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# Search for type-III seesaw heavy leptons in $pp$ collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector

G. Aad *et al.*\*

(ATLAS Collaboration)

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A search for the pair production of heavy leptons ( $N^0, L^\pm$ ) predicted by the type-III seesaw theory formulated to explain the origin of small neutrino masses is presented. The decay channels  $N^0 \rightarrow W^\pm l^\mp$  ( $l = e, \mu, \tau$ ) and  $L^\pm \rightarrow W^\pm \nu$  ( $\nu = \nu_e, \nu_\mu, \nu_\tau$ ) are considered. The analysis is performed using the final state that contains two leptons (electrons or muons), two jets from a hadronically decaying  $W$  boson and large missing transverse momentum. The data used in the measurement correspond to an integrated luminosity of  $20.3 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8$  TeV collected by the ATLAS detector at the LHC. No evidence of heavy lepton pair production is observed. Heavy leptons with masses below 325–540 GeV are excluded at the 95% confidence level, depending on the theoretical scenario considered.

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## I. INTRODUCTION

Experiments show that neutrinos have much smaller masses than charged leptons (see Ref. [1] and references therein). While in the Standard Model (SM) the charged fermions acquire masses by coupling to the Higgs ( $H$ ) boson, the neutrinos may become massive via new physics beyond the SM, e.g. via the introduction of Majorana mass terms [2]. These masses could be small due to the seesaw mechanism [3,4], which relies on new massive states that couple to a charged lepton and the Higgs field. Among different models for the seesaw mechanism, the type-III model [2,5] introduces at least two extra triplets of fermionic fields with zero hypercharge in the adjoint representation of  $SU(2)_L$  that generate neutrino masses and couple to gauge bosons. This model predicts new charged and neutral heavy leptons that could be produced in proton-proton collisions at the LHC.

A search by the CMS experiment [6] excluded the type-III seesaw heavy leptons with masses in the range of 100–210 GeV, depending on theoretical assumptions. A recent search by ATLAS [7] also sets complementary limits on heavy leptons using the three-lepton final state. Similar searches have also been done by the L3 experiment [8] ruling out charged heavy leptons with masses below 100 GeV.

In this paper, a search for heavy leptons predicted by the type-III seesaw mechanism is presented. The search explores the mass region above 100 GeV. A minimal type-III seesaw model [9] is used to optimize the analysis strategy and interpret the search results. The model introduces a triplet with one neutral and two oppositely charged leptons denoted by  $N^0$  and  $L^\pm$ , respectively. The heavy

leptons decay into a SM lepton and a  $W$ ,  $Z$ , or Higgs boson. The heavy leptons are assumed to be degenerate in mass. This assumption does not affect the result because in the case of a small mass splitting due to radiative corrections, the decays within the heavy leptons are highly suppressed [10]. The dominant production mechanism for type-III seesaw heavy leptons in  $pp$  collisions is pair production through the weak coupling to the  $W$  boson propagator,  $pp \rightarrow W^* \rightarrow N^0 L^\pm$ , and the largest branching fraction is the one with two  $W$  bosons in the final state,  $N^0 \rightarrow W^\pm l^\mp$  ( $l = e, \mu, \tau$ ) and  $L^\pm \rightarrow W^\pm \nu$  ( $\nu = \nu_e, \nu_\mu, \nu_\tau$ ). The production cross section does not depend on the mixing angles between the SM leptons and the new heavy lepton states  $V_\alpha$ , ( $\alpha = e, \mu, \tau$ ), which enter only in the expressions for the  $L$  and  $N$  decay widths. The fraction of  $L$  and  $N$  decays to lepton flavor  $\alpha$  is proportional to  $b_\alpha = |V_\alpha|^2 / (|V_e|^2 + |V_\mu|^2 + |V_\tau|^2)$ . The limits obtained may be interpreted in terms of a range of mixing angles and Yukawa couplings [9], allowing tests of a range of models with different couplings to gauge bosons and cross section predictions [11,12]. In the type-III seesaw model considered here, a benchmark point is defined by setting  $V_\tau$  to zero, so that  $b_e$  and  $b_\mu$  are determined only by the ratio  $V_e/V_\mu$ , taken to be 0.87 based on the separately allowed maximum values of  $V_e$  and  $V_\mu$  in Refs. [13–15]. This choice results in values of  $b_e = 0.53$ ,  $b_\mu = 0.43$ , and  $b_\tau = 0$ . The search is performed for the process  $pp \rightarrow N^0 L^\pm \rightarrow W^\pm l^\mp W^\pm \nu$ , where one  $W$  boson decays leptonically and the other  $W$  boson decays hadronically, resulting in a lepton pair in the final state with either the same charge [same sign (SS)] or with the opposite charge [opposite sign (OS)].

## II. DATA SAMPLE AND MONTE CARLO SIMULATION

The analysis uses data from  $\sqrt{s} = 8$  TeV  $pp$  collisions at the LHC that were recorded by the ATLAS detector

\*Full author list given at the end of the article.

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using single-electron and single-muon triggers. A detailed description of the ATLAS detector can be found elsewhere [16]. The data sample corresponds to  $20.3 \pm 0.6 \text{ fb}^{-1}$  [17] of integrated luminosity. Data quality criteria are applied to ensure that events were recorded with stable beam conditions and with all relevant subdetector systems operational. The triggers [18] are fully efficient for leptons with  $p_T > 25 \text{ GeV}$ , where transverse momentum  $p_T$  is defined as the magnitude of the momentum component orthogonal to the beam axis. Events are required to have a reconstructed collision vertex with at least three associated tracks, each with  $p_T > 400 \text{ MeV}$ . In events with multiple vertices, the vertex with the largest  $\sum p_T^2$  of associated tracks is taken as the primary event vertex.

Monte Carlo (MC) samples are used to optimize the event selection and to model the kinematics and normalization of most background processes. Signal samples are generated for heavy lepton masses in the range 100–600 GeV. MADGRAPH 5 [19] is used to calculate the matrix elements for each process, while MADEVENT [20] with the MSTW2008 parton distribution functions (PDFs) set [21] simulates the initial hard scattering and the  $N$  and  $L$  decay. PYTHIA 8.153 [22] is used to simulate the decays of  $W$  bosons and the underlying physics by providing parton showers and hadronization, as well as adding initial- and final-state radiation (ISR and FSR) to the events simulated in MADEVENT. The main background sources arise from the production of a  $Z$  boson in association with jets ( $Z + \text{jets}$ ), single and pair production of top quarks, and diboson production ( $WW, WZ, ZZ$ ). The  $Z + \text{jets}$  and diboson processes are simulated with SHERPA 1.4.1 [23], a generator based on a multileg matrix element calculation matched to the parton shower using the Catani-Krauss-Kuhn-Webber (CKKW) prescription [24], and using the CT10 [25] PDF set. For diboson production, both the electroweak and strong production processes are simulated [26]. Top-quark pair events and single-top-quark events in the  $Wt$ -channel and  $s$ -channel are simulated using MC@NLO 4.03 [27], which is interfaced to HERWIG 6.520 [28] and JIMMY 4.31 [29] with the CT10 PDF set. Top pair production in association with a vector boson,  $t\bar{t} + W/Z$ , is simulated using MADGRAPH with the CTEQ6L1 PDF set [30], interfaced to PYTHIA 8.153 for parton showering and hadronization. Single-top-quark production in the  $t$ -channel is simulated using ACERMC v3.8 [31] with PYTHIA 6.426 [32] and the CTEQ6L1 PDF set.

All samples of simulated events include the effect of multiple  $pp$  interactions in the same and neighboring bunch crossings (pileup) by overlaying simulated minimum-bias events on each generated signal and background event. The number of overlaid events is chosen to match the average number of interactions per  $pp$  bunch crossing observed in the data as it evolved throughout the data-taking period (giving an average of 21 interactions per crossing for the whole data-taking period). The generated

samples are processed through the GEANT4-based detector simulation [33,34] or a fast simulation using a parametrization of the performance of the calorimetry and GEANT4 for the other parts of the detector [35]. The standard ATLAS reconstruction software is used for both simulated and collision data.

### III. OBJECT DEFINITIONS

The reconstructed objects used in this analysis are electrons, muons, jets, and missing transverse momentum. Electrons are reconstructed from clusters of energy depositions in the calorimeter that match a track reconstructed in the inner detector (ID) and satisfy the “tight” criteria defined in Ref. [36]. The electrons are required to have  $p_T > 25 \text{ GeV}$  and pseudorapidity  $|\eta| < 2.47$  [37], excluding the transition region between the barrel and end caps in the liquid argon calorimeter ( $1.37 < |\eta| < 1.52$ ). Muons are reconstructed by combining ID and muon spectrometer tracks that are spatially matched and have consistent curvatures. The muon tracks are required to have  $p_T > 25 \text{ GeV}$  and  $|\eta| < 2.5$ . In addition, leptons are required to be isolated from other tracks and calorimetric activity [38]. To ensure that leptons originate from the interaction point, requirements of  $|d_0|/\sigma_{d_0} < 3$  and  $|z_0 \sin \theta| < 0.5 \text{ mm}$  are imposed on the electrons and muons, where  $d_0$  ( $z_0$ ) is the transverse (longitudinal) impact parameter of the lepton and  $\sigma_{d_0}$  is the uncertainty on the measured  $d_0$ . The lepton impact parameters are measured with respect to the event primary vertex.

Jets are reconstructed from three-dimensional topological clusters of energy depositions in the calorimeter using the anti- $k_r$  algorithm [39] with a radius parameter of  $R = 0.4$ . The energies of jets are calibrated to the hadronic energy scale by correcting for energy losses in passive material, the noncompensating response of the calorimeter, and extra energy due to multiple  $pp$  interactions [40]. The jets are required to have  $p_T > 30 \text{ GeV}$  and  $|\eta| < 2.8$ . For jets with  $p_T < 50 \text{ GeV}$  and  $|\eta| < 2.4$ , the summed scalar  $p_T$  of associated tracks from the reconstructed primary vertex is required to be at least 25% of the summed scalar  $p_T$  of all tracks associated with the jet. In the pseudorapidity range  $|\eta| < 2.5$ , jets containing  $b$ -hadrons are identified using a  $b$ -tagging algorithm [41] with an efficiency of 70% and with a misidentification rate for selecting light-quark or gluon jets of less than 1%. The identification efficiency of the algorithm for jets containing  $c$ -hadrons is 20%. The efficiencies and misidentification rates are determined from  $t\bar{t}$  MC events.

The missing transverse momentum vector (with its magnitude  $E_T^{\text{miss}}$ ) is derived using the calorimeter cell energies within  $|\eta| < 4.9$  and corrected on the basis of dedicated calibrations of the associated physics objects including muons [42]. Calorimeter cells containing energy depositions above noise and not associated with high- $p_T$  physics objects are also included.

#### IV. EVENT SELECTION

Events that contain exactly two reconstructed leptons (electrons or muons), at least two jets, and no  $b$ -tagged jets are selected. One of these leptons is required to match the object upon which the event was triggered. Different sets of optimized selection criteria are used for the events in the OS and SS final states. The optimization is done using simulated heavy lepton pair-production events at a benchmark mass of 300 GeV. For the OS (SS) final state, the leading and next-to-leading lepton candidates are required to have  $p_T$  greater than 100 (70) GeV and 25 (40) GeV, respectively. The invariant mass of the two lepton candidates is required to be larger than 130 (90) GeV in order to suppress background from the production of  $Z + \text{jets}$ . The hadronically decaying  $W$  candidate is formed by combining the two jets with highest  $p_T$ , and the  $p_T$  of the first and second leading jets are, respectively, required to be larger than 60 (40) and 30 (25) GeV, for the OS (SS) final state. The invariant mass of the  $W$  candidate,  $m_{jj}$ , is required to be between 60 and 100 GeV. Events selected in the OS (SS) final state are required to have a  $E_T^{\text{miss}}$  of at least 110 (100) GeV and, for OS events, an angular separation  $\Delta R_{jj} = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} < 2$  between the two jets with highest  $p_T$ , where  $\Delta\eta$  and  $\Delta\phi$  are defined as the differences in pseudorapidity and azimuthal angle between the jets.

