Eq. 4-6

As mentioned before with respect to the styrene/methanol system (see section 4.3.6), the DPT reaction is only feasible if the proton affinity (PA) of the methanol trimer (or now ethanol dimer) is higher than that of the styrene dimer. The reported proton affinity value for the ethanol dimer is 220.4 kcal/mol.[120] This, in turn, lowers the estimated upper limit for the proton affinity of the styrene dimer to  $\leq$  220.4 kcal/mol.

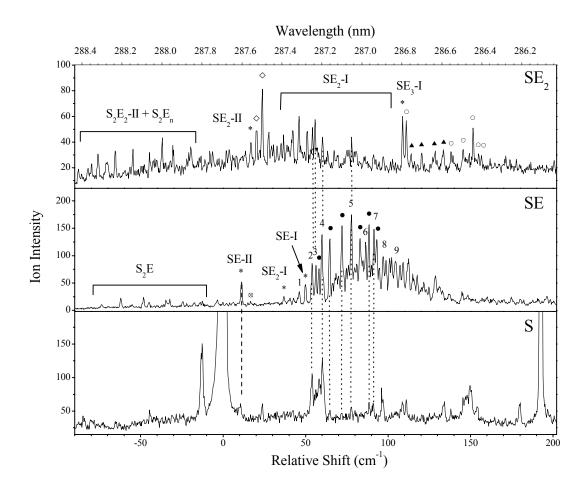


Figure 27: One-color R2PI spectra measured in styrene and styrene-(ethanol)<sub>n</sub> mass channels with n = 1-2, [S, SE and SE<sub>2</sub>] relative to the electronic origin band of the S<sub>1</sub> $\leftarrow$ S<sub>0</sub> transition of the styrene molecule at 34,758.79 cm<sup>-1</sup>. The origin of each cluster isomer is marked with a (\*). The peaks labeled with numbers, letters or symbols following each cluster's origin represent vdW bands associated with the cluster origin. The peaks labeled with ( $\odot$ ) in SE<sub>2</sub> mass channel are assigned to S<sub>2</sub>E<sub>2</sub>-I isomer cluster (see Table 10).

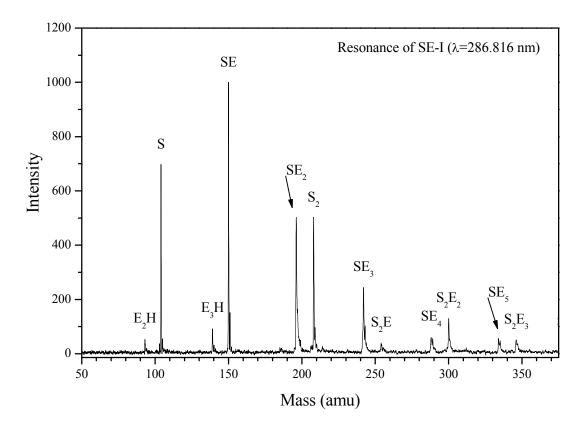


Figure 28: R2PI mass spectra of the styrene-ethanol  $(SE_n)$  cluster beam obtained at the resonance ionization assigned to the origin of SE-I cluster.

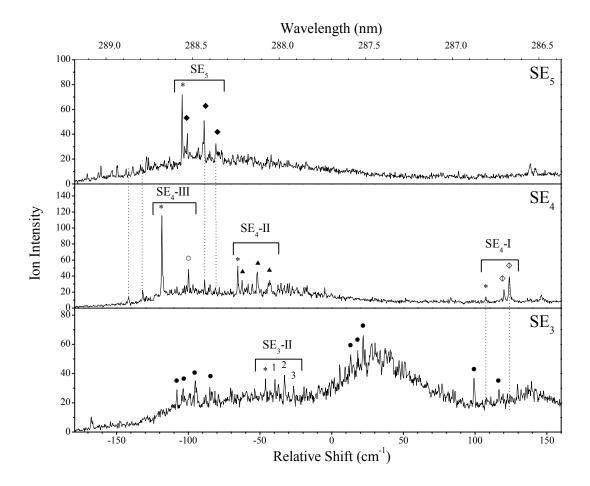


Figure 29: One-color R2PI spectra measured in the styrene-(ethanol)<sub>n</sub> mass channels with n = 3-5, [SE<sub>3</sub>, SE<sub>4</sub> and SE<sub>5</sub>] relative to the electronic origin band of the S<sub>1</sub> $\leftarrow$ S<sub>0</sub> transition of the styrene molecule at 34,758.79 cm<sup>-1</sup>. The origin of each cluster isomer is marked with a (\*). The peaks labeled with numbers, letters or symbols following each cluster's origin represent vdW bands associated with the cluster origin. The peaks labeled with (•) in SE<sub>3</sub> mass channel are assigned to S<sub>2</sub>E<sub>3</sub> cluster (see Table 10).

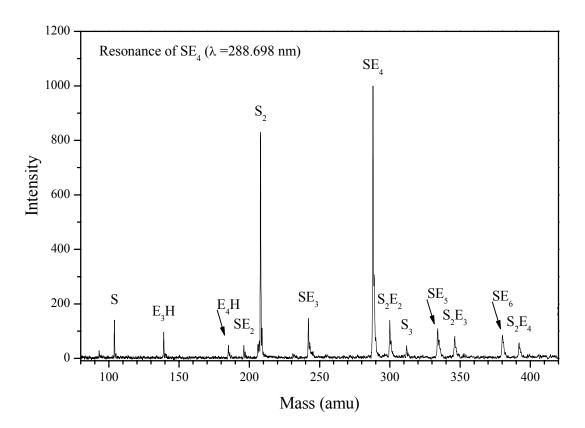


Figure 30: R2PI mass spectra of the styrene-ethanol (SE<sub>n</sub>) cluster beam obtained at the resonance ionization assigned to the origin of  $SE_4$  cluster.

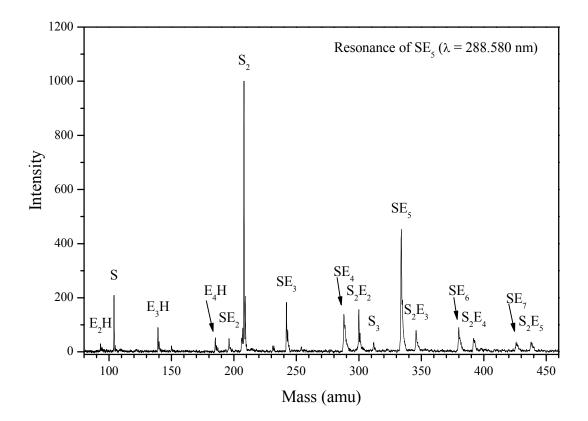


Figure 31: R2PI mass spectra of the styrene-ethanol  $(SE_n)$  cluster beam obtained at the resonance ionization assigned to the origin of  $SE_5$  cluster.

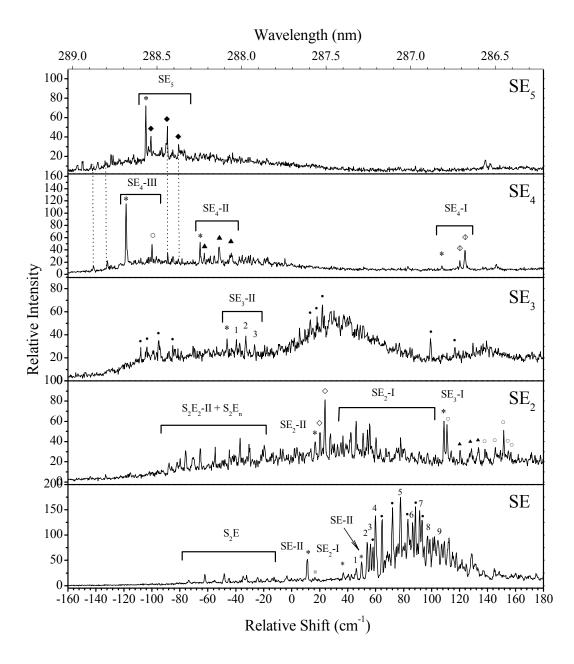


Figure 32: Overall R2PI spectra measured in the styrene-(ethanol)<sub>n</sub> mass channels (SE<sub>n</sub>), n = 1-5, relative to the electronic origin band of the  $S_1 \leftarrow S_0$  transition of the styrene molecule at 34,758.79cm<sup>-1</sup>. The origin of each cluster isomer is marked with a (\*). The peaks labeled with symbols following each cluster's origin represent vdW bands associated with the cluster origin (see Table 8). The peaks labeled with ( $\bullet$ ) and ( $\bigcirc$ ) in SE<sub>2</sub> and SE<sub>3</sub> mass channels are assigned to S<sub>2</sub>E<sub>2</sub> and S<sub>2</sub>E<sub>3</sub> clusters, respectively.

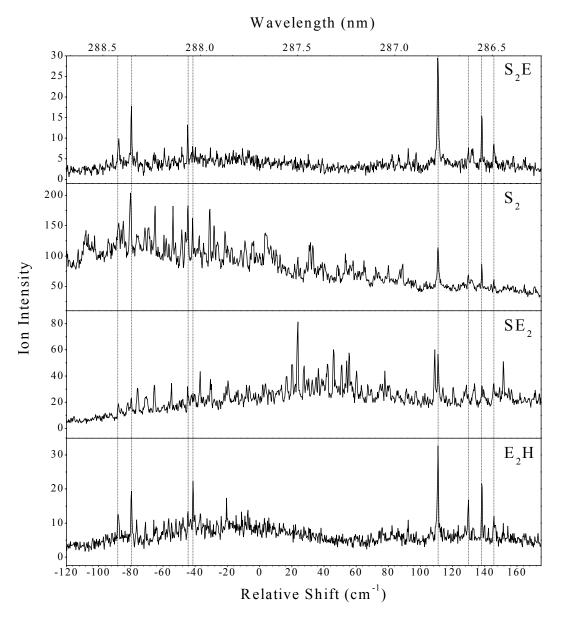


Figure 33: R2PI scans measured in  $E_2H$ ,  $SE_2$ ,  $S_2$  and  $S_2E$  mass channels. The hashed lines point out the spectral features in resonance with those in the proton transfer product mass channel ( $E_2H$ ). These features are assigned to  $S_2E_2$  isomer clusters.

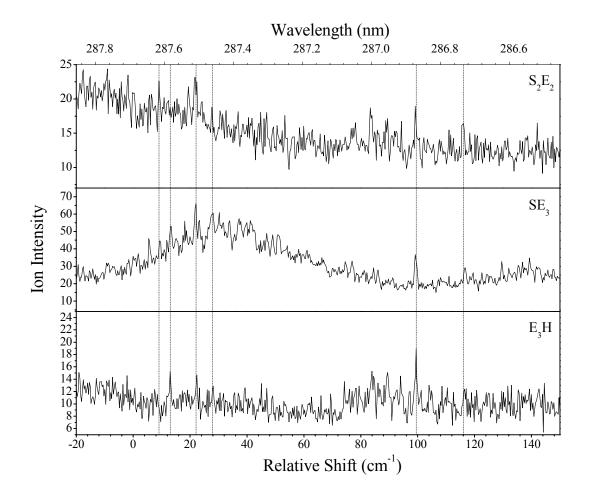


Figure 34: R2PI scans measured in  $E_3H$ ,  $SE_3$ , and  $S_2E_2$  mass channels. The hashed lines point out the spectral features in resonance with those in the proton transfer product mass channel ( $E_3H$ ). These features are assigned to  $S_2E_3$  cluster.

