



Publication Year	2019
Acceptance in OA@INAF	2020-12-23T12:38:15Z
Title	SAPREMO: a simplified algorithm for predicting detections of electromagnetic transients in surveys
Authors	Vinciguerra, S.; Branchesi, M.; CIOLFI, RICCARDO; Mandel, I.; Neijssel, C. J.; et al.
DOI	10.1093/mnras/sty3490
Handle	http://hdl.handle.net/20.500.12386/29155
Journal	MONTHLY NOTICES OF THE ROYAL ASTRONOMICAL SOCIETY
Number	484

SAPREMO: a simplified algorithm for predicting detections of electromagnetic transients in surveys

S. Vinciguerra,^{1,2★} M. Branchesi,^{3,4★} R. Ciolfi,^{5,6} I. Mandel,^{1,7} C. J. Neijssel¹ and G. Stratta^{8,9}

¹*Institute of Gravitational Wave Astronomy and School of Physics and Astronomy, University of Birmingham, B15 2TT Birmingham, UK*

²*Max Planck Institute for Gravitational Physics (Albert Einstein Institute), D-30167 Hannover, Germany*

³*Gran Sasso Science Institute, Viale Francesco Crispi, 7, I-67100 L'Aquila AQ, Italy*

⁴*INFN, Laboratori Nazionali del Gran Sasso, I-67100 Assergi, Italy*

⁵*INAF, Osservatorio Astronomico di Padova, Vicolo dell'Osservatorio 5, I-35122 Padova, Italy*

⁶*INFN-TIFPA, Trento Institute for Fundamental Physics and Applications, Via Sommarive 14, I-38123 Trento, Italy*

⁷*Monash Centre for Astrophysics, School of Physics and Astronomy, Monash University, Clayton, Victoria 3800, Australia*

⁸*University of Urbino 'Carlo Bo', Dipartimento di Scienze di Base e Fondamenti/Physics section, Via Santa Chiara, 27, I-61029 Urbino, Italy*

⁹*INFN, Sezione di Firenze, Sesto Fiorentino, I-50019 Firenze, Italy*

Accepted 2018 December 11. Received 2018 December 7; in original form 2018 September 20

ABSTRACT

The multiwavelength detection of GW170817 has inaugurated multimessenger astronomy. The next step consists in interpreting observations coming from population of gravitational wave sources. We introduce SAPREMO, a tool aimed at predicting the number of electromagnetic signals characterized by a specific light curve and spectrum, expected in a particular sky survey. By looking at past surveys, SAPREMO allows us to constrain models of electromagnetic emission or event rates. Applying SAPREMO to proposed astronomical missions/observing campaigns provides a perspective on their scientific impact and tests the effect of adopting different observational strategies. For our first case study, we adopt a model of spin-down-powered X-ray emission predicted for a binary neutron star merger producing a long-lived neutron star. We apply SAPREMO on data collected by *XMM–Newton* and *Chandra* and during 10^4 s of observations with the mission concept *THESEUS*. We demonstrate that our emission model and binary neutron star merger rate imply the presence of some signals in the *XMM–Newton* catalogues. We also show that the new class of X-ray transients found by Bauer et al. in the *Chandra* Deep Field-South is marginally consistent with the expected rate. Finally, by studying the mission concept *THESEUS*, we demonstrate the substantial impact of a much larger field of view in searches of X-ray transients.

Key words: gravitational waves – methods: data analysis – catalogues – surveys – stars: neutron – X-rays: bursts.

1 INTRODUCTION

GW170817 (Abbott et al. 2017a) has just opened the era of multimessenger astronomy (Abbott et al. 2017b). The first coincident set of gravitational waves (GWs) and electromagnetic (EM) observations has already provided an extraordinary insight into the physics of the binary neutron star (NS) mergers. Among the key results of this revolutionary discovery is the confirmation of the association between the merger of two NSs and (at least some) short

gamma-ray bursts (SGRBs; Abbott et al. 2017b and refs. therein). The last radio VLBI observations demonstrate that a narrow jet was formed and prove the association with a classical SGRB (see Troja et al. 2017; D’Avanzo et al. 2018; Dobie et al. 2018; Ghirlanda et al. 2018; Lyman et al. 2018; Margutti et al. 2018; Mooley et al. 2018a; Mooley et al. 2018b).

The intense multiwavelength follow-ups of gamma-ray bursts in the last decade have revealed new and unexpected features, such as early and late X-ray flares, extended emission, and X-ray plateaus (e.g. Berger 2014 and refs. therein). The challenges posed by this rich astronomical scenario motivated a growing interest of the community in investigating compact binary mergers from both the theoretical and observational points of view. Intensified theoretical

* E-mail: serena.vinciguerra@aei.mpg.de (SV);
marica.branchesi@gmail.com (MB)

efforts have been dedicated to explain these observations and coherently explore these and other possible EM signals generated by these sources.

In order to validate the variety of proposed theoretical scenarios in the context of multimessenger astronomy with compact binary mergers, we present SAPREMO.

We developed SAPREMO, a *Simplified Algorithm for Predicting ElectroMagnetic Observations*, to evaluate how many EM signals, characterized by a specific light curve and spectrum, should be present in a data set given some overall characteristics of the astronomical survey and a cosmological rate of compact binary mergers. Predictions can be used both to validate the theoretical scenarios against already collected data and to critically examine the scientific means of future missions and their observational strategies. While we use compact binary mergers as the prime multimessenger targets, SAPREMO can also be applied to other type of transients (e.g. core-collapse supernovae).

We describe the main features of SAPREMO in Section 2. As first case study, we use SAPREMO to investigate the X-ray emission from Binary NS (BNS) mergers leading to the formation of a long-lived and strongly magnetized NS, following the model of Siegel & Ciolfi (2016a,b; see Section 3.1). We apply SAPREMO to present X-ray surveys, collected by *XMM-Newton* (Jansen et al. 2001; Strüder et al. 2001; Turner et al. 2001) and *Chandra* (Weisskopf & Van Speybroeck 1996; Weisskopf et al. 2000), and study the perspectives of the mission concept *THESEUS* (Amati et al. 2018). Results are reported in Section 3.3 and discussed in Section 4. Finally, in Section 5 we draw our conclusions summarizing our first results and outlining the main features and primary scopes of SAPREMO. Throughout the paper we assume a flat cosmology with: $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_M = 0.3$, and $\Omega_\Lambda = 0.7$.

2 SAPREMO OUTLINE

SAPREMO is a PYTHON algorithm designed to predict how many detectable EM signals, associated with a specific EM emission (*EMe*) model, are present in a survey S . The full code and a short manual are publicly available at https://github.com/sapREMO/source_code.

According to instrument limitations, such as field of view (FoV) and spectral sensitivity, SAPREMO estimates the number of signals whose emission flux F at the observer is above the flux limit F_{lim} of the S survey. Accounting for the energy dependency of the survey sensitivity, we define detections on an instantaneous flux-based criterion: $\exists g, t' \mid F_g(t') > F_{\text{lim}, g}$, where g labels the spectral band of the survey. We therefore simplify our analysis by treating detectability in each band independently, i.e. a source is considered to be detected if and only if it can be independently detected in at least one instrument band. More realistic treatments include flux integration over the observation time and noise simulation (see Carbone et al. 2017 and references therein for a discussion of these data analysis aspects).

SAPREMO does not directly consider the actual sky locations observed by the survey S (even when applied to archival data) and instead focuses on accounting for cosmological distances, relying on the isotropy of the Universe.

2.1 Core analysis

SAPREMO can be applied to any type of EM emission, from transients to continuous sources emitting in any EM spectral range. In this work, we focus on X-ray transients associated with mergers

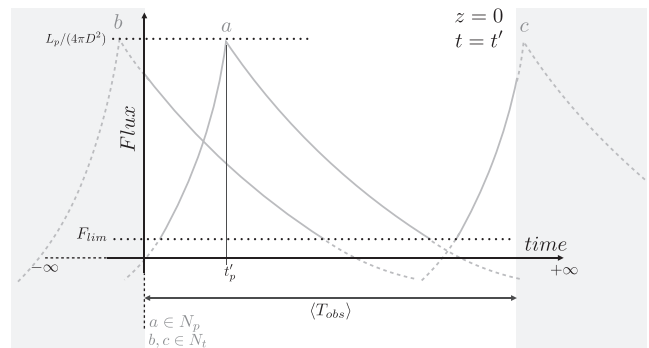


Figure 1. Schematic representation of a peak ($t_p \in \langle T_{\text{obs}} \rangle$) signal (a) and tails (b, c). The solid curves represent the part of signals *EMe* visible during the exposure time at the observer, the dashed components are the missed (because of time or flux restrictions) part of the emissions. The upper dotted line shows the peak flux $F_p = L_p / (4\pi D^2)$, the lower line the flux limit of the survey.

of NS binaries. The expected number of BNS mergers N_{BNS} in the volume enclosed within redshift z_{max} , in a time T at the observer, is

$$N_{\text{BNS}} = T \int_0^{z_{\text{max}}} \frac{R_V(z) dV_C}{1+z} dz, \quad (1)$$

where $R_V(z)$ is the rate of BNS mergers per unit comoving volume, per unit source time. In our case z_{max} is the maximum distance at which the emission following the model of interest *EMe* can be detected. z_{max} is calculated considering both the spectral shift due to the source redshift compared to the instrument energy band $E^I \sim [E_{\text{min}}^I, E_{\text{max}}^I]$ and the maximum luminosity distance, set by the peak luminosity $L_p(E)$ of the *EMe* model and the sensitivity F_{lim} of the survey. We only expect a fraction of N_{BNS} to be observed by a specific instrument, depending on the emission properties and the characteristics of the survey. The number of BNS mergers, detectable by the survey S such that the peak of the considered emission *EMe* falls within the observing time, is given by

$$N_p = \varepsilon \frac{FoV}{4\pi} T \int_0^{z_{\text{max}}} \frac{R_V(z) dV_C}{1+z} dz. \quad (2)$$

The total observing time T can be expressed in terms of the survey S as $T = \langle T_{\text{obs}} \rangle n_{\text{obs}}$, where n_{obs} is the number of the observations and $\langle T_{\text{obs}} \rangle$ is the average exposure time for observation. In equation (2), the *FoV* of the instrument reduces the number N_{BNS} of signals present in the volume enclosed within z_{max} by $FoV/4\pi$. In the specific case of BNS mergers, the efficiency factor ε typically includes the occurrence rate of a specific merger remnant (ε_{sr}), which are expected to generate the emission *EMe*, and source geometry/observational restrictions such as collimation [$\varepsilon_c = 1 - \cos(\theta)$, where θ is the beaming angle], so that $\varepsilon \sim \varepsilon_{\text{sr}} \cdot \varepsilon_c$. We designate as *peaks* the signals included in N_p (see Fig. 1). This contribution only depends on the emission model by the intensity of the light-curve peaks in the energy bands of the survey. This dependence is enclosed in z_{max} .

