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***NuSTAR* observation of the supergiant fast X-ray transient IGR J11215–5952 during its 2017 outburst**

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ABSTRACT

We report on the results of a *NuSTAR* observation of the supergiant fast X-ray transient pulsar IGR J11215–5952 during the peak of its outburst in June 2017. IGR J11215–5952 is the only SFXT undergoing strictly periodic outbursts (every 165 days). *NuSTAR* caught several X-ray flares, spanning a dynamic range of 100, and detected X-ray pulsations at 187.0 s, which is consistent with previous measurements. The spectrum from the whole observation is well described by an absorbed power law (with a photon index of 1.4), which is modified, above ~ 7 keV, by a cutoff with an e-folding energy of ~ 24 keV. A weak emission line is present at 6.4 keV, consistent with K_{α} emission from cold iron in the supergiant wind. The time-averaged flux is $\sim 1.5 \times 10^{-10}$ erg cm⁻² s⁻¹ (3–78 keV, corrected for the absorption), translating into an average luminosity of about 9×10^{35} erg s⁻¹ (1–100 keV, assuming a distance of 6.5 kpc). The *NuSTAR* observation allowed us to perform the most sensitive search for cyclotron resonant scattering features in the hard X-ray spectrum, resulting in no significant detection in any of the different spectral extractions adopted (time-averaged, temporally selected, spin-phase-resolved and intensity-selected spectra). The pulse profile showed an evolution with both the energy (3–12 keV energy range compared with 12–78 keV band) and the X-ray flux: a double-peaked profile was evident at higher fluxes (and in both energy bands), while a single-peaked, sinusoidal profile was present at the lowest intensity state achieved within the *NuSTAR* observations (in both energy bands). The intensity-selected analysis allowed us to observe an anti-correlation of the pulsed fraction with the X-ray luminosity. The pulse profile evolution can be explained by X-ray photon scattering in the accreting matter above magnetic poles of a neutron star at the quasi-spherical settling accretion stage.

Key words. stars: neutron – X-rays: binaries – pulsars: individual: IGR J11215–5952

1. Introduction

Supergiant fast X-ray transients (SFXTs) are a sub-class of high-mass X-ray binaries (HMXBs; [Walter et al. 2015](#); [Martínez-Núñez et al. 2017](#); [Sidoli 2017](#); [Sidoli & Paizis 2018](#)) that show sporadic, short, and recurrent X-ray transient emission ([Sguera et al. 2005, 2006](#)) and are associated with early-type supergiants ([Negueruela et al. 2006](#)). At odds with classical HMXBs hosting a blue supergiant plus a neutron star (NS, hereafter), systems known since the early times of X-ray astronomy ([White et al. 1983](#)), SFXTs are not persistent X-ray emitters. Supergiant fast X-ray transients outbursts span a few days, with a peak usually lasting less than one day, and are punctuated by short X-ray flares whose duration distribution peaks in the range of 100–2000 s ([Sidoli et al. 2019](#)). Their X-ray spectrum is similar to that of classical accreting pulsars, and X-ray pulsations were discovered in a few members of the class (ranging from ~ 20 to 1200 s, [Sidoli 2017](#)), a signature for the presence of a NS. There is no consensus in the literature about the mechanism powering the transient X-ray emission. This is also because some of the crucial NS properties needed to investigate the transition across different regimes,

from quiescence to outburst ([Bozzo et al. 2008](#); [Shakura et al. 2012](#); [Shakura & Postnov 2017](#)), are elusive in SFXTs. It is worth noting that to date there is no SFXT for which both the pulsar spin period and magnetic field are known with certainty.

The X-ray transient source IGR J11215–5952 was discovered in 2005 ([Lubinski et al. 2005](#)) with the INTEGRAL satellite ([Winkler et al. 2003](#)) and later associated to the B0.5 Ia star HD 306414 ([Negueruela et al. 2005](#)). The presence of a supergiant companion together with the short outbursts (a few thousand seconds, as observed by INTEGRAL, [Sidoli et al. 2006](#)) make it a member of the SFXT sub-class. IGR J11215–5952 is also an X-ray pulsar (spin period, P_{spin} , of 186.78 ± 0.3 s, discovered by [Swank et al. 2007](#)) and undergoes an X-ray outburst every ~ 165 days ([Sidoli et al. 2007](#); [Romano et al. 2009](#)), a remarkable and unique characteristic of SFXTs (see [Sidoli 2017](#) for the most recent review focussed on SFXTs).

The periodicity in the occurrence of the outburst is assumed to be due to the orbital period of the system ([Sidoli et al. 2006](#)), making it the supergiant HMXB with the longest orbital period, overlapping with the region occupied by Be/X-ray binary systems in the so-called Corbet diagram ([Corbet 1986](#)). An

extreme orbital eccentricity ($e \gtrsim 0.8$) is suggested by the shortness of its outburst compared with the long orbital period (Romano et al. 2007, 2009; Sidoli et al. 2007; Negueruela et al. 2008; Lorenzo et al. 2014). The extreme orbital geometry is not contradicted by the observed radial velocity curve of the companion HD 306414, which is dominated by effects other than orbital ones, most likely stellar pulsations (Lorenzo et al. 2014).

The observed properties of the X-ray outburst led Sidoli et al. (2007) to propose that the supergiant wind is magnetically compressed along the stellar equatorial plane (similar to the mechanism discussed by ud-Doula & Owocki 2002) in order to produce a short duration outburst when the NS crosses this denser wind component along the orbit. This hypothesis is supported by the detection of a stellar magnetic field of ~ 1 kG in the optical companion (Hubrig et al. 2018). The possible role played by the supergiant magnetic field in triggering SFXT outbursts has been proposed by Shakura et al. (2014) in the quasi-spherical accretion scenario.

The parallax of the star HD 306414 was measured by *Gaia* Data Release 2 (DR2, Gaia Collaboration 2018), translating into a distance of $6.5_{-1.0}^{+1.4}$ kpc (Bailer-Jones et al. 2018). The source distance previously reported in the literature was very close to this value, from $d = 6.2$ kpc (Masetti et al. 2006) to $d \gtrsim 7$ kpc (Lorenzo et al. 2014). HD 306414 was observed at optical and UV wavelengths, with controversial results in terms of the stellar properties (in particular, the wind mass-loss rate and the wind velocity), due to different assumptions made (Lorenzo et al. 2014; Hainich et al. 2020).

A *NuSTAR* observation of IGR J11215–5952 that was performed during an outburst that occurred in 2016 (with a net exposure of 20 ks) led to a hint (at 2.63σ) of the presence of a cyclotron resonant scattering feature (CRSF, hereafter) at 17 keV (Sidoli et al. 2017). Being produced by scattering of X-rays by electrons in quantized energy levels in the strong magnetic field of the pulsar, CRSFs are the only direct way to measure the NS magnetic field strength. In IGR J11215–5952, the suggested feature would imply an NS magnetic field of 2×10^{12} G (Sidoli et al. 2017). These findings make the pulsar IGR J11215–5952 the perfect target for a more in-depth investigation of its hard X-ray spectrum with *NuSTAR*, through which we would seek confirmation of the CRSF.

2. Observation and data reduction

2.1. *NuSTAR*

The Nuclear Spectroscopic Telescope Array (*NuSTAR*, Harrison et al. 2013) carries two co-aligned telescopes that focus X-ray photons onto two independent focal plane modules named A and B (hereafter FPMA and FPMB), with a $12' \times 12'$ field of view (FOV). Each FPM contains four (2×2) solid-state cadmium zinc telluride (CdZnTe) pixel detectors providing a spatial resolution of $18''$ (full width at half maximum) and a spectral resolution of 400 keV (FWHM) at 10 keV, with an effective area calibrated in the energy band 3–78 keV (Madsen et al. 2015).

The *NuSTAR* observation of the source IGR J11215–5952 covered an entire day, starting on 2017 June 21 at 08:58 and ending on 2017 June 22 at 08:31 (TT). It was a fixed-time observation, planned to cover the brightest phase of the outburst, as expected from the known strict periodicity in the outburst recurrence of 164.6 d (Romano et al. 2009). The expectations were confirmed by a *Swift*/XRT monitoring around the expected peak (see Sect. 2.2 and Fig. 1). The low satellite orbit (with data gaps

lasting about 30 min per revolution) reduced the net exposure time to 43.7 ks.

NuSTAR data (Obs.ID 30301010002) were processed using version 1.7.1 of the *NuSTAR* data analysis software (*NuSTAR* DAS). Spectra and light curves were extracted with NUPRODUCTS on the cleaned event files, using circular extraction regions with a radius of $60''$. The background spectrum was extracted away from the point-spread function (PSF) source wings. Since the background was stable and constant along the observation, no further filtering was applied. *NuSTAR* source light curves with arrival times corrected to the solar system barycenter have been extracted with the *NuSTAR* DAS tool NUPRODUCTS and the keyword “barycorr=yes”. When needed (e.g. in pulse phase spectroscopy), good-time intervals (GTIs) were generated using XSELECT and then running NUPRODUCTS using the “usrgtifile” keyword, to correctly extract the temporal selected spectra. The source net count rate in the energy range 3–78 keV from the overall exposure (43.7 ks) was ~ 1.5 counts s^{-1} per single module.

Spectra from FPMA and FPMB in the energy band 3–78 keV were simultaneously fitted in XSPEC (v.12.10.1; Arnaud 1996), adopting cross-calibration constants to take into account calibration uncertainties. All fluxes reported in this paper were estimated using the FPMA response matrix as reference. The same procedure was used when *NuSTAR* spectra were jointly fitted together with the *Swift*/XRT one.

When fitting the spectra in XSPEC, the absorption model TBABS was adopted, assuming the photoelectric absorption cross sections of Verner et al. (1996) and the interstellar abundances of Wilms et al. (2000). The spectra were re-binned to have at least 20 counts per bin in order to apply the χ^2 statistics. All uncertainties in the spectral analysis are given at 90% confidence level, for one interesting parameter. All *NuSTAR* light curves and pulse profiles reported in this paper are background subtracted.

2.2. *Swift*

The *Neil Gehrels Swift* observatory (Gehrels et al. 2004; *Swift* hereafter) monitored the source X-ray flux (0.3–10 keV) during the expected times of the 2017 outburst. We reprocessed the *Swift*/XRT observations (X-ray Telescope XRT; Burrows et al. 2005) using XRPIPELINE in the HEASOFT software package version 6.25, adopting standard procedures. The source flux was low and only photon-counting data (PC) were used.