Cluster	Shift (cm <sup>-1</sup> ) from the $0_0^0$ of Styrene	Observed Mass Channel	Relative Intensity	Assignment	Shift (cm <sup>-1</sup> ) from Cluster's Origin
SE-I	50.0	SE	47	SE-I origin	0
	58.1	"	76	vdW-1	8.1
	64.8	"	131	vdW-2	14.8
	72.1	>>	155	vdW-3	22.1
	83.0	>>	131	vdW-4	33.0
	88.5	>>	157	vdW-5	38.5
	93.0	"	129	vdW-6	43.0
SE-II	11.2	SE	52	SE-II origin	0
	16.6	"	17	vdW	5.4
SE <sub>2</sub> -I	36.9	SE	25	SE <sub>2</sub> -I origin	0
	36.6	$SE_2$	47	"	0
	46.3	SE	34	vdW-1	9.4
	46.0	$SE_2$	61	"	9.4
	53.9	SE	86	vdW-2	17.0
	54.2	$SE_2$	51	"	17.3
	56.3	SE	82	vdW-3	19.4
	56.0	$SE_2$	55	"	19.4
	60.0	SE	138	vdW-4	23.7
	60.3	$SE_2$	44	"	23.7
	77.9	SE	175	vdŴ-5	41.0
	77.9	$SE_2$	44	"	41.3
	86.4	SE	125	vdW-6	49.5
	86.1	$SE_2$	29	"	49.5
	91.5	SE	148	vdW-7	54.6
	91.2	$SE_2$	30	22	54.6
	97.0	SE	98	vdW-8	60.1
	97.0	$SE_2$	28	>>	60.4
	104.6	SE	89	vdW-9	67.7
	104.6	$SE_2$	24	"	68.0
SE <sub>2</sub> -II	16.9	$SE_2$	40	SE <sub>2</sub> -II origin	0
	20.3	>>	49	vdW-1	3.4

Table 10: Spectral features observed in the  $SE_n$  mass channels (S = styrene, E = ethanol, vdW = van der Waals mode).

	23.9	22	81	vdW-2	7.0
SE <sub>3</sub> -I	108.9	$SE_2$	60	SE <sub>3</sub> -I origin	0
	114.0	>>	31	vdW-1	5.1
	120.4	>>	31	vdW-2	11.5
	128.6	"	33	vdW-3	19.7
	133.5	>>	34	vdW-4	24.6
SE <sub>3</sub> -II	-46.2	$SE_3$	37	SE <sub>3</sub> -II origin	0
	-39.6	>>	36	vdW-1	6.6
	-32.9	>>	39	vdW-2	13.3
	-26.6	>>	31	vdW-3	19.6
SE <sub>4</sub> -I	107.4	$SE_3$	22	SE <sub>4</sub> -I origin	0
	107.7	$SE_4$	13	"	0
	120.1	>>	23	vdW-1	12.4
	124.1	$SE_3$	25	vdW-2	16.7
	123.8	$SE_4$	39	>>	16.1
SE <sub>4</sub> -II	-65.2	$SE_4$	52	SE <sub>4</sub> -II origin	0
	-62.5	>>	35	vdW-1	2.7
	-51.9	>>	45	vdW-2	13.3
	-43.2	"	35	vdW-3	22.0
SE <sub>4</sub> -III	-118.4	$SE_4$	115	SE <sub>4</sub> -I origin	0
	-99.8	>>	49	vdW-1	18.6
$SE_5$	-104.3	$SE_5$	72	SE5 origin	0
	-101.0	$SE_4$	28	vdW-1	3.3
	-100.7	$SE_5$	41	>>	3.6
	-88.6	$SE_4$	35	vdW-2	15.1
	-88.9	$SE_5$	51	>>	15.4
	-81.1	$SE_4$	26	vdW-3	23.2
	-80.8	$SE_5$	32	>>	23.5

# 4.5 R2PI of Styrene-Trifluoroethanol Clusters

# 4.5.1 Spectra of $ST_n$ Clusters, n = 1-4

The binary mixed clusters of styrene and trifluoroethanol were investigated by using the R2PI technique. The R2PI spectra of the styrene-(trifluoroethanol)<sub>n</sub> clusters,  $C_8H_8$ -(CF<sub>3</sub>CH<sub>2</sub>OH)<sub>n</sub>, were measured utilizing the  $\tilde{A}^1A' \leftarrow \tilde{X}^1A'$  (0<sup>0</sup><sub>0</sub>) transition of the styrene molecule, with n=1-4. The observed R2PI spectra of styrene-(trifluoroethanol)<sub>n</sub>, ST<sub>n</sub>, mass channels in Figure 35 reveal these spectral characterizations (1) the absence of the R2PI spectrum of ST ion cluster, (2) no evidence of cluster fragmentations was found in that system and (3) no intracluster ion products were obtained for that system under the given experimental conditions.

In liquid phase, reactions of styrene radical cation obtained by the laser excitation with 266 and 308 nm in neat trifluoroethanol alcohol have been observed.[121] Accordingly, one possible explanation for the absence of R2PI spectrum of  $ST^+$  is that, upon photoexcitation of the ST cluster, radical reaction products are formed which are incapable of R2PI detection.

The observed spectral features in  $ST_2$ ,  $ST_3$  and  $ST_4$  mass channels were assigned to 1:2, 1:3 and 1:4 neutral clusters, respectively. The most-red peaks in each mass channel could be assigned to the origins of the corresponding neutral cluster. The origins of 1:2, 1:3 and 1:4 neutral clusters are located at 53.92, 130.82 and 136.30 cm<sup>-1</sup>, respectively. The small intensity peaks to the red side of 1:4 cluster's origin could be due another isomer. In each mass channel, several closely spaced features were observed following the cluster's origin. The relative frequencies and the assignments of the  $ST_n$  clusters, with n =2-4, are tabulated in Table 11.

Shown in Figure 36 through Figure 38 are representative mass spectra measured at the origin of  $ST_n$  clusters. The mass spectra clearly show the mass peaks corresponding to (1) styrene clusters  $S_n$ , with n=1-3, (2) styrene-trifluoroethanol clusters  $ST_n$ , with n=1-5, and (3) styrene dimer-trifluoroethanol clusters  $S_2T_n$ , with n=1-3.

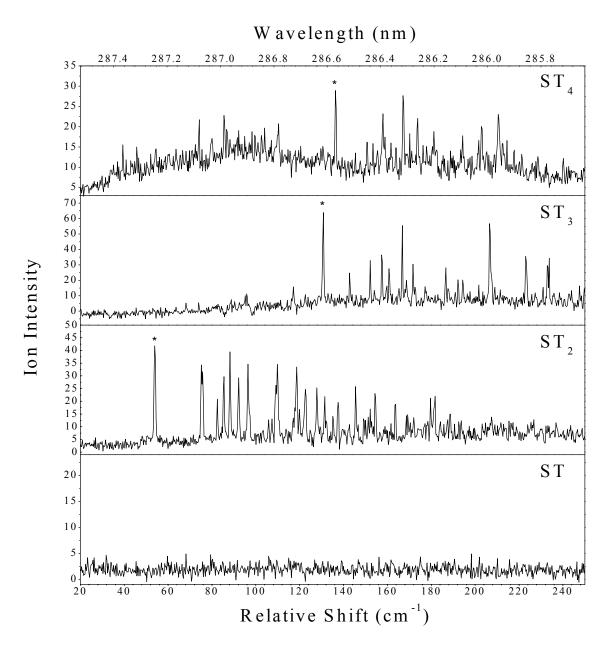


Figure 35: Overall R2PI spectra measured in the styrene-(trifluoroethanol)<sub>n</sub> mass channels  $(ST_n)$ , n = 1-4, relative to the electronic origin band of the S<sub>1</sub> $\leftarrow$ S<sub>0</sub> transition of the styrene molecule at 34,758.79 cm<sup>-1</sup>. The origin of each cluster is marked with a (\*).

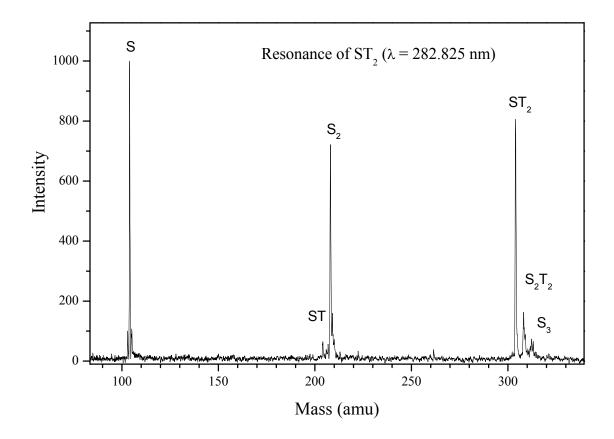


Figure 36: R2PI mass spectra of the styrene-trifluoroethanol  $(ST_n)$  cluster beam obtained at the resonance feature assigned to the origin of  $ST_2$  cluster.

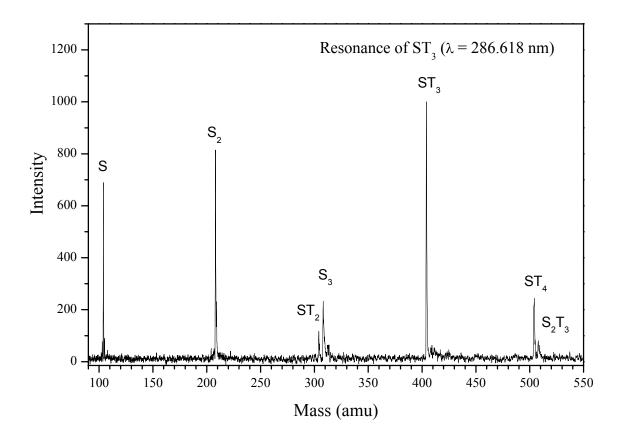


Figure 37: R2PI mass spectra of the styrene-trifluoroethanol  $(ST_n)$  cluster beam obtained at the resonance feature assigned to the origin of  $ST_3$  cluster.

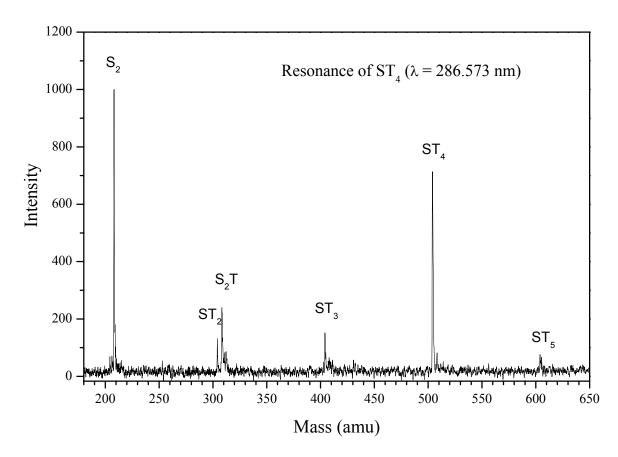


Figure 38: R2PI mass spectra of the styrene-trifluoroethanol  $(ST_n)$  cluster beam obtained at the resonance feature assigned to the origin of  $ST_4$  cluster.

