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Lower bounds for the first eigenvalue of the Laplacian with zero magnetic field in planar domains[☆]



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ABSTRACT

We study the Laplacian with zero magnetic field acting on complex functions of a planar domain Ω , with magnetic Neumann boundary conditions. If Ω is simply connected then the spectrum reduces to the spectrum of the usual Neumann Laplacian; therefore we focus on multiply connected domains bounded by convex curves and prove lower bounds for its ground state depending on the geometry and the topology of Ω . Besides the area, the perimeter and the diameter, the geometric invariants which play a crucial role in the estimates are the fluxes of the potential one-form around the inner holes and the distance between the boundary components of the domain; more precisely, the ratio between its minimal and maximal width. Then, we give a lower bound for doubly connected domains which is sharp in terms of this ratio, and a general lower bound for domains with an arbitrary number of holes. When the inner holes shrink to points, we obtain as a corollary a lower bound for the first eigenvalue of the so-called Aharonov-Bohm operators with an arbitrary number of poles.

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1. Introduction

1.1. Definitions and state of the art

Let Ω be a bounded, open, connected domain with smooth boundary $\partial\Omega$ in a Riemannian manifold (M, g) and let A be a smooth real one-form on Ω , (the *potential* one-form). Define a connection ∇^A on the space of complex-valued functions $C^\infty(\Omega, \mathbf{C})$ as follows:

$$\nabla_X^A u = \nabla_X u - iA(X)u,$$

for all vector fields X on Ω , where ∇ is the Levi-Civita connection of M . The *magnetic Laplacian with potential A* is the operator acting on $C^\infty(\Omega, \mathbf{C})$:

$$\Delta_A = (\nabla^A)^* \nabla^A.$$

In \mathbf{R}^n this gives explicitly, in the usual notation:

$$\Delta_A = (i\nabla + A^\sharp)^2,$$

where A^\sharp is the dual vector field of A , the *vector potential*. The two-form $B = dA$ is the magnetic field; in dimension 2, by duality, B is the scalar field $B = \text{curl}A^\sharp = \partial_{x_1}A_2 - \partial_{x_2}A_1$.

Scope of this paper is to discuss the spectrum of Δ_A for planar domains. Hence in what follows we take $\Omega \subset \mathbf{R}^2$.

The spectrum of the magnetic Laplacian has been studied extensively for Dirichlet boundary conditions ($u = 0$ on $\partial\Omega$), and we denote by $\lambda_1^D(\Omega, A)$ the first eigenvalue. First we remark that, thanks to the diamagnetic inequality, one always has:

$$\lambda_1^D(\Omega, A) \geq \lambda_1^D(\Omega, 0),$$

and in particular $\lambda_1^D(\Omega, A) > 0$. For planar domains and constant magnetic field (that is, $dA = B$ and $|B|$ constant), a Faber-Krahn inequality holds, in the sense that the first eigenvalue of a planar domain is minimized by that of the disk of the same area (see [6]). Estimates for sums of eigenvalues can be found in [5] and [11].

However in this paper we deal with *magnetic Neumann boundary conditions*, that is we impose $\nabla_N^A u = 0$, on the boundary, where N is the inner unit normal to $\partial\Omega$. It is known that then Δ_A admits a discrete spectrum

$$0 \leq \lambda_1^N(\Omega, A) \leq \lambda_2^N(\Omega, A) \leq \dots$$

diverging to $+\infty$. The first eigenvalue has the following variational characterization:

$$\lambda_1^N(\Omega, A) = \min_{u \in H_A^1(\Omega) \setminus \{0\}} \frac{\int_\Omega |\nabla^A u|^2 dx}{\int_\Omega |u|^2 dx}, \tag{1}$$

where $H_A^1(\Omega)$ is the magnetic Sobolev space:

$$H_A^1(\Omega) = \{u \in L^2(\Omega) : \nabla^A u \in L^2(\Omega)\}.$$

When the potential is sufficiently regular, for example $A \in L^\infty(\Omega)$, the magnetic Sobolev space coincides with $H^1(\Omega)$.

For computing lower bounds the diamagnetic inequality is of no use; in fact it gives:

$$\lambda_1^N(\Omega, A) \geq \lambda_1^N(\Omega, 0) = 0,$$

because $\lambda_1^N(\Omega, 0)$ is simply the first eigenvalue of the usual Neumann Laplacian, which is zero (the associated eigenspace being spanned by the constant functions). There are fewer estimates in this regard; let us first discuss the case of a constant magnetic field $|B| = B_0 > 0$ on planar domains. The paper [4] gives a lower bound of $\lambda_1^N(\Omega, A)$ in terms of the inradius of Ω , $\lambda_1^N(\Omega, 0)$ and of course B_0 . Asymptotic expansions as $|B| \rightarrow \infty$ are obtained in [8]. We also mention the paper [7] which investigates the validity of a reverse Faber-Krahn inequality for constant magnetic field B_0 , that is: is it true that $\lambda_1^N(\Omega, A)$ is always bounded above by that of a disk with equal volume? It is proved there that this inequality is true when B_0 is either sufficiently small or sufficiently large, but the general case is still open in the simply connected case.

- In this paper we prove three lower bounds for the first eigenvalue of planar domains under Neumann conditions, when the magnetic field is identically zero. *Since this will be the only boundary condition we consider, from now on we will simply write $\lambda_1(\Omega, A)$ instead of $\lambda_1^N(\Omega, A)$.*

Let us first clarify the circumstances under which the first eigenvalue might be positive even if the magnetic potential is a closed one-form on Ω . This is intimately related to a phenomenon in quantum mechanics predicted in 1959 and known as *Aharonov-Bohm effect*, which has also experimental evidence: a particle traveling a region in the plane might be affected by the magnetic field even if this is identically zero on its path. In fact what the particle “feels” is not the magnetic field but, rather, the magnetic potential A , provided that A is closed but not exact, and that the flux of A around the pole may assume non-integer values (see below for the precise condition).

Let us be more precise. From the definition we see that, if $A = 0$, the spectrum of Δ_A coincides with the spectrum of the usual Laplacian under Neumann boundary conditions. The same is true when $A = df$ is an exact one-form, by the well-known *gauge invariance* of the magnetic Laplacian. This fundamental property states that the spectrum of Δ_{A+df} is the same as the spectrum of Δ_A , for any $f \in C^\infty(\Omega)$, which follows from the identity:

$$\Delta_A e^{-if} = e^{-if} \Delta_{A+df}$$

showing that Δ_A and Δ_{A+df} are unitarily equivalent.

On the other hand, if the magnetic field $B = dA$ is non-zero, then $\lambda_1(\Omega, A)$ is strictly positive. One could then ask if $\lambda_1(\Omega, A)$ has to vanish whenever the magnetic field is zero, that is, whenever A is a *closed* one-form.

To that end, let c be a closed curve in Ω (a loop). The quantity:

$$\Phi_c^A = \frac{1}{2\pi} \oint_c A$$

is called the *flux of A across c* (we assume that c is traveled once, and we will not specify the orientation of the loop; this will not affect any of the statements, definitions or results which we will prove in this paper).

It turns out that

- $\lambda_1(\Omega, A) = 0$ if and only if A is closed and the cohomology class of A is an integer, that is, the flux of A around any loop is an integer.

This was first observed by Shigekawa [14] for closed manifolds, and then proved in [10] for manifolds with boundary. This remarkable feature of the magnetic Laplacian shows its deep relation with the topology of the underlying manifold Ω . In this paper we will focus precisely on the situation where the potential one form is *closed*, and we will then give two lower bounds for the first eigenvalue $\lambda_1(\Omega, A)$.

Let us then recall a few previous results when the magnetic field is assumed to vanish. A lower bound for a general Riemannian cylinder (i.e. the surface $\mathbf{S}^1 \times (0, L)$ endowed with a Riemannian metric) and zero magnetic field has been given in [3], and is somewhat the inspiration of this work: one of two main results here is in fact to improve such bound when Ω is a doubly connected planar domain.

Directly related to the Aharonov-Bohm effect, we mention the papers [1] and [12] which investigate the behavior of the spectrum of a domain with a pole $\Omega \setminus \{a\}$ when the pole a approaches the boundary, for Dirichlet boundary conditions. We remark here that the pole is a distinguished point $a = (a_1, a_2)$ and the potential is the harmonic one-form:

$$A_a = \frac{1}{2} \left(- \frac{x_2 - a_2}{(x_1 - a_1)^2 + (x_2 - a_2)^2} dx_1 + \frac{x_1 - a_1}{(x_1 - a_1)^2 + (x_2 - a_2)^2} dx_2 \right)$$

which has flux $\frac{1}{2}$ across any closed curve enclosing a , giving rise to a magnetic field which is a Dirac distribution concentrated at the pole a (therefore, the magnetic field indeed vanishes on $\Omega \setminus \{a\}$). The magnetic Laplacian Δ_{A_a} acting on $\Omega \setminus \{a\}$ is often called an Aharonov-Bohm operator. One could think to a domain with a pole as a doubly connected domain for which the inner boundary curve shrinks to a point.

We will in fact give a lower bound for the first eigenvalue of Aharonov-Bohm operators with many poles, and Neumann boundary conditions (see Theorem 5).

The Aharonov-Bohm operators play an interesting role in the study of minimal partitions, see chapter 8 of [2].

For Neumann boundary conditions, we mention the paper [10], where the authors study the multiplicity and the nodal sets corresponding to the ground state for non-

simply connected planar domains with harmonic potential. For doubly connected domains, it is shown that $\lambda_1(\Omega, A)$ is maximal precisely when Φ^A is congruent to $\frac{1}{2}$ modulo integers (this fact is no longer true when there are more than two holes). The proof relies on a delicate argument involving the nodal line of a first eigenfunction and the conclusion does not follow from a specific comparison argument, or from an explicit lower bound.

The focus of this paper is on lower bounds for multiply connected planar domains and zero magnetic field defined by the closed potential form A . By what we have just said, it is clear that estimating the first eigenvalue is a trivial problem if Ω is simply connected, because then any closed one-form is automatically exact, and therefore $\lambda_1(\Omega, A) = 0$ by gauge invariance. Therefore, we restrict our study to domains with n holes, with $n \geq 1$.

In this paper we will prove: an improved lower bound for doubly connected domains; a general lower bound for multiply connected domains with an arbitrary number of convex holes; a lower bound for a general convex domain with an arbitrary number of punctures. Let us describe these results in detail.

1.2. A lower bound for doubly connected domains

Let us start from doubly connected domains ($n = 1$) hence domains of type:

$$\Omega = F \setminus \bar{G},$$

where F and G are open and smooth. We assume F and G convex. Let Φ^A be the flux of the closed potential A around the inner boundary curve ∂G : by Shigekawa result, the lower bound is simply zero when Φ^A is an integer. Then, to hope for a positive lower bound, we need to measure how much Φ^A is far from being an integer, and the natural invariant will then be:

$$d(\Phi^A, \mathbf{Z}) = \min\{|\Phi^A - k| : k \in \mathbf{Z}\}.$$

The second important ingredient for our lower bounds is the ratio $\frac{\beta}{B}$ between the minimal width and the maximal width of Ω . To be more precise, let us say that the line segment $\sigma \subset \Omega$ is an *orthogonal ray* if it hits the inner boundary ∂G orthogonally. By definition, the *minimal width* β (resp. *maximal width* B) of Ω is the minimal (resp. maximal) length of an orthogonal ray contained in Ω , see Fig. 1.

