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Orbital characters of bands in the iron-based superconductor BaFe_{1.85}Co_{0.15}As₂

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The unconventional superconductivity in the newly discovered iron-based superconductors is intimately related to its multiband/multiorbital nature. Here we report the comprehensive orbital characters of the low-energy three-dimensional electronic structure in BaFe_{1.85}Co_{0.15}As₂ by studying the polarization and photon-energy dependence of angle-resolved photoemission data. While the distributions of the d_{xz} , d_{yz} , and $d_{3z^2-r^2}$ orbitals agree with the prediction of density functional theory, those of the d_{xy} and $d_{x^2-y^2}$ orbitals show remarkable disagreement with theory. Our results point out the inadequacy of the existing band structure calculations and, more importantly, provide a foundation for constructing the correct microscopic model of iron pnictides.

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I. INTRODUCTION

Unlike the cuprates, the low-energy electronic structure of the iron-based superconductors is dominated by multiple bands and orbitals.¹⁻³ Theoretically, it has been proposed that the Fermi surface sheets with multiple orbitals could result in a strong anisotropy and amplitude variation of the superconducting gaps.^{4,5} Experimentally, a recent study on Ba_{0.6}K_{0.4}Fe₂As₂ shows that the superconducting gap sizes are different at the same Fermi momentum for two bands with different orbital characters.⁶ To construct correct models for the iron-based superconductors, and to understand the unconventional superconductivity, it is thus critical to experimentally identify the orbital characters of the low-energy electronic structure.

There have been inconsistencies over the orbital identities of the bands among theories^{2,3} and a few experiments.⁷⁻¹² Moreover, various physical properties of iron-based superconductors are featured with three-dimensional (3D) characters.¹³ For example, the band structure in the so-called “122” series of iron pnictides is rather 3D,^{10,14} and the gap dependency on the out-of-plane momentum (k_z) has recently been reported.⁶ It has been pointed out that the three-dimensionality of the electronic structure might be essential in inducing both the spin density wave (SDW) and superconductivity.¹⁵ However, the orbital character distribution along the k_z direction in 3D momentum space has not been exposed so far.

Here we report a systematic angle-resolved photoemission spectroscopy (ARPES) study on the orbital character of the electronic structure in an electron-doped “122” compound, BaFe_{1.85}Co_{0.15}As₂. We have observed strong polarization dependency for all the bands near the Fermi energy (E_F) and have obtained a comprehensive picture of the orbital characters of these bands. We confirm previous theoretical findings that the bands with d_{xz} and d_{yz} orbitals form two hole pockets around the zone center and one electron pocket around the zone corner, and the $d_{3z^2-r^2}$ (or d_{z^2} for simplicity) orbital is mostly included in the two bands well below E_F . More importantly, we find that certain bands predicted to be of the

d_{xy} orbital in theory are actually mixed with the $d_{x^2-y^2}$ orbital around the zone center and are dominated by the $d_{x^2-y^2}$ orbital around the zone corner. Furthermore, we have identified the orbital characters of the band structure in 3D momentum space. Our results provide explicit knowledge for constructing the theoretical model of iron-based superconductors.

II. POLARIZATION-DEPENDENT ARPES

BaFe_{1.85}Co_{0.15}As₂ single crystals were synthesized by a self-flux method¹⁶ with a superconducting transition temperature (T_c) of 25 K and no SDW or structural transition. Unconventional superconductivity with multigap behavior has been reported in such material.¹⁷⁻²¹ As this system is optimally doped with electrons, the low-energy band structure is well occupied and free from the complications of the electronic reconstruction in the SDW state.²²⁻²⁶ Because the arsenic ions in the FeAs layer are situated in two inequivalent positions, there are two iron ions per unit cell. The Brillouin zone of BaFe_{1.85}Co_{0.15}As₂ is shown in Figs. 1(a) and 1(b). To compare with theory, we define the k_x and k_y directions to be the Fe-Fe bond directions, as shown in the unfolded Brillouin zone (dashed lines) for one iron ion per unit cell. Two high-symmetry directions, $\Gamma(Z)$ -M(A) and $\Gamma(Z)$ -X(R), are illustrated by red arrows #1 and #2, respectively, in Fig. 1(b).

The polarization sensitivity of the orbitals in ARPES is a powerful tool to identify the orbital characters of band structure.²⁷ The matrix element of the photoemission process can be described by

$$|M_{f,i}^{\mathbf{k}}|^2 \propto |(\phi_f^{\mathbf{k}}|\hat{\mathbf{e}} \cdot \mathbf{r}|\phi_i^{\mathbf{k}})|^2,$$

where $\hat{\mathbf{e}}$ is the unit vector of the electric field of the light.²⁷ For high-kinetic-energy photoelectrons, the final-state wave function $\phi_f^{\mathbf{k}}$ can be approximated by a plane-wave state $e^{i\mathbf{k}\cdot\mathbf{r}}$ with \mathbf{k} in the mirror plane as plotted in Fig. 1(c). Consequently, it is always even with respect to the mirror plane. For the p (or s) experimental geometry in Fig. 1(c), because $\hat{\mathbf{e}}$ is parallel (or perpendicular) to the mirror plane, $\hat{\mathbf{e}} \cdot \mathbf{r}$ is even

