

A Gemini-NIFS view of the merger remnant NGC 34

J. C. Motter[®], ^{1*} R. Riffel[®], ^{1*} T. V. Ricci, ² R. A. Riffel[®], ³ T. Storchi-Bergmann, ¹ M. G. Pastoriza, ¹ A. Rodriguez-Ardila[®], ⁴ D. Ruschel-Dutra[®], ⁵ L. G. Dahmer-Hahn, ^{4,6} N. Z. Dametto⁷ and M. R. Diniz [®]

Accepted 2021 July 6. Received 2021 June 25; in original form 2021 April 30

ABSTRACT

The merger remnant NGC 34 is a local luminous infrared galaxy (LIRG) hosting a nuclear starburst and a hard X-ray source associated with a putative, obscured Seyfert 2 nucleus. In this work, we use adaptive optics assisted near-infrared (NIR) integral field unit observations of this galaxy to map the distribution and kinematics of the ionized and molecular gas in its inner $1.2\,\mathrm{kpc}\times1.2\,\mathrm{kpc}$, with a spatial resolution of 70 pc. The molecular and ionized gas kinematics is consistent with a disc with projected major axis along a mean PA = $-9^{\circ}.2\pm0^{\circ}.9$. Our main findings are that NGC 34 hosts an AGN and that the nuclear starburst is distributed in a circumnuclear star formation ring with inner and outer radii of ≈ 60 and 180 pc, respectively, as revealed by maps of the [Fe II]/Pa β and H₂/Br γ emission-line ratios, and corroborated by PCA tomography analysis. The spatially resolved NIR diagnostic diagram of NGC 34 also identifies a circumnuclear structure dominated by processes related to the stellar radiation field and a nuclear region where [Fe II] and H₂ emissions are enhanced relative to the hydrogen recombination lines. We estimate that the nuclear X-ray source can account for the central H₂ enhancement and conclude that [Fe II] and H₂ emissions are due to a combination of photoionization by young stars, excitation by X-rays produced by the AGN and shocks. These emission lines show nuclear, broad, blueshifted components that can be interpreted as nuclear outflows driven by the AGN.

Key words: galaxies: individual: NGC 34-galaxies: individual: NGC 17-galaxies: individual: Mrk 938-galaxies: ISM-galaxies: nuclei-galaxies: starburst-infrared: galaxies.

1 INTRODUCTION

Local Universe galaxies display a wide range of luminosities, sizes, gas and stellar distribution and kinematics, that are the outcome of a process that lasts \sim 13.8 Gyr. In the hierarchical scenario of galaxy formation and evolution, early-type galaxies form as a result of the merger of spiral galaxies whose discs are disrupted as the merging evolves (e.g.: White & Rees 1978; Hopkins et al. 2008; Deeley et al. 2017). In cases where the interacting galaxies are gasrich ('wet-merger'), the migration of vast amounts of gas towards the innermost regions of the merging system can lead to the onset of starbursts and/or active galactic nuclei (AGN) phenomena. However, this poses several challenges for the understanding of the fate of the interacting gas during these galaxies encounters; at the same time that gas is required either to form stars or to feed the AGN, it can also be expelled from the merger remnant by galactic winds related with these processes. Moreover, it is well known that star formation (SF) on galactic scales may be quenched, suppressed or triggered as a result of the AGN feedback (Harrison 2017), thus raising the

question if the AGN feeding happens simultaneously with the SF (Kawakatu & Wada 2008), follows it during a post-starburst phase (Davies et al. 2007; Riffel et al. 2009) or lack any association with recent SF (Sarzi et al. 2007).

Luminous and ultraluminous infrared galaxies (LIRGs and ULIRGs, respectively) emit the bulk of their energy in the infrared (IR). For LIRGs, IR luminosities are in the range $11.0 \le \log (L_{\rm IR}/L_{\odot})$ \leq 12.0, while for ULIRGs, $\log(L_{\rm IR}/L_{\odot}) \geq$ 12.0 (Sanders et al. 2003). These objects harbor a strong nuclear starburst and/or an AGN (Sanders & Mirabel 1996; Lonsdale, Farrah & Smith 2006), and most of local ULIRGs, and LIRGs with $L_{\rm IR} > 10^{11.5} L_{\odot}$ are often associated with gas-rich merging systems (Hung et al. 2014, for a recent review on the properties of LIRGs, see Pérez-Torres et al. 2021). Therefore, these objects have become increasingly important for galaxy evolution studies since they provide a laboratory for understanding the starburst/AGN feedback processes that shape the evolution of a merger remnant. The main power source of U(LIRGs) is the heating of dust by a starburst and/or an AGN, but the relative contribution of each process to the observed IR luminosities is still under debate. This happens mainly due to difficulties in carrying out high-resolution optical/IR studies of these objects as a consequence of the heavy dust obscuration towards their nuclear regions and to

¹Departamento de Astronomia, Universidade Federal do Rio Grande do Sul, Porto Alegre, RS, 91501-970, Brazil

²Universidade Federal da Fronteira Sul, Campus Cerro Largo, RS, 97900-000, Brazil

³Departamento de Física, Centro de Ciências Naturais e Exatas, Universidade Federal de Santa Maria, Santa Maria, RS, 97105-900, Brazil

⁴Laboratório Nacional de Astrofísica, Itajubá, MG, 37500-000, Brazil

⁵Departamento de Física, Centro de Ciências Físicas e Matemáticas, Universidade Federal de Santa Catarina, Florianópolis, SC, 88040-900, Brazil

⁶Shanghai Astronomical Observatory, Chinese Academy of Sciences, 80 Nandan Road, Shanghai 200030, China

⁷Centro de Astronomía (CITEVA), Universidad de Antofagasta, Antofagasta, 02800, Chile

^{*} E-mail: jujucmotter@gmail.com (JCM); riffel@ufrgs.br (RR)

the fact that they are more numerous at $z \ge 1$ (Sanders & Mirabel 1996; Lonsdale et al. 2006). Therefore, it is clear that in order to map in detail their stellar and gas properties, it is necessary to search for local prototypes of the (U)LIRGs population. Observations indicate that, at least for local LIRGs, the bulk of their IR luminosites is dominated by starburst emission (Alonso-Herrero et al. 2012), and that the AGN contribution to the observed IR luminosities increases with the IR luminosity of the system, and may dominate the energy budget for objects with $L_{\rm IR} > 10^{12.5} L_{\odot}$ (Veilleux et al. 2009; Nardini et al. 2010).

The galaxy NGC 34 (NGC 17, Mrk 938), z = 0.0196 (Rothberg & Joseph 2006) and at a distance of $\approx 82 \,\mathrm{Mpc}$ ($H_0 =$ 67.4 km s⁻¹ Mpc⁻¹), is a local representative of the LIRGs class $(\log(L_{\rm IR}/\rm L_{\odot}) = 11.61$, Chini, Kruegel & Steppe 1992). It hosts a dominant central starburst (e.g.: Riffel, Rodríguez-Ardila & Pastoriza 2006; Schweizer & Seitzer 2007; Riffel et al. 2008b; Dametto et al. 2014) with inferred star formation rates (SFRs) in the range 50-90 $M_{\odot} \text{yr}^{-1}$ (Prouton et al. 2004; Valdés et al. 2005). However, evidence on AGN activity has been extensively debated. Based on nearinfrared (NIR) studies in the range 0.8–2.4 μ m, Riffel et al. (2006) found that NGC 34 displays a poor emission-line spectrum and that the continuum emission is dominated by stellar absorption features. Therefore, they classified NGC 34 as a starburst galaxy. In the optical domain, Mulchaey, Wilson & Tsvetanov (1996) carried out an imaging survey of the $[O III]\lambda 5007$ and $H\alpha + [N II]\lambda\lambda 6548$, 6543 emission lines for a sample of early-type Seyfert galaxies and found that NGC 34 not only is a week [O III] emitter compared to most of the Seyferts in their sample, but also that H α emission is strong over the entire galaxy, indicating that the gas ionization is not related to any AGN. However, Yuan, Kewley & Sanders (2010) and Brightman & Nandra (2011b) used Baldwin, Phillips and Terlevich (BPT) diagnostic diagrams (Baldwin, Phillips & Terlevich 1981; Kewley et al. 2001, 2006) to confirm that NGC 34 hosts a Seyfert 2 nucleus. Despite the highly controversial nature of the NGC 34 nuclear spectrum, X-ray observations provide compelling evidence for the presence of an obscured AGN in its central regions (e.g.: Guainazzi, Matt & Perola 2005; Brightman & Nandra 2011a; Esquej et al. 2012). Esquej et al. (2012) found that the NGC 34 X-ray luminosity in the 2-10 keV energy range $(L_{\rm X} = 1.4^{+0.3}_{-0.2} \times 10^{42} {\rm erg \, s^{-1}})$ is too high to be solely due to the nuclear starburst, and is, therefore, dominated by an AGN.

The detailed study of the optical morphological properties of NGC 34 presented by Schweizer & Seitzer (2007) shows that this galaxy features a red nucleus, with a nuclear starburst confined to a radius ≤ 1 kpc; a young and blue central exponential disc at a position angle (PA) of -9°, a system of young massive star clusters and a pair of unequal tidal tails indicative of the merger of two former gas-rich disc galaxies. In addition, Schweizer & Seitzer (2007) also reported that blueshifted Na I D lines reveal that the inner regions of NGC 34 drive a strong outflow of cool, neutral gas with a mean velocity of $-620\pm60\,\mathrm{km\,s^{-1}}$ that could be due to the concentrated starburst and/or the hidden AGN.

The galaxy NGC 34 has also been probed in the radio and submillimetre regimes. Fernández et al. (2010) used 21 cm Very Large Array (VLA) observations to map the radio continuum emission and the H I distribution and kinematics. They found that the radio continuum emission structure consists of an extended nuclear component that is dominated by the central starburst, and an outer, extra–nuclear diffuse extended component in the shape of two radio lobes that could be evidence for past AGN activity or due to a starburst-driven superwind. The authors detected a broad H I absorption profile with both blueshifted and redshifted velocities that could be explained by the presence of a circumnuclear disc (CND) of neutral and molecular

gas. The existence of the CND was later confirmed by Fernández et al. (2014). They detected a rotating CO disc of $2.1 \,\mathrm{kpc}$ in diameter using Combined Array for Research in Millimeter-wave Astronomy (CARMA) observations of the CO(1–0) transition at 115 GHz. An even more compact molecular rotating disc with a size of 200 pc was later detected by Xu et al. (2014) using Atacama Large Millimeter Array (ALMA) observations of the CO(6–5) emission line (rest-frame frequency = $691.473 \,\mathrm{GHz}$).

Finally, Mingozzi et al. (2018) used archival *Herschel* and ALMA observations of multiple CO transitions along with X-ray data from *NuSTAR* and *XMM*–*Newton* to perform the modelling of the CO spectral line energy distribution (SLED) – that is, the modelling of the luminosities of the CO lines as a function of their upper rotational levels. They found that a combination of a cold and diffuse photodissociation region (PDR) to account for the low-*J* transitions and a warmer and denser X-ray dominated region (XDR) to account for the high-J transitions was necessary to properly fit the CO line luminosities, and concluded that AGN contribution is significant in heating the molecular gas in NGC 34.

In this work, we use adaptive optics (AO) assisted NIR integral field unit (IFU) observations of the galaxy NGC 34 to map the distribution and kinematics of the ionized and molecular gas distributed in the inner $1.2\,\mathrm{kpc}\times1.2\,\mathrm{kpc}$. Our main goal is to investigate the nature of the NGC 34 NIR emission-line spectrum in order to put tighter constraints on the presence of an AGN in its centre. Starburst phenomena, as in the case of NGC 34, are known for residing in dusty systems; therefore, NIR observations are ideally suited for probing the central regions of these objects since they can pierce through highly obscured regions.

This paper is structured as follows. In Section 2, we present the observations and data reduction procedures. In Section 3, we describe our data analysis. Our results are shown in Section 4 followed by the discussion in Section 5. We summarize our findings in Section 6. At the redshift of z=0.0196, 1arcsec corresponds to 411 pc for the adopted cosmology: $H_0=67.4\,\mathrm{km\,s^{-1}\,Mpc^{-1}}$, $\Omega_\mathrm{M}=0.31$, and $\Omega_\Lambda=0.69$ (Planck Collaboration VI 2020). The NGC 34 systemic velocity is $v_\mathrm{sys}=5881\,\mathrm{km\,s^{-1}}$.

2 OBSERVATIONS AND DATA REDUCTION

Data for this work were obtained with the Near-Infrared Integral Field Spectrograph (NIFS) of the Gemini North Telescope under the Gemini Science Programme GN-2011B-Q-71 (PI: Rogério Riffel) using the ALTitude conjugate Adaptive optics for the InfraRed (ALTAIR) system. The NIFS instrument provides a field of view (FoV) of $3.0\times3.0~{\rm arcsec^2}$ that is sampled by spatial pixels (spaxels) with dimensions of $0.103\times0.043~{\rm arcsec^2}$ as a result of its optical and geometrical properties.

Observations of NGC 34 were carried out in 2011 September 9 and 24 in the K_l (1.99–2.40 μ m) and J (1.15–1.33 μ m) bands, respectively. Eight on-target (T) and four sky (S) exposures of 350 s were obtained in each band following the sequence TST. Observations of standard stars for telluric absorption removal and flux calibration, as well as flat-field, darks, Ronchi-flat, and arc-lamp (Ar and ArXe for the J and K_l bands, respectively) were obtained for the data reduction process. The spectral resolution for both bands is $20~{\rm km~s^{-1}}$ as measured by the full width at half maximum (FWHM) of the Ar and ArXe lines.

