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Citation for published version:

Bao, Z, Tseng, YJ, You, W, Zheng, W, Chen, X, Mahlik, S, Lazarowska, A, Lesniewski, T, Grinberg, M, Ma, C, Sun, W, Zhou, W, Liu, RS & Attfield, JP 2020, 'Efficient Luminescence from CsPbBr, Nanoparticles Embedded in Cs, PbBr, ', *The Journal of Physical Chemistry Letters*, vol. 11, no. 18, pp? 7637-7642. https://doi.org/10.1021/acs.jpclett.0c02321

Digital Object Identifier (DOI):

10.1021/acs.jpclett.0c02321

Link:

Link to publication record in Edinburgh Research Explorer

Document Version: Peer reviewed version

Published In: The Journal of Physical Chemistry Letters

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Efficient Luminescence from CsPbBr₃ nanoparticles embedded in Cs₄PbBr₆

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TOC GRAPHICS



Keywords: CsPbBr₃; Cs₄PbBr₆; embedded; superfluorescence; nanocrystals

In recent years, lead halide perovskite materials have become regarded as a new generation of functional materials for solar cells, lasers, photodetectors and light-emitting diodes (LEDs).¹⁻⁶ Perovskite nanocrystals (NCs) are also considered to have great potential for backlight displays because of their excellent photoluminescence (PL) performance and quantum yield.⁷⁻¹¹ Although perovskite NCs have been successfully demonstrated in light-emitting devices, poor stability is the main limitation for commercial applications,^{2, 12} and issues of thermal stability, photostability, water-resistance, and anion exchange need to be overcome.

One approach to solve these problems is to reduce the structural dimension.¹³⁻¹⁷ In zerodimensional perovskite-related Cs₄PbBr₆, strong green PL emission can be observed dependent on particle size.¹⁴ The PL has been attributed to Br defects in the crystals,¹⁸⁻²¹ but other studies²² discovered CsPbBr₃ nanoparticles embedded in Cs₄PbBr₆ and proposed that the emission of green light comes from these rather than the Cs₄PbBr₆ host.^{21, 23} However, further studies have suggested that the embedded CsPbBr₃ phase may tune an impurity mechanism for PL of Cs₄PbBr₆,²⁴ and the mechanism of green emission from Cs₄PbBr₆ remains under debate.²⁵ In this research, we have further explored the emission from Cs₄PbBr₆ and we have also tuned reagent ratios while synthesizing Cs₄PbBr₆ to explore the PL behavior and phase transform process, and we demonstrate that the embedded CsPbBr₃ nanoparticles are the origin of green emission in Cs₄PbBr₆ crystals.

Polycrystalline Cs₄PbBr₆ was synthesized by the antisolvent method with the Cs/Pb precursor ratio of 3.33, and CsPbBr₃ QDs for comparison were synthesized through a hot-injection method. Detailed processes are presented in the Methods section. Powder synchrotron X-ray diffraction (XRD) data in Figure. 1a shows that our Cs₄PbBr₆ crystals are single phase and the

previously reported rhombohedral structure (crystal open database; COD No.1538416) fits the data well with lattice parameters a = 13.72939(15) and c = 17.31874(25) Å (Table S1).

HRTEM, on the other hand, revealed many embedded dots (black spots in each particle) in the Cs₄PbBr₆ crystals, as shown in Figure. 1b. High-quality HRTEM images from this beam sensitive material have been recorded, showing the lattice fringes of the Cs₄PbBr₆ crystals as well as the embedded dots. As exhibited in Figure. 1c, the fringes "A" of the parent crystal have a *d*-spacing of 3.99 Å, corresponding to the (300) planes of the rhombohedral structure of Cs₄PbBr₆, while the fringes "B" in a dark dot have a *d*-spacing of 2.90 Å, which can be indexed to the (200) planes of the cubic CsPbBr₃ structure. The inset of Figure. 1c shows a size distribution of the embedded CsPbBr₃ nanoparticles based on a measurement of randomly selected 200 dots. Almost all these dots are smaller than 10 nm and their average size around 3.8 nm in diameter. The dark contrast of these embedded dots is mainly dominated by the mass contrast due to a relatively higher concentration of Pb in CsPbBr₃.

