

### Overview of FFT-based homogenization techniques from the Galerkin point of view (slides)

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## Overview of FFT-based homogenization techniques from the Galerkin point of view Sébastien Brisard

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Navier Introduction

- Homogenization requires the solution to the so-called "corrector problem"
- Traditional numerical methods (e.g. FEM) can be costly
  - Conforming mesh
  - Large linear system
- Grid-based methods are handy in such situations!
- FFT-based methods first introduced by Moulinec and Suquet (1994)
- Since about 2010, regain of interest for these methods
- Present talk: overview, biased towards a variational point of vew
  - Brief recap on homogenization
  - The Lippmann-Schwinger equation (LS): strong and weak forms
  - Galerkin discretization of LS: consistent and asymptotically consistent discretizations
  - 3D application

## Navier Homogenization in a nutshell



## Navier Computation of the homogenized stiffness

### **Elastic equilibrium of RVE**

<b>_</b>	Heterogeneous
div ( $\mathbf{C}$ : $\mathbf{\varepsilon}$ ) = 0	
$\mathbf{\epsilon} =$ sym grad $\boldsymbol{u}$	
$BC(\boldsymbol{E})$	

### Can be complex!



### **Boundary conditions**

- Ensure that average strain is *E*
- Hill's lemma must hold

## **Example: periodic BCs**

$$\boldsymbol{u}(\boldsymbol{x}) = \boldsymbol{E} \cdot \boldsymbol{x} + \boldsymbol{u}^{\text{per}}(\boldsymbol{x})$$

Periodic

Well-suited to numerical homogenization

Kanit et al. (2003), Int J Sol Struct 40, 3647-3679

**Macroscopic stress** 

$$\Sigma = \overline{\sigma} = C^{\text{eff}} : \overline{\epsilon} = C^{\text{eff}} : E$$

## Navier The Lippmann-Schwinger equation (LS)

### **Reference material**

- Arbitrary, homogeneous stiffness: C<sub>0</sub>
- Interesting additional properties if reference material stiffer/softer than all phases

Hashin and Shtrikman (1962), J Mech Phys Sol 10, 335-342

Willis (1977), J Mech Phys Sol 25, 185-202

### The Green operator for strains

div 
$$(\mathbf{C}_0: \boldsymbol{\varepsilon} + \boldsymbol{\varpi}) = \mathbf{0}$$
  
 $\boldsymbol{\varepsilon} = \text{sym grad } \boldsymbol{u}^{\text{per}}$ 



$$\varepsilon = -\Gamma_0 * \varpi$$

### **The Lippmann-Schwinger equation**

div (**C**:
$$\varepsilon$$
)=**0**  
 $\varepsilon$ =**E**+sym grad **u**<sup>per</sup>

$$(\mathbf{C} - \mathbf{C}_{\mathbf{0}})^{-1}$$
:  $\mathbf{\tau} + \mathbf{\Gamma}_{\mathbf{0}} * \mathbf{\tau} = \mathbf{E}$   
 $\mathbf{\tau} = (\mathbf{C} - \mathbf{C}_{\mathbf{0}})$ :  $\mathbf{\varepsilon}$ 

Korringa (1973), J Math Phys 14, 509-513 Kröner (1974), Topics in Applied Continuum Mechanics, 22-38 Nemat-Nasser et al. (1982), Mech Mat 15, 163-181 Zeller and Dederichs (1973), Physica Status Solidi (B) 55, 831-842

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## Strong form $(\mathbf{C} - \mathbf{C}_0)^{-1} : \mathbf{\tau} + \mathbf{\Gamma}_0 * \mathbf{\tau} = \mathbf{E}$

Weak form: find  $\tau \in V$  such that

$$a(\mathbf{\tau}, \mathbf{\omega}) = f(\mathbf{\omega})$$
 for all  $\mathbf{\omega} \in V$ 

V: space of square integrable, second order, symmetric tensors.

The linear form: 
$$f(\boldsymbol{\varpi}) = \boldsymbol{E} : \int \boldsymbol{\varpi}$$

### The bilinear form

$$a(\mathbf{\tau}, \mathbf{\varpi}) = \mathbf{a}_{diag}(\mathbf{\tau}, \mathbf{\varpi}) + \mathbf{a}_{circ}(\mathbf{\tau}, \mathbf{\varpi})$$
$$a(\mathbf{\tau}, \mathbf{\varpi}) = \mathbf{a}_{diag}(\mathbf{\tau}, \mathbf{\varpi}) + \mathbf{a}_{circ}(\mathbf{\tau}, \mathbf{\varpi})$$
$$a_{circ}(\mathbf{\tau}, \mathbf{\varpi}) = \mathbf{\int} \mathbf{\varpi}(\mathbf{x}) : \mathbf{\Gamma}_{\mathbf{0}}(\mathbf{x} - \mathbf{y}) : \mathbf{\tau}(\mathbf{y}) d\mathbf{x} d\mathbf{y}$$