The data reduction was performed using the IRAF (Image Reduction and Analysis Facility) environment along with standard reduction scripts made available by the Gemini team. The data reduction process includes the trimming of the images, sky subtraction,

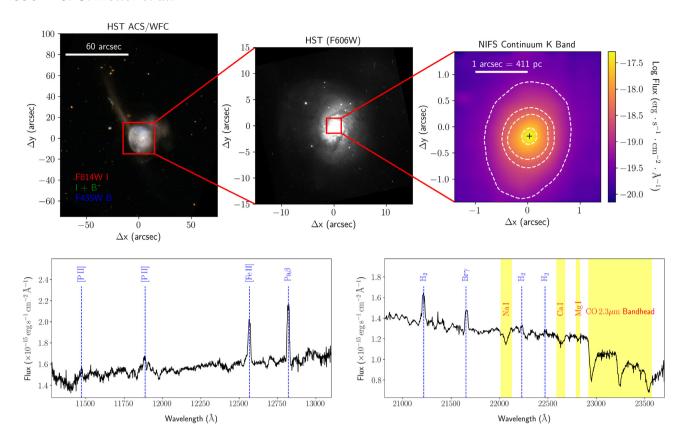


Figure 1. Top left-hand panel: NGC 34 seen by the *HST* ACS/WFC instrument (Original image by NASA, ESA, the Hubble Heritage (STScI/AURA)-ESA/Hubble Collaboration, and A. Evans (University of Virginia, Charlottesville/NRAO/Stony Brook University), axes and text were added by the authors). Top middle panel: optical image of NGC 34 obtained by the *HST* WFPC2 with the filter F606W (Malkan et al. 1998) with the NIFS FoV indicated by the red square. Top right-hand panel: average *K*-band NIFS continuum image. The black cross marks the nucleus of the galaxy. The colour bar shows the fluxes in units of ergs s⁻¹ cm⁻² Å ⁻¹. White, dashed contours are drawn at 1, 5, 10, and 50 per cent of the continuum peak. Bottom panel: integrated NIFS *J*- (left-hand panel) and *K*- (right-hand panel) band spectra over the whole FoV. The blue vertical dashed lines show the main NIR emission lines. The yellow shaded areas show the absorption features.

flat-field, bad pixel and spatial distortion corrections, wavelength calibration, telluric absorption removal, and flux calibration by fitting a blackbody function to the spectrum of the telluric standard star. Individual reduced data cubes were constructed with spaxels of $0.05 \times 0.05 \, \rm arcsec^2$ and were median combined in each band to a single data cube.

Fully calibrated and redshift corrected data cubes were obtained following the treatment procedures for Gemini–NIFS data cubes presented by Menezes, Steiner & Ricci (2014), including the re-sampling of the image to spaxels with dimensions $0.021 \times 0.021 \, \mathrm{arcsec^2}$. High-frequency spatial noise was removed using a Butterworth filter with a cut-off frequency of $0.22F_{\mathrm{NY}}$ (where F_{NY} is the Nyquist frequency) for the *J*-band cube and $0.33F_{\mathrm{NY}}$ for the *K*-band cube, both with a filter order n=2. Then, we removed low-frequency instrumental fingerprints using principal component analysis (PCA) tomography (Steiner et al. 2009).

Our final data cubes have FoVs of $\approx 3.0 \times 3.0 \,\mathrm{arcsec^2}$ corresponding to $1.2 \,\mathrm{kpc} \times 1.2 \,\mathrm{kpc}$ at the galaxy, and spectral ranges of 1.128-1.310 and $2.081-2.370 \,\mu\mathrm{m}$ in the J and K bands, respectively. The spatial resolution corresponding to the FWHM of the brightness profile of the telluric standard star is $0.17 \,\mathrm{arcsec}$ ($\approx 70 \,\mathrm{pc}$) for both the J and K bands.

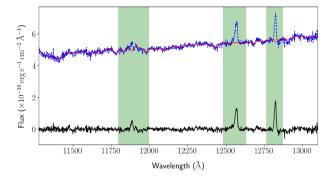
In the top left-hand panel of Fig. 1, we show an image of NGC 34 obtained by the Advance Camera for Surveys/Wide Field Camera (ACS/WFC) on board the *Hubble Space Telescope* (*HST*) with the

filters F814W and F435W. We show in the middle panel of Fig. 1 an HST Wide Field and Planetary Camera 2 (WFPC2) image of NGC 34 obtained with the filter F606W (Malkan, Gorjian & Tam 1998) along with a K-band NIFS continuum image obtained from the average flux in the $2.180-2.195~\mu m$ wavelength range (top right-hand panel). Integrated NIFS J- and K-band spectra over the entire FoV are shown in the bottom panels of Fig. 1. In the J band, the main NIR emission lines that can be seen in our spectra are $[PII]\lambda 11\,470~\text{Å}$, $[PII]\lambda 11\,886~\text{Å}$, $[FeII]\lambda 12\,570~\text{Å}$, and $Pa~\beta$. In the K band, the $H_2~\lambda 21\,218~\text{Å}$, $Br~\gamma$, $H_2~\lambda 22\,230~\text{Å}$, and $H_2~\lambda 22\,470~\text{Å}$ emission lines can be seen, in addition to the Na I $2.20~\mu m$, Ca I $2.26~\mu m$, Mg I $2.28~\mu m$, and CO $2.3~\mu m$ absorption features.

3 DATA ANALYSIS

3.1 Spectral synthesis

In order to study the nature of the gas emission lines present in the observed spectra of active/starburst galaxies, first, one needs to obtain spectra that are free from contamination caused by other components of the galaxy, such as the stellar population, dust content and AGN continuum (accretion disc). This is usually done through the application of the spectral synthesis technique, which requires a base of theoretical and/or empirical stellar spectra models to account for the other galaxy components and a code that quantifies the



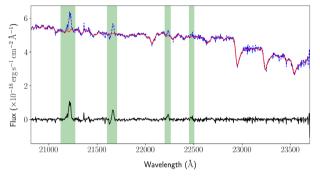


Figure 2. Examples of the PPXF fitting results for J (top panel) and K (bottom panel) bands at the spaxel corresponding to the nucleus of the galaxy. Blue dashed, solid red, and solid black lines indicate the observed, fitted stellar model, and residual (gas) spectra, respectively. The light green, vertical, shaded areas indicate masking intervals for the fit.

contribution of all these elements to the observed spectra. In this work, we used the NASA InfraRed Telescope Facility (IRTF) Spectral Library (Cushing, Rayner & Vacca 2005; Rayner, Cushing & Vacca 2009) and the Penalized Pixel-Fitting (PPXF; Cappellari 2017) code to obtain pure emission-line (gas) spectra for the galaxy NGC 34 at the *J* and *K* bands.

The IRTF Spectral Library (R = 2000) contains observed spectra in the $0.8-5\mu m$ wavelength range for 210 cool stars, with mostly near-solar metallicities, and spectral types between F and M and luminosity classes between I and V, in addition to some AGB, carbon, and S stars (see Riffel et al. 2015, for a similar application). We used this set of observed spectra as an input template for the PPXF package to fit the continuum of the NGC 34 observed spectra in every spaxel of the data cubes. The fitting was performed separately for the J and K bands and the emission lines were masked before the fit. The stellar features were fitted using Gauss-Hermite profiles and we used only multiplicative polynomials to prevent changes in the line strength of the absorption features in the templates as recommended by Cappellari (2017). Examples of the gas spectra obtained following this procedure are shown in Fig. 2 for the J and K bands at the spaxel corresponding to the nucleus of the galaxy.

3.2 Emission-line fitting

We fitted the profiles of the emission lines using the package IFSCUBE (Ruschel-Dutra & de Oliveira 2020), which is a PYTHON based package designed to carry out analysis of data cubes of integral field spectroscopy. IFSCUBE provides a set of tasks that allows the simultaneous fitting by multiple Gaussian curves or Gauss—Hermite series emission-line profiles in velocity space. It also provides a

Monte Carlo implementation to estimate fitting uncertainties pixelby-pixel.

We used Gaussian functions to fit separately, in individual IFSCUBE runs, each one of the emission lines seen in Fig. 1, thus obtaining their amplitudes, radial velocities, velocity dispersion (σ) and integrated fluxes. Pixels where the peak of a given emission line compared to the standard deviation of the adjacent continuum was smaller than five were masked during the fit. We also included a polynomial of order one to account for the contribution of the underlying residual continuum to the observed emission. The [PII] $\lambda 11886 \,\text{Å}$, Pa β , Br γ , H₂ λ 22 230 Å, and H₂ λ 22 470 Å emission lines were well fitted by a single Gaussian component. In the cases of the [Fe II] $\lambda 12570 \,\text{Å}$ and H₂ $\lambda 21218 \,\text{Å}$ emission lines, we started by fitting them with a single Gaussian, however this procedure was not able to account for the presence of broad, blue wings in a region of radius ≈ 0.5 arcsec around the nucleus of the galaxy; therefore, we included a second Gaussian to fit their profiles. We fixed the radial velocities of these components at -420 and -250 km s⁻¹ and their velocity dispersions at 350 and 215 km s⁻¹ for the [Fe II] $\lambda 12570$ Å and H₂ $\lambda 21218$ Å emission lines, respectively, as determined from the fitting of the nuclear spectra, integrated in an aperture of $0.2 \times 0.2 \,\mathrm{arcsec^2}$, also using the package IFSCUBE. However, since these components are fitted simultaneously with their main narrow counterparts, a further threshold needs to be applied to claim their detection (see Section 4.1). We note that we did not detect broad, blueshifted components associated with the other H₂ emission features. This could be due to the fact the H_2 emissions at $\lambda 22\,230\,\text{Å}$ and $\lambda 22\,470\,\text{Å}$ are much weaker compared to the emission at $\lambda 21218 \,\text{Å}$ or a result of different excitation mechanisms of the H₂ molecule, as will be discussed in Section 5.3.

We point out that while fitting the [P II] $\lambda 11$ 886Å emission line, we noticed a neighbouring emission feature at the wavelength 11 910 Å corresponding to a forbidden transition of [Ni II]. For this reason, we fitted these lines simultaneously, although we kept their kinematics independent. The [Ni II] $\lambda 11$ 910Å emission is confined to a region of radius smaller than 0.3 arcsec in the centre of the FoV and it does not affect the fitting of the [P II] $\lambda 11$ 886Å profile.

We show in Fig. 3 the resulting fits for the main emission lines for the spaxel corresponding to the nucleus of the galaxy. From these example fits, one can see that the broad, blueshifted components accompanying the [Fe II] $\lambda12\,570\,\mbox{\normalfont\AA}$ and $H_2\,\lambda21\,218\,\mbox{\normalfont\AA}$ emission lines are clearly distinguished from their main narrow counterparts. Moreover, the panel displaying the resulting fit for the [P II] $\lambda11\,886\,\mbox{\normalfont\AA}$ emission line also shows a second, well-separated feature corresponding to the [N III] $\lambda11\,910\,\mbox{\normalfont\AA}$ emission. We did not attempt to model the [P II] emission at $\lambda11\,470\,\mbox{\normalfont\AA}$ since it is barely seen in the integrated spectra show in Fig. 1 as well as in the gas spectra shown in Fig. 2.

4 RESULTS

We show maps of flux distribution, radial velocities, and σ in the first, second, and third columns, respectively, of Fig. 4 for the Pa β , [PII] λ 11886 Å, and Br γ emission lines as well as for the [FeII] λ 12570 Å and H₂ λ 21218 Å narrow components. In all cases, velocities are relative to the corrected systemic velocity of the galaxy $v_{\rm sys,cor} = 5841\,{\rm km\,s^{-1}}$ (see Section 5.1), and maps of velocity dispersions were corrected for the instrumental broadening. In all panels, north is up, east is to the left, the black cross indicates the nucleus of the galaxy, and pixels in white were not considered in the fit. In Fig. 5, we show the flux distribution of

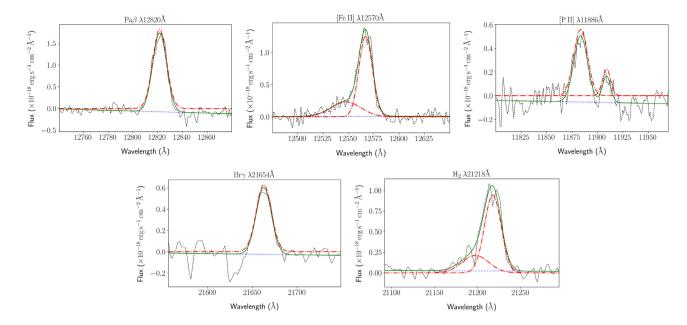


Figure 3. Examples of fitting results using the package IFSCUBE for the spaxel corresponding to the nucleus of the galaxy. In each panel, the black dashed and green solid lines are the observed and modelled spectra, respectively. The fitted Gaussian components are shown by the red dash-dotted lines, and the blue dotted line is the fitted residual continuum. The main fitted feature is indicated at the top of each panel. Note that the second component accompanying the $[P II] \lambda 11886 \text{ Å}$ emission line corresponds to the $[N III] \lambda 11910 \text{ Å}$ emission feature.

the [Fe II] $\lambda 12\,570\,\text{Å}$ and $H_2\,\lambda 21\,218\,\text{Å}$ broad, blueshifted emission. The flux map of the [Ni II] $\lambda 11\,910\,\text{Å}$ emission line is shown in Fig. 6.

Hereafter, we refer to the [Fe II] $\lambda12\,570 \mbox{\normalfone}{\norma$

4.1 Flux distribution of the emission lines

The Pa β , [Fe II], [P II], and Br γ emission features, as well as the H₂ λ 21 218 Å molecular line show spatially resolved emission. The most compact flux distribution is found for [P II] that extends to roughly 0.4 arcsec from the nucleus of the galaxy, while the H₂ λ 21 218 Å emission line extends over almost the entire NIFS FoV. In the *J* band, both the [Fe II] and [P II] flux distributions peak close to the galaxy nucleus and display a slight elongation in the north–west–south-east direction. In the *K* band, H₂ emission peaks at the nucleus and smoothly decreases towards the edges of the FoV.