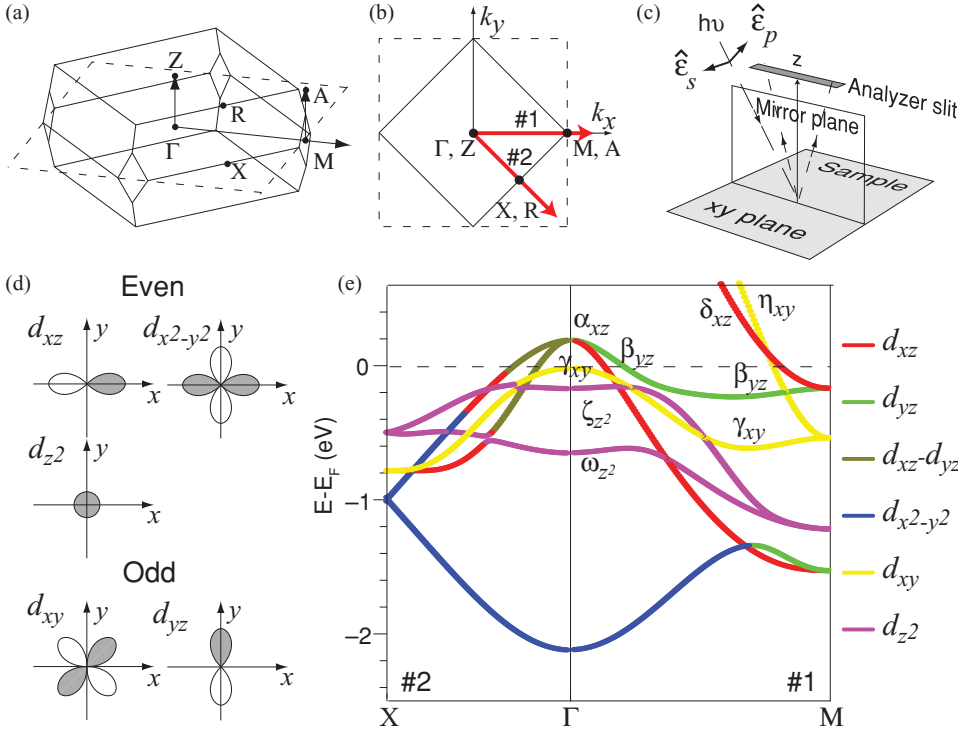


FIG. 1. (Color online) The experimental setup and definitions. (a) The Brillouin zone of $\text{BaFe}_{1.85}\text{Co}_{0.15}\text{As}_2$. (b) Two-dimensional plot of the simplified Brillouin zone (solid line) and the unfolded Brillouin zone (dashed line). (c) Experimental setup for polarization-dependent ARPES. For the p (or s) experimental geometry, the electric field direction of the incident photons $\hat{\epsilon}_p$ (or $\hat{\epsilon}_s$) is parallel (or perpendicular) to the mirror plane defined by the analyzer slit and the sample surface normal. (d) Illustration of the spatial symmetry of the $3d$ orbitals with respect to the mirror plane formed by surface normal and cut #1 in panel (b), that is, the xz plane. (e) A typical orbital assignment of bands of iron pnictide as calculated in Ref. 3.

(or odd). Therefore, to have a finite matrix element, that is, to be observed, the initial state ϕ_i^k has to be even (or odd) in the p (or s) geometry.

Considering the spatial symmetry of the $3d$ orbitals [Fig. 1(d)], when the analyzer slit is along the high-symmetry direction of the sample, the photoemission signal of certain orbitals would appear or disappear by specifying the polarization directions as summarized in Table I. For example, with respect to the mirror plane formed by direction #1 and sample surface normal (or the xz plane), the even orbitals (d_{xz} , d_{z^2} , and $d_{x^2-y^2}$) and the odd orbitals (d_{xy} and d_{yz}) could be only observed in the p and s geometries, respectively. Note that d_{xz} and d_{yz} are not symmetric with respect to the mirror plane defined by direction #2 and surface normal and thus could be observed in both the p and the s geometries.

Figure 1(e) shows the band calculation of iron-based superconductors in a two-dimensional (2D) band model reproduced from Ref. 3, which was confirmed by many other calculations and widely adopted in various models. The bands are labeled with corresponding orbital characters as subscripts. Around the Γ point, the d_{z^2} orbital contributes to the two bands

TABLE I. The possibility to detect $3d$ orbitals along two high-symmetry directions in the p and s geometries by polarization-dependent APRES.

High-symmetry direction	Experimental geometry	$3d$ orbitals				
		d_{xz}	$d_{x^2-y^2}$	d_{z^2}	d_{yz}	d_{xy}
#1 $\Gamma(\text{Z})$ -M(A)	p	✓	✓	✓		
	s				✓	✓
#2 $\Gamma(\text{Z})$ -X(R)	p	✓		✓	✓	✓
	s	✓	✓		✓	

ζ_{z^2} and ω_{z^2} well below E_F . There are three holelike bands near E_F . α_{xz} and β_{yz} cross E_F forming two hole pockets, while the band top of γ_{xy} is just below E_F . The β_{yz} and γ_{xy} disperse to the lower binding energies at M. Two bands, δ_{xz} and η_{xy} , form two electron pockets, and they are degenerate with β_{yz} and γ_{xy} at M, respectively. The $d_{x^2-y^2}$ orbital was found to be irrelevant to the low-energy electronic structure.

ARPES measurements were performed at Beamline 1 of the Hiroshima Synchrotron Radiation Center (HSRC). By rotating the entire photoemission spectroscope around the incoming photon beam, one can switch between the s and the p polarization geometries. All the data were taken with a Scienta electron analyzer at 30 K, the overall energy resolution is 15 meV, and the angular resolution is 0.3° . The sample was cleaved *in situ* and measured under ultrahigh vacuum of 5×10^{-11} torr.