Note that the ratio $\frac{\beta}{B}$ is invariant under homotheties, and reaches its largest value 1 whenever the boundary components are parallel curves.

In Theorem 2 of [3] we prove the lower bound:

$$\lambda_1(\Omega, A) \geq \frac{4\pi^2}{|\partial F|^2} \frac{\beta(\Omega)^2}{B(\Omega)^2} d(\Phi^A, \mathbf{Z})^2. \tag{2}$$

We insist on the fact that if $\frac{\beta}{B}$ is bounded below away from zero we get a positive, uniform lower bound even if β tends to zero. Think for example to a concentric annulus

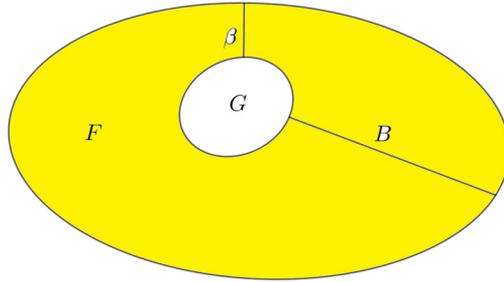


Fig. 1. The minimal width β and the maximal width B of an annulus $\Omega = F \setminus \bar{G}$.

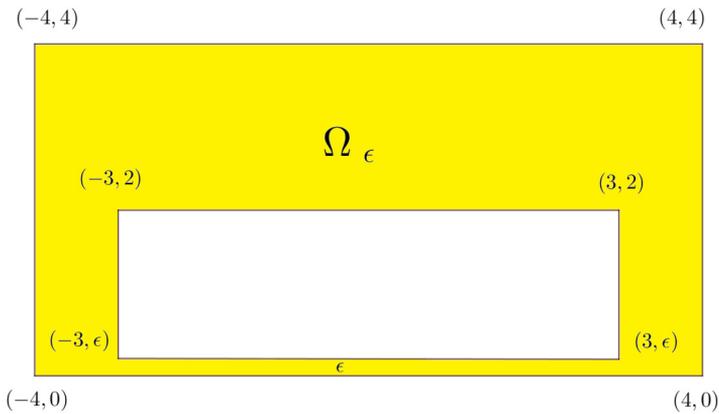


Fig. 2. The domain Ω_ϵ has minimal width ϵ and lowest eigenvalue going to zero proportionally with ϵ .

Ω of radii 1 and $1 + \beta$; then $\frac{\beta}{B} = 1$. Consider the magnetic Laplacian with potential $A = \frac{\gamma}{x_1^2 + x_2^2}(-x_2, x_1)$, for which $\Phi^A = \gamma$. From Proposition 8 of [3] we deduce that, as $\beta \rightarrow 0$, the lower bound will approach $d(\Phi^A, \mathbf{Z})^2$. Note that this number is also the first eigenvalue of the magnetic Laplacian $(-i\frac{d}{d\theta} + A)^2$ on the unit circle (see Remark 2.1.4 in [9]). In particular, (2) reduces to an equality as $\beta \rightarrow 1$.

This means that (for fixed perimeter) in order to get λ_1 small, the ratio $\frac{\beta}{B}$ (and not just β) has to be small.

Sharpness in terms of $\frac{\beta}{B}$. In [3] we showed that if $\frac{\beta}{B}$ is small then the first eigenvalue could indeed be small. We then looked for an example which could show that the dependence on $\frac{\beta^2}{B^2}$ is sharp, and we could not find it. Rather, in Examples 14 and 15 in [3], we constructed examples of domains such that B is bounded below, say by 1, $|\partial F|$ is bounded above, β goes to zero and $\lambda_1(\Omega, A)$ goes to zero proportionally to β , for any non-integral flux. Therefore, if one could replace $\frac{\beta^2}{B^2}$ by the linear factor $\frac{\beta}{B}$ in (2), one would obtain a sharp inequality (with respect to $\frac{\beta}{B}$). See Fig. 2 for the example which shows sharpness.

This is in fact possible, and the theorem which follows should be regarded as the first main theorem of this paper.

Theorem 1. Let $\Omega = F \setminus \bar{G}$ be an annulus in the plane, with F and G convex with piecewise-smooth boundary. Let A be a closed one-form with flux Φ^A around the inner hole G . Then:

a) One has the lower bound:

$$\lambda_1(\Omega, A) \geq \frac{\pi^2}{8} \cdot \frac{|F|^2}{|\partial F|^2 D(F)^4} \cdot \frac{\beta(\Omega)}{B(\Omega)} \cdot d(\Phi^A, \mathbf{Z})^2, \tag{3}$$

where $\beta(\Omega)$ and $B(\Omega)$ are, respectively, the minimal and maximal width of Ω , and $D(F)$ is the diameter of F .

b) If $R_{\text{in}}(F)$ denotes the inradius of F , then (3) takes the form:

$$\lambda_1(\Omega, A) \geq \frac{\pi^2}{32} \cdot \frac{R_{\text{in}}(F)^2}{D(F)^4} \cdot \frac{\beta(\Omega)}{B(\Omega)} \cdot d(\Phi^A, \mathbf{Z})^2. \tag{4}$$

c) If the outer boundary ∂F is smooth, and if $\beta(\Omega)$ is less than the injectivity radius of the normal exponential map of ∂F , then we have the simpler lower bound:

$$\lambda_1(\Omega, A) \geq \frac{\pi^2}{|\partial F|^2} \frac{\beta(\Omega)}{B(\Omega)} d(\Phi^A, \mathbf{Z})^2. \tag{5}$$

- We remark that b) follows immediately from a) by using the inequality $\frac{|F|}{|\partial F|} \geq \frac{R_{\text{in}}(F)}{2}$, which is valid for any convex domain in \mathbf{R}^2 . Thus, the dependence on the geometry in the main lower bound (3) can be expressed solely in terms of the diameter, the inner radius and the minimal and maximal widths.
- Note that, modulo a factor of 4, (5) is formally identical to (2) with β/B replacing β^2/B^2 .

We observe that there is no positive constant c such that

$$\frac{\beta(\Omega)}{B(\Omega)} \geq c \frac{|F|^2}{D(F)^4}$$

for all doubly convex annuli in the plane (otherwise, the lower bound would be independent on the inner hole, and this is impossible). This means that Theorem 1 is not a trivial consequence of (2).

In fact, the proof of Theorem 1 uses a suitable partition of Ω into overlapping annuli for which $\frac{\beta}{B}$ is, so to speak, as small as possible (see Section 2 below, and in particular Fig. 3 for an example). Recall the δ -interior ball condition:

given $x \in \partial F$, there is a ball of radius δ tangent to ∂F at x and entirely contained in F .

Here and for further applications, we say that the injectivity radius of ∂F is $\text{Inj}(\partial F)$ if F satisfies the δ -interior ball condition for any $\delta \leq \text{Inj}(\partial F)$. If ∂F is smooth, its injectivity radius is positive.

Finally we refer to Fig. 2 for a picture of the family of domains Ω_ϵ which realize sharpness. Ω_ϵ is the difference between two rectangles with parallel sides, with boundaries being ϵ units apart. Hence $\beta(\Omega_\epsilon) = \epsilon$ and $B(\Omega_\epsilon)$ is uniformly bounded above by $\sqrt{5}$.

We show in Section 2.8 that

$$\frac{\pi^2}{360\sqrt{5}}d(\Phi^A, \mathbf{Z})^2 \leq \frac{\lambda_1(\Omega_\epsilon, A)}{\epsilon} \leq \frac{1}{10},$$

so that $\lambda_1(\Omega, A)$ goes to zero proportionally to $\epsilon \sim \frac{\beta}{B}$.

1.3. Sharpness in terms of $d(\Phi^A, \mathbf{Z})$

We now examine how precise is the estimate of Theorem 1 with respect to $d(\Phi^A, \mathbf{Z})$, as $d(\Phi^A, \mathbf{Z}) \rightarrow 0$.

Let Ω_a be a concentric annulus, centered at the origin, with radii $a < 1$ and 1. Let $A = \nu(-\frac{x_2}{x_1^2+x_2^2}dx_1 + \frac{x_1}{x_1^2+x_2^2}dx_2)$ be, as usual, the harmonic potential having flux ν around the origin. We assume $0 < \nu \leq \frac{1}{2}$, so that $\nu = d(\Phi^A, \mathbf{Z})$.

We already remarked that estimate (2) is sharp as $a \rightarrow 1$. We show that (5) is order sharp in $\nu = d(\Phi^A, \mathbf{Z})$ for all fixed $a \in (0, 1)$, in the sense that ν^2 is the correct rate of decay of λ_1 as $\nu \rightarrow 0$.

Theorem 2. *In the above notation we have, for all $0 < a < 1$ and $\nu \in (0, \frac{1}{2})$:*

$$\frac{1}{4} \leq \frac{\lambda_1(\Omega_a, A)}{\nu^2} \leq c(a),$$

where $c(a) = \frac{6|\log a|}{1-a^2}$.

- Note that $c(a) \rightarrow \infty$ as $a \rightarrow 1$: thus, $c(a)$ is far from being sharp (we know that in fact $\lambda_1(\Omega_a, A) \rightarrow \nu^2$ as $a \rightarrow 1$).

Proof. The lower bound follows immediately from (5). To prove the upper bound we use Rayleigh’s min-max principle:

$$\lambda_1 \doteq \lambda_1(\Omega_a, A) = \inf \left\{ \frac{\int_\Omega |\nabla^A u|^2}{\int_\Omega u^2} : u \in H_A^1(\Omega) \setminus \{0\} \right\}$$

applied to a real valued, radial function $u = u(r)$, where r is the distance to the origin. As $|A|^2 = \frac{\nu^2}{r^2}$ we see:

$$\begin{aligned} |\nabla^A u|^2 &= |\nabla u|^2 + u^2|A|^2 \\ &= u'^2 + \frac{\nu^2}{r^2}u^2 \end{aligned}$$

We now choose $u(r) = r^\nu$. In polar coordinates (r, θ) we see that $\Omega_a = (a, 1) \times [0, 2\pi)$ and upon integration we get easily:

$$\int_{\Omega} |\nabla^A u|^2 = 2\pi\nu(1 - a^{2\nu}), \quad \text{and} \quad \int_{\Omega} u^2 = 2\pi \cdot \frac{1 - a^{2\nu+2}}{2\nu + 2}$$

so that, by the min-max principle:

$$\frac{\lambda_1}{\nu^2} \leq \frac{2\nu + 2}{\nu} \cdot \frac{1 - a^{2\nu}}{1 - a^{2\nu+2}}.$$

Now $a^{2\nu} = e^{2\nu \log a} \geq 1 - 2\nu|\log a|$ so that $1 - a^{2\nu} \leq 2\nu|\log a|$; moreover $a^{2\nu+2} \leq a^2$ (because $a < 1$), and $\nu \leq \frac{1}{2}$. We end-up with the inequality:

$$\frac{\lambda_1}{\nu^2} \leq \frac{6|\log a|}{1 - a^2} = c(a),$$

as asserted. \square

Remark 3. The situation changes when considering the punctured unit disk with a Aharonov-Bohm singularity at its center. Calculations show in fact that $\lambda_1(\Omega, A)$ behaves asymptotically as $2d(\Phi^A, \mathbf{Z})$ as $d(\Phi^A, \mathbf{Z}) \rightarrow 0$.

