

New magnetic and structural results for uniformly spaced, alternately spaced, and ladder-like copper (II) linear chain compounds

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Results of recent synthetic, structural, and magnetic work on exchange coupled copper (II) ions bridged by nonmagnetic ligands are discussed, including the quasi-linear chain compound bis(dimethyldithiocarbamate)copper (II) in which the copper (II) ions are antiferromagnetically exchange coupled. An unusual correlation exists in the structural and magnetic data for $[\text{Cu}(\text{dmtc})_2]_\infty$ and other known sulfur bridged copper compounds. Data are presented for the alternately spaced linear chain compounds di- μ -chlorobis(4-methylpyridine)copper (II), and a series of dionebis-(thiosemicarbazato)copper (II) compounds. Exchange coupling constants and alternation parameters were obtained from expressions which were empirically generated from calculations on rings of ten $S = 1/2$ ions for $kT/|J| \gtrsim 0.5$ and $J < 0$. Finally, new results for recently discovered spin ladder systems are presented.

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INTRODUCTION

The cluster approach as developed by Bonner and Fisher (1) in their classical work on the Heisenberg chain with $S=1/2$, has received much less attention in applications to exchange coupled systems with lower symmetry geometries. This is a surprising situation since the results of Bonner and Fisher have been used extensively by chemists and physicists for the analysis of magnetic properties of exchange coupled low-dimensional $S=1/2$ chains. A limited number of studies have been carried out. Weng followed Bonner and Fisher's cluster approach in his calculation of the magnetic properties of Heisenberg chains with $S \geq 1$, (2) but this work has not been widely applied since exchange coupled chains with $S > 1/2$ frequently show pronounced anisotropy in the exchange, and one of the anisotropic models, the Ising chain or the XY-chain is required for the analysis of the data.

Bonner and coworkers (3,4) have continued investigating the application of the cluster approach to more complicated spin systems. These include the alternately spaced Heisenberg chain and the interesting spin ladder systems. It is interesting to know that the most thoroughly studied transition metal example of an alternately spaced Heisenberg antiferromagnet is $\text{Cu}(\text{NO}_3)_2 \cdot 2.5\text{H}_2\text{O}$, a compound which has a ladder-like structure at room temperature. (5,6) The initial attempts to analyze the magnetic data for $\text{Cu}(\text{NO}_3)_2 \cdot 2.5\text{H}_2\text{O}$ were in terms of the spin ladder. This is a very good example of a system which exhibits a structure at high temperature which is inconsistent with the observed magnetic properties at low temperatures. There are now several bona fide examples of alternately spaced Heisenberg exchange coupled $S=1/2$ linear chains. These include the pyrazine-bridged binuclear copper acetate chain, (7) $[\text{Cu}_2(\text{OAc})_4\text{pyr}]_\infty$, the low temperature form of tetrathiofulvalenium bis-cis(1,2-perfluoromethylethylene-1,2-dithiolato)copper(II) (8) which results from a spin Peierls transition at 12 K, catena-di- μ -bromo-(N-methylimidazole)copper(II) (9), and a series of diketobis(thiosemicarbazato)copper(II) complexes related to the carcinostatic and carcinolytic agent 3-ethoxy-2-oxobutaldehyde-bis(thiosemicarbazato)copper(II), Cu-KTS. (10) The magnetic properties of several of these alternately spaced Heisenberg exchange-coupled $S = 1/2$ linear chains have been characterized in our laboratory. (11) Basically, because of the general interest in alternating Heisenberg antiferromagnets, and because of the absence of convenient procedures for the analysis of the magnetic properties, we have repeated the calculations reported by Duffy and Barr (12) for such a model system to obtain numerical data

for the generation of an expression for magnetic susceptibility in terms of the exchange coupling constant and the alternation parameter. The generation of the magnetic susceptibility expression and the application of the new expression for the analysis of new magnetic data for several of these compounds will be discussed.

Since bona fide examples $S = 1/2$ ladder systems have not been available for study, there has been little stimulation for theoretical work. However, in recent years, three examples of copper compounds with ladder structures have been reported. These include catena-dibromo(2-aminomethyl)-pyridine-copper (II), (13) catena-dibromo(2-methyl-1,2-diaminopropane)copper(II), (13) and catena-trichlorohydrazinium-copper(II) (14). Weller (15) has carried out calculations of the magnetic properties of ladder systems with a range of exchange coupling constants, and has analyzed the magnetic susceptibility for the known $S=1/2$ ladder systems. Preliminary results of these studies on catena-trichlorohydraziniumcopper(II) are presented here.