III. RESULTS

A. Orbital characters of bands around the zone center

The photoemission data taken near the zone center are shown in Fig. 2. Two bands assigned as ζ and ω could be observed at about 170 and 500 meV below E_F . Both bands only emerge in the p geometry and could not be observed in the s geometry along both #1 and #2 directions [Figs. 2(a)–2(d)]. Based on Table I, these two bands are made of d_{z^2} orbital, assuming they consist of a single orbital, as suggested in Fig. 1(e). The energy distribution curves (EDCs) at $k_{\parallel} = 0 \text{ \AA}^{-1}$ taken with different photon energies are stacked in Fig. 2(f), so that we could track the ζ and ω bands along Γ -Z. The peak position of ω shows a periodic variation if we take the inner potential to be 15 eV [Fig. 2(e)]. Furthermore, the peak intensities of ω and ζ show periodic anticorrelation with the

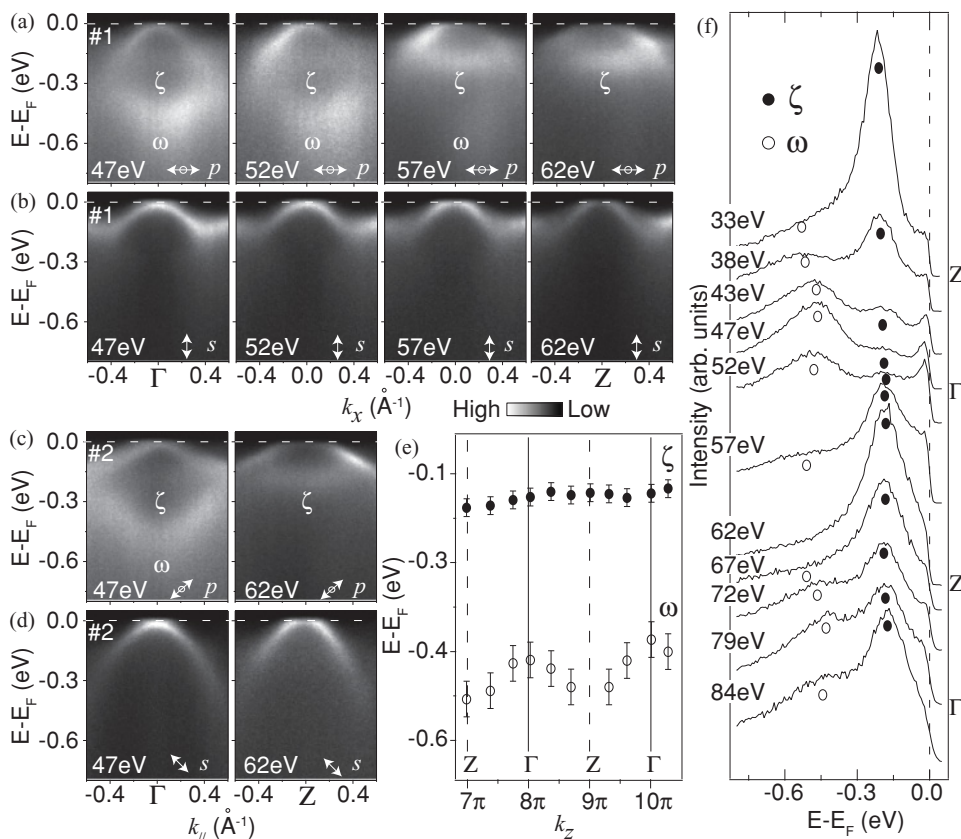


FIG. 2. Photoemission data around the zone center. (a),(b) The photon-energy dependence of the photoemission intensity [$I(k,\omega)$] along Γ (Z)-M(A) (#1) in the p and s geometries, respectively. (c),(d) The photon-energy dependence of $I(k,\omega)$ along Γ (Z)-X(R) (#2) in the p and s geometries, respectively. (e) The k_z dispersion of the ζ and ω bands along Γ -Z. (f) The photon-energy dependence of EDCs sampled along Γ -Z.

photon energy. While ω is at its strongest at Γ (47 and 79 eV), ζ is mostly enhanced at Z (33 and 62 eV). The anticorrelation of intensities of these two bands could be naturally explained by the different wave-function distributions along the z direction for the bonding and antibonding states formed by the d_{z^2} orbitals.²⁸ Summing up these facts, we attribute ζ and ω to the d_{z^2} orbital, which agrees with the assignments in Fig. 1(e).

Near the Fermi energy, three bands (α , β , and γ) could be identified around Γ in Fig. 3. α only shows up in the p geometry, while β only appears in the s geometry, exhibiting opposite spatial symmetries [Figs. 3(a)–3(e)]. The band tops of both α and β are below E_F and degenerate at the Γ point. Note that the bands with d_{xz} and d_{yz} orbitals should be degenerate at the Γ point due to the fourfold symmetry of the sample. Therefore, we attribute the α and β bands to be of the d_{xz} and d_{yz} orbitals, respectively, in this direction. In addition, along the Γ -M direction, α disperses to the binding energy over 200 meV, and β disperses to about 150 meV and then bends over to lower binding energies toward the M point, which is consistent with the behavior of α_{xz} and β_{yz} predicted by theory in Fig. 1(e). The d_{xz} and d_{yz} orbitals have no definite symmetry with respect to the Γ -X direction. The strong polarization dependence of α and β observed along Γ -X direction [Figs. 3(f)–3(j)] is most likely due to the hybridization of the d_{xz} and d_{yz} orbitals. For example, $\frac{d_{xz}+d_{yz}}{\sqrt{2}}$ and $\frac{d_{xz}-d_{yz}}{\sqrt{2}}$ are of even and odd spatial symmetry with respect to the Γ -X direction, respectively. Actually, because of the fourfold symmetry, the d_{xz} component of α along the k_x axis has to become d_{yz} along the k_y axis. Therefore, the

equal mixing of these two orbitals along the Γ -X direction is expected. Similar arguments hold for the β band.

The γ band could be observed in both geometries, indicating that it is a mixture of both odd and even orbitals. With respect to Γ -M, the odd orbital candidates are d_{xy} and d_{yz} . As the d_{yz} orbital contributes to β , the odd orbital in γ is likely d_{xy} , as the γ_{xy} band proposed in Fig. 1(e). The even orbital candidates include d_{xz} , d_{z^2} , and $d_{x^2-y^2}$. Since $\alpha(d_{xz})$, $\zeta(d_{z^2})$, and $\omega(d_{z^2})$ bands contain no odd orbital, the odd orbitals in γ are less likely mixed with the d_{xz} and d_{z^2} orbital in this momentum region. We thus deduce that the most possible even orbital in γ is $d_{x^2-y^2}$. Because both d_{xy} and $d_{x^2-y^2}$ orbitals mainly spread in the FeAs plane, the γ band is expected to be more two dimensional than the α and β bands. Such an assignment is consistent with the polarization dependence of data along Γ -X as well [Figs. 3(f)–3(j)], except $d_{x^2-y^2}$ is the odd orbital and d_{xy} is the even one in this direction.

