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## ON THE SUPERIOR DYNAMIC BEHAVIOR OF A VARIABLE ROTATING SPEED SCROLL COMPRESSOR

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#### ABSTRACT

This study compares the calculated results for the dynamic behavior of a comparatively small cooling-capacity variable-speed scroll compressor with those of a variablespeed rolling-piston rotary compressor of the same cooling capacity. The superior dynamic behavior of the scroll compressor could thus be shown. In order to derive the characteristic of the rotating speed, the rotating-speed fluctuation ratio and the rotatory acceleration, first, the equation of motion of the rotating crankshaft is solved numerically for different mean rotating speeds. Thus, the unbalanced forces of inertia exciting the whole compressor are calculated and the vibration characteristics of the compressor can be shown. In the numerical calculation, the driving torque, the gas-compression load, the frictional coefficients, the mass and the moment of inertia of the compressors are kept constant, thus the calculated results for the two different compressors can be compared.

#### INTRODUCTION

In recent years, many domestic electrical appliances with an inverter system has come into use. The inverter system is a device which converts a commercial alternating current into one with adjustable frequency and voltage. The inverter system is also used in air-conditioning systems to adjust the mean rotating speed of the compressors driven by the induction motor, and thus an air-conditioning system with variable speed compressor is formed. In the conventional air-conditioning systems, the compressors are operated with a constant mean rotating speed and thus the motor is switched off and on to regulate the quantity of cool air which is required for keeping the room temperature at a desired value. The switching off and on of the motor may be accompanied by largeamplitude vibrations of the compressor. As a result, the compressor hits the enclosing housing to make the noisy impulsive sounds in the reciprocating compressors [see Refs. 1-6], and in the rolling-piston rotary compressors [7] it also induces large vibrations of the pipes and panels surrounding the compressor [8]. The variable speed amplitude air-conditioning system has no vibration problems such as those which occur in compressors with constant speed operation, since the variable speed compressor can deliver the optimum quantity of cool air by adjusting the mean speed of operation,

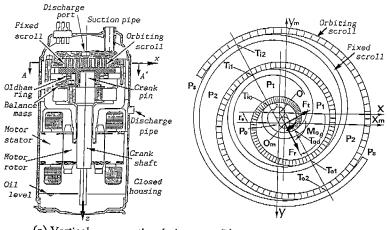
#### NOMENCLATURE

e=rotating radius of rolling piston ro=orbiting radius r<sub>Q</sub>,r<sub>S</sub>=radius of crankshaft and crankpin fi=frictional force at thrust bearing  $f_{xi}, f_{yi}$ =frictional force at Oldham ring [F]=exciting force matrix  $\hat{S}_{x}, \bar{S}_{y}$ =reaction force on crankpin t=scroll thickness Ti=reaction force between Oldham ring and  $F_r$ =radial gas force Ft=tangential gas force fixed guide slot Fti=thrust force on thrust bearing x,y,z=static rectangular coordinate xm, ym=orbiting coordinate i=integer Ic, Io=moment of inertia of crankshaft xy=coordinate of reciprocating blade Ip=moment of inertia of rolling piston [X]=displacement matrix of compressor Ix,Iy,Iz=moment of inertia of compressor  $\alpha$ =rotating-speed fluctuation ratio [K]=spring coefficient matrix  $\beta$ =ratio of tangetial acc. and centrifugail one Lo=frictional torque at crankshaft Ls=frictional torque ar crankpin [y]=damping coefficient matrix Io=height of Oldham ring  $\theta$ =crank angle mo-mass of Oldham ring  $\theta$ =rotating mean speed of crankshaft mo=mass of orbiting scroll my=mass of blade  $\Theta_z$ =rotation around z axis of compressor [M]=mass matrix η<sub>o</sub>=viscosity coefficient of oil N=motor torque O;=reaction force between Oldham ring and µo=frictional coefficient at Oldham ring orbiting scroll µO=frictional coefficient at crankshaft p<sub>1</sub>,p<sub>2</sub>,p<sub>c</sub>=compressed gas pressure ps,p0,pd=suction and discharge pressure µs=frictional coefficient at crankpin Qx, Qy=reaction force on crankshaft

without switching the motor off and on. In addition to this point, the variable airconditioning system has the following advantages: The energy loss in a variable speed air-conditioning system is much less than the switching-on and -off system, and thus it may be said that the variable speed system is an energy-saving system. Furthermore, the variable speed system makes it possible to control the room temperature quickly, for instance, by operating the compressor at a higher speed initially.

