# Towards a volumetric census of close white dwarf binaries – I. Reference samples

K. Inight<sup>®</sup>,<sup>1</sup>\* Boris T. Gänsicke<sup>®</sup>,<sup>1</sup> E. Breedt<sup>®</sup>,<sup>2</sup> T. R. Marsh,<sup>1</sup> A. F. Pala<sup>3</sup> and R. Raddi<sup>4</sup>

<sup>1</sup>Department of Physics, University of Warwick, Coventry CV4 7AL, UK

<sup>2</sup>Institute of Astronomy, University of Cambridge, Cambridge CB3 0HA, UK

<sup>3</sup>European Southern Observatory, Karl Schwarzschild Straße 2, D-85748 Garching, Germany

<sup>4</sup>Departament de Física, Universitat Politécnica de Catalunya, c/Esteve Terrades 5, E-08860 Castelldefels, Spain

Accepted 2021 March 9. Received 2021 February 8; in original form 2020 December 3

#### ABSTRACT

Close white dwarf binaries play an important role across a range of astrophysics, including thermonuclear supernovae, the Galactic low-frequency gravitational wave signal, and the chemical evolution of the Galaxy. Progress in developing a detailed understanding of the complex, multithreaded evolutionary pathways of these systems is limited by the lack of statistically sound observational constraints on the relative fractions of various sub-populations and their physical properties. The available samples are small, heterogeneous, and subject to a multitude of observational biases. Our overarching goal is to establish a volume-limited sample of all types of white dwarf binaries that is representative of the underlying population as well as sufficiently large to serve as a benchmark for future binary population models. In this first paper, we provide an overview of the project, and assemble reference samples within a distance limit of 300 pc of known white dwarf binaries spanning the most common sub-classes: post-common envelope binaries containing a white dwarf plus a main-sequence star, cataclysmic variables, and double-degenerate binaries. We carefully vet the members of these "Gold" samples, which span most of the evolutionary parameter space of close white dwarf binary evolution. We also explore the differences between magnitude and volume limited close white dwarf binary samples, and discuss how these systems evolve in their observational properties across the *Gaia* Hertzsprung–Russell diagram.

**Key words:** stars: evolution – Hertzsprung-Russell and colour-magnitude diagrams – novae, cataclysmic variables – stars: statistics; stars: distances.

#### **1 INTRODUCTION**

Binaries are important in many aspects of astrophysics such as their role as precursors of supernovae (Schaefer & Pagnotta 2012; Livio & Mazzali 2018), detectable sources of gravitational wave radiation (Nelemans 2018), and in the chemical evolution of galaxies (Tsujimoto et al. 1995; De Donder & Vanbeveren 2004). Binary evolution involves complex physical processes including common envelope evolution and orbital angular momentum loss that are not well-understood (Knigge, Baraffe & Patterson 2011; Ivanova et al. 2013). Models describing these physical processes involve a number of free parameters, that require observational calibrations. However, the currently known samples of different types of binaries are severely incomplete and subject to observational biases that are difficult to quantify (Gänsicke 2004; Pretorius, Knigge & Kolb 2007). Consequently, predictions of binary populations have to be considered with some reservations, as the parameters used to encapsulate the physics of binary evolution are not well-calibrated.

Progress requires the identification of representative populations of binaries with well-understood selection effects that sample the entire phase–space of physical parameters, as well as evolutionary stages. The key motivation of the project presented here is to assemble a volume-limited sample that includes a homogeneous representation of all major sub-types of close white dwarf binaries: detached white dwarf plus main-sequence post-common envelope binaries (PCEBs), interacting white dwarf binaries (cataclysmic variables, CVs), and double white dwarfs (DWDs). One example of such a study is the joint population study of all types of white dwarfs, single, and in binaries, within the volume-limited and complete 20 pc sample (Toonen et al. 2017).

The key information required for constructing a truly representative sample of white dwarf binaries is the accurate knowledge of their distances. *Hipparcos* measured parallaxes for  $\simeq 118\,000$ stars brighter than  $V \simeq 11$  (Perryman 1997) and included only a small number of CVs (Duerbeck 1999), as most of these systems are intrinsically faint. Ground-based parallax measurements for a few dozen CVs extended to fainter magnitudes, but were limited in accuracy (Thorstensen 2003; Thorstensen, Lépine & Shara 2008). Precise space-based parallaxes were obtained for a handful of selected CVs using the *Hubble Space Telescope (HST*; e.g. Harrison et al. 2000; Beuermann et al. 2003, 2004; Roelofs et al. 2007), although some of those measurements sparked controversies (Schreiber & Gänsicke 2002; Harrison et al. 2004; Miller-Jones et al. 2013). Accurate distances to faint stars have only now become

Published by Oxford University Press on behalf of Royal Astronomical Society. This is an Open Access article distributed under the terms of the Creative Commons Attribution License (http://creativecommons.org/licenses/by/4.0/), which permits unrestricted reuse, distribution, and reproduction in any medium, provided the original work is properly cited.

<sup>\*</sup> E-mail: keith.inight@gmail.com

<sup>© 2021</sup> The Author(s).

systematically available, thanks to the all-sky astrometry provided by the *Gaia* mission. In Data Release 2<sup>1</sup> (DR2; Gaia Collaboration et al. 2018) *Gaia* provides parallax measurements for ~1.3 billion stars, with a limiting magnitude of  $G \simeq 20$ .

Pala et al. (2020) used the Gaia DR2 data to establish the first truly volume-limited study of CVs within 150 pc. Whilst limited to only 42 systems, this study provided the most robust estimate of the space density of CVs to date ( $\rho_0 = 4.8 \frac{+0.6}{-0.8} \times 10^{-6} \text{ pc}^{-3}$ ), as well as an assessment of the intrinsic make-up of of the CV population, with the surprising conclusion that about one-third of all CVs contain a magnetic white dwarf.

With the *Gaia* astrometry of  $\simeq 1.3$  billion stars in hand, the limitation in extending the study of Pala et al. (2020) to a larger volume and including all types of close white dwarf binaries, is the incompleteness of the known members of the different sub-classes. Pala et al. (2020) estimated a  $\simeq 77$  per cent completeness for the CVs within 150 pc; that value drops for larger distances and is poorly constrained for PCEBs.

The first step within this project is to establish reference samples of each sub-type of white dwarf binaries that are as representative as possible. The definition of these samples and a discussion of their overall properties within the *Gaia* Hertzsprung–Russell (HR) diagram is the focus of this paper. These reference samples will then be used to define algorithms that identify white dwarf binary candidates, combining *Gaia* data with observations extending over a wider wavelength range.

In this paper, we identify the key characteristics of each close white dwarf binary sub-type (Section 2), discuss the limitations of *Gaia* (Section 3), define the reference "Gold" samples for each sub-type (Section 4) including a discussion of the selection effects that affect these samples, and finally use the Gold Samples to discuss the evolution of close white dwarf binaries in the *Gaia* HR diagram (Section 5).

#### 2 CLOSE WHITE DWARF BINARIES

Within this project, we will focus on close white dwarf binaries that have undergone interactions in the past, in the vast majority of cases in the form of a common envelope. Some of these are currently in a mass transferring state, others are in a detached configuration but may interact again in the future. The discussion in this and the following papers will frequently refer to these physical properties, as well as the past and future evolution of the major classes of close white dwarf binaries. A graphical overview of the evolutionary links between the different types of white dwarf binaries discussed below is shown in Fig. 1.

#### 2.1 Post-common envelope binaries (PCEBs)

Between 50 and 80 per cent of stars are formed as part of a binary system (Eggleton & Tokovinin 2008; Raghavan et al. 2010). Both stars will evolve normally through the main sequence until the higher mass primary runs out of hydrogen in its core and ascends the giant branch. If the initial binary separation is sufficiently small, it will eventually fill its Roche lobe, initiating unstable mass transfer on to the lower mass companion (Paczynski 1976; Webbink 1984). The mass transfer time-scale is faster than the rate that the companion can adjust to and the two stars are rapidly engulfed in a common envelope. Friction within this envelope drains energy and angular

<sup>1</sup>Data Release 1 only included parallaxes of  $\simeq 2$  million stars from a joint *Tycho-Gaia* astrometric solution, including 16 CVs (Ramsay et al. 2017).

momentum from the binary orbit, shrinking their separation. The envelope acquires this angular momentum and is ejected leaving a binary consisting of a white dwarf and a main-sequence star.

The orbital shrinkage caused by common envelope evolution implies that the resulting population of white dwarf binaries is bimodal (see fig. 10 in Willems & Kolb 2004) with PCEBs having periods of hours to days and wide white dwarf plus main-sequence star binaries (that avoided a common envelope and evolved as single stars) having periods in excess of a few years. We note that, depending on the details of the initial binary parameters, a small number of PCEBs can emerge with periods of ~100 d (Kruse & Agol 2014; Zorotovic, Schreiber & Parsons 2014). This bi-model distribution has been observationally confirmed with high-resolution *HST* imaging and radial velocity studies of a sample of white dwarf binaries (Farihi, Hoard & Wachter 2010; Ashley et al. 2019).

Identifying PCEBs is subject to two major selection effects: on the one hand, PCEBs in which one component dominates the observed wavelength range will resemble either single white dwarfs or mainsequence stars. On the other hand, with typical distances of a few 100 pc, PCEBs will be indistinguishable from spatially unresolved wide binaries separated by a few to a few tens of au, even at *Gaia*'s space-based resolution. These selection effects will be discussed in more detail in Section 4.1.

#### 2.2 Cataclysmic variables (CVs)

Once orbital angular momentum loss reduces the orbital separation of PCEBs such that the main-sequence star fills its Roche lobe, mass transfer on to the white dwarf is initiated via the inner Lagrangian point between the two stars; the systems morph into CVs, with observational properties that radically differ from those of their detached progenitors.

Most CVs above the period gap have periods  $\leq 14$  h, whilst a small number of systems with donors that have slightly evolved off the main sequence have periods extending up to a few days (Sekiguchi 1992). As a CV evolves and the companion loses mass, the separation reduces and the orbital period decreases. The observed distribution of orbital periods shows a significant dip between 2 and 3 h (Whyte & Eggleton 1980), and this deficiency is known as the period gap. The currently accepted explanation for the period gap is that CVs undergo a major drop in the rate at which they lose angular momentum once the donor stars become fully convective (Rappaport, Verbunt & Joss 1983; Howell, Nelson & Rappaport 2001), at  $\simeq$ 3 h and evolve as detached systems across the period gap (Davis et al. 2008; Zorotovic et al. 2016), until the Roche lobe shrinks sufficiently to, once again, establish a semi-detached configuration, re-starting the mass transfer, albeit at a much lower rate. As mass transfer keeps eroding the mass of the donor star, it eventually becomes a brown dwarf. The result of this change in the internal structure of the donors is that further evolution will drive them back towards longer periods (Paczynski & Sienkiewicz 1983) and hence there is a minimum orbital period,  $P_{\min}$ . As CVs evolve twice through the period range near  $P_{\min}$ , there is an observationally confirmed pile-up of systems with  $P_{\rm min} \simeq 80 \,\rm min$ (Gänsicke et al. 2009).

#### 2.2.1 CV sub-types

Historically, a large number of CV sub-classes have been defined based on observational characteristics, which, ideally, should map on to a small number of fundamental physical properties of these binaries. Three such defining key parameters are the orbital period  $(P_{\rm orb})$  of the system, magnetic field strength (B) of the white dwarf,

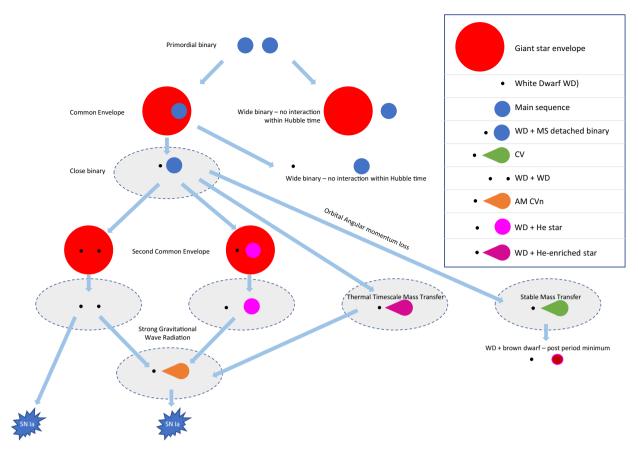


Figure 1. A schematic overview of the evolutionary channels leading to the different sub-classes of close white dwarf binaries discussed in this paper. A number of these channels potentially lead to the ignition of white dwarfs as type Ia supernovae and to the emission of gravitational waves detectable by the *LISA* mission, highlighting the importance of these systems for a wide range of astrophysical problems. The grey ellipses show the sub-classes that are included in our census. Adapted from Toloza et al. (2019).

and the mass-loss rate  $(\dot{M})$  of the companion (or donor) star. These various sub-classes of CVs with different accretion geometries present distinctly different spectroscopic appearances, which will affect our ability to identify them (Section 4.3).

In the absence of a strong magnetic field, an accretion disc forms around the white dwarf. These discs can be subject to a limit cycle based on a thermal instability (Meyer & Meyer-Hofmeister 1981), in which they undergo quasi-periodic large changes in disc temperature and brightness – observationally, CVs that exhibit these disc outbursts are called dwarf novae. Whether or not an accretion disc is subject to this instability depends on both the size of the disc (and hence on the orbital separation, and in turn on  $P_{\rm orb}$ ) and on  $\dot{M}$  (Shafter, Wheeler & Cannizzo 1986; Shafter 1992). If the disc is sufficiently small, and  $\dot{M}$  sufficiently high, the disc will be in a steady, hot, bright state (akin to a dwarf nova permanently in outburst), and such CVs make up the sub-class of nova-like variables, or novalikes for short.

Dwarf novae are the dominant CV sub-type in a volume-limited sample (Pala et al. 2020) and are further divided according to their outburst behaviour. SU UMa systems (Kato et al. 2020) exhibit, in addition to the normal short outbursts lasting a few days, superoutbursts that are somewhat brighter, and can last several weeks. These superoutbursts are triggered when the outer radius of the disc reaches a 3:1 resonance, resulting in tidal interactions with the companion star (Osaki 1989). These interactions force the accretion disc to precess, resulting in a brightness modulation, so-called *superhumps*, on the beat period between the precession period of the disc and the orbital period – i.e. somewhat longer than the orbital period (Patterson et al.

2005). The majority of the SU UMa systems are found below the period gap. WZ Sge systems are a sub-class of dwarf novae defined by having extremely long outburst recurrence times of years to decades (Patterson et al. 2002; Kato 2015), implying that they have low mass transfer rates; most of them have orbital periods in the range  $\simeq$ 75–95 min. Dwarf novae that *do not* undergo superoutbursts are almost exclusively found above the period gap, and further divide into U Gem systems (Smak 1984) that exhibit quasi-periodic outbursts and Z Cam systems (Meyer & Meyer-Hofmeister 1983) that alternate between states where they undergo disc instabilities, and so-called stand-stills during which the disc remains in a hot, stable state, akin to the novalike variables.

If the magnetic field of the white dwarf is sufficiently strong  $(B \gtrsim 10 \text{ MG})$  to control the accretion geometry, essentially such that the Alfvén radius of the white dwarf is larger than its Roche lobe, the formation of an accretion disc is suppressed. The material lost by the donor initially follows a ballistic trajectory, and is then channelled along the magnetic-field lines on to the magnetic poles of the white dwarf. These CVs go by the name *AM Her systems* (after their prototype, Chanmugam & Wagner 1977) or *polars* (because of the strongly polarized cyclotron radiation they emit; Tapia 1977). Interactions of the strong magnetic field of the white dwarf with that of the secondary star lock the white dwarf into synchronous rotation with the orbital period (Lamb et al. 1983). An early overview of the study of polars is given by Cropper (1990).

