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► To cite this version:

Blaise Faugeras. Tokamak plasma boundary reconstruction using toroidal harmonics and an optimal control method. Fusion Science and Technology, Taylor & Francis, 2016, 69 (2), pp.495-504. hal-01227686

HAL Id: hal-01227686

<https://hal.archives-ouvertes.fr/hal-01227686>

Submitted on 11 Nov 2015

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TOKAMAK PLASMA BOUNDARY RECONSTRUCTION USING TOROIDAL HARMONICS AND AN OPTIMAL CONTROL METHOD

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Abstract

This paper proposes a new fast and stable algorithm for the reconstruction of the plasma boundary from discrete magnetic measurements taken at several locations surrounding the vacuum vessel. The resolution of this inverse problem takes two steps. In the first one we transform the set of measurements into Cauchy conditions on a fixed contour Γ_O close to the measurement points. This is done by least square fitting a truncated series of toroidal harmonic functions to the measurements. The second step consists in solving a Cauchy problem for the elliptic equation satisfied by the flux in the vacuum and for the overdetermined boundary conditions on Γ_O previously obtained with the help of toroidal harmonics. It is reformulated as an optimal control problem on a fixed annular domain of external boundary Γ_O and fictitious inner boundary Γ_I . A regularized Kohn-Vogelius cost function depending on the value of the flux on Γ_I and measuring the discrepancy between the solution to the equation satisfied by the flux obtained using Dirichlet conditions on Γ_O and the one obtained using Neumann conditions is minimized. The method presented here has led to the development of a software, called VacTH-KV, which enables plasma boundary reconstruction in any Tokamak.

Keywords: inverse problem, boundary identification, magnetic measurements

I. INTRODUCTION

In order to control the plasma during a fusion experiment in a Tokamak it is important to be able to compute its boundary in the vacuum vessel. This boundary is deduced from the knowledge of the poloidal flux which itself is computed from a set of discrete magnetic measurements of the poloidal magnetic field and flux scattered around the vacuum vessel.

This paper presents a new fast and stable algorithm for the reconstruction of the poloidal flux in the vacuum surrounding the plasma and of the plasma boundary.

Let us first briefly recall the equations modelizing the equilibrium of a plasma in a Tokamak.¹

Assuming an axisymmetric configuration one considers a 2D poloidal cross section of the vacuum vessel in the (r, z) system of coordinates. Throughout the manuscript bold face letters represent vectors or matrices. In this setting the poloidal flux $\psi(r, z)$ is linked to the magnetic field through the relation $\mathbf{B} = (B_r, B_z) = \frac{1}{r}(-\frac{\partial\psi}{\partial z}, \frac{\partial\psi}{\partial r})$ and, as there is no toroidal current density in the vacuum outside the plasma, satisfies the following linear elliptic partial differential equation

$$L\psi = 0$$

where L denotes the elliptic operator

$$L. = -[\frac{\partial}{\partial r}(\frac{1}{r}\frac{\partial.}{\partial r}) + \frac{\partial}{\partial z}(\frac{1}{r}\frac{\partial.}{\partial z})]$$

The unknown plasma boundary Γ_p is determined from the equation $\psi(r, z) = \psi_p$, ψ_p being the value of the flux at an X-point or the value of the flux for the outermost closed flux line inside the limiter (see Figure 1). The reconstruction of this unknown plasma boundary is a Cauchy problem for the operator L . A supplementary difficulty is that one does not have directly access to a complete set of Cauchy data on an external measurement contour Γ_O but only to discrete magnetic measurements from sensors not necessarily positioned on a single contour.

The proposed resolution method for this inverse problem takes two steps. In the first one we transform the set of discrete measurements into Cauchy conditions, ψ and its normal derivative, on a fixed contour Γ_O chosen close to the measurement points. This is done by least square fitting a truncated series of toroidal harmonic functions to the magnetic measurements. In principle this first step can be sufficient to determine the plasma boundary²⁻⁹ (also see the review article¹⁰). Indeed the toroidal harmonics expansion is valid anywhere in the vacuum surrounding the plasma. The method of toroidal harmonics presented in⁹ proved to be very efficient for the WEST Tokamak configuration.¹¹ The equivalent of magnetic measurements were generated with the equilibrium code CEDRES++.¹² Even when noise is added to the synthetic measurements the boundary reconstructions are accurate with a judicious choice of the order of the toroidal harmonics used.

However since then new numerical experiments conducted for other Tokamaks with more elongated plasma shapes than WEST and with real experimental measurements showed that due to the ill-posedness of the inverse problem, the plasma boundary reconstructed after this first step involving toroidal harmonics only can in some cases be inaccurate. Nevertheless the interpolation of the discrete magnetic measurements on a contour Γ_O chosen

to be close to the measurement points is always very accurate even though the extrapolation of the flux towards the plasma boundary can suffer from some perturbations due to the singularity of internal toroidal harmonics. This is the reason why we propose the following second step which can be seen as a regularization procedure for the toroidal harmonics expansion.

The second step consists in solving a Cauchy problem for the elliptic equation $L\psi = 0$ satisfied by the flux and for the overdetermined boundary conditions on Γ_O obtained with the help of toroidal harmonics as explained above. The proposed method consists in a reformulation of the problem as an optimal control problem on a fixed annular domain of external boundary Γ_O and an fictitious inner boundary Γ_I (see Figure 1) located inside the plasma domain. We aim at minimizing a regularized Kohn-Vogelius¹³ cost function depending on the value of the flux on Γ_I and measuring the discrepancy between the solution to the equation satisfied by ψ obtained using Dirichlet conditions on Γ_O and the one obtained using Neumann conditions.