There is also a contribution, which we call *tails* (Fig. 1), from the mergers whose emission is detected only before (first block of equation 3) or after (second block of equation 3) the luminosity peaks (i.e. L_p is outside the observation period). The longer the light curve is above F_{lim} , the higher the probability of it being observed (see Carbone et al. 2017 for a detailed discussion of signal duration in the context of transient detectability and classification). To estimate this contribution, we need to account for the evolution in time of the emission luminosity $L(t')$, which affects the horizon

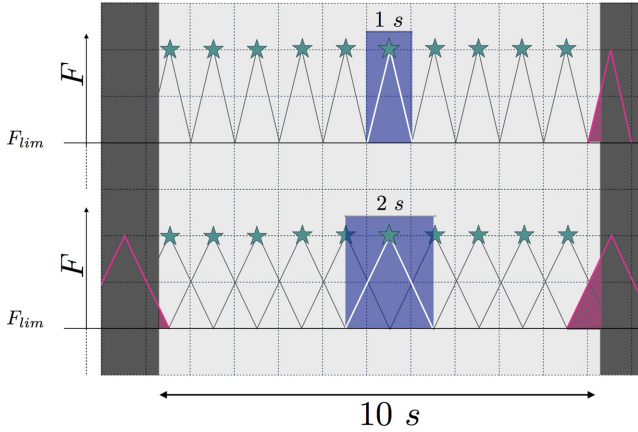


Figure 2. Simplified examples of peaks and tails. In both upper and lower panels we report the detected flux (in arbitrary units on the y axis) as a fraction of time at the observer. We assume a transient class emitting from $z \sim 0$, characterized by the unrealistic rate 1 Hz and a fixed triangular light curve. The flux limit F_{lim} is represented by the horizontal line. The transient light curves above the flux limit are shown in black solid lines. The light grey region highlights the 10 s exposure window. The panels show the peak contribution, i.e. the events whose luminosity peak falls in the observational window (green stars) and the tail contribution, i.e. the events detected only before or after the luminosity peak (fuchsia lines and area). The only difference between upper and lower panels is the duration of the transient light curve above the flux limit, respectively, 1 s and 2 s. Doubling the light-curve duration results in doubling the number of tails, while leaving unchanged the number of peaks.

of the survey:

$$\begin{aligned}
 N_t &= \varepsilon n_{\text{obs}} \frac{FoV}{4\pi} \left[\int_{-\infty}^{t'_p} \int_0^{z_t(L(t'))} \frac{R_V(z)}{1+z} \frac{dVc}{dz} dz dt \right. \\
 &\quad \left. + \int_{t'_p}^{+\infty} \int_0^{z_t(L(t'))} \frac{R_V(z)}{1+z} \frac{dVc}{dz} dz dt \right] \\
 &= \varepsilon n_{\text{obs}} \frac{FoV}{4\pi} \int_{-\infty}^{+\infty} \int_0^{z_t(L(t'))} \frac{R_V(z)}{1+z} \frac{dVc}{dz} dz dt, \quad (3)
 \end{aligned}$$

where t and $t' = t/(1+z)$ are the time, respectively, in observer and source frames and t'_p is the time corresponding to the peak luminosity. $z_t(L(t'))$ represents the horizon of the survey, given the specific intrinsic luminosity of the source $L(t')$. The integration time of equation (3) is practically limited by the duration of the emission above the flux limit. At the moment, SAPREMO does not correctly account for the possibility of detecting multiple times the same event. Multiple detections of the same source might occur if the survey contains repeated observations of the same sky locations and the time interval between the different exposures is shorter than the considered emission EMe . While N_p would be unaffected, in these cases N_t would overestimate the expected number of events by these additional detections. Under these specific conditions, N_t should then be considered as an upper limit. We refer the readers to Carbone et al. 2017 for discussions of transient detectability in the context of multiple images of the same field.

The ratio between the duration of the observable emission and the typical exposure time determines the relative importance of peaks and tails. The trade-off between these two contributions, as well as their different origin, can be understood with Figs 1 and 2. For illustrative purposes, we consider only local events with unphysical rates and a generic triangular light curve. Fig. 1 shows how tail

events b and c can be observed during one exposure of duration $\langle T_{\text{obs}} \rangle$. Fig. 2 shows the impact of the transient observable duration on tails. Upper and lower panels represent exactly the same scenarios (10 s of exposure of transients at $z = 0$ characterized by a 1 Hz rate) except for the transient duration, which is doubled in the lower panel. The number of stars, which represent the peak contribution, is the same in both upper and lower panels, demonstrating that peaks are unaffected by the change of the transient observable duration. However the tail contribution, given by the number of pink triangles, doubles in the lower panel compared to the upper one. On the contrary, extending the exposure to 20 s would double the number of peaks, while leaving unchanged the number of tails. Given a fixed event rate, N_p depends on the observing time, while N_t depends on the duration of the events. We return to this topic in Section 4, when we discuss the results of this study.

While equations (2) and (3) explain the general concept behind the tool, they do not explicitly account for the energy (or wavelength $E = \hbar c/\lambda$) dependence of light curves, instrument sensitivities, and absorption. These effects are particularly important when exploring the Universe at high redshift. We explain how these effects are implemented in SAPREMO analysis in Appendix A.

2.2 Inputs and outputs

We now present inputs and outputs of SAPREMO.

Inputs:

(i) **light curves (LC)**, in terms of luminosity rest frame, of the EM emission EMe in different energy bins (if predicted by the model, also including energies outside the instrument sensitivity band, as they might be redshifted into it, after accounting for cosmology corrections);

(ii) **astrophysical rate** in the source frame $R_V(z)$;

(iii) **efficiency ε** of the EM model EMe (ε should account for source geometry/selection effects, such as collimation, as well as the frequency with which this type of astrophysical source generates the electromagnetic emission EMe);

(iv) **main instrument and survey properties:**

- (a) for each spectral band g of the survey S , minimum and maximum energy included $[E_i, E_s]_g$;
- (b) corresponding flux limits $[F_{\text{lim}}]_g$;
- (c) average exposure time $\langle T_{\text{obs}} \rangle$;
- (d) field of view FoV ¹;
- (e) number of observations $n_{\text{obs}} = 1$.

Outputs:

(i) N_p and N_t : numbers of peaks and tails that are expected in the survey S . The numbers of signals returned by SAPREMO should be interpreted as the expectation value of a Poisson process. Therefore, Poisson statistical errors should be considered in addition to the systematics due to rate and emission model uncertainties;

(ii) dN_p/dz and dN_t/dz : distributions of tail and peak numbers as a function of redshift;

(iii) $dN_p/d\log(D)$ and $dN_t/d\log(D)$: expected distribution of signal observed durations, obtained by convolving the approximate distribution of the survey exposure times P_{obs} with the light-curve span LCS observable at each step in redshift. For each redshift, the LCS represents the total time of the light curve that is above the flux limits at the observer frame (for more details see Appendix A3).

¹Or equivalently covered sky area $f_{\text{sky}} \sim n_{\text{obs}} \frac{FoV}{4\pi}$.

To estimate the distribution of the signal durations, we analytically approximate the exposure time distribution P_{obs} from the average exposure time $\langle T_{\text{obs}} \rangle$ (and standard deviation, when available) with a *Maxwell–Boltzmann* or *Log-normal* function, according to the user input. For each data point saved from the cosmic integrations, we simulate N_{trials} (for both peaks and tails) observation durations T_{obs} and define for each of them the starting time t_s . The starting time is uniformly drawn from a time interval including both the exposure time of the specific trial T_{obs} and the observable emission at observer $(t'_f - t'_i)(1+z)$ (where t'_f and t'_i are, respectively, the last and first LC time at source satisfying our detection criteria at redshift z). If t_s is drawn in the interval $I_p = [t_p - T_{\text{obs}}, t_p]$, where t_p is the time at observer correspondent to the luminosity peak, it contributes to the peak distribution, otherwise it adds up to the tail distribution. For each simulated observation, the total duration is then calculated summing only the contribution of light-curve intervals whose flux is above the limit;

(iv) $dN_p/d\log(F)$ and $dN_t/d\log(F)$: distributions of peak and tail detection numbers as a function of maximum flux. At each step in redshift, necessary to compute the integral (2), SAPREMO also calculates the associated flux. The fluxes are obtained by summing the contribution of each energy and rescaling with the associated luminosity distance. From the same observations simulated for estimating the duration distributions, we obtain the expected distribution of maximum fluxes.

Distributions in redshift are useful to estimate the horizon of the survey to the emission *EMe* and for astrophysical interpretation. They provide a prior on the redshift distribution when a counterpart allowing z measurement is missing, or constrain cosmic rate evolution of BNS mergers when multiwavelength observations yield the source distance. Distributions of fluxes and durations are robust observables, which enable comparisons with real data.²

3 APPLICATION TO SOFT X-RAY EMISSION FROM LONG-LIVED BINARY NEUTRON STAR MERGER REMNANTS

In the following, we consider a specific application of SAPREMO to the case of spin-down-powered X-ray transients from long-lived NS remnants of BNS mergers.

Depending on the involved masses and the NS equation of state (EOS), a BNS merger can either produce a short-lived remnant, collapsing to a black hole (BH) within a fraction of a second, or a long-lived massive NS. The latter can survive for much longer spin-down time-scales (up to minutes, hours, or more) prior to collapse or even be stable forever (Lasky et al. 2014; Lü et al. 2015). After the discovery of NSs with a mass of $\sim 2 M_{\odot}$ (Demorest et al. 2010; Antoniadis et al. 2013), different authors converged to the idea that the fraction of BNS mergers leading to a long-lived NS remnants should range from a few per cent up to more than half (e.g. Piro et al. 2017). Information extracted from the multimessenger observations of the BNS merger event GW170817 did not change this view, although more stringent constraints on the NS EOS were obtained from the GW signal (Abbott et al. 2018a),

from various indications excluding the prompt collapse to a BH, and from the kilonova brightness and the relatively high mass of the merger ejecta (e.g. Bauswein et al. 2017; Margalit & Metzger 2017; Radice et al. 2018; Rezzolla et al. 2018). An additional supporting element in favour of long-lived remnants is given by the observation of long-lasting (approximately minutes to hours) X-ray transients following a significant fraction of SGRBs (e.g. Rowlinson et al. 2013; Gompertz, O’Brien & Wynn 2014; Lü et al. 2015). Given the short accretion time-scale of a remnant disc on to the central BH ($\lesssim 1$ s), such long-lasting emission represents a challenge for the canonical BH-disc scenario of SGRBs while it can be easily explained by alternative scenarios involving a long-lived NS central engine, e.g. the magnetar (Zhang & Mészáros 2001; Metzger, Quataert & Thompson 2008) and the time-reversal (Ciolfi & Siegel 2015; Ciolfi 2018) scenarios. According to this view, the fraction of SGRBs accompanied by long-lasting X-ray transients might reflect the fraction of BNS mergers producing a long-lived NS.