The appropriate spectral redistribution matrices were used, available in the calibration database maintained by the High Energy Astrophysics Science Archive Research Center (HEASARC). Source events were extracted from a circular region with a radius of 20 pixels ($\sim 47''$), while local background events were taken from an annular region, centred on the source, with inner and outer radii of 30 and 60 pixels, respectively. The *Swift*/XRT spectrum was analysed in the energy range 2–8 keV, since there are no net counts either below 2 keV or above 8 keV.

3. Temporal analysis

3.1. Light curves

The source light curve during the 2017 outburst is shown in Fig. 1. It was observed with *Swift*/XRT and *NuSTAR* (in the upper and lower panel, respectively). The *Swift* monitoring around the times of the *NuSTAR* exposure has a large gap of about ten days due to Sun constraints, which prevented additional *Swift*/XRT snapshots, yet it indicates that the *NuSTAR* pointing caught the source at the peak of the outburst.

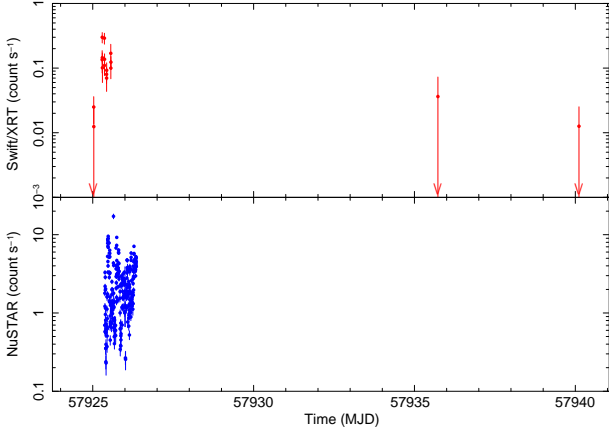


Fig. 1. *Swift*/XRT (0.3–10 keV) light curve of IGR J11215–5952 to monitor the onset of the 2017 outburst (*upper panel*), compared with the times of the *NuSTAR* observation (*lower panel*, 3–78 keV). A time binning of 187 s has been assumed. We note that, since there are gaps in both observations (by *Swift* and *NuSTAR*), unfortunately they were never simultaneous. The down arrows in the *upper panel* indicate upper limits.

Figure 2 reports the *NuSTAR* net source light curve alone, adopting a time bin of 187 s (the known pulsar spin period) in order to avoid intensity variability due to the X-ray pulsations. The three horizontal lines divide the outburst into four intensity regions, which have been adopted to perform the intensity-selected spectroscopy (Sect. 4.3).

In Fig. 3, we show the net light curves extracted from two energy bands, below (S) and above 12 keV (H), together with their hardness ratio (H/S). These two energy ranges were chosen to enable a more direct comparison with previous observations (Sidoli et al. 2017). There is no evidence of a strong trend of the spectral hardness, either with time (Fig. 3, left panel) or with the source intensity (right panel). The huge range of variability of the X-ray flux, spanning a factor of 100 (from 0.15 to 15 count s⁻¹, in the range 3–78 keV) is also evident.

3.2. Pulse profiles

We performed the timing analysis on *NuSTAR* background subtracted light curves (3–78 keV, bin time of 0.1 s), after correcting the arrival times to the solar system barycenter. Using epoch folding techniques, we searched for the known rotational period, measuring a periodicity $P_{\text{spin}} = 187.0 \pm 0.12$ s (1σ), consistent with previous values (Swank et al. 2007; Sidoli et al. 2007, 2017). Figure 4 shows the pulse profile obtained folding the whole exposure time on this periodicity in three energy ranges (3–12 keV, 12–78 keV, 3–78 keV), together with the hardness ratio of 12–78 keV to 3–12 keV (third panel from the top). Variability of the spectral hardness is evident along the rotation of the pulsar.

We calculated the pulsed fraction, PF (defined as $PF \equiv (F_{\text{max}} - F_{\text{min}})/(F_{\text{max}} + F_{\text{min}})$, where F_{max} and F_{min} are the count rates at the maximum and at the minimum of the spin profile, respectively), in three energy ranges, finding a correlation with the energy: $PF = 30 \pm 2\%$ (3–12 keV), $PF = 35 \pm 2\%$ (12–20 keV) and $PF = 43 \pm 4\%$ (20–78 keV). The pulsed fraction measured over the entire *NuSTAR* energy range is $PF = 30.1 \pm 0.9\%$ (3–78 keV).

We investigated the evolution of the pulse shape with the X-ray luminosity, extracting the folded light curves in the four intervals of count rate shown in Fig. 2 and used for

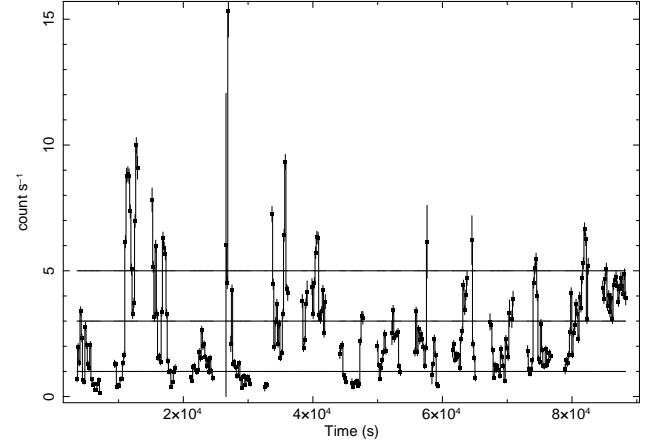


Fig. 2. *NuSTAR* net source light curve (3–78 keV) of the 2017 outburst. A time binning of 187 s has been assumed. The three horizontal lines mark the intensity intervals adopted in the intensity selected spectroscopy (see Sect. 4.3 for details).

the intensity selected spectroscopy (Table 1): count rate interval <1 count s⁻¹, $1-3$ count s⁻¹, $3-5$ count s⁻¹ and count rate >5 count s⁻¹ (3–78 keV, time bin of 187 s). These intensity ranges correspond to the average X-ray luminosities reported in the seventh column of Table 1. In Fig. 5, we report the intensity-selected pulse profiles in three energy bands. The pulse profile also changes with the X-ray luminosity, and an anti-correlation of the pulsed fraction with the X-ray luminosity is observed (Fig. 6).

4. Spectroscopy

We performed spectroscopy extracting *NuSTAR* spectra in different ways (time-averaged, temporally-resolved, intensity-selected, and spin-phase-resolved spectra) to search for evidence of CRSFs, as reported in the following sub-sections.

4.1. Time-averaged spectrum

We first analysed the *NuSTAR* spectra (FPMA and FPMB) extracted from the entire exposure time, fitting them jointly together with the *Swift*/XRT spectrum (but we note that XRT snapshots were never simultaneous with *NuSTAR*). The *Swift*/XRT exposure time is 1.8 ks, and the source average count rate in the 0.3–10 keV energy band is 0.117 ± 0.008 count s⁻¹.

A single power law is not able to model the spectrum appropriately, given the presence of a cutoff at high energies. A faint emission line is evident at ~ 6.4 keV (equivalent width, $EW \sim 50$ eV) consistent with it being produced by neutral iron and modeled with a narrow Gaussian line. Notably, this emission line is detected only in the time-averaged spectrum.

In Table 2, we report the results with two models: the first one (which we refer to as PLCUT hereafter) includes a power law modified by a high-energy cut-off (HIGHECUT in XSPEC), the latter being defined as $M(E) = (\exp[(E_{\text{cut}} - E)/E_{\text{fold}}])$ when $E \geq E_{\text{cut}}$, while $M(E) = 1$ at $E \leq E_{\text{cut}}$; the second model (CUTOFFPL) adopts a cut-off power law (CUTOFFPL in XSPEC), defined as $A(E) = KE^{-\Gamma} \exp[-E/E_{\text{cut}}]$, where Γ is the photon index. The model PLCUT was successfully adopted by Sidoli et al. (2017) to investigate the 2016 *NuSTAR* observation, as a statistically acceptable description. Since the best fit was obtained with PLCUT (Fig. 7), in the following we report results adopting this spectral model only. We show the count spectra

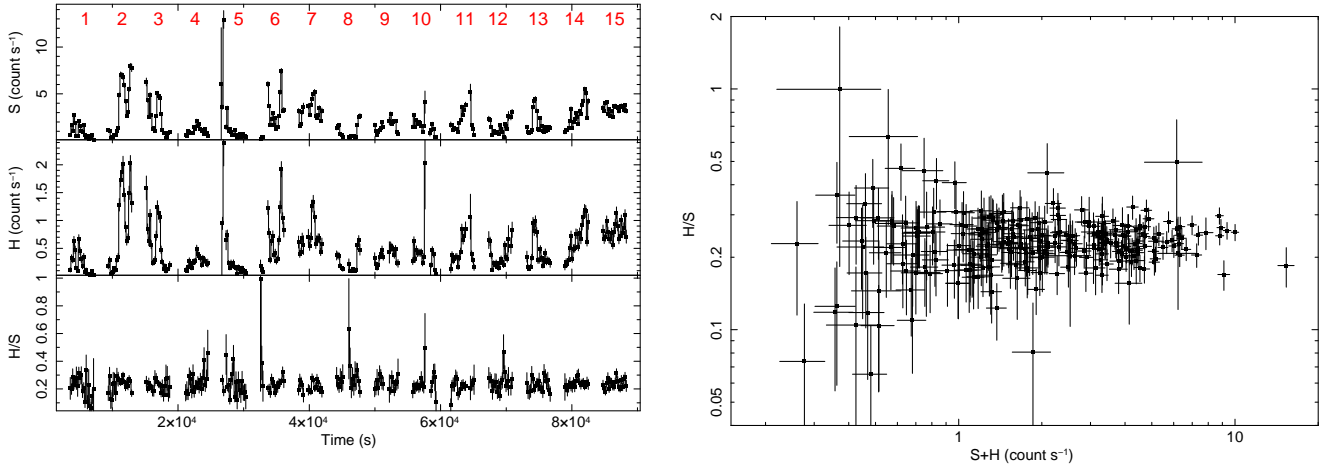


Fig. 3. IGR J11215–5952 light curves in two energy bands ($S = 3\text{--}12\text{ keV}$, $H = 12\text{--}78\text{ keV}$) together with their hardness ratio (HR, on the *left*). Red numbers on the top of the soft X-ray’s light curve indicate the *NuSTAR* satellite revolutions analysed in Table 3. *Right*: same hardness ratio versus the 3–78 keV count rate is shown. In both panels, a bin time of 187 s has been adopted.