#### V. BACKGROUND ESTIMATE

The background in this search can be classified into two categories based on the origin of the charged lepton candidates. The first category of backgrounds consists of events in which two leptons are produced via the decays of  $W$  or  $Z$  bosons and are correctly reconstructed. This category of backgrounds includes the production of  $Z + \text{jets}$ ,  $t\bar{t}$ ,  $Wt$  single-top-quark and diboson events. Smaller contributions originate from  $t\bar{t} + W/Z$  events. Contributions from triboson events, such as  $WWW$  and events containing a Higgs boson, are negligible. The number of events from this background category is estimated using the simulated samples described previously.

The second category corresponds to all other sources, such as events containing at least one particle that is incorrectly identified as a lepton, or a lepton which originates from secondary interactions and decays, which are together denoted as *fakes*, and events with a lepton whose charge is incorrectly determined. Fake electrons originate primarily from jets that have a large electromagnetic energy fraction passing the electron selection requirements, photon conversions, and electrons from semileptonic decays of charm or bottom hadrons. Fake muons include muons arising from semileptonic decays of charm or bottom hadrons, in-flight decays of pions or kaons, or energetic particles that reach the muon spectrometer.

For the OS final state, the background contribution is dominated by events from the first category. For the SS

final state, the expected background is very small and is also primarily from the first category. Events from  $Z + \text{jets}$  production can mimic the SS signal if the charge of one of the electrons is incorrectly determined. The background events from this contribution are modeled by simulation, with correction factors derived using  $Z \rightarrow ee$  events in data. The probability of misidentifying the muon charge is negligible.

The background contribution in the second category (i.e. fake leptons) is estimated using an *in situ* technique [38] that relies minimally on simulation. This is done by reweighting a complementary set of events, selected by changing the electron identification criterion from tight to “loose” [43] and by loosening the muon  $|d_0|/\sigma_{d_0}$  and the electron and muon isolation requirements, while keeping the event selection otherwise identical. The reweighting factors are defined as the ratio of the number of events containing a lepton that satisfies the nominal criteria to the number of events containing a lepton that only fulfills the relaxed criteria. These factors are measured as a function of the candidate  $p_T$  and  $\eta$  in data samples that are enriched in fake leptons [38]. Corrections to the factors due to true leptons from vector boson decay in the background-enriched samples are taken from MC simulation.

Figure 1 shows a comparison of the missing transverse momentum distribution of data, expected backgrounds, and signal predictions when all the selection requirements, except for the missing transverse momentum requirement, are applied. The shape and the rate of the background estimate is in good agreement with the data.

The background estimates are validated by comparing the predicted numbers of events in simulation to those observed in the data in several control regions that have event selection criteria similar to those for the signal region.

The control region for top-quark pairs is defined by selecting events with two  $b$ -tagged jets. In this region, according to MC simulation, all the events are from top-quark pair and single-top-quark production with a negligible contribution from other sources. MC simulation predicts  $26 \pm 3$  (stat.) events, and in the data 32 events are observed. The scale factor, the ratio of the observed and predicted event yields, is found to be consistent with unity.

The diboson control region is obtained using a  $WZ$ -enriched sample of events containing three leptons without a requirement on the missing transverse momentum. In this control region, according to MC simulation, all the events are from diboson production with a negligible contribution from other sources. MC simulation predicts  $11 \pm 1$  events, and in the data nine events are observed. The scale factor is found to be consistent with unity.

#### VI. SYSTEMATIC UNCERTAINTIES

The uncertainties on the rate of top-quark and diboson backgrounds due to potential differences between the data and MC simulation are evaluated using the statistical

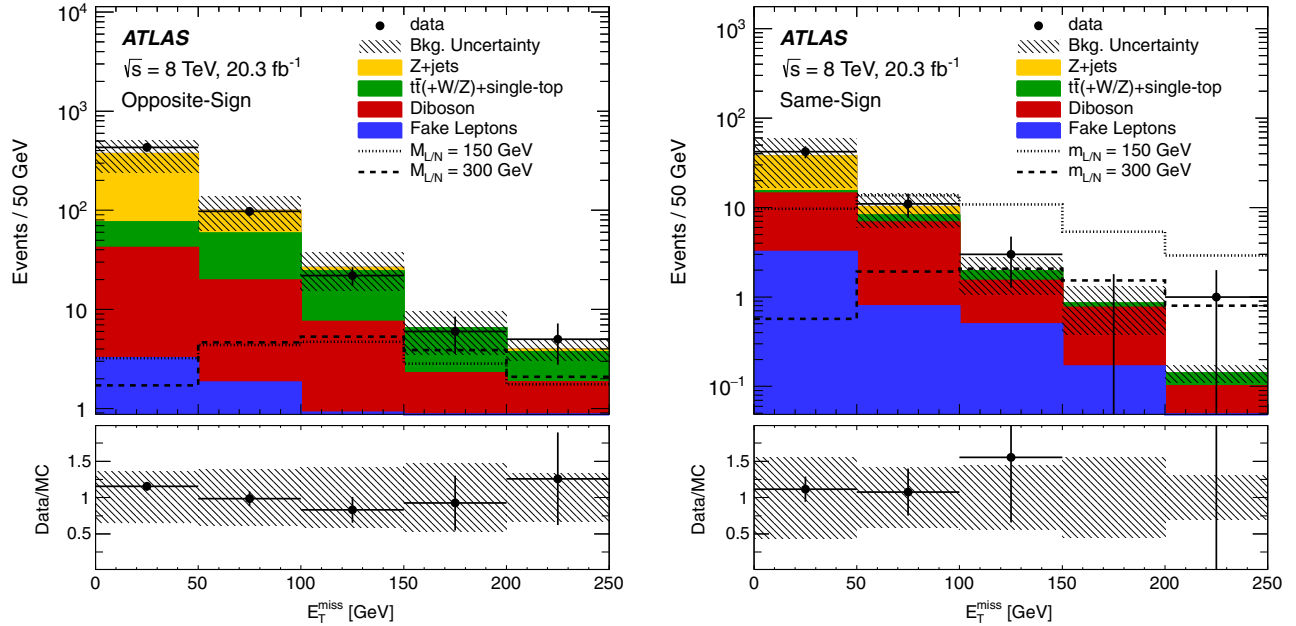


FIG. 1 (color online). The missing transverse momentum distribution of opposite-sign (left) and same-sign (right) events for data and predictions. The dotted and dashed lines show the expected distribution of type-III seesaw lepton pair production with masses of 150 and 300 GeV, respectively. The events shown here are required to pass all selection requirements except that on the missing transverse momentum. The uncertainties shown include both statistical and systematic uncertainties.

uncertainties of the measured scale factors in the control samples. They are the dominant systematic uncertainties ( $\sim 35\%$ ) on the background estimates. For top-quark production, an additional systematic uncertainty is considered to account for a potential difference between scale factors in the control region and the signal region. This is done by comparing the nominal  $t\bar{t}$  sample to alternative  $t\bar{t}$  MC samples. These samples include events that are generated using POWHEG-BOX 1.0 (patch 4) [44–46] and the leading-order, multileg generator ALPGEN v2.13 [47]. The POWHEG-BOX generator is interfaced to the PYTHIA 6.426 showering routines with either CT10 or HeraPDF [48] PDF sets and with the POWHEG  $h_{\text{damp}}$  parameter set to either the mass of the top quark or infinity [49]. ALPGEN is interfaced to HERWIG 6.520 and used to simulate top pair events with up to four additional partons in the matrix element. The uncertainties due to QCD ISR and FSR modeling are estimated with samples generated with ACERMC v3.8 interfaced to PYTHIA 6.426 in which the parton shower parameters are varied in a range consistent with a measurement of additional hadronic activity in  $t\bar{t}$  events [50]. The differences observed in the signal region by using different MC simulations are about 35%.

For the  $Z + n$ -jets ( $n \geq 2$ ) background estimate, the dominant systematic uncertainty in the OS final state is from the uncertainty on its production cross section ( $\sim 50\%$ ) [51]. For the SS final state, the systematic uncertainty is dominated by the statistical uncertainty on the measured electron charge misidentification rate.

Uncertainties on the background estimate due to fake leptons are determined in dedicated studies using a

combination of simulation and data. They account for potential biases in the method used to extract the reweighting factors and for the dependency of the reweighting factors on the event topology.

For both the predicted signal and background event yields, uncertainties resulting from detector effects from jet energy scale and resolution [40], lepton reconstruction and identification efficiencies [36,52], lepton momentum scales and resolutions [52,53], and missing transverse momentum [42] are considered. They are typically small (1%–5%). The theoretical uncertainties on the signal production cross section and acceptance, such as PDF choice and ISR and FSR modeling, are found to be negligible.

The background estimates and their uncertainties are tested in two other regions: a  $Z + \text{jets}$  control region and

TABLE I. Event yields for opposite-sign (OS) and same-sign (SS) selection for predicted backgrounds, data, and type-III seesaw lepton pair production with masses of 150 and 300 GeV. The reported errors include both the statistical and systematic uncertainties.

	OS	SS
Fake leptons	$1.4 \pm 0.9$	$0.67 \pm 0.42$
$Z + \text{jets}$	$2.4 \pm 1.2$	$0.06 \pm 0.23$
$WW/WZ/ZZ$	$9.2 \pm 2.9$	$1.95 \pm 0.58$
$t\bar{t}(+W/Z)$ and single top	$17.9 \pm 6.9$	$0.47 \pm 0.25$
Total	$31.0 \pm 7.7$	$3.15 \pm 0.80$
Data	25	4
Signal $m_{L/N} = 150$ GeV	$9.5 \pm 1.6$	$20.3 \pm 2.3$
Signal $m_{L/N} = 300$ GeV	$12.2 \pm 0.6$	$5.7 \pm 0.5$

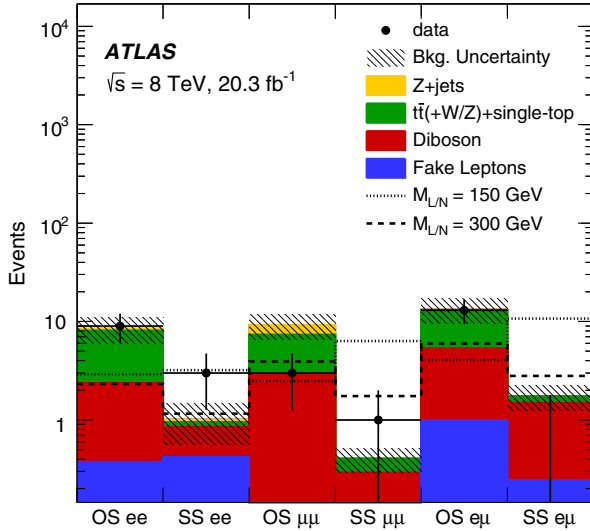


FIG. 2 (color online). Event yields for opposite-sign (OS) and same-sign (SS) selection for  $ee$ ,  $\mu\mu$ , and  $e\mu$  predicted backgrounds, data, and signal events featuring type-III seesaw lepton pair production with masses 150 and 300 GeV. The reported uncertainties include both the statistical and systematic uncertainties.

the hadronic  $W$  sidebands. The  $Z$  + jets control region is selected by requiring the invariant mass of the oppositely charged lepton pair to be consistent with the  $Z$  boson mass. In this region, approximately 70% of the events are from a  $Z$  boson produced in association with jets. The predicted number of events is  $34 \pm 4$ , where the error includes both the statistical and systematic uncertainties, and 32 events are observed in the data. For the hadronic  $W$  sidebands, the analysis was repeated using the same event selection but

requiring the invariant mass of the hadronically decaying  $W$  candidate to be  $35 < m_{jj} < 60$  GeV or  $100 < m_{jj} < 125$  GeV. This selection provides samples dominated by background events with kinematic properties similar to those of the signal candidates. In this region the predicted number of events is  $34 \pm 7$  events, where the error includes both the statistical and systematic uncertainties, and in the data 18 events are observed. The data are in agreement with the predictions within 1.9 standard deviations.