If the merger remnant is a long-lived NS, its spin-down-powered EM emission represents an additional energy reservoir that can potentially result in a detectable transient. Recent studies taking into account the reprocessing of this radiation across the baryon-polluted environment surrounding the merger site have shown that the resulting signal should peak at wavelengths between optical and soft X-rays, with luminosities in the range 10^{43} – 10^{48} erg s⁻¹ and time-scales of minutes to days (e.g. Yu, Zhang & Gao 2013; Metzger & Piro 2014; Siegel & Ciolfi 2016a,b). Besides representing an explanation for the long-lasting X-ray transients accompanying SGRBs, this spin-down-powered emission is a promising counterpart to BNS mergers (e.g. Stratta et al. 2018 and refs. therein), having the advantage of being both very luminous and nearly isotropic.

For our first direct application of SAPREMO, we consider the spin-down-powered transient model by Siegel & Ciolfi 2016a,b (hereafter SC16), described in the next Section 3.1, in which the emission is expected to peak in the soft X-ray band. This model cannot be excluded or constrained by GW170817. The first X-ray observations in the 2–10 keV band were performed by MAXY (Sugita et al. 2018) 4.6 h after the merger with a sensitivity of 8.6×10^{-9} erg s⁻¹ cm⁻², well above the flux that the model predicts at that time after the merger.

In Section 3.2, we briefly describe the model of the BNS merger rate adopted for this first study. Then, in Section 3.3 we present our results referring to three different X-ray satellites: *XMM–Newton*, *Chandra*, and the proposed *THESEUS*. We discuss these results in Section 4.

3.1 Reference emission model

The model proposed by Siegel & Ciolfi (SC16) describes the evolution of the environment surrounding a long-lived NS formed as the result of a BNS merger. The spin-down radiation from the NS injects energy into the system and interacts with the optically thick baryon-loaded wind ejected isotropically in the early post-merger phase, rapidly forming a baryon-free high-pressure cavity or ‘nebula’ (with properties analogous to a pulsar wind nebula) surrounded by a spherical shell of ‘ejecta’ heated and accelerated by the incoming radiation. As long as the ejecta remain optically thick, the non-thermal radiation from the nebula is reprocessed and thermalized before eventually escaping. As soon as the ejecta become optically thin, a signal rebrightening is expected, accompanied by a transition from dominantly thermal to non-thermal spectrum. The

²The reported flux distribution is calculated from the maximum theoretical fluxes of detected events in our simulation (see paragraph on $dN_p/d\log(D)$ output). Quantitative comparisons with actual data would require a more detailed analysis, including the use of the instrument response, a realistic model for noise, the model used to convert photon counts into a light curve, etc. (see e.g. Carbone et al. 2017, who modelled some of these aspects).

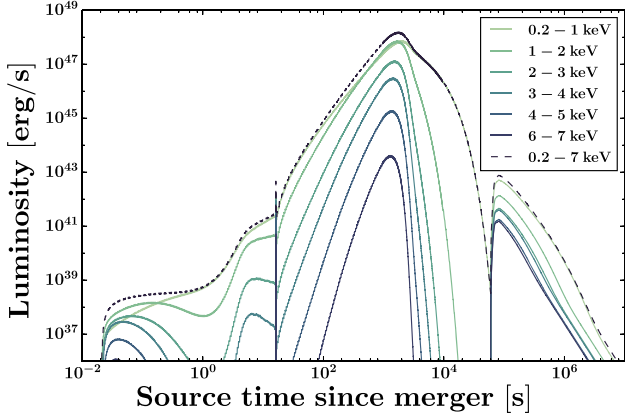


Figure 3. Light curve of the spin-down-powered emission from a long-lived BNS merger remnant according to the model proposed by Siegel & Ciolfi 2016a,b (corresponding to their ‘fiducial’ case; see the text). The solid curves represent the contributions of different energy bands to the total light curve (dashed line).

model can also take into account the collapse of the NS to a BH at any time during the spin-down phase.³

Exploring a wide range of physical parameters, Siegel & Ciolfi found that the escaping spin-down-powered signal has a delayed onset of ~ 10 – 100 s and in most cases peaks ~ 100 – 10^4 s after merger. Furthermore, the emission typically falls inside the soft X-ray band (peaking at ~ 0.1 – 1 keV) and the peak luminosity is in the range 10^{46} – 10^{48} erg s $^{-1}$.

In this work, we consider only one representative case, corresponding to the ‘fiducial case’ of SC16 (SC16f; this model is rescaled for the analysis in Section 3.3.2). The light curve and spectral distribution of this particular model are shown in Fig. 3. The main parameters of the model are as follows. The early baryon-loaded wind ejects mass isotropically at an initial rate of $5 \times 10^{-3} M_{\odot} \text{ s}^{-1}$, decreasing in time with a time-scale of 0.5 s. The dipolar magnetic field strength at the poles of the NS is 10^{15} G and the initial rotational energy of the NS is 5×10^{52} erg (\sim ms initial spin period). Moreover, in this case the remnant evolves without collapsing to a BH. In Fig. 3 we can distinguish two important transitions. The first, around ~ 10 s, marks the end of the early baryon wind phase and the beginning of the spin-down phase. The second, at several times 10^4 s, corresponds to the time when the ejecta become optically thin. While the emission described by the above model is essentially isotropic, allowing us to set $\epsilon_c \sim 1$, only a fraction of BNS mergers ϵ_{LLNS} is expected to generate a long-lived NS. The value of this fraction mainly depends on the unknown NS EOS and distribution of component masses. Here, we assume for simplicity a one-to-one correspondence between the fraction ϵ_{LLNS} and the fraction of SGRBs accompanied by a long-lasting X-ray transient (i.e. extended emission and/or X-ray plateau). Following the analysis presented in Rowlinson et al. 2013, we set ϵ_{LLNS} to 50 per cent.

Once we assign $\epsilon_{\text{sr}} = \epsilon_{\text{LLNS}}$, the resulting total efficiency of the emission is $\epsilon \sim \epsilon_{\text{sr}} \cdot \epsilon_c = 50$ per cent.

³We refer to Siegel & Ciolfi (2016a,b) and Ciolfi (2016) for a detailed discussion of the model and its current limitations.

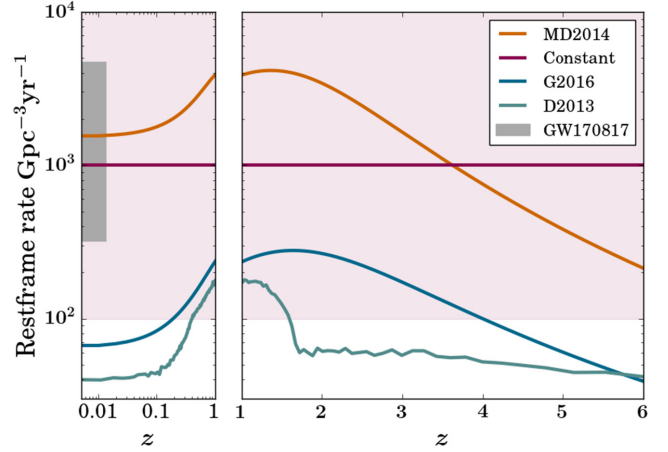


Figure 4. BNS merger rate as a function of redshift for different models: D2013 Dominik et al. (2013) in green, G2016 Ghirlanda et al. (2016) in blue, MD2014 Madau & Dickinson (2014) convolved with $P(t_{\text{del}}) \propto t_{\text{del}}^{-1}$ in orange and our default constant model in violet: the solid line represents the reference rate of $R_V(z) = 1000 \text{ Gpc}^{-3} \text{ yr}^{-1}$, while the light shadowed area includes the whole interval $R_V(z) = [100-10000] \text{ Gpc}^{-3} \text{ yr}^{-1}$. The range in redshift is divided into $z \leq 1$ and $z > 1$ to enhance the visibility of the constraints set by the observation of GW170817 (grey area, 90 per cent probability as reported in Abbott et al. 2017a), which only apply to the local Universe.

3.2 BNS merger rate model

The dependence of the BNS merger rate on redshift is poorly observationally constrained. Several models based on different assumptions have been proposed. For this work we consider four different rate models, a simplified (default) case, and three further astrophysically motivated scenarios:

DEFAULT: a constant BNS merger rate per unit co-moving volume per unit source time in the range $R_V(z) = [100-10000] \text{ Gpc}^{-3} \text{ yr}^{-1}$, extending up to a maximum redshift of $z = 6$;

D2013: the Monte Carlo population synthesis model of Dominik et al. 2013 (their cosmological standard model, high-end metallicity scenario Belczynski et al. 2013);

G2016: the analytic approximation based on SGRB observations described in equation (12) of Ghirlanda et al. 2016 (adopting the average value of the parameter reported for case *a* with an opening angle of 4.5 deg);

MD2014: the analytic prescription for the star formation history proposed by Madau & Dickinson (2014) convolved with a probability distribution of delay times between formation and merger given by the power law $P(t_{\text{del}}) \propto t_{\text{del}}^{-1}$, with a minimum delay time of 20×10^6 yr, normalized to the local BNS merger rate of $1540 \text{ Gpc}^{-3} \text{ yr}^{-1}$, as estimated with GW170817 ($1540_{-1220}^{+3200} \text{ Gpc}^{-3} \text{ yr}^{-1}$ median and uncertainties at 90 per cent probability; Abbott et al. 2017a).

The different BNS merger rates are reported in Fig. 4. We note that D2013 and G2016, as proposed, are inconsistent with the local rate range obtained from a single GW170817 (grey region). However both inferred rate from a single GW observation and population synthesis models rely on poorly constrained astrophysical model assumptions and are therefore highly uncertain. We adopt these rate models for illustrative purposes to test the impact of different redshift-dependent merger rates.

