

Columnar and equiaxed solidification of Al-7wt.%Si alloys in reduced gravity in the framework of the CETSOL project

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Abstract

During casting often a dendritic microstructure is formed, resulting in a columnar or equiaxed grain structure, or leading to a transition from columnar to equiaxed growth (CET). The detailed knowledge of the critical parameters for the CET is important, because the microstructure affects materials properties. To provide unique data for testing of fundamental theories of grain and microstructure formation, solidification experiments in microgravity environment were performed within the ESA MAP project CETSOL. Reduced gravity allows for purely diffusive solidification conditions, i.e., suppressing melt flow and sedimentation and floatation effects. On-board the International Space Station (ISS), Al-7wt.%Si alloys with and without grain refiners were solidified in different temperature gradients and with different cooling conditions. Detailed analysis of the microstructure and the grain structure showed purely columnar growth for non-refined alloys. The CET was detected only for refined alloys,

either as a sharp CET in the case of a sudden increase in the solidification velocity, or as a progressive CET in the case of a continuous decrease of the temperature gradient. The present experimental data were used for numerical modelling of the CET with three different approaches: (i) a front tracking model using an equiaxed growth model, (ii) a 3D cellular automaton – finite element (CAFE) model, (iii) a 3D dendrite needle network (DNN) method. Each model allows predicting the columnar dendrite tip undercooling and the growth rate with respect time. Furthermore, the positions of CET and the spatial extent of the CET, being sharp or progressive, are in reasonably good quantitative agreement with experimental measurements.

Introduction

Materials properties of cast components are determined strongly by their microstructures. During casting often a dendritic microstructure is formed. As a result of competition between the growth of several arrays of dendrites, columnar or equiaxed grain structures, or a transition from columnar to equiaxed growth is observed. The precise investigation on grain formation is hampered by buoyancy-driven flow and by movement of crystals growing in the melt. To allow for pure diffusive solidification conditions, i.e., suppressing melt flow and sedimentation effects, experiments in microgravity environment were performed. These experiments provide unique data for testing fundamental theories of grain and microstructure formation.

Such activities are the task of the research project Columnar-to-Equiaxed Transition in SOLidification Processing (CETSOL) in the framework of the Microgravity Application Promotion (MAP) programme of the European Space Agency (ESA). At present, the CETSOL team consists of seven European scientific partners, three partners from the United States of America and seven partners from industry (see list of co-authors).

Microgravity experiments with Al-7wt.%Si alloys

i) Experimental set-up and process parameters

Several experiments in microgravity were carried out successfully to investigate the columnar-to-equiaxed transition under diffusive conditions for heat and mass transport:

First, within Batch1 (B1) six solidification experiments were performed in the Materials Science Laboratory (MSL) with the Low Gradient Furnace (LGF) module on-board the International Space Station (ISS). Rod-like samples with diameter 7.8 mm and length 245 mm made of Al-7wt.%Si alloy, with (gr) and without grain refiner particles, were used and integrated in a special sample cartridge assembly (SCA) together with 12 thermocouples positioned along the sample (Figure 1). The main parts of the LGF are the ‘cold zone’ which consists of 3 heaters, the ‘hot zone’ which consists of 4 heaters, and an adiabatic zone in between. By controlling the temperatures of the heaters a temperature gradient develops along the sample axis. Melting or solidification of the metallic alloy is realized by movement of the furnace insert along the axis of the fixed sample with a defined speed. Table I summarizes the main parameters of the Batch1 flight experiments. At the beginning of the solidification phase the furnace is in its end position and the set heater temperatures allow for an initial temperature gradient along the sample axis of about 0.9 K/mm. The solidification stage I with $v_1 = 10 \mu\text{m/s}$ for $z_1 = 20 \text{ mm}$ was identical for all experiments (see Table I) and is expected to generate a columnar dendritic growth. A transition to equiaxed growth

should be triggered in the solidification stage II, either by increasing the furnace movement to $v_2 = 200 \mu\text{m/s}$ (B1-FM1 to B1-FM3) or mainly decreasing the temperature gradient by applying a cooling rate of $dT/dt = -4\text{K/min}$ at the heaters of the hot zone (B1-FM4 to B1-FM6).

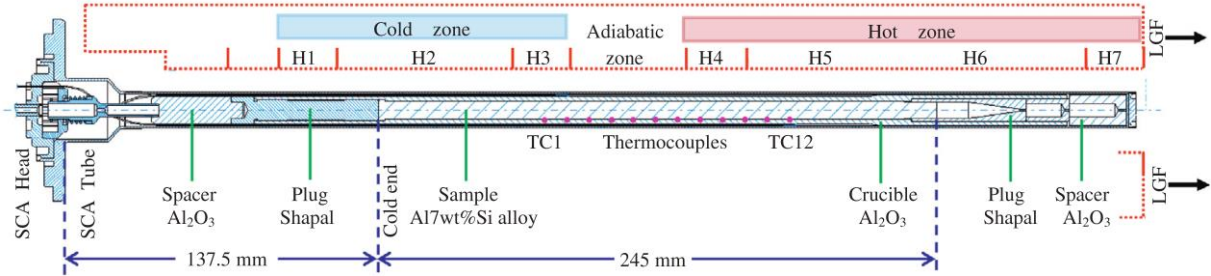


Fig. 1: Schematic setup of the sample cartridge assembly (SCA) inserted in the low gradient furnace (LGF) available in the Materials Science Laboratory (MSL) on board the International Space Station (ISS). The SCA contains the Al-7wt.%Si alloy inserted as a 7.8 mm diameter 245 mm long cylinder. The LGF is sketched at its initial position with respect to the SCA. The positions of the 12 thermocouples TC1 to TC12 regularly distributed from 72.5 mm to 182.5 mm from the cold end of the sample are also provided.

Table I: Process parameters of the CETSOL1 MSL-LGF Batch1 (B1) flight experiments.

