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II. Nitrogen-bearing species

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ABSTRACT

Context. The chemical inventory of planets is determined by the physical and chemical processes that govern the early phases of star formation. Nitrogen-bearing species are of interest as many provide crucial precursors in the formation of life-related matter.
Aims. The aim is to investigate nitrogen-bearing complex organic molecules towards two deeply embedded Class 0 low-mass protostars (Perseus B1-c and Serpens S68N) at millimetre wavelengths with the Atacama Large Millimeter/submillimeter Array (ALMA). Next, the results of the detected nitrogen-bearing species are compared with those of oxygen-bearing species for the same and other sources. The similarities and differences are used as further input to investigate the underlying formation pathways.
Methods. ALMA observations of B1-c and S68N in Band 6 (~ 1 mm) and Band 5 (~ 2 mm) are studied at $\sim 0.5''$ resolution, complemented by Band 3 (~ 3 mm) data in a $\sim 2.5''$ beam. The spectra are analysed for nitrogen-bearing species using the CASSIS spectral analysis tool, and the column densities and excitation temperatures are determined. A toy model is developed to investigate the effect of source structure on the molecular emission.
Results. Formamide (NH_2CHO), ethyl cyanide (C_2H_5CN), isocyanic acid ($HNCO$, $HN^{13}CO$, $DNCO$), and methyl cyanide (CH_3CN, CH_2DCN, and CHD_2CN) are identified towards the investigated sources. Their abundances relative to CH_3OH and $HNCO$ are similar for the two sources, with column densities that are typically an order of magnitude lower than those of oxygen-bearing species. The largest variations, of an order of magnitude, are seen for NH_2CHO abundance ratios with respect to $HNCO$ and CH_3OH and do not correlate with the protostellar luminosity. In addition, within uncertainties, the nitrogen-bearing species have similar excitation temperatures to those of oxygen-bearing species ($\sim 100-300$ K). The measured excitation temperatures are larger than the sublimation temperatures for the respective species.
Conclusions. The similarity of most abundances with respect to $HNCO$ for the investigated sources, including those of CH_2DCN and CHD_2CN, hints at a shared chemical history, especially the high D-to-H ratio in cold regions prior to star formation. However, some of the variations in abundances may reflect the sensitivity of the chemistry to local conditions such as temperature (e.g. NH_2CHO), while others may arise from differences in the emitting areas of the molecules linked to their different binding energies in the ice. The excitation temperatures likely reflect the mass-weighted kinetic temperature of a gas that follows a power law structure. The two sources discussed in this work add to the small number of sources that have been subjected to such a detailed chemical analysis on Solar System scales. Future data from the *James Webb Space Telescope* will allow a direct comparison between the ice and gas abundances of both smaller and larger nitrogen-bearing species.

Key words. Astrochemistry – Stars: low-mass – Stars: protostars – ISM: abundances – Instrumentation: interferometers

1. Introduction

Interstellar molecules with six or more atoms containing carbon atoms as well as hydrogen and oxygen or nitrogen are known as O- or N-bearing complex organic molecules (COMs). Other COMs comprise both oxygen and nitrogen atoms and/or other elements such as sulphur and phosphorus. The Class 0 protostellar stage is the warmest stage during star formation and thus the richest in gas-phase COMs due to the thermal sublimation of ices in hot corinos (van 't Hoff et al. 2020b). Class 0 sources are thus prime targets to unveil their associated chemistry through

the observation of (sub-)millimetre lines. Over the past decades, single dish and interferometric studies on cloud scales have indeed revealed that both low-mass (e.g. van Dishoeck et al. 1995; Cazaux et al. 2003; Bottinelli et al. 2004a; Bottinelli et al. 2004b; Bisschop et al. 2008; Jørgensen et al. 2016; Ceccarelli et al. 2017; Bergner et al. 2017; Bianchi et al. 2019) and high-mass protostars (Blake et al. 1987; Gibb et al. 2000; Nummelin et al. 2000; Fontani et al. 2007; Belloche et al. 2013; Ilee et al. 2016; Bøgelund et al. 2019a; Taniguchi et al. 2020) are rich in COMs. It has also recently been found that planet formation likely starts at earlier stages of low-mass star formation, during the Class 0/I stage (Harsono et al. 2018; Manara et al. 2018; Tychoniec et al. 2018; Tychoniec et al. 2020; Tobin et al. 2020). Therefore, to understand the chemical enrichment of forming planets in complex compounds, one needs to study the formation of such species at these early phases on Solar System scales ($\sim$50 au).

Complex organic molecules start forming at the dawn of star formation, in molecular clouds of gas and dust with temperatures of $\sim$10 − 20 K and initial densities of $\sim$$10^3$ − 10^4 cm^{-3}. Under these conditions, first simple molecules, such as water (H_2O), carbon dioxide (CO_2), and ammonia (NH_3), form on icy grains (Boogert et al. 2015; Linnartz et al. 2015). Once CO freezes out, O-bearing COMs start forming through the hydrogenation of CO, resulting in the formation of formaldehyde (H_2CO) and methanol (CH_3OH; Watanabe & Kouchi 2002; Fuchs et al. 2009). Species as complex as glycolaldehyde ($HC(O)CH_2OH$) and ethylene glycol ($H_2C(OH)CH_2OH$) can form through the recombination of HCO radicals at low temperatures (Fedoseev et al. 2015). It is expected that other radical recombinations result in the formation of even larger O-containing COMs, such as glycerol ($HOCH_2CH(OH)CH_2OH$; Fedoseev et al. 2017). As the temperature increases, more COMs, such as ethanol (CH_3CH_2OH; Öberg et al. 2009a) and propanal (CH_3CH_2CHO; Qasim et al. 2019), can form in the ice, perhaps with the assistance of some UV radiation.

Much less is known about the formation of N-bearing species in ices. Nitrogen-bearing molecules are important as nitrogen is a crucial element in developing biomolecules such as amino acids and nucleobases, and they are thus essential for the emergence of life. N-bearing COMs have been detected in high-mass (e.g. Isokoski et al. 2013; Belloche et al. 2016; Bøgelund et al. 2019b; Csengeri et al. 2019; Ligterink et al. 2020) and low-mass protostars (e.g. Bottinelli et al. 2008; Ligterink et al. 2018; Marcelino et al. 2018; Calcutt et al. 2018; Lee et al. 2019; Belloche et al. 2020). Some are thought to form through the recombination of radicals in the solid state, such as methylamine (CH_3NH_2) from NH_2 and CH_3 produced by the photodissociation of NH_3 and CH_4 ice, although these radicals can also result from hydrogen addition reactions to N and C atoms (Garrod et al. 2008). Other species can form via isocyanic acid (HNCO) in the solid state (e.g. methyl isocyanate, CH_3NCO; Cernicharo et al. 2016; Ligterink et al. 2017). Formamide (NH_2CHO) is thought to have gas-phase formation routes (Barone et al. 2015; Song & Kästner 2016; Codella et al. 2017; Skouteris et al. 2017) as well as two main ice pathways, namely the hydrogenation of HNCO (e.g. Raunier et al. 2004; Haupa et al. 2019) and the radical-radical addition of NH_2 and CHO (Jones et al. 2011, Dulieu et al. 2019; Martín-Doménech et al. 2020). HNCO itself can also form on the surface of interstellar ices via the reaction of CO and NH radicals (Fedoseev et al. 2015). Another molecule put forward as a parent molecule for many N-bearing COMs is methyl cyanide (CH_3CN; Bulak et al. 2021), which may be produced via reactions of CH_3 and CN radicals in ices. One route for the formation of ethyl cyanide (C_2H_5CN) is thought to be a piecewise addition

of its functional groups on dust grains (Belloche et al. 2009). In order to elucidate the formation pathways to N-bearing COMs, more observational constraints on their abundances are needed.

Another clue regarding formation routes comes from observed deuteration fractions. The low temperature in the dense core stage causes the D/H ratio in molecules to increase to higher values than the elemental ratio in the interstellar medium (ISM) of $\sim$$2 \times 10^{-5}$ (Prodanović et al. 2010). The chemical reaction $H_3^+ + HD \rightleftharpoons H_2D^+ + H_2 + \Delta E$ increases the amount of H_2D^+ in the gas phase at low temperatures ($\sim$10 − 20 K) since this reaction has a small activation barrier for its reverse reaction (Watson 1976; Aikawa & Herbst 1999; Tielens 2013; Ceccarelli et al. 2014; Sipilä et al. 2015). Moreover, this process can be enhanced by CO freeze-out as CO is one of the main molecules to destroy H_3^+ and H_2D^+ in the gas phase (Brown & Millar 1989; Roberts et al. 2003). Dissociative recombination of H_2D^+ will enhance the D/H ratio, and, subsequently, D atoms can be transferred onto grains, enriching the ice. Therefore, D/H values provide clues on the temperature at which N-bearing COMs are formed, and thus on their formation history (Taquet et al. 2012, 2014; Furuya et al. 2016).

To assess whether many of the N- and O-bearing COMs are indeed produced in ices, a smoking gun would be to detect their ice features directly at mid-infrared wavelengths. The total number of unambiguously identified ice species is, however, relatively small and so far only comprises one securely identified COM (CH_3OH; Grim et al. 1991; Taban et al. 2003). This is largely a consequence of observational limitations (see the review by Boogert et al. 2015); even with more laboratory ice data for COMs currently available (see e.g. Terwisscha van Scheltinga et al. 2018), only the more abundant molecules in ices can be detected. There are upper limit measurements of some N-bearing COMs in ices towards massive protostars, for example aminomethanol (NH_2CH_2OH; Bossa et al. 2009) and NH_2CHO (Schutte et al. 1999). However, OCN^-, a direct derivative of HNCO (van Broekhuizen et al. 2004; Fedoseev et al. 2016), has been detected in the ISM in ices (Grim & Greenberg 1987; van Broekhuizen et al. 2005; Öberg et al. 2011). Observations of most molecules in the solid state (at near- and mid-infrared) are not possible from the ground because a large part of the wavelength range is blocked by the Earth's atmosphere. Moreover, moderate spectral resolution is needed in the critical 3 − 10 μm wavelength range for observation of most molecules in ices, which the *Spitzer Space Telescope* did not have. The *James Webb Space Telescope* (JWST), with its unique sensitivity and appropriate spectral resolution, will transform the study of COMs in the solid state.

The present work focuses on gas-phase sub-millimetre identifications of N-bearing species and has only become possible because of the superb performance of the Atacama Large Millimeter/submillimeter Array (ALMA; Jørgensen et al. 2020). This is because ALMA has a much higher sensitivity and spatial resolution than the pre-existing telescopes, which enables the study of low-mass protostars on Solar System scales. Moreover, the high sensitivity of ALMA allows the observation of molecule isotopologues. This is especially important for highly abundant molecules that show optically thick emission as their abundances can be measured more accurately using their optically thin isotopologues. ALMA observations provide information on the gas-phase chemical inventory in the hot core, where all ices have sublimated, and these data can then eventually be compared with JWST observations of the ice composition to directly link gas and ice chemistry.

One of the most well-studied low-mass protostars in complex chemistry is IRAS 16293-2422 (hereafter IRAS 16293), which was investigated as part of the ALMA Protostellar Interferometric Line Survey (PILS) programme (Jørgensen et al. 2016). The PILS gives the most complete inventory of nitrogen- and oxygen-bearing COMs in low-mass protostars to date (Coutens et al. 2016; Ligterink et al. 2017; Ligterink et al. 2018; Calcutt et al. 2018; Manigand et al. 2020). Jørgensen et al. (2018) and Manigand et al. (2020) suggest that there are two categories for O-bearing and N-bearing species: Some molecules desorb at temperatures of ~100 K and others at ~300 K, closer to the central protostar. Apart from IRAS 16293, N-bearing COMs have been observed with ALMA towards a handful of low-mass sources: HH 212 (Lee et al. 2019), NGC 1333 IRAS 4A2 (López-Sepulcre et al. 2017), B1b-S (Marcelino et al. 2018), B335 (Imai et al. 2016), and L483 (Oya et al. 2017). This list has also been supplemented using the Plateau de Bure Interferometer (PdBI) and its upgraded version, the Northern Extended Millimeter Array (NOEMA; Taquet et al. 2015; Belloche et al. 2020).

In this paper, ALMA observations are used to study two Class 0 objects: B1-c in the Perseus Barnard 1 cloud and S68N in the Serpens Main star-forming region. These sources are targeted in ALMA Band 6, Band 5, and Band 3. The luminosity of B1-c at its distance of 321 pc is $6.0\,L_\odot$ (Karska et al. 2018; Ortiz-León et al. 2018). S68N has a luminosity of $5.4\,L_\odot$ (Enoch et al. 2011) at its distance of 436 pc (Ortiz-León et al. 2017). Both S68N and B1-c have been observed and studied with ALMA. Very recently, van Gelder et al. (2020) presented observational data for O-bearing COMs towards both sources. In this paper we focus on N-bearing molecules to investigate how similarities and differences between N-bearing and O-bearing species reveal information on the involved chemical processes. Both sources are also targets of the guaranteed time observation (GTO) programme (project ID 1290) of JWST/MIRI (Wright et al. 2015). Therefore, in the near future, it will be possible to directly compare the ice observations of these sources obtained by JWST with what has been done in this work for their gas-phase counterparts in the hot corino, where these ices have sublimated.

The layout of this paper is as follows. Section 2.1 describes the observations. Section 3 presents the methods and the results. In Sect. 4 we discuss our results and put the sources studied here in the context of what has been done so far in the literature. We also compare the measured excitation temperature with the sublimation temperature of each molecule. Moreover, a simple toy model is constructed to understand how source structure may affect abundance ratios. Finally, a summary is given in Sect. 5.

2. Observations and methods

2.1. The data

Two COM-rich protostars (B1-c and S68N) were observed by ALMA (project code: 2017.1.01174.S; principal investigator: E.F. van Dishoeck). The data reduction and first results from these observations are explained in van Gelder et al. (2020). Here we only give a brief overview of the data. B1-c (RA_{J2000}: 03:33:17.88, Dec_{J2000}: 31:09:31.8) and S68N (RA_{J2000}: 18:29:48.08, Dec_{J2000}: 01:16:43.3) were observed during ALMA Cycle 5 at 3 mm (Band 3) and 1 mm (Band 6) using the 12m array. Two bands were used here to be sensitive to both more extended (Band 3, maximum baseline of ~400 m) and hence colder COM emission and the more compact (Band 6, maximum baseline of ~800 m) and thus warmer COM emis-

sion. The targeted N-bearing molecules were originally HNCO and NH_2CHO, species with a likely solid-state formation origin (Fedoseev et al. 2015, 2016). The other N-bearing (complex organic) molecules discussed in this work were serendipitously observed, and hence this work does not aim at a complete inventory of N-bearing molecules. The observational parameters and the lines covered in the data are presented in Appendix D.

The Band 3 data were taken using ALMA configurations C43-2 (S68N) and C43-3 (B1-c) with an angular resolution of ~1.5 − 2.5″. In Band 6 the C43-4 configuration was used with an angular resolution of ~0.45″, corresponding to radii of 72 au and 98 au for B1-c and S68N, respectively. The spectral resolution for most spectral windows is ~0.2 km s⁻¹. A few spectral windows in Band 3 have a spectral resolution of ~0.3 − 0.4 km s⁻¹. The maximum recoverable scales for Band 3 and Band 6 are ~20″ and ~6″, respectively. The line rms is ~0.15 K in the Band 6 data. The absolute flux calibration uncertainty is ≤ 15 %.

Additionally, Band 5 data (project code: 2017.1.01371.S; principal investigator: M.L.R; van 't Hoff) are included in the analysis for B1-c to confirm identifications and get more accurate measurements of the excitation temperatures. The data were reduced using the ALMA pipeline (CASA version 5.1.1), after which line-free regions were carefully selected for continuum subtraction. The data were then imaged using a robust weighting of 0.5. This dataset has a similar angular resolution (~0.45″) to our Band 6 dataset, with a maximum baseline of ~1.3 km. The spectral resolution of this dataset is mostly ~0.1 km s⁻¹, but for some of the data cubes it is ~1.6 km s⁻¹. The line rms ranges from ~0.1 K for spectral windows with ~1.6 km s⁻¹ spectral resolution to ~0.4 K for spectral windows with ~0.1 km s⁻¹ spectral resolution in the Band 5 data. The covered frequency ranges for Bands 5 and 6 are given in Table D.2.

2.2. Spectral modelling

A molecule is considered to be detected when at least three lines are identified at a 3σ level without over-predicting any line emission. It is called 'tentatively detected' when it has one or two lines at a 3σ level in the spectrum. An upper limit is reported when no lines are identified.

We followed the approach in van Gelder et al. (2020) to determine the column densities (N) and excitation temperatures (T_{ex}) of molecules identified in the spectra. First, a grid of N and T_{ex} for each molecule was set. Although the full width half maximums (FWHMs) are fixed for the final fits, they were varied first; however, it was found that either they do not vary significantly for unblended lines or they are non-constrained due to line blending. Therefore, they were fixed to the best-fit value found for clean, single lines. Assuming a single component origin, the simplest assumption is that a single excitation temperature describes the level populations of a molecule. This condition is referred to as 'local thermodynamic equilibrium' (LTE) when the densities are high enough that the excitation temperature approaches the kinetic temperature and one can use the Boltzmann distribution to describe the population of all levels at a single temperature. Jørgensen et al. (2016) found that the assumption of LTE conditions is reasonable on scales of 100 au for low-mass protostars such as IRAS 16293, where densities are $\sim 10^8 - 10^9$ cm⁻³ or higher. Assuming LTE conditions, the corresponding spectrum for each grid point is calculated using the CASSIS[1] (Vastel et al. 2015) spectral analysis tool. For each molecule, we used its corresponding line list from the Jet Propulsion Laboratory

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[1] http://cassis.irap.omp.eu/

(JPL) database (Pickett et al. 1998) and the Cologne Database for Molecular Spectroscopy (CDMS; Müller et al. 2001; Müller et al. 2005). Subsequently, the resulting modelled spectrum of each grid point was overlaid onto the observed spectrum of each source and its χ^2 was computed. In the computation of the best-fit model, we did not include blended or optically thick lines. In fitting all the molecules in the Band 5 and 6 data, we used 20% for the flux uncertainty. The uncertainty used here is a conservative estimate to take into account potential errors caused by continuum contamination due to the line richness of the sources.

We investigated a grid with a large range of column densities, from 10^{12} cm^{-2} to 10^{16} cm^{-2}, with large spacings (for some molecules, a larger initial range was used for the grid). This was to ensure that the full parameter space was covered. Once the range of the final column density was found, a finer grid, with 0.05 spacings in logarithmic scale for the column density, was made. The excitation temperature in our grids mostly ranges from 10 K to 600 K with 10 K spacings on a linear scale. For NH$_2$CHO, C$_2$H$_5$CN, and CHD$_2$CN towards B1-c, the maximum of the temperature grid was larger to guarantee that the resulting excitation temperature is not biased. We only fitted the temperature when there were several lines that covered a range of upper energy levels. Otherwise, we fixed the temperature to 200 K for the Band 5 and 6 data. This is because van Gelder et al. (2020) found 200 K to be a typical excitation temperature for O-bearing species in Band 6. Where it is possible to fit for the temperature, the results from the fits were inspected more closely via the χ^2 plots and by manually changing the temperature for a small range of column densities to make sure that the most accurate excitation temperature is found (see Appendix C). Typically, the FWHM was fixed to 3 km s^{-1} for B1-c and S68N unless broader lines were observed for a certain molecule. In that case, the FWHM was fixed to 4.5 km s^{-1} (HN^{13}CO, C$_2$H$_5$CN, and CH$_2$DCN in S68N).

In this procedure we fixed the source velocities (V_{lsr}) to 6.0 km s^{-1} and 8.5 km s^{-1} for B1-c and S68N, respectively (van Gelder et al. 2020). Furthermore, we followed the same method as in van Gelder et al. (2020) and assumed that the source size for both sources is the same as the beam size for the Band 6 data (0.45″) since the compact emission from the inner envelope is unresolved (see Sect. 3.1). The beam dilution factor is given by $(\theta_b^2 + \theta_s^2)/\theta_s^2$, where θ_b is the beam size and θ_s is the source size. Hence, assuming a source size of 0.45″, the beam dilution factors are ~20 and ~2 in the Band 3 data and the Band 5 and 6 data, respectively. The source size affects the measured column densities but not their ratios, unless the lines become optically thick. This is discussed further in Sect. 3.5.

We derived the uncertainties on the column densities using the χ^2 error calculation of the grid (2σ) or the variation in the column densities when manually fitting for temperature. The 2σ uncertainty on the temperatures using the χ^2 error calculation of the grid is reported when the temperatures derived from both the χ^2 calculations and the inspection by eye are consistent. However, in cases where the χ^2 method is not constraining (e.g. because there are not enough lines covering a large range in upper energy levels), the temperature is reported based on the fit-by-eye method, where we typically find a ~50 K or ~100 K error (see Appendix C for a description of the method).

3. Results

3.1. Line identification and spatial extent

We find four N-bearing molecules and some of their isotopologues towards both sources. Towards B1-c, we detect HNCO, NH$_2$CHO, C$_2$H$_5$CN, CH$_2$DCN, and CHD$_2$CN and tentatively detect HN^{13}CO and DNCO in the Band 6 and 5 data. Moreover, CH$_3$CN is clearly detected in the Band 3 data. Towards S68N, we detect HNCO, C$_2$H$_5$CN, and CH$_2$DCN and tentatively detect NH$_2$CHO and HN^{13}CO with upper limits for DNCO and CHD$_2$CN in the Band 6 data. In addition, CH$_3$CN is identified in the Band 3 data. Among the molecules searched for, we also find upper limits on NH$_2$CN, CH$_3$NH$_2$, CH$_3$NCO, and HOCH$_2$CN for both sources in the Band 5 and 6 data. Given that in the Band 3 data only CH$_3$CN is securely detected, with a tentative detection of HNCO, we focused on the Band 6 data for S68N and the combined Band 5 and 6 data for B1-c for most of the analyses.

A summary of the identified molecules along with the fitted parameters of the models are presented in Table 1 for B1-c and S68N. The temperature was fitted for 50% of the molecules in B1-c and 20% of the species in S68N. The FWHMs found for the N-bearing species are similar to those found by van Gelder et al. (2020) for O-bearing species. The fits to the spectra for NH$_2$CHO and HN^{13}CO towards B1-c are presented in Figs. 1 and 2, respectively, as examples. The fits to the rest of the data are shown in Appendix C.

Figure 3 shows moment zero maps of the SO $6_7 - 5_6$ ($E_{up} = 47.6$ K), HNCO $12_{0,12} - 11_{0,11}$ ($E_{up} = 82.3$ K), and NH$_2$CHO $13_{1,13} - 12_{1,12}$ ($E_{up} = 91.8$ K) lines for B1-c and S68N in Band 6. Both B1-c and S68N are known to have outflows (Jørgensen et al. 2006; Tychoniec et al. 2019). This is seen from the SO moment zero map, where the spatially extended emission traces the outflow (see e.g. Podio et al. 2015). Some more complex molecules, such as NH$_2$CHO and HNCO, can also be detected in the outflow in their low upper energy level lines in the Band 3 data (Tychoniec et al. in prep.), but the analysis in this paper focuses on the compact emission from the inner envelope, and mostly on the Band 5 and 6 data. Figure 3 shows that the compact emission from NH$_2$CHO and HNCO is not spatially resolved and that these molecules show emission that does not extend beyond the continuum. These must be located within ~200 au of the central protostar. The width of the lines of N-bearing species (see Appendix C for the spectra), ~3 − 5 km s^{-1}, is in line with a rotating structure or turbulence in the inner envelope; therefore, the molecules studied in this work likely trace the inner envelope and potential warm disk around the protostar.

3.2. Column densities

A summary of the derived column densities for B1-c and S68N is presented in Table 1. The most abundant N-bearing molecule detected towards B1-c and S68N in the Band 5 and 6 data is HNCO. The emission of this molecule is optically thick in both sources. This is also seen from a comparison of the HN^{13}CO column density with the fitted HNCO column density, where the latter is under-estimated. The column densities of HNCO measured by directly fitting the spectra for HNCO are ~1.8 × 10^{15} cm^{-2} and ~10^{15} cm^{-2} towards B1-c and S68N, respectively. Therefore, the interstellar ratio of ^{12}C/^{13}C ~ 68 (Milam et al. 2005) was used to find HNCO column densities of 1.6 ± 0.7 × 10^{16} cm^{-2} and ~1.2 × 10^{16} cm^{-2} for B1-c and S68N from the HN^{13}CO column densities, assuming the HN^{13}CO lines are optically thin. This assumption can be investigated by searching for HNC^{18}O lines.