THE LINEAR CHAIN MODEL

For an infinite linear chain the Hamiltonian may be written as

$$H = -2J \sum_{i=1}^n \{ \alpha S_{i+1}^z S_i^z + \gamma (S_i^x S_{i+1}^x + S_i^y S_{i+1}^y) \} - g\beta H \sum_{i=1}^n H \cdot S_i \quad (1)$$

If $\alpha = 1$ and $\gamma = 1$, the isotropic Heisenberg model is produced. If $\alpha = 1$ and $\gamma = 0$, the anisotropic Ising model is obtained, and if $\alpha = 0$ and $\gamma = 1$, the anisotropic X-Y model results. Copper(II) systems are usually best described by the Heisenberg model, but in some cases γ -values slightly less than one have been deduced from experimental data.

In 1964, Bonner and Fisher (1) carried out calculations on systems of chains containing from 4 to 12 spins. The results for the 11 and 12 membered chains were used to extrapolate to the infinite-length chain with very good agreement above $kT/|J|$ of 0.5. The maximum in susceptibility is uniquely defined by the following relations:

$$\frac{kT}{|J|} \max = 1.282 \quad \text{and} \quad \frac{\chi_{\max} |J|}{N g^2 \mu_B^2} = 0.07346.$$

Hall (16) noted that the shape of the reduced magnetic susceptibility versus temperature curve obtained by Bonner is very similar to a serpentine curve, and fit the quadratic/cubic function

$$f(x) = (Ax^2 + Bx + C)/(x^3 + Dx^2 + Ex + F)$$

to the theoretically calculated results tabulated by Bonner. (17) A fit, with the greatest deviation between the calculated values for reduced magnetic susceptibility and those obtained by Bonner being less than 0.5%, was obtained with the parameters $A = 0.25$, $B = 0.14995$, $C = 0.30094$, $D = 1.9862$, $E = 0.68854$, and $F = 6.0626$. The magnetic susceptibility of chains of isotropically exchange coupled $S=1/2$ ions is

$$\chi_m = \frac{Ng^2\mu_B^2}{kT} \cdot \frac{A + Bx^{-1} + Cx^{-2}}{1 + Dx^{-1} + Ex^{-2} + Fx^{-3}} \quad (2)$$

The maximum in reduced susceptibility calculated from this expression occurs at

$$\frac{kT}{|J|} \max = 1.297 \quad \text{and} \quad \frac{\chi_{\max}}{Ng^2\mu_B^2} \frac{|J|}{2} = 0.07337.$$

The agreement with the values of Bonner and Fisher (1) which are given above are excellent, and this expression allows data analysis using computer programs rather than graphical methods.

MAGNETIC SUSCEPTIBILITY OF THE LINEAR CHAIN BIS(DIMETHYLDITHIOCARBAMATO)COPPER(II)

Magnetic susceptibility data (18) for a powdered sample of the linear chain compound bis(dimethyldithiocarbamato)copper(II) is plotted as a function of temperature in Figure 1. It may be seen that the susceptibility exhibits a maximum value at about 2.5 K. This behavior of the magnetic susceptibility is indicative of a magnetically condensed substance, and since the compound has a linear chain structure, (19) the data were initially analyzed using the Heisenberg linear chain theory of Bonner and Fisher, (1) and the analytical expression developed by Hall. (16) Since the chains in $[\text{Cu}(\text{dmtc})_2]_{\infty}$ are in rather close contact, substantial interchain interactions were anticipated. The problem was treated as one involving Heisenberg interactions along a chain with a correction for interchain interactions (19a) being accounted for in the expression

