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Modelling Sulphur Clusters for an Understanding of Ultramarine

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ABSTRACT

Ultramarine pigments are aluminosilicate-based and contain sulphur-based chromophores. Self-consistent-field Hartree-Fock and Møller-Plesset second order perturbation theory were applied to determine the relative stability of S_2 , $S_2^{-\bullet}$, S_2^{2-} , and S_3 , $S_3^{-\bullet}$, S_3^{2-} . The singly charged species was found to be the most stable in both sets. The transition from green to blue ultramarine is thought to be the transformation of the doubly charged species to the singly charged species and is known to be exothermic. Modelling studies supported this hypothesis. The open, $C_{2v'}$ isomer was found to be the most stable for the $S_3^{-\bullet}$ molecule, which is the blue chromophore in ultramarine. The closed, $D_{3h'}$ geometry represents a transition state. The S_4 molecule is the most likely chromophore in ultramarine red; however the specific isomer is uncertain. Under the assumption that S_4 was formed by a concerted reaction between $S_3^{-\bullet}$ and $S^{+\bullet}$, a Woodward-Hoffmann analysis of the molecular orbitals of $S_3^{-\bullet}$ and $S^{+\bullet}$ supported the formation of the puckered square S_4 , pyramidal S_4 , double triangle S_4 , and gauche S_4 chain isomers. The gauche S_4 chain is the most likely isomer in ultramarine red.

KEYWORDS

Ultramarine pigments, Woodward-Hoffmann analysis, modelling.

1. Introduction

Ultramarine pigments come in green, blue, greenish blue, reddish blue, violet, pink and red variants. ¹⁻⁷ Several chromophores are present in the aluminosilicate zeolite-type framework of the ultramarine species. S_2^{\bullet} is the yellow chromophore ^{2,5,8-10} in ultramarine and S_3^{\bullet} (Fig. 1) is the blue chromophore. ^{2,5,8,9} Ultramarine green is an ultramarine species in which the yellow and blue chromophores are present in such a ratio as to allow the green colour to be observed.

Ultramarine blue is synthesized in two steps, a reduction step followed by an oxidation step. $^{1,2,5,8,9,11-16}$ The nature of these steps is, however, disputed. Some authors suggested Equation 2,5,8,9 whereas others supported the existence of doubly charged species according to Equations 2 and $^{3,1,11-15}$

$$s_{8} \xrightarrow{\text{Reduction}} s_{2}^{-\bullet} \xrightarrow{\text{Oxidation}} s_{3}^{-\bullet}$$
Yellow Yellow Blue

$$2Na_{8}[Al_{6}Si_{6}O_{24}]S_{2-3} + SO_{2} + O_{2} \rightarrow 2Na_{7}[Al_{6}Si_{6}O_{24}]S_{2-3} + Na_{2}SO_{4}$$
 (2)

$$2 S_{2-3}^{2-}(s) + SO_2(g) + O_2(g) \rightarrow 2 S_{2-3}^{-}(s) + 2 S_4^{2-}$$
 (3)

The oxidation of ultramarine green to ultramarine blue is related to an increasing $S_3^{\bullet,\bullet}/S_2^{\bullet,\bullet}$ ratio, 2,5,8,9,16 in support of Equation 1. Beardsley and Whiting 12 described the formation of ultramarine blue with oxygen as oxidizing agent and introduced the use of SO_2 in the production thereof. 12 Gorshtein 11,14 successfully explained the formation of Na_2SO_4 and the role of SO_2 (Equations 2 and 3), noted by Tarling and co-workers 17 and Gobeltz and co-workers. 8 This explanation was based on the identification of S_3^{2-} and S_2^{-2-} as the yellow, and S_3^{-2-} and S_2^{-2-} as the blue chromophores. Other evidence identified S_3^{-2-} as responsible for the blue and S_2^{-2-} for the yellow colour. 2,5,8,9,18 No spectroscopic evidence of the doubly charged species is available in the litera-

ture.¹6 Köhler and co-workers¹0 synthesized ultramarine yellow from the thiocyanide derivative of sodalite. They confirmed that the yellow chromophore was $S_2^{-\bullet}$ and that oxidation led to the formation of ultramarine green, with the simultaneous formation of $S_3^{-\bullet}$.¹0 The evidence supported the formation of $S_3^{-\bullet}$ from $S_2^{-\bullet}$ rather than from $S_3^{-\bullet}$.

