

1 **Simultaneous observations of a sporadic E layer by Digisonde and SuperDARN**
2 **HF radars at Zhongshan, Antarctica**

3 **Xiangcai Chen¹, Jianjun Liu¹, Michael J. Kosch^{2,3,4}, Zejun Hu¹, Zhiwei Wang^{1,5}, Beichen Zhang¹,**
4 **Huigen Yang¹, and Hongqiao Hu¹**

- 5
- 6 1. MNR Key Laboratory for Polar Science, Polar Research Institute of China, Shanghai, China
 - 7 2. Department of Physics, Lancaster University, Lancaster, United Kingdom
 - 8 3. South African National Space Agency (SANSA), Hermanus, South Africa
 - 9 4. Department of Chemistry and Physics, University of KwaZulu-Natal, Westville, South Africa
 - 10 5. Institute of Space Weather, Nanjing University of Information Science and Technology, Nanjing, China

11

12 *Corresponding author: Xiangcai Chen (chenxiangcai@pric.org.cn) and Jianjun Liu (liujianjun@pric.org.cn)

13

14 **Key points:**

- 15 1. An Es layer formed in the lower F region and descended to the E region, observed simultaneously
16 by Digisonde and SuperDARN HF radars.
- 17 2. The formation and evolution of the Es layer related to the afternoon convection reversal.
- 18 3. SuperDARN HF radar measurements suggest the Es layer is elongated with convection circulation.

20 **ABSTRACT**

21 Sporadic E (Es) layers could be composed of metallic ions and formed, modified, or transported by the
22 action of convective electric fields in the high latitude ionosphere. In this paper, by utilizing
23 simultaneous observations from Digisonde and Super Dual Auroral Radar Network (SuperDARN) HF
24 radars at Zhongshan Station (ZHS, 69.4°S, 76.4°E), Antarctica, a thin Es layer, which initially formed
25 in the lower F region and descended into the lower E region, with wavelike structures, was recorded by
26 Digisonde on 14 November 2019. The Es layer-related concurrent ionospheric irregularities were also
27 detected by the SuperDARN ZHS HF radar. By using a global-scale 2-D convection map, combined
28 with images from the Special Sensor Ultraviolet Spectrographic Imager instruments onboard Defense
29 Meteorological Satellite Program (DMSP) spacecraft, it is proposed that the flow shears associated
30 with the duskside convective circulation are responsible for the evolution of the Es layer. Moreover,
31 using the HF radar elevation angle data to measure the scatter height, it is strongly suggested that the
32 Es layer was elongated with convection circulation. The electrodynamic processes responsible for the
33 formation and evolution of the Es layer are discussed.

34

35 **1. Introduction**

36 A sporadic E layer, denoted as Es layer, is a sudden ionization enhancement in the ionospheric E region.
37 It usually appears as a narrow band in the height ranges of 90–130 km in the mesosphere and lower
38 thermosphere region. Unlike the molecular ions NO^+ and O_2^+ that are formed by the photoionization of
39 N_2 and O_2 , which have lifetimes on the order of seconds, ionization of incoming meteoroid atmospheric
40 ablation, such as Fe^+ , Mg^+ , and Na^+ , produces an Es layer composed of metallic ions that can last
41 several hours and even days (e.g., Plane, 2003). With the suddenly enhanced ionospheric electron
42 density, the Es layer reflects radio waves much more efficiently at higher frequencies and with less
43 absorption. This is sometimes used for high-frequency (HF) radio amateur communications and is also
44 valuable for diagnosing the dynamics, electrodynamics, and ion-chemical processes in the mesosphere
45 and lower thermosphere region (e.g., Bourdillon et al., 1995; Chen et al., 2021; Chisham & Pinnock,
46 2001).

47 Observations and simulation results have shown that the mechanism responsible for Es layer formation

48 at mid-latitudes is wind shear (c.f. Haldoupis, 2011, and references therein), while the combined effects
49 of electric fields and neutral winds (Kirkwood & Nilsson, 2000), as well as the abundance of metallic
50 ions (e.g., Bedey & Watkins, 1997), are more important for the formation of Es layers at high latitudes.
51 The convection electric field in polar regions thus provides an important electromagnetic force for
52 convergence/divergence of ion drift (Bristow & Watkins, 1991; Kirkwood & von Zahn, 1991, 1993;
53 Kirkwood & Nilsson, 2000; Nygrén et al., 1984).

54 According to the convective electric field theory (Kirkwood & Nilsson, 2000), ignoring neutral wind
55 effects, the vertical motion of ions can be expressed as: $v_{iz} = \frac{1}{B_0(1+\rho_i^2)} [E_E + (-)\rho_i E_N] \cos I$, where v_{iz}
56 is the vertical motion of the ions, and ρ_i is the ratio of ion-neutral collision frequency to ion
57 gyrofrequency. E_E and E_N represent the eastward and northward components of the electric field,
58 respectively. I is the magnetic dip angle of the local geomagnetic field B_0 . The sign in parentheses
59 refers to the southern hemisphere. With an assumed dawn-dusk electric field at Cambridge Bay (77°
60 CGM lat.), an obvious ion convergence located around 120 km in the pre-midnight hours was observed
61 by MacDougall & Jayachandran (2005) (see their Figure 3c). The Es appeared to form when the
62 convective reversal passed over the observation site. A two-step mechanism responsible for the
63 formation of the Es layer at this ‘cusp station’ was thus proposed by them. In brief, the Es-related
64 horizontal ionization was firstly concentrated by the convection reversal electric field; then vertical
65 convergence of this enhanced ionization concentration into a moderately narrow Es layers at ~120 km
66 height was achieved by the polar cap dawn-dusk electric field. By investigating the roles of electric
67 field and neutral wind in the generation of Es layers within the polar cap, Nygrén et al. (2008) found
68 that the relative intensity of the electric field comparing to the neutral wind would influence the height
69 of the Es layer and also its density variations. Furthermore, gravity waves may also play an important
70 role in the formation and motion of the Es layer in the vertical direction (e.g., MacDougall et al., 2000a,
71 2000b).

72 It is expected that the Es layer could drift horizontally and vertically, but observations of its evolution
73 process are still rare due to the limited field of view of each instrument. Until now, only a few case
74 studies considered the dynamic process of the Es layer in the horizontal and vertical dimensions
75 simultaneously (e.g., MacDougall et al., 2000a). It is therefore still important to extend observations
76 with instruments to further understand the Es layer features at high latitudes. The motivation for this

77 paper is thus to investigate the formation and evolution process of an Es layer at the Zhongshan station
78 (ZHS, -74.9° CGM lat.), Antarctica, and to investigate its morphological characteristics observed by
79 different instruments, as well as to consider its dynamic process in relation to the convective electric
80 field. An overview of the instruments used in this paper is given in section 2. The Es layer observed by
81 the Digisonde and SuperDARN ZHS radars on 14 November 2019 is described in section 3. The
82 formation process of the Es layer and its related ionization enhancement effects on SuperDARN
83 oblique rays are discussed in section 4. A summary of the results is finally given in section 5.

84

85 **2. Instruments and data sets**

86 **2.1. Digisonde radar (DPS-4D)**

87 An ionosonde is a powerful tool that is extensively being used for monitoring real-time ionospheric
88 conditions all over the world. The Digisonde Portable Sounder (DPS-4D) located at ZHS (69.4°S ,
89 76.4°E , $\text{LT} \approx \text{UT} + 5 \text{ h}$, and $\text{MLT} \approx \text{UT} + 2 \text{ h}$) was developed by the University of Massachusetts
90 Lowell Center for Atmospheric Research (UMLCAR), and has been in operation since 2010. By
91 utilizing one crossed delta antenna for transmission, and four crossed magnetic dipole antennas for
92 reception, the DPS-4D can adopt the multi-beam sounding mode to get six digitally synthesized
93 off-vertical reception beams, as well as one vertical beam that separates O-mode and X-mode echoes
94 (Reinisch et al., 2009). With 0.05 MHz frequency step sweeping from 0.5 MHz to 9.5 MHz, the
95 amplitude of ionospheric echoes at the height where the local cutoff frequency equals the transmitted
96 frequency were recorded in the ionograms. Finally, using SAO software to manually scale the
97 ionograms with a time interval of 7.5 mins and a height resolution of 2.5 km, the main ionospheric
98 parameters such as Es critical frequency ($f_0\text{Es}$) and its virtual height ($h'\text{Es}$) can be obtained.

99 **2.2. SuperDARN HF radar**

100 The ZHS HF radar, as part of the Super Dual Auroral Radar Network (SuperDARN) in the southern
101 hemisphere, consists of a main array of 16 transmitting/receiving antennas (main array) and a passive
102 array (interferometer) of 4 antennas only for receiving. To obtain sufficient backscatter returns at
103 ionospheric altitudes (i.e. transmitted radio waves refracted and orthogonal to the local magnetic field),

104 SuperDARN radars operate at a variable frequency within 8–20 MHz (Greenwald et al., 1985). By
105 transmitting a multi-pulse sequence with a dwell time of 3 or 7 s on each beam, the backscattered
106 signals are sampled/processed to produce multi-lag complex autocorrelation functions (ACF) as a
107 function of range within 1 or 2 mins (Chisham et al., 2007; Greenwald et al., 1995). Finally, the
108 Doppler line-of-sight velocities can be obtained from a least-square fit to the phase of the complex
109 value of the ACF as a function of lag, while the power and spectral width of the backscatter can be
110 estimated from a Lorentzian or Gaussian fit to the decorrelation of the ACF (Hanuise et al., 1993).
111 Moreover, by measuring the phase delay of backscatter signals between the main array and the passive
112 array, the elevation angle (i.e. the vertical arrival angle) for each range gate can be determined using
113 known information about the radar configuration (e.g., Milan et al., 1997).

