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To cite this article: Livia Elena Crica et al 2021 2D Mater. 8 035044

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OPEN ACCESS

RECEIVED

11 January 2021

REVISED

26 April 2021

ACCEPTED FOR PUBLICATION 4 May 2021

PUBLISHED

2 June 2021

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PAPER

A method for the measurement of mass and number of graphene oxide sheets in suspension based on non-spherical approximations

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Keywords: RMM, buoyant mass, particle, 2D materials Supplementary material for this article is available online

Abstract

Currently, particle analysis of 2D materials in suspension is commonly restricted to microscopic techniques in the dry state, and thus does not permit an accurate investigation of colloidal suspensions. Colloids in bulk can be assessed by light scattering and diffraction to investigate features such as their hydrodynamic size, charge and concentration. However, the main drawback of such techniques lies in the application of analytical and computational methods based on models assuming particle sphericity which are not representative for 2D materials. Resonance mass measurement (RMM) is a technique which can enable the analysis of 2D materials in suspension without the assumptions of spherical models. Here, we report the application of RMM to measure particle mass and concentration for three types of graphene oxide (GO) aqueous dispersions. Using micro- and nano-suspended resonating sensors, we were able to decipher gravimetric differences between GO and graphitic materials. Our results support the urge for proper definitions and standardisations of graphene based materials, and offer a new method of characterisation for 2D material colloids in liquid suspension.

1. Introduction

Two-dimensional (2D) materials, starting with graphene-based materials (GBMs), are expected to be incorporated into a vast range of applications from electronics to materials science and from aerospace engineering to biomedicine (Novoselov et al 2012, Dimitrios et al 2013, Randviir et al 2014). However GBMs come in various forms, with different properties: different lateral size and thickness distribution, oxidation degree and purity, often leading to contradictions within the scientific community (Bianco 2013, Wick et al 2014, Kauling et al 2018). There is an urgent need for proper definitions and standardisation for GBMs which are often described with similar product names despite several and significant deviations in physicochemical characteristics, thereby presenting an obstacle to the industrial adoption and translation of these innovative materials (Bianco et al

2013, Bøggild 2018). For this reason, good practice guidelines and standards are under development by various institutions (Bianco et al 2013, ISO 2017, Pollard and Clifford 2017, Kauling et al 2018, Pollard

Currently, microscopic techniques, such as optical microscopy (OM), scanning and transmission electron microscopy (SEM and TEM) and atomic force microscopy (AFM) are the most commonly used methodologies for the analysis of individual GBM sheets (Pollard 2017). However, when 2D sheets are in colloidal suspension, microscopic analysis of GBMs in a dry state is not ideal. This is because sheets can fold, wrinkle, bundle or stack to adopt different conformations, either when freely floating in different fluids or during the drying processes applied, thereby affecting various properties, such as their size, thickness and density or even mass (Whitby et al 2012). Common methods to analyse particles

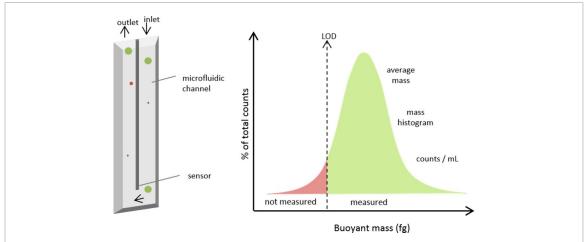


Figure 1. Illustration of the RMM principle: particles flow through the microfluidic channel and change the resonant frequency of the sensor. The frequency shift caused by each particle is translated to buoyant mass. As particles are measured individually in a known volume, the number of counts per millilitre is approximated. The limit of detection (LOD) separates the background noise and undetectable particles from the particle count.

in suspension are dynamic light scattering (DLS) and nanoparticle tracking analysis (NTA), which explore the light scattered by particles in Brownian motion to approximate their size. However, both techniques rely on spherical models which are unsuitable for flat and irregularly shaped nanomaterials such as GBMs.

