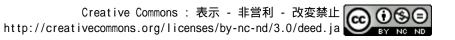


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1 Numerical study of temperature oscillation in loop heat pipe

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7 HIGHLIGHTS

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- 9 Once a two-phase flow is generated in the liquid line, temperature oscillation occurs.
- 10 A low reservoir temperature leads to a two-phase flow in the liquid line.
- When the condensation length oscillates highly, temperature amplitude becomes high.
- 12 Low sink temperature can prevent temperature oscillation.

14 ABSTRACT

15 Loop heat pipes are high-efficient heat transfer devices. Many spacecraft have loop heat pipes to control the 16 temperature of equipment precisely. Temperature oscillation, however, sometimes occurs in the loop heat pipe when 17 the external conditions change. Temperature oscillation may impede the temperature-controlling ability of loop heat 18 pipes. Nevertheless, the cause of temperature oscillation has not been understood yet. To understand the internal 19 flow during temperature oscillation, we developed a transient model that can reproduce the fluid condition in the 20 transport lines. In this study, the reservoir temperature, the liquid line temperature, and the condensation length were 21 focused. Investigating the relation between each parameter, we found that the cause of temperature oscillation is 22 the inflow of two-phase flow into the liquid line, which is caused by decreasing the reservoir temperature. In 23 addition, as the amplitude of the condensation length becomes high, temperature amplitude also becomes high. 24When the sink temperature is low, the condensation length can barely oscillate, and thus, temperature oscilla-25 tion is less likely to occur.

26 Keywords: Loop heat pipe, Temperature oscillation, Transient model, Condensation length

28			
29	A	:	Cross-sectional area (m ²)
30	С	:	Heat capacity (J/K) or parameter in Lockhart-Martinelli model
31	C_p	:	Specific heat (J/(kg·K))
32	d	:	Diameter (m)
33	dt	:	Time step (s)
34	G	:	Thermal conductance (W/K)
35	Gr	:	Grashof number
36	h	:	Enthalpy (J/(kg·K))
37	h_{lv}	:	Latent heat (J/kg)
38	h_{nc}	:	Natural convection heat-transfer coefficient $(W/(m^2 \cdot K))$
39	h_{tc}	:	Heat-transfer coefficient $(W/(m^2 \cdot K))$
40	k	:	Thermal conductivity $(W/(m \cdot K))$
41	L	:	Length (m)
42	m	:	Mass (kg)
43	'n	:	Mass flow rate (kg/s)

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4.4	17		
44	Nu		Nusselt number
45	P		Pressure (Pa)
46			Prandtl number
47	Q	:	Amount of heat transferred between the fluid in the pipe and the ambient or the sink (W)
48	Q _{in}	:	Heat load (W)
49	Re	:	Reynolds number
50	Su	:	Suratman number
51	Т	:	Absolute temperature (K)
52	и	:	flow velocity (m/s)
53	UA/L	: T	hermal conductance per unit length (W/(m·K))
54	V	:	Volume (m ³)
55	x	:	Vapor quality
56	X	:	Lockhart-Martinelli parameter
57	Z.	:	Coordinate in axial direction
58	ε	:	Porosity
59	μ	:	Viscosity (Pa·s)
60	ρ	:	Density (kg/s)
61	σ		Surface tension (N/m)
62	ϕ		Two-phase multiplier
63	T		
64	Subsc	ript	
65			Air
66	amb	:	Ambient
67	bulk	:	Wick material
68	сс	:	Reservoir
69	cond	:	Condenser
70	eff	:	Effective
71	evap	:	Evaporator or evaporation
72	f		Fluid
73	hb	:	Heater block
74	i	:	Node number
75	in	:	Inner
76	int	:	Vapor-liquid interface
77	l		Liquid
78	ll,out	:	Liquid line outlet
79	out		Outer
80	sat		Saturation
81	sink		Sink
82	tp		Two-phase flow
83	v		Vapor
84	vg		Vapor groove
85	wall	:	Wall
86			

87 **1. Introduction**

88 Loop heat pipes (LHPs) are highly-efficient and reliable heat transfer devices [1–3]. Many spacecrafts 89 mount LHPs for thermal control. Fig. 1 shows the schematic of an LHP. The LHP has working fluid and 90 consists of an evaporator, vapor line, condenser, liquid line, and reservoir. The heat applied to the evaporator 91 is absorbed by evaporation, and the generated vapor condenses in the condenser. The LHP can transport a large 92 amount of heat. The evaporator and condenser are connected with simple pipes. The capillary force of a pri-93 mary wick in the evaporator circulates the working fluid. The LHP can precisely control the temperature of 94 equipment with a small amount of electric power. As the LHP has many advantages, there have been several 95 studies on the ground application of the LHP [4–7].

