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# ERRATUM: A CENTER MANIFOLD RESULT FOR DELAYED NEURAL FIELDS EQUATIONS

ROMAIN VELTZ\* AND OLIVIER FAUGERAS<sup>†</sup>

Abstract. Lemma C.1 in [R. Veltz and O. Faugeras, SIAM J. Math. Anal., 45(3) (2013), pp. 1527-1562] is wrong. This lemma is used in the proof of the existence of a smooth center manifold, Theorem 4.4 in [5]. An additional assumption is required to prove this existence. We spell out this assumption, correct the proofs and show that the assumption is satisfied for a large class of delay functions  $\tau$ . We also weaken the general assumptions on  $\tau$ .

Lemma C.1 in [5] is wrong as shown by the counterexample  $\phi(\theta, \bar{\mathbf{r}}) = \mathbb{1}_{\tau(\mathbf{r}_0, \cdot)^{-1}(\{\theta\})}(\bar{\mathbf{r}}), \theta \in [-\tau_m, 0], \bar{\mathbf{r}} \in \Omega$  for some  $\mathbf{r}_0 \in \Omega$  and sufficiently regular delay function  $\tau$ . This lemma is used in the proof of the regularity of  $\mathbf{R}$  in Lemma 4.2 of [5].

1. Corrections to the paper. To correct this problem requires choosing a slightly different functional setup from the one in the paper. We redefine the spaces  $\mathcal{X}^{(q)}$  and  $\mathcal{Y}^{(q)}$  for q > 2 (definition 2.4 of [5]) as

$$\begin{cases} \mathcal{X}^{(q)} \equiv \mathcal{L}^{\infty} \times \mathcal{L}^{q}(-\tau_{\mathrm{m}}, 0; \mathcal{L}^{\infty}), & \mathcal{L}^{\infty} \equiv \mathcal{L}^{\infty}(\Omega, \mathbb{R}^{\mathrm{p}}) \\ \mathcal{Y}^{(q)} \equiv \left\{ u \in \mathcal{L}^{\infty} \times \mathcal{W}^{1, \mathrm{q}}(-\tau_{\mathrm{m}}, 0; \mathcal{L}^{\infty}) \mid \pi_{1} u = (\pi_{2} u)(0) \right\} \end{cases}$$

and keep the original definition for q = 2:

$$\mathcal{X}^{(2)} \equiv \mathrm{L}^2 \times \mathrm{L}^2(-\tau_{\mathrm{m}}, 0; \mathrm{L}^2).$$

This choice does not alter the linear analysis (sections 1-3) in the paper but it affects a) Lemma B.2, b) Lemma 4.2, c) Theorem 4.4 and d) the main text in section 4.1, Lemma C.3 and Proposition C.4 as follows.

- a) Lemma B.2 needs to be proved for the new spaces  $\mathcal{X}^{(q)}$  as shown in section 3.
- b) Lemma 4.2 requires a different proof given in section 4 below. Lemma C.1 needs to be re-written in a way we also explain in section 4.
- c) In Theorem 4.4, the statement  $\Psi \in C^q(\mathcal{X}_c \times \mathbb{R}^{m_{par}}; \mathcal{Y}_h)$  becomes  $\Psi \in C^k(\mathcal{X}_c \times \mathbb{R}^{m_{par}}; \mathcal{Y}_h)$  (where  $S \in C^k(\mathbb{R}^p, \mathbb{R}^p)$ ).
- d) The main text in section 4.1 and Lemma C.3, Proposition C.4 (and their proofs) remain exactly the same modulo the change  $L^q \to L^{\infty}$  (*i.e.*  $\|\cdot\|_{L^q} \to \|\cdot\|_{L^{\infty}}$ ,  $L^q(-\tau_m, 0; L^q) \to L^q(-\tau_m, 0; L^{\infty})$  and  $W^{1,q}(-\tau_m, 0; L^q) \to W^{1,q}(-\tau_m, 0; L^{\infty})$ ).