114 Currently, the SuperDARN ZHS radar operates at a frequency of ~ 10 MHz and scans 16 beams with a
115 dwell time of 3 s for each beam, giving a fan-shaped field-of-view ($\sim 54^\circ$ in azimuth) scanning every 1
116 min. For each beam, with a pulse length of 300 μs corresponding to a range gate length of 45 km, a
117 total of 75 range gates are sampled. The time lag to the first gate is set at 1200 μs , which means that the
118 radar observes theoretically from 180 to 3555 km range (Hu et al., 2013; Liu et al., 2013). Figure 1
119 illustrates the relative configuration of the DPS and SuperDARN radars at ZHS for remote sensing. The
120 oblique rays of SuperDARN radar backscatter are primarily from plasma density irregularities in
121 ionospheric E and F regions (ionospheric scatter) (e.g., Chen et al., 2015; Milan & Lester, 2001),
122 reflected from meteor trails in the lower E region (meteor scatter) (e.g., Chisham & Freeman, 2013;
123 Hall et al., 1997), and from the Earth’s surface after reflection/refraction from the ionosphere (ground
124 scatter) (e.g., Huang et al., 2018), but only ionospheric scatter is valuable for our current study. We,
125 therefore, attempt to filter out ground backscatter by removing data with line-of-sight velocities of less
126 than 30 m/s and spectral widths of less than 35 m/s, as well as spectral power lower than 3 dB
127 (Chisham & Pinnock, 2002).

128 Although the SuperDARN line-of-sight velocities are widely used, more important are the global-scale
129 2-D plasma convection maps (Chisham et al., 2007; Ruohoniemi & Baker, 1998). By applying the
130 so-called Potential Fit analysis (Ruohoniemi & Baker, 1998), the line-of-sight velocities measured by
131 multiple radars are combined into a common data set for each scan. Hence, the vectors of global
132 plasma convection flow can also be inferred. Specifically, the plasma convection flow measurements

133 are firstly median filtered by several neighboring beams and for three consecutive scans. The smoothed
134 velocities are then sorted to fill a uniform grid of magnetic latitude and longitude equivalent to 1° of
135 magnetic latitude. By fitting the observed grid velocities into a statistical model of plasma flow
136 parameterized by the solar wind and IMF conditions and the Earth's magnetic dipole orientation (e.g.,
137 Cousins & Shepherd, 2010), a global-scale 2-D plasma convection map is obtained. As the ionospheric
138 plasma convection is predominantly driven by the electric field in the high latitude regions, its related
139 vectors of electric field and electrostatic potential can also be estimated.

140 The SuperDARN backscatter elevation angle data has recently been utilized for ionospheric diagnostics.
141 Using the elevation angle of ground scatter echoes, André et al. (1998) derived the critical frequencies
142 in the high-latitude E-region. Milan & Lester (2001) applied it to estimate the altitude of ionospheric
143 backscatter. For the classical SuperDARN virtual height model with mapping errors (e.g., Chen et al.,
144 2015, 2016, 2017; Yeoman, et al., 2001, 2008), Chisham et al. (2008) proposed an improved
145 range-finding algorithm based on a statistical analysis of backscatter elevation angle of arrival
146 information. Moreover, Gillies et al. (2009, 2010, 2011) applied the elevation angle for estimating the
147 refractive index of the ionospheric plasma to improve the SuperDARN velocity measurements. After
148 that, the estimated refractive index is further utilized to monitor the F-region maximum electron
149 densities (Ponomarenko et al., 2011). The above mentioned methods for determining the ground range
150 and refractive index of ionospheric scattering volumes were systematically evaluated by Greenwald et
151 al. (2017). Due to the intrinsic technical difficulties with direct calibration of a time offset between the
152 two antenna arrays (Ponomarenko et al., 2015, 2018), getting reliable measurements of echo elevation
153 angles is painstaking. Nevertheless, the relatively straightforward algorithm (e.g., Milan et al., 1997)
154 gives reasonably accurate results, which has recently been evaluated systematically by Shepherd
155 (2017).

156 In this study, both SuperDARN ZHS beam data and convection map in the Southern Hemisphere are
157 used. A snapshot of the global plasma convection map at 16:00 – 16:02 UT is shown on the right panel
158 of Figure 1. The collocated DPS-4D and SuperDARN ZHS radars thus form an especially good
159 geometry for investigating the ionospheric plasma dynamic process.

160 **2.3. Defense Meteorological Satellite Program (DMSP)/SSUSI**

161 With a far-ultraviolet (FUV) imager, i.e. the Special Sensor Ultraviolet Spectrographic Imager (SSUSI),
162 onboard the Defense Meteorological Satellite Program (DMSP) spacecraft, the FUV emissions in 5
163 bands can be recorded by cross-track scanning (Paxton et al., 2002). A series of DMSP satellites
164 operating at an altitude of ~ 840 km, with an inclination of $\sim 98.9^\circ$ and a period of ~ 101 minutes, are
165 sun-synchronous, polar-orbiting around the Earth (Hardy et al., 1984). Each scanning image is built up
166 over 20–30 minutes when the satellite flies over the polar region. During this study, the DMSP F17 and
167 F18 were sequentially (spacing of ~ 4 –5 mins) crossing the Southern Hemisphere. Hence, we only
168 employ the SSUSI LBHS (140–160 nm) emission from DMSP F17, which is primarily produced by
169 precipitating electrons and secondary electrons produced by precipitating protons.

170

171 **3. Observational results**

172 **3.1. Es layer observed by DPS**

173 Figure 2 shows a sequence of 7.5 min interval ionograms for $\sim 15:37$ – $17:30$ UT on 14 November 2019.
174 A thin Es layer initially formed at ~ 260 km (virtual height) and suddenly occurs in the field-of-view of
175 Digisonde shown in Figure 2c. This thin Es layer became stratified at a time after 16:22 UT (see Figure
176 2g), with sometimes discontinuous features in the frequency range (see Figure 2i), echoes spread in
177 height (see Figures 2k and 2l), and a secondary echo such as shown in Figure 2o. The Es layer finally
178 descends from the lower F region into the lower E region.

179 To examine the altitude and temporal evolution of the Es layer in detail, Figure 3 presents the
180 height-time-amplitude of the echoes (i.e. integration of the echo amplitude) recorded by the DPS-4D at
181 sounding frequency ranges of 3.5–3.75 MHz (3a), 6.5–6.75 MHz (3b), and greater than 2 MHz (3c)
182 during the time interval of 13:00–21:00 UT. From Figures 3a and 3c, the ionospheric echoes from the F
183 region are mainly above the virtual height of 250 km. is the F-region was almost completely obscured
184 by the Es layer for time intervals of $\sim 18:15$ – $18:45$, and $\sim 19:45$ – $20:15$ UT. The critical frequency (peak
185 height) of the F2 layer gradually decreased (uplift) from ~ 4 (~ 200) to 2.5 MHz (220 km) during this
186 time interval, which is associated with the diminished photoionization due to sunset. The Es layer
187 suddenly occurred at the lower F region and subsequently descended into the E region (< 140 km) with
188 an average speed of ~ 43.5 (± 7.35) m/s as estimated from Figure 3b. This morphology, called

189 intermediate layer (indicated by arrows in Figure 3b), has been extensively reported and simulated by
190 many studies at mid-latitudes (e.g., Bishop & Earle, 2003; Mathews et al., 2001), and is expected to be
191 modulated by horizontal wind field and tidal winds. Here, we still address this intermediate layer as an
192 Es layer because of its continuous evolution. Subsequently, the Es layer shows ascend/descend
193 wavelike features for about ~ 3.5 h, and finally disappeared out of the field-of-view of the DPS-4D at
194 $\sim 20:15:10$ UT. As shown in Figure 3b, the lowest virtual height of the Es layer is separately observed at
195 17:30:10, 18:30:10, 19:30:10 UT, which appears to be a periodic perturbation of about 1 h. It seems
196 like that the period of the descending Es layer corresponded to strong ionization enhancements (i.e.
197 large integration of the echo amplitude shown in Figure 3c). Nevertheless, the f_0E_s can not be properly
198 estimated at time intervals of $\sim 16:00\text{--}16:15$, $\sim 16:52\text{--}17:15$, and $\sim 18:00\text{--}18:45$ UT, due to the upper
199 limit of the DPS sounding frequency at 9.5 MHz.

200 3.2. Es layer observed by SuperDARN ZHS radar

201 The Es layer-related ionospheric echoes are also detected by the collocated SuperDARN ZHS radar. A
202 range gate versus time plot of Doppler line-of-sight velocities for beams 13(a), 7(b), and 1(c) from
203 13:00 to 21:00 UT is shown in Figure 4. Three intervals with differential backscatter features can be
204 easily identified. Before $\sim 15:10$ UT, seen in Figures 4a and 4b, the HF radar echoes are mainly
205 backscattered from range gates greater than 5 (i.e. slant ranges from 405 to 990 km), which is expected
206 to be ionospheric echoes with 1/2-hop from the F region. Subsequently, the HF radar ionospheric
207 returns are also detected in the near-range gates. These ionospheric echoes show periodic-like flow,
208 which is delineated by the black lines in Figure 4c. The third interval of ionospheric echoes is mainly
209 observed in the far ranges (i.e. range gate > 9) at a time after $\sim 18:00$ UT, which is also expected to be
210 the 1/2-hop backscatter from the ionospheric F region.

211 Although all three panels in Figure 4 display similar ionospheric characteristics, some differences exist,
212 such as the ionospheric echoes were rarely observed in beam 1 before 15:30 UT. This means that the
213 HF radar beams remotely sound different ionosphere regions. To further understand the ionospheric
214 conditions for Es layer formation, four ionospheric plasma convection maps from the SuperDARN
215 network overlaid with the SSUSI LBHS images from DMSP F17 crossing are depicted in Figure 5. To
216 ease the comparison of these two data sets, two color codes have been applied to represent plasma flow
217 (color from reddish black to yellow) and auroral activities (color from green to yellowish-white),

218 respectively. The time for each convection map corresponds, approximately, to the satellite position
219 somewhere in the middle of each plot. Similar to Figure 1b, the field-of-view of the ZHS radar is
220 illustrated by the dashed blue line, and the red star indicates the location of ZHS. In addition, the thick
221 red curve represents the Heppner–Maynard Boundary, which in general represents the equatorward
222 boundary of the auroral oval (e.g., Imber et al., 2013; Wang et al., 2022).