## VII. RESULTS AND INTERPRETATION

Table I shows the predicted numbers of signal and background events and the observed data events in the signal region. The data agree with the background-only hypothesis. Figure 2 shows the flavor composition of the simulated signal and background events and of the observed events in data. In the absence of any significant data excess, upper limits on the production rate of  $pp \rightarrow N^0 L^\pm \rightarrow W^\pm \ell^\mp W^\pm \nu$  at the 95% confidence level (C.L.) are derived as a function of the heavy lepton mass using the  $CL_s$  method [54]. The results of limit calculations combining the observations in the OS and SS final states are shown in Figs. 3 and 4. With the default mixing angles considered here, heavy leptons with masses less than 335 GeV are excluded by the analysis, while masses less than 475 GeV are excluded in the scenario in which heavy leptons can only decay to the  $W\ell$  or  $W\nu$  final states. For comparison, upper limits are also calculated for different theoretical assumptions, such as exclusive coupling between the heavy leptons and muons ( $b_e = 0, b_\mu = 1$ ) or electrons ( $b_e = 1, b_\mu = 0$ ). In the limit calculations for exclusive couplings, events in the two-muon (two-electron) final state

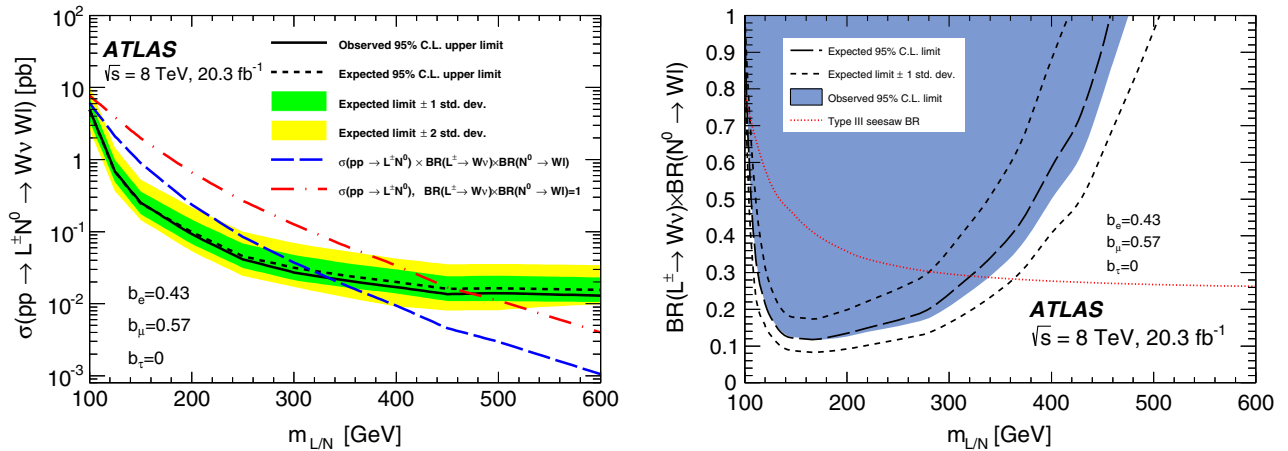


FIG. 3 (color online). Left: Observed (solid line) and expected (dashed line) 95% C.L. upper limits on the type-III seesaw heavy lepton cross section as a function of the heavy lepton  $L$  or  $N$  mass assuming  $b_e = 0.43$ ,  $b_\mu = 0.57$ , and  $b_\tau = 0$ . The bands surrounding the expected limit correspond to one and two standard deviations on the expected limit. The large-dashed (dot-dashed) line shows the theoretical prediction for  $m_{L/N}$ -dependent (maximal) branching fraction for decays to a  $W$  boson. Right: The expected (large-dashed line) and observed (shaded region) upper limits at the 95% C.L. on the  $BR(L^\pm \rightarrow W\nu) \times BR(N^0 \rightarrow W^\pm \ell^\mp)$  vs  $m_{L/N}$ . The dashed line corresponds to one standard deviation around the expected limit. The dotted line shows the nominal mass-dependent branching ratio of the type-III seesaw model.

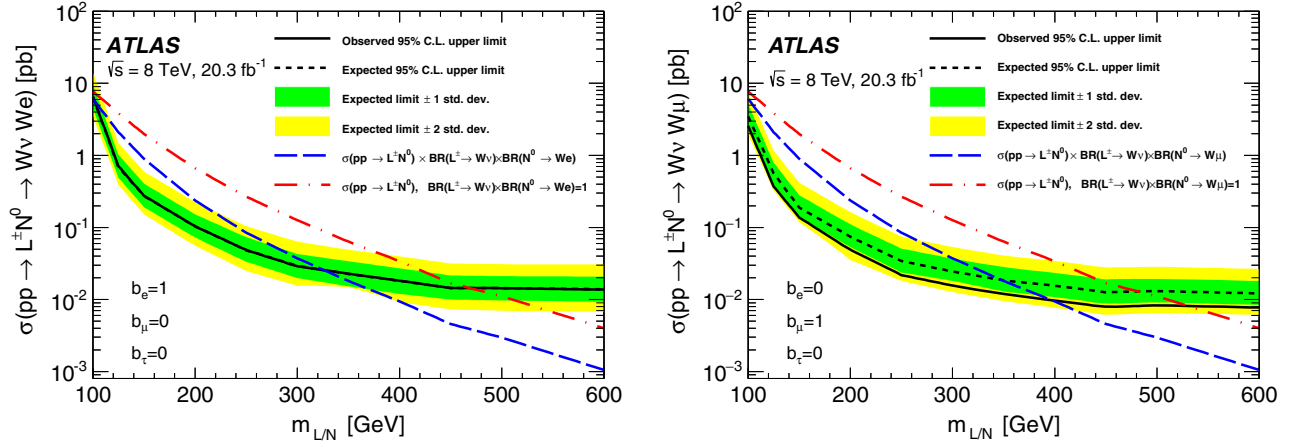


FIG. 4 (color online). Observed (solid line) and expected (dashed line) 95% C.L. upper limits on the type-III seesaw heavy lepton cross section as a function of the heavy lepton  $L$  or  $N$  mass for exclusive coupling to electrons (left) and exclusive coupling to muons (right). In the left plot, the dashed line is covered by the solid line. The large-dashed (dot-dashed) line shows the theoretical prediction for  $m_{L/N}$  dependent (maximal) branching fraction for decays to a  $W$  boson.

are excluded for exclusive electron (muon) coupling. Heavy leptons with a mass below 400 (325) GeV can be ruled out by the data in the case of exclusive coupling to muons (electrons). Masses less than 540 (470) GeV are excluded in the scenario in which heavy leptons can only decay to the  $W\ell$  or  $W\nu$  final state.

### VIII. CONCLUSIONS

A search for the pair production of heavy leptons predicted by the type-III seesaw model is presented. The analysis is performed using a final state that contains two leptons, two jets from a hadronically decaying  $W$  boson, and a large missing transverse momentum. The data used in the search correspond to an integrated luminosity of  $20.3 \text{ fb}^{-1}$  of  $pp$  collisions at  $\sqrt{s} = 8 \text{ TeV}$  collected by the ATLAS detector at the LHC. No evidence of heavy lepton production is observed. Heavy leptons with masses below 325–540 GeV are excluded at the 95% confidence level, depending on the considered theoretical scenarios.

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G. Aad,<sup>85</sup> B. Abbott,<sup>113</sup> J. Abdallah,<sup>151</sup> O. Abidin,<sup>11</sup> R. Aben,<sup>107</sup> M. Abolins,<sup>90</sup> O. S. AbouZeid,<sup>158</sup> H. Abramowicz,<sup>153</sup> H. Abreu,<sup>152</sup> R. Abreu,<sup>30</sup> Y. Abulaiti,<sup>146a,146b</sup> B. S. Acharya,<sup>164a,164b,b</sup> L. Adamczyk,<sup>38a</sup> D. L. Adams,<sup>25</sup> J. Adelman,<sup>108</sup> S. Adomeit,<sup>100</sup> T. Adye,<sup>131</sup> A. A. Affolder,<sup>74</sup> T. Agatonovic-Jovin,<sup>13</sup> J. A. Aguilar-Saavedra,<sup>126a,126f</sup> S. P. Ahlen,<sup>22</sup> F. Ahmadov,<sup>65,c</sup> G. Aielli,<sup>133a,133b</sup> H. Akerstedt,<sup>146a,146b</sup> T. P. A. Åkesson,<sup>81</sup> G. Akimoto,<sup>155</sup> A. V. Akimov,<sup>96</sup>



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 M. Cano Bret,<sup>76</sup> J. Cantero,<sup>82</sup> R. Cantrill,<sup>126a</sup> T. Cao,<sup>40</sup> M. D. M. Capeans Garrido,<sup>30</sup> I. Caprini,<sup>26a</sup> M. Caprini,<sup>26a</sup>

