



UvA-DARE (Digital Academic Repository)

Prospects for Characterizing the Haziest Sub-Neptune Exoplanets with High-resolution Spectroscopy

Hood, C.E.; Fortney, J.J.; Line, M.R.; Martin, E.C.; Morley, C.V.; Birkby, J.L.; Rustamkulov, Z.; Lupu, R.E.; Freedman, R.S.

DOI

[10.3847/1538-3881/abb46b](https://doi.org/10.3847/1538-3881/abb46b)

Publication date

2020

Document Version

Submitted manuscript

Published in

Astronomical Journal

[Link to publication](#)

Citation for published version (APA):

Hood, C. E., Fortney, J. J., Line, M. R., Martin, E. C., Morley, C. V., Birkby, J. L., Rustamkulov, Z., Lupu, R. E., & Freedman, R. S. (2020). Prospects for Characterizing the Haziest Sub-Neptune Exoplanets with High-resolution Spectroscopy. *Astronomical Journal*, 160(5), [198]. <https://doi.org/10.3847/1538-3881/abb46b>

General rights

It is not permitted to download or to forward/distribute the text or part of it without the consent of the author(s) and/or copyright holder(s), other than for strictly personal, individual use, unless the work is under an open content license (like Creative Commons).

Disclaimer/Complaints regulations

If you believe that digital publication of certain material infringes any of your rights or (privacy) interests, please let the Library know, stating your reasons. In case of a legitimate complaint, the Library will make the material inaccessible and/or remove it from the website. Please Ask the Library: <https://uba.uva.nl/en/contact>, or a letter to: Library of the University of Amsterdam, Secretariat, Singel 425, 1012 WP Amsterdam, The Netherlands. You will be contacted as soon as possible.

UvA-DARE is a service provided by the library of the University of Amsterdam (<https://dare.uva.nl>)

Prospects for Characterizing the Haziest Sub-Neptune Exoplanets with High Resolution Spectroscopy

CALLIE E. HOOD ¹, JONATHAN J. FORTNEY ¹, MICHAEL R. LINE ², EMILY C. MARTIN ¹, CAROLINE V. MORLEY ³,
JAYNE L. BIRKBY ⁴, ZAFAR RUSTAMKULOV ^{1,5}, ROXANA E. LUPU ⁶, AND RICHARD S. FREEDMAN^{7,8}

¹*Department of Astronomy & Astrophysics, University of California, Santa Cruz, CA 95064, USA*

²*School of Earth and Space Exploration, Arizona State University, Tempe, AZ 85287, USA*

³*Department of Astronomy, University of Texas, Austin, TX 78712, USA*

⁴*Anton Pannekoek Institute of Astronomy, University of Amsterdam, Science Park 904, Amsterdam, 1098 XH, The Netherlands*

⁵*Department of Earth and Planetary Sciences, Johns Hopkins University, Baltimore, MD, USA*

⁶*BAER Institute / NASA Ames Research Center, Moffett Field, CA 94035, USA*

⁷*SETI Institute, Mountain View, CA 94043, USA*

⁸*NASA Ames Research Center, Moffett Field, CA 94035, USA*

ABSTRACT

Observations to characterize planets larger than Earth but smaller than Neptune have led to largely inconclusive interpretations at low spectral resolution due to hazes or clouds that obscure molecular features in their spectra. However, here we show that high-resolution spectroscopy ($R \sim 25,000$ to $100,000$) enables one to probe the regions in these atmospheres above the clouds where the cores of the strongest spectral lines are formed. We present models of transmission spectra for a suite of GJ1214b-like planets with thick photochemical hazes covering $1 - 5 \mu\text{m}$ at a range of resolutions relevant to current and future ground-based spectrographs. Furthermore, we compare the utility of the cross-correlation function that is typically used with a more formal likelihood-based approach, finding that only the likelihood based method is sensitive to the presence of haze opacity. We calculate the signal-to-noise of these spectra, including telluric contamination, required to robustly detect a host of molecules such as CO, CO₂, H₂O, and CH₄, and photochemical products like HCN, as a function of wavelength range and spectral resolution. Spectra in M band require the lowest S/N_{res} to detect multiple molecules simultaneously. CH₄ is only observable for the coolest models ($T_{\text{eff}} = 412 \text{ K}$) and only in the L band. We quantitatively assess how these requirements compare to what is achievable with current and future instruments, demonstrating that characterization of small cool worlds with ground-based high resolution spectroscopy is well within reach.

Keywords: planets, atmospheres

1. INTRODUCTION

NASA's *Kepler* mission has discovered thousands of exoplanet candidates with sizes between that of Earth and Neptune (Borucki et al. 2011; Batalha et al. 2013; Burke et al. 2014; Mullally et al. 2015; Rowe et al. 2015). These sub-Neptune planets appear to be common, around both M dwarf and Sun-like stars (Petigura et al. 2013; Fressin et al. 2013; Burke et al. 2015; Dressing & Charbonneau 2015). In fact, around a third of Sun-like stars host a planet of this size with orbital pe-

riods less than 100 days (Petigura et al. 2013; Fressin et al. 2013; Burke et al. 2015).

The measured bulk density of these planets could be consistent with a range of compositions (Figueira et al. 2009; Rogers & Seager 2010; Nettelmann et al. 2011). Correspondingly, sub-Neptunes may have a diversity of compositions from rocky to gas-rich as expected from formation and evolution modelling (Fortney et al. 2013; Moses et al. 2013; Lopez & Fortney 2014). A wide range of atmospheres is expected from their bulk compositions (Morley et al. 2017; Kempton et al. 2018). The history of how a planet accreted or outgassed its atmosphere, and its subsequent evolution, may be encoded in the abundances or ratios of molecular species in its atmosphere (Öberg et al. 2011; Booth et al. 2017; Espinoza et al.

2017). Thus, constraining the atmospheric makeup of a sample of sub-Neptune planets may be the best way to understand how and out of what material these objects form. Due to their abundance, more sub-Neptune planets are likely to be found by the TESS mission around nearby, bright M-dwarfs (Ricker et al. 2015; Sullivan et al. 2015; Barclay et al. 2018), providing prime targets for atmospheric characterization with JWST, ARIEL, and large ground-based telescopes (Louie et al. 2018; Kempton et al. 2018; Zellem et al. 2019).

However, these planets have proven hard to characterize; most atmospheric features that are detected are weaker than expected for a solar-metallicity, cloud-free atmosphere (Fraine et al. 2014; Fu et al. 2017; Wakeford et al. 2017; Wakeford et al. 2019). Some measurements have been unable to detect atmospheric features at all (e.g. Knutson et al. 2014a; Kreidberg et al. 2014). Incorporating potentially muted features into yield calculations, Crossfield & Kreidberg (2017) find that the expected yield of *TESS* planets amenable to characterization with *JWST* is up to 7x worse than when assuming cloud-free conditions.

GJ 1214b, a $6.16 \pm 0.91 M_{\oplus}$ and $2.71 \pm 0.24 R_{\oplus}$ planet around a M4.5 star (Charbonneau et al. 2009), is the prototype of this planetary class, and most dramatic example a “difficult” atmosphere. Observations of GJ 1214b taken with ground-based instruments and *HST* are consistent with a flat transmission spectrum (e.g. Bean et al. 2010, 2011; Kreidberg et al. 2014). While early observations were inconclusive, Kreidberg et al. (2014) achieved the signal-to-noise necessary to rule out just a clear but high mean molecular weight atmosphere as the source of the flat transmission spectrum. Instead, significant gray aerosol opacity has been invoked as the source of muted features in transmission spectra, including that of GJ 1214b (e.g. Crossfield et al. 2013; Kreidberg et al. 2014; Knutson et al. 2014a,b; Iyer et al. 2016; Sing et al. 2016).

Aerosols can absorb and scatter light (Heng & Demory 2013), providing an extra opacity source that dampens absorption features in transmission spectra (e.g. Deming et al. 2013; Kreidberg et al. 2014, 2018; Stevenson 2016). Morley et al. (2013) explored two types of aerosols expected to form in GJ 1214b’s atmosphere—clouds from equilibrium chemistry and a photochemical haze layer from the destruction of CH_4 —finding that either aerosol over a range of parameters could flatten the planets transmission spectrum. Further cloud formation work found that KCl and ZnS clouds can only be consistent with observations at high metallicities (1000x solar) and with strong atmospheric mixing ($K_{zz} = 10^{10} \text{ cm}^2 \text{ s}^{-1}$) (Morley et al. 2015; Gao & Benneke 2018). In con-

trast, photochemical hazes could explain the observed *HST* observations with lower metallicities $\sim 50\times$ solar (Morley et al. 2015). If an aerosol is the cause of the observed flat transmission spectra, *JWST* may allow us to characterize sub-Neptunes with its longer wavelength coverage and higher resolution than *HST* (Greene et al. 2016; Mai & Line 2019).

However, another potential avenue for studying these atmospheres is ground-based, high-resolution spectroscopy. Over the past decade, spectroscopy with a resolving power $\geq 25,000$ has been used to characterize the composition, dynamics, and thermal structure of exoplanet atmospheres (see Birkby 2018 for a recent review of the technique and resulting detections). The large variations in the radial velocity of a close-in exoplanet relative to the host star allow for the Doppler-shifted planet spectrum to be disentangled from the relatively static lines of the host star’s spectrum as well as from the spectral absorption features of Earth’s atmosphere. At high spectral resolution, molecular band heads are resolved into unique groups of individual lines allowing for robust detections from matching these lines to theoretical models. Though first suggested by Deming et al. (2000), Brown (2001), and Sparks & Ford (2002), Snellen et al. (2010) was the first robust detection of a molecule (CO) in a planet’s atmosphere using this technique with CRIRES at the VLT. Since then, molecules such as CO and H_2O have been routinely detected for a variety of planets in emission and transmission (e.g. Rodler et al. 2012; Birkby et al. 2013; Brogi et al. 2014). While most of these studies have been of hot Jupiters, a few focused on smaller planets have yielded upper limits on molecular abundances (Crossfield et al. 2011; Esteves et al. 2017; Deibert et al. 2019).

At high spectral resolution, the cores of the strongest molecular lines are formed very high up in the planet’s atmosphere, possibly above whatever cloud or haze deck may obscure features at low-resolution (de Kok et al. 2014; Kempton et al. 2014). Thus, planets whose atmospheres are completely obscured in a low-resolution transmission spectrum may still be successfully characterized at high-resolution (Birkby 2018). Pino et al. (2018) showed that not only is H_2O detectable in the presence of an aerosol for a typical hot Jupiter, but the relative cross-correlation strength across multiple wavelength ranges could be used to detect the aerosol’s presence.

The aim of this paper is to quantitatively study the feasibility of detecting molecular features of the haziest sub-Neptune planets with high-resolution transmission spectroscopy. Is this achievable? And specifically, is this a science case for current instruments, or only for

instruments on upcoming Extremely Large Telescopes? We investigate how a range of observational parameters including signal-to-noise ratio, spectral resolution, and wavelength coverage affect the detection of various molecules.

Furthermore, as high-resolution spectroscopy has become more common, the best way to robustly report detection significances of molecules has been explored. Brogi & Line (2019) proposed a new log likelihood function to use when comparing observed spectra to a model in the place of the traditional cross-correlation function. We will compare the utility of these two metrics, motivating our choice of the log likelihood function for the majority of this work. However, this method may be affected by any “missing” molecules present in the observed atmosphere but absent in the model spectrum. For illustration, we will consider how HCN, a high abundance photochemical product, affects the observed spectra and reported detection significances.

This work is structured as follows. In section 2, we describe how we generate transmission spectra for hazy GJ 1214b analogs and discuss how we quantify the significance of molecular detection with this technique. In Section 3, we give an overview of the prospects for observing CO, CO₂, H₂O, and CH₄ across a range of planetary insolation levels as well as the spectral resolution and wavelength coverage of the data. A discussion of our results are presented in Section 4 and our conclusions are summarized in Section 5.

2. METHODS

2.1. Model Atmosphere and Spectra

We generate high-resolution transmission spectra based on the 1D radiative-convective-photochemical models presented in Morley et al. (2015). The authors assume 50× solar metallicity and use a 1D radiative-convective model to determine a temperature-pressure profile for the atmosphere assuming radiative-convective equilibrium and calculate gas abundances in different layers of the atmosphere assuming chemical equilibrium. They find the total mass of soot precursors from a photochemical model (Miller-Ricci Kempton et al. 2012; results first published in Fortney et al. 2013) at each layer and assume some percentage (f_{haze}) will form a scattering haze at that layer, with f_{haze} and the mode particle size as free parameters. The optical properties of this haze are calculated with Mie theory. For this study, we focus on a particular combination of parameters that reproduce the “flat” Kreidberg et al. (2014) observations at low-resolution: $f_{\text{haze}} = 10\%$ and a mode particle size of 0.1 microns. This haze becomes opaque at a pressure of approximately 10^{-5} bar in the atmosphere. We also

look at models for atmospheres with $0.3\times$ and $3\times$ the insolation of GJ 1214 b with the same haze parameters but not the exact same haze. The atmospheres with $0.3\times$, $1\times$, and $3\times$ the insolation of GJ 1214 b have effective temperatures of 412, 557, and 733 K, respectively (Fortney et al. 2013). The atmosphere models from Morley et al. (2015) go to 10^{-6} bars at the top of the atmosphere, but our highest resolution spectra are sensitive out to $\sim 3\times 10^{-7}$ bars, so we assume an isothermal atmosphere above 10^{-6} bars with constant molecular abundances. The resulting pressure-temperature profiles and molecular abundances we use are shown in Figure 1.