$$\chi_m = \chi_H / (1 - 2zJ'\chi_H / Ng^2\mu_B^2) \quad (3)$$

where χ_H is the magnetic susceptibility of an isolated Heisenberg chain of $S = 1/2$ ions, z is the number of near neighbors in adjacent chains, and J' is the interchain exchange parameter. For the case of $[\text{Cu}(\text{dmtc})_2]_{\infty}$, z was taken to be 4. The magnetic susceptibility expression thus modified was fit to the experimental magnetic susceptibility data, and the best fit yielded the $\langle g \rangle$ value of 2.03, $J = -1.22 \text{ cm}^{-1}$, and $J' = +0.47 \text{ cm}^{-1}$. Since the ratio of J'/J is relatively large, the applicability of the linear chain model to this compound is open to some question. The expression for two exchange-coupled $S = 1/2$ ions gives a reasonable fit to the data also, with the EPR $\langle g \rangle$ value of 2.03 and $J = -1.3 \text{ cm}^{-1}$. However, this latter model is probably not applicable to the present case since the high temperature structure is known to be a uniformly spaced linear chain, and since the magnetic susceptibility data appear to be tending to non-zero values as the temperature approaches zero. Also, there is no anomaly in the dielectric constant versus temperature plot over the temperature range 10-300 K, and structural phase transitions are not indicated. Extensive measurements on single crystals in the low temperature range will be required before an exact description of the magnetic interaction may be had. However, for the purposes of this work, it is clear that the sulfur-bridged copper ions are exchanged coupled. Furthermore, the data are

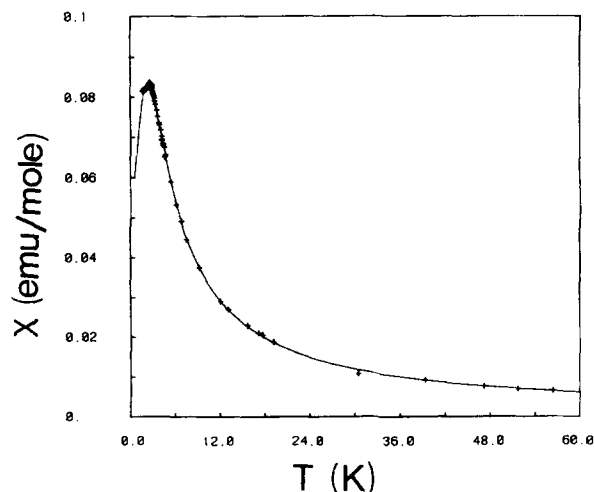


Figure 1. The magnetic susceptibility versus temperature for the linear chain compound $[\text{Cu}(\text{dmtc})_2]_{\infty}$. The solid line is the best fit of the Heisenberg linear chain $S = 1/2$ model to the data with the EPR g value of 2.03 and $J(\text{intra-chain}) = -1.22 \text{ cm}^{-1}$. A correction for an interchain interaction gave $J(\text{interchain}) = 0.47 \text{ cm}^{-1}$.

adequate to permit a correlation between the magnetic properties of this sulfur bridged chain and the geometry of the bridged Cu_2S_2 units in a number of copper complexes. The available structural and magnetic data for the sulfur-bridged copper(II) compounds are compiled in Table I. There are several conclusions which may be drawn from the data. The first and most important point is that even long out-of-plane bridging bond distances provide pathways for appreciable singlet-triplet splittings resulting from superexchange coupling. Next, the short in-plane copper-sulfur bond distances which are involved in the bridge are relatively constant at $2.3 \pm 0.04 \text{ \AA}$, but the out-of-plane bond distances vary from 2.851 \AA in $[\text{Cu}(\text{detc})_2]_2$ to 3.3 \AA in $[\text{Cu}(\text{H}^+\text{TCH})\text{Cl}_2]_2\text{Cl}_2$. In addition, there is a fairly wide range of $\text{Cu-S}_b\text{-Cu}$ angles at the bridging sulfur atom.

Studies on a series of di- μ -hydroxo-bridged copper (II) compounds have established that the singlet-triplet splitting is a linear function of the angle at the bridging oxygen atom. (28) Also, Roundhill, et al. (29) have noted that the singlet states are stabilized as the bridge angle increases in a number of chloro-bridged copper(II) systems. For the purposes of the present study we note that in this particular group of sulfur-bridged compounds there is an additional variable, which also arises in several parallel planar chloro-bridged copper compounds, and that is the out-of-plane bridging bond distances. Since the extent of exchange coupling is known to decrease as the intranuclear separation of exchange coupling sites increases, we were led to investigate the relationship between the exchange coupling constants for these sulfur bridged compounds and the function ϕ/r_o , where ϕ is the $\text{Cu-S}_b\text{-Cu}$ angle and r_o is the out-of-plane copper-sulfur bond distance. This relationship is shown in Figure 2, where it may be seen that a smooth curve may be drawn through the experimental exchange coupling constants J and the function ϕ/r_o .

For the purposes of this work the exchange coupling constant for $[\text{Cu}(\text{detc})_2]_2$ determined earlier in our laboratory was used although there may be some question concerning the magnitude of this number in view of the EPR investigations of Al'tshuler, et al. (22).