To confirm such 2D nature of the γ band and explore the 3D characters of the α and β bands, photon-energy-dependent data around the zone center are shown in Fig. 4. The β band shifts away from the zone center when approaching the Z point [Figs. 4(b)–(d)], which is also illustrated by the shift of the momentum distribution curve (MDC) peaks of β in Fig. 4(e). The k_z dependence of the α band shows similar behavior as β , with its band top below E_F at Γ , but slightly above E_F at Z [Figs. 4(g)–4(j)]. Moreover, the in-plane dispersion of α becomes quite flat and eventually intersects the γ band around Z as shown in Figs. 4(j) and 4(l). The MDC peak positions of γ show almost no photon-energy dependence in Fig. 4(k), and no obvious k_z influence on its in-plane dispersion is

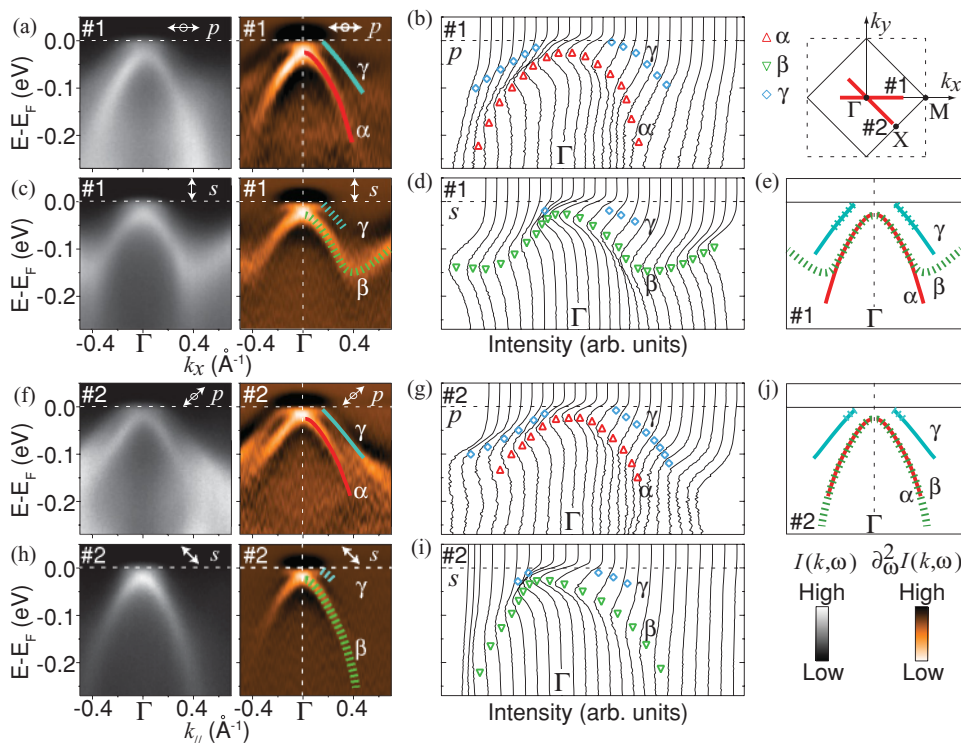


FIG. 3. (Color online) The polarization-dependent APRES data around Γ taken with 47-eV photons. (a) $I(k, \omega)$ map and its second derivative with respect to energy $[\partial^2 I(k, \omega)/\partial \omega^2]$ taken with the p geometry along #1 direction as marked with a red thick solid line in the top-right inset. (b) The EDCs for the data in panel (a). (c), (d) Same as in panels (a) and (b), but taken with the s geometry. (e) The sketch of the band structure in panels (a) and (c). (f)–(j) Same as in panels (a)–(e), but taken along cut #2 in the top-right inset. The band dispersions are identified with the help of second derivative image and EDCs. The solid and dashed sketches of the bands represent the even and odd components of the bands, respectively. The color scale of $I(k, \omega)$ and $\partial^2 I(k, \omega)/\partial \omega^2$ are shown in the bottom-right inset. These color scales and notations are used throughout the paper.

observed as well. Therefore, unlike the strong k_z dependence of α and β , the weak k_z dependence of γ is consistent with the nature of $d_{x^2-y^2}$ and d_{xy} , which are more in-plane than the other $3d$ orbitals. Note that γ is more intensive in both polarization geometries near Z, likely due to matrix element effects related to its distribution along the crystallographic c axis.

The band structure and orbital characters around Z are shown in Fig. 5. The α and β bands cross E_F along both the Z-A and the Z-R directions and thus form two hole pockets around the Z point. The polarization dependence of the spectra near Z are similar to that near Γ . Therefore, around the zone center, all the bands keep the same spatial symmetry at different k_z 's.

B. Orbital characters of bands around the zone corner

The polarization-dependent photoemission data around the zone corner are plotted in Fig. 6. The electronlike δ band is degenerate with β at about 40 meV below E_F at M. The δ band is much more pronounced in the p geometry, while β is very strong in the s geometry. On the other hand, some residual spectral weight of β could still be observed in the p geometry at M, and some infinitesimal δ band spectral weight could be observed in the s geometry at both M and A, indicating a light mixing between β and δ . If we assume that the orbital character of β does not change abruptly, its main composition can be then attributed to the odd d_{yz} orbital, based on its d_{yz} orbital character around Γ . In Fig. 1(e), the δ band is attributed to the even d_{xz} orbital, which is consistent with its mainly even character observed here.