Weller and co-workers, ¹⁹ who also synthesized ultramarine from thiocyanate sodalites, concluded that a colourless polysulphide sodalite was the initial product at high temperatures. This initial product was assumed to undergo a secondary in-cage reaction to produce the coloured polysulphide radical anions *in situ* by heating in SO₂ at about 500°C. ¹⁹ The presence of a doubly-charged species was inferred from the formation of a noselite-like (nosean) phase. ¹⁹ The high ratio of sulphur atoms to unpaired electrons found by Hofmann and co-workers⁶ in the ultramarine green pigment supported the possible existence of doubly-charged, non-radical sulphide species. This conclusion was not generally accepted. ²⁰

The ultramarine red chromophore had been identified as the S_4 molecule (Fig. 2). 1,2,5,21 In order to identify the S_4 isomer it was assumed that S_4 was formed from $S_3^{-\bullet}$ and $S^{+\bullet}$ in a concerted reaction. 22 The feasibility of single-step concerted reaction products could be evaluated by examining the phase symmetry of the molecular orbitals, by the methods of Fukui 23,24 and of Woodward and Hoffmann. $^{25-27}$ The highest occupied molecular orbital (HOMO) and lowest unoccupied molecular orbital

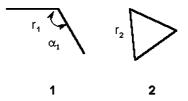


Figure 1. S_3 isomers: (1) open (C_{2v}) , (2) closed (D_{3h})

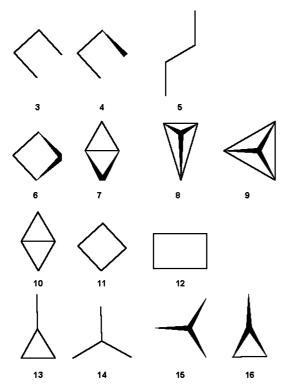


Figure 2. Possible S_4 isomers: (3) cis chain (C_{2v}) , (4) gauche chain (C_2) , (5) trans chain (C_{2h}) , (6) puckered ring (D_{2d}) , (7) butterfly (D_{2d}) , (8) tetrahedral (geometry optimized) (D_{2d}) , (9) tetrahedral (T_d) , (10) double triangle (D_{4h}) , (11) square planar (D_{4h}) , (12) rectangle (D_{2h}) , (13) exocyclic (C_{2v}) , (14) branched chain (D_{3h}) , (15) pyramidal (C_s) , (16) bent exocyclic (C_s)

(LUMO) were assumed to feature in reactions to form new bonds. For a reaction to take place the HOMO (LUMO) of the attacked species and the LUMO (HOMO) of the attacking species need to have the same phase symmetry. This condition allowed regions of the same phase in the molecular orbitals to overlap and to form new bonding and antibonding molecular orbitals provided favourable overlap occurred.

In this study the different possible S_4 isomers were evaluated in a modified Woodward-Hoffmann symmetry correlation analysis, similar to that applied by Carey and Sundberg. The current method modelled the molecular orbitals of the reactants, without modelling the molecular orbitals of the S_4 products. Sannigrahi and Grein used the Woodward-Hoffmann methodology to show that the formation of $S_4^{\,2+}$ from $S_2^{\,+}$ was not allowed.

