

Review

# Atomistic Descriptions of Gas-Surface Interactions on Tin Dioxide

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**Abstract:** Historically, in gas sensing literature, the focus on “mechanisms” has been on oxygen species chemisorbed (ionisorbed) from the ambient atmosphere, but what these species actually represent and the location of the adsorption site on the surface of the solid are typically not well described. Recent advances in computational modelling and experimental surface science provide insights on the likely mechanism by which oxygen and other species interact with the surface of SnO<sub>2</sub>, providing insight into future directions for materials design and optimisation. This article reviews the proposed models of adsorption and reaction of oxygen on SnO<sub>2</sub>, including a summary of conventional evidence for oxygen ionosorption and recent operando spectroscopy studies of the atomistic interactions on the surface. The analysis is extended to include common target and interfering reducing gases, such as CO and H<sub>2</sub>, cross-interactions with H<sub>2</sub>O vapour, and NO<sub>2</sub> as an example of an oxidising gas. We emphasise the importance of the surface oxygen vacancies as both the preferred adsorption site of many gases and in the self-doping mechanism of SnO<sub>2</sub>.

**Keywords:** SnO<sub>2</sub>; sensor; mechanism; operando; spectroscopy; adsorbates; vacancies



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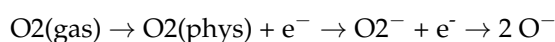
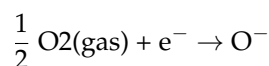
## 1. Introduction

Conductometric gas sensors (CGS) are small, easily fabricated devices that require relatively low power consumption and contain no moving parts. In their simplest form, they consist only of two electrodes interconnected by a layer of gas-sensitive material and can be printed as films directly onto a substrate [1,2]. These characteristics make CGS perfect for mass-scale production and portable device applications. Large scale deployment of such devices could, for example, screen for diabetes or detect imminent electric transformer failure by detecting gases characteristic to the processes [3,4]. However, despite their many advantages, CGS often suffer from cross-sensitivity to interfering gases and performance deterioration over time [5,6], and the main driving force in CGS research targets improvements in sensors' sensitivity, selectivity and response time. Insufficient understanding of the mechanisms of surface reactions makes improvements possible only through empirical exploration.

Spectroscopic investigation of sensors under working conditions is often unattainable due to the UHV requirement of many analytical techniques, for example, X-ray photoelectron spectroscopy (XPS), ion scattering spectroscopy (ISS) or low-energy electron diffraction (LEED). In the catalytic community, this incompatibility is known as the “pressure gap”; this is why early sensor operation models were developed by applying prevailing chemistry principles to explain phenomenological observations. The most prevalent model is the oxygen ionosorption mechanism, which explains the sensor's increased resistance on exposure to oxygen through dissociative adsorption of ambient dioxygen to form ions electrostatically stabilised by the surface [7–13].

It was conjectured that because of its electron affinity, oxygen would withdraw electrons from the conduction band of an n-type semiconductor, and “consequently, there will be no chemisorbed oxygen atoms, but oxygen ions, in the surface” [14]. In this description,

the localisation of negative charge at the surface of SnO<sub>2</sub> grains results in a charge depletion zone below the surface, which constitutes a potential barrier for electrons trying to cross the grain boundary, effectively increasing the sensor's resistance. Usually, the process is only broadly described using the following reaction equation [7,8], which sometimes is extended to include intermediate O<sub>2</sub> physisorbed and O<sub>2</sub><sup>-</sup> chemisorbed species [10,15].



The lack of reference to adsorption sites is not the only problem with ionosorption models, however, which also do not consider whether the oxygen dissociation is symmetric, and if asymmetric, the fate of the other oxygen atom. Moreover, the models often assume that there is no mass transport between the bulk of the sensor and the gas phase [7,10,12], which is not a valid assumption to make, as explained further below. Finally, there is also the question of the existence of O<sup>-</sup> ions adsorbed at the surface, for which there is no conclusive evidence after decades of research [16].

In the absence of adequate analytical techniques, early atomistic insight into the surface processes of gas sensing was provided by computational studies. The chemical state of the surface can be analysed either kinetically, modelling the temporal evolution of a system defined by a set of chemical reactions by optimising an array of parameters, or thermodynamically, estimating the most stable adsorption configurations from first-principles calculations on a model surface, such as a cluster or a slab. The former approach may be used to, for example, estimate the equilibrium density of oxygen adsorbates on O<sub>2</sub> or predict sensor behaviour in varying oxygen pressure [7,15]. The latter is essential in discovering possible adsorbate species and investigating their effects on the surface [17,18]. However, analysis of sensors working under conditions close to ambient has become increasingly more available with the advent of operando spectroscopy. It offers microscopic insight into processes formerly studied only post-mortem in UHV after exposure to a target gas in a preparation chamber, thus presenting new possibilities. For example, changes in the electronic structure of a working sensor may be studied under dynamic pressure and temperature conditions through Near Ambient Pressure (NAP) XPS with simultaneous resistance measurements [19].

In this work, the authors attempt to summarise the proposed mechanisms for the interactions of SnO<sub>2</sub> with oxygen and reducing and oxidising gases as understood from phenomenological and computational studies and highlight those supported by spectroscopic evidence. Since providing an atomistic description of such interactions is the aim of this review, a sensible place to start is by considering the atoms of the surface itself, which is the subject of the next section.

## 2. The Dynamic Surfaces of Tin Dioxide

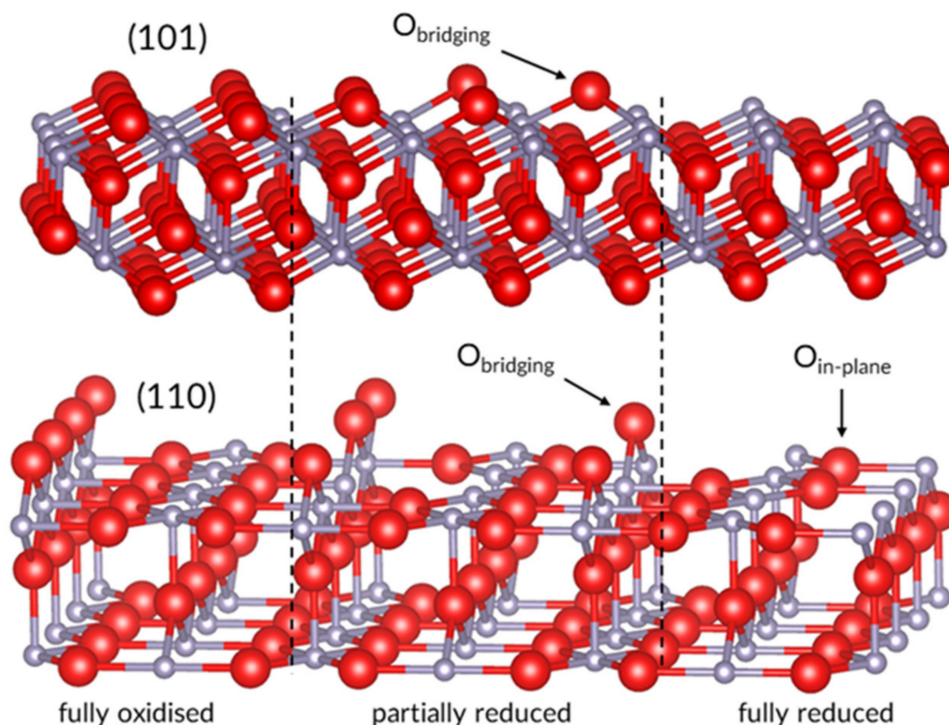
Tin dioxide, SnO<sub>2</sub>, is among the most common sensitive materials used in CGS. It is abundant, inexpensive, non-toxic and shows excellent stability to reducing conditions. Since Taguchi introduced the first commercial SnO<sub>2</sub>-based gas sensors in 1972, and this oxide has been the most widely studied gas-sensitive material, with an initial focus on detecting reducing gases such as CO and volatile organic compounds (VOC). However, recently the focus is shifting towards low-temperature NO<sub>2</sub> detection.

The crystalline form of SnO<sub>2</sub>, known as cassiterite, is a rutile-type tetragonal structure with lattice parameters  $a = b = 4.737 \text{ \AA}$  and  $c = 3.186 \text{ \AA}$  [20]. It is an intrinsic n-type semiconductor with a wide, direct bandgap of 3.6 eV [21], which also makes it transparent and of high interest in touchscreen electronics. Although the exact origin of the intrinsic conductivity is still debated, oxygen vacancies (V<sub>O</sub>) seem to play an important role [22–25]; therefore, a change in the density of V<sub>O</sub> will affect the sensor's resistance.

The surface is the interface of oxygen exchange between the bulk and the ambient and a terminal for interaction with other gases and thus requires due consideration.

Large SnO<sub>2</sub> single crystals, which are notoriously challenging to grow, preferentially expose the (110), (101) and (100) surfaces [26]. Of these three, the (110) was shown to have the lowest energy and is the subject of most computational adsorption studies [17,18,27–31], though some results are also available for the (101) surface [32,33]. Both surfaces feature bridging oxygen sites (O<sub>br</sub>), which protrude outwards from the surface determined by the plane of Sn atoms. The (110) surface has another type of exposed oxygen atom, the in-plane oxygen (O<sub>pl</sub>), which does not have a counterpart on the (101) surface, though the second layer (101) oxygen atoms (first subsurface oxygen) are sometimes referred to in that way.

