



## UvA-DARE (Digital Academic Repository)

### Simultaneous suppression of ferromagnetism and superconductivity in UCoGe by Si substitution

de Nijs, D.E.; Huy, N.T.; de Visser, A.

**DOI**

[10.1103/PhysRevB.77.140506](https://doi.org/10.1103/PhysRevB.77.140506)

**Publication date**

2008

**Published in**

Physical Review B

[Link to publication](#)

**Citation for published version (APA):**

de Nijs, D. E., Huy, N. T., & de Visser, A. (2008). Simultaneous suppression of ferromagnetism and superconductivity in UCoGe by Si substitution. *Physical Review B*, 77, 140506(R-1:4). <https://doi.org/10.1103/PhysRevB.77.140506>

**General rights**

It is not permitted to download or to forward/distribute the text or part of it without the consent of the author(s) and/or copyright holder(s), other than for strictly personal, individual use, unless the work is under an open content license (like Creative Commons).

**Disclaimer/Complaints regulations**

If you believe that digital publication of certain material infringes any of your rights or (privacy) interests, please let the Library know, stating your reasons. In case of a legitimate complaint, the Library will make the material inaccessible and/or remove it from the website. Please Ask the Library: <https://uba.uva.nl/en/contact>, or a letter to: Library of the University of Amsterdam, Secretariat, Singel 425, 1012 WP Amsterdam, The Netherlands. You will be contacted as soon as possible.



## Simultaneous suppression of ferromagnetism and superconductivity in UCoGe by Si substitution

D. E. de Nijs, N. T. Huy, and A. de Visser\*

*Van der Waals–Zeeman Institute, University of Amsterdam, Valckenierstraat 65, 1018 XE Amsterdam, The Netherlands*

(Received 11 February 2008; revised manuscript received 13 March 2008; published 29 April 2008)

We investigate the effect of substituting Si for Ge in the ferromagnetic superconductor UCoGe. dc-magnetization, ac-susceptibility, and electrical resistivity measurements on polycrystalline  $\text{UCoGe}_{1-x}\text{Si}_x$  samples show that ferromagnetic order and superconductivity are progressively depressed with increasing Si content and simultaneously vanish at a critical concentration  $x_{cr} \approx 0.12$ . The non-Fermi-liquid temperature variation in the electrical resistivity near  $x_{cr}$  and the smooth depression of the ordered moment point to a continuous ferromagnetic quantum phase transition. Superconductivity is confined to the ferromagnetic phase, which provides further evidence for magnetically mediated superconductivity.

DOI: 10.1103/PhysRevB.77.140506

PACS number(s): 74.70.Tx, 74.62.Dh, 75.30.Kz

Recently, it was discovered<sup>1</sup> that the intermetallic compound UCoGe belongs to the small group of ferromagnetic superconductors (FMSCs): superconductivity with a transition temperature  $T_s = 0.8$  K coexists with weak itinerant ferromagnetic (FM) order with a Curie temperature  $T_C = 3$  K. Ferromagnetic superconductors attract much interest because in the standard BCS scenario, superconductivity (SC) and FM are incompatible.<sup>2</sup> This is due to the strong depairing effect of the ferromagnetic exchange interaction, which thwarts phonon mediated formation of singlet Cooper pairs. However, an alternative route is offered by spin fluctuation models,<sup>3,4</sup> in which critical magnetic fluctuations associated with a ferromagnetic quantum critical point (FM QCP) mediate SC by pairing the electrons in triplet states. The FMSCs discovered so far are  $\text{UGe}_2$  (under pressure),<sup>5</sup>  $\text{UIr}$  (under pressure),<sup>6</sup>  $\text{URhGe}$ ,<sup>7</sup> and  $\text{UCoGe}$ .<sup>1</sup> The latter two compounds offer the advantage that SC occurs at ambient pressure, which facilitates the use of a wide range of experimental techniques to probe magnetically mediated SC.