The variable speed type air-conditioners began to be produced in 1981 and at present about 50 % of the 2 million heat-pump type air-conditioners on the Japanese market is the variable speed type. This percentage will increase more in the future because of the variable speed air-conditioning system's great advantages. However, it should be noticed that the variable-speed type also has a tendency to induce the following vibration problem. The vibration amplitude of displacement is inversely proportional to the vibration frequency, and hence the compressor vibration displacement will increase as the operat-ing speed decreases. Moreover, the operating speed comes closer to the natural fre-quencies of the compressor supporting system, and thus the amplitude of compressor vibration has a tendency to increase due to a resonance effect. These considerations suggest that the scroll compressors with low-level vibration characteristics will be most suitable for the variable speed air-conditioning system. The scroll compressor has many compression chambers which are simultaneously compressed at comparatively low speed. Thus, the unbalanced moment of inertia about the rotating shaft is fairly small compared with that of the rolling-piston rotary compressor. Therefore, the scroll compressor will be far superior with regard to vibration than the rolling-piston rotary compressor .

Previous studies [9-11] have discussed the superior dynamic behavior of a constant speed scroll compressor. In this study, the dynamic behavior of a variable speed scroll compressor having a small cooling capacity is calculated numerically. Thus, the calculated results of the fluctuating crankshaft speed, the rotating speed fluctuation ratio, the crankshaft rotatory acceleration, the unbalanced forces and moments of inertia and the vibration displacements and accelerations are compared with those of a rolling-piston rotary compressor having the same cooling capacity. In the numerical calculation, the driving torque, the gas-compression load, the frictional coefficients, the mass and the moment of inertia of the compressors are kept constant, in order to compare the calculated results for the two different compressors.



(a) Vertical cross-sectional view (b) A-A' cross-sectional view Fig.1 Scroll Compressor

# EQUATION OF MOTION AND UNBALANCED FORCES OF INERTIA

#### Scroll Compressor

The vertical cross-sectional view of a scroll compressor is shown in Fig. 1a. The compression mechanism is above the motor. The A-A' cross-sectional view of the compression mechanism is shown in Fig. 1b, in which the mating of the two identical scrolls formed by the involutes with the basic-circle radius  $r_b$  can be seen. The unshaded scroll is fixed. The basic circle center O of the fixed scroll is the origin of the static x-y coordinate. z is the downward axis along the crankshaft center. The shaded scroll can move in the x direction, relatively to the Oldham ring which can move in the y direction along a fixed guide slot. Thus, the shaded scroll orbits when the basic-circle center  $O_m$  of the shaded scroll rotates around the fixed point O. The orbiting radius  $r_0$  is given by  $(r_b \pi$ t) where t is the scroll thickness. The clockwise orbiting angle from the x axis is represented by  $\theta$ . The contact points  $T_{ik}$  and  $T_{ok}$  (k=0,1,2...) appear respectively on the two tangents common to the two basic circles. The two scrolls form a series of crescentshaped compression chambers. The pressures  $p_s$ ,  $p_1$ ,  $p_2$  and  $p_0$  acting on the orbiting scroll are reduced to the tangential force  $F_t$  which acts against the clockwise orbiting motion, the radial force  $F_r$  and the clockwise moment  $M_o$  which act on the orbiting scroll center  $O_m$ . They are given by (4)-(6) in the previous study [9]. When the crankshaft is driven by the motor torque N, the equation of motion of the rotating crankshaft is given by the following [9, 11]:

$$(I_0 + m_{\bullet} r_0^2 + m_o r_0^2 \sin^2 \theta) \ddot{\theta} + m_o r_0^2 \sin \theta \cdot \cos \theta \cdot \dot{\theta}^2$$

$$= N - \{P_{\ell}r_{0} + L_{0} + L_{s} + (f_{z1} + f_{z2})r_{0}\sin\theta + (f_{y1} + f_{y2})r_{0}\cos\theta + (f_{\ell 1} + f_{\ell 2})r_{0}\}$$
(1)

where  $I_0$  is the moment of inertia of the crankshaft system comprising the motor rotor,  $m_s$  the mass of the orbiting scroll and  $m_o$  the mass of the Oldham ring.  $L_Q$  and  $L_S$  are the frictional torques on the crankshaft and the crankpin, respectively.  $f_{xi}$  and  $f_{yi}$  (i=1,2) are the frictional forces acting on the Oldham ring. The subscripts x and y represent the directions of the frictional forces.  $f_{ti}$  (i=1,2) are the frictional forces at the thrust bearing. They are given by the following expressions, assuming Coulomb's law of friction.

