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Abstract

Using three independent ice-flow models and several satellite-based datasets, we assess the importance of correctly capturing ice-shelf breakup, shelf thinning, and reduction in basal traction from ungrounding in reproducing observed speed-up and thinning of Thwaites Glacier between 1995 and 2015. We run several transient numerical simulations applying these three perturbations individually. Our results show that ocean-induced ice-shelf thinning generates most of the observed grounding line retreat, inland speed-up, and mass loss, in agreement with previous work. We improve the agreement with observed inland speed-up and thinning by prescribing changes in ice-shelf geometry and a reduction in basal traction over areas that became ungrounded since 1995, suggesting that shelf breakups and thinning-induced reduction in basal traction play a critical role on Thwaites's dynamics, as pointed out by previous studies. These findings suggest that modeling Thwaites's future requires reliable ocean-induced melt estimates in models that respond accurately to downstream perturbations.

Plain Language Summary

Recent observations have shown that Thwaites Glacier, West Antarctica, has been accelerating and thinning over the past decades and its floating part is quickly breaking up. While these observations suggest that warmer ocean currents are the main factor responsible for these changes, it remains unclear which of the following processes are most important to the glacier's dynamics: (i) breakup of its floating section, (ii) ice-shelf thinning, or (iii) changes in grounded-ice area. By employing three ice-sheet models and several satellite-based datasets, we find that thinning induced by ocean melting and the resulting reduction of grounded-ice area explain most of the observed flow acceleration and mass loss of Thwaites, in agreement with other studies. We also find that the breakup of the floating section plays an important role on Thwaites's dynamics. These findings suggest that improved forecasts of Thwaites's future require reliable ocean-induced melt estimates and improved model response to changes in ice-shelf thickness and geometry.

1 Introduction

Thwaites Glacier, one of the largest ice streams in the Amundsen Sea Embayment (Fig. 1), drains a large area of the West Antarctic Ice Sheet (WAIS). Its ice volume holds the equivalent of ~ 0.65 m of sea level (Morlighem et al., 2020), and is resting on deep bedrock, a wide channel below sea level that spreads under WAIS to the Ross Sea Embayment (Holt et al., 2006; Fretwell et al., 2013). The retrograde slope of this channel makes Thwaites potentially vulnerable to marine ice sheet instability (Weertman, 1974; Schoof, 2007; Gudmundsson et al., 2012), a positive feedback of grounding line retreat and increased ice discharge, which may lead ultimately to WAIS's collapse over the coming centuries (Bamber et al., 2009; Joughin et al., 2014; Feldmann & Levermann, 2015; Scambos et al., 2017; Martin et al., 2019). How fast this collapse may happen depends on internal instability mechanisms (e.g., Favier et al., 2014) and on external forcings that could drive significant mass loss of Thwaites Glacier and WAIS (e.g., Gudmundsson et al., 2019).

Recent observations have shown that Thwaites has been accelerating, thinning and experiencing ice-shelf breakups and grounding line retreat since the 1970s (Mouginot et al., 2014; Rignot et al., 2014; Konrad et al., 2018; Shepherd et al., 2019). The pattern of ice thinning in the Amundsen Sea Sector suggests that changes in ocean conditions are likely the main external driver of ocean-induced ice-shelf thinning, increased calving rates, and changes in grounding line positions (Alley et al., 2015; Seroussi et al., 2017; Milillo et al., 2019). Although these changes are not independent, the exact chain of events that led to Thwaites's thinning and acceleration remains unclear. For instance, the increased thinning of the ice shelves may have compromised their mechanical integrity, lead-

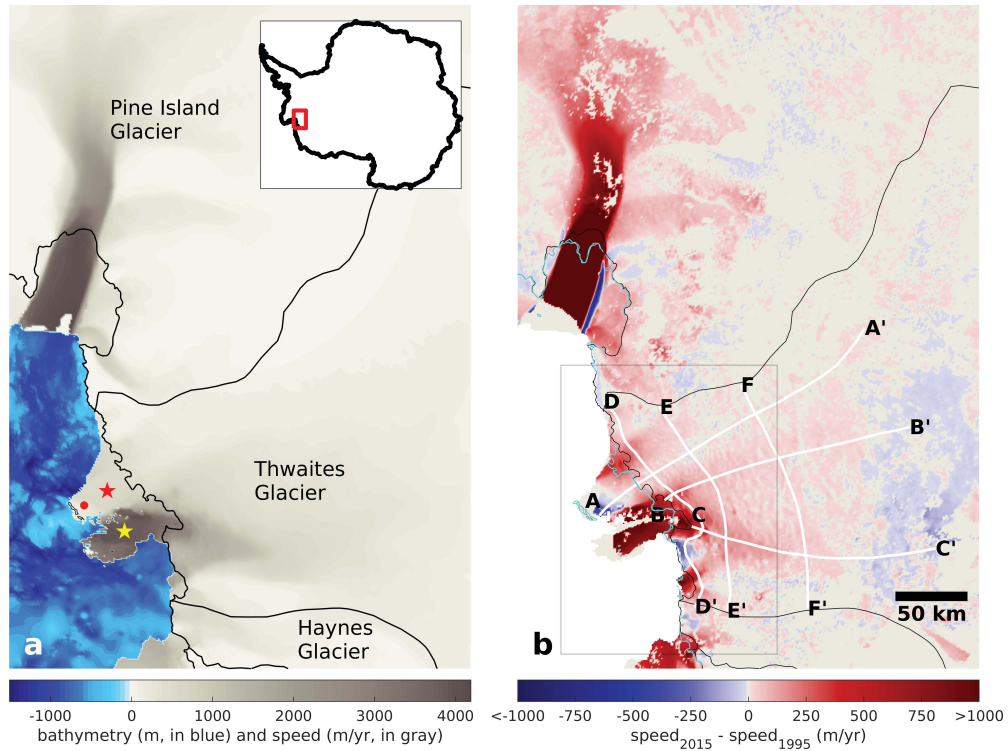


Figure 1. (a) Ocean bathymetry (blue scale in colorbar) and ice speeds (gray scale), and (b) observed speed changes between 1995 and 2015 (Mouginot et al., 2014) of the two largest ice streams of the Amundsen Sea Sector, West Antarctica: Pine Island and Thwaites glaciers. In (a), the red star shows the location of Thwaites Eastern Ice Shelf (TEIS) and the yellow star shows Thwaites Tongue (TT). The red circle shows the location of the pinning point at TEIS's tip. The black lines delineate the drainage basin and the 2015 grounding lines. In (b), the white lines show flow lines and cross sections used in this paper. Transect A-A' is referred to as eastern ice stream, transect B-B' as the main trunk, and transect C-C' as the western ice stream. Lines in light blue are the 1995 grounding line. The box highlights the area shown in Fig. 2.

ing to the partial collapse of Thwaites Eastern Ice Shelf (TEIS) and the complete loss of Thwaites Tongue (TT) (e.g., Scambos et al., 2009; Miles et al., 2020). These changes could have decreased the buttressing provided by the pinning point in TEIS’s tip to the grounded ice (Fig. 1, panel a), inducing glacier speed-up and, as a consequence, the retreat of the grounding line (Fig. 1, panel b).

