

Structure and evolution of the tidally heated hot-Jupiter KELT-9b

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Abstract

The ultra-hot Jupiter KELT-9b is observed to have a strong intrinsic heat flux characterized by $T_{\text{int}} \sim 2400 \pm 800$ K [1], and to be on a polar orbit around its parent star. Obliquity tidal heating can explain the observed high luminosity [1] suggesting parallel stellar and planetary rotation axes. Here, we investigate the possible thermal evolution and internal structure of KELT-9b. We find that its large radius of $1.9 R_E$ at a cooling time of 0.3-0.6 Gyrs requires a strong extra heat source with flat radius-evolution at present, consistent with obliquity tidal heating. However, adjusting the bulk metallicity to the observed planet radius requires that its interior follows a significantly colder adiabat characterized by $T_{\text{int}} \sim 500$ K. Merging both branches yields a tidal heat deposition zone at ~ 1 GPa and $\sim 0.8 R_p$, and consequently a low static Love number $k_2 < 0.1$ [2].

KELT-9b: high intrinsic luminosity (T_{int})

Param	G17	Cs21	used
M_* (M_{Sun})	$2.52^{+0.25}_{-0.20}$	2.26–2.82	2.52
R_* (R_{Sun})	$2.36^{+0.75}_{-0.63}$	2.28–2.54	2.362
T_* (K)	10170		10170
a (AU)	0.0346		0.0346
R_p (R_J)	$1.89^{+0.06}_{-0.05}$		1.89
M_p (M_J)	2.88 ± 0.84		2.88
A_B	—	0–0.23	0–0.25
T_{Day} (K)	—	4800 ± 80	
T_{Night} (K)	—	2770 ± 240	
T_{int}	—	2400 ± 820	800–1000

Table: Stellar and planetary parameters used here.

The energy balance equation used here (L_{eq} : planetary luminosity if in equilibrium with stellar irradiation, L_{int} : intrinsic luminosity, L_{sec} =cooling and contraction), $L_{\text{tid,o}}$: obliquity tidal heating [4]) is

$$L - L_{\text{eq}} = L_{\text{int}} = L_{\text{sec}} + L_{\text{tid,o}} + \dots$$

Using Stefan-Boltzman's law, luminosities are converted to temperatures.

