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# Discrete curvatures motivated from Riemannian geometry and optimal transport 

Bonnet-Myers-type diameter bounds and rigidity

Supanat Kamtue

A thesis submitted for the degree of Doctor of Philosophy<br>Department of Mathematical Sciences<br>Durham University<br>United Kingdom<br>July 2021

# Discrete curvatures motivated from Riemannian geometry and optimal transport 

Bonnet-Myers-type diameter bounds and rigidity

Supanat Kamtue


#### Abstract

This thesis gives an overview of three notions of Ricci curvature for discrete spaces, including Ollivier Ricci curvature (motivated from optimal transport), BakryÉmery curvature (from Bochner's formula in Riemannian geometry) and ErbarMaas entropic Ricci curvature (from optimal transport). The first part of the thesis provides background knowledge in optimal transport theory and Riemannian geometry which is essential to the understanding of generalized Ricci curvatures for metric measure spaces and the mentioned Ricci curvatures for graphs.

For each of the three discrete curvature notions, discussed in their respective part of the thesis, we provide the definition of the curvature and use hypercubes as an example for the curvature calculation. We study various curvature results with an emphasis on upper bounds of diameter and lower bounds of the spectral gap for graphs with positive lower bound on the Ricci curvature. These results can be regarded as discrete analogues of the Bonnet-Myers theorem and the Lichnerowicz theorem in Riemannian geometry. In addition, we deeply investigate into the rigidity results (analogous to Cheng's rigidity) in attempt to classify all graphs which yield the sharp diameter bound in the sense of Ollivier Ricci curvature and Bakry-Émery curvature.


## Preface

Discrete curvatures in this thesis refer to discrete versions of Ricci curvature (which is a geometric concept on manifolds) for discrete spaces such as graphs and discrete Markov chains. Unlike on manifolds, there is no unified notion of curvature on graphs, and each of discrete curvature notions is inspired by different geometric aspects of Ricci curvature. In particular, we focus on three discrete curvature notions, namely,

1. Ollivier Ricci curvature, which is defined via the contraction property of $L^{1}$-transportation distance,
2. Bakry-Émery curvature, which is defined via Bochner's formula in Riemannian geometry and the Laplace operator,
3. Erbar-Maas entropic Ricci curvature, which is defined via the convexity of the entropy functional along $L^{2}$-transportation geodesics.

There are other discrete curvature notions that are beyond the scope of the thesis, for example, combinatorial curvature on tessellated surfaces, Forman-Ricci curvature on simplicial complexes, and sectional-type curvature for Alexandrov spaces.
This thesis is divided into four main parts, where each of Parts II III and IV is dedicated to one of the three discrete curvature notions mentioned above. Part I provides background knowledge in optimal transport theory and Riemannian geometry that is helpful for understanding the discrete curvatures. Outlined below are topics which are covered in each part.

Part I begins with an introduction to optimal transport theory (Kantorovich's problem and its duality) with an emphasis towards geometries of the $L^{2}$-Wasserstein space of probability measures. Otto's formal calculus [Ott01] is used as a framework to view the Wasserstein space as a "Riemannian manifold" and to describe, for example, the transportation distance in fluid dynamics viewpoint (BenamouBrenier formula $[\mathrm{BB} 00]$ ) and the heat flow by a gradient flow structure (Jordan-Kinderlehrer-Otto theorem [JKO98]). We also discuss three famous approaches
by (i) Sturm and Lott-Villani [Stu06, LV09], (ii) Bakry-Émery [BE84], and (iii) Ollivier [Oll09] to generalize Ricci curvature for metric measure spaces. These approaches give rise to the discrete curvatures under discussion. A central theme of the thesis is about the Bonnet-Myers diameter bound theorem [Mye41] and its rigidity (known as Cheng's rigidity [Che75]). They are classical theorems in Riemannian geometry and their analogous results will be re-appearing in the context of each discrete curvature.

In Part II about Ollivier Ricci curvature on graphs, we revisit an optimal transport in the setting of graphs, which prompt us to the curvature's definition. We discuss various properties of Ollivier Ricci curvature, including Cartesian products, a lower bound on the spectral gap (Lichnerowicz-type theorem) and a Bonnet-Myers-type upper bound on diameter. The key result in this part is my own contribution to the rigidity result in collaboration with David Cushing, Jack Koolen, Shiping Liu, Florentin Münch and my supervisor Norbert Peyerimhoff [CKK $\left.{ }^{+} 20\right]$, including the crucial concept of transport geodesics and the classification of all self-centered Bonnet-Myers sharp graphs. The presentation in this thesis deviates somewhat from $\left[\mathrm{CKK}^{+} 20\right]$ by providing an alternative proof of a preliminary rigidity result using only transport geodesics and no properties of the Laplacian. It is also a natural question whether the assumption of self-centeredness can be removed from our classification result. Some partial progress (removal of self-centeredness in diameter three) is provided in [Kam20a].
Part III features Bakry-Émery curvature on weighted graphs and our joint research result [CKLP21] on the curvature reformulation as a minimal eigenvalue of a local matrix. Moreover, we follow the Bakry-Emery theory to provide, in the discrete setting, an equivalent characterization of Bakry-Émery curvature via a gradient estimate of the heat semigroup. In parallel to the previous part, we discuss relevant properties regarding Cartesian products and a lower bound on the spectral gap. We then survey some key ideas of research results by Liu, Münch and Peyerimhoff [LMP18, LMP17] on a Bonnet-Myers-type diameter bound and its rigidity with respect to Bakry-Émery curvature.

In Part IV, we present Erbar and Maas' work [EM12] to define, on discrete Markov chains, a discrete notion of Ricci curvature in the spirit of Sturm and Lott-Villani, called the entropic Ricci curvature. This approach requires a modification of the 2-Wasserstein distance which is motivated from Otto's calculus (in particular, Benamou-Brenier formula). Similarly to Bakry-Émery curvature, we also provide an equivalent reformulation of the entropic Ricci curvature in terms of a Bochner-type formula and a gradient estimate. These tools are crucial to derive a Bonnet-Myers type diameter upper bound for a Markov chain with positive entropic Ricci curvature. This result is a generalization of my paper [Kam20b],
which deals with the special case of simple random walks.
Explicit computation of all three discrete curvatures is given for hypercubes as they represent an important family of graphs which possess positive curvature, a product structure, and (potentially) sharp diameter bounds with respect to all discrete curvatures. Each of the three parts finishes with an outlook chapter discussing some related open questions.
Readers who are particularly interested in one of these discrete curvatures can start reading from its respective part without prior knowledge from Part I. During the discussion in discrete curvatures, there usually are pointers which reconnect us to related topics in the manifold case in Part I. Readers with differential geometry and Riemannian geometry background may find the interplay between Ricci curvature and optimal transport interesting, in particular since the Lott-Sturm-Villani synthetic Ricci curvature notion allows to extend various fundamental curvature results to the more general setting of metric measure spaces, which appear naturally via Gromov-Hausdorff convergence of Riemannian manifolds.

In summary, besides its research results, this thesis serves also as a survey on topics which are motivated from Riemannian geometry and optimal transport with special emphasis on Bonnet-Myers-type diameter bounds and the corresponding rigidity results.

Supanat Kamtue
July 2021

## Declaration

This thesis is the result of the author's own work and includes nothing which is the outcome of work done in collaboration except as declared in the Preface and specified in the text. No substantial part of this thesis has already been submitted, or, is being concurrently submitted for any such degree, diploma or other qualification at Durham University or any other institution except as declared in the Preface and specified in the text.

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## Part I

## Curvature notions on metric measure spaces

## Chapter 1

## Introduction to optimal transport

Let us start with a brief introduction to optimal transport in a general setting of metric measure spaces, where the topics include Monge's problem, Kantorovich's problem, the duality, and the Wasserstein space. In the later section, we will restrict to the setting of Riemannian manifolds, where geometric structures of the Wasserstein space are more enriched. For the contents given in this chapter, we mostly refer to a reader-friendly note by Ambrosio and Gigli [AG13], and occasionally from a book by Santambrogio [San15] and a book by Villani [Vil09].

### 1.1 Setting of metric measure spaces

Let $(X, d)$ and $\left(Y, d^{\prime}\right)$ be Polish spaces (i.e., complete and separable metric spaces), equipped with $\sigma$-Borel algebra. Define $\mathscr{P}(X)$ and $\mathscr{P}(Y)$ to be the sets of all Borel probability measures on $X$ and on $Y$, respectively.
For a given Borel map $T: X \rightarrow Y$ and a measure $\mu \in \mathscr{P}(X)$, the pushforward of $\mu$ by $T$ is the measure denoted by $T_{\#} \mu \in \mathscr{P}(Y)$ such that

$$
T_{\#} \mu(B)=\mu\left(T_{-1}(B)\right),
$$

for all Borel subsets $B \subset Y$. Equivalently, the pushforward satisfies the following change of variable formula:

$$
\int_{Y} f(y) d T_{\#} \mu(y)=\int_{X} f(T(x)) d \mu(x),
$$

for all Borel functions $f: X \rightarrow Y$.
The original optimal transport problem proposed by Monge in [Mon81] is defined as follows.

Definition 1.1. Given a Borel cost function $c: X \times Y \rightarrow \mathbb{R}_{\geq 0}$. Let $\mu \in \mathscr{P}(X)$ and $\nu \in \mathscr{P}(Y)$. Monge's optimal transport problem is the following minimization problem:

$$
\begin{equation*}
\inf _{T}\left\{\operatorname{cost}(T):=\int_{X} c(x, T(x)) d \mu(x) \mid T_{\#} \mu=\nu\right\} . \tag{MP}
\end{equation*}
$$

Any map $T: X \rightarrow Y$ such that $T_{\#} \mu=\nu$ is called a transport map from $\mu$ to $\nu$. A transport map $T$ which minimizes (MP) is called an optimal transport map.
Intuitively, a map $T$ transports the collection of mass distributed as in $\mu$ to mass distributed as in $\nu$ by moving mass from location $x$ to location $T(x)$ (without splitting mass) with the transport cost given by $c(x, T(x))$ per unit mass. Monge's problem is to find the minimum total cost of all such transport maps $T$. One limitation of Monge's problem is about the possibility of non-existence of transport maps. For instance, in a situation that $\mu=\delta_{x}$ is a Dirac measure at $x$ but $\nu$ is not a Dirac measure (e.g. $\nu=\frac{1}{2} \delta_{y}+\frac{1}{2} \delta_{y^{\prime}}$ with $y \neq y^{\prime}$ ), no transport map exists since the splitting of mass at $x$ is not allowed.

In [Kan42], Kantorovich considers a more general optimal transport problem which allows the mass to be split at the source.