Table 1. Average expected values for peaks (N_p) and tails (N_t) for different surveys. *XMM-Newton* PO and SLEW surveys^a, *Chandra* Deep Field - South (CDF-S) and 10 ks of *THESEUS* operation for a single exposure, *case a*, and 10 distinct exposures *case b* considering $N_H = 5 \times 10^{-22} \text{ cm}^{-2}$ ($N_H = 5 \times 10^{-20} \text{ cm}^{-2}$). ^a For the *XMM-Newton* surveys the total observing time was inferred from $T = n_{\text{obs}}(T_{\text{obs}}) = (4\pi f_{\text{sky}} F_o V^{-1})(T_{\text{obs}})$, using the properties reported in tables B1 and B3.

	<i>XMM-Newton</i>		<i>Chandra</i>	<i>THESEUS</i>	
	PO	SLEW	CDF-S	Case a	Case b
N_p	8	0	0.14	5 (4)	3 (2)
N_t	25	165	\times	5 (3)	20 (11)
$F_o V [\text{deg}^2]$	~0.2		0.08	3300	
$T [10^6 \text{s}]$	~160	~1.06	~6.73	0.01	

To investigate the impact of other inputs, we adopt the DEFAULT simplified model of constant cosmological rate as a reference case. We report distributions and results for $R_V(z) = 1000 \text{ Gpc}^{-3} \text{ yr}^{-1}$. Because the considered $R_V(z)$ is constant, the results for the upper (lower) bound of the whole range $R_V(z) = [100 - 10000] \text{ Gpc}^{-3} \text{ yr}^{-1}$, can be obtained by scaling up (down) the output quantities by one order of magnitude. This wide range of the BNS merger rate includes the local rate interval inferred from the detection of GW170817 and is broadly consistent with estimates obtained using Galactic BNS observations and population synthesis models (Abadie et al. 2010; Chruslinska et al. 2018; Paul 2018; Vigna-Gómez et al. 2018).

We use the DEFAULT model in Section 3.3 and discuss the impact of applying different BNS merger rate models in Section 4.

3.3 Results

We now proceed with using SAPREMO and the X-ray transient model described in Section 3.1 to evaluate the expected number of detections of this type of signal in three present and future surveys by *XMM-Newton*, *Chandra*, and *THESEUS*. To emphasize the impact of the survey properties, we initially keep fixed (i) the light curve to the SC16 model described in Section 3.1 (fiducial, SD16f, or fiducial rescaled; see Section 3.3.2); (ii) the assumed astrophysical rate to the DEFAULT case, described in Section 3.2; and (iii) the efficiency $\varepsilon \sim 50$ per cent.

In particular, we consider:

- (i) two present surveys, collected during the decade of operation of *XMM-Newton*, to predict the expected number of detectable signals in these archived data;
- (ii) the *Chandra* Deep Field - South (CDF-S) data set to verify whether the transient class discovered by Bauer et al. (2017) is statistically consistent with the SC16 model;
- (iii) 10 ks of *THESEUS* observations, to explore the sensitivity of this mission concept to transients associated with BNS mergers, such as SC16f.

The main properties of the surveys are summarized in Appendix B.

3.3.1 *XMM-Newton*

We apply SAPREMO to two different collections of data obtained by *XMM-Newton*; we call them SLEW and PO (Pointed Observations), their characteristics are presented in the following. The number of signals predicted by SAPREMO are reported in Table 1.⁴ In the case

of *XMM-Newton* surveys, the sky locations of observations have been used to estimate the impact of the absorption due to the Milky Way (see Appendix A2 for more details on our adopted absorption model).

The PO survey is a collection of pointed observations made between 2000 February 3 and 2016 December 15. The data belong to the *XMM-Newton Serendipitous Source Catalogue (3XMM DR7; Rosen et al. 2016; XMM-Newton SSC Consortium 2017a)*.

PO exposures are longer (typically 10^3 – 10^4 s) compared to the SLEW catalogue (see following paragraph). This implies an extension to lower fluxes (down to 10^{-15} – $10^{-16} \text{ erg s}^{-1} \text{ cm}^{-2}$), as Fig. 5(a) shows. The same figure shows that such low fluxes are however reached only by tails. This is because the luminosity of the model makes the flux higher than the survey flux limits up to $z = 6$, which is the artificial cut of our BNS merger rate. The distribution of source redshift is represented in the bottom graph of Fig. 5(a) and implies that, under the assumption of a constant cosmic BNS merger rate, the median redshift of detectable signals is $z \sim 2$. The double bump in the tail distribution of the PO survey is explained by the blue and purple curves that, respectively, represent the light-curve span above the threshold at a fixed redshift z , $LCS(z)$, and the time-shifted rate of events per unit redshift, $R_V(z) dV_c/dz$. Given our simplified BNS merger rate model, the purple and black solid lines scale like the redshift derivative of the comoving volume, since in both cases only constants multiply the element dV_c/dz . The blue curve has instead a very peculiar behaviour that depends on the specific emission light curve compared to the limit fluxes of the survey. The distinct trends in the blue lines of Fig. 5 are due to features very peculiar to the adopted light curve (Fig. 3). When the flux from the non-thermal second peak drops below the limit, the overall visible duration sharply decreases; this happens at $z \sim 0.5$ and $z \sim 0.05$ for PO and SLEW, respectively. The second turnover at $z = 4$ in the LCS, evident in PO tails (grey area and dashed black line), is instead due to the discretization of the light curves in energy bins and the relatively soft energy spectrum of these transients. In particular the lowest energy bin characterizing the light curve (see Fig. 3) exits the band of the instrument.⁵

The SLEW survey is composed by data collected while changing the target in the sky, according to the *XMM-Newton* observation program (Smale 2017a). The tested observations are collected in the *XMM-Newton Slew Survey Clean Source Catalog, v2.0*.

The SLEW survey is characterized by typical exposure time of only few seconds, and consequent flux limits $\gtrsim 10^{-13} \text{ ergs}^{-1} \text{ cm}^{-2}$. Given the properties of the model SC16, this yields a predominance of tails over peaks (see first point of Section 4). Our results show that SLEW observations, assuming correct identification (see Section 4), could already reveal a population of BNS merger events. The flux limits of the survey determine the distance of most of the X-ray sources at $z < 3$.

Although the data have been collected by the same instrument, PO and SLEW considerably differ in terms of exposure time (and

⁴For both the surveys, the statistical uncertainties due to the assumed Poisson distribution are almost negligible compared to the systematics due to uncertainty in the signal production efficiency ε and the BNS merger rate.

⁵According to Appendix A notation: $z_{\text{exit}} = [E'_{\text{max},h=0}/E'_{\text{min}} - 1] \sim [1 \text{ keV}/0.2 \text{ keV} - 1] = 4$, where z_{exit} is the redshift at which the energy bin of lowest energy, denoted with $h = 0$, exits the instrument band, $E'_{\text{max},h=0} = 1 \text{ keV}$ is the highest energy included in the bin $h = 0$ and $E'_{\text{min}} = 0.2 \text{ keV}$ is the minimum energy included in the instrument band.

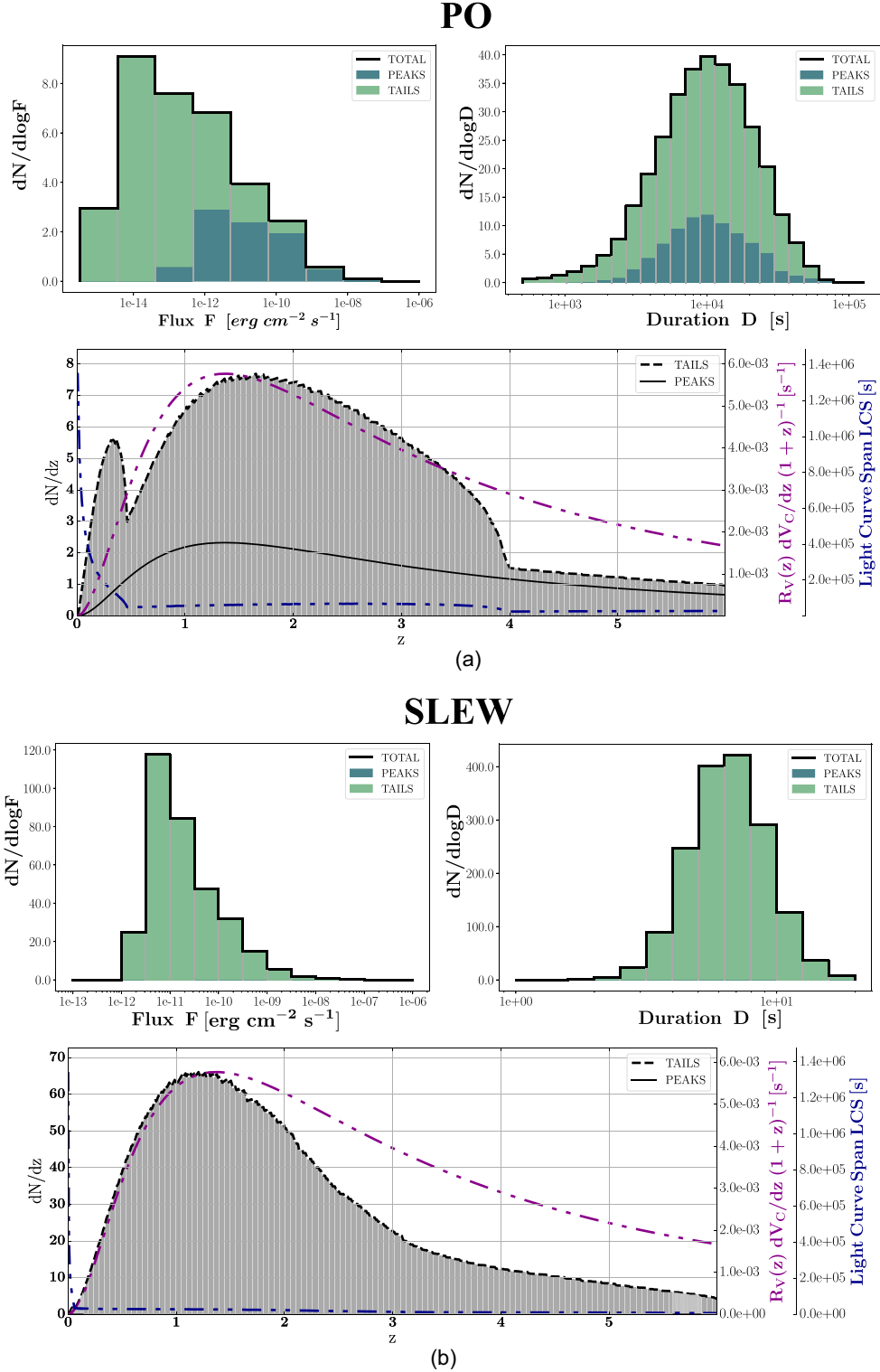


Figure 5. Comparison between expected distributions of detectable signals in PO and SLEW surveys. The upper-left panel shows the total expectations on the observed maximum flux distribution obtained by adding peak (dark green) and tail (green) contributions. The upper-right panel shows the expected distributions of peak and tail durations. The redshift distributions in the bottom panel represent the differential contribution of peaks and tails throughout the scanned comoving volume of the Universe. The violet and blue curves have been added to explain the trend of the tail distribution and represent, respectively, the time-shifted rate of events per unit redshift (scaled argument of the z integral A7) and the light-curve span $LCS(z)$; the duration of the light curve above the flux limits at a fixed z . (a) Expected distributions of observed maximum flux (upper-left), duration of signals above the flux limits (upper-right), and redshift (bottom) of the observations in *XMM-Newton* PO survey. (b) Expected distributions of observed maximum flux (upper-left), duration of signals above the flux limits (upper-right), and redshift (bottom) of the observations in *XMM-Newton* SLEW survey. Because of the low expected number of detections for peaks compared to tails (see Table 1), the black solid line in bottom graph is not visible.

therefore sensitivity), sky coverage, and energy responses (as shown by tables in Appendix B1). The SLEW survey is less sensitive, but it scans a much wider area of the sky compared to the PO survey, so that the total number of expected signals is actually considerably larger (see Table 1).