$$L = L_{\text{Day}} + L_{\text{Night}}$$

$$L_D = 2\pi R_p^2 \sigma_B T_D^4$$

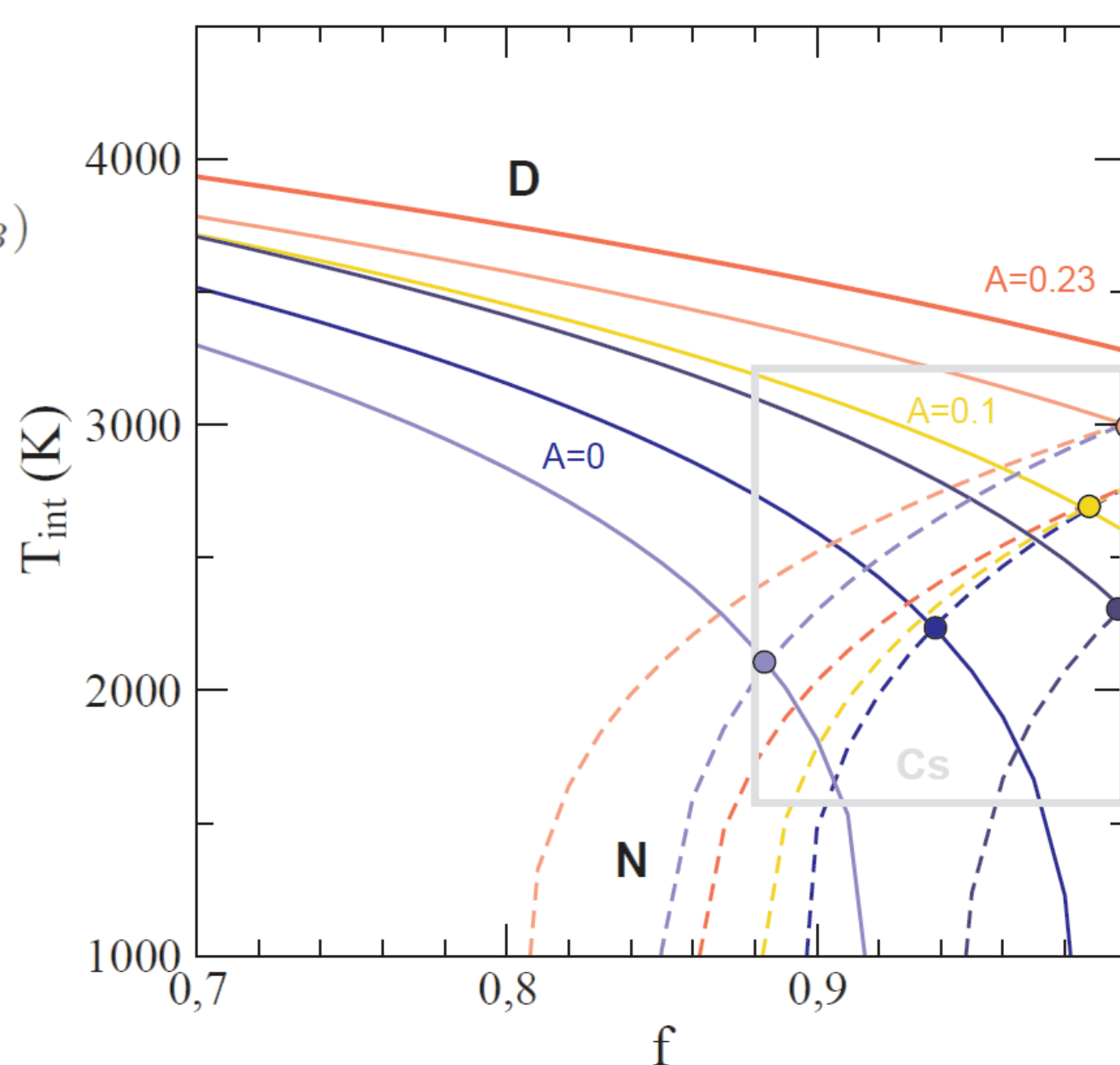
$$L_N = 2\pi R_p^2 \sigma_B T_N^4$$

$$L_{\text{eq}} = 4\pi R_p^2 \sigma_B T_*^4 (R_*/2a)^2 (1 - A_B)$$