We investigate the effect of each of these physical processes on numerical modeling of Thwaites’s dynamics between 1995 and 2015. We perform several numerical simulations using three independent ice sheet models (Úa, ISSM, and STREAMICE) for which we prescribe changes in calving front position, ice-shelf thickness, and basal traction due to grounding line retreat, and we quantify their impacts on upstream flow. The resulting changes in ice speed and thickness are compared with satellite-based measurements. The misfit between modeled and observed ice velocities in each case provides estimates of the relative importance of all those observed changes on the glacier flow.

2 Data and Methods

2.1 Data

The three models are initialized to 1995 conditions and are run forward in time until 2015. The 1995 digital elevation model (DEM) is derived from the European Remote Sensing (ERS-1) radar altimetry (Bamber, 2000). The 2015 DEM is the Reference Elevation Model of Antarctica, REMA (Howat et al., 2019), included in BedMachine Antarctica v2 (Morlighem et al., 2020). We employ the most recent bed elevation product derived from mass conservation and recent survey in Thwaites (Hogan et al., 2020; Jordan et al., 2020; Morlighem et al., 2020). Ice velocities and grounding line positions of the initial (1995) and final (2015) states are derived from interferometric synthetic aperture radar data (InSAR, Mouginot et al., 2014; Rignot et al., 2014).

As described in Sect. 2.3, we impose perturbations based on satellite measurements. Changes in calving-front extension are derived from Landsat imagery (MacGregor et al., 2012). Ice-shelf thinning rates are estimated by radar altimetry (Paolo et al., 2015). Grounding line retreat is measured by InSAR data (Rignot et al., 2014). We compare the modeled velocity change and thinning rates with observations of speed change and ice thinning derived from InSAR and radar/laser altimetry data, respectively (Mouginot et al., 2014; Shepherd et al., 2019; Smith et al., 2020). The basal melting is parameterized by a depth-dependent relationship based on observations and ice-ocean-coupling simulations (Rignot et al., 2013; Seroussi et al., 2017; Milillo et al., 2019; Nakayama et al., 2019). The surface mass balance is derived from the Regional Climate Model (RACMO v2.3, Van Wessem et al., 2014).

Landsat imagery shows a rift propagating between TEIS and TT from the 1980s to 2010/2011, when the main part of TT calved off (MacGregor et al., 2012). Based on the hypothesis of a non-negligible shear stress between TEIS and TT prior to 2006 (Mouginot et al., 2014; Miles et al., 2020), likely due to mélange formation that could act as a granular ice shelf (Burton et al., 2018) into that rifted zone, we start all the experiments in 1995 with TEIS and TT mechanically connected. To set up this connection, we remove this rift from the 1995 Landsat-derived ice-front contour, allowing transfer of stresses across the region where the rift later developed. The model initialization (inversion, see Sect. 2.2) adjusts the rheological parameter of the ice into that rifted zone to model the speed differences between TEIS and TT (Fig. 1, panel a). We keep the 1995 shelf-front position fixed in time, except where otherwise stated (see Sect. 2.3).

2.2 Ice sheet models

To assess the robustness of our conclusions with respect to the use of ice sheet models and ice-flow assumptions, we employ three ice sheet models: \acute{U} a (Gudmundsson, 2020), ISSM (Larour et al., 2012), and STREAMICE (Goldberg & Heimbach, 2013). \acute{U} a employs the two-dimensional Shallow Shelf Approximation (SSA; MacAyeal, 1989), ISSM a three-dimensional High-Order model (HO; Blatter, 1995), and STREAMICE a two-dimensional L1L2-type approximation (Goldberg, 2011; Lipscomb et al., 2019).

The model domains comprise the Amundsen Sea Embayment catchment, including Pine Island, Thwaites, and neighboring glaciers (Haynes, Pope, Smith, and Kohler glaciers). The model domains and meshes are similar to those in Barnes et al. (2021). ISSM and \acute{U} a’s meshes rely on Delaunay triangulation with edge lengths varying according to an interpolation error estimate of the observed ice velocity and on the distance to the grounding line. In ISSM, a 500 m mesh resolution is employed close to the grounding line and 15 km on the far field. ISSM generates the three-dimensional mesh by extruding the two-dimensional triangular mesh. In \acute{U} a, 1 km resolution is used near the grounding line, while a coarser resolution (up to 20 km) is used inland. STREAMICE relies on a quadrilateral-element-type mesh with uniform element’s edge length equal to 1.5 km. These resolutions are sufficient for this type of experiment (e.g., Cornford et al., 2020).

All models employ a Weertman’s sliding law (Weertman, 1957):

$$\boldsymbol{\tau}_b = -\beta^2 \|\mathbf{v}_b\|^{\frac{1}{m}-1} \mathbf{v}_b \quad (1)$$

where $\boldsymbol{\tau}_b$ is the basal drag (the bed-parallel component of the bed traction), β^2 is the drag coefficient, \mathbf{v}_b is the basal velocity, and we here assume the commonly-used value of $m = 3$. We note that the impact of the stress exponent m on modelled changes in ice flow have been studied previously in a number of studies, e.g. a recent study on the drivers of change over Pine Island Glacier by De Rydt et al. (2021).