The analysis of this second step alone was presented by the author in¹⁴ but only tested on a very academic case with synthetic Cauchy data. The algorithm proposed in this paper is a combination of the two previous works⁹ and.¹⁴ It combines the precise data interpolation provided by the toroidal harmonics method to generate Cauchy conditions on Γ_O and the robustness of the optimal control approach to extrapolate these Cauchy conditions towards the plasma and compute the plasma boundary. An implementation for fast numerical computations is also proposed.

The introduction of an optimal control problem on fixed domain appears in^{15,16} where the cost function however is not of the same type as the one proposed in this paper. The advantage of the cost function proposed in this work is that it uses the given Dirichlet and Neumann boundary conditions

on the external contour Γ_O in a symmetric way. For the analysis of section II.B to be valid the control variable has to be the Dirichlet conditions on the internal contour Γ_I .

The use of a fixed inner contour also appears in¹⁷ where a surface current sheet modelizes the plasma and in¹⁸⁻²⁰ in the framework of boundary integral equations which might not be easily applicable in the case of iron core Tokamaks.

The next section presents the two main steps of the proposed method implemented in the code VacTH-KV. Section III then gives details on the numerical methods used and shows some numerical results.

II. OVERVIEW OF THE METHOD

II.A. Step1: interpolation of discrete magnetic measurements with toroidal harmonics

The poloidal flux at any point $\mathbf{x} = (r, z)$ of the vacuum surrounding the plasma is represented by

$$\psi(\mathbf{x}) = \psi_C(\mathbf{x}) + \psi_{TH}(\mathbf{x})$$

The first term ψ_C represents the flux generated by the poloidal field coils. It is computed using Green functions. Each coil is modeled by a number of filaments of current of given intensity.

The second term ψ_{TH} as a solution to the equation $L\psi_{TH} = 0$ can be uniquely decomposed in a series of toroidal harmonics. These toroidal harmonics are found from computations^{8,21} involving a quasi separation of variable technique which lead to define external and internal harmonics. They form a complete set of solutions to this equation in any annular domain.⁸

In particular this implies that this expansion can be used for an iron core Tokamak such as WEST.

The ψ_{TH} term is approximated by the following truncated expansion

$$\left\{ \begin{array}{l} \psi_{TH} = \psi_{ext} + \psi_{int} \\ \psi_{ext} = \frac{r_0 \sinh \zeta}{\sqrt{\cosh \zeta - \cos \eta}} \left[\sum_{n=0}^{n_e} (a_n^e \cos(n\eta) + b_n^e \sin(n\eta)) Q_{n-1/2}^1(\cosh \zeta) \right] \\ \psi_{int} = \frac{r_0 \sinh \zeta}{\sqrt{\cosh \zeta - \cos \eta}} \left[\sum_{n=0}^{n_i} (a_n^i \cos(n\eta) + b_n^i \sin(n\eta)) P_{n-1/2}^1(\cosh \zeta) \right] \end{array} \right.$$

The toroidal harmonic functions are computed in a toroidal coordinate system (ζ, η) which depends on the choice of a pole F_0 . The internal harmonics involve the Legendre functions of the first kind, of degree one and half-integer order, $P_{n-1/2}^1$ which are singular at the pole F_0 . The external harmonics involve the Legendre functions of the second kind, of degree one and half-integer order, $Q_{n-1/2}^1$ which are singular on the axis $r = 0$. If we denote by $\mathbf{u} = (a_0^e, \dots, a_{n_e}^e, b_1^e, \dots, b_{n_e}^e, a_0^i, \dots, a_{n_i}^i, b_1^i, \dots, b_{n_i}^i)$ the vector of the coefficients of the expansion we have an analytic expression for $\psi_{TH}(\mathbf{x}; \mathbf{u})$. Its gradient can be computed analytically as well leading to an expression for the magnetic field in terms of \mathbf{u} , $\mathbf{B}_{TH}(\mathbf{x}; \mathbf{u})$. It remains to identify \mathbf{u} from the set of magnetic measurements available. These discrete magnetic measurements are of three types:

- B probes provide N_B measurements of the poloidal field at points \mathbf{x}_i^B and directions given by unit vectors \mathbf{d}_i such that $B_i^{meas} \approx \mathbf{B}(\mathbf{x}_i^B) \cdot \mathbf{d}_i$ where the dot represents the scalar product
- Flux loops provide N_f flux measurements at points \mathbf{x}_i^f such that $\psi_i^{meas} \approx \psi(\mathbf{x}_i^f)$
- Saddle loops provide N_s flux variations measurements between two

points such that $\delta_i \psi^{meas} \approx \psi(\mathbf{x}_i^1) - \psi(\mathbf{x}_i^2) := \delta\psi(\mathbf{x}_i^1, \mathbf{x}_i^2)$

