# Focused beam dosimetry of short VHEE bunches

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## ABSTRACT

Accelerators driven by 10s TW-class lasers can produce electron bunches with femtosecond-scale duration and energy of 100s of MeV. A potential application of such short bunches is high-dose rate radiotherapy, which could transition to FLASH radiotherapy if a sufficiently large dose is delivered in a single shot. Here we present Monte Carlo simulations to study the bunch length evolution of an electron beam propagating in a water phantom. We show that for electron energies above 100 MeV the bunch lengthens to 1–10 ps duration after interaction with a 30 cm long water phantom, both for a collimated and weakly focused geometry. The corresponding dose rates are on the order of 200 Gy/s per primary electron, much higher than in conventional radiotherapy.

**Keywords:** focused electrons, radiotherapy

#### 1. INTRODUCTION

Laser-wakefield accelerators produce electron beams with energy from tens of MeV to almost  $10\,\text{GeV}^{1-5}$  in centimetre-scale distances. Simulations and experiments have shown that the duration of these beams is on the order of femtoseconds,  $^{6,7}$  much shorter than what is delivered by most conventional radio-frequency accelerators. A possible application of laser-driven accelerators is radiotherapy using very high energy electrons (VHEEs).  $^{8,9}$  Because of the short bunch duration, this results in dose delivery at rates much higher than in conventional radiotherapy.  $^{10}$  Interest in this regime has been renewed by recent developments in FLASH radiotherapy,  $^{11}$  where therapeutical doses are delivered in durations shorter than a few milliseconds, resulting in dose rates in excess of  $40\,\text{Gy/s}$ .

Electron beams propagating in a material are subject to scattering, which results in degradation of the beam quality. Lateral scattering can be reduced by increasing the electron energy and by focusing the beam using magnets, <sup>12</sup>, <sup>14</sup> resulting in high dose concentration at depth. Here we present results of Monte Carlo simulations to study the bunch length evolution in a water phantom for both collimated and focused beams, investigating the dependence on electron energy and initial bunch duration.

## 2. RESULTS

We performed simulations with the Monte Carlo code Geant4<sup>13</sup> to model the propagation of electron beams with different energies and bunch durations in a 300 mm  $\times$  300 mm  $\times$  300 mm water phantom. It has recently been demonstrated theoretically<sup>12</sup> and experimentally<sup>14</sup> that focused VHEE beams can improve dose delivery to deep seated tumours. Therefore, simulations were carried out both for a collimated beam and for a weakly focused beam, using the CLEAR beam line<sup>15</sup> at CERN as a reference. The beam transport was optimised using MAD-X<sup>16</sup> for an electron beam with energy of 200 MeV, 0.1% energy spread, r.m.s. divergence of 1 mrad and normalised emittance of 1  $\pi$  mm mrad passing through two triplets of quadrupole magnets, as shown in Fig. 1. The beam was focused at the centre of the phantom (in vacuum) with f-number  $\approx f/18$ , assuming a beam diameter  $2\sigma_r$ , with  $\sigma_r$  the r.m.s. radius. The magnetic field gradients are presented in table 1. The effective length of all quadrupoles was 226 mm, matching the specifications of the CLEAR beam line, but the aperture radius of the final two quadrupoles (Q515 and Q520) was increased to 250 mm to focus the beam in both planes without clipping. An aperture radius of 150 mm would have produced similar results, but with a charge loss of about 1%. Even with no clipping, the large quadrupole aperture and relatively short length result in a square beam profile in the horizontal plane at the entrance of the phantom. Here, we assumed hard-edged quadrupoles

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with no fringe fields, but simulations performed using GPT<sup>17</sup> with more realistic field profiles produced similar beam shapes. We did not optimise the positions and lengths of the quadrupoles, since these simulations are not intended to model a real experiment, but only to provide insight on the bunch evolution in a general case. In addition to a focusing geometry, we also performed simulations for a collimated beam with an r.m.s. radius of 5 mm.

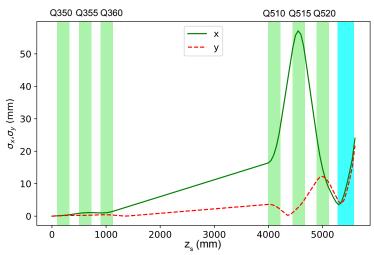


Figure 1: Electron beam transport through a beam line based on the CLEAR facility at CERN. Curves show the r.m.s. beam size in the horizontal (x) and vertical (y) plane for a 200 MeV beam with initial r.m.s. divergence of 1 mrad and normalised emittance of 1  $\pi$  mm mrad focused (in vacuum) at the centre of a water phantom, marked by the blue shaded area. Green shaded areas mark the electromagnetic quadrupoles.