96 One of the problems in the LHP, however, is tem-97 perature oscillation. When the heat load applied to the 98 evaporator or the temperature of the sink attached to the 99 condenser changes, the temperature of the LHP some-100 times oscillates. This may impede the temperature-con-101 trol capabilities of an LHP and may cause equipment 102 failure when the allowable temperature range of the 103 equipment becomes narrow.

104 Various investigations on temperature oscillation 105 have been conducted. Ku et al. observed two types of temperature oscillation in the experiments: high-ampli-106 107 tude/low-frequency oscillation [8] and low-ampli-108 tude/high-frequency oscillation [9]. According to the 109 previous research [8], when the thermal mass attached 110 to the evaporator is high, the heat load applied to the 111 evaporator is kept to a minimum, and the sink tempera-

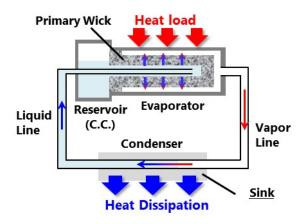


Fig. 1 Schematic of a loop heat pipe

ture is lower than the ambient temperature, high-amplitude/low-frequency oscillation occurs. They also found that the amplitude of temperature becomes high as the condensation length (the length required for vapor to completely condense) oscillates highly [8].

115 On the other hand, the cause of the low-amplitude/high-frequency oscillation has not been understood 116 yet. The proposed experiment mainly focuses on the low-amplitude/high-frequency oscillation. Ku et al. [9] 117 predicted that the temperature oscillation would occur if the condensation length is either longer or shorter 118 than the condenser length. Chen et al. [10] also led to the same conclusion as Ku [9] and elucidated the influ-119 ence of the LHP orientation on the temperature oscillation. Vershinin et al. [11] investigated the favorable 120 conditions for the temperature oscillation. They concluded that the shortage of the working fluid in the LHP 121 and the shortage of the liquid in the reservoir caused the temperature oscillation. However, the cause of the 122 temperature oscillation has been only deduced based on the experimental results.

Numerical simulation is useful to understand the cause of the temperature oscillation. Some researchers developed transient models of LHPs [12–19]. There are only a few transient models related to temperature oscillation. Launay et al. [15] developed a transient model for the temperature oscillation that yielded good accordance with the experimental results. Hoang et al. [16] proposed the stability theory for the LHP based on the analytical model. The stability criteria proposed in their research indicated that the higher the degree of the subcooling in the liquid line, the higher the possibility of temperature oscillation.

However, the relation between the internal flow and temperature oscillation is still unclear. The previous researches have not elucidated the reason why temperature oscillation occurs when the condensation length becomes longer than the condenser length or when the degree of the subcooling in the liquid line is high. In this study, we developed an analytical model to understand the internal flow during temperature oscillation. The model can calculate the vapor quality in the transport lines and investigate the fluid condition.

134 **2. Analytical model**

135 The LHP is modeled, especially focusing on the 136 internal flow of the vapor line, the condenser, and the 137 liquid line. The analytical model is divided into two 138 parts: the vapor line, the condenser and the liquid line; 139 the evaporator and the reservoir. The fluid is assumed 140 to be incompressible and viscous.

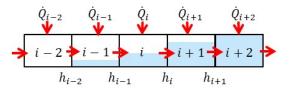


Fig. 2 Discretization of vapor line, condenser, and liquid line

142 **2.1 Vapor line, condenser, and liquid line**

143 The vapor line, condenser, and liquid line are discretized into small-volume nodes (Fig. 2). i expresses 144 the node number and starts from the vapor line inlet. The energy-conservation equation at node i is expressed 145 as

146

$$\left[\rho V \frac{dh}{dt}\right]_{i} = \dot{m}(h_{i-1} - h_{i}) + \dot{Q}_{i} + u_{i}A(P_{i-1} - P_{i+1}) + \frac{kA}{dz}(T_{i-1} - T_{i}) + \frac{kA}{dz}(T_{i+1} - T_{i}).$$
(1)

 \dot{Q}_i in the vapor line and the liquid line is calculated by

$$\dot{Q}_i = G_{f_amb}(T_{amb} - T_i). \tag{2}$$

 G_{f_amb} is given as follows.