The unknown vector \mathbf{u} of coefficients is computed minimizing the least square cost function

$$\left\{ \begin{aligned} J(\mathbf{u}) &= \sum_{i=1}^{N_B} \frac{(B_{TH}(\mathbf{x}_i^B; \mathbf{u}) - \tilde{B}_i^{meas})^2}{\sigma_B^2} \\ &+ \sum_{i=1}^{N_f} \frac{(\psi_{TH}(\mathbf{x}_i^f; \mathbf{u}) - \tilde{\psi}_i^{meas})^2}{\sigma_f^2} + \sum_{i=1}^{N_s} \frac{(\delta\psi_{TH}(\mathbf{x}_i^1, \mathbf{x}_i^2; \mathbf{u}) - \delta_i \tilde{\psi}^{meas})^2}{\sigma_s^2} \end{aligned} \right.$$

Where $B_{TH}(\mathbf{x}_i^B; \mathbf{u}) = \mathbf{B}_{TH}(\mathbf{x}_i^B; \mathbf{u}) \cdot \mathbf{d}_i$ and the known contributions of the poloidal field coils are subtracted from the measurements:

$$\left\{ \begin{aligned} \tilde{B}_i^{meas} &= B_i^{meas} - \mathbf{B}_C(\mathbf{x}_i^B) \cdot \mathbf{d}_i, \quad \text{for } i = 1, \dots, N_B \\ \tilde{\psi}_i^{meas} &= \psi_i^{meas} - \psi_C(\mathbf{x}_i^f), \quad \text{for } i = 1, \dots, N_f \\ \delta_i \tilde{\psi}^{meas} &= \delta_i \psi^{meas} - \delta\psi_C(\mathbf{x}_i^1, \mathbf{x}_i^2), \quad \text{for } i = 1, \dots, N_s \end{aligned} \right.$$

The weights σ_B and σ_f, σ_s correspond to the assumed measurement errors. $J(\mathbf{u})$ is quadratic and is minimized by solving the normal equation giving the optimal set of coefficients \mathbf{u}_{opt} .

Once \mathbf{u}_{opt} is computed an approximation of the flux can be obtained at any point of the vacuum surrounding the plasma by

$$\psi(\mathbf{x}; \mathbf{u}_{opt}) = \psi_{TH}(\mathbf{x}; \mathbf{u}_{opt}) + \psi_C(\mathbf{x})$$

In particular one can evaluate ψ and its normal derivative in order to provide Cauchy boundary conditions on any fixed closed contour Γ_O . In this work Γ_O is a contour chosen to be close to the measurements. If all sensors were located a single contour this sensor or measurement contour would be a natural choice for Γ_O .

II.B. Step 2: an optimal control method

From the method described in the previous paragraph one can compute a complete set of Cauchy conditions, $f = \psi$ on Γ_O and $g = \frac{1}{r} \frac{\partial \psi}{\partial n}$ on Γ_O . It is then possible to employ the optimal control method discussed in¹⁴ which is summed up below.

The poloidal flux satisfies

$$\left\{ \begin{array}{l} L\psi = 0 \quad \text{in } \Omega_X \\ \psi = f \quad \text{on } \Gamma_O \\ \frac{1}{r} \frac{\partial \psi}{\partial n} = g \quad \text{on } \Gamma_O \\ \psi = \psi_p \quad \text{on } \Gamma_p \end{array} \right.$$

In this formulation the annular domain $\Omega_X = \Omega_X(\psi)$ contained between Γ_O and Γ_p (see Fig. 1) is unknown since the free plasma boundary Γ_p is unknown. Moreover the problem is ill-posed as there are two Cauchy conditions on the boundary Γ_O .

In order to compute the flux in the vacuum and to find the plasma boundary we define a new problem approximating the original one. Let us define a fictitious boundary Γ_I fixed inside the plasma. We seek an approximation of the poloidal flux ψ satisfying $L\psi = 0$ in the domain Ω contained between the fixed boundaries Γ_O and Γ_I . The problem becomes one formulated on the fixed annular domain Ω . Freezing the domain to Ω by introducing the fictitious boundary Γ_I enables to remove the nonlinearity of the problem. The plasma boundary Γ_p can still be computed as an iso-flux line and thus is an output of our computations.

Let us insist here on the fact that this problem is an approximation to

the original one since in the domain between Γ_p and Γ_I , ψ should satisfy the Grad-Shafranov equation. The relevance of this approximating model is consolidated by the Cauchy-Kowalewska theorem.²² For Γ_p smooth enough the function ψ can be extended in the sense of $L\psi = 0$ in a neighborhood of Γ_p inside the plasma. Hence the problem formulated on a fixed domain with a fictitious boundary Γ_I not "too far" from Γ_p is an approximation of the free boundary problem. If Γ_I were identical with Γ_p then by the virtual shell principle²³ the quantity $w = \frac{1}{r} \frac{\partial \psi}{\partial n} |_{\Gamma_I}$ would represent the surface current density (up to a factor $\frac{1}{\mu_0}$) on Γ_p for which the magnetic field created outside the plasma by the current sheet is identical to the field created by the real current density spread throughout the plasma.

The shape of the internal contour Γ_I is chosen to be a circle or an ellipse. Although no thorough sensitivity analysis has been conducted concerning the shape and size of Γ_I the numerical results presented in the next section did not prove to be very dependent on these factors.

No boundary condition is known on Γ_I and boundary conditions are over determined on Γ_O . One way to deal with this and to solve such a problem is to formulate it as an optimal control one. We are going to compute a function ψ such that the Dirichlet boundary condition $u = \psi$ on Γ_I is such that the Cauchy conditions on Γ_O are satisfied as nearly as possible in the sense of the Kohn-Vogelius error functional defined below which enables to use Dirichlet and Neumann conditions in a symmetric way.