CETSOL Batch1 Sample No.	Alloy	Homogenization time t_H (min)	Solidification stage I		Solidification stage II			Fast movement stage III
			v_1 ($\mu\text{m/s}$)	z_1 (mm)	v_2 ($\mu\text{m/s}$)	z_2 (mm)	dT/dt (K/min)	v_3 ($\mu\text{m/s}$)
B1-FM1	AlSi7+gr	10	10	20	200	50	-4	3000
B1-FM2	AlSi7+gr	300	10	20	200	50	-4	3000
B1-FM3	AlSi7	300	10	20	200	50	-4	3000
B1-FM4	AlSi7	300	10	20	10	20	-4	3000
B1-FM5	AlSi7+gr	10	10	20	10	20	-4	3000
B1-FM6	AlSi7+gr	300	10	20	10	20	-4	3000

Second, within Batch2a (B2) seven solidification experiments were performed in the MSL with the Solidification and Quenching Furnace (SQF) module on-board the ISS. As with Batch1, rod-like samples of Al-7wt.%Si alloys with the same dimensions made, with (gr) and without grain refiner particles, were used and integrated in a sample

cartridge assembly (SCA) together with 12 thermocouples positioned along the sample. The positions of the 12 thermocouples TC1 to TC12 are regularly distributed from -7.5 mm to 222.5 mm from the cold end of the sample. The SQF consists of a ‘hot’ zone (4 heaters), an actively cooled ‘cold’ zone realized by a Liquid Metal Ring (LMR) and an adiabatic zone in between. By controlling the temperatures of the heater a higher temperature gradient along the sample axis (than in Batch1) was achieved. Melting or solidification of the metallic alloy was realized by movement of the furnace insert along the axis of the fixed sample with a defined speed. Table II summarizes the main parameters of the Batch2a flight experiments. At beginning of the solidification phase the furnace is in its end position and the heater set temperatures allow for an initial temperature gradient along the sample axis of about 4 K/mm and 3 K/mm (B2-FM4 and B2-FM7). The solidification stage I with $v_1 = 20 \mu\text{m/s}$ for $z_1 = 20 \text{ mm}$ or $z_1 = 30 \text{ mm}$ is expected to generate columnar dendritic growth. A transition to equiaxed growth should be triggered in the solidification stage II, either by increasing the furnace movement to $v_2 = 200 \mu\text{m/s}$ (B2-FM1 to B2-FM4) or decreasing the temperature gradient by applying a cooling rate of $dT/dt = -8\text{K/min}$ at the heaters of the hot zone (B2-FM5 to B2-FM7). For the non-refined samples FM2 and FM6 a moderate rotating magnetic field (RMF) with frequency $f = 57 \text{ Hz}$ and magnetic flux density $B = 0.5 \text{ mT}$ was applied to promote eventually fragmentation.

Table II: Process parameters of the CETSOL MSL-LGF Batch2a (B2) flight experiments.

CETSOL Batch2a Sample No.	Alloy	Homogeni- zation time t_H (min)	Solidification stage I		Solidification stage II				Quench stage III
			v_1 ($\mu\text{m/s}$)	z_1 (mm)	v_2 ($\mu\text{m/s}$)	z_2 (mm)	dT/dt (K/min)	RMF	
B2-FM1	AlSi7+gr	240	20	20	200	50	-8	No	Yes
B2-FM2	AlSi7	240	20	20	200	50	-8	Yes	Yes
B2-FM3	AlSi7+gr	60	20	30	200	50	-8	No	Yes
B2-FM4	AlSi7+gr	60	20	30	200	50	-8	No	Yes
B2-FM5	AlSi7+gr	60	20	30	20	50	-8	No	Yes
B2-FM6	AlSi7	60	20	30	20	50	-8	Yes	Yes
B2-FM7	AlSi7+gr	60	20	30	20	50	-8	No	Yes

ii) Experimental Results

Based on the temperature measurements along the sample axis, cooling curves for all Batch1 and Batch2a samples were determined. Also, the average liquidus isotherm velocities and the temperature gradients ahead of the liquidus isotherm were deduced as function of time. It was found that grain refinement and RMF have almost negligible influence on the heat flow.

From axial and longitudinal cross-sections of the samples the microstructure and grain structures were identified. Quantitative evaluation of the microstructure showed the development of eutectic percentage along the sample axis. From electrolytically etched cross-sections observed in a polarizing microscope the different grain structures were identified. Additionally from quantitative EBSD measurements the size (equivalent diameter), the crystallographic orientation and the elongation factor of each grain was determined. The elongation factor was defined by the ratio of the length of the major axis to the length of the minor axis, the two axes intersecting orthogonally at the centroid of the grain.

Some results for the Batch1 benchmark experiments are already published [1 - 4] are summarized here briefly:

- In non-refined alloys (samples B1-FM3 and B1-FM4; see Table I) no CET was observed, but columnar dendritic growth with some fragmentation of dendrite structures. The fragments have a size comparable to the dendrite sidebranch spacing and cannot promote CET during solidification in microgravity environment because the fragments are not transferred by melt convection to the undercooled melt ahead of the columnar front [3].
- CET was obtained only for refined alloys (see Table I). In case of the sudden increase of the solidification velocity (samples B1-FM1 and B1-FM2) a sharp CET within a few mm distance was observed, based on the evaluation of the grain size. In case of a decrease of the temperature gradient, resulting also in a moderate increase of solidification velocity, the formation of elongated and equiaxed grains is found (samples B1-FM5 and B1-FM6). At the end of this transition region of several tenths of millimetres, fully equiaxed grains are detected. Therefore, a progressive CET mode exists.

The evaluation of the Batch2a samples is still in progress. As a recent result, in the following section a comparative analysis of two flight samples – one from Batch1 (B1-FM1) and one from Batch2a (B2-FM1) – is given. Figure 2 shows the microstructure and the eutectic percentage map of the longitudinal central sections situated in the sample at the junction between stage I and II for samples B1-FM1 and B2-FM2. The successive solidification stages (I, II and III) are marked out on the left side which were defined the position of the liquidus isotherm calculated as in our previous work [3].

For both samples, a columnar structure is obtained during stage I. Some quite large eutectic regions (darker areas in Fig. 2) are found in stage I, which are formed due to the solute rejection during solidification and accumulation in the interdendritic regions. The microstructure significantly changes from stage I to stage II in both samples. First, equiaxed and finer (compared to stage I) dendrite grains are obtained in stage II. Moreover, no large eutectic zones are observed anymore. As a consequence, the eutectic distribution is much more uniform, the eutectic percentage is also lower in stage II, but nearly the same for B1-FM1 and B2-FM1 [7]. In comparison of the two samples, finer structures are obtained in B2-FM1, in every stage, than what can be observed in B1-FM1.

For both samples, the increase of furnace velocity results in a sharp increase of the liquidus isotherm velocity. At lower growth velocity in stage I, competition with neighboring dendrites forms large solute-enriched liquid pockets in the interdendritic region that are later transformed into eutectic phase. But in stage II, the velocity increase leads to a larger undercooling at the front of the solidification. In this case, the Si content is higher according to the phase diagram and the growth rate of the dendrite would be lower. As a result, a more homogeneous eutectic distribution and fine dendritic microstructure are generated. Moreover, when comparing the two samples, we can see that for the same growth velocity, higher temperature gradient leads to a shorter mushy zone during

solidification, which accordingly decreases the growth and ripening time of the dendritic microstructure in this region. Therefore, for other parameters comparable, the increase of the temperature gradient leads to a finer microstructure.