The molecular orbitals of S^{+•}, S₂^{-•} and S₃^{-•} were computed in this work in order to use a Woodward-Hoffmann analysis to evaluate the feasibility of the formation of the possible S₄ isomers. S₂, S₂²⁻, S₃ and S₃²⁻ were also modelled in an effort to find the most stable sulphur species. These results might confirm that green ultramarine yielded blue ultramarine by the formation of S₂^{-•}, S₃^{-•} from S₂²⁻, S₃²⁻. ¹¹⁻¹⁵

2. Experimental Data

Naudé^{30,31} studied the rotational spectrum of S_2 and determined the bond length to be 1.893 Å. Based on the absorption band fine structure between 160 and 190 nm, Maeder and Miescher³² calculated a vibrational wavenumber of 724 cm⁻¹ for S_2 . $S_2^{-\bullet}$ had a Raman band at 590 cm⁻¹,²⁵ and S_2^{-} a band at 451 cm⁻¹,³³

Hopkins and co-workers³⁴ observed the S₃ molecule as part of an electric discharge reaction and observed Raman bands at 490, 585 and 650 cm⁻¹. Brabson and co-workers³⁵ recorded the infra-

Table 1. Experimental characteristics of S_3 , S_3^{\bullet} and S_3^{2-} species.

Species	Assignment	Wavenum	Wavenumber/cm ⁻¹		
		Infrared	Raman		
S ₃	v_1		585	[34]	
	v_2 ?	495 or 310	490		
	v_3		650		
	v_1		583	[35]	
	v_2	680		[35,37]	
	Vibrational				
	progression	340		[35,39]	
		450			
	v_1		575	[36]	
	v_2		256		
	v_3		656		
$S_3^{-\bullet}$	v_1		548	[2,5]	
3	v_3	582		[2]	
	v_1	547		[2]	
	v_2		232	[40]	
S_3^{2-}	v_1 or v_3	457	458	[37]	
-	v_2		227/238		
	v_1 or v_3	476	476		
	v_1 and v_3		466	[38]	
	v_2		238		

red spectrum of S_3 in a solid argon matrix (Table 1). They concluded that the structure of S_3 was the $C_{2\nu}$ open geometry, with a bond angle of $116\pm 2^{\circ}.^{35}$ Lenain and co-workers³⁶ studied the Raman spectrum of S_3 and also concluded that the observed species was the $C_{2\nu}$ isomer (Table 1). Clark and Cobbold² observed both the Raman and infrared bands for $S_3^{-\bullet}$ (Table 1). Janz and co-workers^{33,37,38} studied several polysulphides. For Ba S_3 the S_3^{2-} structure was said to be $C_{2\nu}$ with vibrational bands at 227/238 (R), 458 (R,IR) and 476 cm⁻¹ (R,IR).³⁷ For K_2S_3 the same interpretation held, with the stretching modes accidentally degenerate: 238 (R) and 466 cm⁻¹ (R,IR).³⁸

3. Molecular Modelling

HyperChem 41 was used to perform the computations on a Compaq Pressario 1200 Celeron 850 MHz processor with 128 MB RAM

In the work on sulphur compounds the shapes of the molecular orbitals were important because they dictated the reactivity of the species under consideration and were used in the Woodward-Hoffmann analyses. The energies were also important as an estimate of the relative stability of the ionic states of S_2 and S_3 . A potential energy surface specified the classical potential energy as a function of molecular structure. This concept was useful for the comparison of the possible geometrical isomers. The potential energy surface was inspected for S_2 , $S_2^{-\bullet}$, S_2^{-2} , S_3 , $S_3^{-\bullet}$ and S_3^{-2} in an effort to find the most favourable oxidation states of these species, and to obtain the molecular orbitals for use in a Woodward-Hoffmann analysis.

A disadvantage inherent in all the computations was that the data obtained were related to gas phase molecules in vacuum at 0 K. No interaction between molecules within a crystal was taken into account. Therefore disparities between computed and observed quantities could be large.

Geometry optimizations were performed by means of the Fletcher-Reeves conjugate gradient method.

3.1. Semi-empirical ZINDO/1 Modelling Scheme

ZINDO/1 is a semi-empirical method based on a modified version of the intermediate neglect of differential overlap

Table 2. Bond lengths (Å) and angles (degrees) for S_2 and S_3 species.