We consider two separate well-posed sub-problems. In the first one we retain the Dirichlet boundary condition on Γ_O only, assume v is given on Γ_I and seek the solution ψ_D of the boundary value problem:

$$\begin{cases} L\psi_D = 0 & \text{in } \Omega \\ \psi_D = f & \text{on } \Gamma_O \\ \psi_D = v & \text{on } \Gamma_I \end{cases}$$

This solution ψ_D can be decomposed in a part linearly depending on v and a part depending on f only. We have the following decomposition:

$$\psi_D = \psi_D(v, f) = \psi_D(v, 0) + \psi_D(0, f) := \psi_D(v) + \tilde{\psi}_D(f)$$

In the second sub-problem we retain the Neumann boundary condition only and look for ψ_N satisfying the boundary value problem:

$$\begin{cases} L\psi_N = 0 & \text{in } \Omega \\ \frac{1}{r} \frac{\partial \psi_N}{\partial n} = g & \text{on } \Gamma_O \\ \psi_N = v & \text{on } \Gamma_I \end{cases}$$

in which ψ_N can be decomposed in a part linearly depending on v and a part depending on g only. We have the following decomposition:

$$\psi_N = \psi_N(v, g) = \psi_N(v, 0) + \psi_N(0, g) := \psi_N(v) + \tilde{\psi}_N(g)$$

For given Cauchy data $f \in H^{1/2}(\Gamma_O)$ and $g \in H^{-1/2}(\Gamma_O)$ we look for $u \in \mathcal{U} = H^{1/2}(\Gamma_I)$ satisfying $J(u) = \inf_{v \in \mathcal{U}} J(v)$ where J is the regularized error functional defined by

$$J(v) = \frac{1}{2} \int_{\Omega} \frac{1}{r} \|\nabla \psi_D(v, f) - \nabla \psi_N(v, g)\|^2 dx + \frac{\varepsilon}{2} \int_{\Omega} \frac{1}{r} \|\nabla \psi_D(v)\|^2 dx$$

In this expression the first term measures a misfit between the Dirichlet solution and the Neumann solution whereas the second one is a regularization term in which ε is a regularization parameter.

Let $u, v \in H^{1/2}(\Gamma_I)$ and define the bilinear symmetric forms

$$s_D(u, v) = \int_{\Omega} \frac{1}{r} \nabla \psi_D(u) \nabla \psi_D(v) dx$$

$$s_N(u, v) = \int_{\Omega} \frac{1}{r} \nabla \psi_N(u) \nabla \psi_N(v) dx$$

and the linear form l

$$l(v) = - \int_{\Omega} \frac{1}{r} (\nabla \tilde{\psi}_D(f) - \nabla \tilde{\psi}_N(g)) \nabla \psi_D(v) dx$$

as well as the quantity c depending on the Cauchy data only

$$c = \frac{1}{2} \int_{\Omega} \frac{1}{r} \|\nabla \tilde{\psi}_D(f) - \nabla \tilde{\psi}_N(g)\|^2 dx$$

Following the analysis provided in¹⁴ the functional J can be rewritten as

$$J(v) = \frac{1}{2} ((1 + \varepsilon) s_D(v, v) - s_N(v, v)) - l(v) + c$$

and its minimum is given by the first order optimality condition or Euler equation which is linear and reads

$$(J'(u), v) = (1 + \varepsilon) s_D(u, v) - s_N(u, v) - l(v) = 0 \quad \forall v \in \mathcal{U}$$

Solving this equation gives the optimal Dirichlet conditions u_{opt} on Γ_I from which one can compute $\psi_D(u_{opt}, f)$ or $\psi_N(u_{opt}, g)$ and then look for the iso-contour defining the plasma boundary.

III. NUMERICAL IMPLEMENTATION AND RESULTS

III.A. The toroidal harmonics expansion

The Legendre functions of degree 1 and half integer order involved in the toroidal harmonic expansion are computed thanks to the Fortran code provided together with reference.²⁴

In the first step of the method one also needs to provide a pole for the toroidal coordinate system and the number of harmonics to be used.

Because the internal toroidal harmonics functions are singular at the pole F_0 of the coordinate system, the expansion procedure provides meaningful results if the pole lies inside the unknown plasma region and not too close from the boundary. An excellent candidate for the choice of this pole is the current center (r_c, z_c) defined from the moments of the plasma current density.^{10,25}

$$\begin{aligned} I_p &:= \int_{D_\Gamma} j_p dx = \int_\Gamma \frac{1}{\mu_0} B_s ds \\ z_c I_p &:= \int_{D_\Gamma} z j_p dx = \int_\Gamma \frac{1}{\mu_0} (-r \log r B_n + z B_s) ds \\ r_c^2 I_p &:= \int_{D_\Gamma} r^2 j_p dx = \int_\Gamma \frac{1}{\mu_0} (2r z B_n + r^2 B_s) ds \end{aligned}$$

where Γ is any closed contour containing the plasma and D_Γ the domain it defines. These quantities can be directly approximated as weighted sums of magnetic measurements and then precisely recomputed as integrals on the contour Γ_0 at every point of which the flux ψ and the field B can be evaluated.