**2. Preliminaries.** In order to modify and prove Lemma B.2 we need the following measure-theoretical preliminaries. We assume that  $\tau \in L^{\infty}(\Omega^2, \mathbb{R}^+)$ . For each  $\mathbf{x} \in \Omega$  define  $\tau_{\mathbf{x}} : \Omega \to [-\tau_m, 0]$  as  $\tau_{\mathbf{x}}(\mathbf{y}) = -\tau(\mathbf{x}, \mathbf{y})$  for all  $\mathbf{y} \in \Omega$ . We note  $\lambda_p$  the Lebesgue measure on  $\mathbb{R}^p$  and  $\tau_{\mathbf{x}} * \lambda_p$  the pushforward measure<sup>1</sup> of  $\lambda_p$  by  $\tau_{\mathbf{x}}$ , i.e. the measure on  $[-\tau_m, 0]$  such that for each Borelian B of  $[-\tau_m, 0], \tau_{\mathbf{x}} * \lambda_p(B)$  is equal to  $\lambda_p(\tau_{\mathbf{x}}^{-1}(B))$ .

Note that  $\tau_{\mathbf{x}} * \lambda_p([-\tau_m, 0]) \leq \lambda_p(\Omega)$  for all  $\mathbf{x} \in \Omega$  and that for all measurable function f on  $[-\tau_m, 0]$  we have the equality:

$$\int_{\Omega} f(-\tau(\mathbf{x}, \mathbf{y}))) d\lambda_p(\mathbf{y}) = \int_{-\tau_m}^0 f(\theta) d\left(\tau_{\mathbf{x}} * \lambda_p\right)(\theta)$$

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<sup>&</sup>lt;sup>1</sup> Legitimate because  $\tau_{\mathbf{x}}$  is measurable

whenever  $f \circ \tau_{\mathbf{x}}$  is  $\lambda_p$  integrable [2][th. 3.6.1]. As  $\tau_{\mathbf{x}} * \lambda_p$  and  $\lambda_1$  are  $\sigma$ -finite, the Lebesgue-Radon-Nikodým theorem give the following decomposition

$$\forall \mathbf{x} \in \Omega, \quad \frac{1}{\lambda_p(\Omega)} \tau_{\mathbf{x}} * \lambda_p = \mu_{\mathbf{x}}^{(abs)} + \mu_{\mathbf{x}}^{(at)} + \mu_{\mathbf{x}}^{(sing)}$$
(2.1)

where  $\mu_{\mathbf{x}}^{(abs)}$ ,  $\mu_{\mathbf{x}}^{(at)}$  and  $\mu_{\mathbf{x}}^{(sing)}$  are three measures on  $[-\tau_m, 0]$  such that  $\mu_{\mathbf{x}}^{(abs)}$  is absolutely continuous w.r.t  $\lambda_1$ , with density  $g_{\mathbf{x}}$ ,  $\mu_{\mathbf{x}}^{(at)}$  is atomic and  $\mu_{\mathbf{x}}^{(sing)}$  is continuous singular. We have:

$$\mu_{\mathbf{x}}^{(at)} = \sum_{n} a_n(\mathbf{x}) \delta_{-D_n(\mathbf{x})}.$$

If we define  $D_{\mathbf{x}} = \left\{ \theta \in [-\tau_m, 0] \mid \lambda_p\left(\tau_{\mathbf{x}}^{-1}(\theta)\right) > 0 \right\}$ , then  $D_{\mathbf{x}}$  is at most countable [3] [XIII.18.6], hence we write  $D_{\mathbf{x}} = (D_n(\mathbf{x}))_n$ . We make the following hypothesis (justified in section 3.1)

(H1) 
$$\forall \mathbf{x} \in \Omega \quad \mu_{\mathbf{x}}^{(sing)} = 0, \ D_n(\mathbf{x}) = D_n \text{ and } a_n(\mathbf{x}) = a_n.$$