Definition 1.2. Given a Borel cost function $c: X \times Y \rightarrow \mathbb{R}_{\geq 0}$. Let $\mu \in \mathscr{P}(X)$ and $\nu \in \mathscr{P}(Y)$. A transport plan from $\mu$ to $\nu$ is a Borel probability measure $\Gamma \in \mathscr{P}(X \times Y)$ that satisfies the following marginal constraints:

$$
\begin{aligned}
\Gamma(A \times Y) & =\mu(A) \\
\Gamma(X \times B) & =\nu(B)
\end{aligned}
$$

for all Borel $A \subseteq X$ and $B \subseteq Y$. Equivalently, the marginal constraints can be expressed by

$$
\left(\pi_{1}\right)_{\#} \Gamma=\mu \text { and }\left(\pi_{2}\right)_{\#} \Gamma=\nu,
$$

where $\pi_{i}$ denotes the $i$ th-projection: $\pi_{1}(x, y)=x$ and $\pi_{2}(x, y)=y$. Furthermore, we denote by $\Pi(\mu, \nu)$, the set of all transport plans from $\mu$ to $\nu$.
Kantorovich's optimal transport problem (also known as the Monge-Kantorovich problem) is the following minimization problem:

$$
\begin{equation*}
\inf _{\Gamma}\left\{\operatorname{cost}(\Gamma):=\int_{X \times Y} c(x, y) d \Gamma(x, y) \mid \Gamma \in \Pi(\mu, \nu)\right\} . \tag{KP}
\end{equation*}
$$

Any such plan $\Gamma$ which minimizes (KP) is called an optimal transport plan.

Intuitively, $\Gamma(x, y)$ represents the amount of mass that is moved from $x$ to $y$. In contrast to Monge's problem, there always exists a transport plan for Kantorovich's problem (as a trivial one is given by $\Gamma=\mu \times \nu$ ).

It is noteworthy that the support of a transport plan lies in the product of the marginal supports.

Definition 1.3. The support of a measure $\mu \in \mathscr{P}(X)$ is defined as the set

$$
\operatorname{supp}(\mu):=\{x \in X \mid \mu(U)>0 \quad \forall \text { open neighborhood } U \text { of } x\} \subset X .
$$

Proposition 1.4. Let $\mu \in \mathscr{P}(X)$ and $\nu \in \mathscr{P}(Y)$, and let $\Gamma \in \Pi(\mu, \nu)$. Then

$$
\operatorname{supp}(\Gamma) \subset \operatorname{supp}(\mu) \times \operatorname{supp}(\nu)
$$

Proof. Assume that $(x, y) \in \operatorname{supp}(\Gamma)$. Then the marginal constraints implies that, for any open $x \in A \subset X$ and $y \in B \subset Y$, we have $\mu(A)=\Gamma(A \times X)>0$ and $\nu(B)=\Gamma(X \times B)>0$, which means $x \in \operatorname{supp}(\mu)$ and $y \in \operatorname{supp}(\nu)$.

The connection between transport maps and transport plans can be described as in the following proposition.

Proposition 1.5. Given $\mu \in \mathscr{P}(X)$ and two Borel maps $F: X \rightarrow X$ and $G$ : $X \rightarrow Y$. Let the map $(F, G): X \rightarrow X \times Y$ be defined as $x \mapsto(F(x), G(x))$. Then $\Gamma:=(F, G)_{\#} \mu$ is a transport plan from $F_{\#} \mu$ to $G_{\#} \mu$, and its cost is given by

$$
\begin{equation*}
\operatorname{cost}(\Gamma)=\int_{X} c(F(x), G(x)) d \mu(x) \tag{1.1}
\end{equation*}
$$

Proof. For $A \subset X$ and $B \subset Y$, we have

$$
\Gamma(A \times B)=\mu\left((F, G)^{-1}(A \times B)\right)=\mu\left(F^{-1}(A) \cap G^{-1}(B)\right)
$$

In particular, $\Gamma$ satisfies the marginal constraints $\Gamma(A \times X)=\mu\left(F^{-1}(A)\right)=$ $F_{\#} \mu(A)$ and $\Gamma(X \times A)=\mu\left(G^{-1}(B)\right)=G_{\#} \mu(A)$.

The change of variable formula then gives

$$
\operatorname{cost}(\Gamma)=\int_{X \times Y} c(x, y) d(F, G)_{\#} \mu(x, y)=\int_{X} c((F, G) x) d \mu(x)
$$

as desired.

An immediate consequence from Proposition 1.5 is that, for any Borel map $T$ : $X \rightarrow Y$, the pushforward $\Gamma:=(\mathrm{id}, T)_{\#} \mu$ is a transport plan from $\mu$ to $T_{\#} \mu$ with $\operatorname{cost}(\Gamma)=\operatorname{cost}(T)$. This means the problem (KP) can be considered as a relaxation to the problem (MP), and the infimum in (KP) is less than or equal to the infimum in (MP):

$$
\inf _{\Gamma \in \Pi(\mu, \nu)} \operatorname{cost}(\Gamma) \leq \inf _{T_{\# \mu=\nu}} \operatorname{cost}(T) .
$$

When (KP) is regarded as a linear programming problem in $\Gamma$, there occurs naturally a dual problem.

Definition 1.6. Given a Borel cost function $c: X \times Y \rightarrow \mathbb{R}_{\geq 0}$. Let $\mu \in \mathscr{P}(X)$ and $\nu \in \mathscr{P}(Y)$. Consider all pairs of functions $\phi: X \rightarrow \mathbb{R}$ and $\psi: Y \rightarrow \mathbb{R}$ which satisfy

$$
\phi(x)+\psi(y) \leq c(x, y) \quad \forall x \in X, \forall y \in Y
$$

or written shortly as $\phi \oplus \psi \leq c$.
Kantorovich's dual problem is the following maximization problem:

$$
\begin{equation*}
\sup _{\phi, \psi}\left\{J_{\mu, \nu}(\phi, \psi):=\int_{X} \phi(x) d \mu(x)+\int_{Y} \psi(y) d \nu(y) \mid \phi \oplus \psi \leq c\right\} . \tag{DP1}
\end{equation*}
$$

For further discussion on the dual problem, let us recall some relevant concepts of $c$-transform and $c$-concavity.

Definition 1.7. Given a cost function $c: X \times Y \rightarrow \mathbb{R}_{\geq 0}$. Consider functions $\phi: X \rightarrow \mathbb{R}$ and $\psi: Y \rightarrow \mathbb{R}$. The $c$-transform of $\phi$ is a function $\phi^{c}: Y \rightarrow \mathbb{R}$ defined as

$$
\phi^{c}(y):=\inf _{x \in X} c(x, y)-\phi(x) \quad \forall y \in Y .
$$

A function $\psi: Y \rightarrow \mathbb{R}$ is called c-concave if $\psi=\phi^{c}$ for some $\phi: X \rightarrow \mathbb{R}$.
Similarly, the $c$-transform of $\psi$ is a function $\psi^{c}: X \rightarrow \mathbb{R}$ defined as

$$
\psi^{c}(x):=\inf _{y \in Y} c(x, y)-\psi(y) \quad \forall x \in X .
$$

A function $\phi: X \rightarrow \mathbb{R}$ is called $c$-concave if $\phi=\psi^{c}$ for some $\psi: Y \rightarrow \mathbb{R}$. Moreover, we denote by $\operatorname{Lip}^{c}(X)$, the set of all $c$-concave functions on $X$.

Here are some basic properties of $c$-transform and $c$-concave functions.
Proposition 1.8. Let $\phi: X \rightarrow \mathbb{R}$ and $\psi: Y \rightarrow \mathbb{R}$. Then
(i) $\phi \oplus \phi^{c} \leq c$, and $\varphi \leq \phi^{c}$ for all $\varphi: Y \rightarrow \mathbb{R}$ such that $\phi \oplus \varphi \leq c$.
(ii) $\phi^{c c c}=\phi^{c}$, and similarly $\psi^{c c c}=\psi^{c}$.
(iii) $\phi^{c c}=\phi$ if and only if $\phi \in \operatorname{Lip}^{c}(X)$.

Proof. (i) follows directly from the definition. To see (ii), we observe

$$
\begin{aligned}
\phi^{c c c}(y) & =\inf _{x} c(x, y)-\left(\inf _{\tilde{y}} c(x, \tilde{y})-\left(\inf _{\tilde{x}} c(\tilde{x}, \tilde{y})-\phi(\tilde{x})\right)\right) \\
& =\inf _{x} \sup _{\tilde{y}} \inf _{\tilde{x}} \underbrace{c(x, y)-c(x, \tilde{y})+c(\tilde{x}, \tilde{y})-\phi(\tilde{x})}_{=: \Phi(x, y, \tilde{x}, \tilde{y})} .
\end{aligned}
$$

Thus

$$
\begin{aligned}
& \phi^{c c c}(y) \geq \inf _{x} \inf _{\tilde{x}} \Phi(x, y, \tilde{x}, y)=\phi^{c}(y), \\
& \phi^{c c c}(y) \leq \inf _{x} \sup _{\tilde{y}} \Phi(x, y, x, \tilde{y})=\phi^{c}(y),
\end{aligned}
$$

and therefore $\phi^{c c c}=\phi^{c}$. (iii) then follows from (ii).
From the above proposition, one can see that

$$
\sup _{\phi \oplus \psi \leq c} J_{\mu, \nu}(\phi, \psi)=\sup _{\phi \in \operatorname{Lip}^{c}(X)} J_{\mu, \nu}\left(\phi, \phi^{c}\right) .
$$

Definition 1.9. Kantorovich's dual problem can be reformulated as the maximization problem:

$$
\begin{equation*}
\sup _{\phi}\left\{J_{\mu, \nu}\left(\phi, \phi^{c}\right):=\int_{X} \phi(x) d \mu(x)+\int_{Y} \phi^{c}(y) d \nu(y) \mid \phi \in \operatorname{Lip}^{c}(X)\right\} \tag{DP2}
\end{equation*}
$$

Any $\phi \in \operatorname{Lip}^{c}(X)$ that maximizes (DP2) is called an optimal Kantorovich potential with respect to $\mu$ and $\nu$.

The relation between Kantorovich's problem (KP) and its dual problem (DP2) is given by the so-called Kantorovich duality. Moreover, minimizers of (KP) and maximizers of (DP2) exist under a certain condition on the cost function $c$, as given below. We refer to [AG13, Theorem 1.17] for more details and proof.

Theorem 1.10 (Kantorovich duality). Let $\mu \in \mathscr{P}(X), \nu \in \mathscr{P}(Y)$ and let $c$ : $X \times Y \rightarrow \mathbb{R}_{\geq 0}$ be a continuous cost function such that

$$
\begin{equation*}
c(x, y) \leq a(x)+b(y) \tag{1.2}
\end{equation*}
$$

for some $a \in L^{1}(\mu)$ and $b \in L^{1}(\nu)$. Then Kantorovich's problem and its dual problem satisfy the following Kantorovich duality:

$$
\begin{equation*}
\inf _{\Gamma \in \Pi(\mu, \nu)} \operatorname{cost}(\Gamma)=\sup _{\phi \in \operatorname{Lip}^{c}(X)} J_{\mu, \nu}\left(\phi, \phi^{c}\right) . \tag{1.3}
\end{equation*}
$$

Furthermore, there always exist an optimal transport plan $\Gamma$ and an optimal Kantorovich potential $\phi$.

Remark 1.11. 1. The assumption $c \leq a \oplus b$ with $a \in L^{1}(\mu)$ and $b \in L^{1}(\nu)$ guarantees that any transport plan $\Gamma \in \Pi(\mu, \nu)$ has a finite cost:

$$
\int_{X \times Y} c(x, y) d \Gamma(x, y) \leq \int_{X} a(x) d \mu(x)+\int_{Y} b(y) d \nu(y)<\infty .
$$

Without this assumption, the duality (1.3) still holds true in general by allowing infinite values for both sides.
2. If $X=Y$, and the cost function is given by $c=d$, the metric of $X, c$ concave functions are exactly 1-Lipschitz functions, and the $c$-transform of $\phi \in \operatorname{Lip}^{c}(X)$ is given by $\phi^{c}=-\phi$.