Table 1. Column densities and excitation temperatures for B1-c and S68N in a 0.45″ beam.

| Species	Catalogue	B1-c				S68N			
		T_{ex}(K)	N(cm^{-2})	N/CH$_3$OH(%)	N/HNCO(%)	T_{ex}(K)	N(cm^{-2})	N/CH$_3$OH(%)	N/HNCO(%)
HNCO	CDMS	–	$(1.6 \pm 0.7) \times 10^{16}$	0.9 ± 0.5	$\equiv 100.0$	–	$\sim 1.2 \times 10^{16}$	0.9 ± 0.4	$\equiv 100.0$
HN^{13}CO	JPL	$200 \pm 100^*$	$(2.4 \pm 1.1) \times 10^{14}$	0.013 ± 0.007	$\equiv 1.47$	$[200]^*$	$\sim 1.8 \times 10^{14}$	0.013 ± 0.005	$\equiv 1.47$
DNCO	JPL	$[200]^*$	$(3.0 \pm 0.2) \times 10^{14}$	0.016 ± 0.005	1.8 ± 0.8	$[200]$	$< 6.8 \times 10^{13}$	< 0.005	< 0.6
NH$_2$CHO	JPL	> 70	$(5.1 \pm 2.4) \times 10^{14}$	0.03 ± 0.02	3.1 ± 2.1	$[200]^*$	$(2.7 \pm 0.3) \times 10^{14}$	0.019 ± 0.009	2.2 ± 0.3
C$_2$H$_5$CN	CDMS	275 ± 125	$(1.1 \pm 0.3) \times 10^{15}$	0.06 ± 0.02	6.5 ± 3.5	170^{+100}_{-110}	$(1.1 \pm 0.6) \times 10^{15}$	0.08 ± 0.05	8.9 ± 4.7
CH$_3$CN	CDMS	–	$(7.4 \pm 2.2) \times 10^{15}$	0.4 ± 0.2	45.5 ± 24.9	–	$(2.9 \pm 1.0) \times 10^{15}$	0.2 ± 0.1	23.6 ± 8.3
CH$_2$DCN	CDMS	210^{+50}_{-40}	$(2.6 \pm 0.2) \times 10^{14}$	0.014 ± 0.004	1.6 ± 0.7	220^{+60}_{-110}	$(1.0 \pm 0.2) \times 10^{14}$	0.007 ± 0.003	0.8 ± 0.2
CHD$_2$CN	CDMS	> 200	$(1.4 \pm 0.2) \times 10^{14}$	0.007 ± 0.003	0.9 ± 0.4	$[200]$	$< 3.8 \times 10^{13}$	< 0.003	< 0.3
NH$_2$CN	JPL	$[200]$	$< 7.1 \times 10^{13}$	< 0.004	< 0.4	$[200]$	$< 1.0 \times 10^{13}$	< 0.0007	< 0.1
CH$_3$NCO	CDMS	$[200]$	$< 8.8 \times 10^{14}$	< 0.05	< 5.4	$[200]$	$< 3.8 \times 10^{14}$	< 0.03	< 3.1
CH$_3$NH$_2$	JPL	$[200]$	$< 5.1 \times 10^{15}$	< 0.3	< 31.5	$[200]$	$< 1.9 \times 10^{15}$	< 0.1	< 15.6
HOCH$_2$CN	CDMS	$[200]$	$< 2.0 \times 10^{15}$	< 0.1	< 12.2	$[200]$	$< 6.4 \times 10^{14}$	< 0.05	< 5.3

Notes. Band 5 and 6 data are used to fit B1-c, and Band 6 data are used to fit S68N parameters. The column densities for CH$_3$OH are taken from van Gelder et al. (2020), who use the optically thin isotopologue CH$_3^{18}$OH to determine the column density of CH$_3$OH.
The square brackets around the T_{ex} values indicate the molecules for which the temperature was fixed. The column densities for HNCO and CH$_3$CN are found from HN^{13}CO and CH$_2$DCN as their main isotopologues are optically thick. Species with a tentative detection are indicated by a star next to their excitation temperatures.
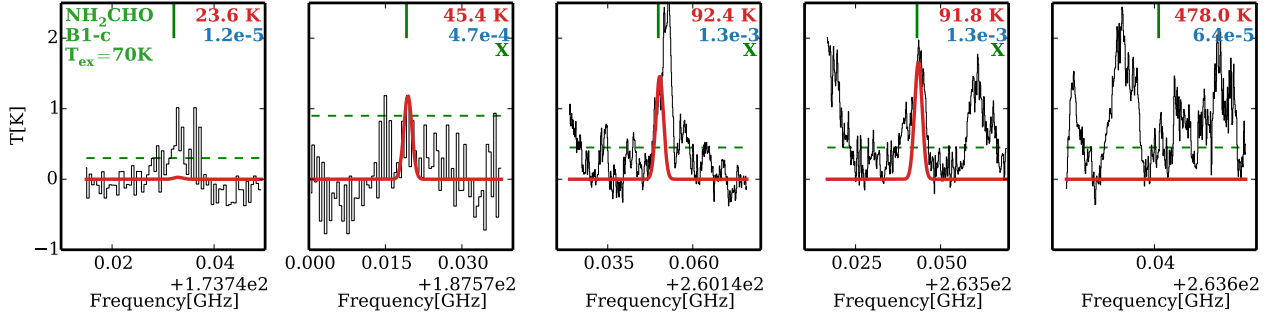
Fig. 1. Best fitted model to combined Band 5 and 6 NH$_2$CHO data for B1-c in red and data in black. Each graph shows one line of NH$_2$CHO with its upper state energy level and the A_{ij} coefficient at the top right in red and blue, respectively. The dashed green line indicates the 3σ level. The lines above the 3σ level that were used in the fitting are indicated by a green X. The lines with upper energy levels above 1000 K and/or A_{ij} below 10^{-5} are not plotted.
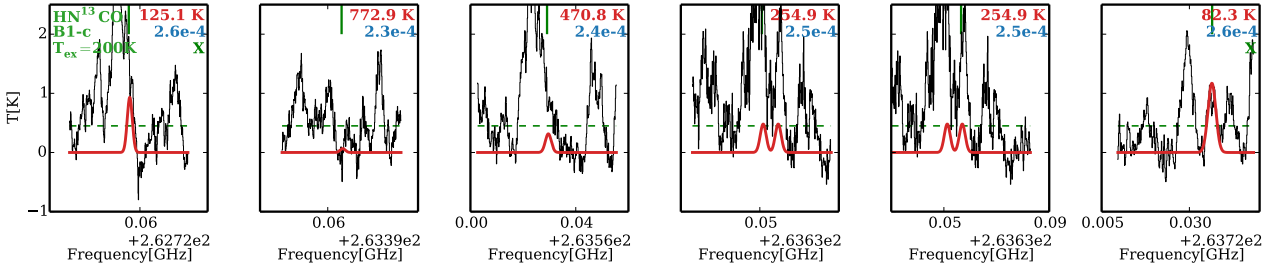

Fig. 2. Same as Fig. 1 but for HN^{13}CO.

The Band 5 data cover HNC^{18}O lines in B1-c. However, this molecule is not detected, and a $\sim 3\sigma$ upper limit column density of $\sim 4 \times 10^{14}$ cm^{-2} at 200 K is found towards B1-c. This upper limit corresponds to a lower limit for the HNCO/HNC^{18}O ratio of $\gtrsim 40$. This lower limit is smaller than the isotope ratio of ^{16}O/^{18}O ~ 560 (Wilson & Rood 1994), implying that HN^{13}CO emission towards B1-c could potentially be marginally optically thick. Using ^{16}O/^{18}O ~ 560 and the upper limit value for HNC^{18}O, an upper limit for HNCO of 2.2×10^{17} cm^{-2} can be found. Therefore, the value for the HNCO column density towards B1-c could potentially differ from the current value by up to an order of magnitude and hence should be taken with care.

In the Band 3 data, HNCO is tentatively detected towards both sources. However, as derived above from our Band 5 and 6 data, it is optically thick, and none of its isotopologues are detected in Band 3. Therefore, we refrained from deriving the HNCO column density from the main isotopologue lines.

Unfortunately, our Band 5 and 6 frequency range covers neither CH$_3$CN nor its ^{15}N or ^{13}C isotopologues. The frequency range covers three rotational CH$_3$CN lines originating from the $v_8 = 1$ excited vibrational level, but they are not very constraining as they give very high upper limits on the column density at a temperature of 200 K ($\lesssim 7 \times 10^{18}$ cm^{-2} and $\lesssim 6 \times 10^{18}$ cm^{-2} for B1-c and S68N, respectively). However, CH$_2$DCN is detected
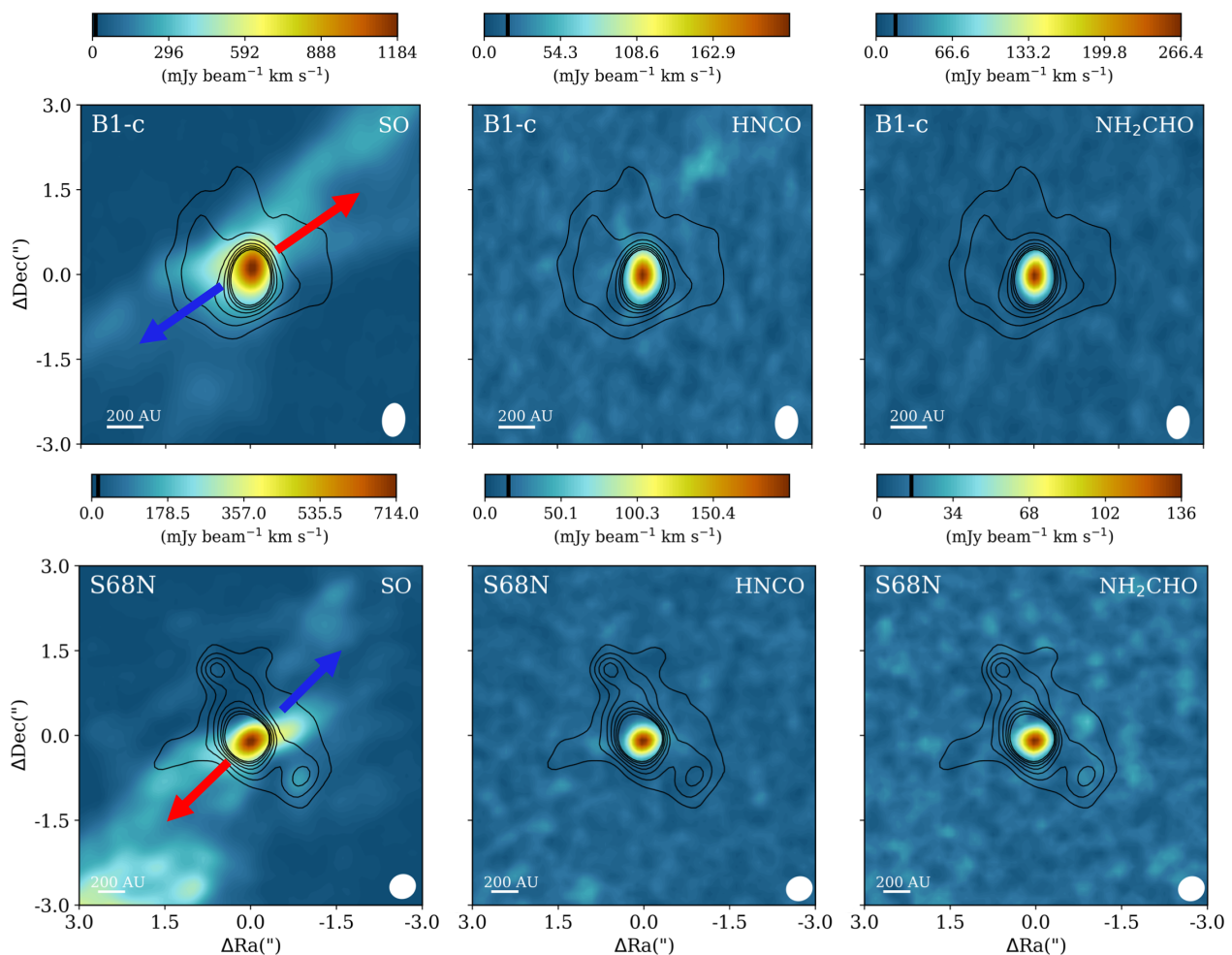
Fig. 3. Moment zero maps of the lines of SO 6_7-5_6 ($E_{\rm up}$ = 47.6 K), HNCO $12_{0,12}-11_{0,11}$ ($E_{\rm up}$ = 82.3 K), and NH$_2$CHO $13_{1,13}-12_{1,12}$ ($E_{\rm up}$ = 91.8 K) (from left to right) for B1-c (top row) and S68N (bottom row) in the Band 6 data. The images are made by integrating over $[-10, 10]$ km s^{-1} with respect to $V_{\rm lsr}$. The black contours show the continuum in levels of $[15, 30, 45, 60, 75, 90, 105]\,\sigma_{\rm cont}$ with a $\sigma_{\rm cont}$ of 0.2 mJy beam^{-1} for B1-c and $[30, 45, 60, 75, 90, 105, 120]\,\sigma_{\rm cont}$ with a $\sigma_{\rm cont}$ of 0.09 mJy beam^{-1} for S68N. The beam size is shown at the right-hand side of each panel. SO traces the extended outflow, while the other two molecules show compact emissions. The approximate directions of the blueshifted and redshifted emission of the outflow are shown with blue and red arrows.

towards both sources. Given that our sources are hot corinos, in each source we used the ratio of CH$_2$DCN/CH$_3$CN = 0.035 from the results of the PILS for IRAS 16293B (Calcutt et al. 2018) to estimate the column density of CH$_3$CN from our value for CH$_2$DCN. This is, however, only an estimate. Interestingly, CHD$_2$CN is also detected towards B1-c, whereas we find an upper limit for this molecule towards S68N. In Band 3, CH$_3$CN is clearly detected towards both sources (see Appendix C), but, given that the lines are optically thick and none of its isotopologues are detected, it is only possible to report lower limits to the actual column densities. These are $\sim10^{15}$ cm^{-2} and $\sim3\times10^{14}$ cm^{-2} at a fixed temperature of 100 K (as van Gelder et al. 2020 found that Band 3 traces colder temperatures) for B1-c and S68N, respectively. These lower limits are consistent with the column densities derived for CH$_3$CN using the Band 6 CH$_2$DCN data towards both sources.

HNCO has the highest column density in both sources. CH$_2$DCN and CHD$_2$CN have the lowest column densities for S68N and B1-c, respectively. The column densities for $\sim85\%$ of the detected and tentatively detected N-bearing species studied here are of the order of $\sim10^{14-15}$ cm^{-2} in a 0.45″ beam for both sources. These values are on average an order of magni-

tude lower than those of the O-bearing molecules studied by van Gelder et al. (2020), which is an interesting observation. Similar to what van Gelder et al. (2020) found, weaker ($T_{\rm b} \lesssim 1$ K) and fewer lines are found towards the more distant S68N source compared to B1-c. On average, the column densities of the species studied in this work are $\sim1.5-3$ times lower in S68N than in B1-c.

3.3. Excitation temperatures

A summary of the derived excitation temperatures is presented in Table 1 for B1-c and S68N. The excitation temperatures found for the N-bearing species span a range between ~100 K and ~300 K. There is no significant difference, within the uncertainties, between the excitation temperatures for the N-bearing COMs and the O-bearing COMs found by van Gelder et al. (2020) (see Sect. 4.1 for a discussion).

For most molecules, a reasonable range of upper state energy level lines is covered in our data, implying that the excitation temperatures fitted here are not biased. The only exception is NH$_2$CHO, where the lines included in the fit only cover upper

energy levels below ~200 K, and thus the excitation temperature derived here is biased towards lower temperatures. Moreover, to derive the excitation temperature of HN^{13}CO towards B1-c, its Band 3 data are used (in addition to the Band 6 data) as an additional constraint, eliminating very low excitation temperatures.

3.4. Column density ratios

Table 1 presents the column density ratios of the studied species with respect to CH$_3$OH and HNCO. The methanol column density is taken from van Gelder et al. (2020). The column density ratios of the species are not calculated with respect to H$_2$ because in low-mass protostars it is difficult to derive accurate values for the warm H$_2$ gas column density from the dust continuum or the CO column density. This is due to the fact that the dust continuum becomes optically thick at $\lesssim$ 100 au scales and may have contributions from a forming, colder disk (Yıldız et al. 2013; Persson et al. 2016; De Simone et al. 2020). Moreover, in low-mass sources not all CO emission comes from the region with warm gas (>100 K) (Yıldız et al. 2013). Therefore, the abundances of the species considered in this work are presented with respect to CH$_3$OH and HNCO instead of H$_2$. Column density ratios with respect to methanol are given in order to compare our results with those of van Gelder et al. (2020) and other studies. The column density ratios with respect to HNCO are calculated to show whether the formation of species discussed here is linked to HNCO or routes resulting in HNCO formation (see Sect. 4.2.2).

The column density ratios for all the detected and tentatively detected molecules with respect to methanol are very similar (within a factor of two) between the two sources. This agrees with what van Gelder et al. (2020) found for most O-bearing COMs. Moreover, the column density ratios for NH$_2$CHO, C$_2$H$_5$CN, HNCO, and NH$_2$CN with respect to methanol agree, within the uncertainties, with the results of Belloche et al. (2020) for S68N observed with the PdBI. It should be noted, though, that our estimated value for CH$_3$CN/CH$_3$OH is ~13.5 times smaller than that derived by Belloche et al. (2020). This discrepancy could originate from the CH$_3$OH lines used in Belloche et al. (2020) being (marginally) optically thick, while the CH$_3$OH column density reported by van Gelder et al. (2020) is derived from its optically thin ^{18}O isotopologue. It should also be noted that we find the CH$_3$CN column density from its deuterated isotopologues, which is another source of uncertainty. Finally, the assumed ^{16}O/^{18}O ratio used to derive the CH$_3$OH column density may differ slightly from the value assumed in van Gelder et al. (2020).

3.5. Source size and optical depth

The emission from our data is not spatially resolved. Therefore, it is not obvious whether the emission is optically thin because the estimated optical depth depends on the assumed size of the emitting region. In this paper it is assumed that the source size is the same as the ALMA Band 6 beam size of 0.45″. A smaller source size will result in larger column densities, and hence the emission can become optically thick. This can be seen by taking the dilution factor into account. The equation $N_1 \theta_{s_1}^2 / (\theta_b^2 + \theta_{s_1}^2) = N_2 \theta_{s_2}^2 / (\theta_b^2 + \theta_{s_2}^2)$ shows how the measured column density can be scaled for different source sizes (where θ_b is the beam size, θ_s is the source size, and the numbers refer to the two assumed source sizes). Although the column densities would become larger for smaller source sizes that depend on the

emitting region of a molecule, column density ratios will stay the same as long as the emission remains optically thin and comes from the same region.

N-bearing species identified in this work show a large range of excitation temperatures (Table 1). It is not unlikely that source sizes are different for molecules with different excitation temperatures. Here, two representative values for the excitation temperatures are considered: one for molecules that are excited at temperatures of ~100 K and one for molecules that are excited at temperatures of ~200 K. Therefore, one can assume two source sizes equal to two radii where temperatures are 100 K and 200 K.

Using the same method as van Gelder et al. (2020), the radius at which the temperature reaches 100 K for a spherically symmetric hot core region can be estimated from $R_{T=100K} \simeq 15.4 \sqrt{L/L_\odot}$ au (Bisschop et al. 2007), where L is the source luminosity. The luminosities of B1-c and S68N are 6.0 L$_\odot$ (Karska et al. 2018) and 5.4 L$_\odot$ (Enoch et al. 2011), respectively. Therefore, the radii at which the temperature reaches 100 K for B1-c and S68N are about 37.7 au (0.23″) and 35.8 au (0.16″), respectively. When these radii are used for the line analysis, NH$_2$CHO remains optically thin in the Band 5 and 6 data towards B1-c even for this smaller source size. Moreover, C$_2$H$_5$CN stays optically thin for S68N at this smaller radius.

To calculate the radius at which the temperature is 200 K, we used a toy model for a spherically symmetric infalling envelope with a power law in temperature and density structure (see Appendix B). For B1-c and S68N, $R_{T=200K}$ is about 6.7 au (0.042″) and 6.3 au (0.029″), respectively. When these radii are used for the line analysis, all the molecules with high T_{ex} (HN^{13}CO, C$_2$H$_5$CN, and CHD$_2$CN towards B1-c and CH$_2$DCN towards both sources) become optically thick in Band 5 and 6 data. Moreover, CH$_3$CN becomes optically thick in Band 3 for both sources, assuming its excitation temperature is similar to that of its deuterated versions (~200 K). It is therefore not possible to derive accurate column densities or column density ratios for such small source sizes.

It is worth noting that some deviation from spherical symmetry due to the presence of a disk can occur at < 100 au scales, and thus our simple toy model runs into limitations. Persson et al. (2016) find that the amount of gas at temperatures above 100 K in low-luminosity sources can vary by more than an order of magnitude depending on the disk size and structure, affecting optical depth. Moreover, the extremely small source sizes predicted by the simple toy model are likely unrealistic. Adopting $R_{T=100K}$ as the source size for both sources, all molecules in Band 5 and 6 remain optically thin and CH$_3$CN remains marginally optically thick ($\tau > 0.1$) in Band 3 for both sources as expected (see Sect. 3.2). For this reason, assuming a source size of 0.45″ in the rest of this paper does not change our results.

4. Discussion

4.1. Excitation versus desorption temperatures

Figure 4 presents a comparison between excitation temperatures of the O-bearing species studied by van Gelder et al. (2020) and the N-bearing species studied in this work towards B1-c and S68N. This figure shows that there is no significant difference in T_{ex} between the O-bearing and N-bearing species; the mean and the scatter of T_{ex} are very similar. Assuming that the excitation temperature is close to the kinetic temperature of the region in which a molecule resides, Fig. 4 suggests that N-bearing species trace a range of temperatures from ~100 K to ~300 K. We note that at the densities considered here some molecules may be sub-
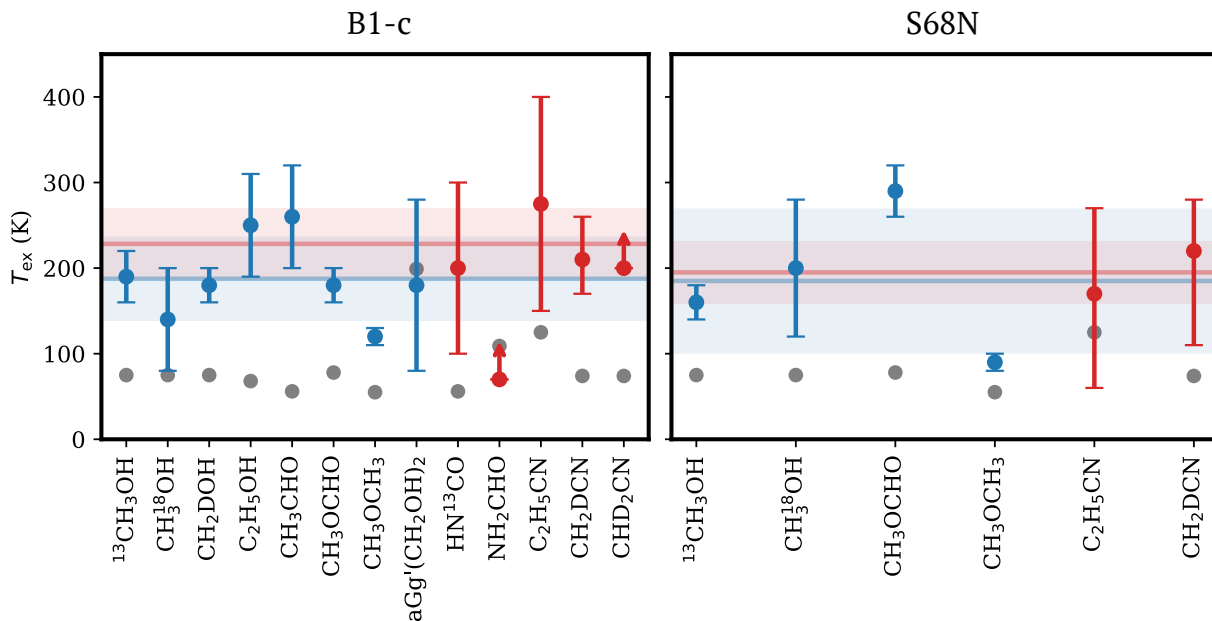
Fig. 4. Excitation temperatures for N-bearing species discussed in this work (red) and O-bearing species discussed in van Gelder et al. (2020) (blue) for B1-c and S68N. Only the species with derived T_{ex} are plotted. The solid red and blue lines show the average values for the red and blue data points, respectively. The shaded red and blue areas show the standard deviation of the data points. The grey points show the sublimation temperatures of the corresponding molecules found by the gas-grain balance model (Hasegawa et al. 1992) and the respective binding energies of each molecule (Penteado et al. 2017; Garrod 2013).

thermally excited, and hence the T_{ex} found in this work may be a lower limit to the kinetic temperature (Jørgensen et al. 2016). Bisschop et al. (2007) also found that there is no difference between the excitation temperatures of N-bearing and O-bearing species in seven high-mass protostars, which is consistent with what is found here.