223 In the early phase of the Es occurrence shown in Figure 5a, three regions of plasma flow within the
224 field-of-view of ZHS radar are observed. Just poleward of the auroral oval around 78° MLAT, the
225 plasma flow with convection reversal is spatially localized between 18 and 21 MLT. This is associated
226 with the ionospheric flow shear observed by beam 13 around 15:36 UT in Figure 4a. An eastward
227 convection flow roughly located at the auroral oval around 21–23 MLT is observed. This is also
228 identified by HF radar beam 1 at range gates of 10–15 at this time in Figure 4c. The observational
229 evidence strongly demonstrates that the ZHS is located under the auroral oval before the appearance of
230 the Es layer. This conclusion could also be verified by the ionospheric echoes from virtual heights of
231 ~ 140 – 190 km shown in Figure 3c before 16:30 UT, which signifies the possible particle precipitation
232 with an energy of \sim keV. Nevertheless, the poleward boundary of the auroral oval is observed by HF
233 radar beams 13 and 7.

234 For the next two satellite crossings shown in Figures 5b and 5c, two spatially localized plasma flows
235 are observed. The large-scale eastward convection flows around midnight are situated at far ranges in
236 the polar cap, while plasma flow near the auroral oval poleward boundary (and/or auroral arc) was
237 observed with low velocities at near ranges. The sheared plasma flow over ZHS associated with aurora
238 emissions may be expected (e.g., Liu et al., 2011), and was observed by the HF radar line-of-sight
239 velocities (indicated by arrows in Figure 3). Finally, for the end phase shown in Figure 5d, ZHS was
240 located in the polar cap region, even though the HMB represents the auroral oval boundary incorrectly.
241 The Es layer had disappeared from the field of view of the DPS. Joint satellite and ground radar
242 observations confirm that the spatial-temporal formation and evolution of the Es layer were associated
243 with the afternoon convection reversal.

244 Line-of-sight plasma drift velocities towards the radar were continuously observed in the near range
245 gates (see Figure 4), which implies that the plasma was convected from the F region to the E region
246 (see the sketch map in Figure 1). To further clarify the Es layer-related ionization enhancement

247 detected by DPS-4D that is also monitored by the SuperDARN ZHS radar, Figure 6a shows the
 248 occurrence distribution of elevation angles from different beams for range gates 0–18, i.e. green, blue,
 249 and black histograms for beams 13, 7 and 1, during this time interval. Assuming spherical propagation
 250 geometry and straight-line propagation, the virtual height of the ionospheric scatter volume from radar
 251 beams can be estimated as

$$252 \quad h = \sqrt{R_E^2 + r^2 + 2R_E r \sin\alpha} - R_E$$

253 where R_E is the Earth's radius, the slant range (r) is determined by the center of the effective scatter
 254 volume at $180 \text{ km} + 45/2 \text{ km} + \text{gate} \times 45 \text{ km}$, and α is the elevation angle for the corresponding range
 255 gate. Figures 6b–6d show the scatterplot of backscatter altitude versus time with slant ranges for three
 256 radar beams. The lowest virtual height of the Es layer and its associated stratified/sublayers observed
 257 by the DPS-4D are also overlaid as red histogram and asterisk, respectively. Additionally, the peak
 258 height of the F2 layer (i.e. hmF2), scaled from ionograms, is plotted as black dots in Figure 6c.
 259 Although the straight-line propagation assumption may not be true, the hmF2 denoted in Figure 6c can
 260 be treated as a demarcation altitude with backscatter from the F region above and from the Es-related E
 261 region below. The ionospheric echoes in the E region (i.e. $< 150 \text{ km}$) from beams 13 and 7 precede the
 262 appearance of Es observed by DPS. The ionospheric echoes from beam 1 are closely associated with
 263 the evolution of the Es layer and its related sublayers (red asterisk) from 260 to 150 km. As the
 264 majority values of elevation angle from beams 13 and 1 are higher than from beam 7, as shown in
 265 Figure 6a, some ambiguity is present near the lowest altitude of Es layer. One reason is due to ignoring
 266 the effects of refraction, which tends to produce an overestimation of the true altitude of the scatter
 267 volume by about 10–20 km (Milan & Lester, 2001). Another reason may be the classic algorithm for
 268 acquiring elevation angle values at ZHS without considering a nonzero X component (see Shepherd,
 269 2017 for more details). Nonetheless, these ambiguities can not cause a serious impact on our height
 270 comparison.

271

272 **4. Discussion**

273 **4.1. Dynamic process of Es layer**

274 We have observed a thin Es layer initially formed in the lower F region followed by a descent into the
275 E region, which is especially different from the normal Es layer usually formed in the E region about
276 90–130 km. With the duration of the Es layer observed by DPS-4D of more than 3 hours, an abundance
277 of metallic ions within the Es layer is expected (MacDougall et al., 2000a), which could be transported
278 by the approximately equatorward directed convective electric field (Bedey & Watkins, 1997; Nygren
279 et al., 1984; Parkinson et al., 1998). The plasma flow in the polar ionosphere is modulated by the
280 large-scale convection electric field. A simple model presented by Bedey & Watkins (1997) has shown
281 that metallic ions could be lifted from the nominal background metallic layer below 100 km into the
282 lower F region on the dayside. With the flow horizontally through the polar cap and toward the
283 nightside, Bedey & Watkins (1997) proposed that the metallic ions could form in a narrow vertical
284 stream (~50 km) and precipitate within a limited band of geomagnetic latitudes (see their Figure 6).
285 Their results clearly showed that the abundance of the metallic ions enhanced only in a small portion of
286 the region where westerly electric field dominated (in the northern hemisphere) and ion precipitation
287 predominantly occurred before local midnight. For our current study of the long-lived thin Es layer at
288 ZHS, Antarctica, we are concerned with the ionospheric E-F region coupling effect caused by the Es
289 layer. A southward electric field associated with eastern plasma flow (see Figure 4) dominates the
290 evolution process of the Es layer from the lower F region to the E region, which means that an
291 anti-sunward plasma drift is occurring over the site. Nevertheless, an ion motion gradient that exists in
292 the F region should be a precondition for the initial convergence of the Es layer.

293 As mentioned in the Introduction, MacDougall & Jayachandran (2005) suggested that a two-step
294 formation mechanism is responsible for the Es formation at a ‘cusp latitude’ station, i.e. firstly
295 horizontal convergence of ionization by the afternoon convection reversal electric field, and then
296 vertical convergence of this ionization by the mainly dawn-dusk electric field in the polar cap. By
297 carefully checking the DMSP/SSUSI data (see Figure 5), we recognize that the ZHS station was shifted
298 from the duskside auroral oval into the polar cap. The initial formation of the Es layer in the lower F
299 region observed by the DPS at ~15:50 UT would correspond to the poleward boundary of the auroral
300 oval. Moreover, the last convection reversal is also clearly observed by the SuperDARN radar beam 13
301 at time ~20:15 UT (see Figure 4a). This implies that the initial formation of the Es layer at the lower F
302 region would be due to the duskside convection reversal passing over the ZHS. It also suggests that

303 once the ionospheric convection reversal condition is suitable for the metallic ions to converge, the Es
304 layer would be formed even in the lower F region.

305 The Es layer descended with an average vertical speed of ~ 43.5 m/s, i.e. estimated from the time
306 interval of $\sim 15:52$ – $16:37$ UT and from 260 to 142.5 km. Considering the convection electric fields
307 within the descending period, we estimate the electric field contributes to the vertical velocity about 15
308 m/s on average, which means the action of gravity waves and/or tidal waves must be responsible for the
309 descending of the Es layer. The stratified features of the Es layer shown in Figure 2 and wave-like
310 ionospheric echoes that were both detected by the DPS (see Figure 3) and SuperDARN radar (see
311 Figure 4c) are strong evidence for the action of gravity waves. Similar features were also observed and
312 explained by MacDougall et al. (2000a). With the Es layer descent into the E region, a combined
313 electric field and neutral wind, as well as a gravity wave (e.g, Kirkwood & Nilsson, 2000) resulting in
314 the observed wave-like motion of the Es layer, are expected. This Es layer contained metallic sodium
315 ions with wave-like motions in the lower E region resulting in the neutralization of sodium atoms that
316 produced a sporadic sodium layer that has recently been reported by Chen et al. (2021).

317 **4.2. SuperDARN backscatter returns refracted by Es layer**

318 The Es layer-related ionization enhancement is accompanied by simultaneous coherent backscatter
319 echoes from the F region at far ranges to the E region in the near ranges in SuperDARN ZHS radar.
320 These ionospheric features strongly indicate the existence of electrical coupling between the E and F
321 regions along the geomagnetic field. Observations of ionospheric irregularities associated with the Es
322 layer were also reported by Ogawa et al. (2009) in the Japanese sector. The SuperDARN ZHS radar
323 monitors the Es layer from the F to E regions as shown in Figure 6, indicating that there is a continuous
324 supply of decameter-scale irregularities associated with the enhanced ionospheric electron density.