Theorem 1.12 (complementary slackness). Assume that $\Gamma$ is an optimal transport plan and $\phi$ is an optimal Kantorovich potential from $\mu$ to $\nu$. Then

$$
\begin{equation*}
\phi(x)+\phi^{c}(y)=c(x, y) \quad \forall(x, y) \in \operatorname{supp}(\Gamma) . \tag{1.4}
\end{equation*}
$$

Proof. Since $\Gamma$ and $\phi$ are a minimizer and a maximizer in the duality, we have

$$
\begin{aligned}
0 & =\operatorname{cost}(\Gamma)-J\left(\phi, \phi^{c}\right) \\
& =\int_{X \times Y} c(x, y) d \Gamma(x, y)-\int_{X} \phi(x) d \mu(x)+\int_{Y} \phi^{c}(y) d \nu(y) \\
& =\int_{X \times Y}\left(c(x, y)-\phi(x)-\phi^{c}(y)\right) d \Gamma(x, y),
\end{aligned}
$$

due to the marginals $d \mu(x)=\int_{y \in Y} d \Gamma(x, y)$ and $d \nu(y)=\int_{x \in X} d \Gamma(x, y)$. Since $c(x, y)-\phi(x)-\phi^{c}(y) \geq 0$, the equality (1.4) must holds $\Gamma$-almost everywhere, and by continuity of $c$, it holds everywhere on $\operatorname{supp}(\Gamma)$. For more details, we refer to [AG13, Remarks 1.15 and 1.18].

From now on, we always consider a situation where $X=Y$ and the cost function $c$ is the power of the distance function $d$, that is, $c=d^{r}$ for some fixed $r \in$
$[1, \infty)$. The minimum value of Kantorovich's problem (KP) defines the so-called $L^{r}$-Wasserstein distance between two measures as

$$
\begin{equation*}
W_{r}(\mu, \nu):=\left(\inf _{\Gamma \in \Pi(\mu, \nu)} \int_{X \times X} d(x, y)^{r} d \Gamma(x, y)\right)^{1 / r} . \tag{1.5}
\end{equation*}
$$

It is important to remark that $W_{r}$ defines a metric on $\mathscr{P}_{r}(X) \subset \mathscr{P}(X)$, the space of all Borel probability measures with finite $r$-moments, that is,

$$
\mathscr{P}_{r}(X):=\left\{\mu \in \mathscr{P}(X) \mid \int_{X} d\left(x_{0}, x\right)^{r} d \mu(x)<\infty \text { for some } x_{0} \in X\right\} .
$$

Proposition 1.13. For $r \in[1, \infty)$, $W_{r}$ is a metric on $\mathscr{P}_{r}(X)$.
A proof to the above proposition can be found in e.g. [AG13, Theorem 2.2] in the case of $W_{2}$, but it can be generalized to any $W_{r}$ for $r \geq 1$. The idea to prove triangle inequality is to apply the gluing lemma to compose transport plans and then apply Minkowski's inequality.
The finite $r$-moment restriction guarantees that $W_{r}\left(\delta_{x_{0}}, \mu\right)<\infty$ for some (and hence for all) $x_{0}$, and therefore $W_{r}(\mu, \nu)<\infty$ for all $\mu, \nu \in \mathscr{P}_{r}(X)$ due to the triangle inequality.

Definition 1.14 (Wasserstein space). For $r \in[1, \infty)$, the $L^{r}$-Wasserstein space is the space of probability measures $\mathscr{P}_{r}(X)$ endowed with the $L^{r}$-Wasserstein metric $W_{r}$.

The following proposition states some basic topological properties of the Wasserstein space $\left(\mathscr{P}_{r}(X), W_{r}\right)$; for details see [Vil09, Theorem 6.18 and Remark 6.19].

Proposition 1.15 (Topological properties of $\mathscr{P}_{r}(X)$ ). Let $(X, d)$ be a complete and separable metric space and $r \in[0, \infty)$. Then $\left(\mathscr{P}_{r}(X), W_{r}\right)$ is also a complete and separable metric space. Moreover, if $X$ is compact, then $\mathscr{P}_{r}(X)$ is also compact. However, if $X$ is only locally compact, then $\mathscr{P}_{r}(X)$ is not locally compact.

Remark 1.16. Note that there is a trivial embedding from $X$ into $\mathscr{P}_{r}(X)$ given by delta distributions. The embedding is isometric, i.e., $W_{r}\left(\delta_{x}, \delta_{y}\right)=d(x, y)$ because transport maps from $\delta_{x}$ to $\delta_{y}$ are those maps $T$ such that $T(x)=y$.

Another important property of the Wasserstein space is about geodesic. Let us recall the definition of geodesics and geodesic spaces.

Definition 1.17 (geodesics in metric spaces). Let $(X, d)$ be a metric space. A curve $\gamma: I \rightarrow X$ (for some interval $I \subset \mathbb{R}$ ) is called a constant speed geodesic
if there is a constant $v \geq 0$ such that $d\left(\gamma_{s}, \gamma_{t}\right)=v|t-s|$ for all $s, t \in I$. Here we write $\gamma_{t}:=\gamma(t)$.
In particular, a curve $\gamma:[0,1] \rightarrow X$ is a constant speed geodesic if for all $s, t \in$ $[0,1]$,

$$
\begin{equation*}
d\left(\gamma_{s}, \gamma_{t}\right)=|t-s| d\left(\gamma_{0}, \gamma_{1}\right) . \tag{1.6}
\end{equation*}
$$

Remark 1.18. Equivalently, one can weaken the condition (1.6) by replacing it with $d\left(\gamma_{s}, \gamma_{t}\right) \leq|t-s| d\left(\gamma_{0}, \gamma_{1}\right)$. For $0 \leq s \leq t \leq 1$, the triangle inequality gives

$$
d\left(\gamma_{0}, \gamma_{1}\right) \leq d\left(\gamma_{0}, \gamma_{s}\right)+d\left(\gamma_{s}, \gamma_{t}\right)+d\left(\gamma_{t}, \gamma_{1}\right) \leq(s+(t-s)+(1-s)) d\left(\gamma_{0}, \gamma_{1}\right)
$$

which makes these inequalities hold with equality, and $d\left(\gamma_{s}, \gamma_{t}\right)=|t-s| d\left(\gamma_{0}, \gamma_{1}\right)$ is recovered.

Definition 1.19 (geodesic spaces). A metric space $(X, d)$ is a geodesic space if for every $x, y \in X$, there exists a constant speed geodesic joining $x$ and $y$, namely $\gamma:[0,1] \rightarrow X$ such that $\gamma_{0}=x$ and $\gamma_{1}=y$ satisfying (1.6).

If there exists a geodesic $\gamma:[0,1] \rightarrow X$ from $x$ to $y$ in $X$, then the curve $t \mapsto \delta_{\gamma(t)}$ is indeed a geodesic in $\mathscr{P}_{r}(X)$ from $\delta_{x}$ to $\delta_{y}$. In fact, if $X$ is a geodesic space, then $\mathscr{P}_{r}(X)$ is also a geodesic space (see e.g. [AG13, Theorem 2.10] for the proof).

Theorem 1.20. Suppose that $(X, d)$ is a geodesic space. Then the Wasserstein space $\left(\mathscr{P}_{2}(X), W_{2}\right)$ is also a geodesic space.

On the other hand, when $X$ is a discrete space (which means $X$ it is not a geodesic space), the Wasserstein space $\left(\mathscr{P}_{2}(X), W_{2}\right)$ is not a geodesic space (see Chapter 9). This fact leads to Erbar-Maas' approach to define Ricci curvature notion on discrete spaces, which is later discussed in Part IV.

Although it is not a focus of this work, we would like to remark another property of the Wasserstein space related to a lower curvature bound in the sense of Alexandrov. We refer to [Oht14, Section 8.2] for the following theorem and to [BBI01] for general knowledge on Alexandrov spaces.

Definition 1.21. A geodesic space $(X, d)$ is an Alexandrov space of curvature $\geq k$ if for every geodesic triangle in $X$ is "thicker" than the triangle of the same side length in $\mathbb{M}^{2}(k)$, the space of constant sectional curvature $k$. More precisely, the space $X$ must satisfy the following property. Let $\triangle x y z$ be any triangle in $X$ and let $\triangle x^{\prime} y^{\prime} z^{\prime}$ be the corresponding triangle in $\mathbb{M}^{2}(k)$, that is, $d(x, y)=\mathrm{d}\left(x^{\prime}, y^{\prime}\right)$ and so on. For any point $p$ on the edge $x y$ and the corresponding point $p^{\prime}$ on $x^{\prime} y^{\prime}$ (that is, $d(x, p)=\mathrm{d}\left(x^{\prime}, p^{\prime}\right)$ ), it must satisfy $d(z, p) \geq \mathrm{d}\left(z^{\prime}, p^{\prime}\right)$.

Note that in particular, $\mathbb{M}^{2}(1)$ is the sphere of radius one, $\mathbb{M}^{2}(0)$ is the Euclidean plane, and $\mathbb{M}^{2}(-1)$ is the real hyperbolic plane.

Theorem 1.22. A compact geodesic space $(X, d)$ is an Alexandrov space of nonnegative curvature if and only if the Wasserstein space $\left(\mathscr{P}_{2}(X), W_{2}\right)$ is an Alexandrov space of nonnegative curvature. Moreover, if $(X, d)$ is not an Alexandrov space of nonnegatively curvature, then $\left(\mathscr{P}_{2}(X), W_{2}\right)$ is not an Alexandrov space of any curvature lower bound.

### 1.2 Setting of Riemannian manifolds

In this section, we focus on geometric properties of the $L^{2}$-Wasserstein space $\left(\mathscr{P}_{2}(M), W_{2}\right)$ in the case that $(M,\langle\cdot, \cdot\rangle)$ is a connected complete Riemannian manifold. The understanding these properties, especially McCann's theorem, will prepare us for later discussion about the synthetic Ricci curvature notion. For a reference on geometric properties of $\left(\mathscr{P}_{2}(M), W_{2}\right)$, we suggest the work by Lott [Lot08].

One of the most remarkable results is about the existence and uniqueness of the optimal transport map on a Riemannian manifold, and it also gives the characterization of such map in terms of a Kantorovich potential. This result is proved by Brenier [Bre87] for Euclidean space $\mathbb{R}^{n}$, and it is generalized by McCann for Riemannian manifolds [McC01]. These results play an important role in the discussion of Otto's calculus in Chapter 2.

Theorem 1.23 (Brenier's theorem). Let $\mu, \nu \in \mathscr{P}_{2}\left(\mathbb{R}^{n}\right)$ with $\mu$ absolutely continuous with respect to Lebesgue measure. Then there exists a unique optimal transport plan from $\mu$ to $\nu$, and this plan is induced by a transport map $T: \mathbb{R}^{n} \rightarrow \mathbb{R}^{n}$ given by

$$
T(x):=x-\operatorname{grad} \varphi(x),
$$

where $\varphi=\frac{1}{2} \phi$ for some optimal Kantorovich potential $\phi$.
Here in the case of Euclidean space, the gradient $\operatorname{grad} \varphi: \mathbb{R}^{n} \rightarrow \mathbb{R}^{n}$ is given at any point $x=\left(x_{1}, \ldots, x_{n}\right) \in \mathbb{R}^{n}$ by $\operatorname{grad} \varphi(x)=\left(\frac{\partial \varphi}{\partial x_{1}}(x), \ldots, \frac{\partial \varphi}{\partial x_{n}}(x)\right)$. The definition of the gradient in a more general setting of manifolds is discussed in Section 5.1.