Each molecule can desorb at a significantly different temperature according to its binding energy, ice, and grain environment (Cuppen et al. 2017), so molecules come off the ice roughly at their respective snow lines depending on their sublimation temperatures. Therefore, one would expect differences in the excitation temperatures of molecules. The idea of an onion-like temperature structure around low-mass protostars was discussed in Jørgensen et al. (2018), in which O-bearing molecules are divided into two categories (see also Manigand et al. 2020): one that includes molecules associated with temperatures of 100-150 K and another that includes molecules associated with temperatures of 250-300 K. Figure 4 shows that C_2H_5CN towards B1-c falls under the hottest category (> 250 K). In addition, three of the N-bearing molecules studied here ($HN^{13}CO$, CH_2DCN, and CHD_2CN) seem to trace the gas with temperatures around ~200 K in B1-c.

The potential relation between the excitation and sublimation temperatures of molecules can be further explored by comparing our results with the sublimation temperatures found from the gas-grain balance model (Hasegawa et al. 1992) and the binding energies of each molecule in the solid state. In this model, the number density of the solid to the gas phase of species i is given by

$$\frac{n_{ice}}{n_{gas}} = \frac{\pi a_d^2 n_d S \sqrt{3k_B T_{gas}/m_i}}{e^{-E_b/T_d} \sqrt{2k_B n_{ss} E_b/(\pi^2 m_i)}}, \tag{1}$$

where a_d is the dust grain size (assumed to be $0.1\,\mu m$), $n_d = 10^{-12} \times n_H$ is the dust number density (with n_H the hydrogen number density, assumed to be $10^7\,cm^{-3}$), S is the sticking coefficient (assumed to be 1), n_{ss} is the number of binding sites per surface area (taken as $8 \times 10^{14}\,cm^{-2}$), E_b is the binding energy of species i in units of kelvin, m_i is the mass of species i, and T_{gas} and T_d are the gas and dust temperatures, respectively. Assuming that the environment is sufficiently dense, the gas would be thermally coupled with the dust and T_d would be equal to T_{gas}. The sublimation temperature of species i is the temperature for which the ice and the gas are in balance for that molecule. In other words, $n_{ice}/n_{gas} = 1$.

Therefore, using Eq. (1) and the binding energies for most molecules from Penteado et al. (2017) (the binding energy for $(CH_2OH)_2$ is taken from Garrod 2013), the desorption temperatures of O-bearing and N-bearing species can be calculated. The grey points in Fig. 4 represent these values. It can be safely assumed that the desorption temperatures for the isotopologues will barely differ from that of the main isotopologue. However, binding energies differ for different molecule interactions in a mixed ice or on different grain surfaces (e.g. Tielens et al. 1991; Collings et al. 2004; Ferrero et al. 2020), and hence in principle desorption temperatures may vary (typically by several tens of kelvin). The values adopted here are mostly for pure ices; if these COMs are mixed in ice layers consisting predominantly of H_2O or CH_3OH, values close to their desorption temperatures (~100 K) are expected.

Figure 4 shows that there is no significant difference between the sublimation temperatures of O-bearing and N-bearing species. This is consistent with what is seen for the excitation temperatures, but with an offset between sublimation temperatures and excitation temperatures for most molecules (except for $(CH_2OH)_2$). Specifically, the excitation temperatures are typically higher than the sublimation temperatures for all molecules

by a factor of $\sim 2 - 4$, except for $(CH_2OH)_2$ where these two values are very similar. This can be interpreted as the excitation temperatures reflecting an average temperature of the region around the protostar, where it is hotter than the sublimation temperature of each molecule. Using the toy model for a spherically symmetric infalling envelope with power law structure in temperature and density (see Appendix B.4), the mass-weighted average temperature can be calculated as $1.36 T_{sub}$. Therefore, if the excitation temperature is measuring an average temperature, it is expected to be larger than the sublimation temperature for each molecule. We note that the factor 1.36 is not large enough to explain the difference seen in Fig. 4. However, this is only using a simple toy model: Developing a more complete model that takes small-scale structures such as disks into account can improve our understanding of the difference between sublimation temperatures and excitation temperatures.

The sublimation temperatures for all species in Fig. 4 (except for $(CH_2OH)_2$) are between ~ 55 and ~ 125 K. While the sublimation temperatures of O-bearing and N-bearing groups of species do not differ significantly, these values for some molecules (i.e. CH_3OCH_3, CH_3CHO, and HNCO) are lower than the rest, and others (i.e. C_2H_5CN and $(CH_2OH)_2$) are higher. Such trends are reflected in the excitation temperatures for some of the molecules with small error bars. For example, CH_3OCH_3 shows very low excitation temperatures in both sources, which agrees with its low sublimation temperature. This may be good evidence for the formation of this molecule in the ice and its desorption at its snow line. On the other hand, a molecule such as CH_3CHO shows a larger excitation temperature compared to the other species, in contradiction with its low sublimation temperature. This could be evidence for other mechanisms (e.g. gas-phase formation pathways or source structures) potentially playing a role in the formation (Garrod et al. 2008; Tideswell et al. 2010; Vazart et al. 2020) and emission of these species. However, given the large error bars on most of the excitation temperatures shown in Fig. 4, all these arguments should be taken with care.

Using *Herschel*-HIFI data, Crockett et al. (2015) find that N-bearing species, C_2H_5CN, NH_2CHO, and CH_3CN, trace hotter gas (~ 250 K) compared with the oxygen-bearing species, CH_3OH, CH_3OCH_3, C_2H_5OH, and CH_3OCHO, (~ 100 K), in the Orion Kleinmann-Low nebula (see Fig. 5). They argue that this could indicate either that N-bearing COMs require higher dust temperatures to desorb from the ice or that more N-bearing molecules are formed through high temperature gas-phase chemistry, reflecting the two possible COM formation scenarios: colder solid-state chemistry and warmer gas-phase chemistry.

Alternatively, at such high temperatures Crockett et al. (2015) could be seeing the carbon grain sublimation into the gas phase that leads to the formation of CH_3CN, an idea put forward by van 't Hoff et al. (2020a). They propose the formation of N-rich COMs inside a 'soot line' at 300 K through carbon grain sublimation and suggest that an excess of hydrocarbons and nitriles with excitation temperatures higher than those of O-bearing species can be a signature of this phenomenon. This process implies that such molecules could be formed through top-down chemistry (from the destruction of larger species) rather than bottom-up chemistry in the solid state or gas phase. van 't Hoff et al. (2020a) assume that no molecules with oxygen atoms are formed through carbon grain sublimation inside of the soot line at ~ 300 K as all the oxygen is locked up in H_2O, CO, and CO_2. Thus, the only two molecules in this work that can be signatures for carbon grain sublimation are C_2H_5CN and CH_3CN. The toy model for a spherically symmetric infalling envelope explained in Appendix B.4 predicts the ~ 300 K radius to be signif-
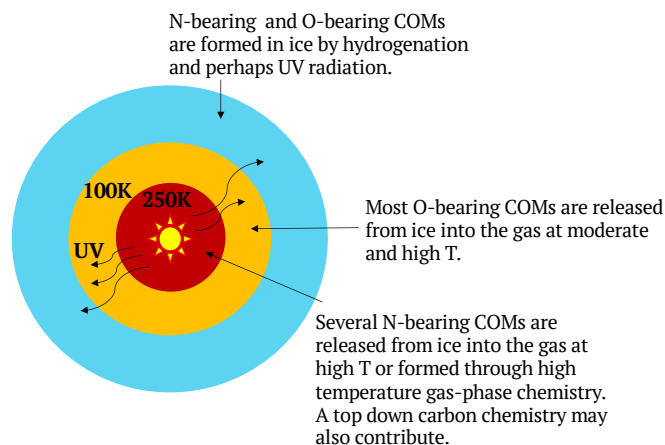
Fig. 5. Cartoon of the temperature structure of protostellar envelopes showing where O-bearing and N-bearing species are most likely to be found in the gas. This is based on the findings of Crockett et al. (2015) for massive protostars, who showed that higher excitation temperatures for N-bearing COMs are found towards several high-mass protostars, and van 't Hoff et al. (2020a), who proposed a top-down chemistry for nitriles. It is not yet clear whether this structure also applies to low-mass protostars.

icantly smaller than the spatial resolution of our data. Therefore, the fact that carbon grain sublimation is not seen for C_2H_5CN and CH_3CN in B1-c or S68N might be due to this effect being concealed at the present spatial resolution.

4.2. Comparison of abundances in different sources

In this section the column density ratios of B1-c and S68N are compared with other sources. In most studies, column density ratios are interpreted as abundances. Therefore, this section also assumes no difference between column density ratios and abundances. The validity of this assumption is further discussed in Sect. 4.3.

4.2.1. Dependence on luminosity

Higher luminosity implies a higher temperature, where chemical reactions can take place more efficiently thereby affecting the production rate of molecules and possibly changing their abundances. In addition, a source with a higher luminosity has more UV radiation, and hence the UV processing on the grains can alter the gas-phase abundances of species, for example through (non-)dissociative photo-desorption but also following UV-induced photochemistry in the ice (Öberg et al. 2009b; Bertin et al. 2013). Therefore, a relation between the luminosity and the abundance of a species with respect to methanol may be present. This is especially important for species with higher desorption temperatures (e.g. NH_2CHO; see Sect. 4.1) that may stay mostly in the ice in sources with lower luminosities.

Figure 6 shows the abundances of NH_2CHO and HNCO with respect to methanol for our sources and other low- and high-mass protostars against the source luminosity (see the caption of Fig. 6 for the full list of sources and reference papers). There is no significant relation between abundances and source luminosity. The same conclusion holds true for the rest of the N-bearing species, which is consistent with the results of van Gelder et al. (2020) for
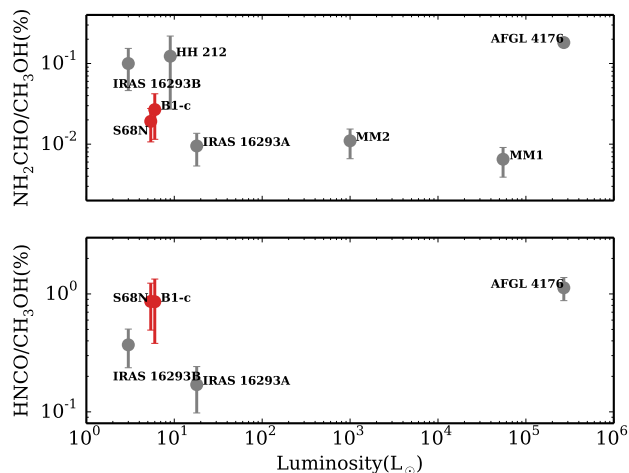
Fig. 6. Abundance ratio of NH₂CHO and HNCO with respect to CH₃OH against source luminosity. Red points show the results for the sources considered in this work. The grey points indicate other low- and high-mass sources. These sources are IRAS 16293A (Ligterink et al. 2017; Manigand et al. 2020), IRAS 16293B (Coutens et al. 2016; Ligterink et al. 2017; Jørgensen et al. 2018), HH 212 (Lee et al. 2019), AFGL 4176 (Bøgelund et al. 2019a), and NGC6334 MM1 and MM2 (Bøgelund et al. 2019b).

the O-bearing COMs. Moreover, our results are consistent with what Belloche et al. (2020) find for protostellar sources observed by the PdBI. Nevertheless, Fig. 6 shows that NH₂CHO/CH₃OH ranges from ~0.01% to ~0.1% (discussed further in Sect. 4.2.2), whereas van Gelder et al. (2020) find that for most O-bearing molecules the column density ratios with respect to methanol are within a factor of a few for various sources. Given that in all these sources the methanol column density is derived from one of its optically thin isotopologues, the difference seen here cannot be due to issues with optical depth.

4.2.2. Abundances with respect to CH₃OH and HNCO

We compared our abundances to values obtained for other low-mass and high-mass sources that have been studied by ALMA. The low-mass sources are IRAS 16293A (Coutens et al. 2016; Ligterink et al. 2017; Calcutt et al. 2018; Manigand et al. 2020), IRAS 16293B (Coutens et al. 2016; Ligterink et al. 2017; Jørgensen et al. 2018; Calcutt et al. 2018; Coutens et al. 2018), HH 212 (Lee et al. 2019), and NGC 1333 IRAS 4A2 (López-Sepulcre et al. 2017). The high-mass sources are AFGL 4176 (Bøgelund et al. 2019a) and NGC 6334 MM1, MM2, and MM3 (Bøgelund et al. 2019b).

Figure 7 shows this comparison for the ratio of the column densities to methanol. At a glance, we can see that these values are generally lower than those of O-bearing species (Fig. 7 in van Gelder et al. 2020), typically by an order of magnitude. Moreover, there is a larger variation (up to an order of magnitude) between the different sources for N-bearing molecules than for O-bearing ones (by a factor of a few).

When comparing the column density ratios of the N-bearing species shown in Fig. 7 among different sources, one can see that some molecules show similar abundance ratios. Most notably, C₂H₅CN is comparable in the low-mass sources and with the high-mass AFGL 4176 (within a factor of ~5).

A larger scatter (around an order of magnitude) is seen for NH₂CHO between all sources (see also Fig. 6). Three cate-

gories seem to be present. First, two of the high-mass sources (NGC6334 MM1 and MM2) seem to agree with three of the low-mass sources (B1-c, S68N, and IRAS 16293A). Second, the high-mass source, AFGL 4176, is similar to the two other low-mass sources, IRAS 16293B and HH212. Third, the high-mass source NGC 6334 MM3 does not show a detection of this COM. Manigand et al. (2020) explain the low abundance ratio of formamide in IRAS 16293A compared with IRAS 16293B by the fact that the offset position used to extract the spectrum for the former seems to be missing the hottest gas and thus most of the emission from this molecule, while this is not the case for the latter. Moreover, in contrast with our results, they find that formamide belongs to a group of species with a high excitation temperature. Another possible explanation for the variation in formamide abundance ratios with respect to methanol across various sources is the fact that there are multiple ways to form formamide in the solid state (Rimola et al. 2018; Haupa et al. 2019; Dulieu et al. 2019; Martín-Doménech et al. 2020) and gas phase (Barone et al. 2015; Codella et al. 2017). Therefore, the amount of this COM formed in the gas or ice could depend on local conditions, such as the presence of a disk, its temperature structure (Jørgensen et al. 2002; Schöier et al. 2002; Crimier et al. 2010), or its exposure to UV (López-Sepulcre et al. 2019).

HNCO and its isotopologues show some scatter in their distributions with respect to methanol. B1-c (0.9 ± 0.5%), S68N (0.9 ± 0.4%), and the high-mass source AFGL 4176 (1.1 ± 0.3%) seem to have a more similar distribution of HNCO with respect to methanol, while the IRAS 16293 sources show a lower abundance ratio (0.17 ± 0.07% and 0.37 ± 0.13% for IRAS 16293A and B, respectively).

Figure 8 shows the abundance ratios of the molecules studied in this work with respect to HNCO. The scatter that was seen in Fig. 7 seems to be less pronounced for most molecules. Among them, DNCO has a more uniform distribution between sources, as expected for a molecule formed directly via the deuteration of HNCO (Noble et al. 2015). Another molecule that shows a much more uniform distribution among sources is CHD₂CN. Ligterink et al. (2020) studied amide molecules towards 12 positions in the high-mass star-forming region NGC 6334I. Their values of NH₂CHO/HNCO measured from the ¹³C isotopologues of these two species span a range between ~10% and ~100%. This range is ~1 − 2 orders of magnitude higher than what is found in this work towards B1-c and S68N, and more similar to IRAS 16293B (Fig. 8).

4.2.3. Deuteration fraction

B1-c and S68N are known to have a high deuteration fraction, as confirmed for methanol by van Gelder et al. (2020) from CH₂DOH. Jørgensen et al. (2018) found that less complex molecules have similar D/H ratios. For example, the D/H ratios for HNCO (~1%; Coutens et al. 2016) and CH₃OH (2%; Jørgensen et al. 2018) in IRAS 16293B are comparable. Therefore, we similarly compared D/H for these two molecules in B1-c and S68N. The deuteration fraction found from this work for B1-c and S68N from the DNCO/HNCO ratio are 1.8 ± 0.8% and < 0.6%. The value for B1-c is in agreement with the 2.7 ± 0.9% methanol value reported in van Gelder et al. (2020), whereas the upper limit found for S68N is somewhat lower than the 1.4±0.6% methanol value reported in the same study.

The D/H ratio can provide a clue as to the temperature history of the environment in which B1-c is forming. This can also be seen in Fig. 8, where DNCO, CHD₂CN, and CH₂DCN have abundance ratios with respect to HNCO that agree, within the
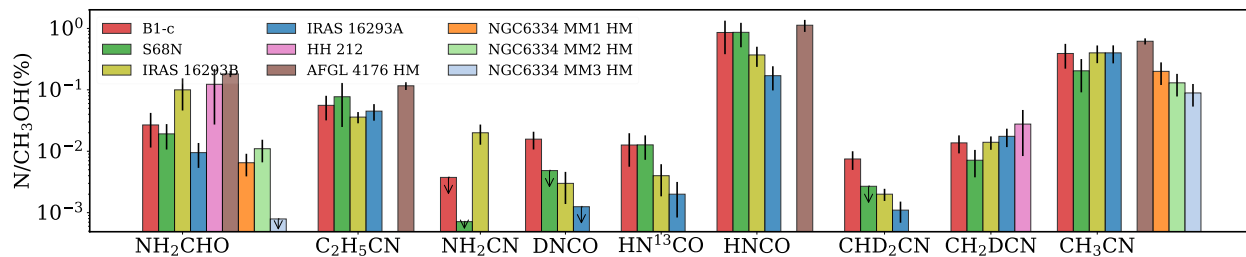
Fig. 7. Column densities of the N-bearing molecules with respect to methanol for sources observed by ALMA. Column densities of methanol for B1-c and S68N are taken from van Gelder et al. (2020). Shown are IRAS 16293A (Coutens et al. 2016; Ligterink et al. 2017; Calcutt et al. 2018; Manigand et al. 2020), IRAS 16293B (Coutens et al. 2016; Ligterink et al. 2017; Jørgensen et al. 2018; Calcutt et al. 2018; Coutens et al. 2018), HH 212 (Lee et al. 2019), AFGL 4176 (Bøgelund et al. 2019a), and NGC 6334 MM1, MM2, and MM3 (Bøgelund et al. 2019b).
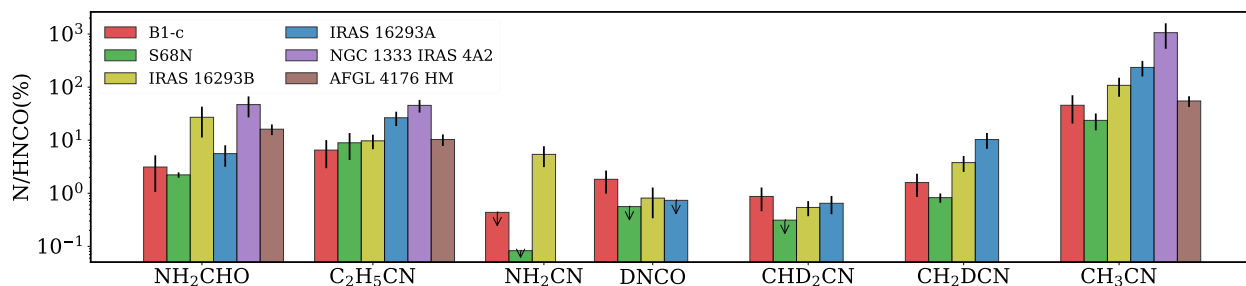
Fig. 8. Column densities of N-bearing species with respect to HNCO as derived from HN^{13}CO. Shown are IRAS 16293A (Coutens et al. 2016; Ligterink et al. 2017; Calcutt et al. 2018; Manigand et al. 2020), IRAS 16293B (Coutens et al. 2016; Ligterink et al. 2017; Calcutt et al. 2018; Coutens et al. 2018), NGC 1333 IRAS 4A2 (López-Sepulcre et al. 2017), and AFGL 4176 (Bøgelund et al. 2019a).

uncertainties, between the low-mass sources. These sources are from different star-forming regions and yet show similar abundance ratios, pointing to the fact that there is a universal mechanism for increasing the D/H ratio prior to star formation in cold environments.

For CH$_2$DCN/CHD$_2$CN, a column density ratio of ~2 is found towards B1-c, which is lower than that found for IRAS 16293A (~16) and IRAS 16293B (~7) by Calcutt et al. (2018). On the other hand, Taquet et al. (2019) find CH$_2$DOH/CHD$_2$OH ~ 1 and ~2 for NGC1333-IRAS2A and NGC1333-IRAS4A, values that are more similar to our ratio for the deuterated versions of methyl cyanide. In general, these similarities and differences are interesting, but, given the small sample of hot corinos with a similar analysis at hand, we did not investigate this further.

4.2.4. Comparison with ices

B1-c is an ice-rich source, but so far CH$_3$OH is the only COM that has been identified along the line of sight towards the neighbouring B1-b source with *Spitzer* (Boogert et al. 2008). Ices have also been detected towards S68N (e.g. Anderson et al. 2013), but they have not yet been analysed in detail. However, OCN$^-$, a direct derivative of HNCO (van Broekhuizen et al. 2004; Fedoseev et al. 2016), is observed towards many low- and high-mass protostars (van Broekhuizen et al. 2005; Öberg et al. 2011). Öberg et al. (2011) find that, for the low-mass sources in their study, the median values for OCN$^-$ and CH$_3$OH ice abundances with respect to water ice are $0.4^{0.4}_{0.3}\%$ and $7^{12}_{5}\%$, respectively. This gives an abundance ratio of OCN$^-$ with respect to CH$_3$OH of ~$6^{8}_{2.5}\%$. This abundance ratio shows a considerable spread from source to source. Our gas-phase abundance ratios of HNCO to CH$_3$OH

are ~0.9±0.5% and ~0.9±0.4% for B1-c and S68N, respectively (Table 1). These values are ~6 times lower than the median ice ratio but close to its lower limit. This comparison assumes that all OCN$^-$ comes off the grains as HNCO upon ice sublimation, but some of it may be converted into other species, such as more refractory salts (Schutte & Khanna 2003; Boogert et al. 2008). This is a reasonable assumption given that N-bearing salts are found to be abundant in comet 67P (Altwegg et al. 2020).

In addition, there are also upper limit estimates for HNCO ice itself, but mostly for high-mass protostars. Taking the upper limit of 0.7% for HNCO/H$_2$O ice from van Broekhuizen et al. (2004, 2005) and dividing it by the median of the CH$_3$OH/H$_2$O ice ratio in high-mass protostars ($8^{16}_{8}\%$) from Öberg et al. (2011) gives an upper limit estimate of 9% for the HNCO/CH$_3$OH ice ratio. Again, this is higher than our observed ratios in low-mass hot corinos.

These differences are interesting, but without the direct observation of the species studied here in ices for the same sources it is not possible to investigate this further. The column density ratios discussed in this section need to be studied in a larger sample of protostars to make progress in this field. More robust conclusions will become possible with the launch of JWST and the direct observation of these molecules in ices. In addition, progress in recording the solid state features of several species will make this comparison more plausible in the near future (Boudin et al. 1998; Terwisscha van Scheltinga et al. 2018; Gerakines & Hudson 2020).

4.3. Effect of emitting areas of molecules on abundance ratios

As the emission from our data is not spatially resolved, variations from source to source or molecule to molecule could be explained by different emitting regions that are possibly related to different binding energies. Using the spherically symmetric infalling envelope toy model (see Appendix B), we can find a relation between the ratio of column densities of two molecules in the same beam with the ratio of their corresponding sublimation temperatures. This is given as

$$\frac{N_1}{N_2} = \frac{X_1}{X_2}\left(\frac{T_{1,\mathrm{sub}}}{T_{2,\mathrm{sub}}}\right)^{-3.75}, \qquad (2)$$

where X is the abundance of a molecule in gas phase with respect to molecular hydrogen, N is the column density, and T_{sub} is the sublimation temperature, with subscripts 1 and 2 indicating the two species. Equation (2) shows that the column density ratios depend not only on the abundance ratios but also on the ratio of sublimation temperatures. This is because the difference in the sublimation temperature results in a difference in the emitting volume.

Figure 4 shows that a typical O-bearing or N-bearing molecule, $\mathcal{M}$, comes off the grains at a temperature of $\sim$100 K. Hence, one can rewrite Eq. (2) for species 1 as

$$f_{\mathrm{sub}} \equiv \frac{X_{\mathcal{M}}}{X_1}\frac{N_1}{N_{\mathcal{M}}} = \left(\frac{T_{1,\mathrm{sub}}}{100}\right)^{-3.75}. \qquad (3)$$

This can be rewritten as

$$\frac{X_1}{X_{\mathcal{M}}} = \frac{1}{f_{\mathrm{sub}}}\frac{N_1}{N_{\mathcal{M}}}. \qquad (4)$$

The column density ratios discussed in this work with respect to methanol, which has a T_{sub} of $\sim$ 100 K, are thus the abundance ratios with respect to methanol times a factor (f_{sub}). The abundance of a molecule is a local quantity that is sensitive to the emitting area, but column density is an integrated quantity within a given beam. In other words, using the expression 'abundance' to refer to a column density ratio is only accurate if the emitting areas are similar.