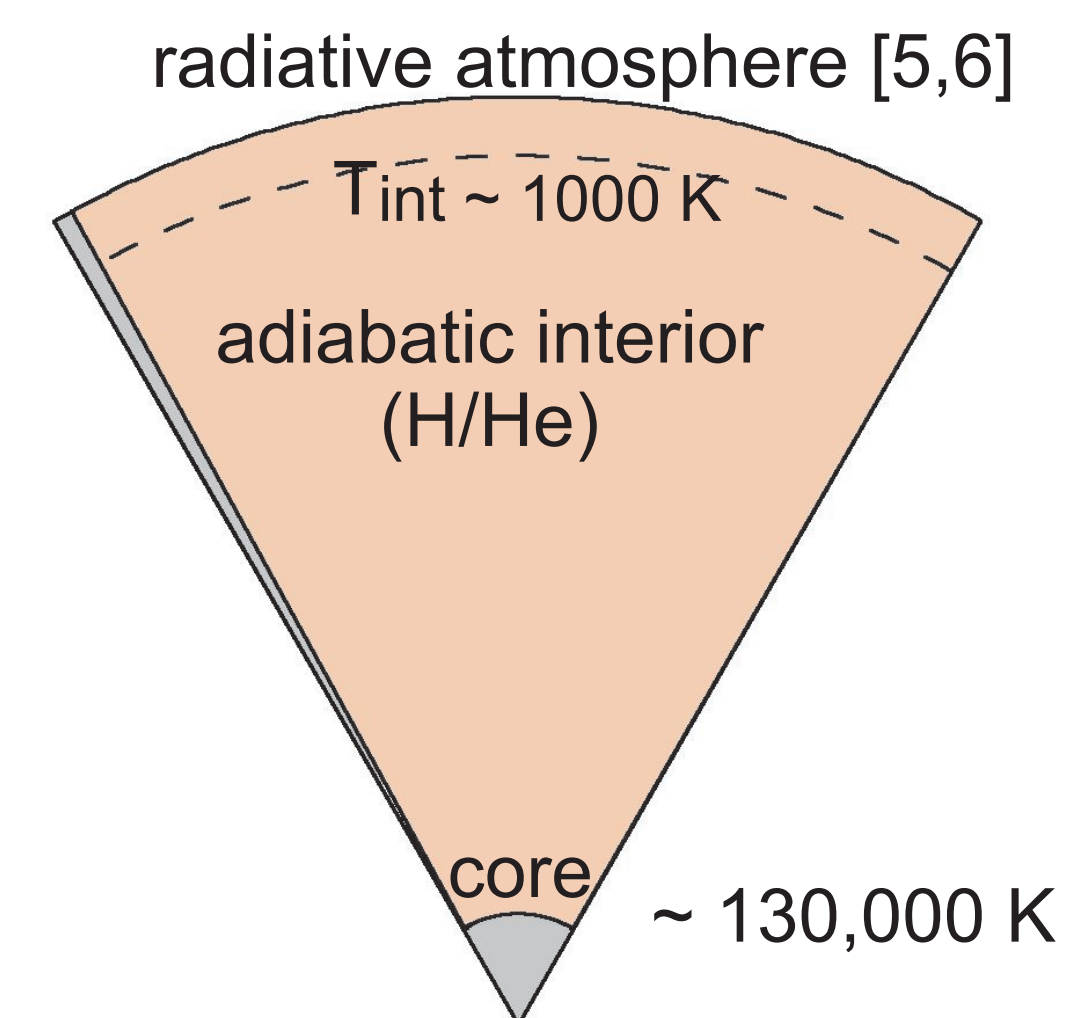
Using the observed T_D and T_N values [1], we can solve for T_{int} as a function of heat distribution factor f and Albedo A . Results are consistent with Csizmadia [1].

