

*Physics**Physics Research Publications*

*Purdue University*Year 2010

Studying the underlying event in
Drell-Yan and high transverse
momentum jet production at the
Tevatron

T. Aaltonen, J. Adelman, B. A. Gonzalez, S. Amerio, D. Amidei, A. Anastassov, A. Annovi, J. Antos, G. Apollinari, A. Apresyan, T. Arisawa, A. Artikov, J. Asaadi, W. Ashmanskas, A. Attal, A. Aurisano, F. Azfar, W. Badgett, A. Barbaro-Galtieri, V. E. Barnes, B. A. Barnett, P. Barria, P. Bartos, G. Bauer, P. H. Beauchemin, F. Bedeschi, D. Beecher, S. Behari, G. Bellettini, J. Bellinger, D. Benjamin, A. Beretvas, D. Berge, A. Bhatti, M. Binkley, D. Bisello, I. Bizjak, R. E. Blair, C. Blocker, B. Blumenfeld, A. Bocci, A. Bodek, V. Boisvert, D. Bortoletto, J. Boudreau, A. Boveia, B. Brau, A. Bridgeman, L. Brigliadori, C. Bromberg, E. Brubaker, J. Budagov, H. S. Budd, S. Budd, K. Burkett, G. Busetto, P. Bussey, A. Buzatu, K. L. Byrum, S. Cabrera, C. Calancha, S. Camarda, M. Campanelli, M. Campbell, F. Canelli, A. Canepa, B. Carls, D. Carlsmith, R. Carosi, S. Carrillo, S. Carron, B. Casal, M. Casarsa, A. Castro, P. Catastini, D. Cauz, V. Cavaliere, M. Cavalli-Sforza, A. Cerri, L. Cerrito, S. H. Chang, Y. C. Chen, M. Chertok, G. Chiarelli, G. Chlachidze, F. Chlebana, K. Cho, D. Chokheli, J. P. Chou, K. Chung, W. H. Chung, Y. S. Chung, T. Chwalek, C. I. Ciobanu, M. A. Ciocci, A. Clark, D. Clark, G. Compostella, M. E. Convery, J. Conway, M. Corbo, M. Cordelli, C. A. Cox, D. J. Cox, F. Crescioli, C. C. Almenar, J. Cuevas, R. Culbertson, J. C. Cully, D. Dagenhart, M. Datta, T. Davies, P. de Barbaro, S. De Cecco, A. Deisher, G. De Lorenzo, M. Dell'Orso, C. Deluca, L. Demortier, J. Deng, M. Deninno, M. d'Errico, A. Di Canto, G. P. di Giovanni, B. Di Ruzza, J. R. Dittmann, M. D'Onofrio, S. Donati, P. Dong, T. Dorigo, S. Dube, K. Ebina, A. Elagin, R. Erbacher, D. Errede, S. Errede, N. Ershaidat, R. Eusebi, H. C. Fang, S. Farrington, W. T. Fedorko, R. G. Feild, M. Feindt, J. P. Fernandez, C. Ferrazza, R. Field, G. Flanagan, R. Forrest, M. J. Frank, M. Franklin, J. C. Freeman, I. Furic, M. Gallinaro, J. Galyardt, F. Garbersson, J. E. Garcia, A. F. Garfinkel, P. Garosi, H. Gerberich, D. Gerdes, A. Gessler, S. Giagu, V. Giakoumopoulou, P.

Giannetti, K. Gibson, J. L. Gimmell, C. M. Ginsburg, N. Giokaris, M. Giordani, P. Giromini, M. Giunta, G. Giurgiu, V. Glagolev, D. Glenzinski, M. Gold, N. Goldschmidt, A. Golossanov, G. Gomez, G. Gomez-Ceballos, M. Goncharov, O. Gonzalez, I. Gorelov, A. T. Goshaw, K. Goulianos, A. Gresele, S. Grinstein, C. Grosso-Pilcher, R. C. Group, U. Grundler, J. G. da Costa, Z. Gunay-Unalan, C. Haber, S. R. Hahn, E. Halkiadakis, B. Y. Han, J. Y. Han, F. Happacher, K. Hara, D. Hare, M. Hare, R. F. Harr, M. Hartz, K. Hatakeyama, C. Hays, M. Heck, J. Heinrich, M. Herndon, J. Heuser, S. Hewamanage, M. Hickman, D. Hidas, C. S. Hill, D. Hirschbuehl, A. Hocker, S. Hou, M. Houlden, S. C. Hsu, R. E. Hughes, M. Hurwitz, U. Husemann, M. Hussein, J. Huston, J. Incandela, G. Introzzi, M. Iori, A. Ivanov, E. James, D. Jang, B. Jayatilaka, E. J. Jeon, M. K. Jha, S. Jindariani, W. Johnson, M. Jones, K. K. Joo, S. Y. Jun, J. E. Jung, T. R. Junk, T. Kamon, D. Kar, P. E. Karchin, Y. Kato, R. Kephart, W. Ketchum, J. Keung, V. Khotilovich, B. Kilminster, D. H. Kim, H. S. Kim, H. W. Kim, J. E. Kim, M. J. Kim, S. B. Kim, S. H. Kim, Y. K. Kim, N. Kimura, L. Kirsch, S. Klimentenko, K. Kondo, D. J. Kong, J. Konigsberg, A. Korytov, A. V. Kotwal, M. Krepis, J. Kroll, D. Krop, N. Krumnack, M. Kruse, V. Krutelyov, T. Kuhr, N. P. Kulkarni, M. Kurata, S. Kwang, A. T. Laasanen, S. Lami, S. Lammel, M. Lancaster, R. L. Lander, K. Lannon, A. Lath, G. Latino, I. Lazzizzera, T. LeCompte, E. Lee, H. S. Lee, J. S. Lee, S. W. Lee, S. Leone, J. D. Lewis, C. J. Lin, J. Linacre, M. Lindgren, E. Lipeles, A. Lister, D. O. Litvintsev, C. Liu, T. Liu, N. S. Lockyer, A. Loginov, L. Lovas, D. Lucchesi, J. Lueck, P. Lujan, P. Lukens, G. Lungu, J. Lys, R. Lysak, D. MacQueen, R. Madrak, K. Maeshima, K. Makhoul, P. Maksimovic, S. Malde, S. Malik, G. Manca, A. Manousakakis-Katsikakis, F. Margaroli, C. Marino, C. P. Marino, A. Martin, V. Martin, M. Martinez, R. Martinez-Ballarín, P. Mastrandrea, M. Mathis, M. E. Mattson, P. Mazzanti, K. S. McFarland, P. McIntyre, R. McNulty, A. Mehta, P. Mehtala, A. Menzione, C. Mesropian, T. Miao, D. Miettlicki, N. Miladinovic, R. Miller, C. Mills, M. Milnik, A. Mitra, G. Mitselmakher, H. Miyake, S. Moed, N. Moggi, M. N. Mondragon, C. S. Moon, R. Moore, M. J. Morello, J. Morlock, P. M. Fernandez, J. Mulmenstadt, A. Mukherjee, T. Muller, P. Murat, M. Mussini, J. Nachtman, Y. Nagai, J. Naganoma, K. Nakamura, I. Nakano, A. Napier, J. Nett, C. Neu, M. S. Neubauer, S. Neubauer, J. Nielsen, L. Nodulman, M. Norman, O. Norriella, E. Nurse, L. Oakes, S. H. Oh, Y. D. Oh, I. Oksuzian, T. Okusawa, R. Orava, K. Osterberg, S. P. Griso, C. Pagliarone, E. Palencia, V. Papadimitriou, A. Papaikonomou, A. A. Paramanov, B. Parks, S. Pashapour, J. Patrick, G. Pauletta, M. Paulini, C. Paus, T. Peiffer, D. E. Pellett, A. Penzo, T. J. Phillips, G. Piacentino, E. Pianori, L. Pinera, K. Pitts, C. Plager, L. Pondrom, K. Potamianos, O. Poukhov, F. Prokoshin, A. Pronko, F. Ptohos, E. Pueschel, G. Punzi, J. Pursley, J. Rademacker, A. Rahaman, V. Ramakrishnan, N. Ranjan, I. Redondo, P. Renton, M. Renz, M. Rescigno, S. Richter, F. Rimondi, L. Ristori, A. Robson, T. Rodrigo, T. Rodriguez, E. Rogers, S. Rolli, R. Roser, M. Rossi, R. Rossin, P. Roy, A. Ruiz, J. Russ, V. Rusu, B. Rutherford, H. Saarikko, A. Safonov, W. K. Sakumoto, L. Santi, L. Sartori, K. Sato, A. Savoy-Navarro, P. Schlabach, A. Schmidt, E. E. Schmidt, M. A. Schmidt, M. P. Schmidt, M. Schmitt, T. Schwarz, L. Scodellaro, A. Scribano, F. Scuri,

A. Sedov, S. Seidel, Y. Seiya, A. Semenov, L. Sexton-Kennedy, F. Sforza, A. Sfyrla, S. Z. Shalhout, T. Shears, P. F. Shepard, M. Shimojima, S. Shiraishi, M. Shochet, Y. Shon, I. Shreyber, A. Simonenko, P. Sinervo, A. Sisakyan, A. J. Slaughter, J. Slaunwhite, K. Sliwa, J. R. Smith, F. D. Snider, R. Snihur, A. Soha, S. Somalwar, V. Sorin, P. Squillacioti, M. Stanitzki, R. S. Denis, B. Stelzer, O. Stelzer-Chilton, D. Stentz, J. Strologas, G. L. Strycker, J. S. Suh, A. Sukhanov, I. Suslov, A. Taffard, R. Takashima, Y. Takeuchi, R. Tanaka, J. Tang, M. Tecchio, P. K. Teng, J. Thom, J. Thome, G. A. Thompson, E. Thomson, P. Tipton, P. Ttito-Guzman, S. Tkaczyk, D. Toback, S. Tokar, K. Tollefson, T. Tomura, D. Tonelli, S. Torre, D. Torretta, P. Totaro, S. Tourneur, M. Trovato, S. Y. Tsai, Y. Tu, N. Turini, F. Ukegawa, S. Uozumi, N. van Remortel, A. Varganov, E. Vataga, F. Vazquez, G. Velez, C. Vellidis, M. Vidal, I. Vila, R. Vilar, M. Vogel, I. Volobouev, G. Volpi, P. Wagner, R. G. Wagner, R. L. Wagner, W. Wagner, J. Wagner-Kuhr, T. Wakisaka, R. Wallny, S. M. Wang, A. Warburton, D. Waters, M. Weinberger, J. Weinelt, W. C. Wester, B. Whitehouse, D. Whiteson, A. B. Wicklund, E. Wicklund, S. Wilbur, G. Williams, H. H. Williams, M. G. Wilson, P. Wilson, B. L. Winer, P. Wittich, S. Wolbers, C. Wolfe, H. Wolfe, T. Wright, X. Wu, F. Wurthwein, A. Yagil, K. Yamamoto, J. Yamaoka, U. K. Yang, Y. C. Yang, W. M. Yao, G. P. Yeh, K. Yi, J. Yoh, K. Yorita, T. Yoshida, G. B. Yu, I. Yu, S. S. Yu, J. C. Yun, A. Zanetti, Y. Zeng, X. Zhang, Y. Zheng, and S. Zucchelli

Studying the underlying event in Drell-Yan and high transverse momentum jet production at the Tevatron

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We study the underlying event in proton-antiproton collisions by examining the behavior of charged particles produced in association with a large transverse momentum jet ($\sim 2.2 \text{ fb}^{-1}$) or with a Drell-Yan lepton pair ($\sim 2.7 \text{ fb}^{-1}$) in the Z-boson mass region [$70 < M(\text{pair}) < 110 \text{ GeV}/c^2$] as measured by CDF at 1.96 TeV center-of-mass energy. We use the direction of the lepton pair or the leading jet in each event to define regions of η - ϕ space that are sensitive to the modeling of the underlying event. The data are corrected to the particle level to remove detector effects and are then compared with several QCD Monte Carlo models.

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I. INTRODUCTION

In order to find physics beyond the standard model at a hadron-hadron collider, it is essential to have Monte Carlo models that accurately simulate QCD hard-scattering events. To do this one must not only have a good model of the hard-scattering part of the process, but also of the beam-beam remnants (BBR) and the multiple-parton interactions (MPI). The underlying event consists of the BBR plus MPI and is an unavoidable background to most collider observables. A good understanding of the underlying event will lead to more precise measurements at the Tevatron and the Large Hadron Collider (LHC). The goal of this analysis is to provide data that can be used to test and improve the QCD Monte Carlo models of the underlying event.

Figure 1 illustrates the way the QCD Monte Carlo models simulate a proton-antiproton collision in which a hard two-to-two parton scattering with transverse momentum, $p_T(\text{hard})$, has occurred. The resulting event contains particles that originate from the two outgoing partons (plus initial and final-state radiation) and particles that come from the breakup of the proton and antiproton. The beam-beam remnants are what is left over after a parton is knocked out of each of the initial two beam hadrons.

They are one of the reasons why hadron-hadron collisions are more complicated than electron-positron annihilations. For the QCD Monte Carlo models the beam-beam remnants are an important component of the underlying event. In addition to the hard two-to-two parton-parton scattering and the beam-beam remnants, sometimes there are additional semihard two-to-two parton-parton scatterings (MPI) that contribute particles to the underlying event. However, on an event-by-event basis these two components cannot be uniquely separated from particles that come from the initial and final-state radiation. Hence, a study of the underlying event inevitably involves a study of the BBR plus MPI plus initial and final-state radiation.