$$L_{q} = \mu_{q} \sqrt{Q_{x}^{2} + Q_{y}^{2}} \cdot r_{q}, \quad L_{s} = \mu_{s} \sqrt{S_{x}^{2} + S_{y}^{2}} \cdot r_{s}$$

$$f_{xi} = \mu_{o} |T_{i}|, \quad f_{yi} = \mu_{o} |O_{i}|, \quad f_{ti} = \mu_{t} |E_{t}|, \quad (i = 1, 2)$$
(2)

where  $Q_x$ ,  $Q_y$  are the reaction forces on the crankshaft,  $S_x$ ,  $S_y$  those on the crankpin,  $T_i$  (i=1,2) those between the Oldham ring and the fixed guide slot and  $O_i$  (i=1,2) those be

tween the orbiting scroll and the Oldham ring. They can be easily derived from the previous study [9]. These reaction forces include the frictional forces of (2) and hence an iterative calculation method should be adopted for calculating the frictional forces and torques of (2).  $F_{ii}$  (i=1,2) are the reaction forces on the thrust bearing and are given by (18) in the previous study [9].

The forces and moments exciting the compressor vibrations are given by the following expressions, namely, unbalanced forces and moments of inertia:

$$F_{x} = -(m_{o}/2)r_{0}\left(\dot{\theta}^{2}\cos\theta + \ddot{\theta}\sin\theta\right), F_{y} = -(m_{o}/2)r_{0}\left(\dot{\theta}^{2}\sin\theta - \ddot{\theta}\cos\theta\right), F_{z} = 0,$$

$$M_{x} = (m_{o}/2)l_{o}r_{0}\left(\dot{\theta}^{2}\sin\theta - \ddot{\theta}\cos\theta\right), M_{y} = (m_{o}/2)l_{o}r_{0}\left(\dot{\theta}^{2}\cos\theta + \ddot{\theta}\sin\theta\right), \qquad (3)$$

$$M_{z} = -(I_{o} + m_{z}r_{o}^{2})\ddot{\theta}$$

in which the best static balancing and the best dynamic balancing have been achieved by choosing optimum values of the balancing masses. However, as shown in Fig.1a, the orbiting scroll is supported from the bottom only, and hence the perfect dynamic balancing cannot be achieved.  $F_x$ ,  $F_y$  and  $F_z$  are the forces acting on the fixed-coordinate origin O.  $M_x$ ,  $M_y$  and  $M_z$  are moments around the x, y and z-axes respectively.

#### Rolling-Piston Rotary Compressor

The vertical cross-sectional view of a rollingpiston rotary compressor is shown in Fig.2. A static rectangular coordinate with the origin at the cylinder center is represented by x, y and zaxes. x axis coincides with the center line of the reciprocating blade. z is the upward axis along the crankshaft center. The equation of motion of the rotating crankshaft has been derived in the previous study [7]. Achieving the best static and dynamic balancing, the unbalanced forces and moments of inertia are given by the following expressions.

$$F_{x} = -m_{v}\dot{x}_{v} - (m_{v}/2)e\left(\dot{\theta}^{2}\cos\theta + \theta\sin\theta\right),$$
  

$$F_{y} = -(m_{v}/2)e\left(\dot{\theta}^{2}\sin\theta - \ddot{\theta}\cos\theta\right), \quad F_{z} = 0,$$
  

$$M_{x} = 0, \quad M_{y} = 0, \quad M_{z} = -(I_{c} + m_{p}e^{2})\ddot{\theta} - I_{p}\dot{\phi} \quad (4)$$

both sides of the cylinder and hence the perfect dy-The crankshaft is supported at namic balancing can be achieved.  $m_v$  and  $x_v$  represent the reciprocating-blade mass and its coordinate [see (2) of ref. 7], respectively. e is the rotating radius of the rolling piston with the mass  $m_p$ .  $I_c$  and  $I_p$  represent the moment of inertia of the crankshaft and the rolling piston, respectively.