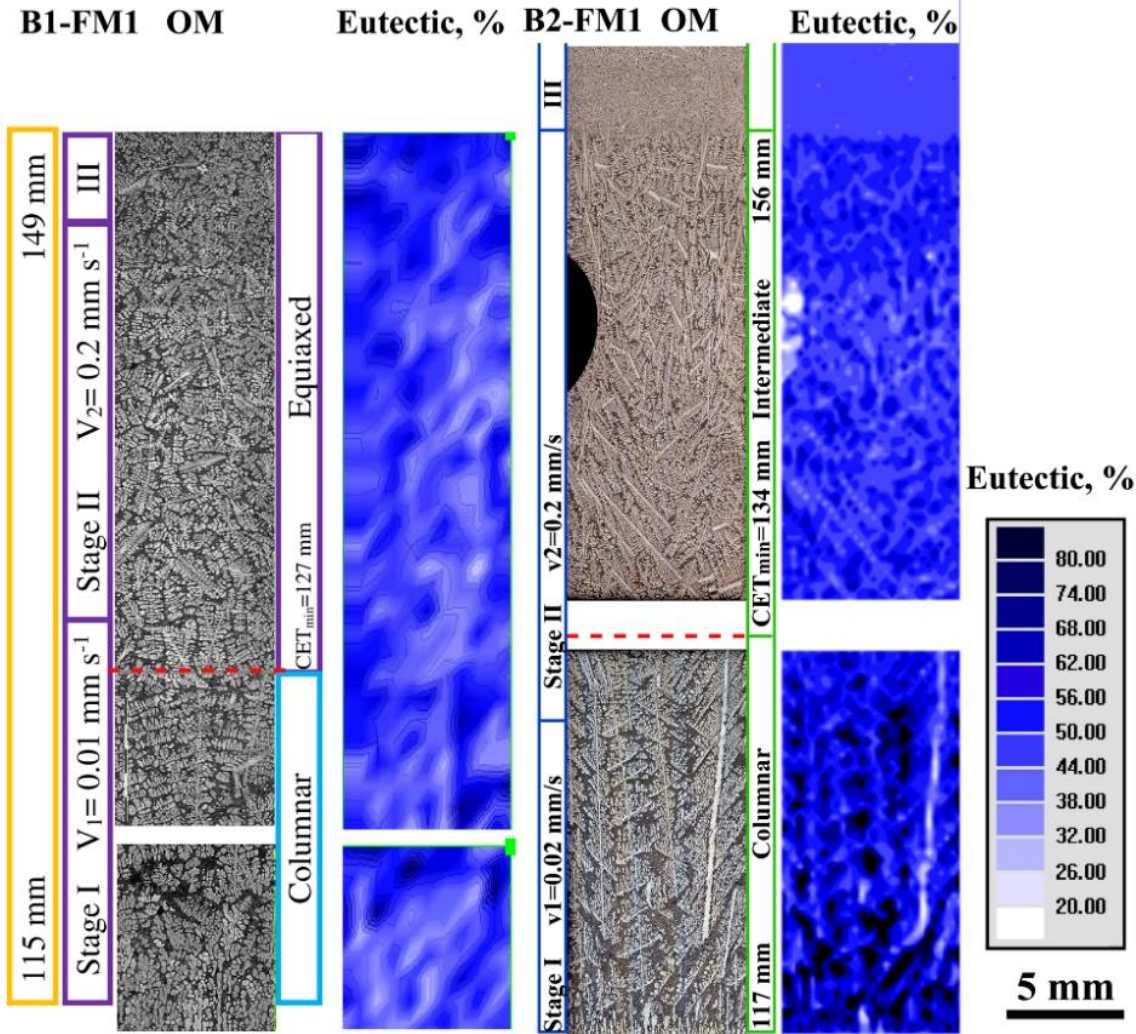


Fig. 2: Longitudinal microstructures determined by optical microscopy OM (left) and eutectic percentage distribution maps (right) for flight samples B1-FM1 and B2-FM1. The blank region represents the material loss during sample preparation.

Considering the grain structure observed by optical microscopy, both samples are columnar during stage I and equiaxed in stage II, which confirms a columnar-to-equiaxed transition in both experiments, triggered by the increase of the solidification velocity. However, the equiaxed grains in B2-FM1 appear to be more elongated than in B1-FM1. The longest length of the columnar zone is defined as the CET inception [5], and noted CET_{min} in Fig. 2. This position is 127 mm for B1-FM1 and 134 mm for B2-FM1.

For quantitative evaluation of the CET, the grain structure is also investigated using EBSD measurements. Figure 3 shows the grain equivalent diameter and elongation factor longitudinal profiles calculated from the EBSD maps. For both samples, the equivalent diameter decreases from stage I to stage II as expected with an increase of growth rate. According to Hunt’s model [6], a critical value of the elongation factor of the grains was deduced to define the occurrence of CET. Grains would be recognized as fully equiaxed below this value. Here, the value “2” is discussed. In the case of B1-FM1, the elongation factor remains constant and lower than 2 except at one isolated position after CET_{min} , indicating an equiaxed zone covering the stage II and a sharp CET. Similarly to B1-FM1, the equivalent diameter of B2-FM1 decreases from stage I to stage II, while the elongation factor mainly stays close above 2 after CET_{min} up to 156 mm. Therefore, the equiaxed grains are more elongated in this region defining a progressive CET. It should be noted that an important deviation of the elongation factor is observed for this sample at about 148 mm, which has to be correlated with the existence of the porosity observed in Fig. 2.

A higher velocity enhances the constitutional undercooling ahead of the solidification front, leading to the CET in both samples. However, due to the higher temperature gradient of 4K/mm in B2-FM1, nucleation takes place in a shorter undercooled region. Therefore, nucleation rate is reduced compared to the experiment with temperature gradient of 0.9K/mm. As a consequence, a mixed zone of some still columnar growing dendrites and new equiaxed dendrites exist, which can be characterized as a progressive CET in B2-FM1, whereas in B1-FM1 a sharp CET is observed [7].