As shown in Fig. 2, Drell-Yan lepton-pair production provides an excellent place to study the underlying event. Here one studies the outgoing charged particles (excluding the lepton pair) as a function of the lepton-pair invariant mass and as a function of the lepton-pair transverse momentum. Unlike high- p_T jet production, for lepton-pair production there is no final-state gluon radiation.

Hard-scattering collider jet events have a distinct topology. A typical hard-scattering event consists of a collection (or burst) of hadrons traveling roughly in the direction of the initial two beam particles and two collections of hadrons (jets) with large transverse momentum. The two large

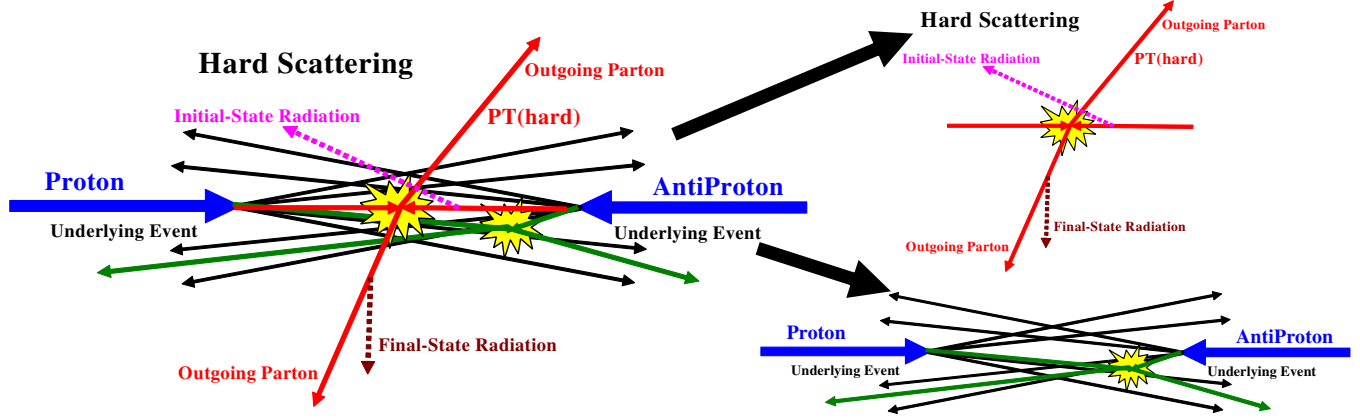


FIG. 1 (color online). Illustration of the way QCD Monte Carlo models simulate a proton-antiproton collision in which a hard two-to-two parton scattering with transverse momentum, $p_T(\text{hard})$, has occurred. The hard-scattering component of the event consists of particles that result from the hadronization of the two outgoing partons (i.e. the initial two jets) plus the particles that arise from initial and final-state radiation (i.e. multijets). The underlying event consists of particles that arise from the beam-beam remnants and from multiple-parton interactions.

transverse momentum jets are roughly back-to-back in azimuthal angle ϕ . One can use the topological structure of hadron-hadron collisions to study the underlying event. We use the direction of the leading (highest p_T) jet in each event to define four regions of η - $\Delta\phi$ space (referred to as leading-jet events). As illustrated in Fig. 3, for leading-jet events $\Delta\phi = \phi - \phi_{\text{jet}\#1}$, where $\phi_{\text{jet}\#1}$ and ϕ are the azimuthal angles of the leading jet and a charged particle, respectively, and $\eta = -\log(\tan(\theta_{\text{cm}}/2))$ is the pseudorapidity, where θ_{cm} is the center-of-mass polar scattering angle of the outgoing charged particles.

As is also shown in Fig. 3 in Drell-Yan lepton-pair production (referred to as Drell-Yan events) $\Delta\phi = \phi - \phi_{\text{pair}}$, where ϕ_{pair} and ϕ are the azimuthal angles of the lepton pair and a charged particle, respectively. On an event-by-event basis, the toward region contains all

charged particles with $|\Delta\phi| < 60^\circ$ and $|\eta| < 1$, while the away region contains all charged particles with $|\Delta\phi| > 120^\circ$ and $|\eta| < 1$. The two transverse regions $60^\circ < -\Delta\phi < 120^\circ$, $|\eta| < 1$ and $60^\circ < \Delta\phi < 120^\circ$, $|\eta| < 1$ are referred to as transverse 1 and transverse 2, respectively. The overall transverse region corresponds to combining the charged particles in the transverse-1 and transverse-2 regions. For Drell-Yan events the two leptons are not included. For leading-jet events, the toward and away regions receive large contributions from the outgoing high- p_T jets, while the transverse region is perpendicular to the plane of the hard two-to-two scattering and is therefore very sensitive to the underlying event. For Drell-Yan events both the toward and the transverse regions are very sensitive to the underlying event, while the away region receives large contributions from the awayside jet from the

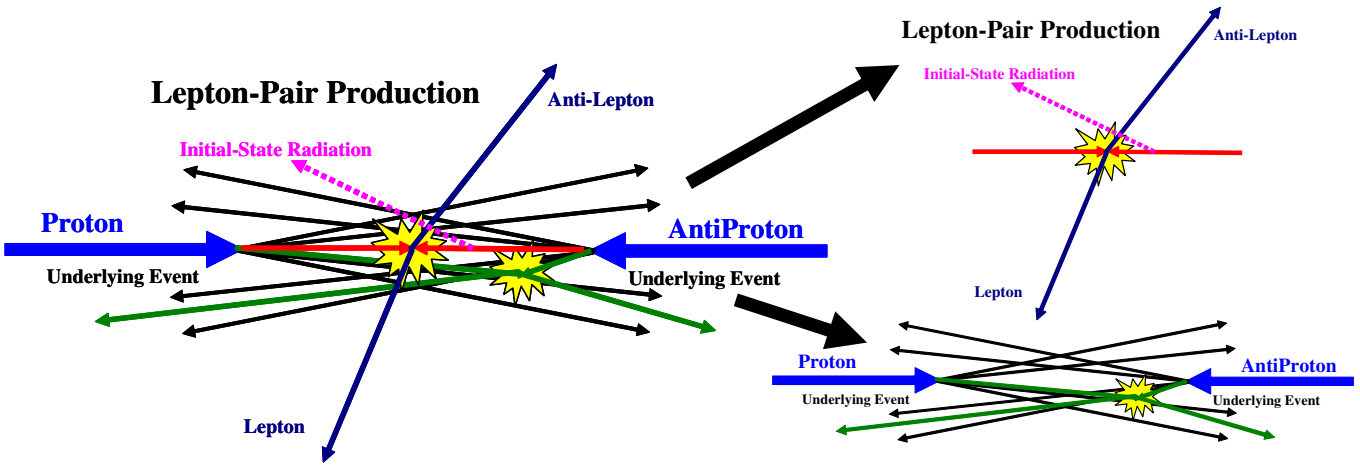


FIG. 2 (color online). Illustration of the way QCD Monte Carlo models simulate Drell-Yan lepton-pair production. The hard-scattering component of the event consists of the two outgoing leptons plus particles that result from initial-state radiation. As in the hard two-to-two parton scattering of Fig. 1, the underlying event consists of particles that arise from the beam-beam remnants and from multiple-parton interactions.

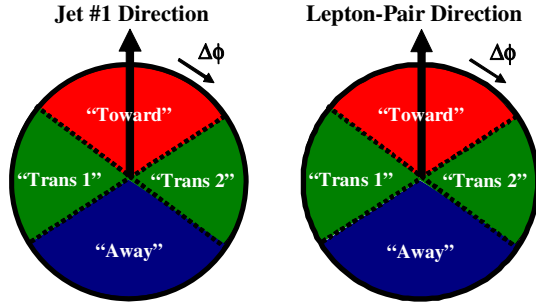


FIG. 3 (color online). Illustration of correlations in azimuthal angle $\Delta\phi$ relative to (left) the direction of the leading jet (highest p_T jet) in the event, jet#1, in high- p_T jet production or (right) the direction of the lepton pair in Drell-Yan production. The angle $\Delta\phi = \phi - \phi_{\text{jet}\#1}$ ($\Delta\phi = \phi - \phi_{\text{pair}}$) is the relative azimuthal angle between charged particles and the direction of jet#1 (lepton pair). The toward region is defined by $|\Delta\phi| < 60^\circ$ and $|\eta| < 1$, while the away region is $|\Delta\phi| > 120^\circ$ and $|\eta| < 1$. The two transverse regions $60^\circ < -\Delta\phi < 120^\circ$, $|\eta| < 1$ and $60^\circ < \Delta\phi < 120^\circ$, $|\eta| < 1$ are referred to as transverse 1 and transverse 2, respectively. Each of the two transverse regions has an area in η - $\Delta\phi$ space of $\Delta\eta\Delta\phi = 4\pi/6$. The overall transverse region corresponds to combining the transverse-1 and transverse-2 regions. The transMAX (transMIN) refer to the transverse region (transverse-1 or transverse-2) containing the largest (smallest) number of charged particles or to the region containing the largest (smallest) scalar p_T sum of charged particles.

subprocesses: $q + \bar{q} \rightarrow l^+ l^- + g$, $q + g \rightarrow l^+ l^- + q$, and $\bar{q} + g \rightarrow l^+ l^- + \bar{q}$.

We study charged particles in the range in the toward, away, and transverse regions. For leading-jet events, we require that the leading jet in the event be in the region $|\eta(\text{jet}\#1)| < 2$; however, charged particles are restricted to the region $|\eta| < 1$. The jets are constructed using the MidPoint algorithm ($R = 0.7$, $f_{\text{merge}} = 0.75$), where R is the jet radius and f_{merge} is the jet splitting and merging

fraction [1]. For Drell-Yan production we require that the invariant mass of the lepton pair be in the mass region of the Z boson, $70 < M(\text{pair}) < 110 \text{ GeV}/c^2$, with $|\eta(\text{pair})| < 6$.

Table I shows the observables that are considered in this analysis as they are defined at the particle level and detector level. The detector level corresponds to the tracks passing good-track criteria, and the particle level corresponds to true charged particles in the event. The particle level can be compared directly with the QCD Monte Carlo models at the generator level. Since we will be studying regions in η - $\Delta\phi$ space with different areas, we construct densities by dividing by the area. For example, the number density corresponds the number of charged particles per unit η - $\Delta\phi$, and the PT_{sum} density corresponds the charged-particle scalar- p_T sum per unit η - $\Delta\phi$.

For both leading-jet and Drell-Yan events we define MAX and MIN transverse regions (transMAX and transMIN) [2]. For the charged-particle density MAX (MIN) refers to the transverse region (transverse-1 or transverse-2) containing the largest (smallest) number of charged particles. For the charged scalar PT_{sum} density MAX (MIN) refers to the transverse region (transverse-1 or transverse-2) containing the largest (smallest) scalar p_T sum of charged particles. For events with large initial or final-state radiation the transMAX region will usually contain the third jet in high- p_T jet production or the second jet in Drell-Yan production while both the transMAX and transMIN regions receive contributions from the beam-beam remnants. Thus, the transMIN region is very sensitive to the beam-beam remnants, while the event-by-event difference between transMAX and transMIN is very sensitive to initial and final-state radiation (transDIF = transMAX - transMIN).

A discussion of the QCD Monte Carlo model is presented in Sec. II. In Sec. III we discuss the data selection, track cuts, and the method we use to correct the data to the

TABLE I. Observables examined in this analysis as they are defined at the particle level and the detector level. Charged tracks are considered good if they pass the track selection criterion given in Sec. III E. The mean charged-particle $\langle p_T \rangle$ is constructed on an event-by-event basis and then averaged over the events. For the average p_T and the PT_{max} , we require that there is at least one charged particle present. Particles are considered stable if $c\tau > 10 \text{ mm}$ (K_s , Λ , Σ , Ξ , and Ω are kept stable).

Observable	Particle level	Detector level
$dN/d\eta d\phi$	Number of stable charged particles per unit η - ϕ ($p_T > 0.5 \text{ GeV}/c$, $ \eta < 1$)	Number of good tracks per unit η - ϕ ($p_T > 0.5 \text{ GeV}/c$, $ \eta < 1$)
$d\text{PT}/d\eta d\phi$	Scalar p_T sum of stable charged particles per unit η - ϕ ($p_T > 0.5 \text{ GeV}/c$, $ \eta < 1$)	Scalar p_T sum of good tracks per unit η - ϕ ($p_T > 0.5 \text{ GeV}/c$, $ \eta < 1$)
$\langle p_T \rangle$	Average p_T of stable charged particles ($p_T > 0.5 \text{ GeV}/c$, $ \eta < 1$) require at least 1 charged particle	Average p_T of good tracks ($p_T > 0.5 \text{ GeV}/c$, $ \eta < 1$) require at least 1 good track
PT_{max}	Maximum p_T stable charged particle ($p_T > 0.5 \text{ GeV}/c$, $ \eta < 1$) require at least 1 charged particle	Maximum p_T good charged tracks ($p_T > 0.5 \text{ GeV}/c$, $ \eta < 1$) require at least 1 good track
Jet	MidPoint algorithm $R = 0.7$ $f_{\text{merge}} = 0.75$ applied to stable particles	MidPoint algorithm $R = 0.7$ $f_{\text{merge}} = 0.75$ applied to calorimeter cells

particle level. Section IV contains the results for leading-jet and Drell-Yan events and comparisons with the QCD Monte Carlo models. Section V is reserved for the summary and conclusions.