O<sub>br</sub> atoms can be removed to form vacancies in the process of surface reduction, facilitated by the dual valency of Sn<sup>2+/4+</sup>. The removal results in either a partially or fully reduced surface, the latter indicating one with all O<sub>br</sub> removed, as presented in Figure 1. The surface vacancies can migrate via in-plane sites into the bulk in a thermally activated process, which becomes noticeable at temperatures above 250 °C [24]. Moreover, the formation energies of these vacancies mutually depend on the defects' densities, as indicated by computational studies [34]. It means that while the initial reduction of the stoichiometric surface will proceed by mainly removing O<sub>br</sub>, at a relative V<sub>br</sub> density of 0.5, the formation energy difference between bridging and in-plane defects is much smaller, so both should coexist at the surface. Therefore, the equilibrium density of surface vacancies should depend not only on the chemical potential of gaseous oxygen but also on the density of subsurface V<sub>O</sub> in bulk SnO<sub>2</sub>, assuming the temperature is high enough for the atoms to reach their equilibrium position. The role of vacancies in sensor operation mechanisms is two-fold; as a self-doping mechanism, which provides intrinsic conductivity in SnO<sub>2</sub>, and as the preferred adsorption site for oxygen, which does not adsorb onto stoichiometric surfaces of SnO<sub>2</sub>, as indicated by many computational studies [17,28,32] and explained in the following section.



**Figure 1.** Progressing reduction of the (110,101) surfaces of SnO<sub>2</sub>. Image prepared in VESTA [35].

### 3. Interactions with Oxygen

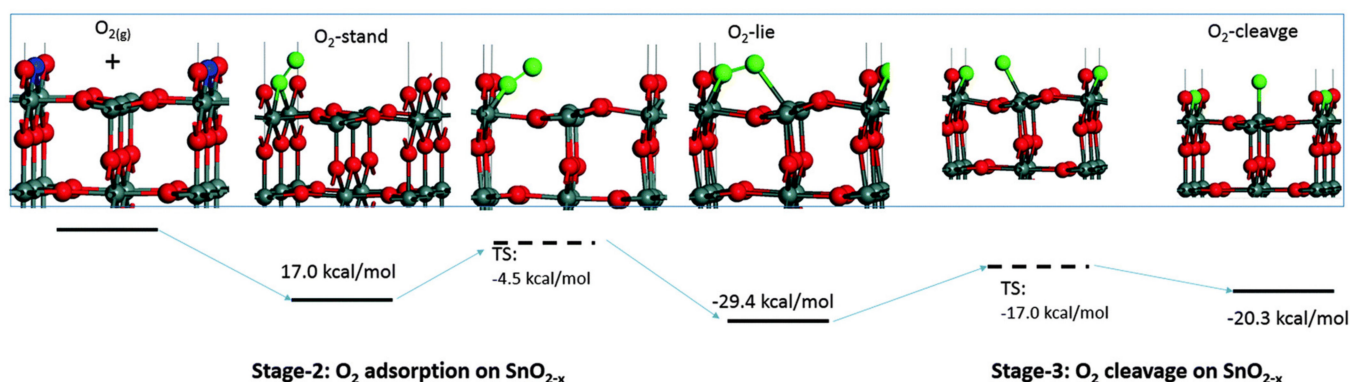
#### 3.1. Insight from Computational Studies

The early computational work by Yamaguchi et al. on the surfaces of SnO<sub>2</sub> explored the adsorption energies of oxygen adsorbates placed on a (110) surface bridging vacancy us-

ing a point charge model [36]. They found that adsorbate interaction with a vacancy yields negative adsorption energies for a range of both mono- and diatomic species. According to their results, a molecule of oxygen can adsorb both in side-on and end-on configurations and can be singly or doubly charged, i.e.,  $O_2^-$  and  $O_2^{2-}$ . The study also shows that an  $O^-$  adsorbate is stabilised by the surface vacancy with an energy similar to that of typical bridging lattice oxygen,  $O^{2-}$  (16.22 vs. 16.97 eV per  $2O^-$ ). However, the authors acknowledge that the investigated configurations of oxygen adsorbates were guessed due to the lack of information on adsorption geometry.

In a follow-up letter, Yamaguchi considers the dissociation of oxygen on a fully reduced surface and concludes that the dissociation of  $O_2^{2-}$  leads to the formation of  $O^-$  monoatomic species, which are coupled with a  $V_{br}$  and immobilised by a rather sizeable potential barrier, and that this adsorption site is more favourable than a neighbouring  $Sn_{5c}$  site [27]. However, a more recent computational study suggests that the range of possible adsorbates is less diverse. Investigation of a similarly reduced surface led to the conclusion that the dissociation of  $O_2^{2-}$  results in the healing of two vacancies, i.e., the formation of  $O^{2-}$  ions indistinguishable from  $O_{br}$  [17]. The same study considers the adsorption of  $O_2$  onto an isolated  $V_{br}$ , in which case they find that the dissociation of  $O_2$  leads to the healing of the vacancy by one of the atoms, while the other becomes an  $O^-$  ion adsorbed onto a neighbouring  $Sn_{5c}$  site, as there are no other nearby vacancies to interact within the model. Therefore, it seems that  $O^-$  should only exist under particular circumstances as a remnant of the final steps of surface oxidation after nearly all vacancies have been healed. For both the reduced surface and an isolated vacancy, the dissociation has a relatively large activation energy (1.32 and 0.54 eV, respectively) [17], which is at variance with the results of Yamaguchi, who reports an activation barrier of only 0.09 eV [27]. Such a small barrier is unlikely as, experimentally, oxidation of the surface was determined to be a thermally activated process [37].

Prior to the dissociation, the molecular oxygen adsorbates may exist in a physisorbed state, where there is no charge transfer between the surface and the molecule, and in a chemisorbed state, where the molecule withdraws electron density from the surface. Currently, there is no agreement on whether the charge of the adsorbate is related to the adsorption geometry. While one DFT study finds that end-on and side-on configurations (Figure 2) correspond to superoxide and peroxide species, respectively [17], a different study only reports a side-on configuration for diatomic oxygen, which they assign to a superoxide [28]. In both studies, the adsorbate is bound to the  $V_{br}$  on one side and  $Sn_{5c}$  on the other, which seems to be the most favourable adsorption geometry; see O-lie in Figure 2.



**Figure 2.** Adsorption of  $O_2$  onto a (110) surface of  $SnO_2$  with an isolated vacancy. From left: a (110) surface with a single (periodic) bridging vacancy (blue), “O-stand” superoxide molecular adsorbate in the end-on configuration, “O-lie” peroxide molecular adsorbate in the side-on configuration and “O-cleavage” monoatomic  $O^-$  adsorbate adjacent to a healed vacancy—now  $O_{br}$ . Reproduced with permission from Reference [17].



The above descriptions were of the (110) surface, for which most computational studies are performed, and there are limited data on the adsorption of O<sub>2</sub> onto other surfaces. An investigation of the (101) surface indicates that O<sub>2</sub> interaction mechanisms are similar to those of the (110) surface [32]. Only weak interactions are found for a stoichiometric surface, and the vacancy is again an essential active site. Upon adsorption onto an isolated vacancy on a partially reduced surface, the molecule transforms into a peroxide-like species with an adsorption energy of 1.02 eV. Unlike for the (110) surface, this adsorption configuration should not lead to dissociation into O<sup>-</sup>, as the peroxide bond cannot break [32]. However, the adsorption of O<sub>2</sub> onto two neighbouring vacancies leads to the adsorbate dissociation and healing of both defects, similar to the (110) surface.

Although the localisation of electrons in the antibonding  $\pi^*$  orbitals of the peroxide adsorbates results in the formation of negatively charged species, this process should not lead to the accumulation of negative charge at the surface or a subsurface depletion layer. It was found that upon adsorption onto a vacancy, the oxidation state of a neighbouring Sn is changed from +2 to +4 [32]; therefore, the charge of the V<sub>br</sub> should be -2 rather than 0 for the surface to remain overall charge-neutral. Because of this, the peroxide species should not contribute to a dipole layer formation, rather than to the localisation of electrons on the vacancy site.