- M. Capua,<sup>37a,37b</sup> R. Caputo,<sup>83</sup> R. Cardarelli,<sup>133a</sup> F. Cardillo,<sup>48</sup> T. Carli,<sup>30</sup> G. Carlino,<sup>104a</sup> L. Carminati,<sup>91a,91b</sup> S. Caron,<sup>106</sup>  
E. Carquin,<sup>32a</sup> G. D. Carrillo-Montoya,<sup>8</sup> J. R. Carter,<sup>28</sup> J. Carvalho,<sup>126a,126c</sup> D. Casadei,<sup>78</sup> M. P. Casado,<sup>12</sup> M. Casolino,<sup>12</sup>  
E. Castaneda-Miranda,<sup>145b</sup> A. Castelli,<sup>107</sup> V. Castillo Gimenez,<sup>167</sup> N. F. Castro,<sup>126a,h</sup> P. Catastini,<sup>57</sup> A. Catinaccio,<sup>30</sup>  
J. R. Catmore,<sup>119</sup> A. Cattai,<sup>30</sup> J. Caudron,<sup>83</sup> V. Cavaliere,<sup>165</sup> D. Cavalli,<sup>91a</sup> M. Cavalli-Sforza,<sup>12</sup> V. Cavasinni,<sup>124a,124b</sup>  
F. Ceradini,<sup>134a,134b</sup> B. C. Cerio,<sup>45</sup> K. Cerny,<sup>129</sup> A. S. Cerqueira,<sup>24b</sup> A. Cerri,<sup>149</sup> L. Cerrito,<sup>76</sup> F. Cerutti,<sup>15</sup> M. Cerv,<sup>30</sup>  
A. Cervelli,<sup>17</sup> S. A. Cetin,<sup>19c</sup> A. Chafaq,<sup>135a</sup> D. Chakraborty,<sup>108</sup> I. Chalupkova,<sup>129</sup> P. Chang,<sup>165</sup> B. Chapleau,<sup>87</sup>  
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S. Chekanov,<sup>6</sup> S. V. Chekulaev,<sup>159a</sup> G. A. Chelkov,<sup>65,i</sup> M. A. Chelstowska,<sup>89</sup> C. Chen,<sup>64</sup> H. Chen,<sup>25</sup> K. Chen,<sup>148</sup> L. Chen,<sup>33d,j</sup>  
S. Chen,<sup>33c</sup> X. Chen,<sup>33f</sup> Y. Chen,<sup>67</sup> H. C. Cheng,<sup>89</sup> Y. Cheng,<sup>31</sup> A. Cheplakov,<sup>65</sup> E. Cheremushkina,<sup>130</sup>  
R. Cherkaoui El Moursli,<sup>135e</sup> V. Chernyatin,<sup>25,a</sup> E. Cheu,<sup>7</sup> L. Chevalier,<sup>136</sup> V. Chiarella,<sup>47</sup> J. T. Childers,<sup>6</sup> G. Chiodini,<sup>73a</sup>  
A. S. Chisholm,<sup>18</sup> R. T. Chislett,<sup>78</sup> A. Chitan,<sup>26a</sup> M. V. Chizhov,<sup>65</sup> K. Choi,<sup>61</sup> S. Chouridou,<sup>9</sup> B. K. B. Chow,<sup>100</sup>  
V. Christodoulou,<sup>78</sup> D. Chromek-Burckhart,<sup>30</sup> J. Chudoba,<sup>127</sup> A. J. Chuinard,<sup>87</sup> J. J. Chwastowski,<sup>39</sup> L. Chytka,<sup>115</sup>  
G. Ciapetti,<sup>132a,132b</sup> A. K. Ciftci,<sup>4a</sup> D. Cinca,<sup>53</sup> V. Cindro,<sup>75</sup> I. A. Cioara,<sup>21</sup> A. Ciocio,<sup>15</sup> Z. H. Citron,<sup>172</sup> M. Ciubancan,<sup>26a</sup>  
A. Clark,<sup>49</sup> B. L. Clark,<sup>57</sup> P. J. Clark,<sup>46</sup> R. N. Clarke,<sup>15</sup> W. Cleland,<sup>125</sup> C. Clement,<sup>146a,146b</sup> Y. Coadou,<sup>85</sup> M. Cobal,<sup>164a,164c</sup>  
A. Coccaro,<sup>138</sup> J. Cochran,<sup>64</sup> L. Coffey,<sup>23</sup> J. G. Cogan,<sup>143</sup> B. Cole,<sup>35</sup> S. Cole,<sup>108</sup> A. P. Colijn,<sup>107</sup> J. Collot,<sup>55</sup> T. Colombo,<sup>58c</sup>  
G. Compostella,<sup>101</sup> P. Conde Muiño,<sup>126a,126b</sup> E. Coniavitis,<sup>48</sup> S. H. Connell,<sup>145b</sup> I. A. Connelly,<sup>77</sup> S. M. Consonni,<sup>91a,91b</sup>  
V. Consorti,<sup>48</sup> S. Constantinescu,<sup>26a</sup> C. Conta,<sup>121a,121b</sup> G. Conti,<sup>30</sup> F. Conventi,<sup>104a,k</sup> M. Cooke,<sup>15</sup> B. D. Cooper,<sup>78</sup>  
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G. Cree,<sup>29</sup> S. Crépe-Renaudin,<sup>55</sup> F. Crescioli,<sup>80</sup> W. A. Cribbs,<sup>146a,146b</sup> M. Crispin Ortuzar,<sup>120</sup> M. Cristinziani,<sup>21</sup> V. Croft,<sup>106</sup>  
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S. D'Auria,<sup>53</sup> M. D'Onofrio,<sup>74</sup> M. J. Da Cunha Sargedas De Sousa,<sup>126a,126b</sup> C. Da Via,<sup>84</sup> W. Dabrowski,<sup>38a</sup> A. Dafinca,<sup>120</sup>  
T. Dai,<sup>89</sup> O. Dale,<sup>14</sup> F. Dallaire,<sup>95</sup> C. Dallapiccola,<sup>86</sup> M. Dam,<sup>36</sup> J. R. Dandoy,<sup>31</sup> N. P. Dang,<sup>48</sup> A. C. Daniells,<sup>18</sup>  
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C. David,<sup>169</sup> T. Davidek,<sup>129</sup> E. Davies,<sup>120,m</sup> M. Davies,<sup>153</sup> P. Davison,<sup>78</sup> Y. Davygora,<sup>58a</sup> E. Dawe,<sup>88</sup> I. Dawson,<sup>139</sup>  
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P. de Jong,<sup>107</sup> H. De la Torre,<sup>82</sup> F. De Lorenzi,<sup>64</sup> L. De Nooij,<sup>107</sup> D. De Pedis,<sup>132a</sup> A. De Salvo,<sup>132a</sup> U. De Sanctis,<sup>149</sup>  
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A. Dell'Acqua,<sup>30</sup> L. Dell'Asta,<sup>22</sup> M. Dell'Orso,<sup>124a,124b</sup> M. Della Pietra,<sup>104a,k</sup> D. della Volpe,<sup>49</sup> M. Delmastro,<sup>5</sup>  
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D. Derendarz,<sup>39</sup> J. E. Derkaoui,<sup>135d</sup> F. Derue,<sup>80</sup> P. Dervan,<sup>74</sup> K. Desch,<sup>21</sup> C. Deterre,<sup>42</sup> P. O. Deviveiros,<sup>30</sup> A. Dewhurst,<sup>131</sup>  
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B. Di Girolamo,<sup>30</sup> A. Di Mattia,<sup>152</sup> B. Di Micco,<sup>134a,134b</sup> R. Di Nardo,<sup>47</sup> A. Di Simone,<sup>48</sup> R. Di Sipio,<sup>158</sup> D. Di Valentino,<sup>29</sup>  
C. Diaconu,<sup>85</sup> M. Diamond,<sup>158</sup> F. A. Dias,<sup>46</sup> M. A. Diaz,<sup>32a</sup> E. B. Diehl,<sup>89</sup> J. Dietrich,<sup>16</sup> S. Diglio,<sup>85</sup> A. Dimitrievska,<sup>13</sup>  
J. Dingfelder,<sup>21</sup> P. Dita,<sup>26a</sup> S. Dita,<sup>26a</sup> F. Dittus,<sup>30</sup> F. Djama,<sup>85</sup> T. Djobava,<sup>51b</sup> J. I. Djuvsland,<sup>58a</sup> M. A. B. do Vale,<sup>24c</sup>  
D. Dobos,<sup>30</sup> M. Dobre,<sup>26a</sup> C. Doglioni,<sup>49</sup> T. Dohmae,<sup>155</sup> J. Dolejsi,<sup>129</sup> Z. Dolezal,<sup>129</sup> B. A. Dolgoshein,<sup>98,a</sup> M. Donadelli,<sup>24d</sup>  
S. Donati,<sup>124a,124b</sup> P. Dondero,<sup>121a,121b</sup> J. Donini,<sup>34</sup> J. Dopke,<sup>131</sup> A. Doria,<sup>104a</sup> M. T. Dova,<sup>71</sup> A. T. Doyle,<sup>53</sup> E. Drechsler,<sup>54</sup>  
M. Dris,<sup>10</sup> E. Dubreuil,<sup>34</sup> E. Duchovni,<sup>172</sup> G. Duckeck,<sup>100</sup> O. A. Ducu,<sup>26a,85</sup> D. Duda,<sup>175</sup> A. Dudarev,<sup>30</sup> L. Duflot,<sup>117</sup>  
L. Duguid,<sup>77</sup> M. Dührssen,<sup>30</sup> M. Dunford,<sup>58a</sup> H. Duran Yildiz,<sup>4a</sup> M. Düren,<sup>52</sup> A. Durglishvili,<sup>51b</sup> D. Duschinger,<sup>44</sup>  
M. Dyndal,<sup>38a</sup> C. Eckardt,<sup>42</sup> K. M. Ecker,<sup>101</sup> R. C. Edgar,<sup>89</sup> W. Edson,<sup>2</sup> N. C. Edwards,<sup>46</sup> W. Ehrenfeld,<sup>21</sup> T. Eifert,<sup>30</sup>  
G. Eigen,<sup>14</sup> K. Einsweiler,<sup>15</sup> T. Ekelof,<sup>166</sup> M. El Kacimi,<sup>135c</sup> M. Ellert,<sup>166</sup> S. Elles,<sup>5</sup> F. Ellinghaus,<sup>83</sup> A. A. Elliot,<sup>169</sup> N. Ellis,<sup>30</sup>  
J. Elmsheuser,<sup>100</sup> M. Elsing,<sup>30</sup> D. Emeliyanov,<sup>131</sup> Y. Enari,<sup>155</sup> O. C. Endner,<sup>83</sup> M. Endo,<sup>118</sup> J. Erdmann,<sup>43</sup> A. Ereditato,<sup>17</sup>  
G. Ernis,<sup>175</sup> J. Ernst,<sup>2</sup> M. Ernst,<sup>25</sup> S. Errede,<sup>165</sup> E. Ertel,<sup>83</sup> M. Escalier,<sup>117</sup> H. Esch,<sup>43</sup> C. Escobar,<sup>125</sup> B. Esposito,<sup>47</sup>  
A. I. Etiennevire,<sup>136</sup> E. Etzion,<sup>153</sup> H. Evans,<sup>61</sup> A. Ezhilov,<sup>123</sup> L. Fabbri,<sup>20a,20b</sup> G. Facini,<sup>31</sup> R. M. Fakhruddinov,<sup>130</sup> S. Falciano,<sup>132a</sup>  
R. J. Falla,<sup>78</sup> J. Faltova,<sup>129</sup> Y. Fang,<sup>33a</sup> M. Fanti,<sup>91a,91b</sup> A. Farbin,<sup>8</sup> A. Farilla,<sup>134a</sup> T. Farooque,<sup>12</sup> S. Farrell,<sup>15</sup>  
S. M. Farrington,<sup>170</sup> P. Farthouat,<sup>30</sup> F. Fassi,<sup>135e</sup> P. Fassnacht,<sup>30</sup> D. Fassouliotis,<sup>9</sup> M. Fauci Giannelli,<sup>77</sup> A. Favareto,<sup>50a,50b</sup>  
L. Fayard,<sup>117</sup> P. Federic,<sup>144a</sup> O. L. Fedin,<sup>123,n</sup> W. Fedorko,<sup>168</sup> S. Feigl,<sup>30</sup> L. Feligioni,<sup>85</sup> C. Feng,<sup>33d</sup> E. J. Feng,<sup>6</sup> H. Feng,<sup>89</sup>  
A. B. Fenjuk,<sup>130</sup> L. Feremenga,<sup>8</sup> P. Fernandez Martinez,<sup>167</sup> S. Fernandez Perez,<sup>30</sup> J. Ferrando,<sup>53</sup> A. Ferrari,<sup>166</sup> P. Ferrari,<sup>107</sup>