To produce high-resolution transmission spectra, we use the flexible radiative transfer code described in the appendix of Morley et al. (2017). This line-by-line code takes in the temperature-pressure profiles, chemical abundance profiles, and haze opacity files from (Morley et al. 2015), as well as a mass and radius of the planet. Using the the line-by-line optical depth calculations and the vectorized method for calculating transmission spectra presented in (Robinson 2017), this code outputs high resolution line-by-line ($R \sim 500,000$) transmission spectra for the planet. We use the cross-section database described in (Freedman et al. 2014). For our calculations we include the opacities of CO (Rothman et al. 2010), CO₂ (Huang et al. 2014, 2013), H₂O (Barber et al. 2006), CH₄ (Yurchenko & Tennyson 2014; Yurchenko et al. 2013), HCN (Harris et al. 2008), and H₂/He collision-induced absorption (Richard et al. 2012). Absorption cross sections for these molecules are shown in Figure 2. A molecule will be easiest to detect where it has the highest cross sections/strongest spectral features, for example H₂O in *J* and *H* bands, CH₄ in *L* band, and CO in *M* band.¹

Example transmission spectra in *K* band for a range of resolutions are shown in Figure 3. Though the low resolution transmission spectrum shows little deviation from a flat line, resolved spectral features from CO and H₂O are visible starting with $R \sim 10,000$ and increase in size as spectral resolution increases. High and low resolution spectra from 1 to 5 μm with and without the haze for all three insolation cases are shown in Figure 4. The haze opacity effectively obscures the molecular features below a certain pressure in the atmosphere, reducing the low resolution spectrum in particular to a mostly flat line. As the stellar insolation (and therefore the effective temperature of the atmosphere) increases, CH₄

¹ Although there have been updates to line lists for certain species since these publications, we do not include them in this study as we are not comparing to observations and are thus internally consistent.

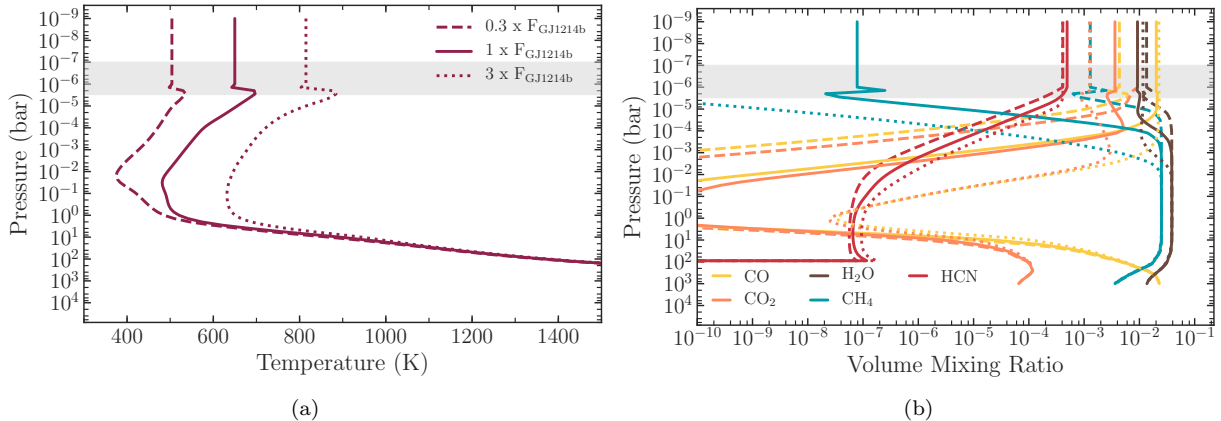


Figure 1. Pressure-temperature profiles and molecular abundances for models with $50\times$ solar metallicity and $0.3\times$ (dashed), $1\times$ (solid), and $3\times$ (dotted) GJ 1214b’s insolation. The high-resolution observations discussed in this paper are most sensitive roughly between 3×10^{-6} - 3×10^{-7} bars, marked by the light grey region in each plot.

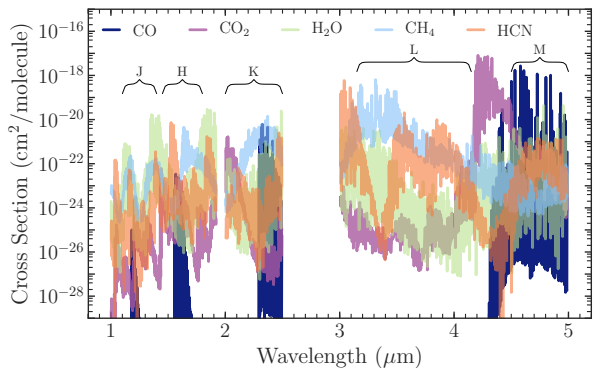


Figure 2. Absorption cross sections for the molecules included in our spectra. These cross sections are calculated at a pressure of 10^{-6} bars and a temperature of 650 K, then smoothed to $R \sim 1000$ for illustrative purposes. Molecules with the strongest features in a particular bandpass, e.g. CO in M band, will be the dominant species in that wavelength range, though this effect is also dependent on the abundance of the molecule (see Figure 1b).

features in L band disappear while CO shows stronger features in K and M band, in accordance with the change in abundances shown in Figure 1b.

Previous analyses of high resolution spectra have involved various methods to remove telluric and stellar contamination of the data which all require normalization of the observed spectra. To approximate an observed and reduced spectrum, we take a $R \sim 500,000$ transmission spectrum model including all opacity sources as described above, Doppler shift it according to the systemic velocity of GJ 1214 (21 km/s), and then to reach the desired spectral resolution smooth with a Gaussian kernel and interpolate the model onto a coarser wavelength grid corresponding to the given

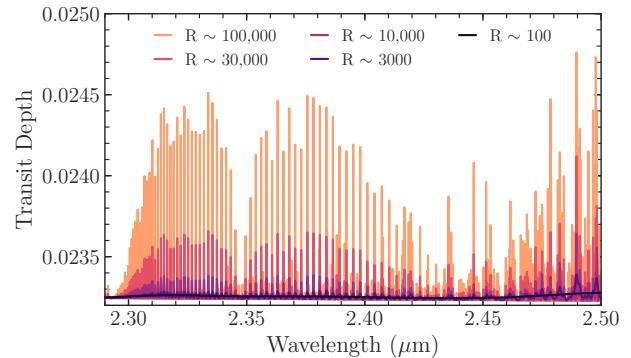
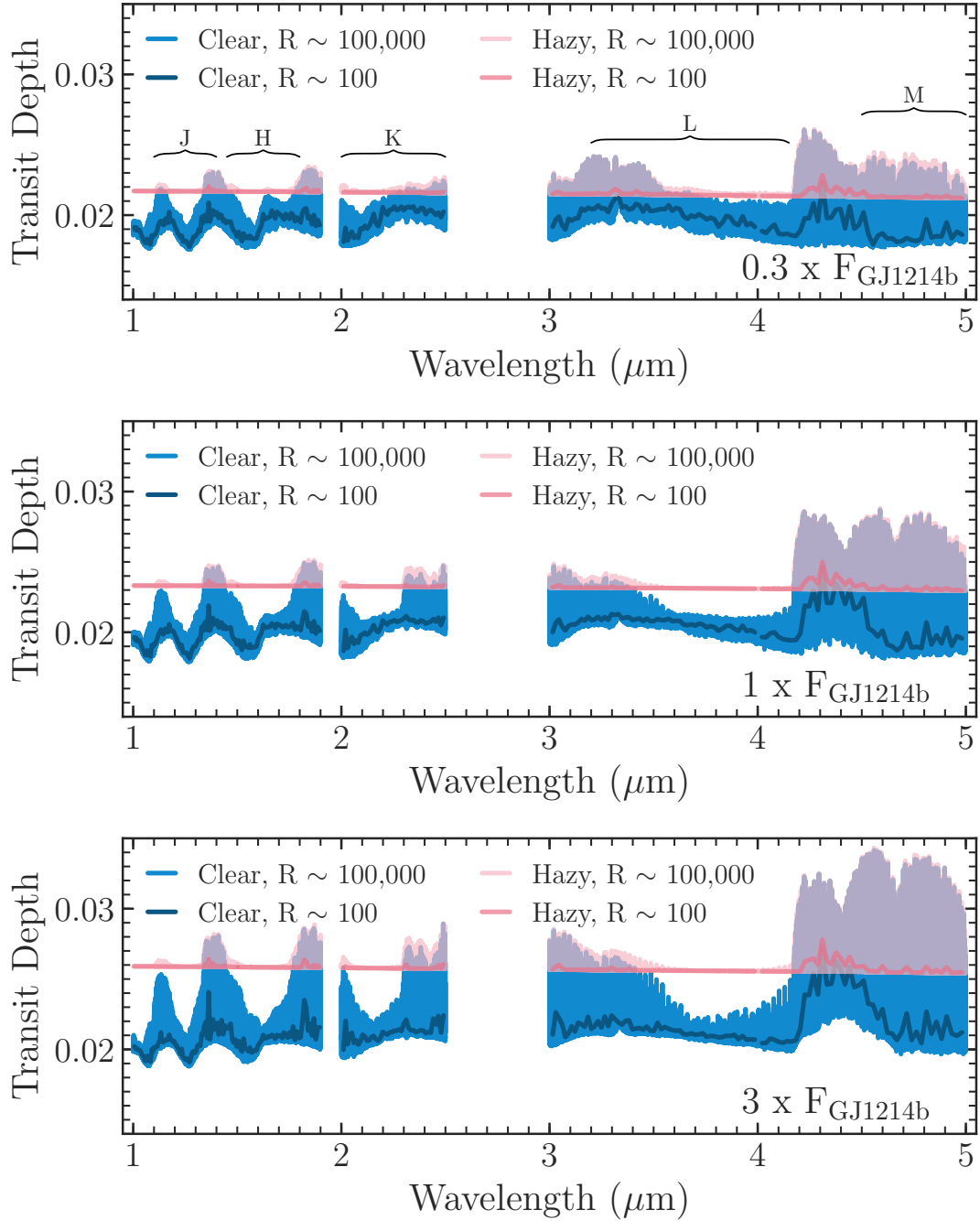


Figure 3. Transmission spectra for a GJ 1214b model at the nominal insolation level across a range of spectral resolutions. At low resolution, the transmission spectrum is essentially a flat line, but spectral features are clearly visible at higher resolutions.

resolution and 2 pixels per resolution element (assuming a Nyquist like sampling of 2 pixels per element). We assume telluric and stellar contamination will be dealt with completely by the data reduction process such that they are removed to the photon noise level. Consequently, we only simulate the planets transmission spectrum, and investigate the effect of the photon noise. While we do not assume a full noise simulator and data processing steps (e.g. Brogi & Line 2019), we approximate the final outcome of such an approach by doing the following. First, to determine the amount of random noise to add to each pixel in our truth spectrum, we choose a particular signal-to-noise ratio on what would have been observed in the stellar spectrum per resolution element (S/N_{res}) as in Pino et al. (2018). For the signal-to-noise per pixel, $S/N_{\text{pix}}(\lambda)$ we multiply this S/N_{res} by the square root of the telluric absorption spectrum $T(\lambda)$ (which has values between 0 and

Figure 4. Clear and hazy transmission spectra for models with $0.3\times$ (top), $1\times$ (middle) and $3\times$ (bottom) GJ 1214b’s insolation. Both low ($R \sim 100$) and high ($R \sim 100,000$) resolution spectra are plotted. The presence of the haze clearly mutes molecular features, particularly in the low resolution spectra. Furthermore, at high resolution the difference between the coolest and hottest models is most pronounced; this difference can be attributed to the larger scale height for the hotter model in addition to differences in abundances.



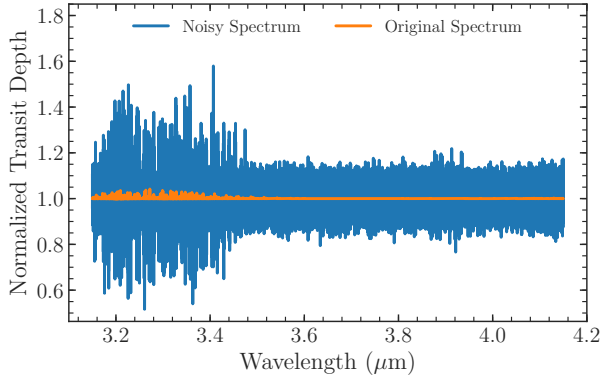


Figure 5. Normalized transmission spectra with $R \sim 100,000$ for a GJ 1214b model before and after the addition of noise. The random noise added to the spectrum has a base S/N_{res} of 1000 (or 1000 ppm of noise), scaled by a model of the Earth’s transmission spectrum.

1), to mimic the reduction in S/N due to telluric extinction, and the square root of the number of pixels per resolution element. We simulate noise by adding a noise value to each pixel drawn from a Gaussian distribution with a standard deviation equal to the reciprocal of the desired signal-to-noise per pixel. Given that we are assuming photon noise, which follows a Poisson distribution, Gaussian distributed noise is an appropriate approximation for the high stellar photon counts. We obtain $T(\lambda)$ using the ESO Skycalc tool based on the Cerro Paranal Sky Model (Noll et al. 2012; Jones et al. 2013). An example of the normalized transmission spectrum in L band before and after the addition of noise is shown in Figure 5. Both the original GJ 1214b transmission spectrum and that of Earth’s atmosphere have stronger features on the bluer end of L band, reflected in the noisier but more prominent features on the shorter wavelength end of the noisy spectrum.

2.2. Quantifying Detection Significance

High-resolution spectra of exoplanet atmospheres are often “self-calibrated” rather than in comparison to a standard star, meaning broadband information and changes in flux at a fixed wavelength over time are removed from the data by fitting a trend with airmass or using a principal component analysis based approach (e.g. Snellen et al. 2010; Birkby et al. 2013). These data processing steps to remove the telluric contamination typically remove any reliable planetary continuum level information. This makes typical data-model “chi-square” comparisons difficult, if not impossible, unless the exact stretching/scaling to the data is known. In light of this, the standard approach is to utilize the cross-correlation function (CCF), which leverages information in individual line ratios to determine planetary

atmosphere information (see Brogi & Line 2019 and Gibson et al. 2020 for a detailed discussion). The cross-correlation function determines the correlation between the data and a model template as a function of Doppler shift. A model perfectly matched to the data will show a “peak”, or maximum correlation, at the planetary velocity (in our case, the systemic velocity). Incorrect models will show no peak or a damped peak relative to the “correct” model. Following the notation of Brogi & Line (2019), we define the variance of the data (s_f^2), the variance of the model (s_g^2), and the cross-covariance $R(s)$ as follows:

$$s_f^2 = \frac{1}{N} \sum_n f^2(n)$$

$$s_g^2 = \frac{1}{N} \sum_n g^2(n-s)$$

$$R(s) = \frac{1}{N} \sum_n f(n)g(n-s)$$

where n is the bin number or spectral channel, s is a bin/wavelength shift, N is the total number of spectral channels, $f(n)$ is an observed spectrum, and $g(n)$ is a model spectrum for comparison, both mean-subtracted.