Also, the geometry of the more tightly bound unit of $[\text{Cu}(\text{KTS})]_{\infty}$ was used to obtain the value of ϕ/r_o for

Table I. Structural and Magnetic Data in Sulfur-Bridged Copper(II) Compounds

Compound	J, cm ⁻¹	Cu-S _b ^a (Å)	Cu-S-Cu (deg)	φ/r _o
[Cu(dmtc) ₂] _∞	-1.22 ^b	3.159(2)	94.2°	29.8 ^c
[Cu(detc) ₂] ₂	12 ^d	2.851	86.9°	30.5 ^c
[Cu(H ⁺ -TCH)Cl ₂] ₂ Cl ₂	6.5±0.5 ^f	3.310(6)	88.4°	26.7 ^h
[Cu(KTS) ₂] _∞	-12.4 ^g	3.101	89.8°	29.0 ^j
	-8 ⁱ	3.312	86.5°	
[Cu(d-n-btc) ₂] ₂	-14.1 ^k	2.899	92.1°	31.8 ^k

^aThis is the long, out-of-plane bridging bond distance; ^bRef. 18; ^cRef. 19; ^dRef. 20; ^eRef. 21; ^fRef. 22; ^gRef. 23; ^hRef. 24; ⁱRef. 25; ^jRef. 26; ^kRef. 27.

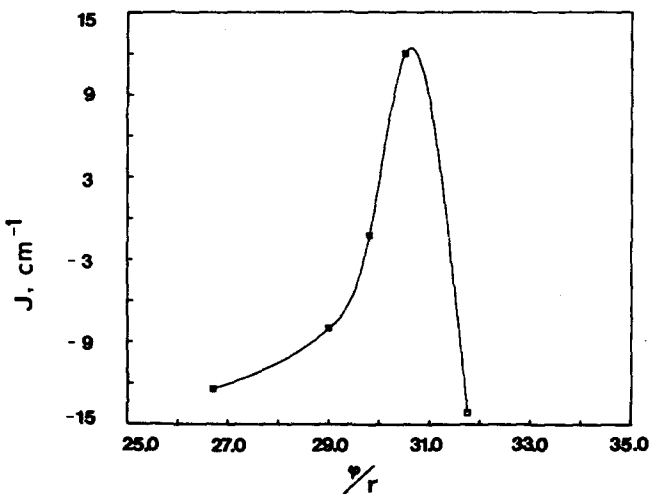


Figure 2. A plot of J versus φ/r_o, where φ is the Cu-S_b-Cu bridging angle in a series of sulfur bridged copper compounds and r_o is the long out-of-plane bond distance.

this alternating chain compound, since this unit may be expected to dominate exchange coupling in the range of Blumberg and Peisach's measurements. (25) It is interesting to point out that at the time that the work on [Cu(dmtc)₂]_∞ was carried out, (18) the data for [Cu(α-d-n-btc)₂]₂ were not available. At that time it was expected that the line in Figure 2 would reach a maximum at some larger value of φ/r_o and then begin to decrease with an increase in φ/r_o. This expectation was anticipated in view of data for chloro-bridged copper(II) complexes. (30)

THE ALTERNATING CHAIN MODEL

The Hamiltonian for the Heisenberg alternating linear chains may be written as

$$H = -2J \sum_{i=1}^{n/2} [\hat{S}_{2i} \cdot \hat{S}_{2i-1} + \alpha \hat{S}_{2i} \cdot \hat{S}_{2i+1}] \quad (4)$$

where J is the exchange integral between a spin and its right neighbor and αJ is the exchange integral between a spin and its left neighbor. The model of most interest is for antiferromagnetic exchange (J less than 0) and for 0 < α < 1. At the extremes, when α = 0, the model reduces to the dimer model with pairwise interactions, and when α = 1, the model reduces to the regular linear chain model.

This model has been studied in detail by Duffy and Barr, (12) as well as Bonner and Friedberg, (3,4) Diederix, et al. (5), and Bonner and coworkers (30a).

For the limiting case of α = 1, Hall (16) has fit Bonner's numerical data (17) to expression (2) given above. For the purposes of this work the reduced magnetic susceptibilities of short alternating rings of up to ten S = 1/2 spins for α = 0, 0.1, 0.2, 0.3, 0.4, 0.6, 0.8, and 1.0 have been calculated using the cluster approach and the programs which have been described earlier. (15,31) The results of these calculations are in excellent agreement with the previously reported results of Duffy and Barr, (12) and the reduced magnetic susceptibilities are displayed as a function of reduced temperature in Figure 3.