II. QCD MONTE CARLO MODELS

QCD Monte Carlo generators such as PYTHIA [3] have parameters which may be adjusted to control the behavior of their event modeling. A specified set of these parameters that has been adjusted to better fit some aspects of the data is referred to as a tune. PYTHIA Tune A was determined by fitting the CDF run 1 underlying-event data [4]. Later it was noticed that Tune A does not fit the CDF run 1 Z-boson p_T distribution very well [5]. PYTHIA Tune AW was tuned to fit the Z-boson p_T distribution as well as the underlying event at the Tevatron [6]. For leading-jet production Tune A and Tune AW are nearly identical. Table II shows the parameters for several tunes for PYTHIA version 6.2. PYTHIA Tune DW is very similar to Tune AW except the setting of one PYTHIA parameter $\text{PARP}(67) = 2.5$, which is the preferred value determined by the D0 Collaboration in fitting their dijet $\Delta\phi$ distribution [8]. $\text{PARP}(67)$ sets the high- p_T scale for initial-state radiation in PYTHIA. It determines the maximal parton virtuality allowed in timelike showers. Tune DW and Tune DWT are identical at 1.96 TeV (the reference point), but Tune DW and DWT extrapolate differently to the LHC. Tune DWT uses the PYTHIA default value for energy dependence of the MPI cutoff [$\text{PARP}(90) = 0.16$], which is the value used in the ATLAS PYTHIA tune [7]. Tune DWT produces more activity in the underlying

event at the LHC than does Tune DW but predicts less activity than Tune DW in the underlying event at energies below 1.96 TeV. Tune DW uses the Tune A value of $\text{PARP}(90) = 0.25$, which was determined by comparing the run 1 data at 1.8 TeV with the CDF underlying-event measurements at 630 GeV [9]. The amount of MPI and hence the tunings depend on the choice of the parton distribution functions. All these tunes use the CTEQ5L [10] parton distribution functions.

The first 9 parameters in Table II tune the MPI. $\text{PARP}(62)$, $\text{PARP}(64)$, and $\text{PARP}(67)$ tune the initial-state radiation, and the last three parameters set the intrinsic transverse momentum of the partons within the incoming proton and antiproton.

Table III shows the computed value of the multiple-parton scattering cross section for the various tunes. The multiple-parton scattering cross section (divided by the total inelastic cross section at the center-of-mass energies of 1.96 and 14 TeV, respectively) determines the average number of multiple-parton collisions per event. The MPI cross section is the same for proton-proton and proton-antiproton collisions.

HERWIG [11] is a QCD Monte Carlo generator similar to PYTHIA except HERWIG employs a cluster fragmentation model while PYTHIA uses a string fragmentation approach. In addition, gluon radiation is modeled differently by the two generators. Also, HERWIG does not include MPI in the underlying event. In HERWIG the underlying event arises solely from the BBR. JIMMY [12] is a multiple-parton interaction model which can be added to HERWIG to improve agreement with the underlying-event observables.

TABLE II. Parameters for several PYTHIA 6.2 tunes. Tune A is the CDF run 1 underlying-event tune. Tune AW and DW are CDF run 2 tunes which fit the existing run 2 underlying-event data and fit the run 1 Z-boson p_T distribution. The ATLAS Tune is the tune used in the ATLAS technical design report [7]. Tune DWT uses the ATLAS energy dependence for the MPI, $\text{PARP}(90)$. The first 9 parameters tune the multiple-parton interactions. $\text{PARP}(62)$, $\text{PARP}(64)$, and $\text{PARP}(67)$ tune the initial-state radiation, and the last three parameters set the intrinsic k_T of the partons within the incoming proton and antiproton.

Parameter	Description	Tune A	Tune AW	Tune DW	Tune DWT	ATLAS
PDF	Parton distribution functions	CTEQ5L	CTEQ5L	CTEQ5L	CTEQ5L	CTEQ5L
MSTP(81)	MPI on	1	1	1	1	1
MSTP(82)	Double Gaussian matter distribution	4	4	4	4	4
PARP(82)	MPI cutoff	2.0	2.0	1.9	1.9409	1.8
PARP(83)	Fraction of matter within core	0.5	0.5	0.5	0.5	0.5
PARP(84)	Core radius	0.4	0.4	0.4	0.4	0.5
PARP(85)	Color connections	0.9	0.9	1.0	1.0	0.33
PARP(86)	Color connections	0.95	0.95	1.0	1.0	0.66
PARP(89)	Reference energy	1800	1800	1800	1960	1000
PARP(90)	MPI energy dependence	0.25	0.25	0.25	0.16	0.16
PARP(62)	Initial-state radiation cutoff	1.0	1.25	1.25	1.25	1.0
PARP(64)	Soft initial-state radiation scale	1.0	0.2	0.2	0.2	1.0
PARP(67)	Hard initial-state radiation scale	4.0	4.0	2.5	2.5	1.0
MSTP(91)	Gaussian intrinsic k_T	1	1	1	1	1
PARP(91)	Intrinsic Gaussian width σ	1.0	2.1	2.1	2.1	1.0
PARP(93)	Intrinsic k_T upper cutoff	5.0	15.0	15.0	15.0	5.0

TABLE III. The computed value of the multiple-parton scattering cross section for the various PYTHIA 6.2 tunes.

Tune	$\sigma(\text{MPI})$ at 1.96 TeV	$\sigma(\text{MPI})$ at 14 TeV
A, AW	309.7 mb	484.0 mb
DW	351.7 mb	549.2 mb
DWT	351.7 mb	829.1 mb
ATLAS	324.5 mb	768.0 mb

To compare with the Drell-Yan data we have constructed a HERWIG tune (with JIMMY MPI) with $\text{JMUEO} = 1$, $\text{PTJIM} = 3.6 \text{ GeV}/c$, $\text{JMRAD}(73) = 1.8$, and $\text{JMRAD}(91) = 1.8$. These parameters govern the MPI activity produced by JIMMY. This tune of JIMMY was arrived at by fitting the data from this analysis on the charged scalar particle PT_{sum} density in the toward region for Drell-Yan production.

In this paper the Monte Carlo model predictions are presented as smooth curves. These curves come from fits to QCD Monte Carlo output with limited statistical accuracy. The curves presented here reproduce the QCD Monte Carlo results (with infinite statistical accuracy) within about 2%.

III. ANALYSIS STRATEGY

A. Data sample and event selection

The CDF run II detector, in operation since 2001, is an azimuthally and forward-backward symmetric solenoidal particle detector [13]. It combines precision charged-particle tracking with fast projective calorimetry and fine-grained muon detection. Tracking systems are designed to detect charged particles and measure their momenta and displacements from the point of collision, termed the primary interaction vertex. The tracking system consists of a silicon microstrip system and an open-cell wire drift chamber, termed the central outer tracker (COT) that surrounds the silicon. Segmented electromagnetic and hadronic sampling calorimeters surround the tracking system and measure the energy of interacting particles. Particles make showers which deposit energy and are sampled via their ionization. The muon system resides beyond the calorimeters. Muons are minimally ionizing particles and, hence, only deposit small amounts of ionization energy in the material. They are the only particles likely to penetrate both the tracking and five pion absorption lengths of calorimeter steel and leave tracks in the muon detection system. At CDF the positive z axis is defined to lie along the incident proton beam direction. The leading-jet data and lepton-pair data correspond to an integrated luminosity of about 2.2 and 2.7 fb^{-1} , respectively. For both data sets we require one and only one primary vertex within the fiducial region $|Z_{\text{vertex}}| \leq 60 \text{ cm}$ centered around the nominal CDF $z = 0$.

B. Jet selection

Jets are selected using the MidPoint cone-based algorithm with a cone size of 0.7 and $f_{\text{merge}} = 0.75$ [1]. For the leading-jet events we require that the highest p_T jet in the calorimeter lie in the range $|\eta| < 2$ or the event is rejected.

C. Lepton selection

Dielectron events are triggered online by either one central ($|\eta| < 1.1$) electron candidate with $E_T > 18 \text{ GeV}$ and a track with $p_T > 18 \text{ GeV}/c$ associated to it or by two electromagnetic clusters with $E_T > 18 \text{ GeV}$ and $|\eta| < 3.2$ where no track association is required. At offline level we consider only electrons with $E_T > 20 \text{ GeV}$ and $|\eta| < 1$ that also have a track matched to the calorimeter cluster. The electrons also have to pass certain quality criteria to verify that they are consistent with the electromagnetic shower characteristics as expected for electrons [14].

Dimuon events are triggered on at least one muon candidate that has a signal in one of the muon chambers with $|\eta| < 1$ and $p_T > 18 \text{ GeV}/c$. The second muon candidate is not required to have a signal in the muon chambers but it must have hits in the COT. At offline level we consider only muon candidates with $p_T > 20 \text{ GeV}$ and $|\eta| < 1$. All muon candidates are required to have calorimeter energy deposits consistent with those expected from a minimum ionizing particle. In addition, we employ a time-of-flight filter to remove cosmic ray muons.

All leptons are required to be isolated from other charged tracks in the event. The lepton is rejected if there is a charged track within the distance of $R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} = 0.4$.

D. Lepton-pair selection

The lepton pairs are formed by oppositely charged leptons, with the requirement that the z positions of the two leptons satisfy $|\Delta z| < 4 \text{ cm}$, to ensure that both leptons came from the same primary collision. For the Drell-Yan data we require that both leptons have $p_T > 20 \text{ GeV}/c$ and $|\eta| < 1$ and that the invariant mass of the lepton pair be in the range $70 < M(\text{pair}) < 110 \text{ GeV}/c^2$, with $|\eta(\text{pair})| < 6$. We chose this lepton-pair mass region because studies have shown that the lepton-pair backgrounds (mostly from events with QCD jets or events with a W boson and jets) are negligible in the region of the Z boson [15].

E. Track selection

We consider charged tracks that have been measured by the COT. The COT [16] is a cylindrical open-cell drift chamber with 96 sense wire layers grouped into eight alternating superlayers of stereo and axial wires. Its active volume covers $40 < r < 137 \text{ cm}$, where r is the radial coordinate in the plane transverse to the z axis, and $|z| < 155 \text{ cm}$, thus providing fiducial coverage in $|\eta| \leq 1.1$ to tracks originating within $|z| \leq 60 \text{ cm}$. We include tracks in

the region $0.5 < p_T < 150$ GeV/ c and $|\eta| < 1$ where COT efficiency is high. At very high p_T the track resolution deteriorates. The upper limit of 150 GeV/ c is chosen to prevent mismeasured tracks with very high p_T from distorting the average charged-particle density and the average charged-particle PT_{sum} density. Tracks are required to hit at least two axial segments with more than 10 total hits and at least two stereo segments with more than 10 total hits in the COT. In addition, the tracks are required to point back to the primary vertex. We consider two track selections: loose and tight. The loose track selection requires $|d_0| < 1.0$ cm and $|z - Z_{\text{vtx}}| < 3$ cm, where d_0 is the beam corrected transverse impact parameter and $z - Z_{\text{vtx}}$ is the distance on the z axis (beam axis) between the track and the primary vertex. The tight track selection requires that $|d_0| < 0.5$ cm and $|z - Z_{\text{vtx}}| < 2$ cm. The loose criterion is similar to the run 1 underlying-event analysis [4].

F. Correcting to the particle level and systematic uncertainties

The raw data at the detector level are corrected to the level of final-state stable particles and are then compared with the QCD Monte Carlo models at the generator level. The particle level corresponds to the event without detector effects. The detector level corresponds to the tracks passing the good-track criterion. We rely on the QCD Monte Carlo models and the CDF detector simulation CDFSIM (parametrized response of the CDF II detector [17,18]) to correct the measured tracks back to the stable charged-particle level. The generator level charged particles have $p_T > 0.5$ GeV/ c , $|\eta| < 1$, and are considered stable if $c\tau > 10$ mm. Hence, to compare the corrected data with QCD Monte Carlo model predictions one must keep the K_{short} meson stable as well as the following baryons: Λ , Σ , Ξ , and Ω .

The QCD Monte Carlo models are used to calculate the observables in Table I at the particle level in bins of particle jet#1 p_T (GEN) and at the detector level in bins of calorimeter jet#1 p_T (CDFSIM). GEN refers to the Monte Carlo model at the generator level and CDFSIM are the GEN particles after detector simulation. The detector-level data in bins of calorimeter jet#1 p_T are corrected by multiplying by the QCD Monte Carlo correction factor, $F_{\text{cor}} = \text{GEN}/\text{CDFSIM}$. This is done bin by bin for every observable. We refer to the ratio $F_{\text{res}} = \text{CDFSIM}/\text{GEN}$ as the response factor for that observable with the correction factor being the reciprocal. Smooth curves are drawn through the QCD Monte Carlo predictions at both the generator level (GEN) and the detector level (CDFSIM) to aid in comparing with the data and also to construct the correction factors. This one-step correction method simultaneously corrects for mismeasurement of the leading-jet transverse momentum (jet energy scale) and for missed and/or fake tracks.

The correction factors are different for every observable and they are different for the tight and loose track selection criteria. The tight track criterion results in fewer tracks than the loose criterion and hence the Monte Carlo corrections factors are different. If the Monte Carlo model described the data perfectly and if CDFSIM were exact, then the corrected observable would be identical regardless of the track selection criterion. Using PYTHIA Tune A for the leading-jet events and PYTHIA Tune AW for the Drell-Yan events, we find that the loose and tight track selections do result in nearly the same particle level result for all the observables presented in this analysis. The differences are used as a source of systematic error and are added in quadrature to the statistical errors.

