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Physical Characterization of Main-Belt Comet (248370) 2005 QN₁₇₃

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ABSTRACT

We report results from new and archival observations of the newly discovered active asteroid (248370) 2005 QN₁₃₇, which has been determined to be a likely main-belt comet based on a subsequent discovery that it is recurrently active near perihelion. From archival data analysis, we estimate g' -, r' -, i' -, and z' -band absolute magnitudes for the nucleus of $H_g = 16.62 \pm 0.13$, $H_r = 16.12 \pm 0.10$, $H_i = 16.05 \pm 0.11$, and $H_z = 15.93 \pm 0.08$, corresponding to nucleus colors of $g' - r' = 0.50 \pm 0.16$, $r' - i' = 0.07 \pm 0.15$, and $i' - z' = 0.12 \pm 0.14$, an equivalent V -band absolute magnitude of $H_V = 16.32 \pm 0.08$, and a nucleus radius of $r_n = 1.6 \pm 0.2$ km (using a V -band albedo of $p_V = 0.054 \pm 0.012$). Meanwhile, we find mean near-nucleus coma colors when 248370 was active of $g' - r' = 0.47 \pm 0.03$, $r' - i' = 0.10 \pm 0.04$, and $i' - z' = 0.05 \pm 0.05$, and similar mean dust tail colors, suggesting that no significant gas coma is present. We find approximate ratios between the scattering cross-sections of near-nucleus dust (within 5000 km of the nucleus) and the nucleus of $A_d/A_n = 0.7 \pm 0.3$ on 2016 July 22, and $1.8 < A_d/A_n < 2.9$ in 2021 July and August. During the 2021 observation period, the coma declined in intrinsic brightness by ~ 0.35 mag (or $\sim 25\%$) in 37 days, while the surface brightness of the dust tail remained effectively constant over the same period. Constraints derived from the sunward extent of the coma suggest that terminal velocities of ejected dust grains are extremely slow (~ 1 m s⁻¹ for 1 μ m particles), indicating

that the observed dust emission may have been aided by rapid rotation of the nucleus lowering the effective escape velocity.

Keywords: Main belt comets — Comets — Main belt asteroids

1. INTRODUCTION

Asteroid (248370) 2005 QN₁₇₃ (hereafter, 248370) was discovered to be active on UT 2021 July 7 in data comprising 120s of total exposure time (Figure 1e) obtained by the Asteroid Terrestrial-Impact Last Alert System (ATLAS; Tonry et al. 2018a) survey telescope (Fitzsimmons et al. 2021). On that date, the object was at a heliocentric distance of $r_h = 2.391$ au and true anomaly of $\nu = 16.0^\circ$, having most recently reached perihelion on UT 2021 May 14. As reported in the discovery announcement, 248370 exhibited a thin, straight dust tail 7'6 in length at a position angle of 245° East of North in confirmation observations obtained by Lowell Observatory's 4.3 m Lowell Discovery Telescope (LDT). Zwicky Transient Facility (ZTF) observations show the presence of the tail as early as UT 2021 June 11 (Kelley et al. 2021).

As of 2021 August 1, 248370 has a semimajor axis of $a = 3.067$ au, eccentricity of $e = 0.226$, and inclination of $i = 0.067^\circ$, according to the JPL Small-Body Database¹, placing it unambiguously in the outer main asteroid belt. 248370's active nature and asteroidal orbit places it among the class of objects known as active asteroids, which exhibit comet-like mass loss yet have dynamically asteroidal orbits (Jewitt et al. 2015). Active asteroids include main-belt comets (Hsieh & Jewitt 2006), for which sublimation of volatile ice is the most likely activity driver, and disrupted asteroids, for which activity is due to other processes such as impacts or rotational destabilization (e.g., Hsieh et al. 2012).

248370 has been previously measured to have a diameter of 3.6 ± 0.2 km and visible geometric albedo of 0.054 ± 0.012 , using $H_V = 16.00$ for the V-band absolute magnitude and $G = 0.15$ (Mainzer et al. 2019). As of July 2021, there are no published rotational lightcurve data available for the object in the Asteroid Lightcurve Photometry Database², nor in the Asteroid Lightcurve Data Base archived by the NASA Planetary Data System (PDS). Similarly, no taxonomic classification for 248370 is available in current PDS catalogs.

Following the discovery of 248370's activity in 2021, Chandler et al. (2021) reported the discovery of activity in archival data from the Dark Energy Camera (DE-

Cam; Flaughner et al. 2015) on the 4 m Victor M. Blanco Telescope (hereafter, Blanco) at Cerro Tololo Interamerican Observatory (CTIO) obtained on UT 2016 July 22, when the object was at a true anomaly of $\nu = 56.5^\circ$, having then most recently passed perihelion on UT 2016 January 3. This discovery of two separate active apparitions of 248370, both near perihelion, is considered a strong indication that sublimation is responsible for the observed activity (e.g., Hsieh et al. 2012, Chandler et al., submitted).