Quadrupole	Gradient (T/m)
Q350	16.4
Q355	-8.3
Q360	14.8
Q510	14.4
Q515	-11.7
Q520	14.9

Table 1: Quadrupole gradients used for the simulations with a focused electron beam.

Fig. 2 shows the full-width-at-half-maximum (FWHM) size of a 200 MeV electron beam as it propagates inside the phantom. For a focused geometry (Fig. 2a), the beam size at the phantom entrance (z=0) is about 10 mm, gradually decreasing to a minimum of about 5 mm at a depth of 100 mm, about 50 mm upstream from the focus position in vacuum. For a collimated beam (Fig. 2b), the size remains approximately constant for the first 100 mm, until scattering causes it to diverge. The corresponding dose profiles are shown in Fig. 3 for a focused geometry and in Fig. 4 for a collimated beam. In both cases, the highest doses are registered in the first half of the phantom, but a focusing geometry enables to concentrate dose deposition in a small volume at larger depths. For the focusing configuration used here, the dose distribution is slightly asymmetric, as shown by the two-dimensional cross-sections in the horizontal (zx) and vertical (zy) planes presented in Fig. 3a and Fig. 3b, respectively. The square beam profile causes sharp edges to appear in the horizontal plane, but these are gradually smoothed away by interaction with water. The on-axis depth-dose distribution calculated in a 2 mm × 2 mm × 300 mm volume is presented in Fig. 3c, which is characterised by a peak at a depth of 81 mm with amplitude of  $8.8 \times 10^{-10}$  Gy per primary electron, corresponding to an enhancement of about 4.4 with respect to the surface dose. The corresponding depth-dose profile generated by a collimated beam (Fig. 4c), on the other hand, is approximately flat for the first 60 mm, with a dose of about  $2 \times 10^{-10}$  Gy per primary electron. The depth-dose profile calculated in a  $100 \,\mathrm{mm} \times 100 \,\mathrm{mm} \times 300 \,\mathrm{mm}$  volume centred on-axis is similar for both

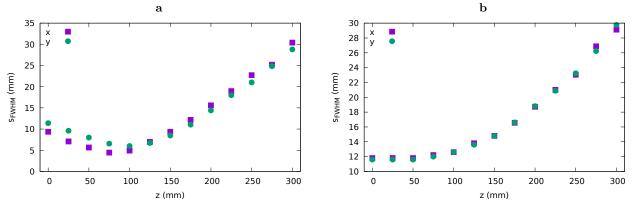


Figure 2: Transverse size (FWHM) versus depth in the phantom of  $\bf a$  a focused and  $\bf b$  a collimated electron beam with energy of 200 MeV.

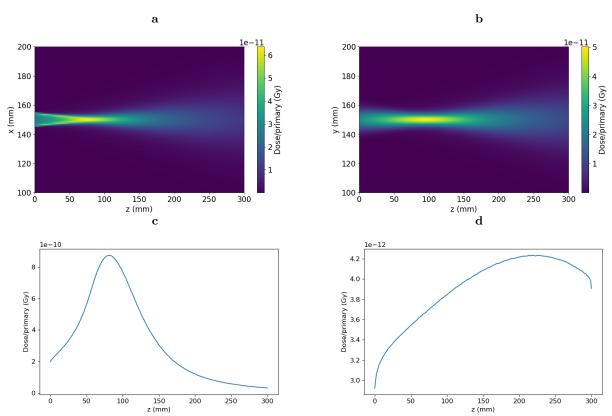


Figure 3: Dose deposited in the phantom by a 200 MeV focused electron beam. **a** Two-dimensional distribution in the zx plane, calculated in a  $2\,\mathrm{mm}\times100\,\mathrm{mm}\times300\,\mathrm{mm}$  volume. **b** Two-dimensional distribution in the zy plane, calculated in a  $100\,\mathrm{mm}\times2\,\mathrm{mm}\times300\,\mathrm{mm}$  volume. **c** On-axis depth-dose profile calculated in a  $2\,\mathrm{mm}\times2\,\mathrm{mm}\times300\,\mathrm{mm}$  volume. **d** Depth-dose profile calculated in a  $100\,\mathrm{mm}\times100\,\mathrm{mm}\times300\,\mathrm{mm}$  volume.