Theorem 1.24 (McCann's theorem). Let $M$ be a connected and closed manifold. Let $\mu, \nu \in \mathscr{P}_{2}(M)$ such that $\mu \ll \mathrm{vol}$. Then there exists a unique optimal transport plan from $\mu$ to $\nu$, and this plan is induced by a transport map $T: M \rightarrow M$ given by

$$
T(x):=\exp _{x}(-\operatorname{grad} \varphi(x)),
$$

where $\varphi=\frac{1}{2} \phi$ for some optimal Kantorovich potential $\phi$.
We note that the factors $\frac{1}{2}$ appear in the above statements (but not in the original statements) because the original ones use the cost function $c(x, y)=\frac{1}{2} d^{2}(x, y)$. Here we only provide the proof for Brenier's theorem to give an intuition how the quadratic cost function play a central role in the argument. For the proof of McCann's theorem, we refer to [ $\mathrm{McC01}$, Theorem 8 and 9].

Proof of Brenier's theorem. In view of Theorem 1.10, there exists an optimal transport plan $\pi \in \Pi(\mu, \nu)$ and an optimal Kantorovich potential $\varphi$. Let $\left(x_{0}, y_{0}\right) \in$ $\operatorname{supp}(\pi)$. By complementary slackness theorem, the map $x \mapsto \phi(x)-c\left(x, y_{0}\right)$ assumes the maximum value of $-\phi^{c}\left(y_{0}\right)$ at $x=x_{0}$. This implies that the gradient of this map vanishes at $x_{0}$ :

$$
\operatorname{grad}\left(\phi(\cdot)-c\left(\cdot, y_{0}\right)\right)\left(x_{0}\right)=0
$$

With the cost function $c(x, y)=h(x-y)$, for some strictly convex function $h$, we have

$$
\operatorname{grad} \phi\left(x_{0}\right)=\operatorname{grad} h\left(x_{0}-y_{0}\right) .
$$

Moreover, $\operatorname{grad} h$ is invertible, so

$$
x_{0}-y_{0}=(\operatorname{grad} h)^{-1}\left(\operatorname{grad} \phi\left(x_{0}\right)\right) .
$$

In the particular case of the quadratic cost $c(x, y)=h(x-y)=|x-y|^{2}$, its gradient satisfies $\operatorname{grad} h(x)=2 x$ and its inverse is simply $(\operatorname{grad} h)^{-1}(x)=\frac{1}{2} x$. Therefore,

$$
y_{0}=x_{0}-\frac{1}{2} \operatorname{grad} \phi\left(x_{0}\right) .
$$

Corollary 1.25. Let $M$ be a connected and closed manifold and let $\mu, \nu \in \mathscr{P}_{2}(M)$ such that $\mu \ll \operatorname{vol}$. Let $T: M \rightarrow M$ be the unique optimal transport map from $\mu$ to $\nu$ given as in McCann's theorem by $T(x)=\exp _{x}(-\operatorname{grad} \phi(x))$ for some function $\phi$. For $t \in[0,1]$, define the map $T_{t}: M \rightarrow M$ by

$$
T_{t}(x):=\exp _{x}(-t \operatorname{grad} \phi(x)),
$$

and define $\mu_{t}:=\left(T_{t}\right)_{\#} \mu \in \mathscr{P}_{2}(M)$ (so $\mu_{0}=\mu$ and $\mu_{1}=\nu$ ). Then the following two statements hold.
(a) $\operatorname{supp}\left(\mu_{t}\right) \subset Z_{t}(\operatorname{supp}(\mu), \operatorname{supp}(\nu))$ where

$$
Z_{t}(X, Y):=\left\{\begin{array}{l|l}
z \in M & \begin{array}{l}
d(x, z)=t d(x, y) \text { and } d(z, y)=(1-t) d(x, y) \\
\text { for some } x \in X, y \in Y
\end{array}
\end{array}\right\}
$$

(b) $\left(\mu_{t}\right)_{t \in[0,1]}$ is a constant speed geodesic in $\left(\mathscr{P}_{2}(M), W_{2}\right)$ from $\mu$ to $\nu$.

Proof. As discussed in the proof of [McC01, Lemma 7], $t \mapsto T_{t}(x)$ is a minimal geodesic for all $x$ where $\phi$ is differentiable (that is $\mu$-a.e.). Therefore, for such $x$, $d\left(T_{s}(x), T_{t}(x)\right)=|t-s| d(x, T(x))$.
In view of Proposition 1.5, we have for $0 \leq s \leq t \leq 1$,

$$
\begin{aligned}
W_{2}\left(\mu_{s}, \mu_{t}\right)^{2} & \leq \int_{M} d^{2}\left(T_{s}(x), T_{t}(x)\right) d \mu(x) \\
& =\int_{M}(t-s)^{2} d^{2}(x, T(x)) d \mu(x)=(t-s)^{2} W_{2}(\mu, \nu)^{2}
\end{aligned}
$$

Therefore, $\left(\mu_{t}\right)_{t \in[0,1]}$ is a constant speed geodesic in $\left(\mathscr{P}_{2}(M), W_{2}\right)$ from $\mu$ to $\nu$ due to Remark 1.18.

Furthermore, the absolute continuity of $\mu=\mu_{0}$ (with respect to vol) implies that $\mu_{t}=\left(T_{t}\right)_{\#} \mu$ is also absolutely continuous for all $t \in[0,1)$, say, with the density $\rho_{t}:=d \mu_{t} / d$ vol. The Jacobian determinant $\operatorname{det} D T_{t}$ satisfies the following MongeAmpère equation:

$$
\begin{equation*}
\rho(x)=\rho_{t}\left(T_{t}(x)\right) \operatorname{det} D T_{t}(x) \tag{1.7}
\end{equation*}
$$

This Monge-Ampère equation is essentially the Jacobian equation for change of variables; for more details see e.g. [San15, Section 1.7.6].
It is important to remark that the uniqueness of optimal transport plans and maps in $L^{2}$-Wasserstein spaces is attributable to the convexity of the quadratic cost function. This is not the case for $L^{1}$-Wasserstein spaces; in fact, there often are plenty of transport maps which give the same optimal transport cost. For an illustration, let us visit the simplest example of transportation on $\mathbb{R}$.

Example 1.26 (linear and quadratic costs in $\mathbb{R}$ ). Let $\mu=U(-1,0)$ and $\nu=$ $U(0,1)$ be the uniform probability measures on $\mathbb{R}$ whose supports are the real intervals $[-1,0]$ and $[0,1]$. All transport maps $T_{\#} \mu=\nu$ can have their domain and range restricted to the support of $\mu$ and $\nu$, respectively. For a linear cost function $c(x, y)=|x-y|$, the cost of any transport map $T:[-1,0] \rightarrow[0,1]$ can be computed as

$$
\begin{aligned}
\int_{-1}^{0}|x-T(x)| d \mu(x) & =\int_{-1}^{0}(T(x)-x) d \mu(x) \\
& =\int_{0}^{1} y d \nu(y)-\int_{-1}^{0} x d \mu(x)=\frac{1}{2}-\left(\frac{-1}{2}\right)=1
\end{aligned}
$$

which means all such maps $T$ have the same (and hence optimal) linear cost. On the other hand, the quadratic cost of $T$ satisfies, by Cauchy-Schwarz inequality,

$$
\int_{-1}^{0}|x-T(x)|^{2} d \mu(x) \geq \frac{\left(\int_{-1}^{0}|x-T(x)| d \mu(x)\right)^{2}}{\int_{-1}^{0} 1 d \mu(x)}=1
$$

and its infimum is attained if and only if $|x-T(x)|$ is constant, i.e., it is the translation map $T_{T}(x)=x+1$. In comparison, the reflection map $T_{R}(x)=-x$ has its quadratic cost equal to $\int_{-1}^{0}|x-(-x)|^{2} d x=\frac{4}{3}>1$. See Figure 1.1 below.

(a) Transport by translation

$$
T_{T}(x)=x+1
$$


(b) Transport by reflection $T_{R}(x)=-x$.

Figure 1.1: Comparison of two transport maps between uniform distributions $U(-1,0)$ and $U(0,1)$. The translation map has linear and quadratic cost of 1 . The reflection map has linear cost of 1 and quadratic cost of $\frac{4}{3}$.

## Chapter 2

## Otto's calculus

In his booklet [Ott01], Otto develops a framework in which the Wasserstein space is studied via a formal Riemannian calculus (where the word "formal" means that arguments are not completely rigorous, especially due to the concerned Riemannian manifold having infinite dimension). A curve in the Wasserstein space is described by a differential equation called the continuity equation, and the Wasserstein distance is the length of a minimizing curve. In this chapter, we explain Otto's calculus in a formal viewpoint of Riemannian geometry. Otto's calculus will reappear in Part IV of this thesis as an important concept to the definition of Erbar-Maas curvature.

### 2.1 Wasserstein space as "Riemannian manifold"

For simplicity, we avoid the discussion with specific regularity and often assume smoothness in most of the setting. Throughout this chapter, we let $\left(M^{n},\langle\cdot, \cdot\rangle\right)$ be a connected complete Riemannian manifold, and we restrict the Wasserstein space of our interest to the space $\mathscr{P}^{\infty}(M)$ consisting of all smooth positive density functions $\rho$ of an absolutely continuous measure $\mu$ with respect to the volume measure (that is, $\rho=d \mu / d \mathrm{vol}$ ):

$$
\mathscr{P}^{\infty}(M):=\left\{\rho \in C^{\infty}(M) \mid \rho>0, \int_{M} \rho=1\right\},
$$

where we identify the distance $W_{2}\left(\rho_{1}, \rho_{2}\right):=W_{2}\left(\mu_{1}, \mu_{2}\right)$ for $d \mu_{i}=\rho_{i} d$ vol. The space $\mathscr{P}^{\infty}(M)$ can be regarded as an $\infty$-dimensional manifold $\mathcal{M}$, and the tangent
space at $\rho \in \mathscr{P}^{\infty}(M)$ is given by

$$
T_{\rho} \mathscr{P}^{\infty}(M)=\left\{(\rho, s) \mid s \in C^{\infty}(M) \text { with } \int_{M} s=0\right\}
$$

where the curve which passes through $\rho$ and represents $s$ is $\rho(t)=\rho_{t}=\rho+t s \in$ $\mathscr{P}^{\infty}(M)$ for small enough $|t|$. The key idea in Otto's calculus is to identify this tangent space by another space, namely

$$
\operatorname{Tan}_{\rho} \mathscr{P}^{\infty}(M)=\left\{(\rho, \operatorname{grad} \phi) \mid \phi \in C^{\infty}(M)\right\},
$$

through the identification map $\operatorname{Tan}_{\rho} \mathscr{P}^{\infty}(M) \rightarrow T_{\rho} \mathscr{P}^{\infty}(M)$ given by

$$
(\rho, \operatorname{grad} \phi) \longmapsto(\rho, s=-\operatorname{div}(\rho \operatorname{grad} \phi)) .
$$