2. OBSERVATIONS

New observations of 248370 were obtained on several nights between UT 2021 July 8 and UT 2021 August 14 with LDT (Levine et al. 2012), Palomar Observatory's 5 m Hale Telescope (hereafter, Palomar), the 2 m Faulkes Telescope North (FTN), and Las Cumbres Observatory (hereafter, LCOGT) 1 m telescopes (Brown et al. 2013) at CTIO and the South African Astronomical Observatory (SAAO) (Table 1). Observations were obtained using the LDT's Large Monolithic Imager (LMI; Bida et al. 2014), Palomar's Wafer-Scale camera for Prime (WaSP; Nikzad et al. 2017) wide field prime focus camera, FTN's Multicolor Simultaneous Camera for studying Atmospheres of Transiting exoplanets (MuSCAT3; Narita et al. 2020), and LCOGT Sinistro cameras. All observations were obtained using Sloan g' -, r' -, i' -, or z' -band filters, and non-sidereal tracking to follow the target's motion.

Multi-filter FTN data were obtained using the simultaneous g' -, r' -, i' -, and z' -band imaging capability of MuSCAT3. Multi-filter Palomar observations were obtained by interspersing filters (i.e., using repeating $r'g'r'i'r'$ or $r'i'g'r'$ sequences) to enable the use of interpolation to approximate simultaneous multi-filter imaging for color computation (i.e., compensating for possible rotational variability in the nucleus brightness between our actual observations in different filters).

Bias subtraction, flat field correction, and cosmic ray removal were performed for LDT and Palomar data using Python 3 code utilizing the `ccdproc` package in Astropy (Astropy Collaboration et al. 2018) and `L.A.Cosmic` code³ (van Dokkum 2001; van Dokkum et al. 2012). FTN and LCOGT 1 m data were pro-

¹ <https://ssd.jpl.nasa.gov/sbdb.cgi>

² <https://minplanobs.org/alcdef/index.php>

³ Written for python by Maltes Tewes (<https://github.com/RyleighFitz/LACosmics>)

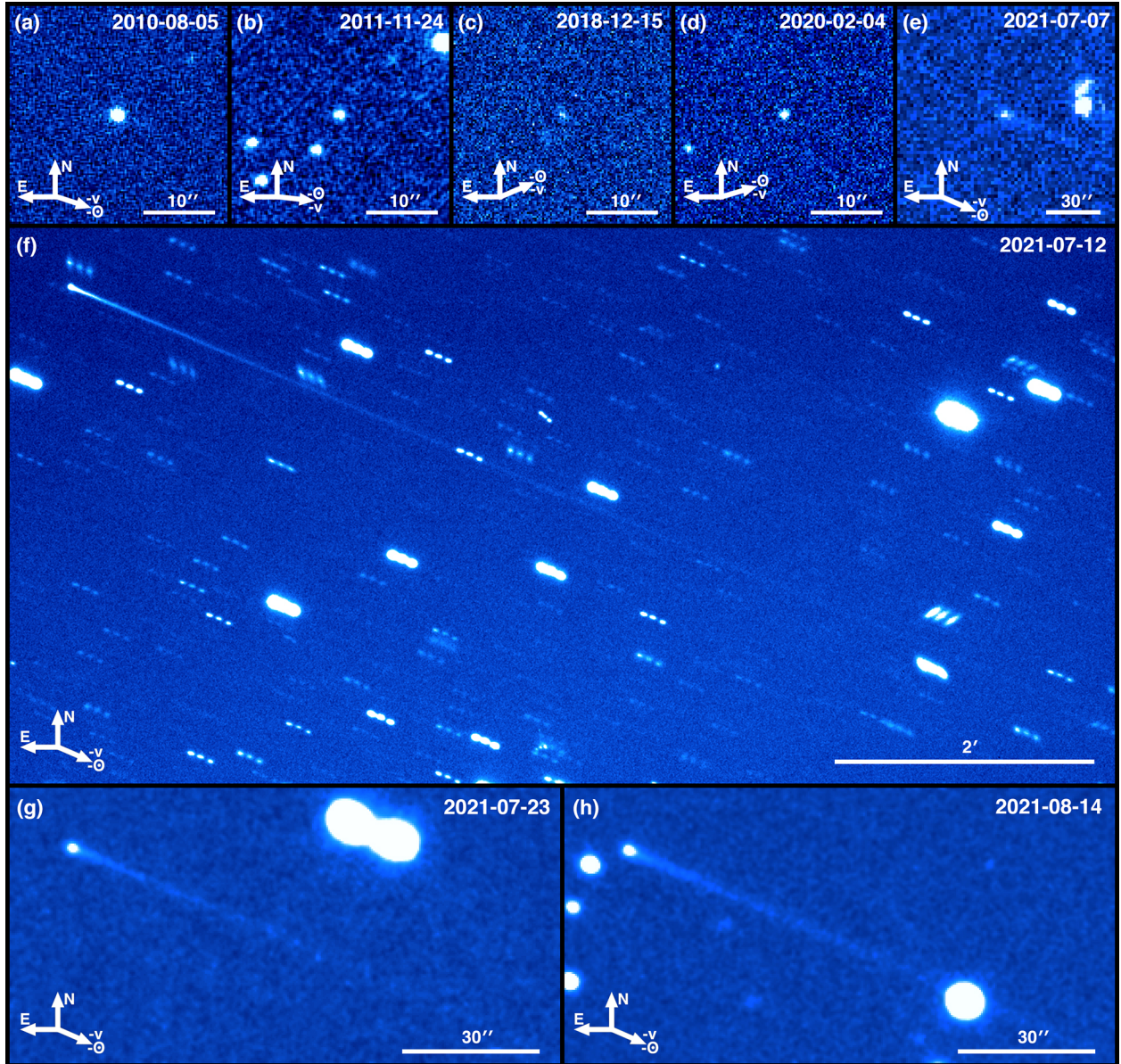


Figure 1. Single or composite images of 248370 for dates indicated in each panel (see Tables 1 and 2 for observation details). All images are in r' -band except for (e) which was obtained using the ATLAS survey’s “cyan” filter (bandpass from 420 nm to 650 nm). Scale bars indicate the size of each panel. North (N), East (E), the antisolar direction ($-\odot$), and the negative heliocentric velocity direction ($-v$) are indicated in each panel. The object is located at the center of panels (a) through (e), while in panels (f) through (h), the object’s nucleus is located in the upper left corner with the tail extending down and to the right, where the latter set of images have been Gaussian-smoothed to enhance the visibility of low surface brightness features.

cessed using standard LCOGT pipeline software (McCully et al. 2018).