325 To quantitatively estimate the Es layer-related electron density that results in the bending of ray paths
326 from the ZHS HF radar, we use Ponomarenko et al.'s (2011) technique to evaluate the maximum
327 electron density in the far range F region and near range E region, respectively. Assuming the
328 ionosphere is spherically stratified and symmetric so that Snell's law applies for the rays launched at an
329 elevation angle (α , at the radar site) and reaching the orthogonality condition in the ionosphere at the
330 height of h_s , the refractive index of the ionospheric plasma can be written as (Gillies et al., 2009):

331
$$n_r = \frac{R_E \cos \alpha}{R_E + h_s \sin I}$$

332 where R_E is the Earth's radius (i.e. 6370 km), I represents the geomagnetic inclination at the backscatter
 333 points, and h_s is the scatter echo height. Since the refraction index of HF radio waves at frequency f_0 is
 334 related to the plasma frequency f_p as

335
$$f_p = f_0 \sqrt{1 - n_r^2}$$

336 the maximum electron density at the Es layer can be evaluated from

337
$$N_{mEs} = \frac{4\pi^2 m \epsilon_0 f_p^2}{e^2}$$

338 where m and e are the electron mass and charge and ϵ_0 is the permittivity of free space.

339 We use the elevation angle data from beam 7, which is closest to a uniform distribution of angle values
 340 (see Shepherd, 2017 Figure 4d). For the initial formation of the Es layer in the lower F region, h_s was
 341 ~ 260 km, and the onset elevation angle value was $\sim 39.95^\circ$, which corresponds to a refractive index n_r
 342 of about ~ 0.76 . The derived maximum electron density (i.e. NmEs) by the ZHS radar is $\sim 0.55 \times 10^6$
 343 cm^{-3} , which is only half of the Digisonde measurement of $\sim 1.12 \times 10^6$ electrons/ cm^3 . For the final Es
 344 echoes recorded by the SuperDARN ZHS radar in the near range gates around 20:00 UT, the elevation
 345 angle of $\sim 20^\circ$ corresponds to a refractive index of 0.97. The peak electron density of 0.76×10^5 cm^3
 346 derived by the SuperDARN radar is ~ 3.8 times smaller than the Digisonde measurement of $\sim 2.9 \times 10^5$
 347 electrons/ cm^3 .

348 Assessing the quality of the peak electron density estimates made from elevation angle measurements
 349 by the SuperDARN Rankin Inlet (RKN) radar was recently carried out by Koustov et al. (2020). By
 350 comparing the SuperDARN RKN results to ionosonde and incoherent scatter radar measurements, the
 351 underestimation of peak electron density by SuperDARN radar (by up to 30% in their cases) could be
 352 mainly due to low background electron densities (Koustov et al., 2020). However, there are other
 353 effects that also very likely contribute to our current measurements. One is the Digisonde that detects
 354 signals from localized regions with the strongest electron density, while the SuperDARN ZHS radar
 355 range gate, with a spatial resolution of 45 km, could smooth out the peak electron densities. This
 356 smoothing effect is expected to be stronger for ionospheric conditions with higher patchiness and

357 poorer propagation conditions (Koustov et al., 2020). Secondly, the critical frequency is manually
358 acquired from Digisonde ionograms containing the extraordinary wave, which has a higher frequency
359 than the ordinary wave, i.e. O-mode wave, by half the electron gyrofrequency. Last but not least, the
360 Digisonde and SuperDARN ZHS radar detected the Es layer from different locations. However, the
361 measurement results will not influence our conclusion that the SuperDARN ZHS radar oblique rays
362 experience significant ionospheric refraction, resulting in bending of radio paths, by the Es layer.

363 4.3. The elongation of the Es layer

364 The Es layer initially forms in the lower F region and descends into the E region, which is detected by
365 both the DPS-4 Digisonde and SuperDARN ZHS radar. As seen from Figure 6c, the onset time of the
366 ionospheric echoes monitored by the SuperDARN ZHS radar at 200 km precedes the DPS detections
367 by about 60 min. The time of the Es layer detected by the DPS at 200 km corresponds to the
368 disappearance of the Es-related ionospheric irregularities at the same height (i.e. ~16:30 UT). As the
369 SuperDARN ZHS radar scans in azimuth to the west of the site, it is expected that the Es layer formed
370 in the lower F region, as detected by the HF radar, will be transported by the convective circulation.
371 The range from the center of the effective scatter volume of the SuperDARN range gate 7 to the
372 field-of-view of the DPS corresponds to about 483 (± 57.7) km at 200 km altitude in the F region.
373 Assuming the Es layer-related plasma flow is always equal to the large-scale convection velocity, i.e.
374 ~174 m/s on average, the Es layer will take ~46 (± 5.5) min to travel 483 km. This means that it is
375 possible the Es layer was formed as an enhanced ionization patch with a size of about 500 km and then
376 transported with the ionospheric convection circulation.

377 Bedey & Watkins (1997) have suggested that the Es layer related metallic ions can be lifted from the
378 background ion layer into the F region, where the ions can then be transported over large distances
379 before “precipitating” in relatively narrow bands in latitude (see their Figure 6). Furthermore, they
380 expected that the Es layer related ions are deposited in the lower E region within a very narrow band of
381 latitudes, 74° – 76° and 76° – 77° MLAT for westerly field strengths of 50 mV/m and 10 mV/m,
382 respectively. During this time interval, the electric field was on average less than ~12 mV/m (see Chen
383 et al., 2021, Figure 1c), the Es layer initially formed at a latitude of ~ 78° MLAT in the lower F region
384 and descended to around 75° MLAT in the lower E region as observed by the SuperDARN radar (see
385 Figure 5). Bedey & Watkins’s (1997) model analysis yields results that are in good agreement with our

386 current observations, which strongly demonstrates that the formation of the Es layer is due to the
387 large-scale redistribution of metallic ions via horizontal transport in the F region (e.g., MacDougall &
388 Jayachandran, 2005) followed by a descent into the lower E region within a narrow band (Bedey &
389 Watkins, 1997).

390

391 **5. Summary**

392 In this paper, we have studied the dynamic process of an Es layer simultaneously observed by the DPS
393 Digisonde and SuperDARN ZHS radars at Zhongshan, Antarctica. As the ionospheric Es layer has a
394 significant impact on radio wave propagation, the traditional techniques that only employ a single
395 instrument, such as ionosonde for Es layer observation, are not sufficient to resolve the morphology
396 and dynamics of the Es layer in the spatial domain. We show a unique example of an Es layer, which
397 initially formed in the lower F region and then descended into the lower E region. This is expected to
398 result from a two-step formation mechanism involving horizontal convergence of ionization by the
399 electric fields of the duskside convection reversal, and then vertical convergence of this enhanced
400 ionization by the electric fields (MacDougall & Jayachandran, 2005). The initial formation of the Es
401 layer-related convection reversal was detected by the SuperDARN line-of-sight velocity, while its
402 subsequent evolution process was related to the auroral activity at the poleward boundary of the auroral
403 oval, which is confirmed by images of SSUSI LHBS emission from the DMSP spacecraft. The
404 Es-related ionization enhancement also gives rise to SuperDARN radar refraction bending the beam
405 into the lower F and E regions. The occurrence of the Es layer has a time lag between the DPS and
406 SuperDARN radar observations, which can be simply explained by the large-scale convection
407 circulation at the lower F region. This implies that the Es layer is elongated as an enhanced ionization
408 patch with a size of about 500 km. Utilizing a multi-instrument to study the evolution processes of the
409 Es layer thus has significant implications for understanding electrodynamics and electrical coupling
410 between the E and F regions in the polar region.

411

412 **Data Availability Statement**

413 The ZHS DPS and SuperDARN data can be download from the Data Centre for Meridian Space Weather

414 Monitoring Project webpage (<http://data.meridianproject.ac.cn>), while the SuperDARN 2-D convection map and
415 the electric field is available from the Virginia Tech portal (<http://vt.superdarn.org>). The DMSP/SSUSI LBHS data
416 is available from [Data Products | SSUSI \(jhuapl.edu\)](https://ssusi.jhuapl.edu/data_products) (https://ssusi.jhuapl.edu/data_products)

417

418 **Acknowledgments**

419 This work was supported by the National Key Research and Development Program of China (Grant
420 NO. 2021YFE0106400 and 2018YFC1407300), the National Natural Science Foundation of China
421 (Grant NO. 42130210, 42120104003, 41704159, and 41831072), and the fund from the Institute of
422 Applied Meteorology. H. G. Yang was supported by Shanghai Science and Technology Innovation
423 Action Plan (No. 21DZ1206100). The author also acknowledges the use of ZHS DPS and HF radar
424 data from the Chinese Meridian Project. SuperDARN is a collection of scientific HF radars funded by
425 the national science funding agencies of Australia, Canada, China, France, Italy, Japan, Norway, South
426 Africa, the United Kingdom, and the United States of America. Thanks to the Johns Hopkins
427 University Applied Physics Laboratory for providing the DMSP/SSUSI auroral FUV data.