The η band disperses over 150 meV below E_F , which is very close to the γ band at about -220 meV. This pair of

bands (γ and η) was previously predicted to be of d_{xy} with odd symmetry [Fig. 1(e)]. However, surprisingly, both of them are actually even and only appear in the p geometry around the zone corner [Figs. 6(a) and 6(f)].

To further settle the orbital characters around the zone corner, Figs. 7(a)–7(f) show the photoemission intensity map around the M and A points by rotating the sample azimuthally. Obviously, multiple disconnected parts of the Fermi surface are either enhanced or suppressed by the p or s geometry, which again demonstrates the multi-orbital nature of the Fermi surface. Although the polarization selection rules of different orbitals in Table I are not strict when the cuts are not exactly along the high-symmetry direction, they still manifest themselves effectively, for example, as strong intensity modulations along Fermi surface sheets composed of sections with opposite spatial symmetries. Combining the Fermi surface parts observed in both the p and s geometries, one obtains two electron pockets around the M and A points [Figs. 7(c) and 7(f)]. The inner electron pocket is contributed by the δ band. If we assume δ to be the δ_{xz} band along Γ -M in Fig. 1(e), its horizontal (k_x) sections are made of the d_{xz} orbital, which could be thus observed in the p geometry [the red solid line in Figs. 7(a) and 7(d)]. Since the fourfold and translational symmetries of the crystal would result in the rotating distribution of the d_{xz} and d_{yz} orbitals, the vertical (k_y) sections of δ are made of the d_{yz} orbital, which could be thus observed in the s geometry [the green dashed line in Figs. 7(b) and 7(e)]. Moreover, the elliptical shape of the δ pocket rotates 90° from M to A, which is consistent with the 3D character of the Brillouin zone [Fig. 7(j)]. As shown in Figs. 7(h) and 7(i), the distance between the two Fermi crossings of the δ band along the vertical (k_y) direction shrinks significantly from M to A. Therefore, the alternating even and odd orbital nature

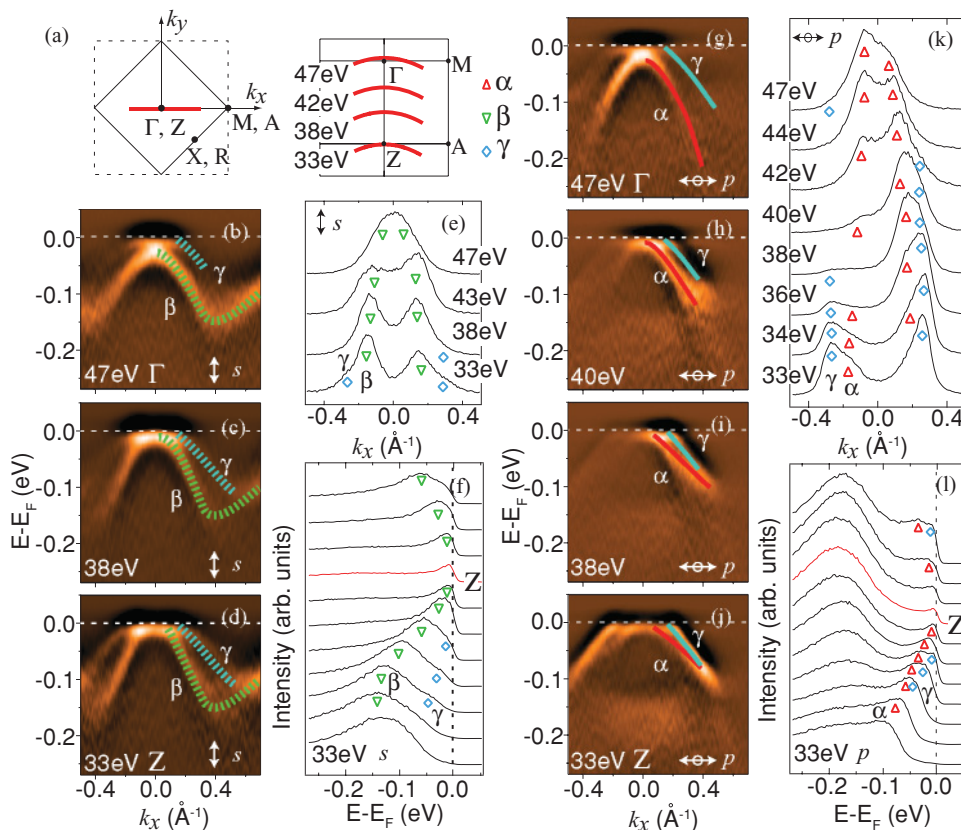


FIG. 4. (Color online) Photon-energy dependence of bands near the zone center. (a) Illustration of the cuts sampled with different photon energies. (b)–(d) $\partial^2 I(k, \omega) / \partial \omega^2$ taken in the s geometry with 47-, 38-, and 33-eV photons, respectively. (e) The photon-energy dependence of the MDCs at 30 meV below E_F in the s geometry. (f) The EDCs for the data in panel (d). (g)–(j) $\partial^2 I(k, \omega) / \partial \omega^2$ taken in the p geometry with 47-, 40-, 38-, and 33-eV photons, respectively. (k) The photon-energy dependence of the MDCs at 30 meV below E_F in the p geometry. (l) The EDCs for the data in panel (j).

and strong k_z dependence of δ electron pocket further supports its d_{xz} and d_{yz} orbital characters.