Figure 9 shows f_{sub} as defined in Eq. (3) against the sublimation temperature for a molecule. This figure shows that, for species with desorption temperatures between 55 K and 125 K (see the grey points in Fig. 4), one would need to multiply the column density ratios of the species with respect to methanol by a factor of between 0.1 and 2 to get the respective abundance ratios. This factor becomes much larger as the sublimation temperature increases. For instance, $(\mathrm{CH_2OH})_2$ has a sublimation temperature of $\sim$200 K, and hence the factor becomes $\sim$14 for this molecule. However, this factor is not applied here because of the simplicity of our model. In reality, one needs to build a model that takes small-scale structures such as disks into account (Harsono et al. 2015; Persson et al. 2016), as discovered by other studies of Class 0 objects (Tobin et al. 2012; Murillo et al. 2013; Martín-Doménech et al. 2019; Maret et al. 2020). This is important because the temperature profile in an envelope of a Class 0 object does not correspond to the temperature profile in the inner regions of the disk (Murillo et al. 2018; van 't Hoff et al. 2020c). Moreover, the opacity of the dust continuum emission can also become relevant at $\lesssim$ 50 au scales (De Simone et al. 2020).
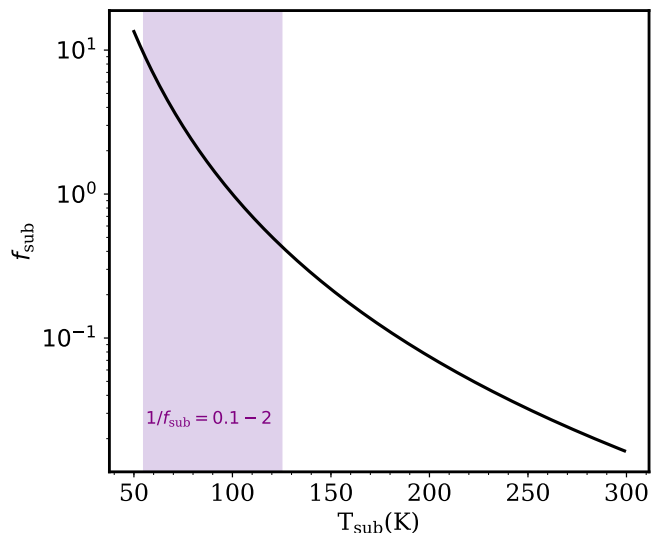
Fig. 9. Correction factor, f_{sub}, introduced in Eq. (4) due to different emitting areas arising from the sublimation temperature. The shaded area shows where most of the species considered in this work fall on this graph (sublimation temperatures between 55 K and 125 K). The abundance ratios of these species with respect to methanol are found by multiplying their column density ratios with respect to methanol by a factor of between 0.1 and 2.

In conclusion, some of the variations seen in abundances between various molecules and between different sources may be related to different emitting regions originating from different binding energies rather than from chemistry.

5. Conclusions

In this work we derived the column densities and excitation temperatures of N-bearing species in two Class 0 objects, B1-c and S68N, using the Band 6, Band 5, and Band 3 of ALMA. The main conclusions are listed below.

- Four N-bearing molecules and their isotopologues are identified towards B1-c and S68N with ALMA. These species are HNCO, HN^{13}CO, DNCO, C$_2$H$_5$CN, NH$_2$CHO, CH$_3$CN, CH$_2$DCN, and CHD$_2$CN. None of these species are spatially resolved, but they are located within 200 au of the central protostar.
- N-bearing and O-bearing species show a similar scatter in their excitation temperatures ($\sim$100$-$300 K), and the average excitation temperature of N-bearing species is roughly the same as the average for the O-bearing species studied by van Gelder et al. (2020) for B1-c and S68N.
- The excitation temperatures for the N-bearing and O-bearing species are larger than the sublimation temperatures of each molecule, which span a range between $\sim$55 and $\sim$125 K. This can be interpreted as the excitation temperature measuring a mass-weighted average temperature of the region that is hotter than the sublimation temperature of the detected molecule.
- The abundances of the N-bearing species with respect to methanol and HNCO are very similar for the two sources studied in this work. Their abundances with respect to methanol are lower than those of the O-bearing species by an order of magnitude.

- Overall, N-bearing species show more uniform abundance ratios with respect to HNCO than relative to methanol between low- and high-mass sources.
- CH_2DCN, CHD_2CN, and DNCO show similar abundance ratios with respect to HNCO between different low-mass sources, suggesting a universal mechanism for increasing the D/H in cold regions prior to the star formation process.
- The NH_2CHO/CH_3OH varies within an order of magnitude between different sources. This can be due either to the different formation pathways suggested for formamide in the solid state or the high sublimation temperature of formamide and the different local temperatures where formamide is detected.
- Variation in abundance ratios between various molecules and sources could be due to the different emitting areas related to the binding energies of the species in spatially unresolved emission.

We emphasise that the sample of COM-rich low-mass protostars that have been subjected to such analyses is still very small. Therefore, reaching robust conclusions on the chemistry of the species discussed here is difficult. An increase in the size of this sample by an order of magnitude will enhance our understanding considerably. Moreover, a broader frequency coverage can help to determine more accurate excitation temperatures and identify other molecular species. Therefore, more ALMA and NOEMA observations of COM-rich Class 0/I objects are necessary for moving this field forward. Moreover, JWST data of the sources discussed here can give a direct comparison between the ice abundances and the gas abundances found in this work. The conclusions presented here can act as a guideline for this future work.

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Appendix A: Spectroscopic data

The spectroscopic data for HNCO are taken from the CDMS (Müller et al. 2001; Müller et al. 2005), which is based on the work by Kukolich et al. (1971), Hocking et al. (1975), Niedenhoff et al. (1995), and Lapinov et al. (2007). DNCO and HN^{13}CO data are based on Hocking et al. (1975) and are taken from the JPL database (Pickett et al. 1998).

The line list for C_2H_5CN is taken from the CDMS. The molecule entries are based on the work by Pearson et al. (1994), Fukuyama et al. (1996), and Brauer et al. (2009). A vibrational factor of 1.8316 is used at 200 K (Heise et al. 1981).

The line list for NH_2CHO is taken from the JPL database. The methods used to analyse the experimental measurements are in Kirchhoff (1972). The measurements are taken from Kurland & Wilson (1957), Costain & Dowling (1960), Kukolich & Nelson (1971), Johnson et al. (1972), and Kirchhoff & Johnson (1973). An upper limit of 1.5 for the vibrational correction factor of NH_2CHO is used at 300 K.

The spectroscopic data used for CH_3CN are taken from the CDMS entry (Bocquet et al. 1988; Koivusaari et al. 1992; Tolonen et al. 1993; Anttila et al. 1993; Gadhi, J et al. 1995; Šimečková et al. 2004; Cazzoli & Puzzarini 2006; Müller et al. 2009; Müller et al. 2015). The line lists of CH_2DCN and CHD_2CN are taken from the CDMS (Nguyen et al. 2013). CH_2DCN has additional data from Le Guennec et al. (1992) and Müller et al. (2009). CHD_2CN has additional data from Halonen & Mills (1978). The vibration factor for CH_3CN is lower than 1.1 up to a temperature of 180 K (Müller et al. 2015) and is thus negligible in this work. It is assumed that the difference between the vibrational factors for the main molecule and its isotopologues are small at the excitation temperatures found in this work.

The spectroscopic data for CH_3NCO are taken from the CDMS (Koput 1986; Cernicharo et al. 2016). The partition function takes vibrational states into account up to 580 K. The line list for NH_2CN is taken from the JPL database (Read et al. 1986). The spectroscopic data for $HOCH_2CN$ are taken from the CDMS (Margulès et al. 2017). The line list for CH_3NH_2 is taken from the JPL database (Ilyushin et al. 2005; Kréglewski & Wlodarczak 1992; Kręglewski et al. 1992; Ohashi et al. 1987; Takagi & Kojima 1971; Nishikawa 1957; Lide 1954; Shimoda et al. 1954; Lide 1957).

Appendix B: Toy model for a spherically symmetric envelope

Appendix B.1: Physical model

A toy model for a spherically symmetric infalling envelope can be developed by assuming a power law structure in density and temperature. The number density of hydrogen is assumed to be

$$n_H(R) = n_{H,0} \left(\frac{R}{R_0} \right)^{-\alpha}, \tag{B.1}$$

where α is 3/2 for a free-falling envelope, R is the radius, and $n_{H,0}$ is the number density of hydrogen at radius R_0. Moreover, assuming that the envelope is passively heated by the central protostar, the temperature profile can be written as

$$T(R) = T_0 \left(\frac{R}{R_0} \right)^{-\beta}, \tag{B.2}$$

where $\beta \simeq 2/5$ (at $R \gtrsim 10$ au) from radiative transfer calculations (Adams & Shu 1985) and T_0 is the temperature at R_0 and is proportional to $L_{bol}^{1/5}$, where L_{bol} is the bolometric luminosity of the source.

In this model we assume that a constant amount of material crosses the snow line and sublimates into the gas phase without further gas-phase reactions. Therefore, each COM comes off the ice at its sublimation temperature, T_{sub}, corresponding to a sublimation radius of $R_{T_{sub}}$. This radius is found by rearranging Eq. (B.2):

$$R(T_{sub}) = R_0 \left(\frac{T_{sub}}{T_0} \right)^{-1/\beta}. \tag{B.3}$$

Appendix B.2: Scaling of COM emission with luminosity

Assuming the emission is optically thin, the total number of a specific COM in the gas phase that comes off the ice at temperatures above its sublimation temperature is proportional to the line flux. The total number of a molecule of a specific COM, $\mathcal{N}_i$, is given by:

$$\mathcal{N}_i = \int_0^{R_{T_{sub}}} 4\pi X_i n_H(R) R^2 dR \tag{B.4}$$

$$= 4\pi X_i \int_0^{R_{T_{sub}}} n_{H,0} \left(\frac{R}{R_0} \right)^{-\alpha} R^2 dR \tag{B.5}$$

$$= \frac{4\pi}{3-\alpha} X_i n_{H,0} R_0^\alpha R_{T_{sub}}^{3-\alpha}, \tag{B.6}$$

where X_i is the abundance of species i with respect to hydrogen atoms.

Dropping the factors that do not depend on the source properties, Eqs. (B.3) and (B.6) yield

$$\mathcal{N}_i \propto n_{H,0} R_{T_{sub}}^{3-\alpha} \propto n_{H,0} T_0^{(3-\alpha)/\beta} \propto n_{H,0} L_{bol}^{(3-\alpha)/(5\beta)}, \tag{B.7}$$

and assuming $\alpha = 3/2$ and $\beta = 2/5$ gives

$$\mathcal{N}_i \propto n_{H,0} L_{bol}^{3/4}. \tag{B.8}$$

Therefore, the toy model predicts that the total number of molecules depends on the density of the envelope at a reference radius, R_0, and on the luminosity of the source.

Appendix B.3: COM emission as a function of T_{sub}

The measured column density of a molecule depends on the assumed source size. If the emission is spatially unresolved and optically thin, one generally assumes a fixed source size for the different species. Hence, the ratio of the measured column densities of two species, 1 and 2, does not depend on the assumed source size. Moreover, the ratio of the measured column densities, assuming the same emitting region, is the same as the ratio of the number of the two molecules in the gas phase. Therefore, the ratio of column densities as found in Sect. 3.2 (N_1/N_2) is the same as the ratio of the total number of molecules for two COMs ($\mathcal{N}_1/\mathcal{N}_2$). For the rest of this section, N_1/N_2 and $\mathcal{N}_1/\mathcal{N}_2$ are the same.

From Eq. (B.6), this ratio is given by

$$\frac{N_1}{N_2} = \frac{X_1}{X_2} \left(\frac{R_{T_{sub,1}}}{R_{T_{sub,2}}} \right)^{3-\alpha}. \tag{B.9}$$

Using Eq. (B.3) and assuming $\alpha = 3/2$ and $\beta = 2/5$, this becomes

$$\frac{N_1}{N_2} = \frac{X_1}{X_2}\left(\frac{T_{\mathrm{sub},1}}{T_{\mathrm{sub},2}}\right)^{-3.75}. \tag{B.10}$$

Therefore, the measured column density ratio for two molecules is not the same as their abundance ratio: The former also depends on the ratio of their sublimation temperatures. In principle, if the thermal structure of the envelope is well constrained, one can recover the abundance ratio from the measured column density ratio by applying the correction factor defined in Eq. 3.

Appendix B.4: Mass-weighted kinetic temperature

Assuming that the excitation temperatures measured in this work probe the mass-weighted kinetic temperature of the environment, one can find a relation between this quantity and the sublimation temperature of each COM. The mass-weighted temperature is given by

$$< T > = \frac{\int_0^{R_{T_{\mathrm{sub}}}} T(R)n_{\mathrm{H}}(R)R^2 dR}{\int_0^{R_{T_{\mathrm{sub}}}} n_{\mathrm{H}}(R)R^2 dR}. \tag{B.11}$$

Using Eqs. (B.1), (B.2), and (B.3), one can find the above average temperature as

$$< T > = 1.36 T_{\mathrm{sub}}. \tag{B.12}$$

Therefore, the measured excitation temperature of a species would be 1.36 times the sublimation temperature of that molecule.

Appendix C: Spectral fitting results

The fitted models are presented in Figs. C.1 to C.25 for B1-c and S68N. To derive the most accurate excitation temperatures, the χ^2 plot for each molecule is examined. Figure C.28 shows the χ^2 plot for one of the molecules (C_2H_5CN) where the χ^2 calculation is most constraining. When the χ^2 method was not constraining or the results from fitting the spectrum by eye were not consistent with the χ^2 results, the spectrum was fitted by eye. The fit by eye was done by fixing the column density to the best fitted value or changing it slightly around the best fitted value while varying the temperature in steps of 10 K to find the best fit. Errors of ~50 K to ~100 K were typically found. All the excitation temperatures that were not fixed were fitted by eye due to line blending. Figures C.1-C.14 and C.15-C.25 present the range of fitted models for B1-c and S68N, including the upper and lower limit temperature fit for each molecule where the fit-by-eye method was used (NH_2CHO, C_2H_5CN, $HN^{13}CO$, CH_2DCN, and CHD_2CN towards B1-c and C_2H_5CN and CH_2DCN towards S68N).

In general, the model, as expected, does not predict a line for transitions with low A_{ij}. This can be seen for the NH_2CHO line with an E_{up} of 23.6 K and an A_{ij} of 1.2×10^{-5} (Figs. C.6 and C.7). Moreover, the line with an E_{up} of 478.0 K has a large error in its frequency in the JPL catalogue and hence is not used in the fit.

In the specific case of C_2H_5CN towards B1-c (Figs. C.4 and C.5), the two lines with upper energy levels of 153.7 K and 233.6 K have the same frequencies and plausible A_{ij} of

4.99×10^{-4} and 4.07×10^{-4}, respectively. Hence it is not possible to say which of them is responsible for the emission line. We produced two models with the best-fit parameters for this molecule, once limiting the E_{up} of lines included in the model to below 160 K and once to above this value to make sure that these two models do not include the two lines at the same time. We found that the two models show very similar line peak intensities, and thus the final best-fit model most likely includes emission from both of these lines.

Appendix D: Additional tables
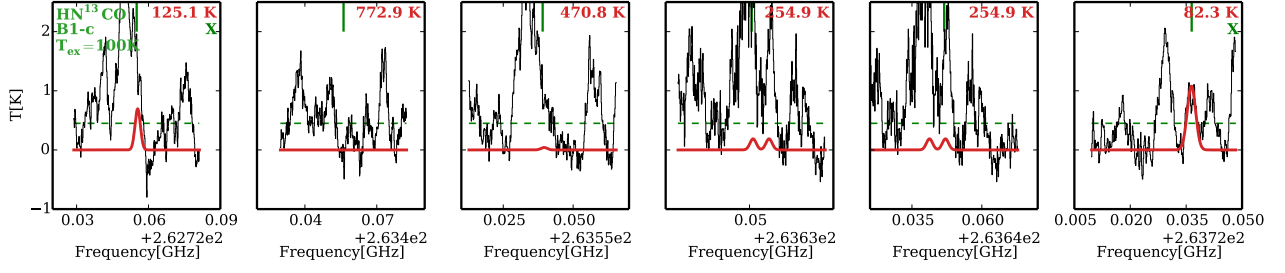
Fig. C.1. Model to combined Band 5 and 6 data of HN^{13}CO for B1-c in red and data in black. The model uses the lower limit on the excitation temperature of HN^{13}CO. Each graph shows one line of HN^{13}CO, indicated by the solid green line at the top of the box, along with its upper energy level, written in red at the top right. The lines with upper energy levels above 1000 K and/or A_{ij} below 10^{-5} are not plotted. The excitation temperature used for the figure is shown at the top left. The dashed green line shows the 3σ level. Cases where a line is seen at the 3σ level or above and is used as part of the fitting are marked with a green X at the top right corner of the box.
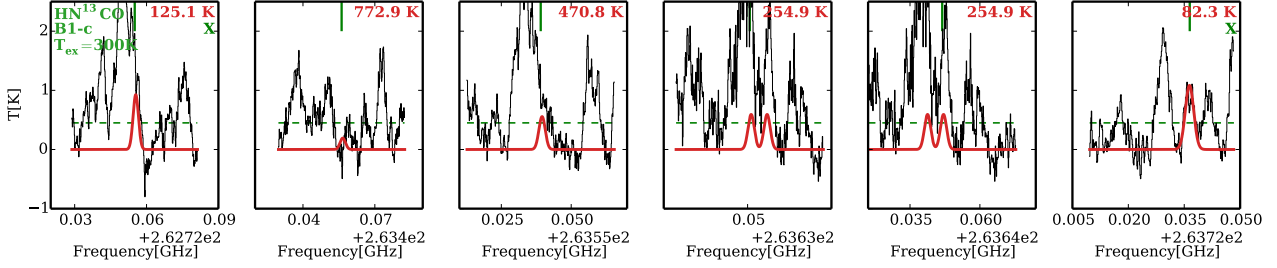
Fig. C.2. Same as Fig. C.1 but for its upper limit on temperature.

Table D.1. Observational parameters of the data from observing programme 2017.1.01174.S.

Parameters	B1-c		S68N	
	Band 3	Band 6	Band 3	Band 6
Configuration	C43-3	C43-4	C43-2	C43-4
Beam ($''$)	2.1×1.4	0.58×0.39	2.8×1.8	0.46×0.42
LAS ($''$)	16	6	22	6
ΔV (km s^{-1})	0.2 / 0.4	0.14	0.17 / 0.32	0.14
Continuum rms (mJy)	14.4	24.4	14.6	9.2
Line rms (K)	0.06	0.14	0.04	0.15

Table D.2. Frequency ranges covered in the Band 5 and 6 data.