The number of harmonics to be used need not be large to obtain precise fit to the measurements.⁹ Typically the order of the harmonics is chosen

to be 4, 5 or 6. Figure 2 shows a WEST example case of boundary reconstruction with the code using toroidal harmonics only (called VacTH). Two reconstructions are performed using harmonics of order 4 (case 1) and 6 (case 2). In both cases the root mean square errors for the fit to the measurements is of $2.0 \times 10^{-3}[T]$ for B sensors and $1.0 \times 10^{-3}[Wb]$ for flux loops. The Cauchy conditions computed in both cases differ from only small values: $max_{\Gamma_O} |\psi_{case1} - \psi_{case2}| = 2.0 \times 10^{-4}$ and $max_{\Gamma_O} |\frac{1}{r} \frac{\partial \psi}{\partial n}_{case1} - \frac{1}{r} \frac{\partial \psi}{\partial n}_{case2}| = 5.0 \times 10^{-4}$. However the plasma boundary reconstructed with VacTH differs quite a lot from one case to the other. In case 1 it is accurate whereas in case 2 it presents unrealistic oscillations. This shows two things. On the one hand the interpolation of the magnetic measurements using toroidal harmonics and the evaluation of Cauchy data on the contour Γ_O is accurate regardless of the order of the harmonics used. On the other hand the reconstructed plasma boundary is quite dependent on the order of the harmonics used and particularly on the order of the internal harmonics. For all WEST cases that have been treated up today the choice of order 4 harmonics always gave accurate results. However it can be understood from this example that the choice of this order might not be easy or even possible for other Tokamak configurations. Figure 3 shows such an issue. We used VacTH to reconstruct the plasma boundary using ASDEX UpGrade experimental data for shot 25374 from the ITM-WPCD database.²⁶ The same choice for the order of the toroidal harmonics gives a good reconstruction at time 1.1s but this is not the case at time 2.5s. This illustrates the need for the optimal control method proposed in this paper and implemented in a code called VacTH-KV.

III.B. The optimal control problem

The resolution of the optimal control problem in VacTH-KV is based on a classical P^1 finite element method.²⁷ Although over choices would have been possible the advantage of this numerical method over finite differences or boundary integral methods is its simplicity when it comes to adapt the mesh to any internal and external contours, and to extract the plasma boundary from the finite dimensional representation of the flux.

Given the fixed domain Ω let us consider the family of triangulation τ_h of Ω , and V_h the finite dimensional subspace of $H^1(\Omega)$ defined by

$$V_h = \{\psi_h \in H^1(\Omega), \psi_h|_T \in P^1(T), \forall T \in \tau_h\}.$$

Let us also introduce the finite element space on Γ_I

$$D_h = \{v_h = \psi_h|_{\Gamma_I}, \psi_h \in V_h\}.$$

Consider $(\phi_i)_{i=1,\dots,N}$ a basis of V_h and assume that the first N_{Γ_I} mesh nodes (and basis functions) correspond to the ones situated on Γ_I . A function $\psi_h \in V_h$ is decomposed as $\psi_h = \sum_{i=1}^N \psi_i \phi_i$ and its trace on Γ_I as $v_h = \psi_h|_{\Gamma_I} = \sum_{i=1}^{N_{\Gamma_I}} \psi_i \phi_i|_{\Gamma_I}$.

Given boundary conditions v_h on Γ_I and f_h, g_h on Γ_O one can compute the approximations $\psi_{D,h}(v_h)$, $\psi_{N,h}(v_h)$, $\tilde{\psi}_{D,h}(f_h)$ and $\tilde{\psi}_{N,h}(g_h)$ with the finite element method. These computations involve two different linear systems. One for the Dirichlet-Dirichlet problem and one for the Dirichlet-Neumann problem with two symmetric positive definite matrices \mathbf{A}_{DD} and \mathbf{A}_{DN} .

In order to minimize the discrete regularized error functional, $J_{\varepsilon,h}(u_h)$

we have to solve the discrete optimality condition which reads

$$(1 + \varepsilon)s_{D,h}(u_h, v_h) - s_{N,h}(u_h, v_h) - l(v_h) = 0 \quad \forall v_h \in D_h$$

which is equivalent to look for the vector \mathbf{u} solution to the linear system

$$\mathbf{S}\mathbf{u} = \mathbf{l}$$

where the $N_{\Gamma_I} \times N_{\Gamma_I}$ matrix \mathbf{S} represents the bilinear form $s_h = (1 + \varepsilon)s_{D,h} - s_{N,h}$ and is defined by

$$\mathbf{S}_{ij} = s_h(\phi_i, \phi_j)$$

and \mathbf{l} is the vector $(l_h(\phi_i))_{i=1, \dots, N_{\Gamma_I}}$.

The matrices which constitutes \mathbf{S} are evaluated by

$$s_{D,h}(\phi_i, \phi_j) = \int_{\Omega} \frac{1}{r} \nabla \psi_{D,h}(\phi_i) \nabla \mathcal{R}(\phi_j) dx$$

and

$$s_{N,h}(\phi_i, \phi_j) = \int_{\Omega} \frac{1}{r} \nabla \psi_{N,h}(\phi_i) \nabla \mathcal{R}(\phi_j) dx$$

where $\mathcal{R}(\phi_j)$ is the trivial extension which coincides with ϕ_j on Γ_I and vanishes elsewhere.