Here, we refer to Section 5.1 for basic knowledge of differential operators (such as grad, div, and the Laplacian $\Delta$ ).
To understand this identification map, note that the geodesic $\rho_{t}$ can be described by $d \mu_{t}=\rho_{t} d \mathrm{vol}=\left(T_{t}\right)_{\#} \rho d \mathrm{vol}$ where $T_{t}$ is the interpolation of an optimal map given via McCann's theorem by $T_{t}(x)=\exp _{x}(\operatorname{tgrad} \phi(x))$ for some function $\phi$ (when $\phi$ here is chosen to be the negative of the function $\phi$ in Corollary 1.25). It suffices to justify at $t=0$ that

$$
\left.\frac{\partial}{\partial t}\right|_{t=0}\left(T_{t}\right)_{\#}(\rho d \operatorname{vol})=-\operatorname{div}(\rho \operatorname{grad} \phi) d \operatorname{vol} .
$$

For a smooth test function $\varphi \in C^{\infty}(M)$ and $x \in M$, we have

$$
\left.\frac{\partial}{\partial t}\right|_{t=0} \varphi\left(T_{t}(x)\right)=\left\langle\operatorname{grad} \varphi(x),\left.\frac{\partial}{\partial t}\right|_{t=0} T_{t}(x)\right\rangle=\langle\operatorname{grad} \varphi(x), \operatorname{grad} \phi(x)\rangle,
$$

and therefore

$$
\begin{aligned}
\left.\frac{\partial}{\partial t}\right|_{t=0} \int_{M} \varphi\left(T_{t}(x)\right) \rho(x) & =\int_{M}\langle\operatorname{grad} \varphi(x), \rho(x) \operatorname{grad} \phi(x)\rangle \\
& =-\int_{M} \varphi(x) \operatorname{div}(\rho(x) \operatorname{grad} \phi(x)),
\end{aligned}
$$

as desired. The above identification $s=-\operatorname{div}(\rho \operatorname{grad} \phi))$ can be summarized as the so-called "continuation equation" for probability measures (or probability densities), which is stated in the following proposition.
Proposition 2.1. Consider any smooth curve $\gamma:[a, b] \rightarrow \mathscr{P}^{\infty}(M)$. Its tangent vector field $\gamma^{\prime}$ along the curve $\gamma$ can be described as $\gamma^{\prime}(t)=\gamma_{t}^{\prime}=\left(\rho_{t}, \operatorname{grad} \phi_{t}\right)$ which satisfies the following continuity equation:

$$
\begin{equation*}
\frac{\partial}{\partial t} \rho_{t}+\operatorname{div}\left(\rho_{t} \operatorname{grad} \phi_{t}\right)=0 \tag{2.1}
\end{equation*}
$$

Definition 2.2. Riemannian metric tensor $g$ on $\mathscr{P}^{\infty}(M)$ is defined as

$$
\begin{equation*}
g_{\rho}\left(s_{1}, s_{2}\right)=\int_{M} \rho\left\langle\operatorname{grad} \phi_{1}, \operatorname{grad} \phi_{2}\right\rangle, \tag{2.2}
\end{equation*}
$$

where $s_{i}=-\operatorname{div}\left(\rho \operatorname{grad} \phi_{i}\right)$ for $i=1,2$.
In principle, with this metric tensor, one can apply Riemannian geometry framework to the Wasserstein spaces and study on e.g. length of curves (which is to be discussed in the next section), Levi-Civita connection, parallel transport, geodesics, as well as Riemann curvature tensor and sectional curvature; for details we refer to the paper by Lott [Lot08]. For example, the following two theorems assert that the nonnegative sectional curvature of $\mathscr{P}^{\infty}(M)$ is induced from the underlying manifold $M$.

Theorem 2.3 ([Ott01]). $\mathscr{P}_{2}\left(\mathbb{R}^{n}\right)$ has nonnegative sectional curvature. Moreover, it is flat if and only if $n=1$.

Theorem 2.4 ([Lot08]). If $M$ has nonnegative sectional curvature, then $\mathscr{P}^{\infty}(M)$ has a nonnegative sectional curvature.

### 2.2 Benamou-Brenier formula

In [BB00], Benamou and Brenier introduce an alternative viewpoint to an optimal transportation problem by reintroducing the time-variable and considering it as a "dynamic problem". It is presented in the paper an influential formulation, later known as Benamou-Brenier formula, which is used for the Wasserstein distance computation. The aim of this section is to discuss the Banamou-Brenier formula in the perspective of formal Riemannian geometry and Otto's calculus.
First, let us recall basic facts from Riemannian geometry about the length of curves. Given a Riemannian manifold $(\mathcal{M}, g)$, the length of a curve $\gamma:[0,1] \rightarrow \mathcal{M}$ is $\ell(\gamma):=\int_{0}^{1}\left|\gamma^{\prime}(t)\right|_{g} d t$, where the norm is induced from Riemannian metric tensor $|v|_{g}:=g(v, v)^{1 / 2}$. The distance of two points $p, q \in M$ is

$$
d(p, q)=\inf _{\gamma} \ell(\gamma)=\inf _{\gamma} \int_{0}^{1}\left|\gamma^{\prime}(t)\right|_{g} d t
$$

where the infimums are taken over all curves $\gamma:[0,1] \rightarrow \mathcal{M}$ with $\gamma(0)=p$ and $\gamma(1)=q$. Such curves $\gamma$ attaining the infimum are called minimal geodesics, and they are not unique since the length $\ell(\gamma)$ is unchanged under any reparametrization. However, the distance of two points can also be described by a curve which minimizes its kinetic energy $E(\gamma):=\int_{0}^{1}\left|\gamma^{\prime}(t)\right|_{g}^{2} d t$. Such a minimizing curve (if exists) is indeed a minimal geodesic with constant speed.

Proposition 2.5. Let $(\mathcal{M}, g)$ be a Riemannian manifold. Given $p, q \in \mathcal{M}$, then

$$
\begin{equation*}
d(p, q)=\inf _{\gamma} \int_{0}^{1}\left|\gamma^{\prime}(t)\right|_{g} d t=\inf _{\gamma}\left(\int_{0}^{1}\left|\gamma^{\prime}(t)\right|_{g}^{2} d t\right)^{\frac{1}{2}} \tag{2.3}
\end{equation*}
$$

where the infimums are taken over all curves $\gamma:[0,1] \rightarrow \mathcal{M}$ with $\gamma(0)=p$ and $\gamma(1)=q$.

Proof. The Cauchy-Schwarz inequality gives

$$
\int_{0}^{1}\left|\gamma^{\prime}(t)\right|_{g} d t \leq\left(\int_{0}^{1}\left|\gamma^{\prime}(t)\right|_{g}^{2} d t\right)^{\frac{1}{2}}\left(\int_{0}^{1} 1 d t\right)^{\frac{1}{2}}=\left(\int_{0}^{1}\left|\gamma^{\prime}(t)\right|_{g}^{2} d t\right)^{\frac{1}{2}}
$$

with the equality holds when $\left|\gamma^{\prime}(t)\right|_{g}$ is constant. Moreover, since the length of curves is unchanged under a reparametrization, the infimum $\inf _{\gamma} \int_{0}^{1}\left|\gamma^{\prime}(t)\right|_{g} d t$ can be restricted to those curves $\gamma$ with constant speed. In the case of constant speed $\gamma$, the above inequality holds with equality, which in turn yields the equality in (2.3).

In view of Otto's calculus, one can apply the above description of distance via minimizing energy to the Wasserstein space $\mathscr{P}^{\infty}(M)$ in order to recover the Wasserstein distance $W_{2}$. A precise statement and its rigorous proof (without formal Riemannian calculus arguments) can be found in original paper by Benamou and Brenier [BB00].

Theorem 2.6 (Benamou-Brenier formula). For $\rho_{0}, \rho_{1} \in \mathscr{P}^{\infty}(M)$,

$$
\begin{equation*}
W_{2}\left(\rho_{0}, \rho_{1}\right)=\inf _{\left(\rho_{t}, \phi_{t}\right) \in C E\left(\rho_{0}, \rho_{1}\right)}\left\{\int_{0}^{1} \int_{M} \rho_{t}(x)\left|\operatorname{grad} \phi_{t}(x)\right|^{2} d \operatorname{vol}(x) d t\right\}^{1 / 2} \tag{2.4}
\end{equation*}
$$

where the infimum is taken over $C E\left(\rho_{0}, \rho_{1}\right)$ the set of all curves $\left(\rho_{t}\right)_{t \in[0,1]} \in$ $\mathscr{P}^{\infty}(M)$ and $\left(\psi_{t}\right)_{t \in[0,1]} \in C^{\infty}(M)$ which satisfy the following continuity equation

$$
\begin{equation*}
\frac{\partial}{\partial t} \rho_{t}+\operatorname{div}\left(\rho_{t} \operatorname{grad} \psi_{t}\right)=0 \quad \forall t \in[0,1], \tag{2.5}
\end{equation*}
$$

with boundary condition $\rho_{t=0}=\rho_{0}$ and $\rho_{t=1}=\rho_{1}$.

### 2.3 Heat flow as gradient flow of entropy

In his original paper [Ott01], Otto uses a gradient flow structure to describe the heat equation. Let us first recall the definition of gradient flow for an abstract Riemannian manifold $(\mathcal{M}, g)$.

Definition 2.7 (gradient flow). Let $(\mathcal{M}, g)$ be a Riemannian manifold, and $E$ : $\mathcal{M} \rightarrow \mathbb{R}$ be a function(al) on $\mathcal{M}$. For a fixed time-interval $I \subset \mathbb{R}$, a curve $\rho: I \rightarrow \mathcal{M}$ (written as $\rho_{t}:=\rho(t)$ ), is called an integral curve of gradient flow of $E$ if the following equation holds true

$$
\begin{equation*}
\frac{\partial}{\partial t} \rho_{t}=-(\operatorname{grad} E) \rho_{t} \tag{2.6}
\end{equation*}
$$

or in other words,

$$
\begin{equation*}
g_{\rho_{t}}\left(\frac{\partial}{\partial t} \rho_{t}, s\right)=-D E\left(\rho_{t}\right)(s), \quad \forall s \in T_{\rho(t)} \mathcal{M} \tag{2.7}
\end{equation*}
$$

Definition 2.8 (entropy). For a complete Riemannian manifold ( $M^{n},\langle\cdot, \cdot\rangle$ ), an entropy functional (with respect to volume measure) $\mathcal{E}_{\text {vol }}: \mathscr{P}(M) \rightarrow \mathbb{R} \cup\{+\infty\}$ is defined as

$$
\mathcal{E}_{\mathrm{vol}}(\mu):= \begin{cases}\int_{M} \rho \log \rho & \text { if } \mu \ll \operatorname{vol} \text { with } \rho=\frac{d \mu}{d \mathrm{vol}}, \\ +\infty & \text { otherwise. }\end{cases}
$$

For simplicity in this chapter, we define the entropy for the space of densities $\mathscr{P}^{\infty}(M)$ as

$$
\mathcal{E}_{\mathrm{vol}}(\rho):=\int_{M} \rho \log \rho
$$

for all $\rho \in \mathscr{P}^{\infty}(M)$.
Definition 2.9 (heat flow). For a fixed time-interval $I \subset \mathbb{R}$, a curve $u: I \rightarrow \mathcal{M}$, is called an integral curve of heat flow if it is a solution of the heat equation: $\frac{\partial}{\partial t} u_{t}=\Delta u_{t}$, where $\Delta:=\operatorname{div}(\mathrm{grad})$.