Structure name Symmetry		Structural parameter ^a	This work	Average from literature ^b	References	
$\overline{S_2}$	D _{∞ h}	r	1.88	1.91 (2)	[48,49,51]	
S_{2}^{-}	$D_{\infty h}$	r	2.01	2.12 (8)	[48,57]	
S ₂ ² -	$\mathrm{D}_{_{\infty}\mathrm{h}}^{^{\mathrm{n}\mathrm{n}}}$	r	2.21 (2.80)	. ,		
Singlet open S ₃	C_{2v}	\mathbf{r}_1	1.90	1.98 (2)	[47,49–56]	
0 1	2.	α_1	117	116.6 (7)	[47,49–56]	
Singlet closed S ₃	D_{3h}	r_2	2.08	2.12 (2)	[47,49–56]	
Triplet open S ₃	C_{2v}^{SR}	r_1	1.99	2.00 (4)	[49,52]	
1 1 0	2.	α_1	93	93 (2)	[49,52]	
Open S ₃	C_{2v}	$\mathbf{r}_{_{1}}$	2.00	2.02 (3)	[47,56,57]	
1 5	24	$\overset{\cdot}{lpha_1}$	115	113 (2)	[47,56,57]	
Closed S ₃ -	D_{3h}	r_2	2.24	2.241 (5)	[47,56]	
Open S ₃ ² -	C_{2v}^{Sit}	$\overset{\mathtt{r}}{\mathbf{r}_{1}}$	2.13	2.08 (8)	[47,57]	
1 0	24	α_1	113	109 (6)	[47,57]	
Closed S ₃ ²⁻	D_{3h}	r_2	2.24	2.244	[47]	

^a As defined in Fig. 1.

(INDO) method. 42 The overlap weighting factors were set to 1 for both $\sigma-\sigma$ and $\pi-\pi$ interactions, as required for geometry optimization. 42

Semi-empirical and *ab initio* methods took into account both molecular geometry and the electron distribution and the results from HyperChem⁴¹ were interpreted in terms of the linear combination of atomic orbitals (LCAO)-molecular orbital (MO) theory.⁴² These molecular orbitals were assumed to describe the motion of the electrons in molecules and the shapes of molecular orbitals to determine the reactivity of bonds in the molecule.⁴²

HyperChem calculated the potential energy as a basis for geometry optimization and the force constants, vibrational modes, charge and spin densities, atomic charges, dipole moments and electrostatic charges. The electronic spectrum could also be calculated when configuration interaction was taken into account.42 Because quantum mechanical computations deal explicitly with the electrons, it was necessary to specify the charges and spin multiplicities of the molecules.⁴² When a singlet state was considered the computation was usually done with restricted Hartree-Fock (RHF) wavefunctions.42 For higher multiplicity computations unrestricted Hartree-Fock (UHF) wave functions which distinguish between spin states were used. 42 Alpha electrons would then be in the highest occupied molecular orbital, because they were assigned first. The remainder of the electrons were assigned to beta spins.42 The excess of alpha over beta electrons was calculated based on the number of electrons in the neutral atoms and the number of electrons lost or gained due to the charge, and on the specified multiplicity.42

Self-consistency was sometimes difficult to obtain, especially when degenerate states were involved.⁴² To avoid this problem, HyperChem was run with a convergence accelerator, the direct inversion in the iterative subspace (DIIS) method, although the computational effort was thereby increased, ⁴² because the analysis of the molecular orbitals was a linear combination of the current result and the results of previous iterations.

Individual molecular orbitals could be represented on a grid in space. This was the key to deductions made regarding the reactivity of molecules, the so-called frontier orbital approach of Woodward and Hoffmann.²³⁻²⁷

The vibrational analysis and infrared spectroscopic data were obtained from the numerical Hessian matrix of second derivatives of the total energy with respect to the nuclear positions, and a normal coordinate analysis, based on massweighted coordinates. ⁴² The analysis was based on data obtained from quantum mechanical computations, but it remained classical. ⁴² The calculation of vibrational frequencies was described by Matsuura and Yoshida. ⁴³

3.2. Ab Initio Modelling Scheme

The Møller-Plesset second-order perturbation (MP2) correlation energy⁴⁴ compensated for some of the assumptions made during the Hartree-Fock computations and was used to correct the calculated energy.⁴² The MP2 energy correction was done only at the energy minimum in a single point computation. This differed from its application in more advanced software packages. The *ab initio* computations in HyperChem expanded the molecular orbitals into a linear combination of atomic orbitals, without any further approximation. Several different basis sets to describe the atomic orbitals exist.⁴² The 6-311G** notation implies a triple valence basis set with polarization functions.⁴² The computations were performed at the self-consistent-field Hartree-Fock level of theory extended by MP2 energy correction