1.4. A general lower bound for multiply connected domains

Now let Ω be an n -holed planar domain, which we write as follows:

$$\Omega = F \setminus (\bar{G}_1 \cup \dots \cup \bar{G}_n) \tag{6}$$

where the inner holes G_1, \dots, G_n are smooth, open and disjoint. We furthermore assume that F, G_1, \dots, G_n are convex. Note that:

$$\partial\Omega = \partial F \cup \partial G_1 \cup \dots \cup \partial G_n.$$

We will call $\partial G_1 \cup \dots \cup \partial G_n$ the *inner boundary* of Ω . The minimal and maximal widths of Ω are defined as in the case $n = 1$, namely β is the minimal length of a line segment contained in Ω and hitting the inner boundary orthogonally, and the maximal length of such line segments is by definition the maximal width B .

It is clear that we could replace $B(\Omega)$ by the diameter of F , and $\beta(\Omega)$ by the invariant:

$$\tilde{\beta}(\Omega) = \min\{d(\partial G_j, \partial G_k), d(\partial G_h, \partial F) : j \neq k, h = 1, \dots, n\}.$$

In this section we give a lower bound of $\lambda_1(\Omega, A)$ when Ω has an arbitrary number of convex holes.

Here is the estimate.

Theorem 4. *Let $\Omega = F \setminus (\bar{G}_1 \cup \dots \cup \bar{G}_n)$ be an n -holed planar domain, where F, G_1, \dots, G_n are smooth, open and convex. Let A be a closed potential having flux Φ_j around the j -th inner boundary curve ∂G_j , for $j = 1, \dots, n$, and let $\gamma = \min_{j=1, \dots, n} d(\Phi_j, \mathbf{Z})$. Then we have:*

$$\lambda_1(\Omega, A) \geq \frac{\pi^2}{2(|\partial F| + 2\pi B(\Omega))^2} \frac{\beta(\Omega)^4}{B(\Omega)^4} \cdot \gamma^2, \tag{7}$$

where $\beta(\Omega)$ and $B(\Omega)$ are, respectively, the minimal and maximal width of Ω .

The proof uses a suitable decomposition of Ω into a finite union of annuli, and a lower bound proved in [3] for annuli whose outer boundary is star-shaped with respect to the inner boundary curve. A stronger estimate is proved when the inner holes are disks of the same radius (see Theorem 15).

1.5. *A lower bound for Aharonov-Bohm operators with many poles*

The power $\frac{\beta^4}{B^4}$ in the previous estimate is probably not sharp; it appears to be there for technical reasons. By shrinking the inner boundary curves to points we obtain an estimate in terms of $\frac{\beta^2}{B^2}$, which has an interesting interpretation in terms of Aharonov-Bohm operators with many poles.

Precisely, we fix a convex domain Ω and choose n points inside it, say $\mathcal{P} = \{p_1, \dots, p_n\}$. Consider the punctured domain $\Omega \setminus \mathcal{P}$. Given a closed one-form A , we define:

$$\lambda_1(\Omega \setminus \mathcal{P}, A) = \liminf_{\delta \rightarrow 0} \lambda_1(\Omega \setminus \mathcal{P}(\delta), A)$$

where $\mathcal{P}(\delta)$ is the δ -neighborhood of \mathcal{P} (it obviously consists of a finite set of disks of radius δ). It is not our scope in this paper to investigate the convergence in terms of δ ; however, what we are looking at could be interpreted as the first eigenvalue of a Aharonov-Bohm operator with poles p_1, \dots, p_n and Neumann boundary conditions. The proof of the theorem in the previous section simplifies, to give a general lower bound in terms of the distance between the poles, and the distance of each pole to the boundary. To that end, define:

$$\begin{cases} \beta(\mathcal{P}) = \min\{d(p_j, p_k), d(p_m, \partial\Omega) : p_j \neq p_k, p_m \in \mathcal{P}\} \\ B(\mathcal{P}) = \max\{d(p_j, p_k), d(p_m, \partial\Omega) : p_j \neq p_k, p_m \in \mathcal{P}\} \end{cases}$$

Of course $B(\mathcal{P})$ could be conveniently bounded above by the diameter of Ω . Let A be as usual a closed one-form having flux Φ_j around the pole p_j . Then we have:

Theorem 5. *Let Ω be a convex domain and $\mathcal{P} = \{p_1, \dots, p_n\}$ a finite set of poles. For the punctured domain $\Omega \setminus \mathcal{P}$ we have the bound:*

$$\lambda_1(\Omega \setminus \mathcal{P}, A) \geq \frac{4\pi^2}{|\partial\Omega|^2} \frac{\beta(\mathcal{P})^2}{B(\mathcal{P})^2} \gamma^2,$$

where $\gamma = \min_{j=1, \dots, n} d(\Phi_j, \mathbf{Z})$, and Φ_j is the flux of the closed potential A around p_j .

- In a forthcoming paper, we will give upper bounds for the Laplacian with zero magnetic field on multiply connected planar domains, which are closely related to the topology (number of holes) of the domain.

The rest of the paper is devoted to the proof of Theorems 1, 4 and 5.

2. Proof of Theorem 1

The proof depends on a suitable way to partition our domain Ω . We first remark the simple fact that the first eigenvalue of a domain is controlled from below by the smallest first eigenvalue of the subdomains of a partition of Ω (Proposition 6). Then, we need to extend inequality (2) to piecewise-smooth boundaries, see Section 2.2. In Section 2.3 we state our main geometric facts, Lemma 8 and Lemma 9, and then we prove Theorem 1 (see Section 2.4). Finally, in Section 2.5, we define the partition and we prove Lemma 8 and Lemma 9.

2.1. A simple lemma

We say that the family of open subdomains $\{\Omega_1, \dots, \Omega_n\}$ is a *partition* of Ω , if $\bar{\Omega} = \bar{\Omega}_1 \cup \dots \cup \bar{\Omega}_n$. Thus, the members of the partition might overlap and some of the intersections $\Omega_j \cap \Omega_k$ could have positive measure. If furthermore $\Omega_j \cap \Omega_k$ is empty for all $j \neq k$ then we say that the partition is *disjoint*. We observe the following standard fact whose proof is easy:

Proposition 6. *Let $\{\Omega_1, \dots, \Omega_n\}$ be a partition of the domain Ω . Let A be any closed potential. Then, there is an index $k = 1, \dots, n$ such that*

$$\lambda_1(\Omega, A) \geq \frac{1}{n} \lambda_1(\Omega_k, A). \tag{8}$$

If the partition is disjoint, then:

$$\lambda_1(\Omega, A) \geq \min_{j=1, \dots, n} \lambda_1(\Omega_j, A). \tag{9}$$

Proof. We start proving (8). Let u be an eigenfunction associated to $\lambda_1(\Omega, A)$. We use it as test-function for $\lambda_1(\Omega_j, A)$ and obtain, for all j :

$$\lambda_1(\Omega_j, A) \int_{\Omega_j} |u|^2 \leq \int_{\Omega_j} |\nabla^A u|^2. \tag{10}$$

Now

$$\int_{\Omega} |u|^2 \leq \sum_{j=1}^n \int_{\Omega_j} |u|^2 \leq n \int_{\Omega_k} |u|^2$$

where the index k is chosen so that $\int_{\Omega_k} |u|^2$ is maximum among all $j = 1, \dots, n$. Then:

$$\begin{aligned} \lambda_1(\Omega_k, A) \int_{\Omega} |u|^2 &\leq n \lambda_1(\Omega_k, A) \int_{\Omega_k} |u|^2 \\ &\leq \int_{\Omega_k} |\nabla^A u|^2 \\ &\leq n \int_{\Omega} |\nabla^A u|^2 \\ &= n \lambda_1(\Omega, A) \int_{\Omega} |u|^2 \end{aligned}$$

That is: $\lambda_1(\Omega_k, A) \leq n \lambda_1(\Omega, A)$, which is the assertion.

For the proof of (9), let $\lambda_{\min} = \min_{j=1, \dots, n} \lambda_1(\Omega_j, A)$. From (10) we have, for all j :

$$\int_{\Omega_j} |\nabla^A u|^2 \geq \lambda_1(\Omega_j, A) \int_{\Omega_j} |u|^2 \geq \lambda_{\min} \int_{\Omega_j} |u|^2$$

We now sum over $j = 1, \dots, n$ and obtain $\int_{\Omega} |\nabla^A u|^2 \geq \lambda_{\min} \int_{\Omega} |u|^2$. As u is a first eigenfunction the left-hand side is precisely $\lambda_1(\Omega, A) \int_{\Omega} |u|^2$, and the inequality follows. \square

2.2. Convex annuli with piecewise-smooth boundary

From now on Ω will be an annulus in the plane with boundary components $\Gamma_{\text{int}}, \Gamma_{\text{ext}}$ which we assume convex and piecewise-smooth. We will write $\Omega = F \setminus \bar{G}$ where F and G are open, convex, with piecewise smooth boundary. In that case $\Gamma_{\text{int}} = \partial G$ and $\Gamma_{\text{ext}} = \partial F$.

Let p be a point of ∂G where ∂G is not smooth (p will then be called a *vertex*). The *normal cone* of G at p is the set

$$N_G(p) = \{x \in \mathbf{R}^2 : \langle x, y - p \rangle \leq 0, \text{ for all } y \in G\}.$$

Then $N_G(p)$ is the closed exterior wedge bounded by the normal lines to the two smooth curves concurring at p , its boundary is the broken line depicted in the figure below. Call α_p its angle at p .

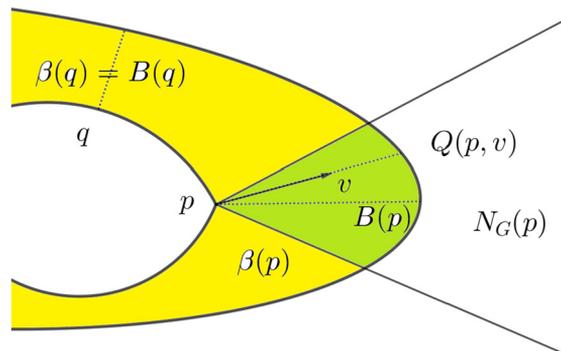


Fig. 3. A vertex p of ∂G and its normal cone $N_G(p)$.