UCoGe crystallizes<sup>8,9</sup> in the orthorhombic  $\text{TiNiSi}$  structure (space group  $P_{nma}$ ). Evidence for the proximity to a FM QCP has been extracted from magnetization and specific heat measurements<sup>1</sup> on polycrystalline samples. The low  $T_C = 3$  K and the small value of the ordered moment  $m_0 = 0.03\mu_B$  reveal that magnetism is weak. Itinerant magnetism is corroborated by the small value of the magnetic entropy<sup>1</sup> (0.3% of  $R \ln 2$ ) associated with the magnetic transition. More recently, the magnetic and SC properties were determined for a single-crystalline sample.<sup>10</sup> Magnetization data reveal that UCoGe is a uniaxial ferromagnet with the ordered moment  $m_0 = 0.07\mu_B \approx 2m_0^{\text{poly}}$  pointing along the  $c$  axis. The electrical resistivity  $\rho(T)$  measured for a current  $I \parallel a$  shows SC below 0.6 K and a sharp kink signaling the Curie temperature  $T_C = 2.8$  K. The temperature variation of the resistivity<sup>10,11</sup> is characteristic<sup>12</sup> for a weak itinerant FM near a critical point, i.e., a Fermi liquid  $\rho \propto T^2$  dependence below  $T_C$  and scattering at critical FM fluctuations  $\rho \propto T^{5/3}$  there above.

In the generic pressure-temperature phase diagram for FMSCs<sup>3,4,13,14</sup> the superconducting phase (the dome) is confined to the magnetic phase and  $T_C$  and  $T_s$  vanish at the same critical pressure. Such a phase diagram has been reported for  $\text{UGe}_2$  (Ref. 5) and  $\text{UIr}$  (Ref. 6) under pressure. In the case of

UCoGe,<sup>1</sup> the analysis of the thermal expansion and specific heat data, using the Ehrenfest relation, shows that  $T_C$  decreases with pressure, whereas  $T_s$  increases. This places UCoGe on the far side of the superconducting dome with respect to the magnetic quantum critical point. Concurrently, under hydrostatic mechanical pressure,  $T_s$  is predicted to go through a maximum before vanishing at the critical point. In this work, we use an alternative route to study the evolution of FM and SC, namely, chemical pressure exerted by replacing Ge by isoelectronic Si. Ferromagnetic UCoGe and paramagnetic<sup>15</sup>  $\text{UCoSi}$  are isostructural.<sup>8,9</sup> The unit cell volume of  $\text{UCoSi}$  is  $\sim 3.5\%$  smaller than the one of  $\text{UCoGe}$ , so chemical pressure is relatively weak. By means of magnetic and transport measurements, we find that FM order and SC are gradually depressed and simultaneously vanish at a critical concentration  $x_{cr} \approx 0.12$ . SC is confined to the FM phase in agreement with the generic phase diagram. This yields further support for magnetically mediated superconductivity.

A series of polycrystalline  $\text{UCoGe}_{1-x}\text{Si}_x$  samples was prepared with  $0 \leq x \leq 0.20$  and  $x = 1$ . The constituents (natural U 3N, Co 4N, Ge 5N, and Si 5N) were weighed according to the nominal composition  $\text{U}_{1.02}\text{Co}_{1.02}\text{Ge}_{1-x}\text{Si}_x$  and arc melted together under a high-purity argon atmosphere in a water-cooled copper crucible. The as-cast samples were annealed for ten days at 875 °C. Samples were cut by spark erosion in a bar shape for transport and magnetic measurements. The phase homogeneity of the annealed samples was investigated by electron probe microanalysis (EPMA). The matrix had the 1:1:1 composition and all samples contained a small amount (2%) of impurity phases. The EPMA technique did, however, not allow for a precise determination of the Ge and Si ratio, and in the following,  $x$  is the nominal concentration. Powder x-ray diffraction patterns at  $T = 300$  K for  $x = 0.0, 0.1, 0.2,$  and  $1.0$  confirmed the  $\text{TiNiSi}$  structure. The measured lattice constants are  $a = 6.864$  Å,  $b = 4.196$  Å, and  $c = 7.261$  Å for  $\text{UCoGe}$  and  $a = 6.876$  Å,  $b = 4.108$  Å, and  $c = 7.154$  Å for  $\text{UCoSi}$ , in good agreement with the literature.<sup>8,9</sup> The unit cell volume  $\Omega$  linearly decreases from 209.5 ( $x = 0$ ) to 202.1 Å<sup>3</sup> ( $x = 1$ ), with the main contraction along the  $b$  and  $c$  axes.