The cross-correlation coefficient $C(s)$ is then:

$$C(s) = \frac{R(s)}{\sqrt{s_f^2 s_g^2}} \quad (1)$$

In the literature, molecules have been detected by reporting a strong signal in the cross-correlation function of observed spectra with a model that contains solely the molecule of interest. When using the CCF, we compare to models that have only one or two opacity sources at a time, e.g. just CO or CO and the haze opacity. The strength of this detection has often been reported as the ratio of the peak of the CCF and the standard deviation of the coefficient around the peak (e.g. Snellen et al. 2010) though some more sophisticated approaches have also been used like the Welch T-test metric (e.g. Birkby et al. 2013). Hawker et al. (2018) find the ratio of peak to standard deviation to be the most conservative metric for evaluating detection significance, so this is how we will report detection significances from a CCF in this work. However, this peak-to-off-peak comparison only determines the S/N within a given model template relative to the on-to-off velocities, making quantitative comparisons amongst differing model templates challenging.

Brogi & Line (2019) proposed a solution to this problem by developing a mapping of the CCF to a log-likelihood function ($\log(L)$) for use in a Bayesian retrieval framework. Using the above definitions, they

related a formal log-likelihood to the CCF (or rather, the cross-covariance) through:

$$\log(L) = -\frac{N}{2} \log(s_f^2 - 2R(s) + s_g^2) \quad (2)$$

where as defined above N is the total number of spectral channels, s_f^2 is the variance of the data, s_g^2 is the variance of the model, and $R(s)$ is the cross-covariance of the data and a model with some wavelength shift s . Before proceeding with our atmospheric analysis, we first investigate the sensitivity of the model comparisons under the CCF and $\log(L)$ assumptions. Since we are not performing a retrieval analysis but instead comparing to forward models, we need a method of mapping the $\log(L)$ value to a detection significance analogous to that obtained with the CCF. To do this, we first calculate $\log(L_1)$ for our “truth” spectrum (e.g., the same underlying model used to generate the simulated data, which again, includes the haze continuum and all of the gases). We then compute $\log(L_2)$ for 5 additional nested models that each lack one of our tested opacity sources, so we can isolate how much that missing opacity source decreases the $\log(L)$. To quantify the detection of each source of opacity, we utilize the change in Bayesian Information Criterion (BIC) (Schwarz 1978):

$$\text{BIC} = p \log N - 2 \log L \quad (3)$$

between the full model and the subset model lacking that opacity source, of which we then approximate $\Delta_{\text{BIC}} \approx 2 * (\log L_1 - \log L_2)$.² We then relate this change in BIC to the Bayes factor, using the formula $\Delta_{\text{BIC}} = 2 * \log B_{12}$ (e.g. Szydlowski et al. 2015). We can then map this Bayes factor to a frequentist p-value using Table 2 from Trotta (2008), which can in turn be converted into a statistical significance. Due to the limited nature of the table, we are only able to report significances smaller than 21.3σ ; anything that would have a stronger significance is reported as this upper limit.

2.3. CCF vs. Log(L) Example: CO in K Band

Here we compare molecular detections from the CCF and $\log(L)$ approaches as a function of the signal-to-noise per resolution element S/N_{res} for a representative K -band (2-2.5 μm) spectrum of our nominal GJ 1214b model.

First, we look at a single truth spectrum with S/N_{res} of 1000 and $R \sim 100,000$, assuming a velocity of 0 km/s

for the planet. Figure 6 shows the CCFs of this spectrum with models that contain either all opacity sources, just CO and the haze, or just CO; the statistical significances of each peak are 9.9σ , 9.7σ , and 9.1σ , respectively³. As expected, the template that contains all of the opacity sources included in the original model gives the highest peak CCF, relative to the off-peak velocity baseline. The template that contains both CO and the haze has a slightly smaller peak, which again decreases when the haze opacity is ignored. However, these decreases are relatively small, suggesting the CCF is not particularly sensitive to the presence of a haze.

Figure 7 similarly shows the $\log(L)$ as a function of velocity for templates that include all opacities, all but CO, and all but the haze. Figure 7a is zoomed in to show how the peak seen when computing $\log(L)$ for the full model disappears when CO is removed, highlighting the necessity of CO to properly match the original spectrum. The decrease in $\log(L)$ at the planet velocity (0 km/s) corresponds to a 10.9σ detection of CO (based on Eqn. 3 and subsequent discussion). We note that the $\log(L)$ for the model without CO has a higher average value than the model with CO; we attribute this effect to the decreased variance (s_g) of the model without the prominent CO lines, which increases the resulting $\log(L)$ as seen in Equation 2. Thus, a clear peak in the $\log(L)$ as a function of velocity should be taken as an indication that a model that correctly matches the truth spectrum rather than just the average value. We assume a retrieval method would identify the correct velocities and so only consider the value at 0 km/s as this is most analogous to what a retrieval detection significance would be. Figure 7b shows the same curves, as well as $\log(L)$ with a template not containing the haze. When the haze opacity is removed, though the peak in $\log(L)$ at 0 km/s is still prominent, $\log(L)$ significantly decreases at this velocity, corresponding to a $> 20\sigma$ detection of the haze. Therefore, the haze is also necessary in addition to CO to match the truth spectrum. However, if one did not test the models containing the haze opacity, they may still detect CO since the $\log(L)$ still shows a peak at 0 km/s but assume the atmosphere was clear. Models including a potential source of continuum opacity should be investigated to maximize the atmospheric information one can learn from the data.

² Effectively, this becomes a likelihood ratio test to compare models as we are not changing the number of free parameters (p in Equation 3). The molecule we remove is not a free parameter as we do not vary its value to fit the data.

³ Typically in the literature this metric is referred to as the S/N of a detection and lacks the sigma symbol (unlike significance values from the Welch T-test metric for example). However, we will include the sigma in this work to aid in comparison to detection significances derived from the $\log(L)$ method.

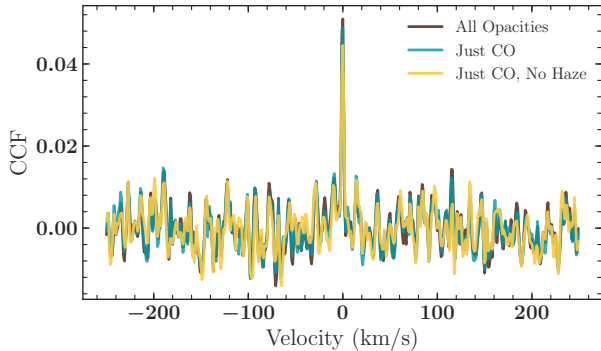


Figure 6. CCFs of transmission spectrum models that include varying opacity sources with a spectrum that includes all opacities and has random noise added to give it a S/N per resolution element of 1000.

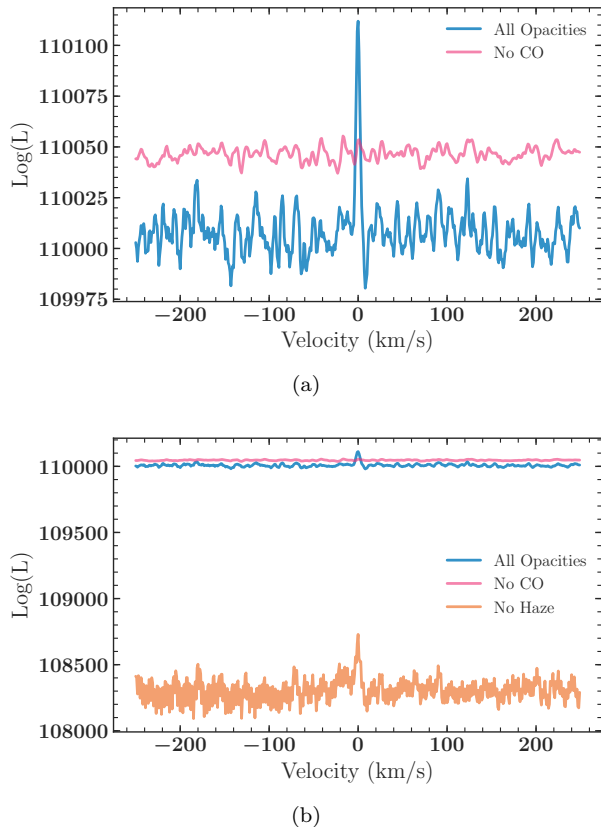


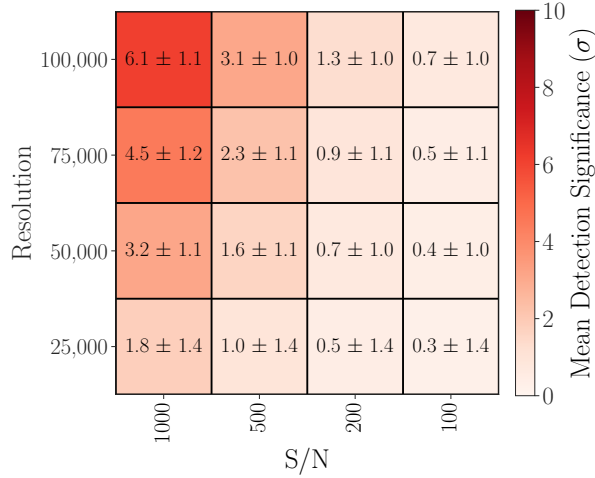
Figure 7. Log(L) for transmission spectrum models that include varying opacity sources with a spectrum that includes all opacities and has random noise added to give it a S/N per resolution element of 1000. In 7a, we see that the peak in log(L) seen at zero relative velocity disappears when CO is removed from the model. In 7b, we see the value of log(L) drastically decreases when the haze opacity is removed from the model.

Next we explore the differences in detection from the CCF vs. $\log(L)$ approach over a small grid of S/N_{res} and resolutions, summarized in Figures 8 and 9. In order to account for the effect of random noise, we repeat the analysis for 25 different noise instances and report the average and standard deviation of the resulting detection strengths. Comparing Figures 8a and 8b, we see that removing the haze opacity does typically slightly lower the average detection strength of CO when using the CCF regardless of spectral resolution or S/N_{res} . However, within their uncertainties, the detection strengths of CO agree between templates that include or ignore the haze opacity. Figures 9a and 9b show the detection strengths for CO and haze, respectively, when using the $\log(L)$ method. We see that 9a resembles 8a and 8b, meaning that the CCF and $\log(L)$ give similar detection strengths for CO. Figure 9b shows we can strongly detect the presence of a haze, even when we cannot robustly detect CO (for example, with $R \sim 25,000$).

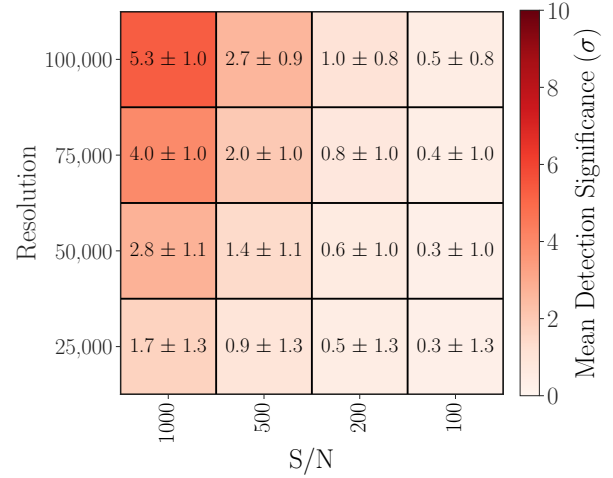
Overall, we find that CCF and $\log(L)$ give similar answers for molecular detections, but the $\log(L)$ method is much more sensitive to the presence of a haze opacity. This suggests the high-resolution spectrum is sensitive to the broadband opacity of the haze due to the loss of a myriad of weaker lines. The increased sensitivity of the $\log(L)$ is attributable to the treatment of the model and data variance terms (s_f and s_g). As noted in Brogi & Line (2019), the $\log(L)$ decreases when s_f and s_g differ significantly, while the CCF is not clearly affected by this discrepancy. The presence of a haze opacity will serve to mute any molecular features (as seen in Figure 4), effectively decreasing the variance of the spectrum. As a result, a model that includes the haze will have a more similar variance to the observed spectrum of a hazy object, leading to a detectable change in $\log(L)$. Thus, for the remainder of this paper, we will only use the $\log(L)$ method to quantify how well we can detect opacity sources.

3. RESULTS

Here we present the detectability given the above metrics for CO, CO₂, H₂O, CH₄, and a haze opacity as a function of spectral band, resolution, and S/N_{res} . The “truth” spectrum is generated by including all of the above opacity sources as well as H₂/He collision-induced absorption from one of three input pressure-temperature profiles, corresponding to 0.3 \times , 1 \times , or 3 \times GJ 1214 b’s stellar insolation (Figure 1). The observing bands we consider and their corresponding wavelengths are shown in Table 1. These bandpasses are a bit wider than are typically defined for ground based observations as at

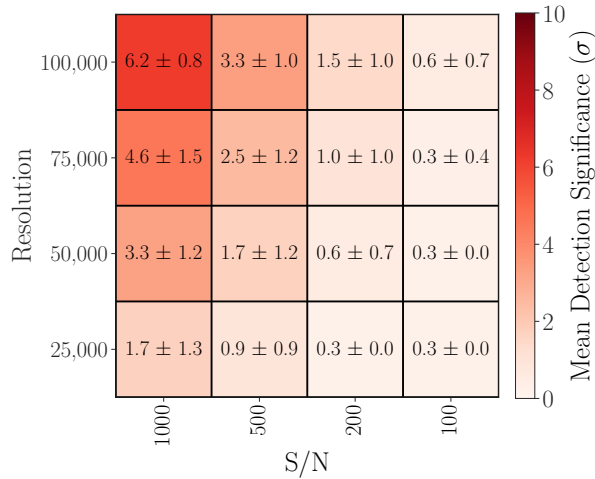


(a) CCF: CO and Haze

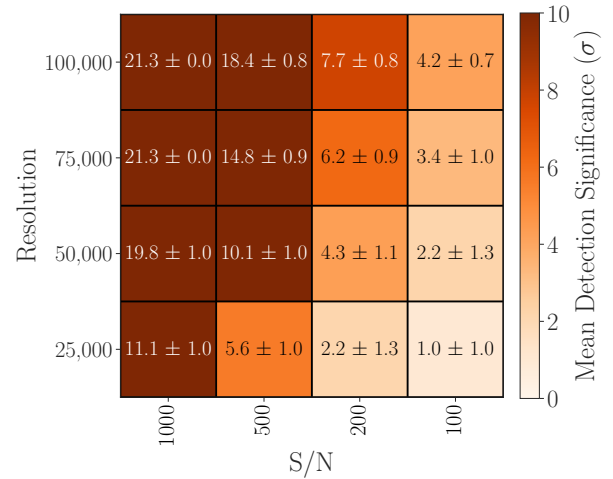


(b) CCF: Just CO

Figure 8. Detection significance in sigma for CO as a function of spectral resolution and S/N per resolution element of the transmission spectrum of nominal GJ 1214b model ($1\times$ insolation). In 8a, we have cross-correlated our input spectra with models including opacity only from CO and a haze. In 8b, we have removed the haze opacity from the model we use for cross-correlation; the detection significances agree with those obtained when including the haze opacity, making it difficult to robustly identify the presence of a haze when using the CCF.