- We remark the obvious fact that $0 < \alpha_p < \pi$.

We now define the minimum and maximum width in the piecewise-smooth case. These are defined in (11) and depicted in the Fig. 3.

For a unit vector v applied in p and pointing inside $N_G(p)$ we let $\gamma_{p,v}(t) = p + tv$ denote the ray exiting p in the direction v , and let $Q(p, v)$ be the intersection of $\gamma_{p,v}$ with $\Gamma_{\text{ext}} = \partial F$. We define:

$$\begin{cases} \beta(p) = \inf_{v \in N_G(p)} d(p, Q(p, v)) \\ B(p) = \sup_{v \in N_G(p)} d(p, Q(p, v)) \end{cases} \tag{11}$$

We notice that at a smooth point q the cone at q degenerates to the normal segment at q . Hence at a smooth point q one has

$$\beta(q) = B(q).$$

We now define

$$\begin{cases} \beta(\Omega) = \inf\{\beta(p) : p \in \partial G\} \\ B(\Omega) = \sup\{\beta(p) : p \in \partial G\} \end{cases} \tag{12}$$

$\beta(\Omega)$ and $B(\Omega)$ will be called the *minimum width* and, respectively, the *maximum width* of Ω . We remark that when the two boundary components are smooth and parallel then $\beta = B$ and the ratio $\frac{\beta}{B}$ assumes its largest possible value, which is 1.

As a first step in the proof of Theorem 1, we extend the inequality (2) to the piecewise-smooth case.

Theorem 7. *Let $\Omega = F \setminus \bar{G}$ be an annulus in the plane whose boundary components are convex and piecewise smooth. Let $\beta = \beta(\Omega)$ and $B = B(\Omega)$ be the invariants defined in*

(12). Then for any closed potential having flux Φ^A around the inner boundary curve one has the lower bound:

$$\lambda_1(\Omega, A) \geq \frac{4\pi^2}{|\partial F|^2} \frac{\beta(\Omega)^2}{B(\Omega)^2} d(\Phi^A, \mathbf{Z})^2,$$

where $|\partial F|$ is the length of the outer boundary.

Proof. First, Ω admits an exhaustion by convex annuli with C^1 -boundary, say $\{\Omega_\epsilon : \epsilon > 0\}$. By that we mean:

- a) $\Omega_\epsilon = F_\epsilon \setminus \bar{G}_\epsilon$ where F_ϵ and G_ϵ are convex and have C^1 -smooth boundary;
- b) $F_\epsilon \subseteq F$ and $G_\epsilon \supseteq G$ so that $\Omega_\epsilon \subseteq \Omega$;
- c) $\Omega = \cup_{\epsilon>0} \Omega_\epsilon$ and in particular $\lim_{\epsilon \rightarrow 0} |\Omega \setminus \Omega_\epsilon| = 0$.

To construct F_ϵ we round off corners at distance ϵ to each of the vertices of ∂F ; to construct G_ϵ we just take the convex domain bounded by the ϵ -neighborhood of G .

Let u be an eigenfunction associated to $\lambda_1(\Omega, A)$; by restriction we obtain a test-function for Ω_ϵ , hence by the min-max principle:

$$\frac{\int_{\Omega_\epsilon} |\nabla^A u|^2}{\int_{\Omega_\epsilon} |u|^2} \geq \lambda_1(\Omega_\epsilon, A).$$

Let $\mathcal{L}(\Omega)$ be the functional:

$$\mathcal{L}(\Omega) = \frac{4\pi^2}{|\partial F|^2} \frac{\beta(\Omega)^2}{B(\Omega)^2} d(\Phi^A, \mathbf{Z})^2$$

We can apply (2) and obtain $\lambda_1(\Omega_\epsilon) \geq \mathcal{L}(\Omega_\epsilon)$ because Ω_ϵ has smooth boundary; then, for all $\epsilon > 0$:

$$\frac{\int_{\Omega_\epsilon} |\nabla^A u|^2}{\int_{\Omega_\epsilon} |u|^2} \geq \mathcal{L}(\Omega_\epsilon).$$

We now pass to the limit as $\epsilon \rightarrow 0$ on both sides; as $\mathcal{L}(\Omega_\epsilon) \rightarrow \mathcal{L}(\Omega)$ (as we can easily see from the definitions in (12)), we obtain the assertion: $\lambda_1(\Omega) \geq \mathcal{L}(\Omega)$. \square

2.3. Preparatory results

In this section we state the two main technical lemmas; the partition of the annulus Ω will be defined in Section 2.5.

So let $\Omega = F \setminus \bar{G}$ be an annulus as above and let $\beta \in (0, \beta(\Omega)]$. We consider the distance functions:

$$\rho_1, \rho_2 : F \rightarrow [0, \infty),$$

where $\rho_1(x) = d(x, G)$ and $\rho_2(x) = d(x, \partial F)$. Fix a parameter $\beta > 0$. As G is convex, with piecewise-smooth boundary, it is well-known that the equidistants $\{\rho_1 = \beta\}$ are C^1 -smooth curves. We say that the parameter β is *regular* if the equidistant $\{\rho_2 = \beta\}$ is a piecewise-smooth curve. Following Appendix 2 in [13], we know that the set of regular parameters has full measure in $(0, \beta(\Omega))$; as a consequence, there exists a sequence of regular parameters $\{\beta_j\} \rightarrow \beta(\Omega)$ as $j \rightarrow \infty$.

- By using an obvious limiting procedure, from now on we take

$$\beta = \beta(\Omega)$$

and can assume that it is a regular parameter, so that $\rho_2 = \beta$ is a piecewise-smooth curve.

Lemma 8. *Let $\Omega = F \setminus \bar{G}$ be an annulus in the plane with F and G convex with piecewise-smooth boundary, and let $\beta = \beta(\Omega)$ (which, by assumption, is a regular parameter). Then Ω admits a partition $\{\Omega_1, \dots, \Omega_n\}$ into (overlapping) subdomains Ω_k with the following properties.*

- a) Ω_k is an annulus bounded by two convex piecewise-smooth curves, that is, $\Omega = F_k \setminus \bar{G}_k$ with F_k and G_k convex, and G_k contains G (see figure in Section 2.5 below).
- b) The number n of annuli in the partition can be taken so that:

$$n \leq \frac{2B(\Omega)}{\beta}.$$

We estimate the ratio $\frac{\beta}{B}$ of each piece as follows.

Lemma 9. *Let $\{\Omega_1, \dots, \Omega_n\}$ be the partition in the previous lemma. For all $k = 1, \dots, n$ one has the following facts.*

- a) $|\partial F_k| \leq |\partial F|$ and $\beta(\Omega_k) = \beta$.
- b) The following estimate holds:

$$\frac{\beta(\Omega_k)}{B(\Omega_k)} \geq \frac{1}{4} \frac{|F|}{D(F)^2}, \tag{13}$$

where $D(F)$ is the diameter of F .

- c) If $\beta(\Omega)$ is less than the injectivity radius of ∂F , then the following simpler lower bound holds for all k :

$$\frac{\beta(\Omega_k)}{B(\Omega_k)} \geq \frac{1}{\sqrt{2}}.$$

The proof of Lemma 8 and Lemma 9 involves rather simple geometric constructions, but there are some delicate points to take care of, and will be done in Section 2.5. In

fact, these two lemmas make it possible to write Ω as a union of subset Ω_k such that the ratio $\frac{\beta(\Omega_k)}{B(\Omega_k)}$ is bounded below, which make the proof of Theorem 1 quite easy, as follows.

2.4. Proof of Theorem 1

We use the partition $\{\Omega_1, \dots, \Omega_n\}$ of Lemma 8. Let A be a closed potential having flux Φ^A around the inner boundary curve ∂G ; then, A has the same flux around the inner component of Ω_k , by Lemma 8a. By Proposition 6a there exists $k \in \{1, \dots, n\}$ such that

$$\lambda_1(\Omega, A) \geq \frac{1}{n} \lambda_1(\Omega_k, A). \tag{14}$$

By Theorem 7 applied to $\Omega = \Omega_k$ we see:

$$\lambda_1(\Omega_k, A) \geq \frac{4\pi^2}{|\partial F_k|^2} \frac{\beta(\Omega_k)^2}{B(\Omega_k)^2} d(\Phi^A, \mathbf{Z})^2.$$

By b) of Lemma 9 we see:

$$\frac{\beta(\Omega_k)^2}{B(\Omega_k)^2} \geq \frac{1}{16} \frac{|F|^2}{D(F)^4}. \tag{15}$$

This, together with the inequality $|\partial F_k| \leq |\partial F|$, gives:

$$\lambda_1(\Omega_k, A) \geq \frac{\pi^2}{4} \cdot \frac{|F|^2}{|\partial F|^2 D(F)^4} \cdot d(\Phi^A, \mathbf{Z})^2.$$

We insert this inequality in (14) and use the inequality $\frac{1}{n} \geq \frac{\beta}{2B(\Omega)}$ (see Lemma 8b) to conclude that

$$\begin{aligned} \lambda_1(\Omega, A) &\geq \frac{\beta}{2B(\Omega)} \lambda_1(\Omega_k, A) \\ &\geq \frac{\pi^2}{8} \cdot \frac{|F|^2}{|\partial F|^2 D(F)^4} \cdot \frac{\beta}{B(\Omega)} \cdot d(\Phi^A, \mathbf{Z})^2. \end{aligned}$$

This proves part a) of Theorem 1.

If $\beta(\Omega)$ is less than the injectivity radius of ∂F we proceed as before, using the lower bound $\frac{\beta(\Omega_k)}{B(\Omega_k)} \geq \frac{1}{\sqrt{2}}$ proved in Lemma 9c. We arrive easily at the inequality:

$$\lambda_1(\Omega, A) \geq \frac{\pi^2}{|\partial F|^2} \frac{\beta(\Omega)}{B(\Omega)} d(\Phi^A, \mathbf{Z})^2,$$

which is Theorem 1b).

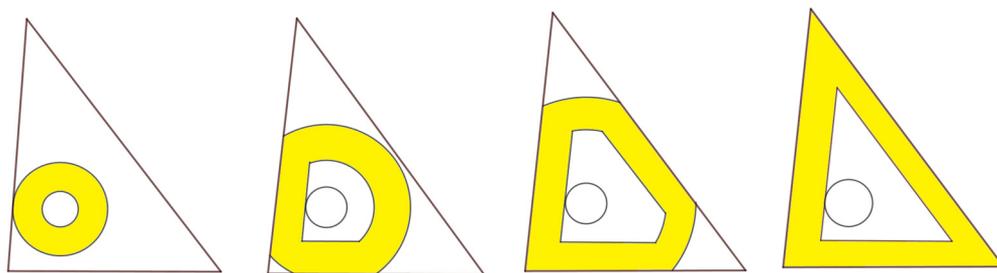


Fig. 4. The pieces $\Omega_1, \Omega_2, \Omega_3$ and the last piece $\Omega_6 = \{\rho_2 < \beta\}$ when the initial domain is the triangle minus the small disk and $\beta = \beta(\Omega)$.