428

429 **References**

- 430 André, D., G. J. Sofko, K. Baker, and J. MacDougall (1998), SuperDARN interferometry: Meteor echoes and electron densities
431 from ground scatter, *Journal of Geophysical Research: Space Physics*, 103(A4), 7003-7015.
- 432 Bedey, D. F., and B. J. Watkins (1997), Large-scale transport of metallic ions and the occurrence of thin ion layers in the polar
433 ionosphere, *Journal of Geophysical Research Space Physics*, 102(A5), 9675-9681, doi:10.1029/96ja03825.
- 434 Bishop, R. L., and G. D. Earle (2003), Metallic ion transport associated with midlatitude intermediate layer development, *Journal*
435 *of Geophysical Research*, 108(A1), doi:10.1029/2002ja009411.
- 436 Bourdillon, A., C. Haldoupis, and J. Delloue (1995), High-frequency Doppler radar observations of magnetic aspect sensitive
437 irregularities in the midlatitude E region ionosphere, *Journal of Geophysical Research: Space Physics*, 100(A11),
438 21503-21521, doi:10.1029/95ja01079.
- 439 Bristow, W. A., and B. J. Watkins (1991), Numerical simulation of the formation of thin ionization layers at high latitudes,
440 *Geophysical Research Letters*, 18(3), 404-407, doi:10.1029/90GL02588.
- 441 Chen, X. C., D. S. Han, D. A. Lorentzen, K. Oksavik, J. I. Moen, and L. J. Baddeley (2017), Dynamic Properties of Throat
442 Aurora Revealed by Simultaneous Ground and Satellite Observations, *Journal of Geophysical Research: Space Physics*, 122,
443 3469-3486, doi:10.1002/2016JA023033.
- 444 Chen, X. C., D. A. Lorentzen, J. I. Moen, K. Oksavik, and L. J. Baddeley (2015), Simultaneous ground-based optical and HF

- 445 radar observations of the ionospheric footprint of the open/closed field line boundary along the geomagnetic meridian,
446 *Journal of Geophysical Research: Space Physics*, 120(11), 9859-9874, doi:10.1002/2015JA021481.
- 447 Chen, X. C., D. A. Lorentzen, J. I. Moen, K. Oksavik, L. J. Baddeley, and M. Lester (2016), F region ionosphere effects on the
448 mapping accuracy of SuperDARN HF radar echoes, *Radio Science*, 51(5), 490-506, doi:10.1002/2016rs005957.
- 449 Chen, X., Huang, W., Ban, C., Kosch, M. J., Murphy, D. J., Hu, Z., et al. (2021). Dynamic properties of a sporadic sodium layer
450 revealed by observations over Zhongshan, Antarctica: A case study. *Journal of Geophysical Research: Space Physics*, 126,
451 e2021JA029787. <https://doi.org/10.1029/2021JA029787>.
- 452 Chisham, G., et al. (2007), A decade of the Super Dual Auroral Radar Network (SuperDARN): scientific achievements, new
453 techniques and future directions, *Surveys in Geophysics*, 28(1), 33-109, doi:10.1007/s10712-007-9017-8.
- 454 Chisham, G., and M. P. Freeman (2013), A reassessment of SuperDARN meteor echoes from the upper mesosphere and lower
455 thermosphere, *Journal of Atmospheric and Solar-Terrestrial Physics*, 102, 207-221, doi:10.1016/j.jastp.2013.05.018.
- 456 Chisham, G., and M. Pinnock (2001), Assessing the contamination of SuperDARN global convection maps by non-F-region
457 backscatter, *Annales Geophysicae*, 20(1), 13-28. <https://doi.org/10.5194/angeo-20-13-2002>
- 458 Chisham, G., T. K. Yeoman, and G. J. Sofko (2008), Mapping ionospheric backscatter measured by the SuperDARN HF radars -
459 Part 1: A new empirical virtual height model, *Annales Geophysicae*, 26(4), 823-841, doi:10.5194/angeo-26-823-2008.
- 460 Cousins, E., and S. Shepherd (2010), A dynamical model of high-latitude convection derived from SuperDARN plasma drift
461 measurements, *Journal of Geophysical Research: Space Physics*, 115(A12), doi:10.1029/2010JA016017.
- 462 Gillies, R. G., G. C. Hussey, G. J. Sofko, K. A. McWilliams, R. A. D. Fiori, P. Ponomarenko, and J. P. St.-Maurice (2009),
463 Improvement of SuperDARN velocity measurements by estimating the index of refraction in the scattering region using
464 interferometry, *Journal of Geophysical Research*, 114(A7), A07305, doi:10.1029/2008ja013967.
- 465 Gillies, R. G., G. C. Hussey, G. J. Sofko, P. V. Ponomarenko, and K. A. McWilliams (2011), Improvement of HF coherent radar
466 line-of-sight velocities by estimating the refractive index in the scattering volume using radar frequency shifting, *Journal of*
467 *Geophysical Research*, 116(A1), A01302, doi:10.1029/2010ja016043.
- 468 Gillies, R. G., G. C. Hussey, G. J. Sofko, D. M. Wright, and J. A. Davies (2010), A comparison of EISCAT and SuperDARN
469 F-region measurements with consideration of the refractive index in the scattering volume, *Journal of Geophysical Research*,
470 115(A6), A06319, doi:10.1029/2009ja014694.
- 471 Greenwald, R. A., et al. (1995), DARN/SuperDARN A global view of the dynamics of high-latitude convection, *Space Science*
472 *Reviews*, 71(1-4), 761-796, doi:10.1007/BF00751350.
- 473 Greenwald, R. A., K. B. Baker, R. A. Hutchins, and C. Hanuise (1985), An HF phased-array radar for studying small-scale
474 structure in the high-latitude ionosphere, *Radio Science*, 20(1), 63-79, doi:10.1029/RS020i001p00063.
- 475 Greenwald, R. A., N. Frissell, and S. de Larquier (2017), The importance of elevation angle measurements in HF radar
476 investigations of the ionosphere, *Radio Science*, 52(3), 305-320, doi:10.1002/2016RS006186.
- 477 Haldoupis, C. (2011), A Tutorial Review on Sporadic E Layers, in *Aeronomy of the Earth's Atmosphere and Ionosphere*, edited
478 by M. A. Abdu and D. Pancheva, pp. 381-394, Springer Netherlands, Dordrecht, doi:10.1007/978-94-007-0326-1_29.
- 479 Hall, G. E., J. W. MacDougall, D. R. Moorcroft, J.-P. St.-Maurice, A. H. Manson, and C. E. Meek (1997), Super Dual Auroral
480 Radar Network observations of meteor echoes, *Journal of Geophysical Research: Space Physics*, 102(A7), 14603-14614,
481 doi:10.1029/97ja00517.

- 482 Hanuise, C., J. Villain, D. Gresillon, B. Cabrit, R. Greenwald, and K. Baker (1993), Interpretation Of HF radar ionospheric
483 Doppler spectra by collective wave scattering-theory, *Annales Geophysicae*, 11, 29-39.
- 484 Hardy, D. A., L. K. Schmitt, M. S. Gussenhoven, F. J. Marshall, and H. C. Yeh (1984), Precipitating electron and ion detectors
485 (SSJ/4) for the block 5D/Flights 6-10 DMSP (Defense Meteorological Satellite Program) satellites: Calibration and data
486 presentation *Rep. Air Force Geophysics Laboratory*
- 487 Hu, H., E. Liu, R. Liu, H. Yang, and B. Zhang (2013), Statistical characteristics of ionospheric backscatter observed by
488 SuperDARN Zhongshan radar in Antarctica, *Advances in Polar Science*, 24(1), 19-31, doi:10.3724/SP.J.1085.2013.00019.
- 489 Huang, D., E. Liu, H. Hu, and J. Liu (2018), Algorithm for the estimation of ionosphere parameters from ground scatter echoes of
490 SuperDARN, *Science China Technological Sciences*, 61(11), 1755-1764, doi:10.1007/s11431-017-9178-4.
- 491 Imber, S. M., S. E. Milan, and M. Lester (2013), The Heppner-Maynard Boundary measured by SuperDARN as a proxy for the
492 latitude of the auroral oval, *Journal of Geophysical Research: Space Physics*, 118(2), 685-697, doi:10.1029/2012ja018222.
- 493 Kirkwood, S., and H. Nilsson (2000), High-latitude Sporadic-E and other Thin Layers – the Role of Magnetospheric Electric
494 Fields, *Space Science Reviews*, 91(3), 579-613, doi:10.1023/A:1005241931650.
- 495 Kirkwood, S., and U. von Zahn (1993), Formation Mechanisms for Low-Altitude Es and Their Relationship with Neutral Fe
496 Layers: Results from the METAL Campaign, *Journal of Geophysical Research*, 98(A12), 21549-21561,
497 doi:10.1029/93ja02644.
- 498 Kirkwood, S., and U. V. Zahn (1991), On the role of auroral electric fields in the formation of low altitude sporadic-E and sudden
499 sodium layers, *Journal of Atmospheric & Terrestrial Physics*, 53(5), 389-407, doi:10.1016/0021-9169(91)90034-5.
- 500 Koustov, A. V., S. Ullrich, P. V. Ponomarenko, R. G. Gillies, D. R. Themens, and N. Nishitani (2020), Comparison of
501 SuperDARN peak electron density estimates based on elevation angle measurements to ionosonde and incoherent scatter
502 radar measurements, *Earth Planets Space*, 72(1), 43, doi:10.1186/s40623-020-01170-w.
- 503 Liu, E. X., H. Q. Hu, K. Hosokawa, R. Y. Liu, Z. S. Wu, and Z. Y. Xing (2013), First observations of polar mesosphere summer
504 echoes by SuperDARN Zhongshan radar, *Journal of Atmospheric and Solar-Terrestrial Physics*, 104, 39-44,
505 doi:10.1016/j.jastp.2013.07.011.
- 506 Liu, J. J., H. Q. Hu, D. S. Han, T. Araki, Z. J. Hu, Q. H. Zhang, H. G. Yang, N. Sato, A. S. Yukimatu, and Y. Ebihara (2011),
507 Decrease of auroral intensity associated with reversal of plasma convection in response to an interplanetary shock as
508 observed over Zhongshan station in Antarctica, *Journal of Geophysical Research: Space Physics*, 116, A03210,
509 doi:10.1029/2010ja016156.
- 510 MacDougall, J. W., and P. T. Jayachandran (2005), Sporadic E at cusp latitudes, *Journal of Atmospheric and Solar-Terrestrial*
511 *Physics*, 67(15), 1419-1426, doi:10.1016/j.jastp.2005.07.011.
- 512 Macdougall, J. W., P. T. Jayachandran, and J. M. C. Plane (2000a), Polar cap Sporadic-E: Part 1, Observations, *Journal of*
513 *Atmospheric and Solar-Terrestrial Physics*, 62(13), 1155-1167, doi:10.1016/S1364-6826(00)00093-6.
- 514 Macdougall, J. W., J. M. C. Plane, and P. T. Jayachandran (2000b), Polar cap Sporadic-E: part 2, modeling, *Journal of*
515 *Atmospheric and Solar-Terrestrial Physics*, 62(13), 1169-1176, doi:10.1016/S1364-6826(00)00092-4.
- 516 Mathews, J. D., D. W. Machuga, and Q. Zhou (2001), Evidence for electrodynamic linkages between spread-F, ion rain, the
517 intermediate layer, and sporadic E: results from observations and simulations, *Journal of Atmospheric and Solar-Terrestrial*
518 *Physics*, 63(14), 1529-1543, doi:https://doi.org/10.1016/S1364-6826(01)00034-7.