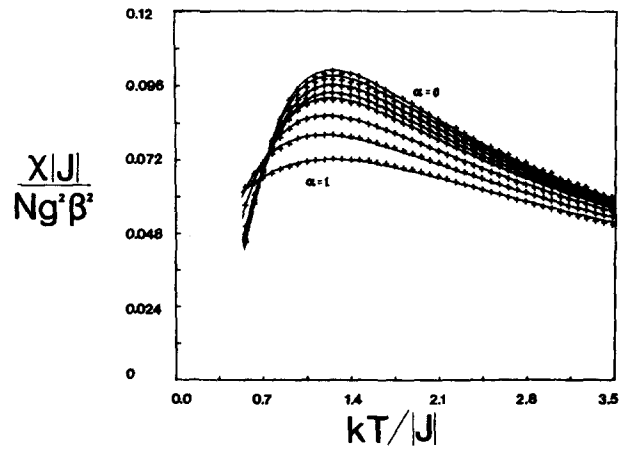


Figure 3. Plot of equation 5 (solid line) in reduced coordinates using coefficients listed in text. Points for 10 membered rings (+) are calculated with H = 1000 Oe, g = 2.2, and α = 0.0, 0.1, 0.2, 0.3, 0.4, 0.6, 0.8, and 1.0.

There reduced magnetic susceptibility curves for 10 member alternating chains with the α values indicated above were fitted to the quadratic-cubic function (5), where the reduced susceptibility χ_r = χ|J|/Ng²μ_B² and the reduced temperature T_r is given by KT/|J|.

$$\chi_r = \frac{AT_r^2 + BT_r + C}{T_r^3 + DT_r^2 + ET_r + F} \quad (5)$$

The fits were truncated below KT/|J| of 0.5, and the constant A in expression 5 was set equal to 0.25 for convergence to the Curie Law at high temperatures.

With the expressions for reduced susceptibility for each of the eight alternating chain cases as a starting point, a unified expression for χ_r as a function of the alternating parameter was developed. The values for the parameters A, B, C, D, E, and F for 0 < α < 0.4 are

$$\begin{aligned} A &= 0.25 \\ B &= -0.12587 + 0.22752\alpha \\ C &= 0.019111 - 0.13307\alpha + 0.509\alpha^2 - 1.3167\alpha^3 + 1.0081\alpha^4 \\ D &= 0.10772 + 1.4192\alpha \\ E &= -0.0028521 - 0.42346\alpha + 2.1953\alpha^2 - 0.82412\alpha^3 \\ F &= 0.37754 - 0.067022\alpha + 5.9805\alpha^2 - 2.1678\alpha^3 + 15.838\alpha^4 \end{aligned}$$

The values for the parameters A-F for 0.4 < α < 1.0 are

$$\begin{aligned} A &= 0.25 \\ B &= -0.13695 + 0.26387\alpha \\ C &= 0.017025 - 0.12668\alpha + 0.49113\alpha^2 - 1.1977\alpha^3 \\ &\quad + 0.87257\alpha^4 \end{aligned}$$

$$D = 0.070509 + 1.3042\alpha$$

$$E = -0.0035767 - 0.40837\alpha + 3.4862\alpha^2 - 0.73888\alpha^3$$

$$F = 0.36184 - 0.065528\alpha + 6.65875\alpha^2 - 20.945\alpha^3 +$$

The expression for the reduced magnetic susceptibility was then converted to magnetic susceptibility in familiar coordinates to give expression 6

$$\chi_m = \frac{Ng^2\mu_B^2}{kT} \cdot \frac{A + Bx + Cx^2}{1 + Dx + Ex^2 + Fx^3} \quad (6)$$

where $x = |J|/kT$. The expression with the two sets of parameters A-F is valid for $kT/|J|$ greater than 0.5 and J less or equal to 0, and as shown in Figure 3, reproduces the calculated magnetic susceptibilities of the 10 membered rings, which should be good approximations for infinite systems for various values of α with an

factor $\sum_n (|\chi_{theo} - \chi_{fit}|^2 / \chi_{theo}^2)^{1/2}$ much better than

0.01. The error in the alternation parameter and the exchange coupling constant which results from the fit is largely that inherent to the assumptions that the magnetic properties of relatively small rings may be extrapolated to the infinite limit. Presently all of our studies with this model has been carried out in the low field region (10 KOe, 1 KOe, 100 Oe); additional studies at higher fields are underway.