The outer squarelike pocket, as shown by the blue solid line in Figs. 7(a) and 7(d), is from the η band. It is even and thus could be only observed in the p geometry. Among the three even orbitals, we could first exclude the d_{xz} orbital in the η pocket; otherwise, one would have observed the corresponding d_{yz} component of η along k_y in the s geometry, due to the fourfold symmetry of the sample. Furthermore, the η Fermi pocket is particularly intense in the Z-A or Z-M direction, as indicated by the thicker blue lines in Figs. 7(a) and 7(d). The rotation of these intense sections from A to M suggests the existence of orbital mixing rather than simple matrix element effect. Therefore, the η band should be the mixture of $d_{x^2-y^2}$ and d_{z^2} orbitals. As shown in Fig. 7(g), the peak intensity of η shows asymmetry with respect to the A point. Similarly, in Fig. 7(d), the intense section of η in the first Brillouin zone is more intensive than that in the extended Brillouin zone on the right side. This contradicts the behavior of the $d_{x^2-y^2}$ orbital, as the matrix element or photoemission intensity of $d_{x^2-y^2}$ should have become stronger as the in-plane momentum became larger. We thus conclude that the intense sections of η are consist of the d_{z^2} orbital over these momentum regions. On the other hand, considering the facts that the η electron pocket shows no k_z dependence as the corresponding MDCs peak positions do not change with the photon energies in Fig. 7(g), the η pocket is mostly made of the more in-plane $d_{x^2-y^2}$ orbital.

The orbital assignment of the η pocket could naturally explain the orbital character of the γ band near the zone corner.

Since γ is actually the same band as η in the unfolded Brillouin zone, they are thus degenerate and have the same orbital characters at M.^{2,3} Therefore, the γ band is also dominated by $d_{x^2-y^2}$ at the zone corner (possibly some d_{z^2}), which indicates that the $d_{x^2-y^2}$ component of γ grows from Γ to M.

IV. DISCUSSION

Based on the above analysis, the orbital characters of the low-energy electronic structure of $\text{BaFe}_{1.85}\text{Co}_{0.15}\text{As}_2$ are summarized in Fig. 8. The bands with d_{xz} and d_{yz} orbitals, including α , β , and δ , are well consistent with the band calculation in Fig. 1(e). The in-plane dispersion of the α band varies much more strongly with k_z than that of the β band. This difference in the k_z dependencies between α and β is consistent with the density function theory calculations in BaFe_2As_2 ²⁹ and previous ARPES studies.^{6,14} In particular, in the 3D model of BaFe_2As_2 ,¹⁵ the sizes of the hole pockets at Z are larger than that at Γ , which quantitatively agrees with our results as well. However, our experimental result seems not to support the orbital character change from Z-A to Z-R predicted in Ref. 15, which states that the α and β would be correspondingly composed of d_{xy} and $d_{x^2-y^2}$ orbitals along the Z-R direction.

The most remarkable findings here are the orbital characters of the γ and η bands. They are dominated by the $d_{x^2-y^2}$ orbital around the zone corner, and γ is a mixture of d_{xy} and $d_{x^2-y^2}$ at the zone center. However, in most band calculations, the γ and η bands are proposed to be purely d_{xy} , and the $d_{x^2-y^2}$ states are far from E_F . Our results disprove this picture. Furthermore,

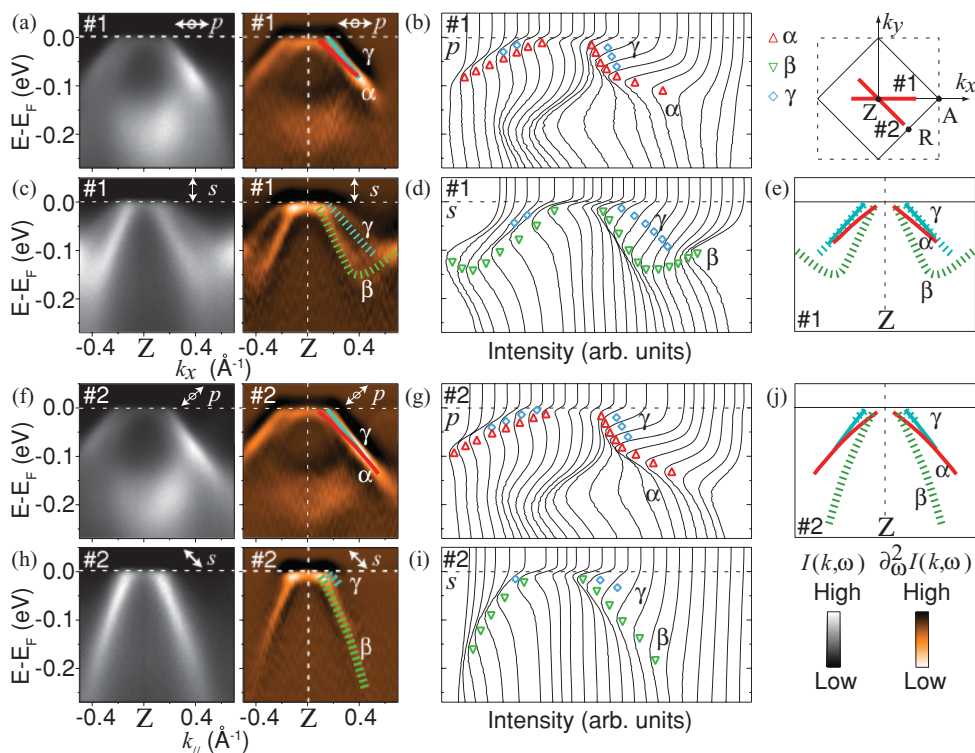


FIG. 5. (Color online) The polarization-dependent APRES data around Z taken with 33-eV photons. (a) $I(k,\omega)$ map and $\partial^2 I(k,\omega)/\partial\omega^2$ taken with the p geometry along #1 direction as marked with a red thick solid line in the top-right inset. (b) The EDCs for the data in panel (a). (c),(d) Same as in panels (a) and (b), but taken with the s geometry. (e) The sketch of the band structure in panels (a) and (c). (f)–(j) Same as in panels (a)–(e), but taken along #2 direction as marked with a red thick solid line in the top-right inset.

the band top of the γ band is found to be higher than those of the α and β at the zone center, which is also inconsistent with the band calculations shown in Fig. 1(e). Considering the 2D character of the γ and η bands, the d_{xy} and $d_{x^2-y^2}$ orbitals should interact less with states outside the Fe layer (such as As $4p_z$ states) than the other $3d$ orbitals. Therefore, our results indicate that the energy of the d_{xy} and $d_{x^2-y^2}$ orbitals may not be calculated correctly, and electron-electron correlations may

be strongly orbital dependent and result in the deviation from a simple local density approximation (LDA) calculation.