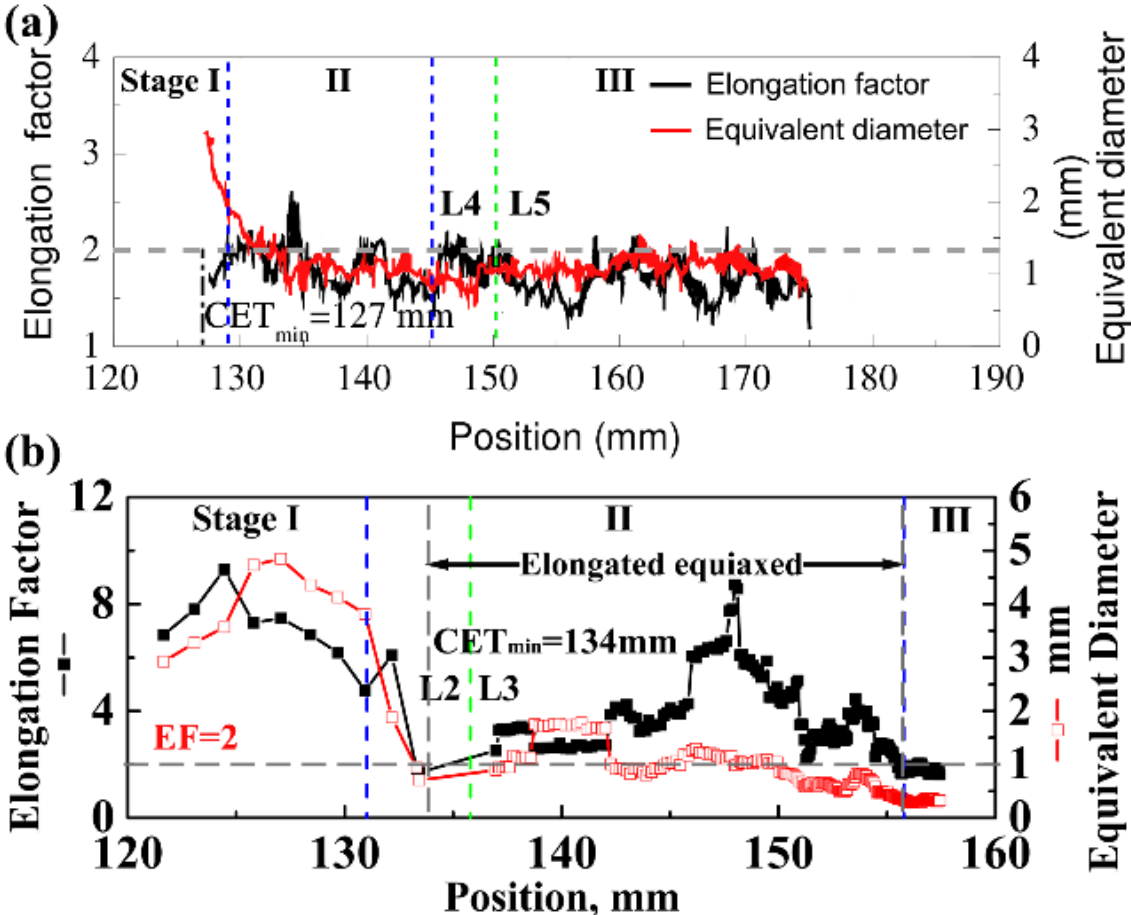


Fig. 3: Longitudinal profiles of the averaged elongation factor and equivalent diameter of the experimental grain contours from EBSD patterns for (a) B1-FM1 and (b) B2-FM1.

Numerical modelling of CET

The flight experiments provide unique data for columnar-to-equiaxed transition in Al-7wt.%Si alloy for pure diffusive solidification conditions, i.e., without melt flow and sedimentation and floatation effects. For this reason, they are used to improve several approaches for numerical modelling of CET, as follows:

- Columnar Front Tracking model with an equiaxed growth model;
- Cellular Automaton – Finite Element (CAFE) method;
- Dendritic Needle Network (DNN) method.

i) CET modelling using front tracking approach

An approach that combines Front Tracking (FT) of columnar growth [8] and an equiaxed volume averaging method [9] was employed to model CET. This model allows to track explicitly the undercooled columnar front and to model equiaxed nucleation from grain refiners. Results of numerical modelling for some of the CETSOL Batch1 experiments are published [1, 4, 10]. For sample B1-FM2 with grain refiner the simulated prediction of the CET position is in very close agreement with that measured experimentally. Also, for B1-FM2 and B1-FM3 samples, the progress of the columnar front was modelled. It was found that the thermal gradient and formation of undercooled zone ahead of the columnar front were similar for both experiments. Therefore, the effect of grain refiners was crucial for the CET in B1-FM2.

Meanwhile, the 1D Bridgman furnace front tracking model (BFFTM) [11, 12] has been significantly modified to simulate columnar and equiaxed growth in the Batch1 flight experiments. The model is suitable for low Biot number ($Bi < 0.1$) applications. We estimate the maximum Biot number for the experiment to be 0.065. Axial and radial heat flow in the crucible has been incorporated using a 2D axisymmetric model. Also, an equiaxed growth model based on the well-established JMAK model [13-15] has been developed and implemented. The equiaxed growth model is currently being redeveloped to improve computational efficiency. The overriding objective of the thermal modelling carried out has been to provide enlightenment regarding experiment design, conditions pertinent to CET, and cross validation in respect of previous modelling efforts.

Fig. 4 illustrates the hybrid of the 2D (annular-shaped control volumes, in black, containing the crucible, $\Delta r = 0.7125$ mm) and the 1D (disc-shaped control volumes, in blue, containing the sample, $\Delta x = 0.5$ mm) axisymmetric model domains. The models are independent but share: (i) a common boundary, (ii) grid resolution in the axial direction, and (iii) their numerical time step, $\Delta t = 0.001$ s.

