

# Improved modelling of alpha-particle emission in nucleon induced reactions

A.Yu. Konobeyev<sup>1</sup>, D. Leichtle<sup>1</sup>, A.J. Koning<sup>2</sup>

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- <sup>1</sup> INR, Karlsruhe Institute of Technology, 76344 Eggenstein-Leopoldshafen, Germany
- <sup>2</sup> Nuclear Data Section, International Atomic Energy Agency, A-1400 Vienna, Austria

#### Impressum

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### Abstract

This report discusses the phenomenological approach proposed to estimate the contribution of direct processes to the emission of  $\alpha$ -particles in nucleon induced reactions. Using available measured energy distributions, the values of the parameters required for the calculations are obtained. The analysis was performed using the TALYS code.

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#### 1. Introduction

The modelling of  $\alpha$ -particle energy distributions in nucleon induced reactions at intermediate primary energies remains a challenge, and it is too early to talk about a final solution. With caution, one can only say about the relative success of modeling the equilibrium and pre-equilibrium  $\alpha$ -emission. The calculation or estimation of the contribution of direct processes in  $\alpha$ -particle spectra is a special task that cannot be easy solved for most nuclei because of the lack of experimental data and complexity of theoretical analysis.

A paragon for calculating the spectra of  $\alpha$ -particles produced in direct interactions is the works of Gadioli et al. [1-3], which combine the application of a rigorous theoretical approach with the use of detailed measured  $\alpha$ -spectra. For most targets, the use of such a method, for obvious reasons, seems difficult.

In this work the question is raised whether it is possible to describe the hard part of the  $\alpha$ -particle spectra using a less rigorous approach than that used in Refs.[1-3], but which, to a certain extent, would give acceptable or close results. Such a method would have to be more suitable for mass calculations and for improving the quality of predictions of existing models and codes. With a certain degree of caution this question is answered positively.

The phenomenological approach discussed and proposed in this work does not claim to be rigorous, but its use provides some advantages and increases, to a certain extent, the agreement between measured and predicted data.

The proposed method of calculation is used within the TALYS code [4,5] and, naturally, will profit from a combination of models already implemented in the code.

The approach is discussed in Section 2. Section 3 presents a comparison of the calculated and measured  $\alpha$ -particle spectra and the discussion.

#### 2. Brief description of method of calculation

The TALYS code implements the Kalbach model [6] for calculating the preequilibrium emission spectra of complex particles. The use of the model does not imply an explicit consideration of direct processes. The present work shows that the situation can be improved to some extent. The simplest attempt to describe direct processes in some way without resorting to a rigorous theory is to use the approach from Refs.[7,8]. In spite of its formulation for a hybrid model [9], and as a possible phenomenological solution, it does not lead to contradictions when using the exciton model [10] to simulate the whole pre-equilibrium emission. Using results of Refs.[7,8] approach, the contribution of direct processes to  $\alpha$ -particle emission can be estimated as follows

$$\frac{d\sigma^{D}}{d\varepsilon_{\alpha}} \sim \beta_{1} \exp(-\beta_{2}(U-\beta_{3})^{2}) \frac{\lambda_{\alpha}^{e}(\varepsilon_{\alpha})}{\lambda_{\alpha}^{e}(\varepsilon_{\alpha}) + \lambda_{\alpha}^{+}(\varepsilon_{\alpha})} g_{\alpha},$$
(1)

where  $\lambda_{\alpha}^{e}(\varepsilon_{\alpha})$  and  $\lambda_{\alpha}^{+}(\varepsilon_{\alpha})$  are emission and absorption rates of  $\alpha$ -particles, which are calculated similarly to Ref.[7,8],  $\beta_{i}$  are parameters, other symbols are conventional.

There are three parameters in Eq.(1), the sensitivity of the calculations to which and which importance, can be checked by comparing the measured and calculated  $\alpha$ -particle spectra.

In the TALYS code, the Kalbach continuum  $\alpha$ -particle spectrum is "mapped" on discrete states, which aptly simulates the effect of real resolution of measurements. The study shows that the change of normalization coefficient of one of the corresponding values "fac1" or "fac2" in the TALYS subroutine "Spectra" together with the use of Eq.(1), can further improve the agreement with the experimental data.