- 519 Milan, S., T. Jones, T. Robinson, E. Thomas, and T. Yeoman (1997), Interferometric evidence for the observation of ground
520 backscatter originating behind the CUTLASS coherent HF radars, *Annales Geophysicae*, 15(1), 29-39.
- 521 Milan, S. E., and M. Lester (2001), A classification of spectral populations observed in HF radar backscatter from the E region
522 auroral electrojets, *Annales Geophysicae*, 19(2), 189-204, doi:10.5194/angeo-19-189-2001.
- 523 Nygrén, T., L. Jalonon, J. Oksman, and T. Turunen (1984), The role of electric field and neutral wind direction in the formation of
524 sporadic E-layers, *Journal of Atmospheric and Terrestrial Physics*, 46(4), 373-381, doi:10.1016/0021-9169(84)90122-3.
- 525 Nygrén, T., M. Voiculescu, and A. T. Aikio (2008), The role of electric field and neutral wind in the generation of polar cap
526 sporadic E, *Annales Geophysicae*, 26, 3757-3763, doi:10.5194/angeo-26-3757-2008.
- 527 Ogawa, T., N. Nishitani, Y. Otsuka, K. Shiokawa, T. Tsugawa, and K. Hosokawa (2009), Medium-scale traveling ionospheric
528 disturbances observed with the SuperDARN Hokkaido radar, all-sky imager, and GPS network and their relation to
529 concurrent sporadic E irregularities, *Journal of Geophysical Research: Space Physics*, 114(A3), A03316,
530 doi:10.1029/2008ja013893.
- 531 Parkinson, M. L., P. L. Dyson, D. P. Monselesan, and R. J. Morris (1998), On the role of electric field direction in the formation
532 of sporadic E -layers in the southern polar cap ionosphere, *Journal of Atmospheric and Solar-Terrestrial Physics*, 60(4),
533 471-491, doi:10.1016/0021-9169(84)90122-3.
- 534 Paxton, L. J., D. Morrison, Y. Zhang, H. Kil, B. Wolven, B. S. Ogorzalek, D. C. Humm, and C.-I. Meng (2002), Validation of
535 remote sensing products produced by the Special Sensor Ultraviolet Scanning Imager (SSUSI): a far UV-imaging
536 spectrograph on DMSP F-16, SPIE.
- 537 Plane, J. M. C. (2003), Atmospheric chemistry of meteoric metals, *Chemical Reviews*, 103(12), 4963-4984,
538 doi:10.1021/cr0205309.
- 539 Ponomarenko, P., N. Nishitani, A. V. Oinats, T. Tsuya, and J.-P. St.-Maurice (2015), Application of ground scatter returns for
540 calibration of HF interferometry data, *Earth, Planets and Space*, 67(1), 138, doi:10.1186/s40623-015-0310-3.
- 541 Ponomarenko, P., J. P. St.-Maurice, and K. A. McWilliams (2018), Calibrating HF Radar Elevation Angle Measurements Using E
542 Layer Backscatter Echoes, *Radio Science*, 53(11), 1438-1449, doi:10.1029/2018rs006638.
- 543 Ponomarenko, P. V., A. V. Koustov, J. P. St.-Maurice, and J. Wiid (2011), Monitoring the F-region peak electron density using HF
544 backscatter interferometry, *Geophysical Research Letters*, 38(21), L21102, doi:10.1029/2011gl049675.
- 545 Reinisch, B. W., I. A. Galkin, G. M. Khmyrov, A. V. Kozlov, K. Bibl, I. A. Lisysyan, G. P. Cheney, X. Huang, D. F. Kitrosser, and
546 V. V. Paznukhov (2009), New Digisonde for research and monitoring applications, *Radio Science*, 44(1), 1-15,
547 doi:10.1029/2008RS004115.
- 548 Ruohoniemi, J. M., and K. B. Baker (1998), Large-scale imaging of high-latitude convection with Super Dual Auroral Radar
549 Network HF radar observations, *Journal of Geophysical Research*, 103(A9), 20797, doi:10.1029/98ja01288.
- 550 Shepherd, S. G. (2017), Elevation angle determination for SuperDARN HF radar layouts, *Radio Science*, 52(8), 938-950,
551 doi:10.1002/2017rs006348.
- 552 Wang, Z., Lu, J., Hu, H., Liu, J., Hu, Z., Wang, M., et al. (2022). HMB Variations Measured by SuperDARN During the
553 Extremely Radial IMFs: Is the coupling function applicable in radial IMF? *Journal of Geophysical Research: Space
554 Physics*, 127, e2021JA029589. <https://doi.org/10.1029/2021JA029589>
- 555 Yeoman, T. K., G. Chisham, L. J. Baddeley, R. S. Dhillon, T. J. T. Karhunen, T. R. Robinson, A. Senior, and D. M. Wright (2008),

- 556 Mapping ionospheric backscatter measured by the SuperDARN HF radars - Part 2: Assessing SuperDARN virtual height
557 models, *Annales Geophysicae*, 26(4), 843-852, doi:10.5194/angeo-26-843-2008.
- 558 Yeoman, T. K., D. M. Wright, A. J. Stocker, and T. B. Jones (2001), An evaluation of range accuracy in the Super Dual Auroral
559 Radar Network over-the-horizon HF radar systems, *Radio Science*, 36(4), 801-813, doi:10.1029/2000rs002558.
560

561 **Figure Captions:**

562 **Figure 1:** (left) Simulation of ray tracing for the SuperDARN radar at Zhongshan Station. Rays are traced for
563 elevation angles from 5° to 55° and group ranges from 0 to 2400 km for beam 7 and at a frequency of 10 MHz.
564 The electron density profiles are provided by the IRI Model for 14 November 2019 at 18:00 UT (\sim 23:10 LT). The
565 red ray represents the last high-angle that can be refracted back to the ground. While the black lines with a 30°
566 off-zenith angle show the maximum field-of-view of Digisonde. (right) An example of the SuperDARN 2-D
567 convection map plotted at 16:00–16:02 UT in MLAT–MLT coordinates in the southern hemisphere. Noon/dawn is
568 toward the top/right. The ionospheric flow measurements are indicated by a dot with a line attached, the color of
569 which gives the magnitude of flow and the line shows the direction. The + and \times symbols are the
570 maximum/minimum of the electrostatic potential distribution, where the positive equipotential contours are shown
571 by black dashed lines and negative contours are solid. The red closed curve shows the Heppner-Maynard Boundary.
572 The field-of-view of the ZHS HF radar is shown by a dashed fan, and directions of beams 13, 7, and 1 are also
573 denoted. The red star represents the ZHS station.

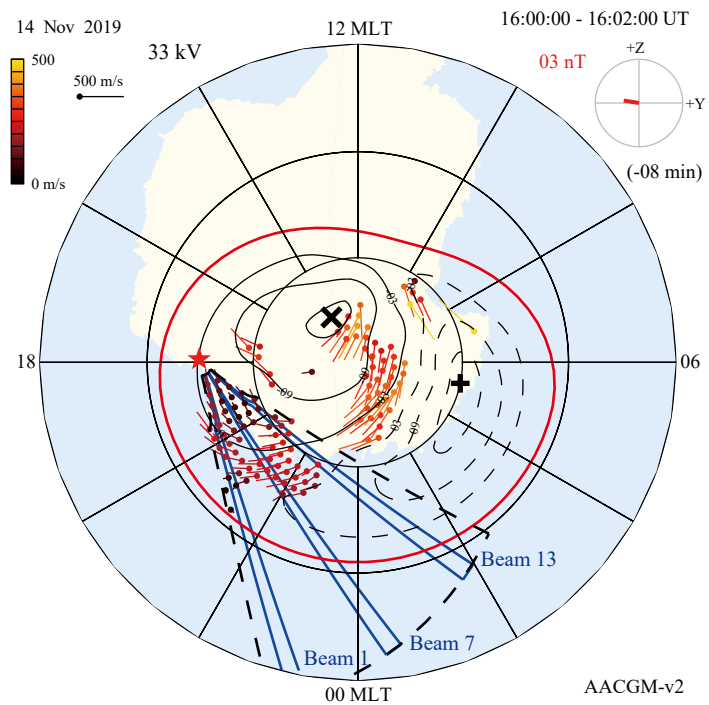
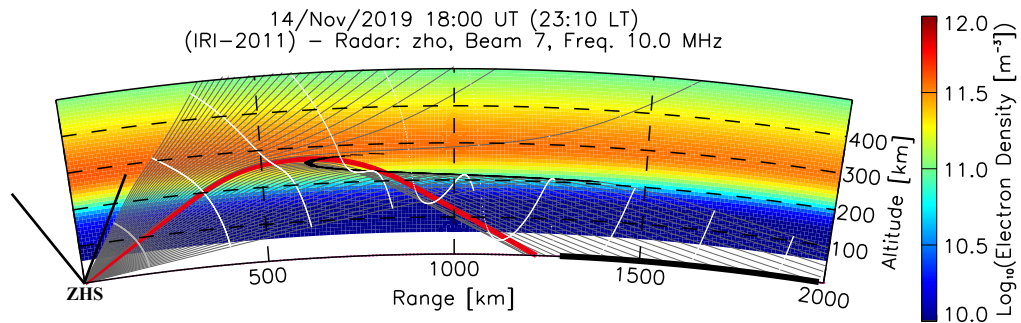
574 **Figure 2:** Sequence of 7.5 min interval ionograms for \sim 15:37 – 17:30 UT on 14 November 2019, which show the
575 initial formation and subsequent evolution of the Es layer.