## Navier Galerkin discretization of the LS equation

Find 
$$\mathbf{\tau} \in \mathbf{V}$$
 such that  $a_{diag}(\mathbf{\tau}, \mathbf{\omega}) + a_{circ}(\mathbf{\tau}, \mathbf{\omega}) = f(\mathbf{\omega})$  for all  $\mathbf{\omega} \in \mathbf{V}$ 

### **Consistent discretization**

Brisard and Dormieux (2010), Comp Mat Sci 49, 663-671

Evaluation over 
$$V^h$$
 remains difficult!  
Find  $\boldsymbol{\tau}^h \in V^h$  such that  $a_{diag}(\boldsymbol{\tau}^h, \boldsymbol{\varpi}^h) + a_{circ}(\boldsymbol{\tau}^h, \boldsymbol{\varpi}^h) = f(\boldsymbol{\varpi}^h)$  for all  $\boldsymbol{\varpi}^h \in V^h$   
Space of cell-wise constant polarization fields

Asymptotically consistent discretization: exact evaluation is not necessary!

Find 
$$\mathbf{\tau}^h \in V^h$$
 such that  $a_{diag}(\mathbf{\tau}^h, \mathbf{\varpi}^h) + \begin{bmatrix} a_{circ}^h, \mathbf{\varpi}^h \end{bmatrix} = f(\mathbf{\varpi}^h)$  for all  $\mathbf{\varpi}^h \in V^h$ 

Brisard and Dormieux (2012), Comp Meth Appl Mech Eng 217-220, 197-212

Asymptotically consistent approximation

# Navier Asymptotically consistent approximations

- Periodic Green operator for strains is in fact given by an infinite Fourier series
- Various estimates of this series for cell-wise constant functions
  - **Truncation** of high frequencies: Moulinec and Suquet (1994, 1998)
  - **Exact** (up to round-off errors): Brisard and Dormieux (2010)
  - Filtering of high frequencies: Brisard and Dormieux (2012)
  - Finite elements approximation: Yvonnet (2012)
  - **Finite differences** approximation: Willot et al. (2014), Willot (2015)
- All these approximations can be fitted in the general framework introduced here!
- If appropriately implemented, they can be switched on-the-fly in a simulation.

Moulinec and Suquet (1994), CR Acad Sci II **318**, 1417-1423 Moulinec and Suquet (1998), Comp Meth Appl Mech Eng **57**, 69-94 Brisard and Dormieux (2010), Comp Mat Sci **49**, 663-671 Brisard and Dormieux (2012), Comp Meth Appl Mech Eng **217-220**, 197-212 Yvonnet (2012), Int J Num Meth Eng **92**, 178-205 Willot et al. (2014), Int J Num Meth Eng **98**, 518-533 Willot (2015), CR Acad Sci Mec **343**, 232-245

## Navier The underlying linear system

**Discrete variational problem results in a linear system** 



### Solving the linear system

- Matrix is not sparse: matrix-free approach
- Use iterative linear solvers
  - Fixed-point iterations: Moulinec and Suquet (1994, 1998)
  - Augmented-Lagrangian: Michel et al. (2001)
  - Conjugate Gradient: Brisard and Dormieux (2010)
- Use FFT to compute matrix-vector products (Moulinec and Suquet, 1994, 1998)

Moulinec and Suquet (1994), CR Acad Sci II 318, 1417-1423 Moulinec and Suquet (1998), Comp Meth Appl Mech Eng 157, 69-94 Michel et al. (2001), Int J Num Meth Eng 52, 139-160 Brisard and Dormieux (2010), Comp Mat Sci 49, 663-671

## Navier Example: 3D microstructure (1/2)



### **Microstructural parameters**

- Flat spheroids (1/8 aspect ratio)
- Dense packing (60%)
- Large model (10000 particles)
- Moderate contrast (inclusions 100 times stiffer than matrix)

## The simulation

- Home-made code
  - Python + Cython + FFTW + MPI
  - Very flexible implementation
  - Soon to be open-sourced (contact me!)
- Simulations run on two servers
  - Intel Xeon X5690, 3.47GHz, 192 Go
  - Intel Xeon E5-2643, 3.30GHz, 762 Go
- Most simulations run on 16 cores

# Navier Example: 3D microstructure (2/2)



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# Navier Conclusion and outlook

- Summary
  - General, unified framework for FFT-based homogenization techniques
  - All avatars of this method (Moulinec & Suquet; Michel, Moulinec and Suquet; Yvonnet; Willot; Monchiet; ...) fit into this unified framework
  - Clear distinction between discretization and iterative solution of the discretized problem: any discrete Green operator can be combined with any iterative linear solver
- Work in progress
  - **A priori** error estimates: with F. Legoll (Navier Laboratory, Ecole des Ponts ParisTech)
  - **A posteriori** error estimates: with L. Chamoin (LMT, ENS Cachan)
- Open questions
  - Matrix-free preconditioners
  - What is the "best" discrete Green operator?
  - What is the "best" reference material?

# Thank you for your attention!

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