An alternative technique for the analysis of suspended particles is resonance mass measurement (RMM). RMM is an innovative and emerging technique used to characterise and count particles within the subvisible and submicron range (Weinbuch et al 2013, Höldrich et al 2017). At the heart of the technology lies a microelectromechanical sensor comprising a resonating cantilever with a microfluidic channel, through which suspended particles are measured under flow (figure 1). When a particle between 50 nm and 5 μ m flows through the cantilever it will cause a shift in the cantilever's resonant frequency, should the density of the particle differ to the fluid in which it is dispersed. This frequency shift indicates the buoyant mass of the particle being measured; denser materials will cause a negative frequency shift and have a positive buoyant mass, whereas less dense materials will cause a positive frequency shift and have a negative buoyant mass. This can be exploited to discriminate between positively and negatively buoyant particles within a sample. The extent of the frequency shift directly relates to the particle's buoyant mass, from which the dry mass can be inferred as well as particle volume and diameter (assuming a spherical shape). Since particles flow through the cantilever individually and are counted in a known volume, RMM can accurately determine particle concentration. Furthermore, given the ability to measure individual particles and the extreme sensitivity of this technique (RMM is able to measure particle mass as low as 10^{-18} g), highresolution particle mass distribution information is inherently obtained. We foresee that characterising

GBMs by mass in this way could be more valuable than reporting a spherical size by light scattering techniques, since these materials are two-dimensional.

So far, RMM has proved useful in the analysis of mammalian and bacterial cells, in identifying protein aggregates, and characterising nanomaterials and nanobubbles (Nittayacharn et al 2018, Krueger et al 2019, Pollo et al 2019). Son et al measured the change in volume and density of cells undergoing mitosis to discover that cells swell and reduce their density as a result of intracellular water regulation (Son et al 2015). Lewis et al used RMM to compare the live and dead bacterial populations and monitor bacterial growth using real time monitoring of cell concentration (2014). The change in mass of nanomaterials upon surface modification was also assessed by RMM. Höldrich et al used this technique to determine the number of pepsin molecules adsorbed on gold nanoparticles, the results of which were in agreement with commercial protein quantification kits (2017). Similarly, Nejadnik et al studied the interaction of bovine serum albumin and monoclonal human antibody with polystyrene nanoparticles and showed a shift in particle mass which could not be determined with other techniques (Nejadnik and Jiskoot 2015). RMM can also be used to differentiate particles based on their buoyancy where size-dependent techniques fail. In pharmaceutical sciences for instance, RMM can be used to detect positively buoyant contaminants such as silicon oil in drug formulations that contain negatively buoyant particles (e.g. protein aggregates) (Weinbuch et al 2013, Pollo et al 2019). In a similar way, ultrafine bubbles can be distinguished from other types of particles in suspensions by measuring their positively buoyant mass (Kobayashi et al 2014).

However, due to its relative novelty, RMM is currently an insufficiently explored technique which could serve sectors where the particle number and

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mass in suspension is of interest but difficult to obtain with conventional techniques. For the characterisation of GBMs in aqueous dispersions in particular, we theorised that RMM could be more precise in measuring the particle mass and concentration of GBMs than other more traditional techniques. Indeed, the uniqueness of this technique stands in its ability to measure individual particles and clusters/aggregates in fluids without using spherical models, a major drawback of many alternative techniques for 2D materials. Herein, we characterised the physicochemical properties, including the particle mass and concentration by RMM, of three types of graphene oxide (GO) aqueous dispersions, including one prepared in-house and two commercial products. To the best of our knowledge, these are the first mass measurements of GO sheets in an aqueous dispersion. We report the added value of measuring buoyant mass using RMM as an extremely sensitive technique that complements conventional microscopy. This new characterisation parameter could be of interest in the undergoing development of standards for the production and quality control of available or emerging 2D materials, where the same product name still serves as an umbrella for several materials despite their different physicochemical properties.

