



Central Limit Theorem and Large Deviation Principle for Continuous Time Open Quantum Walks

Hugo Bringuier

► **To cite this version:**

Hugo Bringuier. Central Limit Theorem and Large Deviation Principle for Continuous Time Open Quantum Walks. 2016. <hal-01376995>

HAL Id: hal-01376995

<https://hal.archives-ouvertes.fr/hal-01376995>

Submitted on 6 Oct 2016

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Central Limit Theorem and Large Deviation Principle for Continuous Time Open Quantum Walks

Hugo BRINGUIER

Institut de Mathématiques de Toulouse ; UMR5219,
Université de Toulouse ; CNRS,
UPS IMT, F-31062 Toulouse Cedex 9, France

October 6, 2016

Abstract

Open Quantum Walks (OQWs), originally introduced in [2], are quantum generalizations of classical Markov chains. Recently, natural continuous time models of OQW have been developed in [24]. These models, called Continuous Time Open Quantum Walks (CTOQWs), appear as natural continuous time limits of discrete time OQWs. In particular they are quantum extensions of continuous time Markov chains. This article is devoted to the study of homogeneous CTOQW on \mathbb{Z}^d . We focus namely on their associated quantum trajectories which allow us to prove a Central Limit Theorem for the "position" of the walker as well as a Large Deviation Principle.

1 Introduction

Open Quantum Walks concern evolution on lattices driven by quantum operations. They describe namely Markovian dynamics influenced by internal degrees of freedom. They have been introduced originally by [2] (see also [15]). These OQWs are promising tools to model physical problems, especially in computer science (see [27]). They can also model a variety of phenomena, as energy transfer in biological systems ([21]).

Continuous time models have been developed as natural continuous time limits of discrete time models [24, 5]. In particular in [5], a natural extension of Brownian

motion called Open Quantum Brownian Motion has been constructed. In this article, we focus on the continuous time open quantum walks (CTOQWs) model presented in [24]. More precisely, we focus on CTOQWs on \mathbb{Z}^d . Briefly speaking, CTOQWs on \mathbb{Z}^d concern the evolution of density operators of the form

$$\mu = \sum_{i \in \mathbb{Z}^d} \rho(i) \otimes |i\rangle\langle i| \in \mathcal{H} \otimes \mathbb{C}^{\mathbb{Z}^d} \quad (1)$$

where the " \mathbb{Z}^d -component" represents the "position" of the walker and \mathcal{H} is a Hilbert space describing the internal degrees of freedom. In particular, if \mathcal{D} denotes the set of density operators of the form (1), CTOQWs are described by a semigroup $\{\phi_t\}$ such that, ϕ_t preserves \mathcal{D} for all $t \geq 0$.

In the context of quantum walks, one is mainly interested in the position of the walker. At time 0, starting with density matrix in \mathcal{D} as (1), the quantum measurement of the "position" gives rise to a probability distribution q_0 on \mathbb{Z}^d , such that, for all $i \in \mathbb{Z}^d$,

$$q_0(i) = \mathbb{P}(\text{"that the walker is in } i\text{"}) = \text{Tr}(\rho(i)) .$$

As well, after evolution, if

$$\mu_t = \phi_t(\mu) = \sum_{i \in \mathbb{Z}^d} \rho^{(t)}(i) \otimes |i\rangle\langle i|$$

then

$$q_t(i) = \mathbb{P}(\text{"that the walker, at time } t, \text{ is in } i\text{"}) = \text{Tr}(\rho^{(t)}(i)) .$$

In [24], it has been shown that usual classical continuous time Markov chains are particular cases of CTOQWs. In particular one can easily construct models where the distribution q_t corresponds to the one of a classical continuous time Markov chain. Contrary to continuous time Markov chains, the distribution q_t of CTOQWs cannot be in general recovered by the knowledge of the initial distribution q_0 . One needs to have access to the full knowledge of the initial state μ . In this sense, this justifies the name *quantum walks*.

Our models of continuous time quantum walks are rather different from the usual models of unitary quantum walks. An essential difference concerns the large time behaviour of the corresponding distribution q_t . Let Q_t be a random variable of law q_t , in the unitary quantum walk theory it has been shown that (Q_t) satisfies a Central Limit Theorem of the type

$$\frac{Q_t}{t} \xrightarrow[t \rightarrow \infty]{} \tilde{Q},$$

where \tilde{Q} has distribution

$$\frac{dx}{\pi\sqrt{1-x^2}}.$$

Note that such behaviour is not usual in classical probability where usually one expects speed in \sqrt{t} and Gaussian law as limit in the Central Limit Theorem (CLT). In our context, the distributions $(q_t)_{t \geq 0}$ express a rather classical behaviour in large time in the sense that a more usual CLT holds. In particular this paper is devoted to show that for CTOQWs one has the following weak convergence

$$\frac{Q_t - m}{\sqrt{t}} \xrightarrow[t \rightarrow \infty]{} \mathcal{N}(0, \sigma^2),$$

where $\mathcal{N}(0, \sigma^2)$ denotes usual Gaussian law. Such phenomena have also been observed in the discrete setting of OQWs [1]. A key point to show this result is the use of the quantum trajectories associated to the CTOQWs. In general, quantum trajectories describe evolutions of quantum system undergoing indirect measurements (see [3] for an introduction). In the context of CTOQWs, quantum trajectories describe the evolution of the states undergoing indirect measurements of the position of the walker. In particular these quantum trajectories appear as solution of jump-type stochastic differential equations called stochastic master equations (see [24] for link between discrete and continuous time models in the context of OQW, one can also consult [5] for such an approach in the context of Open Quantum Brownian Motion). In the physic literature, note that such models appear also naturally in order to describe non-Markovian evolutions. They are called *non-Markov generalization of Lindblad equations* (see [6, 25, 4]).

After establishing the CLT, our next goal is to investigate a Large Deviation Principle (LDP) for the position of the walker. In particular under additional assumptions, one can apply the Gärtner-Ellis Theorem in order to obtain the final result (one can consult [7] for a similar result for discrete time OQWs).

The article is structured as follows. In Section 2, we present the model of CTOQWs on \mathbb{Z}^d . Next we develop the theory of quantum trajectories which describe the continuous measurement of the position. In Section 3, we present the Central Limit Theorem. Section 4 is devoted to the Large Deviation Principle (LDP). Finally in Section 5, we present some examples which illustrate the CLT and the LDP.

2 Continuous Time Open Quantum Walks

2.1 Main setup

The models of Continuous Time Open Quantum Walks have been formalized in [24]. They arise as continuous limits of discrete time OQWs (we do not recall the discrete time models and we refer to [2]). These limits processes are described by

particular types of *Lindblad master equations*. Originally, these equations appear in the "non-Markovian generalization of Lindblad theory" from Breuer [6]. In this article, we focus on nearest neighbors, homogeneous CTOQWs on \mathbb{Z}^d .

In the sequel, \mathcal{H} denotes a finite dimensional Hilbert space and $\mathcal{S}_{\mathcal{H}}$ denotes the space of density matrix on \mathcal{H} :

$$\mathcal{S}_{\mathcal{H}} = \{\rho \in \mathcal{B}(\mathcal{H}) \mid \rho^* = \rho, \rho \geq 0, \text{Tr}(\rho) = 1\}.$$

We put $\mathcal{K}_d = \mathcal{H} \otimes \mathbb{C}^{\mathbb{Z}^d}$ where $\mathbb{C}^{\mathbb{Z}^d}$ stands for the position of a particle while \mathcal{H} corresponds to the internal degree of freedom of this particle. We consider the canonical basis $\{e_1, \dots, e_d\}$ of \mathbb{Z}^d , we set $e_0 = 0_d$ and $e_{d+r} = -e_r$ for all $r \in \{1, \dots, d\}$. The canonical basis of $\mathbb{C}^{\mathbb{Z}^d}$ is denoted by $(|i\rangle)_{i \in \mathbb{Z}^d}$.

As announced we focus on particular diagonal density matrices of \mathcal{K}_d :

$$\mathcal{D} = \left\{ \mu \in \mathcal{B}(\mathcal{K}_d), \mu = \sum_{i \in \mathbb{Z}^d} \rho(i) \otimes |i\rangle\langle i|, \rho(i) \geq 0, \sum_{i \in \mathbb{Z}^d} \text{tr}(\rho(i)) = 1 \right\}.$$

In the sequel we shall consider evolutions on \mathcal{K}_d which preserve \mathcal{D} . To this end we consider a family of operators $\{D_r\}_{r=1, \dots, 2d}$ on $\mathcal{B}(\mathcal{H})$ and we define the operators $\{B_i^r\}_{r=1, \dots, 2d}$ on $\mathcal{B}(\mathcal{K}_d)$ such that $B_i^r = D_r \otimes |i + e_r\rangle\langle i|$.

Now as announced the CTOQWs are generated by particular Lindblad master equations. Let \mathcal{M}_c the following Lindblad operator on $\mathcal{H} \otimes \mathbb{C}^{\mathbb{Z}^d}$,

$$\begin{aligned} \mathcal{M}_c : \mathcal{B}(\mathcal{H} \otimes \mathbb{C}^{\mathbb{Z}^d}) &\rightarrow \mathcal{B}(\mathcal{H} \otimes \mathbb{C}^{\mathbb{Z}^d}) \\ \mu &\mapsto -i[H \otimes I, \mu] + \sum_{i \in \mathbb{Z}^d} \sum_{r=1}^{2d} \left(B_i^r \mu B_i^{r*} - \frac{1}{2} \{B_i^{r*} B_i^r, \mu\} \right) \end{aligned}$$

where H is a self-adjoint operator on \mathcal{H} which is called the Hamiltonian.

Let us introduce the operator

$$D_0 = -iH - \frac{1}{2} \sum_{r=1}^{2d} D_r^* D_r.$$

The next computation shows that \mathcal{M}_c preserves the set \mathcal{D} .

$$\begin{aligned}
 \mathcal{M}_c(\mu) &= \sum_{i \in \mathbb{Z}^d} \left[-i[H, \rho(i)] \otimes |i\rangle\langle i| + \sum_{r=1}^{2d} D_r \rho(i) D_r^* \otimes |i + e_r\rangle\langle i + e_r| \right. \\
 &\quad \left. - \frac{1}{2} \sum_{r=1}^{2d} \{D_r^* D_r, \rho(i)\} \otimes |i\rangle\langle i| \right] \\
 &= \sum_{i \in \mathbb{Z}^d} \left[(D_0 \rho(i) + \rho(i) D_0^*) \otimes |i\rangle\langle i| + \sum_{r=1}^{2d} D_r \rho(i) D_r^* \otimes |i + e_r\rangle\langle i + e_r| \right] \\
 &= \sum_{i \in \mathbb{Z}^d} \left(D_0 \rho(i) + \rho(i) D_0^* + \sum_{r=1}^{2d} D_r \rho(i - e_r) D_r^* \right) \otimes |i\rangle\langle i|,
 \end{aligned}$$

for all $\mu = \sum_{i \in \mathbb{Z}^d} \rho(i) \otimes |i\rangle\langle i|$.

The following proposition describes precisely our model of CTOQWs.