The dc magnetization,  $M(T, B)$ , was measured in a superconducting quantum inference device magnetometer in magnetic fields up to 5 T and temperatures down to 2 K. The low-field ( $B = 10^{-5}$  T) ac susceptibility  $\chi_{ac}$  was measured us-

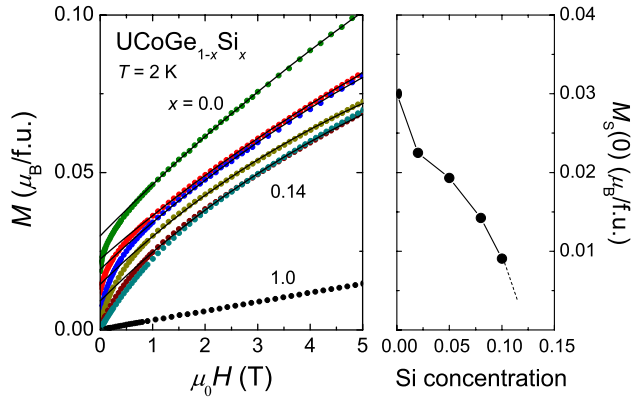


FIG. 1. (Color online) Left panel: magnetization as a function of field for  $\text{UCoGe}_{1-x}\text{Si}_x$  alloys at  $T = 2$  K. Si concentrations are (from top to bottom)  $x = 0, 0.02, 0.05, 0.08, 0.10, 0.14$ , and  $1.0$ . The solid lines represent fits to Eq. (1) for  $x \leq 0.10$ . Right panel: spontaneous magnetization  $M_s(0)$  as a function of Si content.

ing a mutual inductance coil and a phase-sensitive bridge in a  $^3\text{He}$  system with base temperature of 0.23 K or in a dilution refrigerator with base temperature of 0.02 K. Electrical resistivity data  $\rho(T)$  were taken by using a low-frequency ac bridge in a four-point configuration in the same temperature range.

The dc-magnetic susceptibility  $\chi_{dc}(T)$  of the  $\text{UCoGe}_{1-x}\text{Si}_x$  alloys was measured in an applied field of 1 T in the temperature range of 2–300 K. The effect of doping small amounts of Si on  $\chi_{dc}(T)$  is weak. For all  $x \leq 0.20$ , the data for  $T = 50$ –300 K are described by a modified Curie–Weiss law, with a temperature independent susceptibility  $\chi_0 \approx 10^{-8} \text{ m}^3/\text{mol}$  and an effective moment  $p_{eff} \approx 1.6 \pm 0.1 \mu_B/\text{f.u.}$  On the contrary, the effect of doping on the FM transition is large. Measurements of the dc magnetization in a small field ( $B = 0.01$  T) show that upon Si doping, the FM transition is rapidly suppressed to below the low temperature limit of our dc magnetometer (2 K). For  $x = 0.00$  and  $0.02$ , we find  $T_C = 3.0$  and  $2.5$  K, respectively. In Fig. 1, we show the field dependence of the magnetization  $M(H)$  measured at  $T = 2$  K. The gradual increase in  $M(H)$  observed for  $B \geq 1$  T is related to the itinerant nature of the magnetic state. The spontaneous magnetization  $M_s(H=0)$  rapidly drops with increasing Si content. For the ordered compounds, an estimate of  $M_s(0)$  can be made by fitting the data to the empirical expression