We refer to Section 5.5 for details about the heat equation. Now we give a formal proof of the following theorem that the gradient flow of the entropy is in fact the heat flow.

Theorem 2.10. An integral curve of gradient flow of the entropy functional $\mathcal{E}_{\text {vol }}$ on the Riemannian manifold $(\mathcal{M}, g)=\left(\mathscr{P}^{\infty}(M), W_{2}\right)$ is a solution to the heat equation $\frac{\partial}{\partial t} \rho_{t}=\Delta \rho_{t}$.

Proof. Let $\rho_{t}$ be the integral curve of gradient flow of $\mathcal{E}_{\text {vol }}$. For a fixed $t \in I$ and fixed $s \in T_{\rho_{t}} \mathcal{M}$, we first compute $D \mathcal{E}_{\text {vol }}\left(\rho_{t}\right)(s)$ through a curve $c:(-\varepsilon, \varepsilon) \rightarrow \mathcal{M}$ such that $c(0)=\rho_{t}$ and $c^{\prime}(0)=s$. We have

$$
D \mathcal{E}_{\mathrm{vol}}\left(\rho_{t}\right)(s)=\left.\frac{d}{d \tau}\right|_{\tau=0} \mathcal{E}_{\mathrm{vol}}(c(\tau))=\int_{M}(1+\log c(0)) c^{\prime}(0)=\int_{M}\left(1+\log \rho_{t}\right) s
$$

With the identification $s=-\operatorname{div}\left(\rho_{t} \operatorname{grad} \phi\right)$, we have

$$
\begin{aligned}
g_{\rho_{t}}\left(\frac{\partial}{\partial t} \rho_{t}, s\right) & =-D \mathcal{E}_{\mathrm{vol}}\left(\rho_{t}\right)(s)=\int_{M}\left(1+\log \rho_{t}\right) \operatorname{div}\left(\rho_{t} \operatorname{grad} \phi\right) \\
& =-\int_{M} \rho_{t}\left\langle\operatorname{grad}\left(1+\log \rho_{t}\right), \operatorname{grad} \phi\right\rangle \\
& =-\int_{M} \rho_{t}\left\langle\operatorname{grad}\left(\log \rho_{t}\right), \operatorname{grad} \phi\right\rangle .
\end{aligned}
$$

Since the above equality holds for all $s \in T_{\rho_{t}} \mathcal{M}$, in view of (2.2) we deduce that

$$
\frac{\partial}{\partial t} \rho_{t}=\operatorname{div}\left(\rho_{t} \operatorname{grad}\left(\log \rho_{t}\right)\right)=\operatorname{div}\left(\operatorname{grad}\left(\rho_{t}\right)\right)=\Delta\left(\rho_{t}\right)
$$

### 2.4 Jordan-Kinderlehrer-Otto theorem

Theorem 2.10 was proved in a rigorous manner by Jordan-Kinderlehrer-Otto [JKO98] (and hence referred to as JKO theorem), where they consider the following discrete-time minimizing scheme. Fix a functional $E: \mathcal{M} \rightarrow \mathbb{R}$ and fix $\rho_{0} \in \mathcal{M}$. For a given time step size $h>0$, let $\rho^{(0)}=\rho_{0}$ and

$$
\begin{equation*}
\rho^{(k)}=\operatorname{argmin}_{\rho \in \mathcal{M}}\left(\frac{1}{2 h} W_{2}\left(\rho^{(k-1)}, \rho\right)+E(\rho)\right) . \tag{2.8}
\end{equation*}
$$

The theorem states that the step functions $\rho_{h}(t):=\rho^{(k)}$ for $t \in[k h,(k+1) h)$ converges as $h \rightarrow 0$ to the solution of the heat equation $\frac{\partial}{\partial t} \rho=\Delta \rho$ with $\rho_{t=0}=\rho_{0}$.
Here we give a heuristic argument why this minimizing scheme is indeed the discretization of the integral curve. We explain this in a general setting of $(\mathcal{M}, g)$ with the distance function $d$ induced from $g$ and a functional $F: \mathcal{M} \rightarrow \mathbb{R}$.

For $x \in \mathcal{M}$, consider $f_{h}(y):=\frac{1}{2 h} d_{x}(y)^{2}+F(y)$, where $d_{x}(\cdot)$ is the distance from $x$. Let $x_{h}$ solve $\min _{y} f_{h}(y)$. As $h \rightarrow 0$, we know that $x_{h} \rightarrow x$. Now consider a geodesic curve $c:(-\varepsilon, \varepsilon) \rightarrow \mathcal{M}$ such that $c(0)=x$ and $c(h)=x_{h}$. We would like to justify that

$$
c^{\prime}(0)=-\operatorname{grad} F(x) .
$$

Consider this curve $c$ via polar coordinates: $c(t)=\exp _{x}(r(t) v(t))$ where $r(t) \in \mathbb{R}_{\geq 0}$ and $v(t) \in T_{x} M$ is a unit tangent vector. Consider another curve $\gamma(s):=\exp _{x}(s$. $v(t))$, so $\operatorname{grad} d_{x}(c(t))=\gamma^{\prime}(r(t))=D \exp _{x}(r(t) v(t))(v(t))$.

The assumption that $c(h)=x_{h}$ minimizes $f_{h}(\cdot)$ means

$$
0=\left.\operatorname{grad} f_{h}(y)\right|_{y=c(h)}=\frac{1}{h} d_{x}(c(h)) \cdot \operatorname{grad} d_{x}(c(h))+\operatorname{grad} F(c(h)) .
$$

Passing $h \rightarrow 0$ and note that $d_{x}(c(t))=r(t)$, we have

$$
\begin{aligned}
-\operatorname{grad} F(x) & =\lim _{h \downarrow 0} \frac{1}{h} r(h) \operatorname{grad} d_{x}(c(h)) \\
& =\lim _{h \downarrow 0} \frac{r(h)}{h} D \exp _{x}(r(h) v(h))(v(h)) \\
& =r^{\prime}(0) \underbrace{D \exp _{x}\left(0_{x}\right)}_{=\mathrm{Id}}(v(0))=r^{\prime}(0) v(0) .
\end{aligned}
$$

On the other hand, $c^{\prime}(0)=D \exp _{x}\left(0_{x}\right)\left(r^{\prime}(0) v(0)+r(0) v^{\prime}(0)\right)=r^{\prime}(0) v(0)$, and thus $c^{\prime}(0)=-\operatorname{grad} F(x)$ as desired.

## Chapter 3

## Ricci curvature on metric measure spaces

This chapter features a celebrated result from two independent papers by LottVillani [LV09] and Sturm [Stu06] that a weighted Riemannian manifold $M$ has Ricci curvature bounded below by $K \in \mathbb{R}$ if and only if the entropy functional on the Wasserstein space $\left(\mathscr{P}_{2}(M), W_{2}\right)$ is " $K$-convex" along $W_{2}$-geodesics. The latter part of this equivalence does not require differential structure, and hence it allows them to define on metric measure spaces a synthetic notion of Ricci curvature, which is often referred to as Lott-Sturm-Villani (LSV) curvature.

Here we only discuss this main result in the case of nonweighted Riemannian manifold, which was proved earlier by von Renesse and Sturm [vRS05]. The case of weighted manifolds is postponed to Chapter 6.

### 3.1 Ricci lower curvature bound via displacement convexity

Let us start with the definitions of the relative entropy and the displacement convexity on metric measure spaces.

Definition 3.1 (relative entropy). Let $(X, d, m)$ be a metric measure space with a reference measure $m \in \mathscr{P}(X)$. The relative entropy (with respect to $m$ ) is defined on $\mathscr{P}(X)$ as

$$
\mathcal{E}_{m}(\mu):= \begin{cases}\int_{X} \rho \log \rho d m=\int_{X} \log \rho d \mu & \text { if } \mu \ll m \text { with } \rho=\frac{d \mu}{d m}  \tag{3.1}\\ +\infty & \text { otherwise } .\end{cases}
$$

Definition 3.2 (displacement convexity). Let $K \in \mathbb{R}$ and let $(X, d)$ be a metric space. A functional $E: \mathscr{P}_{2}(X) \rightarrow X$ is said to be displacement $K$-convex if for any $\mu_{0}, \mu_{1} \in \mathscr{P}_{2}(X)$, there exists a constant speed geodesic $\left(\mu_{t}\right)_{t \in[0,1]}$ in $\left(\mathscr{P}_{2}(X), W_{2}\right)$ from $\mu_{0}, \mu_{1}$ satisfies

$$
\begin{equation*}
E\left(\mu_{t}\right) \leq(1-t) E\left(\mu_{0}\right)+t E\left(\mu_{1}\right)-\frac{1}{2} K t(1-t) W_{2}\left(\mu_{0}, \mu_{1}\right)^{2} . \tag{3.2}
\end{equation*}
$$

The word "displacement" here indicates that the curve $\left(\mu_{t}\right)_{t \in[0,1]}$ in consideration is a geodesic with respect to the $W_{2}$-metric, and it is not the linear interpolation $\mu_{t}=(1-t) \mu_{0}+t \mu_{1}$. Let us also remark that for the above definition to be valid, the Wasserstein space $\left(\mathscr{P}_{2}(X), W_{2}\right)$ must be a geodesic space. In view of Theorem 1.20 , it is sensible to assume $(X, d)$ to be a geodesic space as well.

Now we are ready to state the main result which characterizes the Ricci lower bound via the displacement convexity of the entropy functional.

Theorem 3.3 (main theorem, [vRS05]). Given a smooth complete Riemannian manifold $(M, g)$ and any $K \in \mathbb{R}$, the following properties are equivalent:
(i) $\operatorname{Ric}(M) \geq K$, that is, $\operatorname{Ric}_{x}(v, v) \geq K|v|^{2}$ for all $x \in M, v \in T_{x} M$.
(ii) The entropy $\mathcal{E}_{\text {vol }}(\cdot)$ is displacement $K$-convex on $\mathscr{P}_{2}(M)$.

Here and henceforth, we use $\operatorname{Ric}(M)$ as the convention for the greatest lower bound of the Ricci curvature on $M$ :

$$
\operatorname{Ric}(M):=\inf \left\{\frac{\operatorname{Ric}_{x}(v, v)}{|v|^{2}}: x \in M, v \in T_{x} M\right\} .
$$

The main theorem gives rise to the following Ricci curvature notion by Lott-SturmVillani on metric measure spaces.

Definition 3.4 (Lott-Sturm-Villani). A metric measure space $(X, d, m)$ is said to have Ricci curvature bound below by $K \in \mathbb{R}$, denoted as $\operatorname{Ric}_{\infty}(X) \geq K$, if the relative entropy $\mathcal{E}_{m}(\cdot)$ is displacement $K$-convex, that is, for any $\mu_{0}, \mu_{1} \in \mathscr{P}_{2}(X)$, there exists a constant speed geodesic $\left(\mu_{t}\right)_{t \in[0,1]}$ in $\mathscr{P}_{2}(X)$ from $\mu_{0}$ to $\mu_{1}$ such that

$$
\begin{equation*}
\mathcal{E}_{m}\left(\mu_{t}\right) \leq(1-t) \mathcal{E}_{m}\left(\mu_{0}\right)+t \mathcal{E}_{m}\left(\mu_{1}\right)-\frac{1}{2} K t(1-t) W_{2}\left(\mu_{0}, \mu_{1}\right)^{2} \tag{3.3}
\end{equation*}
$$

holds true for all $t \in[0,1]$.