In this case, we can write for  $f \tau_{\mathbf{x}} * \lambda_p$ -integrable

$$\forall \mathbf{x}, \quad \frac{1}{\lambda_p(\Omega)} \int_{\Omega} f(-\tau(\mathbf{x}, \mathbf{y})) d\lambda_p(\mathbf{y}) = \int_{-\tau_m}^0 f(\theta) g_{\mathbf{x}}(\theta) d\lambda_1(\theta) + \sum_n a_n f(-D_n).$$

LEMMA 2.1. The two components  $\mu_{\mathbf{x}}^{(abs)}$  and  $\mu_{\mathbf{x}}^{(at)}$  in the decomposition of the measure  $\tau_{\mathbf{x}} * \lambda_p$  satisfy:

•  $\forall \mathbf{x} \in \Omega, \ 0 \le g_{\mathbf{x}} \le 1 \ a.e. \ and \ \sup_{\mathbf{x}} \|g_{\mathbf{x}}\|_1 \le 1.$  It implies  $\sup_{\mathbf{x}} \|g_{\mathbf{x}}\|_q \le 1$  for all

• 
$$0 \le a_n \le 1$$
 and  $\sum_n a_n \le 1$ .

*Proof.* This is a consequence of  $\tau_{\mathbf{x}} * \lambda_p$  being finite and positive.

3. Correction of Lemma B.2 for the new spaces  $\mathcal{X}^{(q)}$ . The domain of  $S_t, T_0(t)$ is changed from  $L^q$  to  $L^\infty$  (see [1]).

LEMMA 3.1. (Lemma B.2 of [5]) Assume that (H1) is satisfied and that  $\mathbf{J} \in \mathrm{L}^{\infty}(\Omega^2, \mathbb{R}^{p \times p})$ . Then for each space  $\mathcal{X}^{(q)}$  with  $2 \leq q < \infty$ , there exists  $Q : \mathbb{R}_+ \to \mathbb{R}_+$  with  $\lim_{t \to 0^+} Q(t) = 0$ such that

$$\forall \begin{bmatrix} x \\ \phi \end{bmatrix} \in D(\mathbf{A}_{(q)}) \quad \int_0^t \|\mathbf{L}_1(S_s x + T_0(s)\phi)\|_{\mathbf{L}^\infty} \, ds \le Q(t) \left\| \begin{bmatrix} x \\ \phi \end{bmatrix} \right\|_{\mathcal{X}^{(q)}} \tag{M}$$

*Proof.* The case q = 2 is proved in the paper. Let us focus on the case q > 2. We first note that for  $\phi \in W^{1,q}(-\tau_m, 0; L^{\infty})$ :

$$\|\phi(s-\tau(\mathbf{x},\mathbf{y}),\mathbf{y})\|_{\mathbb{R}^p} \le \|\phi(s-\tau(\mathbf{x},\mathbf{y}))\|_{\mathrm{L}^{\infty}}.$$
(3.1)

We focus on the second term  $\int_0^t \|\mathbf{L}_1(T_0(s)\phi)\|_{\mathbf{L}^{\infty}} ds$  which is the most difficult to handle. Indeed, the first term  $\int_0^t \|\mathbf{L}_1(S_s x)\|_{\mathbf{L}^{\infty}} ds$  is bounded by  $Kt \|x\|$  for some constant K because the norm of  $\mathbf{L}_1 S_s$  on  $\mathbf{L}^{\infty}$  is bounded by a constant independent of s. By definition of  $\mathbf{L}_1, \mathbf{T}_0$ :

$$\left(\mathbf{L}_{1}(T_{0}(s)\phi)\right)(\mathbf{x}) = \int_{\Omega} \mathbf{1}_{[0,\tau(\mathbf{x},\mathbf{y})]}(s)\mathbf{J}(\mathbf{x},\mathbf{y})\phi(s-\tau(\mathbf{x},\mathbf{y}),\mathbf{y})d\lambda_{p}(\mathbf{y}),$$
2

which gives:

$$\|\left(\mathbf{L}_1(T_0(s)\phi)\right)(\mathbf{x})\|_{\mathbb{R}^p} \le |||\mathbf{J}|||_{\infty} \int_{\Omega} \mathbf{1}_{[0,\tau(\mathbf{x},\mathbf{y})]}(s) \|\phi(s-\tau(\mathbf{x},\mathbf{y}))\|_{\mathbb{R}^p} d\lambda_p(\mathbf{y}).$$