$$T_D^4 = T_{\text{int}}^4 + 2f T_*^4 (R_*/2a)^2 (1 - A)$$

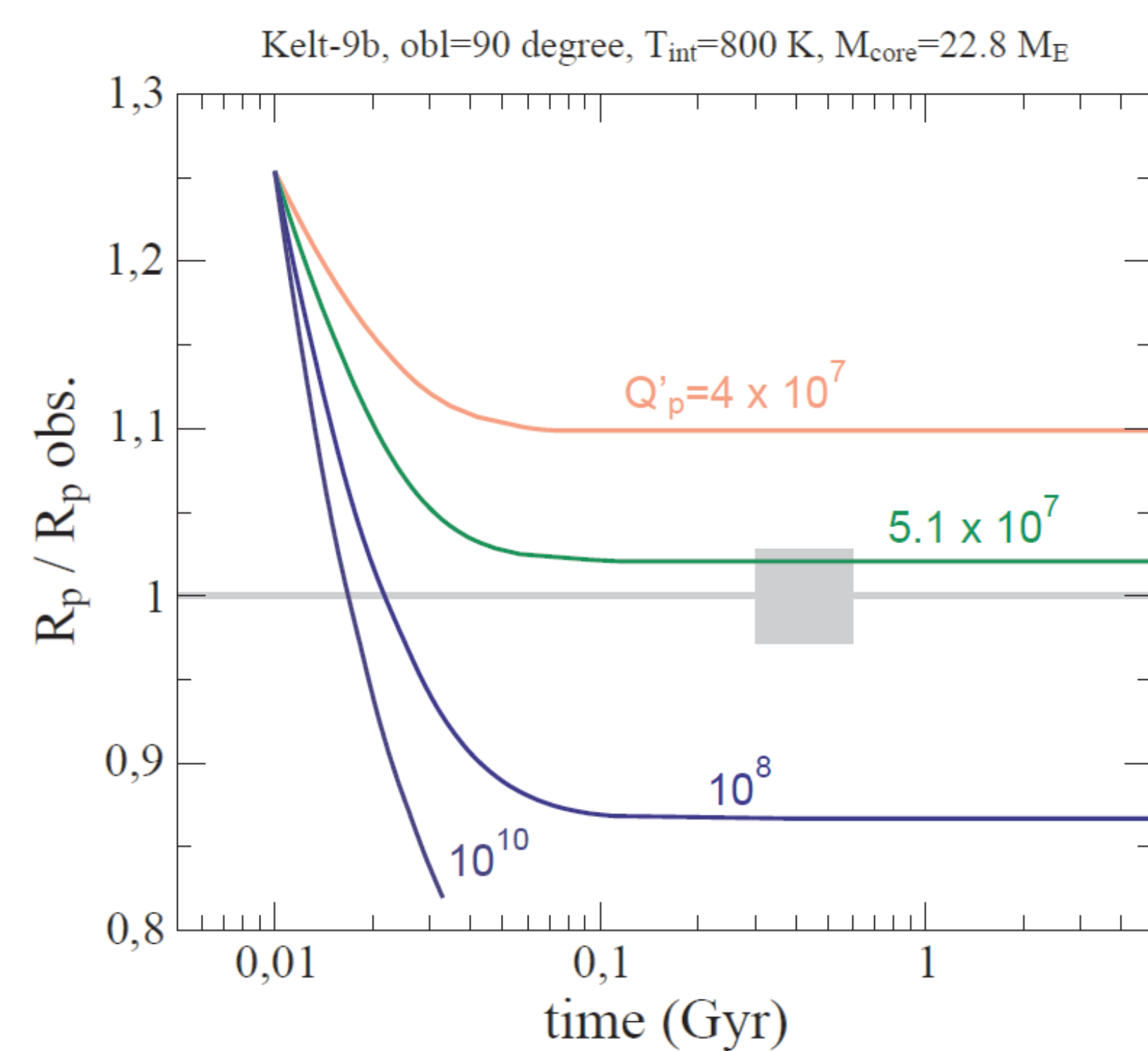
$$T_N^4 = T_{\text{int}}^4 + 2(1-f) T_*^4 (R_*/2a)^2 (1 - A)$$