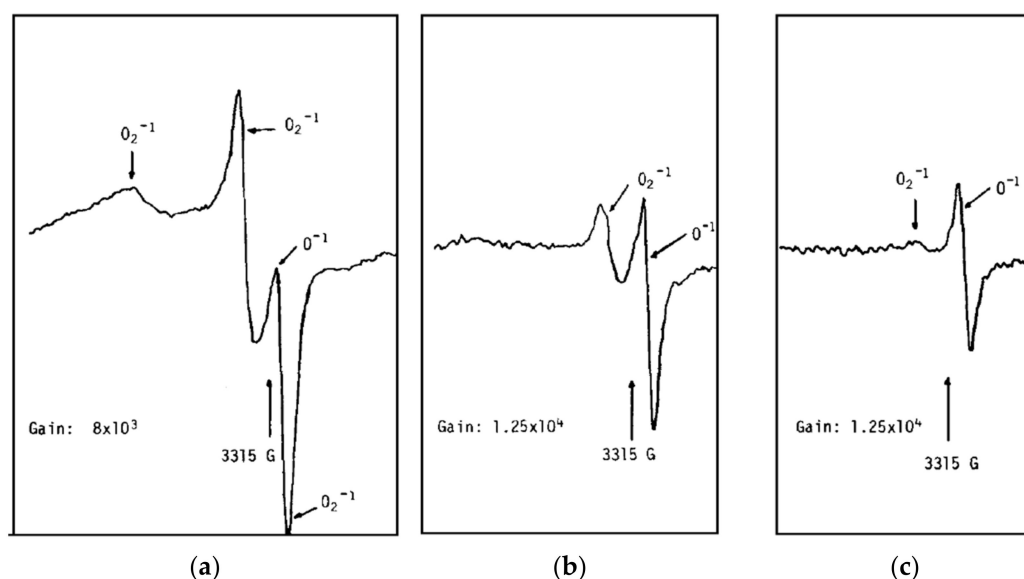
Recently, a different computational description of the adsorption of O<sub>2</sub> onto SnO<sub>2</sub> was proposed, where instead of creating a vacancy, the surface was enriched with extra electrons to enable adsorption. The study shows that an electron added to the system can stabilise it by lowering the energy of one antibonding  $\pi^*$  orbital of O<sub>2</sub><sup>-</sup> to just above the valence band maximum [38]. In such a configuration, the molecule is bound to two neighbouring Sn<sub>5c</sub> sites, and the negative charge of the adsorbate is not compensated by its location, as it lies outside of regular lattice anion sites. Therefore, this process should lead to the accumulation of charge at the surface. However, it is not entirely clear where the electrons come from in a system without oxygen vacancies or dopants, which could significantly alter the results by affecting the electronic structure of the surface [39]. The study also shows that introducing a second electron does not significantly affect the geometry of the adsorbate or the energy of the second  $\pi^*$  orbital. Instead, the second and subsequent electrons occupy the conduction band; therefore, the authors concluded that O<sub>2</sub><sup>2-</sup> should not form on the (110), (100) and (101) surfaces of SnO<sub>2</sub> [38]. This conclusion is at variance with the experimental evidence of peroxide O<sub>2</sub><sup>2-</sup> adsorbate [40], discussed in the following section.

The same study revisits the possibility of monoatomic oxygen adsorbates, pointed out by the authors as the species responsible for the observed resistance change [38]. Several adsorption configurations were found on the surface, which may be divided into two categories, “peroxide-like” and “lattice-like” species. The former is characterised by a peroxide bond to an O<sub>br</sub> in addition to a bond to Sn<sub>5c</sub>; this adsorption configuration is equivalent to the adsorption of O<sub>2</sub><sup>2-</sup> onto a V<sub>br</sub>, and the authors conclude that it is charge neutral with respect to the lattice and should not affect the resistance of the sensor. The latter is bound to either one or two different Sn<sub>5c</sub> sites in a geometry not far removed from the local coordination of lattice oxygen atoms; in this case, stable adsorption configurations were found only on the (100) surface, which features a considerable distortion of the adsorption site. On the (110) and (101) surfaces, all monoatomic adsorbates were found unstable to transformation into O<sub>2</sub> and O<sub>2</sub><sup>-</sup>. Based on these results, the authors concluded that the resistance of SnO<sub>2</sub> is defined by interactions of O<sub>2</sub> with only the (100) surface. However, gas-sensitive phenomena were also observed on other surfaces, for example, an epitaxial (101) thin film [41], which shows that these results might be incomplete.

### 3.2. What Species Were Observed Spectroscopically?

Because of the aforementioned “pressure gap”, the choice of analytical techniques suitable for studying gas sensors is limited, so the evidence is scarce. The most widely accepted evidence of oxygen ionosorption on SnO<sub>2</sub> comes from an ESR investigation

published by Chang in 1980 [42]. The paper describes a reversible transformation of  $O_2^-$  into  $O^-$  witnessed at temperatures between 100 and 150 °C, shown in Figure 3. However, it was already pointed out that some of these assignments are inconsistent with both the theory and other ESR studies [16]; while a singlet was assigned to  $O^-$  in this study, theory shows that the signal should be a triplet [43], which is what was found for many other compounds, including closely related  $TiO_2$  [44,45]. The analysis of ESR spectra is quite complicated, and there is no consensus on the assignment of some peaks; thus, there is no conclusive evidence of paramagnetic monoatomic oxygen species on  $SnO_2$ .

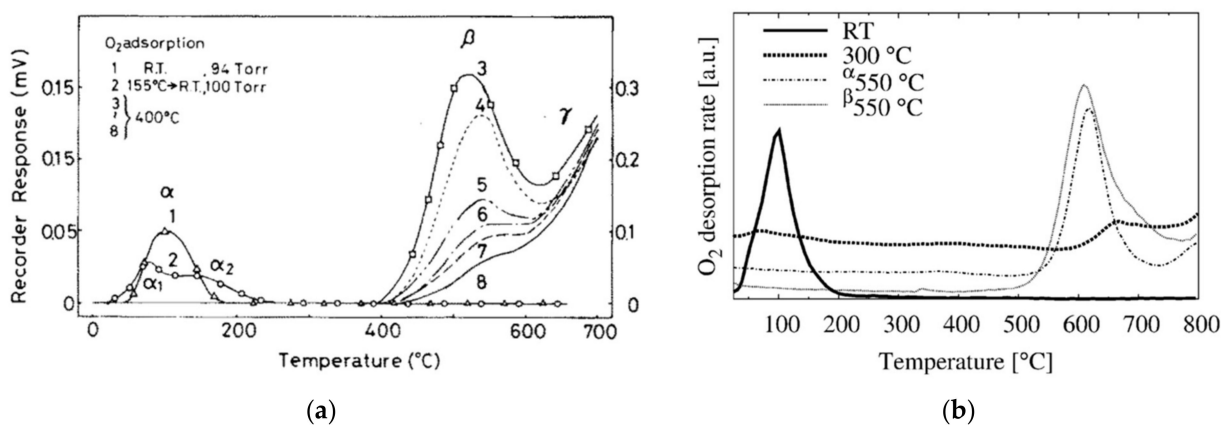


**Figure 3.** ESR spectra of  $SnO_2$  in UHV. The sample was calcined in UHV at 550 °C for 4 h. (a) Sample at 20 °C. (b) The sample was warmed up to 100 °C. (c) The sample was warmed up to 150 °C. Reproduced with permission from Reference [42].

On the other hand, multiple ESR studies show triplet signals consistent with  $O_2^-$  adsorbates [46–49], which was confirmed by the appearance of the hyperfine structure in  $^{17}O$  isotope exchange experiments [49]. This signal appears only on reduced surfaces, which have been heat-treated and cooled in UHV and exposed to  $O_2$  below 150 °C, and its intensity is increased on Pt-doped surfaces [46]. These findings are in agreement with computational models of oxygen adsorption onto  $SnO_2$ , which only show adsorption on reduced surfaces [17,28,32]. However, ESR cannot provide evidence for the other computationally predicted species, the peroxide  $O_2^{2-}$ , which is not paramagnetic.

The existence of such low-temperature molecular adsorbates is also confirmed by TPD studies, which show that oxygen desorption from  $SnO_2$  may be classified into three peaks,  $\alpha$ ,  $\beta$  and  $\gamma$ , as presented in Figure 4. The appearance of the  $\alpha$  peak is correlated with low-temperature adsorption onto a surface preheated in vacuum at 700 and 550 °C (Figure 4a,b, respectively) for several hours and cooled down to below 150 °C before  $O_2$  exposure. However, the peak does not appear if the sample was exposed to  $O_2$  at temperatures above 300 °C; instead, two different desorption events occur, the  $\beta$  and  $\gamma$  peaks, which were assigned to  $O^-$  and lattice oxygen, respectively. However, the assignments are made based solely on the observed desorption temperatures. In one of the studies, the identity of the  $\alpha$  and  $\beta$  adsorbates was investigated by ESR [50]; while  $\alpha$  was confirmed to be  $O_2^-$ , no paramagnetic signal was detected for  $\beta$ . Therefore, the assignment of  $\beta$  as  $O^-$  is again doubtful, and other plausible explanations for the origin of  $\beta$  peak ought to be considered. It should also be pointed out that the onset and intensity of the  $\beta$  desorption peak depend not only on the pressure but also on the temperature at which oxygen dosing was performed [50–52]. Considering this and the fact that the peak coincides with the temperature range at which  $SnO_2$  is known to reduce spontaneously [53], the  $\beta$

peak could be explained by the desorption of  $O_{br}$  (surface reduction). Such assignment would also explain why the  $\beta$  peak does not appear after low-temperature dosing, where the chemisorbed oxygen species have insufficient energy to dissociate and heal the bridging vacancies. However, more data are necessary to make this determination conclusively.



**Figure 4.** TPD chromatograms of  $O_2$  desorption from  $SnO_2$ . (a) Traces 1–2 correspond to the low-temperature adsorption following the evacuation at 700 °C. Traces 3–8 correspond to oxygen dosing at 400 °C and pressures of 99, 48, 16, 6, 3 and 2 Torr, respectively. Reproduced with permission from Reference [50]. (b) TPD of  $SnO_2$  after exposure to  $O_2$ . Trace  $\alpha$  denotes sample cooled in UHV, while  $\beta$  denotes cooling in  $O_2$  background. Reproduced with permission from Reference [54].