Proposition 2.1.1. [24] *Let $\mu^{(0)} = \sum_{i \in \mathbb{Z}^d} \rho^{(0)}(i) \otimes |i\rangle\langle i|$, the equation*

$$\frac{d}{dt} \mu^{(t)} = \mathcal{M}_c(\mu^{(t)}), \tag{2}$$

with initial condition $\mu^{(0)}$ admits a unique solution $(\mu^{(t)})_{t \geq 0}$ with values in \mathcal{D} .

More precisely, $\mu^{(t)}$ is of the form $\mu^{(t)} = \sum_{i \in \mathbb{Z}^d} \rho^{(t)}(i) \otimes |i\rangle\langle i|$ such that:

$$\frac{d}{dt} \rho^{(t)}(i) = D_0 \rho^{(t)}(i) + \rho^{(t)}(i) D_0^* + \sum_{r=1}^{2d} D_r \rho^{(t)}(i - e_r) D_r^*,$$

for all $i \in \mathbb{Z}^d$.

Definition 2.1.2. The evolution (2) is called a Continuous Time Open Quantum Walk on \mathbb{Z}^d .

This definition is justified by the following. The operator B_i^r transcribes the idea that the particle localized in $|i\rangle$ can only jump to one of its nearest neighbors $|i + e_r\rangle$, and in this case, the transformation on \mathcal{H} is governed by D_r . In the case the particle stands still, the evolution on \mathcal{H} is governed by D_0 . It is the exact analogue of the usual OQWs for continuous time evolutions. An interesting fact has been pointed out in [24], usual continuous time classical Markov chains can be realized within this setup.

Now let us describe the probability distributions associated to CTOQWs.

Definition 2.1.3. Let $\mu^{(0)} = \sum_{i \in \mathbb{Z}^d} \rho^{(0)}(i) \otimes |i\rangle\langle i|$. Let $\mu^{(t)} = \sum_{i \in \mathbb{Z}^d} \rho^{(t)}(i) \otimes |i\rangle\langle i|$ be the solution of the equation

$$\frac{d}{dt}\mu^{(t)} = \mathcal{M}_c(\mu^{(t)}).$$

We define

$$q_t(i) = \text{Tr}[\mu^{(t)} I \otimes |i\rangle\langle i|] = \text{Tr}[\rho^{(t)}(i)] \quad (3)$$

and we denote Q_t the random variable on \mathbb{Z}^d of law q_t , that is

$$\mathbb{P}[Q_t = i] = q_t(i),$$

for all $i \in \mathbb{Z}^d$.

As we can see in Section 3 and as it was announced in the introduction, the shape of q_t seems to converge to Gaussian shape. This is exactly the result pointed out by the CLT in Section 3. In order to prove this, we shall need the theory of quantum trajectories for CTOQWs.

2.2 Quantum trajectories

As in the discrete case, quantum trajectories are essential tools for showing the CLT and the LDP. The description of quantum trajectories is less straightforward than the one in OQWs. It makes use of stochastic differential equations driven by jump processes. We refer to [24] for the justification of the below description and the link between discrete and continuous time models. One can also consult [6] where general indirect measurements for non-markovian generalization of Lindblad equations have been developed.

Proposition 2.2.1. *Let $\mu^{(0)} = \sum_{i \in \mathbb{Z}^d} \rho^{(0)}(i) \otimes |i\rangle\langle i|$ be an initial state on $\mathcal{H} \otimes \mathbb{C}^{\mathbb{Z}^d}$.*

The quantum trajectory describing the indirect measurement of the position of the CTOQWs led by \mathcal{M}_c is modeled by a Markov process $(\omega^{(t)} = \rho_t \otimes |X_t\rangle\langle X_t|)_{t \geq 0}$.

This Markov process is valued in the set

$$\mathcal{P} = \{\rho \otimes |i\rangle\langle i|, \rho \in \mathcal{S}_{\mathcal{H}}, i \in \mathbb{Z}^d\}$$

such that

$$\omega^{(0)} = \frac{\rho^{(0)}(i)}{\text{Tr}(\rho^{(0)}(i))} \otimes |i\rangle\langle i| \text{ with probability } \text{Tr}(\rho^{(0)}(i))$$

and such that the following differential equation is satisfied:

$$\begin{aligned} \omega^{(t)} = & \omega^{(0)} + \int_0^t \left(D_0 \rho_{s-} + \rho_{s-} D_0^* - \rho_{s-} \text{Tr}(D_0 \rho_{s-} + \rho_{s-} D_0^*) \right) \otimes |X_{s-}\rangle \langle X_{s-}| \, ds \\ & + \sum_{r=1}^{2d} \int_0^t \int_{\mathbb{R}} \left(\frac{D_r \rho_{s-} D_r^*}{\text{Tr}(D_r \rho_{s-} D_r^*)} \otimes |X_{s-} + e_r\rangle \langle X_{s-} + e_r| \right. \\ & \left. - \rho_{s-} \otimes |X_{s-}\rangle \langle X_{s-}| \right) \mathbb{1}_{0 < y < \text{Tr}(D_r \rho_{s-} D_r^*)} N^r(dy, ds) \end{aligned} \quad (4)$$

where $\{N^r\}_{r \in \{1, \dots, 2d\}}$ are independent Poisson point processes on \mathbb{R}^2 .

In particular the Markov process $(\rho_t, X_t)_{t \geq 0}$ is valued in $\mathcal{S}_{\mathcal{H}} \times \mathbb{Z}^d$ and satisfies

$$\begin{aligned} d\rho_s = & \left(D_0 \rho_{s-} + \rho_{s-} D_0^* - \rho_{s-} \text{Tr}(D_0 \rho_{s-} + \rho_{s-} D_0^*) \right) ds \\ & + \sum_{r=1}^{2d} \int_{y \in \mathbb{R}} \left(\frac{D_r \rho_{s-} D_r^*}{\text{Tr}(D_r \rho_{s-} D_r^*)} - \rho_{s-} \right) \mathbb{1}_{0 < y < \text{Tr}(D_r \rho_{s-} D_r^*)} N^r(dy, ds), \end{aligned} \quad (5)$$

$$dX_s = \sum_{r=1}^{2d} \int_{y \in \mathbb{R}} e_r \mathbb{1}_{0 < y < \text{Tr}(D_r \rho_{s-} D_r^*)} N^r(dy, ds) \quad (6)$$

and $(\rho_0, X_0) = \left(\frac{\rho^{(0)}(i)}{\text{Tr}(\rho^{(0)}(i))}, i \right)$ with probability $\text{Tr}(\rho^{(0)}(i))$.

Remark: The second expression of the description of quantum trajectories is the exact continuous time analogue of the one described in [1] for OQWs. Let us briefly explain how the quantum trajectories evolve in time. To this end we introduce:

$$\forall r \in \{1, \dots, 2d\}, \tilde{N}^r(t) = \int_0^t \int_{\mathbb{R}} \mathbb{1}_{0 < y < \text{Tr}(D_r \rho_{s-} D_r^*)} N^r(dy, ds). \quad (7)$$

The processes \tilde{N}^r are Poisson processes with intensity $\int_0^t \text{Tr}(D_r \rho_{s-} D_r^*) ds$. In particular the processes

$$\tilde{N}^r(t) - \int_0^t \text{Tr}(D_r \rho_{s-} D_r^*) ds$$

are martingales with respect to the filtration induced by $(\rho_t, X_t)_{t \geq 0}$. The evolution described by (4) is deterministic and interrupted by jumps occurring at random time, it is typically a Piecewise Deterministic Markov Process. The jumps are generated by the Poisson processes (7). As we can check from Eq. (4), if $\omega^{(0)} = \rho \otimes |i\rangle \langle i|$ for some $\rho \in \mathcal{S}_{\mathcal{H}}$ and $i \in \mathbb{Z}^d$ (that is $|X_0\rangle = |i\rangle$), the deterministic evolution let the position unchanged until a jump occurs. Since the Poisson processes N^r

are independent, only one Poisson process is involved. If T denotes the time of the first jump and assume the process N^r is involved, the internal degree of freedom is updated by $\rho_T = \frac{D_r \rho_{T-} D_r^*}{\text{Tr}(D_r \rho_{T-} D_r^*)}$ and the position is changed and becomes $|i + e_r\rangle$. This means that the particle has jumped from the position $|i\rangle$ to the position $|i + e_r\rangle$. In other words we have $|X_t\rangle = |i\rangle$, for all $0 \leq t \leq T-$ and $|X_T\rangle = |i + e_r\rangle$. Next, the deterministic evolution starts again with the new initial condition $\rho_T \otimes |i + e_r\rangle\langle i + e_r|$ until a new jump occur and so on.

The following result allows us to make the connection between CTOQWs and their associated quantum trajectories.

Proposition 2.2.2. *Let $\mu^{(t)}$ the OQW defined in Proposition 2.1.1 and $\omega^{(t)}$ the associated quantum trajectory defined in Proposition 2.2.1. Then we have*

$$\forall t \geq 0, \mathbb{E}(\omega^{(t)}) = \mu^{(t)}.$$

Moreover, for all $t \geq 0$, the random variables X_t and Q_t have the same distributions q_t .

Proof. The first part is proved in [24]. For the second part, let ϕ a bounded continuous map on \mathbb{Z}^d , we get:

$$\begin{aligned} \mathbb{E}(\phi(Q_t)) &= \sum_{i \in \mathbb{Z}^d} \phi(i) \text{Tr}(\mu^{(t)}(I \otimes |i\rangle\langle i|)) \\ &= \sum_{i \in \mathbb{Z}^d} \phi(i) \text{Tr}(\mathbb{E}(\omega^{(t)})(I \otimes |i\rangle\langle i|)) \\ &= \sum_{i \in \mathbb{Z}^d} \phi(i) \mathbb{E} \left(\text{Tr}(\omega^{(t)}(I \otimes |i\rangle\langle i|)) \right) \\ &= \sum_{i \in \mathbb{Z}^d} \phi(i) \mathbb{E} \left(\text{Tr}(|X_t\rangle\langle X_t| |i\rangle\langle i|) \right) \\ &= \sum_{i \in \mathbb{Z}^d} \phi(i) \mathbb{E}(\mathbb{1}_{X_t=i}) \\ &= \mathbb{E}(\phi(X_t)), \end{aligned}$$

and the result holds. □

In the next section, we state the CLT.

3 Central Limit Theorem

This section is devoted to prove the Central Limit Theorem for CTOQWs. The result holds under some assumption concerning the Lindblad operator on \mathcal{H} . This

operator is defined below.

$$\begin{aligned} \mathcal{L} : \mathcal{B}(\mathcal{H}) &\rightarrow \mathcal{B}(\mathcal{H}) \\ \rho &\mapsto D_0\rho + \rho D_0^* + \sum_{r=1}^{2d} D_r\rho D_r^*. \end{aligned}$$

Our main assumption for the CLT is the following.

- (H1) There exists a unique density matrix $\rho_{inv} \in \mathcal{S}_{\mathcal{H}}$ such that

$$\mathcal{L}(\rho_{inv}) = 0.$$

In particular $\dim \text{Ker}(\mathcal{L}) = 1$.

Under the condition (H1), we have the following ergodic theorem which is a particular case of the Ergodic Theorem of [19]. In particular this theorem shall be useful in the proof of the CLT.