For any given function  $f \in L^1([-\tau_m, 0], \mathbb{R})$ , we extend it to a function of  $L^1([-\tau_m, \tau_m], \mathbb{R})$  by setting f = 0 on  $[0, \tau_m]$  so that we do not have to worry about the integral bounds. According to (2.1) we have

$$\frac{1}{\lambda_{p}(\Omega)} \int_{\Omega} \mathbf{1}_{[0,\tau(\mathbf{x},\mathbf{y})]}(s) \|\phi(s-\tau(\mathbf{x},\mathbf{y}))\|_{\mathbf{L}^{\infty}} d\lambda_{p}(\mathbf{y}) = \frac{1}{\lambda_{p}(\Omega)} \int_{\Omega} \mathbf{1}_{[0,-\tau_{\mathbf{x}}(\mathbf{y})]}(s) \|\phi(s+\tau_{\mathbf{x}}(\mathbf{y}))\|_{\mathbf{L}^{\infty}} d\lambda_{p}(\mathbf{y}) = \frac{1}{\lambda_{p}(\Omega)} \int_{\Omega} \mathbf{1}_{[0,-\theta]}(s) \|\phi(s+\theta)\|_{\mathbf{L}^{\infty}} d(\tau_{\mathbf{x}} * \lambda_{p})(\theta) = \int_{-\tau_{m}}^{0} \mathbf{1}_{[0,-\theta]}(s) \|\phi(s+\theta)\|_{\mathbf{L}^{\infty}} g_{\mathbf{x}}(s+\theta) d\lambda_{1}(\theta) + \sum_{n} a_{n} \mathbf{1}_{[0,D_{n}]}(s) \|\phi(s-D_{n})\|_{\mathbf{L}^{\infty}}.$$
 (3.2)

From Lemma 2.1

$$\sup_{\mathbf{x}} \int_{-\tau_m}^0 \mathbf{1}_{[0,-\theta]}(s) \|\phi(s+\theta)\|_{\mathbf{L}^{\infty}} g_{\mathbf{x}}(s+\theta) d\lambda_1(\theta) \stackrel{\text{H\"older}}{\leq} \|\phi\|_{\mathbf{L}^{\mathbf{q}}(-\tau_{\mathbf{m}},0;\mathbf{L}^{\infty})} d\lambda_1(\theta) \|_{\mathbf{L}^{\mathbf{q}}(-\tau_{\mathbf{m}},0;\mathbf{L}^{\infty})} d\lambda_1(\theta) \|_{\mathbf{L}^{\mathbf{m}}(-\tau_{\mathbf{m}},0;\mathbf{L}^{\infty})} d\lambda_1(\theta) \|_{\mathbf{L}^{\mathbf{m}}$$

As  $\phi \in D(\mathbf{A}_{(q)})$ , it belongs to  $W^{1,q}(-\tau_m, 0; \mathbf{L}^{\infty})$ . Hence, the function  $\theta \to \|\phi(\theta)\|_{\mathbf{L}^{\infty}}$  is continuous on  $[-\tau_m, 0]$  and its supremum is a max attained at  $\theta = -D_{max}$ . This gives:

$$\sum_{n} a_{n} \mathbf{1}_{[0,D_{n}]}(s) \|\phi(s-D_{n})\|_{\mathbf{L}^{\infty}} \le \|\phi(-D_{max})\|_{\mathbf{L}^{\infty}}$$