3.3.2 Chandra and the new, faint X-ray population

Bauer et al. (2017) have recently claimed the discovery of a X-ray signal, belonging to a new, previously unobserved, transient class. The event was observed within the CDF-S, a deep survey of a 0.11 deg^2 sky region composed of 102 observations collected in different periods during the last decade. Interestingly, the main properties of the event presented by Bauer et al. (2017) are broadly consistent with the SC16 emission model. The maximum luminosity of $\sim 10^{47} \text{ erg s}^{-1}$, the spectral peak around $\sim 2 \text{ keV}$ (source frame), the rise time of $\sim 100\text{s}$, and the overall duration of $\sim 10^4 \text{ s}$ are all in broad agreement with the model predictions. Here we do not attempt to provide convincing evidence for a potential match, but we want to show another interesting case for exploiting the capabilities of SAPREMO. We assume a signal analogous to the SC16f adopted throughout this paper only rescaled to have a maximum luminosity of $\sim 10^{47} \text{ erg s}^{-1}$ (referred to as ‘rescaled SC16 model/signal’ in the following), consistent with the X-ray transient at $z \sim 2.23$ (Bauer et al. 2017). To test the rate consistency between the detected X-ray transient and the rescaled SC16 model, we apply SAPREMO to the CDF-S, adopting a Galactic neutral column density of $n_{H_{\text{tot},\text{MW}}} \approx n_{H,\text{MW}} \sim 8.8 \times 10^{19} \text{ cm}^{-2}$, as reported by Bauer et al. (2017). Given the proposed source redshifts, the shape and fluxes of the detected transient, we assume that the observed maximum flux corresponds to the luminosity peak of our model. We therefore evaluate only the expected number of peaks. SAPREMO predicts an expectation value of ~ 0.14 signals in the CDF-S (see Table 1). Given the adopted constant rate model, the probability of one rescaled SC16 signal being present at its luminosity peak in the $\sim 7\text{Ms}$ of the CDF-S is ~ 12 per cent (with ~ 87 per cent probability of 0 signals). Considering the whole range of allowed BNS merger rates, this value ranges from ~ 1.4 per cent (with ~ 98.6 per cent probability of 0 signals, correspondent to $R_V(z) = 100 \text{ Gpc}^{-3} \text{ yr}^{-1}$) and ~ 35 per cent (with ~ 25 per cent probability of 0 signals, correspondent to $R_V(z) = 10000 \text{ Gpc}^{-3} \text{ yr}^{-1}$).

Despite the broad consistency of the transient revealed by Bauer et al. (2017) with the rescaled SC16 emission model, our analysis shows that a real association between the two is rather disfavoured, although not inconsistent given the uncertainties over rate and emission model. Conversely, assuming that the detected transient is in fact the rescaled SC16 signal adopted in the above calculation, the constant BNS merger rate value is constrained to $6000_{-3700}^{+4000} \text{ Gpc}^{-3} \text{ yr}^{-1}$ (median with 90 per cent credible interval, as in Abbott et al. 2017a), higher than the median inferred from GW170817, though still consistent with the claimed interval (cf. Section 3.2).

3.3.3 Future observations with THESEUS

In the last few years, different wide-FoV X-ray missions have been proposed to monitor the X-ray sky, and specifically to follow up GRBs and GWs (Barcons et al. 2012; Merloni et al. 2012; Yuan

et al. 2015; Zhang et al. 2017). In particular, the mission concept *THESEUS* has been recently selected by ESA for assessment studies (Bauer & Colangeli 2018) to explore the transient high-energy sky and contribute to the multimessenger astronomy (Amati et al. 2018; Frontera et al. 2018; Stratta et al. 2018). We apply SAPREMO to test the sensitivity of the *THESEUS* mission to BNS mergers emitting in the X-ray according to the SC16f model. On the *THESEUS* payload, the Soft X-ray Imager (SXI, O’Brien et al. 2018) would be the instrument sensitive to such emission. SXI flux sensitivities for sources in the Galactic plane ($N_H = 5 \times 10^{22} \text{ cm}^{-2}$) and well outside it ($N_H = 5 \times 10^{20} \text{ cm}^{-2}$) are taken from fig. 4 of Amati et al. (2018). With SAPREMO, we predict detection numbers and properties for two cases of gathering *THESEUS* observations, each having a total observing duration T of 10 ks, acquired with

case a a single exposure of $\langle T_{\text{obs}} \rangle = 10^4 \text{ s}$;

case b 10 exposures of non-overlapping sky regions, each lasting $\langle T_{\text{obs}} \rangle = 10^3 \text{ s}$.

The last two columns of Table 1 show the expectation values for N_p and N_t for both cases *a* and *b*. In less than 3 h of total observing time T , we expect *THESEUS* to detect a number of SC16-like transients comparable to the ones inferred for years-long CDF-S *Chandra* and *PO XMM-Newton* surveys (see also Table 1). The main reason why *THESEUS* is capable of reaching these numbers of detections in such a short observing time T (about four orders of magnitude shorter than for CDF-S and PO/SLEW) is the four orders of magnitude difference between its FoV and the ones in *Chandra* and *XMM-Newton*. Albeit specific for the SC16f emission, our tabulated results demonstrate that the characteristics of the mission concept *THESEUS* suit the search for X-ray emission generated during BNS mergers. We predict that *THESEUS/SXI* with *case a* sensitivity would detect on the order of hundreds to thousands of BNS mergers in only a few years, assuming an emission with peak luminosity in the range $L_p \sim [10^{45} - 10^{48}] \text{ erg cm}^{-2} \text{ s}^{-1}$ and a spectral distribution similar to SC16f. Our analysis therefore shows that proposed large FoV instruments, such as *THESEUS*, offer an incredible opportunity compared to present deep surveys, for the detection of rare but bright transients, such as those of the presented model. *THESEUS/SXI*-like campaigns are expected to detect events generated near the peak of the cosmic star formation and BNS merger rate, at $z \sim 1 - 3$.

The redshift distribution dN_p/dz of cases *a* and *b* in Fig. 6 show that, as expected, longer exposure times decrease the flux limits and therefore probe larger redshift. Meanwhile, multiple shorter exposures at distinct sky locations, as in case *b*, increase the probability of detecting tails. We find that *THESEUS/SXI* would make more detections, during a fixed total observing time, with an observing strategy that increased sky coverage at the cost of shorter exposures.

4 DISCUSSION

With SAPREMO, we tested the sensitivity of different astronomical surveys to the emission model SC16f.

Given a fixed emission model, the tail contribution becomes more important for surveys with shorter exposures $\langle T_{\text{obs}} \rangle$. The analysis has confirmed that when the exposure $\langle T_{\text{obs}} \rangle$ is considerably shorter than the duration of the emission model, as is the case of *XMM-Newton* SLEW observations, detections of pre-peak rises or post-peak decays will be more common than observations of peak flux. Therefore the relation between $\langle T_{\text{obs}} \rangle$ and the typical duration of the EM model sets the relative number of peaks and tails. Given

⁶To estimate the posterior on the constant rate value, we assume a flat prior in the range $[100, 10000] \text{ Gpc}^{-3} \text{ yr}^{-1}$.

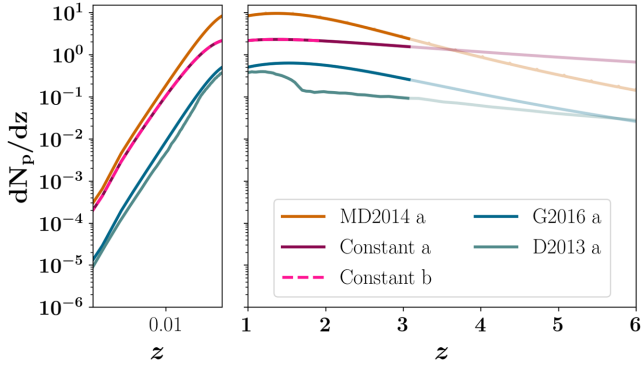


Figure 6. Expected redshift distribution for peak signals in 10 ks of *THESEUS/SXI* observations (opaque lines), assuming $N_H = 5 \times 10^{20} \text{ cm}^{-2}$. In the default configuration of constant BNS merger rates we test two different observing strategies: *case a* and *case b*, as described in Section 3.3.3. They are shown by the solid-violet and dashed-fuchsia lines, which completely overlap for $z \lesssim 2$ and only differ in the maximum achieved redshift, $z \sim 3$ for the former and $z \sim 2$ for the latter. We also report the peak distribution as a function of redshift for other tested BNS merger models: MD2014 in orange, D2013 in green, and G2016 in blue. For comparison with Fig. 4, we show the results in the two different regimes: $z > 1$ in linear and $z \leq 1$ in logarithmic scale. For the same reason, we also add in transparency the distribution of peaks missed because of sensitivity constraints.

Table 2. For *XMM-Newton* PO (top) and *THESEUS* case *a* surveys, comparison of expectation values assuming different BNS merger rate models, from left to right (i) analytic prescription proposed by Madau & Dickinson (2014), assuming a power-law distribution of delay times between formation and merger $P(t_{\text{del}}) \propto t_{\text{del}}^{-1}$, (ii) cosmological rate derived by the population synthesis study Dominik et al. (2013), standard model at high-end metallicity scenario (D2013), and (iii) model based on SGRB statistics Ghirlanda et al. (2016) with assumed opening angle of 4.5 deg (G2016).