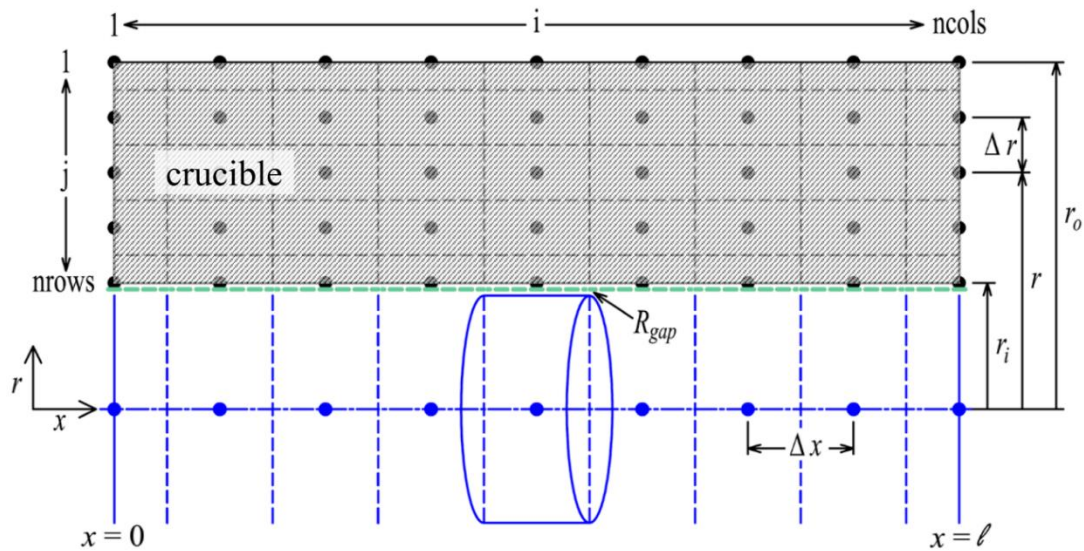


Fig. 4: Axisymmetric model domain of the sample and crucible

The equiaxed growth model applies Johnson-Mehl-Avrami-Kolmogorov (JMAK) theory to independently simulate growth of spherical envelopes of equiaxed mush ahead of a columnar front. The thermal history from twelve thermocouples, located inside grooves on the crucible, is used to drive dynamic boundary conditions in the numerical model. Other key inputs to the model are nuclei density ($N_0=1.6 \times 10^{10} \text{ m}^{-3}$), and mean ($4 \text{ }^\circ\text{C}$) and standard deviation ($0.5 \text{ }^\circ\text{C}$) of nucleation undercooling (see also [20]). It is worth noting that these nucleation parameters are central to the behavior of the model. While a typical maximal value for nucleation undercooling can be determined by examining the magnitude of undercooling ahead of the columnar front (in numerical simulations), the standard deviation of a distribution of nucleation events and the maximum nucleation density are more difficult to estimate. These parameters are selected so that the resulting CET matches that observed in the experiment samples. The current model is computationally inefficient at large nuclei densities ($>2 \times 10^{10} \text{ nuclei/m}^3$), however, a new algorithm is currently in development that solves this problem.

Figure 5 illustrates the mechanism for equiaxed growth ahead of a progressing columnar front within the BFFTM domain. Useful outputs from the model include columnar dendrite tip undercooling, growth rate, temperature gradient with respect to time, and sharp or mixed CET prediction capabilities. Good agreement (within 1 mm) is found between the simulated and experimentally measured CET position in B1-FM2.

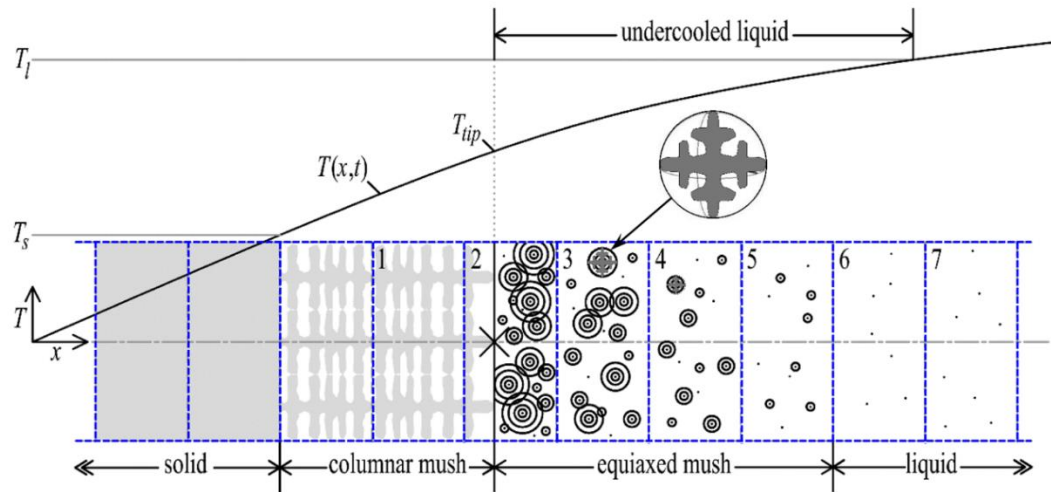


Fig. 5: Spherical equiaxed envelopes growing in undercooled liquid and positive temperature gradient ahead of a columnar front (\times) in the BFFTM domain.

ii) CET modelling using CAFE method

The Cellular Automaton – Finite Element (CAFE) method was proposed to model complex phenomena occurring during solidification processes through the integration of microscopic laws. Resolutions developed at microscopic and macroscopic scales are modelled [16-19]. The conservation equations for energy, solute mass and momentum are solved on a tetragonal FE mesh, which is defined on the domain of interest. A grid of cubic cells is superimposed on to the FE mesh and used by the CA method. Each cell is precisely located in one of the FE element and the macroscopic fields may be interpolated at the cell scale if required.

Results of 2D CAFE modelling of the grain structure evolution and the occurrence of the CET for some of the CETSOL Batch1 experiments are already published [2, 20, 21]. Qualitative agreement was obtained for the position of CET and the CET transition mode, which is sharp in B1-FM1 and progressive in B1-FM5 sample. Further comparisons of the distributions of grain elongation factor and equivalent diameter were conducted and revealed also a fair quantitative agreement.

Now, the CAFE model was extended to the simulation of the B1-FM1 experiment in 3D. Simulation is restricted to the material without considering the whole equipment and its effect on temperature evolution. The sample is virtually divided and simulation of the solidification process is limited to the single domain of interest corresponding to positions from thermocouple TC5 to TC12. Initially, the whole simulation domain is made of a homogeneous liquid of nominal concentration. A Dirichlet boundary condition is used to impose the temperature at the position of the thermocouples. At the macroscopic scale, the energy and solute mass conservations are solved assuming constant densities and fixed solids. Simulation is based on the FE method with about 42,000 isotropic and homogeneous elements and about 9,000 nodes using linear interpolation. At the microscopic scale, the Wang-Beckerman model of microsegregation is used with primary arm spacing $\lambda_1 = 1,000 \mu\text{m}$ [20] and secondary arm spacing $\lambda_2 = 78 \mu\text{m}$ [20]. These values are the mean experimental values determined from sample B1-FM1 [3] and sample B1-FM2 [1], having the same process parameters. Nucleation occurs at the bottom of the simulated part of the sample (TC5). Growth kinetics is approximated with the Ivantsov' solution and the marginal stability criterion using $\sigma = (4\pi^2)^{-1}$. The CA simulation is developed with 3 million cells with characteristic size $80 \mu\text{m}$. The time step is taken equal to 1 s. The parallelized computation took approximately 90 minutes on 40 CPUs. Results of the simulation can be seen on Figure 6 for various solidification times.