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 F. Fiedler,<sup>83</sup> A. Filipčič,<sup>75</sup> M. Filipuzzi,<sup>42</sup> F. Filthaut,<sup>106</sup> M. Fincke-Keeler,<sup>169</sup> K. D. Finelli,<sup>150</sup> M. C. N. Fiolhais,<sup>126a,126c</sup>  
 L. Fiorini,<sup>167</sup> A. Firan,<sup>40</sup> A. Fischer,<sup>2</sup> C. Fischer,<sup>12</sup> J. Fischer,<sup>175</sup> W. C. Fisher,<sup>90</sup> E. A. Fitzgerald,<sup>23</sup> I. Fleck,<sup>141</sup>  
 P. Fleischmann,<sup>89</sup> S. Fleischmann,<sup>175</sup> G. T. Fletcher,<sup>139</sup> G. Fletcher,<sup>76</sup> R. R. M. Fletcher,<sup>122</sup> T. Flick,<sup>175</sup> A. Floderus,<sup>81</sup>  
 L. R. Flores Castillo,<sup>60a</sup> M. J. Flowerdew,<sup>101</sup> A. Formica,<sup>136</sup> A. Forti,<sup>84</sup> D. Fournier,<sup>117</sup> H. Fox,<sup>72</sup> S. Fracchia,<sup>12</sup>  
 P. Francavilla,<sup>80</sup> M. Franchini,<sup>20a,20b</sup> D. Francis,<sup>30</sup> L. Franconi,<sup>119</sup> M. Franklin,<sup>57</sup> M. Frate,<sup>163</sup> M. Fraternali,<sup>121a,121b</sup>  
 D. Freeborn,<sup>78</sup> S. T. French,<sup>28</sup> F. Friedrich,<sup>44</sup> D. Froidevaux,<sup>30</sup> J. A. Frost,<sup>120</sup> C. Fukunaga,<sup>156</sup> E. Fullana Torregrosa,<sup>83</sup>  
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 S. Gadomski,<sup>49</sup> G. Gagliardi,<sup>50a,50b</sup> P. Gagnon,<sup>61</sup> C. Galea,<sup>106</sup> B. Galhardo,<sup>126a,126c</sup> E. J. Gallas,<sup>120</sup> B. J. Gallop,<sup>131</sup>  
 P. Gallus,<sup>128</sup> G. Galster,<sup>36</sup> K. K. Gan,<sup>111</sup> J. Gao,<sup>33b,85</sup> Y. Gao,<sup>46</sup> Y. S. Gao,<sup>143,f</sup> F. M. Garay Walls,<sup>46</sup> F. Garberson,<sup>176</sup>  
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 A. Gaudiello,<sup>50a,50b</sup> G. Gaudio,<sup>121a</sup> B. Gaur,<sup>141</sup> L. Gauthier,<sup>95</sup> P. Gauzzi,<sup>132a,132b</sup> I. L. Gavrilenko,<sup>96</sup> C. Gay,<sup>168</sup> G. Gaycken,<sup>21</sup>  
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 M. H. Genest,<sup>55</sup> S. Gentile,<sup>132a,132b</sup> M. George,<sup>54</sup> S. George,<sup>77</sup> D. Gerbaudo,<sup>163</sup> A. Gershon,<sup>153</sup> H. Ghazlane,<sup>135b</sup>  
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 T. P. S. Gillam,<sup>28</sup> D. Gillberg,<sup>30</sup> G. Gilles,<sup>34</sup> D. M. Gingrich,<sup>3,e</sup> N. Giokaris,<sup>9</sup> M. P. Giordani,<sup>164a,164c</sup> F. M. Giorgi,<sup>20a</sup>  
 F. M. Giorgi,<sup>16</sup> P. F. Giraud,<sup>136</sup> P. Giromini,<sup>47</sup> D. Giugni,<sup>91a</sup> C. Giuliani,<sup>48</sup> M. Giulini,<sup>58b</sup> B. K. Gjelsten,<sup>119</sup> S. Gkaitatzis,<sup>154</sup>  
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 M. Goblirsch-Kolb,<sup>101</sup> J. R. Goddard,<sup>76</sup> J. Godlewski,<sup>39</sup> S. Goldfarb,<sup>89</sup> T. Golling,<sup>49</sup> D. Golubkov,<sup>130</sup> A. Gomes,<sup>126a,126b,126d</sup>  
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 A. Gorišek,<sup>75</sup> E. Gornicki,<sup>39</sup> A. T. Goshaw,<sup>45</sup> C. Gössling,<sup>43</sup> M. I. Gostkin,<sup>65</sup> D. Goujdami,<sup>135c</sup> A. G. Goussiou,<sup>138</sup>  
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 S. Guindon,<sup>2</sup> U. Gul,<sup>53</sup> C. Gumpert,<sup>44</sup> J. Guo,<sup>33e</sup> S. Gupta,<sup>120</sup> G. Gustavino,<sup>132a,132b</sup> P. Gutierrez,<sup>113</sup> N. G. Gutierrez Ortiz,<sup>53</sup>  
 C. Gutschow,<sup>44</sup> C. Guyot,<sup>136</sup> C. Gwenlan,<sup>120</sup> C. B. Gwilliam,<sup>74</sup> A. Haas,<sup>110</sup> C. Haber,<sup>15</sup> H. K. Hadavand,<sup>8</sup> N. Haddad,<sup>135e</sup>  
 P. Haefner,<sup>21</sup> S. Hageböck,<sup>21</sup> Z. Hajduk,<sup>39</sup> H. Hakobyan,<sup>177</sup> M. Haleem,<sup>42</sup> J. Haley,<sup>114</sup> D. Hall,<sup>120</sup> G. Halladjian,<sup>90</sup>  
 G. D. Hallewell,<sup>85</sup> K. Hamacher,<sup>175</sup> P. Hamal,<sup>115</sup> K. Hamano,<sup>169</sup> M. Hamer,<sup>54</sup> A. Hamilton,<sup>145a</sup> G. N. Hamity,<sup>145c</sup>  
 P. G. Hamnett,<sup>42</sup> L. Han,<sup>33b</sup> K. Hanagaki,<sup>118</sup> K. Hanawa,<sup>155</sup> M. Hance,<sup>15</sup> P. Hanke,<sup>58a</sup> R. Hanna,<sup>136</sup> J. B. Hansen,<sup>36</sup>  
 J. D. Hansen,<sup>36</sup> M. C. Hansen,<sup>21</sup> P. H. Hansen,<sup>36</sup> K. Hara,<sup>160</sup> A. S. Hard,<sup>173</sup> T. Harenberg,<sup>175</sup> F. Hariri,<sup>117</sup> S. Harkusha,<sup>92</sup>  
 R. D. Harrington,<sup>46</sup> P. F. Harrison,<sup>170</sup> F. Hartjes,<sup>107</sup> M. Hasegawa,<sup>67</sup> S. Hasegawa,<sup>103</sup> Y. Hasegawa,<sup>140</sup> A. Hasib,<sup>113</sup>  
 S. Hassani,<sup>136</sup> S. Haug,<sup>17</sup> R. Hauser,<sup>90</sup> L. Hauswald,<sup>44</sup> M. Havranek,<sup>127</sup> C. M. Hawkes,<sup>18</sup> R. J. Hawkings,<sup>30</sup> A. D. Hawkins,<sup>81</sup>  
 T. Hayashi,<sup>160</sup> D. Hayden,<sup>90</sup> C. P. Hays,<sup>120</sup> J. M. Hays,<sup>76</sup> H. S. Hayward,<sup>74</sup> S. J. Haywood,<sup>131</sup> S. J. Head,<sup>18</sup> T. Heck,<sup>83</sup>  
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 S. Hellman,<sup>146a,146b</sup> D. Hellmich,<sup>21</sup> C. Helsens,<sup>30</sup> J. Henderson,<sup>120</sup> R. C. W. Henderson,<sup>72</sup> Y. Heng,<sup>173</sup> C. Hengler,<sup>42</sup>  
 A. Henrichs,<sup>176</sup> A. M. Henriques Correia,<sup>30</sup> S. Henrot-Versille,<sup>117</sup> G. H. Herbert,<sup>16</sup> Y. Hernández Jiménez,<sup>167</sup>  
 R. Herrberg-Schubert,<sup>16</sup> G. Herten,<sup>48</sup> R. Hertenberger,<sup>100</sup> L. Hervas,<sup>30</sup> G. G. Hesketh,<sup>78</sup> N. P. Hessey,<sup>107</sup> J. W. Hetherly,<sup>40</sup>  
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 A. Hoecker,<sup>30</sup> M. R. Hoferkamp,<sup>105</sup> F. Hoenig,<sup>100</sup> M. Hohlfield,<sup>83</sup> D. Hohn,<sup>21</sup> T. R. Holmes,<sup>15</sup> M. Homann,<sup>43</sup> T. M. Hong,<sup>125</sup>  
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 J. Howard,<sup>120</sup> J. Howarth,<sup>42</sup> M. Hrabovsky,<sup>115</sup> I. Hristova,<sup>16</sup> J. Hrivnac,<sup>117</sup> T. Hryn'ova,<sup>5</sup> A. Hrynevich,<sup>93</sup> C. Hsu,<sup>145c</sup>  
 P. J. Hsu,<sup>151,q</sup> S.-C. Hsu,<sup>138</sup> D. Hu,<sup>35</sup> Q. Hu,<sup>33b</sup> X. Hu,<sup>89</sup> Y. Huang,<sup>42</sup> Z. Hubacek,<sup>30</sup> F. Hubaut,<sup>85</sup> F. Huegging,<sup>21</sup>  
 T. B. Huffman,<sup>120</sup> E. W. Hughes,<sup>35</sup> G. Hughes,<sup>72</sup> M. Huhtinen,<sup>30</sup> T. A. Hülsing,<sup>83</sup> N. Huseynov,<sup>65,c</sup> J. Huston,<sup>90</sup> J. Huth,<sup>57</sup>  
 G. Iacobucci,<sup>49</sup> G. Iakovidis,<sup>25</sup> I. Ibragimov,<sup>141</sup> L. Iconomidou-Fayard,<sup>117</sup> E. Ideal,<sup>176</sup> Z. Idrissi,<sup>135e</sup> P. Iengo,<sup>30</sup>  
 O. Igonkina,<sup>107</sup> T. Iizawa,<sup>171</sup> Y. Ikegami,<sup>66</sup> K. Ikematsu,<sup>141</sup> M. Ikeno,<sup>66</sup> Y. Ichenko,<sup>31,r</sup> D. Iliadis,<sup>154</sup> N. Ilic,<sup>143</sup> Y. Inamaru,<sup>67</sup>

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Jeanty,<sup>15</sup> J. Jejelava,<sup>51a,s</sup> G.-Y. Jeng,<sup>150</sup> D. Jennens,<sup>88</sup> P. Jenni,<sup>48,t</sup> J. Jentzsch,<sup>43</sup> C. Jeske,<sup>170</sup> S. Jézéquel,<sup>5</sup> H. Ji,<sup>173</sup> J. Jia,<sup>148</sup> Y. Jiang,<sup>33b</sup> S. Jiggins,<sup>78</sup> J. Jimenez Pena,<sup>167</sup> S. Jin,<sup>33a</sup> A. Jinaru,<sup>26a</sup> O. Jinnouchi,<sup>157</sup> M. D. Joergensen,<sup>36</sup> P. Johansson,<sup>139</sup> K. A. Johns,<sup>7</sup> K. Jon-And,<sup>146a,146b</sup> G. Jones,<sup>170</sup> R. W. L. Jones,<sup>72</sup> T. J. Jones,<sup>74</sup> J. Jongmanns,<sup>58a</sup> P. M. Jorge,<sup>126a,126b</sup> K. D. Joshi,<sup>84</sup> J. Jovicevic,<sup>159a</sup> X. Ju,<sup>173</sup> C. A. Jung,<sup>43</sup> P. Jussel,<sup>62</sup> A. Juste Rozas,<sup>12,p</sup> M. Kaci,<sup>167</sup> A. Kaczmarska,<sup>39</sup> M. Kado,<sup>117</sup> H. Kagan,<sup>111</sup> M. Kagan,<sup>143</sup> S. J. Kahn,<sup>85</sup> E. Kajomovitz,<sup>45</sup> C. W. Kalderon,<sup>120</sup> S. Kama,<sup>40</sup> A. Kamenshchikov,<sup>130</sup> N. Kanaya,<sup>155</sup> M. Kaneda,<sup>30</sup> S. Kaneti,<sup>28</sup> V. A. Kantserov,<sup>98</sup> J. Kanzaki,<sup>66</sup> B. Kaplan,<sup>110</sup> A. Kapliy,<sup>31</sup> D. Kar,<sup>53</sup> K. Karakostas,<sup>10</sup> A. Karamaoun,<sup>3</sup> N. Karastathis,<sup>10,107</sup> M. J. Kareem,<sup>54</sup> M. Karneevskiy,<sup>83</sup> S. N. Karpov,<sup>65</sup> Z. M. Karpova,<sup>65</sup> K. Karthik,<sup>110</sup> V. Kartvelishvili,<sup>72</sup> A. N. Karyukhin,<sup>130</sup> L. Kashif,<sup>173</sup> R. D. Kass,<sup>111</sup> A. Kastanas,<sup>14</sup> Y. Kataoka,<sup>155</sup> A. Katre,<sup>49</sup> J. Katzy,<sup>42</sup> K. Kawagoe,<sup>70</sup> T. Kawamoto,<sup>155</sup> G. Kawamura,<sup>54</sup> S. Kazama,<sup>155</sup> V. F. Kazanin,<sup>109,d</sup> M. Y. Kazarinov,<sup>65</sup> R. Keeler,<sup>169</sup> R. Kehoe,<sup>40</sup> J. S. Keller,<sup>42</sup> J. J. Kempster,<sup>77</sup> H. Keoshkerian,<sup>84</sup> O. Kepka,<sup>127</sup> B. P. Kerševan,<sup>75</sup> S. Kersten,<sup>175</sup> R. A. Keyes,<sup>87</sup> F. Khalil-zada,<sup>11</sup> H. Khandanyan,<sup>146a,146b</sup> A. Khanov,<sup>114</sup> A. G. Kharlamov,<sup>109,d</sup> T. J. Khoo,<sup>28</sup> V. Khovanskiy,<sup>97</sup> E. Khramov,<sup>65</sup> J. Khubua,<sup>51b,u</sup> H. Y. Kim,<sup>8</sup> H. Kim,<sup>146a,146b</sup> S. H. Kim,<sup>160</sup> Y. Kim,<sup>31</sup> N. Kimura,<sup>154</sup> O. M. Kind,<sup>16</sup> B. T. King,<sup>74</sup> M. King,<sup>167</sup> S. B. King,<sup>168</sup> J. Kirk,<sup>131</sup> A. E. Kiryunin,<sup>101</sup> T. Kishimoto,<sup>67</sup> D. Kisielewska,<sup>38a</sup> F. Kiss,<sup>48</sup> K. Kiuchi,<sup>160</sup> O. Kivernyk,<sup>136</sup> E. Kladiva,<sup>144b</sup> M. H. Klein,<sup>35</sup> M. Klein,<sup>74</sup> U. Klein,<sup>74</sup> K. Kleinknecht,<sup>83</sup> P. Klimek,<sup>146a,146b</sup> A. Klimentov,<sup>25</sup> R. Klingenberg,<sup>43</sup> J. A. 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Kordas,<sup>154</sup> A. Korn,<sup>78</sup> A. A. Korol,<sup>109,d</sup> I. Korolkov,<sup>12</sup> E. V. Korolkova,<sup>139</sup> O. Kortner,<sup>101</sup> S. Kortner,<sup>101</sup> T. Kosek,<sup>129</sup> V. V. Kostyukhin,<sup>21</sup> V. M. Kotov,<sup>65</sup> A. Kotwal,<sup>45</sup> A. Kourkoumeli-Charalampidi,<sup>154</sup> C. Kourkoumelis,<sup>9</sup> V. Kouskoura,<sup>25</sup> A. Koutsman,<sup>159a</sup> R. Kowalewski,<sup>169</sup> T. Z. Kowalski,<sup>38a</sup> W. Kozanecki,<sup>136</sup> A. S. Kozhin,<sup>130</sup> V. A. Kramarenko,<sup>99</sup> G. Kramberger,<sup>75</sup> D. Krasnopevtsev,<sup>98</sup> M. W. Krasny,<sup>80</sup> A. Krasznahorkay,<sup>30</sup> J. K. Kraus,<sup>21</sup> A. Kravchenko,<sup>25</sup> S. Kreiss,<sup>110</sup> M. Kretz,<sup>58c</sup> J. Kretzschmar,<sup>74</sup> K. Kreutzfeldt,<sup>52</sup> P. Krieger,<sup>158</sup> K. Krizka,<sup>31</sup> K. Kroeninger,<sup>43</sup> H. Kroha,<sup>101</sup> J. Kroll,<sup>122</sup> J. Kroseberg,<sup>21</sup> J. 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Manning,<sup>137</sup> A. Manousakis-Katsikakis,<sup>9</sup> B. Mansoulie,<sup>136</sup> R. Mantifel,<sup>87</sup> M. Mantoani,<sup>54</sup> L. Mapelli,<sup>30</sup> L. March,<sup>145c</sup> G. Marchiori,<sup>80</sup> M. Marcisovsky,<sup>127</sup> C. P. Marino,<sup>169</sup> M. Marjanovic,<sup>13</sup> D. E. Marley,<sup>89</sup> F. Marroquim,<sup>24a</sup> S. P. Marsden,<sup>84</sup> Z. Marshall,<sup>15</sup> L. F. Marti,<sup>17</sup> S. Marti-Garcia,<sup>167</sup> B. Martin,<sup>90</sup> T. A. Martin,<sup>170</sup> V. J. Martin,<sup>46</sup> B. Martin dit Latour,<sup>14</sup> M. Martinez,<sup>12,p</sup> S. Martin-Haugh,<sup>131</sup> V. S. Martoiu,<sup>26a</sup> A. C. Martyniuk,<sup>78</sup> M. Marx,<sup>138</sup> F. Marzano,<sup>132a</sup> A. Marzin,<sup>30</sup> L. Masetti,<sup>83</sup> T. Mashimo,<sup>155</sup> R. Mashinistov,<sup>96</sup> J. Masik,<sup>84</sup> A. L. Maslennikov,<sup>109,d</sup> I. Massa,<sup>20a,20b</sup> L. Massa,<sup>20a,20b</sup> N. Massol,<sup>5</sup> P. Mastrandrea,<sup>148</sup> A. Mastroberardino,<sup>37a,37b</sup> T. Masubuchi,<sup>155</sup> P. Mättig,<sup>175</sup> J. Mattmann,<sup>83</sup> J. Maurer,<sup>26a</sup> S. J. Maxfield,<sup>74</sup> D. A. Maximov,<sup>109,d</sup> R. Mazini,<sup>151</sup> S. M. Mazza,<sup>91a,91b</sup> L. Mazzaferro,<sup>133a,133b</sup> G. Mc Goldrick,<sup>158</sup> S. P. Mc Kee,<sup>89</sup> A. McCarn,<sup>89</sup> R. L. McCarthy,<sup>148</sup> T. G. McCarthy,<sup>29</sup> N. A. McCubbin,<sup>131</sup> K. W. McFarlane,<sup>56,a</sup> J. A. Mcfayden,<sup>78</sup> G. Mchedlidze,<sup>54</sup> S. J. McMahon,<sup>131</sup> R. A. McPherson,<sup>169,l</sup> M. Medinnis,<sup>42</sup> S. Meehan,<sup>145a</sup> S. Mehlhase,<sup>100</sup> A. Mehta,<sup>74</sup> K. Meier,<sup>58a</sup> C. Meineck,<sup>100</sup> B. Meirose,<sup>41</sup> B. R. Mellado Garcia,<sup>145c</sup> F. Meloni,<sup>17</sup> A. Mengarelli,<sup>20a,20b</sup> S. Menke,<sup>101</sup> E. Meoni,<sup>161</sup> K. M. Mercurio,<sup>57</sup> S. 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P. Urquijo,<sup>88</sup> P. Urrejola,<sup>83</sup> G. Usai,<sup>8</sup> A. Usanova,<sup>62</sup> L. Vacavant,<sup>85</sup> V. Vacek,<sup>128</sup> B. Vachon,<sup>87</sup> C. Valderanis,<sup>83</sup> N. Valencic,<sup>107</sup>  
S. Valentinetti,<sup>20a,20b</sup> A. Valero,<sup>167</sup> L. Valery,<sup>12</sup> S. Valkar,<sup>129</sup> E. Valladolid Gallego,<sup>167</sup> S. Vallecorsa,<sup>49</sup> J. A. Valls Ferrer,<sup>167</sup>