Each model performs its own inversion procedure to infer the spatial distributions of the basal drag coefficient β^2 , and an ice rheological parameter, commonly denoted as A , in Glen’s flow law, using 1995 data (DEM and ice velocity; Bamber, 2000; Mougint et al., 2014) and ice temperatures calculated by other studies (Seroussi et al., 2019; Van Liering & Pattyn, 2013). All three models invert for β^2 over grounded ice. \acute{U} a and STREAMICE invert for the ice rheological parameter A over the entire domain (with STREAMICE penalizing variations from a ‘prior’ temperature-based estimate in grounded ice), while ISSM inverts for A only on floating ice. Technical details of the model inversions are described in Barnes et al. (2021). The resulting spatial distributions, i.e., $\beta^2(x, y)$ and $A(x, y)$, are kept constant over the transient runs (except for the experiments where we manually decrease β^2 in specific areas, Sect. 2.3).

The models set the basal traction to a negligible value downstream of the 1995 grounding-line position, which helps to prevent the grounding line from advancing beyond its initial state. The grounding line is based on hydrostatic equilibrium (Seroussi et al., 2014) and is free to migrate in all experiments.

2.3 Numerical experiments

2.3.1 Control experiment

We first run a ‘control’ simulation forced by ice-shelf melting only. None of the observed changes in geometry are imposed during the transient runs, and the ice-shelf thickness and the position of the grounding line are therefore free to evolve in response to this ice-shelf melting. Note that in this control simulation the models may not necessarily

160 reproduce the observed ice-shelf thinning and grounding line retreat since those are un-
 161 constrained in this experimental setup.

Basal melting under floating ice, m_b (in m/yr, positive if melting), is described by a depth-dependent function (see Fig. S1 and S2 in Seroussi et al., 2017):

$$m_b = \begin{cases} 50/500 |z_b|, & \text{if } 0 \geq z_b > -500 \text{ m,} \\ 50, & \text{if } z_b \leq -500 \text{ m,} \end{cases} \quad (2)$$

162 where z_b (in m) is the ice-shelf base depth (negative if below sea level).

163 We apply melt only to elements/cells containing fully floating ice (Seroussi & Morlighem,
 164 2018). The parameters in Eq. 2 are kept fixed during the simulations, although the spa-
 165 tial distribution of basal melting varies in time with the evolving thickness and extent
 166 (due to grounding line migration) of the ice shelf.

167 *2.3.2 Imposed change experiments overview*

168 Since the grounding line and ice-shelf thickness evolve freely in the control exper-
 169 iments, we expect that the agreement with observed speed and thickness change will be
 170 improved if we constrain front geometry, ice-shelf thinning, and loss of basal drag in the
 171 models, all according to observed changes. Prescribing these changes individually allows
 172 their respective impacts on the modeled evolution to be quantified and compared. For
 173 example, if applying the observed change in ice-front position improves the agreement
 174 with observations, it would suggest that calving dynamics played a role in the behav-
 175 ior of Thwaites. Comparing the response of the models will also shed light on the pro-
 176 cesses that have the strongest effect on Thwaites’s dynamics.

177 We run simulations where we apply observed changes to the control setup in (i)
 178 ice-front position, (ii) ice-shelf thickness, and (iii) basal traction downstream of the 2015
 179 grounding line, all based on observations. We force the models to follow these changes
 180 individually by prescribing these observed changes directly in the models. We also run
 181 (iv) an all-external-drivers experiment, where (i), (ii), (iii), and melt are all applied. We
 182 expect the results of experiment (iv) to be more consistent with observations. All ex-
 183 periments employ basal melt as given by Eq. 2, unless otherwise specified.

184 *2.3.3 Ice-front change experiment*

185 The observed retreat and rift propagation on Thwaites’s floating ice are imposed
 186 on a yearly basis at the ice-ocean boundary, following Landsat imagery (MacGregor et
 187 al., 2012). Any dynamic effect from ice-shelf rifting or collapse is captured in this sim-
 188 ulation. We apply these changes only to regions downstream of the 1995 grounding-line
 189 position. We keep TEIS and TT mechanically connected until 2005, following the hy-
 190 pothesis mentioned in Sect. 2.1. To this end, we remove the rift between TEIS and TT
 191 from the Landsat-derived ice-front contours for all years between 1996 and 2005. From
 192 2006 to 2015, we impose the original Landsat-based contours, disconnecting TEIS and
 193 TT (MacGregor et al., 2012). We do not consider any healing of that link after 2006 (Mouginot
 194 et al., 2014; Miles et al., 2020), since TT calved off in 2010/2011 (MacGregor et al., 2012).
 195 The basal melt is applied to all floating ice.

196 *2.3.4 Ice-shelf thinning experiment*

197 The 1995 shelf thickness is manually decreased according to satellite-measured thin-
 198 ning rates (e.g., Paolo et al., 2015; Smith et al., 2020). The 1995 shelf thickness is pro-
 199 portionally changed at each time step from 1995 to 2015. This setup simulates the ef-
 200 fect of decreasing ice-shelf buttressing on grounded ice following the observed shelf thin-
 201 ning. The imposed thinning ‘overrides’ melting except in newly ungrounded ice where

the thinning is not applied, i.e., the melt is only applied to areas upstream of the 1995 grounding line that become ungrounded during the transient runs. Imposing the thinning manually recovers the observed shelf-thickness change, which would probably not be perfectly reproduced by our parameterized basal melt in the control experiment (see Fig. S3 and S6).

2.3.5 *Loss of basal traction experiment*

Due to the lack of spatial and temporal data availability and the technical challenge of preserving hydrostatic equilibrium at the grounding line, we cannot directly prescribe the grounding line positions in transient runs. Instead, we simulate the effect of observed grounded ice retreat by linearly decreasing β^2 with time from its 1995-inverted value to zero between 1995 and 2015. The value of the basal drag coefficient, β^2 , is reduced only in the region that was grounded in 1995 but floating in 2015 (Rignot et al., 2014). This setup simulates a thinning-induced reduction in basal traction: as the ice approaches flotation, the effective pressure declines, reducing the basal traction. Note that the control experiment would not necessarily be reproducing this physical process and the observed grounding line retreat. The basal melt applies to all floating ice.

2.3.6 *All-external-drivers experiment*

The setup imposes the three observed changes together. These changes are the same as those imposed in the experiments described in Sects. 2.3.3, 2.3.4, 2.3.5. The melt is applied only to areas upstream of the 1995 grounding line that become floating over the simulations.