Ab initio-determined vibrational frequencies are usually larger than observed frequencies, and need scaling. 43,45 The scale factor depended on the basis set and the level of theory used. The recommended scale factor for HF/6-311G** was 0.9051. 45 This scale factor was based on 1066 experimental wavenumbers from 122 molecules. The molecules consisted mainly of carbon and hydrogen atoms. Krossing and Passmore, 46 however, suggested that scale factors larger than 1 be used when considering sulphur clusters. Brabson and co-workers suggested a scale factor of 0.87, based on experimental data for S₃ and a self-consistent-field method. This scale factor was used in this work.

HyperChem supported configuration interaction (CI) only in the restricted Hartree-Fock mode. The configuration interaction computations led to the energies of the ground state and singly excited states, which were then used to calculate the electronic spectrum. ⁴² HyperChem could, therefore, only compute an electronic spectrum for molecules with a singlet spin state. ⁴² Configuration interaction was performed in one of two ways, either as a singly excited state or as a microstate. The electrons were allowed to exchange from a limited number of occupied to a limited number of unoccupied orbitals. Configuration interaction was used to obtain a more accurate set of states by taking appropriate linear combinations of these microstates or singly excited states. ⁴²

b The value given is the average from several literature values; the number in brackets denotes one standard deviation of the mean.

Table 3. Relative energies of S₂ species ^a.

Structure number	Structure name	Relative energy/kJ mol ⁻¹
2	S ₂ singlet	146
2	S ₂ triplet	113
2	S_2^-	0
2	S ₂ ² - singlet	427
2	S_2^{2-} singlet S_2^{2-} triplet	632 ^b

^a Energies are reported relative to the most stable isomer, taken as zero.

4. Results

Compared with literature values the geometry optimization gave reasonable results (Table 2). As the charge on the S_2 species became more negative the bond distance increased (Table 2). The geometry of the closed, D_{3h} , $S_3^{\bullet\bullet}$ isomer did not optimize to the desired symmetry and the bond length was therefore set to the average of the values found in the literature (Table 2). The geometry of the closed S_3^{2-} isomer optimized to unrealistically long bond lengths, and the bond distance was set to 2.24 Å.⁴⁷ The singly negative anions of both S_2 and S_3 were found to be the most stable states relative to the neutral and doubly charged states (Tables 3 and 4). Our results at the UHF/6-311G** level of theory indicated that the open, $C_{2\nu}$, structure was the most stable isomer for the $S_3^{\bullet\bullet}$ molecule (Table 4). In fact, at this level, the closed, D_{3h} , $S_3^{\bullet\bullet}$ geometry was shown to be a transition state.

The calculated vibrational wavenumbers for S_2 , $S_2^{\bullet \bullet}$ and S_2^{2-} were in good agreement with the experimental values (Table 5). For S_2^{2-} the unscaled wavenumbers fit best (Table 5). The vibrational wavenumber decreased with an increase in the negative charge on the S_2 molecule as expected from the reverse trend in the bond distances (Table 2). After applying the scale factor of 0.87 to the calculated wavenumbers (Table 6) of the open S_3 isomer, they corresponded well with experimental values (Tables 1 and 6). The scaled vibrational wavenumbers of the closed, D_{3h} , S_3 isomer (Table 7) agreed with the experimental values. The calculated vibrational wavenumbers of $S_3^{\bullet \bullet}$ were less accurate (Table 8). The antisymmetric stretching wavenumber was underestimated, even before scaling. For S_3^{2-} , the vibrational wavenumbers were close to the experimental values without scaling (Table 8).

Molecular modelling of the polysulphide species was done in

Table 4. Relative energies of S₃ species ^a.