$$\frac{1}{G_{f_amb}} = \frac{G_{f_wall} + G_{wall}}{G_{f_wall} \cdot G_{wall}} + \frac{1}{G_{wall_amb}}.$$
(3)

 G_{wall_amb} is calculated using the natural convection heat transfer coefficient. The Nusselt number to determine the natural convection heat transfer coefficient is calculated by

$$Nu = 0.74 (Gr \cdot P_{r_{air}})^{1/15}.$$
 (4)

 \dot{Q}_i in the condenser is expressed as 159

$$\dot{Q}_i = (UA/L)_{cond} \cdot dz \cdot (T_{sink} - T_i).$$
⁽⁵⁾

(UA/L)_{cond} is difficult to calculate theoretically and therefore assumed to be 4 W/(m·K) based on Ref. [20].
 The vapor quality in each node is calculated as follows to investigate fluid condition.

$$x_i = \frac{h_i - h_{l,sat}}{h_{\nu,sat} - h_{l,sat}}.$$
(6)

 h_i is calculated from Eq. (1). The vapor quality basically takes a value from 0 to 1 and shows the vapor mass flow rate. If the vapor quality is higher than 1, it means that the fluid in that node is superheated vapor. If the vapor quality is lower than 0, the fluid in that node is subcooled liquid.

168 When the fluid is superheated vapor or subcooled liquid, the Nusselt number for a laminar and a turbu-169 lent flow is expressed as

$$Nu = \begin{cases} 3.66 & (\text{laminar}) \\ 0.023 \cdot Re^{0.8} \cdot P_r^{0.3} & (\text{turbulant}). \end{cases}$$
(7)

173 The heat transfer coefficient of a two-phase flow is expressed as follows [22].

$$h_{tc_i} = h_{tc,l} \times \left\{ (1 - x_i)^{0.8} + \frac{3.8x^{0.76}(1 - x_i)^{0.04}}{(P_i/P_{cr})^{0.38}} \right\}.$$
(8)

176 The pressure in each node is calculated by

 $P_i = P_{i-1} - \Delta P_i \qquad \text{for } i \neq 1, \tag{9}$

$$P_i = P_{\nu g} - \Delta P_i \qquad \text{for } i = 1. \tag{10}$$

The pressure drop is calculated based on the Darcy-Weisbach equation if the fluid is a single-phase flow.

181 The pressure drop of a two-phase flow is calculated based on the Lockhart-Martinelli method [23]. The pres-182 sure drop of a two-phase flow is calculated by

183

$$\left(\frac{dP}{dz}\right)_{tp} = \phi_l^2 \left(\frac{dP}{dz}\right)_l.$$
(11)

184

185 ϕ_l^2 is calculated by 186

$$\phi_l^2 = 1 + \frac{C}{X} + \frac{1}{X^2},\tag{12}$$

(13)

187

188 where *X* is expressed as 189

 $X^2 = \frac{(dP/dz)_l}{(dP/dz)_n}.$

190

191 To calculate *C* in Eq. (12), Kim's model [21] is used. The value of *C* is expressed as shown in Table 1. Re_{lo} 192 and Su_{vo} in Table 1 are given as follows. 193

 $Re_{lo} = \frac{4\dot{m}}{\pi\mu_l d_{in}},\tag{14}$

$$Su_{vo} = \frac{\rho_v \sigma d_{in}}{\mu_v^2}.$$
(15)

195 196

194

Table 1	Kim's method for calculating C [21]	
		_

Liquid	Vapor	Parameter C
Turbulent	Turbulent	$0.39Re_{lo}^{0.03}Su_{vo}^{0.10}\left(\frac{\rho_l}{\rho_v}\right)^{0.35}$
Laminar	Turbulent	$0.0015 Re_{lo}^{0.59} Su_{vo}^{0.19} \left(\frac{\rho_l}{\rho_v}\right)^{0.36}$
Turbulent	Laminar	$8.7 \times 10^{-4} Re_{lo}{}^{0.17} Su_{vo}{}^{0.50} \left(\frac{\rho_l}{\rho_v}\right)^{0.14}$
Laminar	Laminar	$3.5 \times 10^{-5} Re_{lo}^{0.44} Su_{vo}^{0.50} \left(\frac{\rho_l}{\rho_v}\right)^{0.48}$

197

198 **2.2 Evaporator and reservoir**

177

199 The evaporator model is divided into the heater block, the evaporator wall, the vapor groove, and the 200 primary wick. The energy conservation equation for each part is solved to calculate temperature. The energy 201 conservation equation for the heater block is expressed as