2. Results and discussion

To briefly check the origin of the three materials, we evaluated their physicochemical, structural and morphological properties. Lab GO and Commercial GO-A presented the typical Raman fingerprints, whereas the spectra of Commercial GO-B indicated the existence of graphitic domains by the splitting of the G band to G + D' and by the presence of a sharp 2D peak (figure 2(A)) (Pimenta et al 2007). This is also reflected by the reduced metric of disorder (I_D/I_G ratio) found for the Commercial GO-B ($I_{\rm D}/I_{\rm G}=1.00\pm0.01$) in comparison to the other two samples (Lab GO = 1.23 ± 0.03 ; Commercial GO-A = 1.42 ± 0.01). Furthermore, the UV-VIS absorbance spectra of Commercial GO-B (figure 2(B)) suggests a more graphitic structure, free from functional groups, rather than a GO-specific one (Alam et al 2018). These indications correlate with our observations of the poor colloidal stability of Commercial GO-B in water (figure S1 (available online at stacks.iop.org/2DM/8/ 035044/mmedia)). In contrast, the absorbance bands at 230 nm ($\pi \to \pi^*$ transition of C=C bonds) and 310 nm ($n \to \pi^*$ electron transitions in the C=O groups) shape the fingerprint of highly exfoliated and oxidised GO (Paredes et al 2008, Dhifaf et al 2016), for both Lab GO and Commercial GO-A samples. As a consequence of the oxidised nature, which inhibits platelet re-stacking via electrostatic repulsions (Dreyer et al 2010), both Lab GO and Commercial GO are stable colloids in MilliQ water (figure S1).

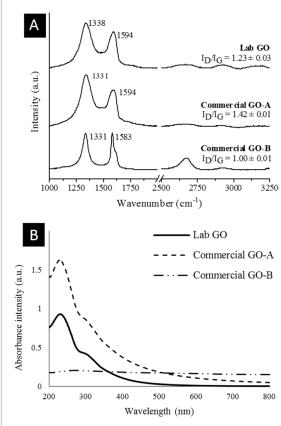


Figure 2. GO material spectroscopic characterisation: structural differences between the lab prepared and commercial GO suspensions, as determined by (A) Raman spectroscopy and (B) UV–VIS spectroscopy. $I_{\rm D}/I_{\rm G}$ values represent the mean value $\pm {\rm SD}$ of three measurements per sample.

Colloidal stability is of great importance for RMM measurements: at high sedimentation rates, a fraction of the sheet population might be inaccurately measured, leading to poor measurements.

Microscopy analysis using TEM and AFM showed predominantly small and thin layers within the Lab GO (figure 3(A)) and larger thin sheets in Commercial GO-A (figure 3(D)), whereas Commercial GO-B contained mostly large and thick agglomerates (figure 3(G)). Overall, Lab GO contained monolayers with lateral sizes up to 1 μ m (figures 3(B) and (C)). Single-layers also dominated in the Commercial GO-A sample, however their lateral size spanned from \sim 0.5 to a few tens of μ m (figures 3(E) and (F)), while Commercial GO-B contained mainly aggregates with >100 nm thickness and micrometric sizes (figures 3(H) and (I)).

To assess the mass via RMM and because the lateral size of the two commercial GO samples exceeded 1 μ m (size limit for the nanosensor), we opted to use only the microsensor, while the Lab GO was analysed using both sensors. Figure 4(A) shows the buoyant mass distribution of the three types of GO measured with the micro-sensor.