In the same way the right hand side \mathbf{l} is evaluated by

$$l_h(\phi_i) = - \int_{\Omega} \frac{1}{r} (\nabla \tilde{\psi}_{D,h}(f_h) - \nabla \tilde{\psi}_{N,h}(g_h)) \nabla \mathcal{R}(\phi_i) dx$$

It should be noticed here that matrices \mathbf{A}_{DD} , \mathbf{A}_{DN} and \mathbf{S} depend on the geometry of the problem only. Therefore an internal contour Γ_I being given they can be computed once as well as their Cholesky decomposition and be used for the resolution of successive problems with varying input data. Only

the righthand side \mathbf{I} has to be recomputed.

An important point however is that in the same way that the pole F_0 of the toroidal coordinate system moves following the plasma current center, the inner contour Γ_I might also need to be displaced. We have adopted the following strategy consisting in defining a number of points in the vacuum vessel. Each of them is the center of a circle defining a contour Γ_I . A different mesh is precomputed for each of these contours as well as the associated finite elements matrices (see Figure 4 for an example). Each time a new set of magnetic measurements is provided the plasma current center is computed in the first step of the algorithm involving the toroidal harmonics expansion. Then one looks for the internal contour whose center is the closest to this point and use the associated and precomputed mesh and matrices in the second step in order to minimize the Kohn-Vogelius functional. This enables fast computations. The computing time is of the order of $1ms$ for one boundary reconstruction in the WEST configuration on a Laptop with two quadcore processors of frequency 2.40 Ghz.

Figure 5 shows the reconstructed plasma boundaries for the same WEST example case as in figure 2 but this time with VacTH-KV. Using toroidal harmonics of order either 4 or 6 for the evaluation of Cauchy conditions on Γ_O leads to the same reconstruction almost superimposed with the reference boundary from the equilibrium code CEDRES++. In addition to the unrealistic results given by VacTH for the two AUG cases with real experimental measurements, figure 3 also shows the excellent results provided by VacTH-KV on these same cases. As a last example illustrating the non tokamak dependance of the method figure 6 shows the plasma boundary reconstructed by VacTH-KV on a JET case (shot 74221 at time 46s). In this case VacTH alone was not able to reconstruct a proper plasma boundary.

In all our numerical experiments the regularization parameter ε is given the value of 5×10^{-4} as suggested by the L-curve method analysis²⁸ conducted in.¹⁴

Finally we would like to point out the fact that the optimal control method presented in section II.B. is completely unchanged if poloidal field coils or passive structures with measured currents are included in the domain Ω . In fact this is the case for the WEST, AUG and JET numerical results presented here.

IV. SUMMARY AND CONCLUSION

In this paper a method for plasma boundary reconstruction is proposed. It is decomposed in two steps. A toroidal harmonics expansion is first used to interpolate discrete magnetic measurements on an external contour. This interpolation using functions which are exact analytic solutions of the equation satisfied by the flux is accurate. On the other hand we have shown that the extrapolation of the flux towards the plasma boundary using toroidal harmonics only can in some cases be inaccurate. This is the reason why we introduce a second step in which a Cauchy problem on a fixed domain is solved by an optimal control method. These two steps together provide an accurate and robust plasma boundary reconstruction method. It is generic and can be used for any Tokamak having ferromagnetic structures or not. An implementation using a finite elements discretization is proposed. The precomputation of several meshes and all associated linear algebra quantities allows fast computation. A code called VacTH-KV has been developed and is available on the ITM-WPCD platform. This paper does not deal with the computation of error bars on the plasma boundary reconstruction results. This is of course of interest as for any ill-posed inverse problem and might

be conducted in the future for example using the epsilon-nets technique²⁹

ACKNOWLEDGMENTS

This work, supported by the European Communities under the contract of Association between EURATOM-CEA was carried out within the framework of the Task Force on Integrated Tokamak Modelling of the European Fusion Development Agreement (ITM-WPCD). The views and opinions expressed herein do not necessarily reflect those of the European Commission.

The author would like to thank his colleagues at the university of Nice, Jacques Blum, Cedric Boulbe and Holger Heumann for many useful discussions. The input magnetic measurements for WEST were provided by colleagues from the IRFM-CEA at Cadarache, Eric Nardon and Remy Nouailletas, and for AUG and JET by the ITM-WPCD project.

Finally the author would like to thank the two anonymous reviewers for their constructive criticism from which the paper has benefitted significantly.

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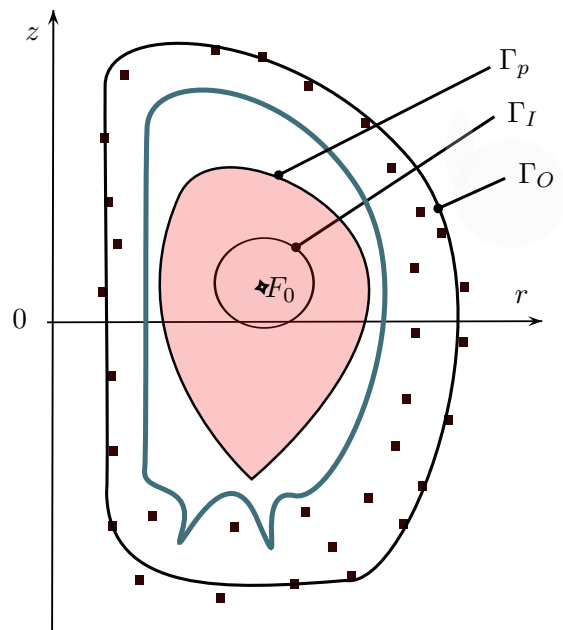


Figure 1: Schematic representation of the poloidal cross section of a Tokamak. The unknown plasma domain of boundary Γ_p and the limiter contour are shown. Magnetic sensors are depicted by squares. The external contour Γ_O , the inner contour Γ_I as well as the current center F_0 are represented.