202

$$C_{hb}\frac{dI_{hb}}{dt} = \dot{Q}_{in} - \dot{Q}_{hb_evap} - \dot{Q}_{hb_amb}.$$
(16)

204 The energy conservation for the evaporator wall is expressed as

205

$$C_{evap,wall} \frac{dT_{evap,wall}}{dt} = \dot{Q}_{hb_evap} - \dot{Q}_{evap_vg} - \dot{Q}_{evap_wick} - \dot{Q}_{evap_cc}.$$
(17)

207 The energy conservation equation for the vapor in the vapor groove is expressed as

208

$$\rho_{v}c_{p_{v}}V_{vg}\frac{dT_{vg}}{dt} = \dot{Q}_{evap_{v}g} - \dot{m}c_{p_{v}}(T_{vg} - T_{int}).$$
(18)

209

The vapor-liquid interface is saturated. The interface temperature is the saturation temperature of the pressure applied in the vapor groove. The pressure in the vapor groove is the sum of the reservoir pressure and the capillary force of the primary wick. The capillary force is equal to the total pressure drop in the LHP. Therefore, the pressure of the vapor groove is expressed as

$$P_{\nu g} = P_{cc} + \Delta P_{total},\tag{19}$$

215

214

216 where ΔP_{total} is the total pressure drop in the LHP.

The primary wick is discretized in the radial direction as shown in Fig. 3. The amount of heat conducted from node *i* to node i - 1 is expressed as

219

$$\dot{Q}_{i_{-}i-1} = \frac{2\pi L_{wick} k_{eff}}{\ln(d_i/d_{i-1})} \left(T_{wick,i} - T_{wick,i-1} \right), \tag{20}$$

220

222

223

226

221 where k_{eff} is calculated as

 $k_{eff} = (1 - \varepsilon)k_{bulk} + \varepsilon k_l. \tag{21}$

When *i* is equal to 1, the amount of heat conducted to the reservoir is expressed as

 $\dot{Q}_{wick_in} = \frac{2\pi L_{wick} k_{eff}}{\ln(d_1/d_{wick,in})} (T_{wick,1} - T_{cc}).$ (22)

The energy conservation equation at the innermost node of the primary wick is expressed as

 $\left(\rho c_{p}\right)_{eff} V_{1} \frac{dT_{wick,1}}{dt} = \dot{Q}_{2,1} - \dot{Q}_{wick_{in}} - \dot{m}c_{p_{l}} \left(T_{wick,1} - T_{cc}\right), \tag{23}$

229

232

234

230 where $(\rho c_p)_{eff}$ is calculated by 231

$$\left(\rho c_p\right)_{eff} = (1-\varepsilon)\left(\rho c_p\right)_{bulk} + \varepsilon\left(\rho c_p\right)_l.$$
(24)

233 The energy conservation equation at the innermost node of the primary wick is expressed as

$$(\rho c_p)_{eff} V_n \frac{dT_{wick,n}}{dt} = \dot{Q}_{evap_{wick}} - \dot{Q}_{n_n-1} - \dot{m}c_{p_l} (T_{int} - T_{wick,n-1}) - \dot{m}_{evap} h_{lv}$$
(25)

235

237

236 \dot{m}_{evap} is calculated by the following equation.

$$\dot{m}_{evap}h_{lv} = h_{evap}A(T_{evap,wall} - T_{int}).$$
⁽²⁶⁾

The evaporation heat transfer coefficient h_{evap} varies with temperature. The heat transfer coefficient, how-ever, is assumed to be constant. The energy conservations at other nodes are expressed as

$$\left(\rho c_{p}\right)_{eff} V_{i} \frac{dT_{wick,i}}{dt} = \dot{Q}_{i+1_{i}} - \dot{Q}_{i_{i}-1} - \dot{m}c_{p_{i}} \left(T_{wick,i} - T_{wick,i-1}\right).$$
(27)

The energy conservation in the reservoir is expressed as

$$m_{cc}c_{p_{cc}}\frac{dT_{cc}}{dt} = \dot{Q}_{wick_in} + \dot{Q}_{evap_cc} - \dot{m}c_{p}(T_{cc} - T_{ll,out}) - \frac{kA}{dz}(T_{cc} - T_{ll,out}) - \dot{Q}_{cc_amb}.$$
 (28)