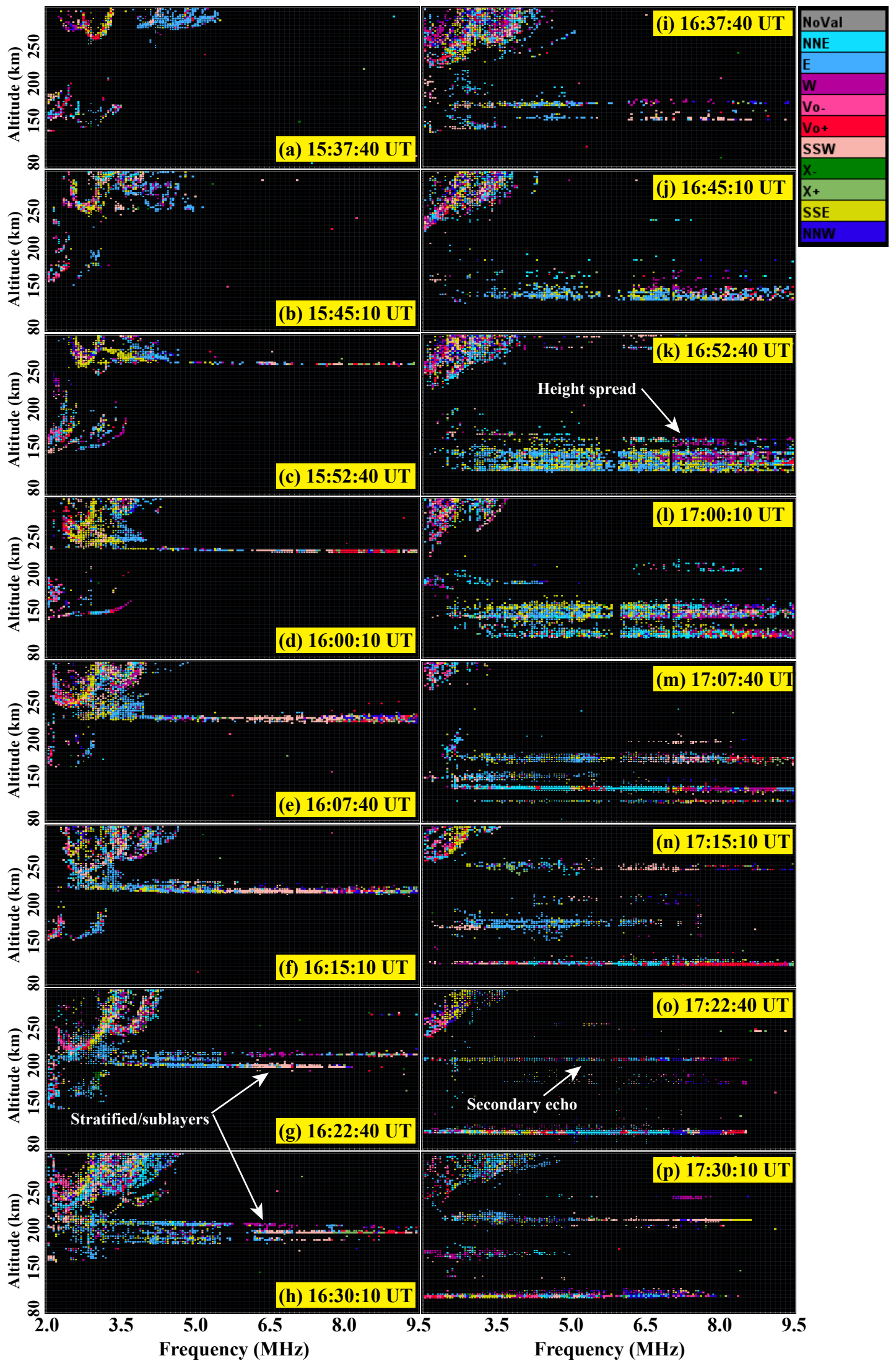
576 **Figure 3:** Height-time-amplitude of the echoes from the DPS at sounding frequency ranges of 3.5–3.75 MHz (a),
577 6.5–6.75 MHz (b), and larger than 2 MHz (c).

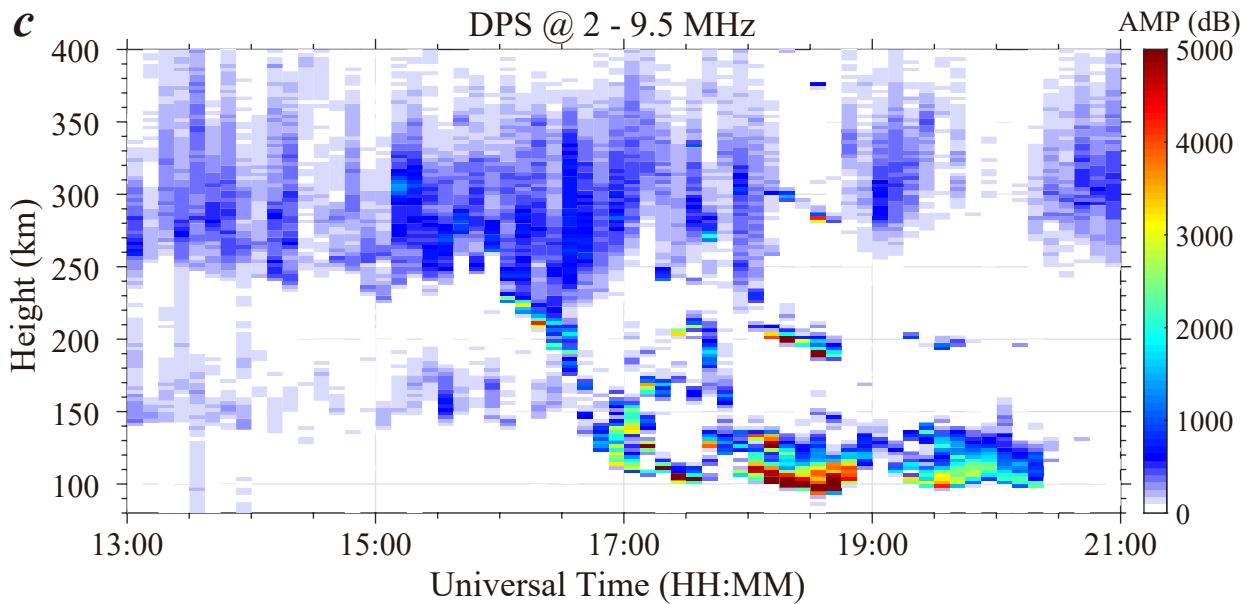
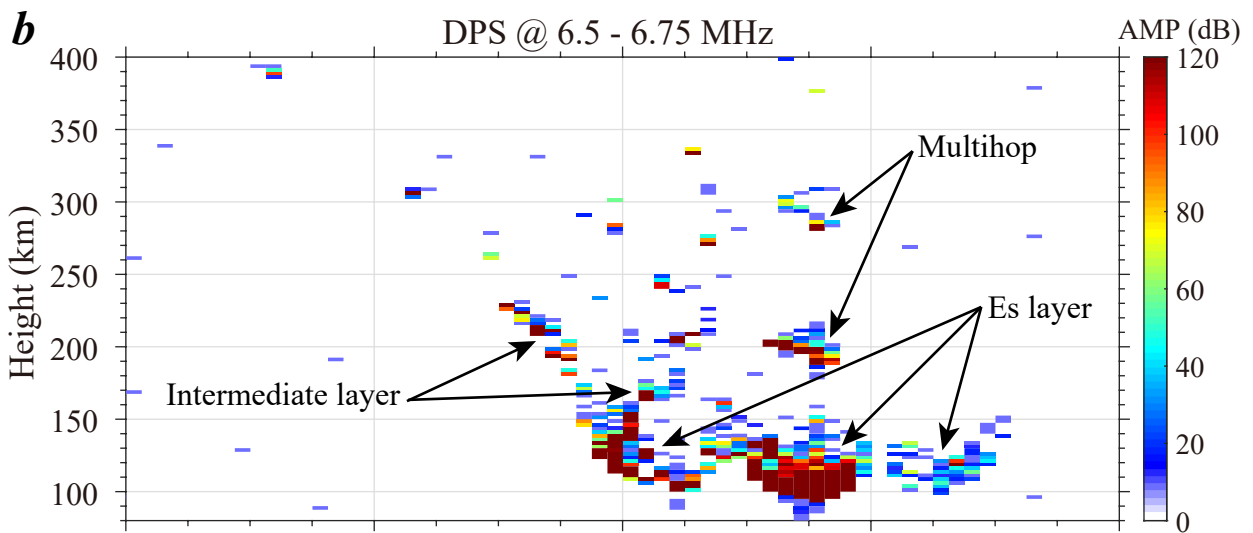
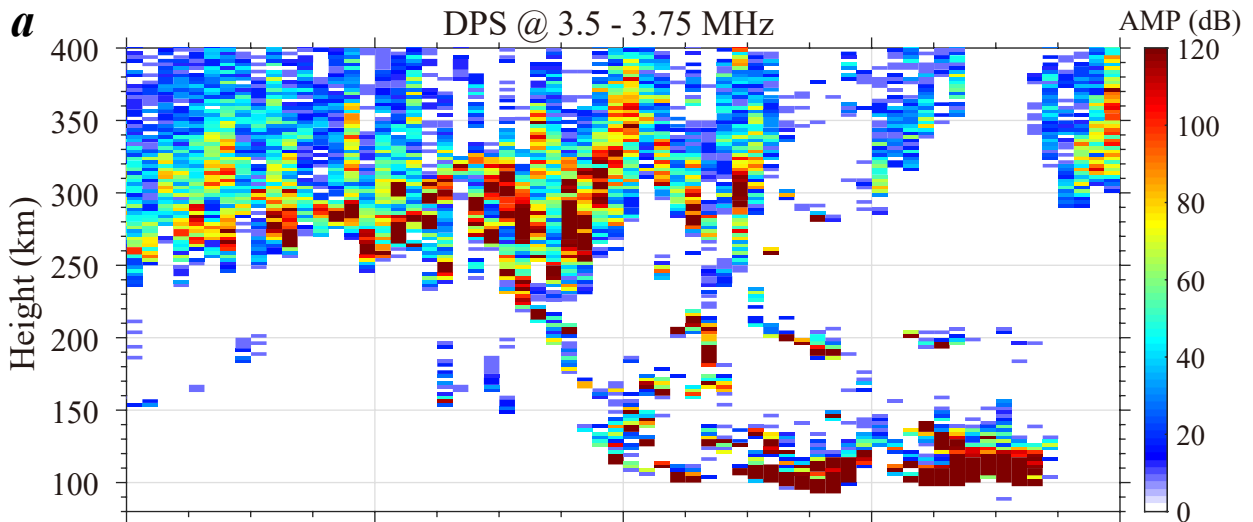
578 **Figure 4:** Range versus time plots from the ZHS HF radar ionospheric line-of-sight velocities for beams 13(a),
579 7(b), and 1(c). Positive (negative) line-of-sight velocities are toward (away from) the radar. The black lines in the
580 bottom panel indicate the Es layer-related ionospheric echoes, while the arrows indicate the shear flow resulting
581 from the auroral arc.