MAGNETIC SUSCEPTIBILITY DATA FOR ALTERNATING CHAIN SYSTEMS

The magnetic susceptibility data (11) for catena-hexanedionebis(thiosemicarbazato)copper(II), Cu-HTS, in the temperature range 1.8 - 120 K are shown in Figure 4. The maximum in the magnetic susceptibility curve provides evidence for antiferromagnetic interactions. Attempts to fit the magnetic data by the uniform Heisenberg chain model described above or by the dimer model failed, but an excellent fit was obtained with the alternating chain model. The solid line in Figure 4 was calculated using the Heisenberg alternating chain expression and the magnetic parameters given in Table II. The inability of the uniform chain and dimer models to describe the magnetic behavior is shown by the "best fit" lines which are calculated from these models and are depicted in Figure 4. Data are also provided for the analogous alternating chain catena-octanedionebis(thiosemicarbazato)copper(II), Cu(OTS).

There are two additional compounds for which an alternating chain model is needed to fit the low temperature magnetic susceptibility data, those being Cu(4-methylpyridine)₂Cl₂ (32) and Cu(N-methylimidazole)₂Br₂. (9) However, both of the com-

pounds have uniformly spaced linear chain structures at room temperature. (33-35) Dielectric measurements on compacted samples of Cu(4-methylpyridine)₂Cl₂ and Cu(N-methylimidazole)₂Br₂ provide evidence for structural phase transformations. As shown in Figure 5, a first order phase transition is clearly indicated by the anomaly in the capacitance data near 50K. There are multiple anomalies in the temperature dependent capacitance data for Cu(N-methylimidazole)₂Br₂. A first order transition near 105K may be a result of a "freezing in" of the expected free rotation of the N-methyl group, and a transition near 50 K is probably comparable to the alternating chain transformation in Cu(4-methylpyridine)₂Cl₂. The magnetic parameters are given in Table II.

The alternating Heisenberg chain model and expressions derived from the cluster approach permit a description of the magnetic susceptibility of the compounds catena-octanedione- and catena-hexanedionebis(thiosemicarbazato)copper(II), catena-di- μ -chlorobis(4-methylpyridine)copper(II), and catena-di- μ -bromobis

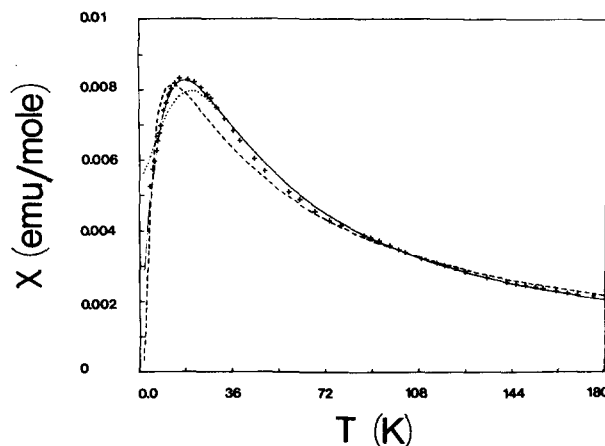


Figure 4. Magnetic susceptibility versus temperature data for catena-hexanedionebis(thiosemicarbazato)copper(II). The solid line is the best fit of the Heisenberg alternating linear chain model with $J = -10.1 \pm 0.2 \text{ cm}^{-1}$, $\alpha = 0.91 \pm 0.02$, and $g = 2.06 \pm 0.02$. For purposes of comparison the "best fits" to the uniform Heisenberg chain (\dots , $J = -10.7 \text{ cm}^{-1}$, $J' = 1.76 \text{ cm}^{-1}$, $g = 2.11$) and dimer ($---$, $2J = -6.4 \text{ cm}^{-1}$, $g = 2.23$, $\theta = -32.5 \text{ K}$) models are also shown.