The Pa β and Br γ flux maps show that the flux distribution is asymmetric and that the peak is off-centred – Pa β and Br γ fluxes peak north-west, ≈ 0.1 arcsec of the nucleus, and decrease towards the south-east. This may indicate either the presence of a gradient in the amount of visual extinction at the galaxy or that these lines are tracing a circumnuclear structure. We note that extended Pa β emission is also seen along the right edge of the FoV.

The flux distributions of the broad, blueshifted components associated with the [Fe II] and $H_2\,\lambda21\,218\mbox{\normalfont A}$ emission lines are shown in the left- and right-hand panels of Fig. 5, respectively. We masked the spaxels where their peaks are smaller than three times the standard deviation of the adjacent continuum. These maps show that the main contribution of these broad, blueshifted components to the observed emission is restricted to a compact region of radius $\approx\!0.25\,\mathrm{arcsec}$ close to the nucleus of the galaxy, and with slightly higher values towards the south-east.

Finally, the [Ni II] emission (Fig. 6) is overall much weaker than that of [P II], although same magnitude – of the order of $10^{-18} {\rm erg \, s^{-1} \, cm^{-2} \, spaxel^{-1}}$ – flux levels can be seen along a patchy distribution close to the galaxy nucleus. As a matter of fact, on average, the [Ni II]/[P II] ratio is ≈ 0.24 , reaching up to 0.3 in some locations. We have not found any report in the literature on the detection of [Ni II] λ 11 910Å emission in NGC 34. However, [Ni II] emission at λ 11 910Å has been found in other AGN in a spectroscopic survey carried out by Lamperti et al. (2017).

4.2 Gas kinematics

The velocity maps for all the fitted emission lines consistently show a rotation signature with a kinematic major axis at the approximate PA of -15° , with redshifted velocities north of the galaxy nucleus, and blueshifted velocities to the south. For the ionized hydrogen, velocities range from -160 to $240~km~s^{-1}$. The ionized gas ([Fe II] and [P II]) shows velocities between -160 and $170~km~s^{-1}$, while the molecular hydrogen has velocities between -190 and $140~km~s^{-1}$. We could not constrain the kinematics of the [Ni II] emission feature due to its patchy distribution.

Velocity dispersions peak close to the nucleus of the galaxy in all cases. The highest values, $160-185~\rm km\,s^{-1}$, are found for the ionized gas. Both the ionized and molecular hydrogen show maximum velocity dispersions in the range $125-145~\rm km\,s^{-1}$.

We recall that we included broad, blueshifted components to the fitting of the [Fe II] and $H_2 \lambda 21 \, 218 \, \mathring{A}$ emission lines, since fitting them with a single Gaussian could not account for the presence of broad, blue wings in a compact region close to the nucleus of the galaxy. These components have fixed radial velocities at -420 and $-250 \, \mathrm{km \, s^{-1}}$, and fixed velocity dispersions of 350 and 215 km s⁻¹ for the [Fe II] and $H_2 \lambda 21 \, 218 \, \mathring{A}$ emission lines, respectively.

We can further investigate the gas kinematics using channel maps, which is the mapping of the emission lines in constant velocity, or equivalently, wavelength bins. We built channel maps of the Pa β ,

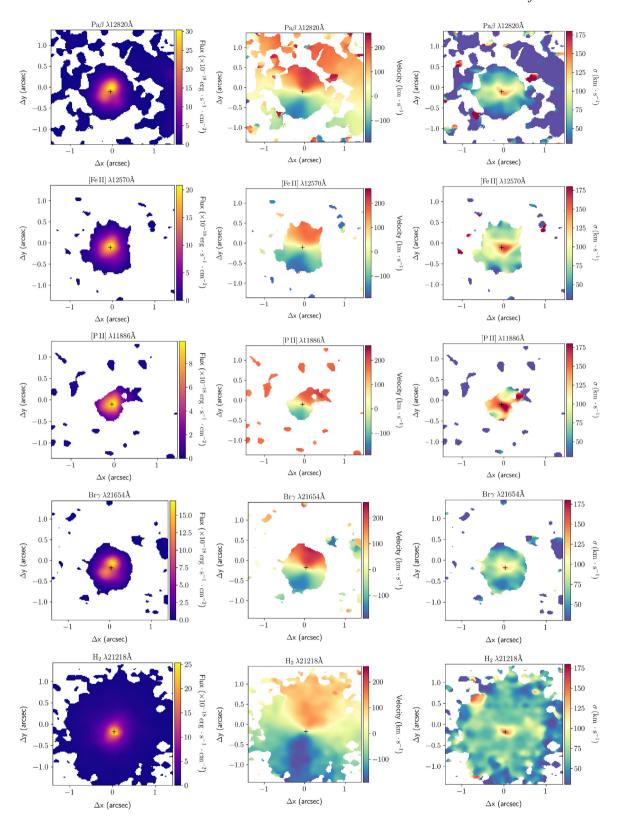


Figure 4. Maps of flux, radial velocities, and σ of the Pa β (first row), narrow [Fe II] $\lambda 12\,570\,\text{Å}$ (second row), [P II] $\lambda 11\,886\,\text{Å}$ (third row), Br γ (fourth row), and narrow H₂ $\lambda 21\,218\,\text{Å}$ (fifth row) emission lines. The range of values are indicated by the accompanying colour bars. The black cross in each panel indicates the nucleus of the galaxy. Pixels in white were not fitted. North is up and east is to the left in all maps.

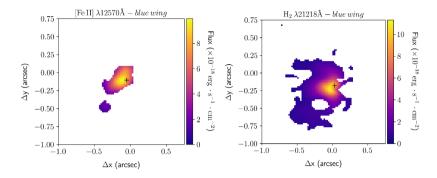


Figure 5. Flux maps of the [Fe II] λ 12 570 Å (left-hand panel) and H₂ λ 21 218 Å (right-hand panel) broad, blueshifted components.

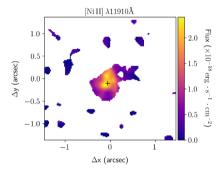


Figure 6. Flux distribution of the [Ni II] $\lambda 11 910 \text{ Å}$ emission line.

[Fe II], and $H_2 \lambda 21218 \, \mathring{A}$ emission lines using the CHANNEL_MAPS MODULE of the IFSCUBE package. This module evaluates the adjacent continuum level and masks spaxels with flux values that after the continuum subtraction fall below a lower threshold defined by the user. Our channel maps are shown in Figs 7–9. Maps were built between -440 and $290 \, \mathrm{km \, s^{-1}}$, with integrated fluxes in velocity bins of $61 \, \mathrm{km \, s^{-1}}$ centred at the velocity shown at the top of each panel. We masked spaxels with fluxes lower than $1.0 \, \mathrm{and} \, 0.1 \times 10^{-19} \mathrm{erg \, s^{-1} \, cm^{-2} \, \mathring{A}^{-1}}$ for the $J \, \mathrm{and} \, K \, \mathrm{bands}$, respectively.

All the channel maps confirm the transition from southern, blueshifted to northern, redshifted velocities. The rotation signature can be seen starting at the $-227\,\mathrm{km\,s^{-1}}$ velocity channel. The channel maps of Pa β (Fig. 7) confirm the asymmetric flux distribution. We can see that lower flux values can be found south of the nucleus, and that as one moves towards redshifted velocities, higher fluxes do not coincide with the nucleus of the galaxy, but instead, the flux distribution seems to encircle it and reaches maximum values to the north. The maps between the 17 and 138 $\mathrm{km\,s^{-1}}$ velocity channels also show the extended emission seen at the right edge of the FoV in the Pa β flux map shown in Fig. 4 (top left-hand panel).

The [Fe II] and $H_2 \lambda 21218$ Å channel maps (Figs 8 and 9) show that for these emission lines, fluxes peak very close to the nucleus of the galaxy. More importantly, these maps confirm the presence of the broad, blueshifted components starting at the $-409 \ \text{km s}^{-1}$ velocity channel. For the [Fe II], this blue wing is very compact and located in the nucleus of the galaxy, while for the $H_2 \lambda 21218$ Å, it is slightly elongated to the south.

4.3 Reddening

To estimate the visual extinction A_v throughout the inner $1.2 \, \text{kpc} \times 1.2 \, \text{kpc}$ of NGC 34, we followed Dametto et al. (2014),

who adopted the ratio of total to selective extinction as $R_V = 4.05 \pm 0.80$ from Calzetti et al. (2000), which has been shown to be the most suitable one when dealing with extinction in starbursts (Calzetti et al. 2000; Fischera, Dopita & Sutherland 2003). In the case B recombination for hydrogen (low-density limit, $T = 10^4$ K; Osterbrock & Ferland 2006), the intrinsic value of the Pa β /Br γ emission-line ratio is 5.88, and the visual extinction A_v is

$$A_{\rm v} = -15.24 \log \left(\frac{1}{5.88} \frac{\text{Pa } \beta}{\text{Br } \gamma} \right). \tag{1}$$

The $A_{\rm v}$ distribution throughout NGC 34 is shown in Fig. 10. The derived values range from \sim 1.9 to 9.0 mag. Higher extinction can be found towards the south of the galaxy. The mean $A_{\rm v}$ in the northern region is 6.78 \pm 0.86 and the median is 7.01, while in the south, the mean $A_{\rm v}$ is 7.74 \pm 0.95 and the median is 7.94. These results are in full agreement with long-slit studies (see table 4 of Dametto et al. 2014).

Alonso-Herrero et al. (2006) used *HST* NICMOS (Near Infrared Camera and Multi-Object Spectrometer) observations to study the NIR and star-forming properties of a sample of local LIRGs. They derived extinctions in the range 2–6 mag to the stars (photometric extinctions), while extinctions to the gas, derived from the H α /Pa α and Pa α /Br γ ratios, tend to be higher, ranging from 0.5 up to 15 mag. The values found for NGC 34 indicate that its inner regions are embedded in a heavily obscured environment, which could potentially affect optical studies of high spatial resolution aiming at investigating the nature of its nuclear activity.

4.4 Emission-line ratios

We show maps of the [Fe II]/Pa β (top panel) and H₂/Br γ (bottom panel) emission-line ratios in Fig. 11 along with the locations of regions identified from A to F that correspond to apertures of 0.2 × 0.2 arcsec². We integrated the gas spectra in each aperture and fitted the emission lines using Gaussian functions and the package IFSCUBE, thus obtaining their amplitudes, σ , and respective uncertainties, which were then used to estimate the uncertainties of the integrated fluxes through error propagation. Fluxes of the emission lines in each region are presented in Table 1 and the integrated spectra are shown in Appendix A.

The [Fe II]/Pa β and H₂/Br γ line-ratios are reddening insensitive and are used to build BPT-like diagrams in the NIR (Larkin et al. 1998; Rodríguez-Ardila et al. 2004; Rodríguez-Ardila, Riffel & Pastoriza 2005; Riffel et al. 2013; Colina et al. 2015; Riffel et al. 2021). The hydrogen recombination lines are tracers of UV-ionizing radiation from young OB-stars or from the AGN. The [Fe II] traces

 $Pa\beta$

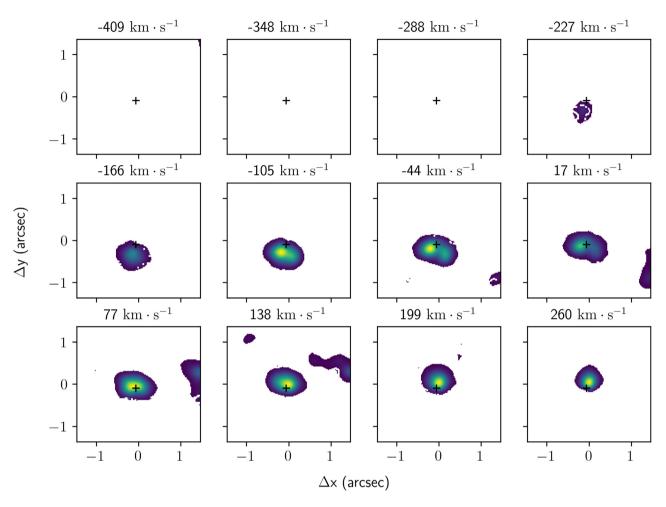


Figure 7. Pa β channel maps. The velocity is shown in top of each panel and the black cross marks the nucleus of the galaxy.

shocked, partially ionized regions and photodissociation regions, while the $H_2\,\lambda 21\,218\mbox{\normalfont\AA}$ traces the hot molecular gas $(T\geq 10^3\mbox{\normalfont\AA})$. Star-forming galaxies (SFGs) have [Fe II]/Pa $\beta<0.6$ and $H_2/\mbox{\normalfont\AA}$ <0.4, AGN-dominated systems have 0.6< [Fe II]/Pa $\beta<2.0$ and 0.4< $H_2/\mbox{\normalfont\AA}$ $\gamma<6.0$, and low ionization nuclear emission regions (LINERs) have [Fe II]/Pa $\beta>2.0$ and $H_2/\mbox{\normalfont﴾}$ >6.0 (e.g.: Riffel et al. 2013).

Our maps of [Fe II]/Pa β and H₂/Br γ for NGC 34 clearly show a ring-like structure with lower values of the ratios around the nucleus of the galaxy. Many locations in the [Fe II]/Pa β map are consistent with pure SF. Along the ring-shaped structure, [Fe II]/Pa β values range from \sim 0.4 to 0.75 and reach values of \sim 1.0 close to the nucleus of the galaxy in region F. In the case of the H₂/Br γ , the obtained values for the emission-line ratios are all consistent with the presence of an AGN, including locations along the ring-shaped structure where values range from \sim 0.7 to 1.4 and reach up to \sim 2.5 in the nucleus.