We also used the Canadian Astronomy Data Centre’s Solar System Object Image Search tool⁴ (Gwyn et al. 2012) and the NASA Planetary Data System (PDS) Small Bodies Node’s Comet Asteroid Telescopic Cat-

alog Hub (CATCH) tool⁵ to identify archival Sloan g' -, r' -, i' -, and z' -band observations of 248370 from 2004 to 2020 (Table 2) from the 1.8 m Panoramic Survey Telescope and Rapid Response System (Pan-STARRS1; hereafter PS1) survey telescope (Chambers et al. 2016; Flewelling et al. 2020), MegaCam (Boulade et al. 2003) on the 3.6 m Canada-France-Hawaii Telescope (CFHT),

⁴ <http://www.cadc-ccda.hia-ihp.nrc-cnrc.gc.ca/en/ssois/>

⁵ <https://catch.astro.umd.edu/>

Table 1. (248370) 2005 QN₁₇₃ Activity Observations

UT Date	Telescope	t^a	Filter	θ_s^b	ν^c	r_h^d	Δ^e	α^f	$m(r_h, \Delta, \alpha)^g$	$m(1, 1, 0)^h$	A_d/A_n^i	$Af\rho^j$	Σ_t^k
2021 Jul 7	ATLAS	120	cyan	5.2	16.0	2.391	1.930	24.4	—	—	—	—	—
2021 Jul 9	Palomar	800	g'	2.1	16.6	2.392	1.909	24.2	19.57±0.01	15.14±0.01	2.9±0.5	15.1±0.7	24.24±0.20
2021 Jul 9	FTN	480	g'	1.5	16.6	2.392	1.909	24.2	19.71±0.01	15.28±0.01	2.4±0.5	13.6±0.8	23.87±0.20
2021 Jul 12	Palomar	300	g'	1.4	17.5	2.394	1.877	23.8	19.52±0.02	15.14±0.02	2.9±0.5	16.4±0.8	23.66±0.20
2021 Jul 13	FTN	480	g'	2.4	17.8	2.395	1.867	23.7	19.52±0.02	15.15±0.02	2.9±0.5	14.5±0.8	24.01±0.20
2021 Jul 15	FTN	480	g'	1.5	18.5	2.396	1.842	23.4	19.55±0.02	15.22±0.02	2.6±0.5	14.6±0.8	23.77±0.20
2021 Jul 8	LDT	2400	r'	1.8	16.3	2.391	1.920	24.3	19.14±0.01	14.69±0.01	2.7±0.4	23.2±0.9	23.45±0.20
2021 Jul 9	Palomar	1300	r'	2.1	16.6	2.392	1.909	24.2	19.11±0.01	14.68±0.01	2.8±0.4	22.8±0.8	23.68±0.20
2021 Jul 9	FTN	480	r'	1.4	16.6	2.392	1.909	24.2	19.19±0.01	14.76±0.01	2.5±0.4	22.2±1.0	23.34±0.20
2021 Jul 10	LCOGT 1m	480	r'	1.8	16.9	2.393	1.898	24.1	19.17±0.02	14.76±0.02	2.5±0.4	21.4±1.0	23.42±0.20
2021 Jul 12	Palomar	900	r'	1.2	17.5	2.394	1.877	23.8	19.12±0.01	14.74±0.01	2.6±0.4	23.2±0.9	23.17±0.20
2021 Jul 13	FTN	480	r'	2.2	17.8	2.395	1.867	23.7	19.05±0.01	14.68±0.01	2.8±0.4	22.6±0.9	23.55±0.20
2021 Jul 14	LCOGT 1m	1130	r'	2.1	18.1	2.395	1.856	23.6	19.11±0.02	14.76±0.02	2.5±0.4	20.8±1.0	23.66±0.20
2021 Jul 15	LCOGT 1m	1130	r'	1.7	18.5	2.396	1.842	23.4	19.05±0.01	14.72±0.01	2.6±0.4	22.8±0.9	23.41±0.20
2021 Jul 15	FTN	480	r'	1.4	18.5	2.396	1.842	23.4	19.13±0.01	14.80±0.01	2.4±0.3	21.2±1.0	23.26±0.20
2021 Jul 18	LCOGT 1m	283	r'	1.7	19.3	2.398	1.815	23.0	19.03±0.06	14.75±0.06	2.5±0.4	22.1±1.9	23.40±0.20
2021 Jul 19	LCOGT 1m	1130	r'	1.9	19.5	2.399	1.805	22.9	19.06±0.02	14.78±0.02	2.4±0.4	20.7±1.0	23.45±0.20
2021 Jul 21	LCOGT 1m	1130	r'	2.2	20.1	2.401	1.785	22.5	19.07±0.03	14.83±0.03	2.3±0.3	18.8±1.1	23.75±0.20
2021 Jul 23	LCOGT 1m	1853	r'	1.9	20.8	2.402	1.762	22.1	19.05±0.03	14.85±0.03	2.2±0.3	19.0±1.1	23.52±0.20
2021 Aug 5	LCOGT 1m	600	r'	2.1	24.5	2.414	1.648	19.2	18.86±0.02	14.89±0.02	2.1±0.3	17.9±1.0	23.55±0.20
2021 Aug 7	LCOGT 1m	600	r'	2.0	25.1	2.415	1.632	18.7	18.80±0.01	14.87±0.01	2.2±0.3	18.7±0.9	23.50±0.20
2021 Aug 14	LCOGT 1m	600	r'	2.0	27.2	2.423	1.577	16.4	18.81±0.01	15.02±0.02	1.8±0.3	15.3±1.0	23.45±0.20
2021 Jul 9	Palomar	900	i'	2.1	16.6	2.392	1.909	24.2	18.98±0.01	14.55±0.01	2.8±0.2	25.7±0.5	23.60±0.20
2021 Jul 9	FTN	480	i'	1.3	16.6	2.392	1.909	24.2	19.16±0.02	14.73±0.02	2.2±0.2	22.2±0.9	23.07±0.20
2021 Jul 12	Palomar	600	i'	1.2	17.5	2.394	1.877	23.8	19.03±0.01	14.65±0.01	2.4±0.2	24.8±0.6	22.94±0.20
2021 Jul 13	FTN	480	i'	2.2	17.8	2.395	1.867	23.7	18.92±0.02	14.55±0.02	2.8±0.2	25.5±0.9	23.40±0.20
2021 Jul 15	FTN	480	i'	1.4	18.5	2.396	1.842	23.4	18.95±0.02	14.62±0.02	2.5±0.2	25.5±0.9	23.07±0.20
2016 Jul 22	Blanco	89	z'	1.1	56.5	2.591	2.571	22.7	20.51±0.10	15.31±0.10	0.7±0.3	7.5±2.4	25.30±0.50
2021 Jul 9	FTN	480	z'	1.2	16.6	2.392	1.909	24.2	19.05±0.04	14.62±0.04	2.2±0.5	24.5±2.1	23.08±0.30
2021 Jul 13	FTN	480	z'	2.1	17.8	2.395	1.867	23.7	18.84±0.03	14.48±0.03	2.6±0.5	27.2±1.9	23.40±0.30
2021 Jul 15	FTN	480	z'	1.3	18.5	2.396	1.842	23.4	18.92±0.03	14.59±0.03	2.2±0.5	25.5±1.9	23.06±0.30