It then follows from the Beppo-Levi's theorem and Lemma 2.1 that

$$\int_{0}^{\min(t,\tau_m)} \sum_{n} a_n \mathbf{1}_{[0,D_n]}(s) \|\phi(s-D_n)\|_{\mathbf{L}^{\infty}} ds = \sum_{n} a_n \int_{0}^{\min(t,\tau_m)} \mathbf{1}_{[0,D_n]}(s) \|\phi(s-D_n)\|_{\mathbf{L}^{\infty}} ds$$

$$\stackrel{\text{Hölder}}{\leq} (\min(t,\tau_m))^{\frac{1}{q}} \|\phi\|_{\mathbf{L}^q(-\tau_m,0;\mathbf{L}^{\infty})},$$

where  $\bar{q}$  is the Hölder conjugate integer of q. Summing up, we have found that:

$$\int_{0}^{t} \|\mathbf{L}_{1}(T_{0}(s)\phi)\|_{\mathbf{L}^{\infty}} ds \stackrel{\text{def. of } T_{0}}{=} \int_{0}^{\min(t,\tau_{m})} \|\mathbf{L}_{1}(T_{0}(s)\phi)\|_{\mathbf{L}^{\infty}} ds \leq \lambda_{p}\left(\Omega\right) |||\mathbf{J}|||_{\infty} \left(t + \min(\tau_{m},t)^{1/\bar{q}}\right) \|\phi\|_{\mathbf{L}^{q}(-\tau_{m},0;\mathbf{L}^{\infty})}$$
(3.3)

which gives

$$\int_0^t ds \left\| \mathbf{L}_1(S_s x + T_0(s)\phi) \right\|_{\mathbf{L}^{\infty}} \le K_2 \max\left(t, \min(\tau_m, t)^{1/\bar{q}}\right) \left\| \begin{bmatrix} x \\ \phi \end{bmatrix} \right\|_{\mathcal{X}^{(q)}}.$$

for some constant  $K_2$ . This ends the proof.  $\Box$ 

**3.1. Example of possible delay functions.** Let us show that the assumptions of lemma 3.1 are satisfied for some realistic delay function. Apart from (H1), the only requirement has been that

$$\tau \in \mathcal{L}^{\infty}(\Omega^2, \mathbb{R}^+).$$

The next lemma shows that (H1) holds for a large class of delay functions that includes a combination of constant and propagation delays,

LEMMA 3.2. Let us consider  $\tau(\mathbf{x}, \mathbf{y}) = D + c\kappa(\mathbf{x}, \mathbf{y})$  with  $c, \kappa \geq 0$ . We assume that  $\forall \mathbf{x} \in \Omega, \kappa(\mathbf{x}, \cdot) \in C^1(\overline{\Omega}, \mathbb{R}^+)$  and that  $\forall \mathbf{x}$ , the gradient of  $\kappa(\mathbf{x}, \mathbf{y})$  w.r.t.  $\mathbf{y}$  is non zero almost everywhere. Then  $\tau$  satisfies (H1).

*Proof.* Straightforward application of integration theory on submanifolds.

#### 4. Correction of Lemmas C.1 and 4.2. The new lemma C.1 reads as follows.

LEMMA 4.1. (Lemmas B.1, C.1 of [5]) Assume that  $\tau \in L^{\infty}(\Omega^2, \mathbb{R}^+)$  and  $\mathbf{J} \in L^{\infty}(\Omega^2, \mathbb{R}^{p \times p})$ , then we have the following results:

1. Define  $\mathbf{J}[s]$  by  $\forall s \in [-\tau_m, 0]$ ,  $J_{ij}(\mathbf{r}, \mathbf{r}')[s] \equiv J_{ij}(\mathbf{r}, \mathbf{r}')H(s + \tau_{ij}(\mathbf{r}, \mathbf{r}'))$  where H is the Heaviside function. Then:

$$\mathbf{L}_{1}\phi = \mathbf{J}\phi(0) - \int_{-\tau_{m}}^{0} \mathbf{J}[s]\dot{\phi}(s)ds \quad \forall \phi \in \mathbf{W}^{1,\mathbf{q}}(-\tau_{m},0;\mathbf{L}^{\infty}), \text{for all } 2 \leq \mathbf{q} < \infty.$$
(4.1)