The more recent evidence of adsorbed species comes from an operando DRIFTS study of  $O_2$  and  $H_2$  on a UHV treated sample, which is expected to have an appreciable initial density of vacancies. The spectra collected in several-minute intervals display emerging absorption bands attributable to superoxide and peroxide species [40]. The bands appear over the first few minutes and stabilise by about 15 min after gas introduction. If the atmosphere is then changed to  $H_2$ , the bands are removed rapidly over the first 5 min and disappear completely after just 10 min. Reintroduction of oxygen leads to the reappearance of the bands, showing the reversible formation of these adsorbates on reduced surfaces of  $SnO_2$ .

The development of near ambient pressure photoelectron analysers opened up the possibility of operando XPS analysis. Conveniently, the thermocouple lead of most NAP XPS systems can be used to measure the resistance of the sample in situ, therefore allowing analysis of a real, working sensor. The recent report of such an investigation does not show changes in the shape of  $O 1s$  emission, which would have confirmed the presence of adsorbed species [19]. However, the spectra were collected at temperatures above 200 °C, so diatomic adsorbates ( $\alpha$  peak in TPD) were unlikely to persist at the surface. On the other hand, a rigid shift in the binding energy (BE) scale was observed for all core levels, suggesting a change in the Fermi level of the sample, which is evidence of band bending. The peaks shift to lower BE in the  $O_2$  atmosphere relative to UHV, which the authors interpret as evidence of upward band bending caused by chemisorbed oxygen. However, another possible interpretation of the observed Fermi level change stems from a different choice of reference. The formation of vacancies in  $SnO_2$  introduces shallow electron donors, which may become thermally ionised and cause downward band bending. This would manifest itself as a shift of peaks to higher BE in photoelectron spectroscopy, which was observed in both UPS and XPS on reduced and oxidised  $SnO_2$  samples [53]. These authors conclude that the flat band condition corresponds to a perfect, stoichiometric surface on which no adsorbates can exist, and the bands bend downward from that, moving peaks in both UPS and XPS spectra to higher binding energies.

#### 4. Interactions with Reducing Gases

##### 4.1. Ionosorption Based Models—Reactions with Preadsorbed Oxygen

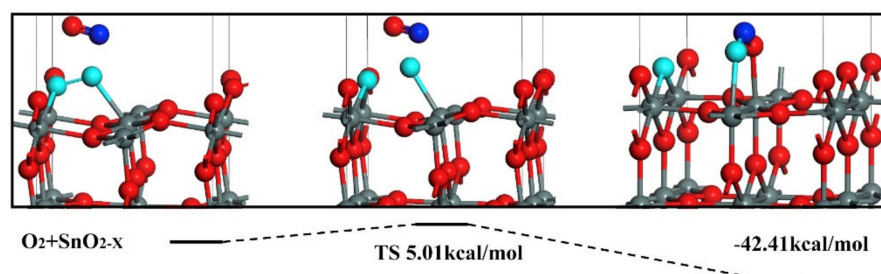
Since the foundations of oxygen ionosorption were established, the detection of reducing gases was understood in terms of the reaction with surface oxygen ions and the release of electrons trapped at the surface into the conduction band. The papers on the detection of common reducing gases, such as CO, H<sub>2</sub> or VOC, describe the surface reactions in a manner presented in Table 1.

**Table 1.** Reactions of common reducing gases with ionosorbed oxygen.

Reaction	References
$\text{CO} + \text{O}^- \rightarrow \text{CO}_2 + \text{e}^-$	[8,55,56]
$\text{H}_2 + \text{O}^- \rightarrow \text{H}_2\text{O} + \text{e}^-$	[50,57–60]
$\text{VOC} + n\text{O}^- \rightarrow x\text{CO}_2 + y\text{H}_2\text{O} + n\text{e}^-$	[4,61,62]

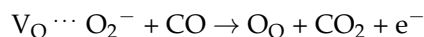
Similar to the core theory behind them, these mechanisms do not refer to the adsorption configurations of the species involved. As discussed earlier, computational studies suggest that O<sup>−</sup> cannot be implicitly assumed to exist, and empirical studies systematically fail to show O<sup>−</sup> on SnO<sub>2</sub>, suggesting that these reactions may be an oversimplification that defines a particular stoichiometry for the surface processes rather than defining the involvement of specific adsorbed species.

In the previous chapter, the adsorption configuration of an oxygen molecule onto a V<sub>br</sub> was discussed. Computational studies show that various combustible gases, such as CO and NO, can adsorb on top of an O<sub>2</sub> molecule and subsequently react with it [17,28,30,32]. On both the (110) and (101) surfaces, one half of the dioxygen heals a V<sub>br</sub>, and the other reacts with the combustible molecule, as shown in Figures 5 and 6.



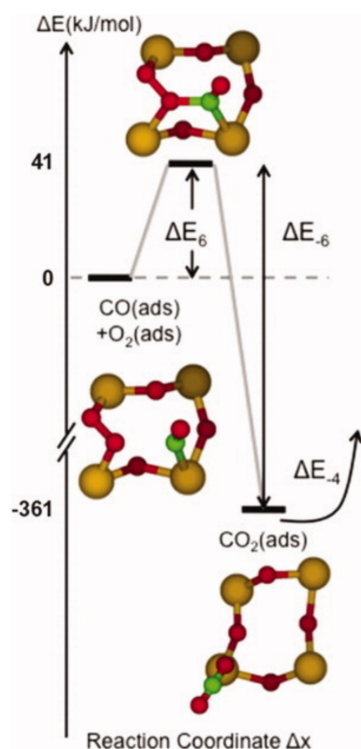
**Figure 5.** Coadsorption of NO and O<sub>2</sub> onto a reduced (110) surface of SnO<sub>2</sub>. Reproduced with permission from Reference [30].

By the example of CO, the overall process occurring at a vacancy can therefore be described as:



Similar to the O<sup>−</sup> based detection in Table 1, this process involves releasing one electron per molecule detected; however, it is described by a species known to exist on SnO<sub>2</sub>, and the calculated transition state energies are reasonably low (0.80 eV on the (110) surface and 0.43 eV on the (101) surface) [17,32]. The products of the reactions are initially adsorbed but may either desorb or further react with the surface. The latter possibility is explored below.





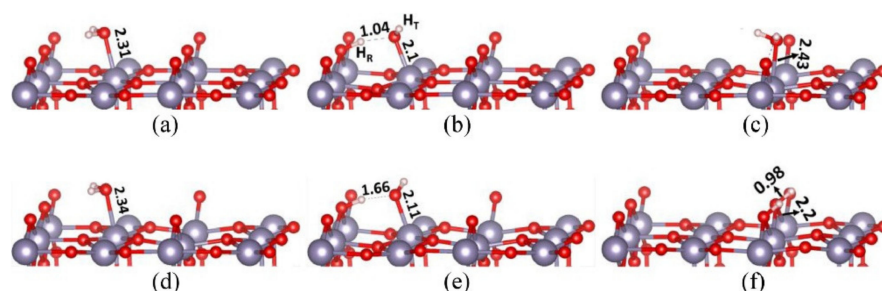
**Figure 6.** Coadsorption of CO and O<sub>2</sub> onto a reduced (101) surface of SnO<sub>2</sub>. CO reacts with O<sub>2</sub> preadsorbed onto a V<sub>O</sub>. Reproduced with permission from reference [32].

While the reaction of CO and NO with preadsorbed oxygen species leads to the healing of a surface oxygen vacancy, the coadsorption of O<sub>2</sub> and H<sub>2</sub> is predicted to look quite different. In this case, one hydrogen atom binds to each of the adsorbed dioxygen atoms to form two distinct hydroxyl groups [63]. Hydrogenation of the oxygen atom bound to the V<sub>O</sub> results in a rooted hydroxyl ((OH)<sub>O</sub>), which occupies the bridging lattice site, while the other oxygen atom becomes a terminal hydroxyl bound to a neighbouring Sn<sub>5c</sub> site (Sn<sub>5c</sub> ⋯ OH). The two hydroxyl groups have dissimilar properties and hence affect the electronic structure of SnO<sub>2</sub> differently [59]. The Sn<sub>5c</sub> ⋯ OH hydroxyl forms a dipole, and it does not contribute to the sensor's resistance other than by removing the oxygen adsorbate during hydroxyl formation. On the other hand, the (OH)<sub>O</sub> is an electron donor (relative to the equivalent O<sub>br</sub> it replaces), which can donate charge carriers to the conduction band and become positively charged; the formation of such hydroxyl surface donors is expected to have a noticeable effect on the sensor's resistance [58,59,64].

Although these mechanisms show how reducing gases may be detected on the surfaces of SnO<sub>2</sub>, they do not provide the complete picture. If a sensor's resistance were regulated solely by the formation and consumption of oxygen adsorbates, then the lowest attainable resistance would be in the absence of oxygen. However, it was shown experimentally that the resistance of a sensor might fall below the no-oxygen baseline during CO detection in humid synthetic air [64]. Therefore, mechanisms extending beyond the interaction with preadsorbed oxygen need to be considered.

#### 4.2. Direct Adsorption of Reducing Gases

Humidity is virtually always present in the atmosphere in varying amounts, which is a potential problem for the reliability of gas sensors, as the adsorption of water could cause the resistance of a sensor to change. Computational studies show that the interaction between the surface and H<sub>2</sub>O is complex; many adsorption configurations are possible, including both molecular and dissociative modes, as shown in Figure 7.