^a Total integration time, in seconds.

^b FWHM seeing, in arcseconds.

^c True anomaly, in degrees.

^d Heliocentric distance, in au.

^e Geocentric distance, in au.

^f Solar phase angle (Sun-object-Earth), in degrees.

^g Mean apparent coma magnitude (measured using 5000 km-radius photometry apertures).

^h Computed absolute magnitude corresponding to measured apparent magnitude assuming a H, G phase function where $G = 0.15$.

ⁱ Inferred ratio of effective scattering cross-sectional areas of coma dust (within 5000 km-radius photometry apertures) and the nucleus, where listed uncertainties only reflect photometric measurement uncertainty and not potential nucleus rotational variability.

^j $A(\alpha = 0^\circ)f\rho$ values, computed using Equation 4 and 5000 km photometry apertures, in cm.

^k Dust tail surface brightness in mag arcsec⁻², as measured in a 1500 km × 10000 km rectangular aperture, as described in the text.

the 1.35 m SkyMapper survey telescope (Wolf et al. 2018), and Blanco. For the purposes of our analysis, PS1 g_{P1} , r_{P1} , i_{P1} , and z_{P1} filters are considered functionally equivalent to their Sloan counterparts. All archival data were pipeline-processed by their respective facilities.

The object was identified in archival images either from its non-sidereal motion when more than one image was available on a particular night, or otherwise from comparison with reference images obtained on other nights when the object was not in the field of view.

3. RESULTS AND ANALYSIS

3.1. Data Analysis

Except for data from 2016 July 22, 248370 had a star-like surface brightness profile in all archival images and

exhibited no other visible indications of activity. Meanwhile, in all 2021 observations, the object exhibited a long, straight dust tail oriented along the coincident antisolar and negative heliocentric velocity vector directions as projected on the sky. In our best composite image from UT 2021 July 12, the tail was seen extending $\sim 9'$ from the nucleus (Figure 1f), corresponding to a physical extent of $\sim 720,000$ km at the geocentric distance of the comet. Minimal coma was present in all images, with full-width at half-maximum (FWHM) measurements of the nucleus's surface brightness profile measured in the direction perpendicular to the dust tail nearly identical to FWHM measurements, θ_s , of field star profiles (listed in Table 1), measured in the direction perpendicular to their trailing due to non-sidereal