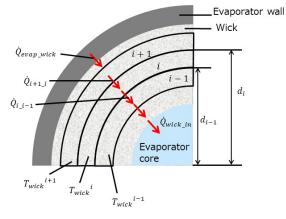


Fig. 3 Discretization of the primary wick

1	able 2 LH	P geometry	
Evaporator		Vapor line	
Length	100.0 mm	Length	800.0 mm
Outer/Inner	19.0/16.0	Outer/Inner	4.85/2.85
diameter	mm	diameter	mm
Primary wick		Liquid line	
Length	67.0 mm	Length	1100 mm
Pore radius	7.0 µm	Outer/Inner	4.85/2.85
role laulus		diameter	mm
Permeability	10^{-13}m^2	10 ⁻¹³ m ² Condenser	
Dorogity	50%	Outer/Inner	6.35/4.35
Porosity		diameter	mm
Outer/Inner	16.0/7.0	Longth	1200 mm
diameter	mm	mm Length	
Groove	2.12/3.22		
height/width mm		Reservoir	
Groove	60 mm	Outer/Inner	45.0/43.0
length	00 11111	diameter	mm
The number	7	Longth	25.0 mm
of grooves	/	Length	23.0 mm

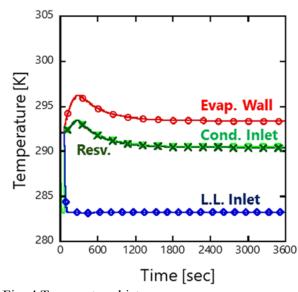
Table 2 I HP geometry

2.3 Calculation Conditions

Table 2 shows the LHP geometry used in this study. The geometry is determined based on the experi-ments of Riehl et al. [24]. The primary wick is made of an ultra-high molecular weight polyethylene, and other parts are made of stainless steel. The working fluid is acetone, and the ambient temperature is 19°C.

3. Results and discussion

255 Fig. 4 shows temperatures at the evaporator wall (Evap. Wall), condenser inlet (Cond. Inlet), liquid line 256 inlet (L.L. Inlet), and reservoir (Resv.) where the heat load is 20 W and the sink temperature is 10 °C. In this case, temperatures converge. In the experiment of Riehl et al. [24], the convergence temperature of the evap-257 258 orator was 320 K when the heat load is 20 W, and the temperature of the cooling bath is -5 °C. The temperature 259 difference between the calculation and the experiment of Riehl at al. [24] is caused by the difference of the 260 evaporation heat transfer coefficient and that of the ability to cool the condenser. The heat transfer coefficient 261 must vary with temperature, but is assumed to be constant in this model. Although the cooling bath temperature 262 is different between the calculation and the experiment of Riehl et al. [24], the temperature of condenser outlet 263 is same at 284 K. However, this LHP model can reproduce the basic operation of the LHP and is qualitatively 264 correct. Fig. 5 shows temperature when the heat load is 40 W. In this case, temperature oscillation occurs, and 265 we will discuss the internal flow based on the calculation results. 266



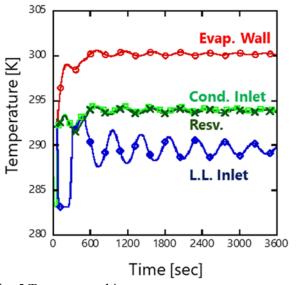


Fig. 4 Temperature history (Heat load is 20 W, sink temperature is 10 °C.)

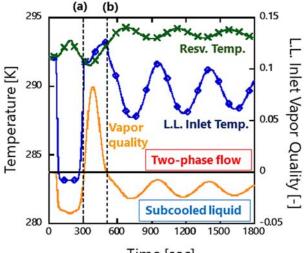
267

Fig. 5 Temperature history (Heat load is 40 W, sink temperature is 10 °C.)

268 **3.1 Internal flow at the start of temperature oscillation**

269 The internal flow at the start of the temperature oscillation is investigated. As shown in Fig. 5, the liquid 270 line inlet temperature rapidly rises before the temperature oscillation starts. To explain that reason, Fig. 6 shows 271 the reservoir temperature, the liquid line inlet temperature, and the vapor quality at the liquid line inlet. When 272 the liquid line temperature rises rapidly, the vapor quality becomes greater than 0 (the line (a) in Fig. 6). It 273 means that a two-phase flow penetrates the liquid line. This causes the rapid rise in the liquid line temperature. 274 After the two-phase flow is generated in the liquid line, the two-phase flow returns to the subcooled liquid (the line (b) in Fig. 6). Now, the liquid line temperature starts to drop, and then the temperature oscillation occurs. 275276 In Ref. [9,10], the authors predicted that the cause of temperature oscillation is the vapor penetrating to the 277 liquid line. This result is in accordance with these previous researches. However, after the two-phase flow in 278 the liquid line return to the subcooled liquid (after the line (b) in Fig. 6), it does not get regenerated in the 279 liquid line.