Overall, the fact that the structure in the shape of a ring around the nucleus of the galaxy can be seen in both maps of emission-line ratios, in addition to many locations displaying [Fe II]/Pa β values consistent with pure SF indicate that the NGC 34 nuclear starburst is distributed in a circumnuclear SF ring. However, larger values found for both line ratios, especially in the case of the H₂/Br γ ,

indicate that additional mechanisms are required to explain the [Fe $\rm II$] emission and the excitation of the $\rm H_2$ molecule, as will be discussed in Section 5.3.

4.5 PCA tomography

We also analysed the data cubes of NGC 34 using PCA tomography (Steiner et al. 2009). This technique uses PCA applied to data cubes in order to extract useful information from a given object (see e.g. Ricci, Steiner & Menezes 2011, 2014; Ricci, Steiner & Giansante 2015; Dahmer-Hahn et al. 2019). It searches for correlations between m spectral pixels across n spaxels. A set of m eigenvectors is built as a combination of the m spectral pixels and they are ordered by their contribution of the variance (obtained with the respective eigenvalues) of the data cube. The tomograms correspond to the projection of the eigenvectors on the data cubes. In practice, the eigenvectors, also called eigenspectra, show the correlations between the wavelengths, while the tomograms reveal where such correlations occur in the spatial dimension.

We applied PCA tomography in both J- and K-band data cubes after the subtraction of the stellar component of NGC 34. In the case of the K band data cube, we used the spectral range 20 814–22 568 Å, which contains all H_2 and the Br γ lines. For the J-band

[Fe II] $\lambda 12570$ Å

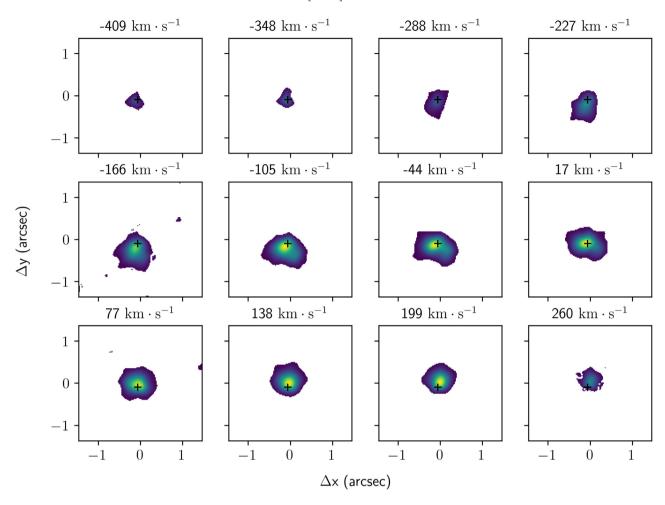


Figure 8. Same as in Fig. 7 for the [Fe II] $\lambda 12\,570\,\text{Å}$ emission line.

data cube, a spectral range of 12 356–12 923 Å was used, containing both [Fe II] and Pa β lines. The most relevant results revealed by the PCA tomography correspond to the fourth eigenspectrum of the K band data cube and the fifth eigenspectrum of the J band data cube, shown in Fig. 12. Eigenspectrum 4 of the K band data cube explains 1.76 per cent of the variance and is characterized by an anticorrelation between the Br γ and the H₂ λ 21 218 Å lines. Its tomogram reveals a nuclear object, which is related to the H₂ λ 21 218 Å line, and a circumnuclear structure, which corresponds to the Br γ line. A similar scheme is seen in the fifth tomogram of the J-band data cube. In this case, the eigenspectrum 5 explains 1.21 per cent of the variance and is characterized by a broad [Fe II] component anticorrelated with the Pa β line. The nuclear object seen in tomogram 5 is related to the broad [Fe II] component, while the circumnuclear structure is associated with the Pa β emission.

5 DISCUSSION

5.1 Nuclear disc of ionized and molecular gas

The maps of radial velocities of the fitted emission lines presented in Section 4.2 clearly show signature of rotation in the ionized and molecular gas phases. In order to compare the kinematics of the gas in NGC 34 with that of the stars, we carried out a new spectral synthesis in the K-band data cube using the code PPXF and the Gemini library of late spectral templates (Winge, Riffel & Storchi-Bergmann 2009), which was especially designed to support stellar kinematics studies in external galaxies. We also followed the recommendations of Cappellari (2017) to include only additive polynomials when using PPXF to fit the stellar kinematics. We then used the code PAFIT, which implements the method presented in appendix C of Krajnović et al. (2006), to determine: (1) a correction ($v_{\rm offset}$) to be applied to the adopted systemic velocity of the galaxy determined from the stellar velocity field ($v_{\rm sys,\,cor} = v_{\rm sys} + v_{\rm offset}$) and (2) the global kinematic PA of the stellar and gas discs.

The NGC 34 stellar velocity field is shown in the left-hand panel of Fig. 13 and the correction to the systemic velocity is $-40~\rm km~s^{-1}$, resulting in $v_{\rm sys,cor}=5841~\rm km~s^{-1}$. In general, one can see that the range of velocities as well as the PA of the rotation signature of the stars are very similar to the ones found for the ionized and molecular gas, indicating that they are under the influence of the same gravitational potential. This can be straightforwardly associated with the presence of a nuclear disc (ND) of stars and gas. For the stars, the kinematic PA of the disc is $-18\%0~\pm~1\%2$. The global kinematic

¹Available at https://pypi.org/project/pafit/.

$H_2 \lambda 21218 \text{\AA}$

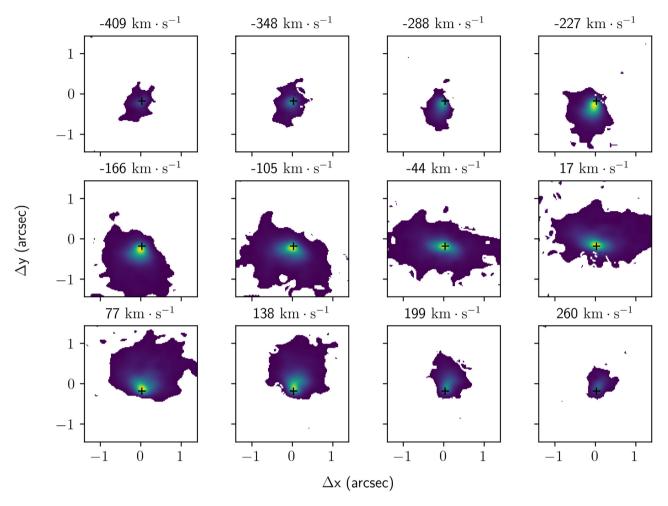


Figure 9. Same as in Fig. 7 for the $H_2 \lambda 21 218 \text{ Å}$ emission line.

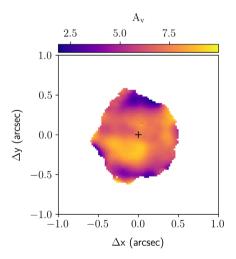


Figure 10. Map of visual extinction A_{v} .

PAs of the ionized and molecular gas discs are $-6^{\circ}0 \pm 0^{\circ}5$ and $-12^{\circ}5 \pm 1^{\circ}8$, respectively, resulting in a mean PA of $-9^{\circ}2 \pm 0^{\circ}9$. The PA of the stellar, ionized, and molecular gas discs are indicated in Fig. 13 by the solid green, blue dash-dotted, and dashed red lines, respectively. The field of stellar velocity dispersion (middle panel) shows higher values in the galaxy nucleus. In the right-hand panel of Fig. 13, we show stellar and gas velocity profiles taken along the kinematic PA of the stellar disc. Velocities rise steeply and reach maximum values in a radius smaller than 0.5 arcsec (≈ 200 pc) from the nucleus of the galaxy, indicating that the disc is compact.

Our results are in agreement with previous works presented by Fernández et al. (2014) and Xu et al. (2014), thus confirming the presence of an ND of ionized and molecular gas in NGC 34 with a northern receding and a southern approaching side. However, Xu et al. (2014) point out that the $\sim\!2$ kpc CO(1–0) disc found by Fernández et al. (2014) is associated with diffuse gas emission, while the much more compact $\sim\!200$ pc CO(6–5) disc detected by them (see fig. 1 in Xu et al. 2014) has the same distribution of the nuclear starburst traced by radio continuum emission. The kinematic PA of the CO(6–5) disc is -15° and the rotation velocity also rises up to a radius of 0.5 arcsec and then flattens, indicating that our NIR IFU

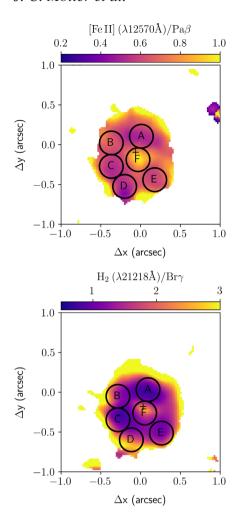


Figure 11. Maps of the [Fe II]/Pa β (top panel) and H₂/Br γ (bottom panel) emission-line ratios. The black, open circles mark regions from A to F where integrated measurements were taken. The black cross marks the galaxy nucleus.

data and the ALMA observations of Xu et al. (2014) share the same kinematics.

5.2 Resolved NIR diagnostic diagram and the NGC 34 power source

As previously mentioned, the nature of the NGC 34 nuclear emission-line spectrum is still under debate, with different works classifying it as due to a starburst (Mulchaey et al. 1996; Riffel et al. 2006) and/or AGN (Yuan et al. 2010; Brightman & Nandra 2011b). In this work, we characterize the nuclear activity in NGC 34 by means of high spatial resolution observations of NIR emission lines, namely the Pa β , [Fe II], Br γ , and H $_2$ $\lambda 21$ 218 Å emission features, which are more suited for probing the dust-embedded environments typically found in (U)LIRGs. While the Pa β and Br γ emissions are mainly due to the UV-ionizing flux produced by young OB-stars (ages $\sim\!10$ Myr) or by the AGN, emissions by the [Fe II] and H $_2$ species have different origin.

Excitation of the H₂ molecule can be due to (1) UV pumping (fluorescence), also regarded as non-thermal emission, where H₂ molecules are electronically excited by the absorption of UV photons in the Lyman–Werner band (918–1108Å) in photodissociation regions (PDRs), followed by rapid transitions to ro-vibrational excited

levels of the ground state (e.g. Black & van Dishoeck 1987); and the *thermal* processes either through (2) direct heating of the molecular gas by shocks (e.g. Hollenbach & McKee 1989) or by (3) X-rays, for example, produced by an AGN (e.g. Maloney, Hollenbach & Tielens 1996). Early works have already shown that for a single object, the observed H₂ spectrum is often the result of multiple excitation mechanisms that are at play (e.g Mouri 1994). Similarly, [Fe II] emission can also be traced to shock-dominated regions, associated either with radio jets, nuclear outflows and/or supernova remnants (SNRs), or to photoionization by a central X-ray source (AGN; Mouri, Kawara & Taniguchi 2000).

In this sense, NIR diagnostic diagrams based on the [Fe II]/Pa β and H₂/Br γ line-ratios provide a progression in the observed values from pure photoionized (such as in H II regions) to pure shock excited regions (e.g. shocks in SNRs), with SF regions displaying low values for both ratios, while regions dominated by the AGN have intermediate values (Larkin et al. 1998; Rodríguez-Ardila et al. 2004, 2005; Riffel et al. 2013; Colina et al. 2015; Riffel et al. 2021).

In Fig. 14, the left-hand panel shows our spatially resolved NIR diagnostic diagram of NGC 34 along with data collected from the literature for other sources (see Appendix B for the complete list of sources, values of the data points and references). The limiting ratio [Fe II]/Pa β < 0.6 from Riffel et al. (2013) for SFGs is indicated by the horizontal, black, dashed line. A zoomed version of the diagram is presented in the right-hand panel along with the values of line-ratios obtained in the integrated regions shown in Fig. 11 and listed in Table 1. We also included the linear relations between the [Fe II]/Pa β and H₂/Br γ line ratios defined by Colina et al. (2015), who investigated the two-dimensional excitation structure of the interstellar medium in a sample of low-z LIRGs and Sevferts. using NIR integral field spectroscopy.² In their work (see section 4.1 of Colina et al. 2015), distinct, linear relations in the [Fe II]/Pa β - $H_2/Br \gamma$ plane were defined for regions identified as (1) young SF regions (SF-young, ages <6 Myr); (2) aged, supernovae dominated regions (SNe-dominated, ages \sim 8–40 Myr); (3) compact regions associated with nuclear AGNs (AGN-compact); and (4) diffuse, extended emitting regions where the AGN radiation field is still detected (AGN-diffuse).

Following the diagnostics presented by Riffel et al. (2013), our results indicate that all the NGC 34 spaxels fall in the AGNdominated region. However, one should note that considering only the [Fe II]/Pa β line ratio, some spaxels fall bellow the 0.6 limiting value for SFGs. This is better illustrated by the integrated measurements taken in the regions A-F (right-hand panel of Fig. 14). Recall that a clear ring-shaped structure around the nucleus of the galaxy can be seen in the maps of the [Fe II]/Pa β and H₂/Br γ line ratios shown in Fig. 11. Regions A–E are all located on this circumnuclear ring. In Fig. 14, regions A, C and D are below the [Fe II]/Pa $\beta = 0.6$ curve, while regions B and E are slightly above it, although considering the measured uncertainties, they could also be below this line. Region F, which is surrounded by the ring structure, is well separated from the other points in the NIR diagnostic diagram and occupies a portion of the diagram where both the [Fe II] and $H_2 \lambda 21 218 \text{ Å}$ emission lines are enhanced. Our interpretation of these findings is that the low [Fe II]/Pa β values found in regions A–E are due to a combination

²Note that the linear relations presented by Colina et al. (2015) are defined on the basis of the [Fe II]1.64 μm/Br γ ratio. Assuming the theoretical ratio [Fe II]1.64 μm/[Fe II]1.26 μm = 0.7646 and case B recombination, we can calculate the [Fe II]1.26 μm/Pa β line ratio using [Fe II]1.64 μm/Br γ = 4.4974 × [Fe II]1.26 μm/Pa β (see Section 2.2 of Colina et al. 2015).