(a) Log(L): CO



(b) Log(L): Haze

Figure 9. Detection significances in sigma for CO and a haze as a function of spectral resolution and S/N per resolution element of the transmission spectrum of nominal GJ 1214b model ($1\times$ insolation). The detection strengths reported in 9a are very similar to those found when detecting CO with the CCF as shown in 8a. However, as this method allows us to probe the presence of a haze directly, we can much more confidently report the detection of a haze using the log(L) as shown in 9b.

Table 1. Observing bands considered in this study.

Observing Band	Wavelength Coverage (microns)
<i>J</i>	1.1 - 1.4
<i>H</i>	1.45 - 1.8
<i>K</i>	2.0 - 2.5
<i>L</i>	3.2 - 4.15
<i>M</i>	4.4 - 5.0

high resolution telluric absorption features can possibly be resolved with usable data in between, and instruments can have varying wavelength coverage. We test a large range of S/N_{res} from 50 to 5000, averaging the detectability over 25 random noise instances (as in Section 2.3). We then find the average detection strength from these 25 noise instances. Tables 3 - 2 report the lowest S/N_{res} required for detecting (threshold for detection set at 5σ) a given opacity source (one table per source) as a function of observing band, insolation (relative to that of GJ 1214b), and spectral resolution. We will go into a more detailed overview of these results in the following sections.

3.1. $1\times$ GJ 1214b Insolation

Three findings are consistently true regardless of wavelength range considered. First, as indicated in Table 2, the haze is always the easiest opacity source to detect since it has the lowest required S/N_{res} for detection, which is 500 or less in all cases considered. Thus, if one achieved the S/N_{res} necessary to detect a molecule like CO in the atmosphere of one of these planets, one would necessarily have the required S/N_{res} to rule out a completely clear atmosphere as well. Again, this is made possible through the s_g and s_f terms in the log-likelihood function and would thus be difficult if not impossible to detect using the classic CCF approach.

Second, increasing the spectral resolution appears to have diminishing returns, i.e. increasing the spectral resolution from $R \sim 25,000$ to 50,000 yields the greatest decrease in the required S/N_{res} (often a factor of 2) but increasing the spectral resolution further does not yield quite as dramatic of a change in the required S/N level. This is because once enough strong spectral lines are resolved to clearly identify the presence of a molecule, adding additional weaker lines from further increased spectral resolution is an increasingly marginal help. However, increasing to higher spectral resolution may still lead important gains in precision on parameters beyond detection of molecules, such as molecular

Table 2. Minimum S/N_{res} required for $\geq 5\sigma$ detection of the haze.

Observing Band	Insolation	Spectral Resolution			
		25,000	50,000	75,000	100,000
(1)	(2)	(3)	(4)	(5)	(6)
<i>J</i>	0.3 \times	400	250	200	150
	1 \times	350	200	150	150
	3 \times	200	150	100	100
<i>H</i>	0.3 \times	350	200	150	150
	1 \times	350	200	150	100
	3 \times	200	100	100	100
<i>K</i>	0.3 \times	400	250	200	150
	1 \times	450	250	200	150
	3 \times	300	150	100	100
<i>L</i>	0.3 \times	300	200	150	100
	1 \times	400	250	200	150
	3 \times	300	200	150	100
<i>M</i>	0.3 \times	350	200	150	150
	1 \times	250	150	100	100
	3 \times	200	150	100	100

Table 3. Minimum S/N_{res} required for $\geq 5\sigma$ detection of CO. There are no $\geq 5\sigma$ detections with $S/N_{res} \leq 5000$ for *J*, *H*, and *L* bands.

Observing Band	Insolation	Spectral Resolution			
		25,000	50,000	75,000	100,000
(1)	(2)	(3)	(4)	(5)	(6)
<i>K</i>	0.3 \times	-	4500	3500	2400
	1 \times	3500	1600	1100	800
	3 \times	1800	1000	700	500
<i>M</i>	0.3 \times	1400	800	600	400
	1 \times	700	400	250	200
	3 \times	400	250	150	150

Table 4. Minimum S/N_{res} required for $\geq 5\sigma$ detection of CO_2 . There are no $\geq 5\sigma$ detections with $S/N_{res} \leq 5000$ for J , H , K , and L bands.

Observing Band	Insolation	Spectral Resolution			
		25,000	50,000	75,000	100,000
(1)	(2)	(3)	(4)	(5)	(6)
M	$0.3\times$	4500	2400	1800	1300
	$1\times$	4500	2400	1700	1300
	$3\times$	4500	2200	1500	1100

Table 5. Minimum S/N_{res} required for $\geq 5\sigma$ detection of H_2O .

Observing Band	Insolation	Spectral Resolution			
		25,000	50,000	75,000	100,000
(1)	(2)	(3)	(4)	(5)	(6)
J	$0.3\times$	-	2600	1900	1200
	$1\times$	-	2800	2200	1300
	$3\times$	3000	1600	1200	700
H	$0.3\times$	-	4500	3000	2200
	$1\times$	-	4000	2800	2000
	$3\times$	3500	1800	1200	900
K	$0.3\times$	-	-	4000	3000
	$1\times$	-	-	4000	3000
	$3\times$	5000	2800	1900	1500
L	$0.3\times$	-	4500	3500	2600
	$1\times$	-	4000	3000	2600
	$3\times$	3000	1500	1300	1100
M	$0.3\times$	5000	2800	1900	1300
	$1\times$	4000	2400	1600	1200
	$3\times$	1800	1000	700	500

abundances, temperature structure, or wind speeds that we do not consider in this work.

Lastly, CH_4 is undetectable across all wavelength ranges and spectral resolutions for models with this insolation level. This inability to detect CH_4 is expected

Table 6. Minimum S/N_{res} required for $\geq 5\sigma$ detection of CH_4 . There are no $\geq 5\sigma$ detections with $S/N_{res} \leq 5000$ for J , H , and M bands.

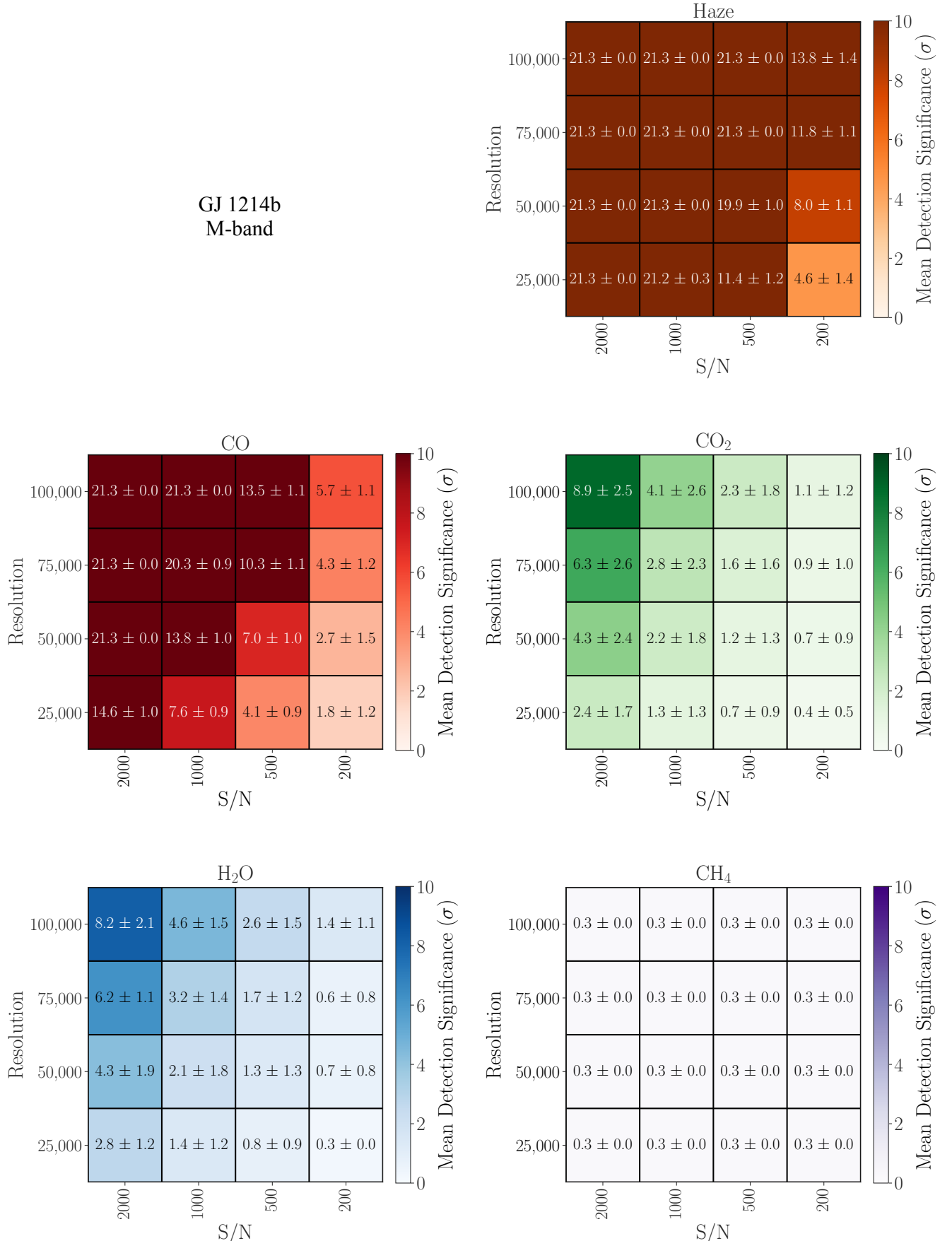
Observing Band	Insolation	Spectral Resolution			
		25,000	50,000	75,000	100,000
(1)	(2)	(3)	(4)	(5)	(6)
K	$0.3\times$	-	3000	2000	1600
	$1\times$	-	-	-	-
	$3\times$	-	-	-	-
L	$0.3\times$	1100	700	450	300
	$1\times$	-	-	-	-
	$3\times$	-	-	-	-

due to the lack of CH_4 above $\sim 10^{-4}$ bar (where these observations are most sensitive) as shown in Figure 1b.

More specifically, here we break-down the band-by-band results:

- J and H Bands: H_2O is the only molecular opacity source detectable in addition to the haze. The required S/N_{res} for detecting H_2O is lower in J band than in H - 2800 and 4000, respectively, for $R \sim 50,000$.
- K Band: CO is detectable for all spectral resolutions, while H_2O is detectable for $R \geq 75,000$. However, H_2O requires much higher S/N spectra. For example, at a $R \sim 75,000$, one needs an effective S/N_{res} of 1100 to detect CO but 4000 for H_2O .
- L Band: H_2O is again the only detectable molecule. A higher S/N_{res} is required to detect H_2O compared to J or H , with a required value of 4000 for $R \sim 50,000$.
- M Band: CO , CO_2 , and H_2O are all potentially detectable. Plots of the detection strength for all considered opacities as a function of selected spectral resolutions and S/N s are shown in Figure 10. Detecting CO is significantly easier than in K ; at a $R \sim 75,000$, one only needs an effective S/N_{res} of 250 to detect CO . In addition, CO_2 is only detectable in this wavelength range, though it does require a higher S/N_{res} than that needed to detect CO (for example, 1700 at a $R \sim 75,000$). H_2O requires comparable S/N_{res} to CO_2 , but with

Figure 10. Detection significances for different opacity sources as function of spectral resolution and S/N_{res} for M band transmission spectra of the model with $1\times$ GJ 1214b's insolation. All tested opacity sources except CH_4 are detectable in M band, though the haze and CO are detectable for a much wider range of combinations of S/N_{res} and spectral resolution than CO_2 and H_2O .



spectra at $R \sim 100,000$, one could detect CO, CO₂, and H₂O with $S/N_{res} \geq 1100$. Thus, the M band is overall the most promising observing band for detecting 2 or more molecular opacity sources at once. However, the high thermal background in M band may make these observations more challenging as discussed in Section 4.1.

In general, the bands in which different molecules are detectable is our work here do reflect previous observational results. H₂O has been detected in J band (Alonso-Floriano et al. 2019), K band (e.g. Hawker et al. 2018), L band (e.g. Birkby et al. 2013, 2017; Piskorz et al. 2018), and data covering 0.95 - 2.45 μm simultaneously (Brogi et al. 2018; Guilluy et al. 2019). CO has been repeatedly detected in K band (e.g. Snellen et al. 2010; Flagg et al. 2019). To date, no molecules have been reported for high-resolution ground-based M band spectra of an exoplanet atmosphere.

de Kok et al. (2014) investigated the optimal wavelength ranges to detect different molecules in exoplanet atmospheres with the 2014-era CRRES instrument on the VLT. Similar to our results presented above, they also find that CO is best detected in K and M bands, while H₂O is detectable in J - M bands. However, they find the region around 3.5 μm as optimal for detecting multiple species (CO₂, CH₄, and H₂O), in contrast to our preference for M band. As we will discuss later in Section 4.1, the thermal background in M band can make it functionally difficult to reach the required S/N_{res} with current instruments. In addition, they assume a truth spectrum with only one molecule at a time and cross-correlate with a model of just that species; they note that CH₄ or H₂O lines could have a shielding effect that make other molecules more difficult to detect. Thus, our inability to also detect CO₂ in L band can be attributed to the higher cross sections of the many lines of CH₄ and H₂O in this region as shown in Figure 2.