280 The generation of the two-phase flow is caused by the low reservoir temperature. The reservoir temper-281 ature is almost the same as the saturation temperature in the LHP because the proportion of reservoir volume to the whole is large. When the reservoir temperature is low, the vapor becomes difficult to condense, and the 282 283 two-phase flow is generated in the liquid line. Hoang et al. [16] suggested that the high degree of a liquid 284subcooling in the liquid line causes temperature oscillation. The calculation result is also in accordance with 285 Hoang's suggestion because the liquid's high degree of the subcooling will lower the reservoir temperature. 286 Therefore, when a subcooled liquid from the liquid line decreases the reservoir temperature, and when a two-287phase flow is generated in the liquid line, the temperature oscillation will start.

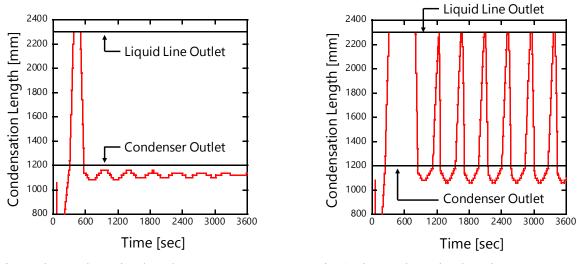


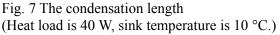
Time [sec]

Fig. 6 The relationship between the fluid condition in the liquid line and temperature in the liquid line and the reservoir

289 **3.2 Internal flow during stable temperature oscillation**

To understand the internal flow during temperature oscillation, we focus on the condensation length, which is the length from the condenser inlet to the position where the vapor is completely condensed. Fig. 7 shows the condensation length during temperature oscillation (Heat load is 40 W, sink temperature is 10 °C), and a two-phase flow is generated in the liquid line only once. However, when the sink temperature becomes 15 °C, the two-phase flow repeatedly penetrates in the liquid line (See Fig. 10).





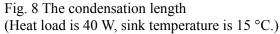
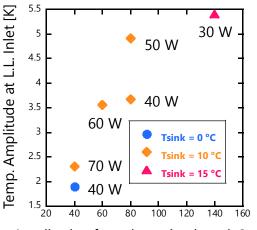
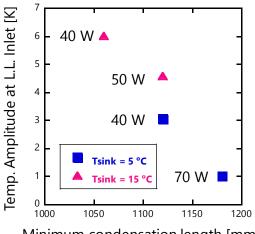


Fig. 9 shows the temperature amplitude in the liquid line inlet against the amplitude of the condensation 295 length when the condensation length oscillates in the condenser like Fig. 7. The temperature amplitude at the 296 297 liquid line inlet is proportional to the amplitude of the condensation length. Fig. 10 shows the temperature 298 amplitude in the liquid line inlet against the minimum condensation length in the case where the condensation 299 length oscillates over the condenser length like Fig. 8. The temperature at the liquid line inlet is affected only 300 by the fluctuation of the condensation length in the condenser, which can be expressed as the difference be-301 tween the condenser length and the local minimum condensation length during steady oscillation. The short 302 minimum condensation length indicates the high amplitude of the condensation length in the condenser. There-303 fore, Figs. 9 and 10 show that the high-amplitude of the condensation length causes the high-temperature 304 amplitude. The high-amplitude of the condensation length indicates that the subcooled region in the condenser 305 fluctuates significantly. The fluid temperature is decreased only in the subcooled region because the heat is 306 released for phase-change in the two-phase region. It means that the long subcooled region reduces the liquid line inlet temperature. Therefore, the fluctuation of the condenser length is important for temperature amplitude. 307 308 According to Ref. [8], the same phenomenon was observed when the high-amplitude/low-frequency oscillation occurred. Our calculation indicates that the relationship between the temperature amplitude and the condensa-309 310 tion length is the same for the high-amplitude/low-frequency oscillation and low-amplitude/high-frequency 311 oscillation.



Amplitude of condensation length [mm]

Fig. 9 The relationship between the temperature amplitude and the amplitude of the condensation length (The condensation length oscillates in the condenser.)