ALMA Band	Frequency range (GHz)
Band 5:	~172.65~172.71
	~173.48~173.54
	~173.65~174.58
	~186.55~188.43
Band 6:	~259.98~260.45
	~261.71~262.18
	~262.70~263.19
	~263.30~263.77
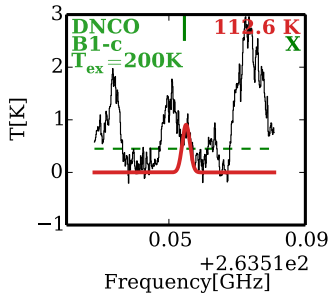
Fig. C.3. Same as Fig. C.1 but for DNCO and the best-fit model where the temperature is fixed.
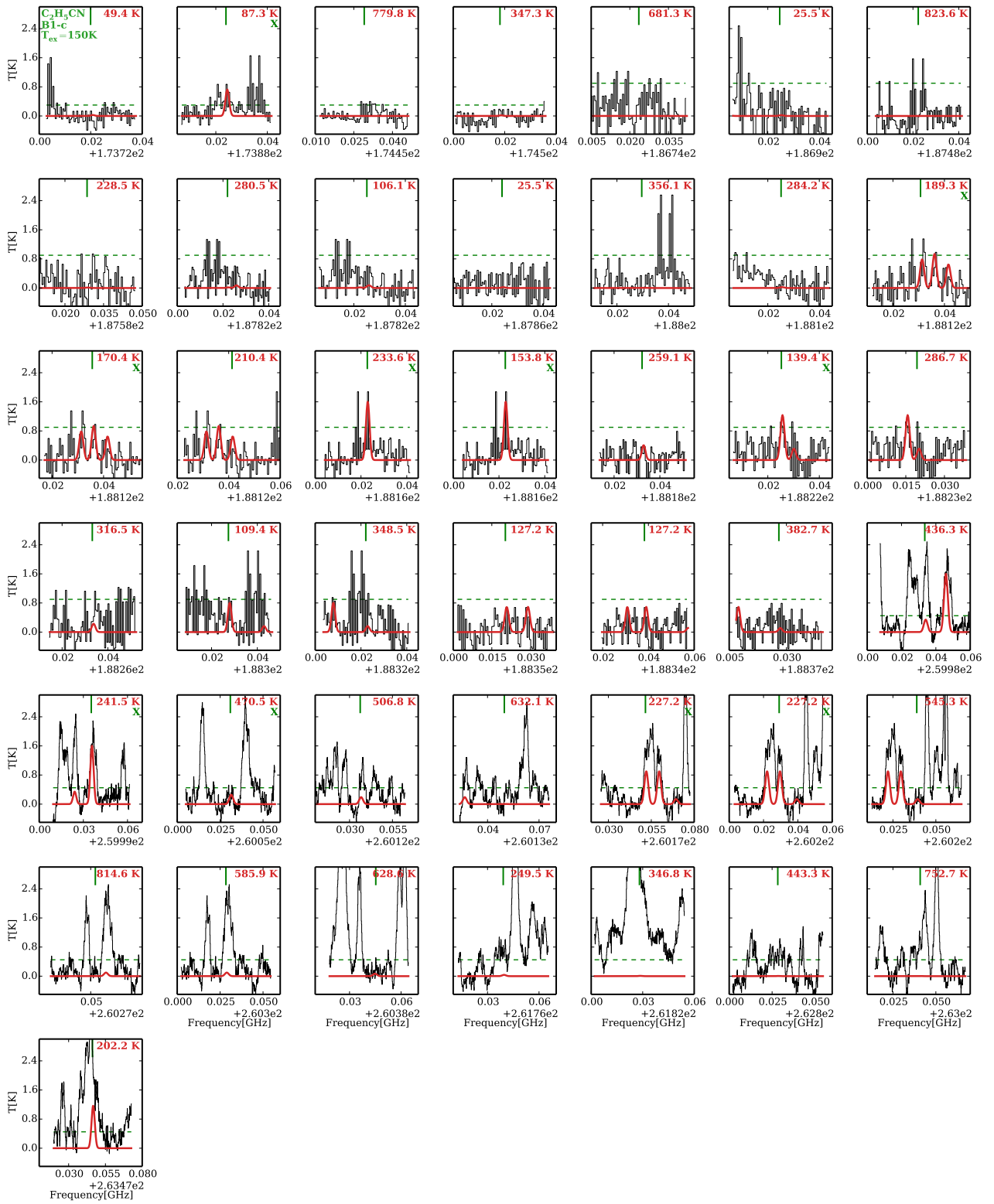
Fig. C.4. Same as Fig. C.1 but for C_2H_5CN.
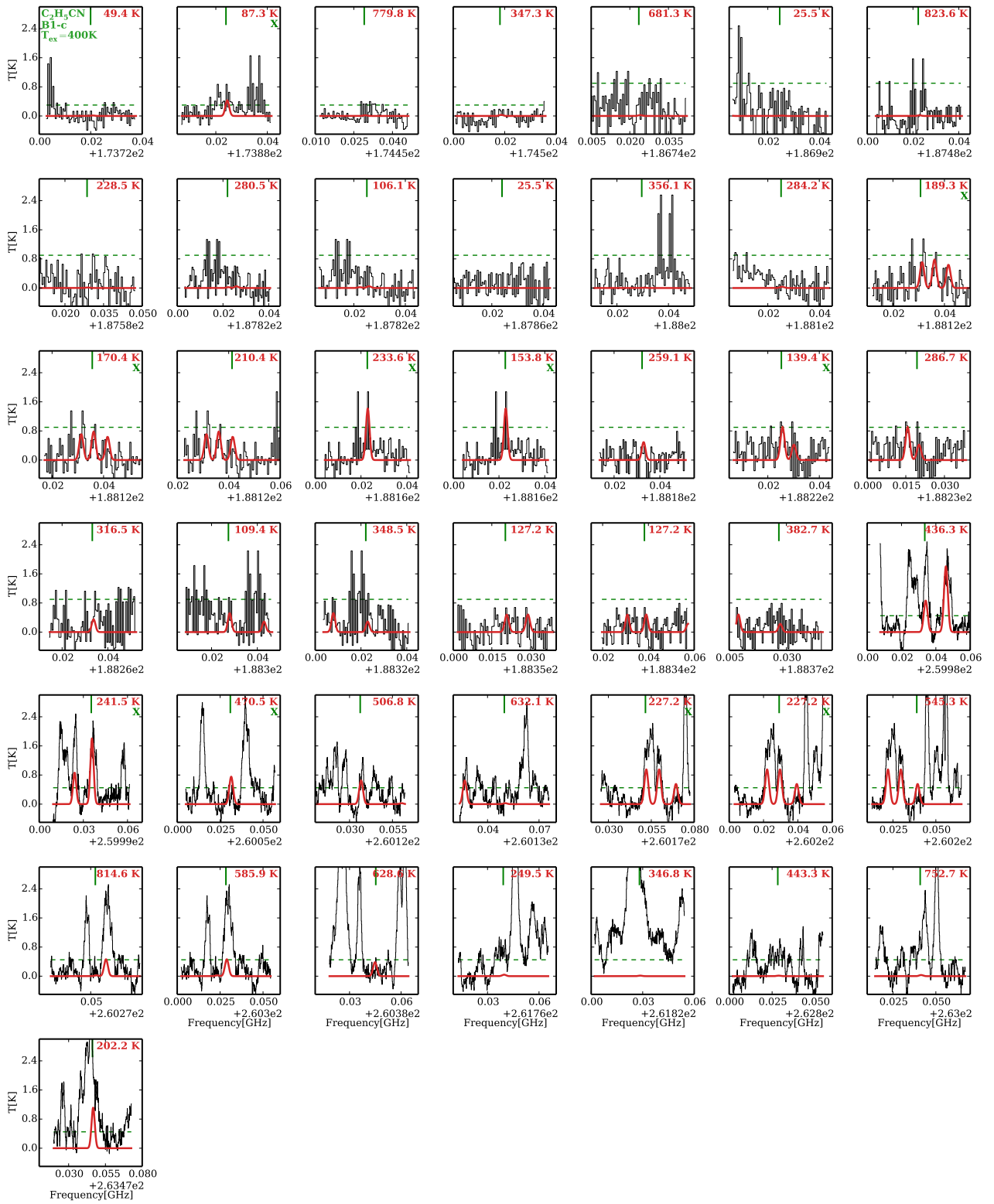
Fig. C.5. Same as Fig. C.1 but for C_2H_5CN and for its upper limit on temperature.
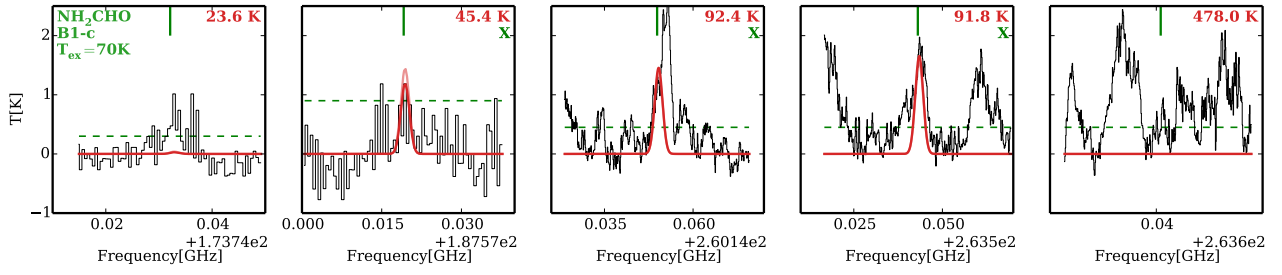
Fig. C.6. Same as Fig. C.1 but for NH_2CHO. The more transparent red line shows the model with a T_{ex} of 50 K to show that it overestimates the line with an E_{up} of 45.4 K.
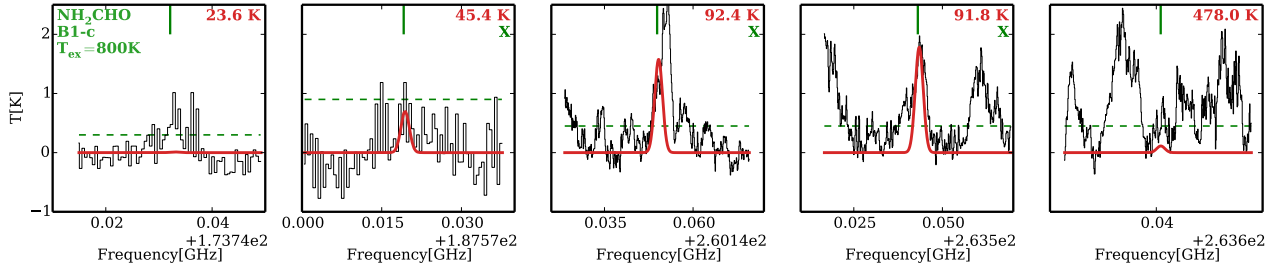
Fig. C.7. Same as Fig. C.1 but for NH_2CHO and a high temperature for the model. This graph is presented to demonstrate that it is not possible to derive an upper limit for the excitation temperature of this molecule because the model used in this plot is very similar to the model in Fig. C.6.
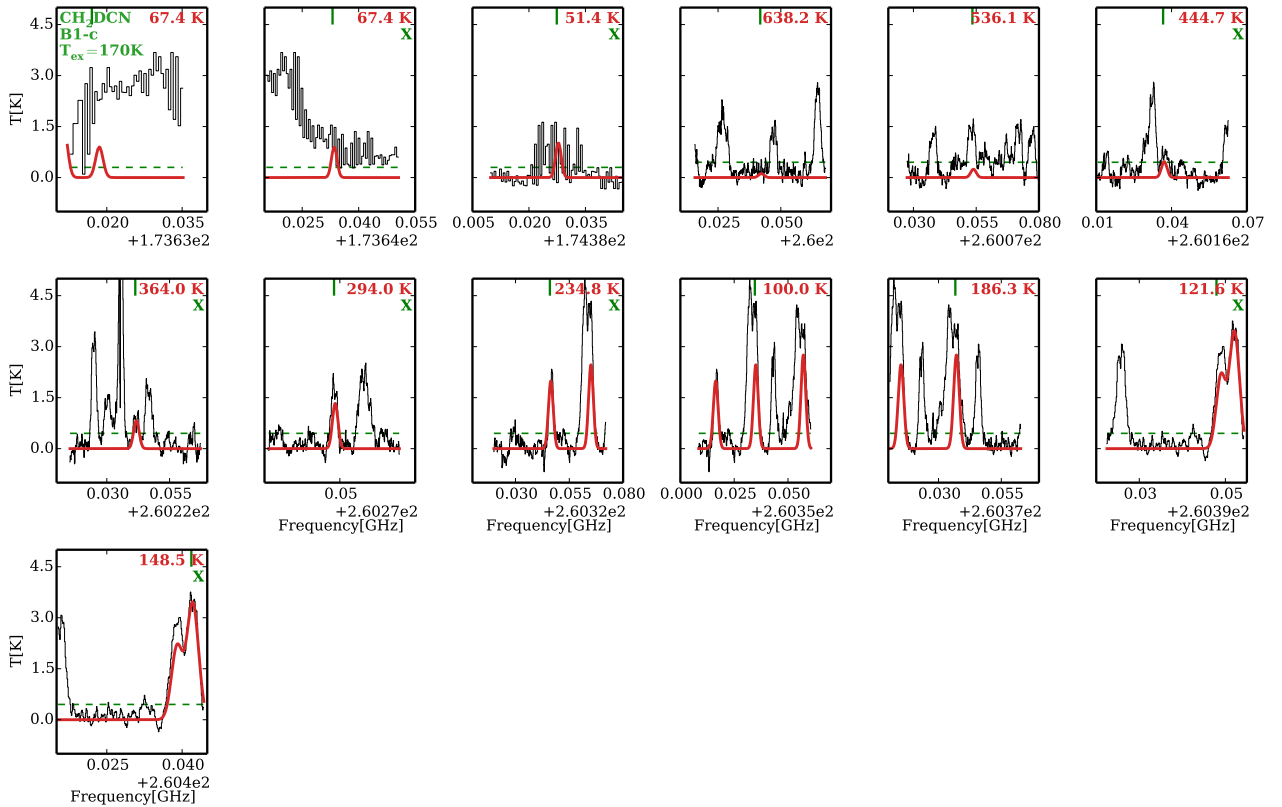
Fig. C.8. Same as Fig. C.1 but for CH_2DCN.
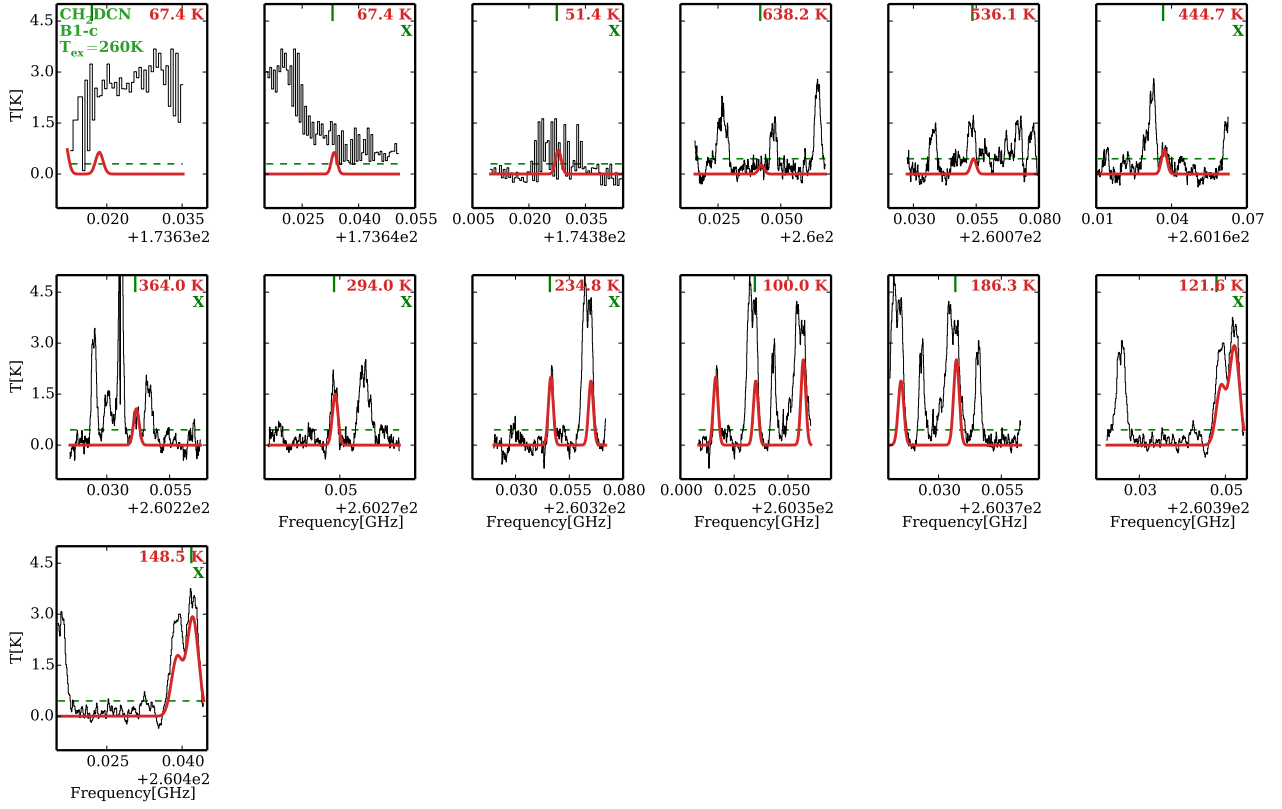
Fig. C.9. Same as Fig. C.1 but for CH$_2$DCN and for its upper limit on temperature.
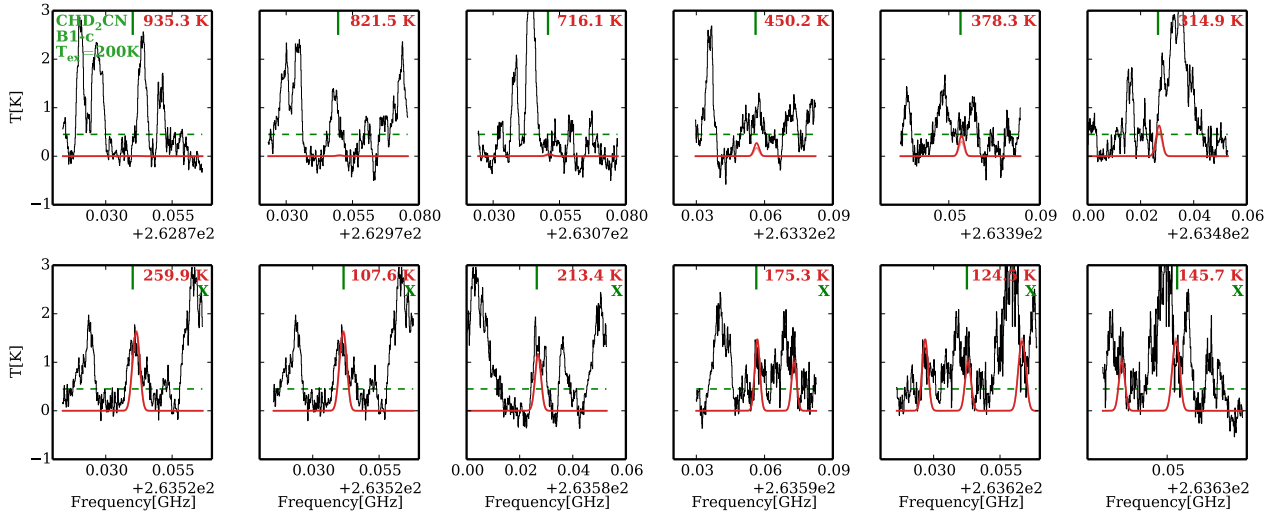
Fig. C.10. Same as Fig. C.1 but for CHD$_2$CN.
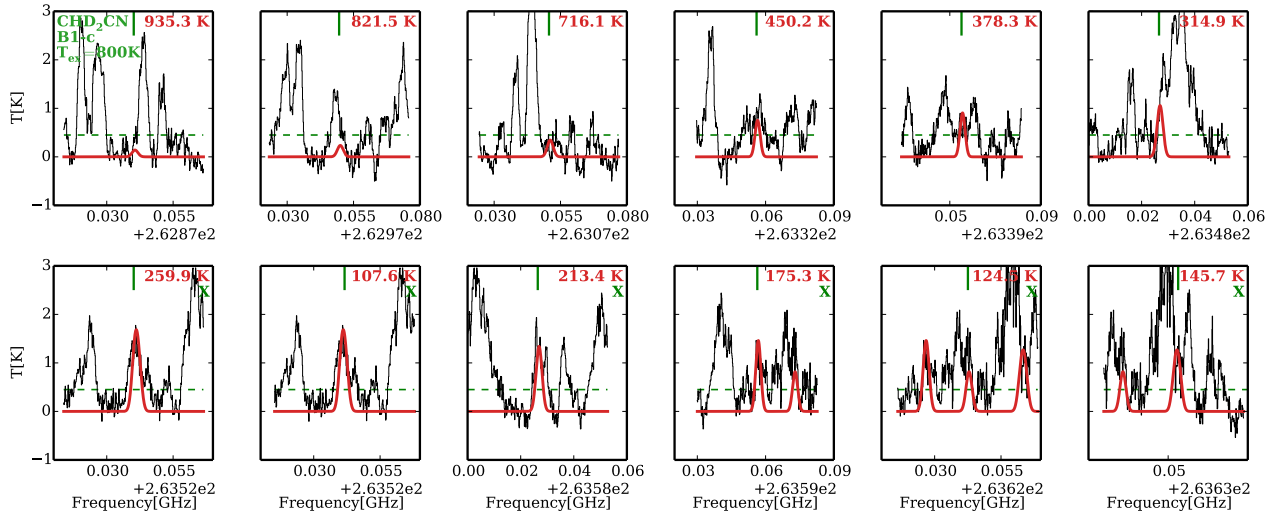
Fig. C.11. Same as Fig. C.1 but for CHD$_2$CN and a high temperature for the model. This graph is presented to demonstrate that it is not possible to derive an upper limit for the excitation temperature of this molecule because the model used in this plot is very similar to the model in Fig. C.10.
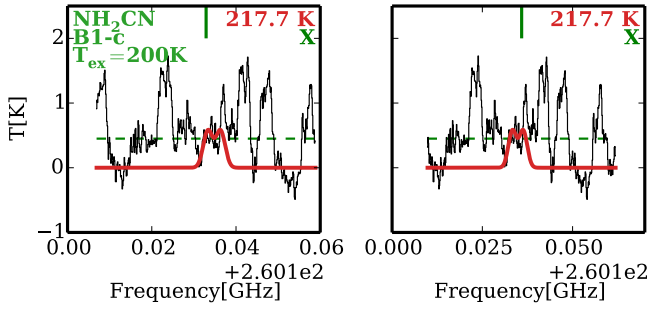
Fig. C.12. Same as Fig. C.1 but for NH$_2$CN and the best-fit model where the temperature is fixed.

Table D.3. Characteristics of the present lines towards B1-c.

Species	$E_{\rm up}$ (K) range[a]	Transitions[b]	FWHM (km/s)[c]
HNCO	82.3-732.9	5	–
HN^{13}CO	82.3-772.9	6	3.0
DNCO	112.6-112.6	1	3.0
NH$_2$CHO	23.6-478.0	6	3.0
C$_2$H$_5$CN	25.5-939.2	65	3.0
CH$_3$CN	–	–	3.0
CH$_2$DCN	51.4-638.2	15	3.0
CHD$_2$CN	28.8-976.3	22	3.0
NH$_2$CN	217.7-217.7	2	3.0
CH$_3$NCO	111.3-395.9	18	3.0
CH$_3$NH$_2$	48.6-843.5	26	3.0

Notes. [a] Upper energy level ranges below 1000 K for the lines present in the less noisy spectral windows of the Band 5 and 6 data. [b] Number of transitions with upper energy levels below 1000 K in the less noisy spectral windows of the Band 5 and 6 data. [c] The FWHM used in the final best-fit model for each molecule.

Table D.4. Characteristics of the present lines towards S68N.

Species	$E_{\rm up}$ (K) range[a]	Transitions[b]	FWHM (km/s)[c]
HNCO	82.3-732.9	5	–
HN^{13}CO	82.3-772.9	6	4.5
DNCO	112.6-112.6	1	3.0
NH$_2$CHO	91.8-478.0	3	3.0
C$_2$H$_5$CN	56.2-814.6	22	4.5
CH$_3$CN	–	–	3.0
CH$_2$DCN	100-638.2	10	4.5
CHD$_2$CN	28.8-935.3	17	3.0
NH$_2$CN	217.7-217.7	2	3.0
CH$_3$NCO	201.1-395.9	7	3.0
CH$_3$NH$_2$	48.6-843.5	10	3.0