3 Results

To assess the relative impact of each imposed change on the models, we compute correlations and root mean square errors (RMSE) between observed and modeled speed and thickness change. The correlations measure the agreement between spatial patterns, while the RMSE quantifies the magnitude of the misfits.

3.1 Ice velocity changes

The observed acceleration of Thwaites Glacier has not been spatially uniform (Mouginot et al., 2014). The ice velocity increased by up to 25 m/yr per year in the vicinity of the grounding line (Fig. 1, panel b). Most of the main trunk and the western ice stream have been accelerating markedly up to 100 km upstream of the glacier's margin. The eastern part of the glacier has not changed significantly, except around the eastern margin of TEIS (see L1 in Fig. 2, panel a1) which accelerated by up to 20 m/yr per year. Most of the regions that sped up coincide with the regions where the grounding line retreated during this period (Fig. 1, panel b). Only a small area at the terminus of the western ice stream (see L2, Fig. 2, panel a1) decelerated between 1995 and 2015. At this location, the grounding line has not changed since the 1990s. The ice flux at the 2011 grounding line of Thwaites increased $30\text{-}33\pm 5$ Gt/yr over the 1994/1996-2013 period (Mouginot et al., 2014) (see Table S2).

The control experiment produces grounding line retreat and inland speed-up with moderate correlations in comparison to other experiments (Fig. 2, panels b1, b2, and b3, and Table 1). The modeled increase in ice flux is also comparable to observations and varies from 20 Gt/yr (ISSM) to 30 Gt/yr (STREAMICE) (Table S2).

Applying observed changes individually produces some differences among the model responses, with some simulations producing substantially better improvements in misfit and correlation for some models (e.g., ice-front change for Úa and ice-shelf thinning

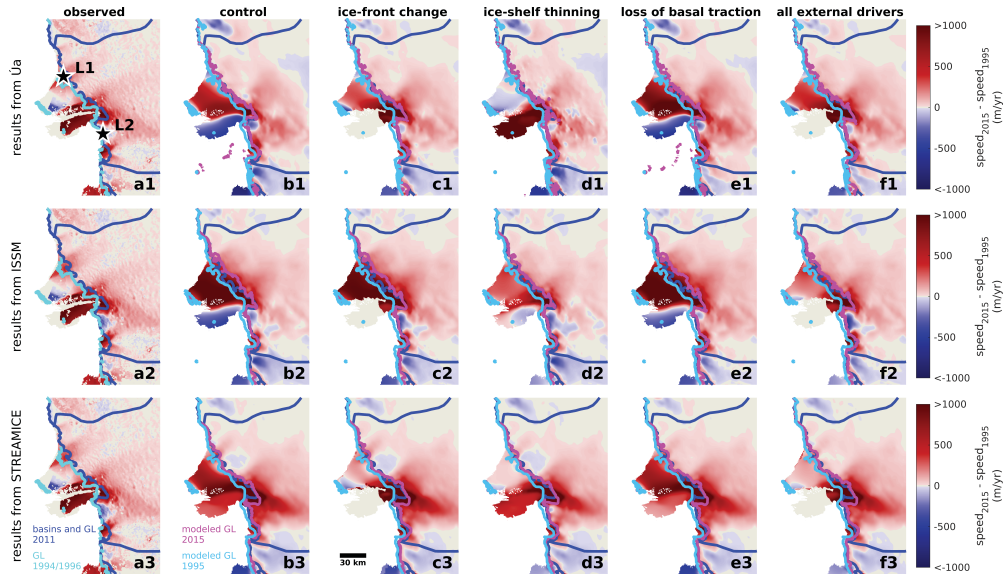


Figure 2. Speed changes of Thwaites Glacier obtained by the transient numerical experiments described in Sect. 2.3. The speed changes are obtained by subtracting the initial speed (1995) from the final speed (2015). Three ice sheet models are used: (b1-f1) \dot{U}_a , (b2-f2) ISSM, and (b3-f3) STREAMICE. (a1-a3) observed speed change (panels a1, a2, and a3 show the same map, for comparison purposes). (b1-b3) control experiment. (c1-c3) ice-front change experiment. (d1-d3) ice-shelf thinning experiment. (e1-e3) loss of basal traction experiment. (f1-f3) all-external-drivers experiment. The panels show the initial and final grounding line positions (from InSAR data and from the models, see the legends portrayed on panels a3 and b3, respectively). In (a1), black stars highlight two regions of localized speed changes (L1 and L2). Location 1 (L1) indicates the acceleration of the eastern margin of TEIS, and location 2 (L2) indicates a region that decelerated in the margin of the western ice stream.

Table 1. Correlations and root mean square errors (RMSE) between modeled and observed speed and thickness changes obtained by the experiments described in Sect. 2.3. The correlations and RMSE for thickness change are calculated using two different data sets: (a) from Shepherd et al. (2019), whose time period is 1992-2017, and (b) from Smith et al. (2020), whose time period is 2003-2019. We apply a Gaussian filter to the modeled thickness changes with kernel size equal to 35 km and $\sigma = 0.20$, following Smith et al. (2020). All correlation coefficients and RMSE are obtained considering Thwaites’s basin and inland extension as given by transect A-A’ portrayed on panel b of Fig. 1. Legend: CR (control experiment), IF (ice-front change experiment), IS (ice-shelf thinning experiment), BT (loss of basal traction experiment), and AD (all-external-drivers experiment).