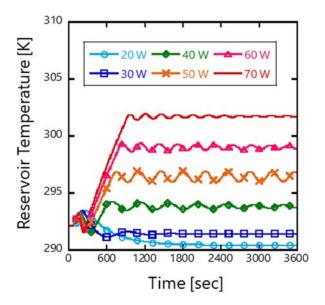
Minimum condensation length [mm]

Fig. 10 The relationship between the temperature amplitude and the minimum condensation length (The condensation length oscillates over the condenser length.)

313 **3.3 Influence of the heat load and the sink temperature**

Fig. 11 shows the reservoir temperature for each heat load. The sink temperature is 10 °C. The temperatures start to oscillate when the heat load is larger than 40 W. When the heat load is high, the degree of subcooling in the liquid line becomes low, and the reservoir temperature does not drop much. Therefore, the temperature oscillation is less likely to occur at a high heat load.

318 Fig. 12 shows the reservoir temperature for each sink temperature. The heat load is 40 W. When the sink temperature is lower than 0 °C, temperatures converge smoothly. The reason why the low sink temperature 319 320 can prevent the temperature from oscillating is related to the condensation length. As mentioned above, the 321 liquid's high degree of the subcooling is necessary to cause the temperature oscillation, and low sink temper-322 ature will lead to a high degree of the subcooling. However, when the sink temperature is low, a large amount 323 of heat is released in the narrow region. At this time, a condensation length is very hard to be changed and 324 does not change enough to cause temperature oscillation. Therefore, even if the two-phase flow is generated 325 in the liquid line due to the low reservoir temperature, temperature oscillation can be prevented by the low sink 326 temperature.



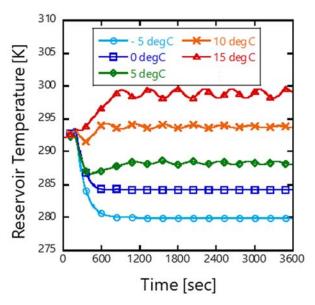


Fig. 11 Reservoir temperature at different heat loads

Fig. 12 Reservoir temperature at different sink temperature

328 4. Conclusions

We investigated the internal flow of an LHP during the temperature oscillation by using the transient model. The calculation results show that the cause of temperature oscillation is the penetration of the twophase flow into the liquid line, which is caused by the low reservoir temperature. After the penetration of the two-phase flow, when the reservoir temperature rises enough, the subcooled liquid flows in the liquid line. Even if the penetration of the two-phase flow in the liquid line occurs only once, temperature oscillation occurs. In addition, we found that when the condensation length starts to oscillate highly, the temperature amplitude becomes high. The condensation length must continue to fluctuate to keep oscillating temperature.

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337 6. References

- J. Ku, Operating Characteristics of Loop Heat Pipes, in: Proc. 29th Int. Conf. Environ. Syst., Denver,
 Colorado, 1999. doi:10.4271/1999-01-2007.
- 340 [2] Y.F. Maydanik, Loop heat pipes, Appl. Therm. Eng. 25 (2005) 635–657.
 341 doi:10.1016/j.applthermaleng.2004.07.010.
- S. Launay, V. Sartre, J. Bonjour, Parametric analysis of loop heat pipe operation: a literature review,
 Int. J. Therm. Sci. 46 (2007) 621–636. doi:10.1016/j.ijthermalsci.2006.11.007.
- V.G. Pastukhov, Y.F. Maidanik, C. V. Vershinin, M.A. Korukov, Miniature loop heat pipes for
 electronics cooling, Appl. Therm. Eng. 23 (2003) 1125–1135. doi:10.1016/S1359-4311(03)00046-2.
- 346[5]J. Li, F. Lin, D. Wang, W. Tian, A loop-heat-pipe heat sink with parallel condensers for high-power347integrated LED chips, Appl. Therm. Eng. 56 (2013) 18–26.
- 348 doi:10.1016/j.applthermaleng.2013.03.016.
- N. Putra, B. Ariantara, R.A. Pamungkas, Experimental investigation on performance of lithium-ion
 battery thermal management system using flat plate loop heat pipe for electric vehicle application,
 Appl. Therm. Eng. 99 (2016) 784–789. doi:10.1016/j.applthermaleng.2016.01.123.