- 2.  $\mathbf{L}_1 \in \mathcal{L} \left( \mathbf{W}^{1,\mathbf{q}}(-\tau_{\mathbf{m}}, 0; \mathbf{L}^{\infty}), \mathbf{L}^{\infty} \right)$  for all  $2 \leq q < \infty$ .
- 3.  $\forall l \in \mathbb{N}, (\phi_1, \cdots, \phi_l) \to \mathbf{L}_1 (\phi_1 \cdots \phi_l)$  is linear continuous from  $(W^{1,q}(-\tau_m, 0; L^\infty))^l$  to  $L^\infty$ .

Proof.

1. Let us consider  $\phi \in W^{1,q}(-\tau_m, 0; L^{\infty})$  with  $2 \leq q < \infty$ . From the definition of Bochner spaces, we have

$$\phi_j(\bar{\mathbf{r}},\theta) = -\int_{\theta}^{0} \dot{\phi}_j(\bar{\mathbf{r}},s) ds + \phi_j(\bar{\mathbf{r}},0) \text{ for almost all } \bar{\mathbf{r}} \in \Omega$$

which gives:

$$\phi_j(\bar{\mathbf{r}},\theta) = -\int_{-\tau_m}^0 \dot{\phi}_j(\bar{\mathbf{r}},s) H(s-\theta) ds + \phi_j(\bar{\mathbf{r}},0) \text{ for almost all } \bar{\mathbf{r}} \in \Omega$$
(4.2)

Moreover,  $\forall \mathbf{r} \in \Omega, \theta \to H(\theta + \tau(\mathbf{r}, \cdot)) \in L^{\infty}(-\tau_{\mathrm{m}}, 0; \mathbb{R}^{+}) \subset L^{\infty}(-\tau_{\mathrm{m}}, 0; \mathrm{L}^{\infty})$ . This shows that  $\forall \mathbf{r} \in \Omega, \theta \to \dot{\phi}_{j}(\cdot, \theta) H(\theta + \tau(\mathbf{r}, \cdot)) \in \mathrm{L}^{q}(-\tau_{\mathrm{m}}, 0; \mathrm{L}^{\infty})$  with  $q \geq 2$ . From (4.2) and the definition of the Bochner integral, it follows that  $\phi(\cdot, -\tau(\mathbf{r}, \cdot)) \in \mathrm{L}^{\infty}$ . As  $\mathbf{J} \in \mathrm{L}^{\infty}(\Omega^{2}, \mathbb{R}^{p \times p})$ , it implies that  $\mathbf{L}_{1}$  is well defined and  $\mathbf{L}_{1}\phi \in \mathrm{L}^{\infty}$ .

Plugging (4.2) in the expression of  $\mathbf{L}_1$  and  $\operatorname{using}^2$  [4][proposition C.4] gives the equality of the lemma with  $\mathbf{J}(\mathbf{r}, \mathbf{r}')[s] \in L^{\infty}(\Omega^2, \mathbb{R}^{p \times p})$ .

- 2. If  $\mathbf{J} \in L^{\infty}(\Omega^2, \mathbb{R}^{p \times p})$ , it is straightforward to show that  $\mathbf{L}_1 \in \mathcal{L}(W^{1,q}(-\tau_m, 0; L^{\infty}), L^{\infty})$  by using (4.1).
- 3. This is a consequence of 2. and the fact that  $W^{1,q}(-\tau_m, 0; L^{\infty})$  is a Banach algebra.

<sup>2</sup>Basically  $\mathbf{J} \int_0^\theta \dot{\phi} = \int_0^\theta \mathbf{J} \dot{\phi}$ 

This allows us to obtain a corrected version of Lemma 4.2 of [5].