Figure 4 shows the response factors F_{res} for the charged-particle density in the toward and transverse regions for leading-jet events. The correction factors ($1/F_{\text{res}}$) are typically small (they differ from one by less than 10%) except in regions where the charged-particle density becomes large, which occurs in the toward and away regions for leading-jet production. The efficiency of detecting charged tracks decreases when the density of tracks becomes large. For this reason we restrict ourselves to the range $p_T(\text{jet}\#1) < 200$ GeV/ c for the toward and away regions, but allow the leading-jet transverse momentum to extend to 400 GeV/ c in the transverse region. For the leading-jet events we have also used HERWIG (without MPI) as well as PYTHIA Tune A to correct the data to the particle level. We use the differences in the corrected data as an additional source of systematic error (added in quadrature). For low $p_T(\text{jet}\#1)$ the correction factors become large due to the uncertainty in the jet energy scale at low energy. Also, the corrections from HERWIG and PYTHIA Tune A differ significantly in this region. This results in very large systematic errors on the data at low leading-jet transverse momentum.

Another important effect and resultant systematic error arises from the uncertainty in the jet energy scale for p_T of the leading jet. The CDF detector simulation does not reproduce perfectly the response of the calorimeters. The overall systematic uncertainty in the CDF jet energy scale is a function of the jet p_T [19]. The uncertainty is about 3% at high p_T and increases to around 8% at low p_T . After correcting the data to the particle level we shift $p_T(\text{jet}\#1)$ up and down by this additional uncertainty with the bin-by-bin differences in the observables in Table I used as another systematic error. The jet energy scale systematic errors are large in the toward and away regions where the observables are varying rapidly with $p_T(\text{jet}\#1)$.

We investigated the dependence of the corrected data to our upper limit of $PT_{\text{max}}(\text{cut}) = 150$ GeV/ c which was applied to all tracks. The sensitivity of the results to this choice of upper limit was checked by changing the upper limit to $PT_{\text{max}}(\text{cut}) = 1.5 \times ET_{\text{max}}(\text{tower})$. Here one looks, on an event-by-event basis, at all the towers in the region

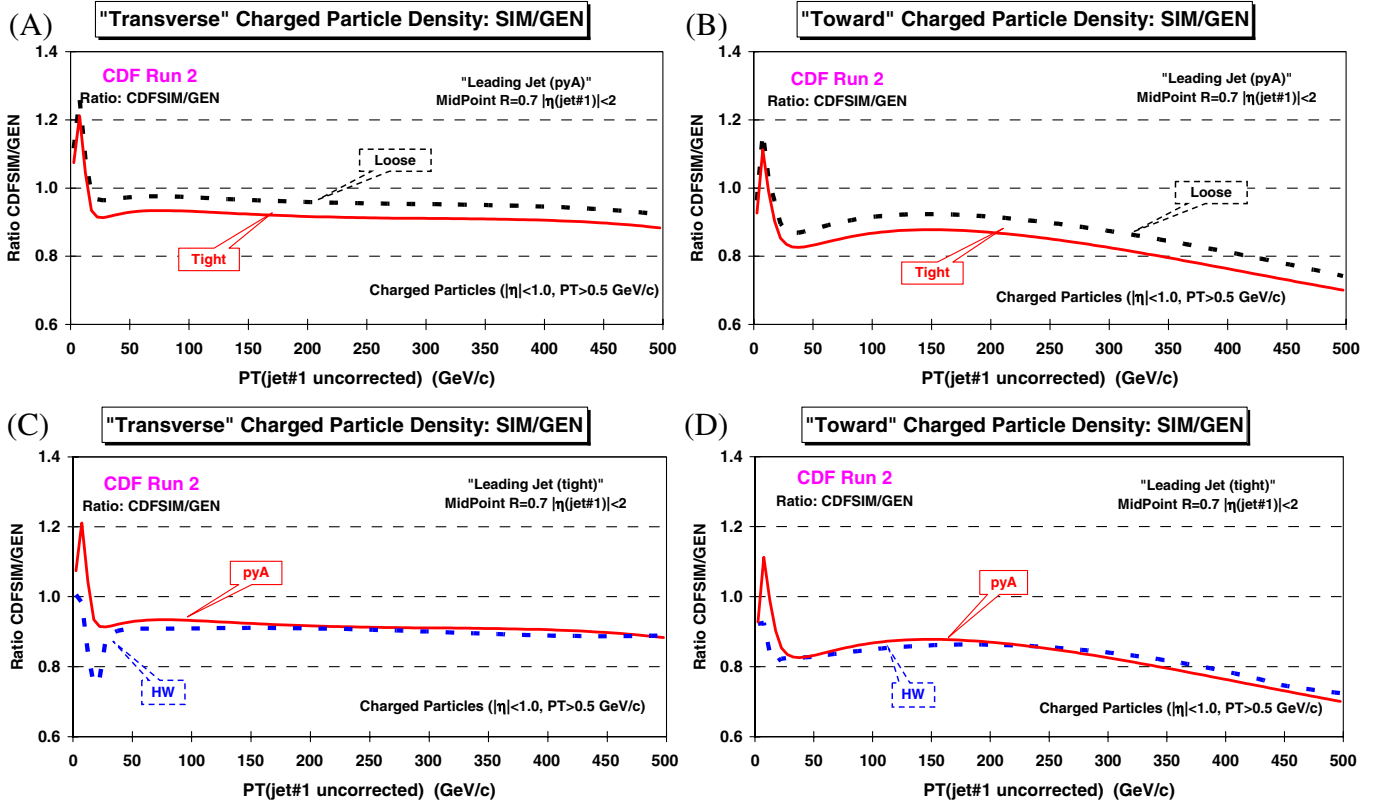


FIG. 4 (color online). (a) The response factors $F_{\text{res}} = \text{CDFSIM}/\text{GEN}$ for PYTHIA Tune A (pyA) with tight and loose track cuts for leading-jet events in the transverse region. (b) The same as (a) for the toward region. (c) Compares the response factors for PYTHIA Tune A (pyA) and HERWIG without MPI (HW) for tight track cuts for leading-jet events in the transverse region. (d) The same as (c) for the toward region. The correction factor is the reciprocal of the response factor ($F_{\text{cor}} = 1/F_{\text{res}}$).

$|\eta| < 1$ and sets the maximum p_T track cut to be equal to 1.5 times the E_T of the tower with the largest transverse energy. High- p_T mismeasured tracks do not deposit energy in the calorimeter. The two maximum p_T track cut methods produce slightly different correction factors; however, after correcting to the particle level the results are nearly identical. For the leading-jet analysis the differences were used as an additional systematic error.

Although we require one and only one high quality vertex, the observables in Table I can still be affected by pileup (more than one proton-antiproton collision in the event). Tracks are required to point back to the primary vertex, but the track observables are affected by pileup when two vertices overlap. Vertices within about 3 cm of each other merge together as one. In the leading-jet analysis we examined the effects of pileup by plotting the transverse charged-particle density and the charged-particle PT_{sum} density versus the instantaneous luminosity (with one and only one vertex). As the instantaneous luminosity increases so does the amount of pileup. We found that these observables did increase slightly with increasing luminosity (roughly linearly). The leading-jet observables in the transverse region are corrected for pileup by extrapolating to the low luminosity limit. To

correct the data, we define a low luminosity region $L_{\text{inst}} < 25 \times 10^{30} \text{ cm}^{-2} \text{ s}^{-1}$ (low) and a high luminosity region $L_{\text{inst}} > 25 \times 10^{30} \text{ cm}^{-2} \text{ s}^{-1}$ (high), where L_{inst} is the instantaneous luminosity. On a bin-by-bin basis, the ratio high/low and all/low was constructed, where all = high + low. The ratio high/low is very close to 1 (usually less than a 1% deviation from 1) and could simply have been absorbed into the overall systematic errors. However, in the leading-jet analysis we corrected the data for pileup by drawing a smooth curve through the ratio all/low and then dividing the data by this ratio. The size of the pileup correction was then taken as the systematic error in making the correction and added in quadrature with the other systematic errors. For the Drell-Yan analysis, the pileup corrections were less than 1% and were simply absorbed into the overall systematic errors.

Figure 5 shows the response factors $F_{\text{res}} = \text{CDFSIM}/\text{GEN}$ for the charged-particle density $dN/d\eta d\phi$ in the transverse region for leading-jet events and for Drell-Yan events. The response factors are similar, but not the same. In the Drell-Yan analysis we required the leptons to be isolated from other particles in the event. This introduces a bias against a very active underlying event which is compensated for by the correction factor.

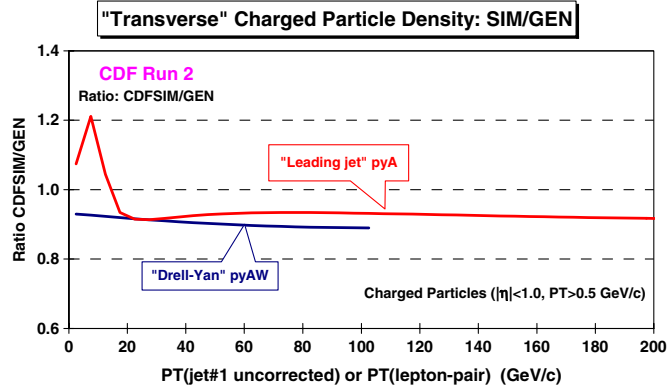


FIG. 5 (color online). The response factors $F_{\text{res}} = \text{CDFSIM/GEN}$ for the charged-particle density $dN/d\eta d\phi$ in the transverse region for leading-jet events and for Drell-Yan events. The plots show the response factors for PYTHIA Tune A (pyA) with tight track cuts (leading-jet) and for PYTHIA Tune AW (pyAW) with tight track cuts (Drell-Yan). The correction factor is the reciprocal of the response factor ($F_{\text{cor}} = 1/F_{\text{res}}$).

IV. RESULTS

A. Leading-jet and Drell-Yan topologies

Figure 6 shows the data on the density of charged particles and the scalar PT_{sum} density, respectively, for the toward, away, and transverse regions for leading-jet and Drell-Yan events. For leading-jet events the densities are plotted as a function of the leading-jet p_T and for Drell-Yan events there are plotted versus the p_T of the lepton pair. The data are corrected to the particle level and are compared with PYTHIA Tune A (leading-jet) and Tune AW (Drell-Yan) at the particle level. For leading-jet events at high $p_T(\text{jet}\#1)$ the densities in the toward and away regions are much larger than in the transverse region because of the towardside and awayside jets. At small $p_T(\text{jet}\#1)$ the toward, away, and transverse densities become equal and go to zero as $p_T(\text{jet}\#1)$ goes to zero. If the leading jet has no transverse momentum, then there can be no particles anywhere. In addition, there are numerous low transverse momentum jets and for $p_T(\text{jet}\#1) < 30$ GeV/c the leading jet is not always the jet resulting from the hard two-to-two scattering. This produces a bump in the transverse density in the range where the toward, away, and transverse densities become similar in size. For Drell-Yan events the

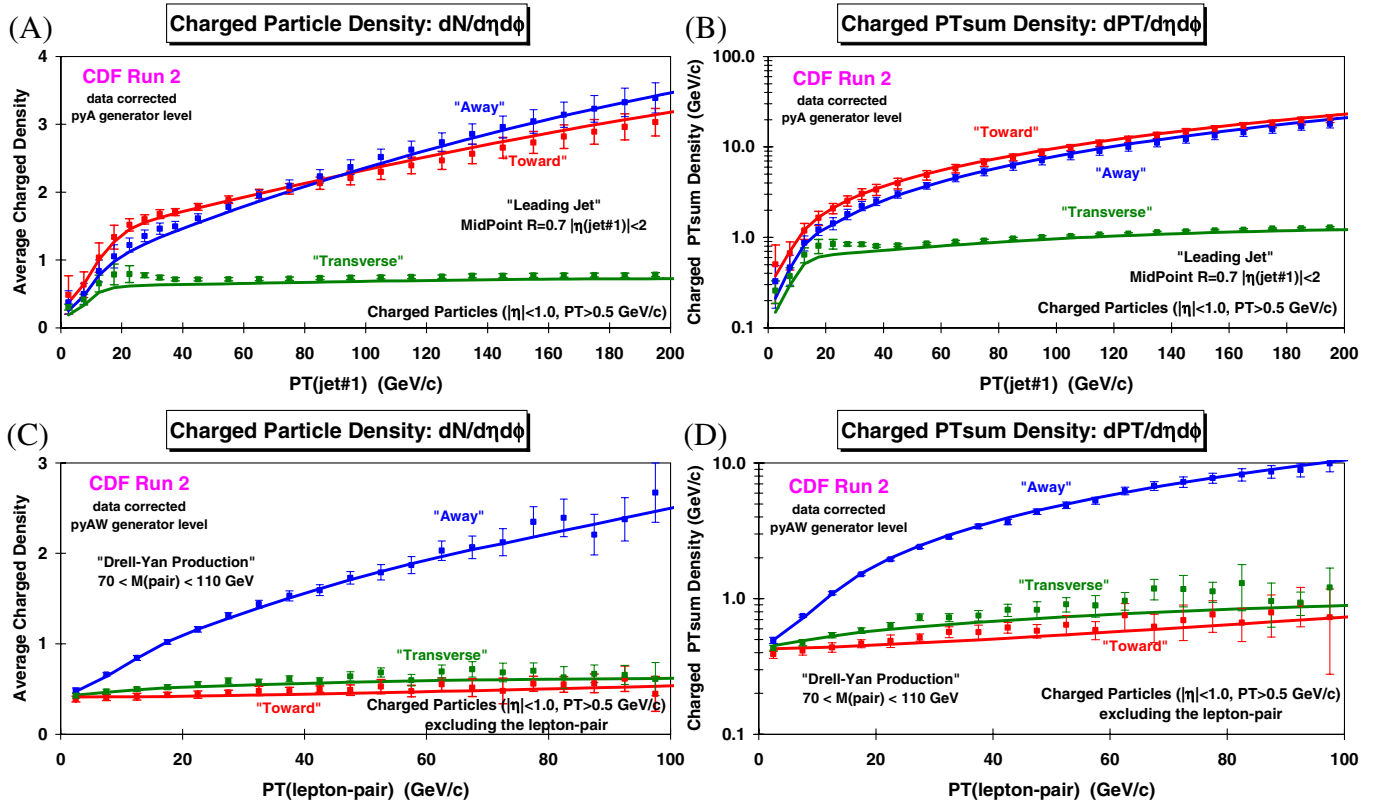


FIG. 6 (color online). (a) CDF data at 1.96 TeV on the density of charged particles in the toward, away, and transverse regions for leading-jet events compared with PYTHIA Tune A (pyA) at the particle level. (b) The same as (a) but for the scalar PT_{sum} density of charged particles. (c) CDF data at 1.96 TeV on the density of charged particles in the toward, away, and transverse regions for Drell-Yan events compared with PYTHIA Tune AW (pyAW) at the particle level. (d) The same as (c) but for the scalar PT_{sum} density of charged particles.