Mass channel	Shift (cm <sup>-1</sup> ) from the $0_0^0$	Relative intensity	Assignments
	of styrene		
$ST_2$	82.5	21	1:2 origin
	85.5	30	1:2
	88.2	39	1:2
	92.2	29	1:2
	96.4	35	1:2
	107.4	13	1:2
	109.8	35	1:2
	118.6	33	1:2
	119.9	17	1:2
	122.6	25	1:2
	127.8	25	1:2
	131.4	22	1:2
	135.1	14	1:2
	137.5	20	1:2
	145.4	26	1:2
	149.1	13	1:2
	154.3	23	1:2
	163.7	19	1:2
	169.2	15	1:2
$ST_3$	130.8	64	1:3 origin
	142.7	25	1:3
	152.1	33	1:3
	157.3	36	1:3
	160.7	28	1:3
	166.8	55	1:3
	171.7	30	1:3
	177.2	16	1:3
	186.6	28	1:3
	192.2	20	1:3
	206.5	57	1:3
	223.1	36	1:3
	233.8	34	1:3
	247.6	16	1:3
$ST_4$	136.3	29	1:4 origin
	145.4	14	1:4
	150.6	16	1:4

Table 11: Spectral features observed in the  $ST_n$  mass channels (S = styrene, T = trifluoroethanol)

101	1	0	1
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157.9	23	1:4
167.1	28	1:4
173.8	22	1:4
181.2	19	1:4
188.8	14	1:4
194.3	18	1:4
201.6	17	1:4
202.9	20	1:4
210.5	23	1:4

#### 4.6 R2PI of Benzene-Ethanol and Benzene Deuterated Ethanol Clusters

The binary mixed clusters of benzene and ethanol were studied by using R2PI technique. The R2PI spectra for benzene-(ethanol)<sub>n</sub>,  $C_6H_6$ -( $C_2H_5OH$ )<sub>n</sub>, were obtained via the  ${}^{1}B_{2u} \leftarrow {}^{1}A_{1g}$  ( $6_0^{1}$ ) transition of the benzene molecule, where n=1-5 (Figure 39). Ion-molecule chemistry yields were observed following the photoionization of benzene-(ethanol)<sub>n</sub> clusters. These are ionized ethanol clusters ( $C_2H_5OH$ )<sub>n</sub>, where n  $\ge 2$ , the protonated ethanol clusters ( $C_2H_5OH$ )<sub>n</sub> $H^+$ , where n  $\ge 3$ , and the ionized products having the formula of ( $C_2H_5OH$ )<sub>n</sub> $-C_2H_5O$ , where n = 1 and 2. Same experiment was conducted on benzene-deuterated ethanol clusters,  $C_6H_6$ -( $C_2H_5OD$ )<sub>n</sub>, to further explore the ion-molecule reactions. The R2PI scan obtained in benzene-(deuterated ethanol)<sub>n</sub>,  $C_6H_6$ -( $C_2H_5OD$ )<sub>n</sub>, where n = 1-6 is displayed in Figure 40. Also, well-resolved R2PI scans were obtained for ( $C_6H_6$ )<sub>2</sub>-( $C_2H_5OD$ )<sub>n</sub> and ( $C_6H_6$ )  $_3$ -( $C_2H_5OD$ )<sub>n</sub> as shown in Figure 41 and Figure 42, respectively.

#### 4.6.1 Intracluster Reactions following Photoionization

In one-color R2PI, the benzene molecule attains about 2730 cm<sup>-1</sup> excess energy above the ionization threshold via its  $6_0^1$  transition. Following the photoionization of the neutral cluster, the ejected electron acquires most of this excess energy as a kinetic energy (Eq. 4-7). However, a considerable part of the excess energy is deposited in the cluster ion and thus increasing its internal energy. This, in turn, causes fragmentation of the cluster ion (Eq. 4-8) and initiates the molecule-ion chemistry. The precursors of the intracluster reactions are the neutral clusters of benzene-(ethanol)<sub>n</sub>, C<sub>6</sub>H<sub>6</sub>-(C<sub>2</sub>H<sub>5</sub>OH)<sub>n</sub>. The observed intracluster reactions are dissociative electron transfer reaction DET (Eq. 4-9), dissociative proton transfer reaction DPT (Eq. 4-10) and hydrogen loss reaction. These reactions are in competition with the cluster fragmentation channel as shown in the following scheme:

R2PI:  

$$C_6H_6 - (C_2H_5OH)_n + 2h\gamma_1 \rightarrow [C_6H_6 - (C_2H_5OH)_n]^{+*} + e^{-}$$
 Eq. 4-7

Fragmentation:  

$$[C_6H_6 - (C_2H_5OH)_n]^{+*} \rightarrow [C_6H_6 - (C_2H_5OH)_{n-1}]^{+} + C_2H_5OH$$
 Eq. 4-8  
DET<sup>-</sup>

$$\left[C_{6}H_{6} - (C_{2}H_{5}OH)_{n}\right]^{+*} \rightarrow \left[(C_{2}H_{5}OH)_{n}\right]^{+} + C_{6}H_{6}$$
 Eq. 4-9

DPT:  

$$[C_6H_6 - (C_2H_5OH)_n]^{+*} \rightarrow [H - (C_2H_5OH)_n]^{+} + C_6H_5$$
 Eq. 4-10

The most efficient pathways for the cluster ions upon the ionization are the fragmentation channel and then the dissociative electron transfer channel. The dissociative electron transfer and dissociative proton transfer channel were observed for  $n \ge 2$  and  $n \ge 3$ , respectively. The exothermicity of the proton transfer channel was tested. Unfortunately, not enough thermochemistry data was obtained to verify the energetic of the electron transfer channel. Assuming the following reaction, one can get the heat of reaction by knowing the heat of formations for  $C_6H_6^+$ ,  $C_2H_5OH$ ,  $C_6H_5$  and  $(C_2H_5OH)_nH^+$ .

$$C_6H_6^+ + nC_2H_5OH \rightarrow (C_2H_5OH)_nH^+ + C_6H_5$$
 Eq. 4-11

For protonated ethanol clusters  $(C_2H_5OH)_nH^+$ , only the heat of formations for n= 1-3 could be obtained. Since the heat of formations of protonated ethanol and protonated methanol clusters are very similar for n=1-3 (Table 12). Therefore, the heat of formation of

protonated methanol for n = 4 and 5, in very good approximation, was used to estimate the exothermicity of the DPT reactions in the analogous protonated ethanol clusters. Summarized in Table 13 is the thermochemical data that was used to calculate the heat of reaction in Eq. 4-11. The dissociative proton transfer DPT channel becomes apparently exothermic at protonated ethanol trimer by 30.3 kcal/mol (see Table 14). Although, the dissociative proton transfer DPT reaction is exothermic at protonated ethanol dimer ( $E_2H^+$ ) by 8.9 kcal/mol, but it may not be enough to dissociate the intermediate complex [ $C_6H_5$ -( $C_2H_5OH_2H^+$ ]. The formation of  $E_nH^+$  becomes only possible if [ $PA(C_6H_5) - PA(E_n) + D(C_6H_5-E_nH^+) < 0$ ], where PA and D are the proton affinity and the dissociation energy, respectively.[122] It seems that this condition is not fulfilled in the case of  $E_2H$  cluster ion but the dissociation energy was estimated for the intermediate complex [ $C_6H_5$ -( $C_2H_5OH_2H^+$ ] to be  $\ge 9.4$  kcal/mol.

# 4.6.2 Spectra of Benzene Deuterated Ethanol (BD<sub>n</sub>) Clusters, n = 1-5

Similar R2PI spectra were obtained for benzene/ethanol and benzene/deuterated ethanol clusters, therefore, only the results from the benzene/deuterated ethanol system were considered to avoid repetitive work.

The assignments for  $BD_n$  clusters, where 1:n is given for the number of benzene-todeuterated ethanol in the neutral cluster responsible for the spectral features, are shown in Figure 40. The observation of intracluster reaction yields led to a clearly identifiable assignment of the neutral precursor cluster responsible for a particular spectral feature. The only exception is the 1:1 assignments in BD mass channel because there is no observable ion-molecule chemistry products for n = 1. However, the fingerprints in the BD mass channel spectrum identical to those in D<sub>2</sub> (m/z = 94) and D<sub>2</sub>-1 (m/z = 93) are assigned to 1:2 neutral cluster and the remaining features are assigned to 1:1 neutral cluster (see Figure 43). The mass spectrum obtained with the laser wavelength adjusted at 1:2 band shows masses 92, 93 and 94, corresponding to D<sub>2</sub>-2, D<sub>2</sub>-1 and D<sub>2</sub>, respectively, arising from hydrogen loss and electron transfer reactions (Figure 44b). These masses confirm that the neutral precursor cluster must contain two C<sub>2</sub>H<sub>5</sub>OD molecules. On the other hand, these masses disappear when the laser is tuned to the feature of 1:1 as shown in Figure 44a.

Spectral characterizations of D<sub>3</sub> mass channel reveal that 1:3 neutral cluster possesses two isomers (Figure 45). Several bands toward the blue side in D<sub>3</sub> mass channel (starting from 123.04 to 161.78 cm<sup>-1</sup>) are due to the first isomer 1:3 (I). While, the two most-red bands at -73.72 and -75.82 cm<sup>-1</sup> are due to the second isomer 1:3 (II). The first isomer 1:3 (I) exhibits, unlike the red shifted isomer 1:3(II), efficient fragmentation into BD<sub>2</sub> mass channel and intense ion-molecule chemistry. The mass spectra obtained at the resonance features of the two isomers support the assignments. The mass spectrum obtained at the resonance feature of 1:3 (I) cluster is displayed in Figure 46a. The mass spectrum shows enhanced ion intensities for D<sub>3</sub>-1 and D<sub>3</sub> masses. While, the mass spectrum recorded at the resonance feature of 1:3 (II) shows small ion intensity for D<sub>3</sub>-1 mass and the BD<sub>2</sub> mass peak is almost disappear (Figure 46b).

As with the 1:2 and 1:3 clusters, the  $D_4$  ion cluster is assigned to the electron transfer reaction taking place in 1:4 neutral precursor cluster. The fingerprints originating at 3.95 cm<sup>-1</sup> in the  $D_4$  mass channel spectrum, which are identical to those in BD<sub>3</sub> and BD<sub>4</sub>

mass channels, are assigned to 1:4 cluster (as shown in Figure 47). The 1:4 cluster is showing a little fragmentation into  $BD_3$  mass channel. The mass spectrum in Figure 48 obtained at the resonance feature of  $BD_4$  shows enhanced intensity for  $BD_4$  and  $D_4$  mass peaks and small ion intensity for  $BD_3$  peak.

 $D_5$  ion cluster is the electron transfer product from the BD<sub>5</sub> neutral cluster. It has spectral features originating at -44.97 cm<sup>-1</sup> and matching some features in BD<sub>4</sub> mass channel spectrum as clearly seen in Figure 49. These features are assigned to 1:5 cluster. The mass spectrum obtained at the resonance feature of 1:5 cluster is displayed in Figure 50. The relative frequencies and careful assignments of the BD<sub>n</sub> clusters are recorded in Table 15.