2.5. The partition of Ω and the proof of Lemma 8

We start by showing the partition on a particular example, see Fig. 4. The initial domain is a triangle F minus a disk G and $\beta = \beta(\Omega)$. We draw the first three pieces $\Omega_1, \Omega_2, \Omega_3$ and then the last one, which is Ω_6 and which coincides with the β -neighborhood of the exterior boundary ∂F (this is always the case). Note that the pieces overlap, hence the partition is not disjoint.

We now proceed to construct the partition in general. Let then Ω be convex annulus $\Omega = F \setminus \bar{G}$ as above, and consider the distance functions:

$$\rho_1, \rho_2 : F \rightarrow [0, \infty),$$

where $\rho_1(x) = d(x, G)$ and $\rho_2(x) = d(x, \partial F) = d(x, F^c)$.

- At step 1, we let $F_1 = \{\rho_1 < \beta\}$, $G_1 = G$ and $\Omega_1 = F_1 \setminus \bar{G}_1$. That is, Ω_1 is simply the subset of F at distance less than β to G .
- At step 2, we let $F_2 = \{\rho_1 < 2\beta\}$ and $G_2 = \{\rho_1 < \beta\} \cap \{\rho_2 > \beta\}$ and define $\Omega_2 = F_2 \setminus \bar{G}_2$.
- At the arbitrary step k , we let $F_k = \{\rho_1 < k\beta\}$ and $G_k = \{\rho_1 < (k-1)\beta\} \cap \{\rho_2 > \beta\}$, and define:

$$\Omega_k = F_k \setminus \bar{G}_k.$$

Observe that for any choice of positive numbers a, b the sets $\{\rho_1 < a\}$ and $\{\rho_2 < b\}$ are convex. Therefore, both F_k and G_k are convex; moreover $G_k \subset F_k$ and then Ω_k is an annulus. This proves part a) of Lemma 8.

For b) we first prove the following fact:

Fact. Let $n \doteq \left\lceil \frac{B(\Omega)}{\beta} \right\rceil$ (the smallest integer greater than or equal to $\frac{B(\Omega)}{\beta}$). Then $F_n = F$ and $G_n = \{\rho_2 > \beta\}$. In particular, $\Omega_n = \{\rho_2 < \beta\}$ and then, starting from n , the sequence Ω_n stabilizes: $\Omega_n = \Omega_{n+1} = \dots$

For the proof we first observe that, from the definition of $B(\Omega)$, we have $F \subseteq \{\rho_1 < B(\Omega)\}$; then, if we fix $n \geq \frac{B(\Omega)}{\beta}$ we have by definition $F \subseteq F_n$ hence $F = F_n$. To show that $G_n = \{\rho_2 > \beta\}$ it is enough to show:

$$\{\rho_2 > \beta\} \subseteq \{\rho_1 < (n - 1)\beta\}. \tag{16}$$

In fact, if not, there would be a point $x \in F$ such that $d(x, \partial F) = \rho_2(x) > \beta$ and $\rho_1(x) \geq (n - 1)\beta$. Let $y \in \partial G$ be a point at minimum distance to x , and prolong the segment from y to x till it hits $\Gamma_{\text{ext}} = \partial F$ at the point z . It is clear that then $d(y, z) = d(y, x) + d(x, z) > n\beta$. By definition of $B(\Omega)$ we then have:

$$B(\Omega) \geq d(y, z) > n\beta,$$

which contradicts the definition of n . Hence (16) holds.

We now prove part b) of Lemma 8. Observe that $\bar{\Omega} = \cup_{k=1}^n \bar{\Omega}_k$ and (by the definition of n) $\frac{B(\Omega)}{\beta} \geq n - 1$. Since $n - 1 \geq \frac{n}{2}$ for all $n \geq 2$ we see that, for all n :

$$\frac{B(\Omega)}{\beta} \geq \frac{n}{2},$$

which gives the assertion.

2.6. Proof of Lemma 9

We now study the typical piece $\Omega_k = F_k \setminus \bar{G}_k$ in the partition. Observe that

$$\partial F_k = \partial_1 F_k \cup \partial_2 F_k, \quad \partial G_k = \partial_1 G_k \cup \partial_2 G_k,$$

where

$$\begin{cases} \partial_1 F_k = \{\rho_1 = k\beta\} \cap \bar{F} \\ \partial_2 F_k = \{\rho_1 \leq k\beta\} \cap \partial F \end{cases} \quad \begin{cases} \partial_1 G_k = \{\rho_1 = (k - 1)\beta\} \cap \{\rho_2 \geq \beta\} \\ \partial_2 G_k = \{\rho_1 \leq (k - 1)\beta\} \cap \{\rho_2 = \beta\} \end{cases}$$

(some of these boundary pieces may be empty). As the equidistants $\{\rho_1 = r\}$ are C^1 -smooth, and the equidistant $\{\rho_2 = \beta\}$ is piecewise smooth, we see that ∂F_k and ∂G_k are both piecewise smooth, hence

- $\partial\Omega_k$ is piecewise smooth.

The inner boundary is ∂G_k , it is piecewise smooth with vertices in the set

$$S = \{\rho_1 = (k - 1)\beta\} \cap \{\rho_2 = \beta\}.$$

Now we have to estimate the ratio $\beta(\Omega_k)/B(\Omega_k)$ for a fixed $k = 1, \dots, n$. Recall the function

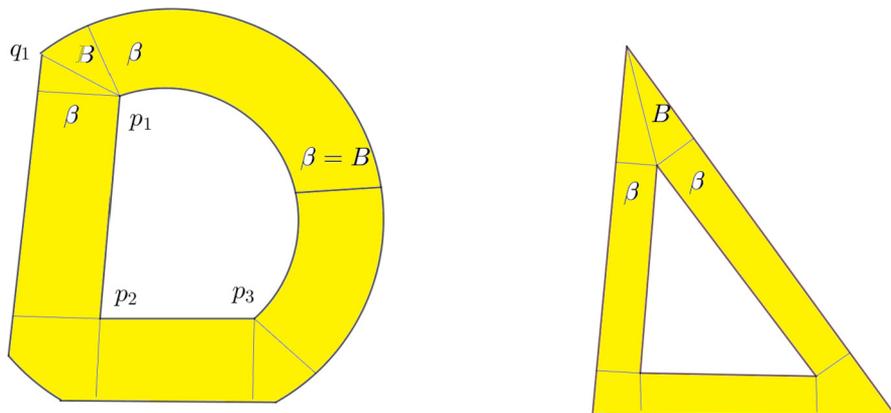


Fig. 5. On the left: the piece Ω_2 and its wedges. On the right: the last piece $\Omega_6 = \{\rho_2 < \beta\}$; the ratio $\frac{\beta}{B}$ is small at the upper wedge, because the angle of the wedge at its break point is near π .

$$\beta(x) : \partial G_k \rightarrow \mathbf{R}$$

defined in (11). First notice that the regular parts

$$\partial_{1,\text{reg}} F_k = \{\rho_1 = k\beta\} \cap F \quad \text{and} \quad \partial_{1,\text{reg}} G_k = \{\rho_1 = (k - 1)\beta\} \cap \{\rho_2 > \beta\}$$

are parallel, at distance β to each other. Hence

$$\frac{\beta(x)}{B(x)} = 1 \tag{17}$$

at all points $x \in \partial_{1,\text{reg}} G_k$. Similarly, the regular sets

$$\partial_{2,\text{reg}} F_k = \{\rho_1 < k\beta\} \cap \partial F \quad \text{and} \quad \partial_{2,\text{reg}} G_k = \{\rho_1 < (k - 1)\beta\} \cap \{\rho_2 = \beta\}$$

are parallel at distance β and $\frac{\beta(x)}{B(x)} = 1$ on $\partial_{2,\text{reg}} G_k$.

Therefore, it only remains to control the ratio $\frac{\beta(x)}{B(x)}$ at the vertices of ∂G_k , which are finite, say p_1, \dots, p_m .

Each break point p_j gives rise to a corresponding wedge $W(p_j) \doteq N_G(p_j) \cap \Omega$. In Fig. 5 we enlarge the domain Ω_2 relative to the partition of Fig. 4 and we show its set of wedges. In other words, every annulus Ω_k is made up of strips of constant width β and wedges, and we need to control $\frac{\beta(x)}{B(x)}$ only at the wedges.

Typically, the ratio $\frac{\beta}{B}$ is small when there are small angles; nevertheless, this ratio is controlled from below by the diameter and the volume of the outer domain, as we will see in the next section.

As F_k is a convex subset of F , we see that $|\partial F_k| \leq |\partial F|$. Now it is clear from the construction that $\beta(x) \geq \beta$ for all $x \in \partial F_k$; moreover, the inequality is attained. Therefore

$$\beta(\Omega_k) = \beta$$

for all k . This proves part a) of Lemma 9.

2.7. End of proof of Lemma 9

The estimate $\frac{\beta}{B}$ on the wedges of the generic piece Ω_k will be a consequence of Lemma 10 below.

We recall that the cut-locus of ∂F is the closure of the set of all points which can be joined to ∂F by at least two minimizing segments; moreover, the injectivity radius of ∂F is the minimum distance of ∂F to the cut-locus. If ∂F is smooth, its injectivity radius is positive. Finally the distance function $d(\cdot, \partial F)$ is smooth outside the cut-locus.

Then, we fix a piece Ω_k and recall that $\beta(\Omega_k) = \beta$. For simplicity we write $B(\Omega_k) = B$ and recall that, by definition, we have $B \geq \beta$.

Let $\{p_1, \dots, p_n\}$ be the vertices of ∂G_k . For p in this set, write $p = \gamma_1 \cap \gamma_2$, where γ_1, γ_2 are the arcs concurring at p . Note that either γ_j is an equidistant to ∂G , that is, is a subset of $\rho_1 = (k - 1)\beta$, and in that case we say that γ_j is *parallel to ∂G* , or γ_j is an equidistant to ∂F , that is, is a subset of $\rho_2 = \beta$, and in that case we say that γ_j is *parallel to ∂F* . There are two possibilities:

Type 1. The vertex $p = \gamma_1 \cap \gamma_2$ where γ_1 is parallel to ∂G and γ_2 is parallel to ∂F ;

Type 2. γ_1 and γ_2 are both parallel to ∂F .

Note that the second type corresponds to the situation where the vertex p is a point of the cut-locus of ∂F . The situation where γ_1 and γ_2 are both parallel to ∂G does not occur, because then p would belong to the cut locus of ∂G ; however the cut-locus of a convex domain is always contained in the interior of the domain; as p is outside G this is impossible.

- For the partition in the example and its piece Ω_2 (see Fig. 5), the vertices p_1 and p_3 are of type 1, while the vertex p_2 is of type 2.