Despite all materials having their maximum located at a similar buoyant mass, the percentage of the

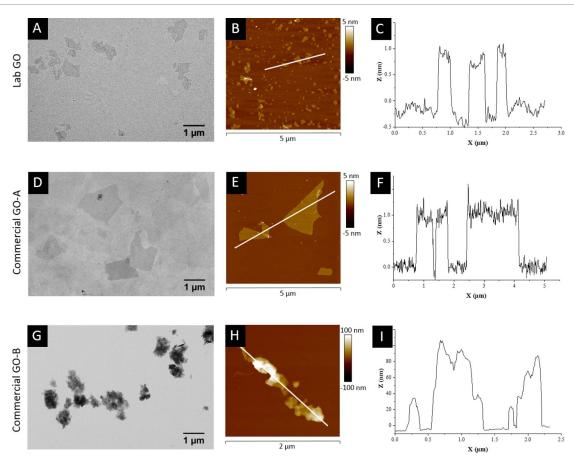


Figure 3. GO material morphology: aspect, lateral size and thickness determined by TEM and AFM for (A)–(C) Lab GO, (D)–(F) Commercial GO-A and (G)–(I) Commercial GO-B. Figures (C), (F) and (I) show the size and thickness of the flakes sectioned within the AFM images (B), (E) and (H), respectively.

total counts was higher for the Lab GO. Moreover, Lab GO spanned from 6.18 fg (LOD micro) to 84 fg, while the commercial GO-A and -B contained particles with buoyant mass up to 968 and 12 400 fg, respectively. The narrower mass distribution of Lab GO suggested a lower sample polydispersity. When measured with the nanosensor, the mass distribution of Lab GO ranged from 0.75 to 12.3 fg (figure 4(B)). In these gravimetric ranges, the average buoyant mass of the particles was 11.2 and 1.2 fg for Lab GO (microsensor vs nanosensor), 24 fg for Commercial GO-A and 28 fg for Commercial GO-B (figure 4(C)).

Next, the particle number n (#/mg) reported in figure 4(D) was calculated using the following formula:

$$n = \frac{C_{\rm p} \left(\frac{\#}{\rm ml}\right)}{C_{\rm m} \left({\rm mg \ ml}^{-1}\right)}$$

where C_p (#/ml) is the particle concentration and C_m (mg ml⁻¹) is the mass concentration of the GO stock. By using this formula instead of simply dividing 1 mg by the average dry mass, we avoid an overestimation, i.e. assuming that all particles in 1 mg of GO would be detectable by the instrument (figure S3).

Based upon previous work (Dhifaf et al 2016, Artur Filipe et al 2018), we expected that Lab GO

contained thin particles with a lateral size down to a few tens of nm (20-50 nm by AFM). A fraction of these sheets would lie below the detection limit of both the micro- and nano-sensors of the RMM instrument. In agreement with this, the low particle count of Lab GO measured using the microsensor suggested that the highest fraction of sheets was indeed below the detection limit. This was then confirmed using the nanosensor, which is more sensitive to lighter sheets. As expected, when using this sensor, Lab GO contained a higher number of particles in 1 mg compared to Commercial GO-A and GO-B. All RMM results correlated with the AFM and TEM data, where Lab GO contained thin small sheets, Commercial GO-A had thin large sheets and Commercial GO-B contained larger and thicker graphitic aggregates.

Based on these results, the authors disagree with defining the Commercial B sample as 'graphene oxide' and would highlight the importance in consulting the good practice guides and ISO standard referenced here before labelling such materials (Pollard and Clifford 2017, Pollard 2017, ISO 2017). We believe that the average mass per sheet and total number of sheets in dispersion are complementary information that can be used for a more comprehensive characterisation of GBMs. RMM could also serve as a quality control technique in the production