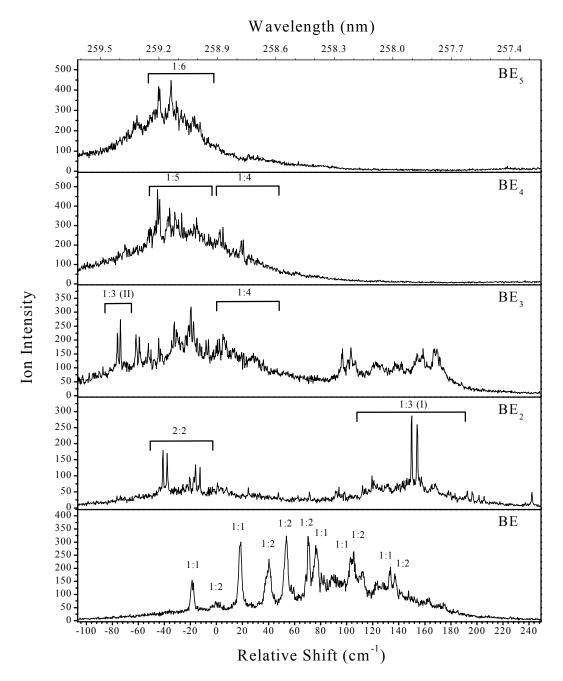


Figure 39 Overall R2PI spectra measured in the benzene  $(ethanol)_n$  mass channels  $(BE_n)$ , n = 1-5, relative to the  $6_0^1$  transition of the benzene molecule at 38,624.95 cm<sup>-1</sup>. The peak labels (m:n) stand for the number of benzene:ethanol in the neutral cluster responsible for the spectral features.

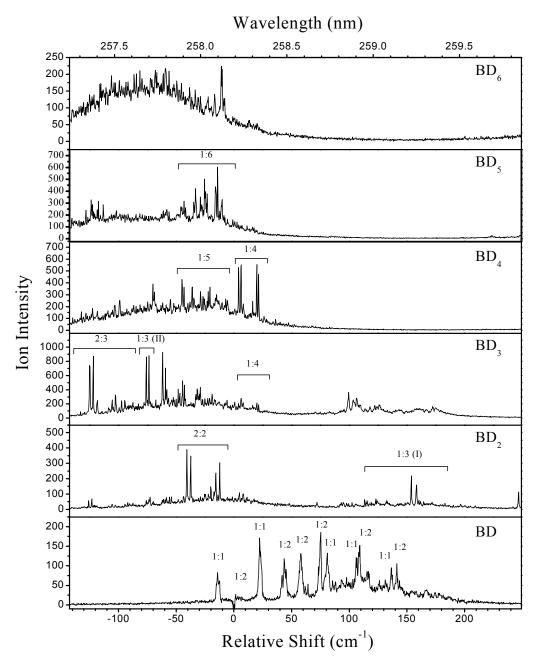


Figure 40: Overall R2PI spectra measured in the benzene (deuterated ethanol)<sub>n</sub> mass channels (BD<sub>n</sub>), n = 1-6, relative to the  $6_0^1$  transition of the benzene molecule at 38,624.95 cm<sup>-1</sup>. The peak labels (m:n) stand for the number of benzene:deuterated ethanol in the neutral cluster responsible for the spectral features (see Table 15).

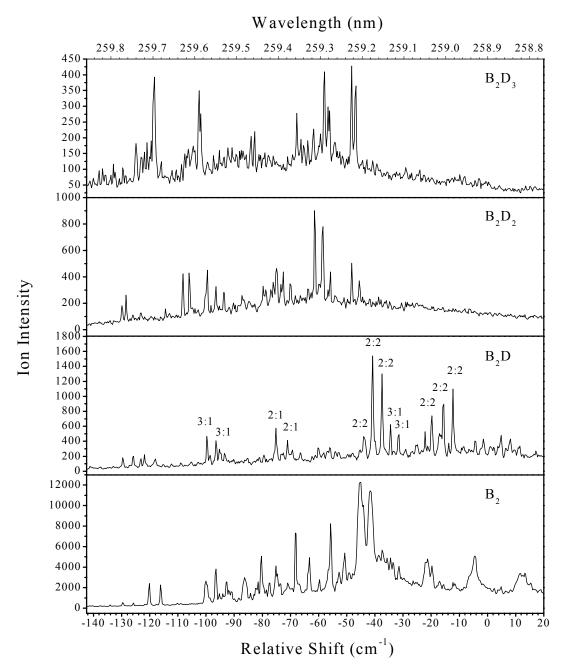


Figure 41: One-color R2PI spectra measured in  $(benzene)_2$ , B<sub>2</sub>, and  $(benzene)_2$  -  $(deuterated ethanol)_n$  mass channels  $(B_2D_n)$ , n = 1-3, relative to the  $6_0^1$  transition of the benzene molecule at 38,624.95 cm<sup>-1</sup>.

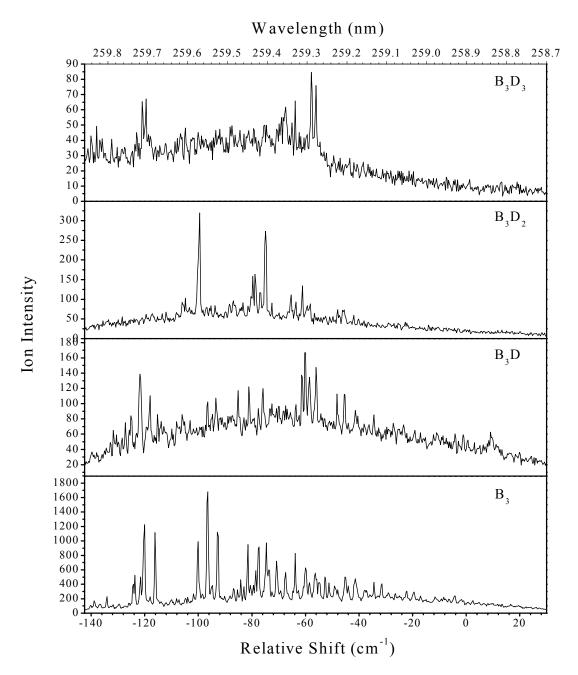


Figure 42: One-color R2PI spectra measured in (benzene)<sub>3</sub>, B<sub>3</sub>, and (benzene)<sub>3</sub>-(deuterated ethanol)<sub>n</sub> mass channels (B<sub>3</sub>D<sub>n</sub>), n = 1-3, relative to the  $6_0^1$  transition of the benzene molecule at 38,624.95 cm<sup>-1</sup>.

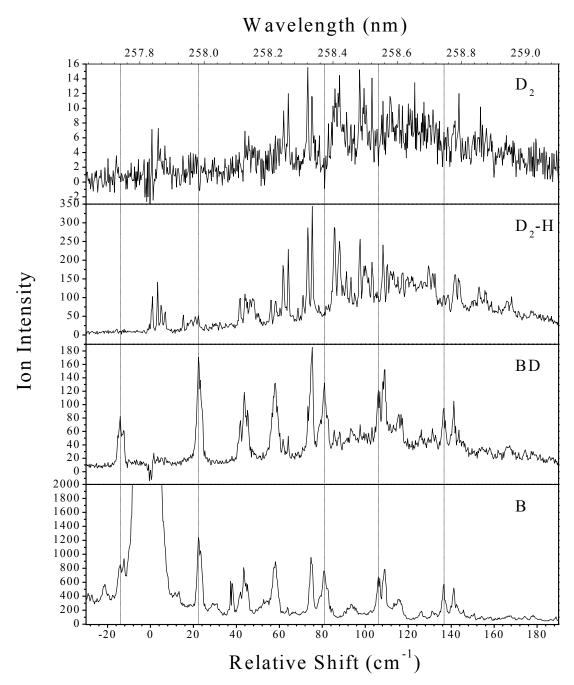


Figure 43: One-color R2PI spectra measured in benzene, B, deuterated ethanol, BD,  $(C_2H_5OD)(C_2H_4OD)$ , D<sub>2</sub>-H, and  $(C_2H_5OD)_2$ , D<sub>2</sub>, mass channels relative to the  $6_0^1$  transition of the benzene molecule at 38,624.95 cm<sup>-1</sup>. The dashed lines indicate the peaks that assigned to 1:1 cluster. Notice the disappearance of these peaks from the D<sub>2</sub> and D<sub>2</sub>-H mass channels (see the text).

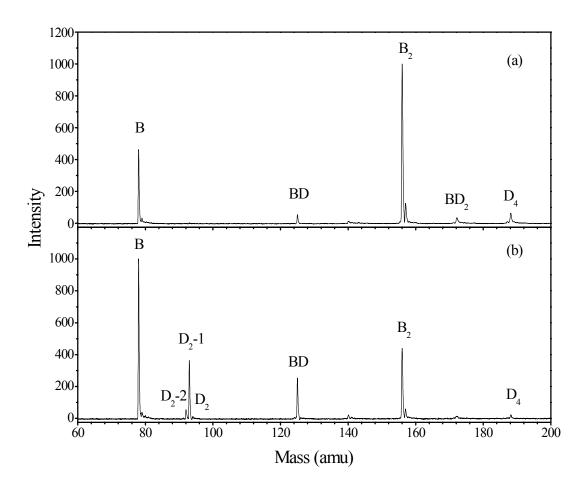


Figure 44: R2PI mass spectra of benzene-deuterated ethanol (BD) cluster beam obtained at the resonance wavelengths assigned to 1:1 and 1:2 neutral clusters at (a) 258.994 nm and (b) at 258.396 nm, respectively.

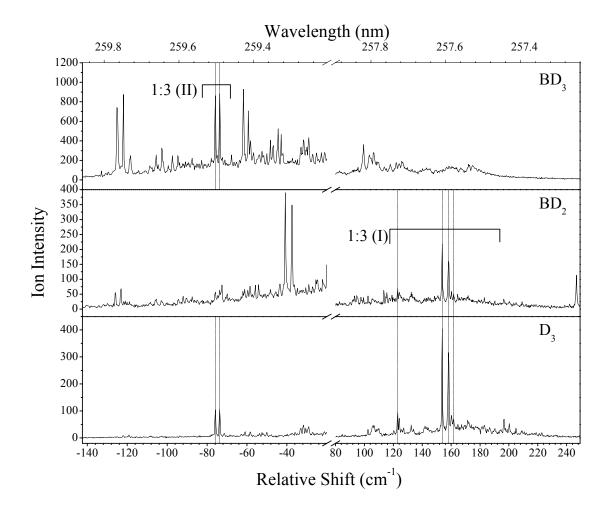


Figure 45: One-color R2PI spectra measured in  $(C_2H_5OD)_3$ , D<sub>3</sub>, benzene-(deuterated ethanol)<sub>2</sub>, BD<sub>2</sub>, benzene-(deuterated ethanol)<sub>3</sub>, BD<sub>3</sub>, and mass channels relative to the  $6_0^1$  transition of the benzene molecule at 38,624.95 cm<sup>-1</sup>. The dashed lines indicate the peaks that assigned to the blue shifted 1:3 (I) and the red shifted 1:3 (II) clusters in BD<sub>2</sub> and BD<sub>3</sub> mass channels, respectively.