Notes. [a] Upper energy level ranges below 1000 K for the lines present in the Band 6 data. [b] Number of transitions with upper energy levels below 1000 K in the Band 6 data. [c] The FWHM used in the final best-fit model for each molecule.
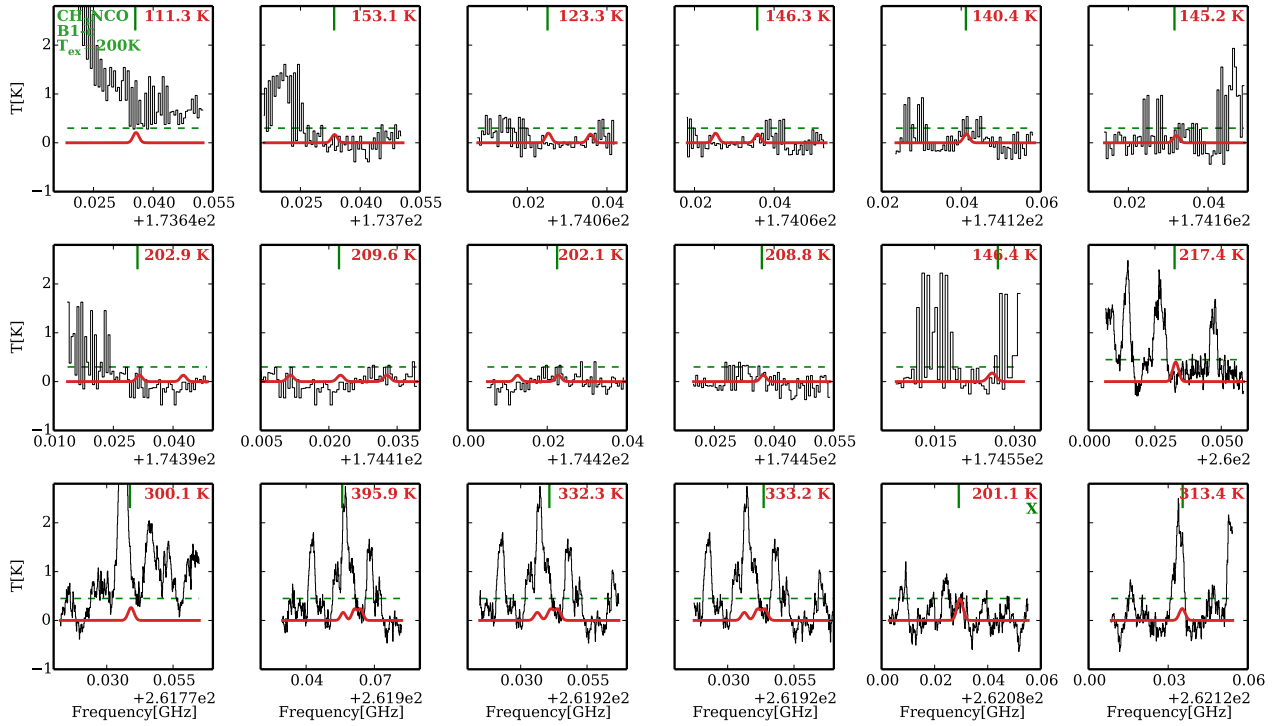
Fig. C.13. Same as Fig. C.1 but for CH₃NCO and the best-fit model where the temperature is fixed.
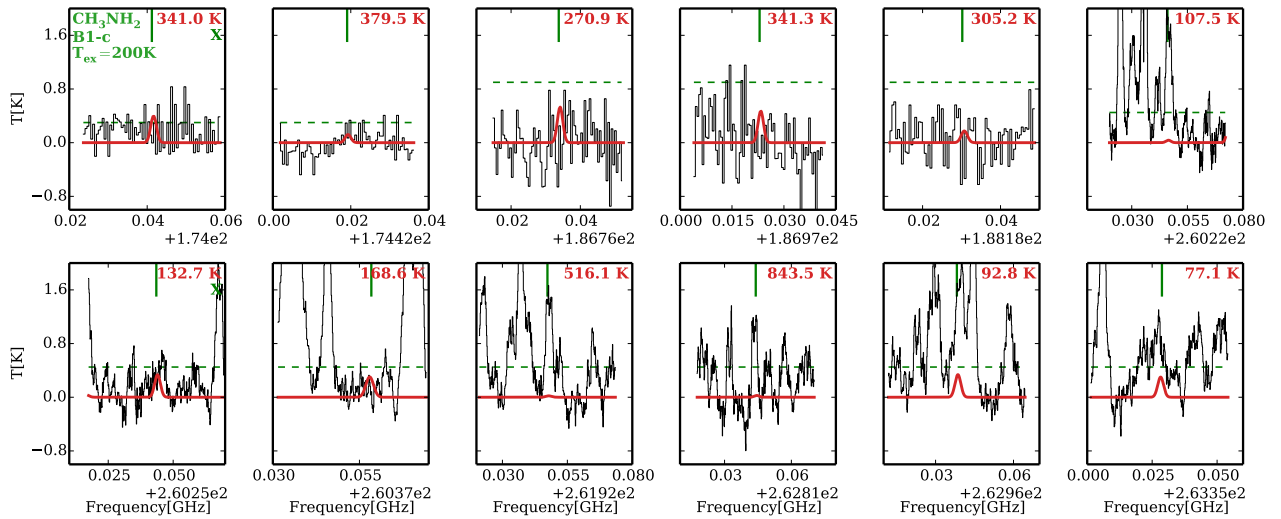
Fig. C.14. Same as Fig. C.1 but for CH₃NH₂ and the best-fit model where the temperature is fixed.
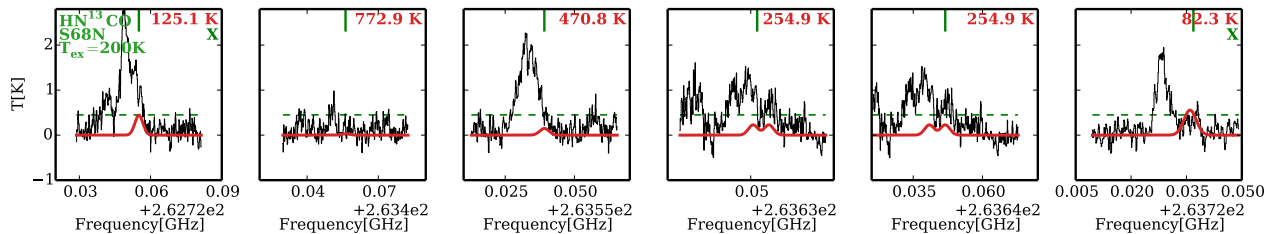
Fig. C.15. Model to Band 6 data of HN¹³CO for S68N in red and data in black. The model uses an excitation temperature that is fixed for HN¹³CO. Each graph shows one line of HN¹³CO, indicated by the solid green line at the top middle along with its upper energy level at the top right. The lines with upper energy levels above 1000 K and/or A_{ij} below 10^{-5} are not plotted. The excitation temperature used for this figure is shown at the top left. The dashed green line shows the 3σ level. Cases where a line is seen at the 3σ level or above and is used as part of the fitting are marked with a green X at the top right corner of the box.
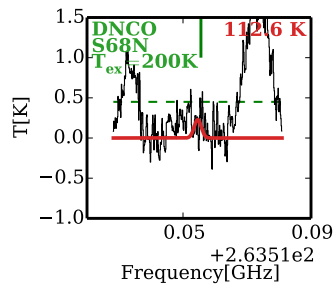
Fig. C.16. Same as Fig. C.15 but for DNCO.
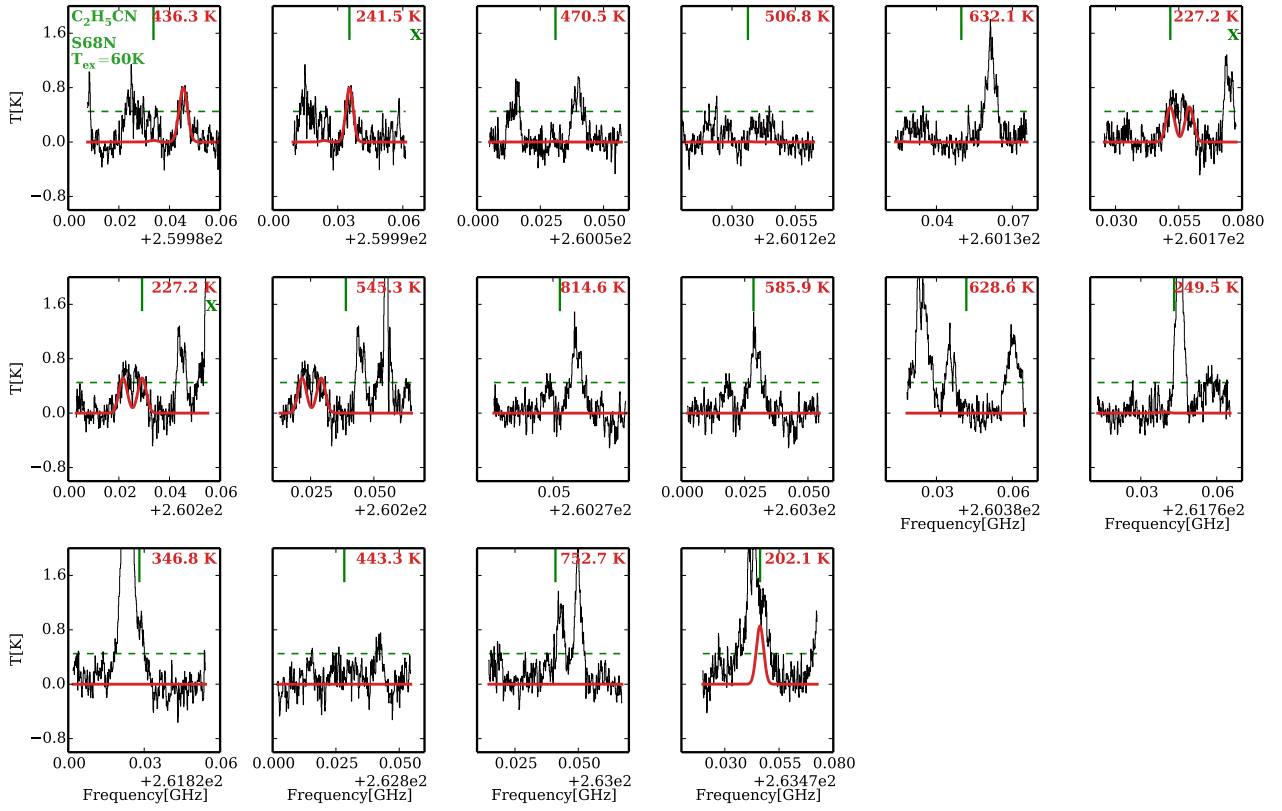
Fig. C.17. Same as Fig. C.15 but for C_2H_5CN and the excitation temperature used is the lower limit for C_2H_5CN.
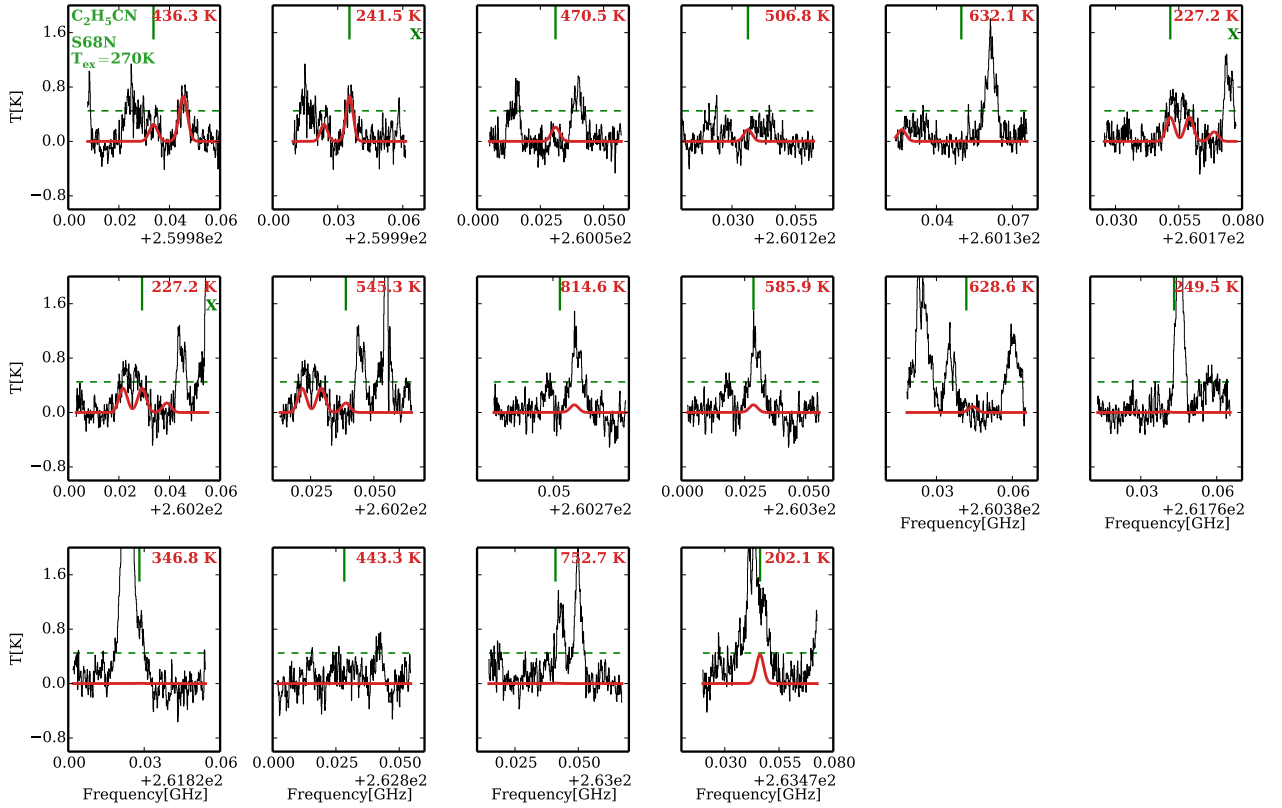
Fig. C.18. Same as Fig. C.15 but for C_2H_5CN and the excitation temperature used is the upper limit for C_2H_5CN.
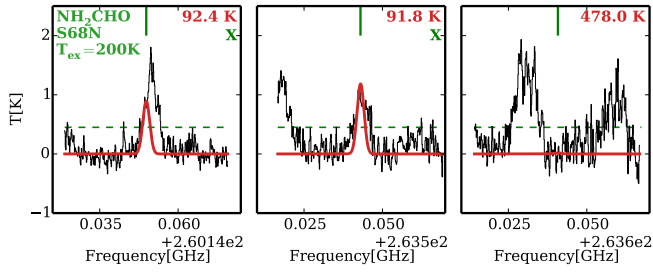
Fig. C.19. Same as Fig. C.15 but for NH$_2$CHO.
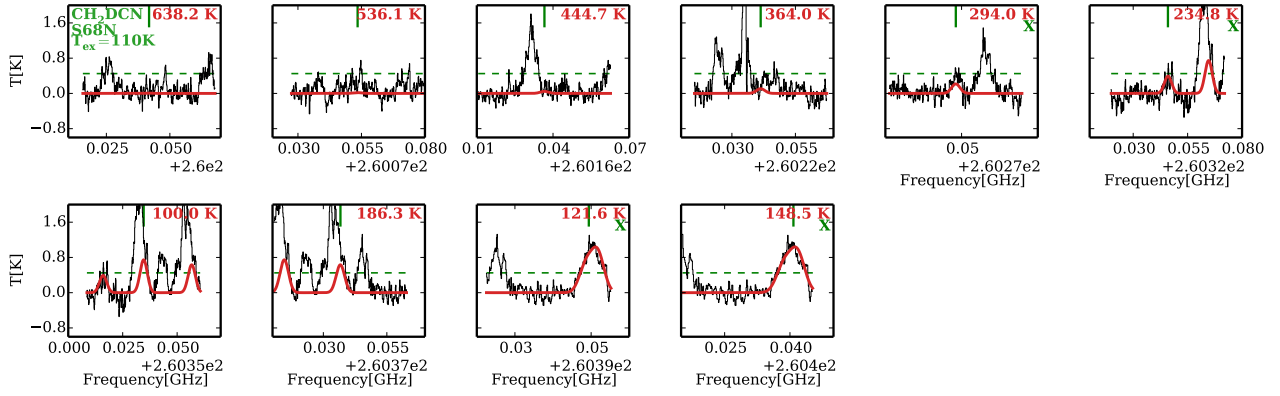
Fig. C.20. Same as Fig. C.15 but for CH$_2$DCN and the excitation temperature used is the lower limit for CH$_2$DCN.
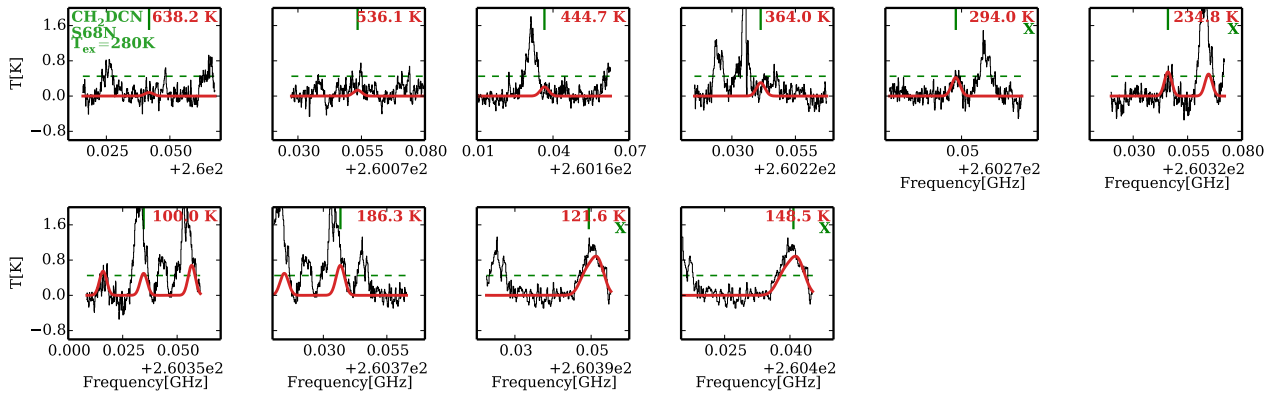
Fig. C.21. Same as Fig. C.15 but for CH$_2$DCN and the excitation temperature used is the upper limit for CH$_2$DCN.
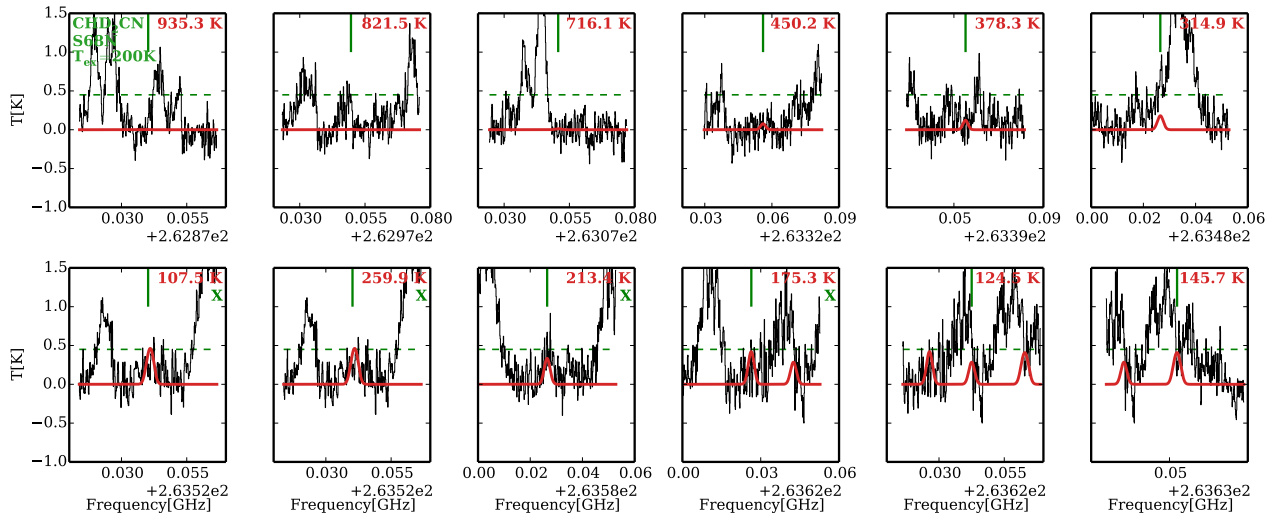
Fig. C.22. Same as Fig. C.15 but for CHD$_2$CN.
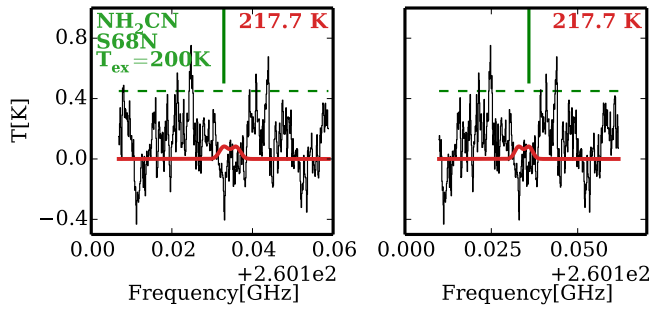
Fig. C.23. Same as Fig. C.15 but for NH$_2$CN.
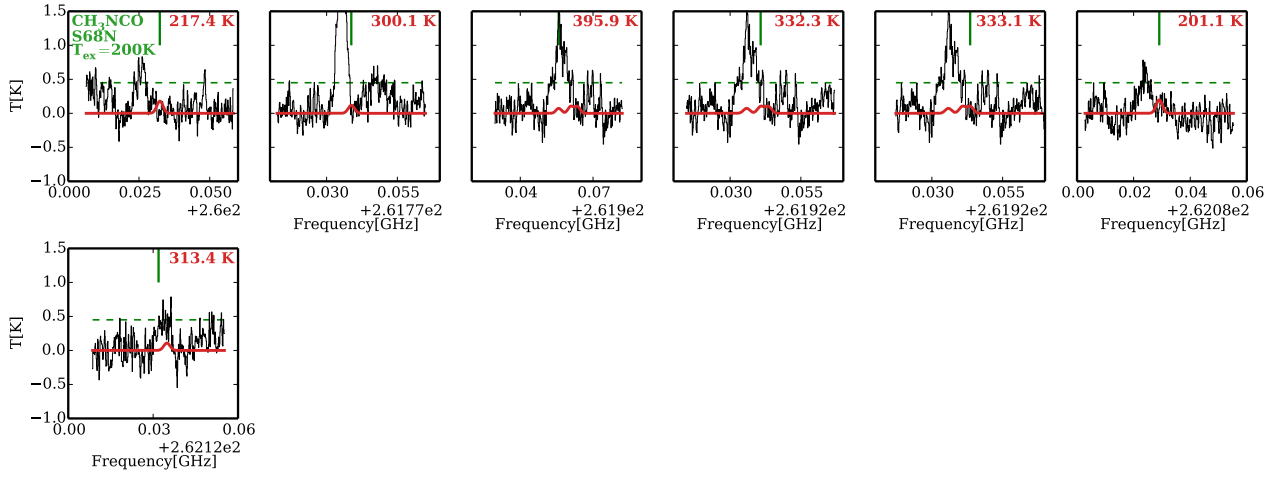
Fig. C.24. Same as Fig. C.15 but for CH$_3$NCO.
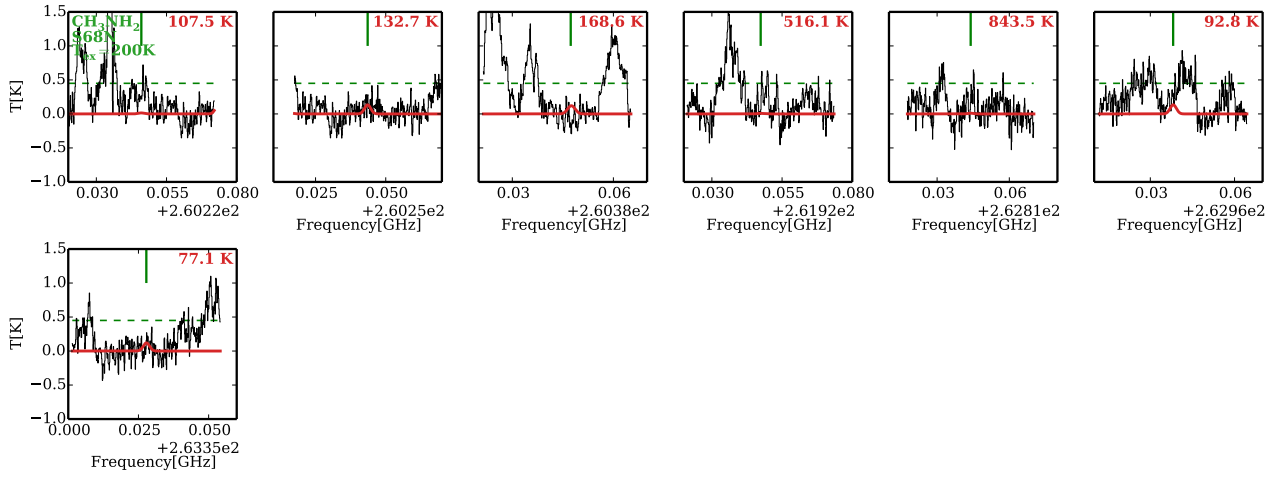
Fig. C.25. Same as Fig. C.15 but for CH$_3$NH$_2$.
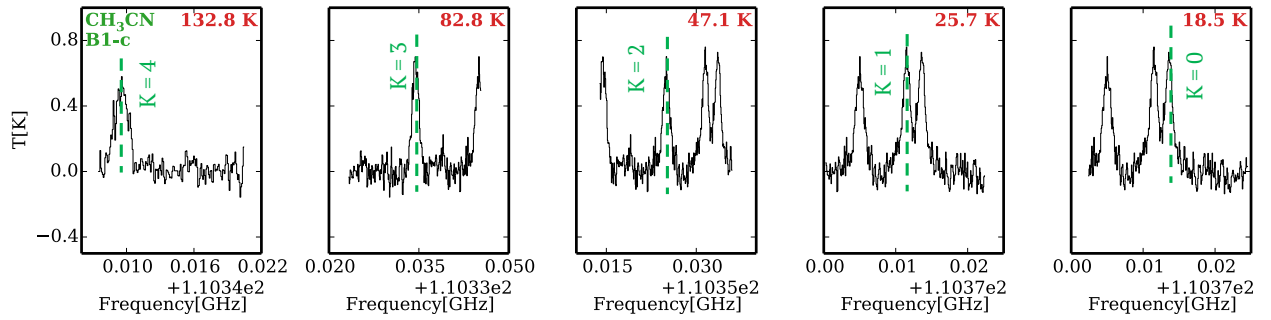
Fig. C.26. Spectrum for CH$_3$CN towards B1-c in the Band 3 data. The upper energy levels of each line are shown at the top right of each panel.
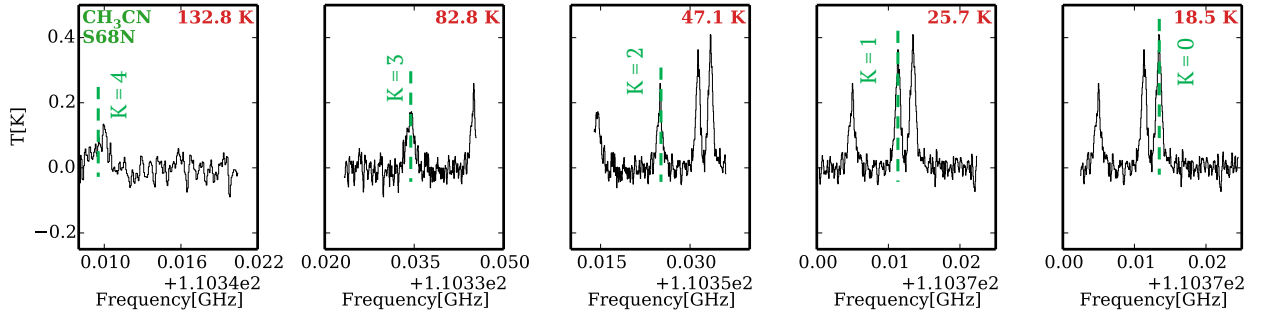
Fig. C.27. Spectrum for CH$_3$CN towards S68N in the Band 3 data. The upper energy levels of each line are shown at the top right of each panel.
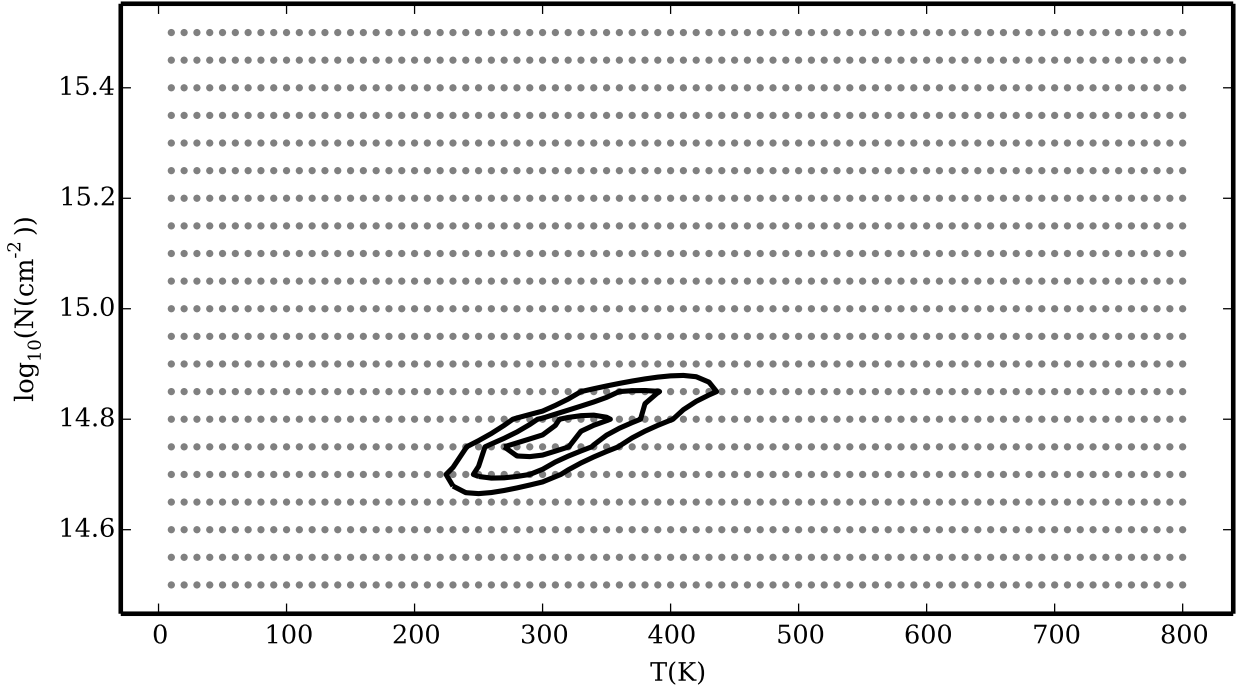
Fig. C.28. χ^2 plot for C$_2$H$_5$CN for B1-c in the Band 5 and 6 data. The dots in the background show the grid used for this molecule, and the contours show the 3σ, 2σ, and σ levels of the best fit. It is seen that T$_{ex}$ is constrained to between ~200 K and ~400 K, whereas the column density is well determined at $(5.8 \pm 1.7) \times 10^{14}$. We note that the vibrational correction factor is not taken into account for this plot.