**Figure 7.** Adsorption of H<sub>2</sub>O onto (a,b) stoichiometric and (c–f) reduced (110) surfaces of SnO<sub>2</sub>. Adsorption energies are given for each configuration. (a) Molecular adsorption onto Sn<sub>5c</sub>, −1.29 eV. (b) Dissociative adsorption on Sn<sub>5c</sub>, −1.85 eV. (c) Molecular adsorption onto V<sub>br</sub>, −0.81 eV. (d) Molecular adsorption onto Sn<sub>5c</sub> with a neighbouring V<sub>br</sub>, −1.32 eV. (e) Dissociative adsorption on Sn<sub>5c</sub> with a neighbouring V<sub>br</sub>, −1.92 eV. (f) Dissociative adsorption on V<sub>br</sub>, −2.48 eV. Reproduced with permission from Reference [29].

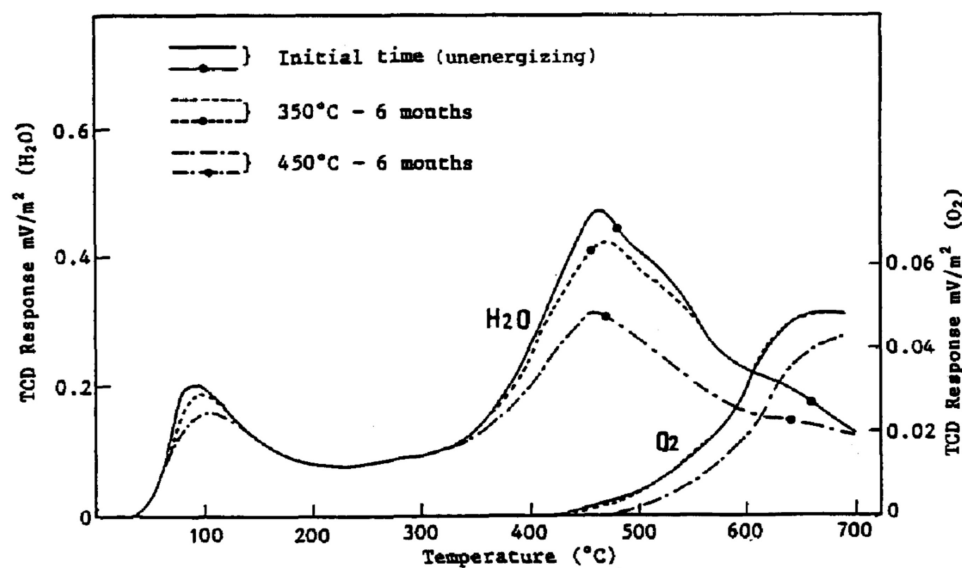
On a stoichiometric surface, H<sub>2</sub>O can adsorb onto an Sn<sub>5c</sub> site [29]. The adsorption energy is relatively small compared to other possible configurations but can be further lowered by the dissociation of the H<sub>2</sub>O molecule, where one of the H atoms coordinates to a neighbouring O<sub>br</sub>. The remaining OH coordinated to the Sn<sub>5c</sub> becomes a terminal hydroxyl (Sn<sub>5c</sub> ··· OH), and the reaction of O<sub>br</sub> with H atom transforms it into a rooted hydroxyl ((OH)<sub>O</sub>); an example of such an interaction is shown in Figure 7b. In fact, the result of dissociative adsorption of H<sub>2</sub>O on a stoichiometric surface is quite similar to the reaction of H<sub>2</sub> with preadsorbed O<sub>2</sub>, as discussed above.

The formation of an oxygen vacancy on the surface strongly influences the adsorption of H<sub>2</sub>O; the adsorption energy of H<sub>2</sub>O onto an Sn<sub>5c</sub> is more negative on a surface with a V<sub>br</sub> in the vicinity, both in the molecular and dissociative configurations [29]. However, the vacancy is also a new adsorption site for H<sub>2</sub>O, leading to more stable species. Although molecular adsorption onto V<sub>br</sub> is the least stable of all configurations, the dissociative mode on V<sub>br</sub>, as shown in Figure 7f, is the opposite, being the most favourable adsorption configuration of all. This configuration is also unique in that it produces two (OH)<sub>O</sub> instead of one of each type of hydroxyl. Therefore, it should affect the resistance more strongly as two electron donors are formed per molecule detected instead of one [29,65].

Although dissociative adsorption is more stable than its molecular counterpart on both adsorption sites, H<sub>2</sub>O cannot always spontaneously dissociate upon adsorption. The process was shown to be thermally activated in an operando DRIFTS study, where the appearance of bands corresponding to OH was observed only after dosing H<sub>2</sub>O at temperatures above 200 °C [66].

The preferred adsorption mode of H<sub>2</sub>O, therefore, depends strongly on the surface of SnO<sub>2</sub>, which determines the range of possible configurations based on the availability of adsorption sites, for example neighbouring V<sub>br</sub> and O<sub>br</sub>, to allow dissociative adsorption. This may not necessarily be a common occurrence, as the bridging vacancies tend to group together, for example, in the proposed structure of the 2 × 1 reconstruction of the (110) surface, where every other row of bridging oxygen atoms is removed. This suggests that depending on the preparation, the surfaces should behave differently in the presence of humidity, which was shown experimentally in an H<sub>2</sub>O/D<sub>2</sub>O exchange study using DRIFTS [67]. The study also confirms that SnO<sub>2</sub> surface hydroxylation is a reversible process in a dynamic equilibrium with the gas phase. Additional TPD data show that H<sub>2</sub>O desorbs from the surface over a wide range of temperatures, suggesting multiple adsorption configurations with different adsorption energies [51,52]. The TPD chromatograms (Figure 8) show two broad peaks, one around 100 °C, which corresponds to molecularly adsorbed water, and another above 400 °C, which results from various dissociatively adsorbed species. Since the adsorption energy decreases with increasing

OH coverage, it takes progressively more energy to desorb H<sub>2</sub>O from SnO<sub>2</sub>, which is why temperatures of at least 600 °C are required for complete dehydroxylation [66].

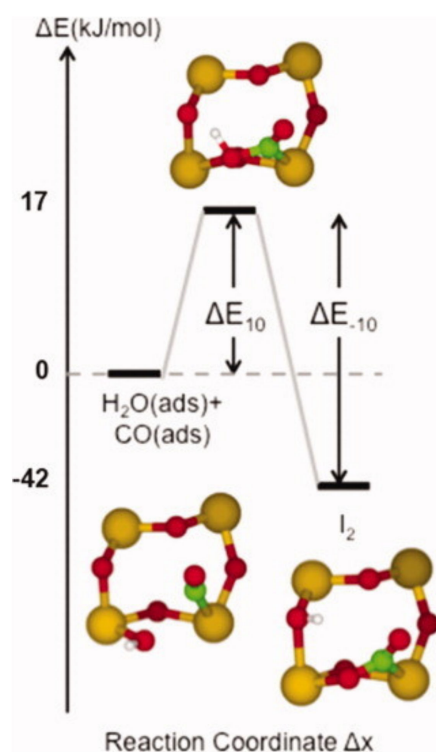


**Figure 8.** TPD chromatograms of H<sub>2</sub>O and O<sub>2</sub> desorption from SnO<sub>2</sub>. Reproduced with permission from Reference [51].

An operando DRIFTS study with simultaneous resistance measurements reveals that H<sub>2</sub>O does not always induce a reducing response. [67]. Two differently prepared samples of SnO<sub>2</sub> were exposed to increasingly humid synthetic air. While one of the samples shows a strong response, the effect of humidity on the other sample is negligible. However, spectra reveal that the former is accompanied by relatively small changes in the OH bands, while intense OH formation is observed for the latter. Therefore, experimental data suggest that there must be two adsorption mechanisms, one which causes a reducing type response, likely linked to the formation of (OH)<sub>O</sub> electron donors, and one which is relatively electroneutral, which could be dominated by the formation of Sn<sub>5c</sub> ··· OH [59].

Although ionosorption is the prevalent sensor operation mechanism discussed in the literature, it has been known for a long time that CO elicits a response on SnO<sub>2</sub>-based sensors in the absence of oxygen [56]. Since no CO<sub>2</sub> was observed during CO adsorption in the absence of O<sub>2</sub>, the change in resistance was attributed to the direct adsorption of CO onto SnO<sub>2</sub>, which produces a surface donor state that can become ionised by donating an electron into the conduction band, increasing the sensor's resistance. The authors support their argument with the study of Heiland et al., who studied the interactions of SnO<sub>2</sub> with acetic acid by TPD [68]. According to the TPD results, the reaction products desorb from the surface primarily as CO instead of CO<sub>2</sub>, which Hahn interprets as a lack of reaction between CO and the surface. However, these results are at odds with the findings of Yamazoe et al., who recorded TPD chromatograms after dosing CO, rather than acetone, onto SnO<sub>2</sub> powder. They found only CO<sub>2</sub> desorbing from the surface and concluded that CO irreversibly adsorbs onto SnO<sub>2</sub> and desorbs as CO<sub>2</sub> upon heating, suggesting the reaction follows a Mars-van Krevelen mechanism [69]. Therefore, it is likely that multiple adsorption pathways for CO are possible.