In the case of an intermediate magnetic-field strength ( $1 \leq B \leq 10$  MG), a vestigial disc may form, disrupted at the Alfvén radius,

from where the material is again funnelled on to the magnetic poles of the white dwarf. The magnetic-field strengths in these *intermediate polars* (sometimes called *DQ Her stars*), are poorly known, as usually neither Zeeman-split photospheric absorption lines from the photosphere, nor cyclotron emission from the accretion funnel, are observed. The white dwarf spin periods in intermediate polars are always shorter than the orbital periods, and some of them may evolve into polars as their period shortens and mass transfer decreases (Norton, Wynn & Somerscales 2004; Norton et al. 2008). An overview of the properties of intermediate polars is given by Patterson (1994).

A sub-class of CVs with distinct physical properties, and evolutionary histories, are the AM CVn stars: binaries containing a white dwarf accreting from a hydrogen-deficient donor star that is fully or partially degenerate (Deloye, Bildsten & Nelemans 2005), with orbital periods of  $\simeq$ 5–68 min (Levitan et al. 2015; Ramsay et al. 2018; Green et al. 2020). These systems are important in the context of gravitational wave physics, as they will be easily detected by the *LISA* space mission (Kupfer et al. 2018).

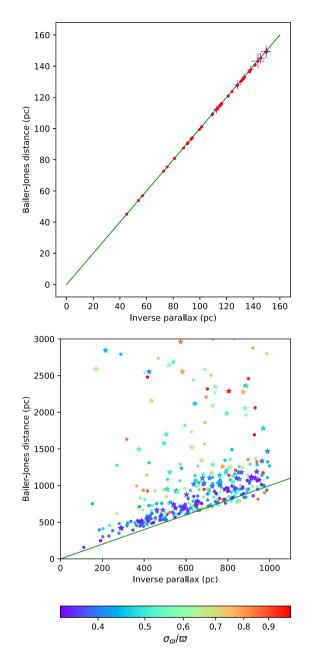
Finally, there are a handful of exotic, and nearby, white dwarf binaries such as the radio-pulsating AR Sco (Marsh et al. 2016) that fall outside the current evolutionary models.

#### 2.3 Double white dwarfs (DWDs)

DWDs (Marsh, Dhillon & Duck 1995) are often formed as the result of a second common envelope event. Others may have formed the first white dwarf in a period of stable mass transfer (Algol-like) to the companion. DWDs lose angular momentum through gravitational radiation and their separation will decrease - eventually leading to a merger. If their combined mass is sufficiently large, this merger may result in a thermonuclear supernova (Toonen, Nelemans & Portegies Zwart 2012; Shen et al. 2018). DWDs are the most common, and have the best-determined space density, of all close white dwarf binaries, thanks to the detailed studies of the local 20 pc sample: Toonen et al. (2017) lists one confirmed and six candidate DWDs (see their table 1) within 20 pc. Four of those candidates were confirmed as DWDs based on their Gaia parallaxes, with two remaining possible candidates (Hollands et al. 2018), implying that  $\simeq$ 3–5 per cent of all apparently single (i.e. spatially unresolved) white dwarfs are actually DWDs, with a space density of  $\simeq (1.4 - 2.3) \times 10^{-4} \text{ pc}^{-3}$ . Most of these DWDs are cool and would be very hard to identify at larger distances. A complete volumetric sample therefore will need to encompass a much smaller volume than that of CVs and PCEBs.

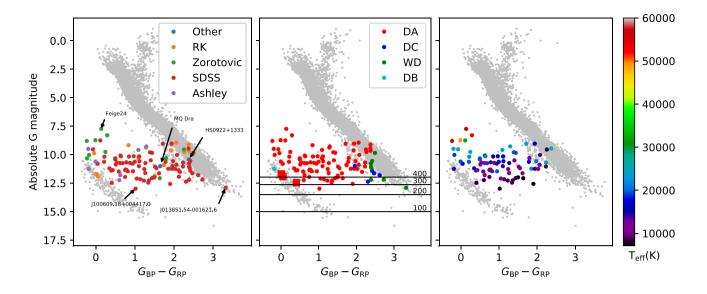
#### **3 FROM PARALLAXES TO DISTANCES**

The astrometry of the *Gaia* space mission (Gaia Collaboration et al. 2016, 2018) is revolutionizing Galactic and stellar astronomy. *Gaia* has vastly increased the number of stars with accurate parallaxes,  $\varpi$ , however, a large fraction of the objects within DR2 have substantial parallax uncertainties, such that  $d = 1/\varpi$  is no longer a reliable estimate of their distances. Bailer-Jones (2015) has used a probabilistic model using the known geometric distributions of stars to derive distance estimates that avoid the errors in traditional methods, such as spectro-photometric estimates which are subject to assumptions regarding interstellar extinction and stellar effective temperatures. We illustrate (Fig. 2) the difference between  $d = 1/\varpi$  and the distances of Bailer-Jones et al. (2018) using the sample of 42 CVs (top panel) known within 150 pc (Pala et al. 2020), all of which have  $\sigma_{\varpi} \leq 0.24$  mas and  $\varpi/\sigma_{\varpi} > 27.7$ . This contrasts with the 374 CVs and CV candidates with distances beyond 1000 pc with



**Figure 2.** Comparison between  $d = 1/\varpi$  and the distances of Bailer-Jones et al. (2018). Top panel: 42 known CVs within 150 pc (Pala et al. 2020), with  $\sigma_{\varpi} \leq 0.24$  mas and  $\varpi/\sigma_{\varpi} > 27.7$ , illustrating that for accurate parallaxes, the inverse of the parallax is a very good approximation for the distance. Bottom panel: 374 CVs and CV candidates with  $\sigma_{\varpi} > 0.33$ mas and  $1 < \varpi/\sigma_{\varpi} < 3$ . A star indicates that the object is outside the Galactic plane ( $|b| > 10^{\circ}$ ). It is evident that for large parallax uncertainties,  $1/\varpi$  often significantly underestimates the likely distance of the systems. In both panels, the solid green line indicates a 1:1 correspondence.

 $\varpi > 1$  mas and  $1 < \varpi/\sigma_{\varpi} < 3$ . Fig. 2 demonstrates that  $1/\varpi$  results in substantial underestimates with respect to the distances of Bailer-Jones et al. (2018) for data with large parallax uncertainties. We note from Lindegren et al. (2018) that faint sources or those in crowded areas of the sky are most susceptible to errors in the five-parameter solution (and the parallax in particular). We have split the sources in Fig. 2 according to their location on sky, with those in crowded areas (defined as the galactic plane,  $|b| < 10^\circ$ ) shown as circles and those in



**Figure 3.** The WD+M Gold Sample in the HR diagram. Left-hand panel: distribution of the sub-samples. Middle panel: distribution of the white dwarf spectral types. The concentration of systems with no spectroscopic classification (WD, green symbols) near the main sequence indicates that the companion dominates the optical flux of these systems. Systems with brown dwarf companions are indicated by squares; their location within the white dwarf cooling sequence is consistent with the low luminosity of the brown dwarfs. The minimum absolute magnitude to be detectable at an apparent magnitude of G = 20 is indicated by grey lines for distances of 200, 300, and 400 pc. Right-hand panel: The white dwarf effective temperatures of the WD+M Gold Sample plotted on the HR diagram. The increase in absolute magnitude with decreasing effective temperature is evident.

uncrowded areas as stars. It appears that for CVs and CV candidates, there is no clear correlation between the magnitude of the discrepancy between  $1/\varpi$  and the distances of Bailer-Jones et al. (2018) and whether a source is found within a crowded or uncrowded region. Bailer-Jones et al. (2018) acknowledge that there are shortcomings in the relatively simple prior used and indeed have been developing more sophisticated models for Gaia EDR3 (Bailer-Jones et al. 2020). Nevertheless, the Bailer-Jones et al. (2018) estimates (r\_est) are a significant improvement on simple inverse parallax estimates and we therefore have adopted them throughout this paper.

One major shortcoming of Gaia DR2 for the construction of volumetric samples is that 361 009 408 (21.3 per cent) of the objects in Gaia DR2 do not have parallaxes and hence their distances are unknown. The most common reason is that those sources are too faint for Gaia to obtain a parallax. Whilst this is not a major problem for the purpose of this paper (the definition of Gold Samples of various classes of white dwarf binaries), it is a concern for the future papers of this project, as several sub-classes (such as low  $\dot{M}$  CVs, DWDs, and binaries containing cool white dwarfs with very lowmass companions) are intrinsically faint. Hence, the incompleteness of Gaia parallaxes is likely to reduce the size of our census and may introduce biases. However, we expect the situation to improve with Gaia DR3. Here, we assess the significance of the lack of parallaxes of a significant fraction of the Gaia DR2 sources by establishing, for each white dwarf binary subclass individually, the fraction of previously known systems without a parallax that are likely to be within 300 pc.

#### **4 REFERENCE SAMPLES**

The objective of this section is to define a prototypical Gold Sample for each type of close white dwarf binary based on the known systems that can then be used to define selection criteria for an all-sky search. The ultimate goal of this search is to establish a view of the Galactic white dwarf binary populations that is homogeneous, representative, and as complete as possible within a given volume limit. We emphasize that each Gold Sample will inevitably be subject to the selection effects underlying the known populations of white dwarf binaries. However, our goal here is not to achieve completeness at any level, but to map out the parameter space occupied by these white dwarf binaries.

For the goal of completeness within a given volume, two aspects limit the maximum distance for which it will be possible to identify white dwarf binaries, i.e. the requirement for (1) a *Gaia* parallax, and (2) spectroscopic confirmation. Within Gaia DR2, parallaxes require a five-parameter astrometric solution; which exists for  $\simeq 83$  per cent of targets with  $G \leq 20$ , dropping to  $\simeq 16$  per cent at  $G \leq 21$  (interpolated from fig. B1 in Lindegren et al. 2018). The required follow-up spectroscopy of a large number ( $\simeq 100\ 000$ ) of candidate white dwarf binaries can only be achieved via multiobject spectroscopic surveys, such as DESI (DESI Collaboration et al. 2016a,b), WEAVE (Dalton et al. 2012), 4MOST (de Jong et al. 2016), and SDSS-V (Kollmeier et al. 2017). These experiments will deliver spectroscopy with a sufficient quality for a reliable identification for sources as faint as  $G \simeq 20$ 

We therefore adopt a magnitude limit of G = 20, and using  $m = 20 - 5\log(d/10)$ , we determine the intrinsically faintest systems (assuming minimal extinction) that we can detect for a given distance limit, as illustrated in Figs 3, 8, and 11 for the different classes of white dwarf binaries. The overall volume limit of our census is a compromise between sample size and completeness and we have chosen

$$r_{est} < 300 \, pc \tag{1}$$

where  $r_{\text{est}}$  is the distance estimate of Bailer-Jones et al. (2018), with the exception of the DWDs (Section 4.4). This choice will introduce well-quantifiable selection effects (some types of white dwarf binaries are inherently fainter than others), which we explore further in the context of each individual sub-class.

Accurate astrometry and photometry are needed to avoid misleading conclusions when we use the samples to evaluate potential cuts. Following Lindegren et al. (2018) and Lindegren (2018), we filter the astrometry by

$$\varpi/\sigma_{\varpi} > 10 \tag{3}$$

where ruwe is the "re-normalized unit weight errors", and  $\varpi$  and  $\sigma_{\varpi}$  are the parallaxes and its uncertainties.

To ensure that the precision of the red and blue photometry is within 10 per cent ( $\simeq 0.1$  mag), we again followed Lindegren et al. (2018) and filtered on

$$phot_rp_mean_flux_over_error > 10$$
 (5)

Known CVs that failed these quality cuts are listed in Table A4. These filters do not just exclude dim stars. Two examples closer than 100 pc are U Gem (excluded because ruwe = 1.48) and AM Her (excluded because phot\_rp\_mean\_flux\_over\_error = 6.04). Although it may seem counter-intuitive to omit well-known systems, it must be remembered that the purpose of a Gold Sample is not to be complete, but to be clean, with good *Gaia* data. The significance of this will become apparent in the context of target selection in Paper II.

The accurate astrometry enables the derivation of an important parameter of the various white dwarf binary classes: their space density. The space density makes assumptions about the age and hence the scale height of systems and can be calculated using the approximation

$$\rho = \rho_0 \,\mathrm{e}^{-|z|/h} \tag{6}$$

where h is the scale height. The effective volume for a given radius (R) is then calculated by integration

$$V_{\rm eff} = \int_0^R \int_{-\sqrt{R^2 - x^2}}^{\sqrt{R^2 - x^2}} (2\pi x) \left( \exp^{-|z|/h} \right) dz \, dx \tag{7}$$

Using this expression, the space density of the considered type of white dwarf binary is then given by the size of the sample divided by  $\rho = N/V_{\text{eff}}$ . This requires an assumption for the value of *h*, which depends upon the age of the stars in the sample (Pretorius et al. 2007). Below, we follow Pala et al. (2020) and calculate the space densities of the Gold Samples for three different values of *h* (100, 280, 500 pc). The estimates of  $\rho_0$  will represent lower limits to the true space densities of the respective white dwarf binary populations as the Gold Samples are not complete.

#### 4.1 PCEBs: White dwarfs with M-type companions

#### 4.1.1 Establishing the Gold Sample

There has been a continuing effort over many years to build a comprehensive catalogue of binaries consisting of a white dwarf and a main-sequence M-type star (hereafter WD+M; see Rebassa-Mansergas et al. 2007, 2010, 2012, 2013, 2016). This catalogue contains WD+M binaries identified from the detection of both stellar components in their SDSS spectroscopy. Whereas these systems are unresolved in the SDSS imaging data, given the typical spatial resolution  $\simeq 1.5$  arcsec, and distances of several 100 pc, they can be either short-period (hours to days) PCEBs or wide binaries with separations of tens of au that never interacted. Extensive radial velocity studies of the SDSS WD+M binary sample showed that

 $\simeq$ 21–24 per cent are PCEBs, with a period distribution peaking at  $\simeq$ 10.3 h (Nebot Gómez-Morán et al. 2011). In the online catalogue of SDSS WD+M binaries,<sup>2</sup> PCEBs are defined as systems where a  $\ge$ 3 $\sigma$  radial velocity variation has been detected in the spectroscopy, and we select these "SDSS" systems for our Gold Sample.

To these, we added confirmed WD+M binaries from Farihi, Hoard & Wachter (2006) and Farihi et al. (2010), some of which were later followed up by Ashley et al. (2019).

These WD+Ms (hereafter, the "Ashley sample") were identified from catalogues of white dwarfs due to their infrared excess and were unresolved in high-resolution *HST* imaging. Ashley et al. (2019) obtained radial velocity measurements of these systems, and corroborated the bi-modality of the orbital separations referred to earlier, i.e. that we can expect WD+M binaries to be either close (with periods of hours to days) or wide (with periods > years).

We also added confirmed PCEBs from appendix A2 of Zorotovic et al. (2010) (hereafter the "Zorotovic sample") as well as the confirmed PCEBs from the catalogue of Ritter & Kolb (2003) that were not contained in any of the other published lists (hereafter, the "RK sample").

The process for assembling the Gold Sample is illustrated in Table 1. We first match each system to the closest *Gaia* source within 2 arcsec and then removed those with distances ( $r\_est$  from Bailer-Jones et al. 2018) greater than 300 pc. We then removed any systems that fail the data quality criteria (2)–(5) and finally rejected wrongly classified systems, e.g. those with an A to K class companion (which we incorporate in the sample discussed in the next Section), or those considered to be unreliable.

The WD+M Gold Sample includes four systems containing strongly magnetic white dwarfs in which accretion of the wind from the M-dwarf companion results in the emission of cyclotron radiation: MQ Dra (Szkody et al. 2008), WX LMi (Reimers, Hagen & Hopp 1999), SDSS J030308.35+005444.1 (Parsons et al. 2013), and HS0922+1333 (Reimers & Hagen 2000). These systems are considered to be progenitors of magnetic CVs, and hence are called pre-polars, or PREPs.