Lemma 10. a) If p is of type 1 then the interior angle of G_k at p is larger than or equal to $\frac{\pi}{2}$, hence the angle of the wedge $W(p)$ at p is at most $\frac{\pi}{2}$. Consequently,

$$\frac{\beta}{B(p)} \geq \frac{1}{\sqrt{2}}.$$

b) If p is of type 2, then p is in the cut-locus of ∂F and one has:

$$\frac{\beta}{B(p)} \geq \frac{1}{4} \frac{|F|}{D(F)^2}.$$

c) If $\beta = \beta(\Omega)$ is less than the injectivity radius of ∂F then the estimate in a) will hold at all vertices of the decomposition.

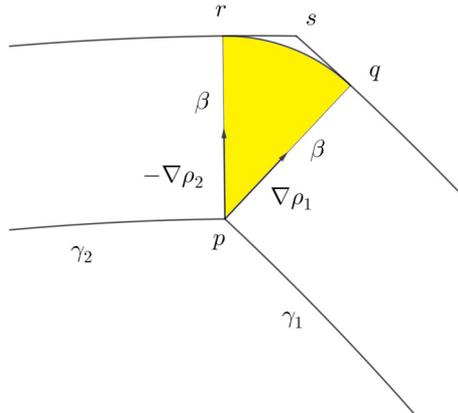


Fig. 6. Proof of Lemma 10a.

Since the lower bound in b) is always weaker than that in a), we have b) at all vertices of ∂G_k . It is clear that Lemma 10 completes the proof of Lemma 9.

Proof of Lemma 10a). If p is of type 1 then p is not on the cut-locus of ∂F , hence $\nabla \rho_2$ exists and is a well-defined unit vector in a neighborhood of p . Note that $\nabla \rho_2(p)$ points in the direction where the distance to ∂F increases (obviously an analogous observation holds for $\nabla \rho_1$). Now observe that the angle of the wedge $W(p)$ is the angle between the vectors $\nabla \rho_1$ and $-\nabla \rho_2$ (see Fig. 6).

Hence, it is enough to show that the quantity

$$\langle \nabla \rho_1(p), \nabla \rho_2(p) \rangle = c(p)$$

is non-positive. Assume on the contrary that $c(p) > 0$. We let $\alpha(t)$ denote the segment which minimizes distance from p to ∂G (parametrized by arc-length); hence $\alpha'(t) = -\nabla \rho_1(\alpha(t))$. We let $f(t)$ be the function which measures distance from $\alpha(t)$ to ∂F , so that:

$$f(t) = \rho_2(\alpha(t)).$$

Now $f'(t) = \langle \nabla \rho_2(\alpha(t)), \alpha'(t) \rangle = -\langle \nabla \rho_2(\alpha(t)), \nabla \rho_1(\alpha(t)) \rangle$. In particular,

$$f'(0) = -c(p) < 0.$$

As $f(0) = \beta$, this means that for small t one has $\rho_2(\alpha(t)) < \beta$, but this impossible because $\alpha(t) \in G_k$, and all points of G_k are, by definition, at distance at least β to ∂F .

Hence $c_p \leq 0$ and the angle of the wedge at p is at most $\frac{\pi}{2}$. Now, the wedge $W(p)$ is contained in the polygon with vertices p, q, s, r as in the picture, hence $B(p) \leq d(p, s) \leq \sqrt{2}\beta$ because the angle at p is at most $\frac{\pi}{2}$, the angles at r and s are $\frac{\pi}{2}$, and $d(p, r) = d(p, q) = \beta$.

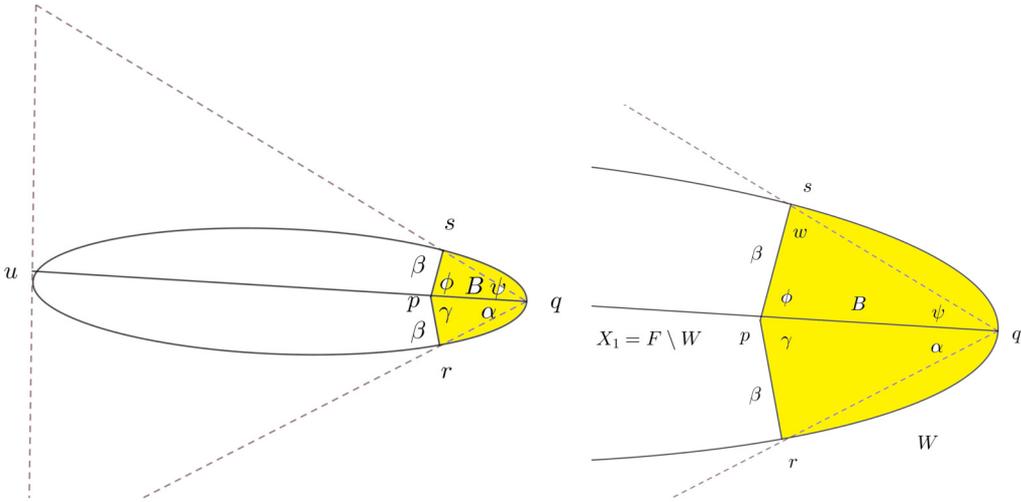


Fig. 7. Estimate of β/B on a typical wedge.

Proof of b). Let p be a vertex of type 2: then, the two arcs concurring at p are parallel to ∂F , and p belongs to the cut locus of ∂F . The boundary of the wedge $W(p)$ is made of two distinct segments of the same length β minimizing distance to ∂F . Then, $W(p)$ is contained in a wedge of the last member of the partition, that is, $\{\rho_2 < \beta\}$. As $B(p)$ depends only on $W(p)$, we could as well estimate the ratio $\frac{\beta}{B(p)}$ by estimating the corresponding ratio for the wedges of $\{\rho_2 < \beta\}$, which will allow to express $\frac{\beta}{B(p)}$ in terms of the geometry of $\{\rho_2 < \beta\}$, hence, in terms of the geometry of F .

The relevant picture is Fig. 7, in which we evidence such an edge W (dark shadowed in the picture): it has its vertex in the point p of $\rho_2 = \beta$; we let $q \in W$ be a point such that $d(p, q) = B$. We omit to draw the inner boundary as it will play no role in the proof.

Let T be the triangle with dotted boundary, with a vertex in q and such that $F \setminus W \subseteq T$. As $\phi + \gamma$ is the exterior angle at a vertex of the piecewise-smooth curve $\rho_2 = \beta$, we see that $\phi + \gamma \leq \pi$.

Each of the angles ψ and α is less than $\frac{\pi}{2}$. Consider the circle with center p and radius β . If q is inside this circle then $B < \beta$ which is impossible. Then q is outside the circle; α and ψ are, each, less than the corresponding angles at the vertex q' obtained by projecting q on the circle. It is clear that each of these two angles is less than $\pi/2$.

Let w be the angle at the vertex s . Then

$$\frac{\beta}{\sin \psi} = \frac{B}{\sin w} \geq B, \tag{18}$$

and similarly $\frac{\beta}{\sin \alpha} \geq B$. If $\psi \geq \frac{\pi}{4}$ or $\alpha \geq \frac{\pi}{4}$ then $\frac{\beta}{B} \geq \frac{1}{\sqrt{2}}$ and we are finished because $\frac{1}{\sqrt{2}} > \frac{|F|}{4D(F)^2}$.

- Hence we can assume from now on $\alpha \leq \psi \leq \frac{\pi}{4}$.

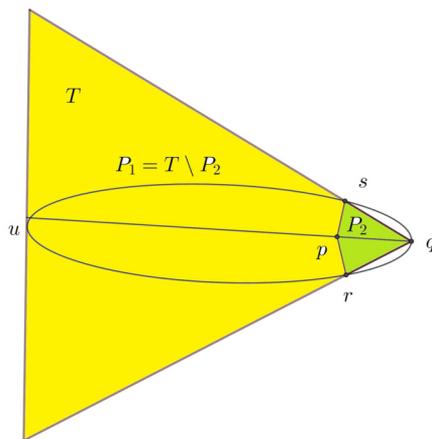


Fig. 8. The polygons P_1 and P_2 .

Lemma 11. *In the above notation we have $|T| \geq \frac{\sqrt{2}}{4}|F|$.*

Proof. We first remark that $\frac{\pi}{4} \leq \phi \leq \frac{3\pi}{4}$. In fact, we have $\psi \leq \frac{\pi}{4}$ and $w \leq \frac{\pi}{2}$ so that $\phi \geq \frac{\pi}{4}$. On the other hand, the same argument applies to γ , that is, $\gamma \geq \frac{\pi}{4}$. Therefore, as $\phi + \gamma \leq \pi$ we conclude $\phi \leq \frac{3\pi}{4}$. The same bounds are satisfied by γ .

As W is contained in the union of two parallelograms with sides β and B we see:

$$|W| \leq 2B\beta. \tag{19}$$

We set $X_1 = F \setminus W$; we also let P_2 be the convex polygon with vertices p, q, r, s and $P_1 = T \setminus P_2$ (see Fig. 8).

We have disjoint unions:

$$\begin{cases} F = X_1 \cup W \\ T = P_1 \cup P_2 \end{cases}$$

We will show that

$$\begin{cases} |P_1| \geq |X_1| \geq \frac{\sqrt{2}}{4}|X_1| \\ |P_2| \geq \frac{\sqrt{2}}{4}|W| \end{cases}$$

and the assertion will follow by summing up the two inequalities. Now the first inequality is obvious, because $X_1 \subseteq P_1$. By the bounds for ϕ and γ we see that $\sin \phi$ and $\sin \gamma$ are both, at least, $\frac{1}{\sqrt{2}}$. Then:

$$\begin{aligned}
 |P_2| &= \frac{1}{2}B\beta \sin \phi + \frac{1}{2}B\beta \sin \gamma \\
 &\geq \frac{\sqrt{2}}{2}B\beta
 \end{aligned}$$

Combining the two estimates we see that $|P_2| \geq \frac{\sqrt{2}}{4}|W|$ as asserted. \square

End of proof of Lemma 9. Refer to Fig. 8. We can assume that $\alpha \leq \psi \leq \frac{\pi}{4}$. We let δ be the length of the segment joining q and u (which meets the side opposite to q orthogonally, by definition), so that:

$$|T| = \frac{1}{2}\delta^2(\tan \alpha + \tan \psi) \leq \delta^2 \tan \psi.$$

The assumptions give $\tan \psi \leq \sqrt{2} \sin \psi$, so that $|T| \leq \sqrt{2}\delta^2 \sin \psi$. Using the lower bound for $|T| \geq \frac{\sqrt{2}}{4}|F|$ proved before, we have $\sin \psi \geq \frac{|F|}{4\delta^2}$ and then, from (18):

$$\frac{\beta}{B} \geq \sin \psi \geq \frac{|F|}{4\delta^2} \geq \frac{|F|}{4D(F)^2},$$

the last inequality holding because evidently $\delta \leq D(F)$. This proves Lemma 10b and, with it, Lemma 9 is completely proved.