<i>XMM-Newton</i> PO			
Rate model	MD2014	D2013	G2016
N_p	20	1	2
N_t	65	2	5
FoV[deg ²]		~ 0.25	
$n_{\text{obs}} (T_{\text{Obs}})[s]$		$\sim 234 \times 10^6$	
<i>THESEUS</i> Case <i>a</i>			
Rate model	MD2014	D2013	G2016
N_p	17 (15)	0.54 (0.46)	1.2 (1.0)
N_t	16 (10)	0.54 (0.34)	1.0 (0.6)

a fixed amount of total observing time, the number of expected tails increases with the number of pointings of different sky positions. This is shown, for example, by the two cases tested with *THESEUS* (see Tables 1 and 2).

Inferences on BNS merger models

(i) **A few detections can already constrain the lower limit of cosmic BNS merger rate.** For example, assuming the emission model proposed by Siegel & Ciolfi (2016b), the probability of detecting more than three peaks in *XMM-Newton* PO assuming a constant merger rate of $100 \text{ Gpc}^{-3} \text{ yr}^{-1}$ is < 1 percent; thus, a few detections could set a lower limit on the BNS merger rate. The proposed emission model also offers the unique possibility of exploring mergers in the high-redshift Universe: with peak luminosity as high as $\sim 10^{48} \text{ erg s}^{-1}$, *XMM-Newton* could detect signals generated at redshifts as high as $z \sim 15$ with PO sensitivities.

Constraints on BNS merger rate can also be obtained by assuming the association of the *Chandra* X-ray transient with the SC16 emission model. In this case our analysis can put a lower limit on the BNS merger rate of $\sim 2300 \text{ Gpc}^{-3} \text{ yr}^{-1}$ at 90 per cent confidence interval, assuming a constant rate up to $z \sim 6$. The peak luminosity of the rescaled SC16 emission is indeed bright enough to actually be detectable by *Chandra* up to $z \sim 6$, even after accounting for the Milky Way absorption.

(ii) **Larger PO-like or *THESEUS* surveys will probe $R_V(z)$ and likely also constrain EM emission models.** In the case of bright emission, as expected by SC16, *SAPREMO* predicts that the PO campaign can detect all BNS merger peaks in the Universe, localized within *XMM-Newton*'s FoV. However, Table 1 shows that to probe a statistically significant population of BNS mergers, PO-like sensitivities must be obtained with longer time window T and/or larger FoV. Large FoV instrument, such as *THESEUS/SXI*, can also detect many of these events. Although limited to smaller redshift, cosmological distances could still be achieved if we assume bright emission models. This is shown for the case of SC16 in Fig. 6. These campaigns should therefore yield the BNS merger rate distribution as a function of redshift, providing that multiwavelength observations will allow redshift associations. These observations would then enable us to constrain the proposed BNS merger rate scenarios, which indeed predict different z distributions (see Fig. 6).

(iii) **Redshift measurement play a fundamental role for breaking the degeneracy between emission parameters and rate models.** We apply *SAPREMO* to the PO and *THESEUS* cases, *case a*, for the three scenarios introduced in Section 3.2: D2013, G2016, and MD2014 (see Fig. 4). The absolute expectation values reported in Table 2 and Fig. 6 reflect the trend of the rate models reported in Fig. 4. The overall results show that, without perfect knowledge of light curve and spectrum of the emission, measurements of source distances are necessary to constrain the redshift dependence of the BNS merger rate.

Considerations on *SAPREMO*'s results. In this paragraph we highlight some general considerations to realistically interpreting *SAPREMO* results. The main output of the analysis consists in the number of peaks and tails expected for a specific emission model in a selected survey of data. However, depending on the purpose of the analysis, other information, such as more accurate requirements for detectability and classification, should be taken into account. In the following we give some examples.

Challenges for detectability: despite this does not concern the results presented in the previous session, short transients in long exposure observations can be lost in the integrated background flux. To overcome this issue, targeted analyses might be required (see for e.g. the work of EXTraS group De Luca 2014; De Luca et al. 2016 for detection of $\sim 10^2$ s lasting transients in PO).

Challenges for classifications: because of their definition, we generally expect durations and fluxes of tails to extend to lower values compared to the peak ones (as shown in Fig. 5a). This generally worsen the performances of signal classifications and identification among more common phenomena. In particular for our analyses, given the shape of SC16f's light curve, tails should mostly appear as simple decaying signals, which can be challenging to distinguish from other X-ray transients (e.g. tails of tidal disruption events Lodato & Rossi 2011 or supernovae Dwarkadas & Gruszko 2012). The correct classification of X-ray events can also be challenged by short exposure times. This is, for example, the case of the SLEW survey, where observations typically last only few seconds. Indeed some emission models, as SC16, predict a long-scale time evolution

of the emission properties that would likely result in detections of dissimilar signals, challenging their association to a common origin. Campaigns characterized by typically longer observations, such as the *XMM–Newton* PO and *Chandra* CDF-S, are less affected by classification problems. The extension of the typical exposure time to thousands of seconds and improved spectral resolution, allow for the acquisition of more informative data, simplifying transient identification.

5 SUMMARY AND OUTLOOK

In this study we showed some applications of our tool SAPREMO; we applied it on few present and possible future surveys, assuming a specific emission model and cosmological BNS merger rate $R_V(z)$. In terms of multimessenger astronomy, our results show that the luminosities predicted by the SC16 emission model can be detected up to cosmological distances that extend much further than the horizon of present (Aasi et al. 2016; Abbott et al. 2018b) and future GW detectors (Sathyaprakash et al. 2012) to binary NS mergers, both in the cases of current surveys, such as CDF-S and *XMM–Newton* PO and SLEW, and proposed missions, such as *THESEUS*.

SAPREMO provides theoretical predictions allowing us

(i) *to compare predictions with actual data*. E.g. we proved that some signals consistent with the model could already be detected in present surveys of data such as *XMM–Newton* PO and SLEW;

(ii) *to test potential associations*. E.g. we proved that the new transient found by Bauer et al. (2017) is marginally consistent with the model;

(iii) *to assess the effectiveness of proposed mission concepts for a specific type of signal*. We illustrate the utility of SAPREMO for evaluating proposed missions with a case study of *THESEUS*. We demonstrate that, with few years of operation, the large FoV of *THESEUS/SXI* could allow for the detection of up to thousands of SC16-like signals, enabling considerable constrains on both the BNS merger rate and the emission models;

(iv) *and to compare different observational strategies*. SAPREMO can be used to determine advantages and disadvantages compared to a particular emission of adopting different observational strategies. With the case of *THESEUS*, we indeed demonstrate that SAPREMO can compare observations characterized by different values of typical exposure time ($\langle T_{\text{obs}} \rangle$) and point out the main properties of the relative detections. In general, given a total observing time T , different observational strategies can be applied, increasing the exposure time to increase the sensitivity or decreasing the exposure time to enlarge the sky coverage. The effect of adopting different exposures depends on several parameters, including both source and instrumental properties (such as rate, luminosity, flux limit dependence on exposure time, etc). In this paper, we specifically prove that, given *THESEUS/SXI* sensitivity as a function of exposure time, 10 observations of disjoint sky areas lasting 1 ks would encapture more SC16f-like transients than an extended single pointing of 10 ks.

In general SAPREMO allows us to test both survey and astrophysical properties. This study has mainly focused on the former, exploring the impact of different trade-offs among such properties (including exposure time, sky localization, and spectral sensitivity), assuming a single light-curve model from SC16. However SAPREMO can also test (and be used for inference on) astrophysical quantities such as emission duration, peak luminosity, and spectra.

Once the design sensitivity of Advanced GW interferometers is achieved, GW detections of EM bright sources, such as GW170817, will occur more and more often and very likely at lower SNRs. In this context of multimessenger astronomy, SAPREMO can be used to optimize the analysis by identifying specific emission model. We conclude remarking that the flexibility of the implemented methodology allows considerations of emission model spanning the whole EM spectrum (e.g. kilonovae models can also be tested). Moreover our analysis include no priors on nature of EM sources, so that it can be applied to a wide range of astrophysical phenomena. With its analysis dedicated to treat high redshift effects, SAPREMO particularly suits studies on emission of cosmological origin.

ACKNOWLEDGEMENTS

The authors thank Lorenzo Amati for his assistance with the case of *THESEUS* and for the useful comments. We thank Andrea Belfiore, Andrea De Luca, Martino Marelli, David Salvietti, and Andrea Tiengo for the help in understanding and interpreting *XMM–Newton* data. We thank Ruben Salvaterra for the useful suggestions and discussions. The research leading to these results has received funding from the People Programme (Marie Curie Actions) of the European Union’s Seventh Framework Programme FP7/2007-2013/ (PEOPLE-2013-ITN) under Research Executive Agency (REA) grant agreement no. 606176. This paper reflects only the authors’ view and the European Union is not liable for any use that may be made of the information contained therein. Giulia Stratta acknowledges European Gravitational Observatory (EGO) support through a VESF fellowship (EGO-DIR-133-2015). This research has made use of data obtained from XMMSL2, the Second *XMM–Newton* Slew Survey Catalogue, produced by members of the XMM SOC, the EPIC consortium, and using work carried out in the context of the EXTraS project (‘Exploring the X-ray Transient and variable Sky’, funded from the European Union’s Seventh Framework Programme under grant agreement no. 607452). This research has made use of data obtained from the 3XMM *XMM–Newton* serendipitous source catalogue compiled by the 10 institutes of the *XMM–Newton* Survey Science Centre selected by European Space Agency (ESA).