	CR	IF	IS	BT	AD	CR	IF	IS	BT	AD
Speed change	Correlation					RMSE (m/yr)				
\dot{U}_a	0.58	0.74	0.55	0.67	0.77	17.53	10.47	17.39	17.14	10.29
ISSM	0.58	0.69	0.61	0.69	0.74	14.46	14.17	12.19	15.35	10.37
STREAMICE	0.70	0.72	0.76	0.72	0.73	25.36	26.98	18.91	24.72	24.08
Thickness change (a)	Correlation					RMSE (m)				
\dot{U}_a	0.79	0.85	0.71	0.84	0.88	4.20	4.50	4.33	4.53	4.81
ISSM	0.84	0.86	0.81	0.90	0.89	5.47	5.95	3.83	6.72	5.00
STREAMICE	0.82	0.81	0.80	0.86	0.84	6.67	7.81	6.69	7.49	7.94
Thickness change (b)	Correlation					RMSE (m)				
\dot{U}_a	0.78	0.87	0.71	0.83	0.89	4.86	3.53	5.69	4.55	3.48
ISSM	0.79	0.82	0.80	0.86	0.88	4.71	4.80	3.75	5.10	3.55
STREAMICE	0.86	0.86	0.85	0.89	0.88	4.53	5.53	4.59	5.10	5.52

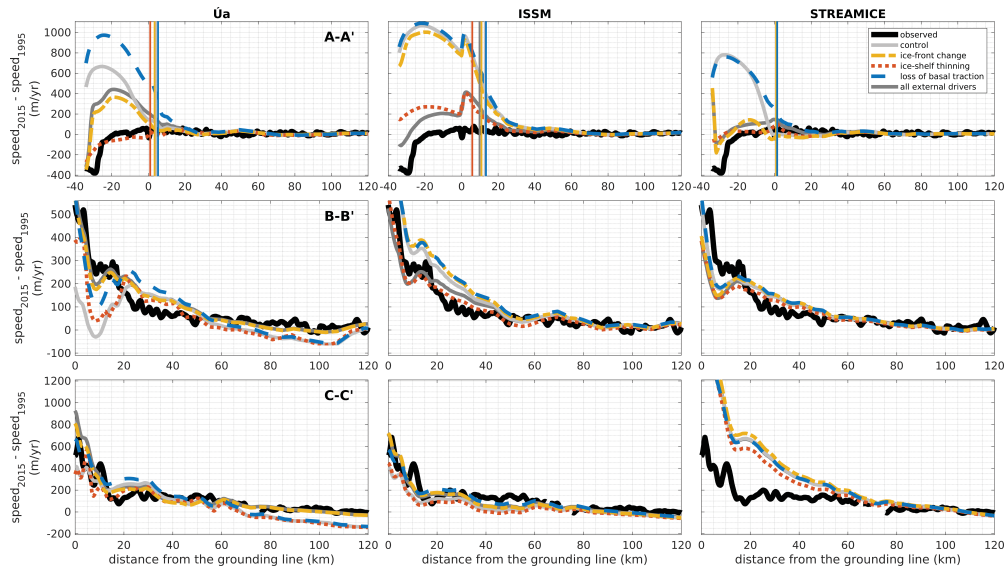


Figure 3. Speed changes at the end of the transient experiments (2015) over flow lines: A-A', B-B', and C-C'. The experiments are described in Sect. 2.3. The InSAR-derived speed change is in black. The distance from the grounding line refers to the 2015-observed grounding line position. Vertical lines shown on transect A-A' panels are the grounding line positions at the end of the experiments (2015). All flow lines' locations are shown in Fig. 1.

248 for STREAMICE) than for others. In $\acute{U}a$, prescribing ice-shelf thinning appears to in-
 249 troduce numerical inconsistencies at the boundary of the area for which thickness is en-
 250 forced, which explains the differences in the shape of the grounding line compared to other
 251 experiments (Fig. 2, panel d1).

252 As expected, the all-external-drivers experiment reproduces the overall pattern of
 253 observed speed change with the least error, including the localized changes at L1 and
 254 L2 (Fig. 2, panels f1, f2, and f3). Changes in L2 are not captured by STREAMICE be-
 255 cause the grounding line retreats in this region with this model. Overall, the final mod-
 256 eled grounding line positions obtained with the all-external-drivers experiment are also
 257 in good agreement with observations. STREAMICE overestimates grounding line retreat
 258 along the western side of the grounding line 'bight' in front of western Thwaites (location
 259 'A' of Milillo et al., 2019). This overestimated retreat coincides with overestimated ac-
 260 celeration in the western part of Thwaites (Fig. 2, panels from b3 to f3, and Fig. 3, tran-
 261 sect C-C'), which contributes to the higher RMSE and glacier flux compared to $\acute{U}a$ and
 262 ISSM (Tables 1 and S2, respectively). The increase in glacier flux varies from 15 Gt/yr
 263 (ISSM) to 30 Gt/yr (STREAMICE) (Table S2).

264 3.2 Ice thinning

265 The observed thinning of Thwaites followed the observed ice speed-up, extending
 266 tens of kilometers over the interior of the glacier (Fig. S9, panel a1). The margins of the
 267 main trunk and western ice streams as well as the eastern part of TEIS (L1) thinned the
 268 most (up to 45 m over the 2003-2019 period; Smith et al., 2020). To compare with our
 269 results, we use two different datasets of observed thickness changes (Shepherd et al., 2019;
 270 Smith et al., 2020), since they employ an acquisition period different from the period con-
 271 sidered here (see Table 1).

272 The three ice sheet models produce patterns of ice thinning similar to observations
 273 in all experiments, as seen by the relatively high correlations (Table 1), although some
 274 thickening appears in regions around the glacier’s basin likely due to thickness adjust-
 275 ments during the transient runs. The inland thickening produced by Úa in the ice-shelf
 276 experiment is likely caused by the numerical inconsistencies described above.

277 The control experiment generates a pattern of thickness change comparable to ob-
 278 servations, as noted by the relatively high correlation of ~ 0.81 . The thinning rates ob-
 279 tained with all-external-drivers and loss of basal traction experiments show the highest
 280 correlation coefficients (~ 0.88 and ~ 0.86 , respectively), followed by ice-front change (~ 0.85)
 281 and ice-shelf thinning (~ 0.78) experiments.

282 4 Discussion

283 The control experiment produces upstream acceleration, thinning, and grounding
 284 line migration comparable to observations. Our parameterized melt is not necessarily
 285 in balance with the ice-shelf flow (e.g., Rignot et al., 2013), which produces some shelf
 286 thinning (Fig. S6). As a consequence, the melt sustains thinning of newly-ungrounded
 287 ice upstream of the 1995 grounding line as downstream changes (i.e., shelf thinning) in-
 288 duce a loss of buttressing on upstream flow (e.g., in the vicinity of the grounding line
 289 along the main trunk and the western ice stream), causing inland speed-up, thinning,
 290 and grounding line retreat. This mechanism may be enhanced by local reverse-slope bedrock,
 291 where the grounding line may retreat faster (Schoof, 2007; Joughin et al., 2014; Rignot
 292 et al., 2014; Seroussi et al., 2017; Morlighem et al., 2020). The observed mass loss over
 293 the last decades in the Amundsen Sea Sector is therefore likely associated with increas-
 294 ing ocean-induced melt (e.g., Pritchard et al., 2012; Joughin et al., 2014; Jenkins et al.,
 295 2016; Seroussi et al., 2017; Martin et al., 2019; Milillo et al., 2019; Hoffman et al., 2019;
 296 Robel et al., 2019).