The [111] direction of the CsPbBr₃ nanoparticles observed by HRTEM is coincident with [001] of Cs₄PbBr₆ host crystals and the hexagonal symmetries of these phases match well in the (001) planes of Cs₄PbBr₆ with the [100] direction of Cs₄PbBr₆ parallel to [110] of CsPbBr₃. Consequently, the ideal inter-axes angle between the [100] of Cs₄PbBr₆ and the [100] of CsPbBr₃ is 65.9 degrees. The corresponding measured angle between the marked fringes in Figure. 1c is 63 degrees. The structural projection along the [001] zone axis of Cs₄PbBr₆ and the [111] axis of CsPbBr₃ (as shown in Figure. 1d) with the [100] direction of Cs₄PbBr₆ parallel to the [110] of CsPbBr₃ gives a good lattice match (a mismatch of +4.7% between d₁₀₀(Cs₄PbBr₆) = 11.89 Å and 3d₁₁₀(CsPbBr₃) = 12.45 Å). Similarly in the [001] axis direction of Cs₄PbBr₆, the d₀₀₁(Cs₄PbBr₆) = 17.32 Å has only a -2.1% mismatch to 5d₁₁₁(CsPbBr₃) = 16.945 Å. Although the lattice matching

is good, the slight mismatch will create a strain field in the host material that likely acts to limit the size of the CsPbBr₃ particles and keep them well separated, which is important to the optical properties as discussed below. Thus they can be embedded epitaxially into Cs₄PbBr₆ crystals.



Figure 1. Structural characterization of the Cs₄PbBr₆ crystals. a, Refined synchrotron XRD (beamline energy 20 keV) pattern of the specimen. b, TEM image of Cs₄PbBr₆ crystals with embedded CsPbBr₃ crystals as dark dots. c, HRTEM image of a Cs₄PbBr₆ crystal with embedded CsPbBr₃ dots. The fringes marked A (d = 3.99 Å) are in the parent Cs₄PbBr₆ crystal region and

that marked B (d = 2.90 Å) in an embedded dot can be indexed to the CsPbBr₃ structure. The inset of (c) shows a size distribution of the embedded CsPbBr₃ dots, matching to a log-normal distribution curve. The arrow points to an area of slight crystal damage. d, Model showing how a CsPbBr₃ nanoparticle is embedded epitaxially within the Cs₄PbBr₆ host lattice, projected on the (001) plane of the latter.

To explore the optical behavior of Cs_4PbBr_6 crystals, we carried out temperature-dependent steady-state PL spectroscopy measurements (10-300 K). As shown in Figure. 2a, the crystals exhibited an intense green emission band in the spectral region from 490 nm to 550 nm upon above-bandgap excitation at 274 nm. When the temperature rose from 10 K to 300 K, the intensity of the emission band decreased greatly by a factor 6.7, along with an obvious blue-shift in peak wavelength from 530.5 nm to 520.0 nm. Such temperature-dependent PL behavior coincides very well with that of CsPbBr₃ quantum dots (QDs) previously reported,²⁶⁻²⁷ suggesting that the green PL of Cs₄PbBr₆ crystals is dictated by the emission of CsPbBr₃ embedded in the lattice of Cs₄PbBr₆. Figure. 2b shows the temperature-dependent PL excitation spectra of Cs₄PbBr₆ crystals by monitoring the green emission at 524 nm, which displays broad excitation bands from 250 nm to 500 nm with "hole burning" at around 310 nm. The excitation band with the energy below 3.82 eV and above 4.13 eV are ascribed to the absorption of CsPbBr₃, while the temperature-dependent "hole burning" at 310 nm agrees well the localized exciton absorption of Cs₄PbBr₆ crystals.²⁸ This can be further verified by the temperature-dependent PL emission spectra of Cs4PbBr₆ crystals upon excitation at 310 nm, as shown in Figure. 2c. As the temperature fell below 200 K, the green emission of CsPbBr₃ was detected along with an ultra-violet emission with a central wavelength at 375 nm. The emission band at 375 nm can be assigned to the ${}^{3}P_{0} \rightarrow {}^{1}S_{0}$ transition of Pb²⁺ ion in isolated [PbBr₆]⁴⁻ octahedra.²⁹⁻³¹ Meanwhile, it was found that both the CsPbBr₃ and Pb²⁺