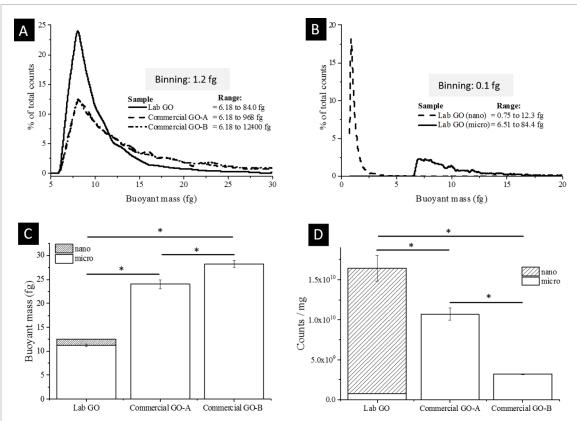


Figure 4. RMM analysis: (A) and (B) buoyant mass distribution of the three types of GO measured with the microsensor and nanosensor, respectively. The nanosensor provides better resolution for measuring the Lab GO as compared with the microsensor. (C) and (D) Average buoyant mass (\pm SD) and counts/mg (\pm SD) of Lab GO compared with Commercial GO-A and -B. Each sample was measured in triplicate, with 1500 counts per replicate. Statistical significance was assessed using one-way ANOVA, with * <0.01.

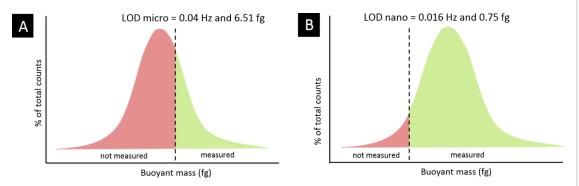


Figure 5. Impact of the LOD on particle measurement: the illustrations exemplify, in an oversimplified manner, why (A) the Lab-GO was poorly measured with the microsensor (high fraction of the flakes below the LOD), while (B) the nanosensor improved the measurement.

of different batches of materials. This tool could be also used to evaluate the mass changes in GBMs upon reduction, functionalisation and adsorption, or to investigate the behaviour of GBMs in different solvents. The main advantages that RMM brings are the possibility to perform measurements directly in the desired solvents (as opposed to microscopy in dry state) and applying non-spherical models (in comparison to DLS or NTA). The limitations of the technique rely on the dependence of detection limit primarily on the density and size of the material to be tested. The graphic in figure 5 illustrates how the LOD impacted on the RMM measurement of our Lab-GO.

The development of suspended resonating sensors with higher sensitivity towards less dense and smaller sheets would be of great benefit for more accurate measurements of the mass and number of graphene-based platelets in dispersions.

Examples where RMM might be useful are nanomedicine and nanosafety, where, similar to the lateral size (Ma *et al* 2015, Vranic *et al* 2018), the mass or number of sheets are complementary information for understanding the effects of 2D materials on biological systems. Although most often the effects of GBMs in cell culture and animal models are expressed as a function of mass concentration, we have seen

here that, when comparing different materials, the same mass can equate to different numbers of sheets. The buoyant mass and particle count could be used to better bridge the gap between the physicochemical properties of 2D materials and their effects on living systems.

3. Conclusions

Firstly, this study emphasises the potential of RMM as a novel technique for the characterisation of 2D materials, providing unique information such as the mass and number of individual particles or aggregates in suspension. RMM offers an alternative for the quantification of 2D materials in dispersion when their optical properties cannot be explored.

Secondly, this paper highlights the importance of thoroughly describing and understanding different types of GBMs defined by similar nomenclature. Here, we compared a lab-prepared GO with two different commercially available GO aqueous dispersions. We report for the first time that the mass and concentration of GO sheets can be measured using suspended resonating sensors. Significant differences in terms of average buoyant mass and particle counts were identified between the three different types of GO, results which were in agreement with particle size and thickness by AFM and TEM.

4. Experimental

Materials: an aqueous dispersion of small GO sheets prepared as previously described (Dhifaf *et al* 2016, Artur Filipe *et al* 2018) was compared with two commercially available GO dispersions. Samples were labelled as Lab GO, Commercial GO-A and Commercial GO-B.

Raman spectroscopy: Raman spectroscopy was performed on drop-casted GO ($100~\mu g\, ml^{-1}$) on glass slides using a calibrated micro-Raman spectrometer (Thermo Scientific, UK), with laser of 633 nm and 0.4 mW. The objective and pinhole aperture were $50\times$ (NA 0.75) and $50~\mu m$, respectively. Three measurements of different spots were acquired for each sample. The raw data was processed by fluorescence correction and further normalised by the intensity of the G band.