Theorem 3.0.3 ([19]). *Assume (H1). Let $(\rho_t, X_t)_{t \geq 0}$ the Markov process defined in Proposition 2.2.1, therefore*

$$\frac{1}{t} \int_0^t \rho_s ds \xrightarrow{a.s.} \rho_{inv}.$$

Now, our strategy to show the CLT consists in reducing the problem to a CLT for martingales with the help of the solution of the Poisson equation. To this end let us introduce the generator of the process $(\rho_t, X_t)_{t \geq 0}$.

We denote \mathcal{A} the Markov generator of the process $(\rho_t, X_t)_{t \geq 0}$ and $\mathbb{D}(\mathcal{A})$ his domain. For all $f \in \mathbb{D}(\mathcal{A})$, $\rho \in \mathcal{S}_{\mathcal{H}}$ and $x \in \mathbb{Z}^d$, we get

$$\begin{aligned} \mathcal{A}f(\rho, x) &= D_{(\rho, x)}f(\mathcal{F}(\rho), 0) \\ &\quad + \sum_{r=1}^{2d} \left[f\left(\frac{D_r\rho D_r^*}{\text{Tr}(D_r\rho D_r^*)}, x + e_r\right) - f(\rho, x) \right] \text{Tr}(D_r\rho D_r^*) \end{aligned} \quad (8)$$

where $\mathcal{F}(\rho) = D_0\rho + \rho D_0^* - \rho \text{Tr}(D_0\rho + \rho D_0^*)$ for all $\rho \in \mathcal{S}_{\mathcal{H}}$ and where $D_{(\rho, x)}f$ denotes the differential of f at (ρ, x) .

Remark: Note that in the sequel we do not need to make precise the exact domain of \mathcal{A} . Actually we shall apply the Markov generator on \mathcal{C}^1 functions.

We shall also need the following quantity,

$$m = \sum_{r=1}^{2d} \text{Tr}(D_r\rho_{inv} D_r^*) e_r.$$

The following lemma shall be used in the proof.

Lemma 3.0.4. *For all $u \in \mathbb{R}^d$, the equation*

$$\mathcal{L}^*(J_u) = -\left(\sum_{r=1}^{2d}(e_r \cdot u)D_r^*D_r - (m \cdot u)I\right) \quad (9)$$

admits a solution and the difference between any couple of solutions of (9) is a multiple of the identity.

Proof. First, let us remark that

$$\mathrm{Tr}\left(\rho_{inv}\left(\sum_{r=1}^{2d}(e_r \cdot u)D_r^*D_r - (m \cdot u)I\right)\right) = \sum_{r=1}^{2d}\mathrm{Tr}(D_r\rho_{inv}D_r^*)(e_r \cdot u) - (m \cdot u)\mathrm{Tr}(\rho_{inv}) = 0,$$

which implies that $-\left(\sum_{r=1}^{2d}(e_r \cdot u)D_r^*D_r - (m \cdot u)I\right) \in \{\rho_{inv}\}^\perp$. But by hypothesis, we have $\{\rho_{inv}\}^\perp = \mathrm{Ker}(\mathcal{L})^\perp$. Moreover, since $\mathrm{Ker}(\mathcal{L})^\perp = \mathrm{Im}(\mathcal{L}^*)$, we finally get that

$$-\left(\sum_{r=1}^{2d}(e_r \cdot u)D_r^*D_r - (m \cdot u)I\right) \in \mathrm{Im}(\mathcal{L}^*)$$

which proves the existence of the lemma. Now we prove the second part. To this end consider J_u and J'_u two solutions of (9) and set $H_u = J_u - J'_u$. It is then clear that

$$\mathcal{L}^*(H_u) = 0.$$

Therefore $H_u \in \mathrm{Ker}(\mathcal{L}^*)$. Since $\dim \mathrm{Ker}(\mathcal{L}) = 1$, we get $\dim \mathrm{Ker}(\mathcal{L}^*) = 1$ and since $\mathcal{L}^*(I) = D_0^* + D_0 + \sum_{r=1}^{2d}D_r^*D_r = 0$, the operator H_u is necessarily a multiple of the identity. \square

From now on, for $u \in \mathbb{R}^d$, we denote J_u the unique solution of (9) such that $\mathrm{Tr}(J_u) = 0$. Moreover, if $u = e_r$, then we simply write $J_u = J_r$. Using the linearity of \mathcal{L}^* , one can notice that:

$$J_u = \sum_{r=1}^d u_r J_r,$$

for all $u = (u_1, \dots, u_d) \in \mathbb{R}^d$.

The next lemma concerns the Poisson equation in our context (see [20] for more details on the Poisson equation).

Lemma 3.0.5. *For all $(\rho, x) \in \mathbb{S} \times \mathbb{Z}^d$ and $u \in \mathbb{R}^d$, let set*

$$f_u(\rho, x) = \mathrm{Tr}(\rho J_u) + x \cdot u. \quad (10)$$

Then f_u is solution of the Poisson equation:

$$\mathcal{A}f_u(\rho, x) = m \cdot u. \quad (11)$$

Proof. For all $(\rho, x) \in \mathbb{S} \times \mathbb{Z}^d$ and $u \in \mathbb{R}^d$, we complete the following computation:

$$\begin{aligned}
 \mathcal{A}f_u(\rho, x) &= \text{Tr}(\mathcal{F}(\rho)J_u) \\
 &+ \sum_{r=1}^{2d} \left[\text{Tr} \left(\frac{D_r \rho D_r^*}{\text{Tr}(D_r \rho D_r^*)} J_u \right) + x \cdot u + e_r \cdot u - \text{Tr}(\rho J_u) - x \cdot u \right] \text{Tr}(D_r \rho D_r^*) \\
 &= \text{Tr} \left(D_0 \rho J_u + \rho D_0^* J_u - \text{Tr}(D_0 \rho + \rho D_0^*) \rho J_u \right) \\
 &+ \sum_{r=1}^{2d} \text{Tr}(D_r \rho D_r^* J_u) + \text{Tr}(D_r \rho D_r^*) (e_r \cdot u) - \text{Tr} \left(\text{Tr}(D_r \rho D_r^*) \rho J_u \right) \\
 &= \text{Tr} \left(\rho \left[J_u D_0 + D_0^* J_u + \sum_{r=1}^{2d} D_r^* J_u D_r + \sum_{r=1}^{2d} D_r^* D_r (e_r \cdot u) \right] \right) \\
 &= \text{Tr} \left(\rho \left[\mathcal{L}^*(J_u) + \sum_{r=1}^{2d} D_r^* D_r (e_r \cdot u) \right] \right) \\
 &= \text{Tr} \left((m \cdot u) \rho \right) \\
 &= m \cdot u,
 \end{aligned}$$

so f_u is solution of the Poisson equation (11). \square

Now we have found the solution of the Poisson equation, we express the CLT for martingales that we shall use.

Theorem 3.0.6 ([10]). *Let $(M_t)_{t \geq 0}$ be a real, càdlàg, and square integrable martingale. Suppose the following conditions:*

$$\lim_{t \rightarrow \infty} \mathbb{E} \left(\frac{1}{\sqrt{t}} \sup_{0 \leq s \leq t} |\Delta M_s| \right) = 0 \tag{12}$$

and

$$\lim_{t \rightarrow \infty} \frac{[M, M]_t}{t} = \sigma^2 \tag{13}$$

for some $\sigma \geq 0$, then

$$\frac{M_t}{\sqrt{t}} \xrightarrow[t \rightarrow +\infty]{\mathcal{L}} \mathcal{N}(0, \sigma^2).$$

We shall also use the following lemma which is a straightforward consequence of the law of large numbers for martingales (see [26]).

Lemma 3.0.7. *Let Z_t a real, càdlàg, and square integrable martingale which satisfies $\langle Z, Z \rangle_t \leq Kt$ for a constant K , then*

$$\frac{Z_t}{t} \xrightarrow[t]{a.s.} 0.$$

The last lemma below shall be useful in this part as well as in the next one.

Lemma 3.0.8. *For all $t \geq 0$ and all $u \in \mathbb{R}^d$, we have*

$$\begin{aligned} \mathbb{E} \left[\sup_{0 \leq s \leq t} |X_s - X_0| \right] &\leq (2d)t \quad \text{and} \\ \mathbb{E} \left[\sup_{0 \leq s \leq t} e^{u \cdot (X_s - X_0)} \right] &\leq \exp \left((2d)t(e^{|u|} - 1) \right). \end{aligned}$$

Proof. Let $t \geq 0$ and $u \in \mathbb{R}^d$,

$$\begin{aligned} \mathbb{E} \left[\sup_{0 \leq s \leq t} e^{u \cdot (X_s - X_0)} \right] &\leq \mathbb{E} \left[e^{|u| \int_0^t |dX_s|} \right] \\ &\leq \mathbb{E} \left[\exp \left(|u| \sum_{r=1}^{2d} \int_{v=0}^t \int_{y \in \mathbb{R}} \mathbb{1}_{0 < y < \text{Tr}(D_r \rho_v D_r^*)} N^r(dy, dv) \right) \right] \\ &\leq \mathbb{E} \left[\exp \left(|u| \sum_{r=1}^{2d} \int_0^t \int_{y=0}^1 N^r(dy, dv) \right) \right]. \end{aligned}$$

Since N^r are independent Poisson point processes on \mathbb{R}^2 , we get

$$\begin{aligned} \mathbb{E} \left[\sup_{0 \leq s \leq t} e^{u \cdot (X_s - X_0)} \right] &\leq \mathbb{E} \left[\exp \left(|u| \int_0^t \int_{y=0}^1 N^1(dy, dv) \right) \right]^{2d} \\ &\leq \exp \left((2d)t(e^{|u|} - 1) \right). \end{aligned}$$

In the same way, one can prove that

$$\mathbb{E} \left[\sup_{0 \leq s \leq t} |X_s - X_0| \right] \leq \mathbb{E} \left[\sum_{r=1}^{2d} \int_{v=0}^t \int_{y=0}^1 N^r(dy, dv) \right] \leq (2d)t.$$

□

Now, we are in the position to state the main result of this section.

Theorem 3.0.9. *Assume (H_1) holds. Let $(\rho_t, X_t)_{t \geq 0}$ the Markov process defined in Proposition 2.2.1 then*

$$\frac{X_t - tm}{\sqrt{t}} \xrightarrow[t \rightarrow +\infty]{\mathcal{L}} \mathcal{N}(0, V),$$

where $V \in \mathcal{M}_d(\mathbb{R})$ such that for all $r, q \in \{1, \dots, d\}$,

$$\begin{aligned} V_{rq} &= -m_q \text{Tr}(\rho_{inv} J_r) - m_r \text{Tr}(\rho_{inv} J_q) \\ &\quad + \delta_{rq} \left(\text{Tr}(D_r \rho_{inv} D_r^*) + \text{Tr}(D_{r+d} \rho_{inv} D_{r+d}^*) \right) \\ &\quad + \text{Tr}(D_q \rho_{inv} D_q^* J_r) + \text{Tr}(D_r \rho_{inv} D_r^* J_q) \\ &\quad - \text{Tr}(D_{q+d} \rho_{inv} D_{q+d}^* J_r) - \text{Tr}(D_{r+d} \rho_{inv} D_{r+d}^* J_q). \end{aligned}$$

Remark: Proposition 2.2.2 implies then the CLT for the process $(Q_t)_{t \geq 0}$ as it holds for $(X_t)_{t \geq 0}$.