Remark $3.5(C D(K, N)$ spaces). Such a metric measure space $(X, d, m)$ is usually referred to as a $C D(K, \infty)$ space, where $C D$ stands for curvature-dimension. A more general convention of $C D(K, N)$, where $N \in(0, \infty]$, means $K$ is a lower bound of Ricci curvature, and $N$ is "an upper bound of the dimension" (in case $X$ is a manifold). The definition of $C D(K, N)$ spaces where the $N<\infty$ requires the Rényi entropy $S_{N}(\mu):=-\int_{X} \rho^{1-1 / N} d m$ to satisfy a more complicated $K$ -displacement-convex inequality than (3.2). However, the case $N<\infty$ is not our focus, and we refer readers to a concise summary by Ohta [Oht14, Section 7.4.2] and references therein.

We come back to a proof sketch of the main theorem, which we follow from [Vil09]. Here we only prove the forward implication, and we ignore all regularity concerns. For readers' interest, we will present in Chapter 8 different ideas of proof from the original papers [vRS05] and Cordero-Erausquin, McCann, and Schmuckenschläger's [CEMS01] (the latter of which proves in the case of $\operatorname{Ric}(M) \geq 0)$.

Proof sketch of Theorem 3.3. First suppose that $\mu_{0}, \mu_{1} \ll$ vol otherwise the right-hand-side of (3.3) is infinity. By Theorem 1.24, there exists a unique geodesic $\left(\mu_{t}\right)_{t \in[0,1]}$ in $\mathscr{P}_{2}(M)$ connecting $\mu_{0}$ and $\mu_{1}$, namely $\mu_{t}=\left(T_{t}\right)_{\#} \mu_{0}$ where $T_{t}(x)=$ $\exp _{x}(-t \operatorname{grad} \phi(x))$ for some $\phi: M \rightarrow \mathbb{R}$. From the viewpoint that $t \mapsto T_{t}(x)$ is a minimal constant speed geodesic (see Corollary 1.25), we may denote this geodesic by $\gamma_{x}:[0,1] \rightarrow M$, that is,

$$
\gamma_{x}(t):=T_{t}(x)
$$

For all $t, \mu_{t} \ll \operatorname{vol}$, so it can be written as $d \mu_{t}(x)=\rho_{t}(x) d \operatorname{vol}(x)$ with a finite probability density $\rho_{t}(x)$. Moreover, $J_{x}(t):=D T_{t}(x)$ satisfies the Monge-Ampère equation (1.7):

$$
\begin{equation*}
\rho_{0}(x)=\rho_{t}\left(F_{t}(x)\right) \operatorname{det} J_{x}(t) . \tag{3.4}
\end{equation*}
$$

The pushforward relation and Monge-Ampère equation give

$$
\begin{align*}
\mathcal{E}_{\mathrm{vol}}\left(\mu_{t}\right)=\int_{M} \log \rho_{t}(x) d \mu_{t}(x) & =\int_{M} \log \left(\frac{\rho_{0}(x)}{\operatorname{det} J_{x}(t)}\right) d \mu_{0}(x) \\
& =\mathcal{E}_{\mathrm{vol}}\left(\mu_{0}\right)-\int_{M} \log \left(\operatorname{det} J_{x}(t)\right) d \mu_{0}(x) \tag{3.5}
\end{align*}
$$

To derive the second derivatives of $\mathcal{E}_{\text {vol }}\left(\mu_{t}\right)$, we compute that of the term $\log \left(\operatorname{det} J_{x}(t)\right)$ for a fixed $x$ as follows:

$$
\begin{aligned}
\frac{d}{d t} \log \left(\operatorname{det} J_{x}(t)\right) & =\frac{1}{\operatorname{det} J_{x}(t)} \cdot \operatorname{det} J_{x}(t) \operatorname{tr}\left(J_{x}^{\prime}(t) J_{x}^{-1}(t)\right) \\
& =\operatorname{tr}\left(J_{x}^{\prime}(t) J_{x}^{-1}(t)\right)
\end{aligned}
$$

and

$$
\begin{align*}
\frac{d^{2}}{d t^{2}} \log \left(\operatorname{det} J_{x}(t)\right) & =\operatorname{tr}\left(J_{x}^{\prime \prime}(t) J_{x}^{-1}(t)-\left(J_{x}^{\prime}(t) J_{x}^{-1}(t)\right)^{2}\right) \\
& \leq \operatorname{tr}\left(J_{x}^{\prime \prime}(t) J_{x}^{-1}(t)\right) \tag{3.6}
\end{align*}
$$

Remark that the derivative of an exponential map is a matrix consisting of Jacobi fields, that is, for any $u, v \in T_{x} M$, by considering the geodesic variation $F(s, t):=$ $\exp _{x}(s u+t v)$, we infer that $\left.\frac{\partial}{\partial s}\right|_{s=0} F(s, t)=D \exp _{x}(t v)(u)$ is a Jacobi field. Here, it means for any fixed $x \in M, J_{x}(t)=D T_{t}(x)$ is a matrix whose entries are Jacobi fields along $\gamma_{x}(t)$, so it satisfies the following Jacobi equation:

$$
J_{x}^{\prime \prime}(t)+A_{x}(t) J_{x}(t)=0,
$$

where the map $A_{x}(t): T_{\gamma_{x}(t)} M \rightarrow T_{\gamma_{x}(t)} M$ is given by

$$
A_{x}(t)(v):=R\left(\gamma_{x}^{\prime}(t), v\right) \gamma_{x}^{\prime}(t),
$$

and $R$ is Riemann curvature tensor. Moreover, the trace of $A_{t}(x)$ can be expressed in terms of Ricci curvature:

$$
\begin{equation*}
-\operatorname{tr}\left(J_{x}^{\prime \prime}(t) J_{x}^{-1}(t)\right)=\operatorname{tr}\left(A_{x}(t)\right)=\operatorname{Ric}_{\gamma_{x}(t)}\left(\gamma_{x}^{\prime}(t), \gamma_{x}^{\prime}(t)\right) \tag{3.7}
\end{equation*}
$$

The Ricci curvature assumption $\operatorname{Ric}(M) \geq K$ implies $\operatorname{Ric}_{\gamma_{x}(t)}\left(\gamma_{x}^{\prime}(t), \gamma_{x}^{\prime}(t)\right) \geq$ $K\left|\gamma_{x}^{\prime}(t)\right|$. Furthermore, since the curve $\gamma_{x}(t)=T_{t}(x)$ is a minimal constant speed geodesic, we have $\left|\gamma_{x}^{\prime}(t)\right|=d\left(x, T_{1}(x)\right)$ for all $t \in[0,1]$. It then follows from (3.5), (3.6) and (3.7) that

$$
\frac{d^{2}}{d t^{2}} \mathcal{E}_{\mathrm{vol}}\left(\mu_{t}\right) \geq \int_{M} K d\left(x, T_{1}(x)\right) d \mu_{0}(x)=K W_{2}\left(\mu_{0}, \mu_{1}\right)^{2} .
$$

Let $\tilde{E}(t):=\mathcal{E}_{\text {vol }}\left(\mu_{t}\right)+\frac{1}{2} K t(1-t) W_{2}\left(\mu_{0}, \nu_{1}\right)^{2}$. Note that $\tilde{E}(0)=\mathcal{E}_{\text {vol }}\left(\mu_{0}\right), \tilde{E}(1)=$ $\mathcal{E}_{\text {vol }}\left(\mu_{1}\right)$ and

$$
\frac{d^{2}}{d t^{2}} \tilde{E}(t)=\frac{d^{2}}{d t^{2}} \mathcal{E}_{\text {vol }}\left(\mu_{t}\right)-K W_{2}\left(\mu_{0}, \mu_{1}\right)^{2} \geq 0
$$

Therefore $\tilde{E}(t)$ is a convex function and it satisfies $\tilde{E}(t) \leq(1-t) \tilde{E}(0)+t \tilde{E}(1)$, which means

$$
\mathcal{E}_{\mathrm{vol}}\left(\mu_{t}\right) \leq(1-t) \mathcal{E}_{\mathrm{vol}}\left(\mu_{0}\right)+t \mathcal{E}_{\mathrm{vol}}\left(\mu_{1}\right)-\frac{1}{2} K t(1-t) W_{2}\left(\mu_{0}, \mu_{1}\right)^{2} .
$$

Example 3.6 (Euclidean space $\mathbb{R}^{n}$ ). For all $n \in \mathbb{N}$, the Euclidean space $\mathbb{R}^{n}$ has zero (sectional and Ricci) curvature everywhere, and it indeed satisfies $\operatorname{Ric}\left(\mathbb{R}^{n}\right) \geq$ 0 . To see that the relative entropy $\mathcal{E}_{\text {vol }}$ (with respect to the Lebesgue/volume measure) is displacement 0 -convex, one can argue similarly to the proof given above. From Brenier's theorem (and Corollary 1.25), a constant speed $W_{2}$-geodesic $\left(\mu_{t}\right)_{t \in[0,1]}$ can be described as $\mu_{t}=\left(T_{t}\right)_{\#} \mu_{0}$, where $T_{t}$ is the interpolation of the optimal transport map given in an explicit form of $T_{t}(x)=x-t \nabla \phi(x)$, for some $\phi$. The Jacobian matrix is $J_{x}(t)=D T_{t}(x)=\operatorname{Id}-t \operatorname{Hess} \phi(x)$, so its second derivative in $t$ vanishes (or alternatively, one may see $J_{x}^{\prime \prime}(t)=0$ from the Jacobi equation with zero Riemann curvature tensor). It is then implies by (3.6) that $\log \left(\operatorname{det} J_{x}(t)\right)$ is concave in $t$, and by (3.5) that $\mathcal{E}_{\text {vol }}\left(\mu_{t}\right)$ is convex, i.e., $\mathcal{E}_{\text {vol }}\left(\mu_{t}\right) \leq$ $(1-t) \mathcal{E}_{\text {vol }}\left(\mu_{0}\right)+t \mathcal{E}_{\text {vol }}\left(\mu_{1}\right)$.

One consequence of the displacement convexity is to provide an alternative proof of the following variation of Brunn-Minkowski inequality on $\mathbb{R}^{n}$ (see also [Vil09, Theorem 18.5]).