3.2. 0.3 \times and 3 \times Insolation

Since we would like to be able to observe a wider range of planets than just GJ 1214b, we also investigate how changing the stellar insolation effects our results, looking at models with 0.3 \times and 3 \times the true insolation of GJ 1214b. The primary effect is the change over of the dominant carbon bearing species with temperature: CH₄ dominating in the 0.3 \times case and CO/CO₂ in the 3 \times scenario, with an overlap/transition for the 1 \times case (Figure 1b).

3.2.1. 0.3 \times Insolation

As above, we list our results by observing band, but to simplify the discussion, only summarize the differences

from the 1 \times insolation case. In some cases the haze opacity requires a slightly higher S/N_{res} to detect than the nominal insolation case, but it still remains the easiest opacity source to detect across all wavelength bands and resolutions. Overall, the decrease in temperature for the models with 0.3 \times GJ 1214 b's insolation makes CO, CO₂, and H₂O slightly harder to detect, while these are the only models where CH₄ is detectable. Though the volume mixing ratio of H₂O is actually slightly higher in the cooler case than that of the nominal insolation, this increase in difficulty may be attributed to the decrease in scale height with decreasing temperature which makes spectral features smaller. However, interference from the many CH₄ lines may also make H₂O harder to detect. To test this idea, we took a 0.3 \times insolation model in L band (which has the strongest CH₄ features) without CH₄ but including all of our other opacity sources as the "truth" model and computed the required S/N_{res} to detect H₂O in this case. When CH₄ lines are not included, the required S/N_{res} to detect H₂O decreases for all tested spectral resolutions, indicating that interference from CH₄ is indeed a source of the increased difficulty in detecting H₂O for these cooler models.

- J and H Bands: The required S/N_{res} for detecting H₂O is slightly lower in J band and higher in H band when compared to the 1 \times results. For example, with $R \sim 50,000$ spectra, detecting H₂O would require $S/N_{res} \geq 2600$ in J band and $S/N_{res} \geq 4500$ in H band, compared to 2800 and 4000 for the 1 \times insolation case, respectively.
- K Band: Detecting CO requires a much higher S/N_{res} due to its decreased abundance, and is not detectable at $R \leq 25,000$ for any of our noise scenarios. CH₄, which is not detectable in the nominal 1 \times models, is more readily detectable than CO and H₂O in this cooler scenario (i.e. with $S/N_{res} \geq 2000$ at $R \sim 75,000$, compared to 3500 for CO and 4000 for H₂O). For the highest resolution case ($R \sim 100,000$), one could detect CO, H₂O, and CH₄ with a $S/N_{res} \geq 3000$.
- L Band: CH₄ and H₂O are both detectable. Due to the proximity to the v3 band, detecting CH₄ requires a much lower S/N_{res} than in K , with a lower limit of 300 in the $R \sim 100,000$ case. A slightly higher S/N_{res} is required to detect H₂O than in the 1 \times insolation case; $S/N_{res} \geq 3500$ rather than 3000 for $R \sim 75,000$.
- M Band: In all cases the required S/N_{res} to detect CO, CO₂, and H₂O is higher than for the nominal

models. For $R \sim 100,000$, CO, CO₂, and H₂O require $S/N_{res} \geq 400, 1300,$ and 1300 , respectively (compared to 200, 1300, and 1200 for the $1 \times$ insolation case).

3.2.2. $3 \times$ Insolation

The hotter $3 \times$ insolation models show qualitatively similar detection behaviour to the $1 \times$ case. In most cases the haze opacity requires a slightly lower S/N_{res} to detect than the nominal insolation case, but the change is often small. The biggest difference is that CO, CO₂ and H₂O are all easier to detect, sometimes by up to a factor of 2 decrease in the required S/N_{res} .

- *J* and *H* Bands: The required S/N_{res} for detecting H₂O is approximately $2 \times$ lower than in the $1 \times$ insolation case for all resolutions in both bands, likely due to the increase in the size of the features in these bands compared to the nominal insolation case as seen in Figure 4.
- *K* Band: Both CO and H₂O are detectable, but again require lower S/N_{res} for a 5σ detection. This decrease in required S/N_{res} is $\sim 40\%$ for CO, while closer to $\sim 50\%$ for H₂O. Notably, in this hotter scenario, H₂O is detected at $R \leq 50,000$ for all explored S/N cases, in contrast to the nominal $1 \times$ insolation case where it is only detectable at higher resolutions.
- *L* Band: The required S/N_{res} for H₂O detection is again roughly 50% lower than needed in the $1 \times$ insolation case across spectral resolutions.
- *M* Band: The required S/N_{res} to detect CO, CO₂, and H₂O is lower than in the nominal insolation case, with the largest effect for H₂O. With spectra at $R \sim 100,000$, one could detect CO, CO₂, and H₂O simultaneously with $S/N_{res} \geq 1100$ (compared to 1300 in the $1 \times$ insolation case).

3.3. Observing Multiple Bands

Modern instruments are now able to observe multiple atmospheric windows simultaneously, often some selection of *J*, *H*, and *K* bands, including CARMENES (Quirrenbach et al. 2016), NIRPS (Wildi et al. 2017), IGRINS (Park et al. 2014), GIANO (Origlia et al. 2014), and SPIRou (Artigau et al. 2014). Access to multiple bands in one exposure allows for more spectral lines of a molecule to be observed, strengthening the signal of that molecule, as the molecular detection S/N scales as the $\sqrt{N_{lines}}$. Over these wavelength ranges, H₂O in particular has millions of spectral lines in each band. Figure 11 demonstrates the effect of increasing wavelength

range on H₂O detection significance. Figure 11b shows a clear improvement when combining *J* and *H* bands over just observing *J* band alone (Figure 11a). Furthermore, H₂O is more difficult to observe in *K* band, as shown in Figure 11c and expected since H₂O has larger cross sections in *J* and *H* bands as shown in Figure 2. However, H₂O is much more detectable at fixed spectral resolution or S/N_{res} when *J* and *H* band are also observed (Figure 11d). Thus, instruments that can observe *J-K* band simultaneously will more readily detect H₂O and CO (which is detectable in *K* band but neither of the other two bands; see Table 3 with the same observations).

Looking to future instrumentation, GMTNIRS (Jaffe et al. 2016) is a proposed high-resolution spectrograph for the GMT that would cover 1 to 5 μm in one exposure. Table 7 shows the minimum S/N_{res} required to detect each opacity source when observing all bands (*J-M*) simultaneously. While the required S/N_{res} to detect H₂O is substantially less than when looking at any single band as shown in Table 5, other opacity sources are not as benefited by simultaneous wavelength coverage, due to their narrower span or dominance by H₂O over most bands. In particular, the minimum S/N_{res} required to detect CO and CO₂ is in most cases identical to that required for when considering solely *M* band spectra, so adding in other bands does not make either molecule easier to detect. Similarly, the minimum S/N_{res} reported for CH₄ in Table 7 are identical to those for analyzing solely *L* band spectra from Table 6. Thus, H₂O is the opacity source, under these specific atmospheric conditions, most benefited by an instrument with wide instantaneous wavelength coverage. In scenarios in which other broad-band absorbing molecules dominate, say, higher metallicity where CO/CO₂ are more prominent, or much cooler where CH₄ dominates the full near infrared, the influence of multiple bands on specific molecular detections would undoubtedly change. We leave this detailed analysis to a future study.

3.4. Photochemical Products

Though our results so far have only included four molecules (CO, CO₂, H₂O, and CH₄), other molecules could be important opacity sources for these planets in the near infrared. In particular, photochemical products could affect the high-resolution transmission spectra at these wavelengths and their detection could provide an avenue to distinguish between a photochemical haze and equilibrium condensate clouds. However, not all potential molecules have high fidelity line lists at the temperatures and pressures necessary for generating these models. Hawker et al. (2018) and Cabot et al. (2019) have

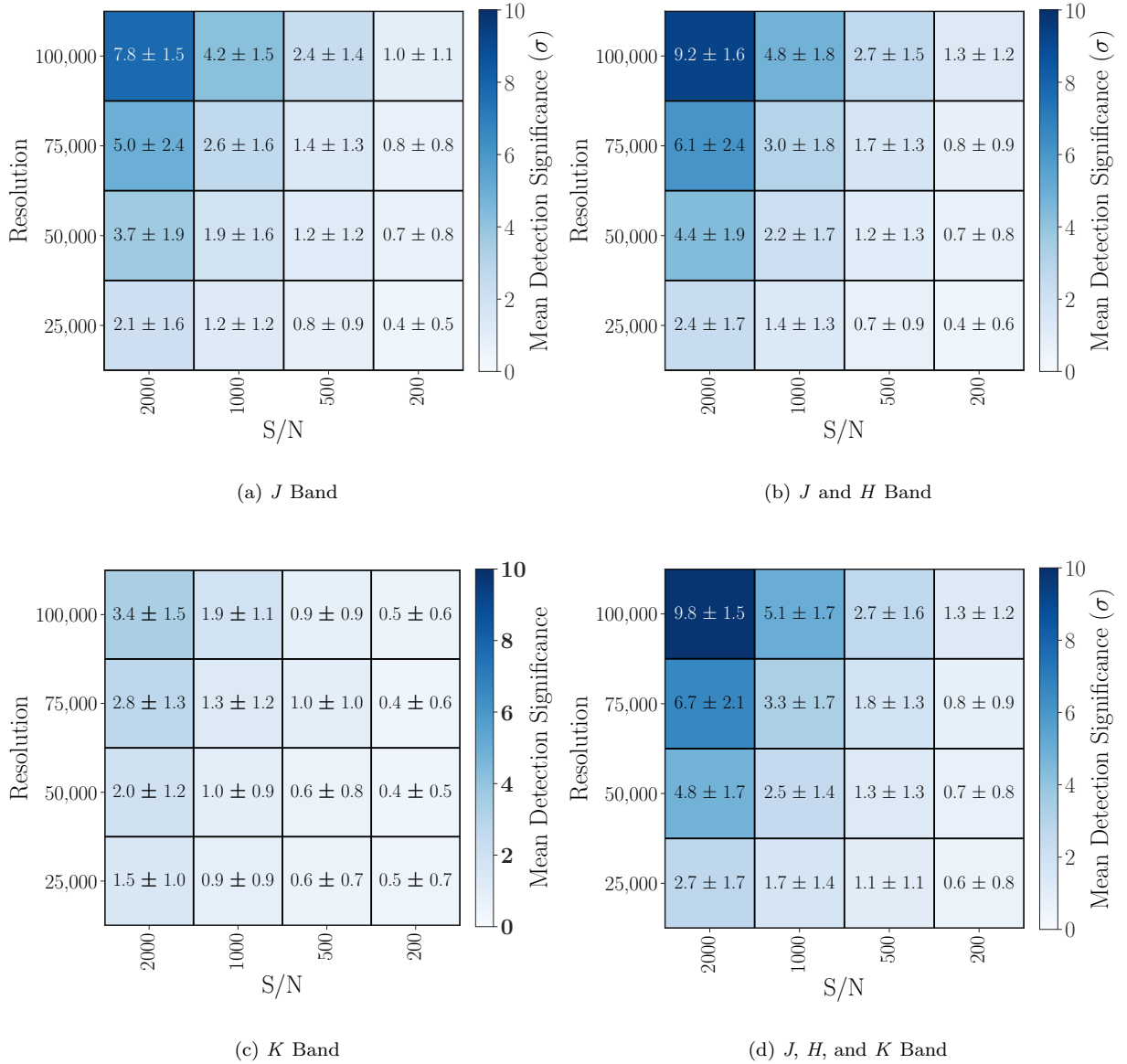


Figure 11. Detection significances for H_2O for different observing bands as a function of spectral resolution and S/N per resolution element of the transmission spectrum. Observing *J* and *H* bands simultaneously increased the detection significance of H_2O as shown in 11a and 11b. Similarly, while the H_2O detection in *K* band is marginal in 11c, adding *J* and *H* bands allows for a much stronger detection of H_2O .

both presented evidence of HCN in the atmospheres of hot Jupiters using high-resolution spectroscopy, indicating the possibility of detecting this molecule with current line lists. Thus, to test whether photochemical products would be detectable with these kinds of observations, we focus on HCN as an illustrative example using the line list from Harris et al. (2008).

As in Section 2.1, we create a new “truth” spectrum that includes HCN in addition to our other opacity sources. We use the results of the photochemical model cited in Section 2.1 to determine the HCN abundance.

We can then use the $\log(L)$ method described in Section 2 to compare to the models without HCN and quantify our ability to detect the molecule as a function of S/N_{res} and spectral resolution. We find that *L* band is the only observing band where HCN is detectable; so far the only spectra used to detect HCN in Hawker et al. (2018) and Cabot et al. (2019) covered 3.18 - 3.27 μm . The minimum S/N_{res} needed to detect HCN is listed in Table 8. HCN is only detectable for the hottest models and highest spectral resolutions. However, as shown in Figure 2, HCN has its strongest features between 3 and

simultaneously.

Table 7. Minimum S/N_{res} required for $\geq 5\sigma$ detection of each opacity source for transmission spectra that cover observing bands J - M

Opacity Source	Insolation	Spectral Resolution			
		25,000	50,000	75,000	100,000
(1)	(2)	(3)	(4)	(5)	(6)
CO	0.3×	1400	800	600	400
	1×	700	350	250	200
	3×	400	250	150	150
CO ₂	0.3×	4500	2400	1800	1200
	1×	4500	2400	1700	1300
	3×	4500	2200	1500	1100
H ₂ O	0.3×	2800	1500	1100	800
	1×	2600	1400	1100	800
	3×	1200	700	500	350
CH ₄	0.3×	1100	700	450	300
	1×	-	-	-	-
	3×	-	-	-	-
Haze	0.3×	150	100	100	100
	1×	150	100	100	50
	3×	100	100	50	50

3.2 μm . If spectra starting at 3.1 μm instead of 3.15 μm can be obtained, HCN could be much easier to detect - for example, the required S/N_{res} for 3× insolation and $R \sim 100,000$ decreases from 3500 to 1000.