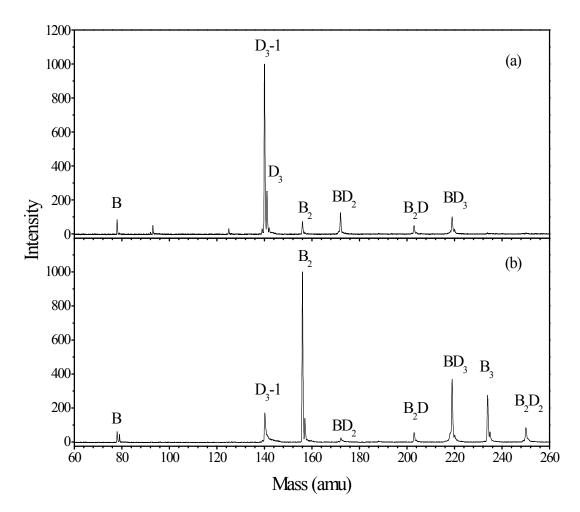


Figure 46: R2PI mass spectra of benzene-deuterated ethanol (BD) cluster beam obtained at the resonance wavelengths assigned to the 1:3 (I) and 1:3 (II) neutral clusters at (a) 257.842 nm and (b) at 259.395 nm, respectively.

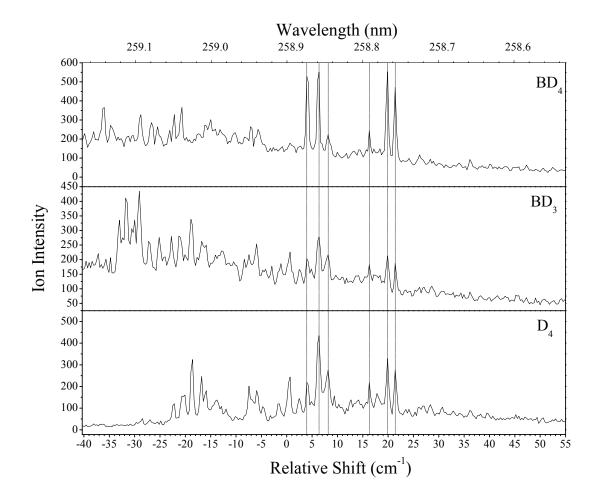


Figure 47: One-color R2PI spectra measured in  $(C_2H_5OD)_4$ , D<sub>4</sub>, benzene-(deuterated ethanol)<sub>3</sub>, BD<sub>3</sub>, benzene-(deuterated ethanol)<sub>4</sub>, BD<sub>4</sub>, and mass channels relative to the  $6_0^1$  transition of the benzene molecule at 38,624.95 cm<sup>-1</sup>. The dashed lines indicate the peaks that assigned to the spectral features of 1:4 neutral cluster in BD<sub>3</sub> and BD<sub>4</sub> mass channels.

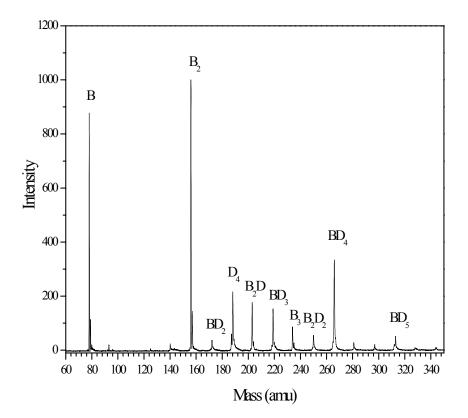


Figure 48: R2PI mass spectra of benzene-deuterated ethanol (BD) cluster beam obtained at the resonance wavelength assigned to the 1:4 neutral cluster at 258.857 nm.

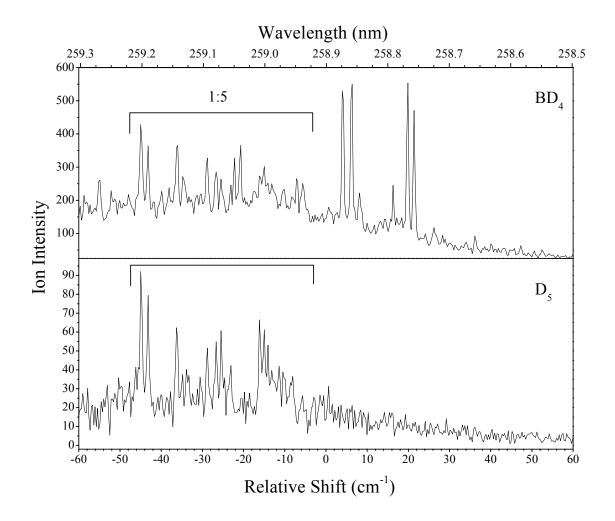


Figure 49: One-color R2PI spectra measured in  $(C_2H_5OD)_5$ , D<sub>5</sub>, and benzene-(deuterated ethanol)<sub>4</sub>, BD<sub>4</sub>, mass channels relative to the  $6_0^1$  transition of the benzene molecule at 38,624.95 cm<sup>-1</sup>. The same spectral features that are in resonance in both channels assigned to 1:5 cluster.

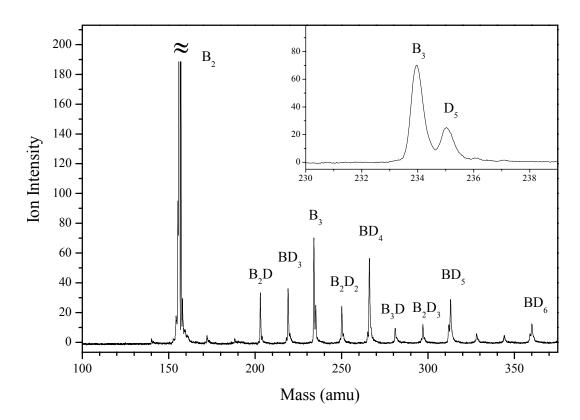


Figure 50: R2PI mass spectra of benzene-deuterated ethanol (BD) cluster beam obtained at the resonance wavelength assigned to the 1:5 neutral cluster at 259.202 nm.

(KOII) <sub>n</sub> II		
(n-1), n	Heat of reaction, <sup>a</sup> kcal/mol	Heat of reaction, <sup>b</sup> kcal/mol
	R=CH <sub>3</sub>	$R=C_2H_5$
0, 1	-180.3 (-182)	-188.3
1, 2	-32.6	-32.1
2, 3	-21.2	-21.4
3, 4	-14.0	
4, 5	-11.3	
<sup>a</sup> Ref [82, 123]		
<sup>b</sup> Ref $\begin{bmatrix} 120 \end{bmatrix}$		

Table 12: Thermochemistry of the clustering reactions:  $(ROH)_{n-1} + H^+ + ROH \rightarrow (ROH)_n H^+$ 

<sup>b</sup> Ref [120]

Table 13: Thermochemical data relevant to the BE<sub>n</sub> clusters (B =  $C_6H_6$ , E =  $C_2H_5OH$ )

Species	Heat of formation	Proton affinity	
	kcal/mol	kcal/mol	
$\mathrm{H}^{+}$	356.7		
$B^+$	233.2		
$C_6H_5$	79	211	
Е	-56.2	188.3	
$E_2H^+$	32.9 <sup>a</sup>	220.4	
$E_3H^+$	$-44.7^{a}$	241.8	
$E_4H^+$	-114.9 <sup>a</sup>		
$E_5H^+$	-182.4 <sup>a</sup>		

<sup>a</sup>Calculated from the thermochemical data in Table 12. <sup>b</sup> Ref [120].

Table 14: Exothermicity of the proton transfer reactions (PT)

finieity of the proton transfer reactions (1.1)		
n	$\Delta H^0$	
	kcal/mol	
1	87.6	
2	-8.9	
3	-30.3	
4	-44.3	
5	-55.6	

Mass channel	Shift (cm <sup>-1</sup> ) from the $6_0^1$	Relative intensity	Assignments
	of benzene		
BD	141.20	105.27	1:2
	136.66	94.80	1:1
	108.22	134.36	1:2
	106.11	121.33	1:1
	81.06	132.39	1:1
	75.32	185.42	1:2
	58.14	131.85	1:2
	43.67	118.51	1:2
	22.30	170.84	1:1
	3.36	22.52	1:2
	-14.06	82.60	1:1
BD <sub>2</sub>	161.78	51.27	1:3 (I)
	160.28	59.39	1:3 (I)
	158.16	159.78	1:3 (I)
	153.92	218.29	1:3 (I)
	124.25	56.34	1:3 (I)
	123.04	68.92	1:3 (I)
	-12.26	304.62	2:2
	-15.57	236.80	2:2
	-20.07	148.16	2:2
	-22.17	99.21	2:2
	-37.47	348.47	2:2
	-40.77	389.97	2:2
	-73.72	63.90	1:3 (II)
	-75.82	57.06	1:3 (II)
$BD_3$	21.39	187.81	1:4
	19.90	212.86	1:4
	16.28	182.98	1:4
	8.16	216.95	1:4
	6.36	277.58	1:4
	3.95	203.55	1:4
	-73.72	883.01	1:3 (II)
	-75.82	860.96	1:3 (II)
	-77.91	221.81	2:3
	-82.71	201.65	2:3
	-87.49	222.30	2:3

Table 15: Spectral features observed in the  $BD_n$  mass channels (B = benzene, D = deuterated ethanol, n = 1-5).