2.8. Example showing sharpness

This example is taken from [3], we repeat it below for the sake of clarity. Its scope is to show that the inequality of Theorem 1 is sharp in $\frac{\beta}{B}$.

We take F to be the rectangle $[-4, 4] \times [0, 4]$, $G_\epsilon = [-3, 3] \times [\epsilon, 2]$ and consider the doubly connected domain:

$$\Omega_\epsilon = F \setminus \bar{G}_\epsilon.$$

We refer to the picture in the Introduction. We let A be any closed 1-form. As a direct consequence of the gauge invariance of the magnetic Laplacian, it is proved in [3] that, for any planar domain Ω one has:

$$\lambda_1(\Omega, A) \leq \nu_1(D), \tag{20}$$

where D is any closed, simply connected subdomain of $\bar{\Omega}$, and where $\nu_1(D)$ denotes the first eigenvalue of the usual Laplacian with Neumann boundary conditions on $\partial D \cap \partial\Omega$ and with Dirichlet boundary conditions on $\partial D \cap \Omega$.

Given our choice of Ω_ϵ , we remove from it the rectangle $(-1, 1) \times (0, \epsilon)$ to get the simply connected subdomain called D_ϵ . We estimate $\nu_1(D_\epsilon)$ by taking the test-function as follows:

$$\phi(x, y) = \begin{cases} 1 & \text{on the complement of } [-2, -1] \times [0, \epsilon] \cup [1, 2] \times [0, \epsilon] \\ x - 1 & \text{on } [1, 2] \times [0, \epsilon] \\ 1 - x & \text{on } [-2, -1] \times [0, \epsilon] \end{cases}$$

It is readily seen that $\int_{D_\epsilon} |\nabla \phi|^2 = 2\epsilon$, while $\int_{D_\epsilon} \phi^2 \geq C > 0$ (note that $C > 20$). Therefore:

$$\nu_1(D_\epsilon) \leq \frac{\epsilon}{10}.$$

Given (20) we conclude that:

$$\lambda_1(\Omega_\epsilon, A) \leq \frac{\beta(\Omega_\epsilon)}{10} = \frac{\epsilon}{10} \tag{21}$$

Applying our lower bound in Theorem 1 to this case, we have

$$\beta = \epsilon; B = \sqrt{5}; |F| = 32; |\partial F| = 24; D(F) = 4\sqrt{5},$$

and we obtain

$$\lambda_1(\Omega_\epsilon, A) \geq \frac{\pi^2}{360\sqrt{5}} d(\Phi^A, \mathbf{Z})^2 \epsilon. \tag{22}$$

We now observe that the minimum width $\beta(\Omega_\epsilon) = \epsilon$, by construction, and that $B(\Omega_\epsilon)$ is bounded above by 4. Taking into account (21) and (22) we see that $\lambda_1(\Omega_\epsilon)$ goes to zero proportionally to $\epsilon \sim \frac{\beta(\Omega_\epsilon)}{B(\Omega_\epsilon)}$.

3. Proof of Theorem 4

Let Ω be an n -holed domain, which we write: $\Omega = F \setminus (\bar{G}_1 \cup \dots \cup \bar{G}_n)$, with F, G_1, \dots, G_n smooth, open and convex.

From now on we denote $\Gamma_j = \partial G_j$. The idea is to use a suitable partition of Ω by annuli Ω_j whose boundary is either a piece of ∂F or is an equidistant curve from two interior boundary curves; each Ω_j is an annulus $\Omega_j = F_j \setminus G_j$ with piecewise-smooth exterior boundary ∂F_j which is star-shaped with respect to $\Gamma_j = \partial G_j$. We can then apply a theorem in [3] and obtain the uniform bound, valid for all j :

$$\lambda_1(\Omega_j, A) \geq \frac{2\pi^2}{9(|\partial F| + 2\pi B(\Omega))^2} \frac{\beta(\Omega)^4}{B(\Omega)^4} \cdot \gamma^2. \tag{23}$$

As the bound holds for all subdomains of a disjoint partition it holds a fortiori for Ω , thanks to Proposition 6.

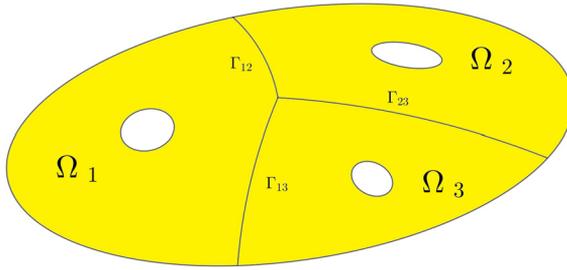


Fig. 9. The partition $\{\Omega_1, \Omega_2, \Omega_3\}$ of a domain Ω with three holes. The curves Γ_{jk} are equidistant sets.

3.1. The partition of Ω

We start by giving in Fig. 9 the picture of the partition $\{\Omega_1, \Omega_2, \Omega_3\}$ when Ω has three holes. The inner boundary of each piece Ω_j is made of equidistant sets from two suitable holes.

Here is the construction. For each $j = 1, \dots, n$ we consider the non-empty open set:

$$F_j = \{x \in F : d(x, G_j) < d(x, G_k) \text{ for all } k \neq j\}.$$

If we set:

$$H_{jk} = \{x \in \mathbf{R}^2 : d(x, G_j) < d(x, G_k)\} \tag{24}$$

we see that we can write

$$F_j = \bigcap_{k \neq j} (H_{jk} \cap F).$$

It is clear that $\bar{F} = \bigcup_{j=1}^n \bar{F}_j$. We remark that ∂H_{jk} is the equidistant set from G_j and G_k :

$$\partial H_{jk} = \{x \in \mathbf{R}^2 : d(x, G_j) = d(x, G_k)\}. \tag{25}$$

We have the following general fact.

Lemma 12. *Let G_1 and G_2 be disjoint smooth convex domains. Then the equidistant set ∂H_{12} as above is a smooth curve.*

Proof. Let $A = \mathbf{R}^2 \setminus (G_1 \cup G_2)$ and let ρ_j be the distance function to G_j , $j = 1, 2$. The convexity of G_j implies that ρ_j is smooth on the complement of G_j , so that ρ_j is smooth on A . Let $f = \rho_1 - \rho_2$, so that ∂H_{12} is the zero set of f . One has $\nabla f = \nabla \rho_1 - \nabla \rho_2$ and it is enough to show that ∇f has no critical points on $\{f = 0\}$. Let p be a point in $\partial H_{12} = \partial H_{21}$ and let γ_1 be the line segment which minimizes the distance from p to G_1 . One has: $\gamma_1(t) = p - t \nabla \rho_1(p)$. The corresponding minimizing segment from p to G_2 is then

$\gamma_2(t) = p - t\nabla\rho_2(p)$. If $\nabla\rho_1(p) = \nabla\rho_2(p)$ then $\gamma_1(t) = \gamma_2(t)$ and, as $d(p, G_1) = d(p, G_2)$, the two minimizing segments would have the same foot q , which would then belong to both G_1 and G_2 : but this impossible because G_1 and G_2 are disjoint.

Hence, on $\{f = 0\}$ one has $\nabla f \neq 0$ which proves smoothness. \square

As G_1, \dots, G_k are disjoint we see that $G_j \subset F_j$ and then we can introduce the annulus

$$\Omega_j \doteq F_j \setminus \bar{G}_j,$$

that is:

$$\Omega_j = \{x \in \Omega : d(x, G_j) < d(x, G_k) \text{ for all } k \neq j\}.$$

The family $\{\Omega_1, \dots, \Omega_n\}$ gives rise to a disjoint partition of Ω , as the next lemma shows.

Lemma 13. *The following properties hold:*

- a) $\bar{\Omega} = \cup_{j=1}^n \bar{\Omega}_j$.
- b) For $j \neq k$ one has that $\Omega_j \cap \Omega_k = \emptyset$ and $\bar{\Omega}_j \cap \bar{\Omega}_k$ is a smooth curve (eventually empty).
- c) Ω_j is an annulus with smooth inner boundary G_j and piecewise smooth outer boundary ∂F_j . Moreover:

$$\partial F_j = \left(\cup_{k \neq j} \Gamma_{jk} \right) \cup (\partial F \cap \bar{F}_j),$$

where $\Gamma_{jk} = \bar{F}_j \cap \bar{F}_k = \bar{\Omega}_j \cap \bar{\Omega}_k$ is contained in the equidistant curve ∂H_{jk} from G_j and G_k .

Note that actually $\bar{\Omega}_j \cap \bar{\Omega}_k = \bar{F}_j \cap \bar{F}_k$. The proof of the lemma is clear from the definitions.

3.2. Estimate of $\lambda_1(\Omega_j, A)$

As the partition of Lemma 13 is disjoint, from Proposition 6 we have:

$$\lambda_1(\Omega, A) = \min_{j=1, \dots, n} \lambda_1(\Omega_j, A).$$

Therefore, in this section, we estimate the first eigenvalue of the generic member of the partition. To that end, recall a relevant theorem from [3]. Let $\Omega_1 = F_1 \setminus G_1$ be an annulus with inner boundary curve $\Gamma_1 = \partial G_1$, where G_1 is smooth and convex. For $x \in \Gamma_1$ and $t \geq 0$, consider the segment $\gamma(t) = x + tN(x)$ where $N(x)$ is the unit normal to Γ_1 oriented outside G_1 . Let $Q(x)$ be the first intersection of $\gamma_x(t)$ with the outer boundary curve ∂F_1 , and let θ_x be the angle between γ'_x and the outer normal ν to F_1 at $Q(x)$. We set:

$$m(\Omega_1) \doteq \min_{x \in \Gamma_1} \cos \theta_x.$$

We recall that Ω_1 is said to be *starlike with respect to Γ_1* if, for any $y \in F_1$, the segment minimizing distance from y to Γ_1 is entirely contained in Ω_1 .

We also set:

$$\begin{cases} \beta(\Omega_1) \doteq \min\{d(x, Q(x)) : x \in \Gamma_1\} \\ B(\Omega_1) \doteq \max\{d(x, Q(x)) : x \in \Gamma_1\} \end{cases}$$

which are called, respectively, the minimum and maximum width of Ω_1 .