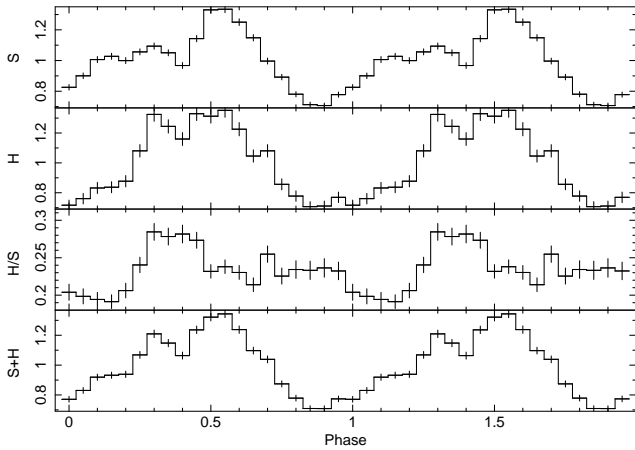


Fig. 4. IGR J11215–5952 pulse profile folded on the spin period of 187 s, assuming MJD 57432 as epoch of phase zero. The whole *NuSTAR* exposure time has been considered. *From top to bottom*: soft profile ($S = 3\text{--}12\text{ keV}$), hard one ($H = 12\text{--}78\text{ keV}$), their hardness ratio (H/S) and profile in the total band ($S+H = 3\text{--}78\text{ keV}$), respectively. Each pulse profile has been normalised by dividing by the average source intensity in the considered energy band.

fitted with this model in Fig. 7. We found no evidence of statistically significant absorption features.

We note that the analysis of the *NuSTAR* spectra alone led to spectral parameters consistent with the ones reported in Table 2.

4.2. Temporally resolved spectroscopy

In Table 3, we report on the temporally selected spectroscopy, adopting the natural segmentation of the *NuSTAR* observation into 15 satellite orbits, as is evident from Fig. 2 (marked with red numbers). We did not find any evidence of evolution in the spectral parameters with time, nor with the X-ray flux, within the uncertainties. No evidence of CRSFs is present.

4.3. Intensity-selected spectroscopy

Although there is no strong evidence of a trend in the hardness ratio (Fig. 3, $H/S = 12\text{--}78\text{ keV}/3\text{--}12\text{ keV}$) with the source flux,

we nevertheless performed an intensity-selected spectroscopy to investigate the eventual presence of absorption features compatible with CRSFs. We extracted intensity selected spectra, after binning the net-source light curve on the pulsar spin period, to avoid source variability due to the X-ray pulsations. Four X-ray intensity ranges were considered (to get comparable statistics) along with net source count rates in the following intervals (shown in Fig. 2): below 1 count s^{-1} , $1\text{--}3\text{ count s}^{-1}$, $3\text{--}5\text{ count s}^{-1}$, and above 5 count s^{-1} , resulting in net exposure times of 13.5 ks, 17 ks, 9.7 ks, and 3.5 ks, respectively, in the four spectra. In Table 1, we list the spectral parameters obtained with the PLCUT model, confirming no variability with the X-ray flux, within the uncertainties. No significant absorption lines were present.

4.4. Spin-phase-resolved spectroscopy

The source shows evidence of variability of the spectral hardness along the pulse profile (Fig. 4). This behaviour led us to perform a spin-phase resolved spectroscopy. The results are reported in Table 4 and shown in Fig. 8. Harder X-ray emission is found at spin phases $\Delta\phi = 0.3\text{--}0.5$, preceding the main peak of the 3–12 keV profile, while the softest emission appears in the range $\Delta\phi = 0.0\text{--}0.2$ (Fig. 4). The spectral analysis indicates that this hardening can be explained by two different reasons (Fig. 8): in the spin phase interval $\Delta\phi = 0.3\text{--}0.4$, the power law extends up to very high energies (high E_{cut} compared with other spin phases), while in the interval $\Delta\phi = 0.4\text{--}0.5$, the flattest power law is obtained.

4.5. Spin-phase-resolved spectroscopy of intensity-selected spectra

In Table 5 and Fig. 9, we report on the spin-phase-selected spectroscopy performed within the four intensity states, adopting the same model and ten spin-phase intervals per period. Given the shorter exposure time compared with previous spectroscopy, the energy range covered by these spectra is limited to the band 3–50 keV, with the only exception of the lowest luminosity state at the minimum of the pulse profile, where the spectrum extends up to 30 keV. No evidence of significant absorption features was obtained. In a few cases, the parameters of the high-energy cut-off component (E_{cut} and E_{fold}) were unconstrained (and are

Table 1. Intensity-selected spectroscopy (see Sect. 4.3).

Count rate range (count s ⁻¹)	N_{H} (10 ²² cm ⁻²)	Γ ^(a)	E_{cut} (keV)	E_{fold} (keV)	$UF_{1-100\text{keV}}$ ^(b) (erg cm ⁻² s ⁻¹)	$L_{1-100\text{keV}}$ (10 ³⁶ erg s ⁻¹)	$\chi^2_{\nu}/\text{d.o.f.}$	Net exposure (ks)
<1	5.3 ^{+2.2} _{-2.1}	1.44 ^{+0.15} _{-0.16}	6.3 ^{+1.1} _{-1.1}	23 ⁺⁷ ₋₅	0.70×10^{-10}	0.35 ^{+0.17} _{-0.11}	1.048/538	13.5
1–3	4.6 ^{+1.4} _{-1.4}	1.45 ^{+0.11} _{-0.10}	6.5 ^{+1.0} _{-0.5}	25 ⁺⁵ ₋₄	1.31×10^{-10}	0.66 ^{+0.32} _{-0.20}	0.920/746	17.0
3–5	6.6 ^{+1.4} _{-1.2}	1.50 ^{+0.11} _{-0.08}	6.5 ^{+1.0} _{-0.4}	27 ⁺⁵ ₋₃	2.90×10^{-10}	1.47 ^{+0.70} _{-0.46}	1.027/800	9.7
> 5	5.0 ^{+1.6} _{-1.6}	1.33 ^{+0.12} _{-0.10}	6.5 ^{+0.8} _{-0.5}	21 ⁺⁴ ₋₃	4.75×10^{-10}	2.40 ^{+1.15} _{-0.74}	1.019/681	3.5

Notes. ^(a)Power-law photon index. ^(b)Flux corrected for the absorption.

indicated by vertical lines in Fig. 9), and a single absorbed power law was an acceptable deconvolution of the X-ray continuum emission.

5. Discussion

We report the X-ray properties shown by the SFXT pulsar IGR J11215–5952 during the periodic outburst occurred in June, 2017 observed with the *NuSTAR* satellite (3–78 keV) at its maximum brightness. Monitoring observations with *Swift*/XRT (2–8 keV) confirmed the rise to the outburst peak and the shortness of the flaring activity phase, in agreement with previous ones, where the X-ray luminosity was above 10³⁴ erg s⁻¹ for fewer than 10 days (Romano et al. 2007, 2009).

The *NuSTAR* observation resulted in a net exposure time of about 44 ks, but since *NuSTAR* operates in a low orbit with several observational gaps, the total temporal coverage of the X-ray activity was about one day, a time span corresponding to the typical duration of the peak of the outburst in IGR J11215–5952.

The *NuSTAR* light curve showed a flaring activity with an ample range of variability, reaching a factor of 100. Although the light curve is segmented in 15 satellite revolutions (so that the temporal evolution of the X-ray variability cannot be entirely followed), it is possible to distinguish a few short (~hundred seconds) X-ray flares within single satellite orbits. This huge flaring variability in the outburst is the typical X-ray activity characterising the SFXTs as a class.

5.1. X-ray spectrum

The X-ray spectrum was well modeled with an absorbed power law modified by a high-energy cut-off, a phenomenological deconvolution frequently used to describe the spectra of accreting pulsars (e.g. Coburn et al. 2002), together with a faint emission line at 6.4 keV, consistent with neutral iron, which is usually observed in HMXBs and due to fluorescent emission in the companion wind (Torrejón et al. 2010; Giménez-García et al. 2015).

This is the second time *NuSTAR* has observed IGR J11215–5952 during an outburst. The difference this time is the doubled exposure time (i.e. ~44 ks) with respect to the 2016 outburst reported by Sidoli et al. (2017). In 2016, the outburst was followed also by *XMM-Newton*, simultaneously with *NuSTAR*, implying a very sensitive coverage of the soft region of the spectrum (below 3 keV), which showed a soft excess that could be modelled equally well either with a hot black body or with an additional absorption model, covering only a fraction of the power-law emission at soft energies. In the reported 2017 outburst, the *Swift*/XRT observations were too short to detect the source below 2 keV, and the 2–78 keV

spectrum did not show evidence of a more complex model than a single absorbed power law with a high-energy rollover.

The *NuSTAR* spectra collected in 2016 and 2017, extracted from the whole exposures, were very similar. The re-analysis of the whole 2016 exposure (20 ks), showed a slightly less absorbed ($(2.4 \pm 0.7) \times 10^{22}$ cm⁻²) harder power-law model ($\Gamma = 1.26^{+0.10}_{-0.07}$), with a slightly higher flux of $F = 1.6 \times 10^{-10}$ erg cm⁻² s⁻¹ (3–78, not corrected for the absorption) than in 2017. The cut-off and e-folding energies, together with the parameters of the faint and narrow emission line from neutral iron, were compatible in the two outbursts.

During the spectroscopy of a flare of the 2016 outburst, we found (Sidoli et al. 2017) a hint (at 2.63 σ) of an absorption feature at 17 keV, probably a variable CRSF, which needed confirmation. Unfortunately, we did not find any evidence of CRSFs in any of the spectra extracted from the 2017 observation. We indeed searched for significant absorption features by adopting many different spectral extractions: a time-averaged spectrum, temporally selected, spin-phase-selected, and intensity-selected spectra (together with a spin-phase-resolved spectroscopy at different X-ray intensity states). All these spectra were well fitted with the same model, with the exception of the faint K α iron emission line, which was clearly detected only in the longest exposure, time-averaged spectrum. In none of them did we find significant absorption features compatible with the presence of a CRSF.