$\begin{array}{c ccccccccccccccccccccccccccccccccccc$				
$BD_4 = \begin{array}{ccccccccccccccccccccccccccccccccccc$		-89.29	177.58	2:3
$BD_4 = \begin{bmatrix} -102.74 & 327.15 & 2.3 \\ -105.43 & 256.09 & 2.3 \\ -118.29 & 249.82 & 2.3 \\ -121.86 & 876.17 & 2.3 \\ -125.15 & 741.25 & 2.3 \\ BD_4 = \begin{bmatrix} 21.39 & 470.97 & 1.4 \\ 19.90 & 553.93 & 1.4 \\ 16.28 & 245.23 & 1.4 \\ 8.16 & 222.28 & 1.4 \\ 6.36 & 550.34 & 1.4 \\ 3.95 & 529.97 & 1.4 \\ -5.65 & 250.97 & 1.5 \\ -7.16 & 265.03 & 1.5 \\ -20.67 & 366.00 & 1.5 \\ -20.67 & 366.00 & 1.5 \\ -20.67 & 366.00 & 1.5 \\ -26.67 & 285.52 & 1.5 \\ -28.77 & 327.53 & 1.5 \\ -34.77 & 271.97 & 1.5 \\ -35.98 & 365.34 & 1.5 \\ -44.97 & 428.52 & 1.5 \\ BD_5 & -10.16 & 329.85 & 1.6 \\ -29.07 & 351.10 & 1.6 \\ -41.97 & 260.77 & 1.6 \\ -43.17 & 317.29 & 1.6 \\ -56.06 & 227.09 & 1.6 \\ \end{bmatrix}$		-94.67	249.66	2:3
$BD_4 = \begin{array}{ccccccccccccccccccccccccccccccccccc$		-97.37	247.54	2:3
$BD_4 = \begin{array}{ccccccccccccccccccccccccccccccccccc$		-102.74	327.15	2:3
$BD_4 = \begin{array}{ccccccccccccccccccccccccccccccccccc$		-105.43	256.09	2:3
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$		-118.29	249.82	2:3
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$		-121.86	876.17	2:3
$BD_5 = \begin{array}{ccccccccccccccccccccccccccccccccccc$		-125.15	741.25	2:3
$BD_5 = \begin{array}{ccccccccccccccccccccccccccccccccccc$	BD₄	21 39	470 97	1.4
$BD_5 \qquad \begin{array}{ccccccccccccccccccccccccccccccccccc$	т			
$BD_5 = \begin{array}{ccccccccccccccccccccccccccccccccccc$				
$BD_5 \qquad \begin{array}{ccccccccccccccccccccccccccccccccccc$				
$BD_5 = \begin{array}{ccccccccccccccccccccccccccccccccccc$				
$BD_5 \qquad \begin{array}{ccccccccccccccccccccccccccccccccccc$				
$\begin{array}{cccccccccccccccccccccccccccccccccccc$		-5.65	250.97	1:5
$\begin{array}{cccccccccccccccccccccccccccccccccccc$		-7.16	265.03	1:5
$\begin{array}{cccccccccccccccccccccccccccccccccccc$		-14.97	302.24	1:5
$\begin{array}{cccccccccccccccccccccccccccccccccccc$		-20.67	366.00	1:5
$\begin{array}{cccccccccccccccccccccccccccccccccccc$		-26.67	285.52	1:5
$\begin{array}{cccccccccccccccccccccccccccccccccccc$		-28.77	327.53	1:5
$\begin{array}{cccccccccccccccccccccccccccccccccccc$		-34.77	271.97	1:5
$\begin{array}{cccccccccccccccccccccccccccccccccccc$		-35.98	365.34	1:5
-29.07351.101:6-41.97260.771:6-43.17317.291:6-56.06227.091:6		-44.97	428.52	1:5
-29.07351.101:6-41.97260.771:6-43.17317.291:6-56.06227.091:6	BD <sub>5</sub>	-10.16	329.85	1:6
-41.97260.771:6-43.17317.291:6-56.06227.091:6	5		351.10	
-56.06 227.09 1:6				
-56.06 227.09 1:6				
			235.40	1:6

# 4.7 R2PI of Benzene-Trifluoroethanol

The binary mixed clusters of benzene and trifluoroethanol were investigated by using the R2PI method. The R2PI spectra for benzene-(trifluoroethanol)<sub>n</sub>, C<sub>6</sub>H<sub>6</sub>-(CF<sub>3</sub>CH<sub>2</sub>OH)<sub>n</sub>, were achieved by utilizing the <sup>1</sup>B<sub>2u</sub>  $\leftarrow$  <sup>1</sup>A<sub>1g</sub> (6<sup>1</sup><sub>0</sub>) transition of benzene molecule, where n=1-7 (Figure 51). Well-resolved R2PI scans were obtained for (C<sub>6</sub>H<sub>6</sub>)<sub>2</sub>-(CF<sub>3</sub>CH<sub>2</sub>OH)<sub>n</sub>, where n = 1-5, as shown in Figure 52. Intracluster reactions were observed following the photoionization of benzene-(trifluoroethanol)<sub>n</sub> clusters. The protonated trifluoroethanol clusters (CF<sub>3</sub>CH<sub>2</sub>OH)<sub>n</sub>H<sup>+</sup>, where n ≥ 2, were detected. The R2PI technique was carried out on deuterated benzene-trifluoroethanol clusters, C<sub>6</sub>D<sub>6</sub>-(CF<sub>3</sub>CH<sub>2</sub>OH)<sub>n</sub>, to investigate the proton transfer reactions. The R2PI scan measured in C<sub>6</sub>D<sub>6</sub>-(CF<sub>3</sub>CH<sub>2</sub>OH)<sub>n</sub>, where n = 1-7, with respect to the 6<sup>1</sup><sub>0</sub> of C<sub>6</sub>D<sub>6</sub> molecule is shown in Figure 53. The frequency of the 6<sup>1</sup><sub>0</sub> transition was detected at 38,805.64 cm<sup>-1</sup>. The spectral characterization of the R2PI scans obtained of the C<sub>6</sub>D<sub>6</sub>/trifluoroethanol system revealed identical spectral features as those of C<sub>6</sub>H<sub>6</sub>/trifluoroethanol system.

In the following sections, the spectral assignments of benzene-trifluoroethanol  $(BT)_n$  clusters, where n = 1-5, and the intracluster reactions upon photoionization of  $BT_n$  neutral clusters are discussed.

# 4.7.1 Spectra of Benzene Trifluoroethanol (BT<sub>n</sub>) Clusters, n = 1-2

The R2PI absorption spectra obtained by monitoring the benzene, B, and benzene- $(CF_3CH_2OH)_n$ ,  $BT_n$ , (n = 1 and 2) mass channels are represented in Figure 54. The assignments have been carefully given in BT and  $BT_2$  mass channels to the observed

spectral features considering the possible fragmentations from higher mass channels. The BT mass channel shows several features, which appear in the BT<sub>2</sub> mass channel with very small ion intensities (as indicated by dashed lines in Figure 54). These features are originating at 109.25 cm<sup>-1</sup> relative to the  $6_0^1$  transition of benzene and assigned to 1:2 cluster. Three peaks are found in the mass channel corresponding to BW (m/z = 96) at the same frequencies as they appear in the BT channel (as shown in Figure 55). These features are assigned to the multi-component cluster (BTW) and they appear in BT channel as a result of the fragmentations of BTW cluster upon photoionization into B<sup>+</sup>T and B<sup>+</sup>W. The observation of benzene-trifluoroethanol-water clusters could be attributed to the contamination of traces of water in the gas line. Several weak features in BT channel could be assigned to 1:1 cluster with no traces in the higher mass channels. The  $6_0^1$  origin of 1:1 cluster is blue-shifted by 65.47 cm<sup>-1</sup>. The mass spectra obtained at the resonance features of 1:1 and 1:2 are displayed in Figure 56. The remaining features in BT mass could not be assigned due to their complexity.

# 4.7.2 Spectra of Benzene Trifluoroethanol (BT<sub>n</sub>) Clusters, n = 3-5

Besides the spectral features described above in  $BT_2$  channel, two clusters were identified to be responsible for the remaining features. Four peaks are present in  $BT_2$ channel arising from the fragmentation of  $B^+T_2W$  [ $B^+T_2W \rightarrow B^+T_2 + W$ ]; these peaks were assigned to the  $BT_2W$  neutral cluster (as indicated by dashed lines in Figure 57). The bands assigned to 1:3 in  $BT_2$  mass channel arise from the fragmentation of  $B^+T_3$  cluster are clearly seen in Figure 57. The spectral shift for the first band is a blue-shift by 179.48 cm<sup>-1</sup>. This cluster was the most blue shifted cluster, which is indicative of the extent of the hydrogen bonding between the OH group of trifluoroethanol and the styrene  $\pi$ -system.. The efficient fragmentation exhibited by the 1:3 cluster is consistent with the  $\pi$ -hydrogen bonding that may be incorporated in the 1:3 structure. The TOF mass spectrum when the laser is tuned to the resonance feature of BT<sub>3</sub> cluster is shown in Figure 58. It is interesting to see how the fragmentation process affects the shape of the BT<sub>2</sub> mass peak. Since B<sup>+</sup>T<sub>3</sub> breaks up in the acceleration region, in a time scale of sub-microsecond, the B<sup>+</sup>T<sub>2</sub> peak will have a tail toward B<sup>+</sup>T<sub>3</sub> mass peak.

The assignments of both 1:4 and 1:5 clusters were afforded by the observation of their fingerprints in T<sub>2</sub>H mass channel. The peaks marked with the dashed lines in T<sub>2</sub>H channel and linked with those in either BT<sub>4</sub> or BT<sub>5</sub> channels were designated to 1:4 or 1:5, respectively, as illustrated in Figure 59. The mass spectrum recorded at the resonance feature of BT<sub>4</sub> cluster has a T<sub>2</sub>H mass peak, which is the product of intracluster reaction (Figure 60). The  $6_0^1$  origins of 1:4 and 1:5 are located at 78.14 and 35.94 cm<sup>-1</sup>, respectively. Although, the spectral shift reverses direction in the 1:4 cluster, all BT<sub>n</sub> clusters (n = 1-5) possess blue-shifted  $6_0^1$  origins relative to that of benzene molecule. The relative frequencies and careful assignments of the BT<sub>n</sub> clusters are tabulated in Table 16.

### 4.7.3 Intracluster Reactions following Photoionization

The  $T_nH^+$  intracluster products were observed for  $n \ge 2$ . The R2PI spectrum of  $T_2H$  channel carries signatures for 1:4 and 1:5 neutral clusters, while the R2PI spectrum of  $T_3H$  channel (in poor resolution) has bands from 1:n, where n = 5-7. To examine the

mechanism of the intracluster reactions, the R2PI experiments were carried out using  $C_6D_6$ . Only the reaction products of protonated trifluoroethanol clusters  $T_nH^+$  ( $n \ge 2$ ) were observed. The R2PI spectra of these products have similar spectral features as those obtained using  $C_6H_6$  (for example, compare Figure 61 with Figure 59). The observation of the  $T_nH^+$  protonated trifluoroethanol clusters (not  $T_nD^+$  clusters) in the R2PI experiments using  $C_6D_6$  suggests that the dissociative proton transfer reactions DPT are taking place in the trifluoroethanol sub-clusters. The proposed mechanism involves a dissociative electron transfer reaction via the neutral precursor 1:n, where  $n \ge 4$ , to produce  $T_n^+$  cluster followed by an evaporative loss (or losses) of trifluoroethanol monomer to form either  $T_3^+$  or  $T_4^+$ , and then the later undergo a dissociative proton transfer reaction to form  $T_2H^+$  or  $T_3H^+$ , respectively. The proposed mechanism is the following:

$$C_6H_6^+ - T_n \xrightarrow{\text{DET}} T_n^+ + C_6H_6$$
 Eq. 4-12

$$T_n^+ \xrightarrow{\text{evaporative loss(es)}} T_3^+ \xrightarrow{\text{DPT}} T_2 H^+ + CF_3 CH_2 O \quad n = 4-5$$
 Eq. 4-13

or

$$T_{n}^{+} \xrightarrow{\text{evaporative loss(es)}} T_{4}^{+} \xrightarrow{\text{DPT}} T_{3}H^{+} + CF_{3}CH_{2}O \quad n \ge 5$$
 Eq. 4-14

The proposed mechanism is supported by these facts: first, the ionization potential of the trifluoroethanol molecule is about 2.25 ev above that of the benzene molecule and it is reasonably to deduce that the DET channel is endothermic for 1:1, 1:2 and 1:3 clusters. It seems that three trifluoroethanol molecules are involved in stabilizing the  $(CF_3CH_2OH)^+$  and hence lowering the ionization potential of the  $(CF_3CH_2OH)_4^+$  cluster to be equal or less

than that of the benzene molecule (IP of benzene is 9.24 ev). Compared that with the benzene/ethanol system, the DET channel is exothermic for 1:2 cluster and the IP(ethanol) is above that of benzene by 1.24 ev. Second, no fragmentation processes were observed for the cluster ions above the 1:3 ion cluster. Third, the evaporative loss of the monomer constituents is the preferential process for cluster fragmentation in many clusters.[124]