REFERENCES

- Aasi J. et al., 2016, *Living Rev. Relativ.*, 19, 1
 Abadie J. et al., 2010, *Class. Quantum Gravity*, 27, 173001
 Abbott B. P. et al., 2017a, *Phys. Rev. Lett.*, 119, 161101
 Abbott B. P. et al., 2017b, *ApJ*, 848, L12
 Abbott B. P. et al., 2018a, *Phys. Rev. Lett.*, 121, 161101
 Abbott B. P. et al., 2018b, *Living Rev. Relativ.*, 21, 3
 Amati L. et al., 2018, *Adv. Space Res.*, 62, 191
 Antoniadis J. et al., 2013, *Science*, 340, 448
 Barcons X. et al., 2012, preprint (arXiv:1207.2745)
 Bauer M., Colangeli L., 2018, Space for Europe. Available at: https://www.esa.int/Our_Activities/Space_Science/ESA_selects_the_new_mission_concepts_for_study
 Bauer F. E. et al., 2017, *MNRAS*, 467, 4841
 Bauswein A., Just O., Janka H. T., Stergioulas N., 2017, *ApJ*, 850, L34
 Belczynski K. et al., 2013, Synthetic Universe. Available at: <https://www.syntheticuniverse.org>
 Berger E., 2014, *ARA&A*, 52, 43
 Carbone D., van der Horst A. J., Wijers R. A. M. J., Rowlinson A., 2017, *MNRAS*, 465, 4106

- Chengalur J. N., Kanekar N., Roy N., 2013, *MNRAS*, 432, 3074
- Chruslinska M., Belczynski K., Klencki J., Benacquista M., 2018, *MNRAS*, 474, 2937
- Ciolfi R., 2016, *ApJ*, 829, 72
- Ciolfi R., 2018, *Int. J. Mod. Phys. D*, 27, 1842004
- Ciolfi R., Siegel D. M., 2015, *ApJ*, 798, L36
- D’Avanzo P. et al., 2018, *A&A*, 613, L1
- De Luca A., 2014, Exploring the X-ray Transients and the variable Sky. Available at: <http://www.extras-fp7.eu/index.php/scientific-community/the-project>
- De Luca A. et al., 2016, in Napolitano N. R., Longo G., Marconi M., Paolillo M., Iodice E., eds, *The Universe of Digital Sky Surveys*. Springer International Publishing, Switzerland, p. 291
- Demorest P. B., Pennucci T., Ransom S. M., Roberts M. S. E., Hessels J. W. T., 2010, *Nature*, 467, 1081
- Dobie D. et al., 2018, *ApJ*, 858, L15
- Dominik M., Belczynski K., Fryer C., Holz D. E., Berti E., Bulik T., Mandel I., O’Shaughnessy R., 2013, *ApJ*, 779, 72
- Dwarkadas V. V., Gruszko J., 2012, *MNRAS*, 419, 1515
- ESA, 2017, THE XMM-NEWTON SLEW SURVEY SOURCE CATALOGUE: XMMSL2. Available at: <https://www.cosmos.esa.int/web/xmm-newton/xmmsl2-ug>
- Frontera F. et al., 2018, *J. Ital. Astron. Soc.*, 89, 157
- Ghirlanda G. et al., 2016, *A&A*, 594, A84
- Ghirlanda G. et al., 2018, preprint (arXiv:1808.00469)
- Gompertz B. P., O’Brien P. T., Wynn G. A., 2014, *MNRAS*, 438, 240
- Jansen F. et al., 2001, *A&A*, 365, L1
- Kalberla P. M. W., Burton W. B., Hartmann D., Arnal E. M., Bajaja E., Morras R., Pöppel W. G. L., 2005, *A&A*, 440, 775
- Lasky P. D., Haskell B., Ravi V., Howell E. J., Coward D. M., 2014, *Phys. Rev. D*, 89, 047302
- Lehmer B. D. et al., 2005, *ApJS*, 161, 21
- Lodato G., Rossi E. M., 2011, *MNRAS*, 410, 359
- Lü H.-J., Zhang B., Lei W.-H., Li Y., Lasky P. D., 2015, *ApJ*, 805, 89
- Luo B. et al., 2017, *ApJS*, 228, 2
- Lyman J. D. et al., 2018, *Nature Astron.*, 2, 751
- Madau P., Dickinson M., 2014, *ARA&A*, 52, 415
- Margalit B., Metzger B. D., 2017, *ApJ*, 850, L19
- Margutti R. et al., 2018, *ApJ*, 856, L18
- Merloni A. et al., 2012, preprint (arXiv:1209.3114)
- Metzger B. D., Piro A. L., 2014, *MNRAS*, 439, 3916
- Metzger B. D., Quataert E., Thompson T. A., 2008, *MNRAS*, 385, 1455
- Mooley K. et al., 2018a, *Nature*, 554, 207
- Mooley K. et al., 2018b, *Nature*, 561, 355
- Morrison R., McCammon D., 1983, *ApJ*, 270, 119
- O’Brien P., Bozzo E., Willingale R., Hutchinson I., Osborne J., Amati L., Götz D., 2018, *J. Ital. Astron. Soc.*, 89, 130
- Paul D., 2018, *MNRAS*, 477, 4275
- Piro A. L., Giacomazzo B., Perna R., 2017, *ApJ*, 844, L19
- Radice D., Perego A., Zappa F., Bernuzzi S., 2018, *ApJ*, 852, L29
- Read A. M., Saxton R. D., 2016, EPIC-pn SLEW-specific PSF parameterisation. Available at: <https://www.cosmos.esa.int/web/xmm-newton/ccf-release-notes>
- Rezzolla L., Most E. R., Weih L. R., 2018, *ApJ*, 852, L25
- Rosen S. R. et al., 2016, *A&A*, 590, A1
- Rowlinson A., O’Brien P. T., Metzger B. D., Tanvir N. R., Levan A. J., 2013, *MNRAS*, 430, 1061
- Sathyaprakash B. et al., 2012, *Class. Quantum Gravity*, 29, 124013
- Saxton R. et al., 2008, *A&A*, 480, 611
- Saxton R. D. et al., 2017, Optical loading in the XMM-Newton slew source catalogue. Available at: <http://xmm2.esac.esa.int/docs/documents/CAL-TN-0210-0-1.ps.gz>
- Siegel D., Ciolfi R., 2016a, *ApJ*, 819, 14
- Siegel D., Ciolfi R., 2016b, *ApJ*, 819, 15
- Smale A. P., 2017a, the Astrophysics Science Division at NASA/GSFC t. H. E. A. D. o. t. S. A. O., XMM-SLEWCLN - XMM-Newton Slew Survey Clean Source Catalog, v2.0). Available at: <https://heasarc.gsfc.nasa.gov/w3browse/all/xmmslewcln.html>
- Smale A. P., 2017b, the Astrophysics Science Division at NASA/GSFC t. H. E. A. D. o. t. S. A. O., CHANDFS7MS - Chandra Deep Field-South 7-Megasecond X-Ray Source Catalog. Available at: <https://heasarc.gsfc.nasa.gov/W3Browse/chandra/chandfs7ms.html>
- Smale A. P., 2017c, the Astrophysics Science Division at NASA/GSFC t. H. E. A. D. o. t. S. A. O., XMMSSC - XMM-Newton Serendipitous Source Catalog (3XMM DR7 Version). Available at: <https://heasarc.gsfc.nasa.gov/W3Browse/xmm-Newton/xmmsc.html>
- Stratta G. et al., 2018, *Adv. Space Res.*, 62, 662
- Strüder L. et al., 2001, *A&A*, 365, L18
- Sugita S. et al., 2018, *PASJ*, 70, 81
- Troja E. et al., 2017, *Nature*, 551, 71
- Turner M. J. et al., 2001, *A&A*, 365, L27
- Vigna-Gómez A. et al., 2018, *MNRAS*, 481, 4009
- Weisskopf M. C., Van Speybroeck L., 1996, in Hoover R. B., Walker A. B. C., eds, *Proc. SPIE Conf. Ser. Vol. 2805, Multilayer and Grazing Incidence X-Ray/EUV Optics III*. SPIE, Bellingham, p. 2
- Weisskopf M. C., Tananbaum H. D., Van Speybroeck L. P., O’Dell S. L., 2000, in Truemper J. E., Aschenbach B., eds, *Proc. SPIE Conf. Ser. Vol. 4012, X-Ray Optics, Instruments, and Missions III*. SPIE, Bellingham, p. 2
- Willingale R. et al., 2013, *MNRAS*, 431, 394
- XMM-Newton SSC Consortium, 2017a, XMMSSC - XMM-Newton Serendipitous Source Catalog (3XMM DR7 Version). Available at: <https://heasarc.gsfc.nasa.gov/w3browse/all/xmmsc.html>
- XMM-Newton SSC Consortium, 2017b, Table 2.1: A list of all 9710 observations and exposures used for the 3XMM-DR7 source detection. Available at: http://xmmsc.irap.omp.eu/Catalogue/3XMM-DR7/3xmmdr7_obslist.html
- XMM-Newton SSC Consortium, 2017c, 3XMM-DR7. Available at: http://xmmsc.irap.omp.eu/Catalogue/3XMM-DR7/3XMM_DR7.html
- Yu Y., Zhang B., Gao H., 2013, *ApJ*, 776, L40
- Yuan W. et al., 2015, preprint (arXiv:1506.07735)
- Zhang B., Mészáros P., 2001, *ApJ*, 552, L35
- Zhang S. N. et al., 2017, in den Herder J. W. A., Takahashi T., Bautz M., eds, *Proc. SPIE Conf. Ser. Vol. 9905, Space Telescopes and Instrumentation 2016: Ultraviolet to Gamma Ray*. SPIE, Bellingham, p. 9905

APPENDIX A: SAPREMO FEATURES

A1 Redshift / K correction

To evaluate if the emission is visible in the g band of the instrument, we need to calculate the fraction of source light curve that contributes to the flux in the g band at the observer. Here we denote each energy band $[E'_{\min}, E'_{\max}]_g$ with the label g ; we use $[E'_{\min}, E'_{\max}]$ without subscripts for the whole range of operation.

We assume a set of light curves $[L_h(t')]_{h=0}^{h_{\max}}$, where each element $L_h(t')$ represents the emission in the source frame within a fixed h energy bin $[E'_{\min}, E'_{\max}]_h$. The redshifted energy bin h , associated to the light-curve element $L_h(t')$, might only partially overlap with an instrumental energy bin g . To consider only the part of the light curve that contributes to the emission visible in the g energy bin, we calculate the fraction of the h energy bin that falls into the g band and assume the energy is uniform across its intrinsic spectrum. Therefore for each numerical step in z , we calculate the emission contribution to each g band:

$$L_g(t', z) = \sum_{h=0}^{h_{\max}} L_h(t') w_{hg}(z). \quad (\text{A1})$$

When the redshifted h bin and the g band of the instrument overlap, the $L_h(t')$ emission contributes to the total observable emission in the g band $L_g(t', z)$ with a weight defined by the ratio between the amount of overlap and the width of the light-curve energy bin:

$$w_{hg}(z) = \begin{cases} \alpha & \text{if } \alpha > 0 \\ 0 & \text{otherwise} \end{cases}, \quad (\text{A2})$$

where

$$\alpha = \frac{\min((1+z)^{-1} E'_{\max,h}, E'_{\max,g}) - \max((1+z)^{-1} E'_{\min,h}, E'_{\min,g})}{(1+z)^{-1} [E'_{\max,h} - E'_{\min,h}]}. \quad (\text{A3})$$

An increased resolution in the energy bins of the emission model will result in more precise estimates.