emissions weakened significantly with the temperature rise and nearly vanished when the temperature was higher than 200 K, due to the increased thermal quenching at higher temperatures. 10 K PL excitation spectrum of the crystals by monitoring the Pb²⁺ emission at 375 nm shows a strong excitation band around 310 nm, ascribed to the ${}^{1}S_{0}\rightarrow{}^{3}P_{1}$ transition of Pb²⁺ ion in isolated [PbBr₆]⁴⁻ octahedra in Cs₄PbBr₆ (Figure. 2d). In summary, the temperature-dependent PL evolution for the green emission of Cs₄PbBr₆ crystals is very similar to that of CsPbBr₃ QDs.



Figure 2. Photoluminescence of the Cs₄PbBr₆ sample. a, Temperature-dependent PL emission spectra of Cs₄PbBr₆ upon excitation at 274 nm in the temperature range of 10–300 K. b, Normalized PL excitation spectra of Cs₄PbBr₆ as a function of temperature by monitoring the CsPbBr₃ emission at 524 nm. c, Temperature-dependent PL emission spectra of Cs₄PbBr₆ upon excitation at 310 nm in the temperature range of 10–300 K. d, 10 K PL excitation spectrum of Cs₄PbBr₆ by monitoring the Pb²⁺ emission at 375 nm.

To further examine the type of luminescence centers in the green-emitting Cs₄PbBr₆ material, the time-resolved emission spectra of Cs₄PbBr₆ crystals and dried CsPbBr₃ QD powders (for comparison) were compared in the temperature range of 10-300 K. Figure. 3 presents timeresolved emission spectra (streak images) of Cs₄PbBr₆ crystals (Figures. 3a, b, c) and CsPbBr₃ QDs (Figures. 3d, e, f) obtained at a time range of 1 ns. The room temperature streak images of both Cs4PbBr6 crystals and CsPbBr3 QDs systems show an emission band located around the same central wavelength (520 nm), however with different FWHM (~10 nm for Cs₄PbBr₆ crystals and ~30 nm for CsPbBr₃ QDs). As the temperature decreases the FWHM of both systems decreases, most significantly for CsPbBr₃ QDs which is accompanied by a shift of the emission center towards red (green curves in Figures. 3d, e, f). CsPbBr₃ QDs at the temperature lower than 250 K exhibit an additional type of luminescence, possessing emission wavelength slightly shifted towards the red (red curves in Figures. 3d, e). Under pulsed excitation, the intensity of the additional luminescence is much greater than the regular emission and the decay time of the luminescence is at least an order of magnitude shorter than the regular luminescence of CsPbBr₃ QDs. The precise determination of the decay time was not possible due to the finite duration of the excitation pulse (FWHM ~30 ps, see the pink curve in Figures. 3 a-f).

We interpret the additional luminescence as superfluorescence²⁸ – luminescence due to the collective emission of multiple QDs located close to each other (at distances below the wavelength of the emitted light) so light emission from a QD can induce synchronous emission from nearby QDs. As a result, the QDs emit light collectively due to induced emission, with a much greater radiative rate, which corresponds to a much shorter decay time (of the order of picoseconds). The streak images of low-temperature emission in CsPbBr₃ QDs (Figure 3d, e) show intense, short-time superfluorescence from QDs remaining in proximity (coupled QDs) as well as residual

luminescence from the QDs that are too far from other QDs to emit collectively (uncoupled QDs). The Cs₄PbBr₆ sample does not exhibit superfluorescence at any temperature which indicated that the QDs embedded in the Cs₄PbBr₆ matrix are too far apart to exhibit coupling. In contrast, the free CsPbBr₃ QDs tend to agglomerate easily leading to coupling and superfluorescence. High-pressure studies also confirmed that the green PL of Cs₄PbBr₆ originates from the emission of CsPbBr₃ embedded in the lattice of Cs₄PbBr₆ (shown in Supplementary Information). Hence, this study demonstrates that embedded CsPbBr₃ anoparticles in Cs₄PbBr₆ show different emission properties to aggregates of free CsPbBr₃ QDs.