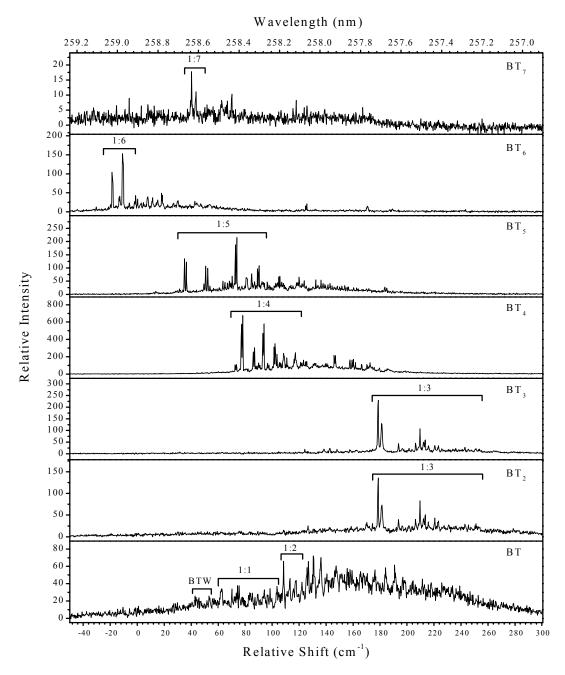


Figure 51: Overall R2PI spectra measured in the benzene (trifluoroethanol)<sub>n</sub> mass channels (BT<sub>n</sub>), n = 1-6, relative to the  $6_0^1$  transition of the benzene molecule at 38,624.95 cm<sup>-1</sup>. The peak designations (m:n) stand for the number of benzene:trifluoroethanol in the neutral cluster which is responsible for the spectral features.

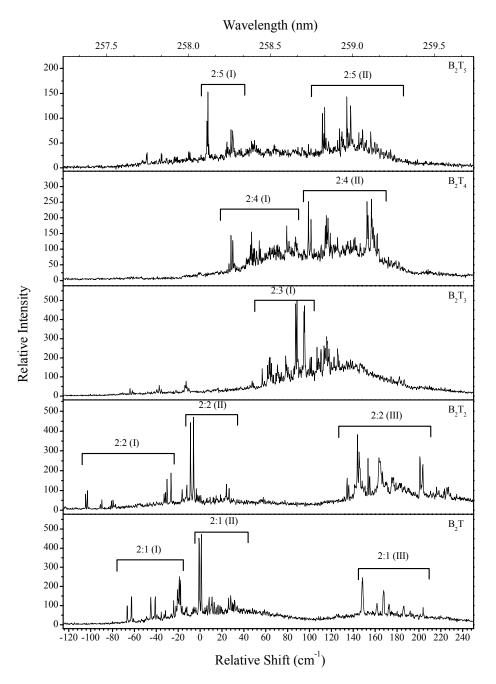


Figure 52: One-color R2PI spectra measured in and  $(benzene)_2$ -(trifluoroethanol)<sub>n</sub> mass channels (B<sub>2</sub>T<sub>n</sub>), n = 1-5, relative to the  $6_0^1$  transition of the benzene molecule at 38,624.95 cm<sup>-1</sup>.

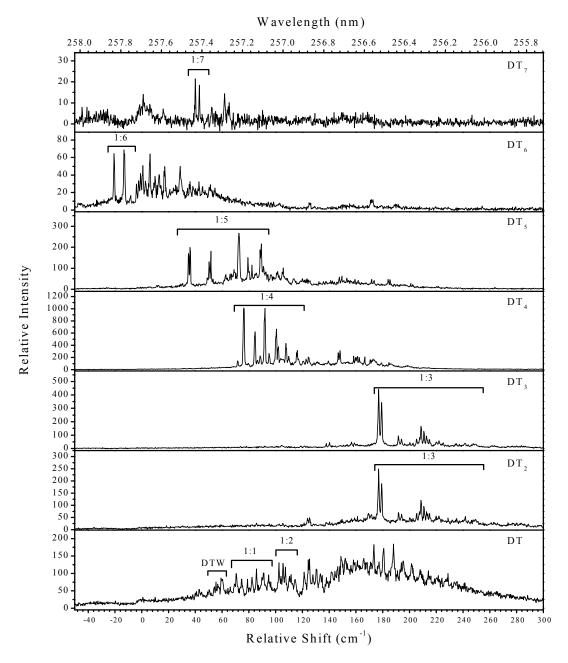


Figure 53: Overall R2PI spectra measured in the  $C_6D_6$ -(trifluoroethanol)<sub>n</sub> mass channels (DT<sub>n</sub>), n = 1-7, relative to the  $6_0^1$  transition of  $C_6D_6$  molecule at 38,805.64 cm<sup>-1</sup>. The peak designations (m:n) stand for the number of  $C_6D_6$ :trifluoroethanol in the neutral cluster which is responsible for the spectral features.

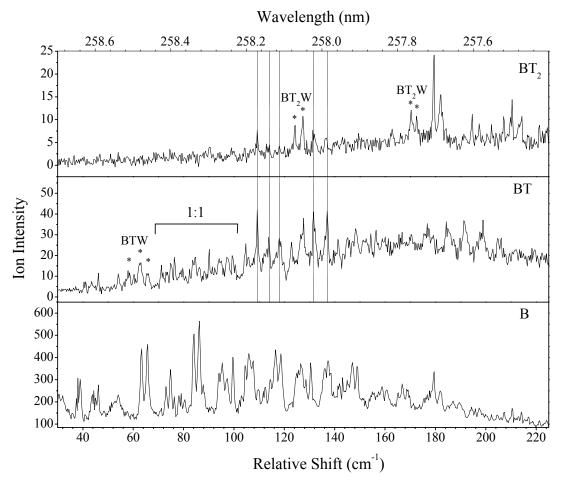


Figure 54: One-color R2PI spectra measured in benzene, B, benzene-trifluoroethanol, BT, and benzene-(trifluoroethanol)<sub>2</sub>, BT<sub>2</sub>, mass channels relative to the  $6_0^1$  transition of the benzene molecule at 38,624.95 cm<sup>-1</sup>. The dashed lines indicate the peaks that assigned to 1:2 cluster. BT<sub>n</sub>W donate to ternary clusters of benzne-(trifluoroethanol)<sub>n</sub>-water.

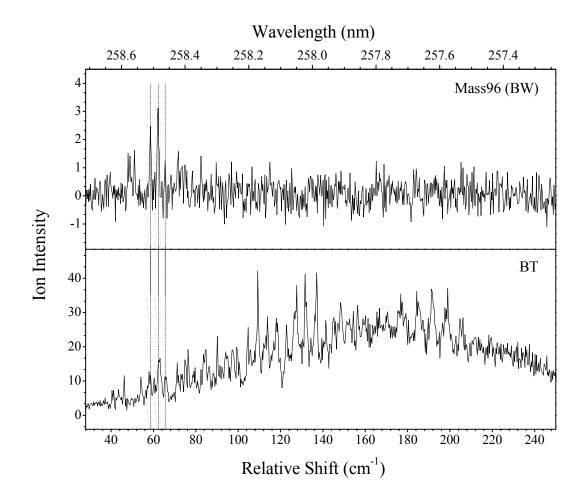


Figure 55: One-color R2PI spectra measured in benzene-trifluoroethanol, BT, and benzene-water, BW, mass channels relative to the  $6_0^1$  transition of the benzene molecule at 38,624.95 cm<sup>-1</sup>. The dashed lines indicate the peaks that assigned to the ternary BTW cluster.

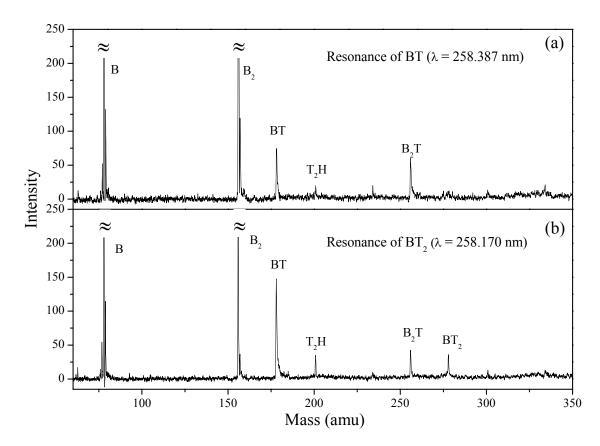


Figure 56: R2PI mass spectra of the benzene-trifluorothanol  $(BT_n)$  cluster beam obtained at the resonance features assigned to (a) BT and (b) BT<sub>2</sub> clusters.

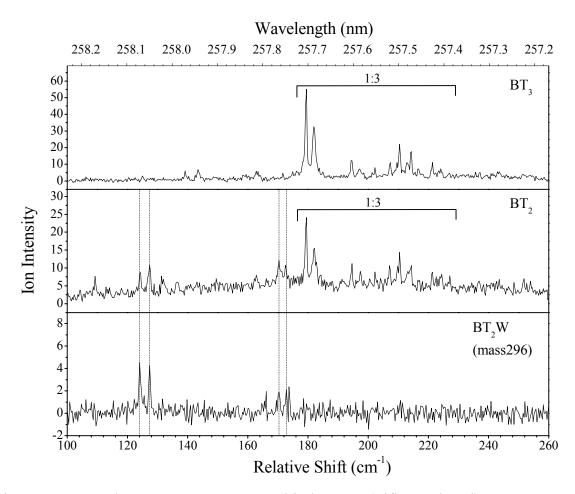


Figure 57: One-color R2PI spectra measured in benzene-(trifluoroethanol)<sub>2</sub>-water,  $BT_2W$ , benzene-(trifluoroethanol)<sub>2</sub>,  $BT_2$ , and benzene-(trifluoroethanol)<sub>3</sub>,  $BT_3$ , mass channels relative to the  $6_0^1$  transition of the benzene molecule at 38,624.95 cm<sup>-1</sup>. The dashed lines indicate the peaks that assigned to the ternary  $BT_2W$  cluster.

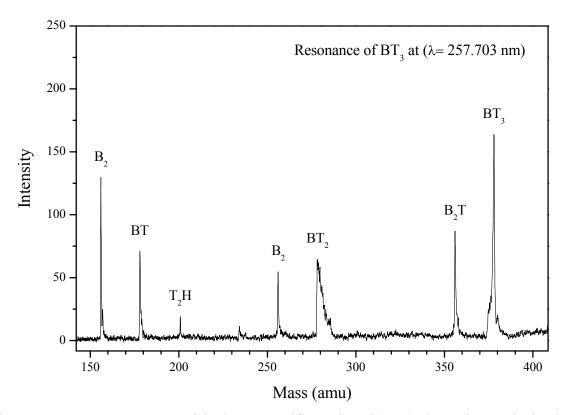


Figure 58: R2PI mass spectra of the benzene-trifluorothanol  $(BT_n)$  cluster beam obtained at the resonance feature assigned to  $BT_3$  cluster.