The pairing in the superconducting state can be strongly tied to the simultaneous existence of both Fermi surface sheets between the zone center and zone corner in iron-based superconductors. It has been suggested that the pairing is the strongest when the nesting condition between these two Fermi surface sheets are reached.^{30,31} However, in Fig. 8, α and β

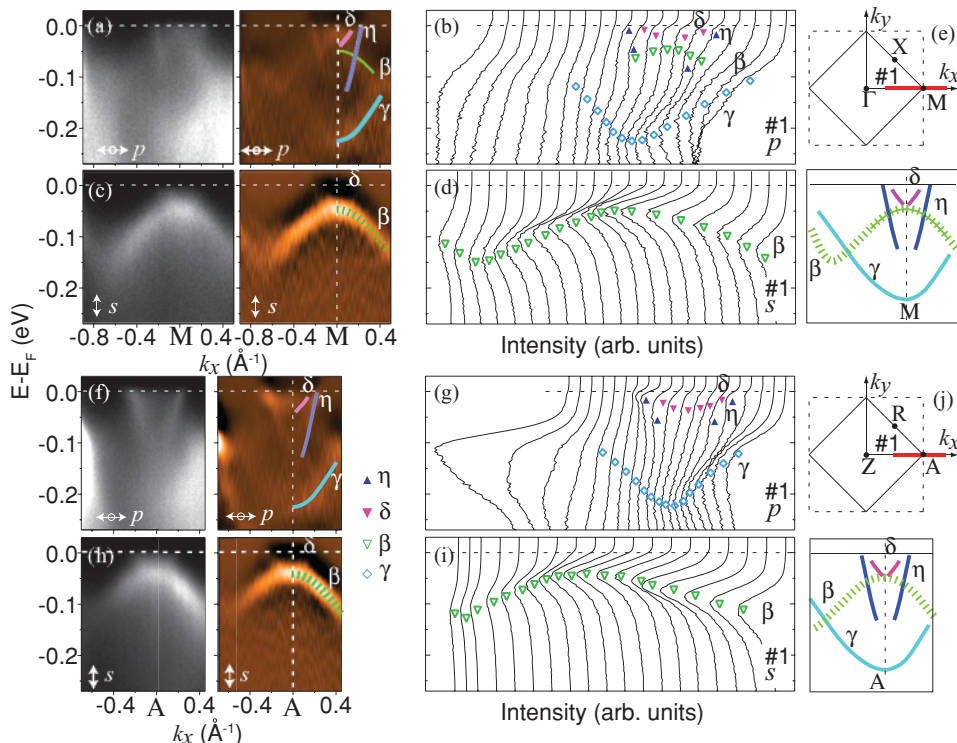


FIG. 6. (Color online) The polarization-dependent APRES data around the zone corner. (a) $I(k,\omega)$ and $\partial^2 I(k,\omega)/\partial\omega^2$ map taken with the p geometry and 51-eV photons across M along the red thick line in the inset of panel (e). (b) The EDCs for the data in panel (a). (c),(d) Same as in panels (a) and (b), but taken under the s geometry. (e) The sketch of the band structure in panels (a) and (c). (f)–(j) Same as in panels (a)–(e), but taken with 38-eV photons across A along the red thick line in the inset of panel (j).