#### NUMERICAL CALCULATIONS

#### Frictional coefficients

When a small cooling capacity rolling-piston rotary compressor was driven by a motor with shaft power of about 600 W, the measured mechanical efficiency was 90.3 % and the measured rolling piston mean speed was 32 rpm under a suction pressure of 0.54 MPa and discharge pressure of 2.06 MPa [12]. On the basis of these measured values, the lubrication at each pair of machine elements in the mechanism was evaluated theoretically, by using those into three major categories of lubrication: a fluid lubrication between the rolling piston and the crankpin, a comparatively well-conditioned boundary lubrication between the crankshaft and the crank journal, and a comparatively poorly-conditioned boundary lubrication between the blade and the rolling piston or the cylinder guide slot [13]. It was concluded from this study that the frictional coefficient at the crankshaft is 0.013 and the frictional coefficient at the blade is 0.083 when the oil viscosity coefficient is about 2.076 mPa.s. In order to compare the calculated results for the two different compressors, it is assumed that the frictional coefficients in the scroll compressor are as follows: The frictional coefficients  $\mu_Q$  and  $\mu_S$  at the comparatively well-lubricated crankshaft and crankpin are 0.013, and

 $\mu_0$  and  $\mu_t$  at the comparatively poorly-lubricated Oldham ring and thrust bearing are 0.083.

#### Motor Torque and Gas Pressure

The motor torques for five synchronous speeds from 30 to 150 Hz are shown by dotted lines in

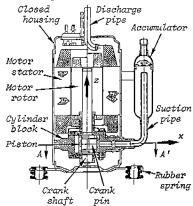


Fig.2 Rolling-piston rotary compressor

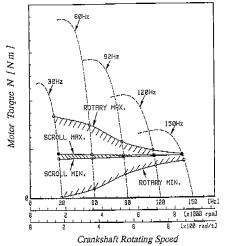


Table 1 Major mechanical constants

Scroll Comp.	Rotary comp.
$I_{o}=0.0422 Ncms^{2}$ $m_{o}=0.051 Kg$ $m_{s}=0.277 Kg$ $r_{o}=0.25 cm$ $r_{Q},r_{S}=1.0 cm$	e=0.326  cm $I_c=0.0422 \text{ Ncms}^2$ $I_p=0.0015 \text{ Ncms}^2$ $m_p=0.0741 \text{ Kg}$ $m_v=0.0113$
$I_{X}=3.519 \ Ncms^{2}$ $I_{Y}=3.904 \ Ncms^{2}$ $I_{Z}=0.996 \ Ncms^{2}$ $M=8.7 \ Kg$ $(x_{Gr}y_{Gr}z_{G})=(6.54 \ cm, \ 0, \ 0)$	

#### Fig.3 Motor torque and the range of torque fluctuation

Fig. 3. In the numerical calculations, these motor torque curves are approximated by a 4th-order polynomial expression for the rotating speed  $\theta$ . Assuming that the gas compression process is subjected to the adiabatic change of specific heat ratio 1.32 which is the value of the superheated Freon (R-22) vapor at a pressure of 1.32 MPa and a temperature of 55.6 °C, the pressures in the compression chambers are calculated. The suction pressure is 0.617 MPa and the discharge pressure is 2.17 MPa.

#### Mechanical Constants

Table 1 shows the major mechanical constants of the two compressors having a small cooling capacity. The maximum suction volume is  $10.26 \text{ cm}^3$ . The crankshaft moments of inertia  $I_0$  and  $I_c$  are the same. As shown in (3) and (4), the unbalanced forces and moments of inertia except for  $M_Z$  consist of the inertia terms of the reciprocating elements: the Oldham ring in the scroll compressor and the blade in the rolling-piston rotary compressor. Thus, it should be noticed that the mass  $m_0$  of the Oldham ring is about 5 times larger than the mass  $m_V$  of the blade. This difference is caused by the structural difference between the scroll compressor and the rolling-piston rotary compressor.