It is well known that humidity has a substantial impact on the detection of CO in SnO<sub>2</sub>-based gas sensors. One of the possible contributory processes is the transformation of O<sub>br</sub> into (OH)<sub>O</sub> when a CO molecule reacts with a terminal OH, which was found on a (101) surface [32]. The initial and final adsorbate configurations and the transition state are shown in Figure 9. Although the process forms a donor state, which will influence the sensor's resistance, it permanently removes important adsorption sites for CO, as explained in the next section.



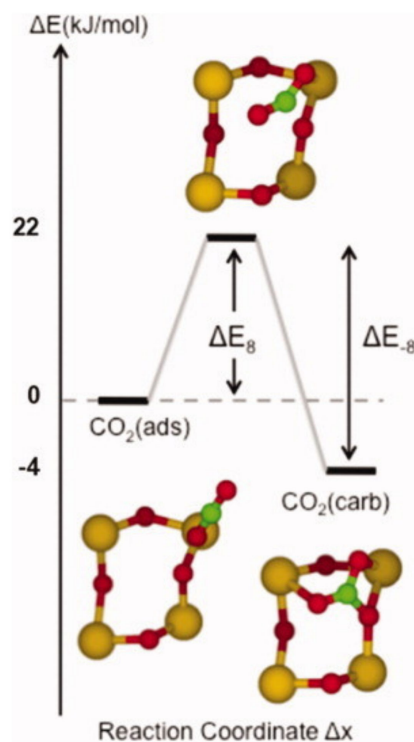
**Figure 9.** Adsorption of CO onto a (101) surface of SnO<sub>2</sub> with a terminal hydroxyl on an adjacent Sn<sub>5c</sub> site. Reproduced with permission from Reference [32].

Computational investigation of NO adsorption onto SnO<sub>2</sub> reveals that various geometries are possible, but the most stable ones are linked to nitrogen (N<sup>-</sup>) down configurations [30]. Contrary to H<sub>2</sub>O adsorption, the bonding interaction between NO and Sn<sub>5c</sub> is more favourable on a stoichiometric surface than on a reduced one. However, when present at the surface, the V<sub>br</sub> is again the preferred adsorption site. Although a different study reports no O-down geometries [18], both studies agree that the most stable adsorption configuration is N-down onto a V<sub>br</sub> with an adsorption energy of about 1.2 eV. Analysis of the charge density transfer reveals that, in terms of the influence on the electronic structure, NO can act not only like a reducing gas but also like an oxidising one [18]. While adsorption onto O<sub>br</sub> results in a donor-like NO<sup>+</sup> state, the same molecule adsorbing onto a V<sub>br</sub> results in electron density being withdrawn from the surface and the formation of an NO<sup>-</sup>. Effectively, NO adsorption will produce a reducing gas response on a stoichiometric surface but an oxidising gas response on a fully reduced surface. For intermediate surfaces, a mixture of the adsorbates is expected at the surface; such behaviour complicates NO detection.

The detection mechanisms of many gases involve their transformation into reaction products such as H<sub>2</sub>O (from H<sub>2</sub>), NO<sub>2</sub> (from NO) and CO<sub>2</sub> (from CO). Of these, the formation of H<sub>2</sub>O is usually lost in the background humidity, but the ambient concentrations of CO<sub>2</sub> and NO<sub>2</sub> are respectively small and negligible, which means that their interference should be taken into consideration. NO<sub>2</sub> is considered an oxidising gas, which strongly increases the resistance of the sensor. Its adsorption is discussed in the next section.

On the other hand, CO<sub>2</sub> seems to have little effect on the sensor's resistance compared to other gases; considerable concentration changes produce a relatively small response. A computational study found that CO<sub>2</sub> does not adsorb onto a stoichiometric (110) surface of SnO<sub>2</sub>, as only very weak interactions were found [70]. Moreover, on a reduced (110) surface with preadsorbed V<sub>O</sub> ··· O<sub>2</sub>, the CO<sub>2</sub> does not interact with O<sub>2</sub>, as the two molecules are localised far apart from each other in their equilibrium position. These two conclusions imply that CO<sub>2</sub> should not substantially affect the sensor's resistance, as they eliminate two

main detection mechanisms. However, a different study suggests the possibility of CO<sub>2</sub> adsorption on a stoichiometric (101) SnO<sub>2</sub> surface, with the molecule coordinating to an O<sub>br</sub> and two Sn<sub>5c</sub> sites in a tridentate carbonate form, as shown in Figure 10 [32]. A charge transfer occurs from the surface, with each forming carbonate gaining 0.120 e<sup>-</sup>, determined from population analysis. Moreover, it was shown that the (110) surface could also be activated towards CO<sub>2</sub> detection; an investigation of CO<sub>2</sub> adsorption onto a hydroxylated surface shows a possible reaction leading to the formation of a bidentate carbonate ion adsorbed onto two Sn<sub>5c</sub>, with the remaining H atom coordinating to a neighbouring O<sub>br</sub> to form an (OH)<sub>O</sub>. It appears that the presence of Sn<sub>5c</sub> ⋯ OH also activates neighbouring O<sub>br</sub> to react with CO<sub>2</sub> forming a bidentate CO<sub>3</sub><sup>2-</sup> between an Sn<sub>5c</sub> and V<sub>br</sub>. These findings show how humidity should enhance the sensing properties of SnO<sub>2</sub> towards CO<sub>2</sub>, which has been demonstrated experimentally [70]. However, because of the poor performance of pure SnO<sub>2</sub> in CO<sub>2</sub> detection, most studies involve modification of the surface, for example, by lanthanum doping [71,72], and relatively few studies investigate the interaction of CO<sub>2</sub> with undoped SnO<sub>2</sub>.

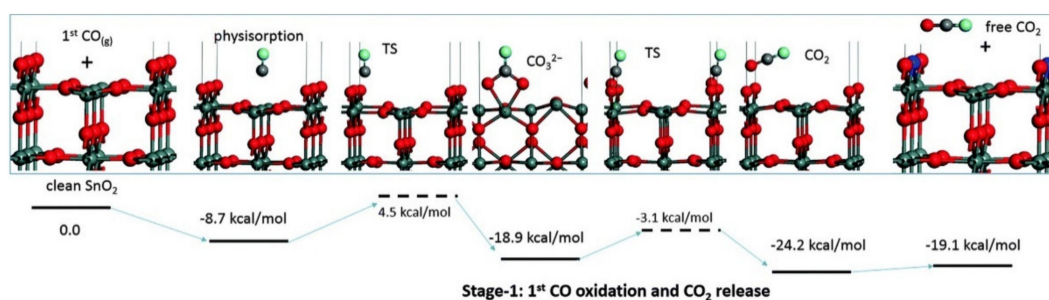


**Figure 10.** Adsorption of CO<sub>2</sub> onto a stoichiometric (101) surface of SnO<sub>2</sub>. The formed adsorbate is a tridentate carbonate-like species. Reproduced with permission from Reference [32].

#### 4.3. Reduction of the Surface

There is another reaction pathway for detecting CO, which recently has gained more recognition within the gas sensing community. As mentioned before, the direct reduction of the SnO<sub>2</sub> surface by CO was initially discounted as no CO<sub>2</sub> was detected in the outflowing gas [56]. However, such behaviour would vary significantly from many other oxides, on which CO oxidation proceeds via a Mars-van Krevelen (MvK) mechanism, which features a cyclic reduction/oxidation of the adsorption sites [73,74]. Indeed, DFT studies on various SnO<sub>2</sub> surfaces have shown repeatedly that the direct reduction of the surface by CO should also be possible [17,28,32,75]. Several mechanisms were proposed, through which the CO molecule may react with O<sub>br</sub>, for example, through a carbonate intermediary, as shown in Figure 11.





**Figure 11.** Direct reduction of the (110) surface of SnO<sub>2</sub> during CO adsorption. Reproduced with permission from Reference [17].

Although the formation of a carbonate species on the surface is often associated with CO detection, it is likely a secondary process resulting from the adsorption of the produced CO<sub>2</sub>. Experimentally, it was found that carbonates do not always appear during CO dosing on SnO<sub>2</sub> and that identical features in the DRIFTS spectra may be produced by dosing CO<sub>2</sub> instead of CO [76]. Indeed, the formation of surface carbonates seems to be related to the immediate sample pre-treatment, with the species appearing after preparation at 450 °C but not after 1000 °C [67,76].

Whether the carbonates are formed during or after the interaction of SnO<sub>2</sub> with CO, the predicted reduction of the surface by CO was confirmed experimentally in operando UV-Vis diffuse reflectance spectroscopy [77]. The analysis revealed strong absorption in the visible region upon dosing with CO, consistent with the formation of in-gap electronic states of oxygen vacancies, making the typically transparent semiconductor opaque. Additional evidence comes from a TPD study, which shows that CO desorbs from the SnO<sub>2</sub> surface almost exclusively as CO<sub>2</sub>, which indicates that CO reacts with the surface upon adsorption [69]. Moreover, the amount of CO<sub>2</sub> desorbed after CO dosing depends strongly on the degree of surface reduction, and therefore the density of O<sub>br</sub> (versus V<sub>br</sub>) on the surface. This indicates that surface oxygen atoms are important adsorption sites and that non-reductive adsorption of CO onto SnO<sub>2</sub> is limited. Therefore, there is strong experimental evidence supporting this mechanism of gas–surface interactions.