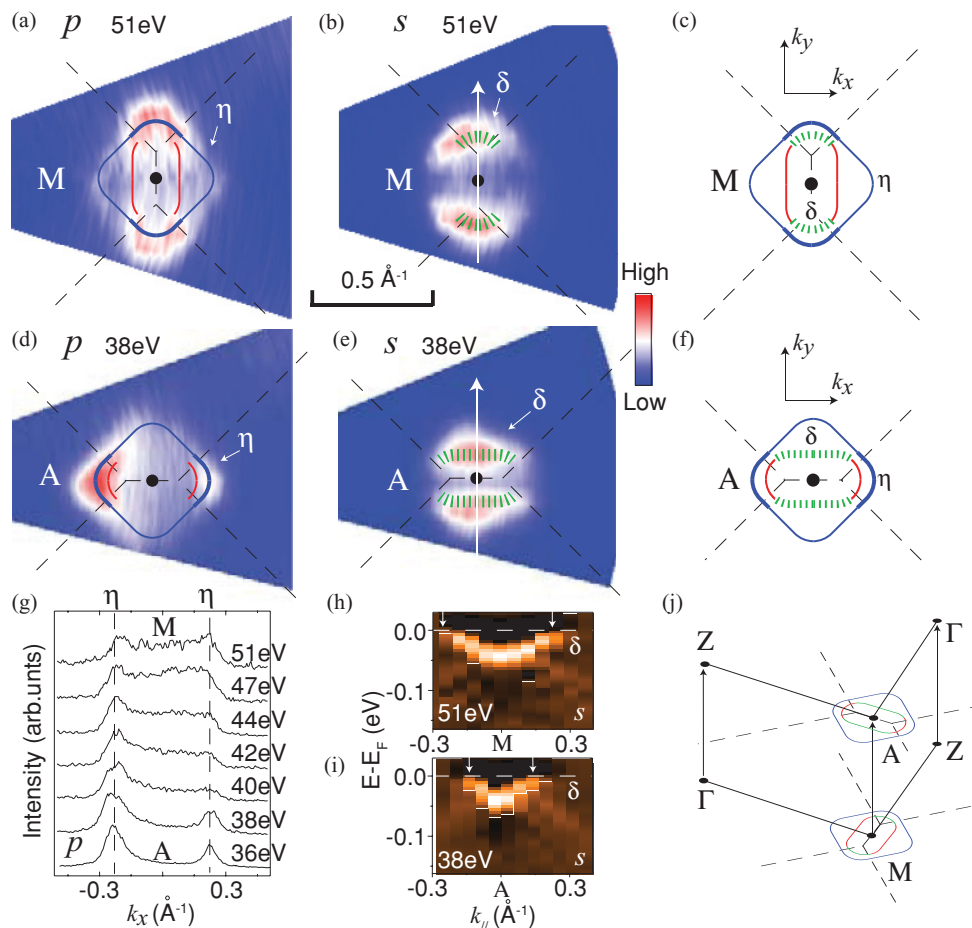


FIG. 7. (Color online) Fermi surface around the zone corner. (a),(b) Photoemission intensity maps taken with 51-eV photons around M in the p and s geometries, respectively. (c) The sketch of the Fermi surface sheets observed in panels (a) and (b). (d)–(f) Same as in panels (a)–(c), but taken with 38-eV photons around A. (g) The photon-energy dependence of the MDCs at E_F in the p geometry along the Γ (Z)-M(A) direction. (h),(i) $\partial^2 I(k, \omega)/\partial \omega^2$ map obtained by merging selected EDCs in the s geometry along the white arrows in panels (b) and (e), respectively. (j) The illustration of the Fermi surface at A and M in the 3D Brillouin zone.

show Fermi crossings with the distance about 0.18 \AA^{-1} at the Z point, which is still smaller than the 0.2 \AA^{-1} of the δ band at the M point. Therefore, the α and β bands do not nest well with the δ band. The γ and η bands form one hole pocket and one electron pocket respectively with a cylinderlike shape in 3D Brillouin zone due to their 2D character and they are

clearly not nested to each other as demonstrated by the black dashed circles in Fig. 8(c). If we include the k_z dependence, the size of the δ electron pocket changes from 0.2 \AA^{-1} to 0.36 \AA^{-1} , passing the 0.34 \AA^{-1} of the γ bands. Therefore, the only possible nesting is between γ and δ within a small k_z momentum window. It is hard to argue that this type of nesting

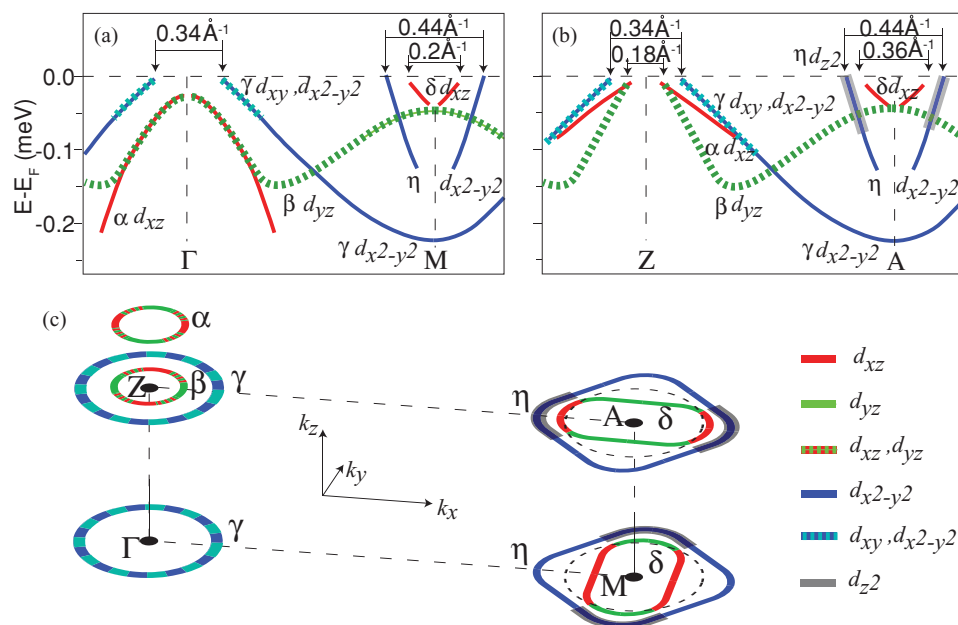


FIG. 8. (Color online) The summary of the orbital characters of low-energy electronic structure. (a),(b) The orbital characters of low-energy electronic structure along Γ -M and Z-A, respectively. The solid and dashed lines represent the bands observed in the p and the s geometries, respectively. (c) The illustration of the orbital characters on the Fermi surface sheets at Γ , M, Z, and A. The almost overlapping α and β Fermi surface sheets at Z are separated for a better illustration. The black dashed circles around M and A are the duplicates of the γ Fermi surface sheets shifted from Γ and Z. The slight mixing between δ and β near the zone corner is not shown though.

can effectively enhance pairing since it takes place at a small region of k_z . In summary, our result does not support that the nesting condition (at least the strict nesting condition) plays a strong role in forming Cooper pairs.

V. SUMMARY

To summarize, we have carried out a systematic investigation on orbital characters of $\text{BaCo}_{0.15}\text{Fe}_{1.85}\text{As}_2$ and analyzed all the possible orbital characters of the multibands based on their strong polarization-dependent photoemission response. As recapitulated in Fig. 8, we found that although the main band structure is qualitatively consistent with the prediction of the present 2D band model, there are important differences between the current theories and our experimental results. The distribution of the d_{xz} , d_{yz} , and d_{z^2} orbital agrees well with the theories, but the bands predicted to be d_{xy} orbital show a strong mixture with $d_{x^2-y^2}$ orbital around the zone center

and almost pure $d_{x^2-y^2}$ orbital around the zone corner. The discrepancy between the calculated band structures and our experimental results suggests that orbital-dependent electron-electron correlations play an important role. We also study the 3D character of the electronic structure and show the strong k_z dependence of the Fermi surfaces and in-plane dispersion of certain bands, which proves that the 3D character is also strongly orbital dependent. Our results lay out a comprehensive picture of the orbital identities of the multiband electronic structures and provide explicit ingredients for constructing the theory of iron-based superconductors.

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