We also explore how unaccounted for photochemical products could affect the detectability of the major molecular species. Our “truth” spectrum includes our standard set of opacities (CO, CO₂, H₂O, CH₄, and H₂/He CIA) plus HCN, but compare to models without HCN (e.g., an “incorrect” model). We find that the detection S/N values do not change in any case, even for L band where HCN is detectable (Table 9). This consistency suggests that our above results are robust against missing absorbers. However, since HCN is not easily detectable, the minimum S/N_{res} for a detection may be slightly higher than reported here if there are unaccounted-for molecules that would more significantly affect the high-resolution transmission spectrum.

Table 8. Minimum S/N_{res} required for $\geq 5\sigma$ detection of HCN with L band spectra.

Insolation	Spectral Resolution			
	25,000	50,000	75,000	100,000
(1)	(2)	(3)	(4)	(5)
0.3×	-	-	-	-
1×	-	-	-	-
3×	-	-	4000	3500

Table 9. Minimum S/N_{res} required for $\geq 5\sigma$ detection of each opacity source for transmission spectra in L band when HCN is included in the observed spectrum but not the comparison models. CO and CO₂ are not observable with $S/N_{res} \leq 5000$.

Opacity Source	Insolation	Spectral Resolution			
		25,000	50,000	75,000	100,000
(1)	(2)	(3)	(4)	(5)	(6)
H ₂ O	0.3×	-	4500	3500	2600
	1×	-	4000	3000	2600
	3×	3000	1500	1300	1100
CH ₄	0.3×	1100	700	450	300
	1×	-	-	-	-
	3×	-	-	-	-
Haze	0.3×	300	200	150	100
	1×	400	250	200	150
	3×	300	200	100	100

4. DISCUSSION

4.1. Observability with Current and Future Instrumentation

In Section 3 we presented the required S/N per resolution element to detect molecules over a range of stellar insolation levels ($T_{\text{eff}} = 412, 557, \text{ and } 733 \text{ K}$) and observational parameters ($R \sim 25,000 - 100,000$, $50 \leq S/N_{res} \leq 5000$, and $J - M$ bands). However, the exposure times required to reach these S/N_{res} will vary depending on host star brightness, sky background, telescope aperture

size, and instrument sensitivity. While a detailed instrument/observational investigation is outside the scope of this work, we can focus on GJ 1214b as an example of how these factors might affect observations.

We used a simplified noise model for estimating instrumental S/N_{res} as a function of exposure time to determine the observability of molecules as a function of source brightness, telescope size, resolving power, and wavelength. The estimated measured S/N per resolution element depends upon the number of photons received from the source (which itself depends on exposure time, throughput, collecting area, brightness, and resolution), thermal background (which depends on emissivity, throughput, etendue $A\Omega$, and temperature), and instrumental quantities such as total throughput ($\tau \sim 0.05$), emissivity ($\epsilon = 0.3$), and noise properties of the detector (dark current, read noise), which all combine into:

$$S/N = \frac{S \times T}{\sqrt{S \times T + N_{exp} \times (T_{exp} \times (BKGD + DC) + RN^2)}} \quad (4)$$

where S is the signal, T is the total exposure time, N_{exp} is the number of exposures, T_{exp} is the time for a single exposure (assumed to be 600 seconds), $BKGD$ is thermal background, DC is the dark current, and RN is read noise.

Using this equation, it is then possible to estimate the S/N_{res} expected for a given source magnitude, resolving power, and exposure time. For each molecule, we determined the minimum exposure time required to detect the molecule in the $1\times$ insolation scenario (using GJ 1214b-like planet/star parameters) for varying source brightnesses. We simulated the exposure time vs. source brightness relation for three different scenarios: 1) $R \sim 25,000$ instrument on a 10-m telescope, seeing-limited (e.g. NIRSPEC on Keck, McLean et al. 1998, Martin et al. 2018); 2) $R \sim 50,000$ instrument on an 8-m telescope, seeing-limited (e.g. IGRINS on Gemini, Park et al. 2014, Mace et al. 2018); 3) $R \sim 100,000$ instrument on a 30-m telescope, behind Adaptive Optics (AO), such as the proposed METIS on ELT (Brandl et al. 2018) or MODHIS (Mawet et al. 2019) on TMT. The dark current and read noise for the IGRINS-like and ELT instruments were assumed to be typical Teledyne values⁴; the NIRSPEC-like curve used measured values from that instrument⁵. To easily normalize for

each type of instrument, we assumed a “seeing disk” of 20 pixels per resolution element (assuming the resolution element in the imaging and dispersion directions was 5 pixels by 4 pixels). We assume exposures with the maximum practical length one would want to take in a particular band due to sky lines and thermal background: 600 seconds for K band and 30 seconds for M band. We do not assume any overhead time that might occur, for example due to detector readout, which would increase the required amount of telescope time for these observations. Each resolving power requires a different minimum S/N_{res} to detect a molecule. For example, in K band, to detect CO at 5σ significance for $1\times$ insolation, an $R\sim 25,000$ instrument requires $S/N_{res} > 3500$, an $R\sim 50,000$ instrument needs $S/N_{res} > 1600$, and an $R\sim 100,000$ instrument needs $S/N_{res} > 800$.

Figure 12 shows the results of our noise observability analysis. For K band, shown in Figure 12a, an instrument similar to NIRSPEC would take about 8 hours while the IGRINS-like instrument would take about 4 hours to detect CO. The mean transit duration of GJ 1214b is $52.73^{+0.49}_{-0.35}$ minutes, which translates to ~ 8 transits with NIRSPEC and ~ 4 transits with IGRINS. An ELT instrument behind AO with $R\sim 100,000$ would reduce this required time to around 6 minutes, so observing one full transit will be more than enough time. In contrast, for M band observations, both current instruments would take greater than 100 hours to detect CO due to the much higher background noise in this wavelength range. An ELT instrument behind AO would take about 8 minutes to detect CO in this wavelength range as shown in Figure 12b, so again one full transit would be more than enough time. However, we note that these observing times assume good observing conditions and could readily double in the case of low seeing or slit losses.

We performed similar observability calculations for H_2O and CO_2 . For H_2O , we looked at H and M bands. For H band spectra, we did not detect H_2O with $S/N_{res} \leq 5000$ for $R\sim 25,000$ (see Table 5) so Figure 13a only shows the exposure times for the IGRINS-like and TMT instruments. H band spectra with the IGRINS-like instrument would about 12 hours of total integration for H_2O to be detected. Detecting H_2O in M band spectra, however, would take an inordinate amount of observing time (> 100 hours) with current instruments but again should be more easily detectable with about 7 hours on a 30-m telescope with AO for GJ 1214b. Similarly, only the $R\sim 100,000$ instrument on a 30-m telescope can detect CO_2 within any reasonable amount of time, ~ 8 hours for GJ 1214b as seen in Figure 14.

⁴ <http://www.teledyne-si.com/products/Documents/H2RG%20Brochure%20-%20September%202017.pdf>

⁵ <https://www2.keck.hawaii.edu/inst/nirspec/Specifications.html>

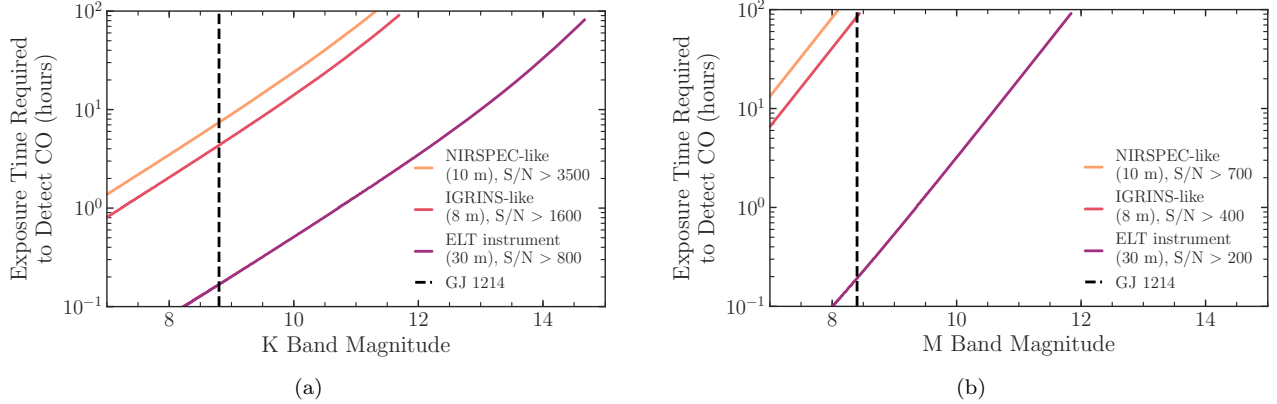


Figure 12. Estimated exposure time required to detect CO as a function of host star magnitude for different instruments in K band (Figure 12a) and M band (Figure 12b). The magnitude of GJ 1214 is marked by the dashed line. We consider three different types of instruments- a NIRSPEC-like $R \sim 25,000$ spectrograph on a 10-m telescope, an IGRINS-like $R \sim 50,000$ spectrograph on an 8-m telescope, and a proposed $R \sim 100,000$ spectrograph behind AO on a 30-m telescope like the TMT. We find that although detecting CO requires a lower S/N_{res} in M band than K band, K -band observations are actually much more feasible, particularly for current instruments. However, a potential high-resolution spectrograph on a 30-m telescope would be able to detect CO in GJ 1214b in M band with one transit.

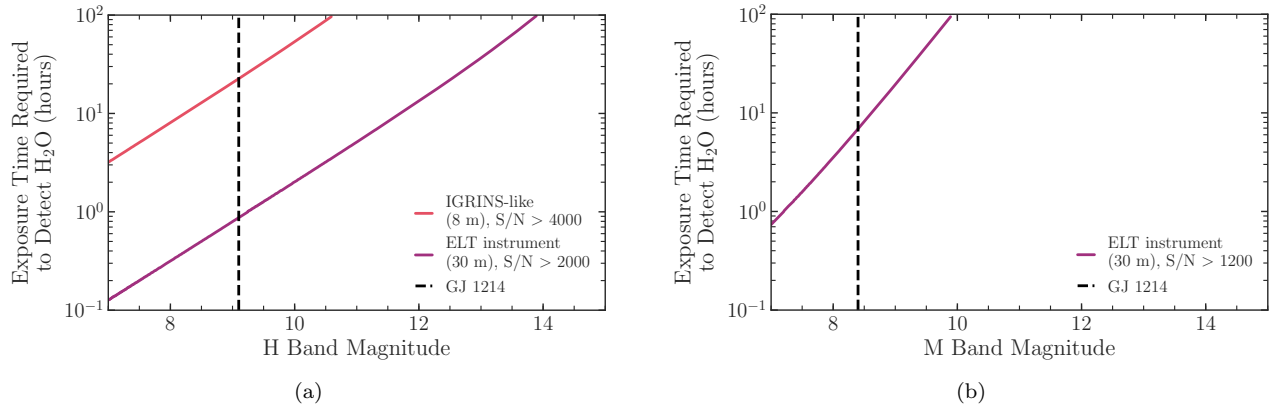


Figure 13. Estimated exposure time required to detect H_2O as a function of host star magnitude for different instruments in H band (Figure 13a) and M band (Figure 13b). The magnitude of GJ 1214 is marked by the dashed line. In H band we did not have any H_2O detections with $S/N_{res} \leq 5000$ with $R \sim 25,000$, so in Figure 13a we just consider an IGRINS-like $R \sim 50,000$ spectrograph on an 8-m telescope and a proposed $R \sim 100,000$ spectrograph behind AO on a 30-m telescope like the TMT. Similar to Figure 12, we find that although detecting H_2O requires a lower S/N_{res} in M band than H band, H -band observations are more feasible, particularly for current instruments. However, a potential high-resolution spectrograph on a 30-m telescope would be able to detect H_2O in GJ 1214b in H band with ~ 7 hours of integration time.

4.2. Additional Caveats

There are other caveats to consider when using the S/N_{res} values in Tables 2-9 to plan observations for GJ 1214b or any other hazy sub-Neptune. First, when computing these values, we assumed all spectra covered the full wavelength range stated in Table 1, regardless of resolution, as the exact wavelength coverage varies between instruments. However, increasing spectral resolution may come with decreased wavelength coverage, leading to fewer spectral lines in the data, which would make these molecules more difficult to detect.

In addition, GJ 1214b has a relatively slow radial velocity change over the course of one transit (12 km s^{-1} ; Crossfield et al. 2011), compared to the hot Jupiters for which this technique has successfully been applied. Thus, it may be difficult to remove the quasi-stationary contamination from our atmosphere and the star while preserving the planet spectrum.

Finally, we are assuming no contamination from stellar lines in our spectra (see Section 2.1). This is a good approximation for host stars with minimal spectral features in the observed wavelengths. However, for planets around M-dwarfs like GJ 1214, a myriad of stellar lines

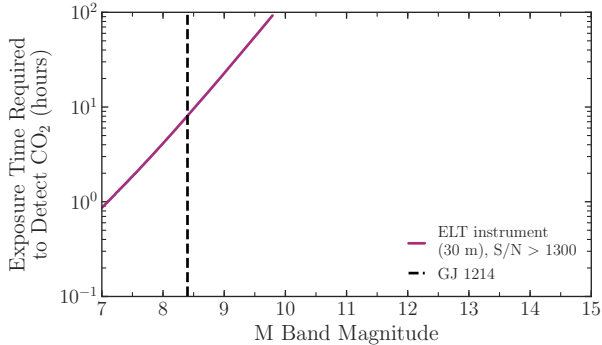


Figure 14. Estimated exposure time required to detect CO₂ as a function of host star magnitude for different instruments in *M* band. The magnitude of GJ 1214 is marked by the dashed line. Of our three investigated instrument types, only a proposed $R \sim 100,000$ spectrograph behind AO on a 30-m telescope like the TMT can detect CO₂ in a reasonable amount of observing time (~ 8 hours for GJ 1214b).

may make proper removal of the telluric and stellar contamination challenging. Brogi et al. (2016) and Schwarz et al. (2016) successfully modelled and removed stellar lines before removing the tellurics, but it remains to be seen if this technique can be successfully applied to M-dwarf spectra as more complex stellar models and data analysis are needed to remove overlapping stellar-planet molecular features (e.g., Chiavassa & Brogi 2019). Stellar variability may also pose a challenge, particularly for combining data from multiple nights, especially for M-dwarfs which are known to have high levels of magnetic activity (e.g., Newton et al. 2016). As a result, a conservative approach may be to focus on detection and characterization of species unlikely to be abundantly present in M-dwarf photospheres (e.g., CH₄, NH₃, HCN, etc.).