Proof. As announced, the proof is a combination of Lemma 3.0.5 and Theorem 3.0.6. Let $u \in \mathbb{R}^d$ and f_u the \mathcal{C}^1 function defined in Lemma 3.0.5. Since \mathcal{A} is the generator of $(\rho_t, X_t)_{t \geq 0}$, following the theory of problem of martingale, the process $(M_t)_{t \geq 0}$ defined by

$$\begin{aligned} M_t &= f_u(\rho_t, X_t) - f_u(\rho_0, X_0) - \int_0^t \mathcal{A} f_u(\rho_{s-}, X_{s-}) ds \\ &= \text{Tr}(\rho_t J_u) - \text{Tr}(\rho_0 J_u) + X_t \cdot u - X_0 \cdot u - (m \cdot u)t \end{aligned}$$

is a local martingale with respect to the filtration \mathcal{F} associated to $(\rho_t, X_t)_{t \geq 0}$ (see [26, 13] for more details on problem of martingale). In order to apply Theorem 3.0.6, we shall show that (M_t) is a true martingale. To this end it is sufficient to show that $\mathbb{E} \left(\sup_{0 \leq s \leq t} |M_s| \right) < \infty$ (see [13] for more details). This way, since $|\text{Tr}(\rho J_u)| \leq \|J_u\|_\infty$ for all $\rho \in \mathcal{S}_{\mathcal{H}}$, one can check with the help of Lemma 3.0.8 that

$$\mathbb{E} \left(\sup_{0 \leq s \leq t} |M_s| \right) \leq 2\|J_u\|_\infty + 2d|u|t + |m \cdot u|t.$$

Now we shall see that (M_t) fulfills the conditions of Theorem 3.0.6. The first one is the easiest one. Indeed,

$$|\Delta M_s| \leq |\text{Tr}(\Delta \rho_s J_u)| + |\Delta X_s \cdot u| \leq 2\|J_u\|_\infty + |u|.$$

This shows that ΔM_s is bounded independently of s and thus the condition (12) holds. Now, we check that (M_t) satisfies Equation (13). The bracket $[M, M]_t$

satisfies:

$$\begin{aligned}
 d[M, M]_s &= d[u.X, u.X]_s + 2d[u.X, \text{Tr}(\rho J_u)]_s + d[\text{Tr}(\rho J_u), \text{Tr}(\rho J_u)]_s \\
 &= \sum_{r=1}^{2d} (e_r \cdot u)^2 \tilde{N}^r(ds) + 2 \sum_{r=1}^{2d} (e_r \cdot u) \text{Tr} \left(\frac{D_r \rho_s - D_r^*}{\text{Tr}(D_r \rho_s - D_r^*)} J_u \right) \tilde{N}^r(ds) \\
 &\quad - 2 \sum_{r=1}^{2d} (e_r \cdot u) \text{Tr}(\rho_s - J_u) \tilde{N}^r(ds) + \sum_{r=1}^{2d} \text{Tr} \left(\frac{D_r \rho_s - D_r^*}{\text{Tr}(D_r \rho_s - D_r^*)} J_u \right)^2 \tilde{N}^r(ds) \\
 &\quad - 2 \sum_{r=1}^{2d} \text{Tr} \left(\frac{D_r \rho_s - D_r^*}{\text{Tr}(D_r \rho_s - D_r^*)} J_u \right) \text{Tr}(\rho_s - J_u) \tilde{N}^r(ds) + \sum_{r=1}^{2d} \text{Tr}(\rho_s - J_u)^2 \tilde{N}^r(ds) \\
 &= \sum_{r=1}^{2d} \left(\text{Tr} \left(\frac{D_r \rho_s - D_r^*}{\text{Tr}(D_r \rho_s - D_r^*)} J_u \right)^2 - \text{Tr}(\rho_s - J_u)^2 \right) \tilde{N}^r(ds) \\
 &\quad - 2 \text{Tr}(\rho_s - J_u) \sum_{r=1}^{2d} \left(\text{Tr} \left(\frac{D_r \rho_s - D_r^*}{\text{Tr}(D_r \rho_s - D_r^*)} J_u \right) - \text{Tr}(\rho_s - J_u) + (e_r \cdot u) \right) \tilde{N}^r(ds) \\
 &\quad + \sum_{r=1}^{2d} (e_r \cdot u)^2 \tilde{N}^r(ds) + 2 \sum_{r=1}^{2d} (e_r \cdot u) \text{Tr} \left(\frac{D_r \rho_s - D_r^*}{\text{Tr}(D_r \rho_s - D_r^*)} J_u \right) \tilde{N}^r(ds).
 \end{aligned}$$

Now we shall make the martingales $Y^r(t) = \tilde{N}^r(t) - \int_0^t \text{Tr}(D_r \rho_s - D_r^*) ds$ appear in the first and the last term of the above expression. Concerning the second term, we recognize $d\text{Tr}(\rho_s J_u)$ and $d(X_s \cdot u)$ to get

$$\begin{aligned}
 d[M, M]_s &= \sum_{r=1}^{2d} \left(\text{Tr} \left(\frac{D_r \rho_s - D_r^*}{\text{Tr}(D_r \rho_s - D_r^*)} J_u \right)^2 - \text{Tr}(\rho_s - J_u)^2 \right) Y^r(ds) \\
 &\quad + \sum_{r=1}^{2d} \left(\text{Tr} \left(\frac{D_r \rho_s - D_r^*}{\text{Tr}(D_r \rho_s - D_r^*)} J_u \right)^2 - \text{Tr}(\rho_s - J_u)^2 \right) \text{Tr}(D_r \rho_s - D_r^*) ds \\
 &\quad - 2 \text{Tr}(\rho_s - J_u) \left(d\text{Tr}(\rho_s J_u) + d(X_s \cdot u) - \text{Tr}(\mathcal{F}(\rho_s) J_u) ds \right) \\
 &\quad + \sum_{r=1}^{2d} (e_r \cdot u)^2 Y^r(ds) + 2 \sum_{r=1}^{2d} (e_r \cdot u) \text{Tr} \left(\frac{D_r \rho_s - D_r^*}{\text{Tr}(D_r \rho_s - D_r^*)} J_u \right) Y^r(ds) \\
 &\quad + \sum_{r=1}^{2d} (e_r \cdot u)^2 \text{Tr}(D_r \rho_s - D_r^*) ds + 2 \sum_{r=1}^{2d} (e_r \cdot u) \text{Tr}(D_r \rho_s - D_r^* J_u) ds. \quad (14)
 \end{aligned}$$

One can remark that for $h(\rho, x) := \text{Tr}(\rho J_u)^2$, we get for all $(\rho, x) \in \mathcal{S}_{\mathcal{H}} \times \mathbb{Z}^d$:

$$\begin{aligned} \mathcal{A}h(\rho, x) &= 2\text{Tr}(\mathcal{F}(\rho)J_u)\text{Tr}(\rho J_u) \\ &\quad + \sum_{r=1}^{2d} \left(\text{Tr} \left(\frac{D_r \rho D_r^*}{\text{Tr}(D_r \rho D_r^*)} J_u \right)^2 - \text{Tr}(\rho J_u)^2 \right) \text{Tr}(D_r \rho D_r^*). \end{aligned} \quad (15)$$

Since $h \in \mathcal{C}^1$, the process $(S_t^h)_{t \geq 0}$ defined by:

$$S_t^h = \text{Tr}(\rho_t J_u)^2 - \text{Tr}(\rho_0 J_u)^2 - \int_0^t \mathcal{A}h(\rho_{s-}, X_{s-}) ds$$

is a local martingale. Besides, since $|\text{Tr}(\rho J_u)| \leq \|J_u\|_{\infty}$ for all $\rho \in \mathcal{S}_{\mathcal{H}}$, one has

$$\mathbb{E} \left(\sup_{0 \leq s \leq t} |S_t^h| \right) \leq \alpha + \beta t < \infty$$

for all $t \geq 0$, so S_t^h is actually a true martingale.

Now using Equation (15) in the second line of (14) and recognizing dM_t in the third one, we have

$$\begin{aligned} d[M, M]_s &= \sum_{r=1}^{2d} \left(\text{Tr} \left(\frac{D_r \rho_{s-} D_r^*}{\text{Tr}(D_r \rho_{s-} D_r^*)} J_u \right)^2 - \text{Tr}(\rho_{s-} J_u)^2 \right) Y^r(ds) \\ &\quad + \mathcal{A}h(\rho_{s-}, X_{s-}) ds - 2\text{Tr}(\mathcal{F}(\rho_{s-})J_u)\text{Tr}(\rho_{s-} J_u) ds \\ &\quad - 2\text{Tr}(\rho_{s-} J_u) \left(dM_s + (m \cdot u) ds - \text{Tr}(\mathcal{F}(\rho_{s-})J_u) ds \right) \\ &\quad + \sum_{r=1}^{2d} (e_r \cdot u)^2 Y^r(ds) + 2 \sum_{r=1}^{2d} (e_r \cdot u) \text{Tr} \left(\frac{D_r \rho_{s-} D_r^*}{\text{Tr}(D_r \rho_{s-} D_r^*)} J_u \right) Y^r(ds) \\ &\quad + \sum_{r=1}^{2d} (e_r \cdot u)^2 \text{Tr}(D_r \rho_{s-} D_r^*) ds + 2 \sum_{r=1}^{2d} (e_r \cdot u) \text{Tr}(D_r \rho_{s-} D_r^* J_u) ds \\ &= \sum_{r=1}^{2d} \left(\text{Tr} \left(\frac{D_r \rho_{s-} D_r^*}{\text{Tr}(D_r \rho_{s-} D_r^*)} J_u \right)^2 - \text{Tr}(\rho_{s-} J_u)^2 + (e_r \cdot u)^2 \right. \\ &\quad \left. + 2(e_r \cdot u) \text{Tr} \left(\frac{D_r \rho_{s-} D_r^*}{\text{Tr}(D_r \rho_{s-} D_r^*)} J_u \right) \right) Y^r(ds) \\ &\quad - dS_s^h - 2\text{Tr}(\rho_{s-} J_u) dM_s + d(\text{Tr}(\rho_s J_u)^2) \\ &\quad + \text{Tr} \left(\rho_{s-} \left[-2(m \cdot u) J_u + \sum_{r=1}^{2d} (e_r \cdot u)^2 D_r^* D_r + 2 \sum_{r=1}^{2d} (e_r \cdot u) D_r^* J_u D_r \right] \right) ds. \end{aligned} \quad (16)$$