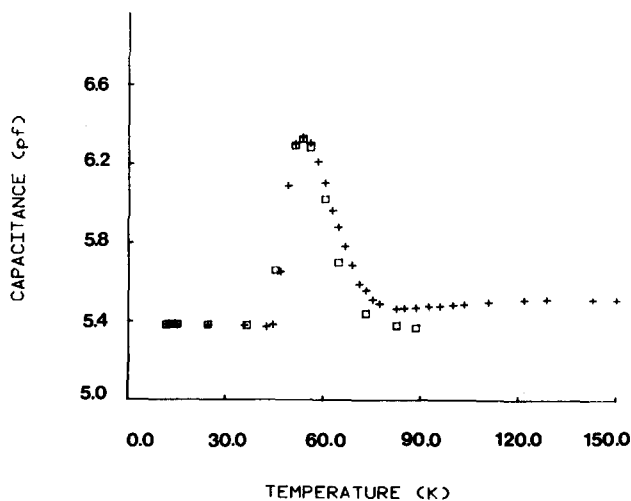


Figure 5. Temperature dependent capacitance data for a pressed pellet sample of catena-di- μ -chlorobis(4-methylpyridine)copper(II) (\square , temperature increasing, + temperature decreasing).

Table II. Magnetic Parameters for the Alternating Chain Compounds

Compound	$J, \text{ cm}^{-1}$	α	g
Cu-HTS	-10.1	0.91	2.06
Cu-OTS	-12.9	0.90	2.11
Cu(4MP) ₂ Cl ₂ ^a	-9.6	0.67	2.17
Cu(NMI) ₂ Br ₂ ^b	-7.2	0.4	2.14

^a4MP = 4-methylpyridine

^bNMI = N-methylimidazole

(N-methylimidazole)copper(II). It is reasonable to suggest that the sulfur bridged compounds have alternating chain structures since the analogous compound Cu-KTS has been shown by X-ray diffraction studies (26) to have an alternating chain structure. X-ray data for twinned crystals of $\text{Cu}(4\text{-methylpyridine})_2\text{Cl}_2$ (33,34) are consistent with a uniform chain structure at room temperature, but an anomaly in the dielectric behavior signals a structural phase transition presumably to an alternating chain as suggested by the low temperature magnetic susceptibility data. Multiple anomalies were seen in the dielectric behavior of $\text{Cu}(\text{N-methylimidazole})_2\text{Br}_2$ and we suggest that a structural phase transition from the uniform chain at room temperature (35) has occurred, thus permitting an explanation of the low temperature magnetic and magnetization data. The neutron diffraction results (36) for $\text{Cu}(\text{N-methylimidazole})_2\text{Br}_2$ at 4 K suggest that this problem merits additional attention.

LADDER SYSTEMS

Ladder systems have also been investigated using the cluster approach. Although it had been thought for a long time that $\text{Cu}(\text{NO}_3)_2 \cdot 2.5\text{H}_2\text{O}$ was an example of a spin ladder, in view of the ladder-like structure at room temperature, (6) recent work has shown that the magnetic data are consistent with an alternating chain. Weller (15) has carried out calculations for short ladder-like chains and rings. For the purposes of illustration, the energy levels resulting from the calculation for a four member chain are plotted in Figure 6. There are several important points to consider. First, the ground state of the system is a quintet when both exchange coupling constants are positive, a singlet when both exchange coupling constants are negative or when $J_{12} < 0$ and $J_{13} > 0$, and there is a ground state spin crossover at $\alpha \approx 0.3$ when $J_{12} > 0$ and $J_{13} < 0$. It is clear that a wide variety of behavior is to be expected from ladder compounds and this expectation provides the stimulation for extensive synthetic and structural studies.

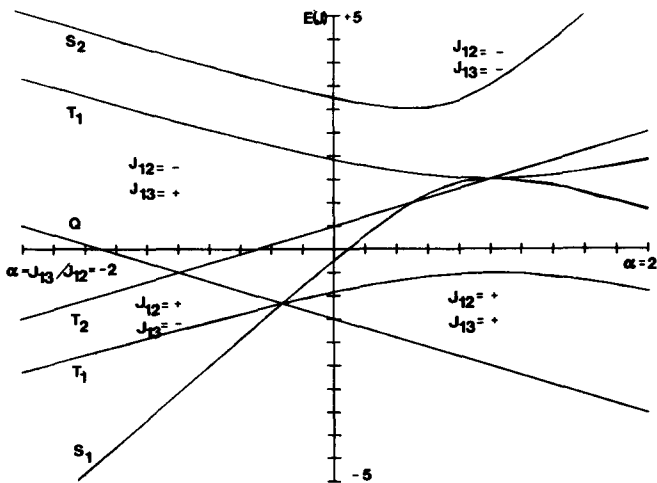


Figure 6. Energy level diagram for a spin ladder with four copper(II) ions.