UV–Visible (UV–VIS) spectroscopy: UV–VIS absorbance spectra were recorded using a Cary 50 Bio UV/Vis spectrophotometer (Varian Inc., Agilent Technologies, UK). Serial dilutions of GO (2.5–20 μg ml⁻¹) were prepared and their absorbance at 200–800 nm was recorded. All measurements were performed at room temperature, using MilliQ water (Merck-Millipore, UK) as blank.

Atomic force microscopy (AFM): AFM samples were prepared on 0.01% poly-L-lysine (PLL, Sigma Aldrich, UK) coated mica substrates (Agar Scientific, UK). PLL was adsorbed on freshly cleaved mica for 1–2 min, then the substrate was gently washed with

MilliQ water and the excess was removed using filter paper (WhatmanTM cat. no. 1001-070, grade no. 1). Next, GO suspension (100 μ g ml⁻¹) was added and left to adsorb for 1–2 min, followed again by a washing step and removal of the excess. The ared samples were then dried overnight at 37 °C before imaging. The measurements were done in ScanAsyst mode using ScanAsyst Air probes on a Multimode 8 instrument (Bruker, UK). Data post-processing was done with Nanoscope Analysis version 1.4 (Bruker) dedicated software.

Transmission electron microscopy (TEM): samples for TEM were prepared on 300-mesh carbon-coated copper grids (C300Cu, EM resolutions, UK). A FEI Tecnai 12 BioTwin microscope (FEI, Eindhoven, NL), operating at 100 kV and equipped with an AMT digital camera (Gatan, UK) was used.

Resonant mass measurements (RMM): RMM was performed on an Archimedes instrument (Malvern Panalytical, UK) operating with Archimedes software version 1.21. Based on the specifications of the instrument, the target concentration for best measurements would be 8×10^6 particles ml⁻¹ for the microsensor and 2×10^8 particles ml⁻¹ for the nano-sensor. The LOD was optimised to 0.04 Hz for the microsensor and to 0.016 Hz for the nanosensor. The rationale in setting the LOD was to enable the measurement of lightest possible GO particles without significant background noise. A low background noise would normally be seen as zero to few particle counts in the water reference sample, which was run before each measurement. In addition, the coincidence factor, which is a measurement indicator referring to measurement duplicates (e.g. particle clusters) was monitored to be maintained below 3%, where possible. The particle density was fixed to 2.26 g cc^{-1} , i.e. the density of crystal graphite (Dismukes et al 1993), and samples were measured in water until 1500 counts were reached. Results are reported as particle mass distribution, particle average mass and counts per mg.

Statistical analysis: the buoyant mass and number counts per milligram were expressed as means with standard deviation (\pm SD) and evaluated for statistical differences by one-way ANOVA using OriginPro SR2 (b9.4.2.380, Academic).

Data availability statement

All data that support the findings of this study are included within the article (and any supplementary files).

Acknowledgments

This work was supported through a studentship offered by the EPSRC Centre for Doctoral Training in the Science and Applications of Graphene and Related Nanomaterials (GrapheneNOWNANO CDT; EP/L01548X/1) within the Nanomedicine Lab

at the Faculty of Biology, Medicine and Health (FBMH) and the National Graphene Institute, The University of Manchester, UK. The microscopy facilities at the FBMH EM Core Facility and the Bioimaging Facility and the Wellcome Trust were accessed in this study. The authors would also like to acknowledge Artur Filipe Rodrigues and Cyrill Bussy for their scientific advice and for critically reading this manuscript. Malvern Panalytical UK funded and offered technical support for all resonance mass measurements.

Authors contribution

L E C synthesized the GO, characterised all samples, designed the study and wrote the manuscript. T J D and E A G offered support and advice for Archimedes use and RMM data interpretation. K K designed and directed the study, offered intellectual support and contributed to writing the manuscript.

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