Structure number	Structure name	Relative energy/kJ mol ⁻¹	
3	Open S ₃ singlet	197	
4	Closed S ₃ singlet	234	
3	Open S ₃ triplet	297	
4	Closed S ₃ triplet	473 ^b	
3	Open S ₃	0	
4	Closed S ₃ -	192 ь	
4	Closed S ₃ ²⁻ singlet	720 ^b	
3	Open S ₃ ² - singlet	331	
4	Closed S ₃ ²⁻ triplet	682 ^b	
3	Open S ₃ ² - triplet	577 ^ь	

^a Energies are reported relative to the most stable isomer, taken as zero.

Table 5. Vibrational wavenumbers (cm⁻¹) of S₂ stretching vibration.

Species	This work	Scaled values	Experimental	References
Singlet S ₂	791	687	724	[32]
Triplet S_2 Doublet $S_2^{-\bullet}$	790 615	687 535	590	[2,5]
Singlet S ₂ ²⁻	426	371	451	[33]



Figure 3. Highest occupied and lowest unoccupied molecular orbitals of S_2^{\bullet}

order to obtain a molecular orbital picture of $S_2^{\bullet\bullet}$ (Fig. 3), and $S^{+\bullet}$ and $S_3^{\bullet\bullet}$ (Fig. 4). HyperChem⁴¹ was used to model the $S_2^{\bullet\bullet}$ and $S_3^{\bullet\bullet}$ structures by the ZINDO/1 semi-empirical method, and self-consistent-field unrestricted Hartree-Fock simulation using a 6-311G** basis set extended by MP2 correlation energy. The ZINDO/1 results confirmed the character of the highest occupied molecular orbital for both open (b₁) and closed (a₂) $S_3^{\bullet\bullet}$ of a

Table 6. Vibrational wavenumbers (cm⁻¹) of open S₃ (1).

Mode	Description	Symmetry	This work	Scaled values ^a	Experimental values b	Spectroscopic activity ^c
v_1 v_2 v_3	Symmetric stretching Bending Antisymmetric stretching	$egin{aligned} \mathbf{a_1} \\ \mathbf{a_1} \\ \mathbf{b_2} \end{aligned}$	654 (627) 294 (211) ^d 737 (490)	569 256 641	575 256 656	R, IR R, IR R, IR

^a Scaled with a scale factor of 0.87 (see ref. 35).

Table 7. Vibrational wavenumbers (cm⁻¹) of closed S₃ (2).

Mode	Description	Symmetry	This work	Scaled values ^a	Experimental values	Spectroscopic activity b
$\begin{matrix} \nu_1 \\ \nu_2 \end{matrix}$	Symmetric stretching Deformation	$a_{1}^{\ \prime}\\e^{\prime}$	662 500	576 435	583 ³⁵ , 585 ³⁴ , 575 ³⁶ 450 ^{35,39} , 490 ³⁴	R R, IR

^a Scaled with a scale factor of 0.87 (reference 35).

^b Species has an imaginary frequency.

^b Species has an imaginary frequency.

^b Reference 36.

^c R: Raman active; IR: infrared active.

^d Values in brackets are for the triplet state.

^b R: Raman active; IR: infrared active

Table 8. Vibrational wavenumbers (cm⁻¹) of negatively charged open S₃ (1).

Mode	Description	Symmetry	$S_3^{-\bullet}$ doublet	S₃ ^{-•} singlet	S ₃ ² -	open	Spectroscopic activity ^a
			Calculated	Experimental	Calculated (ref. 37)	Experimental	
$\begin{array}{c} v_1 \\ v_2 \\ v_3 \end{array}$	Symmetric stretching Bending Antisymmetric stretching	a ₁ a ₁ b ₂	558 250 449	548 ² 232 ⁴⁰ 582 ²	459 202 496	458 227/238 476	R, IR R, IR R, IR

^a R: Raman active; IR: infrared active

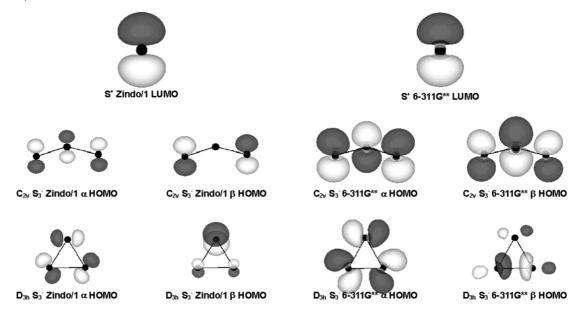


Figure 4. Molecular orbital diagram of S^{+•} (ZINDO/1 and 6-311G** LUMO) and S₃^{-•} (ZINDO/1 and 6-311G** α and β HOMO)

previous pictorial model²² and the work of Koch and coworkers⁴⁷ (Fig. 4).