The routines providing calculations using Eq.(1) from Refs.[11,12] were added to the TALYS code [4,5].

The next Section discusses the comparison of calculated  $\alpha$ -particle spectra with experimental data, the use of parameters, and the problems associated with calculations.

#### 3. Results and discussion

The available measured energy distributions of  $\alpha$ -particles in neutron induced reactions for 38 target nuclei and in proton induced reactions for 43 target nuclei from <sup>12</sup>C to <sup>238</sup>U [1,3,13-72] have been analysed.

Not all experimental data, or rather a smaller part of them, are suitable for studying the contribution of direct processes to the  $\alpha$ -particle energy distributions. Such a study requires detailed information about the hard part of the spectra, which was not always obtained in the measurements.

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The parameters were chosen to achieve the best visible agreement with the experimental data. If possible, the parameters were taken to be the same for different targets or, practically, were not used, as in the case of the  $\beta_3$  parameter.

The study shows that in most cases, when suitable experimental information is available, the use of discussed phenomenological model for calculations significantly improves the agreement between the calculated and measured data.

Figures 1-68 show examples of calculated α-particle energy distributions together with measured data. The results for neutron induced reactions are shown in Figs.1-51 and for proton induced reactions in Figs.52-68.

The resulting value of the  $\beta_2$  parameter for all reactions except (n,x $\alpha$ ) on <sup>95</sup>Mo, <sup>142,143,144</sup>Nd, and <sup>144,149</sup>Sm was taken equal to 0.3, assuming that subroutines from Refs.[11,12] are used for calculations with Eq.(1). The  $\beta_3$  value is equal to zero for all considered reactions except n+<sup>60,61,62,63</sup>Ni. The value of  $\beta_1$  varies from 10<sup>-5</sup> to 10<sup>-2</sup> with a general trend that heavier target nuclei have lower  $\beta_1$  values.

The normalization coefficient  $C_f$  for "fac2" ranges from 1 (no changes), to 10. The number of cases of coefficient value 2 is 25%, value 5 is 47%, and 10 is 10%.

All parameters  $\beta_i$  and  $C_f$  were treated as independent of incident nucleon energy.

The results indicate the principal possibility of using the discussed phenomenological approach to improve the agreement between the calculated and experimental  $\alpha$ -particle energy distributions. In most cases, as mentioned above, it is necessary to know only the values of  $\beta_1$  and  $C_f$  parameters. Of course, in the absence of necessary experimental information, it is possible to use global or systematic values of these parameters.

Obviously, the use of discussed method cannot replace the rigorous theoretical calculations of the  $\alpha$ -particle spectra. However, even in this case the experimental information is necessary to obtain the absolute values of the spectra [1].

The described approach can apparently be extended to describe the emission of other complex particles using the models implemented in the TALYS code.

Hopefully, new measurements of α-particle spectra will provide more information necessary to improve and refine this approach, and will create the possibility of a detailed theoretical analysis of the spectra.

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Fig.1 Comparison between experimental  $\alpha$ -particle energy distributions and calculations using the TALYS code [4,5] without (blue line) and using (red line) the discussed phenomenological model for the estimation of contribution of direct processes to the  $\alpha$ -particle emission for the n+<sup>12</sup>C reaction at the primary neutron energy equal to 50 MeV.



Fig.2 The same as in Fig.1 but for the  $n+^{12}C$  reaction at  $E_n=62.7$  MeV.



Fig.3 The same as in Fig.1 but for the  $n+^{12}C$  reaction at  $E_n=95.6$  MeV.



Fig.4 The same as in Fig.1 but for the  $n+{}^{16}O$  reaction at  $E_n=37.5$  MeV.



Fig.5 The same as in Fig.1 but for the  $n+{}^{16}O$  reaction at  $E_n=62.7$  MeV.



Fig.6 The same as in Fig.1 but for the  $n+{}^{16}O$  reaction at  $E_n=95.6$  MeV.



Fig.7 The same as in Fig.1 but for the  $n+^{27}AI$  reaction at  $E_n=31.5$  MeV.



Fig.8 The same as in Fig.1 but for the  $n+^{27}AI$  reaction at  $E_n=45$  MeV.