W. Van Den Wollenberg,<sup>107</sup> P. C. Van Der Deijl,<sup>107</sup> R. van der Geer,<sup>107</sup> H. van der Graaf,<sup>107</sup> R. Van Der Leeuw,<sup>107</sup> N. van Eldik,<sup>152</sup> P. van Gemmeren,<sup>6</sup> J. Van Nieuwkoop,<sup>142</sup> I. van Vulpen,<sup>107</sup> M. C. van Woerden,<sup>30</sup> M. Vanadia,<sup>132a,132b</sup> W. Vandelli,<sup>30</sup> R. Vanguri,<sup>122</sup> A. Vaniachine,<sup>6</sup> F. Vannucci,<sup>80</sup> G. Vardanyan,<sup>177</sup> R. Vari,<sup>132a</sup> E. W. Varnes,<sup>7</sup> T. Varol,<sup>40</sup> D. Varouchas,<sup>80</sup> A. Vartapetian,<sup>8</sup> K. E. Varvell,<sup>150</sup> F. Vazeille,<sup>34</sup> T. Vazquez Schroeder,<sup>87</sup> J. Veatch,<sup>7</sup> L. M. Veloce,<sup>158</sup> F. Veloso,<sup>126a,126c</sup> T. Velz,<sup>21</sup> S. Veneziano,<sup>132a</sup> A. Ventura,<sup>73a,73b</sup> D. Ventura,<sup>86</sup> M. Venturi,<sup>169</sup> N. Venturi,<sup>158</sup> A. Venturini,<sup>23</sup> V. Vercesi,<sup>121a</sup> M. Verducci,<sup>132a,132b</sup> W. Verkerke,<sup>107</sup> J. C. Vermeulen,<sup>107</sup> A. Vest,<sup>44</sup> M. C. Vetterli,<sup>142,e</sup> O. Viazlo,<sup>81</sup> I. Vichou,<sup>165</sup> T. Vickey,<sup>139</sup> O. E. Vickey Boeriu,<sup>139</sup> G. H. A. Viehhauser,<sup>120</sup> S. Viel,<sup>15</sup> R. Vigne,<sup>62</sup> M. Villa,<sup>20a,20b</sup> M. Villaplana Perez,<sup>91a,91b</sup> E. Vilucchi,<sup>47</sup> M. G. Vincter,<sup>29</sup> V. B. Vinogradov,<sup>65</sup> I. Vivarelli,<sup>149</sup> F. Vives Vaque,<sup>3</sup> S. Vlachos,<sup>10</sup> D. Vladoiu,<sup>100</sup> M. Vlasak,<sup>128</sup> M. Vogel,<sup>32a</sup> P. Vokac,<sup>128</sup> G. Volpi,<sup>124a,124b</sup> M. Volpi,<sup>88</sup> H. von der Schmitt,<sup>101</sup> H. von Radziewski,<sup>48</sup> E. von Toerne,<sup>21</sup> V. Vorobel,<sup>129</sup> K. Vorobev,<sup>98</sup> M. Vos,<sup>167</sup> R. Voss,<sup>30</sup> J. H. Vosseveld,<sup>74</sup> N. Vranjes,<sup>13</sup> M. Vranjes Milosavljevic,<sup>13</sup> V. Vrba,<sup>127</sup> M. Vreeswijk,<sup>107</sup> R. Vuillermet,<sup>30</sup> I. Vukotic,<sup>31</sup> Z. Vykydal,<sup>128</sup> P. Wagner,<sup>21</sup> W. Wagner,<sup>175</sup> H. Wahlberg,<sup>71</sup> S. Wahrmund,<sup>44</sup> J. Wakabayashi,<sup>103</sup> J. Walder,<sup>72</sup> R. Walker,<sup>100</sup> W. Walkowiak,<sup>141</sup> C. Wang,<sup>151</sup> F. Wang,<sup>173</sup> H. Wang,<sup>15</sup> H. Wang,<sup>40</sup> J. Wang,<sup>42</sup> J. Wang,<sup>33a</sup> K. Wang,<sup>87</sup> R. Wang,<sup>6</sup> S. M. Wang,<sup>151</sup> T. Wang,<sup>21</sup> X. Wang,<sup>176</sup> C. Wanotayaroj,<sup>116</sup> A. Warburton,<sup>87</sup> C. P. Ward,<sup>28</sup> D. R. Wardrope,<sup>78</sup> M. Warsinsky,<sup>48</sup> A. Washbrook,<sup>46</sup> C. Wasicki,<sup>42</sup> P. M. Watkins,<sup>18</sup> A. T. Watson,<sup>18</sup> I. J. Watson,<sup>150</sup> M. F. Watson,<sup>18</sup> G. Watts,<sup>138</sup> S. Watts,<sup>84</sup> B. M. Waugh,<sup>78</sup> S. Webb,<sup>84</sup> M. S. Weber,<sup>17</sup> S. W. Weber,<sup>174</sup> J. S. Webster,<sup>31</sup> A. R. Weidberg,<sup>120</sup> B. Weinert,<sup>61</sup> J. Weingarten,<sup>54</sup> C. Weiser,<sup>48</sup> H. Weits,<sup>107</sup> P. S. Wells,<sup>30</sup> T. Wenaus,<sup>25</sup> T. Wengler,<sup>30</sup> S. Wenig,<sup>30</sup> N. Wermes,<sup>21</sup> M. Werner,<sup>48</sup> P. Werner,<sup>30</sup> M. Wessels,<sup>58a</sup> J. Wetter,<sup>161</sup> K. Whalen,<sup>116</sup> A. M. Wharton,<sup>72</sup> A. White,<sup>8</sup> M. J. White,<sup>1</sup> R. White,<sup>32b</sup> S. White,<sup>124a,124b</sup> D. Whiteson,<sup>163</sup> F. J. Wickens,<sup>131</sup> W. Wiedenmann,<sup>173</sup> M. Wielers,<sup>131</sup> P. Wienemann,<sup>21</sup> C. Wigglesworth,<sup>36</sup> L. A. M. Wiik-Fuchs,<sup>21</sup> A. Wildauer,<sup>101</sup> H. G. Wilkens,<sup>30</sup> H. H. Williams,<sup>122</sup> S. Williams,<sup>107</sup> C. Willis,<sup>90</sup> S. Willocq,<sup>86</sup> A. Wilson,<sup>89</sup> J. A. Wilson,<sup>18</sup> I. Wingerter-Seez,<sup>5</sup> F. Winklmeier,<sup>116</sup> B. T. Winter,<sup>21</sup> M. Wittgen,<sup>143</sup> J. Wittkowski,<sup>100</sup> S. J. Wollstadt,<sup>83</sup> M. W. Wolter,<sup>39</sup> H. Wolters,<sup>126a,126c</sup> B. K. Wosiek,<sup>39</sup> J. Wotschack,<sup>30</sup> M. J. Woudstra,<sup>84</sup> K. W. Wozniak,<sup>39</sup> M. Wu,<sup>55</sup> M. Wu,<sup>31</sup> S. L. Wu,<sup>173</sup> X. Wu,<sup>49</sup> Y. Wu,<sup>89</sup> T. R. Wyatt,<sup>84</sup> B. M. Wynne,<sup>46</sup> S. Xella,<sup>36</sup> D. Xu,<sup>33a</sup> L. Xu,<sup>33b,11</sup> B. Yabsley,<sup>150</sup> S. Yacoob,<sup>145b,mm</sup> R. Yakabe,<sup>67</sup> M. Yamada,<sup>66</sup> Y. Yamaguchi,<sup>118</sup> A. Yamamoto,<sup>66</sup> S. Yamamoto,<sup>155</sup> T. Yamanaka,<sup>155</sup> K. Yamauchi,<sup>103</sup> Y. Yamazaki,<sup>67</sup> Z. Yan,<sup>22</sup> H. Yang,<sup>33e</sup> H. Yang,<sup>173</sup> Y. Yang,<sup>151</sup> W-M. Yao,<sup>15</sup> Y. Yasu,<sup>66</sup> E. Yatsenko,<sup>5</sup> K. H. Yau Wong,<sup>21</sup> J. Ye,<sup>40</sup> S. Ye,<sup>25</sup> I. Yeletsikh,<sup>65</sup> A. L. Yen,<sup>57</sup> E. Yildirim,<sup>42</sup> K. Yorita,<sup>171</sup> R. Yoshida,<sup>6</sup> K. Yoshihara,<sup>122</sup> C. Young,<sup>143</sup> C. J. S. Young,<sup>30</sup> S. Youssef,<sup>22</sup> D. R. Yu,<sup>15</sup> J. Yu,<sup>8</sup> J. M. Yu,<sup>89</sup> J. Yu,<sup>114</sup> L. Yuan,<sup>67</sup> A. Yurkewicz,<sup>108</sup> I. Yusuff,<sup>28,nn</sup> B. Zabinski,<sup>39</sup> R. Zaidan,<sup>63</sup> A. M. Zaitsev,<sup>130,cc</sup> J. Zalieckas,<sup>14</sup> A. Zaman,<sup>148</sup> S. Zambito,<sup>57</sup> L. Zanello,<sup>132a,132b</sup> D. Zanzi,<sup>88</sup> C. Zeitnitz,<sup>175</sup> M. Zeman,<sup>128</sup> A. Zemla,<sup>38a</sup> K. Zengel,<sup>23</sup> O. Zenin,<sup>130</sup> T. Ženiš,<sup>144a</sup> D. Zerwas,<sup>117</sup> D. Zhang,<sup>89</sup> F. Zhang,<sup>173</sup> H. Zhang,<sup>33c</sup> J. Zhang,<sup>6</sup> L. Zhang,<sup>48</sup> R. Zhang,<sup>33b</sup> X. Zhang,<sup>33d</sup> Z. Zhang,<sup>117</sup> X. Zhao,<sup>40</sup> Y. Zhao,<sup>33d,117</sup> Z. Zhao,<sup>33b</sup> A. Zhemchugov,<sup>65</sup> J. Zhong,<sup>120</sup> B. Zhou,<sup>89</sup> C. Zhou,<sup>45</sup> L. Zhou,<sup>35</sup> L. Zhou,<sup>40</sup> N. Zhou,<sup>163</sup> C. G. Zhu,<sup>33d</sup> H. Zhu,<sup>33a</sup> J. Zhu,<sup>89</sup> Y. Zhu,<sup>33b</sup> X. Zhuang,<sup>33a</sup> K. Zhukov,<sup>96</sup> A. Zibell,<sup>174</sup> D. Zieminska,<sup>61</sup> N. I. Zimine,<sup>65</sup> C. Zimmermann,<sup>83</sup> S. Zimmermann,<sup>48</sup> Z. Zinonos,<sup>54</sup> M. Zinser,<sup>83</sup> M. Ziolkowski,<sup>141</sup> L. Živković,<sup>13</sup> G. Zobernig,<sup>173</sup> A. Zoccoli,<sup>20a,20b</sup> M. zur Nedden,<sup>16</sup> G. Zurzolo,<sup>104a,104b</sup> and L. Zwalinski<sup>30</sup>