Table D.3. Lines of the species in the Band 3 data.

Species	Transition J K L (M)	Frequency (MHz)	A_{ij} (s^{-1})	E_{up} (K)
HNCO	5 4 1 - 4 4 0	109 778.752	5.20×10^{-6}	666.6
	5 4 2 - 4 4 1	109 778.752	5.20×10^{-6}	666.6
	5 0 5 - 4 0 4	109 905.749	1.75×10^{-5}	15.8
CH$_3$CN	6 4 0 - 5 4 0	110 349.471	6.17×10^{-5}	132.8
	6 3 0 - 5 -3 0	110 364.354	8.33×10^{-5}	82.8
	6 -3 0 - 5 3 0	110 364.354	8.33×10^{-5}	82.8
	6 2 0 - 5 2 0	110 374.989	9.88×10^{-5}	47.1
	6 1 0 - 5 1 0	110 381.372	1.08×10^{-4}	25.7
	6 0 0 - 5 0 0	110 383.5	1.11×10^{-4}	18.5

Table D.3. Lines of the species in the Band 5 and 6 data.

Species	Transition J K L (M)	Frequency (MHz)	A_{ij} (s^{-1})	E_{up} (K)
HNCO	12 1 12 - 11 1 11	262 769.477	2.48×10^{-4}	125.3
	12 4 8 - 11 4 7	263 449.24	1.87×10^{-4}	732.9
	12 4 9 - 11 4 8	263 449.24	1.87×10^{-4}	732.9
	12 3 10 - 11 3 9	263 580.924	2.16×10^{-4}	457.2
	12 3 9 - 11 3 8	263 580.928	2.16×10^{-4}	457.2
	12 2 11 - 11 2 10	263 672.912	2.37×10^{-4}	252.5
	12 2 10 - 11 2 9	263 678.709	2.37×10^{-4}	252.5
	12 0 12 - 11 0 11	263 748.625	2.56×10^{-4}	82.3
HN^{13}CO	12 1 12 11 - 11 1 11 11	262 774.075	1.95×10^{-6}	125.1
	12 1 12 13 - 11 1 11 12	262 775.057	2.58×10^{-4}	125.1
	12 1 12 12 - 11 1 11 11	262 775.062	2.57×10^{-4}	125.1
	12 1 12 11 - 11 1 11 10	262 775.064	2.56×10^{-4}	125.1
	12 1 12 12 - 11 1 11 12	262 775.968	1.79×10^{-6}	125.1
	12 4 8 11 - 11 4 7 11	263 455.652	1.76×10^{-6}	772.9
	12 4 9 11 - 11 4 8 11	263 455.652	1.76×10^{-6}	772.9
	12 4 9 11 - 11 4 8 10	263 456.155	2.31×10^{-4}	772.9
	12 4 8 11 - 11 4 7 10	263 456.155	2.31×10^{-4}	772.9
	12 4 9 13 - 11 4 8 12	263 456.155	2.33×10^{-4}	772.9
	12 4 8 13 - 11 4 7 12	263 456.155	2.33×10^{-4}	772.9
	12 4 8 12 - 11 4 7 11	263 456.198	2.31×10^{-4}	772.9
	12 4 9 12 - 11 4 8 11	263 456.198	2.31×10^{-4}	772.9
	12 4 8 12 - 11 4 7 12	263 456.659	1.62×10^{-6}	772.9
	12 4 9 12 - 11 4 8 12	263 456.659	1.62×10^{-6}	772.9
	12 3 10 11 - 11 3 9 11	263 588.502	1.86×10^{-6}	470.8
	12 3 9 11 - 11 3 8 11	263 588.505	1.86×10^{-6}	470.8
	12 3 10 13 - 11 3 9 12	263 589.13	2.46×10^{-4}	470.8
	12 3 10 11 - 11 3 9 10	263 589.132	2.44×10^{-4}	470.8
	12 3 9 13 - 11 3 8 12	263 589.132	2.46×10^{-4}	470.8
	12 3 9 11 - 11 3 8 10	263 589.135	2.44×10^{-4}	470.8
	12 3 10 12 - 11 3 9 11	263 589.154	2.44×10^{-4}	470.8
	12 3 9 12 - 11 3 8 11	263 589.157	2.44×10^{-4}	470.8
	12 3 10 12 - 11 3 9 12	263 589.731	1.71×10^{-6}	470.8
	12 3 9 12 - 11 3 8 12	263 589.734	1.71×10^{-6}	470.8
	12 2 11 11 - 11 2 10 11	263 680.119	1.93×10^{-6}	254.9
	12 2 11 13 - 11 2 10 12	263 680.835	2.56×10^{-4}	254.9
	12 2 11 11 - 11 2 10 10	263 680.839	2.54×10^{-4}	254.9
	12 2 11 12 - 11 2 10 11	263 680.847	2.54×10^{-4}	254.9
	12 2 11 12 - 11 2 10 12	263 681.508	1.77×10^{-6}	254.9
	12 2 10 11 - 11 2 9 11	263 685.825	1.93×10^{-6}	254.9
	12 2 10 13 - 11 2 9 12	263 686.54	2.56×10^{-4}	254.9
	12 2 10 11 - 11 2 9 10	263 686.544	2.54×10^{-4}	254.9
	12 2 10 12 - 11 2 9 11	263 686.552	2.54×10^{-4}	254.9

Table D.3. continued.

Species	Transition J K L (M)	Frequency (GHz)	A_{ij} (s^{-1})	$E_{\rm up}$ (K)
	12 2 10 12 - 11 2 9 12	263 687.21	1.77×10^{-6}	254.9
	12 0 12 11 - 11 0 11 11	263 755.182	1.99×10^{-6}	82.3
	12 0 12 11 - 11 0 11 12	263 755.909	3.45×10^{-9}	82.3
	12 0 12 13 - 11 0 11 12	263 755.97	2.63×10^{-4}	82.3
	12 0 12 12 - 11 0 11 11	263 755.972	2.61×10^{-4}	82.3
	12 0 12 11 - 11 0 11 10	263 755.975	2.61×10^{-4}	82.3
	12 0 12 12 - 11 0 11 12	263 756.699	1.83×10^{-6}	82.3
DNCO	13 1 13 12 - 12 1 12 12	263 563.661	1.67×10^{-6}	112.6
	13 1 13 14 - 12 1 12 13	263 564.594	2.62×10^{-4}	112.6
	13 1 13 13 - 12 1 12 12	263 564.598	2.60×10^{-4}	112.6
	13 1 13 12 - 12 1 12 11	263 564.599	2.60×10^{-4}	112.6
	13 1 13 13 - 12 1 12 13	263 565.464	1.55×10^{-6}	112.6
NH$_2$CHO	6 1 6 5 - 5 0 5 4	173 772.286	1.22×10^{-5}	23.6
	6 1 6 7 - 5 0 5 6	173 772.398	1.27×10^{-5}	23.6
	6 1 6 6 - 5 0 5 5	173 773.239	1.23×10^{-5}	23.6
	19 2 17 20 - 18 3 16 19	174 482.455	4.46×10^{-6}	209.7
	9 0 9 10 - 8 0 8 9	187 589.055	4.75×10^{-4}	45.4
	9 0 9 8 - 8 0 8 7	187 589.059	4.68×10^{-4}	45.4
	9 0 9 9 - 8 0 8 8	187 589.158	4.69×10^{-4}	45.4
	12 2 10 - 11 2 9	260 189.848	1.25×10^{-3}	92.4
	13 1 13 - 12 1 12	263 543.025	1.33×10^{-3}	91.8
	29 3 26 29 - 29 2 27 29	263 640.492	6.36×10^{-5}	478.0
	29 3 26 30 - 29 2 27 30	263 640.985	6.37×10^{-5}	478.0
	29 3 26 28 - 29 2 27 28	263 641.002	6.37×10^{-5}	478.0
C$_2$H$_5$CN	67 4 63 - 66 6 60	173 479.725	3.59×10^{-7}	1006.9
	68 10 58 - 69 8 61	173 481.563	1.22×10^{-7}	1119.8
	53 3 51 - 54 0 54	173 496.547	8.35×10^{-8}	619.7
	44 13 31 - 45 12 34	173 505.22	6.02×10^{-6}	613.1
	44 13 32 - 45 12 33	173 505.22	6.02×10^{-6}	613.1
	10 3 8 - 10 1 9	173 739.951	1.99×10^{-7}	33.7
	14 2 13 - 13 1 12	173 740.541	1.64×10^{-5}	49.4
	19 2 17 - 18 2 16	173 904.151	4.37×10^{-4}	87.3
	8 6 3 - 9 5 4	174 114.528	1.84×10^{-6}	55.5
	8 6 2 - 9 5 5	174 114.533	1.84×10^{-6}	55.5
	57 4 53 - 56 6 50	174 241.737	5.62×10^{-7}	738.0
	58 6 52 - 58 5 53	174 479.058	3.53×10^{-5}	779.8
	38 5 33 - 38 4 34	174 518.373	2.90×10^{-5}	347.4
	64 4 60 - 63 6 57	186 317.795	5.52×10^{-7}	922.2
	38 7 31 - 39 5 34	186 652.738	1.10×10^{-7}	373.3
	54 6 48 - 54 5 49	186 763.733	4.01×10^{-5}	681.3
	22 9 14 - 23 8 15	186 905.217	5.85×10^{-6}	198.8
	22 9 13 - 23 8 16	186 905.217	5.85×10^{-6}	198.8
	8 3 6 - 7 2 5	186 924.942	2.24×10^{-5}	25.5
	60 4 56 - 59 6 53	187 192.928	6.78×10^{-7}	814.6
	58 7 52 - 57 8 49	187 379.431	9.51×10^{-6}	791.6
	32 1 31 - 32 1 32	187 437.247	2.18×10^{-6}	228.5
	60 5 55 - 60 4 56	187 502.511	3.90×10^{-5}	823.6
	32 1 31 - 32 0 32	187 608.572	1.25×10^{-5}	228.5
	56 2 54 - 57 1 57	187 643.268	9.43×10^{-8}	689.0
	10 2 8 - 9 0 9	187 663.517	6.94×10^{-7}	28.2
	35 3 33 - 35 2 34	187 842.025	2.26×10^{-5}	280.5
	22 0 22 - 21 1 21	187 845.577	4.68×10^{-5}	106.1
	90 13 78 - 89 14 75	187 874.315	9.72×10^{-6}	1945.3
	90 13 77 - 89 14 76	187 880.086	9.72×10^{-6}	1945.3
	8 3 5 - 7 2 6	187 884.105	2.26×10^{-5}	25.5
	63 8 55 - 62 9 54	187 931.933	9.70×10^{-6}	939.2

Table D.3. continued.

Species	Transition J K L (M)	Frequency (GHz)	A_{ij} (s^{-1})	E_{up} (K)
	39 4 35 - 38 5 34	188 030.172	1.01×10^{-5}	356.1
	27 10 18 - 28 9 19	188 061.915	6.51×10^{-6}	273.6
	27 10 17 - 28 9 20	188 061.915	6.51×10^{-6}	273.6
	34 5 29 - 34 4 30	188 125.196	3.39×10^{-5}	284.2
	21 9 12 - 20 9 11	188 151.024	4.58×10^{-4}	189.3
	21 9 13 - 20 9 12	188 151.024	4.58×10^{-4}	189.3
	21 8 13 - 20 8 12	188 155.826	4.80×10^{-4}	170.4
	21 8 14 - 20 8 13	188 155.826	4.80×10^{-4}	170.4
	21 10 11 - 20 10 10	188 161.248	4.34×10^{-4}	210.4
	21 10 12 - 20 10 11	188 161.248	4.34×10^{-4}	210.4
	21 11 10 - 20 11 9	188 182.647	4.07×10^{-4}	233.6
	21 11 11 - 20 11 10	188 182.647	4.07×10^{-4}	233.6
	21 7 15 - 20 7 14	188 182.647	4.99×10^{-4}	153.8
	21 7 14 - 20 7 13	188 182.647	4.99×10^{-4}	153.8
	21 12 9 - 20 12 8	188 212.59	3.78×10^{-4}	259.1
	21 12 10 - 20 12 9	188 212.59	3.78×10^{-4}	259.1
	21 6 16 - 20 6 15	188 245.402	5.16×10^{-4}	139.4
	21 6 15 - 20 6 14	188 245.402	5.16×10^{-4}	139.4
	56 3 54 - 57 0 57	188 247.954	9.52×10^{-8}	689.0
	21 13 8 - 20 13 7	188 249.82	3.47×10^{-4}	286.7
	21 13 9 - 20 13 8	188 249.82	3.47×10^{-4}	286.7
	21 14 7 - 20 14 6	188 293.433	3.13×10^{-4}	316.5
	21 14 8 - 20 14 7	188 293.433	3.13×10^{-4}	316.5
	21 3 19 - 20 3 18	188 327.62	5.51×10^{-4}	109.4
	21 15 6 - 20 15 5	188 342.625	2.76×10^{-4}	348.5
	21 15 7 - 20 15 6	188 342.625	2.76×10^{-4}	348.5
	29 6 24 - 30 4 27	188 354.152	9.34×10^{-8}	227.2
	21 5 17 - 20 5 16	188 370.494	5.31×10^{-4}	127.2
	21 5 16 - 20 5 15	188 378.708	5.31×10^{-4}	127.2
	21 16 5 - 20 16 4	188 397.003	2.36×10^{-4}	382.7
	21 16 6 - 20 16 5	188 397.003	2.36×10^{-4}	382.7
	63 4 59 - 62 6 56	188 398.285	6.07×10^{-7}	894.7
	84 4 80 - 85 3 83	188 416.325	2.91×10^{-7}	1554.4
	26 5 22 - 27 2 25	188 422.769	1.20×10^{-7}	178.8
	68 18 50 - 69 17 53	188 436.437	8.17×10^{-6}	1366.8
	68 18 51 - 69 17 52	188 436.437	8.17×10^{-6}	1366.8
	29 15 14 - 28 15 13	260 013.644	1.09×10^{-3}	436.3
	29 15 15 - 28 15 14	260 013.644	1.09×10^{-3}	436.3
	29 7 23 - 28 7 22	260 025.426	1.40×10^{-3}	241.5
	29 7 22 - 28 7 21	260 025.426	1.40×10^{-3}	241.5
	29 16 13 - 28 16 12	260 081.02	1.04×10^{-3}	470.5
	29 16 14 - 28 16 13	260 081.02	1.04×10^{-3}	470.5
	29 17 12 - 28 17 11	260 156.297	9.80×10^{-4}	506.8
	29 17 13 - 28 17 12	260 156.297	9.80×10^{-4}	506.8
	53 3 50 - 53 2 51	260 179.828	6.18×10^{-5}	632.1
	29 6 24 - 28 6 23	260 221.658	1.43×10^{-3}	227.2
	29 6 23 - 28 6 22	260 229.158	1.43×10^{-3}	227.2
	44 7 38 - 45 4 41	260 232.74	3.29×10^{-7}	480.6
	29 18 11 - 28 18 10	260 238.936	9.19×10^{-4}	545.3
	29 18 12 - 28 18 11	260 238.936	9.19×10^{-4}	545.3
	60 4 56 - 60 3 57	260 322.752	7.43×10^{-5}	814.6
	29 19 10 - 28 19 9	260 328.45	8.54×10^{-4}	585.9
	29 19 11 - 28 19 10	260 328.45	8.54×10^{-4}	585.9
	29 20 9 - 28 20 8	260 424.396	7.86×10^{-4}	628.6
	29 20 10 - 28 20 9	260 424.396	7.86×10^{-4}	628.6
	31 7 24 - 32 5 27	261 749.397	1.87×10^{-7}	267.8
	33 2 31 - 32 3 30	261 799.173	5.81×10^{-5}	249.5

Table D.3. continued.

Species	Transition J K L (M)	Frequency (GHz)	A_{ij} (s^{-1})	$E_{\rm up}$ (K)
	29 12 17 - 30 11 20	261 848.139	1.53×10^{-5}	346.8
	29 12 18 - 30 11 19	261 848.139	1.53×10^{-5}	346.8
	75 21 54 - 76 20 57	261 866.013	2.11×10^{-5}	1711.9
	75 21 55 - 76 20 56	261 866.013	2.11×10^{-5}	1711.9
	53 5 49 - 53 3 50	262 101.305	8.70×10^{-6}	644.7
	34 13 21 - 35 12 24	262 828.348	1.68×10^{-5}	443.3
	34 13 22 - 35 12 23	262 828.348	1.68×10^{-5}	443.3
	87 11 77 - 86 12 74	262 862.696	2.67×10^{-5}	1779.1
	11 5 7 - 12 3 10	262 873.8	3.81×10^{-8}	56.2
	98 13 86 - 97 14 83	262 947.313	2.68×10^{-5}	2268.9
	21 6 15 - 22 4 18	262 953.486	1.16×10^{-7}	139.4
	98 8 90 - 98 8 91	263 004.792	1.02×10^{-5}	2170.5
	80 22 58 - 81 21 61	263 007.909	2.16×10^{-5}	1926.5
	80 22 59 - 81 21 60	263 007.909	2.16×10^{-5}	1926.5
	31 7 25 - 32 5 28	263 026.766	1.89×10^{-7}	267.8
	57 6 52 - 57 5 53	263 040.95	9.13×10^{-5}	752.7
	30 2 29 - 29 2 28	263 516.223	1.54×10^{-3}	202.1
CH$_2$DCN	10 8 2 - 9 8 1	173 480.998	1.60×10^{-4}	390.6
	10 8 3 - 9 8 2	173 480.998	1.60×10^{-4}	390.6
	10 7 3 - 9 7 2	173 524.046	2.27×10^{-4}	309.9
	10 7 4 - 9 7 3	173 524.046	2.27×10^{-4}	309.9
	10 0 10 - 9 0 9	173 638.56	4.46×10^{-4}	45.8
	10 3 8 - 9 3 7	173 641.17	4.06×10^{-4}	94.4
	10 3 7 - 9 3 6	173 641.17	4.06×10^{-4}	94.4
	10 2 9 - 9 2 8	173 648.22	4.28×10^{-4}	67.4
	10 2 8 - 9 2 7	173 673.22	4.28×10^{-4}	67.4
	37 1 36 - 37 0 37	173 837.52	3.54×10^{-7}	593.4
	16 0 16 - 15 1 15	174 131.5	2.20×10^{-7}	113.3
	10 1 9 - 9 1 8	174 407.53	4.47×10^{-4}	51.4
	15 10 5 - 14 10 4	260 041.811	8.46×10^{-4}	638.2
	15 10 6 - 14 10 5	260 041.811	8.46×10^{-4}	638.2
	15 9 6 - 14 9 5	260 123.472	9.75×10^{-4}	536.1
	15 9 7 - 14 9 6	260 123.472	9.75×10^{-4}	536.1
	15 8 7 - 14 8 6	260 196.625	1.09×10^{-3}	444.7
	15 8 8 - 14 8 7	260 196.625	1.09×10^{-3}	444.7
	15 7 8 - 14 7 7	260 261.337	1.19×10^{-3}	364.0
	15 7 9 - 14 7 8	260 261.337	1.19×10^{-3}	364.0
	15 6 9 - 14 6 8	260 317.716	1.28×10^{-3}	294.0
	15 6 10 - 14 6 9	260 317.716	1.28×10^{-3}	294.0
	15 5 10 - 14 5 9	260 365.984	1.36×10^{-3}	234.8
	15 5 11 - 14 5 10	260 365.984	1.36×10^{-3}	234.8
	15 0 15 - 14 0 14	260 384.626	1.53×10^{-3}	100.0
	15 4 12 - 14 4 11	260 406.697	1.42×10^{-3}	186.3
	15 4 11 - 14 4 10	260 406.699	1.42×10^{-3}	186.3
	15 2 14 - 14 2 13	260 438.549	1.50×10^{-3}	121.6
	15 3 13 - 14 3 12	260 441.387	1.47×10^{-3}	148.5
	15 3 12 - 14 3 11	260 441.972	1.47×10^{-3}	148.5
CHD$_2$CN	49 1 48 - 49 1 49	186 743.167	5.09×10^{-7}	976.3
	34 0 34 - 33 2 31	187 059.069	4.89×10^{-8}	469.9
	26 1 25 - 25 2 23	187 083.112	1.29×10^{-7}	283.1
	45 2 44 - 45 1 44	187 826.836	3.12×10^{-7}	834.8
	25 4 21 - 26 3 23	188 372.962	9.57×10^{-8}	324.8
	10 2 9 - 10 1 9	260 444.239	6.53×10^{-7}	60.4
	8 2 7 - 8 1 7	261 931.581	6.55×10^{-7}	45.4
	35 1 35 - 34 2 33	261 979.861	2.60×10^{-7}	499.9
	16 14 2 - 15 14 1	262 910.198	3.69×10^{-4}	935.3
	16 14 3 - 15 14 2	262 910.198	3.69×10^{-4}	935.3

Table D.3. continued.

Species	Transition J K L (M)	Frequency (GHz)	A_{ij} (s^{-1})	$E_{\rm up}$ (K)
	16 13 3 - 15 13 2	263 019.578	5.36×10^{-4}	821.5
	16 13 4 - 15 13 3	263 019.578	5.36×10^{-4}	821.5
	6 2 5 - 6 1 5	263 107.386	6.48×10^{-7}	33.6
	16 12 4 - 15 12 3	263 120.686	6.91×10^{-4}	716.1
	16 12 5 - 15 12 4	263 120.686	6.91×10^{-4}	716.1
	16 9 7 - 15 9 6	263 376.04	1.08×10^{-3}	450.2
	16 9 8 - 15 9 7	263 376.04	1.08×10^{-3}	450.2
	16 8 8 - 15 8 7	263 445.162	1.19×10^{-3}	378.3
	16 8 9 - 15 8 8	263 445.162	1.19×10^{-3}	378.3
	16 7 9 - 15 7 8	263 506.504	1.28×10^{-3}	314.9
	16 7 10 - 15 7 9	263 506.504	1.28×10^{-3}	314.9
	16 6 10 - 15 6 9	263 560.166	1.36×10^{-3}	259.9
	16 6 11 - 15 6 10	263 560.166	1.36×10^{-3}	259.9
	16 0 16 - 15 0 15	263 561.61	1.59×10^{-3}	107.6
	5 2 4 - 5 1 4	263 578.106	6.38×10^{-7}	28.8
	16 5 11 - 15 5 10	263 606.473	1.43×10^{-3}	213.4
	16 5 12 - 15 5 11	263 606.473	1.43×10^{-3}	213.4
	53 3 50 - 52 4 48	263 614.777	3.08×10^{-7}	1169.4
	16 4 13 - 15 4 12	263 646.356	1.49×10^{-3}	175.3
	16 4 12 - 15 4 11	263 646.356	1.49×10^{-3}	175.3
	16 2 15 - 15 2 14	263 662.527	1.57×10^{-3}	124.5
	16 3 14 - 15 3 13	263 682.153	1.53×10^{-3}	145.7
	16 3 13 - 15 3 12	263 683.596	1.53×10^{-3}	145.7
NH$_2$CN	13 3 11 0 - 12 3 10 0	260 132.895	1.77×10^{-3}	217.7
	13 3 10 0 - 12 3 9 0	260 135.878	1.77×10^{-3}	217.7
CH$_3$NCO	20 2 19 0 - 19 2 18 0	173 483.466	2.44×10^{-4}	111.3
	20 2 18 0 - 19 2 17 0	173 675.392	2.45×10^{-4}	111.3
	20 -3 0 1 - 19 -3 0 1	173 733.277	2.41×10^{-4}	153.1
	20 2 0 1 - 19 2 0 1	174 085.007	2.44×10^{-4}	123.3
	20 -1 0 2 - 19 -1 0 2	174 095.57	2.46×10^{-4}	146.3
	20 0 0 2 - 19 0 0 2	174 161.002	2.47×10^{-4}	140.4
	19 3 0 1 - 18 3 0 1	174 191.731	2.17×10^{-4}	145.2
	20 0 0 -3 - 19 0 0 -3	174 421.255	2.47×10^{-4}	202.9
	20 1 0 -3 - 19 1 0 -3	174 432.265	2.51×10^{-4}	209.6
	20 0 0 3 - 19 0 0 3	174 442.476	2.47×10^{-4}	202.1
	20 1 0 3 - 19 1 0 3	174 486.974	2.51×10^{-4}	208.8
	20 1 0 2 - 19 1 0 2	174 575.538	2.47×10^{-4}	146.4
	30 2 29 0 - 29 2 28 0	260 032.391	8.32×10^{-4}	217.4
	30 -3 0 2 - 29 -3 0 2	261 808.735	8.34×10^{-4}	300.1
	30 0 0 4 - 29 0 0 4	261 955.843	8.43×10^{-4}	395.9
	30 2 0 3 - 29 2 0 3	261 960.792	8.43×10^{-4}	332.3
	30 2 0 -3 - 29 2 0 -3	261 963.534	8.43×10^{-4}	333.1
	30 1 29 0 - 29 1 28 0	262 109.072	8.55×10^{-4}	201.1
	30 -1 0 -3 - 29 -1 0 -3	262 154.856	8.26×10^{-4}	313.4
CH$_3$NH$_2$	10 1 7 - 9 2 7	173 654.159	1.41×10^{-6}	120.0
	13 2 5 - 12 3 4	173 922.286	1.37×10^{-6}	209.9
	24 5 4 - 23 6 4	173 958.808	7.10×10^{-6}	735.2
	17 2 4 - 17 1 5	174 041.351	2.90×10^{-5}	341.0
	10 1 4 - 9 2 5	174 135.742	2.33×10^{-6}	120.2
	22 0 7 - 21 3 7	174 146.048	2.49×10^{-7}	535.3
	40 -3 1 - 40 2 0	174 158.953	3.19×10^{-5}	1776.0
	18 1 6 - 18 0 7	174 198.267	1.07×10^{-7}	371.0
	24 5 6 - 23 6 6	174 397.684	7.13×10^{-6}	735.6
	18 -2 1 - 18 1 0	174 438.993	3.08×10^{-5}	379.5
	17 3 4 - 16 4 4	186 769.592	9.24×10^{-6}	361.1
	15 2 4 - 15 1 5	186 793.84	3.29×10^{-5}	270.9
	39 9 5 - 38 10 5	186 802.979	8.44×10^{-6}	1969.6

Table D.3. continued.