Due to the local temperature history and the nucleation parameters, only a columnar grain structure can first develop, the dendrite growth undercooling being too small to favor nucleation of equiaxed grains. After 2,000 s (see Table I), velocity is increased and so is the columnar front undercooling. Liquid cells exist ahead of the columnar front containing a nucleation site sufficiently undercooled to give birth to equiaxed grains. As can be seen by comparison of the simulation and experimental grain structure, qualitative comparison is reached. Moreover, the CET position is accurately and precisely predicted. The CET is measured at 133.2 mm using the same criterion as in reference 3 with 3D data inducing a larger error (Figure 6). The relative error with the experiments is less than 5%. Previous exhaustive explanations [20, 21] of the CET appearance remain valid. The velocity jump increases the constitutional undercooling leading to heterogeneous nucleation of equiaxed grains blocking the columnar front.

To further improve the present computations, the methodology developed within the DNN model (see next chapter) for computation of the dendrite tip growth kinetics is presently under implementation in the CAFE model.

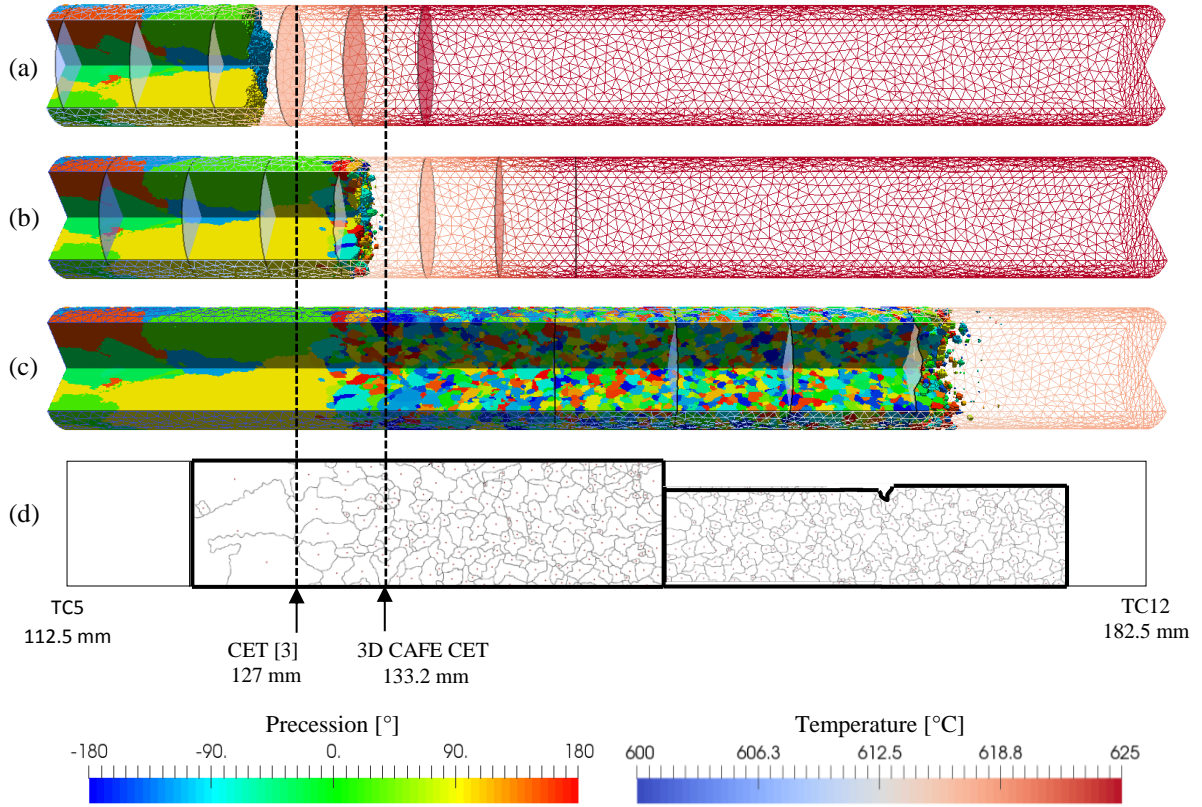


Fig. 6: CAFE simulation of the solidification grain structure of the B1-FM1 sample between thermocouples TC5 (112.5 mm) and TC12 (182.5 mm). Isothermal surfaces from 602 °C to 626 °C with a step of 4 °C are shown. A quarter of the cylinder was removed by post-processing in order to offer an exploded view at various times of (a) columnar growth at $t = 1,800$ s, (b) nucleation and growth of equiaxed grains triggered by a sudden increase of the withdrawal rate at $t = 2,140$ s and (c) the equiaxed structure forming at $t = 2,400$ s. The position of the CET favourably compare with the (d) experimental grain structure observed in a longitudinal metallographic cross section [3].

iii) CET modelling using DNN method

To model the CET in microgravity experiments, we use a three-dimensional (3D) multi-scale Dendritic Needle Network (DNN) method that tracks the diffusion-controlled dynamics of branches of the hierarchically structured dendritic network to model the solidification of dendritic alloys on the grain scale. This approach models explicitly the growth kinetics of all branches of the hierarchical dendritic network (including primary, secondary, and tertiary branches) interacting via the solutal field under both non-steady-state and steady-state growth conditions. The DNN model describes the dendritic tip as a paraboloid [22, 23] and the dynamic of each dendrite is prescribed by two conditions that determine the evolution of its tip velocity V and radius ρ [24].

We model the two CETSOL experiments B1-FM1 and B1-FM5 (see Table I). To investigate the CET, the DNN model is discretized in space on a finite difference cubic grid and the code is implemented on graphics processing

units (GPUs) using the Compute Unified Device Architecture (CUDA). To keep the simulations computationally tractable, we study the CET in a 3D rectangular slab geometry of dimensions $L_x \times L_y \times L_z$ with $L_y = 7$ mm (~experimental sample diameter), L_x chosen large enough to include the CET, and the slab thickness $L_z = 1$ mm chosen much smaller than L_y or L_z but large enough to accurately model 3D dendrite growth kinetics.

The tip selection parameter $\sigma = 0.058$ is used in this study [23, 24]. We model the growth of randomly oriented grains ahead of the columnar front by incorporating nucleation on foreign inoculant particles randomly located in the melt. We choose a nucleation undercooling $\Delta T_n = 4$ K [2] and an effective density $1.25/\text{mm}^3$ [1]; the values are determined from the experiments. We use the measured temperatures at the 12 thermocouples along the sample as input data into the model. The initial condition consists of an array of columnar dendrites [23] with an undercooling ΔT_c^{ini} that was varied in the range of 2 K to 8 K; the results are independent of the choice of initial undercooling within this range.