Finally, whilst we refer to the set binaries assembled here as the WD+M Gold Sample, we have also included four systems with brown dwarf companions, as their properties and evolution are very similar to those of white dwarfs with low-mass Mdwarf companions: SDSS J013532.97+144555.9 (= NLTT 5306, Steele et al. 2013), WD 0137–3457, (Burleigh et al. 2006), SDSS J141126.20+200911.1 (Casewell et al. 2018), and GD1400, (Farihi, Zuckerman & Becklin 2005a).

As mentioned above, we defined 300 pc as the limit for the WD+M Gold Sample yielding 114 systems, which are listed in Table A1.

#### 4.1.2 Selection effects

The distribution of the WD+M binaries in the *Gaia* HR diagram is illustrated in Fig. 3, where they define a horizontal "bridge" between the white dwarf cooling sequence and the main sequence. The evolutionary interpretation of this bridge is discussed in more detail in Section 5.1. We discuss a few individual systems to illustrate the effect of the white dwarf and M-dwarf spectral type on the location in the HR diagram. Feige 24 is a well-studied PCEB (Kawka et al. 2008) that contains a particularly hot ( $\simeq$  56 000 K; Vennes et al. 2000) white dwarf, consistent with having the smallest absolute magnitude among the WD+M Gold Sample.

<sup>2</sup>http://www.sdss-wdms.org see (Rebassa-Mansergas et al. 2012a)

System type		WD+M			W	WD+AFGK		CV		DWD		
Source	SDSS	Ashley sample	Zorotovic sample	RK sample	Other	Ren sample	RK sample	Other		20 pc sample	40 pc sample	Other
Source systems	3287	16	25	190	2	23	25	7	5192	5	24	198
Duplicates	-1	-6	-3	-131				-2			-3	-2
Systems with no Gaia match	-213	-1		-2								
Systems with no parallax	-429			-3					-1150			-6
Systems $> 300  \text{pc}$	-2049		-4	-26		-10	-23		-3737			-129
Systems failing quality criteria <sup><i>a</i></sup>	-272	-2	-7	-5		-3	-1	-1	-104	-1	-10	-9
Unreliable/invalid systems				-8					-50			
Wide binaries	-125											
Possible wide binaries	-119											
Final Gold Sample	79	7	11	15	2	10	1	4	151	4	11	52
			114				15		151		67	

 Table 1. Assembling the Gold samples. Positive numbers indicate the initial sample size, negative numbers indicate the effects of the individual cuts that were subsequently applied. The requirement for a complete set of accurate attributes dramatically reduces the final sample size.

<sup>a</sup>see equations (2)–(5).

In contrast SDSS J100609.18+004417.0 contains a cold (7819 K) white dwarf and a late-type (M7), low-mass ( $0.12 M_{\odot}$ ) secondary (Nebot Gómez-Morán et al. 2011). Both stellar components are intrinsically faint and the system has one of the largest absolute magnitudes in the Gold Sample. Despite the low white dwarf temperature, the system is located relatively close to the white dwarf cooling track, a consequence of the feeble luminosity of its M-dwarf companion. Finally, SDSS J013851.54-001621.6 contains an ultraccol (3570 K) white dwarf and an M5 companion (Parsons et al. 2012). The white dwarf contributes hardly anything to the optical flux of the system, and it is hence located within the low-mass end of the main sequence. These two systems illustrate that the limit for a truly volumetric sample is  $\simeq 250$  pc to ensure faint systems such as SDSS J100609.18+004417.0 are fully represented (see Fig. 3, left-hand panel).

A global selection effect that affects the WD+M sample is that it has been primarily constructed from systems identified by SDSS. Consequently, the restricted sky coverage of SDSS has to be taken into account. Based on the sky distribution of  $\simeq$ 5.8 million optical spectra contained within DR16 (Ahumada et al. 2020), we estimate that SDSS spectroscopy has been obtained over  $\simeq$ 29 per cent of the sky.

The known population of WD+M binaries will be subject to additional selection effects that will affect both the completeness and distribution in the HR diagram parameter space of the Gold Sample defined here. As these systems are, for the vast majority, selected from optical spectroscopy, a key criterion necessary for their identification is that both stellar components contribute noticeable amounts of flux in the visual wavelength range. That immediately implies that a WD+M in which one of the components strongly dominates the optical spectrum will be difficult if not impossible to identify (see Section 4.2). One example consists of hot white dwarfs with very late type M-dwarf companions. The detection of emission lines from the irradiated companion may mitigate the low overall flux contribution of the companion, although that introduces another selection effect, i.e. the binary separation. Very cool white dwarfs with relatively early M-dwarf companions form another example. In addition to these intrinsic limitations in the identification of WD+M binaries, there will be additional biases related to the specific observations that were used in a given study, such as the wavelength range covered by the spectroscopy, and brightness limits.

The following discussion focuses on the SDSS WD+M sample, which was identified and analysed in a homogeneous fashion, and

bias, that we have mitigated by including the systems from Ritter & Kolb (2003), Ashley et al. (2019), and Zorotovic et al. (2010). More subtle and hard-to-quantify effects arise from the fact that the SDSS sample was obtained using the 640-fibre SDSS and the 1000-fibre BOSS spectrographs sampling a  $\simeq 7 \text{ deg}^2$  field-of-view (York et al. 2000; Smee et al. 2013). The allocation of these fibres is subject to a complex selection algorithm, that has evolved over the course of the SDSS, and will result in both incompleteness and magnitude and colour-dependent selection effects. Rebassa-Mansergas et al. (2010, 2012, 2013, 2016) reported spectral types for the white dwarf and M-dwarf components, as well as effective temperatures for a sub-set of the white dwarfs. We discuss the distributions of those parameters in turn. *White dwarf spectral type*: The white dwarfs within the SDSS WD+M sample were classified, according to the key features in their optical spectra, into DAs (clearly detected Balmer lines), DB (clearly

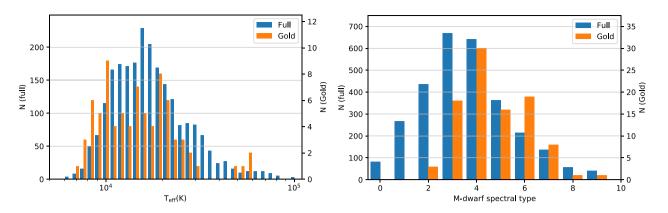
which makes up the bulk of the Gold Sample. Specifically, the SDSS WD+M sample has a bright-limit of  $g \simeq 15$  to avoid saturation

of the CCDs and contamination of the spectra of fainter objects

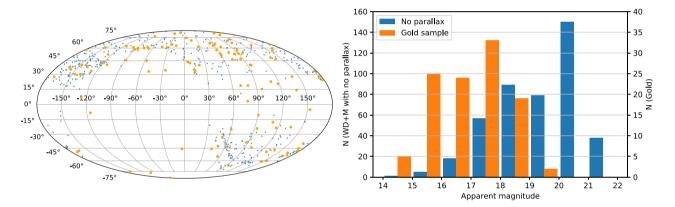
on adjacent fibres (Strauss et al. 2002). This limit will introduce a

WD+M sample were classified, according to the key features in their optical spectra, into DAs (clearly detected Balmer lines), DB (clearly detected helium lines), DC (featureless continua), and WD (definite blue excess over the M-dwarf, but insufficient to assess the intrinsic spectrum of the white dwarf); the respective numbers of these types are listed in Table 2. The vast majority of the WD+M binaries contain DA white dwarfs, and that fraction may be somewhat higher than among single field white dwarfs in a magnitude-limited sample (e.g. Kleinman et al. 2013). Accretion of the wind of the companion (Debes 2006) would relatively rapidly build up a sufficient amount of hydrogen to convert an initially helium-dominated white dwarf into a DA. Fig. 3 (centre panel) illustrates the distribution of white dwarf spectral types within the *Gaia* HR diagram. As expected, the WD+M systems with barely recognizable contributions from the white dwarf are mostly localized close to or in the main sequence.

White dwarf temperature: The temperature distribution of the white dwarfs within the Gold WD+M sample in the HR diagram is shown in Fig. 3 (right-hand panel). Again, as expected, the WD+M binaries with the lowest temperatures have the largest absolute magnitudes. Comparing the white dwarf temperatures in the Gold WD+M sample ( $d \leq 300 \text{ pc}$ ) with those in the full sample (Fig. 4) reveals a bias against hot white dwarfs within the Gold Sample. This bias is a consequence of the brightness limit of the SDSS spectroscopy – whilst the actual value of that limit varied throughout the different installments of the survey, and its diverse sub-programs,



**Figure 4.** Left-hand panel: The distribution of the effective temperatures (where known) of the white dwarfs in the WD+M Gold Sample contrasted with the full (not volume-limited) set from from Rebassa-Mansergas et al. (2016). Right-hand panel: The distribution of the spectral types (where known) of the white dwarf companions. The WD+M Gold Sample is contrasted with the full (not volume-limited) set from Rebassa-Mansergas et al. (2016) and shows that the WD+M Gold Sample has a higher proportion of late companions.



**Figure 5.** Left-hand panel: The distribution of the WD+M Gold Sample (orange) and candidates without a *Gaia* parallax (blue) in Galactic coordinates. Right-hand panel: The *G*-band magnitude distribution of the WD+M Gold Sample compared with the WD+M systems without a parallax. It is evident that most of the systems without a parallax are fainter than the Gold Sample – although the parallax incompleteness is already rising when  $G \gtrsim 17$ , which is several magnitudes brighter than *Gaia*'s magnitude limit.

Table 2. The number of each WD sub-type (where known) in the WD+M Gold Sample.

Sub-type		n
DA	Only hydrogen features	94
DC	Featureless continuum	5
WD	No WD spectral type	8
DB	Only helium features	1

very few stars brighter than  $g \simeq 16$  were observed with the SDSS or BOSS spectrographs, and the hottest white dwarfs will violate that limit at d < 300 pc.

White dwarf temperature: Rebassa-Mansergas et al. (2010) provide the spectral types of the M-dwarf companions of their sample. Fig. 4 shows that the distribution of the WD+M Gold Sample is biased in favour of later stars. This again shows the difference between a volume-limited and a magnitude-limited sample, with the intrinsically fainter late spectral types being underrepresented in the latter. A more subtle selection effect is the correlation between donor spectral type and white dwarf temperature: for the WD+M to be detected as such, the white dwarf will need to be brighter, and hence hotter, for earlier type donors. It follows that a magnitudelimited sample will favour intrinsically bright white dwarfs and hence earlier type donors.

#### 4.1.3 The Gold Sample as a fraction of the total population

The space density of the WD+M Gold Sample is reported in Table 3 for the three adopted values of the scale height. Remembering the selection effect caused by the SDSS sky coverage these space densities should be multiplied by ~3.4 before making comparisons with the intrinsic space density of WD+M binaries. Few published studies on observational estimates of the space density of WD+M binaries exist, Schreiber & Gänsicke (2003)<sup>3</sup> quote  $6 - 30 \times 10^{-6} \text{ pc}^{-3}$  which is substantially higher than the values that we derive here – however, these authors accounted for the selection effects introduced by the dimming of the white dwarf throughout the PCEB evolution. The WD+M Gold Sample clearly forms a small subset of the underlying population, and is biased in favour of systems with hot white dwarfs and early-type M-dwarfs.

<sup>3</sup>The sample of Schreiber & Gänsicke (2003) includes a small number of PCEBs with companion spectral types earlier than M.

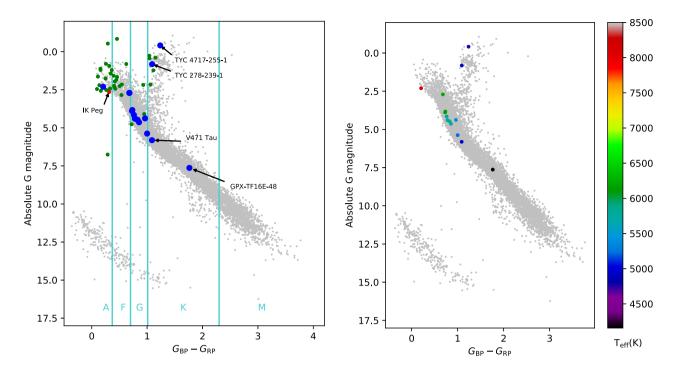


Figure 6. Left-hand panel: The WD+AFGK Gold Sample in the HR diagram contrasting the WD+AFGK PCEBs over 300 pc (green), IK Peg (red), and the Gold Sample (blue). The boundaries between spectral types are shown in turquoise. Right-hand panel: The distribution of the effective temperatures of the main-sequence (and sub-giant) companions of the WD+AFGK Gold Sample.

**Table 3.** The estimated effective volume ( $V_{eff}$ ) and space density of the Gold Samples as a function of scale height (h). These ignore selection effects such as SDSS sky coverage (see text for details).

h (pc)	$V_{eff}$ (10 <sup>6</sup> pc <sup>3</sup> )		$\rho_0$ (10 <sup>-6</sup> pc <sup>-3</sup>	)	
	· • •	WD+M	WD+AFGK	CV	DWD
100	46.5	2.47	0.37	3.27	1.35
280	65.4	1.76	0.26	2.23	1.04
500	93.3	1.23	0.18	1.63	0.73

#### 4.1.4 Candidates with no parallax

Of the 3520 candidates (Table 1), 432 did not have a Gaia parallax. We consider the distribution of these 432 systems across the sky area and magnitude range (Fig. 5). Whereas the sky coverage of the SDSS footprint is evident (see Blanton et al. 2017), there is no apparent spatial pattern among the systems without parallax. However, the magnitude distribution reveals a clear drop in the fraction of WD+M candidates with parallaxes towards fainter systems. The parallax incompleteness is already very noticeable for  $G \gtrsim 17$ , which is substantially above the *Gaia* magnitude limit. The reason for the apparent difficulty of Gaia in measuring parallaxes of WD+M systems is not clear, but we speculate that some of those systems might be wide binaries that exhibit sufficient astrometric accelerations to foil the DR2 five-parameter astromeric solution. This hypothesis is supported by the large number of systems found between the white dwarf cooling track and the main sequence that have excessive ruwe values (see fig. 1 in Belokurov et al. 2020).

#### 4.2 PCEBs - white dwarf + A/F/G/K stars

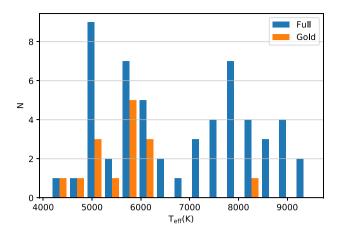
PCEBs consisting of a white dwarf and a star of spectral type A, F, G or K (hereafter WD+AFGK) present very different observational issues from those with M-dwarf companions as the optical flux of the companion star is so much greater than that of even hot and young white dwarfs. Nevertheless, the process of establishing systems as close WD+AFGK binaries is the same as for the WD+M ones, i.e. firstly identifying candidate systems and secondly establishing whether they are wide binaries or PCEBs via the appropriate follow-up observations.

#### 4.2.1 Establishing the Gold Sample

WD+AFGK binary candidates have been identified by searching for systems with an ultraviolet excess – which in most case signals the presence of the white dwarf. As PCEBs have short periods their confirmation requires the detection of radial velocity variations related to the orbital motion. Given that the companions in these systems totally dominate the optical flux, their spectra are rich in sharp photospheric features, and they are typically relatively bright compared with WD+M systems.

The identification of WD+AFGK binary candidates has been pioneered by Parsons et al. (2016) and Rebassa-Mansergas et al. (2017) who cross-matched AFGK-type stars spectroscopically confirmed by the RAVE and LAMOST surveys respectively, with the *GALEX* ultraviolet photometry. Parsons et al. (2016) obtained *HST* spectroscopy of nine of these potential WD+AFGK binaries, all of which contained white dwarfs (see also Werner et al. 2020).<sup>4</sup>

<sup>&</sup>lt;sup>4</sup>This does leave the possibility that some systems are triples. Lagos et al. (2020b) estimated this fraction to be less than 10 per cent.