It was assumed that $S^{+\bullet}$ (Fig. 4) attacked the $S_3^{-\bullet}$ (Fig. 4) species in a concerted reaction to form S_4 (Fig. 2). The validity of different isomers as potential products was tested by means of the frontier orbital symmetry theory of Woodward and Hoffmann. ^{25–27} The HOMO of $S_3^{-\bullet}$ (Fig. 4) was placed near the LUMO of $S_3^{+\bullet}$ (Fig. 4) in such a way as to ensure that lobes of the same phase overlapped to form σ -orbitals. Other orientations of the $S_3^{+\bullet}$ molecular orbital were possible. The charge distribution on the sulphur atoms of the $S_3^{-\bullet}$ isomers indicated that the central sulphur atom was positively charged (Fig. 5), therefore attack of $S_3^{+\bullet}$ was unlikely to occur at this atom. To test whether the electrostatic charges would be unfavourable to this form of attack the distance between the fourth sulphur atom and the central sulphur atom, in the branched chain isomer, was increased (Fig. 5). The

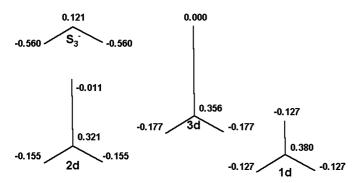


Figure 5. Charge distribution in $S_3^{\bullet\bullet}$ and branched chain S_4 at three times (3d), twice (2d) and at its normal distance (1d).

induced charge distribution is similar to the original charge distribution in the $S^{3-\bullet}$ moiety, in that the central atom is positive relative to the terminal sulphur atoms, therefore the positive charge on the central sulphur atom in $S_3^{-\bullet}$ is not necessarily a problem in the formation of the branched chain S_4 from $S_3^{-\bullet}$ and $S_3^{+\bullet}$

A Woodward-Hoffmann analysis at the ZINDO/1 level showed that the puckered S_4 ring (6), pyramidal branched chain S_4 (15), gauche S_4 chain (4), double triangle (10) and bent exocyclic S_4 (16) isomers were possible products of the concerted reaction of $S^{+\bullet}$ and $S_3^{-\bullet}$. Using the 6-311 G^{**} basis set, the highest occupied molecular orbital of the closed, D_{3h} , $S_3^{-\bullet}$ isomer precluded the formation of the bent exocyclic S_4 (16) isomer.

Furthermore, the symmetry of the highest occupied and lowest unoccupied molecular orbitals of S_2^{\bullet} (Fig. 3) predicted that the closed, D_{3h} , S_3^{\bullet} isomer could not be formed from S_2^{\bullet} in a concerted way.

5. Discussion

Computational results should not be accepted at face value because of the approximations made during the analysis. Strictly speaking, *ab initio* computations are only relevant to gas phase molecules in vacuum at 0 K. The environment created by the sodium counterions and the aluminosilicate framework was not taken into account. The structure of the aluminosilicate framework did ensure that the sulphur chromophores acted independently as molecules and did not form part of an extended sulphur-based 'molecule'.