Let us denote by H_s^r the term in front of $Y^r(ds)$. Now, we shall apply Lemma 3.0.7. To this end, recall that for all $\rho \in \mathcal{S}_{\mathcal{H}}$, $|\text{Tr}(\rho J_u)| \leq \|J_u\|_{\infty}$, this implies the following estimates:

$$\begin{aligned} \left\langle \int_0^\cdot H_s^r dY_s^r, \int_0^\cdot H_s^r dY_s^r \right\rangle_t &\leq (2\|J_u\|_{\infty}^2 + |u|^2 + 2|u|\|J_u\|_{\infty})^2 \|D_r^* D_r\|_{\infty} t, \\ \left\langle \int_0^\cdot -2\text{Tr}(\rho_{s-} J_u) dM_s, \int_0^\cdot -2\text{Tr}(\rho_{s-} J_u) dM_s \right\rangle_t &\leq 4\|J_u\|_{\infty}^2 (|u| + 2\|J_u\|_{\infty})^2 \left(\sum_{r=1}^{2d} \|D_r^* D_r\|_{\infty} \right) t, \\ \langle S^h, S^h \rangle_t &\leq 64\|J_u\|_{\infty}^4 \left(\sum_{r=1}^{2d} \|D_r^* D_r\|_{\infty} \right) t. \end{aligned}$$

Lemma 3.0.7 shows that only the last term of (16) contributes to $\lim_{t \rightarrow \infty} \frac{[M, M]_t}{t}$. Applying Theorem 3.0.3, we get

$$\begin{aligned} \lim_{t \rightarrow \infty} \frac{[M, M]_t}{t} &= \lim_{t \rightarrow \infty} \frac{1}{t} \int_0^t \text{Tr} \left(\rho_{s-} \left[-2(m \cdot u) J_u + \sum_{r=1}^{2d} (e_r \cdot u)^2 D_r^* D_r \right. \right. \\ &\quad \left. \left. + 2 \sum_{r=1}^{2d} (e_r \cdot u) D_r^* J_u D_r \right] \right) ds \\ &= \text{Tr} \left(\rho_{inv} \left[-2(m \cdot u) J_u + \sum_{r=1}^{2d} (e_r \cdot u)^2 D_r^* D_r + 2 \sum_{r=1}^{2d} (e_r \cdot u) D_r^* J_u D_r \right] \right). \end{aligned}$$

$$\text{Now defining } \sigma_u^2 = \text{Tr} \left(\rho_{inv} \left[-2(m \cdot u) J_u + \sum_{r=1}^{2d} (e_r \cdot u)^2 D_r^* D_r + 2 \sum_{r=1}^{2d} (e_r \cdot u) D_r^* J_u D_r \right] \right),$$

Theorem 3.0.6 states that:

$$\frac{M_t}{\sqrt{t}} = \frac{X_t \cdot u - (m \cdot u)t + \text{Tr}(\rho_t J_u) - \text{Tr}(\rho_0 J_u) - X_0 \cdot u}{\sqrt{t}} \xrightarrow[t \rightarrow +\infty]{\mathcal{L}} \mathcal{N}(0, \sigma_u^2).$$

Now since, since $\left(\text{Tr}(\rho_t J_u) - \text{Tr}(\rho_0 J_u) - X_0 \cdot u \right)$ is bounded independently of t , one can obviously deduce that for all $u = (u_1, \dots, u_d) \in \mathbb{R}^d$, one has

$$\frac{X_t \cdot u - (m \cdot u)t}{\sqrt{t}} \xrightarrow[t \rightarrow +\infty]{\mathcal{L}} \mathcal{N}(0, \sigma_u^2),$$

where

$$\begin{aligned} \sigma_u^2 &= -2 \sum_{r,q=1}^d u_r u_q m_q \text{Tr}(\rho_{inv} J_r) + \sum_{r=1}^d u_r^2 \left(\text{Tr}(D_r \rho_{inv} D_r^*) + \text{Tr}(D_{r+d} \rho_{inv} D_{r+d}^*) \right) \\ &\quad + 2 \sum_{r,q=1}^d u_r u_q \left(\text{Tr}(D_q \rho_{inv} D_q^* J_r) - \text{Tr}(D_{q+d} \rho_{inv} D_{q+d}^* J_r) \right). \end{aligned}$$

Finally we can check that $\sigma_u^2 = \sum_{r,q=1}^d u_r u_q V_{rq}$ for all $u = (u_1, \dots, u_d) \in \mathbb{R}^d$, which ends the proof. \square

We finish this section by specifying the case $d = 1$. This is the simpler case where the walker can only jump to the right or the left. The Markov process $(\rho_t, X_t)_{t \geq 0}$, with values in $\mathcal{S}_{\mathcal{H}} \times \mathbb{Z}$, is defined by the following differential equations:

$$\begin{aligned} d\rho_s &= \left(D_0 \rho_{s-} + \rho_{s-} D_0^* - \rho_{s-} \text{Tr}(D_0 \rho_{s-} + \rho_{s-} D_0^*) \right) ds \\ &\quad + \left(\frac{D_1 \rho_{s-} D_1^*}{\text{Tr}(D_1 \rho_{s-} D_1^*)} - \rho_{s-} \right) \tilde{N}^1(ds) + \left(\frac{D_2 \rho_{s-} D_2^*}{\text{Tr}(D_2 \rho_{s-} D_2^*)} - \rho_{s-} \right) \tilde{N}^2(ds) \end{aligned}$$

$$\text{and } dX_s = \tilde{N}^1(ds) - \tilde{N}^2(ds),$$

where $D_0, D_1, D_2 \in \mathcal{B}(\mathcal{H})$ such that $D_0 + D_0^* + D_1^* D_1 + D_2^* D_2 = 0$.

Theorem 3.0.10. *Suppose that the Lindblad operator*

$$\mathcal{L}(\rho) = D_0 \rho + \rho D_0^* + D_1 \rho D_1^* + D_2 \rho D_2^*$$

admits a unique density matrix ρ_{inv} such that $\mathcal{L}(\rho_{inv}) = 0$.

Set $m = \text{Tr}(D_1 \rho_{inv} D_1^) - \text{Tr}(D_2 \rho_{inv} D_2^*)$, and let J the unique solution of $\mathcal{L}^*(J) = -D_1^* D_1 + D_2^* D_2 + mI$ such that $\text{Tr}(J) = 0$.*

Then, we have the following CLT

$$\frac{X_t - tm}{\sqrt{t}} \xrightarrow[t \rightarrow +\infty]{\mathcal{L}} \mathcal{N}(0, \sigma^2)$$

where $\sigma^2 = \text{Tr}(\rho_{inv}[-2mJ + D_1^ D_1 + D_2^* D_2 + 2D_1^* J D_1 - 2D_2^* J D_2])$.*

4 Large Deviation Principle

Here, we study a Large Deviation Principle (LDP) for CTOQWs. Our proof is inspired by strategies developed in [17, 7] which are essentially based on the application of the Gärtner-Ellis Theorem ([12]). In the following, one can notice that the Perron-Frobenius Theorem for positive maps ([14]) is the main tool to apply the Gärtner-Ellis Theorem.

In order to prove the LDP, we shall use a deformed Lindblad operator. From now on, we define for all $u \in \mathbb{R}^d$, the operators $D_r^{(u)} = e^{\frac{u \cdot e_r}{2}} D_r$, $r \in \{0, \dots, 2d\}$, and we denote $\mathcal{L}^{(u)}$ the deformed Lindblad operator associated to the operators $D_r^{(u)}$, that is,

$$\mathcal{L}^{(u)}(\rho) = D_0 \rho + \rho D_0^* + \sum_{r=1}^{2d} e^{u \cdot e_r} D_r \rho D_r^*,$$

for all $\rho \in \mathcal{S}_{\mathcal{H}}$.

Now, defining

$$\phi^{(u)}(\rho) = \sum_{r=1}^{2d} e^{u \cdot e_r} D_r \rho D_r^*,$$

for all $\rho \in \mathcal{S}_{\mathcal{H}}$ and all $u \in \mathbb{R}^d$, we can see that $\mathcal{L}^{(u)}$ is written in the usual Lindblad form, that is $\mathcal{L}^{(u)}(\rho) = D_0 \rho + \rho D_0^* + \phi^{(u)}(\rho)$. This way, the semi-group $\{e^{t\mathcal{L}^{(u)}}\}_{t \geq 0}$ is a completely positive (CP) semi-group (see [9] for the proof). In a same way, we write \mathcal{L} as:

$$\mathcal{L}(\rho) = D_0 \rho + \rho D_0^* + \phi(\rho)$$

for all $\rho \in \mathcal{S}_{\mathcal{H}}$ where $\phi(\rho) = \sum_{r=1}^{2d} D_r \rho D_r^*$.

In this part, the notion of irreducibility is required. This notion was originally defined in [11]. There are several equivalent definitions that the reader can find in [7, 8].

Definition 4.0.11. The CP map $\phi : \rho \mapsto \sum_{r=1}^{2d} D_r \rho D_r^*$ is called irreducible if for any non-zero $x \in \mathcal{H}$, the set $\mathbb{C}[D]x$ is dense in \mathcal{H} , where $\mathbb{C}[D]$ is the set of polynomials in D_r , $r \in \{1, \dots, 2d\}$.

In the sequel, we need $\{e^{t\mathcal{L}^{(u)}}\}_{t \geq 0}$ to be positivity improving. This means that for all $t > 0$ and for all $A \geq 0$, $A \in \mathcal{B}(\mathcal{H})$, one has $e^{t\mathcal{L}^{(u)}}(A) > 0$. The following lemma provides an effective criterion for verifying that $\{e^{t\mathcal{L}^{(u)}}\}_{t \geq 0}$ is positivity improving.

Lemma 4.0.12. *If ϕ is irreducible, then $\phi^{(u)}$ is irreducible by a direct application of Definition 4.0.11 and therefore $\{e^{t\mathcal{L}^{(u)}}\}_{t \geq 0}$ is positivity improving ([17]).*

The next two lemmas are relevant in the proof of the main theorem of this part. In particular, the following lemma describes the largest eigenvalue associated to the deformed Lindblad semigroup $\{e^{t\mathcal{L}^{(u)}}\}_{t \geq 0}$.

Lemma 4.0.13. *Let $t \geq 0$, suppose that ϕ is irreducible. Set*

$$l_u = \max\{\operatorname{Re}(\lambda), \lambda \in \operatorname{Sp}(\mathcal{L}^{(u)})\}.$$

Then e^{tl_u} is an algebraically simple eigenvalue of $e^{t\mathcal{L}^{(u)}}$, and the associated eigenvector V_u is strictly positive (which can be normalized to be in $\mathcal{S}_{\mathcal{H}}$). Besides, the map $u \mapsto l_u$ can be extended to be analytic in a neighbourhood of \mathbb{R}^d .

Proof. The first part has been proved in [17], the proof is based on the Perron-Frobenius Theorem for CP maps ([14]). In particular, using such result, one get that the multiplicity of e^{tl_u} is geometrically simple and that V_u is strictly positive. It remains to show the algebraic simplicity. To this end, we introduce:

$$\Psi(X) = V_u^{-\frac{1}{2}} e^{t(\mathcal{L}^{(u)} - l_u)} (V_u^{\frac{1}{2}} X V_u^{\frac{1}{2}}) V_u^{-\frac{1}{2}}.$$

Then $\Psi(I) = I$, hence $\|\Psi\| = 1$ and Ψ^* is trace preserving. Moreover it is clear that Ψ is positivity improving, then 1 is a geometrically simple eigenvalue of Ψ , therefore this holds for Ψ^* too. By applying Theorem 2.5. of [14], one get that the associated eigenvector X_1 is positive. Assume by contradiction that 1 is not algebraically simple for Ψ^* . Then the Jordan decomposition shows that there exists X_2 such that $\Psi^*(X_2) = X_1 + X_2$. Then, since Ψ^* is trace preserving, $\operatorname{Tr}(X_1) = 0$, hence $X_1 = 0$ which is impossible. This implies that e^{tl_u} is algebraically simple for $e^{t\mathcal{L}^{(u)}}$. The analyticity of $u \mapsto l_u$ is a simple application of perturbation theory for matrix eigenvalues (see Chapter II in [18]). \square

The next lemma describes the link between the moment generating function of $(X_t - X_0)$ and the deformed Lindblad semigroup.