Figure 3. Time-resolved emission spectra (streak images) at a time range of 1 ns. Cs_4PbBr_6 sample at 10K (a), 150 K (b), 293 K (c). $CsPbBr_3$ QDs at 10K (d), 150 K (e), 293 K (f). The green curve presents an integrated emission spectrum of light emitted in the time interval denoted by the green dashed lines. The red curve denotes the integrated emission spectrum of superfluorescence. The pink curve denotes the temporal shape of the excitation laser pulse.

It is important to find how synthesis conditions affect the formation of embedded CsPbBr₃ in Cs₄PbBr₆. Hence we have synthesized a series of materials with the ratio of Cs/Pb reactants varied in the range from 1 to 4.5 using the same synthesis procedure as for the original Cs_4PbBr_6 sample. From the XRD patterns, as shown in Figure. 4a, polycrystalline Cs₄PbBr₆ remains the only phase observed by XRD with the Cs/Pb ratio between 4.5 and 2.5. Peaks from the secondary phase of CsPbBr₃ begin to emerge at a Cs/Pb ratio of 2.25, and increase in intensity as the Cs/Pb ratio decreases to 1. Correspondingly, the external quantum efficiency (EQE) of the samples rises smoothly up to 28.4% as the Cs/Pb ratio decreased from 4.5 to 2.5, as the concentration of embedded CsPbBr₃ nanoparticles increases with falling Cs/Pb ratio, (Figure. 4b) but sharply decreases from 28.4% to 14.7% when the Cs/Pb precursor ratio reached 2.25, where secondary CsPbBr₃ phase starts to emerge (Figure. 4c). CsPbBr₃ has high photoluminescence if particle size is smaller than ~ 10 nm, but for Cs/Pb ratios below 2.5, the unconstrained secondary phase particles aggregate to larger sizes leading to their appearence in XRD patterns and the observed drop in overall EQE. These results thus confirm that the high-quality green PL emission of the Cs₄PbBr₆ is from embedded CsPbBr3 nanoparticles, while any secondary CsPbBr3 phase luminescences with a much lower EQE.



Figure 4. Properties of Cs₄PbBr₆ samples when tuning the Cs/Pb precursor ratio. a, XRD patterns, b, Normalized PL emission spectra, and c, EQE, of Cs₄PbBr₆ samples synthesized with different Cs/Pb precursor ratios, showing single and mixed phase regions as observed by XRD.

In summary, our results demonstrate that epitaxially embedded CsPbBr₃ nanoparticles are responsible for the efficient green emission from Cs₄PbBr₆ crystals. CsPbBr₃ impurity phase may also be present if samples are synthesized with a Cs/Pb precursor ratio lower than 2.5, but these have lower emission intensity. Our study shows that the properties of optically active nanoparticles are significantly altered by embedding them in a suitable inert host lattice, and this mechanism may enable further new efficient luminescent composites to be developed.

ASSOCIATED CONTENT

Supporting Information. The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsenergylett.XXXXXXX.

Complete experimental section, optical properties, further temperature-dependent PL properties, and pressure-dependent PL properties.

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Notes

The authors declare no competing financial interest.

ACKNOWLEDGMENT

This work was financially supported by the "Advanced Research Center of Green Materials Science and Technology" from The Featured Area Research Center Program within the framework of the Higher Education Sprout Project by the Ministry of Education (107L9006) and the Ministry of Science and Technology in Taiwan (MOST 107-2113-M-002-008-MY3, MOST 107-2923-M-002-004-MY3 and MOST 107-3017-F-002-001), the National Centre for Research and Development Poland Grant (No. PL-TW/V/1/2018), the Strategic Priority Research Program of the Chinese Academy of Sciences (XDB2000000), the CAS/SAFEA International Partnership Program for Creative Research Teams, and the NSFC (nos. U1805252 and 11774345). J.P.A. acknowledges financial support from EPSRC, UK.

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