toward and transverse densities are both small and almost equal. The away density is large due to the awayside jet. The toward, away, and transverse densities become equal as p_T of the lepton pair goes to zero, but unlike the leading-jet case the densities do not vanish at $p_T(\text{lepton-pair}) = 0$. For Drell-Yan events with $p_T(\text{lepton-pair}) = 0$ the hard scale is set by the lepton-pair mass which is in the region of the Z boson, whereas in leading-jet events the hard scale goes to zero as transverse momentum of the leading jet goes to zero.

Figure 7 compares the data for leading-jet events with the data for Drell-Yan events for the density of charged

particles and the scalar PT_{sum} density, respectively, in the transverse region. The data are compared with PYTHIA Tune A (leading-jet), Tune AW (Drell-Yan), and HERWIG (without MPI). For large $p_T(\text{jet}\#1)$ the transverse densities are similar for leading-jet and Drell-Yan events as one would expect. HERWIG (without MPI) does not produce enough activity in the transverse region for either process. HERWIG (without MPI) disagrees more with the transverse region of Drell-Yan events than it does with the leading-jet events. This is because there is no final-state radiation in Drell-Yan production so that the lack of MPI becomes more evident.

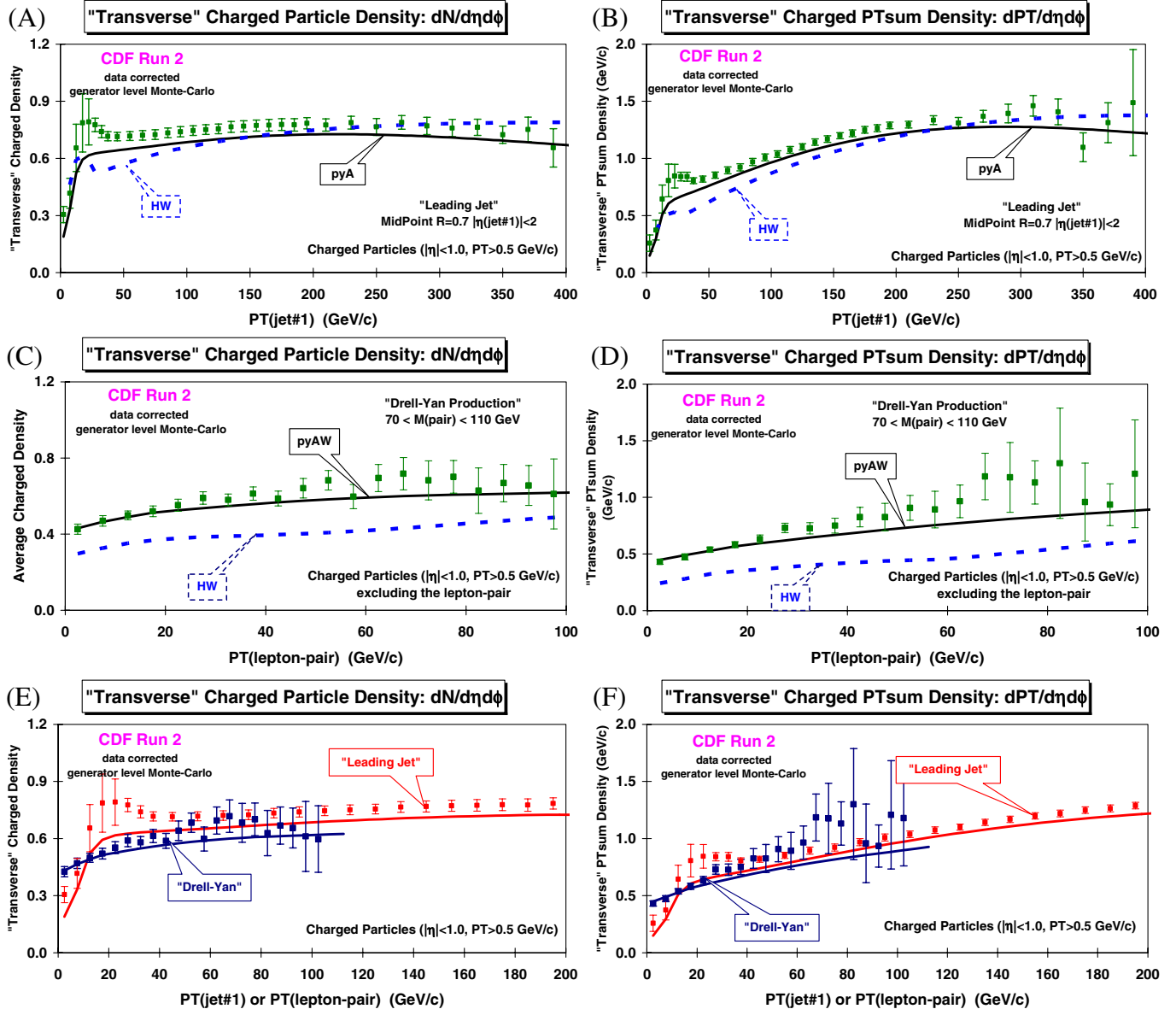


FIG. 7 (color online). (a) CDF data at 1.96 TeV on the density of charged particles in the transverse region for leading-jet events compared with PYTHIA Tune A (pyA) and HERWIG without MPI (HW). (b) The same as (a) but for the scalar PT_{sum} density of charged particles. (c) CDF data at 1.96 TeV on the density of charged particles in the transverse region for Drell-Yan events compared with PYTHIA Tune AW (pyAW) and HERWIG without MPI (HW). (d) The same as (c) but for the scalar PT_{sum} density of charged particles. (e) Compares (a) with (c) without the HERWIG curves. (f) Compares (b) with (d) without the HERWIG curves.

Figure 8 compares the data for leading-jet events with the data for Drell-Yan events for the average charged-particle p_T and the average maximum charged-particle p_T , respectively, in the transverse region. The data are compared with PYTHIA Tune A (leading-jet), Tune AW (Drell-Yan), and HERWIG (without MPI). MPI provides a hard component to the underlying event, and for HERWIG (without MPI) the p_T distributions in the transverse region for both processes are too soft, resulting in an average p_T and average PT_{\max} that are too small.

Figure 9 compares the data for leading-jet events with the data for Drell-Yan events for the density of charged particles and the scalar PT_{sum} density, respectively, for the transMAX and transMIN regions. The data are compared with PYTHIA Tune A (leading-jet), Tune AW (Drell-Yan), and HERWIG (without MPI). For events with large initial-state or final-state radiation the transMAX region would contain the third jet in high- p_T jet production or the second jet in Drell-Yan production. Thus, the transMIN region is very sensitive to the modeling of the underlying event.

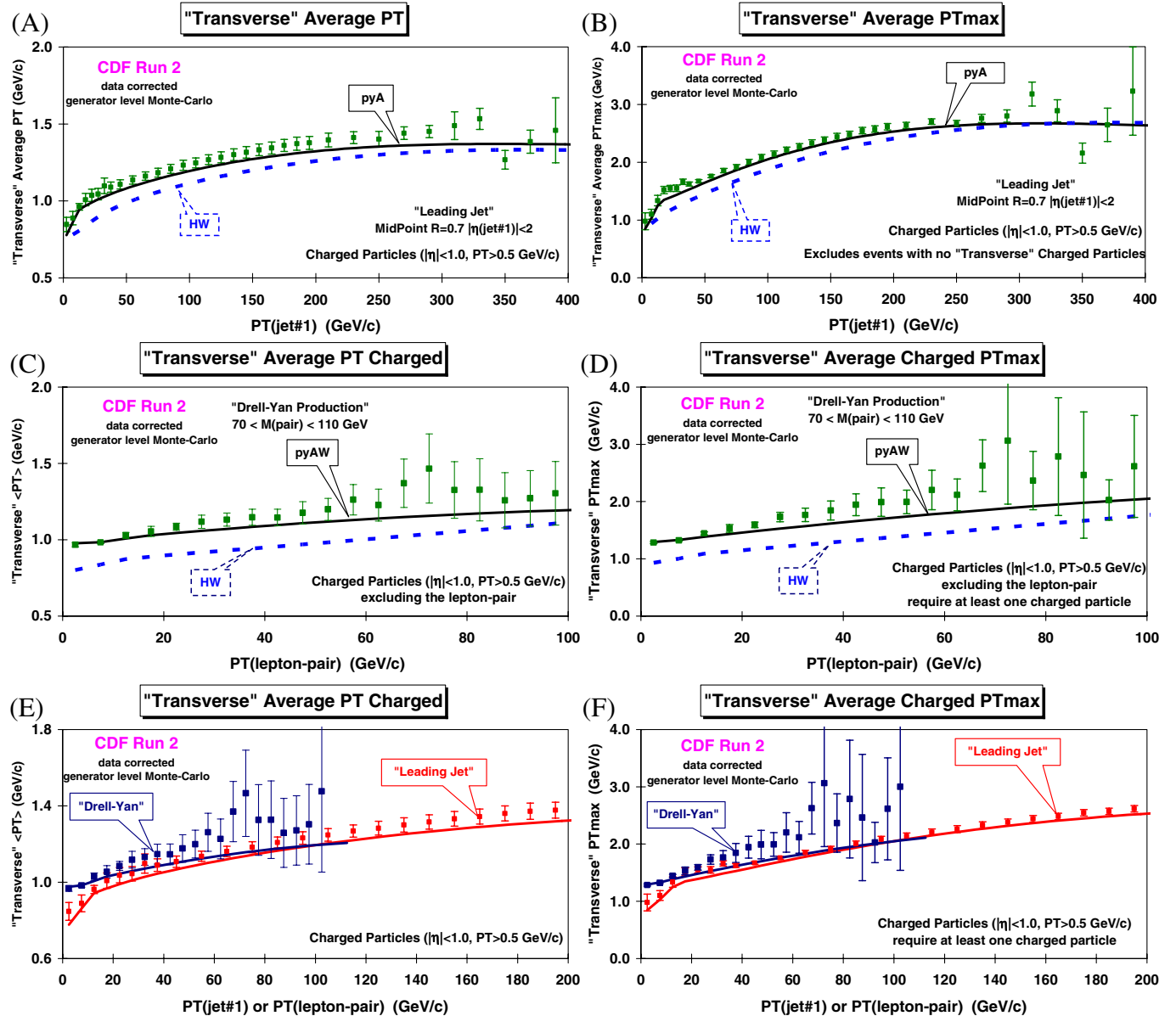


FIG. 8 (color online). (a) CDF data at 1.96 TeV on the average charged-particle transverse momentum in the transverse region for leading-jet events compared with PYTHIA Tune A (pyA) and HERWIG without MPI (HW). (b) The same as (a) but for the average maximum charge particle p_T . (c) CDF data at 1.96 TeV on the average charged-particle transverse momentum in the transverse region for Drell-Yan events compared with PYTHIA Tune AW (pyAW) and HERWIG without MPI (HW). (d) The same as (c) but for the average maximum charge particle p_T . (e) Compares (a) with (c) without the HERWIG curves. (f) Compares (b) with (d) without the HERWIG curves.

Figure 10 compares the data for leading-jet events with the data for Drell-Yan events for the density of charged particles and the scalar PT_{sum} density for $\text{transDIF} = \text{transMAX} - \text{transMIN}$. The data are compared with PYTHIA Tune A (leading-jet) and Tune AW (Drell-Yan). The transDIF region is sensitive to the hard initial and final-state radiation and is predicted to be very similar in the two processes. Figure 10 also compares the data for leading-jet events with the data for Drell-Yan events for the density of charged particles and the scalar PT_{sum} density in

the away region. The away-side jet pseudorapidity distribution and type (quark or gluon) is different for leading-jet and Drell-Yan events, so we do not expect the away region to be the same and it is not. However, PYTHIA Tune A and Tune AW describe the data very well.