Lemma 4.0.14. *For all $t \geq 0$ and all $u \in \mathbb{R}^d$, one has*

$$\mathbb{E}\left(e^{u \cdot (X_t - X_0)}\right) = \operatorname{Tr}\left(e^{t\mathcal{L}^{(u)}}(\mathbb{E}[\rho_0])\right).$$

Proof. The idea of the proof consists in rewriting $\mathbb{E}\left(e^{u \cdot (X_t - X_0)}\right)$ with the help of a Dyson expansion. From now, we set $u \in \mathbb{R}^d$ and $f : (\rho, x) \mapsto e^{u \cdot x} \in \mathcal{C}^1$. Since \mathcal{A} is also the generator of $(\rho_t, X_t - X_0)_{t \geq 0}$, the process $(M_t^f)_{t \geq 0}$ defined by

$$M_t^f = f(\rho_t, X_t - X_0) - f(\rho_0, 0) - \int_0^t \mathcal{A}f(\rho_{t_1-}, X_{t_1-} - X_0) dt_1$$

is a local martingale. Due to Lemma 3.0.8, one has the following upper bound.

$$\begin{aligned} \mathbb{E}\left(\sup_{0 \leq s \leq t} |M_s^f|\right) &\leq \exp\left((2d)t(e^{|u|} - 1)\right) + 1 \\ &\quad + (e^{|u|} + 1) \left(\sum_{r=1}^{2d} \|D_r^* D_r\|_{\infty}\right) \exp\left((2d)t(e^{|u|} - 1)\right). \end{aligned}$$

Then $\mathbb{E}\left(\sup_{0 \leq s \leq t} |M_s^f|\right) < \infty$, for all $t \geq 0$ which implies that $(M_t^f)_{t \geq 0}$ is a true martingale. This leads to

$$\mathbb{E}\left(f(\rho_t, X_t - X_0)\right) = \mathbb{E}\left(f(\rho_0, 0)\right) + \mathbb{E}\left(\int_0^t \mathcal{A}f(\rho_{t_1}, X_{t_1} - X_0) dt_1\right).$$

This way, we can develop $\mathbb{E}\left(e^{u \cdot (X_t - X_0)}\right)$. For all $t \geq 0$,

$$\begin{aligned} \mathbb{E}\left(e^{u \cdot (X_t - X_0)}\right) &= 1 + \mathbb{E}\left(\int_0^t \sum_{r=1}^{2d} \left[e^{u \cdot (X_{t_1} - X_0 + e_r)} - e^{u \cdot (X_{t_1} - X_0)}\right] \text{Tr}(D_r \rho_{t_1} D_r^*) dt_1\right) \\ &= 1 + \mathbb{E}\left(\int_0^t e^{u \cdot (X_{t_1} - X_0)} \left[\text{Tr}(\mathcal{L}^{(u)}(\rho_{t_1})) \right. \right. \\ &\quad \left. \left. - \text{Tr}(D_0 \rho_{t_1} + \rho_{t_1} D_0^* + \sum_{r=1}^{2d} D_r \rho_{t_1} D_r^*) \right] dt_1\right) \\ &= 1 + \mathbb{E}\left(\int_0^t e^{u \cdot (X_{t_1} - X_0)} \text{Tr}(\mathcal{L}^{(u)}(\rho_{t_1})) dt_1\right) \\ &= 1 + \int_0^t \text{Tr}\left(\mathcal{L}^{(u)}\left(\mathbb{E}\left[e^{u \cdot (X_{t_1} - X_0)} \rho_{t_1}\right]\right)\right) dt_1. \end{aligned} \quad (17)$$

In a similar way, we want to develop $\mathbb{E}\left[e^{u \cdot (X_{t_1} - X_0)} \rho_{t_1}\right]$. Let $g : (\rho, x) \mapsto e^{u \cdot x} \rho \in \mathcal{C}^1$, the process $(M_t^g)_{t \geq 0}$ defined by

$$M_t^g = g(\rho_t, X_t - X_0) - g(\rho_0, 0) - \int_0^t \mathcal{A}g(\rho_{t_2-}, X_{t_2-} - X_0) dt_2$$

is a local martingale. One can also check that $\mathbb{E}\left(\sup_{0 \leq s \leq t} |M_s^g|\right) < \infty$ for all $t \geq 0$,

which implies that $(M_t^g)_{t \geq 0}$ is a true martingale. And therefore, one has

$$\begin{aligned}
 \mathbb{E} \left[e^{u \cdot (X_{t_1} - X_0)} \rho_{t_1} \right] &= \mathbb{E}(\rho_0) + \mathbb{E} \left(\int_0^{t_1} e^{u \cdot (X_{t_2} - X_0)} \mathcal{F}(\rho_{t_2}) dt_2 \right) \\
 &\quad + \mathbb{E} \left(\int_0^{t_1} \sum_{r=1}^{2d} \left[e^{u \cdot (X_{t_2} - X_0 + e_r)} \frac{D_r \rho_{t_2} D_r^*}{\text{Tr}(D_r \rho_{t_2} D_r^*)} \right. \right. \\
 &\quad \quad \left. \left. - e^{u \cdot (X_{t_2} - X_0)} \rho_{t_2} \right] \text{Tr}(D_r \rho_{t_2} D_r^*) dt_2 \right) \\
 &= \mathbb{E}(\rho_0) + \mathbb{E} \left(\int_0^{t_1} e^{u \cdot (X_{t_2} - X_0)} \left[\text{Tr}(\mathcal{L}^{(u)}(\rho_{t_2})) \right. \right. \\
 &\quad \quad \left. \left. - \rho_{t_2} \text{Tr} \left(D_0 \rho_{t_2} + \rho_{t_2} D_0^* + \sum_{r=1}^{2d} D_r \rho_{t_2} D_r^* \right) \right] dt_2 \right) \\
 &= \mathbb{E}(\rho_0) + \int_0^{t_1} \text{Tr} \left(\mathcal{L}^{(u)} \left(\mathbb{E} \left[e^{u \cdot (X_{t_2} - X_0)} \rho_{t_2} \right] \right) \right) dt_2. \quad (18)
 \end{aligned}$$

We plug (18) into (17) and we get, for all $t \geq 0$,

$$\begin{aligned}
 \mathbb{E} \left(e^{u \cdot (X_t - X_0)} \right) &= 1 + t \text{Tr} \left(\mathcal{L}^{(u)}(\mathbb{E}[\rho_0]) \right) \\
 &\quad + \int_0^t \int_0^{t_1} \text{Tr} \left((\mathcal{L}^{(u)})^2 \left(\mathbb{E} \left[e^{u \cdot (X_{t_2} - X_0)} \rho_{t_2} \right] \right) \right) dt_2 dt_1.
 \end{aligned}$$

By iterating this procedure, we obtain

$$\begin{aligned}
 \mathbb{E} \left(e^{u \cdot (X_t - X_0)} \right) &= \text{Tr} \left(\sum_{k=0}^j \frac{t^k}{k!} (\mathcal{L}^{(u)})^k \left(\mathbb{E}[\rho_0] \right) \right) \\
 &\quad + \int_{0 < t_j < \dots < t} \text{Tr} \left((\mathcal{L}^{(u)})^{j+1} \left(\mathbb{E} \left[e^{u \cdot (X_{t_{j+1}} - X_0)} \rho_{t_{j+1}} \right] \right) \right) dt_{j+1} \dots dt_1.
 \end{aligned}$$

for all $j \in \mathbb{N}$. Now it is obvious that the first term converges to $\text{Tr} \left(e^{t \mathcal{L}^{(u)}} \left(\mathbb{E}[\rho_0] \right) \right)$ when j goes to infinity. In order to conclude it remains to prove that the second terms converges to zero. Let us estimate its norm.

$$\begin{aligned}
 &\left\| \int_{0 < t_j < \dots < t} \text{Tr} \left((\mathcal{L}^{(u)})^{j+1} \left(\mathbb{E} \left[e^{u \cdot (X_{t_{j+1}} - X_0)} \rho_{t_{j+1}} \right] \right) \right) dt_{j+1} \dots dt_1 \right\|_1 \\
 &\leq \frac{t^{j+1}}{(j+1)!} \left\| \mathcal{L}^{(u)} \right\|_1^{j+1} \sup_{0 \leq s \leq t} \mathbb{E} \left[e^{u \cdot (X_s - X_0)} \right].
 \end{aligned}$$

Finally, thanks to Lemma 3.0.8 and Jensen's inequality,

$$\begin{aligned} & \left\| \int_{0 < t_j < \dots < t} \text{Tr} \left((\mathcal{L}^{(u)})^{j+1} \left(\mathbb{E} \left[e^{u \cdot (X_{t_{j+1}} - X_0)} \rho_{t_{j+1}} \right] \right) \right) dt_{j+1} \dots dt_1 \right\|_1 \\ & \leq \frac{t^{j+1}}{(j+1)!} \|\mathcal{L}^{(u)}\|_1^{j+1} e^{(2d)t(e^{|u|}-1)} \end{aligned}$$

which converges to 0 when j goes to infinity. \square

Now, we can state the main result of this part.

Theorem 4.0.15. *Let $(\rho_t, X_t)_{t \geq 0}$ the Markov process defined in Proposition 2.2.1. Assume that ϕ is irreducible. The process $\left(\frac{X_t - X_0}{t} \right)_{t \geq 0}$ satisfies a Large Deviation Principle with a good rate function Λ^* .*

Explicitly there exists a lower semicontinuous mapping $\Lambda^ : \mathbb{R}^d \mapsto [0, +\infty]$ with compact level sets $\{x | \Lambda^*(x) \leq \alpha\}$, such that, for all open set G and all closed set F with $G \subset F \subset \mathbb{R}^d$, one has:*

$$\begin{aligned} - \inf_{x \in G} \Lambda^*(x) & \leq \liminf_{t \rightarrow +\infty} \frac{1}{t} \log \mathbb{P} \left(\frac{X_t - X_0}{t} \in G \right) \\ & \leq \limsup_{t \rightarrow +\infty} \frac{1}{t} \log \mathbb{P} \left(\frac{X_t - X_0}{t} \in F \right) \leq - \inf_{x \in F} \Lambda^*(x). \end{aligned}$$

Moreover, Λ^* can be expressed explicitly,

$$\Lambda^* : x \mapsto \sup_{u \in \mathbb{R}^d} (u \cdot x - l_u)$$

where l_u is defined in Lemma 4.0.13.