Table 1. Observed emission-line fluxes for the regions indicated in Fig. 11.

Region	[P II] (11 886 Å)	[Fe II] (12 570 Å)	Pa β (12 818 Å)	H ₂ (21 218 Å)	H ₂ (22 230 Å)	H ₂ (22 470 Å)	Br γ (21 654 Å)
A	0.56 ± 0.00	1.24 ± 0.07	2.15 ± 0.09	1.03 ± 0.11	0.26 ± 0.02	0.16 ± 0.01	1.09 ± 0.06
В	0.27 ± 0.02	0.56 ± 0.03	0.86 ± 0.03	0.55 ± 0.04	0.18 ± 0.01	0.08 ± 0.00	0.40 ± 0.02
C	0.24 ± 0.01	0.61 ± 0.03	1.07 ± 0.04	0.63 ± 0.03	0.18 ± 0.01	0.10 ± 0.00	0.63 ± 0.02
D	0.11 ± 0.01	0.27 ± 0.01	0.50 ± 0.01	0.39 ± 0.03	0.13 ± 0.01	0.05 ± 0.00	0.23 ± 0.01
E	0.14 ± 0.01	0.43 ± 0.02	0.70 ± 0.02	0.37 ± 0.02	0.11 ± 0.01	0.06 ± 0.00	0.31 ± 0.01
F	0.82 ± 0.02	1.81 ± 0.13	1.98 ± 0.08	2.25 ± 0.43	0.55 ± 0.09	0.22 ± 0.03	1.08 ± 0.05

Notes. Fluxes are in units of $\times 10^{-15}$ erg cm⁻² s⁻¹. Each region corresponds to an aperture of 0.2×0.2 arcsec².

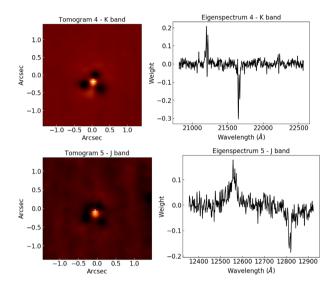


Figure 12. PCA tomography results. In both tomograms (left-hand panels), the white cross marks the position of the centre of the bulge of NGC 34. The tomograms identify two distinct structures in the central region of NGC 34: a nuclear component that is related to the [Fe II] and $H_2 \lambda 21218$ Å emission lines; and a circumnuclear structure associated with the Pa β and Br γ emission. The nuclear and circumnuclear structures appear anticorrelated in the eigenspectra (right-hand panels).

of excitation mainly caused by the stars with some contribution of the AGN, and that the nuclear starburst in NGC 34 is currently distributed in the circumnuclear star-forming ring (CNSFR) traced by these regions. The higher ratio found in region F reveals the presence of the AGN. Results presented by Rodríguez-Ardila et al. (2005) had already placed NGC 34 in the AGN-dominated region of the diagram (yellow circle in Fig. 14). However, since these results were based on one-dimensional spectra, they could not isolate the CNSFR. In addition, when considering the results found by Colina et al. (2015), we can see that region F, which we associate with the AGN, does follow the linear relation defined for AGN-dominated regions. The other regions in the galaxy follow the linear relation found for regions where the radiation field of the AGN acts as a contributor to the observed emission.

Our interpretation is corroborated by the results of the PCA tomography presented in Section 4.5, which clearly identify two distinct structures in the central region of NGC 34: a nuclear component which is related to the [Fe II] and $H_2 \, \lambda 21 \, 218 \, \mathring{A}$ emission lines; and a circumnuclear structure associated with the Pa β and Br γ emission.

5.3 The nature of the [Fe II] and H₂ λ21 218 Å emission lines

As can be seen by the NIR diagnostic diagram in Fig. 14, emission by [Fe II] and H_2 is detected in objects displaying several degrees of nuclear activity. In starburst galaxies, [Fe II] emission is enhanced by supernova-driven shocks. For this reason, as the starburst ages, the level of [Fe II] is expected to rise due to the increasing SN activity, while the levels of the hydrogen recombination lines decrease following the aging of the stellar population. In Seyfert objects, the X-ray emission produced by the central engine is able to penetrate deeper in the molecular clouds, thus creating extensive partially ionized zones where both [Fe II] and H_2 emissions are enhanced. Moreover, shocks associated with nuclear outflows also contribute as an additional excitation mechanism (Mouri et al. 2000).

Larkin et al. (1998) point out that while H_2 can be easily destroyed, 98 per cent of the iron is tied up in dust grains. Therefore, in order to invoke a single mechanism to power both lines, it must not destroy the H_2 molecules and must free up iron through dust grain destruction, both of which can be achieved in partially ionized zones created by the nuclear X-ray emission.

The centre of NGC 34 hosts a hard X-ray source associated with an obscured Seyfert 2 nucleus (Esquej et al. 2012). The fact that region F not only shows higher [Fe II]/Pa β and H₂/Br γ values but also clearly lies in a different position in the NIR diagnostic diagram in Fig. 14 can be explained by an enhancement of [Fe II] and H₂ due to the AGN X-ray emission. Moreover, although contribution from the ionizing flux related to the circumnuclear stellar activity likely plays a role in powering both the [Fe II] and H₂ emissions in the locations where the [Fe II]/Pa β values are low, an additional mechanism is necessary mainly to explain the enhanced H₂ emission in these same locations. Such an additional excitation component could also be due to the X-rays emitted by the central engine.

Following Riffel et al. (2008a) and Zuther et al. (2007), we can use the models of Maloney et al. (1996) to verify if X-ray heating by a source with intrinsic luminosity L_X is a viable H_2 excitation mechanism. Considering a cloud at a distance d from a hard X-ray source and with gas density n, the emergent $H_2 \lambda 21218 \text{Å}$ flux can be obtained from fig. 6 in Maloney et al. (1996) through the determination of the effective ionization parameter ξ_{eff} :

$$\xi_{\text{eff}} = 1.26 \times 10^{-4} \frac{f_x}{n_5 N_{22}^{0.9}},$$
(2)

where f_x is the incident X-ray flux at the distance d, n_5 is total hydrogen gas density in units of 10^5 cm⁻³, and N_{22} is the attenuating column density between the AGN and the gas cloud in units of 10^{22} cm⁻². The value of f_x can be determined using $L_x/4\pi d^2$. We calculated $\xi_{\rm eff}$ for three different distances from the AGN (40, 80 and 120 pc) adopting the intrinsic hard X-ray luminosity $L_X = 1.4 \times 10^{42} {\rm erg \ s^{-1}}$ (Esquej et al. 2012), $N_{\rm H} = 4.2 \times 10^{20} {\rm cm^{-2}}$ (Fernández

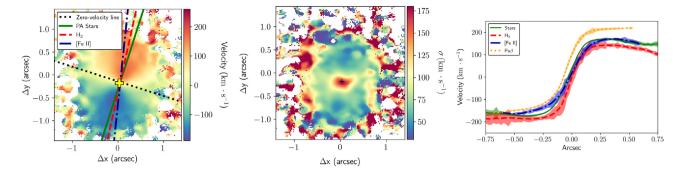


Figure 13. Left-hand panel: map of the stellar velocity field. The solid green, blue dash–dotted, and dashed red lines show the global kinematic PA of the stellar, ionized, and molecular gas discs, respectively. The black dotted line indicates the zero-velocity line and the yellow cross marks the galaxy nucleus. Middle panel: field of stellar velocity dispersion. In both fields, white pixels correspond to regions where uncertainties in the velocities are $> 50 \, \mathrm{km \, s^{-1}}$. Right-hand panel: velocities profiles of the gas and stars along the kinematic PA of the stellar disc. The shaded areas are 3σ uncertainties.

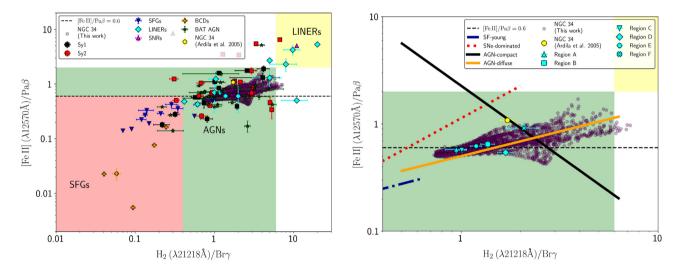


Figure 14. NIR diagnostic diagram based on the [Fe II]/Pa β and H₂/Br γ emission-line ratios. In both panels, the horizontal, dashed line corresponds to [Fe II]/Pa β = 0.6 from Riffel et al. (2013). The shaded areas divide the diagram into the regions dominated by SFGs, AGNs, and LINERs following Riffel et al. (2013). Left-hand panel: The light purple circles are the values obtained from this work for the spaxels where the four emission lines were fitted. Data from the literature for the Seyfert galaxies, types 1 and 2 (Sy1 and Sy2, respectively), are from Rodríguez-Ardila et al. (2004, 2005), Riffel et al. (2006), Riffel et al. (2013), Reunanen, Kotilainen & Prieto (2002), Knop et al. (2001), and Riffel et al. (2021). Data for the BAT AGN Spectroscopic Survey are from Lamperti et al. (2017). Data for SFGs are from Rodríguez-Ardila et al. (2005), Riffel et al. (2013), Dale et al. (2004), and Larkin et al. (1998). For LINERs, data are from Larkin et al. (1998) and Riffel et al. (2013). SNRs are from Larkin et al. (1998) and blue compact dwarfs (BCDs) data points are from Izotov & Thuan (2011). Right-hand panel: zoomed version of the left-hand panel showing our resolved NIR diagnostic diagram for NGC 34 and the results obtained from the integrated regions of Fig. 11 along with the linear relations defined by Colina et al. (2015) for young SF regions (SF-young), SNe-dominated regions, AGN-dominated, and AGN-diffuse regions.

et al. 2014) and for a gas density³ of $n=10^5$ cm⁻³. With the aid of fig. 6(a) in Maloney et al. (1996), we determined the emergent H₂ flux in erg cm⁻² s⁻¹ for an aperture of 0.021×0.021 arcsec² (the dimensions of a spaxel in our data cubes), which corresponds to a solid angle of 1.03×10^{-14} sr. Our results are shown in Table 2. The observed H₂ λ 21 218Å fluxes were determined from three different apertures centred in region F as follows: (1) at d=40 pc, $\log(F_{\rm H_2})$ corresponds to the mean flux value in a spaxel located in a circular aperture of radius 40 pc, (2) at d=80 pc, $\log(F_{\rm H_2})$ corresponds to the mean flux value in a spaxel located in a ring with inner and outer radii of 60 and 100 pc, respectively, and (3) at d=120 pc, $\log(F_{\rm H_2})$

³The models of Maloney et al. (1996) of X-ray irradiated molecular gas were calculated for $n = 10^3$ and 10^5 cm⁻³. However, the former does not provide predictions on the observed H₂ fluxes in every region of interest of the source. Therefore, for completeness, we chose $n = 10^5$ to perform our calculations.

Table 2. Observed and expected $H_2 \lambda 21218 \text{ Å}$ fluxes according to the X-ray heating models of Maloney et al. (1996) for an aperture of 0.021 \times 0.021 arcsec².

	d	Observed	n = 10	0^5cm^{-3}
	(pc)	$\log(F_{\rm H_2})$	$\log \xi_{\rm eff}$	$\log(F_{\rm H_2})$
(1)	40	-16.7	-1.8	-16.2
(2)	80	-17.0	-2.4	-18.7
(3)	120	-17.2	-2.8	-18.5

Note. Fluxes are in units of erg cm $^{-2}$ s $^{-1}$.

corresponds to the mean flux value in a spaxel located in a ring with inner and outer radii of 100 and 140 pc, respectively.

Our estimates indicate that X-ray heating can fully account for the H_2 emission seen in region F, which is probably associated

with the AGN. For larger distances, although X-rays contribute with some amount to the observed $H_2\,\lambda 21\,218\,\text{Å}$ fluxes, they are not the dominant excitation mechanism. Therefore, we conclude that, in NGC 34, the observed [Fe II] and H_2 emissions are due to a combination of photoionization by young stars, excitation by X-rays produced by the AGN and shocks. Mingozzi et al. (2018), when modelling the CO spectral line energy distribution, did not find that shocks contribute significantly to the heating of the molecular gas in NGC 34. However, our results provide kinematic signatures of shocks, as evidenced by the broad [Fe II] and H_2 components, indicating that shocks do contribute to the observed emission.

Note that our calculations using equation (2) were carried out assuming a column density derived from radio observations of the atomic hydrogen (HI at 21 cm; Fernández et al. 2014). The column densities obtained by Esquej et al. (2012) from X-rays and by Xu et al. (2014) using ALMA CO(6-5) observations are of the order of 10²³cm⁻². In this case, the predicted H₂ intensities at d = 40 pc would be two orders of magnitude lower than the ones shown in Table 2, indicating that even at the nucleus of NGC 34, the AGN would have a minor role in the observed H₂ emission. This is unlikely considering the discussion presented in Section 5.2. Rodríguez-Ardila et al. (2004) also estimated the H₂ fluxes for a list of sources using the models of Maloney et al. (1996) and found that the calculations using $N_{\rm H}$ derived from X-rays result in a poorer fit compared to the calculations using $N_{\rm H}$ derived from radio observations. They concluded that this is because $N_{\rm H}$ derived from X-rays probes obscuring material in the innermost regions of active galaxies, while H₂ and also [Fe II] emissions must arise farther out in the narrow-line region, thus justifying the adoption of a column density obtained from 21 cm observations.