5.2. X-ray pulsations

X-ray pulsations were clearly detected at a periodicity $P_{\text{spin}} = 187.0 \pm 0.12$ s, consistent with past observations (Swank et al. 2007; Sidoli et al. 2007, 2017). This shows a remarkable stability of the pulsar rotation since its discovery, which perhaps suggests that the equilibrium period has been reached.

The pulsed fraction showed an evolution with both the energy and the X-ray luminosity: increasing at harder energies and decreasing at higher luminosity. The changing pulse shape is typically observed in accreting pulsars and has been investigated by many authors, especially in transient Be X-ray binaries, which span a large range of X-ray luminosity (Parmar et al. 1989). An evolution in the pulse profile indicates a change in the radiation beam pattern, and is thought to be due to a changing accretion regime onto magnetised NS. In particular, it has been investigated in the most luminous and transient X-ray pulsars crossing the critical luminosity threshold at $\sim 10^{37}$ erg s⁻¹ (Becker et al. 2012). However, in IGR J11215–5952, the luminosity reached during the outburst is always sub-critical, at $\sim 10^{36}$ erg s⁻¹ (1–100 keV). Nevertheless, we observed a changing pulse pattern from a single, sinusoidal peak at low luminosity (at an average $L_{1-100\text{keV}} = 3.5 \times 10^{35}$ erg s⁻¹, see Table 1) and in

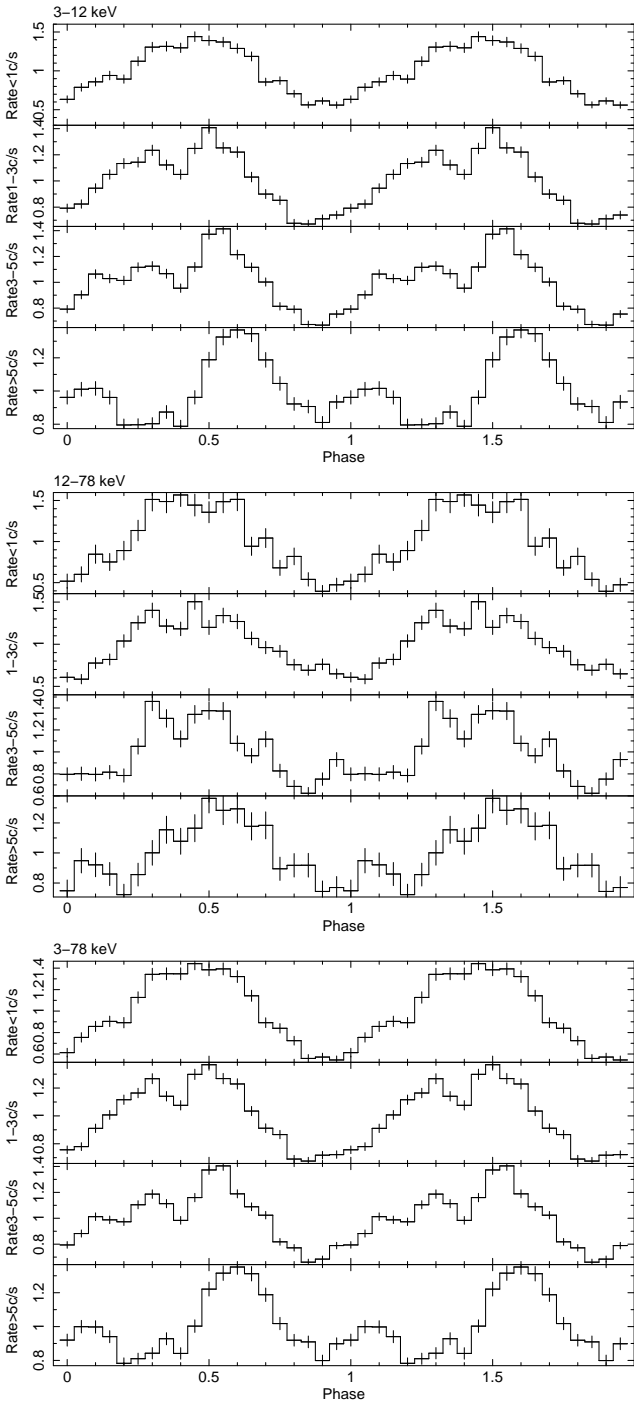


Fig. 5. IGR J11215–5952 pulse profiles in three energy bands (from top to bottom, 3–12 keV, 12–78 keV and 3–78 keV). For each energy band, we report pulse profiles in four intensity-selected intervals. Each profile has been normalised by dividing by the average source intensity in the considered energy band.

all energy ranges (3–12 keV, 12–78 keV and 3–78 keV) and up to a double-peaked pattern at the highest luminosity range (at the peak of the brightest flares, $L_{1-100\text{keV}} = 2.4 \times 10^{36} \text{ erg s}^{-1}$). It was composed of a main and a secondary peak, and the main peak was shifted with respect to the single peak of the lowest luminosity state by about $\Delta\phi = +0.15$.

A more complex pulse profile is observed at intermediate luminosities (Fig. 5). In particular, at soft energies (3–12 keV), a notch appears at spin phase $\phi = 0.4$ in the profile extracted from

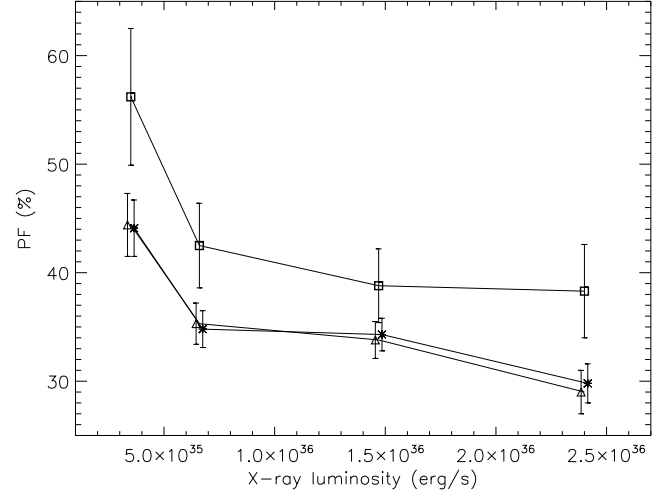


Fig. 6. Pulsed fraction dependence on the average X-ray luminosity (1–100 keV) measured from the four intensity-selected spectra (Table 1). Results in three energy bands are shown with the following symbols: triangles mark the pulsed fractions in the energy band 3–12 keV, squares in the band 12–78 keV, stars in the total band 3–78 keV. A small shift on the x-axis was applied for visualisation purposes.

Table 2. Broad-band (*Swift*/XRT together with *NuSTAR*) time-averaged spectrum (0.3–78 keV; see Sect. 4.1).

Parameter	Model PLCUT	Model CUTOFFPL
N_{H} (10^{22} cm^{-2})	$5.28^{+0.72}_{-0.76}$	$6.28^{+0.65}_{-0.65}$
Γ ^(a)	$1.46^{+0.05}_{-0.07}$	$1.33^{+0.05}_{-0.05}$
E_{cut} (keV)	$7.24^{+0.46}_{-0.66}$	$20.8^{+1.5}_{-1.4}$
E_{fold} (keV)	$23.6^{+1.9}_{-1.7}$	–
E_{line} (keV)	$6.32^{+0.08}_{-0.03}$	$6.36^{+0.04}_{-0.09}$
σ_{line} (keV)	<0.8	<0.8
Flux _{line} ($10^{-5} \text{ ph cm}^{-2} \text{ s}^{-1}$)	4.2 ± 0.1	4.5 ± 0.1
EW _{line} ^(b) (eV)	53 ± 13	57 ± 13
$F_{3-78\text{keV}}$ ($\text{erg cm}^{-2} \text{ s}^{-1}$) ^(c)	1.39×10^{-10}	1.37×10^{-10}
UF _{3-78keV} ($\text{erg cm}^{-2} \text{ s}^{-1}$) ^(d)	1.47×10^{-10}	1.47×10^{-10}
UF _{1-100keV} ($\text{erg cm}^{-2} \text{ s}^{-1}$) ^(d)	1.77×10^{-10}	1.78×10^{-10}
$L_{1-100\text{keV}}$ (erg s^{-1}) ^(e)	$(8.9^{+4.3}_{-2.7}) \times 10^{35}$	$(9.0^{+4.3}_{-2.8}) \times 10^{35}$
$\chi^2_{\nu}/\text{d.o.f.}$	1.045/1131	1.077/1132

Notes. ^(a)Power-law photon index. ^(b)Equivalent width. ^(c)Flux not corrected for the absorption. ^(d)Fluxes corrected for the absorption. ^(e)The luminosity is calculated assuming a source distance of $6.5^{+1.4}_{-1.1} \text{ kpc}$. The quoted error on the luminosity was calculated from the uncertainty on the source distance only.

the count rate range 1–3 count s^{-1} (average $L_{1-100\text{keV}} = 6.6 \times 10^{35} \text{ erg s}^{-1}$, Table 1), becoming deeper and broader towards higher luminosity states. At hard energies (12–78 keV), the pulse shape undergoes a less clear evolution with the luminosity, but with a similar final result: the profile changes from a single-peaked (low luminosity) to a double-peaked profile at the highest luminosity state (average $L_{1-100\text{keV}} = 2.4 \times 10^{36} \text{ erg s}^{-1}$). However, this is with a broader main peak than the one observed in the energy range 3–12 keV.