**Figure 7.** The distribution of the effective temperatures of the systems in the WD+AFGK Gold Sample contrasted with the full (not volume-limited) set. The volume-limited sample has a significantly lower average effective temperature.

Follow-up spectroscopy of TGAS-selected WD+AFGK candidates led to the confirmation of 23 PCEBs (Ren et al. 2020, hereafter the "Ren sample"). Four additional PCEBs were confirmed by Parsons et al. (2015) and Hernandez et al. (2020, submitted). We included also the well-studied PCEBs V741 Tau (Nelson & Young 1970) and IK Peg (Wonnacott, Kellett & Stickland 1993), as well as the recently discovered GPX-TF16E-48 (Krushinsky et al. 2020).

We also consider in this sample the EL CVn binaries: short-period ( $\simeq 1-3$  d) binaries containing A/F-type stars with extremely lowmass (ELM) white dwarf companions. These systems are identified as eclipsing binaries (Maxted et al. 2014; van Roestel et al. 2018). EL CVn stars represent a different evolutionary channel from the WD+M and WD+AFGK binaries discussed so far, forming via dynamically stable mass transfer that initiates once the initially more massive star in a main-sequence binary enters the sub-giant branch (Chen et al. 2017; Sun & Arras 2018). There is strong evidence that most EL CVn systems are inner binaries of hierarchical triples (Lagos et al. 2020a). We therefore include the EL CVn stars listed by Ritter & Kolb (2003; hereafter the "RK sample"). As before, systems without a *Gaia* match or parallax were dropped and quality filters (equations 1–5) were applied (see Table 1).<sup>5</sup> The WD+AFGK Gold Sample contains 15 systems (Table A2).

#### 4.2.2 HR diagram

The distribution of the WD+AFGK binaries in the *Gaia* HR diagram is shown in Fig. 6. As the size of this Gold Sample is small (Table 1), we include in this diagram, purely for the purpose of visualization, systems beyond the 300 pc volume limit; these systems are not included in the Gold Sample.

WD+A binaries are located at the top of the main sequence and the intrinsic brightness of the main-sequence companions in these systems results in a significant selection effect: all but one are EL CVn binaries with a low mass pre-He white dwarf formed by stable mass transfer rather than in the common envelope phase. These systems have been almost exclusively discovered because of

<sup>5</sup>In the case of WD+AFGK binaries, the main-sequence star is intrinsically so bright that it will be reliably detected to well over 300 pc. However, for compatibility with the other Gold Samples, we use the same limiting volume.

mutual eclipses of the two stellar components. The Gold Sample contains one EL CVn star: the prototype itself.

The one non-EL CVn system among the WD+A binaries is IK Peg. Whilst it is not included in the Gold Sample because of *Gaia* quality issues, it deserves mention as it further underlines the difficulty in finding WD binaries with luminous companions. IK Peg had long been known as a single-lined spectroscopic binary with a period of 21.7 d (Harper 1928), but the white dwarf component was discovered only much later via ultraviolet spectroscopy (Wonnacott et al. 1993). The fact that IK Peg is nearby,  $\simeq$ 47 pc, suggests that WD+A PCEBs may be relatively common.

Whilst the Gold Sample contains a number of white dwarf binaries with F- and G-type companions, there is a distinct lack of systems with K-type companions. Noticeable exceptions are V471 Tau (with a K2-type companion; Kamiński et al. 2007) and GPX-TF16E-48 (with a K7-type companion; Krushinsky et al. 2020). This is probably related to selection effects caused by a gap between the two different approaches used to identify white dwarf plus main-sequence binaries (i.e. using either optical spectroscopy or ultraviolet excess detections).

Within the HR diagram (Fig. 6), the majority of the WD+AFGK Gold Sample occupies the expected well-confined area within the main sequence, albeit two Gold systems are located within the giant branch: TYC 4717-255-1 is a K0 III star (Houk & Swift 1999) at 299 pc and TYC 278-239-1 is also a K0 III star at 159 pc (Fabricius et al. 2002).

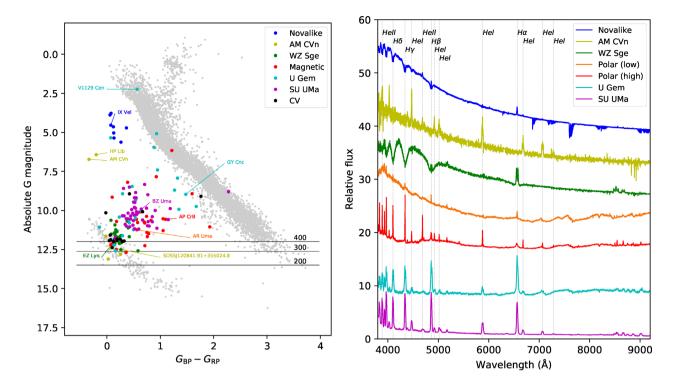
Given that the contribution of the white dwarf to the optical fluxes of the WD+AFGK binaries is negligible, the effective temperatures (taken from Gaia DR2 where available) are essentially those of the main-sequence or sub-giant stars. The Gold sample is dominated by systems with cooler, G/K-type companions (Fig. 7), with only a single WD+A system, i.e. the prototypical EL CVn. This reflects the difficulty in identifying white dwarf binaries with early type companions, the exception being the EL CVn systems which are identified photometrically via the detection of eclipses.

#### 4.2.3 The Gold sample as a fraction of the total population

As before (Section 4.1.3), we calculate the space density of this Gold Sample for three different values of h (100, 280, 500 pc; see Table 3). These figures compare with  $\rho_0 \simeq 4 - 10 \times 10^{-6} \text{ pc}^{-3}$  for the ELCVn stars alone (Chen et al. 2017) and a lower limit on WD+FGK binaries of  $\rho_0 > 1.9 \times 10^{-6} \text{ pc}^{-3}$  (Ren et al. 2020), both of which suggest that the WD+AFGK Gold Sample forms a small subset of the underlying population. Moreover, this subset is heavily biased in favour of systems where the white dwarf signal can be detected, i.e. namely hot white dwarfs and we conclude that the population of known WD+AFGK binaries is likely the most incomplete out of all the close white dwarf binary sub-types discussed in this paper.

#### 4.3 Cataclysmic variables

Historically CVs, or more specifically CV candidates, were primarily discovered because of their large-amplitude optical variability, either by visual observers, or from the comparison of photographic plates taken at different epochs (Breedt et al. 2014). Examples of early discoveries include the dwarf nova U Gem (Pogson 1857) and SS Cyg (Pickering & Fleming 1896). Variability is still a major source of new CV candidate discoveries, with time-domain surveys such as CRTS (Drake et al. 2014), ASAS-SN (Shappee et al. 2014), ZTF



**Figure 8.** Left-hand panel: The HR diagram of the CV Gold Sample shows that the different sub-types cluster in preferred regions of parameter space. There is insufficient data to classify the objects labelled "CV". The location of these clusters arises from the different SEDs of the various sub-classes, which are a function of the properties of both stellar components and the accretion flow. The minimum absolute magnitude to be detectable at an apparent magnitude of G = 20 is indicated by grey lines for distances of 200, 300, and 400 pc. Right-hand panel: Typical examples of the optical spectra for each sub-type. The spectra have been normalized and stacked applying appropriate vertical offsets. The location of each individual CV is identified in the HR diagram. The spectrum of the novalike (=IX Vel) was obtained using X-Shooter (Vernet et al. 2011). The spectra of the AM CVn (=SDSS J120841.91+355024.8), two polars (AR UMa in a low state and AP CrB in a high state), and three dwarf novae of the sub-types U Gem, SU UMa, and WZ Sge (=GY Cnc, BZ UMa, and EZ Lyn, respectively) were obtained from SDSS DR16 (Ahumada et al. 2020).

(Bellm et al. 2019), and *Gaia* alerts (Hodgkin et al. 2013) continually generating new candidates for investigation. Given that a number of phenomena can be mistaken for CVs, such as minor planets (Schmadel, Schmeer & Börngen 1996) or large-amplitude pulsating stars (Dupke 1992), the unambiguous confirmation of CVs typically requires follow-up spectroscopy (Zwitter & Munari 1994; Szkody et al. 2020).

CV candidates discovered because of their optical variability are biased towards systems that undergo frequent, and/or largeamplitude brightness changes. Alternative characteristic hallmarks of CVs are their blue colour, Balmer and helium emission, and X-ray emission, which allow their identification in multicolour photometric surveys (e.g. Green, Schmidt & Liebert 1986), spectroscopic surveys (e.g. Aungwerojwit et al. 2006; Szkody et al. 2009), and X-ray surveys (Schwope et al. 2002; Landi et al. 2009). The relative contributions of the different CV identification methods to the overall population of known CVs and their affinity with different CV subclasses are discussed by Gänsicke et al. (2005). Despite all efforts, CVs that mimic either single white dwarfs, or main-sequence stars, and have low accretion rates remain very difficult to discover (e.g. Rebassa-Mansergas et al. 2014; Teyssier 2019; Yu et al. 2019).

#### 4.3.1 Establishing the Gold Sample

An extensive literature review was conducted as part of Pala et al. (2020) to obtain a comprehensive inventory of both confirmed CVs and CV candidates. This list has been updated for new CV candidates that have been discovered since that study and contains 5192 systems

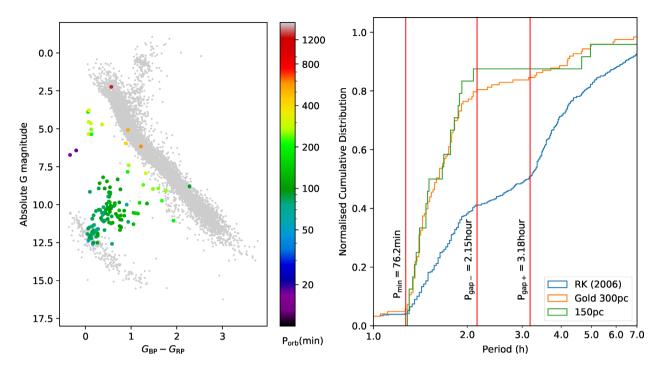
of which 4042 have a *Gaia* parallax (see Table 1). Applying the distance limit, we were left with 305 CV candidates. We then applied the quality filters outlined in Section 3 resulting in 201 systems.

As discussed by Pala et al. (2020), a fraction of CV candidates reported in the literature are contaminants, including flaring Mdwarfs, young stellar objects, single white dwarfs, and detached binaries being mistakenly identified as CVs. In order to establish the CV Gold Sample, we therefore proceeded to validate the CV nature of these 201 candidates.

We began by adopting the CV classification for the 42 systems with  $d \leq 150$  pc, that were already vetted by Pala et al. (2020), as well as 139 with 150  $\leq d \leq 300$  pc that were in the catalogue of bona-fide CVs compiled by Ritter & Kolb (2003). We then extracted light curves, where available, for the 124 remaining candidates from ZTF (Masci et al. 2019), ASAS-SN (Jayasinghe et al. 2019; Masci et al. 2019), and CRTS (Masci et al. 2019) together with spectra from SDSS (Smee et al. 2013) and LAMOST (Zhao et al. 2012). We found a total of 152 CVs with reliable spectra and a history of variability and these are henceforth referred to as the CV Gold Sample (Table A3).

The CVs in the Gold Sample have reliable parallax measurements with an average distance uncertainty (based on the estimated errors in Bailer-Jones et al. 2018) of 2.7 per cent.

In passing, we note that we have excluded  $\simeq 50$  potential CV candidates from the Gold Sample primarily due to a lack of spectroscopy. Follow-up observations are planned and will be reported in a subsequent paper. We also expect that a number of the CV candidates rejected due to poor data quality will be confirmed by Gaia EDR3 astrometry.



**Figure 9.** Left-hand panel: An HR diagram showing the distribution of CVs in the Gold Sample with their measured orbital periods colour-coded (in minutes). As CVs evolve towards shorter periods, they move from the main sequence towards the white dwarf cooling track. Right-hand panel: Comparison of cumulative orbital period distributions of the 150 pc sample (green, Pala et al. 2020), the 300 pc CV Gold Sample (orange), and the Ritter & Kolb (2003) sample (blue – we use V7.6 which enables a direct comparison with fig. 4 in Knigge 2006), which is a heterogeneous mix of CVs identified by all identified selection methods. The period gap starting at 2 h is very evident in each case.

#### 4.3.2 The CV Gold Sample in the Gaia HR diagram

The distribution of the CV Gold Sample within the *Gaia* HR diagram is shown in Fig. 8. We re-iterate that the CV Gold Sample is not intended to be complete, but to establish the parameter space that CVs inhabit and to assess how their location within this parameter space links to their physical properties (the evolution of CVs across the HR diagram is discussed in more detail in Section 5). We review how well the CV Gold Sample meets those goals in four ways – first, by checking that the variety of sub-types are well-represented, secondly, by checking that there is a spread in evolutionary states (for which orbital period is a proxy), thirdly, by comparing the number of CVs in the Gold Sample with the estimated size for a complete 300 pc CV sample, and lastly by investigating whether the CV candidates from our initial list that do not have a parallax introduce some major bias in the CV Gold Sample.

*CV subtypes*: The CV Gold Sample contains examples of all six major different sub-types, their respective numbers are listed in Table 4. The different sub-types form clearly distinct clusters in the *Gaia* HR diagram (Fig. 8, left-hand panel);<sup>6</sup> and the location of these clusters closely links to the spectro-photometric properties of each sub-type (Fig. 8, right-hand panel).

The novalike variables are characterized by stable high mass transfer rates and hot accretion discs that dominate the optical emission of these systems. Correspondingly, their cluster is located at small absolute magnitudes and blue colours. Two AM CVn stars,

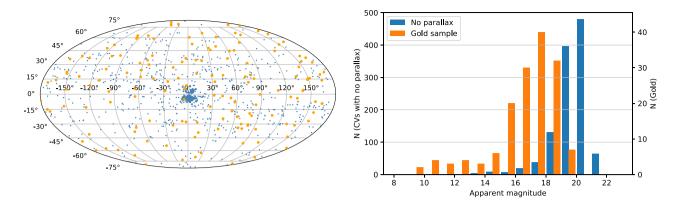
<sup>6</sup>A similar HR diagram has been shown by Abril et al. (2020), presenting overall similar features. However, these authors included a much larger number of CVs with much looser selection criteria than those defining the Gold Sample. Consequently, the clustering in their diagram is less pronounced.

**Table 4.** The number of each type in the CV GoldSample together with the range of their orbital periods.

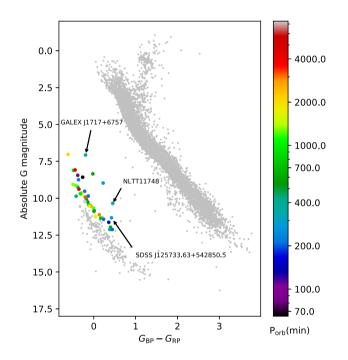
Туре	Ν	Period (h)				
		Median	Minimum	Maximum		
Novalike	9	4.42	1.51	5.88		
AM CVn	6	0.31	0.29	0.88		
WZ Sge	26	1.37	1.12	1.82		
Magnetic	28	1.89	1.37	9.88		
U Gem	23	2.09	0.99	21.43		
SU UMa	45	1.63	1.07	2.81		
CV	14	1.44	1.3	4.26		

HPLib and the prototype AMCVn itself (Faulkner, Flannery & Warner 1972) are novalike variables with stable, hot helium discs that are somewhat bluer and fainter than their cousins with hydrogen-rich discs.