We further investigate the nature of the H_2 emission using the ratio between the $H_2 \lambda 22470$ Å and $H_2 \lambda 21218$ Å emission lines to distinguish thermal ($\sim 0.1 - 0.2$) from non-thermal (~ 0.55) excitation mechanisms of the H_2 molecule (Mouri 1994). In NGC 34, $F_{2.247\mu m}/F_{2.122\mu m} < 0.2$ in all regions, thus confirming that the K-band H_2 lines are predominantly excited by thermal processes (shocks and/or X-ray heating). We can not, however, rule out the contribution of UV-fluorescence to the observed H_2 spectrum since the H_2 emission at 21218Å is itself much more sensitive to *thermal* processes than the $H_2 \lambda 22470$ Å line, which is a tracer of *fluorescence* (Black & van Dishoeck 1987).

The predominance of thermal processes in the excitation of the H₂ molecule is further illustrated by the fact that a broad, blueshifted component is only detected in the H₂ emission at 21 218 Å. The velocity and line width of this blue wing, as well as of that seen in [Fe II], indicate that the central regions of NGC 34 also drive a nuclear outflow of molecular and ionized gas. Although the outflow could, in principle, be driven either by the AGN, by winds associated with the circumnuclear starburst or by SNdriven shocks, we favour the first one since (1) the highest fluxes of these broad components coincide, or almost coincide, with the nucleus of the galaxy and (2) we do not observe the high [Fe II]/Pa β values expected according to the linear relation defined by Colina et al. (2015) for SNe-dominated regions, where SNdriven shocks are responsible for enhancing the [Fe II] emission. In the analysis carried out by Colina et al. (2015) for the composite LIRG NGC 5135 for example, they found a value larger than 1.5 for the [Fe II]/Pa β ratio at the location of [Fe II] peak associated with an SNe-dominated circumnuclear SF clump. Therefore, we conclude that in NGC 34, an AGN-driven nuclear outflow is the most likely primary source of shock excitation of the H_2 and [Fe $\scriptstyle\rm II]$ species.

Finally, we can also estimate the H_2 vibrational excitation temperature $T_{\rm vib}$ using $T_{\rm vib} \simeq 5600/\ln(1.355 \times F_{2.122\,\mu{\rm m}}/F_{2.247\,\mu{\rm m}})$ (Reunanen et al. 2002). We find a mean value of 2504 ± 174 K, which is in the range of values found for the H_2 thermal components in other sources (1800–2700 K; Reunanen et al. 2002). Rodríguez-Ardila et al. (2005) derived an upper limit of $T_{\rm vib} < 1800$ K for NGC 34. However, this result was based on one-dimensional spectra and includes contribution from the broad H_2 component. If we include the contribution of the broad, blueshifted component to the observed H_2 $\lambda 21$ 218 Å flux, we find $T_{\rm vib} = 2196 \pm 154$ K. The values found for $T_{\rm vib}$ support the assumption $T_{\rm vib} = 2000$ K used in equation (5) to estimate the mass of hot H_2 .

5.4 Circumnuclear SF ring

Our maps of emission-line ratios and the J- and K-band PCA tomography show that the nuclear starburst in NGC 34 is distributed in a CNSFR. We can use the Br γ luminosity ($L_{\rm Br}_{\gamma}$) to estimate the SFR in the ring adopting the following relation (Kennicutt 1998):

$$SFR(M_{\odot} \text{ yr}^{-1}) = 8.2 \times 10^{-40} L_{Br\gamma} (\text{erg s}^{-1}).$$
 (3)

The integrated, extinction-corrected Br γ flux in the continuous ring A–E–A, with inner and outer radii of 60 and 180 pc, respectively, is $F_{\rm Br}{}_{\gamma} \approx 20 \times 10^{-15} {\rm erg \ cm^{-2} \ s^{-1}}$ and $L_{\rm Br}{}_{\gamma} \approx 1.6 \times 10^{40} {\rm erg \ s^{-1}}$. From equation (3), we derive SFR $\approx 13 \ {\rm M}_{\odot} {\rm \ yr^{-1}}$. Also based on the Br γ luminosity, Valdés et al. (2005) estimated an SFR = 26.7 ${\rm \ M}_{\odot} {\rm \ yr^{-1}}$, and Dametto et al. (2014) found SFR = 9.63 ${\rm \ M}_{\odot} {\rm \ yr^{-1}}$.

Valdés et al. (2005) carried out NIR medium-resolution spectroscopy of a sample of (U)LIRGs and found that the Br γ luminosity provides SFRs which are on average 60 per cent of those derived from the far-infrared luminosity. In the case of NGC 34, they found $SFR(L_{FIR}) = 49.9 \text{ M}_{\odot} \text{ yr}^{-1}$. They stated that although the AGN might contribute to the infrared emission, this contribution would have to exceed 80 per cent to explain this large discrepancy. Alonso-Herrero et al. (2006) studied the NIR and SF properties of a sample of local LIRGs and also reported that the SFRs derived from the number of ionizing photons are on average 0.2–0.3 dex lower than those inferred from the total IR luminosity. This is because the luminosities of hydrogen recombination lines trace the most recent SFR, and there is a tendency for the measured IR luminosity to include some contribution from older stars. In addition, Alonso-Herrero et al. (2012) performed the spectral decomposition of a complete-volume-limited sample of local LIRGs and found that for the majority of them, the total AGN bolometric contribution to the observed IR luminosities has an upper limit of 5 per cent. Therefore, most of their IR luminosities are related to SF processes.

We point out that the SFR of NGC 34 derived by us is a lower limit since our observations provide a smaller aperture compared to previous works, and also due to the fact that the circumnuclear ring may extend to a larger radius. CNSFRs have been detected in a variety of LIRGs with diameters ranging from 0.7 to 2 kpc (e.g Alonso-Herrero et al. 2006).

5.5 Mass of ionized and molecular gas

We can estimate the mass of ionized and molecular gas in NGC 34 using integrated, extinction-corrected fluxes of Br γ and H₂ λ 21 218Å . Following Storchi-Bergmann et al. (2009) and assuming an electron

temperature of $T = 10^4$ K, the mass of ionized hydrogen, in solar masses, is given by

$$M_{\rm H\,II} \approx 3 \times 10^{19} \left(\frac{F_{\rm Br\gamma}}{{\rm erg\,cm}^{-2}\,{\rm s}^{-1}} \right) \left(\frac{D}{{\rm Mpc}} \right)^2 \left(\frac{N_{\rm e}}{{\rm cm}^{-3}} \right)^{-1},$$
 (4)

where $F_{\rm Br\,\gamma}$ is the Br γ integrated flux, D is the distance to NGC 34 and we have assumed an electron density $N_{\rm e}=160\,{\rm cm}^{-3}$, which is the mean value calculated by Kakkad et al. (2018) using the [S II] $\lambda\lambda$ 6716, 6731 lines for a sample of nearby (z<0.02) Seyfert 1s, Seyfert 2s, and AGN–Starburst composite systems.

The mass of hot H_2 in solar masses can be estimated as follows (Scoville et al. 1982):

$$M_{\rm H_2} = \frac{2 \, m_{\rm p} \, F_{\rm H_2 \lambda 21 \, 218} \, 4 \, \pi \, d^2}{f_{(\nu=1, J=3)} \, A_{S(1)} \, h \, \nu}$$
$$= 5.0776 \times 10^{13} \left(\frac{F_{\rm H_2 \lambda 21 \, 218}}{\rm erg \, cm^{-2} \, s^{-1}} \right) \left(\frac{D}{\rm Mpc} \right)^2, \tag{5}$$

where $m_{\rm p}$ is the proton mass, $F_{\rm H_2\lambda21218}$ is the integrated flux of the $\rm H_2~\lambda21~218~\mathring{A}$ emission line, D is the distance to the galaxy, h is the Planck constant, and ν is the frequency of the $\rm H_2$ line. Adopting a typical vibrational excitation temperature of $T_{\rm vib}=2000~\rm K$, which is consistent with the values we have obtained in Section 5.3, the population fraction is $f_{(\nu=1,J=3)}=1.22\times10^{-2}$ and the transition probability is $A_{S(1)}=3.47\times10^{-7}\rm s^{-1}$ (Turner, Kirby-Docken & Dalgarno 1977; Scoville et al. 1982).

We carried out the mass estimates for the inner $1.0 \times 1.0 \, \mathrm{arcsec^2}$ of the galaxy, since this corresponds to the region where we have derived visual extinction values to be applied to the observed fluxes. The integrated, extinction-corrected Br γ and H₂ $\lambda 21\, 218\, \mathring{\rm A}$ fluxes are $F_{\rm Br}_{\gamma} \approx 26 \times 10^{-15}$ and $F_{\rm H_2}_{\lambda 21}_{218} \approx 34 \times 10^{-15} \, \mathrm{erg} \, \mathrm{cm^{-2}} \, \mathrm{s^{-1}}$, respectively, resulting in $M_{\rm H\,II} \approx 3 \times 10^7$ and $M_{\rm H_2} \approx 11\, 570 \, \mathrm{M_{\odot}}$. Although $M_{\rm H\,II}$ is nearly 10^3 times larger than $M_{\rm H_2}$, we note that the later refers only to the hot molecular gas that gives rise to the observed NIR emission lines. The value derived by us for the hot molecular gas mass is similar to the one presented by Rodríguez-Ardila et al. (2005) who found $M_{\rm H_2} = 12\,600\,\mathrm{M_{\odot}}$.

Following Mazzalay et al. (2013), we can also estimate the mass of cold molecular gas using

$$\frac{M_{\rm H_2cold}}{\rm M_{\odot}} \approx 1174 \times \left(\frac{L_{\rm H_2\lambda21218}}{\rm L_{\odot}}\right),\tag{6}$$

where $L_{\rm H_2\lambda21218}$ is the luminosity of the $\rm H_2$ line. We find $M_{\rm H_2cold} \approx 8 \times 10^9 \rm M_{\odot}$ for NGC 34, which is roughly six orders of magnitude higher than that derived for the hot $\rm H_2$. This is in agreement with previous works that estimated the mass of cold molecular gas in NGC 34 from observations of the CO molecule ($M_{\rm H_2cold} \approx 7 \times 10^9 \rm M_{\odot}$; Kruegel, Chini & Steppe 1990; Chini et al. 1992; Kandalyan 2003).

6 CONCLUSIONS

The galaxy NGC 34 is a local LIRG whose nature of its emission-line features has been explained either as due to a pure starburst or due to a starburst–AGN composite source. In this work, we used AO-assisted IFU observations carried out with the NIFS instrument in the J and K bands to map the inner $1.2 \, \text{kpc} \times 1.2 \, \text{kpc}$ of the galaxy in order to investigate the excitation mechanisms of its NIR spectrum. We summarize our main findings as follows:

(i) The NGC 34 NIR spectra are characterized by the $[P II] \lambda 11 470 \text{Å}$, $[P II] \lambda 11 886 \text{Å}$, $[F e II] \lambda 12 570 \text{Å}$ and Pa β emission features in the J band, and by the $H_2 \lambda 21 218 \text{Å}$, Br γ , $H_2 \lambda 22 230 \text{Å}$, and $H_2 \lambda 22 470 \text{Å}$ emission lines in the K band.

- (ii) We report the detection of [Ni II] emission at λ11 910 Å.
- (iii) The fluxes of all emission features but Pa β and Br γ peak at the nucleus of the galaxy. The Pa β and Br γ flux distributions are asymmetric and peak north-west. This asymmetry is confirmed by the channel maps and indicates that the hydrogen recombination lines are tracing a circumnuclear structure.
- (iv) Through the analysis of the gas kinematics, we confirm the presence of an ND of ionized and molecular gas in NGC 34, with a northern-receding and a southern-approaching side. We determined that the mean kinematic PA of the gas disc is -9.2 ± 0.9 ;
- (v) We find that the NGC 34 nuclear starburst is distributed in a CNSFR with approximate inner and outer radii of 60 and 180 pc, respectively, as revealed by maps of the [Fe II]/Pa β and H₂/Br γ emission-line ratios. These maps clearly show a ring-like structure with lower values around the nucleus of the galaxy, especially for the [Fe II]/Pa β ratios, where many locations display values that are consistent with pure SF ([Fe II]/Pa $\beta \leq 0.6$, Riffel et al. 2013). The higher values found for both ratios indicate that multiple excitation mechanisms are at play in NGC 34.
- (vi) The presence of a CNSFR is corroborated by the PCA tomography, which shows a nuclear object associated with the H_2 $\lambda 21\,218\,\mathring{A}$ and [Fe II] emission lines, and a circumnuclear structure associated with the Pa β and Br γ lines.
- (vii) The resolved NIR diagnostic diagram based on the [Fe II]/Pa β and H₂/Br γ emission-line ratios show that all the NGC 34 spaxels fall in the AGN-dominated region. Integrated measurements taken in different regions A-F also confirm that the inner regions of NGC 34 are characterized by a nuclear and circumnuclear structures. Regions A-E trace the CNSFR since the [Fe II]/Pa β ratios lie bellow or slightly above the limiting value of 0.6 for pure SF defined by Riffel et al. (2013). Region F, in the nucleus of the galaxy, clearly occupies a portion of the diagram where both the [Fe II] and H₂ λ21 218Å lines are enhanced relative to the hydrogen recombination lines. Therefore, we associate region F with the AGN. NGC 34 hosts a hard X-ray source associated with an obscured Seyfert 2 nucleus (Esquej et al. 2012), and whose incident flux on the gas clouds could contribute to their enhancement. Moreover, the line ratios observed in region F follow the linear relation found by Colina et al. (2015) for AGN-dominated regions;
- (viii) Using the models of Maloney et al. (1996), we concluded that X-ray heating can fully account for the $H_2 \lambda 21218$ Å emission in the nucleus of the galaxy (region F), but it is not the predominant excitation mechanism at larger distances. Therefore we state that, in NGC 34, emission by [Fe II] and H_2 is due to a combination of photoionization by young stars, excitation by X-rays produced by the AGN and shocks as evidenced by the kinematic signatures of outflows in the ionized and molecular gas phases;
- (ix) We detected broad, blueshifted components associated with the [Fe II] and $H_2 \, \lambda 21 \, 218 \, \mathring{A}$ emission lines at -420 and -250 km \cdot s⁻¹, respectively, that can be interpreted as a nuclear outflow of ionized and molecular gas. Our results indicate that the outflow is AGN-driven and is the most likely primary source of shock excitation of the H_2 and [Fe II] species;
- (x) From the Br γ luminosity of the CNSFR we estimated a lower limit of SFR $\approx 13 \, M_{\odot} \, \text{yr}^{-1}$;
- (xi) The mass of ionized hydrogen in NGC 34 is $M_{\rm H{\scriptstyle II}} \approx 3 \times 10^7 {\rm M}_{\odot}$, and the mass of cold molecular gas is $M_{\rm H_2cold} \approx 8 \times 10^9 {\rm M}_{\odot}$, in agreement with previous works.