We note that a pulse-profile evolution was already observed in IGR J11215–5952 using *XMM-Newton* to investigate the 2006 outburst (Sidoli et al. 2007). On that occasion, the change

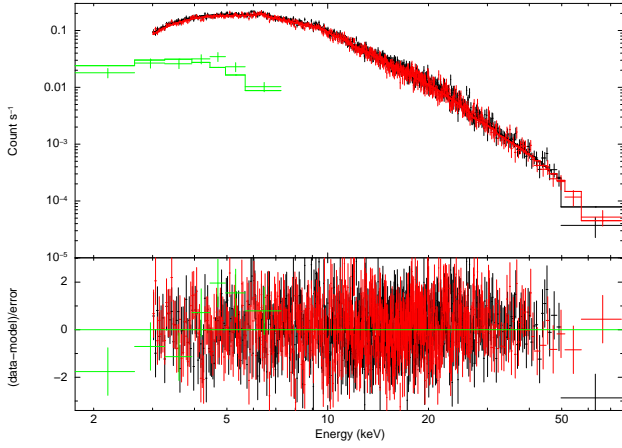


Fig. 7. *Swift*/XRT and *NuSTAR* time-averaged spectra, fitted with an absorbed power law modified by a high-energy cutoff, together with a Gaussian line in emission, accounting for iron line at ~ 6.4 keV (see Table 2, model PLCUT, for the best fit parameters). *Lower panel:* residuals are shown in terms of standard deviation.

from a single-peaked (at low luminosity) to a double-peaked profile (at high luminosity) resulted in a larger phase shift of the main peak with respect to the single-peak at lower luminosity. However, it is worth noting that the low-luminosity state probed by *XMM-Newton* was fainter than the lowest luminosity state probed by *NuSTAR*: indeed, rescaling luminosities in the same band (0.5–10 keV) and at the same distance of 6.5 kpc, we found that in 2006, the single-peaked pulse profile of the faint state was at a luminosity of 0.9×10^{34} erg s $^{-1}$ (0.5–10 keV), while with *NuSTAR* it was observed at an average luminosity of 2×10^{35} erg s $^{-1}$ (0.5–10 keV). Thus, we can conclude that the phase shift was larger in 2006, because the pulse profile was extracted at a lower luminosity with *XMM-Newton* than it was in 2017 using *NuSTAR*.

5.3. X-ray pulse-profile evolution in IGR J11215–5952 with luminosity: a physical explanation

The analysis of the pulse profile properties (time evolution with X-ray flux and time-resolved spectroscopy) suggests the following features, which should be explained by a physical model:

1. The 3–12 keV pulse-profile shape changes with X-ray flux, from a single-peak sine-line form at low luminosities (a few 10^{35} erg s $^{-1}$), to more complex double-peak profile at high luminosities (a few 10^{36} erg s $^{-1}$), while the 12–78 keV pulse shape almost does not change (see Figs. 4–9).
2. The hardness ratio HR changes along with the pulse (Fig. 9) but keeps, on average, constant at different X-ray luminosities (Fig. 3).
3. The pulsed fraction *PF* increases below $\sim 6 \times 10^{35}$ erg s $^{-1}$ (Fig. 6).
4. No CRSF has been found in the 3–78 keV spectral range (Fig. 7).

At low and moderate accretion rates, the X-ray emission is generated by converting the accreting plasma kinetic energy into heat in a thin layer near the NS surface (Zel’dovich & Shakura 1969; Lamb et al. 1973). The shape of the X-ray pulse profile and its evolution with luminosity are determined by the physical conditions above the NS polar cap (magnetic field strength, plasma density, and temperature) and are insensitive to the mechanism triggering the source flaring activity.

Table 3. Temporally resolved spectroscopy (see Sect. 4.2).

Rev.	N_{H} (10^{22} cm $^{-2}$)	Γ ^(a)	E_{cut} (keV)	E_{fold} (keV)	$UF_{1-100\text{keV}}$ ^(b) (erg cm $^{-2}$ s $^{-1}$)	$\chi^2_{\nu}/\text{d.o.f.}$
1	7.3 $^{+5.0}_{-4.4}$	1.40 $^{+0.33}_{-0.27}$	6.6 $^{+2.9}_{-6.6}$	22 $^{+18}_{-7}$	0.82×10^{-10}	0.891/195
2	6.1 $^{+2.2}_{-2.7}$	1.34 $^{+0.14}_{-0.21}$	7.2 $^{+1.1}_{-1.3}$	22 $^{+5}_{-5}$	2.89×10^{-10}	0.873/545
3	5.2 $^{+2.6}_{-2.6}$	1.51 $^{+0.15}_{-0.18}$	7.0 $^{+0.9}_{-1.0}$	26 $^{+9}_{-6}$	1.97×10^{-10}	0.975/418
4	6.4 $^{+4.1}_{-4.0}$	1.38 $^{+0.49}_{-0.29}$	6.1 $^{+1.1}_{-6.6}$	25 $^{+24}_{-8}$	0.95×10^{-10}	0.734/216
5	6.6 $^{+3.9}_{-6.8}$	1.51 $^{+0.20}_{-0.33}$	9.7 $^{+2.3}_{-4.8}$	24 $^{+11}_{-6}$	0.91×10^{-10}	0.981/195
6	8.7 $^{+3.9}_{-2.5}$	1.62 $^{+0.61}_{-0.19}$	6.5 $^{+5.7}_{-6.5}$	36 $^{+41}_{-11}$	3.13×10^{-10}	0.944/407
7	6.2 $^{+2.6}_{-2.5}$	1.54 $^{+0.15}_{-0.16}$	6.3 $^{+1.2}_{-0.8}$	28 $^{+11}_{-6}$	2.69×10^{-10}	1.055/429
8	2.6 $^{+5.5}_{-4.1}$	1.36 $^{+0.36}_{-0.33}$	6.6 $^{+8.0}_{-6.6}$	25 $^{+27}_{-12}$	0.78×10^{-10}	0.927/152
9	7.7 $^{+4.1}_{-3.7}$	1.67 $^{+0.23}_{-0.25}$	6.8 $^{+3.0}_{-6.8}$	45 $^{+71}_{-27}$	1.57×10^{-10}	0.872/223
10	5.9 $^{+4.2}_{-4.1}$	1.43 $^{+0.43}_{-0.30}$	6.1 $^{+7.5}_{-6.1}$	22 $^{+27}_{-7}$	1.27×10^{-10}	0.793/224
11	5.7 $^{+3.1}_{-3.3}$	1.64 $^{+0.19}_{-0.19}$	6.1 $^{+1.7}_{-6.1}$	35 $^{+25}_{-12}$	1.74×10^{-10}	0.991/268
12	4.7 \pm 4.0	1.49 $^{+0.24}_{-0.23}$	5.8 $^{+2.5}_{-5.8}$	28 $^{+33}_{-9}$	1.11×10^{-10}	0.868/215
13	3.0 $^{+4.1}_{-3.0}$	1.27 $^{+0.30}_{-0.22}$	7.0 $^{+2.4}_{-1.0}$	19 $^{+9}_{-4}$	1.38×10^{-10}	0.963/325
14	5.6 $^{+2.8}_{-2.3}$	1.49 $^{+0.20}_{-0.17}$	6.8 $^{+1.9}_{-1.2}$	26 $^{+12}_{-6}$	2.36×10^{-10}	0.958/471
15	5.0 $^{+2.3}_{-2.0}$	1.38 $^{+0.23}_{-0.14}$	6.3 $^{+2.6}_{-0.8}$	23 $^{+8}_{-4}$	2.93×10^{-10}	0.971/546

Notes. ^(a)Power-law photon index. ^(b)Flux corrected for the absorption.

5.3.1. Photon scattering in accreting plasma above one NS magnetic pole

Firstly, let us suppose that in IGR J11215–5952 we are observing only X-ray pulsations from one NS magnetic pole, as suggested by the 12–78 keV pulse profile. This is possible for a given range of viewing angles of the NS spin and magnetic axes. In this case, we can propose the following physical explanation to the observed gross properties of the X-ray pulse evolution with luminosity.

Shakura et al. (2013) discussed, in the context of the “off-states” of the canonical wind-accreting X-ray pulsar Vela X-1, the pulse profile changes in low- and moderate-luminosity X-ray pulsars with luminosity. The pulse-profile evolution is expected, due to the possible transformation of a pencil X-ray beam produced by ordinary (O) photons at the NS polar caps into a more spherically symmetric pattern, or even a fan beam at high accretion rates, due to photon scatterings in the accreting matter and possible vacuum polarisation effects. This transition can occur above the X-ray luminosity

$$L_{\dagger} \sim 3 \times 10^{35} \mu_{30}^{-3/10} \text{ erg s}^{-1}.$$

Here, $\mu = 10^{30} \mu_{30} \text{ G cm}^3$ is the NS dipole magnetic moment normalised to a canonical NS surface magnetic field of 10^{12} G.

In this model, the change in the X-ray beam for photons above the cyclotron energy E_{cyc} , from pencil-like to fan-like as the X-ray luminosity increases, leads to the X-ray pulse phase shifting by about 90° , as indeed is the case for Vela X-1 with $E_{\text{cyc}} \approx 20$ keV (Doroshenko et al. 2011).

In IGR J11215–5952, however, only a slight pulse-phase shift has been detected (Sidoli et al. 2007, the *XMM-Newton* observations, and the present paper). The unchanged mean hardness ratio HR (property 2 of Sect. 5.3) suggests that the observed absorption feature in the X-ray pulse observed at high luminosities is due to scatterings of O-photons generated near the polar cap of the magnetised NS. The smallness in the X-ray pulse phase shift in the range 3–78 keV observed by *NuSTAR* can be interpreted by the lack of transformation of the pencil X-ray diagram in IGR J11215–5952 with luminosity. This is possible if there is no conversion of ordinary to extraordinary (X) photons in the strong magnetic field expected for photons below E_{cyc} (see the discussion in Shakura et al. 2013), likely pointing

Table 4. Spin-phase-resolved spectroscopy of the overall *NuSTAR* spectrum (see Sect. 4.4).