Table 2. (248370) 2005 QN₁₇₃ Nucleus Observations

UT Date	Tel. ^a	N ^b	t ^c	Filter	ν^d	r_h^e	Δ^f	α^g	$m(r_h, \Delta, \alpha)^h$	$m(1, 1, 0)^i$
2010 Aug 6	PS1	1	43	<i>g'</i>	354.7	2.389	1.377	2.0	19.17±0.03	16.35±0.03
2010 Sep 6	PS1	2	86	<i>g'</i>	3.8	2.388	1.463	12.2	20.08±0.05	16.63±0.05
2011 Nov 24	PS1	2	86	<i>g'</i>	109.4	3.150	2.180	4.1	21.20±0.13	16.64±0.13
2011 Dec 1	PS1	2	86	<i>g'</i>	110.6	3.164	2.180	1.4	20.98±0.10	16.60±0.10
<i>Median^j</i>	<i>g'</i>	16.62±0.13
2010 Aug 5	PS1	1	40	<i>r'</i>	354.4	2.389	1.378	2.5	18.81±0.03	15.95±0.03
2011 Nov 24	PS1	2	80	<i>r'</i>	109.4	3.150	2.180	4.0	20.74±0.10	16.19±0.10
2018 Dec 15	Blanco	1	45	<i>r'</i>	191.7	3.733	3.508	15.2	22.63±0.31	16.20±0.31
2020 Feb 4	Blanco	2	81	<i>r'</i>	248.9	3.165	3.059	18.1	21.91±0.11	16.04±0.11
<i>Median^j</i>	<i>r'</i>	16.12±0.10
2004 Jul 8	CFHT	3	540	<i>i'</i>	287.5	2.740	2.028	17.7	20.72±0.03	16.07±0.03
2010 Aug 2	PS1	1	45	<i>i'</i>	353.5	2.390	1.383	3.9	19.01±0.05	16.05±0.05
2010 Aug 31	PS1	2	90	<i>i'</i>	2.1	2.388	1.429	9.8	19.23±0.04	15.93±0.04
2011 Nov 30	PS1	2	90	<i>i'</i>	110.4	3.162	2.179	1.9	20.26±0.09	15.84±0.09
2015 Aug 18	SkyMapper	1	100	<i>i'</i>	320.3	2.483	1.888	21.8	20.72±0.26	16.31±0.26
<i>Median^j</i>	<i>i'</i>	16.05±0.11
2010 Jun 14	PS1	2	60	<i>z'</i>	339.3	2.416	1.728	21.1	20.07±0.13	15.93±0.13
2010 Oct 30	PS1	5	150	<i>z'</i>	19.6	2.413	2.018	23.8	20.48±0.23	15.93±0.13
2020 Feb 10	Blanco	1	199	<i>z'</i>	250.0	3.152	2.960	18.2	21.59±0.14	15.80±0.14
<i>Median^j</i>	<i>z'</i>	15.93±0.08

^a Telescope used^b Number of exposures.^c Total integration time, in seconds.^d True anomaly, in degrees.^e Heliocentric distance, in au.^f Geocentric distance, in au.^g Solar phase angle (Sun-object-Earth), in degrees.^h Mean apparent magnitude in specified filter.ⁱ Computed absolute magnitude corresponding to measured apparent magnitude assuming IAU *H, G* phase function behavior where *G* = 0.15.^j Median values of computed absolute magnitudes, where standard deviations are used as uncertainties.**Table 3.** (248370) 2005 QN₁₇₃ Colors

UT Date	Telescope	Coma			Tail		
		<i>g' - r'</i>	<i>r' - i'</i>	<i>i' - z'</i>	<i>g' - r'</i>	<i>r' - i'</i>	<i>i' - z'</i>
2021 Jul 9	Palomar	0.47±0.01	0.13±0.01	—	0.56±0.30	0.08±0.30	—
2021 Jul 9	FTN	0.52±0.02	0.03±0.03	0.14±0.04	0.53±0.30	0.27±0.30	-0.02±0.40
2021 Jul 12	Palomar	0.42±0.02	0.08±0.01	—	0.49±0.30	0.23±0.30	—
2021 Jul 13	FTN	0.47±0.02	0.13±0.03	0.05±0.04	0.45±0.30	0.16±0.30	-0.01±0.40
2021 Jul 15	FTN	0.50±0.02	0.10±0.03	0.03±0.04	0.51±0.30	0.19±0.30	0.01±0.40
<i>Median^a</i>	...	0.47±0.03	0.10±0.04	0.05±0.05	0.51±0.04	0.19±0.07	-0.01±0.01

^a Median values of computed colors, where standard deviations are used as uncertainties.

tracking. We did however find the half-width at half-maximum (HWHM) of the object's profile measured along the sunward direction directly opposite the dust tail to be $\sim 10\%$ larger than stellar HWHM values.

To maximize signal-to-noise ratios (S/N) for sets of observations where more than one image was obtained in the same filter in a night, we constructed composite images by shifting and aligning individual images in each filter on the object's photocenter using linear interpolation, and adding them together. Representative single or composite images are shown in Figure 1.

For photometry of all data, measurements of 248370 and 10-30 nearby reference stars were performed using IRAF software (Tody 1986, 1993), with absolute cali-

bration performed using field star magnitudes in Sloan bandpasses derived from the RefCat2 all-sky catalog (which uses the PS1 photometric system; Tonry et al. 2018b). Nucleus or near-nucleus coma photometry of 248370 was performed using circular apertures with sizes chosen using curve-of-growth analyses when the object appeared inactive, or circular apertures with fixed radii equivalent to 5000 km at the geocentric distance of the object when it was active. For the latter, photometry aperture radii, θ_{obs} , were determined from convolving the projected angular equivalent, θ_0 , of the desired intrinsic distance (i.e., 5000 km) at the geocentric distance of the comet with the FWHM seeing, θ_s , on a given night

using

$$\theta_{\text{obs}} = (\theta_0^2 + \theta_s^2)^{1/2} \quad (1)$$

where $\theta_{\text{obs}} \sim 4''$ for most of our observations. Background statistics for comet photometry were measured in nearby regions of blank sky to avoid dust contamination from the object or nearby field stars.