It seems fortuitous that CO can reduce the surface of SnO<sub>2</sub> directly by reacting with O<sub>br</sub>; the V<sub>br</sub> are essential adsorption sites for oxygen that enable the secondary reaction of CO with the preadsorbed oxygen, reforming the O<sub>br</sub> lost in the reduction. This cyclic process results in a dynamically changing density of O<sub>br</sub> on the surface. The steady-state coverage, which eventually determines the electronic properties of the surface, will depend on the rates of reduction and oxidation, and therefore the concentration of CO, assuming the ambient oxygen concentration is constant. Because of the strong evidence supporting the reaction, it is more reasonable to perceive the detection mechanism of CO in terms of plausible reactions of well-established species rather than an elusive monoatomic adsorbate; such descriptions of CO detection have been proposed recently [76], which foreshadows a new era in gas-sensor modelling.

Although the reactant geometry of the reduction of SnO<sub>2</sub> by NO is similar to CO, the reaction is less likely in the case of NO. DFT calculation determined that this process is endothermic, making it less likely to occur at noticeable rates [30]. Nevertheless, there is a possibility of a direct surface reduction by NO, which results in the formation of NO<sub>2</sub> adsorbed onto a bridging vacancy. The production of this interfering gas is somewhat problematic, as it is an oxidising gas that causes a significant increase in the sensor's resistance. The adsorption mechanisms of NO<sub>2</sub> are the subject of the following section.

## 5. Interactions with Oxidising Gases—NO<sub>2</sub>

### 5.1. Adsorption onto Stoichiometric Surfaces

Even though oxygen lends its name to this category of target gases, it rarely is considered one, as it defines the baseline resistance of sensors. Therefore, oxidising gases are

those that increase the sensor's resistance relative to an ambient oxygen background, for example,  $\text{NO}_2$  or  $\text{SO}_2$ . The current research focuses on the former, and there is very little research on other oxidising gases, with  $\text{SO}_2$  sometimes included as a frequent interfering gas to  $\text{NO}_2$ .

Conventionally,  $\text{NO}_2$  detection is considered in terms similar to oxygen ionosorption; the adsorbates withdraw electrons from the conduction band of  $\text{SnO}_2$ . However, the formed states are deeper than those of oxygen, causing stronger band bending and a more significant resistance increase [78]. Even though this description is fairly straightforward, the reactions of  $\text{NO}_2$  on the surface of  $\text{SnO}_2$  seem to be quite complex. In contrast to carbon, which only forms two oxides, nitrogen can form a variety of oxides, such as  $\text{N}_2\text{O}$ ,  $\text{NO}$ ,  $\text{NO}_2$ ,  $\text{N}_2\text{O}_4$  and the  $\text{NO}_3^-$  ion, many of which are formed during parasitic reactions on the surface of the sensor [78] as evidenced by IR and Raman spectroscopic analyses [66,79]. Long sensor recovery time, during which conductivity of the sensing layer returns to its baseline value after removal of the target gas, has been linked to the formation of the  $\text{NO}_3^-$  ion, which has significantly larger adsorption energy than  $\text{NO}_2$ , and therefore desorbs at a slower rate [80]. In order to improve the performance of sensors, it is crucial to understand the formation mechanisms of various nitrogen oxides on the surfaces of  $\text{SnO}_2$  and to design sensitive layers that could hinder parasitic reactions.

On stoichiometric  $\text{SnO}_2$  (101) surfaces,  $\text{NO}_2$  can adsorb in a monodentate or bidentate bridging geometry with one or two neighbouring  $\text{Sn}_{5c}$  sites [81]. The adsorption energy between the two configurations differs by only about 0.13 eV. It means that even at room temperature, the molecule should convert easily between the two states, which would result in a "random walk" along the  $\text{Sn}_{5c}$  rows. This could eventually lead to two  $\text{NO}_2$  molecules happening upon one another at the surface, resulting in them reacting; one of the molecules abstracts an oxygen atom from the other to become an  $\text{NO}_3^-$  ion, while the other is reduced to  $\text{NO}$ , which may adsorb onto a neighbouring site or desorb from the surface. The occurrence of this reaction is supported by the observation of an intermediate  $\text{NO}_2$  dimer, denoted as  $\text{ONONO}_2$ , using in-situ Raman spectroscopy [79].

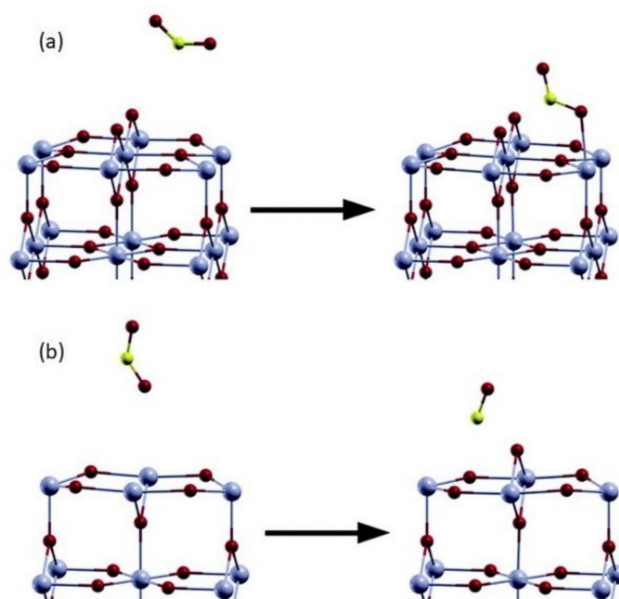
Whereas the interaction of  $\text{SnO}_2$  with the produced  $\text{NO}$  is relatively weak, the  $\text{NO}_3^-$  adsorbs much more strongly than  $\text{NO}_2$ . Contrary to  $\text{NO}_2$ , the energy difference between the monodentate and bridging geometries of  $\text{NO}_3^-$  is considerable [81]; therefore, a random walk along the Sn paths is unlikely. Moreover, because of the stronger adsorption of  $\text{NO}_3^-$ , temperatures of up to 700 °C are necessary to remove the adsorbates altogether. In terms of their electronic effects, the Bader charge trapped by the  $\text{NO}_3^-$  is over twice as large as that trapped by an adsorbed  $\text{NO}_2$  molecule, which suggests it should affect the sensor's resistance much more strongly [80]. As a result, the formation of  $\text{NO}_3^-$  may lead to sensor poisoning and long recovery times.

### 5.2. Adsorption onto Reduced Surfaces

Although a "walking"  $\text{NO}_2$  molecule is interesting to picture, the preferable adsorption site for this gas is a surface oxygen vacancy, as shown in independent computational studies [18,31]. The molecule interacts with a vacancy via the oxygen atoms. Both studies determined that  $\text{NO}_2$  adsorption onto a bridging vacancy is stronger than at an in-plane vacancy.  $\text{NO}_2$  adsorbed onto  $V_{br}$  is most stable in a monodentate configuration, though bidentate geometries coordinated to a nearby  $\text{Sn}_{5c}$  or another  $V_{br}$  are also possible. In all such geometries, the molecule is immobile, which lowers its chance of reacting with another  $\text{NO}_2$  to form nitrates. Many studies highlight the importance of oxygen vacancies in the low-temperature detection of  $\text{NO}_2$  on  $\text{SnO}_2$ -based gas sensors, although they sparsely comment on the atomistic origin of the superior sensitivity [82–84].

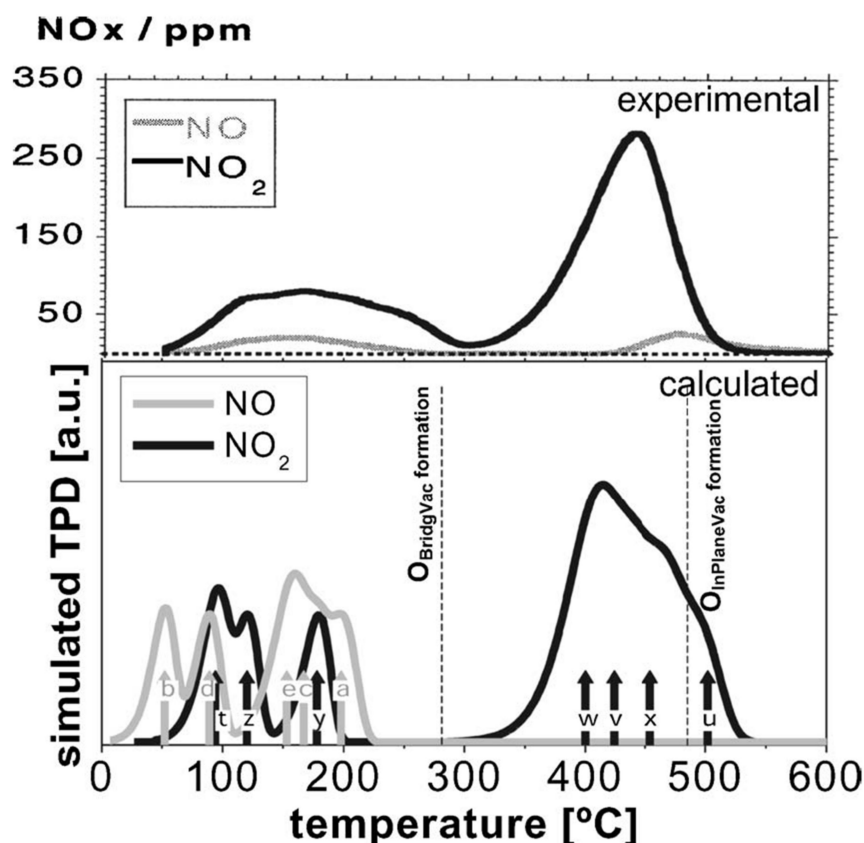
Adsorption onto an oxygen-deficient surface offers a reaction pathway that is unavailable on a stoichiometric one, which is the dissociation of  $\text{NO}_2$ , as shown in Figure 12. We have already considered the mechanism by which reducing gases affect the sensor's resistance; the current process can be thought of as the direct opposite. Similar to oxygen, the adsorption of  $\text{NO}_2$  at a surface vacancy can lead to the oxidation of the surface, as