297 In the experiments where we impose observed changes instead of letting the model
 298 evolve freely, we find that forcing the geometry of the ice shelf and basal traction increases
 299 the correlations in both speed change and ice thinning. These findings suggest that rift-
 300 ing propagation between TEIS and TT and thinning-induced reduction in basal trac-
 301 tion play an important role on Thwaites’s dynamics, as pointed out by previous stud-
 302 ies (e.g., Mouginot et al., 2014; Miles et al., 2020; Joughin et al., 2014; Nias et al., 2016).
 303 Given its importance, the evolution of basal drag as the grounding line retreats may there-
 304 fore need to be further improved in ice-sheet models (Nias et al., 2016; De Rydt et al.,
 305 2021). For instance, we employ here Weertman’s sliding law with an exponent $m = 3$
 306 and we invert for the drag coefficient (β^2). The resulting spatial distribution of β^2 is then
 307 kept fixed in all simulations (except for the loss of basal traction setup). Other sliding
 308 laws that reduce the basal traction as the grounding line migrates could potentially gen-
 309 erate a different upstream response to ice thinning/front retreat (Brondex et al., 2017,
 310 2019; Joughin et al., 2019; De Rydt et al., 2021). Also, it remains unclear whether a ‘me-
 311 chanical link’ between TEIS and TT could be reinstated in the future, or whether the
 312 mechanical integrity of TEIS will be compromised due to structural weakening (e.g., Miles
 313 et al., 2020). Thus, enhanced calving dynamics may also improve the accuracy of nu-
 314 merical simulations of Thwaites (e.g., Crawford et al., 2021).

315 Our parameterized basal melt is based on observations and ocean simulations, and
 316 similar parameterizations have been used in other studies of Thwaites Glacier (Depoorter
 317 et al., 2013; Rignot et al., 2013; Milillo et al., 2019; Seroussi et al., 2017; Nakayama et
 318 al., 2018, 2019; Joughin et al., 2014; Hoffman et al., 2019). At the end of the control ex-
 319 periments, the integrated melt is, on average, 110 Gt/yr, which is slightly greater than
 320 satellite-based (97.5 ± 7 Gt/yr, Rignot et al., 2013) and simulation-based (80-120 Gt/yr,
 321 Seroussi et al., 2017) estimates (see Table S1 and Fig. S10). Depth-dependent melt pa-
 322 rameterizations tend to overestimate grounding line retreat in comparison to ice-ocean

323 simulations in longer runs (Seroussi et al., 2017). The parameters in Eq. 2 were kept con-
 324 stant over the transient runs, although it is likely that ocean conditions have changed
 325 over the last decades, which could have affected the response of the models. For instance,
 326 rerunning the control experiment in ISSM with the parameterized melt multiplied by 4,
 327 the model overestimates the inland acceleration (the resulting RMSE is 28.58 m/yr. See
 328 also Fig. S5, panel b) and the integrated melt along the entire transient run (Fig. S10),
 329 although the spatial pattern of the response is similar to observations (correlation of 0.78).
 330 To improve the forecast of Thwaites’s future, reliable estimates of melt rates are required,
 331 especially close to the grounding line, where thinning-induced reduction in basal trac-
 332 tion is critical.

333 The differences between the models’ results may be caused by several factors: stress
 334 balance approximation, inversion procedure, mesh resolution, numerical issues caused
 335 by imposed forcings, etc. (e.g., Cornford et al., 2020; Barnes et al., 2021). STREAMICE
 336 had more extensive retreat in the western part of the Thwaites grounding line than that
 337 of Úa or ISSM, which might be the reason for larger acceleration in this region and re-
 338 sulting higher RMSE. The difference may arise from a differing treatment of a small ice
 339 rise in TT arising from a topographic high in the bathymetry data (see the supporting
 340 information), or from resolution in the vicinity of the grounding line which may be too
 341 coarse (Cornford et al., 2020). A deceleration can be seen upstream in some Úa results,
 342 particularly evident across transect F-F’ in Fig. S8, due to slight differences in the in-
 343 verted basal sliding and rate factor fields compared to the other models. These factors
 344 all play an important role in transient simulations and shall be investigated in future work.
 345 Our results also illustrate how challenging reproducibility is in the field of ice sheet mod-
 346 eling (e.g., Seroussi et al., 2020), which calls for further numerical developments and model
 347 inter-comparison initiatives.

348 Uncertainties in the data and inversion procedures may have an impact on our re-
 349 sults. For example, the mass loss observed from over the last decades could be part of
 350 an already existing dynamic imbalance prior to 1995, and our inversions were not able
 351 to capture this early loss trend. Also, using a previous bed elevation version (BedMa-
 352 chine v1), artificial ‘bumps’ downstream of the 2015 grounding line (and close to the 1995
 353 grounding line) prevented inland acceleration and grounding line retreat in most of the
 354 experiments. These results highlight the need for further improvements in bed topog-
 355 raphy data, as noted by others (e.g., Durand et al., 2011; Nias et al., 2016).

356 5 Conclusions

357 By conducting time-dependent numerical simulations of Thwaites Glacier between
 358 1995 and 2015 with three independent ice sheet models and several satellite-based datasets,
 359 we find that thinning induced by ocean melting and the resulting grounding line retreat
 360 explain much of the observed speed-up of Thwaites. The models also suggest that changes
 361 in the ice-shelf geometry, especially the rifting propagation between the Eastern Ice Shelf
 362 and Thwaites Tongue, improve the agreement with observations. The results suggest that
 363 improved forecasts of Thwaites’s future require reliable ocean-induced melt estimates and
 364 improved model response to downstream perturbations, particularly thinning-induced
 365 reduction in basal traction.