Our results indicate that the merger remnant NGC 34 is a gasrich system that hosts an AGN surrounded by a ring of SF. The AGN and the CNSFR are embedded in a compact, highly obscured environment making it difficult for optical studies to probe such inner, dusty regions. A common, evolutionary picture for (U)LIRGs is that mergers of gas-rich galaxies are responsible for funneling large amounts of molecular gas towards the centre of the merger, thus providing fuel to the starburst and AGN, both of which act on heating the surrounding dust. As the starburst declines, and the combined effects of SN explosions and AGN feedback clear out the nuclear dust, these objects become optically selected quasars (e.g. Sanders et al. 1988). In the case of NGC 34, we, therefore, conclude that we may be witnessing the early, evolutionary stage of a dust-enshrouded AGN.

ACKNOWLEDGEMENTS

JCM thanks Coordenação de Aperfeiçoamento de Pessoal de Nível Superior (CAPES) for the financial support under the grant 88882.316156/2019-01. RR and RAR thank Conselho Nacional de Desenvolvimento Científico e Tecnológico (CNPq), CAPES, and Fundação de Amparo à Pesquisa do Estado do Rio Grande do Sul (FAPERGS) for the financial support. TVR also thanks CNPq for the financial support under the grant 306790/2019-0. NZD acknowledges partial support from FONDECYT through project 3190769. We thank Dr João Evangelista Steiner†(1950-2020) for enlightening us with comments and suggestions before embarking on his journey to the stars. The authors also thank the anonymous referee for his/her careful revision that helped us to improve the quality of this paper. This paper is based on observations obtained at the international Gemini Observatory, which is managed by the Association of Universities for Research in Astronomy (AURA) under a cooperative agreement with the National Science Foundation, on behalf of the Gemini Observatory partnership: the National Science Foundation (United States), National Research Council (Canada), Agencia Nacional de Investigación y Desarrollo (Chile), Ministerio de Ciencia, Tecnología e Innovación (Argentina), Ministério da Ciência, Tecnologia, Inovações e Comunicações (Brazil), and Korea Astronomy and Space Science Institute (Republic of Korea). This research made use of observations made with the NASA/ESA Hubble Space Telescope, and obtained from the Hubble Legacy Archive, which is a collaboration between the Space Telescope Science Institute (STScI/NASA), the Space Telescope European Coordinating Facility (ST-ECF/ESAC/ESA), and the Canadian Astronomy Data Centre (CADC/NRC/CSA). This research made use of ASTROPY, ⁴ a community-developed core PYTHON package for astronomy (Astropy Collaboration et al. 2013; Price-Whelan et al. 2018); PHOTUTILS, an ASTROPY package for detection and photometry of astronomical sources (Bradley et al. 2021), the NUMPY (Harris et al. 2020), and MATPLOTLIB (Hunter 2007) PYTHON libraries.

DATA AVAILABILITY

The data underlying this paper are available in the Gemini Observatory Archive at https://archive.gemini.edu/searchform, and can be accessed with the Program ID GN-2011B-Q-71.

REFERENCES

Alonso-Herrero A., Rieke G. H., Rieke M. J., Colina L., Pérez-González P. G., Ryder S. D., 2006, ApJ, 650, 835

Alonso-Herrero A., Pereira-Santaella M., Rieke G. H., Rigopoulou D., 2012, ApJ, 744, 2

Astropy Collaboration et al., 2013, A&A, 558, A33

Baldwin J. A., Phillips M. M., Terlevich R., 1981, PASP, 93, 5

Black J. H., van Dishoeck E. F., 1987, ApJ, 322, 412

Bradley L. et al., 2021, astropy/photutils: 1.0.2, Zenodo, http://doi.org/10.5 281/zenodo.4453725

Brightman M., Nandra K., 2011a, MNRAS, 413, 1206

Brightman M., Nandra K., 2011b, MNRAS, 414, 3084

Calzetti D., Armus L., Bohlin R. C., Kinney A. L., Koornneef J., Storchi-Bergmann T., 2000, ApJ, 533, 682

Cappellari M., 2017, MNRAS, 466, 798

Chini R., Kruegel E., Steppe H., 1992, A&A, 255, 87

Colina L. et al., 2015, A&A, 578, A48

Cushing M. C., Rayner J. T., Vacca W. D., 2005, ApJ, 623, 1115

Dahmer-Hahn L. G. et al., 2019, MNRAS, 489, 5653

Dale D. A. et al., 2004, ApJ, 601, 813

Dametto N. Z., Riffel R., Pastoriza M. G., Rodríguez-Ardila A., Hernandez-Jimenez J. A., Carvalho E. A., 2014, MNRAS, 443, 1754

Davies R. I., Müller Sánchez F., Genzel R., Tacconi L. J., Hicks E. K. S., Friedrich S., Sternberg A., 2007, ApJ, 671, 1388

Deeley S. et al., 2017, MNRAS, 467, 3934

Esquej P., Alonso-Herrero A., Pérez-García A. M., Pereira-Santaella M., Rigopoulou D., Sánchez-Portal M., Castillo M. et al., 2012, MNRAS, 423, 185

Fernández X., van Gorkom J. H., Schweizer F., Barnes J. E., 2010, AJ, 140, 1965

Fernández X., Petric A. O., Schweizer F., van Gorkom J. H., 2014, AJ, 147, 74

Fischera J., Dopita M. A., Sutherland R. S., 2003, ApJ, 599, L21

Guainazzi M., Matt G., Perola G. C., 2005, A&A, 444, 119

Harrison C. M., 2017, Nat. Astron., 1, 0165

Harris C. R. et al., 2020, Nature, 585, 357

Hollenbach D., McKee C. F., 1989, ApJ, 342, 306

Hopkins P. F., Hernquist L., Cox T. J., Kereš D., 2008, ApJS, 175, 356

Hung C.-L. et al., 2014, ApJ, 791, 63

Hunter J. D., 2007, Comput. Sci. Eng., 9, 90

Izotov Y. I., Thuan T. X., 2011, ApJ, 734, 82

Kakkad D. et al., 2018, A&A, 618, A6

Kandalyan R. A., 2003, A&A, 398, 493

Kawakatu N., Wada K., 2008, ApJ, 681, 73

Kennicutt Robert C. J., 1998, ARA&A, 36, 189

Kewley L. J., Heisler C. A., Dopita M. A., Lumsden S., 2001, ApJS, 132, 37

Kewley L. J., Groves B., Kauffmann G., Heckman T., 2006, MNRAS, 372,

Knop R. A., Armus L., Matthews K., Murphy T. W., Soifer B. T., 2001, AJ, 122, 764

Krajnović D., Cappellari M., de Zeeuw P. T., Copin Y., 2006, MNRAS, 366, 787

Kruegel E., Chini R., Steppe H., 1990, A&A, 229, 17

Lamperti I. et al., 2017, MNRAS, 467, 540

Larkin J. E., Armus L., Knop R. A., Soifer B. T., Matthews K., 1998, ApJS, 114, 59

Lonsdale C. J., Farrah D., Smith H. E., 2006, inMason J. W., ed., Ultraluminous Infrared Galaxies, Astrophysics Update 2, Springer, Berlin Heidelberg

Malkan M. A., Gorjian V., Tam R., 1998, ApJS, 117, 25

Maloney P. R., Hollenbach D. J., Tielens A. G. G. M., 1996, ApJ, 466, 561

Mazzalay X. et al., 2013, MNRAS, 428, 2389

Menezes R. B., Steiner J. E., Ricci T. V., 2014, MNRAS, 438, 2597

Mingozzi M. et al., 2018, MNRAS, 474, 3640

Mouri H., 1994, ApJ, 427, 777

Mouri H., Kawara K., Taniguchi Y., 2000, ApJ, 528, 186

Mulchaey J. S., Wilson A. S., Tsvetanov Z., 1996, ApJS, 102, 309

Nardini E., Risaliti G., Watabe Y., Salvati M., Sani E., 2010, MNRAS, 405, 2505

⁴http://www.astropy.org.

- Osterbrock D. E., Ferland G. J., 2006, Astrophysics of Gaseous Nebulae and Active Galactic Nuclei University Science Books, Sausalito, CA
- Pérez-Torres M., Mattila S., Alonso-Herrero A., Aalto S., Efstathiou A., 2021, A&AR, 29, 2
- Planck Collaboration et al., 2020, A&A, 641, A6
- Price-Whelan A. M. et al., 2018, AJ, 156, 123
- Prouton O. R., Bressan A., Clemens M., Franceschini A., Granato G. L., Silva L., 2004, A&A, 421, 115
- Rayner J. T., Cushing M. C., Vacca W. D., 2009, ApJS, 185, 289
- Reunanen J., Kotilainen J. K., Prieto M. A., 2002, MNRAS, 331, 154
- Ricci T. V., Steiner J. E., Menezes R. B., 2011, ApJ, 734, L10
- Ricci T. V., Steiner J. E., Menezes R. B., 2014, MNRAS, 440, 2419
- Ricci T. V., Steiner J. E., Giansante L., 2015, A&A, 576, A58
- Riffel R. et al., 2015, MNRAS, 450, 3069
- Riffel R., Rodríguez-Ardila A., Pastoriza M. G., 2006, A&A, 457, 61
- Riffel R., Pastoriza M. G., Rodríguez-Ardila A., Maraston C., 2008b, MNRAS, 388, 803
- Riffel R., Pastoriza M. G., Rodríguez-Ardila A., Bonatto C., 2009, MNRAS, 400, 273
- Riffel R., Rodríguez-Ardila A., Aleman I., Brotherton M. S., Pastoriza M. G., Bonatto C., Dors O. L., 2013, MNRAS, 430, 2002
- Riffel R. A., Storchi-Bergmann T., Winge C., McGregor P. J., Beck T., Schmitt H., 2008a, MNRAS, 385, 1129
- Riffel R. A., Bianchin M., Riffel R., Storchi-Bergmann T., Schönell A. J., Dahmer-Hahn L. G., Dametto N. Z., Diniz M. R., 2021, MNRAS, 503, 5161
- Rodríguez-Ardila A., Pastoriza M. G., Viegas S., Sigut T. A. A., Pradhan A. K., 2004, A&A, 425, 457
- Rodríguez-Ardila A., Riffel R., Pastoriza M. G., 2005, MNRAS, 364, 1041 Rothberg B., Joseph R. D., 2006, AJ, 131, 185
- Ruschel-Dutra D., de Oliveira B. D., 2020, IFSCube, Zenodo, https://ifscube.readthedocs.io/en/latest/
- Sanders D. B., Mirabel I. F., 1996, ARA&A, 34, 749

- Sanders D. B., Soifer B. T., Elias J. H., Madore B. F., Matthews K., Neugebauer G., Scoville N. Z., 1988, ApJ, 325, 74
- Sanders D. B., Mazzarella J. M., Kim D. C., Surace J. A., Soifer B. T., 2003, AJ, 126, 1607
- Sarzi M., Shields J. C., Pogge R. W., Martini P., 2007, in Ho L. C., Wang J. W., eds, ASP Conf. Ser. Vol. 373, The Central Engine of Active Galactic Nuclei. Astron. Soc. Pac., San Francisco, p. 643
- Schweizer F., Seitzer P., 2007, AJ, 133, 2132
- Scoville N. Z., Hall D. N. B., Ridgway S. T., Kleinmann S. G., 1982, ApJ, 253, 136
- Steiner J. E., Menezes R. B., Ricci T. V., Oliveira A. S., 2009, MNRAS, 395, 64
- Storchi-Bergmann T., McGregor P. J., Riffel R. A., Simões Lopes R., Beck T., Dopita M., 2009, MNRAS, 394, 1148
- Turner J., Kirby-Docken K., Dalgarno A., 1977, ApJS, 35, 281
- Valdés J. R., Berta S., Bressan A., Franceschini A., Rigopoulou D., Rodighiero G., 2005, A&A, 434, 149
- Veilleux S. et al., 2009, ApJS, 182, 628
- White S. D. M., Rees M. J., 1978, MNRAS, 183, 341
- Winge C., Riffel R. A., Storchi-Bergmann T., 2009, ApJS, 185, 186
- Xu C. K. et al., 2014, ApJ, 787, 48
- Yuan T. T., Kewley L. J., Sanders D. B., 2010, ApJ, 709, 884
- Zuther J., Iserlohe C., Pott J. U., Bertram T., Fischer S., Voges W., Hasinger G., Eckart A., 2007, A&A, 466, 451

APPENDIX A: SPECTRA OF REGIONS A-F

We show in Fig. A1 the integrated spectra of regions A–F indicated in Fig. 11 corresponding to apertures of $0.2 \times 0.2 \, arcsec^2$.