The results of the DNN simulations for samples B1-FM1 and B1-FM5 are shown in Fig. 7. Figs. 7a-b show the grain structures for the FM1 (a) and FM5 (b) samples. Figs. 7c-d show the corresponding temperature-time curves collected from the 12 thermocouples (thin black dashed lines) and the temperature of the most advanced columnar front T_c (solid red lines) in FM1 (c) and FM5 (d) samples. We measured the starting CET position (CET_{min}) and the ending CET position (CET_{max}) in simulations, where CET_{min} is defined as the position along the sample where the first nucleant particle is initiated, and CET_{max} is defined as the final position of the tip of the most advanced surviving columnar dendrite. In the CETSOL experiments [1-3], a sharp CET is observed in FM1 ($\text{CET}_{\text{min}} = 125$ mm and $\text{CET}_{\text{max}} = 129$ mm) and a more progressive CET is observed in FM5 ($\text{CET}_{\text{min}} = 130$ mm and $\text{CET}_{\text{max}} = 156$ mm). In reasonably good quantitative agreement with those measurements, the simulated grain structures predict that the CET is relatively sharp in FM1 (Fig. 7a with $\text{CET}_{\text{min}} = 129$ mm and $\text{CET}_{\text{max}} = 133$ mm) and more progressive in FM5 (Fig. 7b with $\text{CET}_{\text{min}} = 146$ mm and $\text{CET}_{\text{max}} = 160$ mm). In the insets of Fig. 7c and Fig. 7d, we show zoomed in views of the front undercooling ΔT_c versus time profiles (solid red lines) compared to the nucleation undercooling $\Delta T_n = 4$ K (horizontal solid black lines) in FM1 and FM5 samples, respectively. Only when $\Delta T_c > \Delta T_n$ ($t > 2150$ s in FM1 and $t > 2900$ s in FM5), nucleation is able to take place ahead of the columnar front, thereby allowing equiaxed grains to develop and block columnar growth.

Figure 8 further shows the time evolution of the velocity of the most advanced columnar front in simulations. The growth velocity increases initially during a transient growth regime occurring early in stage I ($t < 1000$ s), which is associated with a coarsening of the primary spacing. After this transient growth regime, the columnar front reaches a steady-state growth regime with a constant velocity for the rest of stage I. Upon faster cooling during stage II, the front velocity increases again leading to a transient growth regime with a non-uniformity of the primary spacing before the onset of the CET as clearly seen in Fig. 7a. Two types of CET are observed in simulations as mentioned above (a sharp CET in FM1 and a progressive CET in FM5). We attribute this difference to the different cooling processes applied at stage II ($t = 2000$ - 2250 s in FM1 and $t = 2000$ - 4000 s in FM5). For all times after the start of stage II, the front undercooling ΔT_c is always larger in FM1 than FM5 because of the faster movement of the furnace during this stage, which boosts the nucleation ahead of the columnar front and results in a sharper CET.

Summarizing, DNN modeling predictions for the positions and spatial extent of the CET are in reasonably good quantitative agreement with experimental measurements in two Batch1 samples (B1-FM1 and B1-FM5) that exhibit sharp or progressive CET depending on the thermal conditions. As a next step in this project, DNN

simulations will be extended to cylindrical bulk samples. This will allow us to compare simulated 3D grain structures to measurements of grain size and orientation using serial sectioning [26]. Furthermore, the DNN method will be used to model the growth of equiaxed dendrites solidified in microgravity in the transparent model alloy Neopentylglycol-(d)Camphor during the sounding rocket mission MASER-13.