$$M(H) = M_s(0) + \Delta M(1 - e^{-\mu_0 H/B_0}), \quad (1)$$

where the parameter  $B_0$  probes the magnetic interaction strength of the fluctuating moments. In the high-field limit  $M(H=\infty) = M_s(H=0) + \Delta M$ , Eq. (1) describes the experimental data well for  $B \geq 1$  T (solid lines in Fig. 1). The intercepts of the fits with the vertical axis yield the fit parameters  $M_s(H=0)$  in the limit  $T \rightarrow 0$ . The deviations for  $B < 1$  T are due to the finite temperature at which the data are taken (the ordered moment is not fully developed yet). For  $x = 0.00$ ,  $M_s(0) \approx 0.029 \mu_B$  ( $T \rightarrow 0$ ), in agreement with previous results,<sup>1</sup> while for  $x = 0.02$ ,  $M_s(0) \approx 0.022 \mu_B$ . For the samples with  $x = 0.05, 0.08$ , and  $0.10$ , the data have been taken at  $T$

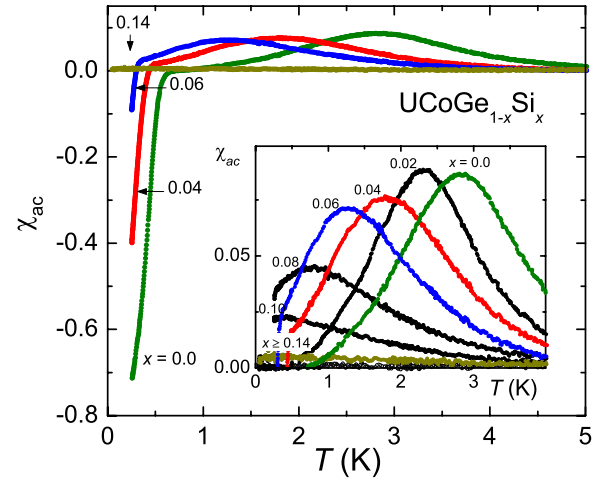


FIG. 2. (Color online) Temperature variation of the ac susceptibility (SI units) of  $\text{UCoGe}_{1-x}\text{Si}_x$  alloys for  $x = 0.00, 0.04, 0.06$ , and  $0.14$ . The inset shows  $\chi_{ac}$  around the ferromagnetic transition for  $0.00 \leq x \leq 0.20$ . The data for  $x = 0.20$  fall on the horizontal axis.

$> T_C$ . Nevertheless, a rough estimate of  $M_s(0)$  can be obtained, as the magnetic transition shows a large temperature broadening in applied fields  $B > 1$  T. The resulting values of  $M_s(0)$  are traced in the right panel of Fig. 1. We conclude that  $M_s(0)$  smoothly goes to zero in the concentration range  $0.10 < x < 0.14$ .

The suppression of  $T_C$  was studied in more detail by the ac-susceptibility technique. The data, taken down to 0.23 K ( $0.00 \leq x \leq 0.10$ ) and down to 0.02 K ( $0.14 \leq x \leq 0.20$ ), are shown in Fig. 2. The maximum in  $\chi_{ac}$  locates the Curie temperature, which equals 2.8 and 2.3 K, for  $x = 0.00$  and  $0.02$ , respectively. These values compare well with those extracted from the dc magnetization. With increasing Si content, the transition becomes weaker and broadens (see inset of Fig. 2) and for  $x \geq 0.14$ , a maximum in  $\chi_{ac}$  can no longer be identified. This confirms that magnetism vanishes in the concentration range  $0.10 < x < 0.14$ . The large diamagnetic signal measured for  $x = 0.00, 0.04$ , and  $0.06$  down to 0.23 K signals bulk SC. SC is progressively depressed and is no longer observed for  $x = 0.14$  (at least down to 0.02 K).