5. CONCLUSIONS

In this work, we have investigated the feasibility of detecting molecules in the atmospheres of hazy sub-Neptunes with ground-based, high-resolution spectroscopy. To do so, we generated high-resolution transmission spectra of GJ 1214b analogs with a photochemical haze that matches the featureless low resolution transmission spectrum. We considered two different metrics from the literature to quantify our detection significances: the cross correlation function (CCF) and a log likelihood function ($\log(L)$) derived by Brogi & Line (2019). While both metrics produced similar detection significances for molecules, only the $\log(L)$ was sensitive to the presence of the hazy opacity due to the additional terms that track the spectral variance relative to the data variance. Thus, we used the $\log(L)$ for the remainder of this work. However, our method relies on measuring the change in $\log(L)$ as we remove one opac-

ity source at a time, indicating more care may need to be taken when determining which opacity sources to include in one’s model.

We have calculated the minimum signal-to-noise (S/N) required for a $> 5\sigma$ detection of each opacity source (CO, CO₂, H₂O, CH₄, and the haze) as a function of stellar insolation, spectral resolution, and wavelength range. Our key results are as follows.

1. High resolution infrared transmission spectrum observations for hazy GJ 1214b analogs probe pressures around $1 \mu\text{bar}$ and numerous molecular features can be detected for spectra that otherwise appear “featureless” at $R \sim 100\text{-}1000$.
2. The haze is always the easiest opacity source to detect and observable in all observing bands. Thus, achieving the S/N per resolution element required to detect a molecule will also allow one to rule out a completely clear atmosphere.
3. H₂O is detectable with $S/N_{res} \leq 5000$ for almost all combinations of spectral resolution and wavelength coverage, and is the only molecule observable in *J* and *H* bands, but always requires a higher S/N_{res} than any other molecules observable in that band. In contrast, CO is only observable in *K* and *M* bands, but is the easiest molecule to detect. CO₂ is only observable in *M* band; in fact, *M* band is the observing band that requires the lowest S/N_{res} to detect two or more molecules at once (although in practice requires long observing time due to thermal background noise).
4. In general, increasing the stellar insolation of the model lowers the required S/N_{res} for molecular detections. However, CH₄ is only detectable in the *L* band spectra of the coldest models.
5. HCN is detectable with *L* band spectra for high resolution spectra of the hottest models. Detecting HCN, along with other potential photochemical products, could be a way to distinguish between equilibrium cloud and photochemical haze opacity in a planet’s atmosphere.

To further investigate the observability of these ultra-hazy sub-Neptunes, we used a simple model in Section 4 to determine the feasibility of observing high-resolution transmission spectroscopy of planets with current and future instruments. We found that detecting CO and H₂O for GJ 1214b with current instruments observing in *K* and *H* bands, respectively, requires on the order of 10 hours of observing time. Furthermore, though a lower S/N_{res} is required to detect these molecules in

M band, current instruments would need an unreasonable investment of observing time due to the high background. However, such an observation would be trivial behind AO on an ELT class telescope. In addition, a high resolution spectrograph on an actively cooled space telescope could observe in M band without issues from thermal background.

As discussed in Section 4, more detailed, instrument-specific simulations for particular targets may be needed for careful observation planning. Such a study for ELT instrument concepts could help inform design decisions to maximize the information we can learn about these objects. Furthermore, an analogous study could be conducted of thermal emission spectra, which will likely require the S/N achievable with ELT instruments. A number of sub-Neptunes will be observable with *JWST*, albeit with lower spectral resolution. Analysis of simulated joint *JWST* and ground-based high-resolution observations could pinpoint optimal observing strategies to take full advantage of these complementary data

sets. As more sub-Neptunes around bright, nearby stars are discovered (as expected from NASA's TESS mission; Barclay et al. 2018), considering the best application of observational techniques is essential to maximizing the scientific return on this abundant class of planets.

ACKNOWLEDGMENTS

We are grateful to Dr. Gregory Mace for his assistance and advice on the instrumental observability section. Posting of this manuscript on the arXiv was coordinated with S. Ghandi et al. ECM is supported by an NSF Astronomy and Astrophysics Postdoctoral Fellowship under award AST-1801978. MRL and JJF acknowledge the support of NASA XRP grant 80NSSC19K0293. JLB acknowledges funding from the European Research Council (ERC) under the European Unions Horizon 2020 research and innovation program under grant agreement No 805445.

REFERENCES

- Alonso-Floriano, F. J., Sánchez-López, A., Snellen, I. A. G., et al. 2019, *A&A*, 621, A74, doi: [10.1051/0004-6361/201834339](https://doi.org/10.1051/0004-6361/201834339)
- Artigau, É., Kouach, D., Donati, J.-F., et al. 2014, *Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series*, Vol. 9147, SPIROu: the near-infrared spectropolarimeter/high-precision velocimeter for the Canada-France-Hawaii telescope, 914715, doi: [10.1117/12.2055663](https://doi.org/10.1117/12.2055663)
- Barber, R. J., Tennyson, J., Harris, G. J., & Tolchenov, R. N. 2006, *MNRAS*, 368, 1087, doi: [10.1111/j.1365-2966.2006.10184.x](https://doi.org/10.1111/j.1365-2966.2006.10184.x)
- Barclay, T., Pepper, J., & Quintana, E. V. 2018, *ApJS*, 239, 2, doi: [10.3847/1538-4365/aae3e9](https://doi.org/10.3847/1538-4365/aae3e9)
- Batalha, N. M., Rowe, J. F., Bryson, S. T., et al. 2013, *ApJS*, 204, 24, doi: [10.1088/0067-0049/204/2/24](https://doi.org/10.1088/0067-0049/204/2/24)
- Bean, J. L., Miller-Ricci Kempton, E., & Homeier, D. 2010, *Nature*, 468, 669, doi: [10.1038/nature09596](https://doi.org/10.1038/nature09596)
- Bean, J. L., Désert, J.-M., Kabath, P., et al. 2011, *ApJ*, 743, 92, doi: [10.1088/0004-637X/743/1/92](https://doi.org/10.1088/0004-637X/743/1/92)
- Birkby, J. L. 2018, arXiv e-prints, arXiv:1806.04617, <https://arxiv.org/abs/1806.04617>
- Birkby, J. L., de Kok, R. J., Brogi, M., et al. 2013, *MNRAS*, 436, L35, doi: [10.1093/mnrasl/slt107](https://doi.org/10.1093/mnrasl/slt107)
- Birkby, J. L., de Kok, R. J., Brogi, M., Schwarz, H., & Snellen, I. A. G. 2017, *AJ*, 153, 138, doi: [10.3847/1538-3881/aa5c87](https://doi.org/10.3847/1538-3881/aa5c87)
- Booth, R. A., Clarke, C. J., Madhusudhan, N., & Ilee, J. D. 2017, *MNRAS*, 469, 3994, doi: [10.1093/mnras/stx1103](https://doi.org/10.1093/mnras/stx1103)
- Borucki, W. J., Koch, D. G., Basri, G., et al. 2011, *ApJ*, 736, 19, doi: [10.1088/0004-637X/736/1/19](https://doi.org/10.1088/0004-637X/736/1/19)
- Brandl, B. R., Absil, O., Agócs, T., et al. 2018, in *Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series*, Vol. 10702, Proc. SPIE, 107021U, doi: [10.1117/12.2311492](https://doi.org/10.1117/12.2311492)
- Brogi, M., de Kok, R. J., Albrecht, S., et al. 2016, *ApJ*, 817, 106, doi: [10.3847/0004-637X/817/2/106](https://doi.org/10.3847/0004-637X/817/2/106)
- Brogi, M., de Kok, R. J., Birkby, J. L., Schwarz, H., & Snellen, I. A. G. 2014, *A&A*, 565, A124, doi: [10.1051/0004-6361/201423537](https://doi.org/10.1051/0004-6361/201423537)
- Brogi, M., Giacobbe, P., Guilluy, G., et al. 2018, *A&A*, 615, A16, doi: [10.1051/0004-6361/201732189](https://doi.org/10.1051/0004-6361/201732189)
- Brogi, M., & Line, M. R. 2019, *AJ*, 157, 114, doi: [10.3847/1538-3881/aaffd3](https://doi.org/10.3847/1538-3881/aaffd3)
- Brown, T. M. 2001, *ApJ*, 553, 1006, doi: [10.1086/320950](https://doi.org/10.1086/320950)
- Burke, C. J., Bryson, S. T., Mullally, F., et al. 2014, *ApJS*, 210, 19, doi: [10.1088/0067-0049/210/2/19](https://doi.org/10.1088/0067-0049/210/2/19)
- Burke, C. J., Christiansen, J. L., Mullally, F., et al. 2015, *ApJ*, 809, 8, doi: [10.1088/0004-637X/809/1/8](https://doi.org/10.1088/0004-637X/809/1/8)
- Cabot, S. H. C., Madhusudhan, N., Hawker, G. A., & Gandhi, S. 2019, *MNRAS*, 482, 4422, doi: [10.1093/mnras/sty2994](https://doi.org/10.1093/mnras/sty2994)
- Charbonneau, D., Berta, Z. K., Irwin, J., et al. 2009, *Nature*, 462, 891, doi: [10.1038/nature08679](https://doi.org/10.1038/nature08679)