Remark: Moreover, if $\mathbb{E}(e^{u \cdot X_0}) < \infty$ then the LDP holds for $(X_t)_{t \geq 0}$ and not only for $(X_t - X_0)_{t \geq 0}$. In this case, Proposition 2.2.2 allows us to have the LDP for $(Q_t)_{t \geq 0}$.

Proof. The main tool of the proof is the Gärtner-Ellis Theorem (GET) (see [12]). We focus on the moment generating function which is involved in the GET. Let $t \geq 0$ and $u \in \mathbb{R}^d$, Lemma 4.0.14 implies that

$$\mathbb{E}(e^{u \cdot (X_t - X_0)}) = \text{Tr} \left(e^{t\mathcal{L}^{(u)}} \left(\mathbb{E}[\rho_0] \right) \right) = \sum_{i \in \mathbb{Z}^d} \text{Tr} \left(e^{t\mathcal{L}^{(u)}} \left(\rho^{(0)}(i) \right) \right).$$

Set $0 < \epsilon < t$. Due to Lemma 4.0.12, $e^{\epsilon\mathcal{L}^{(u)}}$ has the property of positivity improving, therefore $e^{\epsilon\mathcal{L}^{(u)}}(\rho^{(0)}(i))$ is strictly positive for all $i \in \mathbb{Z}^d$.

If we set $r_{u,i} = \inf \left(\text{Sp} \left[e^{\epsilon\mathcal{L}^{(u)}}(\rho^{(0)}(i)) \right] \right) > 0$ and $s_{u,i} = \frac{\text{Tr} \left(e^{\epsilon\mathcal{L}^{(u)}}(\rho^{(0)}(i)) \right)}{\inf \text{Sp}(V_u)}$ then $\text{Sp} \left[e^{\epsilon\mathcal{L}^{(u)}}(\rho^{(0)}(i)) - r_{u,i}V_u \right] \subset \mathbb{R}^+$ and $\text{Sp} \left[s_{u,i}V_u - e^{\epsilon\mathcal{L}^{(u)}}(\rho^{(0)}(i)) \right] \subset \mathbb{R}^+$, and thus

$$r_{u,i}V_u \leq e^{\epsilon\mathcal{L}^{(u)}}(\rho^{(0)}(i)) \leq s_{u,i}V_u.$$

Since $e^{(t-\epsilon)\mathcal{L}^{(u)}}$ preserves the positivity, we get

$$r_{u,i}e^{(t-\epsilon)l_u}V_u \leq e^{t\mathcal{L}^{(u)}}(\rho^{(0)}(i)) \leq s_{u,i}e^{(t-\epsilon)l_u}V_u.$$

Taking the trace, Lemma 4.0.14 yields

$$e^{(t-\epsilon)l_u} \sum_{i \in \mathbb{Z}^d} r_{u,i} \leq \mathbb{E} \left(e^{u \cdot (X_t - X_0)} \right) \leq e^{(t-\epsilon)l_u} \sum_{i \in \mathbb{Z}^d} s_{u,i}.$$

The sums are finite and positive, then we have

$$\lim_{t \rightarrow +\infty} \frac{1}{t} \log \left(\mathbb{E} \left(e^{u \cdot (X_t - X_0)} \right) \right) = l_u.$$

Now define $\Lambda_t : u \mapsto \log \left(\mathbb{E} \left(e^{u \cdot \frac{X_t - X_0}{t}} \right) \right)$ the logarithm of the moment generating function of $\frac{X_t - X_0}{t}$, and $\Lambda : u \mapsto l_u$. We have shown that

$$\lim_{t \rightarrow +\infty} \frac{1}{t} \Lambda_t(tu) = \Lambda(u).$$

Since Λ is analytic (Lemma 4.0.13), Gärtner-Ellis Theorem can be applied, this proves the LDP and the associated good rate function is

$$\Lambda^* : x \mapsto \sup_{u \in \mathbb{R}^d} (u \cdot x - l_u).$$

□

5 Examples

Let us illustrate the results above with some examples.

In the case $d = 1$, we obtain a random walk on \mathbb{Z} where the walker can stand still some random time, jump to the right or jump to the left, these transformations are respectively governed by D_0, D_1 or D_2 . We must have:

$$D_0 + D_0^* + D_1^*D_1 + D_2^*D_2 = 0.$$

1. The constraint above is respected in this concrete example:

$$D_0 = -\frac{1}{2}I, \quad D_1 = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix} \quad \text{and} \quad D_2 = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 0 \\ -1 & 1 \end{pmatrix}.$$

This example falls within the scope of the CLT, in fact (H1) is checked with $\rho_{inv} = \frac{1}{2}I$. We get

$$m = \text{Tr}(D_1 \rho_{inv} D_1^*) - \text{Tr}(D_2 \rho_{inv} D_2^*) = 0; \quad J = \frac{1}{6} \begin{pmatrix} -5 & 2 \\ 2 & 5 \end{pmatrix}; \quad \sigma^2 = \frac{8}{9}.$$

Then $\frac{X_t}{\sqrt{t}} \xrightarrow[t \rightarrow +\infty]{\mathcal{L}} \mathcal{N}(0, \frac{8}{9})$. This result is illustrated with numerical simulations in Figure 1.

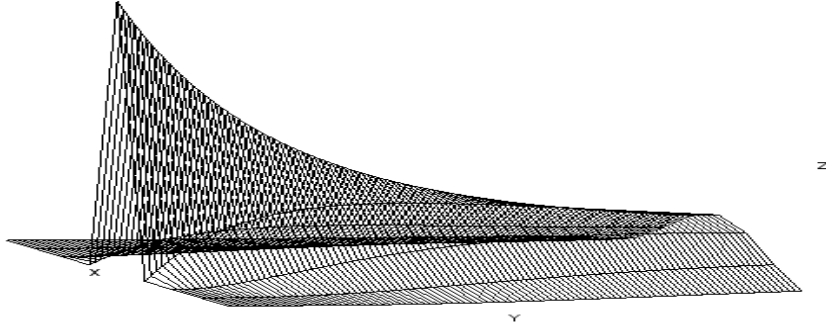
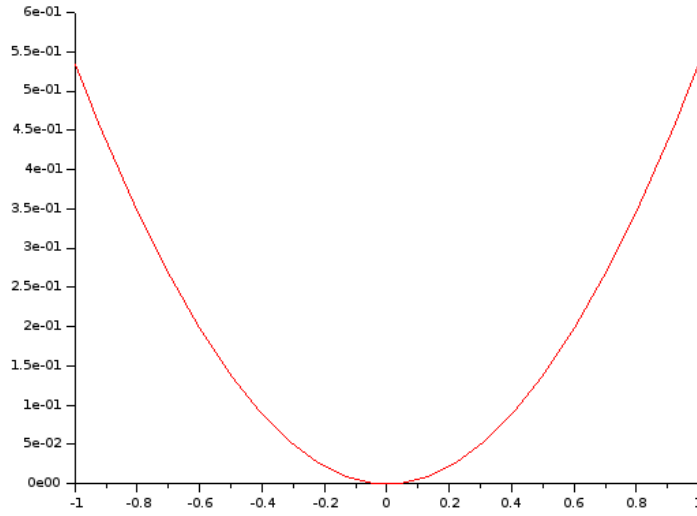


Figure 1: Evolution of the distribution of the CTOQW starting from a state localized in 0. The X -axis stands for the position $|i\rangle$ on \mathbb{Z} , the Y -axis stands for the time t and the Z -axis returns the distribution $q_t(i)$.

One can check that ϕ is irreducible. Hence the process $\left(\frac{X_t - X_0}{t}\right)_{t \geq 0}$ satisfies a LDP with a good rate function $\Lambda^* : x \mapsto \sup_{u \in \mathbb{R}^d} (u \cdot x - l_u)$ (see Figure 2 for numerical computations).

2. We focus on an other example on \mathbb{Z} :

$$D_0 = \begin{pmatrix} -\frac{3}{8} & 0 \\ 0 & -\frac{1}{4} \end{pmatrix}, \quad D_1 = \begin{pmatrix} 0 & \frac{1}{2} \\ \frac{1}{2} & 0 \end{pmatrix} \quad \text{and} \quad D_2 = \begin{pmatrix} 0 & \frac{1}{2} \\ \frac{1}{\sqrt{2}} & 0 \end{pmatrix}.$$

Figure 2: Λ^* for the first example.

We get

$$\mathcal{L}(\rho) = 0 \iff \rho = \begin{pmatrix} \frac{2}{5} & 0 \\ 0 & \frac{3}{5} \end{pmatrix}.$$

This implies that (H_1) is satisfied with

$$\rho_{inv} = \begin{pmatrix} \frac{2}{5} & 0 \\ 0 & \frac{3}{5} \end{pmatrix}.$$

In particular, we have here

$$m = \text{Tr}(D_1 \rho_{inv} D_1^*) - \text{Tr}(D_2 \rho_{inv} D_2^*) = -\frac{1}{10}; \quad J = \frac{1}{10} \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}; \quad \sigma^2 = \frac{73}{125}.$$

The Central Limit Theorem yields:

$$\frac{X_t + \frac{t}{10}}{\sqrt{t}} \xrightarrow[t \rightarrow +\infty]{\mathcal{L}} \mathcal{N}\left(0, \frac{73}{125}\right).$$

The reader can easily check that ϕ is irreducible with the help of Definition 4.0.11. Hence the process $\left(\frac{X_t - X_0}{t}\right)_{t \geq 0}$ satisfies a LDP with a good rate function $\Lambda^* : x \mapsto \sup_{u \in \mathbb{R}^d} (u \cdot x - l_u)$ where $l_u = \max\{\Re(\lambda), \lambda \in \text{Sp}(\mathcal{L}^{(u)})\}$.

Here $\mathcal{L}^{(u)}$ has a simple form in the canonical basis of $\mathcal{M}_2(\mathbb{C})$:

$$\mathcal{L}^{(u)} = \frac{1}{8} \begin{pmatrix} -6 & 0 & 0 & 2(e^u + e^{-u}) \\ 0 & -5 & 2(e^u + \sqrt{2}e^{-u}) & 0 \\ 0 & 2(e^u + \sqrt{2}e^{-u}) & -5 & 0 \\ 2(e^u + 2e^{-u}) & 0 & 0 & -4 \end{pmatrix}.$$

Hence, tedious computations show that $l_u = \frac{1}{32} \left(20 + \sqrt{208 + 64e^{2u} + 128e^{-2u}} \right)$. Figure 3 displays the good rate function Λ^* .

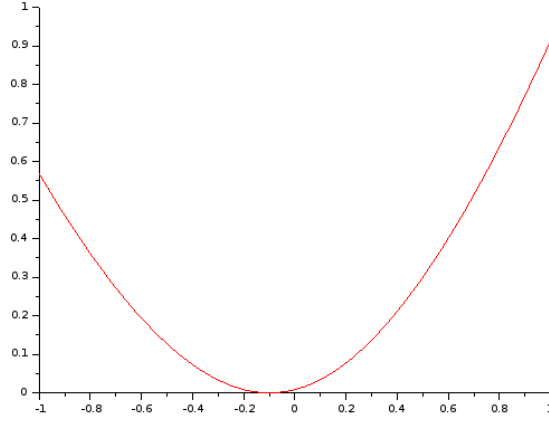


Figure 3: Λ^* for the second example.