Adiabatic models

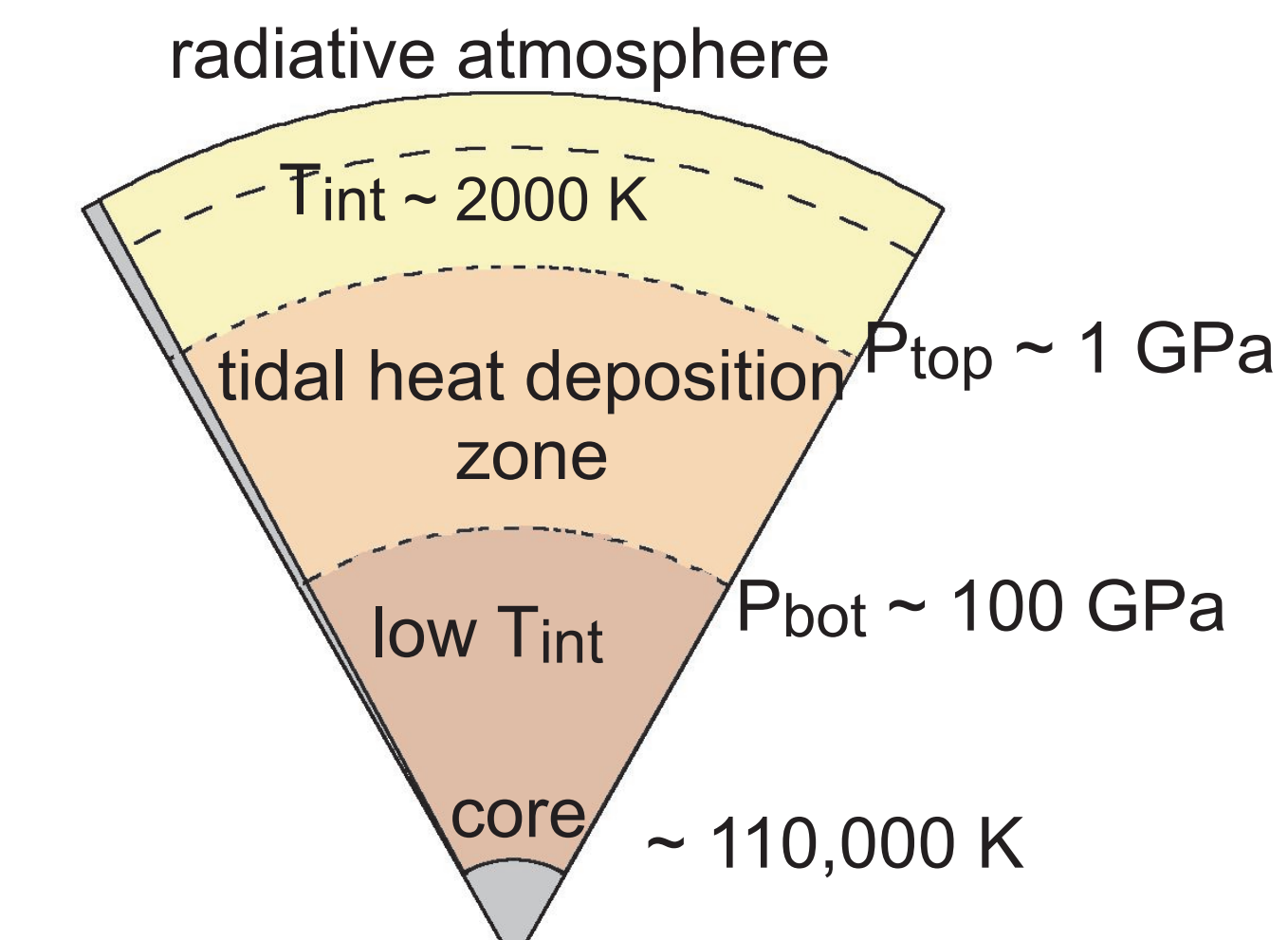


Thermal evolution



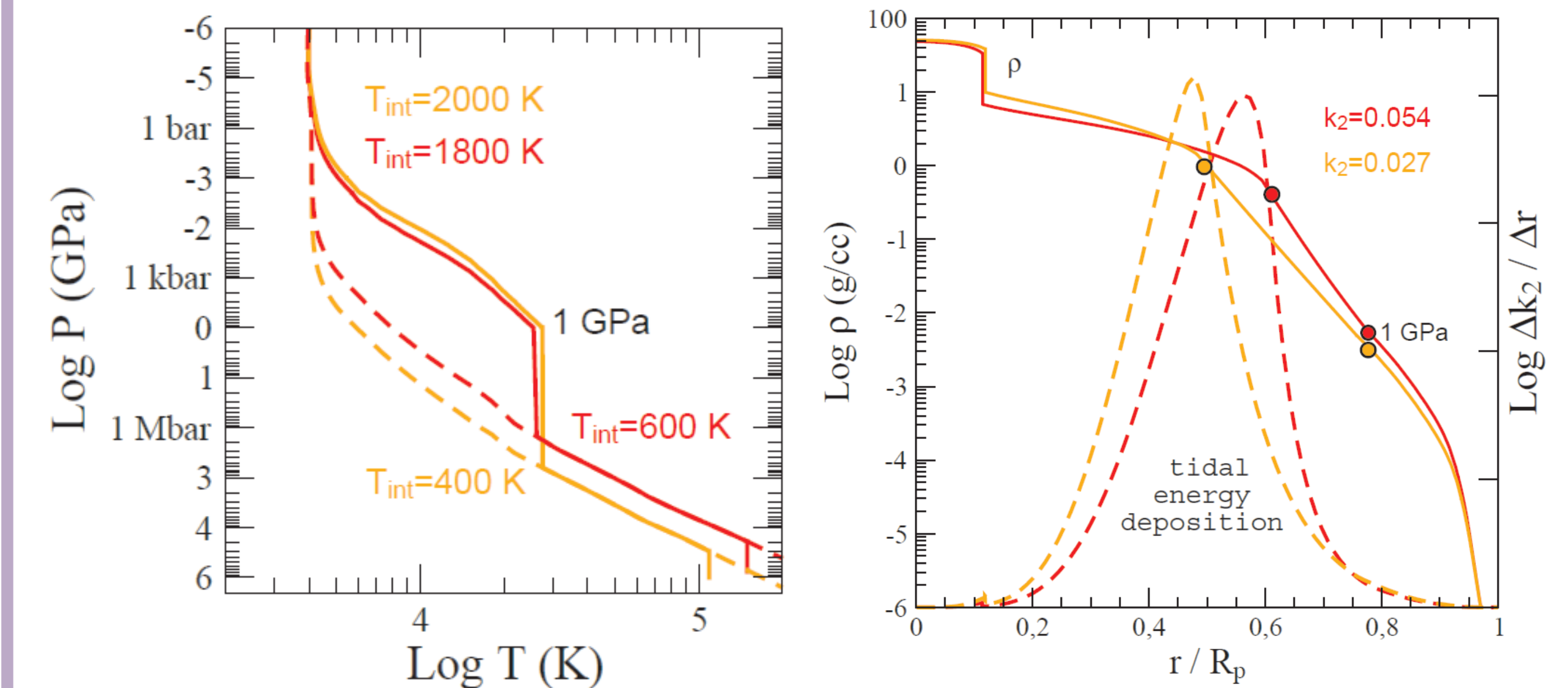
Strong extra heating that is required to obtain the observed radius after ~ 300 Mio cooling (grey box), leading to halted contraction (flat R_p (time)).

Structure models with tidal heating in a shell



ASSUMPTIONS:

- i) outer P-T profile above " P_{top} " follows solution for $T_{\text{int}} = 1800-2000$ K
 - ii) interior P-T profile below " P_{bottom} " follows cool adiabat for $T_{\text{int}} \sim 500$ K
- RESULT:**
 $P_{\text{top}} \sim 1$ GPa
 $P_{\text{bottom}} \sim 1$ Mbar (near maximum of static tidal response, dashed)
 low $k_2 < 0.1$