(ATLAS Collaboration)

<sup>1</sup>Department of Physics, University of Adelaide, Adelaide, Australia

<sup>2</sup>Physics Department, SUNY Albany, Albany New York, USA

<sup>3</sup>Department of Physics, University of Alberta, Edmonton Alberta, Canada

<sup>4a</sup>Department of Physics, Ankara University, Ankara, Turkey

<sup>4b</sup>Istanbul Aydin University, Istanbul, Turkey

<sup>4c</sup>Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey

<sup>5</sup>LAPP, CNRS/IN2P3 and Université Savoie Mont Blanc, Annecy-le-Vieux, France

<sup>6</sup>High Energy Physics Division, Argonne National Laboratory, Argonne Illinois, USA

<sup>7</sup>Department of Physics, University of Arizona, Tucson Arizona, USA

<sup>8</sup>Department of Physics, The University of Texas at Arlington, Arlington Texas, USA

<sup>9</sup>Physics Department, University of Athens, Athens, Greece

<sup>10</sup>Physics Department, National Technical University of Athens, Zografou, Greece

<sup>11</sup>Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan

<sup>12</sup>Institut de Física d'Altes Energies and Departament de Física de la Universitat Autònoma de Barcelona, Barcelona, Spain



- <sup>13</sup>*Institute of Physics, University of Belgrade, Belgrade, Serbia*
- <sup>14</sup>*Department for Physics and Technology, University of Bergen, Bergen, Norway*
- <sup>15</sup>*Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley California, USA*
- <sup>16</sup>*Department of Physics, Humboldt University, Berlin, Germany*
- <sup>17</sup>*Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland*
- <sup>18</sup>*School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom*
- <sup>19a</sup>*Department of Physics, Bogazici University, Istanbul, Turkey*
- <sup>19b</sup>*Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey*
- <sup>19c</sup>*Department of Physics, Dogus University, Istanbul, Turkey*
- <sup>20a</sup>*INFN Sezione di Bologna, Italy*
- <sup>20b</sup>*Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna, Italy*
- <sup>21</sup>*Physikalisches Institut, University of Bonn, Bonn, Germany*
- <sup>22</sup>*Department of Physics, Boston University, Boston Massachusetts, USA*
- <sup>23</sup>*Department of Physics, Brandeis University, Waltham Massachusetts, USA*
- <sup>24a</sup>*Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro, Brazil*
- <sup>24b</sup>*Electrical Circuits Department, Federal University of Juiz de Fora (UFJF), Juiz de Fora, Brazil*
- <sup>24c</sup>*Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei, Brazil*
- <sup>24d</sup>*Instituto de Fisica, Universidade de Sao Paulo, Sao Paulo, Brazil*
- <sup>25</sup>*Physics Department, Brookhaven National Laboratory, Upton New York, USA*
- <sup>26a</sup>*National Institute of Physics and Nuclear Engineering, Bucharest, Romania*
- <sup>26b</sup>*National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj Napoca, Romania*
- <sup>26c</sup>*University Politehnica Bucharest, Bucharest, Romania*
- <sup>26d</sup>*West University in Timisoara, Timisoara, Romania*
- <sup>27</sup>*Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina*
- <sup>28</sup>*Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom*
- <sup>29</sup>*Department of Physics, Carleton University, Ottawa ON, Canada*
- <sup>30</sup>*CERN, Geneva, Switzerland*
- <sup>31</sup>*Enrico Fermi Institute, University of Chicago, Chicago Illinois, USA*
- <sup>32a</sup>*Departamento de Física, Pontificia Universidad Católica de Chile, Santiago, Chile*
- <sup>32b</sup>*Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile*
- <sup>33a</sup>*Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China*
- <sup>33b</sup>*Department of Modern Physics, University of Science and Technology of China, Anhui, China*
- <sup>33c</sup>*Department of Physics, Nanjing University, Jiangsu, China*
- <sup>33d</sup>*School of Physics, Shandong University, Shandong, China*
- <sup>33e</sup>*Department of Physics and Astronomy, Shanghai Key Laboratory for Particle Physics and Cosmology, Shanghai Jiao Tong University, Shanghai, China*
- <sup>33f</sup>*Physics Department, Tsinghua University, Beijing 100084, China*
- <sup>34</sup>*Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Clermont-Ferrand, France*
- <sup>35</sup>*Nevis Laboratory, Columbia University, Irvington New York, USA*
- <sup>36</sup>*Niels Bohr Institute, University of Copenhagen, Kobenhavn, Denmark*
- <sup>37a</sup>*INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati, Italy*
- <sup>37b</sup>*Dipartimento di Fisica, Università della Calabria, Rende, Italy*
- <sup>38a</sup>*AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland*
- <sup>38b</sup>*Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland*
- <sup>39</sup>*Institute of Nuclear Physics Polish Academy of Sciences, Krakow, Poland*
- <sup>40</sup>*Physics Department, Southern Methodist University, Dallas Texas, USA*
- <sup>41</sup>*Physics Department, University of Texas at Dallas, Richardson Texas, USA*
- <sup>42</sup>*DESY, Hamburg and Zeuthen, Germany*
- <sup>43</sup>*Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany*
- <sup>44</sup>*Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany*
- <sup>45</sup>*Department of Physics, Duke University, Durham NC, USA*
- <sup>46</sup>*SUPA - School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom*
- <sup>47</sup>*INFN Laboratori Nazionali di Frascati, Frascati, Italy*
- <sup>48</sup>*Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany*
- <sup>49</sup>*Section de Physique, Université de Genève, Geneva, Switzerland*

- <sup>50a</sup>*INFN Sezione di Genova, Italy*
- <sup>50b</sup>*Dipartimento di Fisica, Università di Genova, Genova, Italy*
- <sup>51a</sup>*E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi, Georgia*
- <sup>51b</sup>*High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia*
- <sup>52</sup>*II Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany*
- <sup>53</sup>*SUPA - School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom*
- <sup>54</sup>*II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany*
- <sup>55</sup>*Laboratoire de Physique Subatomique et de Cosmologie, Université Grenoble-Alpes, CNRS/IN2P3, Grenoble, France*
- <sup>56</sup>*Department of Physics, Hampton University, Hampton VA, USA*
- <sup>57</sup>*Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge Massachusetts, USA*
- <sup>58a</sup>*Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany*
- <sup>58b</sup>*Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany*
- <sup>58c</sup>*ZITI Institut für technische Informatik, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany*
- <sup>59</sup>*Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan*
- <sup>60a</sup>*Department of Physics, The Chinese University of Hong Kong, Shatin, N.T., Hong Kong, China*
- <sup>60b</sup>*Department of Physics, The University of Hong Kong, Hong Kong, China*
- <sup>60c</sup>*Department of Physics, The Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China*
- <sup>61</sup>*Department of Physics, Indiana University, Bloomington Indiana, USA*
- <sup>62</sup>*Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria*
- <sup>63</sup>*University of Iowa, Iowa City IA, USA*
- <sup>64</sup>*Department of Physics and Astronomy, Iowa State University, Ames Iowa, USA*
- <sup>65</sup>*Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia*
- <sup>66</sup>*KEK, High Energy Accelerator Research Organization, Tsukuba, Japan*
- <sup>67</sup>*Graduate School of Science, Kobe University, Kobe, Japan*
- <sup>68</sup>*Faculty of Science, Kyoto University, Kyoto, Japan*
- <sup>69</sup>*Kyoto University of Education, Kyoto, Japan*
- <sup>70</sup>*Department of Physics, Kyushu University, Fukuoka, Japan*
- <sup>71</sup>*Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina*
- <sup>72</sup>*Physics Department, Lancaster University, Lancaster, United Kingdom*
- <sup>73a</sup>*INFN Sezione di Lecce, Italy*
- <sup>73b</sup>*Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy*
- <sup>74</sup>*Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom*
- <sup>75</sup>*Department of Physics, Jožef Stefan Institute and University of Ljubljana, Ljubljana, Slovenia*
- <sup>76</sup>*School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom*
- <sup>77</sup>*Department of Physics, Royal Holloway University of London, Surrey, United Kingdom*
- <sup>78</sup>*Department of Physics and Astronomy, University College London, London, United Kingdom*
- <sup>79</sup>*Louisiana Tech University, Ruston LA, USA*
- <sup>80</sup>*Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France*
- <sup>81</sup>*Fysiska institutionen, Lunds universitet, Lund, Sweden*
- <sup>82</sup>*Departamento de Física Teórica C-15, Universidad Autónoma de Madrid, Madrid, Spain*
- <sup>83</sup>*Institut für Physik, Universität Mainz, Mainz, Germany*
- <sup>84</sup>*School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom*
- <sup>85</sup>*CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France*
- <sup>86</sup>*Department of Physics, University of Massachusetts, Amherst Massachusetts, USA*
- <sup>87</sup>*Department of Physics, McGill University, Montreal QC, Canada*
- <sup>88</sup>*School of Physics, University of Melbourne, Victoria, Australia*
- <sup>89</sup>*Department of Physics, The University of Michigan, Ann Arbor Michigan, USA*
- <sup>90</sup>*Department of Physics and Astronomy, Michigan State University, East Lansing Michigan, USA*
- <sup>91a</sup>*INFN Sezione di Milano, Italy*
- <sup>91b</sup>*Dipartimento di Fisica, Università di Milano, Milano, Italy*
- <sup>92</sup>*B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Republic of Belarus*
- <sup>93</sup>*National Scientific and Educational Centre for Particle and High Energy Physics, Minsk, Republic of Belarus*
- <sup>94</sup>*Department of Physics, Massachusetts Institute of Technology, Cambridge Massachusetts, USA*
- <sup>95</sup>*Group of Particle Physics, University of Montreal, Montreal QC, Canada*
- <sup>96</sup>*P.N. Lebedev Institute of Physics, Academy of Sciences, Moscow, Russia*
- <sup>97</sup>*Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia*