We also measured the surface brightnesses of 248370's dust tail on each night by rotating composite images to make the dust tail horizontal in each image frame, measuring net fluxes in rectangular apertures placed along the length of each tail, and converting those fluxes to surface brightnesses in mag arcsec⁻² using the measured mean magnitudes of the nucleus for data comprising each composite image for absolute photometric calibration. We chose rectangular apertures that extended 750 km above and below the tail's central axis in the vertical direction (i.e., $\sim 2''$ in total height for most of our observations) and from 5000 km to 15000 km (i.e., from $\sim 4''$ to $\sim 12''$ for most of our observations) from the nucleus in the horizontal direction, where the angular sizes of these apertures were computed in the same manner described earlier for near-nucleus photometry apertures. The details of this method of measuring surface brightnesses were chosen to maximize signal-to-noise while also minimizing nucleus flux contribution by focusing on the bright central core of the tail, and measuring close, but not too close, to the nucleus where the tail is brightest.

Photometric results for data obtained when 248370 appeared active and inactive are shown in Tables 1 and 2, respectively. Colors computed by comparing coma magnitudes or tail surface brightnesses in different filters for nights on which multi-filter data were obtained are shown in Table 3.

3.2. Nucleus Properties

Using measured apparent magnitudes of 248370 from archival data, we derive magnitudes normalized to $r_h = \Delta = 1$ au and $\alpha = 0^\circ$, or $m(1, 1, 0)$, by assuming inverse-square-law fading and IAU H, G phase function behavior (Bowell et al. 1989) where $G = 0.15$ (Table 2). We then take medians of these computed $m(1, 1, 0)$ values to estimate absolute magnitudes in each filter. We estimate 248370's absolute magnitudes to be $H_g = 16.62 \pm 0.13$, $H_r = 16.12 \pm 0.10$, $H_i = 16.05 \pm 0.11$, and $H_z = 15.93 \pm 0.08$ (Table 2), corresponding to nucleus colors of $g' - r' = 0.50 \pm 0.16$, $r' - i' = 0.07 \pm 0.15$, and $i' - z' = 0.12 \pm 0.14$, which within uncertainties, are effectively solar (e.g., Holmberg et al. 2006). These colors are consistent with a C-type taxonomic classification (see DeMeo & Carry 2013), which is the most likely classification expected for an outer main belt as-

teroid like 248370, but large uncertainties on the colors derived here from sparse archival data mean that other taxonomic types cannot necessarily be excluded at this time. Using $V = g' - 0.565(g' - r') - 0.016$ (Jordi et al. 2006), we find an equivalent V -band absolute magnitude of $H_V = 16.32 \pm 0.10$.

Using

$$r_n = \left(\frac{2.24 \times 10^{22}}{p_V} \times 10^{0.4(m_{\odot, V} - H_V)} \right)^{1/2} \quad (2)$$

where we use $p_V = 0.054 \pm 0.012$ (Mainzer et al. 2019) for the object's V -band albedo and $m_{\odot, V} = -26.71 \pm 0.03$ for the apparent V -band magnitude of the Sun (Hardorp 1980), we find an effective nucleus radius of $r_n = 1.6 \pm 0.2$ km, or slightly smaller than the radius computed by Mainzer et al. (2019).

The ranges in computed absolute magnitudes in each filter are $\Delta m_g = 0.29$ mag, $\Delta m_r = 0.25$ mag, $\Delta m_i = 0.23$ mag, and $\Delta m_z = 0.13$. These values are not particularly meaningful given the small number of data points used to derive them, but in the present absence of better measurements, they suggest that 248370's photometric range due to rotation is $\Delta m \gtrsim 0.3$ mag, implying a minimum axis ratio of $a/b = 1.3$.

3.3. Activity Properties

3.3.1. Dust Composition

We find mean coma colors of $g' - r' = 0.47 \pm 0.03$, $r' - i' = 0.10 \pm 0.04$, and $i' - z' = 0.05 \pm 0.05$ and mean dust tail colors of $g' - r' = 0.51 \pm 0.04$, $r' - i' = 0.19 \pm 0.07$, and $i' - z' = -0.01 \pm 0.01$ (Table 3). Within uncertainties, coma and dust tail colors are comparable to one another, indicating that both are dominated by dust of similar composition with no apparent color gradient with distance from the nucleus that might indicate the presence of a significant near-nucleus gas coma. The apparent compositional similarity of coma and tail dust also means that we see no evidence of grain fragmentation or loss of icy grains to sublimation that could cause overall color changes to observed dust. The colors of both are also similar within uncertainties to colors found for the bare nucleus (Section 3.2), suggesting that the dust coma and tail are compositionally similar to the nucleus's surface regolith.