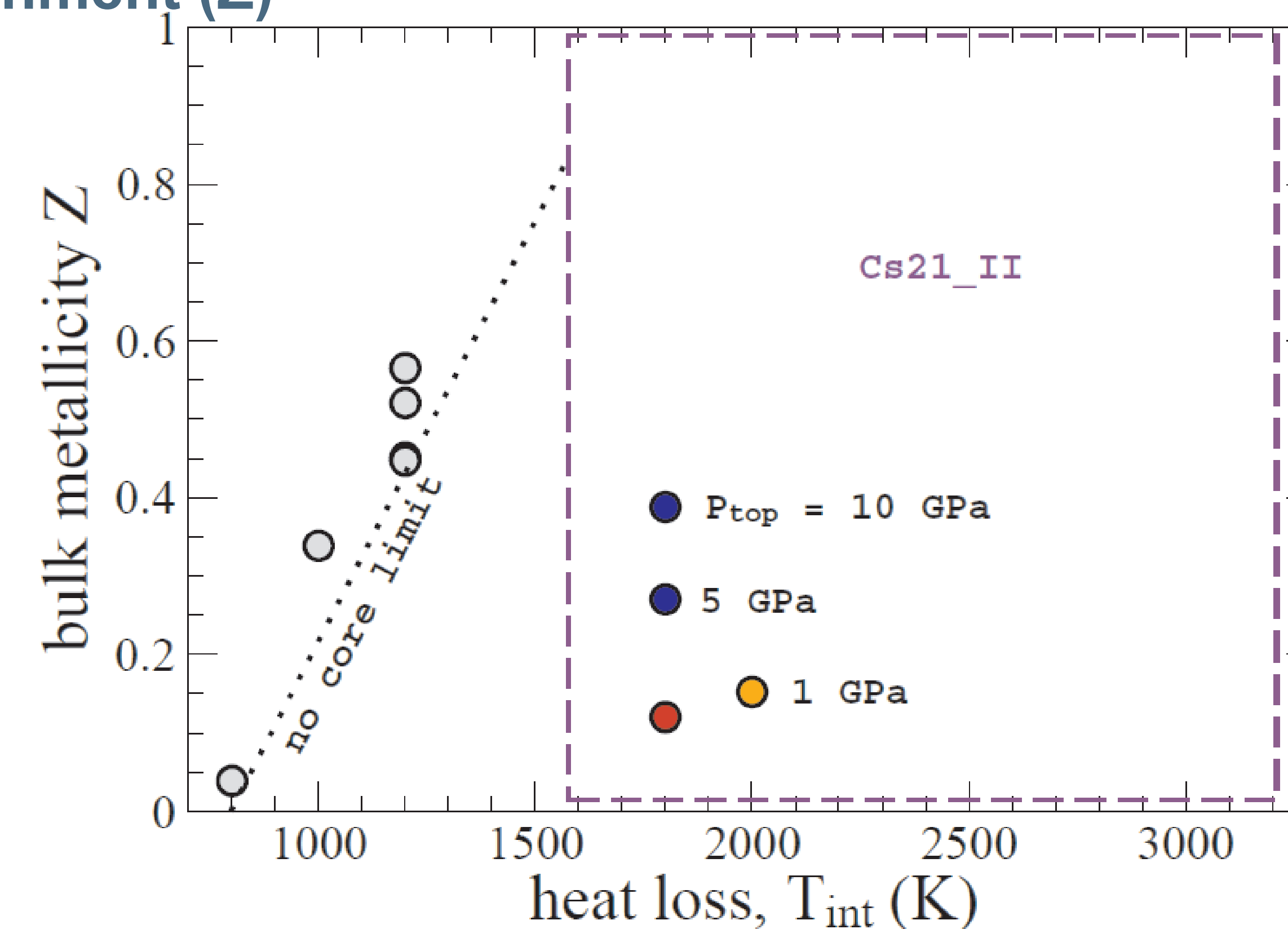
Heavy element enrichment (Z)

ADIABATIC MODELS (grey): possible T_{int} values range from ~ 800 K (zero Z) to ~ 1600 K ($Z=100\%$), inconsistent with observations.

MODELS w SHELL HEATING (color):

$P_{\text{top}} \leq 1$ GPa yields $Z \leq 20\%$ (good!)

models with $T_{\text{int}} \sim 2000$ K are possible, while $\gg 2000$ K still too hot (too high Z!)



References

- [1] Csizmadia, Smith, et al (submitted)
- [2] Nettelmann, et al (in prep)
- [3] Gaudi BS et al (2017), Nature Astr.
- [4] Millholland S (2019), ApJ
- [5] Guillot T (2010), A&A
- [6] Poser AJ et al (2019), Atmosphere

Acknowledgement

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