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 371 codes used in this study are publicly available. The Ice-sheet and Sea-level System Model
 372 can be accessed at <https://issm.jpl.nasa.gov> (we used version 4.18). STREAMICE

373 is a module of MITgcm and can be download at <https://mitgcm.org/source-code/>.
 374 The source code of Ua can be downloaded at <http://doi.org/10.5281/zenodo.3706623>.
 375 BedMachine Antarctica is available at NSIDC (<http://nsidc.org/data/nsidc-0756>).
 376 The 1995 Antarctic 5 km DEM is also available at NSIDC (<https://nsidc.org/data/nsidc-0076>).
 377 InSAR-Based ice velocity of the Amundsen Sea Embayment is found at
 378 NSIDC (<https://nsidc.org/data/NSIDC-0545>). Coastal and terminus history of Thwaites
 379 (shapefile format) is found at NSIDC (<https://nsidc.org/data/NSIDC-0522>). Ground-
 380 ing line positions are found at NSIDC (<https://nsidc.org/data/NSIDC-0498>). Maps
 381 of thickness changes are found at Nasa’s EarthData (https://sealevel.nasa.gov/data/dataset/?identifier=SLCP_ice_shelf_dh_v1_1),
 382 University of Washington’s digital repository (<https://digital.lib.washington.edu/researchworks/handle/1773/45388>),
 383 and at CPOM (<http://www.cpom.ucl.ac.uk/csopr>). The Antarctic surface mass bal-
 384 ance (RACMO 2.3) is available at [https://www.projects.science.uu.nl/iceclimate/](https://www.projects.science.uu.nl/iceclimate/models/antarctica.php)
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Figure 1.

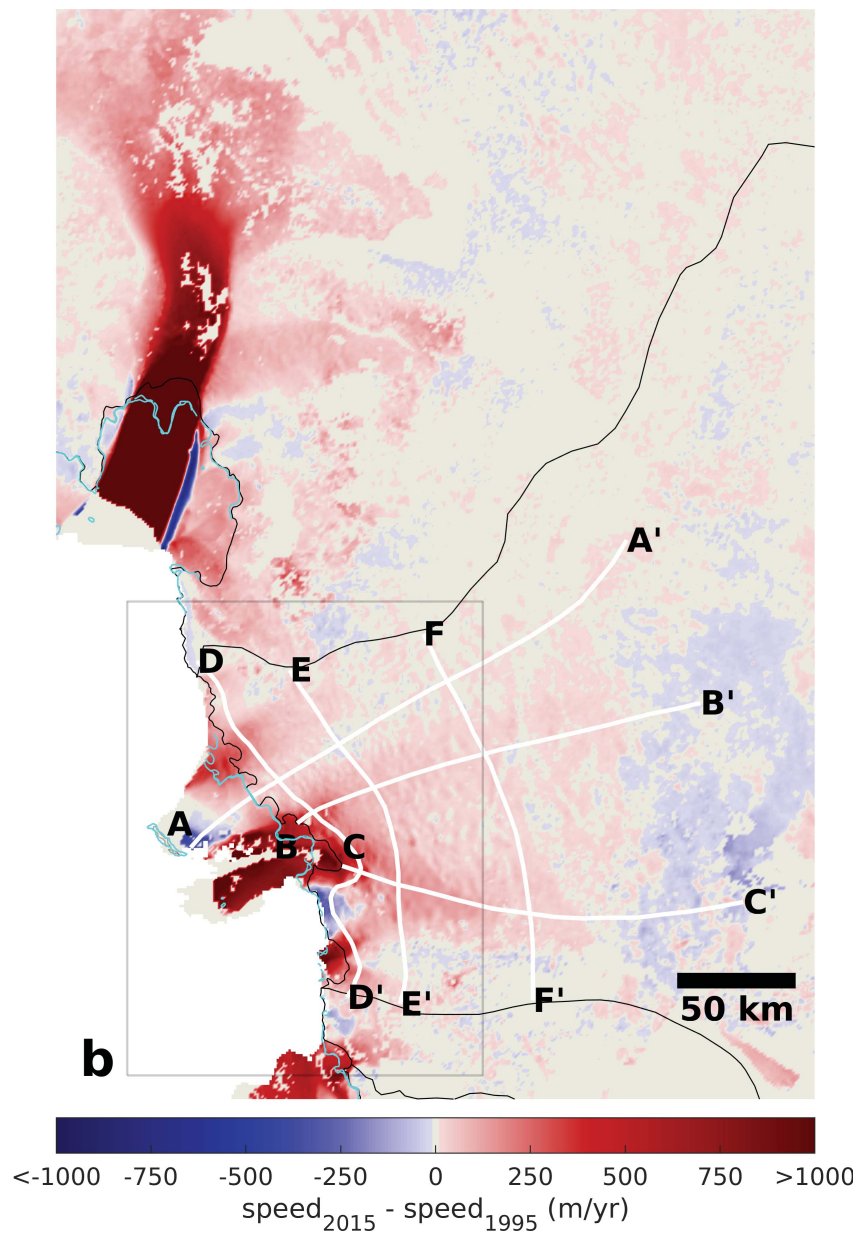
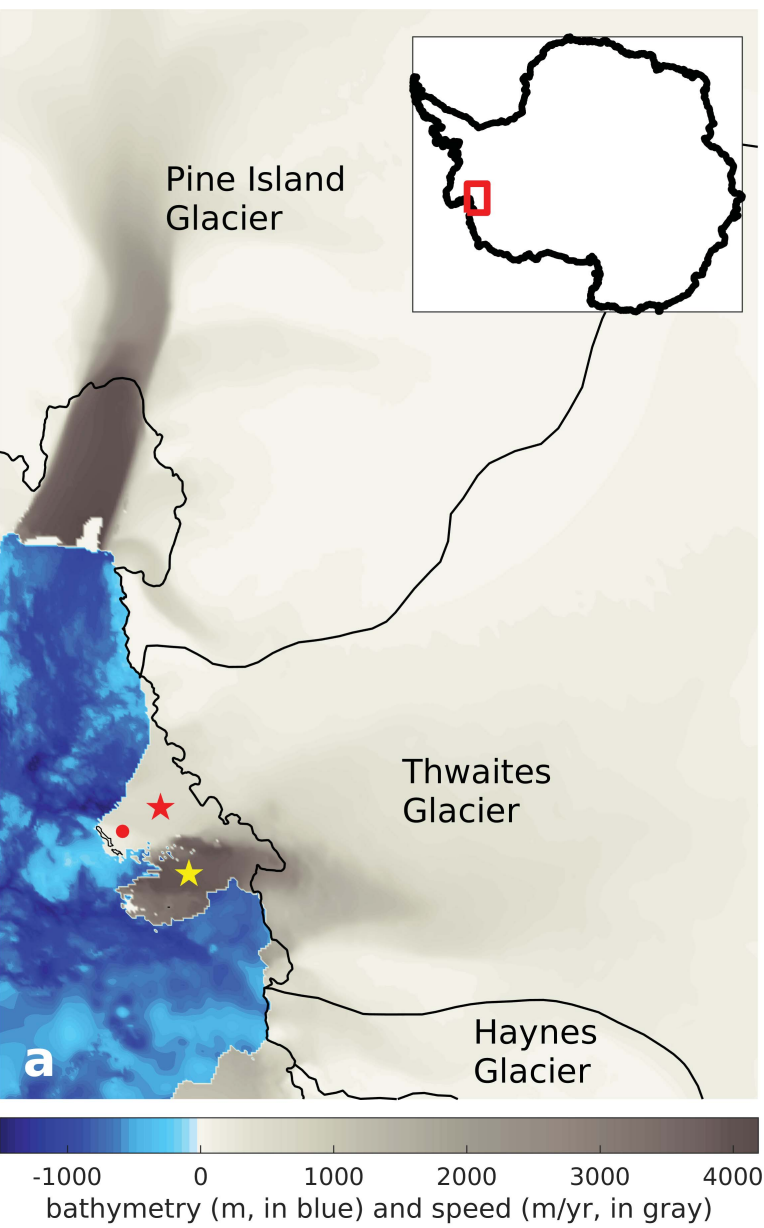
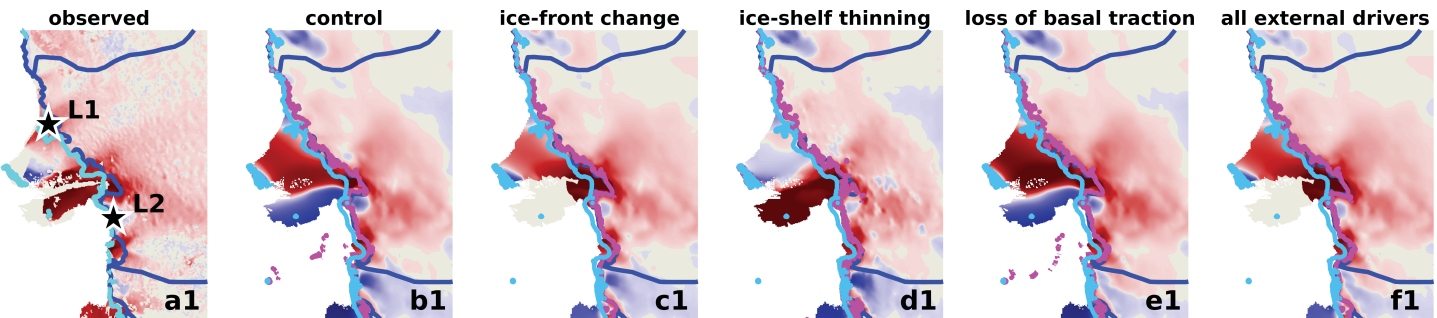
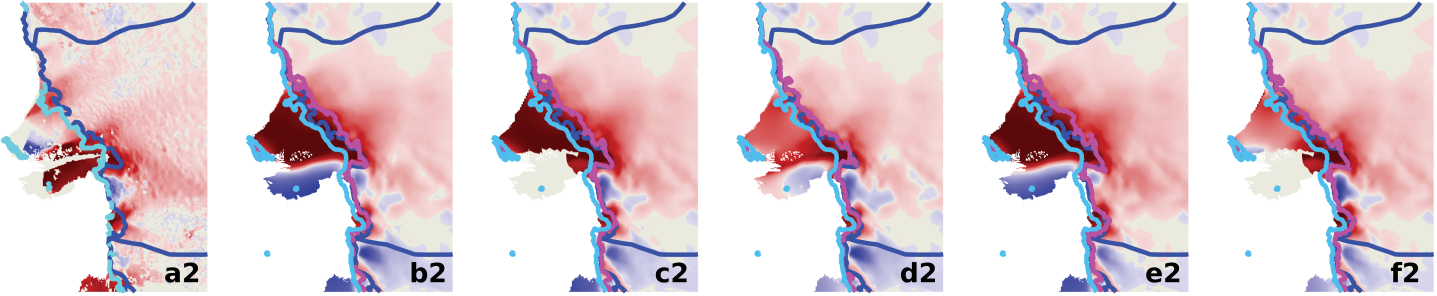