Fig.9 The same as in Fig.1 but for the  $n+^{27}AI$  reaction at  $E_n=62.7$  MeV.



Fig.10 The same as in Fig.1 but for the  $n+^{28}Si$  reaction at  $E_n=21$  MeV.



Fig.11 The same as in Fig.1 but for the  $n+^{28}$ Si reaction at E<sub>n</sub>=34 MeV.



Fig.12 The same as in Fig.1 but for the  $n+^{28}Si$  reaction at  $E_n=95.6$  MeV.



Fig.13 The same as in Fig.1 but for the n+50 Cr reaction at E<sub>n</sub>=14.8 MeV.



Fig.14 The same as in Fig.1 but for the n+Fe reaction at  $E_n=14.1$  MeV.



Fig.15 The same as in Fig.1 but for the n+Fe reaction at  $E_n=41$  MeV.



Fig.16 The same as in Fig.1 but for the n+Fe reaction at  $E_n=45$  MeV.



Fig.17 The same as in Fig.1 but for the n+Fe reaction at  $E_n=49$  MeV.



Fig.18 The same as in Fig.1 but for the n+Fe reaction at  $E_n=53.5$  MeV.



Fig.19 The same as in Fig.1 but for the  $n+^{54}$ Fe reaction at E<sub>n</sub>=14.8 MeV.



Fig.20 The same as in Fig.1 but for the  $n+{}^{56}$ Fe reaction at  $E_n=14.8$  MeV.



Fig.21 The same as in Fig.1 but for the  $n+^{59}$ Co reaction at  $E_n=37.5$  MeV.



Fig.22 The same as in Fig.1 but for the  $n+^{59}$ Co reaction at  $E_n=53.5$  MeV.



Fig.23 The same as in Fig.1 but for the n+58Ni reaction at E<sub>n</sub>=14.8 MeV.



Fig.24 The same as in Fig.1 but for the  $n+^{60}Ni$  reaction at  $E_n=14.1$  MeV.



Fig.25 The same as in Fig.1 but for the  $n+^{65}$ Cu reaction at  $E_n=14.8$  MeV.



Fig.26 The same as in Fig.1 but for the  $n+^{89}$ Y reaction at  $E_n=14.8$  MeV.



Fig.27 The same as in Fig.1 but for the  $n+^{90}Zr$  reaction at  $E_n=14.3$  MeV.



Fig.28 The same as in Fig.1 but for the  $n+^{90}$ Zr reaction at E<sub>n</sub>=18.1 MeV.



Fig.29 The same as in Fig.1 but for the  $n+^{91}Zr$  reaction at  $E_n=14.3$  MeV.



Fig.30 The same as in Fig.1 but for the  $n+^{91}Zr$  reaction at  $E_n=18.1$  MeV.



Fig.31 The same as in Fig.1 but for the  $n+^{92}$ Mo reaction at  $E_n=14.8$  MeV.



Fig.32 The same as in Fig.1 but for the  $n+^{96}Mo$  reaction at  $E_n=14.8$  MeV.



Fig.33 The same as in Fig.1 but for the  $n+^{142}Nd$  reaction at  $E_n=14.3$  MeV.



Fig.34 The same as in Fig.1 but for the  $n+^{142}Nd$  reaction at  $E_n=18.1$  MeV.



Fig.35 The same as in Fig.1 but for the  $n+^{143}Nd$  reaction at  $E_n=12.3$  MeV.



Fig.36 The same as in Fig.1 but for the  $n+^{143}Nd$  reaction at  $E_n=14.1$  MeV.



Fig.37 The same as in Fig.1 but for the  $n+^{143}Nd$  reaction at  $E_n=18.2$  MeV.



Fig.38 The same as in Fig.1 but for the  $n+^{143}Nd$  reaction at  $E_n=20.6$  MeV.

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Fig.39 The same as in Fig.1 but for the  $n+^{144}Nd$  reaction at  $E_n=14.3$  MeV.



<sup>144</sup>Nd(n,x $\alpha$ ) E<sub>n</sub>=18.1 MeV

Fig.40 The same as in Fig.1 but for the  $n+^{144}Nd$  reaction at  $E_n=18.1$  MeV.