3. Now, we study an example of CTOQWs on \mathbb{Z}^2 , this example is defined by the operators below:

$$D_0 = \begin{pmatrix} -\frac{1}{2} & 0 \\ 0 & -\frac{3}{8} \end{pmatrix}, \quad D_1 = \frac{1}{\sqrt{6}} \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}, \quad D_2 = \frac{1}{2\sqrt{2}} \begin{pmatrix} 0 & 1 \\ 0 & 1 \end{pmatrix},$$

$$D_3 = \frac{1}{\sqrt{6}} \begin{pmatrix} 1 & 0 \\ -1 & 1 \end{pmatrix} \quad \text{and} \quad D_4 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}.$$

Again, we check that:

$$D_0 + D_0^* + D_1^* D_1 + D_2^* D_2 + D_3^* D_3 + D_4^* D_4 = 0.$$

Figure 4 represents the evolution of the distribution of the CTOQW with

$$\mu^{(0)} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \\ 0 & \frac{1}{2} \end{pmatrix} \otimes |(0, 0)\rangle\langle(0, 0)| + \begin{pmatrix} \frac{1}{2} & 0 \\ 0 & 0 \end{pmatrix} \otimes |(0, 5)\rangle\langle(0, 5)|$$

as initial state of the system.

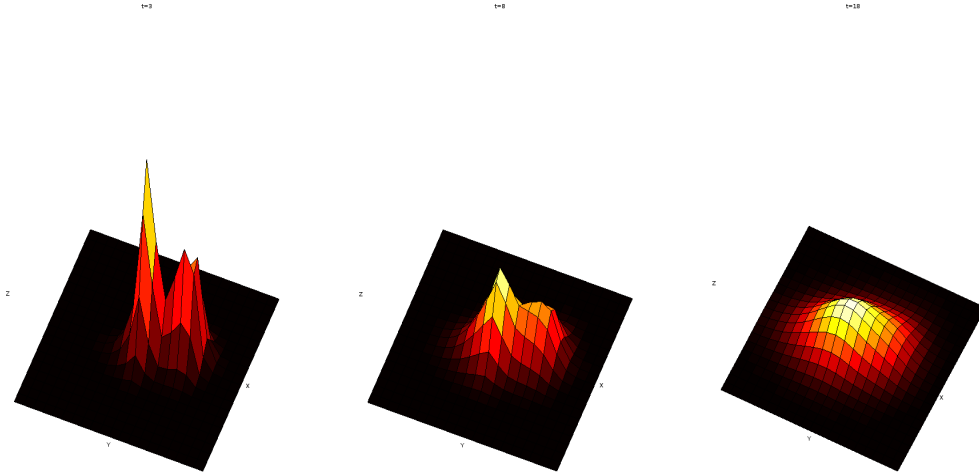


Figure 4: Distribution of the CTOQW on \mathbb{Z}^2 at time $t = 3, 8$ and 18 .

One can check that (H1) holds with $\rho_{inv} = \begin{pmatrix} \frac{7}{11} & 0 \\ 0 & \frac{4}{11} \end{pmatrix}$ and then the CLT is satisfied. The characteristics of this example are

$$m = \left(-\frac{1}{22}, -\frac{5}{22} \right),$$

$$J_1 = \frac{4}{33} \begin{pmatrix} -5 & 2 \\ 2 & 5 \end{pmatrix} \text{ and } J_2 = \frac{3}{77} \begin{pmatrix} -13 & -8 \\ -8 & 13 \end{pmatrix}; V = \frac{1}{23958} \begin{pmatrix} 10651 & -414 \\ -414 & 14661 \end{pmatrix}.$$

The Central Limit Theorem states then

$$\frac{X_t - mt}{\sqrt{t}} \xrightarrow[t \rightarrow +\infty]{\mathcal{L}} \mathcal{N}(0, V).$$

One can check easily that ϕ is irreducible. Therefore the process $\left(\frac{X_t - X_0}{t}\right)_{t \geq 0}$ satisfies a LDP with a good rate function $\Lambda^* : x \mapsto \sup_{u \in \mathbb{R}^d} (u \cdot x - l_u)$. We can plot an numerical approximation of Λ^* (see Figure 5) with the help of $\mathcal{L}^{(u)}$ developed in the canonical basis thereafter:

$$\mathcal{L}^{(u)} = \frac{1}{24} \begin{pmatrix} 4(-6+e^{u_1}+e^{-u_1}+3e^{-u_2}) & 4e^{u_1} & 4e^{u_1} & 4e^{u_1}+3e^{u_2} \\ -4e^{-u_1} & -21+4(e^{u_1}+e^{-u_1}) & 0 & 4e^{u_1}+3e^{u_2} \\ -4e^{-u_1} & 0 & -21+4(e^{u_1}+e^{-u_1}) & 4e^{u_1}+3e^{u_2} \\ 4e^{-u_1} & -4e^{-u_1} & -4e^{-u_1} & -18+4e^{-u_1}+3e^{u_2} \end{pmatrix}.$$

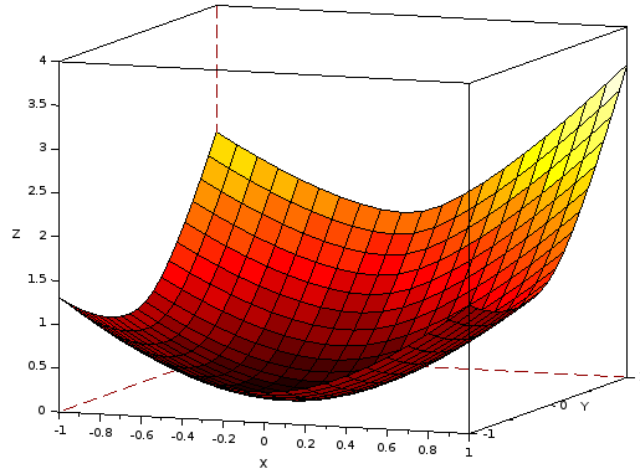


Figure 5: Λ^* for the last example.

Acknowledgments. The author is grateful to Y. Pautrat for interesting discussions around the topics and examples of the present manuscript. This work is financially supported by the ANR project STOQ: ANR-14-CE25-0003. The author wants to thank the organizers of the summer school "stochastic methods in Quantum Mechanics" in Autrans (2016) where this works have been completed.

References

- [1] S. Attal, N. Guillotin-Plantard, and C. Sabot. Central limit theorems for open quantum random walks and quantum measurement records. *Ann. Henri Poincaré*, 16(1):15–43, 2015.
- [2] S. Attal, F. Petruccione, C. Sabot, and I. Sinayskiy. Open quantum random walks. *J. Stat. Phys.*, 147(4):832–852, 2012.
- [3] A. Barchielli and M. Gregoratti. *Quantum trajectories and measurements in continuous time: the diffusive case*, volume 782. Springer, 2009.
- [4] A. Barchielli and C. Pellegrini. Jump-diffusion unravelling of a non-Markovian generalized Lindblad master equation. *J. Math. Phys.*, 51(11):112104, 18, 2010.
- [5] M. Bauer, D. Bernard, and A. Tilloy. The open quantum Brownian motions. *J. Stat. Mech. Theory Exp.*, (9):p09001, 48, 2014.
- [6] H.-P. Breuer. Non-Markovian generalization of the Lindblad theory of open quantum systems. *Phys. Rev. A (3)*, 75(2):022103, 9, 2007.
- [7] R. Carbone and Y. Pautrat. Homogeneous open quantum random walks on a lattice. *J. Stat. Phys.*, 160(5):1125–1153, 2015.
- [8] R. Carbone and Y. Pautrat. Open quantum random walks: reducibility, period, ergodic properties. *Ann. Henri Poincaré*, 17(1):99–135, 2016.
- [9] E. Christensen and D. Evans. Cohomology of operator algebras and quantum dynamical semigroups. *J. London Math. Soc. (2)*, 20(2):358–368, 1979.
- [10] I. Crimaldi and L. Pratelli. Convergence results for multivariate martingales. *Stochastic Process. Appl.*, 115(4):571–577, 2005.
- [11] E. B. Davies. Quantum stochastic processes. II. *Comm. Math. Phys.*, 19:83–105, 1970.
- [12] A. Dembo and O. Zeitouni. *Large deviations techniques and applications*, volume 38 of *Stochastic Modelling and Applied Probability*. Springer-Verlag, Berlin, 2010.
- [13] S. Ethier and T. Kurtz. *Markov processes*. Wiley Series in Probability and Mathematical Statistics: Probability and Mathematical Statistics. John Wiley & Sons, Inc., New York, 1986.

REFERENCES

- [14] D. Evans and R. Høegh-Krohn. Spectral properties of positive maps on C^* -algebras. *J. London Math. Soc. (2)*, 17(2):345–355, 1978.
- [15] S. Gudder. Quantum Markov chains. *J. Math. Phys.*, 49(7):072105, 14, 2008.
- [16] P. Hall and C. C. Heyde. *Martingale limit theory and its application*. Academic Press, Inc. [Harcourt Brace Jovanovich, Publishers], New York-London, 1980.
- [17] V. Jakšić, C.-A. Pillet, and M. Westrich. Entropic fluctuations of quantum dynamical semigroups. *J. Stat. Phys.*, 154(1-2):153–187, 2014.
- [18] T. Kato. *Perturbation theory for linear operators*. Classics in Mathematics. Springer-Verlag, Berlin, 1995.
- [19] B. Kümmerer and H. Maassen. A pathwise ergodic theorem for quantum trajectories. *J. Phys. A*, 37(49):11889–11896, 2004.
- [20] A. Makowski and A. Shwartz. The Poisson equation for countable Markov chains: probabilistic methods and interpretations. In *Handbook of Markov decision processes*, volume 40 of *Internat. Ser. Oper. Res. Management Sci.*, pages 269–303. Kluwer Acad. Publ., Boston, MA, 2002.
- [21] A. Marais, I. Sinayskiy, A. Kay, F. Petruccione, and A. Ekert. Decoherence-assisted transport in quantum networks. *New J. Phys.*, 15(January):013038, 18, 2013.
- [22] R. Motwani and P. Raghavan. *Randomized algorithms*. Cambridge University Press, Cambridge, 1995.
- [23] C. Pellegrini. Existence, uniqueness and approximation of a stochastic Schrödinger equation: the diffusive case. *Ann. Probab.*, 36(6):2332–2353, 2008.
- [24] C. Pellegrini. Continuous time open quantum random walks and non-Markovian Lindblad master equations. *J. Stat. Phys.*, 154(3):838–865, 2014.
- [25] C. Pellegrini and F. Petruccione. Non-Markovian quantum repeated interactions and measurements. *J. Phys. A*, 42(42):425304, 21, 2009.
- [26] D. Revuz and M. Yor. *Continuous martingales and Brownian motion*, volume 293 of *Grundlehren der Mathematischen Wissenschaften*. Springer-Verlag, Berlin, third edition, 1999.
- [27] I. Sinayskiy and F. Petruccione. Efficiency of open quantum walk implementation of dissipative quantum computing algorithms. *Quantum Inf. Process.*, 11(5):1301–1309, 2012.