3.3.2. Activity Strength and Evolution

From our calculations of 248370's absolute magnitudes (Section 3.2), we find that the near-nucleus region of the object was ~ 0.5 mag brighter than expected for the inactive nucleus on 2016 July 22 and ~ 1 mag brighter than expected in 2021 (Table 1). We also compute the ratios,

A_d/A_n , of the scattering cross-sections of ejected near-nucleus dust within our 5000 km photometry apertures and the underlying nucleus when 248370 was active using

$$A_d/A_n = \frac{1 - 10^{0.4(m(1,1,0)-H)}}{10^{0.4(m(1,1,0)-H)}} \quad (3)$$

(e.g., Hsieh et al. 2021). We find $A_d/A_n = 0.7 \pm 0.3$ on 2016 July 22, and $1.8 < A_d/A_n < 2.9$ in 2021 (Table 1).

Plotting $m(1,1,0)$ and A_d/A_n as functions of time, we see that the coma faded during our 2021 observations (Figures 2a and 2b), declining in intrinsic brightness by ~ 0.35 mag (or $\sim 25\%$) in 37 days. Increasing activity strength would suggest ongoing dust production, and therefore the action of a prolonged, possibly sublimation-driven, emission event. However, declining activity strength does not necessarily rule out a sublimation-driven emission event, especially at the relatively gradual rate (~ 0.01 mag/day) seen for 248370, similar to the rate of fading of the coma of confirmed recurrently active MBC 259P/Garradd (Hsieh et al. 2021) of ~ 0.015 mag/day observed after its discovery in 2008 (Jewitt et al. 2009).

Despite the fading of 248370's coma, the dust tail remained relatively consistent in brightness during our observations (Table 1; Figure 2c), suggesting that the tail may consist of larger particles on average than the coma. Larger particles in the tail would be dissipated by radiation pressure more slowly than presumably smaller particles in the coma, which would explain the slower fading of the tail to apparently weakening dust production from the nucleus.

For reference, we also compute $A(\alpha = 0^\circ)f\rho$ values (hereafter, $Af\rho$), which are nominally independent of photometry aperture sizes for observations of comae with r^{-1} radial profiles, and are given by

$$A(\alpha = 0^\circ)f\rho = \frac{(2r_h\Delta)^2}{\rho} 10^{0.4[m_\odot - m_d(r_h, \Delta, 0)]} \quad (4)$$

(A'Hearn et al. 1984), where r_h is in au, Δ is in cm, ρ is the physical radius in cm of the photometry aperture at the geocentric distance of the comet, m_\odot is the Sun's apparent magnitude in the specified filter (using $m_{g,\odot} = -26.60$, $m_{r,\odot} = -27.05$, $m_{i,\odot} = -27.17$, and $m_{z,\odot} = -27.21$; Hardorp 1980; Jordi et al. 2006; Holmberg et al. 2006), and $m_d(r_h, \Delta, 0)$ is the phase-angle-normalized (to $\alpha = 0^\circ$) magnitude of the excess dust mass of the comet (i.e., with the flux contribution of the nucleus subtracted out). These results are tabulated in Table 1, where we see fading behavior similar to that seen for $m(1,1,0)$ and A_d/A_n .

3.3.3. Dust Ejection Parameter Constraints

Dust grains ejected sunward with a terminal ejection velocity of v_g are turned back by solar radiation pressure on a distance scale given by

$$X_R \sim \frac{v_g^2 r_h [\text{au}]^2}{2\beta_d g_\odot} \quad (5)$$

(Jewitt & Meech 1987), where r_h [au] is the heliocentric distance in au, $g_\odot = 0.006 \text{ m s}^{-2}$ is the gravitational acceleration to the Sun at 1 au, and β_d is the ratio of the acceleration experienced by a particle due to solar radiation pressure to the local acceleration due to solar gravity (Burns et al. 1979). Comet dust modeling analyses commonly use β_d to represent particle sizes, where $a_d \approx \beta_d^{-1}$ gives approximate corresponding dust particle radii, a_d , in μm .

On UT 2021 July 12, we measure a HWHM value for the sunward portion of the coma of $\theta_{\text{obs}}/2 = 0''.68$, while nearby field stars had HWHM values of $\theta_s/2 = 0''.61$. Using an analogous form of Equation 1 to compute an inferred intrinsic half-width of the coma, $\theta_0/2$, in the absence of atmospheric seeing, we find $\theta_0/2 = 0''.3$, or ~ 400 km at the geocentric distance of the object. Using $X_R = 400$ km, Equation 5 gives $v_g \sim 0.9\beta_d^{1/2} \text{ m s}^{-1}$, or about half the ejection velocities found for 133P/Elst-Pizarro (Jewitt et al. 2014), another MBC very similar in morphology to 248370.

Analyzing the composite image for UT 2021 July 12, we measure a median FWHM of $\theta_{\text{obs}} = 1''.47$ for the dust tail over the $2'$ of the tail closest to the nucleus, as measured in intervals of 20 pixels ($3''.5$) along the tail. The tail's FWHM flares slightly from $\theta_{\text{obs}} = 1''.45$ close to the nucleus to $\theta_{\text{obs}} = 1''.50$ at $2'$ from the nucleus. Using Equation 1, we find an intrinsic median tail FWHM of $\theta_0 = 0''.8 \pm 0''.3$, corresponding to a physical projected width of ~ 1100 km in the plane of the sky. The narrowness of the tail is consistent with low dust ejection velocities, and similar to that found for 133P (Hsieh et al. 2004).