- <sup>98</sup>*National Research Nuclear University MEPhI, Moscow, Russia*
- <sup>99</sup>*D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia*
- <sup>100</sup>*Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany*
- <sup>101</sup>*Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany*
- <sup>102</sup>*Nagasaki Institute of Applied Science, Nagasaki, Japan*
- <sup>103</sup>*Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan*
- <sup>104a</sup>*INFN Sezione di Napoli, Italy*
- <sup>104b</sup>*Dipartimento di Fisica, Università di Napoli, Napoli, Italy*
- <sup>105</sup>*Department of Physics and Astronomy, University of New Mexico, Albuquerque New Mexico, USA*
- <sup>106</sup>*Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands*
- <sup>107</sup>*Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands*
- <sup>108</sup>*Department of Physics, Northern Illinois University, DeKalb Illinois, USA*
- <sup>109</sup>*Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia*
- <sup>110</sup>*Department of Physics, New York University, New York New York, USA*
- <sup>111</sup>*Ohio State University, Columbus Ohio, USA*
- <sup>112</sup>*Faculty of Science, Okayama University, Okayama, Japan*
- <sup>113</sup>*Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman Oklahoma, USA*
- <sup>114</sup>*Department of Physics, Oklahoma State University, Stillwater Oklahoma, USA*
- <sup>115</sup>*Palacký University, RCPTM, Olomouc, Czech Republic*
- <sup>116</sup>*Center for High Energy Physics, University of Oregon, Eugene OR, USA*
- <sup>117</sup>*LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France*
- <sup>118</sup>*Graduate School of Science, Osaka University, Osaka, Japan*
- <sup>119</sup>*Department of Physics, University of Oslo, Oslo, Norway*
- <sup>120</sup>*Department of Physics, Oxford University, Oxford, United Kingdom*
- <sup>121a</sup>*INFN Sezione di Pavia, Italy*
- <sup>121b</sup>*Dipartimento di Fisica, Università di Pavia, Pavia, Italy*
- <sup>122</sup>*Department of Physics, University of Pennsylvania, Philadelphia Pennsylvania, USA*
- <sup>123</sup>*National Research Centre “Kurchatov Institute” B.P.Konstantinov Petersburg Nuclear Physics Institute, St. Petersburg, Russia*
- <sup>124a</sup>*INFN Sezione di Pisa, Italy*
- <sup>124b</sup>*Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy*
- <sup>125</sup>*Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh Pennsylvania, USA*
- <sup>126a</sup>*Laboratório de Instrumentação e Física Experimental de Partículas - LIP, Lisboa, Portugal*
- <sup>126b</sup>*Faculdade de Ciências, Universidade de Lisboa, Lisboa, Portugal*
- <sup>126c</sup>*Department of Physics, University of Coimbra, Coimbra, Portugal*
- <sup>126d</sup>*Centro de Física Nuclear da Universidade de Lisboa, Lisboa, Portugal*
- <sup>126e</sup>*Departamento de Física, Universidade do Minho, Braga, Portugal*
- <sup>126f</sup>*Departamento de Física Teórica y del Cosmos and CAFPE, Universidad de Granada, Granada, Spain*
- <sup>126g</sup>*Dep Física and CEFITEC of Faculdade de Ciências e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal*
- <sup>127</sup>*Institute of Physics, Academy of Sciences of the Czech Republic, Praha, Czech Republic*
- <sup>128</sup>*Czech Technical University in Prague, Praha, Czech Republic*
- <sup>129</sup>*Faculty of Mathematics and Physics, Charles University in Prague, Praha, Czech Republic*
- <sup>130</sup>*State Research Center Institute for High Energy Physics, Protvino, Russia*
- <sup>131</sup>*Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom*
- <sup>132a</sup>*INFN Sezione di Roma, Italy*
- <sup>132b</sup>*Dipartimento di Fisica, Sapienza Università di Roma, Roma, Italy*
- <sup>133a</sup>*INFN Sezione di Roma Tor Vergata, Italy*
- <sup>133b</sup>*Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy*
- <sup>134a</sup>*INFN Sezione di Roma Tre, Italy*
- <sup>134b</sup>*Dipartimento di Matematica e Fisica, Università Roma Tre, Roma, Italy*
- <sup>135a</sup>*Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies-Université Hassan II, Casablanca, Morocco*
- <sup>135b</sup>*Centre National de l’Energie des Sciences Techniques Nucleaires, Rabat, Morocco*
- <sup>135c</sup>*Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech, Morocco*
- <sup>135d</sup>*Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda, Morocco*
- <sup>135e</sup>*Faculté des sciences, Université Mohammed V-Agdal, Rabat, Morocco*

- <sup>136</sup>*DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l'Univers), CEA Saclay (Commissariat à l'Energie Atomique et aux Energies Alternatives), Gif-sur-Yvette, France*
- <sup>137</sup>*Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz California, USA*
- <sup>138</sup>*Department of Physics, University of Washington, Seattle Washington, USA*
- <sup>139</sup>*Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom*
- <sup>140</sup>*Department of Physics, Shinshu University, Nagano, Japan*
- <sup>141</sup>*Fachbereich Physik, Universität Siegen, Siegen, Germany*
- <sup>142</sup>*Department of Physics, Simon Fraser University, Burnaby BC, Canada*
- <sup>143</sup>*SLAC National Accelerator Laboratory, Stanford California, USA*
- <sup>144a</sup>*Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava, Slovak Republic*
- <sup>144b</sup>*Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic*
- <sup>145a</sup>*Department of Physics, University of Cape Town, Cape Town, South Africa*
- <sup>145b</sup>*Department of Physics, University of Johannesburg, Johannesburg, South Africa*
- <sup>145c</sup>*School of Physics, University of the Witwatersrand, Johannesburg, South Africa*
- <sup>146a</sup>*Department of Physics, Stockholm University, Sweden*
- <sup>146b</sup>*The Oskar Klein Centre, Stockholm, Sweden*
- <sup>147</sup>*Physics Department, Royal Institute of Technology, Stockholm, Sweden*
- <sup>148</sup>*Departments of Physics & Astronomy and Chemistry, Stony Brook University, Stony Brook New York, USA*
- <sup>149</sup>*Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom*
- <sup>150</sup>*School of Physics, University of Sydney, Sydney, Australia*
- <sup>151</sup>*Institute of Physics, Academia Sinica, Taipei, Taiwan*
- <sup>152</sup>*Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel*
- <sup>153</sup>*Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel*
- <sup>154</sup>*Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece*
- <sup>155</sup>*International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan*
- <sup>156</sup>*Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan*
- <sup>157</sup>*Department of Physics, Tokyo Institute of Technology, Tokyo, Japan*
- <sup>158</sup>*Department of Physics, University of Toronto, Toronto ON, Canada*
- <sup>159a</sup>*TRIUMF, Vancouver BC, Canada*
- <sup>159b</sup>*Department of Physics and Astronomy, York University, Toronto ON, Canada*
- <sup>160</sup>*Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan*
- <sup>161</sup>*Department of Physics and Astronomy, Tufts University, Medford Massachusetts, USA*
- <sup>162</sup>*Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia*
- <sup>163</sup>*Department of Physics and Astronomy, University of California Irvine, Irvine California, USA*
- <sup>164a</sup>*INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine, Italy*
- <sup>164b</sup>*ICTP, Trieste, Italy*
- <sup>164c</sup>*Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy*
- <sup>165</sup>*Department of Physics, University of Illinois, Urbana Illinois, USA*
- <sup>166</sup>*Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden*
- <sup>167</sup>*Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNM), University of Valencia and CSIC, Valencia, Spain*
- <sup>168</sup>*Department of Physics, University of British Columbia, Vancouver BC, Canada*
- <sup>169</sup>*Department of Physics and Astronomy, University of Victoria, Victoria BC, Canada*
- <sup>170</sup>*Department of Physics, University of Warwick, Coventry, United Kingdom*
- <sup>171</sup>*Waseda University, Tokyo, Japan*
- <sup>172</sup>*Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel*
- <sup>173</sup>*Department of Physics, University of Wisconsin, Madison Wisconsin, USA*
- <sup>174</sup>*Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany*
- <sup>175</sup>*Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany*
- <sup>176</sup>*Department of Physics, Yale University, New Haven Connecticut, USA*
- <sup>177</sup>*Yerevan Physics Institute, Yerevan, Armenia*
- <sup>178</sup>*Centre de Calcul de l'Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France*

<sup>a</sup>Deceased.

<sup>b</sup>Also at Department of Physics, King's College London, London, United Kingdom.

<sup>c</sup>Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.

<sup>d</sup>Also at Novosibirsk State University, Novosibirsk, Russia.

<sup>e</sup>Also at TRIUMF, Vancouver BC, Canada.

<sup>f</sup>Also at Department of Physics, California State University, Fresno CA, USA.

<sup>g</sup>Also at Department of Physics, University of Fribourg, Fribourg, Switzerland.

<sup>h</sup>Also at Departamento de Física e Astronomia, Faculdade de Ciências, Universidade do Porto, Portugal.

<sup>i</sup>Also at Tomsk State University, Tomsk, Russia.

<sup>j</sup>Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France.

<sup>k</sup>Also at Università di Napoli Parthenope, Napoli, Italy.

<sup>l</sup>Also at Institute of Particle Physics (IPP), Canada.

<sup>m</sup>Also at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom.

<sup>n</sup>Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg, Russia.

<sup>o</sup>Also at Louisiana Tech University, Ruston LA, USA.

<sup>p</sup>Also at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain.

<sup>q</sup>Also at Department of Physics, National Tsing Hua University, Taiwan.

<sup>r</sup>Also at Department of Physics, The University of Texas at Austin, Austin TX, USA.

<sup>s</sup>Also at Institute of Theoretical Physics, Ilia State University, Tbilisi, Georgia.

<sup>t</sup>Also at CERN, Geneva, Switzerland

<sup>u</sup>Also at Georgian Technical University (GTU), Tbilisi, Georgia.

<sup>v</sup>Also at Ochadai Academic Production, Ochanomizu University, Tokyo, Japan.

<sup>w</sup>Also at Manhattan College, New York NY, USA.

<sup>x</sup>Also at Hellenic Open University, Patras, Greece.

<sup>y</sup>Also at Institute of Physics, Academia Sinica, Taipei, Taiwan.

<sup>z</sup>Also at LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France.

<sup>aa</sup>Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan.

<sup>bb</sup>Also at School of Physics, Shandong University, Shandong, China.

<sup>cc</sup>Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia.

<sup>dd</sup>Also at Section de Physique, Université de Genève, Geneva, Switzerland.

<sup>ee</sup>Also at International School for Advanced Studies (SISSA), Trieste, Italy.

<sup>ff</sup>Also at Department of Physics and Astronomy, University of South Carolina, Columbia SC, USA.

<sup>gg</sup>Also at School of Physics and Engineering, Sun Yat-sen University, Guangzhou, China.

<sup>hh</sup>Also at Faculty of Physics, M.V.Lomonosov Moscow State University, Moscow, Russia.

<sup>ii</sup>Also at National Research Nuclear University MEPhI, Moscow, Russia.

<sup>jj</sup>Also at Department of Physics, Stanford University, Stanford CA, USA.

<sup>kk</sup>Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary.

<sup>ll</sup>Also at Department of Physics, The University of Michigan, Ann Arbor MI, USA.

<sup>mm</sup>Also at Discipline of Physics, University of KwaZulu-Natal, Durban, South Africa.

<sup>nn</sup>Also at University of Malaya, Department of Physics, Kuala Lumpur, Malaysia.