The calculated wavenumber of S_2 at 793 cm⁻¹, scaled (0.87) to 690 cm⁻¹, was in reasonable agreement with the experimental value of 724 cm⁻¹ (Table 5). The bond length of 1.88 Å was in close

agreement with the experimental 1.893 Å. 30,31 Heinemann and co-workers 48 calculated the vibrational wavenumber of S $_2$ $^{\bullet}$ to be 582 cm $^{-1}$, close to the experimental value (Table 5). The current value was not far from the expected value (Table 5). Heinemann and co-workers 48 calculated that the triplet S $_2$ molecule was less energetic than the S $_2$ $^{\bullet}$ molecule. The current results confirm this

Rau and co-workers⁵⁸ suggested that sulphur allotrope tautomers with an energy difference less than 33 kJ mol⁻¹ contributed appreciably to the vapour composition of that particular species at temperatures between 823 and 1273 K. This was suggested on the basis of the work of Pauling,⁵⁹ who mentioned the estimate of tautomeric occurrence at 41 kJ mol⁻¹. Therefore, sulphur allotropic isomers found, in this work, to have an energy difference of less than 41 kJ mol⁻¹ relative to the most stable isomer should be present in appreciable amounts.

The bond angle calculated for the open S_3 isomer was 117° , which is within the experimental range. The question of whether the open, C_{2v} , 47,49,52,53,55,56 or closed, D_{3h} , 48,55,56,60,61 S_3 isomer was the more stable seemed to be an open question, since several methods gave contradictory results, depending on the methodology applied to the problem. The open, C_{2v} , S_3 isomer was calculated to be the more stable in this work. The computed closed, D_{3h} , S_3 isomer had an energy only 37 kJ mol^{-1} higher than that of the open isomer, therefore both these isomers should be observable (Table 4) and probably were (Tables 6 and 7).

From the literature it seemed as though the open, $C_{2v'}^{47,54}$ $S_3^{\bullet\bullet}$ isomer was deemed to be more stable than the closed, $D_{3h'}$ isomer. This was confirmed in the current work. The scale factor of 0.87, suggested by Brabson and co-workers, 35 worked well for the neutral sulphur molecules (Tables 5–7), but failed for the ionized species (Tables 5 and 8). Since the $S_3^{\bullet\bullet}$ molecule is the blue chromophore in ultramarine blue, $^{1,2,5,8,9}_{-}$ the geometry of this important stable radical polysulphide was calculated, and found to be C_{2v} . This geometry could now be used in modelling the ultramarine pigment itself, 62 and could assist in crystallographic interpretation. 17

The reaction that transforms ultramarine green to blue is exothermic. ¹² The calculated results (Tables 3 and 4) supported the proposal that this reaction involves the transition from S_2^{2-} , S_3^{-1} to S_2^{-1} , S_3^{-1} . ^{1,11–15}

The assumption that the S_4 molecules were formed by concerted reactions of $S_3^{-\bullet}$ and $S^{+\bullet}$ was theoretically useful, although the existence of $S^{+\bullet}$ was unlikely, even though it had been observed in photoionization experiments.⁶³

Based on the fact that the tetrahedral geometry of S_4 has an open-shell electronic configuration, Salahub and co-workers⁵⁴ ruled out the tetrahedral geometry as the ground state of S_4 by $X\alpha$ -scattered wave calculations. This conclusion was also reached by Landman and De Waal based on Woodward-Hoffmann selection rules,²² and was verified in the current work. The puckered S_4 ring (6), and especially the double triangle S_4 (10) would probably not be the most stable isomers due to internal ring strain. The pyramidal branched chain S_4 isomer (15) also seemed unlikely to be the most stable isomer, based on the close proximity of the sulphur atoms. The gauche S_4 chain (4) was likely to be the energy minimum for the S_4 isomers, due to the strain and interaction-free geometry.

6. Conclusion

At the UHF/6-311G** level of theory the open, C_{2v} , structure was found to be the more stable isomer for the $S_3^{-\bullet}$ molecule, which was the blue chromophore in ultramarine. The closed,

 D_{3h} , geometry was calculated to be a transition state (Table 4). For both S_2 and S_3 the singly negative anion was calculated to be the most stable state, in relation to the neutral and doubly-charged states (Tables 3 and 4), and supported the proposal that this reaction involved the transition from S_2^{2-} , S_3^{-2-} to S_2^{-4-} , $S_3^{-4-1,11-15}$

Woodward-Hoffmann orbital analyses and steric interactions indicated that the gauche chain S_4 (4) was the most probable S_4 isomer in ultramarine red.

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