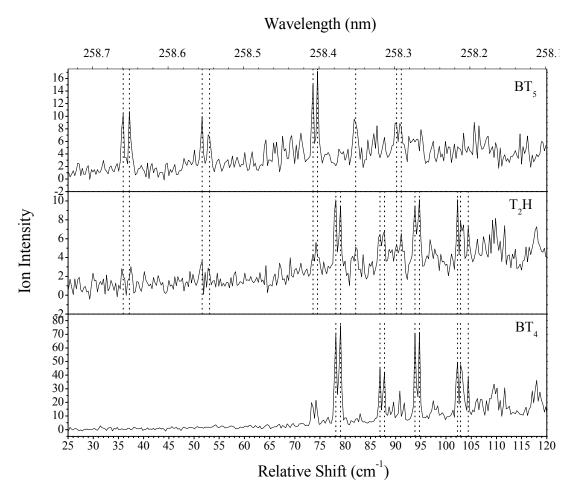


Figure 59: One-color R2PI spectra measured in benzene-(trifluoroethanol)<sub>4</sub>, BT<sub>4</sub>, (trifluoroethanol)<sub>2</sub>H, T<sub>2</sub>H, and benzene-(trifluoroethanol)<sub>5</sub>, BT<sub>5</sub>, mass channels relative to the  $6_0^1$  transition of the benzene molecule at 38,624.95 cm<sup>-1</sup>.

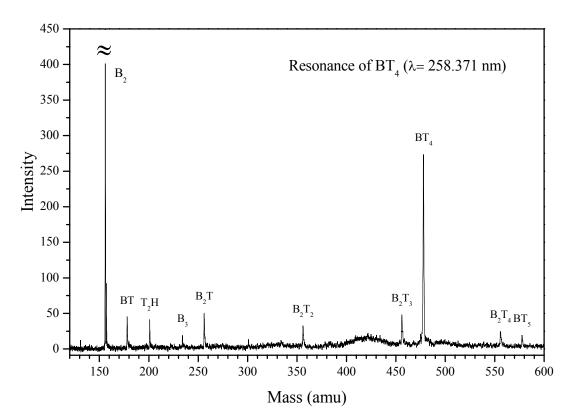


Figure 60: R2PI mass spectra of the benzene-trifluorothanol  $(BT_n)$  cluster beam obtained at the resonance feature assigned to  $BT_4$  cluster.

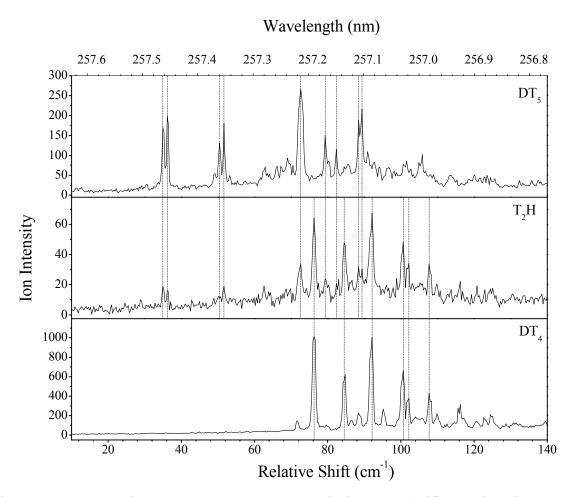


Figure 61: One-color R2PI spectra measured in  $C_6D_6$ -(trifluoroethanol)<sub>4</sub>, DT<sub>4</sub>, (trifluoroethanol)<sub>2</sub>H, T<sub>2</sub>H, and  $C_6D_6$ -(trifluoroethanol)<sub>5</sub>, DT<sub>5</sub>, mass channels relative to the  $6_0^1$  transition of the  $C_6D_6$  molecule at 38,805.64 cm<sup>-1</sup>.

Mass channel	Shift (cm <sup>-1</sup> ) from the $6_0^1$	Relative intensity	Assignments	
	of benzene			
BT	137.07	41.61	1:2	
	131.61	41.14	1:2	
	118.01	28.34	1:2	
	113.78	28.80	1:2	
	109.25	42.14	1:2	
	99.27	19.97	1:1	
	97.16	19.16	1:1	
	94.14	18.51	1:1	
	86.29	14.32	1:1	
	84.78	19.22	1:1	
	83.58	17.74	1:1	
	78.44	11.66	1:1	
	76.33	19.20	1:1	
	74.83	16.21	1:1	
	73.02	12.06	1:1	
	71.20	15.42	1:1	
	65.47	11.45	1:1	
	63.06	16.65	BTW	
	57.94	12.73	BTW	
BT2	221.37	17.13	1:3	
	214.39	21.66	1:3	
	210.44	30.06	1:3	
	207.41	20.16	1:3	
	202.23	14.07	1:3	
	197.98	11.23	1:3	
	194.34	16.83	1:3	
	182.21	32.52	1:3	
	179.48	47.92	1:3	
	172.50	11.69	$BT_2W$	
	170.38	13.36	$\overline{\mathrm{BT}_{2}\mathrm{W}}$	
	136.77	5.97	1:2	
	131.31	7.01	1:2	
	127.38	11.45	$BT_2W$	
	124.96	7.68	1:2	
	124.36	9.25	$BT_2W$	

Table 16: Spectral features observed in the  $BT_n$  mass channels (B = benzene, T = trifluoroethanol, n = 1-5).

	118.10	4.29	1:2
	112.86	6.65	1:2
	109.25	7.06	1:2
BT <sub>3</sub>	221.37	15.82	1:3
	214.09	22.35	1:3
	210.44	28.39	1:3
	207.41	14.67	1:3
	202.23	12.22	1:3
	197.38	9.62	1:3
	194.65	16.84	1:3
	181.91	42.29	1:3
	179.48	70.21	1:3
	163.41	8.18	1:3
BT4	104.40	39.47	1:4
	102.90	47.64	1:4
	102.29	49.57	1:4
	94.74	71.55	1:4
	93.83	70.01	1:4
	87.80	41.60	1:4
	86.89	44.73	1:4
	79.05	76.45	1:4
	78.14	70.93	1:4
BT5	92.63	6.80	1:5
	90.82	8.82	1:5
	81.76	9.57	1:5
	74.51	17.13	1:5
	73.62	14.95	1:5
	52.81	7.07	1:5
	51.61	10.00	1:5
	37.15	10.76	1:5
	35.94	10.33	1:5

## **Chapter 5 Conclusions**

Well-resolved R2PI spectra of the styrene-methanol clusters were described. The spectra reveal a rapid increase in complexity with the number of methanol molecules in the cluster, associated with van der Waal modes and isomeric forms. While only single cluster origins were found for the SM, SM<sub>4</sub> and SM<sub>8</sub> clusters, two distinct isomers were identified for each of the SM<sub>n</sub> clusters with n = 2, 3, 5-7, and 9. The progressive addition of methanol molecules to the SM complex led to the formation of stable cyclic and bicyclic methanol sub-clusters within the SM<sub>4</sub> and SM<sub>8</sub> clusters, respectively. The spectral shifts of the cluster origins reflect the nature of the intermolecular interactions within the binary clusters. Blue shifts were observed for the SM, SM<sub>2</sub> and SM<sub>3</sub> clusters consistent with hydrogen bonding interactions between the OH groups and the styrene  $\pi$ -system. A remarkable switch in the spectral shift from blue to red was observed at the SM<sub>4</sub> cluster consistent with the ring structure of the methanol tetramer. Evidence was provided for intracluster dissociative proton transfer reactions within the  $S_2M_n^+$  clusters with  $n\geq 3$  to generate protonated methanol clusters. It was proposed that the interaction of the methanol molecules with the olefin center of styrene promotes the proton transfer from the styrene dimer to the hydrogen-bonded methanol cluster. These reactions may explain the strong inhibition effects exerted by small concentrations of methanol on the cationic polymerization of styrene.

Well-resolved R2PI spectra of the small styrene-water clusters were reported, and compared with the spectra of the analogous benzene containing clusters. The styrene-water complex exhibited very different spectral shift and structure as compared to the benzene-water complex. The small blue shift of 21.5 cm<sup>-1</sup>, relative to the styrene  $0_0^0$  transition, was consistent with the predominant interaction between water and the ethylene group of styrene. The favorable interactions of water and methanol with the olefin group of styrene may explain the strong inhibition effects of exerted by small concentrations of water on the cationic polymerization of styrene. It was proposed that water molecules tend to aggregate around the olefin center, which promotes the transfer of the charge from the propagating chain to the hydrogen-bonded water cluster.

The R2PI spectra of the styrene-ethanol clusters (SE<sub>n</sub>) were obtained in the region of the S<sub>1</sub>-S<sub>0</sub> transition of styrene molecule. The spectra showed an increase in complexity with the number of ethanol molecules in the cluster, associated with intermolecular vibration modes and isomeric structures. Two distinct isomers were recognized for each of the SE<sub>n</sub> clusters with n = 1, 2, 3 and 5, and three isomers were identified for SE<sub>4</sub> cluster. Evidence was provided for the dissociative proton transfer reactions within the S<sub>2</sub>E<sub>n</sub><sup>+</sup> clusters with n≥2 to generate protonated ethanol clusters.

The binary mixed clusters of styrene and trifluoroethanol were investigated by using the R2PI technique. The R2PI spectra of the  $ST_n$  clusters (with n = 1-4) reveal these spectral characterizations (1) the absence of the R2PI spectrum of ST ion cluster, (2) no evidence of cluster fragmentations is found in that system and (3) No intracluster reaction was observed for styrene-trifluoroethanol  $(ST_n)$  system under the experimental condition. One isomer was identified for  $ST_2$ ,  $ST_3$  and  $ST_4$  clusters.

The binary mixed clusters of benzene and ethanol were studied by using R2PI technique. The R2PI spectra for benzene-(ethanol)<sub>n</sub>, BE<sub>n</sub>, clusters were obtained through the  $6_0^1$  transition of the benzene molecule, where n=1-5. Both the dissociative electron transfer and the dissociative proton transfer reactions were observed within the BE<sub>n</sub> clusters and intracluster ion products were detected following the photoionization of BE<sub>n</sub> clusters. These products were the ionized ethanol clusters (C<sub>2</sub>H<sub>5</sub>OH)<sub>n</sub>, where n  $\geq$  2, and the protonated ethanol clusters (C<sub>2</sub>H<sub>5</sub>OH)<sub>n</sub>H<sup>+</sup>, where n  $\geq$  3. The thermochemistry of the intracluster reactions was discussed.

The R2PI spectra of the benzene-trifluoroethanol clusters  $(BT_n)$  were obtained in the region of the  $6_0^1$  transition of the benzene molecule. The R2PI have been analyzed and the assignments for  $(BT_n)$  clusters have been given, where n = 1-5. This system showed different ion-molecule chemistry. Proton transfer reactions were observed following dissociative electron transfer reactions within  $(BT_n)$  clusters, where n = 4. The mechanisms for the ion-molecule reactions involved evaporative losses of neutral trifluoroethanol molecules. Literature Cited

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## VITA

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