The electrical resistivity was measured in the temperature interval of 0.23–10 K for  $x \leq 0.08$  and in the range of 0.02–10 K for  $0.10 \leq x \leq 0.20$ . For  $x = 0.00$ , the residual resistivity  $\rho_0 = 26 \mu\Omega \text{ cm}$ . Upon alloying,  $\rho_0$  linearly increases at least up to  $x = 0.08$  at the fast rate of  $12 \mu\Omega \text{ cm/at. \% Si}$ . This shows all Si substitutes for Ge. Concurrently, the residual resistance ratio (RRR) =  $R(300 \text{ K})/R(1 \text{ K})$ , which amounts to 27 for  $x = 0$ , drops to  $\sim 5$  for  $x = 0.08$ . For  $x \geq 0.10$ , however, the RRR levels off at a value of  $\sim 4$ . The strong doping sensitivity of  $\rho_0$  is possibly related to an enhanced site inversion  $\text{Ge, Si} \leftrightarrow \text{Co}$ . Notice that the  $\text{TiNiSi}$  structure is an ordered variant of the  $\text{CeCu}_2$  structure<sup>16</sup> (for UTX compounds, crystallizing in the latter structure the transition metal atoms  $T$  and group IV atoms  $X$  are randomly distributed over the  $8h$  Cu sites).

The FM transition appears as a broad hump in  $\rho(T)$  for pure  $\text{UCoGe}$ .<sup>1</sup> Upon alloying, the hump shifts to lower temperatures at the same rate as the maximum in  $\chi_{ac}$ . In Fig. 3,

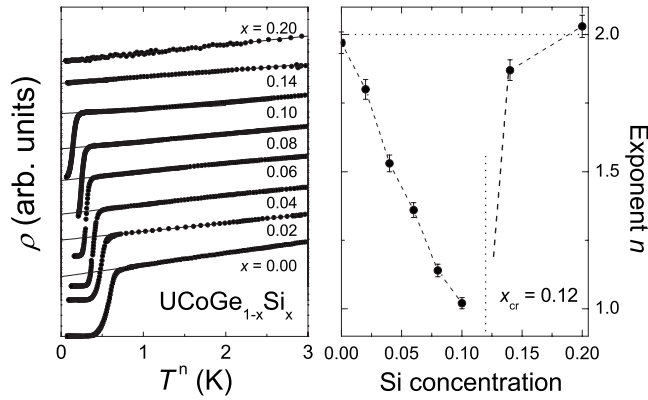


FIG. 3. Left panel: the electrical resistivity  $\rho$  (arbitrary units) plotted versus  $T^n$  of  $\text{UCoGe}_{1-x}\text{Si}_x$  alloys for  $x$  as indicated. The curves are shifted along the vertical axis for clarity. The straight solid lines represent fits  $\rho \sim T^n$  (see text). Right panel: exponent  $n$  versus Si concentration. The dashed line serves as a guide for the eyes. The vertical dotted line locates  $x_{cr}$ . The horizontal dotted line indicates  $n=2$ .

we show the low-temperature part of the resistivity data in a plot of  $\rho$  versus  $T^n$ . Here,  $n$  is determined by fitting  $\rho \sim T^n$  for  $T_s < T < T_C$ . For each  $x$ , the best value of  $n$  was obtained by fitting over a larger and larger temperature range, while keeping  $n$  constant and the error small. In the magnetic phase ( $x \leq 0.10$ ), the exponent shows a quasilinear decrease from  $n=2$  for  $x=0.00$  to the non-Fermi-liquid value  $n \approx 1$  for  $x=0.10$  (see Fig. 3). Close to the critical point, the temperature range for the fit becomes very small and the values of  $n$  should be interpreted with care. Nevertheless, the decreasing trend is evident. For  $x \geq 0.14$ , the Fermi liquid value  $n=2$  is recovered. The SC transition is depressed with increasing Si content and no SC has been observed down to 0.02 K for  $x=0.14$ .

Having determined the evolution of the FM and SC phases in the  $\text{UCoGe}_{1-x}\text{Si}_x$  alloys by magnetic and transport measurements, we construct the phase diagram shown in Fig. 4.  $T_C$  is quasilinearly depressed, at least until  $x=0.08$ , at a rate  $dT_C/dx = -0.25$  K/at. % Si. By extrapolating  $T_C(x) \rightarrow 0$ , we arrive at a critical Si concentration for the suppression of FM order  $x_{cr}^{FM} = 0.11$ . For  $x > 0.08$ , a tail appears, and the data extrapolate to  $x_{cr}^{FM} \approx 0.12$ .  $T_s$ , resistively determined by the midpoint of the transition, is depressed somewhat faster than linear, initially at a rate  $dT_s/dx = -0.06$  K/at. % Si. By smoothly extrapolating  $T_s(x) \rightarrow 0$ , we obtain a critical Si concentration for the suppression of SC  $x_{cr}^{SC} \approx 0.12$ . The  $T_s(x)$  values measured by  $\chi_{ac}(T)$  for  $x \leq 0.06$  signal the onset of bulk<sup>1</sup> SC and follow the same trend. Notice that  $T_s(x)$  bulk extrapolates to a slightly lower  $x_{cr}$ , i.e., close to the value  $x_{cr} = 0.11$  obtained by the linear extrapolation of  $T_C(x)$ .