The three sub-classes of dwarf novae exhibit clear trends in their location within the HR Diagram. The UGem type dwarf novae have lower mass transfer rates than the novalike variables, their quiescent discs are therefore less luminous and their donor stars can contribute noticeably to the optical flux of the systems, which is reflected in the number of these systems being located close to the main sequence. The donors and accretion discs in the SU UMa systems are less luminous than those in UGem-type dwarf novae, hence they are concentrated closer to the white dwarf cooling sequence. Finally, the WZ Sge type dwarf novae have the dimmest accretion discs and donor stars; in fact, some of them contain brown dwarf companions (GD552 Unda-Sanzana et al. 2008; QZ Lib (Pala et al. 2018), 1RXS J105010.3–140431 (Mennickent et al. 2001), SDSS J102905.21+485515.2 (Thorstensen, Alper & Weil 2016),



**Figure 10.** Left-hand panel: The distribution of the CV Gold Sample (orange) and the CV candidates without a valid parallax (blue) in Galactic coordinates. The concentration of systems with no parallax close to the Galactic centre is very noticeable, with  $\simeq 44$  per cent contained within the area bounded by -12 < l < 12 and -8 < b < 5. Most ( $\simeq 79$  per cent) of this concentration were identified by the OGLE survey (Mróz et al. 2015). Right-hand panel: The distribution of apparent magnitudes of the Gold Sample contrasted with the CVs without a parallax. It is evident that most of the CVs without a parallax have apparent magnitudes  $\gtrsim 19$ , i.e. near the limiting magnitude of *Gaia*.



**Figure 11.** The Gold Sample of DWDs with known orbital periods (colourcoded) are located slightly above the general white dwarf cooling sequence in the HR diagram, reflecting the combined luminosity of the two white dwarfs. The lack of DWDs at the faint end of the white dwarf sequence is an observational selection effect.

SDSS J143317.78+101123.3 (Hernández Santisteban et al. 2016), and SSS J122221.7–311525 (Neustroev et al. 2017)). WZ Sge are therefore inherently the faintest CVs and merge into the white dwarf cooling sequence.

A noticeable outlier among the dwarf novae is V1129 Cen, a long-period ( $\simeq$ 21.4 h) CV that exhibits low-amplitude ( $\simeq$ 0.6 mag) outbursts about once per year, but no emission lines are detected in its spectrum (Bruch 2017). The donor is a luminous F-type+ star, explaining the location in the HR diagram.

Magnetic CVs have a larger spread due to their more varied spectral appearances. During states of low mass transfer they resemble the detached WD+M binaries whilst during high states, their optical spectra can contain cyclotron emission lines (Wickramasinghe 1988) that will affect their colour and absolute magnitude.

Overall, the CV Gold Sample contains a representative mix of the different CV sub-types that well sample the spread of the underlying CV population in the HR diagram.

Distribution of periods: As outlined earlier, the evolution of CVs is driven by orbital angular momentum loss, which results in their periods decreasing down to the period minimum before reversing to longer periods as the donor stars become degenerate. The orbital periods of the Gold Sample (Fig. 9) span the typical range occupied by CVs,  $\simeq 80$  to  $\simeq 600$  min, and the distribution of the periods broadly resembles that of the volume-limited 150 pc sample (Pala et al. 2020). However, both the distributions of the 150 and 300 pc Gold Sample differ markedly from that of the overall population of known CVs, drawn from the Ritter & Kolb (2003) catalogue, which displays a steep increase in the number of CVs with periods in the range 3-4 h. This difference highlights the importance of using a volume-limited sample for population studies, as intrinsically bright novalikes (with mass transfer rates that, in fact, exceed the predictions of the standard CV evolution model – see Townsley & Gänsicke 2009; Knigge et al. 2011) are detected out to much larger distances than other CVs and dominate the period distribution of the Ritter & Kolb sample.

We note, in passing, that the upper edge of the canonical period gap at  $\simeq 3$  h is much less pronounced in the volume-limited sample when compared to the overall population of CVs in the RK catalogue, which is subject to a multitude of selection effects. This underlines once more that great care has to be taken when using heterogeneous CV samples to underpin generic CV evolution models.

The Gold Sample as a fraction of the total population: In the same manner as the other Gold Samples, we compute the space density of the Gold CVs for three different values of h (100, 280, 500 pc), resulting in  $\rho_0 \simeq 1.6 - 3.3 \times 10^{-6} \text{ pc}^{-3}$  (see Table 3). These values should be compared with  $\rho_0 = 4.8 \times 10^{-6} \text{ pc}^{-3}$  for the volume-limited 150 pc sample (Pala et al. 2020), and suggest that the CV Gold Sample represents a substantial fraction ( $\simeq$ 50 per cent) of the underlying population within 300 pc.

#### 4.3.3 Candidates with no parallax

Among the 5192 CV candidates that we began with 1150 have no *Gaia* parallax and an important question is whether any of those are

within 300 pc. A characteristic of *Gaia* is that the *G*-band magnitude of variable systems with only a two-parameter solution (i.e. no parallax) will tend to be the (brighter) outburst magnitude not the (fainter) quiescent one. Fig. 10.10 in Antoja et al. (2018) illustrates this effect and upon closer inspection of a sample, we have found that the majority of CVs with *G*-band magnitudes in the range  $\simeq 17$ – 20 (Fig. 10, left-hand panel), and no parallax have indeed fainter quiescent magnitudes,  $\gtrsim 20$ , in SDSS and PanSTARRS1 imaging, suggesting that the *Gaia* detections were obtained during outburst(s) or high states.

For an assumed distance  $\lesssim 300 \text{ pc}$ , these very faint quiescent apparent magnitudes would imply that these systems are intrinsically extremely faint with low mass transfer rates, cool white dwarfs, and very low-mass donors. Whilst we cannot rule out that a small number of intrinsically faint WZ Sge type dwarf novae with distances near the volume-limit fall into this category, inspection of the spatial distribution of the 1150 CVs without parallaxes shows a strong concentration towards the Galactic Centre (Fig. 10, right-hand panel). CVs within a 300 pc volume limit are expected to be approximately isotropic in their sky distribution and we conclude that the majority of CVs without a parallax are beyond our adopted distance limit.

#### 4.4 Double white dwarfs

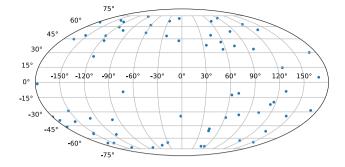
Over recent years, many DWDs have been documented in the literature. Volumetric surveys of white dwarfs encompassing 20 pc (Hollands et al. 2018) and 40 pc (McCleery et al. 2020) contain a number of confirmed DWDs and these were supplemented by a collection of systems from various other publications, which include most of the known DWDs. As before (see table 1), we remove duplicates and systems for which there is no *Gaia* parallax or which fail our quality criteria outlined in Section 3. The remaining 67 form the Gold Sample (Table A5).

#### 4.4.1 The DWD Gold Sample in the Gaia HRD

A system consisting of two white dwarfs typically has a similar colour to a single white dwarf but is somewhat (up to  $\simeq 0.75$  mag depending upon their temperatures and to a lesser extent masses) brighter and therefore appears slightly higher in the HR diagram (Fig. 11). The four outliers bounded by NLTT11748 (Kawka, Vennes & Vaccaro 2010; Kaplan et al. 2014) and GALEX J1717+6757 (Vennes et al. 2011) include an extremely low mass (ELM) white dwarf (Brown et al. 2020; Kosakowski et al. 2020). ELM white dwarfs (typically  $\sim 0.2 M_{\odot}$ ) have greater luminosities than more massive white dwarfs because of their larger radii ( $R \propto M^{-1/3}$ ). This increased luminosity therefore causes them to appear above the white dwarf cooling sequence in the HR diagram.

The lack of DWDs at the cool end ( $T_{\rm eff} \leq 7000 \, {\rm K}$ ) of the white dwarf sequence is a selection effect. The vast majority of DWDs were identified as spectroscopic binaries (e.g. Morales-Rueda et al. 2005; Napiwotzki et al. 2020); only a very small number were photometrically detected because of eclipses (e.g. Steinfadt et al. 2010) or because of inconsistencies in their spectroscopic features (Bergeron, Greenstein & Liebert 1990). At low temperatures, the Balmer lines weaken and eventually disappear, therefore the confirmation of DWD candidates (typically selected because they are overluminous; e.g. Giammichele, Bergeron & Dufour 2012; Hollands et al. 2018) becomes very difficult for cool white dwarfs.

Only 38 of the DWDs in the Gold Sample have an orbital period measurement and there is no apparent correlation between period and position in the HR diagram (Fig. 11).



**Figure 12.** The distribution of the DWD Gold Sample in Galactic coordinates. They appear to be evenly distributed away from the Galactic plane. The lack of DWDs near the Galactic plane is a selection effect due to the difficulty in identifying white dwarf candidates in crowded regions.

The DWD Gold Sample appears to be evenly distributed away from the Galactic plane (Fig. 12). The lack of DWDs near the Galactic plane is a selection effect due to the difficulty in identifying white dwarf candidates in crowded regions

#### 4.4.2 The Gold Sample as a fraction of the total population

As before (Section 4.1.3), we calculate the space density for three different values of h (100, 280, 500 pc; see Table 3). These figures compare with  $> 0.62 \times 10^{-3} \, \text{pc}^{-3}$  from Holberg et al. (2016) which demonstrates that the DWD Gold Sample forms a tiny subset of the underlying population. This is partly a reflection of the inherently low luminosity of these systems – Holberg et al. (2016) derived their estimates from a volume-limited survey of 25 pc – and partly, the difficulty in distinguishing DWDs from single white dwarfs.

#### 4.4.3 Candidates with no parallax

Only 6 of the 167 candidates had no *Gaia* parallax and in each case, their *G*-band magnitude was larger than any of the members of the Gold Sample. It is therefore likely that they are all at a distance of more than 300 pc.

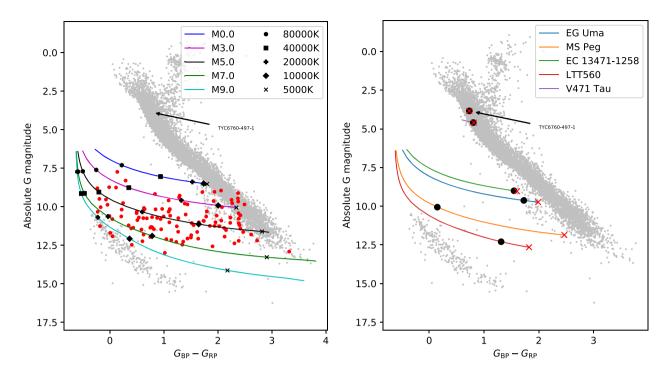
This volume-limited sample is weighted towards normal mass white dwarfs rather than the rare, but more luminous ELM DWDs that are detectable out to larger distances, again highlighting the difference between a volume-limited survey and a magnitude-limited one.

# 5 EVOLUTION OF WHITE DWARF BINARIES WITHIN THE HR DIAGRAM

The discussion, so far, has focused on establishing and validating the various white dwarf binary sub-samples, and to discuss their observational properties. Here, we will briefly investigate the evolution of white dwarf binaries across the HR diagram, as well as the effects of systems moving from one sub-sample into another one.

### 5.1 PCEBs

A PCEB emerges from the common envelope as a detached binary consisting of a hot ( $\simeq 100\,000\,\text{K}$ ) white dwarf and a main-sequence star companion with a spectral type A–M (and a small number of brown dwarfs). These systems have typical orbital periods of hours to days, are hence always spatially unresolved, and *Gaia* will consequently detect and characterize them based on the combined



**Figure 13.** Left-hand panel: The evolution of WD+M PCEBs across the *Gaia* HR diagram is illustrated by the coloured tracks showing a range of companion spectral types; the black markers indicate a sequence of white dwarf temperatures. White dwarfs cool very fast initially and then gradually slow down (the cooling ages for a  $0.6 M_{\odot}$  white dwarf for the five temperatures annotated in the left-hand panel are  $\simeq 3.4 \times 10^5$ ,  $4.2 \times 10^6$ ,  $6.2 \times 10^7$ ,  $6.3 \times 10^8$ ,  $6.3 \times 10^9 \text{ yr}^8$ . This results in a selection effect as most (i.e. except the youngest) WD+M0 PCEBs will be seen near the main sequence whilst white dwarfs with less massive companions of the same age will be found in the area between the main sequence and the white dwarf cooling sequence. As the white dwarfs cool and the flux contribution of the M-dwarfs remains constant, the systems fade and become redder, moving to the right in the HR diagram, until eventually merging into the main sequence. The WD+M Gold Sample is shown as red dots. Right-hand panel: Evolutionary paths for five PCEBs adapted from Schreiber & Gänsicke (2003) and Zorotovic, Schreiber & Gänsicke (2011). Their current locations in the HR diagram are indicated by the black dot; and the positions where they will evolve into CVs once the companion star becomes Roche lobe filling are shown by the red crosses. TYC 6760-497-1 (Parsons et al. 2015) is a WD+F type. Due to its high mass ratio, this system will undergo a short phase of thermal-time-scale mass transfer during which the white dwarf will accrete at a rate that is sufficiently high to sustain nuclear shell burning (Fujimoto 1982; Iben 1982) – and will be a super-soft X-ray source, and correspondingly bright and blue (Kahabka & van den Heuvel 1997).

flux from both stars. In the optical passbands of *Gaia*, PCEBs with K-type or earlier companions are indistinguishable from single mainsequence stars even in the early stages of their evolution (Section 4.2).

As angular momentum losses drive the evolution of these binaries towards shorter orbital periods, the white dwarf gradually cools, fades, and becomes redder in colour – a process that takes  $\simeq$ 415 Myr–2.6 Gyr to reach 10 000K for white dwarfs with masses of 0.4 – 1.2 M<sub>☉</sub>. In contrast, companion stars with masses  $\lesssim 0.9 M_{\odot}$  will not significantly evolve; and their flux contribution consequently remains constant throughout the evolution of the binary.

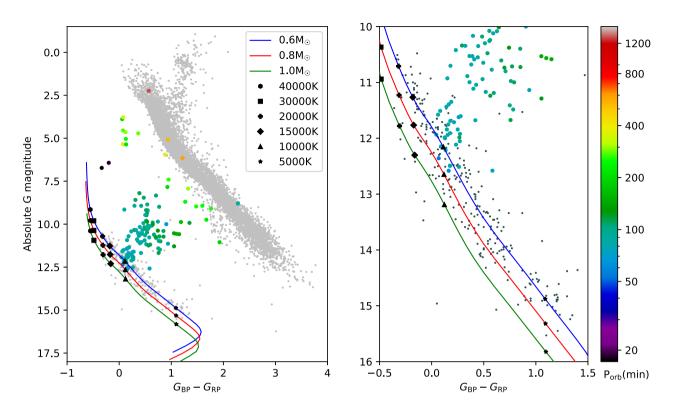
We have computed the evolutionary tracks of a range of WD+M binaries by combining the synthetic *G*,  $G_{\rm BP}$ , and  $G_{\rm RP}$  magnitudes for a typical 0.6 M<sub>☉</sub> white dwarf (Bergeron et al. 2011; Tremblay, Bergeron & Gianninas 2011), with the cooling models provided by Pierre Bergeron.<sup>7</sup> The M-dwarf companions were modelled using the Phoenix (Husser et al. 2013) and ATLAS9 (Castelli & Kurucz 2003) grids.

Comparison between these evolutionary tracks and the WD+M Gold Sample (Fig. 13, left-hand panel) shows reasonably good agreement, highlighting that WD+M PCEBs containing early or late M-type companions are relatively rare (see also fig. 9 in Rebassa-Mansergas et al. 2010). This is likely to be a combination of the slope of the initial mass function of low-mass stars, which drops off for very low stellar masses (Kroupa, Tout & Gilmore 1993)<sup>8</sup> and the selection effects discussed in Section 4.1. Similarly, WD+M PCEBs with very hot white dwarfs are relatively rare in the Gold Sample (again, see fig. 9 in Rebassa-Mansergas et al. 2010), which is a consequence of the fact that white dwarfs cool quasi-exponentially together with the selection effects examined in Section 4.1.