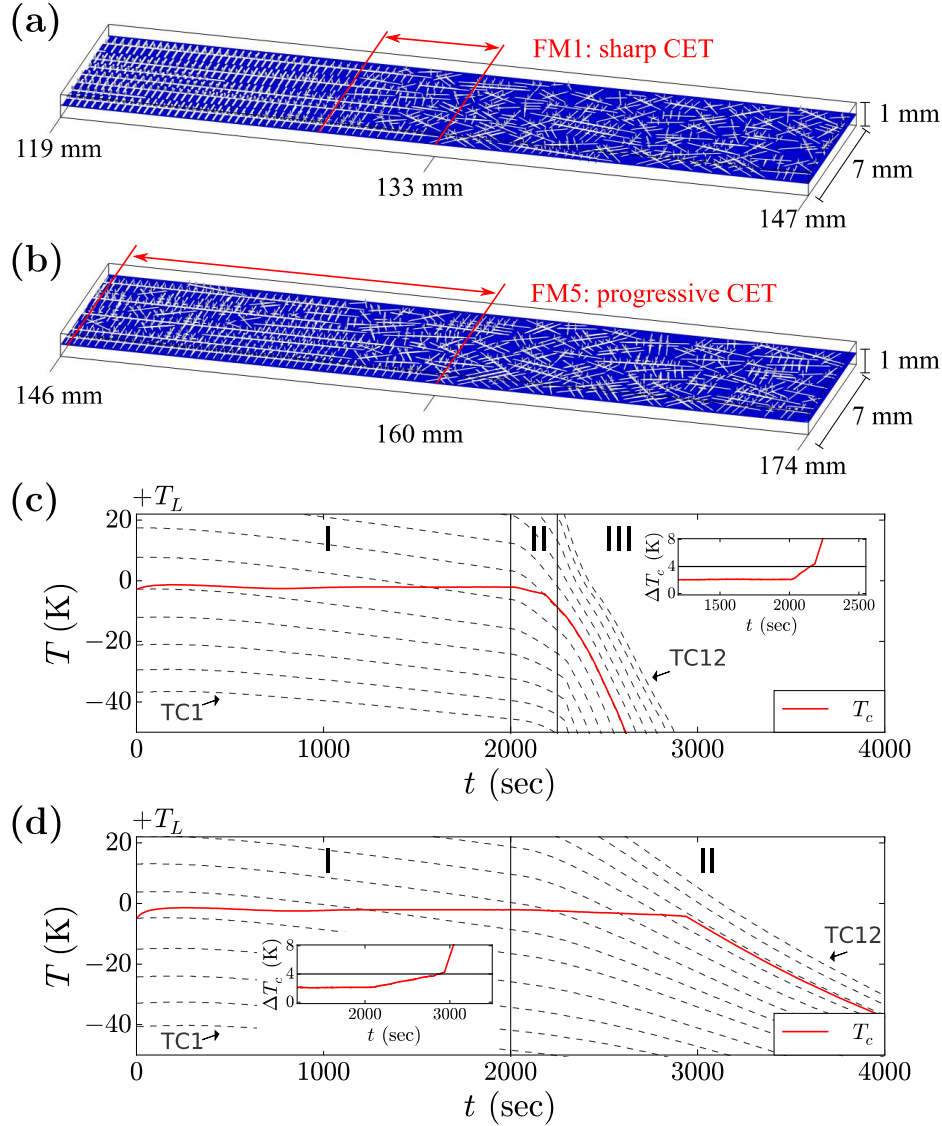


Fig. 7: DNN simulations of dendritic network and grain structures for Al-7wt.%Si samples for B1-FM1 (a) and B1-FM5 (b) samples (see Table I). A sharp CET is observed in FM1 ($CET_{\min} = 125$ mm and $CET_{\max} = 129$ mm) and a more progressive CET is observed in FM5 ($CET_{\min} = 130$ mm and $CET_{\max} = 156$ mm) [2]. (c) and (d) show the temperature curves recorded from 12 thermocouples (thin black dashed lines labelled TC1 to TC12) and the temperature of the most advanced columnar front T_c (solid red lines) in FM1 (c) and FM5 (d). The insets of (c) and (d) show zoomed in views of the time evolution of the columnar front undercooling ΔT_c . The solid line represents the nucleation undercooling $\Delta T_n = 4$ K, which is used in simulations.

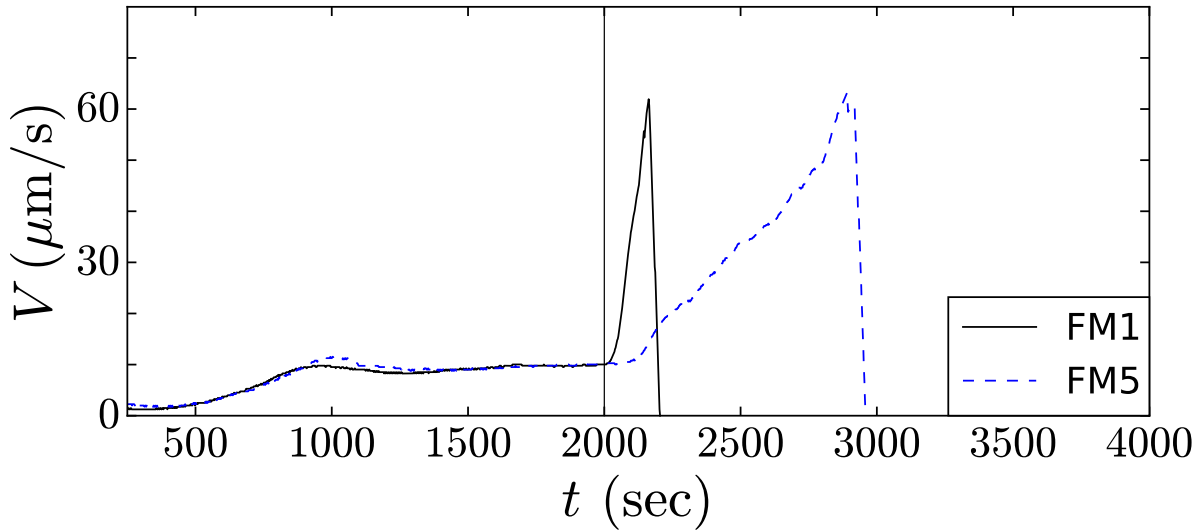


Fig. 8: Time evolution of the velocity of the most advanced columnar front for Batch1 FM1 and FM5 in DNN simulations. The front velocity increases during the growth competition of columnar dendrites leading to an increase of primary spacing ($t < 1000$ s). After this transient regime, the columnar front grows at a constant velocity. The front velocity increases again upon faster cooling during stage II and drops abruptly when columnar growth is blocked by the nucleated equiaxed grains.

Conclusion

This paper summarizes results obtained so far from solidification experiments on-board the International Space Station ISS by the CETSOL team. In the Materials Science Laboratory Al-7wt.%Si alloys with and without grain refiners were solidified in diffusive conditions. Detailed analysis of the microstructure and the grain structure showed columnar growth with some fragments in case of non-refined alloy. CET was detected only for refined alloys. In case of the sudden increase of the solidification velocity, a sharp CET within a few mm distance was observed. A progressive CET mode exists in case of a continuous decrease of the temperature gradient. This mode is characterized by a longer transition region with elongated equiaxed grains followed by fully equiaxed growth.

The unique experimental data basis was used for numerical modelling of CET with three different methods:

- (i) A front tracking model using an equiaxed growth model which applies JMAK theory, predicts the columnar dendrite tip undercooling, growth rate, and temperature gradient with respect time, and sharp or progressive CET. Good agreement is found between the simulated and experimentally measured CET position (B1-FM2).
- (ii) Using a CAFE model both segregation and grain structures, as well as CET, were numerically modelled in 3D. The CET transition mode, be it sharp or progressive, is retrieved. Distributions of grain elongation factor and equivalent diameter are fairly reproduced. Moreover, the CET positions are predicted accurately and precisely.
- (iii) To predict CET a 3D multi-scale Dendrite Needle Network (DNN) method was used. The model was applied to two Batch1 flight samples with sharp and progressive CET (B1-FM1 and B1-FM5), using a simplified slab-like

geometry. The DNN modeling predictions for the positions and spatial extent of the CET are in reasonably good quantitative agreement with experimental measurements in both cases.

Outlook

The experimental data basis for description of CET under diffusive conditions will be extended in the coming years. Therefore, within the CETSOL project solidification experiments with three different hypo-eutectic AlCu alloys, with and without grain refiner, will be performed in reduced gravity on the ISS. Applying different cooling conditions may result in different types of CET. Selected parts of these samples and of the already processed Al-7wt%Si samples will additionally be used for a 3D-reconstruction of the grain structure by serial sectioning [25, 26]. Such 3D features of the structures will be investigated to improve the knowledge of the grain competition, microstructure formation, and CET. In addition, these results will be coupled with further numerical simulations, especially in 3D.

To investigate the dynamics of columnar and equiaxed growth as well as of the CET, further use will be made of the transparent organic alloy Neopentylglycol-(d)Camphor system that solidifies like a metal. As part of the CETSOL project, the sounding rocket experiment MEDI (Multiple Equiaxed Dendrite Interaction) was performed on the MASER-13 sounding rocket flight to observe in-situ free dendrite growth and solutal interaction between close crystals in this alloy under microgravity conditions without sedimentation of the equiaxed crystals and melt flow [27]. This experiment has already provided valuable benchmark data for modelling equiaxed growth at different length scales. This transparent alloy system will be used for future experiments in the Materials Science Glovebox on-board the ISS. Therefore, on behalf of ESA, the TRANSPARENT facility is under development [28]. In experiments with different (d)Camphor compositions, especially the time-dependent evolution of CET for different cooling conditions will be investigated using optical CCD cameras. Again, the results will be used to improve modelling of dendrite growth and interaction.

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