In order to compare the effect of chemical and hydrostatic pressure, we calculate from the difference in unit cell volume of  $\text{UCoGe}$  and  $\text{UCoSi}$  that 1 at. % Si is equivalent to 0.35 kbar (here, we assume the isothermal compressibility  $\kappa \approx 10^{-11}$  Pa<sup>-1</sup>). Concurrently, the measured doping-induced depression of  $T_C$  (Fig. 4) translates to  $dT_C/dp = -0.71$  K/kbar, which is about a factor of 3 larger than the value of  $-0.25$  K/kbar calculated<sup>1</sup> via the Ehrenfest relation.

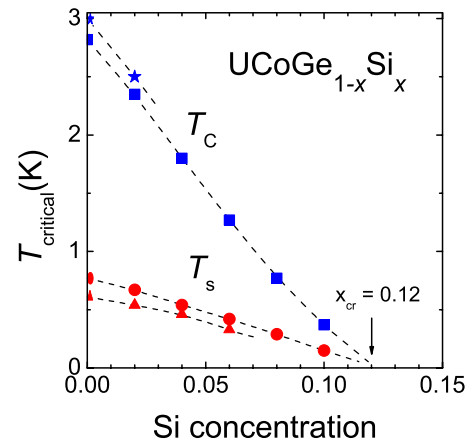


FIG. 4. (Color online) Curie temperature, determined by  $M(T)$  ( $\star$ ) and  $\chi_{ac}(T)$  ( $\blacksquare$ ), and superconducting transition temperature, determined by  $\rho(T)$  ( $\bullet$ ) and  $\chi_{ac}(T)$  ( $\blacktriangle$ ), as a function of  $x$  for  $\text{UCoGe}_{1-x}\text{Si}_x$  alloys. The dashed lines serve as a guide for the eyes. Superconductivity and ferromagnetism both vanish at  $x_{cr} \approx 0.12$ .

This indicates that Si does not merely exert chemical pressure. Indeed, hybridization phenomena in UTX alloys are, in general, strongly anisotropic.<sup>16</sup> As regards to the SC transition, Si doping obviously has a different effect than hydrostatic pressure. The measured doping-induced depression of  $T_s$  (Fig. 4) translates to  $dT_s/dp = -0.17$  K/kbar, while the Ehrenfest relation shows that  $T_s$  increases at a rate of  $dT_s/dp = 0.02$  K/kbar.<sup>17</sup>

The suppression of magnetic order in the  $\text{UCoGe}_{1-x}\text{Si}_x$  alloys can be understood in terms of a simple Doniach picture:<sup>18</sup> by doping the smaller Si atoms, the  $5f$ - $3d$  hybridization strength increases, which leads to a loss of magnetism. The rapid suppression of FM order provides further evidence that  $\text{UCoGe}$  is close to a FM QCP. This is corroborated by the steady decrease in the non-Fermi-liquid exponent  $n$  of the resistivity measured in the FM phase (see Fig. 3). The itinerant nature of the FM state suggests that the critical point is of the Moriya-Hertz-Millis<sup>12,19,20</sup> type. The extracted exponent  $n \approx 1$  near  $x_{cr}^{FM}$  is much smaller than the value  $n=5/3$  predicted for a clean FM QCP. A similar observation was made for the doping-induced FM QCP in  $\text{URh}_{1-x}\text{Ru}_x\text{Ge}$  alloys:<sup>21</sup> at  $x_{cr} = 0.38$ ,  $n \approx 1.2$ . Clearly, disorder reduces<sup>22</sup>  $n$ . The smooth depression of  $M_s(0)$  indicates that the ferromagnetic-to-paramagnetic transition at  $T=0$  K is a continuous phase transition. Additional experiments, e.g., specific heat, are required to set the evidence for a FM QCP at  $x_{cr} \approx 0.12$  on firm footing.