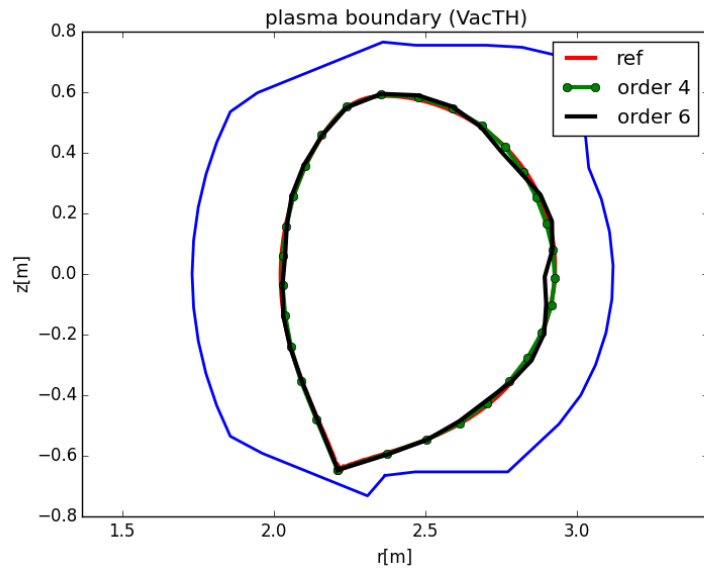


Figure 2: A WEST equilibrium case. The reference plasma boundary from the equilibrium code CEDRES++ and the boundary reconstructed using harmonics of order 4 in VacTH (bullets) are almost superimposed. The boundary reconstructed using harmonics of order 6 presents some unrealistic oscillations.

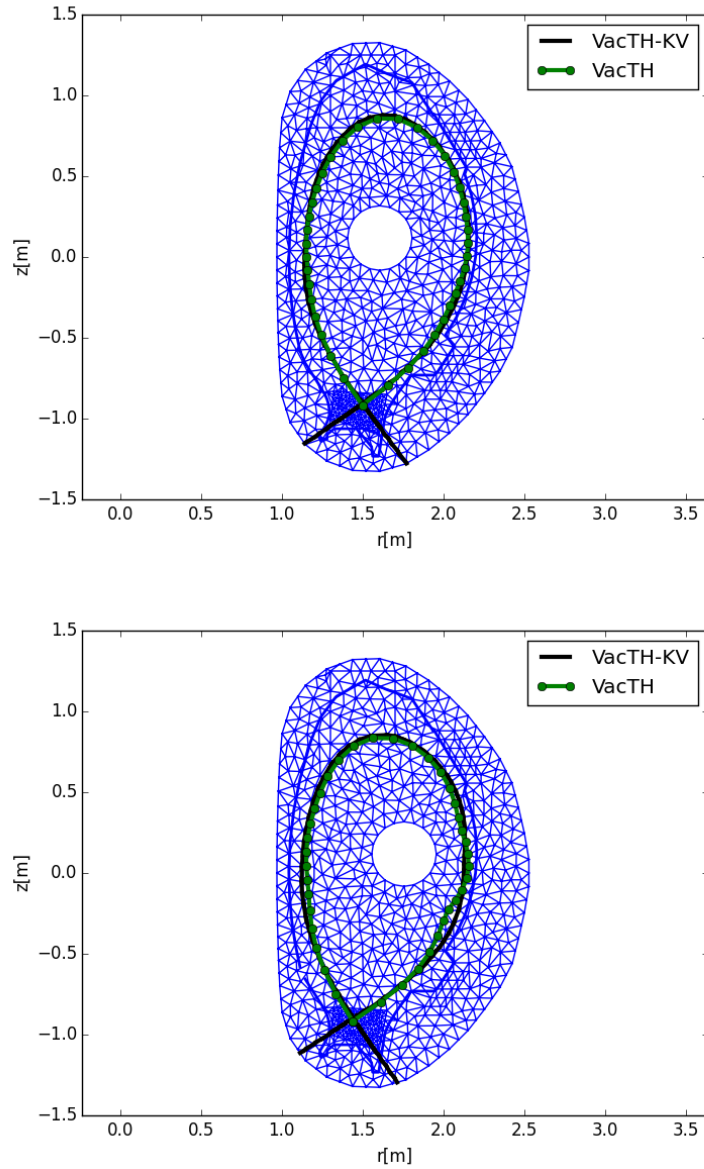


Figure 3: AUG shot 25374 plasma boundary reconstructions at time 1.1s (top) and time 2.5 (bottom). The plasma boundary computed from toroidal harmonics of order 4 in VacTH (bullets) is satisfactory at 1.1s but presents unrealistic oscillations at 2.5s. The plasma boundaries computed with the new proposed method VacTH-KV are shown with continuous lines.