Species	Transition J K L (M)	Frequency (GHz)	A_{ij} (s^{-1})	$E_{\rm up}$ (K)
	19 1 6 - 19 0 7	186 980.617	1.10×10^{-7}	411.7
	17 2 2 - 17 -1 3	186 993.005	3.68×10^{-5}	341.3
	22 8 0 - 23 -7 1	187 323.266	6.99×10^{-6}	786.6
	22 -8 1 - 23 7 0	187 323.267	6.99×10^{-6}	786.6
	22 8 5 - 23 7 5	187 339.428	7.00×10^{-6}	786.3
	16 -2 1 - 16 1 0	188 210.397	3.67×10^{-5}	305.2
	6 4 7 - 7 3 7	260 266.034	1.07×10^{-5}	107.5
	10 2 4 - 10 1 4	260 293.536	2.26×10^{-5}	132.7
	25 0 5 - 24 3 4	260 303.993	8.54×10^{-7}	686.4
	12 1 4 - 11 2 4	260 427.523	2.05×10^{-5}	168.6
	17 7 7 - 18 6 7	261 967.364	1.73×10^{-5}	516.1
	9 2 6 - 9 1 7	262 121.321	1.37×10^{-6}	111.5
	26 5 4 - 25 6 4	262 853.982	2.58×10^{-5}	843.5
	37 8 4 - 36 9 4	262 858.032	2.44×10^{-5}	1741.1
	8 2 0 - 8 -1 1	262 998.165	6.98×10^{-5}	92.8
	31 -1 1 - 30 4 0	263 356.252	7.76×10^{-7}	1063.8
	8 0 7 - 7 1 6	263 377.826	5.77×10^{-5}	77.1
	6 1 4 - 5 1 4	263 431.705	8.98×10^{-6}	48.6
	37 8 0 - 36 -9 1	263 436.944	2.47×10^{-5}	1740.8
	37 -8 1 - 36 9 0	263 436.963	2.47×10^{-5}	1740.8
HOCH$_2$CN	58 6 53 0 - 57 7 50 0	172 653.345	8.21×10^{-6}	809.1
	3 3 1 1 - 2 2 0 1	172 657.358	3.66×10^{-5}	20.6
	46 11 36 1 - 47 10 37 1	172 661.411	7.70×10^{-6}	651.7
	46 11 35 1 - 47 10 38 1	172 661.415	7.70×10^{-6}	651.7
	60 8 53 1 - 59 9 51 0	172 661.689	8.77×10^{-6}	904.5
	3 3 0 1 - 2 2 1 1	172 662.849	3.66×10^{-5}	20.6
	77 9 69 1 - 76 10 66 1	172 681.141	8.84×10^{-6}	1446.6
	54 4 51 0 - 53 5 48 0	172 689.37	3.60×10^{-6}	680.1
	21 7 15 0 - 22 6 16 0	172 698.2	6.49×10^{-6}	170.3
	21 7 14 0 - 22 6 17 0	172 698.426	6.49×10^{-6}	170.3
	17 2 16 0 - 16 0 16 1	172 716.049	2.20×10^{-6}	73.1
	55 4 52 1 - 54 5 49 1	173 518.066	3.03×10^{-6}	709.3
	72 7 65 0 - 71 8 63 1	173 520.841	9.70×10^{-6}	1234.4
	36 2 34 1 - 35 4 32 0	173 629.603	1.09×10^{-6}	310.1
	27 3 24 1 - 26 4 22 0	173 647.439	9.29×10^{-6}	186.2
	56 4 52 0 - 56 4 53 0	173 712.043	3.10×10^{-6}	737.3
	35 3 32 0 - 35 1 34 1	173 722.216	2.26×10^{-6}	294.7
	16 3 13 0 - 15 2 13 1	173 742.253	1.52×10^{-5}	72.8
	62 5 58 1 - 62 4 58 0	173 760.542	1.15×10^{-5}	905.9
	20 0 20 1 - 19 1 19 1	173 836.595	4.12×10^{-5}	96.7
	54 4 51 1 - 53 5 48 1	173 873.415	3.06×10^{-6}	685.1
	52 12 41 0 - 53 11 42 0	173 931.489	7.96×10^{-6}	809.7
	52 12 40 0 - 53 11 43 0	173 931.491	7.96×10^{-6}	809.7
	19 2 18 1 - 18 2 17 1	173 962.325	1.61×10^{-4}	94.7
	74 7 68 0 - 73 8 66 1	173 987.188	1.08×10^{-5}	1298.5
	46 3 43 0 - 46 3 44 0	173 996.613	3.03×10^{-6}	497.9
	19 2 18 0 - 18 2 17 0	174 009.939	1.55×10^{-4}	89.4
	77 9 69 0 - 76 10 66 0	174 038.606	8.93×10^{-6}	1441.9
	50 4 46 0 - 49 5 44 1	174 040.287	1.91×10^{-5}	592.8
	80 8 73 0 - 79 9 71 1	174 056.834	1.04×10^{-5}	1524.3
	15 3 13 0 - 14 2 13 1	174 177.206	1.55×10^{-5}	65.7
	18 3 16 0 - 18 0 18 1	174 309.545	1.30×10^{-6}	88.3
	52 4 48 0 - 52 3 49 0	174 312.887	5.08×10^{-5}	639.1
	12 0 12 1 - 11 1 10 0	174 442.231	1.66×10^{-5}	39.6
	47 3 44 0 - 46 4 42 1	186 298.941	1.88×10^{-5}	519.0
	20 1 19 0 - 19 1 18 0	186 316.378	1.93×10^{-4}	96.0

Table D.3. continued.

Species	Transition J K L (M)	Frequency (GHz)	A_{ij} (s^{-1})	E_{up} (K)
	43 4 39 0 - 42 5 38 0	186 329.612	1.16×10^{-5}	444.3
	68 9 60 1 - 67 10 58 0	186 335.61	1.11×10^{-5}	1156.0
	55 7 49 1 - 54 8 47 0	186 729.585	1.13×10^{-5}	755.6
	7 5 3 0 - 8 4 4 0	186 743.896	3.27×10^{-6}	47.1
	7 5 2 0 - 8 4 5 0	186 744.026	3.27×10^{-6}	47.1
	32 9 24 1 - 33 8 25 1	186 825.606	9.03×10^{-6}	351.6
	32 9 23 1 - 33 8 26 1	186 825.623	9.03×10^{-6}	351.6
	62 12 50 1 - 63 11 52 0	186 844.69	9.73×10^{-6}	1069.1
	62 12 51 1 - 63 11 53 0	186 844.733	9.73×10^{-6}	1069.1
	29 4 25 0 - 29 2 27 1	186 942.085	1.50×10^{-5}	215.3
	35 4 31 1 - 34 5 29 0	186 942.311	1.17×10^{-5}	307.5
	73 7 66 0 - 72 8 64 1	186 964.419	1.21×10^{-5}	1267.0
	54 4 50 0 - 54 3 51 0	186 998.042	5.38×10^{-5}	687.4
	10 2 9 1 - 10 1 9 0	187 020.922	3.09×10^{-5}	35.3
	27 4 23 1 - 27 3 24 1	187 029.983	3.65×10^{-5}	195.2
	7 5 3 1 - 8 4 4 1	187 053.965	3.29×10^{-6}	52.6
	7 5 2 1 - 8 4 5 1	187 054.088	3.29×10^{-6}	52.6
	21 1 21 0 - 20 1 20 0	187 114.251	2.00×10^{-4}	100.5
	19 5 14 1 - 20 4 16 0	187 135.775	8.93×10^{-6}	124.3
	30 3 28 1 - 30 2 28 0	187 174.006	4.29×10^{-6}	223.8
	27 2 26 1 - 26 3 24 0	187 227.832	8.46×10^{-6}	177.1
	21 1 21 1 - 20 1 20 1	187 229.632	1.87×10^{-4}	105.9
	75 4 72 0 - 74 5 70 1	187 235.57	9.23×10^{-6}	1277.0
	19 5 15 1 - 20 4 17 0	187 320.422	8.95×10^{-6}	124.3
	58 6 52 1 - 57 7 51 1	187 325.576	1.09×10^{-5}	814.7
	51 5 46 0 - 51 4 47 0	187 343.315	2.59×10^{-5}	624.7
	29 2 27 0 - 28 3 26 0	187 359.604	1.55×10^{-5}	201.4
	55 7 48 1 - 54 8 46 0	187 446.036	1.14×10^{-5}	755.6
	31 1 30 1 - 31 0 31 1	187 524.072	1.51×10^{-5}	227.8
	63 14 49 0 - 64 13 52 0	187 664.457	1.01×10^{-5}	1163.7
	63 14 50 0 - 64 13 51 0	187 664.457	1.01×10^{-5}	1163.7
	51 14 37 0 - 52 13 39 1	187 682.587	8.07×10^{-6}	858.5
	51 14 38 0 - 52 13 40 1	187 682.587	8.07×10^{-6}	858.5
	20 2 18 1 - 19 2 17 1	187 815.959	2.11×10^{-4}	105.0
	21 0 21 0 - 20 0 20 0	187 877.118	2.03×10^{-4}	100.3
	52 5 47 1 - 52 4 48 1	187 921.804	4.73×10^{-5}	653.0
	15 0 15 1 - 14 1 13 0	187 985.123	1.80×10^{-5}	57.9
	31 3 29 0 - 31 2 30 0	188 047.392	3.27×10^{-5}	232.2
	20 2 18 0 - 19 2 17 0	188 060.556	2.02×10^{-4}	99.7
	38 10 29 0 - 39 9 30 0	188 170.704	9.51×10^{-6}	466.9
	38 10 28 0 - 39 9 31 0	188 170.709	9.51×10^{-6}	466.9
	36 3 33 1 - 35 4 32 1	188 171.403	1.24×10^{-5}	316.0
	68 13 55 1 - 69 12 57 0	188 176.019	9.95×10^{-6}	1277.2
	68 13 56 1 - 69 12 58 0	188 176.029	9.95×10^{-6}	1277.2
	30 1 30 1 - 29 2 28 0	188 209.049	5.83×10^{-6}	205.7
	21 0 21 1 - 20 0 20 1	188 283.514	1.59×10^{-4}	105.7
	31 3 29 1 - 30 4 27 0	188 303.563	1.12×10^{-5}	237.5
	31 3 29 1 - 31 2 30 1	188 394.528	1.02×10^{-5}	237.5
	76 5 72 0 - 75 6 69 0	188 399.299	5.32×10^{-6}	1329.6
	46 3 43 0 - 46 2 44 0	188 408.722	4.36×10^{-5}	497.9
	50 5 45 0 - 50 4 46 0	188 418.654	7.90×10^{-6}	601.9
	46 2 44 0 - 46 2 45 0	259 949.966	5.25×10^{-6}	488.9
	78 7 71 1 - 78 6 72 1	259 964.069	1.25×10^{-4}	1441.1
	21 5 17 1 - 21 4 18 1	259 967.536	8.82×10^{-5}	142.4
	77 7 70 0 - 77 6 71 0	259 970.889	8.69×10^{-5}	1401.8
	18 5 13 0 - 18 4 14 0	259 977.031	8.61×10^{-5}	110.5

Table D.3. continued.

Species	Transition J K L (M)	Frequency (GHz)	A_{ij} (s^{-1})	$E_{\rm up}$ (K)
	28 25 3 0 - 27 25 2 0	260 045.93	1.09×10^{-4}	1033.3
	28 25 4 0 - 27 25 3 0	260 045.93	1.09×10^{-4}	1033.3
	28 26 2 1 - 27 26 1 1	260 047.45	7.53×10^{-5}	1107.7
	28 26 3 1 - 27 26 2 1	260 047.45	7.53×10^{-5}	1107.7
	23 3 21 0 - 22 2 21 1	260 050.159	4.54×10^{-5}	134.8
	18 5 14 0 - 18 4 15 0	260 057.55	8.61×10^{-5}	110.5
	24 9 16 0 - 25 8 17 0	260 073.689	1.95×10^{-5}	245.3
	24 9 15 0 - 25 8 18 0	260 073.69	1.95×10^{-5}	245.3
	17 5 12 0 - 17 4 13 0	260 110.383	8.53×10^{-5}	102.5
	20 5 16 1 - 20 4 17 1	260 116.183	8.76×10^{-5}	133.1
	19 5 14 1 - 19 4 15 1	260 131.715	8.68×10^{-5}	124.3
	17 5 13 0 - 17 4 14 0	260 161.917	8.53×10^{-5}	102.5
	42 4 39 0 - 42 2 40 0	260 162.5	5.40×10^{-6}	422.9
	45 4 42 1 - 44 5 40 0	260 194.78	2.87×10^{-5}	486.5
	83 7 77 0 - 82 8 75 1	260 208.253	3.98×10^{-5}	1613.5
	28 26 2 0 - 27 26 1 0	260 213.351	7.44×10^{-5}	1101.5
	28 26 3 0 - 27 26 2 0	260 213.351	7.44×10^{-5}	1101.5
	16 5 11 0 - 16 4 12 0	260 218.203	8.43×10^{-5}	95.0
	28 27 1 1 - 27 27 0 1	260 220.963	3.84×10^{-5}	1178.5
	28 27 2 1 - 27 27 1 1	260 220.963	3.84×10^{-5}	1178.5
	19 5 15 1 - 19 4 16 1	260 247.467	8.69×10^{-5}	124.3
	16 5 12 0 - 16 4 13 0	260 250.329	8.44×10^{-5}	95.0
	18 5 13 1 - 18 4 14 1	260 285.695	8.61×10^{-5}	115.9
	13 3 10 1 - 13 2 12 0	260 291.934	5.81×10^{-5}	58.2
	6 6 0 0 - 5 5 0 1	260 295.415	1.29×10^{-4}	59.3
	6 6 1 0 - 5 5 1 1	260 295.415	1.29×10^{-4}	59.3
	40 1 39 1 - 40 0 40 1	260 297.087	3.26×10^{-5}	368.9
	15 3 12 1 - 15 2 14 0	260 297.601	5.67×10^{-5}	71.1
	39 4 36 0 - 39 2 37 0	260 302.383	4.94×10^{-6}	368.3
	15 5 10 0 - 15 4 11 0	260 304.408	8.33×10^{-5}	87.9
	15 5 11 0 - 15 4 12 0	260 323.792	8.33×10^{-5}	87.9
	18 5 14 1 - 18 4 15 1	260 361.653	8.62×10^{-5}	115.9
	48 7 42 0 - 48 6 42 1	260 362.51	6.58×10^{-5}	589.1
	34 4 31 1 - 34 3 31 0	260 370.704	5.21×10^{-5}	291.4
	14 5 9 0 - 14 4 10 0	260 372.309	8.20×10^{-5}	81.2
	14 5 10 0 - 14 4 11 0	260 383.545	8.20×10^{-5}	81.2
	28 27 1 0 - 27 27 0 0	260 386.776	3.80×10^{-5}	1172.2
	28 27 2 0 - 27 27 1 0	260 386.776	3.80×10^{-5}	1172.2
	5 4 1 1 - 6 3 3 0	260 408.16	6.97×10^{-6}	34.3
	17 5 12 1 - 17 4 13 1	260 410.819	8.53×10^{-5}	107.9
	5 4 2 1 - 6 3 4 0	260 411.948	6.97×10^{-6}	34.3
	13 5 8 0 - 13 4 9 0	260 424.892	8.05×10^{-5}	75.0
	13 5 9 0 - 13 4 10 0	260 431.158	8.05×10^{-5}	75.0
	17 5 13 1 - 17 4 14 1	260 459.444	8.54×10^{-5}	107.9
	44 5 40 1 - 44 4 41 1	261 739.468	1.08×10^{-4}	479.1
	18 2 16 1 - 18 1 18 0	261 767.085	2.02×10^{-5}	87.4
	29 2 27 1 - 28 3 25 0	261 797.89	5.47×10^{-5}	206.3
	82 7 76 1 - 82 6 76 0	261 833.407	1.81×10^{-4}	1581.4
	40 1 39 0 - 40 1 40 0	261 835.122	3.23×10^{-6}	363.6
	16 3 13 1 - 16 2 15 0	261 835.599	6.36×10^{-5}	78.2
	43 5 38 0 - 43 3 40 1	261 888.854	1.74×10^{-6}	454.8
	50 7 43 0 - 50 6 45 1	261 922.423	6.56×10^{-5}	633.1
	58 4 54 0 - 57 5 52 1	261 937.931	5.31×10^{-5}	789.0
	40 1 39 0 - 40 0 40 0	261 939.987	3.97×10^{-5}	363.6
	60 7 53 0 - 60 6 55 1	261 942.087	5.76×10^{-5}	879.4
	54 3 51 1 - 54 2 52 1	261 971.845	7.83×10^{-5}	683.3

Table D.3. continued.

Species	Transition J K L (M)	Frequency (GHz)	A_{ij} (s^{-1})	E_{up} (K)
	76 9 67 1 - 75 10 65 0	261 983.481	3.28×10^{-5}	1412.5
	28 3 25 1 - 27 3 24 1	262 020.774	5.56×10^{-4}	198.8
	14 3 11 1 - 14 2 13 0	262 027.209	3.59×10^{-5}	64.4
	30 10 21 0 - 31 9 22 0	262 032.199	2.19×10^{-5}	344.6
	30 10 20 0 - 31 9 23 0	262 032.199	2.19×10^{-5}	344.6
	51 2 50 1 - 50 3 48 0	262 074.734	1.89×10^{-5}	587.0
	45 7 39 0 - 45 6 39 1	262 094.475	6.80×10^{-5}	526.6
	49 7 42 0 - 49 6 44 1	262 157.573	6.62×10^{-5}	610.9
	11 5 6 1 - 12 4 8 0	262 714.706	1.60×10^{-5}	69.4
	11 5 7 1 - 12 4 9 0	262 718.054	1.60×10^{-5}	69.4
	13 7 7 0 - 13 6 7 1	262 736.376	5.96×10^{-5}	108.3
	13 7 6 0 - 13 6 8 1	262 736.376	5.96×10^{-5}	108.3
	30 10 20 1 - 31 9 23 1	262 762.028	2.21×10^{-5}	350.1
	30 10 21 1 - 31 9 22 1	262 762.028	2.21×10^{-5}	350.1
	81 6 75 0 - 81 5 76 0	262 765.813	1.42×10^{-4}	1532.1
	30 12 18 0 - 31 11 20 1	262 773.324	1.59×10^{-5}	405.5
	30 12 19 0 - 31 11 21 1	262 773.324	1.59×10^{-5}	405.5
	28 3 25 0 - 27 3 24 0	262 797.496	5.36×10^{-4}	193.8
	14 7 8 0 - 14 6 8 1	262 825.871	6.20×10^{-5}	114.5
	14 7 7 0 - 14 6 9 1	262 825.871	6.20×10^{-5}	114.5
	37 4 34 0 - 37 2 35 0	262 846.614	4.71×10^{-6}	334.1
	46 7 39 0 - 46 6 41 1	262 856.464	6.80×10^{-5}	547.0
	43 7 37 0 - 43 6 37 1	262 877.193	6.92×10^{-5}	487.2
	15 7 9 0 - 15 6 9 1	262 918.607	6.39×10^{-5}	121.2
	15 7 8 0 - 15 6 10 1	262 918.607	6.39×10^{-5}	121.2
	45 5 41 1 - 45 4 42 1	262 927.287	1.12×10^{-4}	499.1
	27 3 24 0 - 26 2 24 1	262 967.477	4.05×10^{-5}	181.1
	45 5 41 0 - 45 4 42 0	262 995.615	1.02×10^{-4}	494.0
	16 7 10 0 - 16 6 10 1	263 014.239	6.55×10^{-5}	128.3
	16 7 9 0 - 16 6 11 1	263 014.239	6.55×10^{-5}	128.3
	57 5 52 0 - 57 4 54 1	263 043.568	4.87×10^{-5}	771.6
	45 7 38 0 - 45 6 40 1	263 071.806	6.85×10^{-5}	526.6
	17 7 11 0 - 17 6 11 1	263 111.982	6.67×10^{-5}	135.8
	17 7 10 0 - 17 6 12 1	263 111.982	6.67×10^{-5}	135.8
	42 7 36 0 - 42 6 36 1	263 181.856	6.97×10^{-5}	468.1
	15 2 13 1 - 14 2 13 0	263 186.129	1.51×10^{-5}	64.5
	18 7 12 0 - 18 6 12 1	263 211.256	6.78×10^{-5}	143.8
	18 7 11 0 - 18 6 13 1	263 211.256	6.78×10^{-5}	143.8
	62 7 55 0 - 62 6 57 1	263 304.782	5.54×10^{-5}	934.0
	19 7 13 0 - 19 6 13 1	263 311.186	6.87×10^{-5}	152.2
	19 7 12 0 - 19 6 14 1	263 311.186	6.87×10^{-5}	152.2
	55 16 39 0 - 56 15 41 1	263 311.961	2.12×10^{-5}	1035.9
	55 16 40 0 - 56 15 42 1	263 311.961	2.12×10^{-5}	1035.9
	20 7 14 0 - 20 6 14 1	263 410.999	6.95×10^{-5}	161.0
	20 7 13 0 - 20 6 15 1	263 410.999	6.95×10^{-5}	161.0
	17 3 14 1 - 17 2 16 0	263 422.063	6.50×10^{-5}	85.7
	40 2 39 1 - 40 1 40 1	263 422.917	3.38×10^{-5}	369.0
	41 7 35 0 - 41 6 35 1	263 437.433	7.02×10^{-5}	449.5
	43 7 36 0 - 43 6 38 1	263 457.612	6.95×10^{-5}	487.2
	54 12 42 1 - 55 11 44 0	263 490.25	2.54×10^{-5}	862.2
	54 12 43 1 - 55 11 45 0	263 490.252	2.54×10^{-5}	862.2
	63 7 57 1 - 62 8 55 0	263 491.36	3.37×10^{-5}	966.7
	21 7 15 0 - 21 6 15 1	263 509.996	7.01×10^{-5}	170.3
	21 7 14 0 - 21 6 16 1	263 509.996	7.01×10^{-5}	170.3
	52 2 51 1 - 51 3 49 0	263 555.984	1.88×10^{-5}	609.3
	38 4 35 1 - 38 2 36 1	263 581.59	1.52×10^{-6}	356.2

Table D.3. continued.

Species	Transition J K L (M)	Frequency (GHz)	A_{ij} (s^{-1})	E_{up} (K)
	22 7 16 0 - 22 6 16 1	263 607.225	7.06×10^{-5}	180.1
	22 7 15 0 - 22 6 17 1	263 607.225	7.06×10^{-5}	180.1
	42 7 35 0 - 42 6 37 1	263 624.45	6.99×10^{-5}	468.1
	40 7 34 0 - 40 6 34 1	263 649.421	7.06×10^{-5}	431.3
	48 4 44 0 - 47 5 43 0	263 656.627	4.52×10^{-5}	547.9
	29 2 28 0 - 28 2 27 0	263 680.233	5.59×10^{-4}	196.6
	23 7 17 0 - 23 6 17 1	263 701.584	7.10×10^{-5}	190.3
	23 7 16 0 - 23 6 18 1	263 701.948	7.10×10^{-5}	190.3