B. The underlying event in Drell-Yan production

Figure 11 compares the data in the toward region with the data in the transverse region for Drell-Yan events for the density of charged particles, the scalar PT_{sum} density,

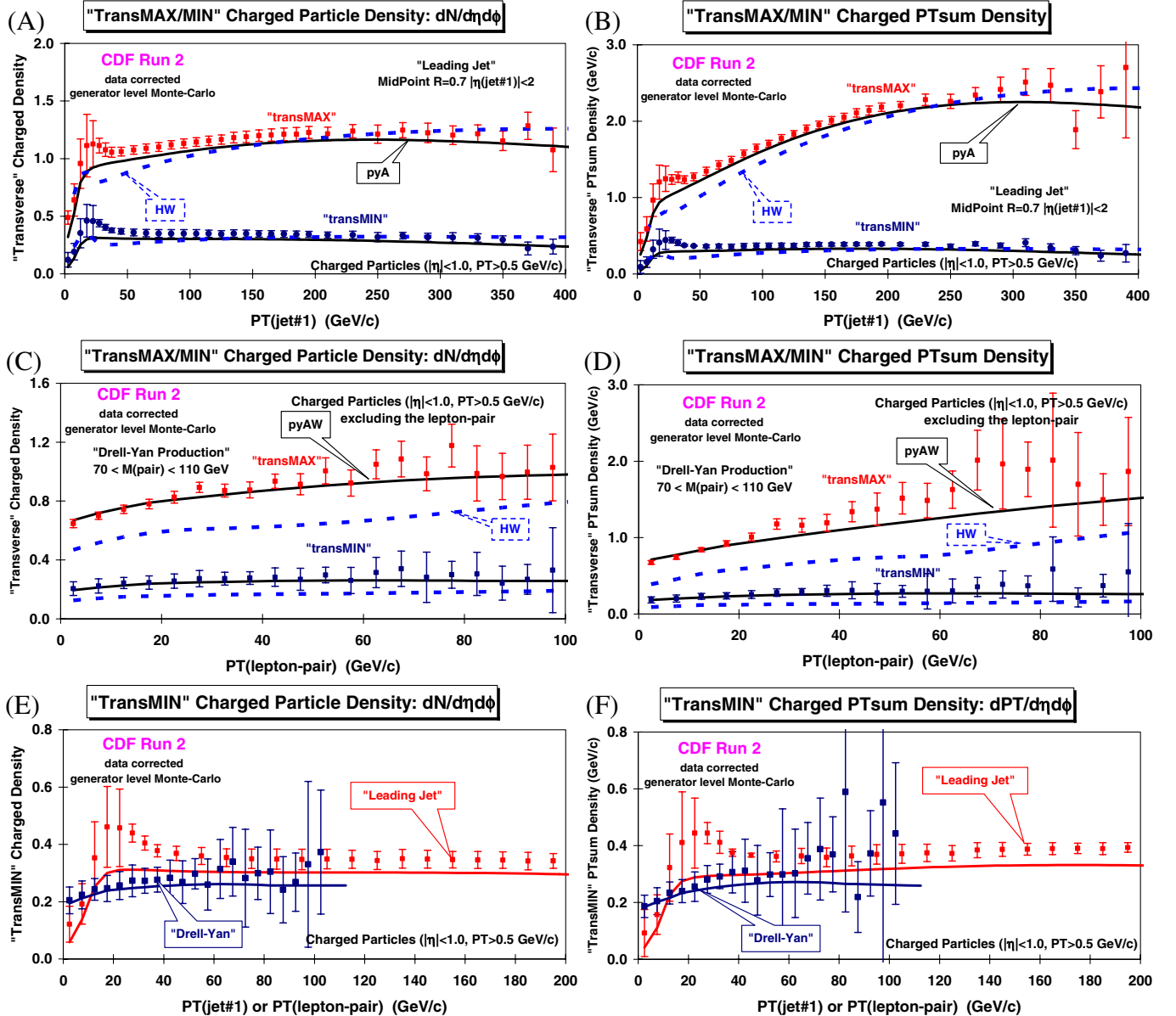


FIG. 9 (color online). (a) CDF data at 1.96 TeV on the transMAX and transMIN density of charged particles for leading-jet events compared with PYTHIA Tune A (pyA) and HERWIG without MPI (HW). (b) The same as (a) but for the scalar PT_{sum} density of charged particles. (c) CDF data at 1.96 TeV on the transMAX and transMIN density of charged particles for Drell-Yan events compared with PYTHIA Tune AW (pyAW) and HERWIG without MPI (HW). (d) The same as (c) but for the scalar PT_{sum} density of charged particles. (e) Compares the transMIN regions from (a) with (c) without the HERWIG curves. (f) Compares the transMIN regions from (b) with (d) without the HERWIG curves.

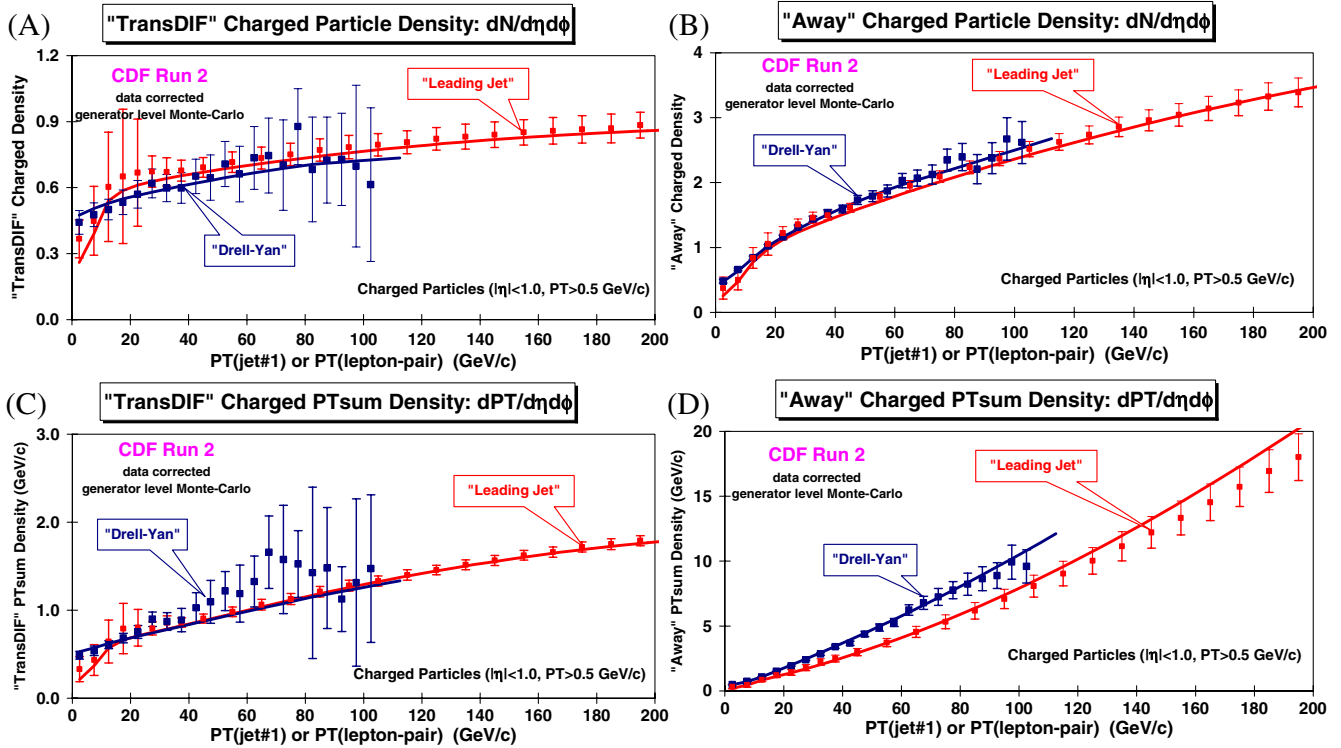


FIG. 10 (color online). (a) CDF data at 1.96 TeV on the transDIF density of charged particles for leading-jet and Drell-Yan events (transDIF = transMAX – transMIN). The Drell-Yan data are compared with PYTHIA Tune AW and the leading-jet data are compared with PYTHIA Tune A. (b) The same as (a) for density of charged particles in the away region. (c) The same as (a) for the transDIF scalar PT_{sum} density. (d) The same as (b) for the scalar PT_{sum} density.

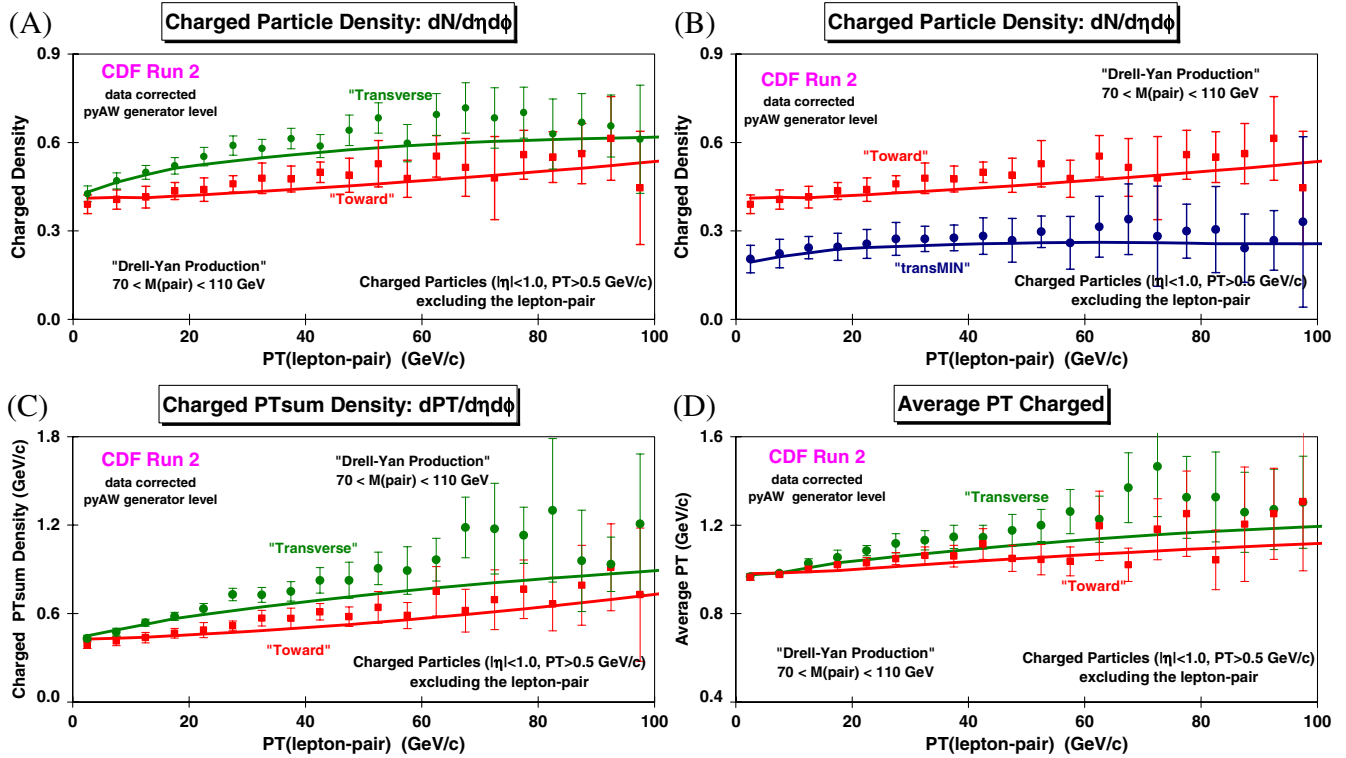


FIG. 11 (color online). (a) Compares the CDF data at 1.96 TeV on the charged-particle density in the toward and transverse regions for Drell-Yan events. Also shows the predictions of PYTHIA Tune AW. (b) Compares the CDF data on the charged-particle density in the toward and transMIN regions for Drell-Yan events and also shows the predictions of PYTHIA Tune AW. (c) The same as (a) for the scalar PT_{sum} density. (d) The same as (a) for the average charged-particle p_T of charged particles.

and the average charged-particle p_T . The data are compared with PYTHIA Tune AW. For high transverse momentum lepton-pair production, particles from initial-state radiation are more likely to populate the transverse region than the toward region and hence the densities are slightly larger in the transverse region. PYTHIA Tune AW describes this very nicely.