The magnetic and SC phase diagram (Fig. 4) presents compelling evidence that superconductivity is confined to the FM phase. Moreover, by smoothly extrapolating  $T_C(x)$  and  $T_s(x)$ , we arrive at a most important conclusion, namely,  $x_{cr}^{FM} = x_{cr}^{SC} \approx 0.12$ . This shows that FM order and SC are closely tied together. The simultaneous suppression of FM order and SC yields strong support for triplet SC mediated by FM spin fluctuations.<sup>3,4,13,14</sup> Evidence for triplet SC is furnished by the absence of Pauli limiting in the upper critical field  $B_{c2}$ .<sup>10</sup> Moreover, the observed anisotropy in  $B_{c2}$  provides evidence for an axial SC gap with nodes along the direction of the ordered moment, as calculated<sup>23</sup> for the A

phase of an orthorhombic FMSC. On the other hand, it is recognized<sup>3,24</sup> that triplet SC is extremely sensitive to scattering at nonmagnetic impurities and defects. Therefore, it is surprising that SC survives until doping concentrations of  $\sim 12$  at. % Si. For our polycrystalline UCoGe samples, with  $RRR \sim 30$ , we calculate<sup>1</sup> an electron mean free path,  $\ell \approx 500$  Å, in excess of the SC coherence length  $\xi \approx 150$  Å, a necessary condition for unconventional SC. Upon replacing Ge by Si, the residual resistance increases, leading to a corresponding decrease in  $\ell$ . Unconventional SC therefore would require a strong doping-induced reduction in  $\xi$  as well. The depression of non- $s$ -wave SC by nonmagnetic impurities can be modeled using a generalized form<sup>25,26</sup> of the Abrikosov-Gor'kov pair-breaking theory. A recent example is provided by the defect-driven depression of  $p$ -wave SC in the paramagnet  $\text{Sr}_2\text{RuO}_4$ .<sup>27</sup> In the case of the  $\text{UCoGe}_{1-x}\text{Si}_x$  alloys, however, the defect-driven depression of  $T_s$  is partly compensated by  $T_s$  increasing due to chemical pressure. Also, one may speculate that upon the approach of the FM

QCP, FM fluctuations stimulate triplet SC even stronger. Obviously, more experiments are needed to unravel the different pairing and depairing contributions to  $T_s$ .

In summary, magnetic and transport measurements on a series of polycrystalline  $\text{UCoGe}_{1-x}\text{Si}_x$  samples show that ferromagnetic order and superconductivity are both depressed and vanish at the same critical concentration  $x_{cr} \approx 0.12$ . The non-Fermi-liquid exponent in the resistivity near  $x_{cr}$  and the smooth depression of the ordered moment point to a continuous FM quantum phase transition. Superconductivity is confined to the ferromagnetic phase, which provides further evidence for magnetically mediated superconductivity. These results offer a unique route to investigate the emergence of superconductivity near a FM QCP at ambient pressure.

This work was part of the research program of FOM (Dutch Foundation for Fundamental Research of Matter) and COST Action P16 ECOM.