#### Calculated Results for Crankshaft Dynamic Behavior

The crankshaft dynamic behavior of the scroll compressor can be calculated numerically from the expressions (1) and (2), as follows: The initial values of  $\theta$  and  $\dot{\theta}$ , and the crankshaft mean rotating speed which is a little smaller than the synchronous speed are given first, and thus  $\theta$ ,  $\dot{\theta}$  and  $\ddot{\theta}$  over one period are calculated. The initial value of  $\dot{\theta}$  is modified if  $\theta$  does not increase in one period by one revolution  $2\pi$ , and the crankshaft mean rotating speed is modified if  $\dot{\theta}$  after one period does not coincide with the initial value, to repeat the same numerical calculation over one period again. Thus, the periodical dynamic behavior of the crankshaft is obtained. The range of fluctuation of the motor torque is shown in *Fig.3*, in which the calculated results for the same cooling capacity rolling-piston rotary compressor are shown. The torque fluctuation of the scroll compressor is far smaller than the crankshaft rotating speed fluctuation ratio  $\alpha$  of the scroll compressor is far smaller than the rolling-piston rotary compressor, as shown in *Fig.4*. As the synchronous speed decreases,  $\alpha$  for both compressors increase, but the increasing ratio for the rotary compressor is far larger than the scroll compressor.

The increase of  $\alpha$  for the decreasing synchronous speed is caused by the following character of the crankshaft rotatory acceleration  $\ddot{\theta}$ : Fig. 5 shows the peak to peak value of the fluctuating acceleration  $\ddot{\theta}$ . The peak to peak value has a tendency to decrease, as the synchronous speed decrease, but the decrease is quite small for both compressors. As seen from (1),  $\ddot{\theta}$  depends mainly upon the fluctuating gas load (F<sub>t</sub>r<sub>o</sub>) which is kept constant for each synchronous speed, and thus  $\ddot{\theta}$  is fundamentally constant independently of the synchronous speed. The small decrease of  $\ddot{\theta}$  for the decreasing synchronous speed is caused mainly by a decrease of frictional torques and forces which are partially af-

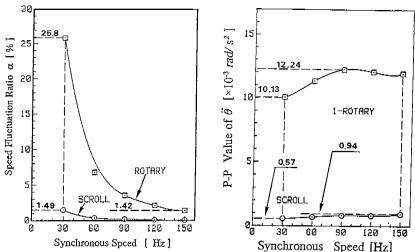


Fig.4 Crankshaft speed fluctuation ratio Fig.5 Peak-peak value of rotatoryacceleration fected by the centrifugal acceleration. The centrifugal acceleration decreases, of course, as the synchronous speed decreases.

#### Compressor Vibration

The compressor vibration is subjected to the following matrix equation of motion:

$$[M][\ddot{X}] + [\gamma][\dot{X}] + [K][X] = [E][F]$$
(5)

where [X] represents the displacements of the compressor gravity center, which consist of X, Y, Z,  $\Theta_X$ ,  $\Theta_Y$  and  $\Theta_Z$ . [M] is the mass matrix consisted of M,  $I_X$ ,  $I_Y$  and  $I_Z$  which are given in Table 1. [ $\gamma$ ] is the damping coefficient matrix and [K] the spring constant matrix. [E] is a transfer matrix which consists of the gravity coordinate  $x_G$ ,  $y_G$  and  $z_G$  (see Table 1) and is given by (28) in the previous study [9]. Assume that the damping effect of coiled springs can be neglected and the undamped natural frequencies are 8.6 Hz in the x and y directions, 15 Hz in the z direction, 15.1 Hz around the x axis, 11.9 Hz around the y axis and 13.6 Hz around the z axis. Thus, the compressor vibrations can be calculated.

As shown in (3) and (4), the unbalanced forces and moments except for  $M_z$  consist of both accelerations  $\dot{\theta}^2$  and  $\ddot{\theta}$ . It is important to find that even when the synchronous speed is 30 Hz, the maximum value of the centrifugal acceleration  $\dot{\theta}^2$  is about 7 times larger than that of the rotatory acceleration  $\ddot{\theta}$ , that is, the exciting forces and moments except for  $M_z$  depend mainly upon the centrifugal forces of the unbalanced mass, which vary proportionally to the quadratic rotating speed. On the other hand, the vibration displacements vary inversely proportional to the quadratic rotating speed. Thus, the amplitude of the vibration displacements is almost constant for any synchronous speed. As an instance, the vibration displacement  $X_G$  in the x direction is shown in Fig.6, in which the horizontal axis is the crank angle  $\theta$ . As shown in Table I, the mass  $m_0$  of the Oldham ring is larger than the mass  $m_V$  of the blade. Thus, the amplitude of the scroll compressor is larger than the rolling-piston rotary compressor, but the maximum amplitude is quite small 7.37  $\mu m$ .