Once the orbital separation reduces sufficiently for the companion to fill its Roche lobe, the system will morph into a CV, and the onset of mass transfer will drastically alter its properties within the HR diagram (Sections 4.3 and 5.2). We have used the parameters of five PCEBs in the WD+M and WD+AFGK Gold Samples to compute their tracks within the HR diagram, illustrating both their past evolution as well as the location at which they will transition into CVs and hence disappear from the PCEB population (Fig. 13, right-hand panel). Except for LTT 560, which contains a relatively cool M-dwarf, the onset of mass transfer will occur close to the main sequence, when the white dwarf has cooled below 10 000 K

<sup>8</sup>Farihi, Becklin & Zuckerman (2005b) already noted (see their figs 6 and 7) the very similar distribution in M-dwarf spectral types among a 20 pc volumelimited sample of field M-dwarfs and their sample of WD+M binaries.

<sup>&</sup>lt;sup>7</sup>http://www.astro.umontreal.ca/~bergeron/CoolingModels (Fontaine, Brassard & Bergeron 2001;Holberg & Bergeron 2006; Kowalski & Saumon 2006; Tremblay et al. 2011; Bédard et al. 2020).



**Figure 14.** The evolutionary channel from the main sequence is clearly evident with a wide range of long period CVs near the main-sequence funnelling in as their periods fall towards a small locus in the white dwarf area of the HR diagram (left-hand panel with expanded view in the right-hand panel). By this time, CVs have become dominated by their white dwarfs, following prolonged mass transfer. The effective temperature of the CV is essentially that of the white dwarf and is focused towards a region with  $T_{\rm eff} \simeq 10\,000 - 15\,000$  K. Subsequent evolution follows the cooling track of a white dwarf – the cooling tracks of the white dwarf component are shown for white dwarf masses of 0.6, 0.8, and  $1.0 M_{\odot}$  (Bergeron et al. 2011).

(Schreiber & Gänsicke 2003; Zorotovic et al. 2010). V471 Tau illustrates that PCEBs with a K-type companion exit the common envelope very close to the main sequence, and rapidly descend into it (even though the white dwarf in this system is still fairly hot,  $\simeq 35\,000$  K; O'Brien, Bond & Sion 2001).

Although detached white dwarf plus main-sequence binaries are mostly formed via common envelope evolution, the standard theory of CV evolution predicts that some WD+M binaries with periods of  $\simeq 2-3$  h are actually CVs in which the donor detached from its Roche lobe (Rappaport et al. 1983; Davis et al. 2008; Zorotovic et al. 2016). Angular momentum loss will still drive the evolution of these systems towards shorter orbital periods until the donor fills its Roche lobe again at  $P_{\rm orb} \simeq 2$  h, when they will transit back into the CV population. Abril et al. (2020) have suggested that these "interlopers" among the WD+M PCEBs population may be identified because of their, on average, higher white dwarf mass (Zorotovic et al. 2011), resulting in larger absolute magnitudes compared to the WD+M PCEBs that have not yet reached the first onset of mass transfer.

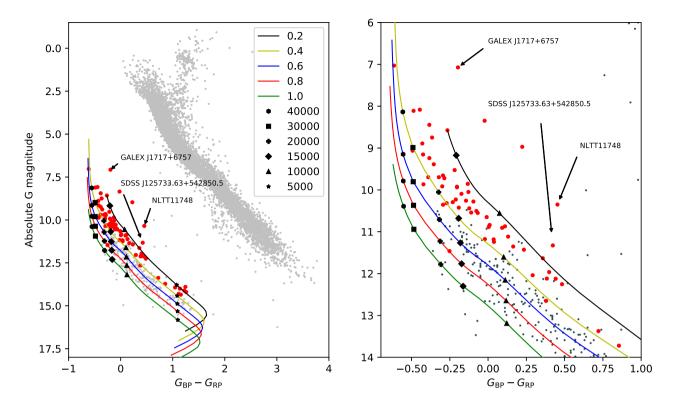
A final note concerns those white dwarfs with companions that are sufficiently massive ( $\gtrsim 0.9 \, M_{\odot}$ ) to evolve off the main sequence within a Hubble time. In the *Gaia* HR diagram, those systems will emerge from the common envelope within the main-sequence. Given their relatively small orbital separations, they will initiate mass transfer before the companion evolves significantly into the giant branch, and will hence transform into CVs "*in situ*". An interesting example is the WD+F binary TYC 6760-497-1 (Parsons et al. 2015): because of the high mass ratio,  $M_2/M_{wd} \simeq 2$ , this system will undergo a short phase of thermal-time-scale mass transfer during which the white dwarf will accrete at a rate that is sufficiently high to sustain nuclear shell burning (Fujimoto 1982; Iben 1982) – and will be a super-soft X-ray source, and correspondingly bright and blue (Kahabka & van den Heuvel 1997). There is no super-soft source known within the 300 pc CV Gold Sample.

#### 5.2 CVs

CVs commence mass transfer resulting in the formation of a luminous accretion flow and heating of the surface of the white dwarf (Townsley & Bildsten 2003). As discussed in Section 4.3, the location of a CV within the HR diagram depends on the relative contributions of the white dwarf, the donor star, and the accretion flow to the overall flux.

As shown in the previous section, the location of most CVs within the HR diagram at the onset of mass transfer is expected to be close to, or within, the main sequence. Once a companion in a PCEB fills its Roche lobe, the system will very quickly move blue-wards, and, if the mass transfer rate is within the regime of dwarf novae, sit slightly below the main sequence (Fig. 8). The distinct cluster of novalike CVs (see Fig. 8) with bright, blue accretion discs and very similar orbital periods of  $\simeq 3-4$  h, i.e. sandwiched between the longer period U Gem type dwarf novae and the shorter period SU UMa type dwarf novae, is noteworthy: these systems appear to have mass transfer rates that exceed the predictions of the standard evolution model (Townsley & Gänsicke 2009), and very likely reflect a peculiar phase in CV evolution, and/or a specific set of initial parameters.

As CVs evolve the mass, radius, and effective temperature of the secondary decrease, and so does the mass transfer rate, and hence the evolutionary curve moves back towards the white dwarf



**Figure 15.** The position of the DWD Gold Sample on the HR diagram together with the cooling tracks of single white dwarfs for masses of 0.2, 0.4, 0.6, 0.8, and  $1.0 \, M_{\odot}$  (Bergeron et al. 2011). The colours of the DWDs will be the weighted average of the colours of its constituents whilst the absolute magnitude will be the sum of their fluxes. The position will therefore be similar to that of a single white dwarf but  $\simeq 0.75 \, \text{mag}$  brighter – in practice, the more massive white dwarf will have formed first and hence have progressed further down the cooling track leaving the secondary dominant. The clumping between the 0.2 and  $0.4 \, M_{\odot}$  cooling curves indicates that the brighter star in these systems typically has a low mass ( $\lesssim 0.4 \, M_{\odot}$ ), highlighting possible observational selection effects. *GALEX* J1717+6757 and NLTT11748 are ELMs with  $M_{wd} \lesssim 0.2 \, M_{\odot}$ , which, despite containing relatively cool white dwarfs ( $\simeq 10\,000 - 15\,000 \, \text{K}$ ) are substantially more luminous, thanks to their large radii (Kawka et al. 2010; Steinfadt et al. 2010; Vennes et al. 2011; Hermes et al. 2014).

cooling sequence. This explains the funnel-shaped distribution of CVs in Fig. 9 (see also Abril et al. 2020), which clearly illustrates the decrease of  $P_{\rm orb}$  as the systems come closer to the white dwarf cooling sequence. An interesting question is whether the 2–3 h period gap, i.e. when CV donors detach from their Roche lobes and the systems turn into detached WD+M binaries (see previous section), will leave a detectable imprint in the population of CVs within the *Gaia* HR diagram. We would argue that, at the moment, the number of CVs within a well-defined sample with accurate *Gaia* data is not sufficiently large to assess this question; and we caution that the intrinsic variability of CVs might smear out the effect of the period gap within the HR diagram.

As the mass of the donors continues to decrease, they will eventually change into brown dwarfs.<sup>9</sup> These are inherently faint systems, and their optical appearance is entirely dominated by the accretion-heated white dwarf, hence their location lies very close to or within the white dwarf cooling track. It is noticeable that the evolutionary path of CVs ends in a small area of the HR diagram corresponding to white dwarf temperatures of  $\simeq 10\,000$  K (Fig. 14, right-hand panel). The mass transfer rates implied by these accretion

rates,  $\dot{M} \simeq 5 \times 10^{-11} \,\mathrm{M_{\odot} yr^{-1}}$  are consistent with the predictions for systems that have reached the period minimum, however, as CVs evolve back towards longer periods (Howell et al. 2001), the mass transfer rates are predicted to drop by an order of magnitude – however the known "period bouncers" have mass accretion rates much higher than these predictions. This exhibits the problem of the "missing period bouncers" (Howell et al. 2001; Pala et al. 2020).

A final note concerns the evolution of AMCVn stars within the Gaia HR diagram. As they may be formed via three different channels, their locations at the onset of mass transfer may radically differ. In the case of a double-degenerate progenitor (Paczyński 1967; Deloye et al. 2007), they will start mass transfer in, or close to, the white dwarf cooling sequence. In the case of a white dwarf plus a helium-burning star (Savonije, de Kool & van den Heuvel 1986; Yungelson 2008), they would probably start mass transfer within the region occupied by subdwarfs ( $M_G \simeq 4, G_{\rm BP} - G_{\rm RP}$  $\simeq$  -0.5; Geier 2020). Finally, if they descend from CVs with nuclearly evolved donor stars (Podsiadlowski, Han & Rappaport 2003; Goliasch & Nelson 2015), they would begin mass transfer within the main sequence, and TYC 6760-497-1 (Parsons et al. 2015) may be a relevant example (Fig. 13). The AM CVn stars within the Gold Sample appear bi-modal, with two luminous blue "high state" systems, and the remaining four "low state" systems close to the white dwarf cooling sequence. This may be a result of small number statistics - the larger sample in fig. 2 of Ramsay et al. (2018) exhibits a larger spread - however, that sample is very likely subject to

<sup>&</sup>lt;sup>9</sup>The CV Gold Sample contains at least five potentially "period bouncer" systems with brown dwarf donors: 1RXS J105010.3-140431 (Mennickent et al. 2001), GD 552 (Unda-Sanzana et al. 2008), SDSS J102905.21+485515.2 (Thorstensen et al. 2016), SDSS J143317.78+101123.3 (Hernández Santisteban et al. 2016), and QZ Lib (Pala et al. 2018).

observational biases. The two luminous AM CVn stars within the Gold Sample (AM CVn with a period of 17.14 min and HP Lib with a period of 18.38 min) have high mass transfer rates that maintain their accretion discs in a hot, stable state – in other words, they are miniature novalike variables. They appear bluer and less luminous than the "normal" novalikes with typical periods of 3–4 h, reflecting their much more compact size with their spectral energy distribution (SED) dominated by very small and hot accretion discs.

#### 5.3 DWDs

Double white dwarfs tend to sit towards, or just above, the bright edge of the white dwarf cooling sequence. Just as single white dwarfs, they will evolve along the cooling track, though not necessarily follow any single-mass track, as the relative contribution of the two stars gradually changes (Fig. 15). With the exception of the extremely low mass white dwarfs mentioned earlier (well above even the  $0.2 M_{\odot}$  cooling track), the majority of the Gold Sample sits around the cooling curves of  $0.2 - 0.4 \,\mathrm{M_{\odot}}$  single white dwarfs. This implies that the brighter dwarf is of low mass - the companion may be an older, and cooler, higher mass white dwarf (see Napiwotzki et al. 2020). This almost certainly represents a selection effect, as the brighter low-mass white dwarfs in these systems are easier to find. There are a few DWDs scattered within the cooling curves of higher mass white dwarfs, which suggests that the Gold Sample captures the range of mass ratios found among DWDs, though may not be representative. An exception is SDSS J125733.63+542850.5.

Fig. 15 inBours et al. (2015) which paradoxically has a low mass white dwarf that is older ( $\simeq$ 5 Gyr) than the more massive one ( $\sim$ 1 Gyr), highlights the potentially complex intricacies of the formation of DWDs. As noted in Section 4.4.1, the Gold Sample has a sharp cut-off at  $G_{\rm BP} - G_{\rm RP} \simeq 0.6$ , corresponding to  $\simeq$ 7000 K due to the disappearance of the Balmer lines.

#### 6 CONCLUSIONS

We have defined "Gold" reference samples within  $d < 300 \,\mathrm{pc}$  of the most common populations of close white dwarf binaries: CVs. PCEBs, and DWDs. We have used those samples, which encompass most of the parameter space of physical properties (component masses, orbital periods, ages) spanned by close white dwarf binaries, to map out their distribution in the Gaia HR diagram. CVs are inherently easier to discover than PCEBs and DWDs. Whereas, the CV Gold Sample is likely to be  $\simeq$ 50 per cent complete, this figure drops to  $\sim 10$  per cent for PCEBs (with a large uncertainty) and less than 0.2 per cent for DWDs. We have furthermore explored how close white dwarf binaries evolve through the Gaia HR diagram, criss-crossing between the white dwarf cooling track and the main sequence and occupying similar parameter spaces during different evolutionary phases. We have investigated the differences between magnitude and volume-limited close white dwarf binary samples and we have demonstrated the importance of a volume-limited sample to further improve our understanding of the evolution of these systems. We will use the Gold Samples established here to develop selection functions that will reliably identify close white dwarf binary candidates for spectroscopic and photometric confirmation, with the overarching goal to assemble a volumetric 300 pc sample that is representative of the underlying population of these systems, and sufficiently large to provide accurate statistical insight into the relative numbers, and physical properties of the various subtypes.

#### ACKNOWLEDGEMENTS

BTG and TRM were supported by the UK STFC grant ST/T000406/1 and by a Leverhulme Research Fellowship. BTG also acknowledges the support of a Leverhulme Research Fellowship. RR has received funding from the postdoctoral fellowship programme Beatriu de Pinós, funded by the Secretary of Universities and Research (Government of Catalonia) and by the Horizon 2020 programme of research and innovation of the European Union under the Maria Skłodowska-Curie grant agreement no. 801370. EB was supported by UK Research and Innovation grant ST/S000623/1. Based on observations made with ESO Telescopes at the La Silla Paranal Observatory under programme ID 094.D-0344

This work has made use of data from the European Space Agency (ESA) mission *Gaia* (https://www.cosmos.esa.int/gaia), processed by the Gaia Data Processing and Analysis Consortium (DPAC, https: //www.cosmos.esa.int/web/Gaia/dpac/consortium). Funding for the DPAC has been provided by national institutions, in particular the institutions participating in the Gaia Multilateral Agreement. We acknowledge ESA *Gaia*, DPAC, and the Photometric Science Alerts Team (http://gsaweb.ast.cam.ac.uk/alerts)

Funding for the Sloan Digital Sky Survey IV has been provided by the Alfred P. Sloan Foundation, the U.S. Department of Energy Office of Science, and the Participating Institutions. SDSS-IV acknowledges support and resources from the Center for High-Performance Computing at the University of Utah. The SDSS web site is www.sdss.org.