A2 Absorption

SAPREMO can account for both host and Galactic absorptions. The host-galaxy absorption is included by substituting $L_h(t')$ with $\tilde{L}_h(t') = L_h(t')e^{-n_{H,h}\sigma_h}$, where $n_{H,h}$ is a typical value of the effective hydrogen column density and σ_h is the average of the absorption cross-section in the h energy band in the source frame. Both of these quantities may depend on the type of the host-galaxy. Similarly, the Milky Way absorption is accounted adopting $\tilde{L}_g(t', z) = L_g(t', z)e^{-n_{H,MW}\sigma_g}$. We estimate an effective hydrogen column density as a function of the observed sky locations (Galactic latitudes), adopting the sky map of HI emission-line brightness T_b released by Kalberla et al. (2005). For each position in the sky, the Galactic column density along the line of sight $n_{H,MW}$ is calculated adopting equation (4) of Chengalur, Kanekar & Roy (2013; valid for negligible total opacity). These values then have to be averaged along Galactic longitudes l , $\langle n_{H,MW} \rangle_l$, and finally associated to the relative frequency of observations in the survey to calculate an effective column density n_H :

$$n_{H,MW}[\text{cm}^{-2}] = \frac{1}{n_{\text{obs}}} \sum_{j=1}^{n_{\text{obs}}} \langle n_{H,MW} \rangle_{l_j}(b_j, l_j), \quad (\text{A4})$$

where b is the Galactic latitude.

To establish the detectability of the light curves $\tilde{L}_g(t', z)$, we calculate the corresponding fluxes:

$$F_g(t', z) = \frac{\tilde{L}_g(t', z)}{4\pi D_L^2(z)}. \quad (\text{A5})$$

A2.1 X-ray absorption model

Different specific absorption models can be implemented accordingly to the energy range of interest. In this paper we only consider the effect of X-ray absorption at the observer. In Willingale et al. (2013), the authors investigate hundreds of GRB afterglows detected by Swift to model the effective total Galactic column density in X-ray $n_{H_{\text{tot}},MW}$. Atomic and molecular hydrogen represents the dominant components of $n_{H_{\text{tot}},MW} \approx n_{H,MW} + 2n_{H_2,MW}$. To estimate the molecular hydrogen component from the atomic contribution, we adopt the model proposed by Willingale et al. (2013):

$$n_{H_2,MW} = n_{H_2\text{max}} \left[1 - \exp\left(\frac{n_{H,MW}}{n_c}\right) \right]^\alpha, \quad (\text{A6})$$

where $n_{H_2\text{max}} = 7.5 \times 10^{20} \text{ molecules cm}^{-2}$, $n_c = 2.37 \times 10^{21} \text{ atoms cm}^{-2}$, and $\alpha = 2$. We apply this effective total Galactic column density $n_{H_{\text{tot}},MW}$ to both *XMM-Newton* surveys, PO and SLEW.

The effective cross-section of the interstellar medium, for each of the energy bands in the range 0.03–10 keV, is analytically estimated as a function of energy E , following Morrison & McCammon (1983).

A3 Detailed integrations over cosmical scales

To estimate the peak contribution, we need to account for the redshift dependence of all the quantities involved in the cosmic integration of equation (2). At each step in z , the portion of light curve still present (after the redshift) in each of the instrumental energy band is calculated, the absorption is applied, and the new maximum redshift $z_{\max,g}$ is computed. The contribution of the correspondent step in redshift z is considered if at least on one energy band g , $z_{\max,g} > z$.

Similarly, we compute the number of tails with

$$N_t = \varepsilon \tilde{n}_{\text{obs}} \frac{F_o V}{4\pi} \int_0^{\bar{z}_{\max}} R_V(z) \frac{dV_C}{dz} LCS dz, \quad (\text{A7})$$

where \bar{z}_{\max} is the maximum redshift that contributes to peak count and $LCS(z) = \bigcup_g [dt'_{\text{vis}}(z)]_g$ represents the union of source time-intervals of the light curve $dt'_{\text{vis},i}$ which are visible at z in at least one g band, i.e.

$$\forall g \quad dt'_{\text{vis},i} = t'_{i+1} - t'_i \mid [F_g(t'_i, z) + F_g(t'_{i+1}, z)] > 2F_{\text{lim},g} \quad (\text{A8})$$

This contribution does not include peaks, since the peak duration can be consider infinitesimal. The time dependence for tail calculation is only set by the signal duration, while it is completely unaffected by the survey exposures, which instead determine the peak contribution.

APPENDIX B: SURVEY PROPERTIES

B1 XMM-Newton parameters

We report the parameters adopted to apply SAPREMO to *XMM-Newton* surveys: PO and SLEW in Tables B1, B2, B3, and B4. The data used to characterize PO and SLEW are collected in source catalogues. From there, relative papers (Saxton et al. 2008; Rosen et al. 2016) and webpages (Read & Saxton 2016; ESA 2017; Saxton et al. 2017; *XMM-Newton* SSC Consortium 2017a), we extract the general properties necessary to apply SAPREMO. In the case of PO and SLEW we estimate the impact of absorption from the locations of the sources collected in the catalogues

Table B1. PO: general characteristics of pointed observations contributing to the *XMM-Newton* Serendipitous Source Catalog. The data are part of the 3XMM-DR7 catalogue (Rosen et al. 2016); we adopted the fit file *3xmm-dr7_obslist.fits* available at *XMM-Newton* SSC Consortium (2017b). ^a: From clean observations (OBS_CLASS < 3); ^b: Excluding overlaps.

Parameter	Value
Minimum energy	0.2 keV
Maximum energy	12 keV
$\langle T_{\text{obs}} \rangle^a$	19 000 s
$\sigma_{T_{\text{obs}}}^a$	17 900 s
Covered Sky area	1750(1032 ^b) deg ²

Table B2. PO: spectral bands of pointed observations contributing to the *XMM–Newton* Serendipitous Source Catalog. Medians in each band from catalogue as reported in website (Smale 2017c), catalogue available at *XMM–Newton* SSC Consortium (2017c) catalogue: *3XMM_DR7* cleaned with the same criteria used for calculating average and variance of exposure times.

Energy band (keV)	Sensitivity ($\text{erg cm}^{-2} \text{s}^{-1}$)
0.2–0.5	5.8×10^{-16}
0.5–1.0	1.7×10^{-15}
1.0–2.0	2.7×10^{-15}
2.0–4.5	3.8×10^{-15}
4.5–12.0	6.6×10^{-15}

Table B3. SLEW: general characteristics of slew data contributing to the *XMM–Newton* Slew Survey Catalogue. ^a: from clean observations according to *xmmsl2_clean.fits* file at (Smale 2017a); ^b:percentage when overlaps are excluded from Smale (2017a).

Parameter	Value
Minimum energy	0.2 keV
Maximum energy	12 keV
$\langle T_{\text{obs}} \rangle^a$	6.9 s
$\sigma_{T_{\text{obs}}}^a$	2.4 s
Covered Sky area ^b	84 per cent of the sky

Table B4. SLEW: spectral bands of slew data contributing to the *XMM–Newton* Slew Survey Catalogue. Flux limits from ESA (2017). The energy range in SLEW catalogue is divided just in two bands.

Energy band (keV)	Sensitivity ($\text{erg cm}^{-2} \text{s}^{-1}$)
0.2–2.0	0.57×10^{-12}
2.0–12.0	3.7×10^{-12}

(Rosen et al. 2016).⁷ Since no flux limits have been quoted for PO, for that case we use as a proxy the medians of the fluxes available from the catalogue.

B2 CHANDRA

To apply SAPREMO to the CDF-S we use data from Smale (2017b), Lehmer et al. (2005), and Luo et al. (2017). As quoted in the same references, we adopted $n_{\text{H,MW}} = 8.8 \times 10^{19} \text{ cm}^{-2}$.

The adopted survey properties are reported in the Tables B5 and B6. Tables B5 and B6 report conservative values extracted from figs 2, 28, and 29 of Luo et al. (2017), approximating the main characteristic to a region of the sky described by roughly homogeneous properties.⁸ More details on the set of observations are available at Smale (2017b).

⁷We adopt clean data, requiring: CLEAN = OBS_CLASS < 3. Data concerning average and standard deviation of observation durations are based on total band (*Exp_Map_B8*). Locations adopted for the absorption model are inferred from source locations.

⁸To account for homogeneity, we decide to only consider a 285 arcmin² region, out of the 484.2 arcmin² of the entire survey, which corresponds to the ACIS FoV and roughly to the region observed by at least 6 ms (see Fig. 2 of Luo et al. 2017). A similar sky area is also characterized by flux limits in the total energy band < $2 \times 10^{-16} \text{ erg cm}^{-2} \text{ s}^{-1}$.

Table B5. CDF-S: general properties of CDF-S (roughly homogeneous region) from Lehmer et al. (2005), and Luo et al. (2017). ^a: the maximum cleaned exposure is $6.727 \times 10^6 \text{ s}$.

Parameter	Value
Minimum energy	0.5 keV
Maximum energy	7 keV
FoV	285 arcmin ²
$\langle T_{\text{obs}} \rangle^a$	$\sim 6 \times 10^6 \text{ (s)}$
n_{obs}	1

Table B6. CDF-S: conservative estimates of flux limits from figs 28 and 29 of Luo et al. (2017; at ~ 50 per cent completeness).

Energy bands (keV)	Sensitivity ($\text{erg cm}^{-2} \text{s}^{-1}$)
0.5–2.0	$\sim 6 \times 10^{-17}$
2.0–7.0	$\sim 4 \times 10^{-16}$

Table B7. THESEUS: general properties of the Soft X-ray Imager (SXI) from Amati et al. (2018).

Parameter	Value
Minimum energy	0.3 keV
Maximum energy	6 keV
FoV	$110 \times 30 \text{ deg}^2$

Table B8. THESEUS : properties of the two considered scenarios for a total of 10 ks observations. Sensitivities extrapolated from fig. 4 of Amati et al. (2018) assuming column density typical of regions outside (inside) the Galactic plane.

	$\langle T_{\text{obs}} \rangle$ [s]	n_{obs}	SXI sensitivity ($\text{erg cm}^{-2} \text{s}^{-1}$)
Case a.	10^4	1	7.82×10^{-12} (1.93×10^{-11})
Case b.	10^3	10	3.20×10^{-11} (7.87×10^{-11})

B3 THESEUS

The data concerning *THESEUS* have been extrapolated from Amati et al. (2018). We analyse 10 ks of exposure collected with two different strategies:

- (a) 1 single exposure of 10 ks;
- (b) 10 distinct exposures of 1 ks each.

We report in Tables B7 and B8 the properties adopted for the results presented in this paper.

This paper has been typeset from a $\text{\TeX}/\text{\LaTeX}$ file prepared by the author.