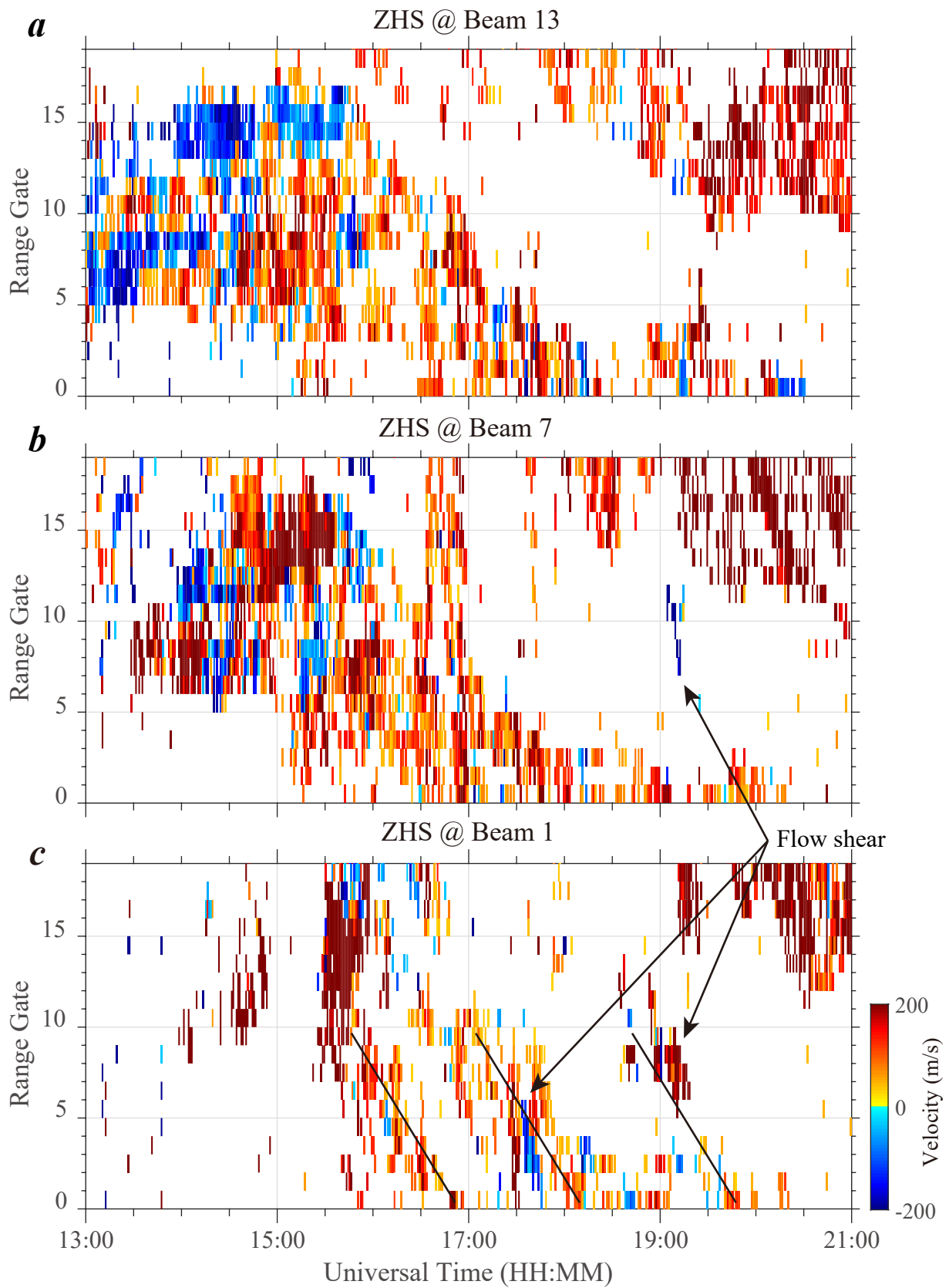
582 **Figure 5:** Four plots of SuperDARN 2-D convection maps overlaid with images recorded by the SSUSI instrument
583 in the LBHS wavelength when the DMSP F17 crossing.

584 **Figure 6:** (a) Histograms of the elevation angle values from HF radar beams 13 (green), 7 (blue), and 1 (black) at
585 range gates of 0–18, during the time interval of 13:00–21:00 UT. (b–d) Scatterplots of backscatter altitude versus
586 time for three radar beams. The overlaid red histogram and asterisks represent the lowest virtual height of the Es
587 layer and its stratified/sublayers observed by the DPS. The black dots show the peak height of the F2 layer scaled
588 from ionograms.

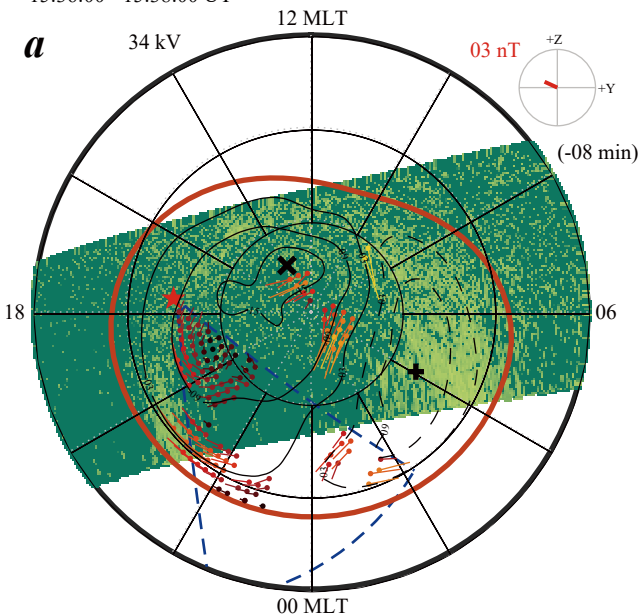




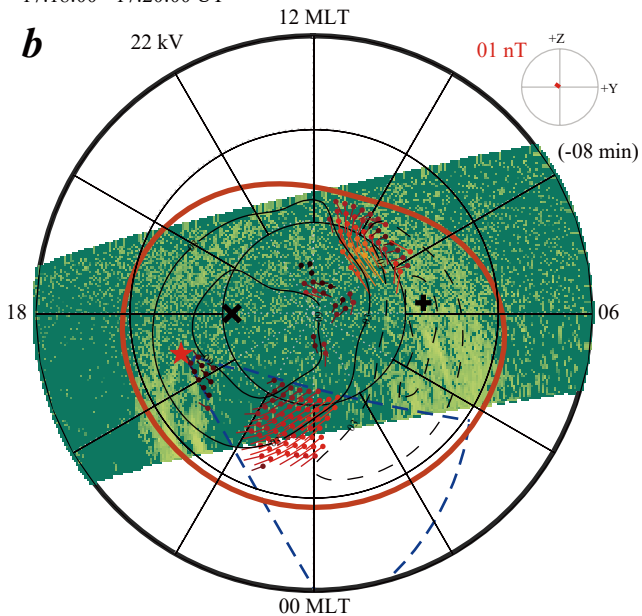




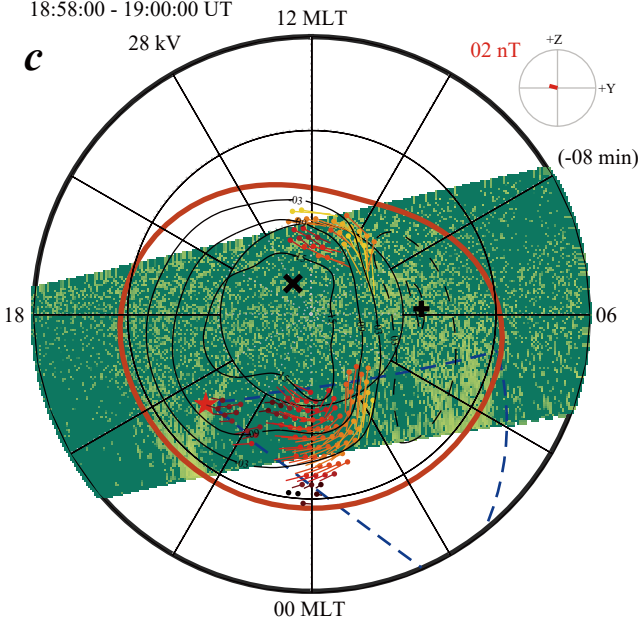
DMSP F17
15:36:00 - 15:38:00 UT



DMSP F17
17:18:00 - 17:20:00 UT



DMSP F17
18:58:00 - 19:00:00 UT



DMSP F17
20:38:00 - 20:40:00 UT