- J. Esarte, J.M. Blanco, A. Bernardini, J.T. San-José, Optimizing the design of a two-phase cooling
 system loop heat pipe: Wick manufacturing with the 3D selective laser melting printing technique and
 prototype testing, Appl. Therm. Eng. 111 (2017) 407–419.
- doi:10.1016/j.applthermaleng.2016.09.123.
- J. Ku, J. Rodriguez, Low Frequency High Amplitude Temperature Oscillations in Loop Heat Pipe
 Operation, in: Proc. 33rd Int. Conf. Environ. Syst., Vancouver, Canada, 2003. doi:10.1360/zd-201343-6-1064.
- J. Ku, L. Ottenstein, M. Kobel, P. Rogers, T. Kaya, Temperature Oscillation in Loop Heat Pipe
 Operation, in: AIP Conf. Proc., 2001: pp. 255–262.
- [10] Y. Chen, M. Groll, R. Mertz, Y.F. Maydanik, S. V. Vershinin, Steady-state and transient performance
 of a miniature loop heat pipe, Int. J. Therm. Sci. 45 (2006) 1084–1090.
 doi:10.1016/j.ijthermalsci.2006.02.003.
- 364[11]S. V. Vershinin, Y.F. Maydanik, Investigation of pulsations of the operating temperature in a365miniature loop heat pipe, Int. J. Heat Mass Transf. 50 (2007) 5232–5240.
- 366 doi:10.1016/j.ijheatmasstransfer.2007.06.024.
- T. Hoang, J. Ku, Transient Modeling of Loop Heat Pipes, in: 1st Int. Energy Convers. Eng. Conf.,
 Portsmouth, Virginia, 2003. doi:10.2514/6.2003-6082.
- T. Kaya, J. Goldak, Numerical analysis of heat and mass transfer in the capillary structure of a loop
 heat pipe, Int. J. Heat Mass Transf. 49 (2006) 3211–3220.
- doi:10.1016/j.ijheatmasstransfer.2006.01.028.
- M. Nishikawara, H. Nagano, T. Kaya, Transient Thermo-Fluid Modeling of Loop Heat Pipes and
 Experimental Validation, J. Thermophys. Heat Transf. 27 (2013) 641–647. doi:10.2514/1.T3888.
- S. Launay, V. Platel, S. Dutour, J.-L. Joly, Transient Modeling of Loop Heat Pipes for the Oscillating
 Behavior Study, J. Thermophys. Heat Transf. 21 (2007) 487–495. doi:10.2514/1.26854.
- T.T. Hoang, R.W. Baldauff, J.R. Maxwell, Stability Theory for Loop Heat Pipe Design, Analysis and
 Operation, in: Proc. 45th Int. Conf. Environ. Syst., Bellevue, Washington, 2015.
- T.T. Hoang, R.W. Baldauff, C.E. Tiu, Verification of Loop Heat Pipe Stability Theory Part I Low Frequency/High-Amplitude Oscillations, in: Proc. 45th Int. Conf. Environ. Syst. Int. Conf. Environ.
 Syst., Bellevue, Washington, 2015.
- [18] V. V. Vlassov, R.R. Riehl, Mathematical model of a loop heat pipe with cylindrical evaporator and
 integrated reservoir, Appl. Therm. Eng. 28 (2008) 942–954.
- 383 doi:10.1016/j.applthermaleng.2007.07.016.
- L. Bai, G. Lin, D. Wen, Modeling and analysis of startup of a loop heat pipe, Appl. Therm. Eng. 30
 (2010) 2778–2787. doi:10.1016/j.applthermaleng.2010.08.004.
- T. Kaya, J. Ku, A Parametric Study of Performance Characteristics of Loop Heat Pipes, in: Proc. 29th
 Int. Conf. Environ. Syst., Denver, Colorado, 1999.
- S.M. Kim, I. Mudawar, Universal approach to predicting two-phase frictional pressure drop for
 adiabatic and condensing mini/micro-channel flows, Int. J. Heat Mass Transf. 55 (2012) 3246–3261.
 doi:10.1016/j.ijheatmasstransfer.2012.02.047.

- M.M. Shah, A general correlation for heat transfer during film condensation inside pipes, Int. J. Heat
 Mass Transf. 22 (1979) 547–556. doi:10.1016/0017-9310(79)90058-9.
- R.W. Lockhart, R.C. Martinelli, Proposed Correlation of Data for Isothermal Two-Phase, Two Component Flow in Pipes, Chem. Eng. Prog. 45 (1949) 39–48.
- R.R. Riehl, T. Dutra, Development of an experimental loop heat pipe for application in future space
 missions, Appl. Therm. Eng. 25 (2005) 101–112. doi:10.1016/j.applthermaleng.2004.05.010.