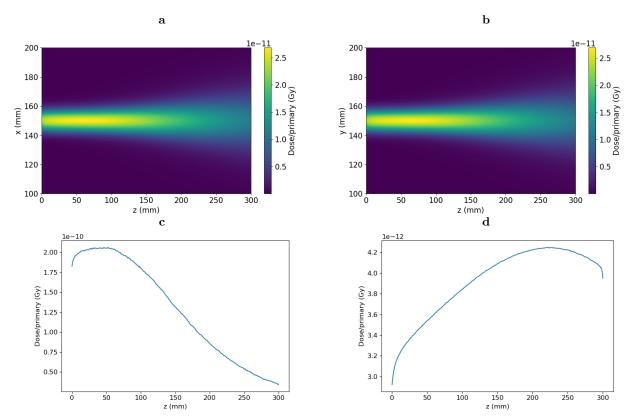


Figure 4: Dose deposited in the phantom by a 200 MeV collimated electron beam with r.m.s. size of 5 mm. a Two-dimensional distribution in the zx plane, calculated in a  $2\,\mathrm{mm}\times100\,\mathrm{mm}\times300\,\mathrm{mm}$  volume. b Two-dimensional distribution in the zy plane, calculated in a  $100\,\mathrm{mm}\times2\,\mathrm{mm}\times300\,\mathrm{mm}$  volume. c On-axis depth-dose profile calculated in a  $2\,\mathrm{mm}\times2\,\mathrm{mm}\times300\,\mathrm{mm}$  volume. d Depth-dose profile calculated in a  $100\,\mathrm{mm}\times100\,\mathrm{mm}\times300\,\mathrm{mm}$  volume.

geometries, as shown in Fig. 3d and Fig. 4d.

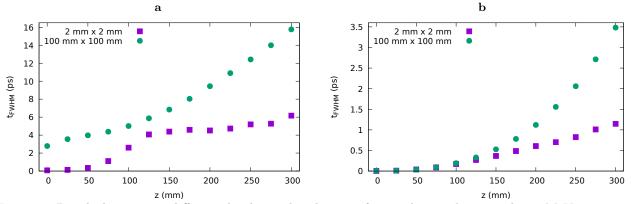


Figure 5: Bunch duration at different depths in the phantom for an electron beam with 200 MeV energy and initial r.m.s. bunch duration of 1 fs. a Focused geometry. b A collimated beam with 5 mm r.m.s. size.

The evolution of the bunch duration inside the phantom for an electron beam with 200 MeV energy and initial r.m.s. duration of 1 fs and 1 ps is shown in Fig. 5 and Fig. 6, respectively. Only electrons propagating in the forward direction with momentum  $p_z > 1 \,\mathrm{MeV/c}$  are included. For a focused geometry, electrons converge

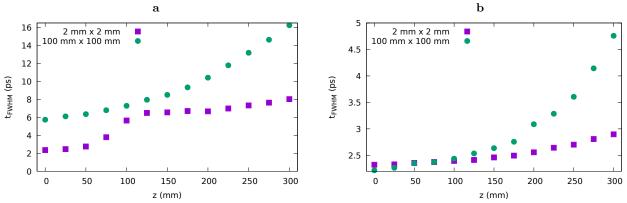


Figure 6: Bunch duration at different depths in the phantom for an electron beam with 200 MeV energy and initial r.m.s. bunch duration of 1 ps. **a** Focused geometry. **b** A collimated beam with 5 mm r.m.s. size.

towards the centre of the phantom after travelling along different trajectories, resulting in larger path differences and longer bunch durations than for a collimated beam. For an initial r.m.s. bunch duration of 1 fs (2.35 fs FWHM), the FWHM bunch duration at the centre of the phantom ( $z=150\,\mathrm{mm}$ ) calculated in a 2 mm  $\times$  2 mm cross-sectional area centred on-axis is about 4.5 ps, in contrast to a 0.35 ps duration obtained for a collimated beam. The bunch duration calculated in a 100 mm  $\times$  100 mm area is about 7 ps for a focused beam and 0.5 ps for a collimated beam. If the initial r.m.s. bunch duration is 1 ps (2.35 ps FWHM), on the other hand, the FWHM bunch duration at the centre of the phantom is about 6.5 ps for a focused beam and 2.5 ps for a collimated beam when using a 2 mm  $\times$  2 mm area, growing to 8.5 ps and 2.6 ps, respectively, in a 100 mm  $\times$  100 mm area.