The estimate in Theorem 2 of [3] says that:

$$\lambda_1(\Omega_1, A) \geq \frac{4\pi^2}{|\partial F_1|^2} \frac{\beta(\Omega_1)m(\Omega_1)}{B(\Omega_1)} d(\Phi^A, \mathbf{Z})^2. \tag{26}$$

We will apply (26) to each annulus Ω_j in the above partition of Ω . We start from:

Lemma 14. *Let $\Omega = F \setminus (G_1 \cup \dots \cup G_n)$ and let $\Omega_j = F_j \setminus G_j$ be a piece in the partition defined above. Then:*

a) Ω_j is an annulus which is starlike with respect to ∂G_j , and moreover:

$$m(\Omega_j) \geq \frac{\beta(\Omega_j)}{2B(\Omega_j)}$$

b) One has the estimate:

$$|\partial F_j| \leq \frac{2B(\Omega_j)}{\beta(\Omega_j)} (|\partial G_j| + 2\pi B(\Omega_j))$$

Lemma 14 allows to prove Theorem 4 as follows. We apply (26) to Ω_j and get:

$$\lambda_1(\Omega_j, A) \geq \frac{\pi^2}{2(|\partial G_j| + 2\pi B(\Omega_j))^2} \frac{\beta(\Omega_j)^4}{B(\Omega_j)^4} \cdot \gamma^2.$$

To make the lower bound independent on j , it is enough to observe that $\beta(\Omega_j) \geq \beta(\Omega)$, $B(\Omega_j) \leq B(\Omega)$ and $|\partial G_j| \leq |\partial F|$. Then we get:

$$\lambda_1(\Omega_j, A) \geq \frac{\pi^2}{2(|\partial F| + 2\pi B(\Omega))^2} \frac{\beta(\Omega)^4}{B(\Omega)^4} \cdot \gamma^2,$$

which is the final step of the proof.

Then, it remains to prove Lemma 14.

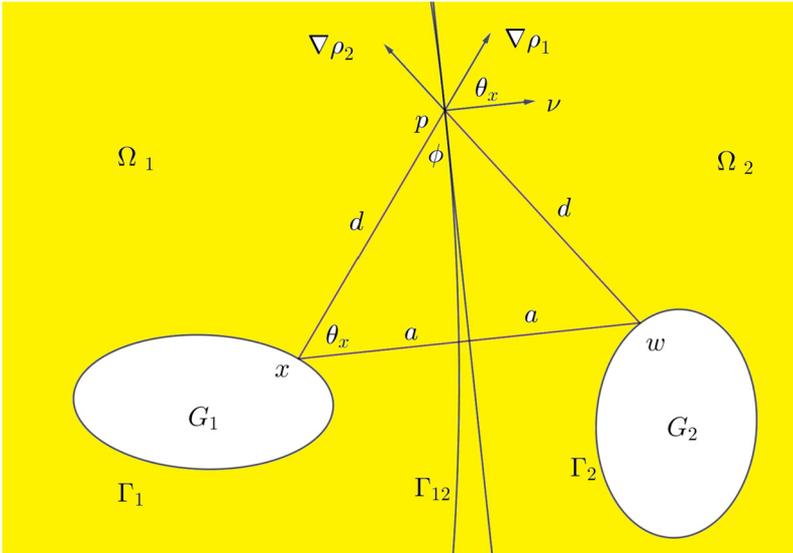


Fig. 10. The curve is the equidistant Γ_{12} from Γ_1 and Γ_2 . The tangent to Γ_{12} at $p = Q(x)$ is the line through p orthogonal to ν . It cuts the angle between $\nabla\rho_1$ and $\nabla\rho_2$ in half.

3.3. Proof of Lemma 14a

It is enough to prove it for $j = 1$. We first prove that $\Omega_1 = F_1 \setminus \bar{G}_1$ is star shaped with respect to $\Gamma_1 = \partial G_1$. Let $y \in \partial F_1$ and let σ be the segment starting at y and minimizing distance to Γ_1 : let $x \in \Gamma_1$ be the foot of σ .

- Note that, as $y \in \partial F_1$, we must have $d(y, \Gamma_1) \leq d(y, \Gamma_k)$ for all $k \neq 1$.

We prove that σ is entirely contained in Ω_1 . Assume by contradiction that there is $z \in \sigma$ such that $z \notin \Omega_1$. Then $z \in \Omega_h$ for some $h \neq 1$, and there exists $q \in \Gamma_h$ with $d(z, q) < d(z, x)$. But then:

$$\begin{aligned} d(y, x) &= d(y, z) + d(z, x) \\ &> d(y, z) + d(z, q) \\ &\geq d(y, q) \end{aligned}$$

that is, $d(y, x) > d(y, q)$ and this means that $d(y, \Gamma_1) > d(y, \Gamma_h)$, which contradicts the assumption. Hence Ω_1 is star shaped.

We now estimate $\cos \theta_x$, and for convenience we refer to Fig. 10.

Let $x \in \Gamma_1$ and draw the segment $\gamma_x(t) = x + tN_x$ where N_x is the unit normal vector to Γ_1 pointing outside G_1 . It hits ∂F_1 at the point $p = Q(x)$. If $p \in \partial F$ we proceed as in [3] (because ∂F is convex) and get

$$\cos \theta_x \geq \frac{\beta(\Omega_1)}{B(\Omega_1)}.$$

If $p \notin \partial F$ (as in the picture) then $p \in \Gamma_{1k}$ for some $k \neq 1$; we can assume that $k = 2$. Let w be the point in Γ_2 such that $d(p, x) = d(p, w) = d$; observe that $\cos \theta_x = \langle \nabla \rho_1, \nu \rangle$ where ν is the normal to Γ_{12} at p pointing away from Γ_1 . Observe that ν is the unit vector in the direction of $\nabla \rho_1 - \nabla \rho_2$, and that $\nabla \rho_1 + \nabla \rho_2$ is tangent to Γ_{12} at p . If 2ϕ is the angle between $\nabla \rho_1$ and $\nabla \rho_2$ then we see that $2\phi + 2\theta_x = \pi$, that is

$$\theta_x = \frac{\pi}{2} - \phi.$$

Consider the triangle with vertices x, w, p ; it is isosceles on the basis xw , (whose length is denoted $2a$); its height is part of the tangent line to the equidistant at p . One sees that

$$d \cos \theta_x = d \sin \phi = a$$

hence

$$\cos \theta_x = \frac{a}{d}.$$

Now $2a = d(x, w) \geq \beta(\Omega_1)$ by definition of $\beta(\Omega_1)$; as the segment joining x and p is entirely contained in Ω_1 we see that $d \leq B(\Omega_1)$. Hence

$$\cos \theta_x \geq \frac{\beta(\Omega_1)}{2B(\Omega_1)}$$

as asserted.

3.4. Proof of Lemma 14b

Recall that the typical piece of the decomposition is $\Omega_j = F_j \setminus G_j$. We need to estimate $|\partial F_j|$; this is a bit more difficult now because F_j is no longer convex (there are circumstances under which each F_j is actually convex - for example, when all holes are disks of the same radius - and we will discuss this case in the next section, to obtain a simpler final estimate).

Set $j = 1$ for concreteness. We apply Green formula to the function $\rho_1(x) = d(x, \partial G_1)$. Note that $\Delta \rho_1(x)$ is the curvature at x of the equidistant to ∂G_1 through x ; as ∂G_1 is convex one has $\Delta \rho_1 \leq 0$ on the complement of G_1 . By Green formula:

$$\int_{\Omega_1} \Delta \rho_1 = \int_{\partial \Omega_1} \langle \nabla \rho_1, N \rangle,$$

where N is the inner unit normal. We let $D(G_1, B)$ denote the B -neighborhood of G_1 , so that $F_1 \subseteq D(G_1, B)$ by the definition of B . Since $\Delta \rho_1 \leq 0$:

$$\int_{\Omega_1} \Delta \rho_1 \geq \int_{D(G_1, B)} \Delta \rho_1.$$

By co-area formula:

$$\int_{D(G_1, B)} \Delta \rho_1 = \int_0^B \int_{\rho_1=r} \Delta \rho_1 \, dr = -2\pi B$$

because $\int_{\rho_1=r} \Delta \rho_1 = -2\pi$ for all r (we are integrating the opposite of the curvature of a closed curve, and we always obtain -2π). Therefore:

$$\int_{\Omega_1} \Delta \rho_1 \geq -2\pi B. \tag{27}$$

On the other hand $\partial\Omega_1 = \partial G_1 \cup \partial F_1$. Hence:

$$\int_{\partial\Omega_1} \langle \nabla \rho_1, N \rangle = \int_{\partial G_1} \langle \nabla \rho_1, N \rangle + \int_{\partial F_1} \langle \nabla \rho_1, N \rangle.$$

The first piece is $|\partial G_1|$. On the outer boundary ∂F_1 we see that:

$$\langle \nabla \rho_1, N \rangle = -\cos \theta_x \leq -\frac{\beta}{2B},$$

where θ_x is as in the proof of part a), and the inequality then follows from part a). Then:

$$\int_{\Omega_1} \langle \nabla \rho_1, N \rangle \leq |\partial G_1| - \frac{\beta}{2B} |\partial F_1|,$$

and given (27) we obtain $-2\pi B \leq |\partial G_1| - \frac{\beta}{2B} |\partial F_1|$, that is:

$$|\partial F_1| \leq \frac{2B}{\beta} (|\partial G_1| + 2\pi B),$$

which gives the assertion.

4. Proof of Theorem 5

About the partition $\{\Omega_1, \dots, \Omega_n\}$ of the previous section for domains with n holes, we remark that if the inner holes G_1, \dots, G_n are disks of the same radius r , then the equidistant set between any pair of them is simply a straight line, and therefore each $\Gamma_{jk} = \bar{\Omega}_j \cap \bar{\Omega}_k$ is a line segment; moreover the subdomains F_1, \dots, F_n are all convex: see Fig. 11 which illustrates the partition when all holes shrink to a point.

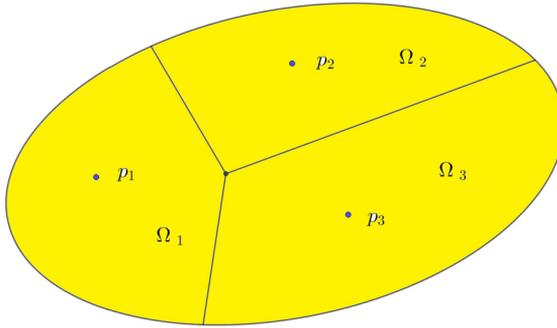


Fig. 11. The partition $\{\Omega_1, \Omega_2, \Omega_3\}$ for a domain Ω punctured at the points p_1, p_2, p_3 .

We can directly apply (2) to each Ω_j and obtain:

$$\lambda_1(\Omega_j, A) \geq \frac{4\pi^2}{|\partial F_j|^2} \frac{\beta(\Omega_j)^2}{B(\Omega_j)^2} \cdot d(\Phi_j, \mathbf{Z})^2$$

As F_j is convex, we have $|\partial F_j| \subseteq |\partial F|$ and therefore we arrive at the following estimate.

Theorem 15. *Let $\Omega = F \setminus (G_1 \cup \dots \cup G_n)$ with F convex and G_1, \dots, G_n being disjoint disks of center, respectively, p_1, \dots, p_n and common radius $r > 0$. Then:*

$$\lambda_1(\Omega, A) \geq \frac{4\pi^2}{|\partial F|^2} \frac{\beta(\Omega)^2}{B(\Omega)^2} \cdot \gamma^2$$

with $\gamma = \min_{j=1, \dots, n} d(\Phi_j, \mathbf{Z})$.

We remark that if we let $r \rightarrow 0$ in Theorem 15 we get the lower bound for the punctured domain $\Omega \setminus \{p_1, \dots, p_n\}$ as in Theorem 5.

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