$\Delta\phi$	N_{H} (10^{22} cm^{-2})	Γ ^(a)	E_{cut} (keV)	E_{fold} (keV)	$F_{3-78 \text{ keV}}$ ^(b) ($\text{erg cm}^{-2} \text{ s}^{-1}$)	$\chi^2_{\nu}/\text{d.o.f.}$
0.0–0.1	$6.6^{+2.6}_{-2.9}$	$1.63^{+0.18}_{-0.21}$	$7.0^{+1.2}_{-0.8}$	22^{+8}_{-6}	1.14×10^{-10}	0.994/408
0.1–0.2	9.8 ± 2.5	$1.79^{+0.15}_{-0.18}$	$7.7^{+1.5}_{-1.4}$	31^{+12}_{-8}	1.43×10^{-10}	1.117/458
0.2–0.3	$8.2^{+2.4}_{-2.2}$	$1.63^{+0.15}_{-0.14}$	$6.3^{+1.5}_{-1.0}$	36^{+17}_{-9}	1.67×10^{-10}	1.107/483
0.3–0.4	$7.6^{+1.5}_{-1.4}$	$1.72^{+0.05}_{-0.05}$	30^{+2}_{-5}	16^{+10}_{-5}	1.71×10^{-10}	0.896/506
0.4–0.5	$2.6^{+1.2}_{-2.2}$	$1.05^{+0.14}_{-0.15}$	$6.2^{+0.6}_{-0.8}$	18^{+2}_{-3}	1.84×10^{-10}	1.033/550
0.5–0.6	$6.1^{+2.2}_{-2.1}$	$1.40^{+0.17}_{-0.14}$	$6.4^{+1.3}_{-0.7}$	22^{+6}_{-4}	1.91×10^{-10}	0.926/550
0.6–0.7	$4.7^{+2.8}_{-2.3}$	$1.37^{+0.21}_{-0.22}$	$6.6^{+1.7}_{-1.7}$	22^{+8}_{-5}	1.68×10^{-10}	0.942/507
0.7–0.8	$6.1^{+2.7}_{-2.9}$	$1.49^{+0.18}_{-0.21}$	$7.2^{+1.7}_{-1.2}$	25^{+9}_{-6}	1.35×10^{-10}	0.858/443
0.8–0.9	$3.6^{+2.7}_{-2.9}$	$1.43^{+0.21}_{-0.20}$	$6.5^{+2.0}_{-1.1}$	24^{+11}_{-6}	1.15×10^{-10}	0.875/372
0.9–1.0	$3.7^{+3.0}_{-2.9}$	$1.38^{+0.20}_{-0.20}$	$6.6^{+1.2}_{-0.8}$	21^{+8}_{-5}	1.13×10^{-10}	1.047/374

Notes. ^(a)Power-law photon index. ^(b)Flux not corrected for the absorption.

to a value of E_{cyc} below three keV, meaning in the observed X-ray band 3–78 keV we observe only ordinary photons. As E_{cyc} in IGR J11215–5952 has not been measured, it may seem difficult to make robust conclusions. However, we can make use of the fact that the X-ray pulse period of IGR J11215–5952 stays constant at $P_* = 187$ s, suggesting its equilibrium value. In this case, to eliminate the unknown NS magnetic field, we can use the expression for the equilibrium period for quasi-spherical settling accretion (which appears to be the case for SFXTs, see Shakura et al. 2014 for more details),

$$P_{\text{eq}} \approx 1000 F(e) \mu_{30}^{12/11} (P_{\text{b}}/10 \text{ d}) \dot{M}_{16}^{-4/11} v_8^4 \text{ s}.$$

Here, P_{b} is the binary orbital period, $v = 10^8 v_8 \text{ cm s}^{-1}$ is the optical stellar wind velocity near the NS orbital location, and $\dot{M} = 10^{16} \dot{M}_{16} \text{ g s}^{-1}$ is the mass accretion rate related to the X-ray accretion luminosity as $L_{\text{X}} = 0.1 \dot{M} c^2$ (c is the velocity of light). The factor $F(e) < 1$ takes into account the possible reduction in the value of P_{eq} due to the orbital eccentricity. For e up to ~ 0.8 , this factor can be $F(e) \sim 0.1$ – 0.01 (Yungelson et al. 2019), depending on the (unknown in IGR J11215–5952) wind radiation acceleration parameters and orbital eccentricity.

The scattering of O-photons generated in the NS polar cap at low accretion rates mostly occurs in the freely falling matter canalised by the NS magnetic field. For a NS dipole field, the optical depth in the accreting matter for ordinary photons above E_{cyc} with the Thomson scattering cross-section is $\tau_{\text{v}} \approx 3(R_{\text{A}}/10^9 \text{ cm})^{1/2} \dot{M}_{16}$ (Lamb et al. 1973). When inserting the expression for the Alfvén radius, $R_{\text{A}} \sim 1.4 \times 10^9 \mu_{30}^{6/11} \dot{M}_{16}^{-2/11} \text{ cm}$, for the quasi-spherical settling accretion (Shakura et al. 2012), and eliminating the unknown NS magnetic field through the equilibrium spin period, we arrive at the optical depth

$$\tau_{\text{v}} \approx 3.6 (P_*/10^3 \text{ s})^{1/4} F(e)^{-1/4} (P_{\text{b}}/10 \text{ d})^{-1/4} \dot{M}_{16} v_8^{-1}.$$

For IGR J11215–5952, $P_* = 187$ s, and the binary orbital period is $P_{\text{b}} \approx 165$ days, so we get $\tau_{\text{v}} \approx 1.2 \dot{M}_{16} F(e)^{-1/4} v_8^{-1}$. We note that the uncertain reduction of P_{eq} due to orbital eccentricity $F(e) < 1$ only increases this estimate.

As the wind velocity can hardly be higher than 1000 km s^{-1} (especially considering the likely origin of the outbursts in this source near the periastron passages, i.e. closer to the optical star where the wind velocity is lower), we can conclude that at

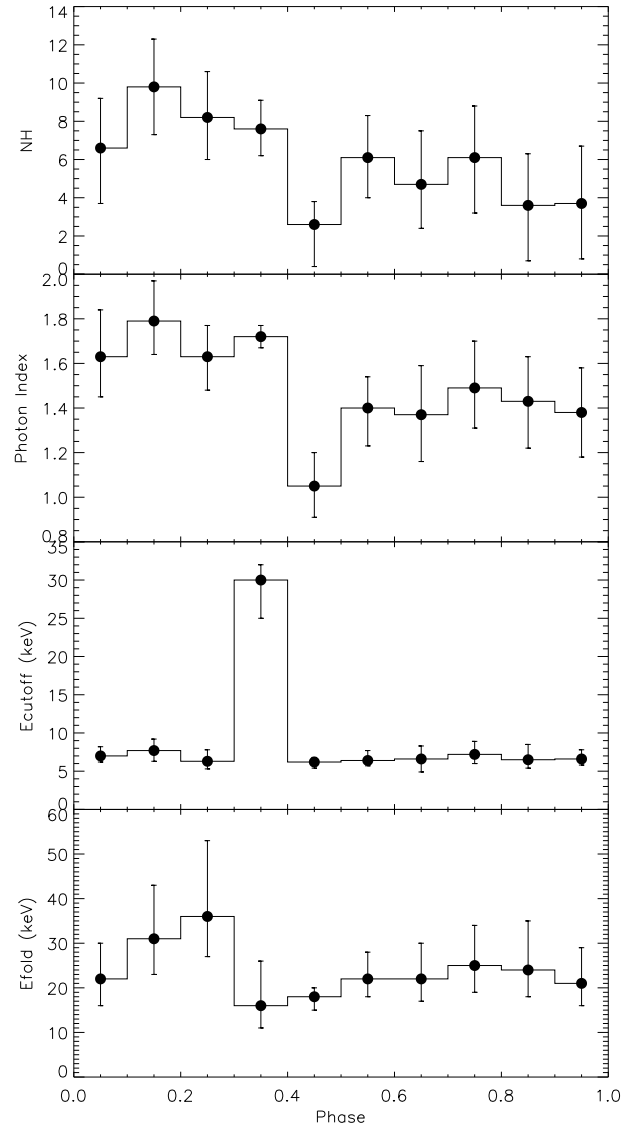


Fig. 8. Spin-phase-resolved spectroscopy over the pulse, obtained folding the whole *NuSTAR* observation on the spin period (the parameters are reported in Table 4). *Top panel:* the absorption column density is in units of 10^{22} cm^{-2} .

Table 5. Results of the spin-phase-resolved spectroscopy of intensity-selected spectra (see Sect. 4.5).