Performing a simple (zero-ejection velocity) dust modeling analysis using the online Comet Toolbox⁶, we find that particles with $\beta = 1$ (or $a \sim 1 \mu\text{m}$) would take ~ 20 days to reach an apparent angular separation from the nucleus of $\sim 9'$ (the visible length of the tail on UT 2021 July 12; Figure 1f). Meanwhile, particles with $\beta = 0.1$, $\beta = 0.01$, and $\beta = 0.001$ (or $a \sim 10 \mu\text{m}$, $a \sim 100 \mu\text{m}$, and $a \sim 1 \text{ mm}$), which span the range of particle sizes found for other MBCs (e.g., Hsieh et al. 2009; Licandro et al. 2013; Jewitt et al. 2014), would take 60, 150, and 430 days, respectively, to reach the

⁶ <http://comet-toolbox.com/FP.html>

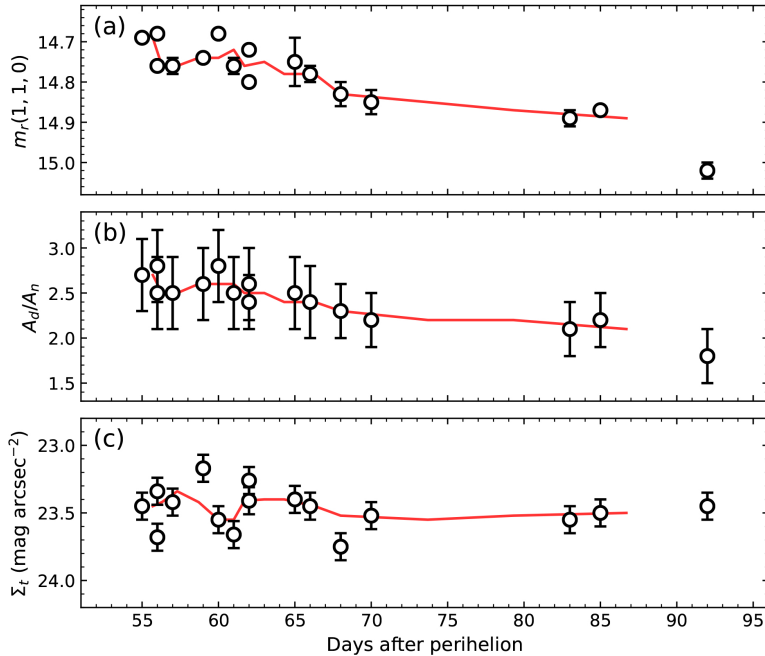


Figure 2. Plots of (a) r' -band near-nucleus magnitudes normalized to $r_h = \Delta = 1$ au and $\alpha = 0^\circ$, (b) inferred ratios of dust scattering cross-sections to nucleus scattering cross-sections measured within 5000 km-radius photometry apertures, and (c) surface brightnesses in mag arcsec^{-2} of a fixed portion of the dust tail, computed for 248370 from UT 2021 July 8 to 2021 August 14 as a function of days after perihelion (UT 2021 May 14), where red lines in each panel show moving medians (computed for groups of three data points each) for each quantity.

same apparent separation. Without additional particle size constraints at the present time, however, we are unable to meaningfully constrain the likely ejection times of the most distant dust grains in 248370’s tail. We note that if activity began when 248370 was at $\nu = 300^\circ$ (the earliest activation point confirmed to date for a MBC; Hsieh & Sheppard 2015), which the object passed on 2020 October 22, particles as large as $a \sim 400 \mu\text{m}$ ($\beta = 0.0025$) would have been able to reach a $9'$ separation from the nucleus by 2021 July 12.

3.4. Future Work

The discovery that 248370 is recurrently active near perihelion strongly suggests that sublimation is a primary driver of its activity, although it does not rule out other processes that could also contribute to the current observed activity or may have triggered it. In particular, the slow terminal velocities inferred for ejected dust grains (Section 3.3.3) suggest that rapid rotation, nucleus elongation, or both could be acting to reduce or negate the effective gravity felt by dust particles at certain locations on the nucleus surface, allowing even extremely large particles to be ejected, similar to what may be occurring on 133P (Jewitt et al. 2014). As such, measurement of 248370’s rotational period and nucleus shape, as well as its taxonomic type, should be considered a high priority once its current activity ends. Con-

tinued monitoring of 248370’s current activity is also highly encouraged to enable further characterization of the object’s fading behavior, which can help constrain the dust grain size distribution.

A detailed dynamical analysis of 248370 is outside the scope of this paper, but should also be performed in the near future. Issues to consider include whether the object can be linked to any dynamical asteroid families (e.g., Hsieh et al. 2018), its long-term dynamical stability and whether it may be an implanted object (e.g., Hsieh & Haghhighipour 2016), and whether it follows dynamical trends found for previously discovered MBCs (Kim et al. 2018).

In the long term, 248370 will be well-placed for monitoring during the approach to its next perihelion passage on UT 2026 September 3. It becomes observable from the Southern Hemisphere in 2026 February at $\nu \sim 300^\circ$, i.e., the earliest activation point confirmed to date for a MBC, as discussed earlier. Monitoring during this time will be extremely valuable for further confirming the recurrent nature of 248370’s activity, constraining the orbital range over which activity occurs (with implications for constraining ice depth on the object, as well as its active lifetime), measuring initial dust production rates, and comparing the object’s activity levels from one orbit to another as well as to other MBCs.

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Facilities: Blanco (DECam), CFHT (MegaCam), PS1, LDT (LMI), Hale (WASP), LCOGT, FTN (MuSCAT3), Skymapper

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