SDSS-IV is managed by the Astrophysical Research Consortium for the Participating Institutions of the SDSS Collaboration including the Brazilian Participation Group, the Carnegie Institution for Science, Carnegie Mellon University, the Chilean Participation Group, the French Participation Group, Harvard-Smithsonian Center for Astrophysics, Instituto de Astrofísica de Canarias, The Johns Hopkins University, Kavli Institute for the Physics and Mathematics of the Universe (IPMU) / University of Tokyo, the Korean Participation Group, Lawrence Berkeley National Laboratory, Leibniz Institut für Astrophysik Potsdam (AIP), Max-Planck-Institut für Astronomie (MPIA Heidelberg), Max-Planck-Institut für Astrophysik (MPA Garching), Max-Planck-Institut für Extraterrestrische Physik (MPE), National Astronomical Observatories of China, New Mexico State University, New York University, University of Notre Dame, Observatário Nacional / MCTI, The Ohio State University, Pennsylvania State University, Shanghai Astronomical Observatory, United Kingdom Participation Group, Universidad Nacional Autónoma de México, University of Arizona, University of Colorado Boulder, University of Oxford, University of Portsmouth, University of Utah, University of Virginia, University of Washington, University of Wisconsin, Vanderbilt University, and Yale University.

The CSS survey is funded by the National Aeronautics and Space Administration under grant no. NNG05GF22G issued through the Science Mission Directorate Near-Earth Objects Observations Program. The CRTS survey is supported by the U.S. National Science Foundation under grants AST-0909182.

Guoshoujing Telescope (the Large Sky Area Multi-Object Fiber Spectroscopic Telescope LAMOST) is a National Major Scientific Project built by the Chinese Academy of Sciences. Funding for the project has been provided by the National Development and Reform Commission. LAMOST is operated and managed by the National Astronomical Observatories, Chinese Academy of Sciences.

Based on observations obtained with the Samuel Oschin 48inch Telescope at the Palomar Observatory as part of the Zwicky Transient Facility project. ZTF is supported by the National Science Foundation under grant no. AST-1440341 and a collaboration includ-

# 2438 K. Inight et al.

ing Caltech, IPAC, the Weizmann Institute for Science, the Oskar Klein Center at Stockholm University, the University of Maryland, the University of Washington, Deutsches Elektronen-Synchrotron and Humboldt University, Los Alamos National Laboratories, the TANGO Consortium of Taiwan, the University of Wisconsin at Milwaukee, and Lawrence Berkeley National Laboratories. Operations are conducted by COO, IPAC, and UW.

# DATA AVAILABILITY

The data underlying this article are available from the sources referenced in the text. The results of our analysis are provided and in the online supplementary material.

## REFERENCES

- Abril J., Schmidtobreick L., Ederoclite A. r., López-Sanjuan C., 2020, MNRAS, 492, L40
- Ahumada R. et al., 2020, ApJS, 249, 3
- Antoja T. et al., 2018, Gaia DR2 documentation Chapter 10: Catalogue validation. European Space Agency
- Ashley R. P., Farihi J., Marsh T. R., Wilson D. J., Gänsicke B. T., 2019, MNRAS, 484, 5362
- Aungwerojwit A. et al., 2006, A&A, 455, 659
- Bailer-Jones C. A. L., 2015, PASP, 127, 994
- Bailer-Jones C. A. L., Rybizki J., Fouesneau M., Demleitner M., Andrae R., 2021, AJ, 161, 147
- Bailer-Jones C. A. L., Rybizki J., Fouesneau M., Mantelet G., Andrae R., 2018, AJ, 156, 58
- Bellm E. C. et al., 2019, PASP, 131, 018002
- Belokurov V. et al., 2020, MNRAS, 496, 1922
- Bergeron P., Greenstein J. L., Liebert J., 1990, ApJ, 361, 190
- Bergeron P. et al., 2011, ApJ, 737, 28
- Beuermann K., Harrison T. E., McArthur B. E., Benedict G. F., Gänsicke B. T., 2003, A&A, 412, 821
- Beuermann K., Harrison T. E., McArthur B. E., Benedict G. F., Gänsicke B. T., 2004, A&A, 419, 291
- Blanton M. R. et al., 2017, AJ, 154, 28
- Bours M. C. P. et al., 2015, MNRAS, 450, 3966
- Breedt E. et al., 2014, MNRAS, 443, 3174
- Brown W. R. et al., 2020, ApJ, 889, 49
- Bruch A., 2017, New A, 57, 51
- Burleigh M. R., Hogan E., Dobbie P. D., Napiwotzki R., Maxted P. F. L., 2006, MNRAS, 373, L55
- Bédard A., Bergeron P., Brassard P., Fontaine G., 2020, ApJ, 901, 93
- Casewell S. L., Littlefair S. P., Parsons S. G., Marsh T. R., Fortney J. J., Marley M. S., 2018, MNRAS, 481, 5216
- Castelli F., Kurucz R. L., 2003, in Piskunov N., Weiss W. W., Gray D. F., eds, IAU Symp. 210, Modelling of Stellar Atmospheres. Kluwer, Dordrecht, p. A20
- Chanmugam G., Wagner R. L., 1977, ApJ, 213, L13
- Chen X., Maxted P. F. L., Li J., Han Z., 2017, MNRAS, 467, 1874
- Cropper M., 1990, Space Sci. Rev., 54, 195
- Dalton G. et al., 2012, Proc. SPIE Conf. Ser. Vol. 4, Ground-Based and Airborne Instrumentation for Astronomy IV . SPIE, Bellingham, p. 84460P
- Davis P. J., Kolb U., Willems B., Gänsicke B. T., 2008, MNRAS, 389, 1563
- Debes J. H., 2006, ApJ, 652, 636
- De Donder E., Vanbeveren D., 2004, New A Rev., 48, 861
- de Jong R. S. et al., 2016, Proc. SPIE Conf. Ser. Vol. 6, Ground-Based and Airborne Instrumentation for Astronomy VI. SPIE, Bellingham, p. 990810
- Deloye C. J., Bildsten L., Nelemans G., 2005, ApJ, 624, 934
- Deloye C. J., Taam R. E., Winisdoerffer C., Chabrier G., 2007, MNRAS, 381, 525
- DESI Collaboration et al., 2016a, preprint (arXiv:1611.00036)

- DESI Collaboration et al., 2016b, preprint (arXiv:1611.00037) Drake A. J. et al., 2014, MNRAS, 441, 1186
- Drake A. J. et al., 2014, MNRAS, 441, 1180
- Duerbeck H. W., 1999, Inf. Bull. Var. Stars, 4731, 1 Dupke R. A., 1992, Inf. Bull. Var. Stars, 3749, 1
- Eggleton P. P., Tokovinin A. A., 2008, MNRAS, 389, 869
- Fabricius C., Makarov V. V., Knude J., Wycoff G. L., 2002, A&A, 386, 709
- Farihi J., Becklin E. E., Zuckerman B., 2005b, ApJS, 161, 394
- Farihi J., Hoard D. W., Wachter S., 2006, ApJ, 646, 480
- Farihi J., Hoard D. W., Wachter S., 2010, ApJS, 190, 275
- Farihi J., Zuckerman B., Becklin E. E., 2005a, AJ, 130, 2237
- Faulkner J., Flannery B. P., Warner B., 1972, ApJ, 175, L79
- Fontaine G., Brassard P., Bergeron P., 2001, PASP, 113, 409
- Fujimoto M. Y., 1982, ApJ, 257, 767
- Gaia Collaboration et al., 2016, A&A, 595, A1
- Gaia Collaboration et al., 2018, A&A, 616, A1
- Geier S., 2020, A&A, 635, A193
- Giammichele N., Bergeron P., Dufour P., 2012, ApJS, 199, 29
- Goliasch J., Nelson L., 2015, ApJ, 809, 80
- Green M. J. et al., 2020, MNRAS, 496, 1243 Green R. F., Schmidt M., Liebert J., 1986, ApJS, 61, 305
- Gänsicke B. T., 2004, in Tovmassian G., Sion E., eds, Rev. Mex. Astron. Astrofis. Conf. Ser., 20, 152
- Gänsicke B. T., Szkody P., Howell S. B., Sion E. M., 2005, ApJ, 629, 451
- Gänsicke B. T. et al., 2009, MNRAS, 397, 2170
- Harper W. E., 1928, Publ. Dom. Astrophys. Obs. Victoria, 4, 171
- Harrison T. E., Johnson J. J., McArthur B. E., Benedict G. F., Szkody P., Howell S. B., Gelino D. M., 2004, AJ, 127, 460
- Harrison T. E., McNamara B. J., Szkody P., Gilliland R. L., 2000, AJ, 120, 2649
- Hermes J. J. et al., 2014, MNRAS, 444, 1674
- Hernández Santisteban J. V. et al., 2016, Nature, 533, 366
- Hodgkin S. T., Wyrzykowski L., Blagorodnova N., Koposov S., 2013, Phil. Trans. R. Soc. London Ser. A, 371, 20120239
- Holberg J. B., Bergeron P., 2006, AJ, 132, 1221
- Holberg J. B., Oswalt T. D., Sion E. M., McCook G. P., 2016, MNRAS, 462, 2295
- Hollands M. A., Tremblay P. E., Gänsicke B. T., Gentile-Fusillo N. P., Toonen S., 2018, MNRAS, 480, 3942
- Houk N., Swift C., 1999, Michigan Spectral Surv., 5, 0
- Howell S. B., Nelson L. A., Rappaport S., 2001, ApJ, 550, 897
- Husser T. O., Wende-von Berg S., Dreizler S., Homeier D., Reiners A., Barman T., Hauschildt P. H., 2013, A&A, 553, A6
- Iben I. J., 1982, ApJ, 259, 244
- Ivanova N. et al., 2013, A&A Rev., 21, 59
- Jayasinghe T. et al., 2019, MNRAS, 485, 961
- Kahabka P., van den Heuvel E. P. J., 1997, ARA&A, 35, 69
- Kamiński K. Z. et al., 2007, AJ, 134, 1206
- Kaplan D. L. et al., 2014, ApJ, 780, 167
- Kato T., 2015, PASJ, 67, 108
- Kato T. et al., 2020, PASJ, 72, 14
- Kawka A., Vennes S., Dupuis J., Chayer P., Lanz T., 2008, ApJ, 675, 1518
- Kawka A., Vennes S., Vaccaro T. R., 2010, A&A, 516, L7
- Kleinman S. J. et al., 2013, ApJS, 204, 5
- Knigge C., 2006, MNRAS, 373, 484
- Knigge C., Baraffe I., Patterson J., 2011, ApJS, 194, 28
- Kollmeier J. A. et al., 2017, preprint (arXiv:1711.03234)
- Kosakowski A., Kilic M., Brown W. R., Gianninas A., 2020, ApJ, 894, 53
- Kowalski P. M., Saumon D., 2006, ApJ, 651, L137
- Kroupa P., Tout C. A., Gilmore G., 1993, MNRAS, 262, 545
- Kruse E., Agol E., 2014, Science, 344, 275
- Krushinsky V. et al., 2020, MNRAS, 493, 5208
- Kupfer T. et al., 2018, MNRAS, 480, 302
- Lagos F., Schreiber M. R., Parsons S. G., Gänsicke B. T., Godoy N., 2020a, MNRAS, 499, L121
- Lagos F. et al., 2020b, MNRAS, 494, 915
- Lamb F. K., Aly J. J., Cook M. C., Lamb D. Q., 1983, ApJ, 274, L71
- Landi R., Bassani L., Dean A. J., Bird A. J., Fiocchi M., Bazzano A., Nousek J. A., Osborne J. P., 2009, MNRAS, 392, 630

- Levitan D., Groot P. J., Prince T. A., Kulkarni S. R., Laher R., Ofek E. O., Sesar B., Surace J., 2015, MNRAS, 446, 391
- Lindegren L., 2018, Re-normalising the astrometric chi-square in Gaia DR2, GAIA-C3-TN-LU-LL-124, Lund Observatory, Available at: http://www. rssd.esa.int/doc\_fetch.php?id=3757412
- Lindegren L. et al., 2018, A&A, 616, A2
- Livio M., Mazzali P., 2018, Phys. Rep., 736, 1
- Marsh T. R., Dhillon V. S., Duck S. R., 1995, MNRAS, 275, 828
- Marsh T. R. et al., 2016, Nature, 537, 374
- Masci F. J. et al., 2019, PASP, 131, 018003
- Maxted P. F. L. et al., 2014, MNRAS, 437, 1681
- McCleery J. et al., 2020, MNRAS, 499, 1890
- Mennickent R. E., Diaz M., Skidmore W., Sterken C., 2001, A&A, 376, 448
- Meyer F., Meyer-Hofmeister E., 1981, A&A, 104, L10
- Meyer F., Meyer-Hofmeister E., 1983, A&A, 121, 29
- Miller-Jones J. C. A., Sivakoff G. R., Knigge C., Körding E. G., Templeton M., Waagen E. O., 2013, Science, 340, 950
- Morales-Rueda L., Marsh T. R., Maxted P. F. L., Nelemans G., Karl C., Napiwotzki R., Moran C. K. J., 2005, MNRAS, 359, 648
- Mróz P. et al., 2015, AcA, 65, 313
- Napiwotzki R. et al., 2020, A&A, 638, A131
- Nebot Gómez-Morán A. et al., 2011, A&A, 536, A43
- Nelemans G., 2018, preprint (arXiv:1807.01060)
- Nelson B., Young A., 1970, PASP, 82, 699
- Neustroev V. V. et al., 2017, MNRAS, 467, 597
- Norton A. J., Butters O. W., Parker T. L., Wynn G. A., 2008, ApJ, 672, 524
- Norton A. J., Wynn G. A., Somerscales R. V., 2004, ApJ, 614, 349
- Osaki Y., 1989, PASJ, 41, 1005
- O'Brien M. S., Bond H. E., Sion E. M., 2001, ApJ, 563, 971
- Paczynski B., 1976, in Eggleton P., Mitton S., Whelan J., eds, Proc. IAU Symp. 73, Structure and Evolution of Close Binary Systems. Kluwer, Dordrecht, p. 75
- Paczynski B., Sienkiewicz R., 1983, ApJ, 268, 825
- Paczyński B., 1967, AcA, 17, 287
- Pala A. F., Schmidtobreick L., Tappert C., Gänsicke B. T., Mehner A., 2018, MNRAS, 481, 2523
- Pala A. F. et al., 2020, MNRAS, 494, 3799
- Parsons S. G., Marsh T. R., Gänsicke B. T., Schreiber M. R., Bours M. C. P., Dhillon V. S., Littlefair S. P., 2013, MNRAS, 436, 241
- Parsons S. G., Rebassa-Mansergas A., Schreiber M. R., Gänsicke B. T., Zorotovic M., Ren J. J., 2016, MNRAS, 463, 2125
- Parsons S. G. et al., 2012, MNRAS, 426, 1950
- Parsons S. G. et al., 2015, MNRAS, 452, 1754
- Patterson J., 1994, PASP, 106, 209
- Patterson J. et al., 2002, PASP, 114, 721
- Patterson J. et al., 2005, PASP, 117, 1204
- Perryman M., ed., 1997, The HIPPARCOS and TYCHO catalogues. Astrometric and photometric star catalogues derived from the ESA HIPPAR-COS Space Astrometry Mission, ESA Special Publication Vol. 1200. European Space Agency, Publications Division, ESTEC, Noordwijk, The Netherlands
- Pickering E. C., Fleming W. P., 1896, ApJ, 4, 369
- Podsiadlowski P., Han Z., Rappaport S., 2003, MNRAS, 340, 1214
- Pogson N., 1857, MNRAS, 17, 200
- Pretorius M. L., Knigge C., Kolb U., 2007, MNRAS, 374, 1495
- Raghavan D. et al., 2010, ApJS, 190, 1
- Ramsay G., Schreiber M. R., Gänsicke B. T., Wheatley P. J., 2017, A&A, 604, A107
- Ramsay G. et al., 2018, A&A, 620, A141
- Rappaport S., Verbunt F., Joss P. C., 1983, ApJ, 275, 713
- Rebassa-Mansergas A., Agurto-Gangas C., Schreiber M. R., Gänsicke B. T., Koester D., 2013, MNRAS, 433, 3398
- Rebassa-Mansergas A. et al., 2012 MNRAS, 419, 806
- Rebassa-Mansergas A., Gänsicke B. T., Rodríguez-Gil P., Schreiber M. R., Koester D., 2007, MNRAS, 382, 1377
- Rebassa-Mansergas A., Gänsicke B. T., Schreiber M. R., Koester D., Rodríguez-Gil P., 2010, MNRAS, 402, 620