The most sensitive regions to the underlying event in Drell-Yan production are the toward and the transMIN regions, since these regions are less likely to receive contributions from the awayside jet and from initial-state

radiation. Figure 12 shows the data for Drell-Yan events for the density of charged particles and the scalar PT_{sum} density, respectively, in the toward and transMIN regions. The data are compared with PYTHIA Tune AW, Tune DW, the PYTHIA ATLAS tune, HERWIG (without MPI), and HERWIG (with JIMMY MPI). The densities are smaller in the transMIN region than in the toward region, and this is described well by PYTHIA Tune AW. Comparing HERWIG (without MPI) with HERWIG (with JIMMY MPI) clearly shows the importance of MPI in these regions. Tune AW and Tune DW are very similar. The ATLAS tune and

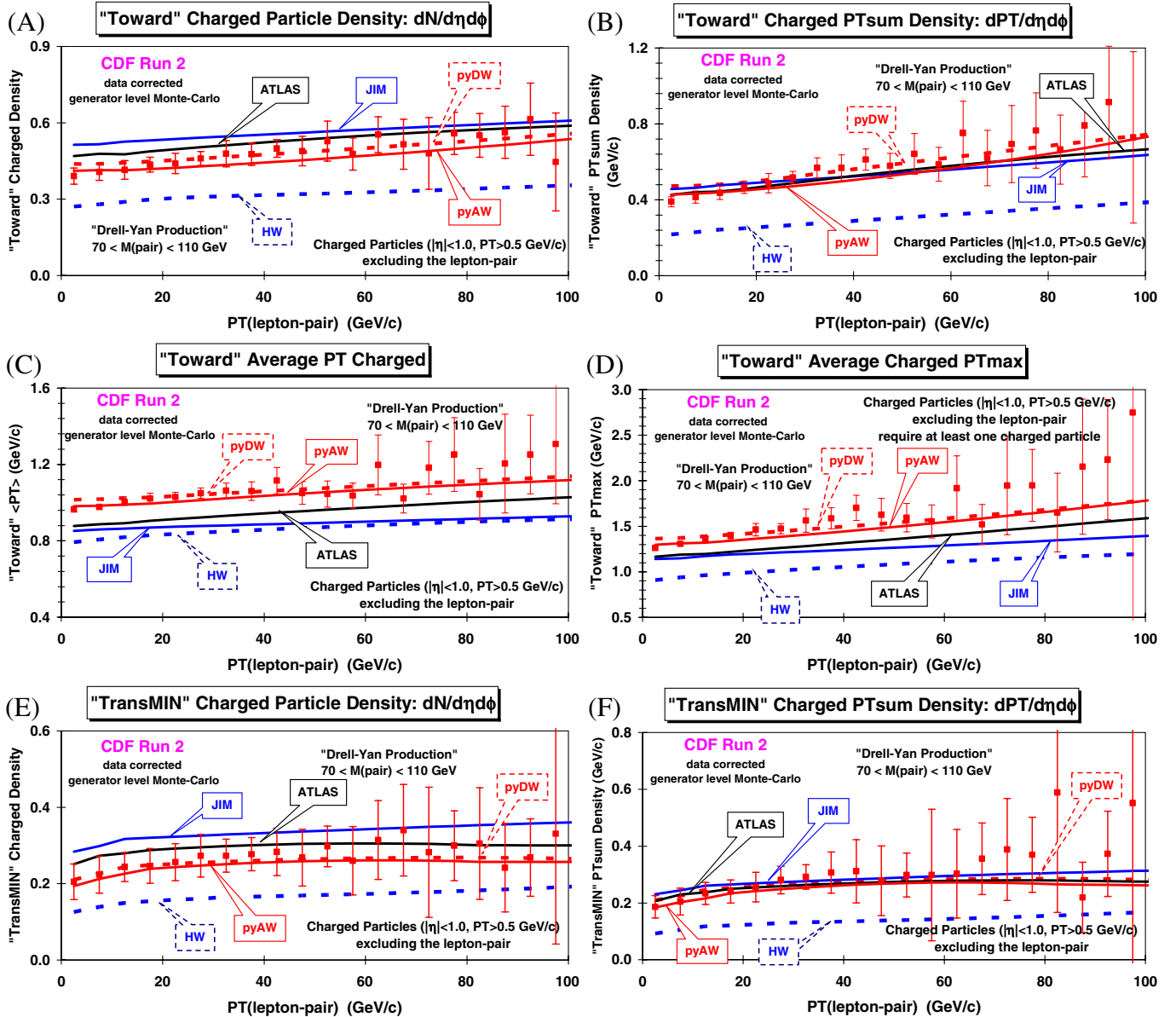


FIG. 12 (color online). (a) CDF data at 1.96 TeV on the density of charged particles in the toward region for Drell-Yan events. The data are compared with HERWIG without MPI (HW), HERWIG with JIMMY MPI (JIM), and three PYTHIA Tunes (pyAW, pyDW, and ATLAS). (b) The same as (a) but for the scalar PT_{sum} density of charged particles. (c) The same as (a) for the toward average charge particle p_T . (d) The same as (a) for the toward average maximum charge particle p_T . (e) The same as (a) for the transMIN charged particle density. (f) The same as (a) for the transMIN charged PT_{sum} density.

HERWIG (with JIMMY MPI) agree with Tune AW for the scalar PT_{sum} density in the toward and transMIN regions. However, both the ATLAS tune and HERWIG (with JIMMY MPI) produce too much charged-particle density in these regions. The ATLAS tune and HERWIG (with JIMMY MPI) fit the PT_{sum} density, but they do so by producing too many charged particles. They both have too soft a p_T spectrum in these regions. This can be seen clearly in Fig. 12, which shows the data for Drell-Yan events on the average charged-particle p_T and the average maximum charged-particle p_T , in the toward region compared with the QCD Monte Carlo models.

C. Extrapolating to the LHC

Figure 13 shows the extrapolation of PYTHIA Tune DW, Tune DWT, and HERWIG (without MPI) to 14 TeV (LHC) for the density of charged particles and the average transverse momentum of charged particles within the towards region of Drell-Yan production. The underlying-event activity is the same for proton-proton and proton-antiproton collisions. For HERWIG (without MPI) the toward region of Drell-Yan production does not change much in going from

the Tevatron to the LHC. Figure 13 also shows the extrapolation of PYTHIA Tune DW and Tune DWT to 14 TeV (LHC) for the transverse density of charged particles for leading-jet events. Models with multiple-parton interactions like PYTHIA Tune DW and Tune DWT predict that the underlying event will become much more active (with larger $\langle p_T \rangle$) at the LHC. PYTHIA Tune DW predicts about a factor of 2 increase in the activity of the underlying event as measured by the charged-particle density in the towards region of Drell-Yan production and the transverse region in leading-jet events. Tune DWT used the default value for PARP(90) and predicts an even greater increase in the activity of the underlying event at the LHC. However, Tune DWT produces less activity than Tune DW in the underlying event at energies below 1.96 TeV and the CDF data at 630 GeV [9] favor Tune DW over Tune DWT.

D. $\langle p_T \rangle$ versus the multiplicity: min-bias and Drell-Yan events

The total proton-antiproton cross section is the sum of the elastic and inelastic components: $\sigma_{\text{tot}} = \sigma_{\text{EL}} + \sigma_{\text{IN}}$. The inelastic cross section consists of three terms: single

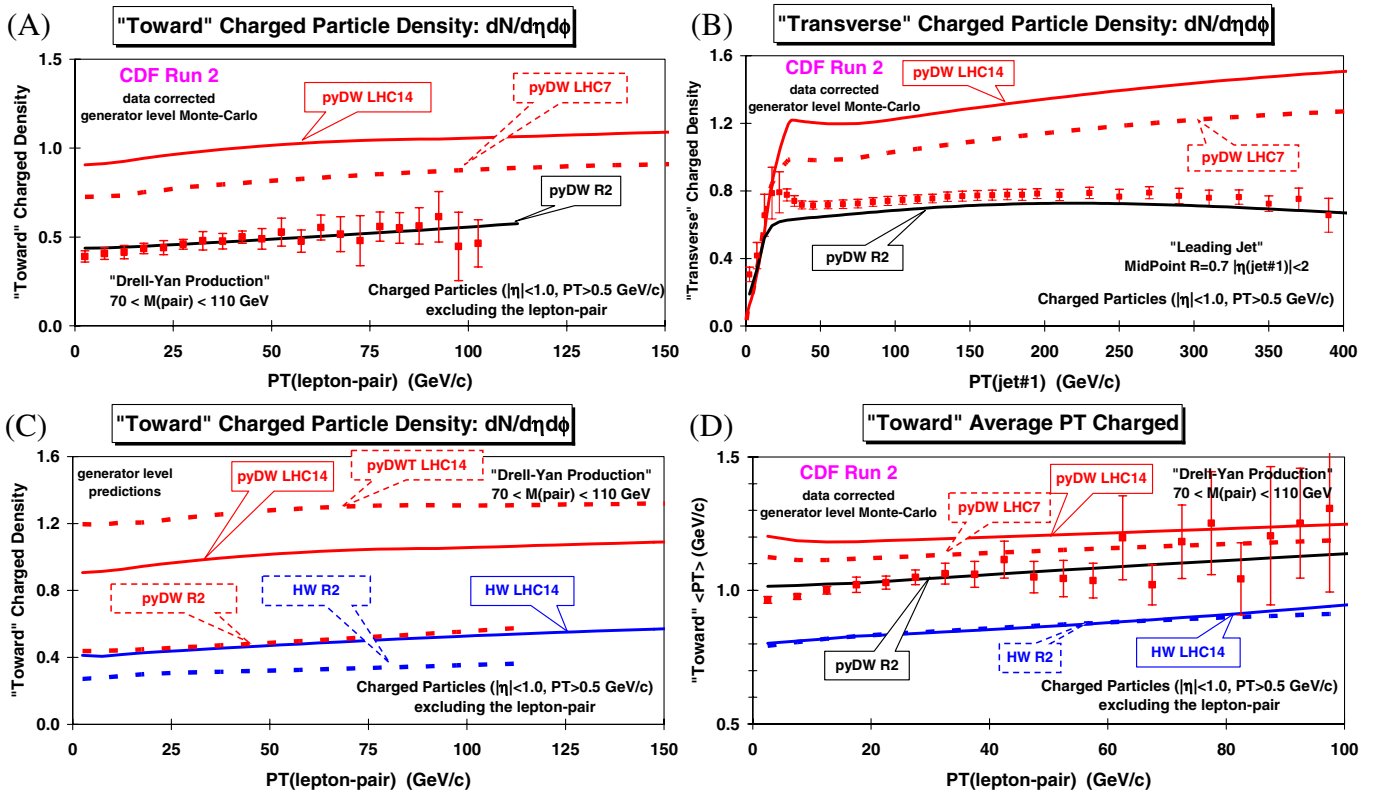


FIG. 13 (color online). (a) CDF data on the density of charged particles for Drell-Yan events in the toward region compared with Tune DW (pyDW R2). Also shows the predictions of Tune DW at 7 (LHC7) and 14 TeV (LHC14). (b) CDF data at the density of charged particles for leading-jet events in the transverse region compared with Tune DW (R2). Also shows the predictions of Tune DW at 7 (LHC7) and 14 TeV (LHC14). (c) Predictions of Tune DW and HERWIG without MPI (HW) for the density of charged particles for Drell-Yan events in the toward region at 1.96 TeV (R2) and at 14 TeV (LHC14). Also shows the prediction of Tune DWT at 14 TeV. (d) CDF data at on the average p_T of charged particles for Drell-Yan events in the toward region compared with Tune DW (pyDW R2) and HERWIG without MPI (HW R2). Also shows the predictions of Tune DW at 7 (LHC7) and 14 TeV (LHC14).

diffraction, double diffraction, and everything else (referred to as the hard core), $\sigma_{\text{IN}} = \sigma_{\text{SD}} + \sigma_{\text{DD}} + \sigma_{\text{HC}}$. For elastic scattering neither of the beam particles break apart. For single and double diffraction one or both of the beam particles are excited into a high mass color singlet state (i.e. N^* states) which then decays. Single and double diffraction also correspond to color singlet exchange between the beam hadrons. When color is exchanged, the outgoing remnants are no longer color singlets and one has a separation of color resulting in a multitude of quark-antiquark pairs being pulled out of the vacuum. The hard core component σ_{HC} involves color exchange and the separation of color. However, the hard core contribution has both a soft and a hard component. Most of the time the color exchange between partons in the beam hadrons occurs through a soft interaction with no high transverse momentum and the two beam hadrons ooze through each other producing lots of soft particles with a uniform distribution in rapidity and many particles flying down the beam pipe. Occasionally there is a hard scattering among the constituent partons producing outgoing particles and jets with high transverse momentum.

Minimum bias (min-bias) is a generic term which refers to events that are selected with a loose trigger that accepts a large fraction of the inelastic cross section. All triggers produce some bias and the term min-bias is meaningless until one specifies the precise trigger used to collect the data. The CDF min-bias trigger consists of requiring at least one charged particle in the forward region $3.2 < \eta < 5.9$ and simultaneously at least one charged particle in the backward region $-5.9 < \eta < -3.2$. Monte Carlo studies show that the CDF min-bias trigger collects most of the σ_{HC} contribution plus small amounts of single and double diffraction [20].

Minimum bias collisions are a mixture of hard processes (perturbative QCD) and soft processes (nonperturbative QCD) and are, hence, very difficult to simulate. Min-bias collisions contain soft beam-beam remnants, hard QCD two-to-two parton-parton scattering, and multiple-parton interactions (soft and hard). To correctly simulate min-bias collisions one must have the correct mixture of hard and soft processes together with a good model of the multiple-parton interactions. We have seen that multiple-parton interactions are a significant component of the underlying event in high- p_T jet production and in Drell-Yan lepton-pair production. Multiple-parton interactions are also an important component in min-bias collisions. Min-bias collisions are not the same as the underlying event in a hard-scattering process, since the rate at which MPI occurs is different, but they are related. Selecting a hard-scattering process such as high- p_T jet production or a lepton pair in the mass region of the Z boson corresponds to selecting a small fraction of min-bias collisions that are very central; the initial proton and antiproton collide with a small impact parameter. For these central collisions the probability of

additional parton-parton collisions is higher than it is for an average min-bias event.