MAGNETIC DATA FOR THE LADDER CATENA-TRICHLOROHYDRAZINIUMCOPPER(II)

The structure and magnetic properties of $(\text{HzH})\text{CuCl}_3$ have been reported by Brown, et al. (14) The ladder-like structure of the compound is shown in Figure 7. The coordination about the copper ion is the familiar 4+2 form with the equatorial plane being

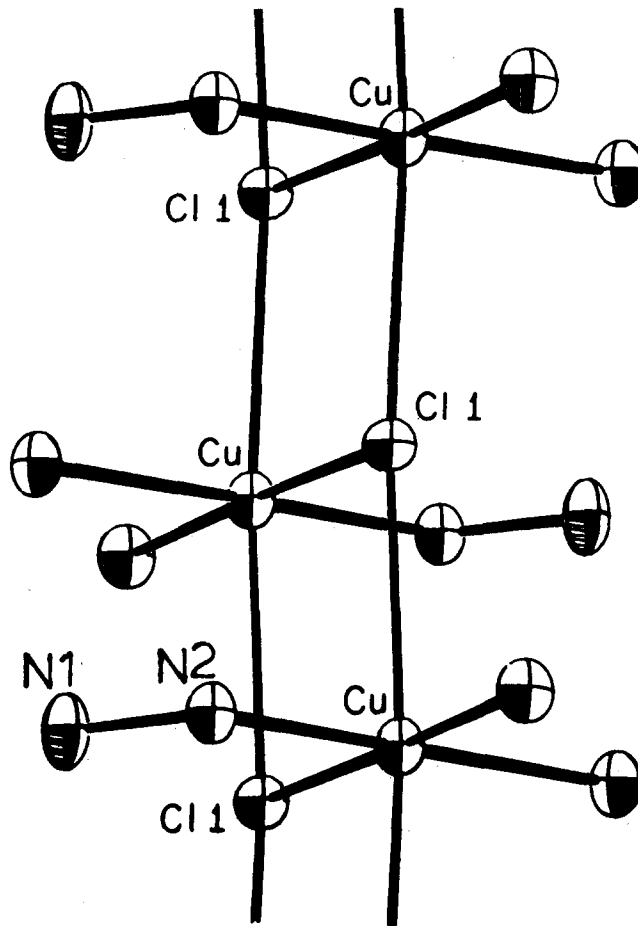


Figure 7. The structure of the ladder compound hydrazinium copper(II) chloride.

formed by three chloride ions and the monodentate hydrazinium ion. The apical positions in the coordination sphere are occupied by chloride ions from adjacent units. The Cu-Cl in-plane bridging distance is 2.297(1) Å and the out-of-plane bridging distance is 2.856 Å.

The magnetic susceptibility data reflect an antiferromagnetic interaction with a maximum in the magnetic susceptibility at about 12 K. The data above 25 K may be fit by the Curie-Weiss law to yield a g value of 2.196, while the entire data set may be fit by an eight membered chain to yield the same g value, J_{12} (the 90° interaction) of -7.5 cm^{-1} , and J_{13} (the 180° interaction) of -7.0 cm^{-1} . The agreement factor $R = [\sum(\chi_{\text{obs}} - \chi_{\text{calc}})^2 / \sum(\chi_{\text{obs}})^2]^{1/2}$ for the fit is 0.018.

A similar calculation was carried out for an eight membered ring to yield $J_{12} = -5.4 \text{ cm}^{-1}$, $J_{13} = -5.3 \text{ cm}^{-1}$, $g = 2.14$, and $R = 0.013$. The calculations firmly support the conclusion that the superexchange interactions are antiferromagnetic by both exchange pathways, but high field magnetization measurements will have to be made before the magnitude of the coupling constants can be known with certainty.

CONCLUSION

The cluster approach to the problem of exchange interactions has been reasonably successful in accounting for the magnetic behavior of extended linear chains, alternately spaced linear chains, and ladders. With use of the cluster approach it has been possible to define exactly the indivi-

dual exchange interactions and obtain detailed analyses of magnetic properties. There is a problem in that small clusters which are not realistic models may yield fits to experimental data which are nearly as good as the fits with larger clusters, but the magnetic parameters are usually unacceptable. As research in this area progresses, an adequate body of data will be generated so that reasoned judgements may be made about the acceptability of magnetic parameters from model calculations.

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