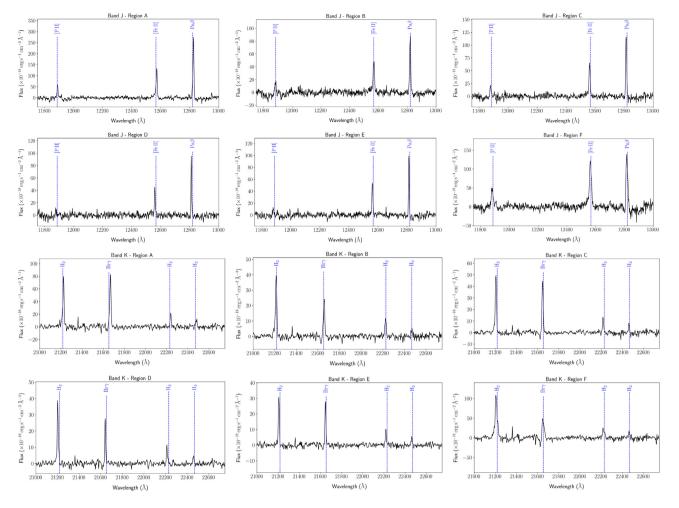


Figure A1. Integrated gas spectra of regions A–F corresponding to apertures of 0.2×0.2 arcsec². *J*-band spectra are shown in the first and second rows. *K*-band spectra are shown in the third and fourth rows. The NIR emission lines are indicated by the blue, dashed, vertical lines. *J*-band emission lines: [P II] λ 11 886Å, [Fe II] λ 12 570 Å, and Pa β . *K*-band emission lines: H₂ λ 21 218Å, Br γ , H₂ λ 22 230 Å, and H₂ λ 22 470 Å.

APPENDIX B: DATA FROM THE LITERATURE

In Table B1, we list the data points collected from the literature that were used to build the NIR diagnostic diagram shown in Fig. 14. The columns are as follows: (1) object type: Seyfert galaxies,

types 1 and 2 (Sy1 and Sy2, respectively), Burst Alert Telescope AGN Spectroscopic Survey (BAT AGN), SFGs, LINERs, SNRs, and BCDs; (2) object name; (3) $H_2/Br \gamma$ emission-line ratio; (4) [Fe II]/Pa β emission-line ratio; (5) reference.

Table B1. Data from the literature used to build the NIR diagnostic diagram shown in Fig. 14.

Object type (1)	Object name (2)	$H_2/Br \gamma$ (3)	[Fe II]/Paβ (4)	Reference (5)
Sy 1	Mrk 334	0.57 ± 0.06	0.57 ± 0.04	Riffel et al. (2006)
•	NGC 7469	1.93 ± 0.10	0.39 ± 0.05	Riffel et al. (2006)
	NGC 3227	2.58 ± 0.76	1.76 ± 0.14	Rodríguez-Ardila et al. (2004)
	NGC 4151	0.63 ± 0.05	0.57 ± 0.03	Rodríguez-Ardila et al. (2004)
	Mrk 766	0.32 ± 0.04	0.28 ± 0.03	Rodríguez-Ardila et al. (2004)
	NGC 4748	1.14 ± 0.26	0.84 ± 0.06	Rodríguez-Ardila et al. (2004)
	NGC 5548	0.74 ± 0.15	0.45 ± 0.09	Rodríguez-Ardila et al. (2004)
	PG 1612+261	1.01 ± 0.30	1.10 ± 0.10	Rodríguez-Ardila et al. (2004)
	NGC 1386	1.14 ± 0.10	1.59 ± 0.14	Reunanen et al. (2002)
	NGC 6814	4.09 ± 2.52	1.31 ± 0.26	Riffel et al. (2013)
	NGC 3227	1.44 ± 0.25	2.71 ± 0.35	Riffel et al. (2021)
	NGC 3516	4.19 ± 0.73	1.9 ± 0.24	Riffel et al. (2021)
	NGC 5506	0.22 ± 0.05	0.18 ± 0.02	Riffel et al. (2021)
Sy 2	ESO428-G014 NGC 591	0.92 ± 0.02 1.81 ± 0.08	0.67 ± 0.02 1.14 ± 0.06	Rodríguez-Ardila et al. (2005) Rodríguez-Ardila et al. (2005)
	Mrk 573	0.64 ± 0.09	0.61 ± 0.02	Rodríguez-Ardila et al. (2005)
	Mrk 1066	0.04 ± 0.05 0.94 ± 0.02	0.01 ± 0.02 0.71 ± 0.01	Rodríguez-Ardila et al. (2005)
	NGC 2110	3.36 ± 0.15	5.45 ± 0.14	Rodríguez-Ardila et al. (2005)
	NGC 7682	5.04 ± 0.13	0.46 ± 0.11	Rodríguez-Ardila et al. (2005)
	NGC 7674	0.68 ± 0.08	1.04 ± 0.10	Rodríguez-Ardila et al. (2005)
	NGC 5929	2.13 ± 0.29	1.09 ± 0.05	Rodríguez-Ardila et al. (2005)
	Mrk 1210	0.33 ± 0.03	0.50 ± 0.03	Rodríguez-Ardila et al. (2004)
	NGC 5728	2.97 ± 0.27	0.69 ± 0.13	Rodríguez-Ardila et al. (2004)
	NGC 4945	3.13 ± 0.06	0.86 ± 0.03	Reunanen et al. (2002)
	NGC 5128	2.04 ± 0.10	3.42 ± 0.16	Reunanen et al. (2002)
	NGC 4388	0.88 ± 0.08	0.40 ± 0.03	Knop et al. (2001)
	Mrk 3	0.31 ± 0.04	1.24 ± 0.04	Knop et al. (2001)
	Mrk 993	5.31 ± 0.00	0.34 ± 0.11	Rodríguez-Ardila et al. (2005)
	NGC 5953	1.28 ± 0.00	3.50 ± 0.00	Rodríguez-Ardila et al. (2005)
	NGC 1144	6.71 ± 0.00	6.50 ± 0.00	Rodríguez-Ardila et al. (2005)
	NGC 788	0.81 ± 0.14	0.23 ± 0.03	Riffel et al. (2021)
	Mrk 607	0.69 ± 0.13	0.26 ± 0.03	Riffel et al. (2021)
	NGC 5899	2.94 ± 0.51	1.75 ± 0.23	Riffel et al. (2021)
BAT AGN	33	1.20 ± 0.17	1.07 ± 0.02	Lamperti et al. (2017)
	308	3.95 ± 0.37	5.14 ± 0.16	Lamperti et al. (2017)
	382	1.68 ± 0.18	0.46 ± 0.03	Lamperti et al. (2017)
	404	1.25 ± 0.10	0.75 ± 0.01	Lamperti et al. (2017)
	517 533	1.24 ± 0.16 1.83 ± 0.17	0.72 ± 0.02 0.87 ± 0.04	Lamperti et al. (2017) Lamperti et al. (2017)
	585	2.62 ± 0.23	0.87 ± 0.04 0.17 ± 0.04	Lamperti et al. (2017)
	586	1.10 ± 0.15	0.42 ± 0.02	Lamperti et al. (2017)
	588	3.67 ± 1.23	0.59 ± 0.07	Lamperti et al. (2017)
	590	0.68 ± 0.05	0.69 ± 0.03	Lamperti et al. (2017)
	592	1.22 ± 0.82	1.54 ± 0.14	Lamperti et al. (2017)
	595	0.64 ± 0.04	0.93 ± 0.01	Lamperti et al. (2017)
	608	0.28 ± 0.02	0.27 ± 0.01	Lamperti et al. (2017)
	615	1.00 ± 0.02	0.38 ± 0.01	Lamperti et al. (2017)
	616	0.84 ± 0.06	0.42 ± 0.01	Lamperti et al. (2017)
	635	1.09 ± 0.13	0.68 ± 0.03	Lamperti et al. (2017)
	641	0.76 ± 0.15	0.73 ± 0.03	Lamperti et al. (2017)
	712	0.22 ± 0.03	0.38 ± 0.01	Lamperti et al. (2017)
	717	$0.75\ \pm\ 0.27$	0.70 ± 0.04	Lamperti et al. (2017)
	723	0.72 ± 0.03	0.53 ± 0.01	Lamperti et al. (2017)
	738	0.85 ± 0.08	0.59 ± 0.00	Lamperti et al. (2017)
	739	4.75 ± 2.78	0.81 ± 0.05	Lamperti et al. (2017)
	774	1.00 ± 0.60	0.55 ± 0.03	Lamperti et al. (2017)
	1117	2.50 ± 0.34	0.51 ± 0.03	Lamperti et al. (2017)
	1133	2.44 ± 0.11	0.62 ± 0.01	Lamperti et al. (2017)

Table B1 - continued

Object type	Object name	$H_2/Br \gamma$	[Fe II]/Pa β	Reference	
(1)	(2)	(3)	(4)	(5)	
	1157	0.30 ± 0.11	0.14 ± 0.01	Lamperti et al. (2017)	
	1158	3.29 ± 0.49	0.93 ± 0.03	Lamperti et al. (2017)	
	1161	0.69 ± 0.05	0.69 ± 0.02	Lamperti et al. (2017)	
	1182	1.09 ± 0.13	0.39 ± 0.01	Lamperti et al. (2017)	
	1198	5.41 ± 0.64	0.42 ± 0.02	Lamperti et al. (2017)	
	1287	3.19 ± 0.16	0.84 ± 0.01	Lamperti et al. (2017)	
SFGs	NGC 34	1.72 ± 0.14	1.08 ± 0.09	Rodríguez-Ardila et al. (2005)	
	NGC 7714	0.21 ± 0.02	0.29 ± 0.02	Rodríguez-Ardila et al. (2005)	
	NGC 1614	0.13 ± 0.01	0.27 ± 0.02	Rodríguez-Ardila et al. (2005)	
	NGC 3310	0.14 ± 0.03	0.33 ± 0.03	Rodríguez-Ardila et al. (2005)	
	NGC 1797	0.60 ± 0.05	0.42 ± 0.02	Riffel et al. (2013)	
	NGC 7678	0.35 ± 0.06	0.36 ± 0.06	Riffel et al. (2013)	
	NGC 6835	0.29 ± 0.05	0.42 ± 0.05	Riffel et al. (2013)	
	NGC 1222	0.12 ± 0.02	0.22 ± 0.01	Riffel et al. (2013)	
	NGC 2388	0.56 ± 0	0.26 ± 0	Dale et al. (2004)	
	NGC 6946	0.95 ± 0	0.49 ± 0	Dale et al. (2004)	
	M 82	0.15 ± 0	0.19 ± 0	Larkin et al. (1998)	
	II Zw 040	0.09 ± 0	0.15 ± 0	Larkin et al. (1998)	
	NGC 5253	0.07 ± 0	0.14 ± 0	Larkin et al. (1998)	
LINERs	NGC 5194	8 ± 3	2.3 ± 0.6	Larkin et al. (1998)	
	NGC 7743	5 ± 2	1.3 ± 0.4	Larkin et al. (1998)	
	NGC 660	0.42 ± 0.05	0.48 ± 0.02	Riffel et al. (2013)	
	NGC 1204	0.61 ± 0.05	0.43 ± 0.02	Riffel et al. (2013)	
	NGC 1266	9.79 ± 2.28	4.21 ± 0.82	Riffel et al. (2013)	
	NGC 7465	1.04 ± 0.20	1.23 ± 0.34	Riffel et al. (2013)	
	NGC 7591	1.18 ± 0.09	0.70 ± 0.03	Riffel et al. (2013)	
	UGC 12150	1.39 ± 0.08	0.60 ± 0.05	Riffel et al. (2013)	
	NGC 404	5 ± 0	2.7 ± 0.3	Larkin et al. (1998)	
	NGC 3998	2 ± 0	0.6 ± 0.3	Larkin et al. (1998)	
	NGC 4826	1 ± 0	0.7 ± 0.4	Larkin et al. (1998)	
	NGC 7479	11 ± 4	0.5 ± 0	Larkin et al. (1998)	
	NGC 4736	20 ± 0	5.3 ± 0	Larkin et al. (1998)	
SNRs	M 2	0.3 ± 0	8.5 ± 0	Larkin et al. (1998)	
	IC 0443	11 ± 0	5 ± 0	Larkin et al. (1998)	
BCDs	I Zw 040	0.0405 ± 0.0031	0.0226 ± 0.0018	Izotov & Thuan (2011)	
	Mrk 71	0.0940 ± 0.0076	0.0055 ± 0.0005	Izotov & Thuan (2011)	
	Mrk 930	0.2462 ± 0.0183	0.1633 ± 0.0076	Izotov & Thuan (2011)	
	Mrk 996	0.1746 ± 0.0139	0.0760 ± 0.0065	Izotov & Thuan (2011)	
	SbS	0.0583 ± 0.0068	0.0229 ± 0.0053	Izotov & Thuan (2011)	

Notes. (1) Object type: Seyfert galaxies, types 1 and 2 (Sy1 and Sy2, respectively), Burst Alert Telescope AGN Spectroscopic Survey (BAT AGN), SFGs, LINERs, SNRs, and BCDs; (2) object name; (3) $H_2/Br \gamma$ emission-line ratio; (4)[Fe II]/Pa β emission-line ratio; (5) reference.

This paper has been typeset from a $\ensuremath{\text{TEX/IATeX}}$ file prepared by the author.