- Chiavassa, A., & Brogi, M. 2019, *A&A*, 631, A100, doi: [10.1051/0004-6361/201936566](https://doi.org/10.1051/0004-6361/201936566)
- Crossfield, I. J. M., Barman, T., & Hansen, B. M. S. 2011, *ApJ*, 736, 132, doi: [10.1088/0004-637X/736/2/132](https://doi.org/10.1088/0004-637X/736/2/132)
- Crossfield, I. J. M., Barman, T., Hansen, B. M. S., & Howard, A. W. 2013, *A&A*, 559, A33, doi: [10.1051/0004-6361/201322278](https://doi.org/10.1051/0004-6361/201322278)
- Crossfield, I. J. M., & Kreidberg, L. 2017, *AJ*, 154, 261, doi: [10.3847/1538-3881/aa9279](https://doi.org/10.3847/1538-3881/aa9279)
- de Kok, R. J., Birkby, J., Brogi, M., et al. 2014, *A&A*, 561, A150, doi: [10.1051/0004-6361/201322947](https://doi.org/10.1051/0004-6361/201322947)
- Deibert, E. K., de Mooij, E. J. W., Jayawardhana, R., et al. 2019, *AJ*, 157, 58, doi: [10.3847/1538-3881/aaf56b](https://doi.org/10.3847/1538-3881/aaf56b)
- Deming, D., Wiedemann, G., & Bjoraker, G. 2000, *Astronomical Society of the Pacific Conference Series*, Vol. 212, *Prospects for Direct Spectral Detection and Characterization of "Hot Jupiters"*, ed. C. A. Griffith & M. S. Marley, 308
- Deming, D., Wilkins, A., McCullough, P., et al. 2013, *ApJ*, 774, 95, doi: [10.1088/0004-637X/774/2/95](https://doi.org/10.1088/0004-637X/774/2/95)
- Dressing, C. D., & Charbonneau, D. 2015, *ApJ*, 807, 45, doi: [10.1088/0004-637X/807/1/45](https://doi.org/10.1088/0004-637X/807/1/45)
- Espinoza, N., Fortney, J. J., Miguel, Y., Thorngren, D., & Murray-Clay, R. 2017, *ApJL*, 838, L9, doi: [10.3847/2041-8213/aa65ca](https://doi.org/10.3847/2041-8213/aa65ca)
- Esteves, L. J., de Mooij, E. J. W., Jayawardhana, R., Watson, C., & de Kok, R. 2017, *AJ*, 153, 268, doi: [10.3847/1538-3881/aa7133](https://doi.org/10.3847/1538-3881/aa7133)
- Figueira, P., Pont, F., Mordasini, C., et al. 2009, *A&A*, 493, 671, doi: [10.1051/0004-6361/20078951](https://doi.org/10.1051/0004-6361/20078951)
- Flagg, L., Johns-Krull, C. M., Nofi, L., et al. 2019, *ApJL*, 878, L37, doi: [10.3847/2041-8213/ab276d](https://doi.org/10.3847/2041-8213/ab276d)
- Fortney, J. J., Mordasini, C., Nettelmann, N., et al. 2013, *ApJ*, 775, 80, doi: [10.1088/0004-637X/775/1/80](https://doi.org/10.1088/0004-637X/775/1/80)
- Fraine, J., Deming, D., Benneke, B., et al. 2014, *Nature*, 513, 526, doi: [10.1038/nature13785](https://doi.org/10.1038/nature13785)
- Freedman, R. S., Lustig-Yaeger, J., Fortney, J. J., et al. 2014, *ApJS*, 214, 25, doi: [10.1088/0067-0049/214/2/25](https://doi.org/10.1088/0067-0049/214/2/25)
- Fressin, F., Torres, G., Charbonneau, D., et al. 2013, *ApJ*, 766, 81, doi: [10.1088/0004-637X/766/2/81](https://doi.org/10.1088/0004-637X/766/2/81)
- Fu, G., Deming, D., Knutson, H., et al. 2017, *ApJL*, 847, L22, doi: [10.3847/2041-8213/aa8e40](https://doi.org/10.3847/2041-8213/aa8e40)
- Gao, P., & Benneke, B. 2018, *ApJ*, 863, 165, doi: [10.3847/1538-4357/aad461](https://doi.org/10.3847/1538-4357/aad461)
- Gibson, N. P., Merritt, S., Nugroho, S. K., et al. 2020, *MNRAS*, 493, 2215, doi: [10.1093/mnras/staa228](https://doi.org/10.1093/mnras/staa228)
- Greene, T. P., Line, M. R., Montero, C., et al. 2016, *ApJ*, 817, 17, doi: [10.3847/0004-637X/817/1/17](https://doi.org/10.3847/0004-637X/817/1/17)
- Guilluy, G., Sozzetti, A., Brogi, M., et al. 2019, *A&A*, 625, A107, doi: [10.1051/0004-6361/201834615](https://doi.org/10.1051/0004-6361/201834615)
- Harris, G. J., Larner, F. C., Tennyson, J., et al. 2008, *MNRAS*, 390, 143, doi: [10.1111/j.1365-2966.2008.13642.x](https://doi.org/10.1111/j.1365-2966.2008.13642.x)
- Hawker, G. A., Madhusudhan, N., Cabot, S. H. C., & Gandhi, S. 2018, *ApJL*, 863, L11, doi: [10.3847/2041-8213/aac49d](https://doi.org/10.3847/2041-8213/aac49d)
- Heng, K., & Demory, B.-O. 2013, *ApJ*, 777, 100, doi: [10.1088/0004-637X/777/2/100](https://doi.org/10.1088/0004-637X/777/2/100)
- Huang, X., Freedman, R. S., Tashkun, S. A., Schwenke, D. W., & Lee, T. J. 2013, *JQSRT*, 130, 134, doi: [10.1016/j.jqsrt.2013.05.018](https://doi.org/10.1016/j.jqsrt.2013.05.018)
- Huang, X., Gamache, R. R., Freedman, R. S., Schwenke, D. W., & Lee, T. J. 2014, *JQSRT*, 147, 134, doi: [10.1016/j.jqsrt.2014.05.015](https://doi.org/10.1016/j.jqsrt.2014.05.015)
- Iyer, A. R., Swain, M. R., Zellem, R. T., et al. 2016, *ApJ*, 823, 109, doi: [10.3847/0004-637X/823/2/109](https://doi.org/10.3847/0004-637X/823/2/109)
- Jaffe, D. T., Barnes, S., Brooks, C., et al. 2016, *Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series*, Vol. 9908, *GMTNIRS: progress toward the Giant Magellan Telescope near-infrared spectrograph*, 990821, doi: [10.1117/12.2232994](https://doi.org/10.1117/12.2232994)
- Jones, A., Noll, S., Kausch, W., Szyszka, C., & Kimeswenger, S. 2013, *A&A*, 560, A91, doi: [10.1051/0004-6361/201322433](https://doi.org/10.1051/0004-6361/201322433)
- Kempton, E. M. R., Perna, R., & Heng, K. 2014, *ApJ*, 795, 24, doi: [10.1088/0004-637X/795/1/24](https://doi.org/10.1088/0004-637X/795/1/24)
- Kempton, E. M. R., Bean, J. L., Louie, D. R., et al. 2018, *PASP*, 130, 114401, doi: [10.1088/1538-3873/aadf6f](https://doi.org/10.1088/1538-3873/aadf6f)
- Knutson, H. A., Benneke, B., Deming, D., & Homeier, D. 2014a, *Nature*, 505, 66, doi: [10.1038/nature12887](https://doi.org/10.1038/nature12887)
- Knutson, H. A., Dragomir, D., Kreidberg, L., et al. 2014b, *ApJ*, 794, 155, doi: [10.1088/0004-637X/794/2/155](https://doi.org/10.1088/0004-637X/794/2/155)
- Kreidberg, L., Line, M. R., Thorngren, D., Morley, C. V., & Stevenson, K. B. 2018, *ApJL*, 858, L6, doi: [10.3847/2041-8213/aabfce](https://doi.org/10.3847/2041-8213/aabfce)
- Kreidberg, L., Bean, J. L., Désert, J.-M., et al. 2014, *Nature*, 505, 69, doi: [10.1038/nature12888](https://doi.org/10.1038/nature12888)
- Lopez, E. D., & Fortney, J. J. 2014, *ApJ*, 792, 1, doi: [10.1088/0004-637X/792/1/1](https://doi.org/10.1088/0004-637X/792/1/1)
- Louie, D. R., Deming, D., Albert, L., et al. 2018, *PASP*, 130, 044401, doi: [10.1088/1538-3873/aaa87b](https://doi.org/10.1088/1538-3873/aaa87b)
- Mace, G., Sokal, K., Lee, J.-J., et al. 2018, in *Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series*, Vol. 10702, *Proc. SPIE*, 107020Q, doi: [10.1117/12.2312345](https://doi.org/10.1117/12.2312345)
- Mai, C., & Line, M. R. 2019, *ApJ*, 883, 144, doi: [10.3847/1538-4357/ab3e6d](https://doi.org/10.3847/1538-4357/ab3e6d)
- Martin, E. C., Fitzgerald, M. P., McLean, I. S., et al. 2018, in *Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series*, Vol. 10702, *Proc. SPIE*, 107020A, doi: [10.1117/12.2312266](https://doi.org/10.1117/12.2312266)

- Mawet, D., Fitzgerald, M., Konopacky, Q., et al. 2019, in *BAAS*, Vol. 51, 134. <https://arxiv.org/abs/1908.03623>
- McLean, I. S., Becklin, E. E., Bendiksen, O., et al. 1998, in *Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series*, Vol. 3354, Proc. SPIE, ed. A. M. Fowler, 566–578, doi: [10.1117/12.317283](https://doi.org/10.1117/12.317283)
- Miller-Ricci Kempton, E., Zahnle, K., & Fortney, J. J. 2012, *ApJ*, 745, 3, doi: [10.1088/0004-637X/745/1/3](https://doi.org/10.1088/0004-637X/745/1/3)
- Morley, C. V., Fortney, J. J., Kempton, E. M. R., et al. 2013, *ApJ*, 775, 33, doi: [10.1088/0004-637X/775/1/33](https://doi.org/10.1088/0004-637X/775/1/33)
- Morley, C. V., Fortney, J. J., Marley, M. S., et al. 2015, *ApJ*, 815, 110, doi: [10.1088/0004-637X/815/2/110](https://doi.org/10.1088/0004-637X/815/2/110)
- Morley, C. V., Kreidberg, L., Rustamkulov, Z., Robinson, T., & Fortney, J. J. 2017, *ApJ*, 850, 121, doi: [10.3847/1538-4357/aa927b](https://doi.org/10.3847/1538-4357/aa927b)
- Moses, J. I., Line, M. R., Visscher, C., et al. 2013, *ApJ*, 777, 34, doi: [10.1088/0004-637X/777/1/34](https://doi.org/10.1088/0004-637X/777/1/34)
- Mullally, F., Coughlin, J. L., Thompson, S. E., et al. 2015, *ApJS*, 217, 31, doi: [10.1088/0067-0049/217/2/31](https://doi.org/10.1088/0067-0049/217/2/31)
- Nettelmann, N., Fortney, J. J., Kramm, U., & Redmer, R. 2011, *ApJ*, 733, 2, doi: [10.1088/0004-637X/733/1/2](https://doi.org/10.1088/0004-637X/733/1/2)
- Newton, E. R., Irwin, J., Charbonneau, D., Berta-Thompson, Z. K., & Dittmann, J. A. 2016, *ApJL*, 821, L19, doi: [10.3847/2041-8205/821/1/L19](https://doi.org/10.3847/2041-8205/821/1/L19)
- Noll, S., Kausch, W., Barden, M., et al. 2012, *A&A*, 543, A92, doi: [10.1051/0004-6361/201219040](https://doi.org/10.1051/0004-6361/201219040)
- Öberg, K. I., Murray-Clay, R., & Bergin, E. A. 2011, *ApJL*, 743, L16, doi: [10.1088/2041-8205/743/1/L16](https://doi.org/10.1088/2041-8205/743/1/L16)
- Origlia, L., Oliva, E., Baffa, C., et al. 2014, *Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series*, Vol. 9147, High resolution near IR spectroscopy with GIANO-TNG, 91471E, doi: [10.1117/12.2054743](https://doi.org/10.1117/12.2054743)
- Park, C., Jaffe, D. T., Yuk, I.-S., et al. 2014, *Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series*, Vol. 9147, Design and early performance of IGRINS (Immersion Grating Infrared Spectrometer), 91471D, doi: [10.1117/12.2056431](https://doi.org/10.1117/12.2056431)
- Petigura, E. A., Howard, A. W., & Marcy, G. W. 2013, *Proceedings of the National Academy of Science*, 110, 19273, doi: [10.1073/pnas.1319909110](https://doi.org/10.1073/pnas.1319909110)
- Pino, L., Ehrenreich, D., Allart, R., et al. 2018, *A&A*, 619, A3, doi: [10.1051/0004-6361/201832986](https://doi.org/10.1051/0004-6361/201832986)
- Piskorz, D., Buzard, C., Line, M. R., et al. 2018, *AJ*, 156, 133, doi: [10.3847/1538-3881/aad781](https://doi.org/10.3847/1538-3881/aad781)
- Quirrenbach, A., Amado, P. J., Caballero, J. A., et al. 2016, *Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series*, Vol. 9908, CARMENES: an overview six months after first light, 990812, doi: [10.1117/12.2231880](https://doi.org/10.1117/12.2231880)
- Richard, C., Gordon, I. E., Rothman, L. S., et al. 2012, *JQSRT*, 113, 1276, doi: [10.1016/j.jqsrt.2011.11.004](https://doi.org/10.1016/j.jqsrt.2011.11.004)
- Ricker, G. R., Winn, J. N., Vanderspek, R., et al. 2015, *Journal of Astronomical Telescopes, Instruments, and Systems*, 1, 014003, doi: [10.1117/1.JATIS.1.1.014003](https://doi.org/10.1117/1.JATIS.1.1.014003)
- Robinson, T. D. 2017, *ApJ*, 836, 236, doi: [10.3847/1538-4357/aa5ea8](https://doi.org/10.3847/1538-4357/aa5ea8)
- Rodler, F., Lopez-Morales, M., & Ribas, I. 2012, *ApJL*, 753, L25, doi: [10.1088/2041-8205/753/1/L25](https://doi.org/10.1088/2041-8205/753/1/L25)
- Rogers, L. A., & Seager, S. 2010, *ApJ*, 712, 974, doi: [10.1088/0004-637X/712/2/974](https://doi.org/10.1088/0004-637X/712/2/974)
- Rothman, L. S., Gordon, I. E., Barber, R. J., et al. 2010, *JQSRT*, 111, 2139, doi: [10.1016/j.jqsrt.2010.05.001](https://doi.org/10.1016/j.jqsrt.2010.05.001)
- Rowe, J. F., Coughlin, J. L., Antoci, V., et al. 2015, *ApJS*, 217, 16, doi: [10.1088/0067-0049/217/1/16](https://doi.org/10.1088/0067-0049/217/1/16)
- Schwarz, G. 1978, *Annals of Statistics*, 6, 461
- Schwarz, H., Ginski, C., de Kok, R. J., et al. 2016, *A&A*, 593, A74, doi: [10.1051/0004-6361/201628908](https://doi.org/10.1051/0004-6361/201628908)
- Sing, D. K., Fortney, J. J., Nikolov, N., et al. 2016, *Nature*, 529, 59, doi: [10.1038/nature16068](https://doi.org/10.1038/nature16068)
- Snellen, I. A. G., de Kok, R. J., de Mooij, E. J. W., & Albrecht, S. 2010, *Nature*, 465, 1049, doi: [10.1038/nature09111](https://doi.org/10.1038/nature09111)
- Sparks, W. B., & Ford, H. C. 2002, *ApJ*, 578, 543, doi: [10.1086/342401](https://doi.org/10.1086/342401)
- Stevenson, K. B. 2016, *ApJL*, 817, L16, doi: [10.3847/2041-8205/817/2/L16](https://doi.org/10.3847/2041-8205/817/2/L16)
- Sullivan, P. W., Winn, J. N., Berta-Thompson, Z. K., et al. 2015, *ApJ*, 809, 77, doi: [10.1088/0004-637X/809/1/77](https://doi.org/10.1088/0004-637X/809/1/77)
- Szydlowski, M., Krawiec, A., Kurek, A., & Kamionka, M. 2015, *European Physical Journal C*, 75, 5, doi: [10.1140/epjc/s10052-014-3236-1](https://doi.org/10.1140/epjc/s10052-014-3236-1)
- Trotta, R. 2008, *Contemporary Physics*, 49, 71, doi: [10.1080/00107510802066753](https://doi.org/10.1080/00107510802066753)
- Wakeford, H. R., Wilson, T. J., Stevenson, K. B., & Lewis, N. K. 2019, *Research Notes of the AAS*, 3, 7, doi: [10.3847/2515-5172/aafc63](https://doi.org/10.3847/2515-5172/aafc63)
- Wakeford, H. R., Stevenson, K. B., Lewis, N. K., et al. 2017, *ApJL*, 835, L12, doi: [10.3847/2041-8213/835/1/L12](https://doi.org/10.3847/2041-8213/835/1/L12)
- Wildi, F., Blind, N., Reshetov, V., et al. 2017, in *Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series*, Vol. 10400, Proc. SPIE, 1040018, doi: [10.1117/12.2275660](https://doi.org/10.1117/12.2275660)
- Yurchenko, S. N., & Tennyson, J. 2014, *MNRAS*, 440, 1649, doi: [10.1093/mnras/stu326](https://doi.org/10.1093/mnras/stu326)
- Yurchenko, S. N., Tennyson, J., Barber, R. J., & Thiel, W. 2013, *Journal of Molecular Spectroscopy*, 291, 69, doi: [10.1016/j.jms.2013.05.014](https://doi.org/10.1016/j.jms.2013.05.014)
- Zellem, R. T., Swain, M. R., Cowan, N. B., et al. 2019, *PASP*, 131, 094401, doi: [10.1088/1538-3873/ab2d54](https://doi.org/10.1088/1538-3873/ab2d54)