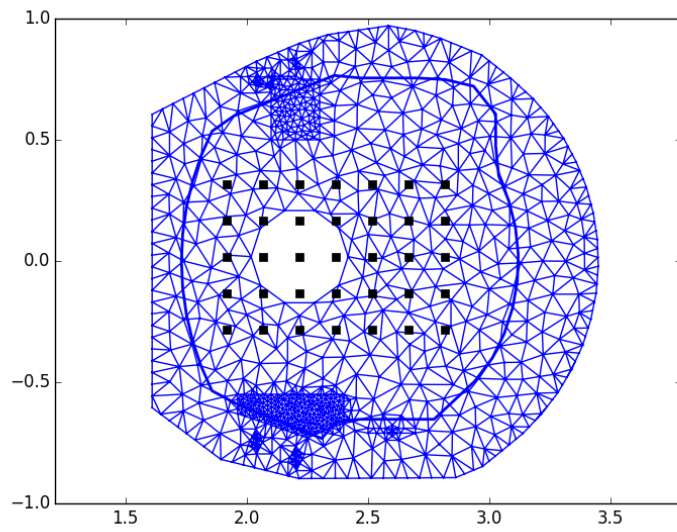
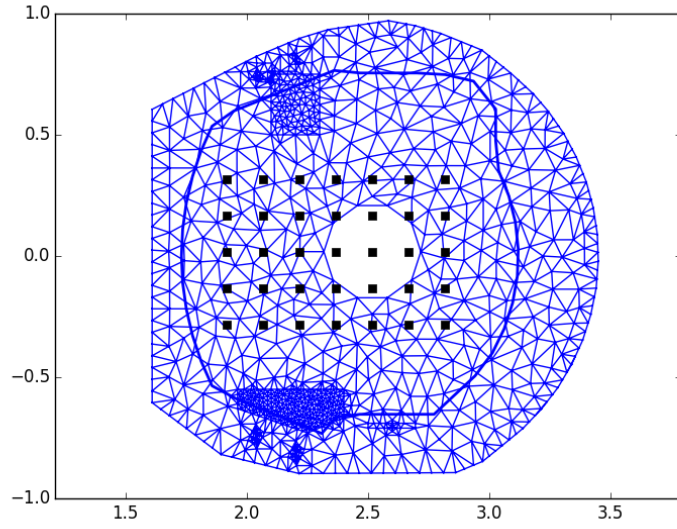


Figure 4: Examples of precomputed meshes for WEST. The limiter is represented. The mesh is refined in lower and upper X-point regions. Each small square inside the limiter region is the center of a circle defining an internal contour Γ_I .

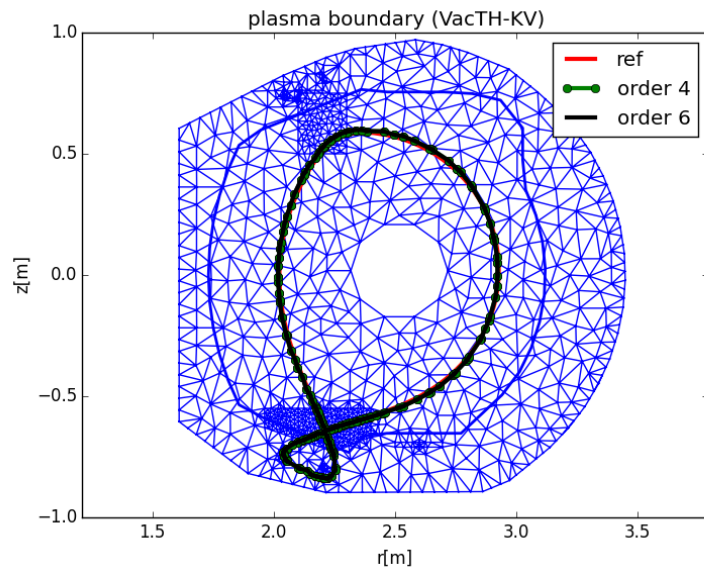


Figure 5: Same WEST test case as in Figure 2. The boundaries reconstructed with the new optimal control proposed method, VacTH-KV, using toroidal harmonics of order either 4 or 6 are almost superimposed with the reference boundary from CEDRES++

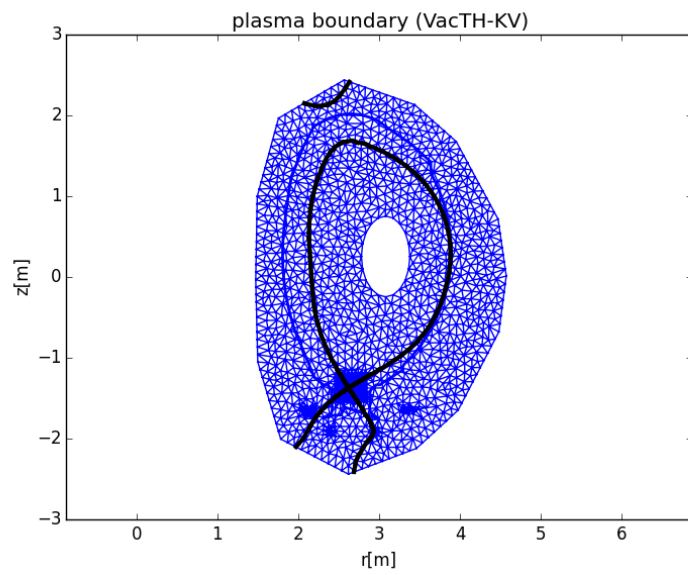


Figure 6: VacTH-KV boundary reconstruction for a JET test case example (shot 74221 at 46s). The method using toroidal harmonics only, VacTH, could not reconstruct a proper boundary.