LEMMA 4.2. (Lemma 4.2 of [5]) If  $\mathbf{J} \in \mathcal{L}^{\infty}(\Omega^2, \mathbb{R}^{p \times p})$ ,  $\mathbf{S} \in \mathcal{C}^k(\mathbb{R}^p, \mathbb{R}^p)$  and  $2 < q < \infty$ . Then

$$\begin{array}{lll} \mathbf{A}_{(q)} & \in & \mathcal{L}(\mathcal{Y}^{(q)}, \mathcal{X}^{(q)}), \\ \mathbf{R} & \in & \mathrm{C}^{k}(\mathcal{Y}^{(q)} \times \mathbb{R}^{m_{par}}, \mathcal{X}^{(q)}), \end{array}$$

and

$$D_{u}^{l}\mathbf{R}(u_{0},\mu)[u_{1},\cdots,u_{l}] = \begin{bmatrix} \mathbf{L}_{1}(\mu)\mathbf{S}^{(l)}(\mathbf{V}^{f}+\pi_{2}u_{0})\pi_{2}(u_{1}\cdots u_{l})\\ 0 \end{bmatrix}, \ l = 1,\cdots,k$$

where  $u_1 \cdots u_l$  is the component-wise product of the functions  $u_i$  in  $\mathcal{Y}^{(q)}$ .

Proof.

**Case of A**<sub>(q)</sub>. The only "difficulty" is showing that  $\phi \to \mathbf{L}_1 \cdot (D\mathbf{S}(\mathbf{V}^f)\phi)$  belongs to  $\mathcal{L}(W^{1,q}(-\tau_m, 0; \mathbf{L}^\infty), \mathbf{L}^\infty)$ . This was done in Lemma 4.1.2.

**Case of R.** Recall that  $\pi_2 \mathbf{R} = 0$  while

$$\pi_1 \mathbf{R}(u,\mu) \stackrel{def}{=} \mathbf{L}_1(\mu) \cdot \left[ \mathbf{S} \left( \mathbf{V}^f + \pi_2 u \right) - \mathbf{S} \left( \mathbf{V}^f \right) \right] - \mathbf{L}_1(\mu_c) \cdot D\mathbf{S} \left( \mathbf{V}^f \right) \pi_2 u$$

We focus on the differentiability at u = 0 and ignore the differentiability w.r.t. the parameter  $\mu$ . The differentiability at  $u \neq 0$  follows from the same argument. It is easy to see from the definition of **R** that, since **S** is  $C^k$  and **L**<sub>1</sub> is bounded (lemma 4.1.3)

$$\pi_1 D^l \mathbf{R}(0,\mu)[u_1,\cdots,u_l] = \begin{cases} (\mathbf{L}_1(\mu) - \mathbf{L}_1(\mu_c)) \cdot \mathbf{S}^{(1)}(\mathbf{V}^f) \pi_2 u_1 & l = 1 \\ \mathbf{L}_1(\mu) \cdot \mathbf{S}^{(l)}(\mathbf{V}^f) \pi_2(u_1\cdots u_l) & l = 2,\cdots,k \end{cases}$$

The proof that  $\pi_1 R(u, \mu)$  is  $C^k$  at u = 0 then follows from the fact that **S** is  $C^k$  and **L**<sub>1</sub> is a bounded operator:

$$\pi_1 \mathbf{R}(h,\mu) - \sum_{l=1}^k \frac{1}{l!} \pi_1 D^l \mathbf{R}(0,\mu) \pi_2 h^l = \mathbf{L}_1(\mu) \cdot \left( \mathbf{S}(\mathbf{V}^f + \pi_2 h) - \mathbf{S}(\mathbf{V}^f) - \sum_{l=1}^k \frac{\mathbf{S}^{(l)}(\mathbf{V}^f)}{l!} \pi_2 h^l \right)$$

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