Fig.41 The same as in Fig.1 but for the  $n+^{144}$ Sm reaction at  $E_n=14.3$  MeV.



Fig.42 The same as in Fig.1 but for the  $n+^{144}$ Sm reaction at  $E_n=18.1$  MeV.



Fig.43 The same as in Fig.1 but for the  $n+^{147}$ Sm reaction at E<sub>n</sub>=12.4 MeV.



Fig.44 The same as in Fig.1 but for the  $n+^{147}$ Sm reaction at  $E_n=14.1$  MeV.



Fig.45 The same as in Fig.1 but for the  $n+^{147}$ Sm reaction at  $E_n=18.2$  MeV.



Fig.46 The same as in Fig.1 but for the  $n+^{147}$ Sm reaction at  $E_n=19.5$  MeV.



Fig.47 The same as in Fig.1 but for the  $n+^{149}$ Sm reaction at  $E_n=12.3$  MeV.



Fig.48 The same as in Fig.1 but for the  $n+^{149}$ Sm reaction at  $E_n=14.1$  MeV.



Fig.49 The same as in Fig.1 but for the  $n+^{149}$ Sm reaction at  $E_n=18.2$  MeV.



Fig.50 The same as in Fig.1 but for the  $n+^{238}U$  reaction at  $E_n=37.5$  MeV.



Fig.51 The same as in Fig.1 but for the  $n+^{238}U$  reaction at  $E_n=62.7$  MeV.



Fig.52 The same as in Fig.1 but for the  $p+^{12}C$  reaction at  $E_p=61$  MeV.



Fig.53 The same as in Fig.1 but for the  $p+^{16}O$  reaction at  $E_p=61$  MeV.



Fig.54 The same as in Fig.1 but for the  $p+^{27}AI$  reaction at  $E_p=61.7$  MeV.



Fig.55 The same as in Fig.1 but for the  $p+^{28}Si$  reaction at  $E_p=26$  MeV.



Fig.56 The same as in Fig.1 but for the  $p+^{54}$ Fe reaction at  $E_p=28.8$  MeV.



Fig.57 The same as in Fig.1 but for the  $p+^{54}$ Fe reaction at  $E_p=61.5$  MeV.



Fig.58 The same as in Fig.1 but for the  $p+^{56}$ Fe reaction at  $E_p=61.5$  MeV.



Fig.59 The same as in Fig.1 but for the  $p+^{93}Nb$  reaction at  $E_p=24.6$  MeV.



Fig.60 The same as in Fig.1 but for the  $p+^{96}Mo$  reaction at  $E_p=18$  MeV.



Fig.61 The same as in Fig.1 but for the  $p+^{98}Mo$  reaction at  $E_p=18$  MeV.



Fig.62 The same as in Fig.1 but for the  $p+^{106}Pd$  reaction at  $E_p=18$  MeV.



Fig.63 The same as in Fig.1 but for the  $p+^{112}Cd$  reaction at  $E_p=18$  MeV.



Fig.64 The same as in Fig.1 but for the  $p+^{118}$ Sn reaction at  $E_p=18$  MeV.



Fig.65 The same as in Fig.1 but for the  $p+^{118}C$  reaction at  $E_p=44.3$  MeV.



Fig.66 The same as in Fig.1 but for the  $p+^{120}$ Sn reaction at  $E_p=18$  MeV.



Fig.67 The same as in Fig.1 but for the  $p+^{120}$ Sn reaction at  $E_p=61.5$  MeV.



Fig.68 The same as in Fig.1 but for the  $p+^{128}$ Te reaction at  $E_p=18$  MeV.

## 4. Conclusion

In this work, a simple phenomenological model is proposed for calculating the contribution of direct processes to the production of  $\alpha$ -particles in reactions induced by nucleons. The calculated and experimental data for 38 target nuclei irradiated by neutrons and 44 target nuclei from <sup>12</sup>C to <sup>238</sup>U irradiated by protons are compared. The study shows that relative good agreement with the experimental data, in most cases, can be achieved using only two parameters. It is hoped that the proposed method will improve the quality of mass calculations of the spectra of  $\alpha$ -particles.

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