$\Delta\phi$	N_{H} (10^{22} cm^{-2})	Γ ^(a)	E_{cut} (keV)	E_{fold} (keV)	$\text{UF}_{1-100\text{keV}}$ ^(b) ($\text{erg cm}^{-2} \text{ s}^{-1}$)	$\chi^2_{\nu}/\text{d.o.f.}$
Rate < 1 count s ⁻¹						
0.0–0.1	12.6 ^{+4.8} _{-4.4}	2.15 ^{+0.14} _{-0.14}	–	–	0.90×10^{-10}	1.017/43
0.1–0.2	4.9 ^{+11.1} _{-4.9}	1.44 ^{+0.60} _{-0.35}	7.1 ^{+3.0} _{-7.1}	17 ⁺²⁵ ₋₁₂	0.71×10^{-10}	0.615/52
0.2–0.3	7.9 ^{+4.7} _{-4.3}	1.88 ^{+0.14} _{-0.14}	–	–	1.37×10^{-10}	1.017/74
0.3–0.4	4.1 ^{+7.9} _{-4.1}	1.40 ^{+0.50} _{-0.22}	6 ⁺¹³ ₋₆	29 ⁺¹⁶⁰ ₋₁₄	1.16×10^{-10}	0.977/97
0.4–0.5	10.3 ^{+5.1} _{-4.7}	1.72 ^{+0.15} _{-0.07}	23 ⁺¹⁴ ₋₁₃	15 ⁺⁴² ₋₁₃	1.21×10^{-10}	0.918/102
0.5–0.6	2.6 ^{+6.9} _{-2.6}	1.20 ^{+0.45} _{-0.34}	6.4 ^{+4.1} _{-6.4}	19 ⁺²² ₋₆	1.09×10^{-10}	0.938/107
0.6–0.7	10.7 ^{+2.9} _{-2.7}	1.98 ^{+0.08} _{-0.08}	–	–	1.30×10^{-10}	1.167/82
0.7–0.8	3.6 ^{+15.3} _{-3.6}	1.21 ^{+0.85} _{-0.62}	6.1 ^{+11.7} _{-6.1}	16 ⁺⁴⁵ ₋₆	0.58×10^{-10}	1.044/56
0.8–0.9	8.8 ^{+10.9} _{-8.8}	1.78 ^{+0.37} _{-1.14}	12.7 ^{+10.4} _{-12.7}	20 ⁺³⁵⁰ ₋₁₇	0.58×10^{-10}	0.998/41
0.9–1.0	2.1 ^{+11.0} _{-2.1}	1.38 ^{+0.63} _{-0.77}	7.4 ^{+5.1} _{-7.4}	17 ⁺¹⁰⁸ ₋₁₄	0.46×10^{-10}	0.831/34
Rate 1–3 count s ⁻¹						
0.0–0.1	7.0 ^{+5.2} _{-6.2}	1.72 ^{+0.32} _{-0.46}	7.6 ^{+2.4} _{-1.9}	22 ⁺²⁶ ₋₁₀	1.07×10^{-10}	1.029/130
0.1–0.2	5.6 ^{+5.4} _{-4.8}	1.52 ^{+0.39} _{-0.35}	6.7 ^{+3.3} _{-1.6}	19 ⁺¹⁸ ₋₆	1.40×10^{-10}	1.090/173
0.2–0.3	12.9 ^{+3.0} _{-2.8}	2.01 ^{+0.09} _{-0.09}	–	–	2.56×10^{-10}	1.106/193
0.3–0.4	4.2 ^{+2.9} _{-2.9}	1.64 ^{+0.20} _{-0.45}	14.7 ^{+26.5} _{-4.8}	47 ⁺⁴⁷ ₋₄₇	1.96×10^{-10}	0.849/194
0.4–0.5	4.3 ^{+4.5} _{-4.3}	1.06 ^{+0.29} _{-0.29}	4.2 ^{+8.5} _{-4.2}	16 ⁺⁸ ₋₄	1.69×10^{-10}	0.904/205
0.5–0.6	3.6 ^{+4.6} _{-3.6}	1.27 ^{+0.28} _{-0.30}	5.6 ^{+1.6} _{-5.6}	19 ⁺¹¹ ₋₆	1.66×10^{-10}	0.905/189
0.6–0.7	5.7 ^{+4.8} _{-5.7}	1.44 ^{+0.27} _{-0.56}	8.0 ^{+2.2} _{-8.0}	24 ⁺¹⁸ ₋₁₀	1.52×10^{-10}	0.919/168
0.7–0.8	4.9 ^{+6.7} _{-4.9}	1.40 ^{+0.44} _{-0.31}	6.5 ^{+4.9} _{-6.5}	23 ⁺⁴⁵ ₋₁₁	1.24×10^{-10}	0.720/135
0.8–0.9	2.1 ^{+6.6} _{-2.1}	1.49 ^{+0.42} _{-0.41}	6.8 ^{+12.2} _{-6.8}	30 ⁺⁷² ₋₁₃	0.99×10^{-10}	1.008/110
0.9–1.0	5.9 ^{+5.7} _{-5.9}	1.51 ^{+0.29} _{-0.65}	8.4 ^{+2.8} _{-8.4}	22 ⁺²⁰ ₋₁₂	1.01×10^{-10}	1.009/117
Rate 3–5 count s ⁻¹						
0.0–0.1	5.7 ^{+4.8} _{-4.7}	1.55 ^{+0.31} _{-0.32}	6.7 ^{+1.7} _{-1.8}	20 ⁺¹⁷ ₋₇	2.34×10^{-10}	0.872/166
0.1–0.2	13.3 ^{+3.7} _{-3.9}	2.07 ^{+0.20} _{-0.24}	8.0 ^{+5.0} _{-5.0}	70 ⁺¹⁰⁰ ₋₅₀	3.84×10^{-10}	0.998/192
0.2–0.3	10.1 ^{+4.0} _{-4.6}	1.72 ^{+0.23} _{-0.21}	7.5 ^{+3.3} _{-7.5}	40 ⁺⁷⁰ ₋₂₀	3.70×10^{-10}	1.032/201
0.3–0.4	6.9 ^{+3.0} _{-2.9}	1.69 ^{+0.09} _{-0.09}	30 ⁺¹⁰ ₋₃₀	20 ⁺⁶⁰⁰ ₋₂₀	3.41×10^{-10}	0.959/201
0.4–0.5	2.1 ^{+4.5} _{-2.1}	1.03 ^{+0.33} _{-0.23}	6.5 ^{+2.2} _{-1.1}	18.6 ^{+5.4} _{-4.4}	3.23×10^{-10}	0.951/220
0.5–0.6	9.2 ^{+3.6} _{-3.5}	1.65 ^{+0.20} _{-0.23}	7.9 ^{+1.9} _{-1.7}	34 ⁺²⁷ ₋₁₁	4.00×10^{-10}	1.054/243
0.6–0.7	8.4 ^{+4.9} _{-4.2}	1.53 ^{+0.35} _{-0.31}	6.3 ^{+5.6} _{-6.3}	26 ⁺²⁸ ₋₉	3.19×10^{-10}	0.717/205
0.7–0.8	6.0 ^{+4.0} _{-5.0}	1.55 ^{+0.26} _{-0.35}	7.6 ^{+2.6} _{-1.9}	27 ⁺²² ₋₁₀	2.58×10^{-10}	0.794/169
0.8–0.9	4.9 ^{+5.3} _{-4.9}	1.49 ^{+0.34} _{-0.34}	6.6 ^{+2.9} _{-6.6}	26 ⁺²⁹ ₋₁₁	2.09×10^{-10}	0.912/142
0.9–1.0	3.9 ^{+4.7} _{-3.9}	1.43 ^{+0.28} _{-0.18}	6.0 ^{+2.2} _{-6.0}	27 ⁺¹⁸ ₋₁₀	2.25×10^{-10}	0.960/153
Rate > 5 count s ⁻¹						
0.0–0.1	6.6 ^{+4.7} _{-6.6}	1.66 ^{+0.24} _{-0.67}	9.0 ^{+2.3} _{-9.0}	26 ⁺²⁵ ₋₉	4.58×10^{-10}	0.993/130
0.1–0.2	7.7 ^{+5.5} _{-6.7}	1.59 ^{+0.28} _{-0.46}	8.1 ^{+2.6} _{-2.1}	29 ⁺⁴⁶ ₋₁₅	4.62×10^{-10}	1.098/131
0.2–0.3	13.8 ^{+4.2} _{-3.8}	1.94 ^{+0.12} _{-0.12}	–	–	6.04×10^{-10}	0.906/114
0.3–0.4	5.7 ^{+5.4} _{-5.7}	1.41 ^{+0.33} _{-0.39}	8.7 ^{+4.9} _{-8.7}	40 ⁺⁷⁰ ₋₂₀	5.37×10^{-10}	1.098/120
0.4–0.5	5.0 ^{+5.9} _{-5.0}	1.21 ^{+0.41} _{-0.35}	6.4 ^{+12.4} _{-6.4}	30 ⁺⁶⁰ ₋₁₂	5.81×10^{-10}	0.991/143
0.5–0.6	7.1 ^{+5.1} _{-4.8}	1.40 ^{+0.38} _{-0.31}	6.3 ^{+4.9} _{-6.3}	23 ⁺²² ₋₇	6.56×10^{-10}	0.933/178
0.6–0.7	5.2 ^{+3.9} _{-3.9}	1.51 ^{+0.17} _{-0.76}	11.2 ^{+1.7} _{-11.2}	24 ⁺¹² ₋₆	6.33×10^{-10}	1.084/174
0.7–0.8	6.0 ^{+5.3} _{-5.7}	1.38 ^{+0.35} _{-0.45}	7.0 ^{+2.4} _{-1.8}	19 ⁺¹⁵ ₋₇	4.61×10^{-10}	1.009/148
0.8–0.9	8.5 ^{+5.6} _{-4.2}	1.63 ^{+0.29} _{-0.56}	8.4 ^{+5.2} _{-8.4}	36 ⁺⁵⁷ ₋₂₀	4.53×10^{-10}	0.962/125
0.9–1.0	3.1 ^{+6.2} _{-3.1}	1.30 ^{+0.36} _{-0.33}	7.0 ^{+1.6} _{-1.6}	16 ⁺⁹ ₋₄	3.66×10^{-10}	0.958/119

Notes. ^(a)Power-law photon index. ^(b)Flux corrected for the absorption.