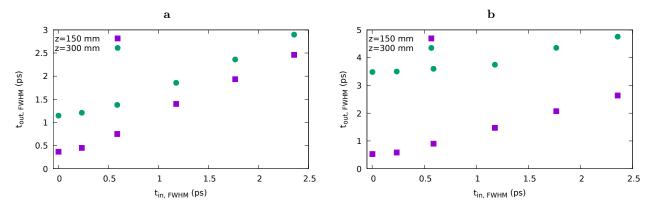


Figure 7: Bunch duration at the centre ( $z=150\,\mathrm{mm}$ ) and exit ( $z=300\,\mathrm{mm}$ ) of the phantom calculated **a** in a  $2\,\mathrm{mm}\times2\,\mathrm{mm}$  cross-sectional area and **b** in a  $100\,\mathrm{mm}\times100\,\mathrm{mm}$  cross-sectional area. The electron beam is collimated and has energy of  $200\,\mathrm{MeV}$ , initial r.m.s. size of  $5\,\mathrm{mm}$  and varying initial bunch duration.

The dependence of bunch lengthening on the initial bunch duration is shown in Fig. 7 for a collimated electron beam with energy of  $200\,\mathrm{MeV}$ . When the initial bunch duration is on the order of femtoseconds or tens of femtoseconds, bunch lengthening in water is large. For example, a bunch with r.m.s duration of 1 fs (2.35 fs FWHM) lengthens by a factor of about 150 after 150 mm propagation in water and by a factor of 500 after 300 mm, when considering a  $2\,\mathrm{mm} \times 2\,\mathrm{mm}$  area. For bunch durations on the order of 100s of femtoseconds or longer, on the other hand, bunch lengthening is small. For example, a bunch with r.m.s. duration of 250 fs (590 fs FWHM) lengthens by a factor of 1.3 after 150 mm and 2.3 after 500 mm.

The dependence of bunch lengthening on the electron beam energy is shown in Fig. 8 for a collimated beam with initial r.m.s. duration of 1 fs. The bunch length increases steeply at low electron energies. For example, for an energy of  $100 \, \text{MeV}$  the bunch lengthens by a factor of about  $350 \, \text{after} \, 150 \, \text{mm}$  propagation in water and by a factor of  $2400 \, \text{at}$  the exit of the phantom, when considering a  $2 \, \text{mm} \times 2 \, \text{mm}$  cross-sectional area. In

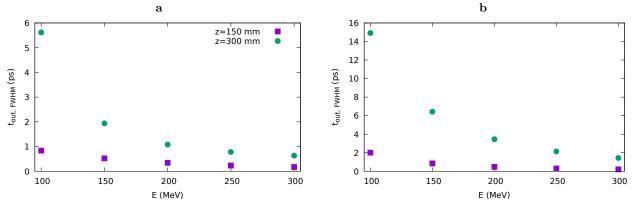


Figure 8: Bunch duration at the centre ( $z=150\,\mathrm{mm}$ ) and exit ( $z=300\,\mathrm{mm}$ ) of the phantom calculated  $\mathbf a$  in a  $2\,\mathrm{mm}\times2\,\mathrm{mm}$  cross-sectional area and  $\mathbf b$  in a  $100\,\mathrm{mm}\times100\,\mathrm{mm}$  cross-sectional area. The electron beam is collimated and has energy between  $100\,\mathrm{MeV}$  and  $300\,\mathrm{MeV}$ , initial r.m.s. size of 5 mm and initial r.m.s. bunch duration of 1 fs.

a  $100\,\mathrm{mm} \times 100\,\mathrm{mm}$  area, the bunch lengthens by a factor of 850 and 6300, respectively. For energies above  $200\,\mathrm{MeV}$ , on the other hand, the bunch lengthening slowly decreases with energy.

#### 3. DISCUSSION

We investigated the propagation of VHHE beams in water for collimated and focused geometries. We have shown that the typical bunch durations in the middle of a  $300\,\mathrm{mm}\times300\,\mathrm{mm}\times300\,\mathrm{mm}$  phantom are close to 1 ps for electron beams with  $200\,\mathrm{MeV}$  energy and initial bunch duration between 1 fs and 1 ps. Assuming an accelerator operating at a repetition rate of 1 Hz, the corresponding dose rates are about  $200\,\mathrm{Gy/s}$  per primary electron, which is many orders of magnitude higher than in conventional radiotherapy. We have shown that focusing the electron beam results in an additional bunch lengthening by a factor between 3 and 10, but the peak dose at depth is enhanced. Typically, laser-driven accelerators produce electron beams with charges of the order of 10s or 100s of pC, resulting in doses of 10s of mGy. These values are expected to be too low to observe FLASH effect, which requires doses of at least 5–10 Gy, based on current knowledge. Nevertheless, these beams allow to explore regimes that have not been extensively studied so far.

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