The most significant component of the compressor vibrations is the rotation  $\Theta_z$ around the z axis. As given by the last expressions in (3) and (4) respectively, the unbalanced moment of inertia  $M_z$  which excites the vibration  $\Theta_z$  consists mainly of the rotatory acceleration  $\ddot{\theta}$ . As shown in *Fig.5*, the amplitude change of the rotatory acceleration  $\ddot{\theta}$  for the synchronous speed is very small compared with a change proportional to the quadratic rotating speed, that is, the exciting moment  $M_z$  is considered to be almost constant for any synchronous speed. Thus, the amplitude of the rotation  $\Theta_z$  increases inversely proportional to the decreasing quadratic synchronous speed, as shown

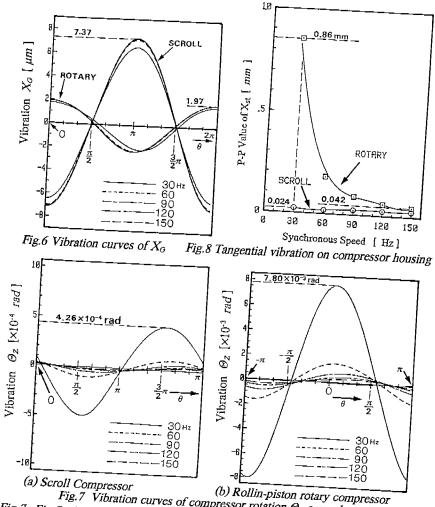


Fig.7 Vibration curves of compressor rotation  $\Theta_Z$  around z axis in Fig. 7. Fig. 7a shows the calculated results for the scroll compressor and Fig. 7b those for the rolling-piston rotary compressor. Both figures show the same tendency that the rotation amplitude increases inversely proportional to the quadratic synchronous speed. However, the rotation amplitude of the scroll compressor is far smaller than the rollingpiston rotary compressor. The rotation amplitude of the scroll compressor at the synchronous speed 30 Hz is quite small  $4.26 \times 10^{-4}$  rad which is about 25 dB smaller than 7.8×10<sup>-3</sup> rad of the rotary compressor.

The amplitude characteristics of the tangential vibration  $X_{st}$  on the compressor housing with the radius 55 mm is shown in Fig.8. It is a matter of course that the main factor inducing the tangential vibration  $X_{st}$  is the rotation  $\Theta_Z$  shown in Fig. 7. Thus, as the synchronous speed decreases, the peak to peak value of  $X_{st}$  for the rolling-piston rotary compressor increase rapidly. But  $X_{st}$  for the scroll compressor is almost independent of the synchronous speed and its peak to peak value is quite small 0.024 mm which is about 31 dB smaller than the maximum value 0.86 mm at the synchronous speed 30

Hz for the rolling-piston rotary compressor.

### CONCLUSIONS

The dynamic behavior of a small cooling capacity variable-speed scroll compressor with a suction volume of  $10.26 \text{ cm}^3$  which is operated at a suction pressure of 0.617MPa and a discharge pressure of 2.17 MPa was calculated for 5 different synchronous speed from 30 to 150 Hz, and the calculated results were compared with those of a variable-speed rolling-piston rotary compressor having the same cooling capacity. It is concluded from this study that the variable speed scroll compressors have the following dynamic behavior which is far superior than the variable-speed rolling-piston rotary compressors: The crankshaft speed fluctuation ratio of the scroll compressor is very small compared with the rolling-piston rotary compressor. For instance, it is 1.49 % at the synchronous speed 30 Hz, which is about one seventeenth of 25.8 % for the rollingpiston rotary compressor. The compressor vibration displacements except for the component around the rotating crankshaft have almost the same amplitude for any synchronous speed. However, the amplitude of the compressor rotation around the rotating crankshaft increases, as the synchronous speed decreases. Especially, the amplitude for the rolling-piston rotary compressor increases up to a large value which is 25 dB larger than the scroll compressor, when the synchronous speed is 30 Hz. As a result, the peak to peak value of the tangential vibration on the compressor housing of the scroll compressor is quite small 0.024 mm which is about 31 dB smaller than the rolling-piston rotary compressor, even when the compressor is driven by the motor with the synchronous speed 30 Hz.

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