- Rebassa-Mansergas A., Parsons S. G., Copperwheat C. M., Justham S., Gänsicke B. T., Schreiber M. R., Marsh T. R., Dhillon V. S., 2014, ApJ, 790, 28
- Rebassa-Mansergas A., Ren J. J., Parsons S. G., Gänsicke B. T., Schreiber M. R., García-Berro E., Liu X. W., Koester D., 2016, MNRAS, 458, 3808
- Rebassa-Mansergas A. et al., 2012, MNRAS, 423, 320
- Rebassa-Mansergas A. et al., 2017, MNRAS, 472, 4193
- Reimers D., Hagen H. J., 2000, A&A, 358, L45
- Reimers D., Hagen H. J., Hopp U., 1999, A&A, 343, 157
- Ren J. J. et al., 2020, ApJ, 905, 38
- Ritter H., Kolb U., 2003, A&A, 404, 301
- Roelofs G. H. A., Groot P. J., Benedict G. F., McArthur B. E., Steeghs D., Morales-Rueda L., Marsh T. R., Nelemans G., 2007, ApJ, 666, 1174
- Savonije G. J., de Kool M., van den Heuvel E. P. J., 1986, A&A, 155, 51
- Schaefer B. E., Pagnotta A., 2012, Nature, 481, 164
- Schmadel L. D., Schmeer P., Börngen F., 1996, A&A, 312, 496
- Schreiber M. R., Gänsicke B. T., 2002, A&A, 382, 124
- Schreiber M. R., Gänsicke B. T., 2003, A&A, 406, 305
- Schwope A. D., Brunner H., Buckley D., Greiner J., Heyden K. v. d., Neizvestny S., Potter S., Schwarz R., 2002, A&A, 396, 895
- Sekiguchi K., 1992, Nature, 358, 563
- Shafter A. W., 1992, ApJ, 394, 268
- Shafter A. W., Wheeler J. C., Cannizzo J. K., 1986, ApJ, 305, 261
- Shappee B. J. et al., 2014, ApJ, 788, 48
- Shen K. J., Kasen D., Miles B. J., Townsley D. M., 2018, ApJ, 854, 52
- Smak J., 1984, PASP, 96, 5
- Smee S. A. et al., 2013, AJ, 146, 32
- Steele P. R. et al., 2013, MNRAS, 429, 3492
   Steinfadt J. D. R., Kaplan D. L., Shporer A., Bildsten L., Howell S. B., 2010, ApJ, 716, L146
- Strauss M. A. et al., 2002, AJ, 124, 1810
- Sun M., Arras P., 2018, ApJ, 858, 14
- Szkody P., Linnell A. P., Campbell R. K., Plotkin R. M., Harrison T. E., Holtzman J., Seibert M., Howell S. B., 2008, ApJ, 683, 967
- Szkody P. et al., 2009, AJ, 137, 4011
- Szkody P. et al., 2020, AJ, 159, 198
- Tapia S., 1977, ApJ, 212, L125
- Teyssier F., 2019, Astron. Telegram, 12936, 1
- Thorstensen J. R., 2003, AJ, 126, 3017
- Thorstensen J. R., Alper E. H., Weil K. E., 2016, AJ, 152, 226
- Thorstensen J. R., Lépine S., Shara M., 2008, AJ, 136, 2107
- Toloza O. et al., 2019, BAAS, 51, 168
- Toonen S., Hollands M., Gänsicke B. T., Boekholt T., 2017, A&A, 602, A16
- Toonen S., Nelemans G., Portegies Zwart S., 2012, A&A, 546, A70
- Townsley D. M., Bildsten L., 2003, ApJ, 596, L227
- Townsley D. M., Gänsicke B. T., 2009, ApJ, 693, 1007
- Tremblay P. E., Bergeron P., Gianninas A., 2011, ApJ, 730, 128
- Tsujimoto T., Nomoto K., Yoshii Y., Hashimoto M., Yanagida S., Thielemann F. K., 1995, MNRAS, 277, 945
- Unda-Sanzana E. et al., 2008, MNRAS, 388, 889
- van Roestel J. et al., 2018, MNRAS, 475, 2560
- Vennes S., Polomski E. F., Lanz T., Thorstensen J. R., Chayer P., Gull T. R., 2000, ApJ, 544, 423
- Vennes S. et al., 2011, ApJ, 737, L16
- Vernet J. et al., 2011, A&A, 536, A105
- Webbink R. F., 1984, ApJ, 277, 355
- Werner K., Reindl N., Löbling L., Pelisoli I., Schaffenroth V., Rebassa-Mansergas A., Irawati P., Ren J., 2020, A&A, 642, A228
- Whyte C. A., Eggleton P. P., 1980, MNRAS, 190, 801

Willems B., Kolb U., 2004, A&A, 419, 1057

York D. G. et al., 2000, AJ, 120, 1579

Wickramasinghe D. T., 1988, Cyclotron radiation from AM Herculis type systems, Univ. Arizona Press, Tucson, Arizona, p. 199

MNRAS 504, 2420-2442 (2021)

Wonnacott D., Kellett B. J., Stickland D. J., 1993, MNRAS, 262, 277

# 2440 K. Inight et al.

Yungelson L. R., 2008, Astron. Lett., 34, 620
Yu Z. et al., 2019, MNRAS, 489, 1023
Zhao G., Zhao Y., Chu Y., Jing Y., Deng L., 2012, Research in Astronomy and Astrophysics, 12, 723
Zorotovic M., Schreiber M. R., Gänsicke B. T., 2011, A&A, 536, A42
Zorotovic M., Schreiber M. R., Gänsicke B. T., Nebot Gómez-Morán A., 2010, A&A, 520, A86
Zorotovic M., Schreiber M. R., Parsons S. G., 2014, A&A, 568, L9
Zorotovic M. et al., 2016, MNRAS, 457, 3867
Zwitter T., Munari U., 1994, A&AS, 107, 503

# SUPPORTING INFORMATION

The Gold Samples in CSV format are available at MNRAS online.

CV appendix.csv

CV qfail appendix.csv

DWD appendix.csv WD + AFGK appendix.csv WD + M appendix.csv

Please note: Oxford University Press is not responsible for the content or functionality of any supporting materials supplied by the authors. Any queries (other than missing material) should be directed to the corresponding author for the article.

#### **APPENDIX A: THE GOLD SAMPLES**

We provide the Gold samples in the Supplementary Information as data files in CSV format. Missing information for individual properties of some systems are left as blank entries in these CSV files.

Table A1. The WD+M Gold Sample.

Field	Units	Description
name		Identifier
source_id		Gaia DR2 unique source identifier
references		References to literature provided as ADS bibcodes
ra	deg	Barycentric right ascension (ICRS) at Epoch 2015.5
dec	deg	Barycentric declination (ICRS) at Epoch 2015.5
type		Type of white dwarf and companion where known (PREP = pre-polar,BD = brown dwarf)
teffwd	Κ	White dwarf effective temperature where known (blank otherwise)
logg	$\log (\mathrm{cm}\mathrm{s}^{-2})$	Stellar log(surface gravity) where known (blank otherwise)
porb	h	Orbital period where known (blank otherwise)
sp		Spectral type of (M-type) companion $(0-9)$ where known, blank or $-1$ otherwise
parallax	mas	Absolute stellar parallax
parallax_error	mas	Uncertainty in parallax
pmra	mas yr <sup>-1</sup>	Proper motion in right ascension direction (pmRA*cosDE)
pmdec	mas yr <sup>-1</sup>	Proper motion in declination direction
phot_g_mean_mag		G-band mean magnitude (Vega)
phot_g_mean_mag_error		Uncertainty in G-band mean magnitude
phot_bp_mean_mag		Integrated $G_{\rm BP}$ mean magnitude (Vega)
phot_bp_mean_mag_error		Uncertainty in G <sub>BP</sub> -band mean magnitude
phot_rp_mean_mag		Integrated $G_{\rm RP}$ mean magnitude (Vega)
phot_rp_mean_mag_error		Uncertainty in G <sub>RP</sub> -band mean magnitude
1	deg	Galactic longitude
b	deg	Galactic latitude
ruwe		Gaia renormalized unit weight error
r_est	pc	Estimated distance
r_est_error	pc	Uncertainty in estimated distance
galex_coverage	-	1 if this location is present in one or more GALEX tiles, 0 otherwise
objid		GALEX identifier for the source
fuv_mag		GALEX FUV calibrated magnitude in AB system
fuve		Uncertainty in GALEX FUV
nuv_mag		GALEX NUV calibrated magnitude in AB system
nuve		Uncertainty in GALEX NUV

Field	Units	Description
name		Name
source_id		Gaia DR2 unique source identifier
references		References to literature
ra	deg	Barycentric right ascension (ICRS) at $Ep = 2015.5$
dec	deg	Barycentric declination (ICRS) at $Ep = 2015.5$
teff	Κ	Stellar effective temperature where known (blank otherwise)
porb	h	Orbital period where known (blank otherwise)
parallax	mas	Absolute stellar parallax
parallax_error	mas	Uncertainty in parallax
pmra	mas yr <sup>-1</sup>	Proper motion in right ascension direction (pmRA×cosDE)
pmdec	mas yr <sup>-1</sup>	Proper motion in declination direction
phot_g_mean_mag		G-band mean magnitude (Vega)
phot_g_mean_mag_error		Uncertainty in G-band mean magnitude
phot_bp_mean_mag		Integrated BP mean magnitude (Vega)
phot_bp_mean_mag_error		Uncertainty in BP-band mean magnitude
phot_rp_mean_mag		Integrated <i>RP</i> mean magnitude (Vega)
phot_rp_mean_mag_error		Uncertainty in RP-band mean magnitude
1	deg	Galactic longitude
b	deg	Galactic latitude
ruwe		Gaia renormalized unit weight error
r_est	parsec	Estimated distance
r_est_error	parsec	Uncertainty in estimated distance
galex_coverage		1 if this location is present in one or more GALEX tiles, 0 otherwise
objid		GALEX identifier for the source
fuv_mag		GALEX FUV calibrated magnitude in AB system
fuv_magerr		Uncertainty in GALEX FUV
nuv_mag		GALEX NUV calibrated magnitude in AB system
nuv_magerr		Uncertainty in GALEX NUV

#### Table A2. The WD+AFGK Gold Sample:.

#### Table A3. CV Gold Sample.

Field	Units	Description
name		Identifier
source_id		Gaia DR2 unique source identifier
references		References to literature
ra	deg	Barycentric right ascension (ICRS) at $Ep = 2015.5$
dec	deg	Barycentric declination (ICRS) at $Ep = 2015.5$
cv_type		CV type (Novalike, AM CVn, WZ Sge, Magnetic, U Gem or SU UMa) or CV if type is unknown
porb	h	Orbital period where known (blank otherwise)
parallax	mas	Absolute stellar parallax
parallax_error	mas	Uncertainty in parallax
pmra	mas yr <sup>-1</sup>	Proper motion in right ascension direction (pmRA*cosDE)
pmdec	$mas yr^{-1}$	Proper motion in declination direction
phot_g_mean_mag		G-band mean magnitude (Vega)
phot_g_mean_mag_error		Uncertainty in G-band mean magnitude
phot_bp_mean_mag		Integrated BP mean magnitude (Vega)
phot_bp_mean_mag_error		Uncertainty in BP-band mean magnitude
phot_rp_mean_mag		Integrated RP mean magnitude (Vega)
phot_rp_mean_mag_error		Uncertainty in RP-band mean magnitude
1	deg	Galactic longitude
b	deg	Galactic latitude
ruwe		Gaia renormalized unit weight error
r_est	pc	Estimated distance
r_est_error	pc	Uncertainty in estimated distance
galex_coverage		1 if this location is present in one or more GALEX tiles, 0 otherwise
objid		GALEX identifier for the source
fuv_mag		GALEX FUV calibrated magnitude in AB system
fuv_magerr		Uncertainty in GALEX FUV
nuv_mag		GALEX NUV calibrated magnitude in AB system
nuv_magerr		Uncertainty in GALEX NUV

# 2442 *K. Inight et al.*

**Table A4.** Confirmed CVs with d < 300 pc excluded from the Gold sample because of data quality issues:.

Field	Units	Description
name		Identifier
source_id		Gaia DR2 unique source identifier
references		References to literature
ra	deg	Barycentric right ascension (ICRS) at $Ep = 2015.5$
dec	deg	Barycentric declination (ICRS) at $Ep = 2015.5$
cv_type		CV type (Novalike, AM CVn, WZ Sge, Magnetic, U Gem or SU Uma) or CV if type is unknown
porb	h	Orbital period where known (0 otherwise)
parallax	mas	Absolute stellar parallax
parallax_error	mas	Uncertainty in parallax
pmra	mas yr <sup>-1</sup>	Proper motion in right ascension direction (pmRA*cosDE)
pmdec	mas yr <sup>-1</sup>	Proper motion in declination direction
phot_g_mean_mag		G-band mean magnitude (Vega)
phot_g_mean_mag_error		Uncertainty in G-band mean magnitude
phot_bp_mean_mag		Integrated BP mean magnitude (Vega)
phot_bp_mean_mag_error		Uncertainty in BP-band mean magnitude
phot_rp_mean_mag		Integrated RP mean magnitude (Vega)
phot_rp_mean_mag_error		Uncertainty in RP-band mean magnitude
1	deg	Galactic longitude
b	deg	Galactic latitude
ruwe		Gaia renormalized unit weight error
r_est	pc	Estimated distance
r_est_error	pc	Uncertainty in estimated distance
galex_coverage		1 if this location is present in one or more GALEX tiles, 0 otherwise
objid		GALEX identifier for the source
fuv_mag		GALEX FUV calibrated magnitude in AB system
fuv_magerr		Uncertainty in GALEX FUV
nuv_mag		GALEX NUV calibrated magnitude in AB system
nuv_magerr		Uncertainty in GALEX NUV

 Table A5.
 The DWD Gold Sample.

Field	Units	Description
name		Identifier
source_id		Gaia DR2 unique source identifier
references		References to literature
ra	deg	Barycentric right ascension (ICRS) at $Ep = 2015.5$
dec	deg	Barycentric declination (ICRS) at $Ep = 2015.5$
porb	h	Orbital period where known (blank otherwise)
parallax	mas	Absolute stellar parallax
parallax_error	mas	Uncertainty in parallax
pmra	mas yr <sup>-1</sup>	Proper motion in right ascension direction (pmRA*cosDE)
pmdec	mas yr <sup>-1</sup>	Proper motion in declination direction
phot_g_mean_mag		G-band mean magnitude (Vega)
phot_g_mean_mag_error		Uncertainty in G-band mean magnitude
phot_bp_mean_mag		Integrated BP mean magnitude (Vega)
phot_bp_mean_mag_error		Uncertainty in BP-band mean magnitude
phot_rp_mean_mag		Integrated RP mean magnitude (Vega)
phot_rp_mean_mag_error		Uncertainty in RP-band mean magnitude
1	deg	Galactic longitude
b	deg	Galactic latitude
ruwe		Gaia renormalized unit weight error
r_est	pc	Estimated distance
r_est_error	pc	Uncertainty in estimated distance
galex_coverage		1 if this location is present in one or more GALEX tiles, 0 otherwise
objid		GALEX identifier for the source
fuv_mag		GALEX FUV calibrated magnitude in AB system
fuv_magerr		Uncertainty in GALEX FUV
nuv_mag		GALEX NUV calibrated magnitude in AB system
nuv_magerr		Uncertainty in GALEX NUV

This paper has been typeset from a  $T_EX/I \Delta T_EX$  file prepared by the author.