\*devisser@science.uva.nl

- <sup>1</sup>N. T. Huy, A. Gasparini, D. E. de Nijs, Y. Huang, J. C. P. Klaasse, T. Gortenmulder, A. de Visser, A. Hamann, T. Görlach, and H. v. Löhneysen, *Phys. Rev. Lett.* **99**, 067006 (2007).
- <sup>2</sup>N. F. Berk and J. R. Schrieffer, *Phys. Rev. Lett.* **17**, 433 (1966).
- <sup>3</sup>D. Fay and J. Appel, *Phys. Rev. B* **22**, 3173 (1980).
- <sup>4</sup>G. G. Lonzarich, in *Electron: A Centenary Volume*, edited by M. Springford (Cambridge University Press, Cambridge, 1997), Chap. 6.
- <sup>5</sup>S. S. Saxena, K. Ahilan, P. Agarwal, F. M. Grosche, R. K. Haselwimmer, M. Steiner, E. Pugh, I. R. Walker, S. R. Julian, P. Monthoux, G. G. Lonzarich, A. D. Huxley, I. Sheikin, D. Braithwaite, and J. Flouquet, *Nature (London)* **406**, 587 (2000).
- <sup>6</sup>T. Akazawa, H. Hidaka, T. Fujiwara, T. C. Kobayashi, E. Yamamoto, Y. Haga, R. Settai, and Y. Ōnuki, *J. Phys.: Condens. Matter* **16**, L29 (2004).
- <sup>7</sup>D. Aoki, A. D. Huxley, E. Ressouche, D. Braithwaite, J. Flouquet, J. P. Brison, E. Lhotel, and C. Paulsen, *Nature (London)* **413**, 613 (2001).
- <sup>8</sup>B. Lloret, Ph.D. thesis, University of Bordeaux I, 1988.
- <sup>9</sup>F. Canepa, P. Manfrinetti, M. Pani, and A. Palenzona, *J. Alloys Compd.* **234**, 225 (1996).
- <sup>10</sup>N. T. Huy, D. E. de Nijs, Y. K. Huang, and A. de Visser, *Phys. Rev. Lett.* **100**, 077002 (2008).
- <sup>11</sup>N. T. Huy, Ph.D. thesis, University of Amsterdam, 2008.
- <sup>12</sup>T. Moriya, *Spin Fluctuations in Itinerant Electron Magnetism* (Springer, Berlin, 1985).
- <sup>13</sup>T. R. Kirkpatrick, D. Belitz, T. Vojta, and R. Narayanan, *Phys. Rev. Lett.* **87**, 127003 (2001).
- <sup>14</sup>P. Monthoux, D. Pines, and G. G. Lonzarich, *Nature (London)* **450**, 1177 (2007).
- <sup>15</sup>R. Troć and V. H. Tran, *J. Magn. Magn. Mater.* **73**, 389 (1988).
- <sup>16</sup>V. Sechovský and L. Havela, *Handbook of Magnetic Materials*, edited by K. H. J. Buschow (North-Holland, Amsterdam, 1998), Vol. 11, pp. 1–289.
- <sup>17</sup>This value corrects the estimate  $dT_s/dp=0.048$  K/kbar given in Ref. 1.
- <sup>18</sup>S. Doniach, *Physica B & C* **91**, 231 (1977).
- <sup>19</sup>J. Hertz, *Phys. Rev. B* **14**, 1165 (1976).
- <sup>20</sup>A. J. Millis, *Phys. Rev. B* **48**, 7183 (1993).
- <sup>21</sup>N. T. Huy, A. Gasparini, J. C. P. Klaasse, A. de Visser, S. Sakarya, and N. H. van Dijk, *Phys. Rev. B* **75**, 212405 (2007).
- <sup>22</sup>See, e.g., C. Pfleiderer, S. R. Julian, and G. G. Lonzarich, *Nature (London)* **414**, 427 (2001).
- <sup>23</sup>V. P. Mineev and T. Champel, *Phys. Rev. B* **69**, 144521 (2004).
- <sup>24</sup>I. F. Foulkes and B. L. Gyorffy, *Phys. Rev. B* **15**, 1395 (1977).
- <sup>25</sup>P. J. Hirschfeld, P. Wölfle, and D. Einzel, *Phys. Rev. B* **37**, 83 (1988).
- <sup>26</sup>A. J. Millis, S. Sachdev, and C. M. Varma, *Phys. Rev. B* **37**, 4975 (1988).
- <sup>27</sup>A. P. Mackenzie, R. K. W. Haselwimmer, A. W. Tyler, G. G. Lonzarich, Y. Mori, S. Nishizaki, and Y. Maeno, *Phys. Rev. Lett.* **80**, 161 (1998).