shown in a computational study of the (110) surface [18]. The adsorption of  $\text{NO}_2$  onto  $V_{\text{br}}$  leads to the dissociation of the molecule and healing the vacancy; the authors comment on the likeness between the final structure achieved here and during the adsorption of  $\text{NO}$  onto a stoichiometric surface. Therefore, apart from the electron-withdrawing properties of  $\text{NO}_2$  adsorption, the removal of  $V_{\text{O}}$  self-doping in  $\text{SnO}_2$  must be considered. Moreover, the molecule preferentially heals the bridging oxygen site on a surface where in-plane vacancies are also available, in agreement with the preferable adsorption onto  $O_{\text{br}}$ . This selectivity is inverted in the case of  $\text{SO}_2$ , which preferentially adsorbs onto  $V_{\text{pl}}$  [31]. Therefore, in theory, one can regulate the selectivity of a sensor by selectively creating bridging vacancies.



**Figure 12.** Adsorption of  $\text{NO}_2$  onto (110) surface of  $\text{SnO}_2$ . (a) Stoichiometric surface, monodentate adsorption onto  $\text{Sn}_{5c}$ . (b) Reduced surface, dissociative adsorption onto a  $V_{\text{br}}$ , leading to the healing of the vacancy and an adsorbed  $\text{NO}$  molecule. Reproduced with permission from Reference [18].

Adsorption of  $\text{NO}_2$  onto  $\text{SnO}_2$  with various distributions of both bridging and in-plane vacancies was studied via first-principles thermodynamics, and the results show clearly that bridging vacancies are the more favourable of the two [85]. Based on these results and the Redhead equation, a simulated TPD chromatogram was computed and compared to experimental data, as shown in Figure 13. The match is reasonable, although experimental data show some high-temperature desorption of  $\text{NO}$ , which was not predicted by DFT. Moreover, the empirical TPD study that benchmarks simulated desorption shows  $\text{NO}$  desorption only at temperatures above  $200\text{ }^\circ\text{C}$  on a differently prepared surface [66]. This could be explained by a temperature-activated dissociation of  $\text{NO}_2$  at a surface vacancy, resulting in  $\text{NO}$  desorbing at temperatures above typical, which the DFT study did not consider. However, it should be mentioned that while the desorption is dominated by  $\text{NO}_2$  in this study, an earlier TPD study found only  $\text{NO}$  and some  $\text{N}_2\text{O}$  desorbing from  $\text{SnO}_2$  powder after adsorption at room temperature [69]. This interesting difference, if confirmed, could indicate the importance of the atomistic state of the surface in determining its interactions with gas adsorbates.



**Figure 13.** Comparison of empirical and simulated NO<sub>x</sub> desorption in TPD of SnO<sub>2</sub> after exposure to NO<sub>2</sub>. Reproduced with permission from Reference [85].

## 6. Oxygen Vacancies versus Gas Sensing Mechanism

The previous chapters discussed the importance of oxygen vacancies as adsorption sites, their formation during the interaction with reducing gases, and their healing via the complex interactions with oxygen or oxidising gases. However, as mentioned before, oxygen vacancies are also a crucial self-doping mechanism of SnO<sub>2</sub> and largely determine its intrinsic conductivity [24]. Although a detailed description of the involvement of V<sub>O</sub> in the transducer function of gas sensors is beyond the scope of this review, there are several points worth considering.

Even though oxygen vacancies have long been recognised as the source of conductivity in bulk SnO<sub>2</sub>, recent computational studies show these electronic states to be deep in the bandgap, thus unlikely to contribute to conductivity [23]. Moreover, the defects are now considered negative-U type, where the attractive interaction between two electrons bound at a vacancy means that the system's energy with a singly occupied vacancy can be lowered by capturing a conduction electron at the defect [23]. By implication, singly ionised vacancies are not thermodynamically stable for any Fermi level position [86]. This computational finding contradicts the well-established evidence of paramagnetic (singly occupied) oxygen vacancies observed in ESR studies [22,46], which could hint at a difference between the nature of bulk and near-surface vacancies. The effect of vacancies on the electronic properties of SnO<sub>2</sub> could depend on their depth below the surface, with bulk vacancies being energetically deep, negative-U and likely neutral, and surface vacancies singly or doubly ionised depending on the temperature. The singly charged vacancies, which could emerge in the substoichiometric coordination of a reduced surface, should then lead to the formation of a surface conductivity layer, while the bulk of the material remains an insulator (relative to the surface). Indeed, such surface conductivity layer models have been proposed recently for In<sub>2</sub>O<sub>3</sub>, a gas-sensitive material analogous to SnO<sub>2</sub>,

whose conductivity as a thick particulate film is several times larger than bulk conductivity models can account for [87].

The existence of singly ionised vacancies is important to acknowledge, as it reveals certain parallelism between ionosorption and surface reduction. If the formed oxygen vacancy is a singly ionised donor, the conduction band has gained one electron per lattice oxygen atom lost, similar to the formation of  $O^-$ . Moreover, such interpretation of the meaning of phenomenological “ionisorbed oxygen” could explain why some studies observe macroscopic changes consistent with the appearance of ionisorbed  $O^{2-}$  at elevated temperatures, even though such species are known to be unstable outside of the stabilising influence of the Madelung potential on oxygen lattice sites.

## 7. Conclusions

The insight provided by recent molecular modelling studies, backed by emerging operando spectroscopy, offers a new perspective on the gas-sensitive phenomena of  $SnO_2$ -based gas sensors, and therefore our understanding of the sensing mechanisms needs to be re-evaluated. Although the prevalent model of oxygen ionosorption can explain some of the observed macroscopic behaviour of the sensors, it is based on species thus far unconfirmed. The systematic lack of evidence after so many years of research indicates that new mechanisms need to be considered.

The only convincing evidence of oxygen adsorbates on the  $SnO_2$  surface is linked to diatomic species, whose effect on the sensor's resistance is limited according to ionosorption theory, which requires  $O^-$  to work. Although  $O^-$  may exist under special circumstances, most interactions between the surface,  $O_2$  and target gases are mediated by the diatomic oxygen adsorbates. It is important to mention that both the dissociation of  $O_2$  and the reaction of a target gas with preadsorbed oxygen results in the healing of a surface vacancy, which should affect the equilibrium density of vacancies, providing a mechanism for resistance change.

Furthermore, direct evidence of the reduction of the  $SnO_2$  surface was finally obtained, which settles the debate of whether CO can react with the  $O_{br}$ , as predicted for some years based on DFT calculations. The lack of  $CO_2$  observed during the dosing of CO on  $SnO_2$  in the absence of oxygen is then somewhat puzzling but could be explained by a slow rate of reaction and the depletion of  $O_{br}$ , which cannot be replenished by ambient oxygen. This is supported by the strong dependence on the amount of desorbed species after CO dosing based on the degree of surface reduction.

The importance of bridging sites is further highlighted in the adsorption mechanisms of  $H_2O$ .  $O_{br}$  is a reaction partner for the dissociation of  $H_2O$  on both reduced and stoichiometric surfaces, and the dissociation on a  $V_{br}$  leads to the formation of an additional electron donor, thus causing a larger resistance change. Moreover, the transformation of  $O_{br}$  into  $(OH)_O$  is the most likely candidate for the interference in CO detection, as the essential adsorption site is removed.

Finally, the bridging sites are heavily involved in the detection of both NO and  $NO_2$ , which transform into one another in a network of complex interactions with the surface. In the process, they either form or heal bridging vacancies, which should strongly affect the baseline conductivity of  $SnO_2$ . However, these two gases also adsorb readily onto various other sites, which should also have a noticeable effect on the sensor's resistance. The mutual coexistence of the absorbed species in these processes is evidenced by the results of the adsorption of one or the other, which leads to a similar final state, as noticed independently by both computational and empirical scientists.

In conclusion, the majority of evidence points to the mechanisms centred around oxygen vacancies. The surface of  $SnO_2$  is a dynamic assembly, which can support a large density of vacancies both at the surface (bridging and in-plane) and in several atomic layers below the surface. As evidenced by experimental studies, the in-plane oxygen vacancies play an important role in determining its conductivity and reactivity. Therefore, the surface of  $SnO_2$  is an exchange zone between bulk and gas-phase oxygen, and the sensors response



can be seen as resulting from the change in the position of the exchange equilibrium, on top of which the effects of direct adsorption may manifest as additional resistance change.

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