Figure 2.

results from Ua



results from ISSM



results from STREAMICE

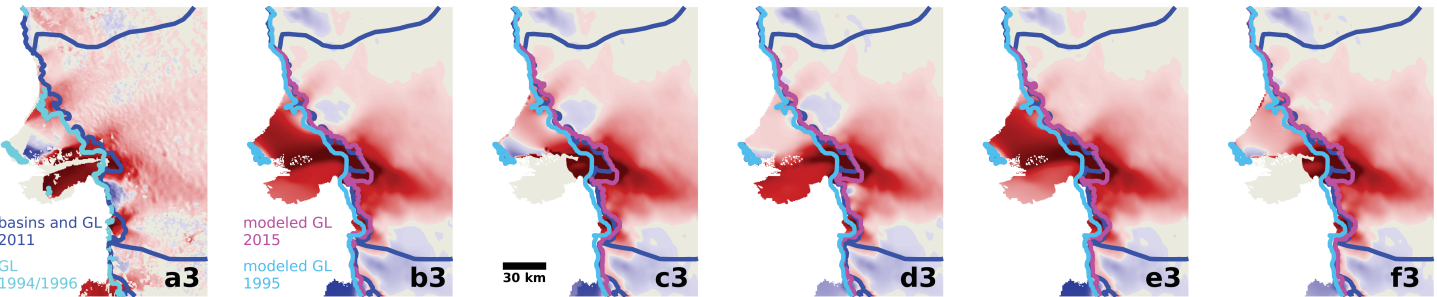


Figure 3.