The first model that roughly described min-bias collisions at CDF was PYTHIA Tune A. However, Tune A was not tuned to fit min-bias collisions. It was tuned to fit the activity in the underlying event in high transverse momentum jet production [4]. However, PYTHIA uses the same p_T cutoff for the primary hard two-to-two parton-parton scattering and for additional MPI. Hence, fixing the amount of multiple-parton interactions by setting the p_T cutoff allows one to run the hard two-to-two parton-parton scattering all the way down to $p_T(\text{hard}) = 0$ without hitting a divergence. For PYTHIA the amount of hard scattering in min-bias is, therefore, related to the activity of the underlying event in hard-scattering processes. Neither HERWIG (without MPI) or HERWIG (with JIMMY MPI) can be used to describe min-bias events since they diverge as $p_T(\text{hard})$ goes to zero.

Figure 14 shows CDF min-bias data corrected to the particle level at 1.96 TeV on the average p_T of charged particles, $\langle p_T \rangle$, versus the multiplicity for charged particles with $p_T > 0.4$ GeV/ c and $|\eta| < 1$ from Ref. [20]. The data are compared with PYTHIA Tune A, the PYTHIA ATLAS tune, and PYTHIA Tune A without MPI (pyAnoMPI). The average p_T is an important observable. The rate of change of $\langle p_T \rangle$ versus charged multiplicity is a measure of the amount of hard versus soft processes contributing to min-bias collisions and it is sensitive to the modeling of the multiple-parton interactions [21]. If only the soft beam-beam remnants contributed to min-bias collisions, then $\langle p_T \rangle$ would not depend on charged multiplicity. If one has two processes contributing, one soft (beam-beam remnants) and one hard (hard two-to-two parton-parton scattering), then demanding large multiplicity will preferentially select the hard process and lead to a high $\langle p_T \rangle$. However, we see that with only these two processes $\langle p_T \rangle$ increases much too rapidly as a function of multiplicity (see pyAnoMPI). Multiple-parton interactions provide another mechanism for producing large multiplicities that are harder than the beam-beam remnants but not as hard as the primary two-to-two hard scattering. PYTHIA Tune A gives a fairly good description of the $\langle p_T \rangle$ versus multiplicity, although not perfect. PYTHIA Tune A does a better job describing the data than the ATLAS tune. Both Tune A and the ATLAS tune include multiple-parton interactions, but with different choices for the color connections [22].

Figure 14 also shows the data at 1.96 TeV on the average p_T of charged particles versus the multiplicity for charged particles with for Drell-Yan events from this analysis. HERWIG (without MPI) predicts the $\langle p_T \rangle$ to rise too rapidly as the multiplicity increases. This is similar to the pyAnoMPI behavior in min-bias collisions. For HERWIG (without MPI) large multiplicities come from events with a high- p_T lepton pair and hence a large p_T awayside jet. This

can be seen clearly in Fig. 14, which also shows the average p_T of the lepton pair versus the charged multiplicity. Without MPI the only way of getting large multiplicity is with high- p_T (lepton-pair) events. For the models with MPI one can get large multiplicity either from high- p_T (lepton-pair) events or from MPI, and hence

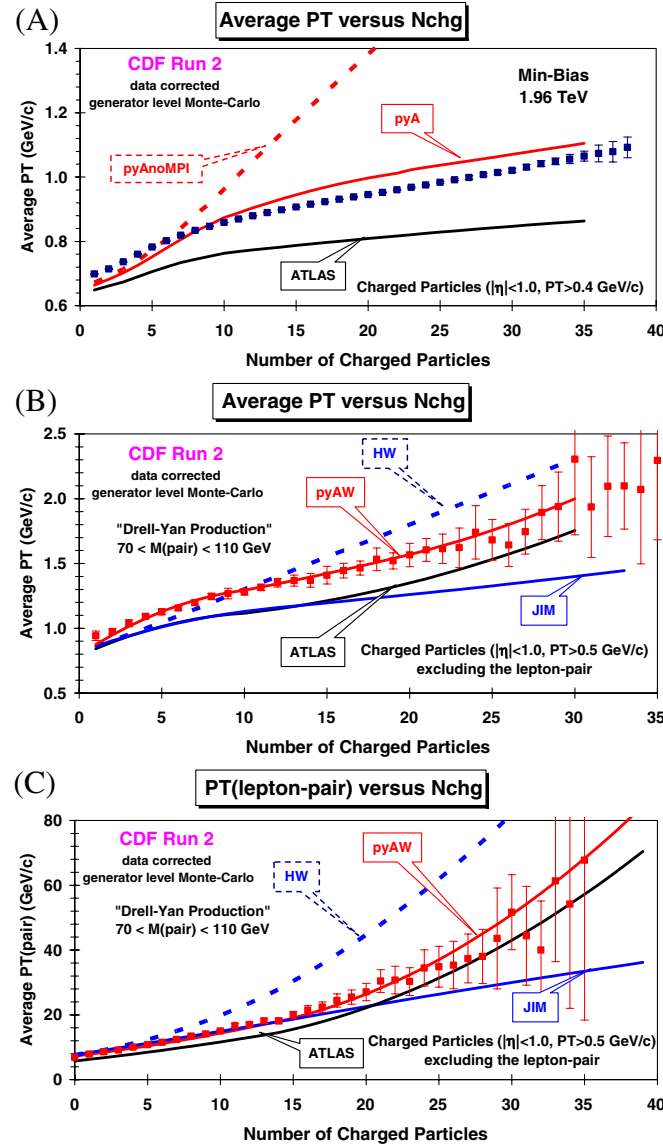


FIG. 14 (color online). (a) CDF min-bias data on the average p_T of charged particles versus the multiplicity for charged particles with $p_T > 0.4 \text{ GeV}/c$ and $|\eta| < 1$ from Ref. [20]. The data are compared with PYTHIA Tune A (pyA), the PYTHIA ATLAS tune, and PYTHIA Tune A without MPI (pyAnoMPI). (b) CDF data on the average p_T of charged particles versus the multiplicity for charged particles with $p_T > 0.5 \text{ GeV}/c$ and $|\eta| < 1$ for Drell-Yan events. The data are compared with PYTHIA Tune AW, the PYTHIA ATLAS tune, HERWIG without MPI (HW), and HERWIG with JIMMY MPI (JIM). (c) The same as (b) for the average p_T of the lepton pair versus the multiplicity for charged particles.

$\langle p_T(\text{lepton-pair}) \rangle$ does not rise as sharply with multiplicity in accord with the data. PYTHIA Tune AW describes the Drell-Yan data fairly well.

Figure 15 shows the data at 1.96 TeV on the average p_T of charged particles versus the multiplicity for Drell-Yan events in which $p_T(\text{lepton-pair}) < 10 \text{ GeV}/c$. We see that $\langle p_T \rangle$ still increases as the multiplicity increases although not as fast. If we require $p_T(\text{lepton-pair}) < 10 \text{ GeV}/c$, then HERWIG (without MPI) predicts that the $\langle p_T \rangle$ decreases slightly as the multiplicity increases. This is because without MPI and without the high- p_T awayside jet which is suppressed by requiring low p_T of the lepton pair, large multiplicities come from events with a lot of initial-state radiation and the particles coming from initial-state radiation are soft. PYTHIA Tune AW describes the behavior of $\langle p_T \rangle$ versus the multiplicity fairly well even when we select $p_T(\text{lepton-pair}) < 10 \text{ GeV}/c$.

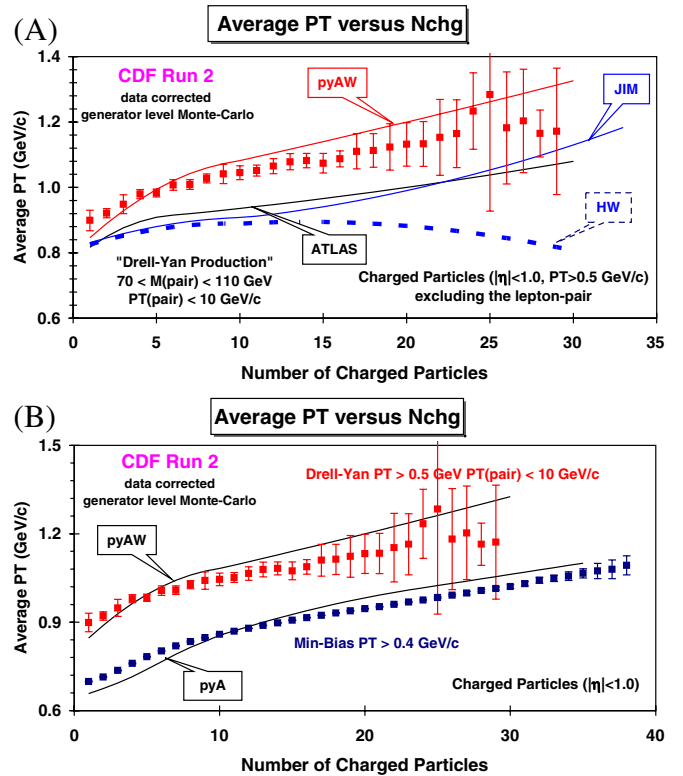


FIG. 15 (color online). (a) CDF data at 1.96 TeV on the average p_T of charged particles versus the multiplicity for charged particles for Drell-Yan events in which $p_T(\text{lepton-pair}) < 10 \text{ GeV}/c$. The data are compared with PYTHIA Tune AW (pyAW), the PYTHIA ATLAS tune, HERWIG without MPI (HW), and HERWIG with JIMMY MPI (JIM). (b) Comparison of the average p_T of charged particles versus the charged multiplicity for min-bias events from Ref. [20] with the Drell-Yan events with $p_T(\text{lepton-pair}) < 10 \text{ GeV}/c$ from this analysis. The min-bias data require $p_T > 0.4 \text{ GeV}/c$ and are compared with PYTHIA Tune A (pyA), while the Drell-Yan data require $p_T > 0.5 \text{ GeV}/c$ and are compared with PYTHIA Tune AW (pyAW).

Figure 15 also shows a comparison of the average p_T of charged particles versus the charged multiplicity for min-bias events from Ref. [20] with the Drell-Yan events with $p_T(\text{lepton-pair}) < 10 \text{ GeV}/c$. There is *a priori* no reason for the min-bias to behave like the Drell-Yan events with $p_T(\text{lepton-pair}) < 10 \text{ GeV}/c$. However, data have remarkably similar shape and are described fairly well by PYTHIA Tune A and Tune AW, respectively. This strongly suggests that MPI are playing an important role in both these processes.

V. SUMMARY AND CONCLUSIONS

Observables that are sensitive to the underlying event in high transverse momentum jet production (leading-jet events) and Drell-Yan lepton-pair production in the mass region of the Z boson (Drell-Yan events) have been presented and compared with several QCD Monte Carlo model tunes. The data are corrected to the particle level and compared with the Monte Carlo models at the particle level. The underlying event is similar for leading-jet and Drell-Yan events as one would expect. This analysis provides data that can be used to test and improve the QCD Monte Carlo models of the underlying event that are used to simulate hadron-hadron collisions. The data presented here are also important for tuning the new QCD Monte Carlo MPI models [20,21].

PYTHIA Tune A and Tune AW do a good job in describing the data on the underlying-event observables for leading-jet and Drell-Yan events, respectively, although the agreement between predictions and data is not perfect. The leading-jet data show slightly more activity in the underlying event than PYTHIA Tune A. PYTHIA Tune AW is essentially identical to Tune A for leading-jet events. All the tunes with MPI agree better than HERWIG without MPI. This is especially true in the toward region in Drell-Yan production. Adding JIMMY MPI to HERWIG greatly improves the agreement with data, but HERWIG with JIMMY MPI produces a charged-particle p_T spectrum that is considerably softer than the data. The PYTHIA ATLAS tune also produces a charged-particle p_T spectra that is considerably softer than the data.

The behavior of the average charged-particle p_T versus the charged-particle multiplicity is important. The rate of change of $\langle p_T \rangle$ versus charged multiplicity is a measure of the amount of hard versus soft processes contributing, and it is sensitive the modeling of the multiple-parton interactions. PYTHIA Tune A and Tune AW do a good job in describing the data on $\langle p_T \rangle$ versus multiplicity for min-bias and Drell-Yan events, respectively, although again the

agreement between the models and data is not perfect. The behavior of $\langle p_T \rangle$ versus multiplicity is remarkably similar for min-bias events and Drell-Yan events with $p_T(\text{lepton-pair}) < 10 \text{ GeV}/c$, suggesting that MPI are playing an important role in both these processes.

Models with multiple-parton interactions like PYTHIA Tune DW predict that the underlying event will become much more active (with larger $\langle p_T \rangle$) at the LHC. For HERWIG (without MPI) the toward region of Drell-Yan production does not change much in going from the Tevatron to the LHC. PYTHIA Tune DW predicts about a factor of 2 increase in the activity of the underlying event in going from the Tevatron to the LHC as measured by the charged-particle density in the towards region of Drell-Yan production and the transverse region in leading-jet events. Tune DWT predicts an even greater increase in the activity of the underlying event at the LHC. However, Tune DWT produces less activity than Tune DW in the underlying event at energies below 1.96 TeV. Tune DW does a better job in fitting the CDF underlying-event data at 630 GeV [9] and is hence favored over Tune DWT. At present, PYTHIA tunes with PARP(90) around the value of Tune AW and Tune DW (≈ 0.25) seem to be preferred. We will learn a lot about the energy dependence of MPI by comparing the Tevatron results with the early LHC measurements and precise measurements at the LHC require good modeling of the underlying event.

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