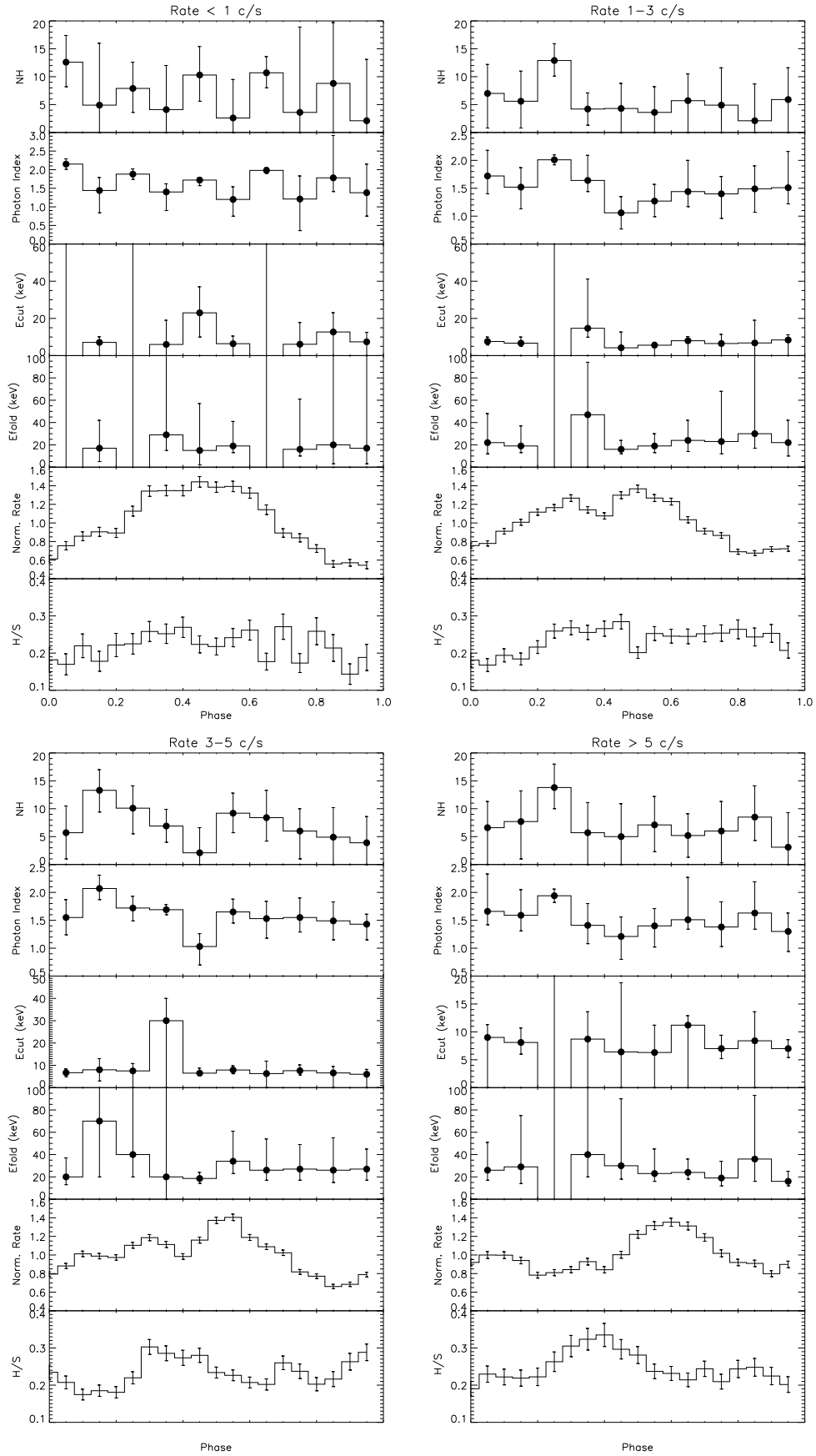


Fig. 9. Spin-phase-selected spectroscopy in four intensity states, as reported in Table 5: from top to bottom, from left to right, the intensity states are below 1 count s^{-1} , 1–3 count s^{-1} , 3–5 count s^{-1} , and above 5 count s^{-1} . Bottom panels: the intensity-selected pulse profiles (3–78 keV) are reported, together with the hardness ratio (H/S) of hard pulse profile ($H = 12$ –78 keV) to the soft pulse profile ($S = 3$ –12 keV).

high X-ray luminosities of IGR J11215–5952 observed (a few 10^{36} erg s $^{-1}$), the optical depth in the accreting matter above the NS polar caps is indeed substantial for the O-photon scatterings.

Therefore, the properties 1 and 2 (Sect. 5.3) of the observed pulse profiles in IGR J11215–5952 listed above can be naturally explained by this model. We note that the pulse profile evolution found in the previous *XMM–Newton* observations (see Fig. 3 in Sidoli et al. 2007) can also be explained by the scattering in the accreting material as described above. Indeed, in those observations, the mean hardness ratio HR does not change at different X-ray luminosity states either, and the pulse shape demonstrates a strong scattering absorption feature at the bright state of the source.

In the proposed model, we eliminated the NS magnetic field by expressing it through the observed NS spin and binary orbital period. The non-observation of the CRSF feature in the *NuSTAR* spectrum (property 4 above) may imply either a rather strong NS magnetic field, $\mu_{30} \gtrsim 10$, or, oppositely, a weak field, $\mu_{30} < 1$. The low NS magnetic field seems more likely because, for the high field, the vacuum polarisation effects leading to the O-X photon transformation can be significant. The scattering cross-section of the X-photons is reduced below E_{cyc} : $\sigma_{\perp} \sim \sigma_{\text{T}}(E/E_{\text{cyc}})^2$, and so is the optical depth in the accreting matter above the NS polar caps. However, the scattering absorption feature in the pulse profiles signaling a large scattering optical depth at high luminosities has been observed in both *XMM–Newton* (Sidoli et al. 2007) and the *NuSTAR* pulse profiles (this paper).

Scattering of the O-photons in the optically thick accreting matter above the polar caps also explains why the pulse fraction increases with decreasing X-ray luminosity (see Fig. 6, property 3 listed in Sect. 5.3). In terms of our model, we can make use of the fact the *PF* stays almost constant down to the X-ray luminosity $L_{\text{PF}} \sim 6 \times 10^{35}$ erg s $^{-1}$ (Fig. 6). By identifying this luminosity with the critical L_{\dagger} corresponding to $\tau_{\text{v}} \sim 1$ for scattering, we can evaluate the required NS magnetic field $\mu_{30} \sim (3 \times 10^{35} \text{ erg s}^{-1} / L_{\text{PF}})^{10/3} \approx 0.1$. This estimate supports our conjecture that the NS magnetic field in IGR J11215–5952 is indeed lower than the canonical NS value 10^{12} G. We also note that the expected NS equilibrium period at the settling quasi-spherical accretion for this low NS magnetic field could be exactly in the right range of a few 100 s for the binary period of 165 days, the typical X-ray luminosities, and stellar wind velocities¹.

We conclude that all significant properties 1–4 of the gross X-ray pulse profile evolution of IGR J11215–5952 revealed by the X-ray spectroscopy of the *NuSTAR* observations, might be explained coherently by the scattering of O-photons, which are mostly produced near the NS surface at low accretion rates in the accreting plasma above one NS polar cap. However, in this case, a detailed pulse profile change with X-ray luminosity shown in Fig. 5 remains unexplained.

5.3.2. Photon scatterings above two NS magnetic poles

The X-ray pulse evolution of IGR J11215–5952 is most clearly visible in Fig. 5. The most striking feature is that the hard 12–78 keV pulse shape changes rather little with X-ray luminosity (the middle panel) compared to the 3–12 keV pulse (the upper panel). The latter reveals a more complicated behaviour,

¹ An attempt to estimate the NS magnetic field using the NS equilibrium spin period only, $P_{\text{eq}} \sim \mu^{12/11} v_w^4$, would fail due to an unknown value of the stellar wind velocity.

from a sine-like shape at low luminosities transforming through a two-hump form at mediate fluxes to an almost single-pulse narrower shape at the highest luminosities. Here, the left peak at the phase ~ 0.3 almost disappears, leaving only the right peak at the phase ~ 1.5 –1.6.

This picture may suggest that, in fact, we are observing two peaks from two NS magnetic poles (we call them south (S) and north (N) poles corresponding to the left and right peaks, respectively), which is generally more likely. In the hard 12–78 keV range, these two peaks almost merge to form the apparently one sine-like peak, although traces of two S and N peaks could be still discerned.

The suppression of the low-energy S peak with increasing X-ray luminosity suggests that the 3–12 keV emission gets scattered in the accreting matter, as described above, which is possible if the O-X photon transformation occurs below 12 keV. This is possible if the magnetic field of the S pole is about 10^{12} G corresponding to CRFS at $E_{\text{cyc}}^{\text{S}} \sim 12$ keV in the spectrum. The scattering of X-photons from the S-beam below E_{cyc} explains the enhanced hardness ratio at the phase 0.3–0.4 at the highest luminosities (Fig. 9, the bottom right panel).

As for the N-beam that remains visible and non-scattered at all luminosities (the upper panel of Fig. 5), it may be due to a lower magnetic field of the N pole, such that $E_{\text{cyc}}^{\text{N}} < 3$ keV, meaning that all generated O-photons are above the CRSF energy. As no strong transformation of the N-pole has occurred up to luminosities $\sim 10^{36}$ erg s $^{-1}$, we can apply the same model as in the previous sub-section to identify the critical luminosity L_{\dagger} , corresponding to the unit optical depth, to estimate the NS magnetic field at this pole to be a few times 10^{10} G.

The different magnetic field at NS poles is not unusual for X-ray pulsars due to the possible non-dipole structure of the NS magnetic field (see the discussion of Her X-1 in Shakura et al. 1991; Postnov et al. 2013). If this is true for IGR J11215–5952, it can indeed be difficult to find the CRSF signature in the X-ray spectrum, because at high luminosities, the O-photons from the N pole with a lower magnetic field dominate. To verify this model, searches for the putative 12–17 keV CRSF should be made only at low luminosities, which is challenging. Here, the future X-ray polarisation observations may be crucial.

6. Conclusions

The *NuSTAR* observation of the SFXT pulsar IGR J11215–5952, performed at the expected peak of the June 2017 outburst, showed an ample range of variability of the flaring activity, reaching two orders of magnitude in the energy range 3–78 keV. This observation allowed us to perform the most sensitive search of CRSFs in the hard X-ray spectrum of IGR J11215–5952 to date, with null results. We were not able to confirm the hint of a 17 keV CRSF obtained during a *NuSTAR* observation of the 2016 outburst, at a similar X-ray flux (Sidoli et al. 2017). Thus, the issue of the direct measurement of the NS magnetic field in this SFXT pulsar remains open. However, indirect inference regarding the possible value of the NS magnetic field in this source can be made from the analysis of the spectroscopic properties of X-ray pulses and their change with X-ray luminosity (see Sect. 5.3).

The pulsar spin period displayed no significant variability with respect to any previous measurements, with a profile evolving with both the energy and the luminosity. The pulsed fraction increased with the energy and decreased with the source X-ray intensity, within the range of luminosities probed during

this observation (from an average $L_{1-100\text{keV}} = 3.5 \times 10^{35} \text{ erg s}^{-1}$ to $L_{1-100\text{keV}} = 2.4 \times 10^{36} \text{ erg s}^{-1}$).

The long spin period and the low X-ray luminosity in the outburst led us to discuss the source properties within the framework of the quasi-spherical settling accretion model as the physical mechanism driving the SFXT phenomenology (see Shakura et al. 2014 for more details). The evolution of the pulse profile at different X-ray luminosities is a further observational fact supporting this model. Indeed, the observed 187 s NS spin period in IGR J11215–5952, perceived as the equilibrium spin period at the quasi-spherical settling accretion, implies a rather moderate NS magnetic field. This conjecture is simultaneously supported by an insignificant pulse phase shift, the constant hardness ratio of the X-ray light curve at different luminosities, the increase in the pulse fraction at low luminosity, and the null detection of CRSF. These properties can be naturally explained by the scattering of ordinary X-ray photons in the accreting matter above NS polar caps. The estimated surface NS magnetic field in this model is about 10^{11} G.

Alternatively (see Sect. 5.3.2), the detailed pulse profile evolution can be understood in terms of accretion onto NS magnetic poles with the surface magnetic field strength differing by one order of magnitude (e.g. due to a non-dipoles magnetic field structure). In this case, the strongest NS magnetic field was estimated to be around the standard NS field value $\sim 10^{12}$ G. The null detection CRSF in our observations at the brightest state may be due to the most of photons generated at this state being ordinary photons above the cyclotron frequency from the pole with lower magnetic field. The model can be further tested by searches for CRSF at low energies and by future X-ray polarisation observations.

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