Theorem 3.7 (Brunn-Minkowski). Let vol be the Lebesgue measure of $\mathbb{R}^{n}$. Then for compact subsets $A_{0}, A_{1} \subset \mathbb{R}^{n}$, the following inequality holds

$$
\operatorname{vol}\left((1-t) A_{0}+t A_{1}\right) \geq \operatorname{vol}\left(A_{0}\right)^{1-t} \operatorname{vol}\left(A_{1}\right)^{t}
$$

Proof. For each $\in\{0,1\}$, let $\mu_{i}$ be the uniform probability distribution whose support is $A_{i}$, that is, $\mu_{i}(E)=\frac{\operatorname{vol}\left(E \cap A_{i}\right)}{\operatorname{vol}\left(A_{i}\right)}$ for any measurable set $E \subset \mathbb{R}^{n}$. Its density function satisfies $\rho_{i}:=\frac{d \mu_{i}}{d \mathrm{vol}}=\frac{1}{\operatorname{vol}\left(A_{i}\right)} \mathbb{1}_{A_{i}}$, and the entropy can be computed as

$$
\mathcal{E}_{\mathrm{vol}}\left(\mu_{i}\right)=\int_{\mathbb{R}^{n}} \log \rho_{i} d \mu_{i}=\int_{A_{i}} \log \left(\frac{1}{\operatorname{vol}\left(A_{i}\right)}\right) d \mu_{0}=-\log \left(\operatorname{vol}\left(A_{i}\right)\right) .
$$

Let $\mu_{t}=\left(T_{t}\right)_{\#} \mu_{0}$ be the constant speed $W_{2}$-geodesic from $\mu_{0}$ to $\mu_{1}$, obtained from Corollary 1.25. For $t \in(0,1)$, let $A_{t}:=\operatorname{supp}\left(\mu_{t}\right)$. Then Jensen's inequality applied to the convex function $s \mapsto s \log s$ yields

$$
\begin{aligned}
\mathcal{E}_{\mathrm{vol}}\left(\mu_{t}\right) & =\int_{\mathbb{R}^{n}} \rho_{t} \log \rho_{t} d \operatorname{vol} \\
& =\operatorname{vol}\left(A_{t}\right) \int_{A_{t}} \rho_{t} \log \rho_{t} \frac{d \operatorname{vol}}{\operatorname{vol}\left(A_{t}\right)} \\
& \geq \operatorname{vol}\left(A_{t}\right)\left(\int_{A_{t}} \rho_{t} \frac{d \operatorname{vol}}{\operatorname{vol}\left(A_{t}\right)}\right) \log \left(\int_{A_{t}} \rho_{t} \frac{d \operatorname{vol}}{\operatorname{vol}\left(A_{t}\right)}\right) \\
& =\operatorname{vol}\left(A_{t}\right) \cdot \frac{1}{\operatorname{vol}\left(A_{t}\right)} \cdot \log \left(\frac{1}{\operatorname{vol}\left(A_{t}\right)}\right)=-\log \left(\operatorname{vol}\left(A_{t}\right)\right)
\end{aligned}
$$

Since $\mathcal{E}_{\text {vol }}$ is displacement 0 -convex, we have

$$
\begin{aligned}
-\log \left(\operatorname{vol}\left(A_{t}\right)\right) \leq \mathcal{E}_{\mathrm{vol}}\left(\mu_{t}\right) & \leq(1-t) \mathcal{E}_{\mathrm{vol}}\left(\mu_{0}\right)+t \mathcal{E}_{\mathrm{vol}}\left(\mu_{1}\right) \\
& \leq-(1-t) \log \left(\operatorname{vol}\left(A_{0}\right)\right)-t \log \left(\operatorname{vol}\left(A_{1}\right)\right)
\end{aligned}
$$

which implies $\operatorname{vol}\left(A_{t}\right) \geq \operatorname{vol}\left(A_{0}\right)^{1-t} \operatorname{vol}\left(A_{1}\right)^{t}$.
Finally we note that $\operatorname{vol}\left((1-t) A_{0}+t A_{1}\right) \geq \operatorname{vol}\left(A_{t}\right)$; this is due to that fact that $\operatorname{supp}\left(A_{t}\right)$ lies inside $\left((1-t) A_{0}+t A_{1}\right)$ as $\mu_{t}=\left(T_{t}\right)_{\#} \mu_{0}$ where $t \mapsto T_{t}(x)$ is a constant speed geodesic for all $x$.

## Chapter 4

## Bonnet-Myers theorem and diameter rigidity

The Bonnet-Myers diameter bound theorem [Mye41] is a classical theorem in Riemannian geometry. It states that a connected complete manifold $\left(M^{n}, g\right)$ with Ricci curvature bounded from below by $K>0$ is a compact manifold, and its diameter is bounded from above by the diameter of the round sphere of dimension $n$ with Ricci curvature $K$. Furthermore, there is a rigidity theorem, known as Cheng's rigidity [Che75], which asserts that the diameter upper bound is attained if and only if $M^{n}$ is isometric to an $n$-dimensional sphere.

The Bonnet-Myers diameter bound theorem is the main theme of this thesis, and its analogous result will be reoccurring in the context of every curvature notion for discrete spaces.
Theorem 4.1 (Bonnet-Myers). Let $\left(M^{n}, g\right)$ be a connected and complete Riemannian manifold. Suppose there exists $r>0$ such that the Ricci curvature satisfies

$$
\operatorname{Ric}_{x}(v) \geq \frac{n-1}{r^{2}}>0
$$

for all $x \in M$ and unit vector $v \in T_{x} M$. Then $M$ is compact and its diameter is bounded from above by $\operatorname{diam}(M) \leq \pi r$.

Equivalently, the diameter bound can be written as

$$
\operatorname{diam}(M) \leq \pi \sqrt{\frac{n-1}{K}}
$$

given that $K:=\operatorname{Ric}(M)>0$. Here we provide a standard proof of this theorem via the second variation formula of the energy (see e.g. the book by do Carmo [dC92]).

Proof. Let $x, y \in M$ be two arbitrary points, and let $a:=d(x, y)$. By Hopf-Rinow theorem, the completeness of $M$ implies that there exists a minimal constant-speed geodesic $c:[0, a] \rightarrow M$ such that $c(0)=x, c(a)=y,\left|c^{\prime}(t)\right|=1$ for all $t \in[0, a]$. It suffices to prove that $a \leq \pi r$. The compactness of $M$ would then follow from $M$ being complete and bounded .
First, we construct proper variations of $c$ as follows. Choose unit vectors $e_{i} \in T_{x} M$ such that $\left\{e_{1}, e_{2}, \ldots, e_{n-1}, c^{\prime}(0)\right\}$ forms an orthonormal basis of $T_{x} M$. For each $1 \leq i \leq n-1$, let $V_{i}$ be a parallel vector field along $c$ such that $V_{i}(0)=e_{i}$. Note that

$$
\frac{d}{d t}\left\langle V_{i}(t), V_{j}(t)\right\rangle=\left\langle\frac{D}{d t} V_{i}(t), V_{j}(t)\right\rangle+\left\langle V_{i}(t), \frac{D}{d t} V_{j}(t)\right\rangle=0
$$

since $\frac{D}{d t} V_{i}(t)=\frac{D}{d t} V_{j}(t)=0$ from being parallel. It means $\left\langle V_{i}(t), V_{j}(t)\right\rangle$ is constant and $\left\langle V_{i}(t), V_{j}(t)\right\rangle=\left\langle e_{i}, e_{j}\right\rangle=\delta_{i j}$.
For each $i$, define $X_{i}(t):=\sin \left(\frac{\pi t}{a}\right) V_{i}(t)$, and let $F_{i}:(-\varepsilon, \varepsilon) \times[0, a] \rightarrow M$ be a variation of $c$ whose variational vector field is $X_{i}$, that is

$$
F_{i}(0, t)=c(t) \quad \text { and } \quad \frac{\partial}{\partial s} F_{i}(s, t)=X_{i}(t) .
$$

Since $X_{i}(0)=X_{i}(a)=0$, it means that for every $s \in(-\varepsilon, \varepsilon)$ the curve $F_{i}(s, \cdot)$ has the same endpoints as the curve $c$, that is, $F_{i}$ is a proper variation of $c$.

The energy for the curve $F_{i}(s, \cdot)$ is defined by

$$
E_{i}(s):=\frac{1}{2} \int_{0}^{a}\left|\frac{\partial}{\partial t} F_{i}(s, t)\right|^{2} d t
$$

The second variation formula of energy gives

$$
\begin{align*}
E_{i}^{\prime \prime}(0) & =\int_{0}^{a}\left|\frac{D}{d t} X_{i}(t)\right|^{2}-\left\langle X_{i}(t), R\left(c^{\prime}(t), X_{i}(t)\right) c^{\prime}(t)\right\rangle d t \\
& =\int_{0}^{a}\left|\frac{\pi}{a} \cos \left(\frac{\pi t}{a}\right) V_{i}(t)\right|^{2}-\sin ^{2}\left(\frac{\pi t}{a}\right)\left\langle V_{i}(t), R\left(c^{\prime}(t), V_{i}(t)\right) c^{\prime}(t)\right\rangle d t \\
& =\int_{0}^{a} \frac{\pi^{2}}{a^{2}} \cos ^{2}\left(\frac{\pi t}{a}\right)-\sin ^{2}\left(\frac{\pi t}{a}\right) K\left(c^{\prime}(t), V_{i}(t)\right) d t \tag{4.1}
\end{align*}
$$

where $K\left(c^{\prime}(t), V_{i}(t)\right)$ is sectional curvature of the two-dimensional plane spanned by $c^{\prime}(t)$ and $V_{i}(t)$.

Summing (4.1) over $i$ and recalling $\sum_{i=1}^{n-1} K\left(c^{\prime}(t), V_{i}(t)\right)=\operatorname{Ric}\left(c^{\prime}(t)\right) \geq \frac{n-1}{r^{2}}$, we obtain

$$
\begin{align*}
\sum_{i=1}^{n-1} E_{i}^{\prime \prime}(0) & \leq(n-1) \int_{0}^{a} \frac{\pi^{2}}{a^{2}} \cos ^{2}\left(\frac{\pi t}{a}\right)-\frac{1}{r^{2}} \sin ^{2}\left(\frac{\pi t}{a}\right) d t \\
& =(n-1)\left(\frac{\pi^{2}}{a^{2}}-\frac{1}{r^{2}}\right) \frac{a}{2} \tag{4.2}
\end{align*}
$$

Since $c$ is a minimal constant speed geodesic, $E_{i}(0)=\inf _{s \in(-\varepsilon, \varepsilon)} E_{i}(s)$ and thus $E_{i}^{\prime \prime}(0) \geq 0$ for all $i$. Equation (4.2) then implies that $a \leq \pi r$ as desired.

Theorem 4.2 (Cheng's rigidity [Che75]). A connected and complete Riemannian manifold $M$ with $\operatorname{Ric}(M) \geq K>0$ and diameter $\operatorname{diam}(M)=\sqrt{\frac{n-1}{K}}$ is isometric to an $n$-dimensional round sphere $S_{r}^{n}\left(\right.$ where $\operatorname{Ric}\left(S_{r}^{n}\right)=\frac{n-1}{r^{2}}$ and $\left.\operatorname{diam}\left(S_{r}^{n}\right)=\pi r\right)$.

Remark 4.3. It is crucial to emphasize the necessity of the assumption of the positive lower Ricci curvature bound $\operatorname{Ric}(M) \geq K>0$ in the Bonnet-Myers theorem. One cannot relax this assumption to "everywhere positive Ricci curvature". As a counterexample, we consider an elliptic paraboloid $x_{3}=x_{1}^{2}+x_{2}^{2}$. This paraboloid has everywhere positive Gaussian curvature $K=\frac{4}{\left(1+4 x_{1}^{2}+4 x_{2}^{2}\right)^{2}}$; however, the infimum of curvature is zero so the Bonnet-Myers theorem is not applicable, and the paraboloid is indeed noncompact.

In [Oht07a], Ohta provides a generalized Bonnet-Myers diameter bound for metric measure spaces, where the lower bound of Ricci curvature is replaced by the Lott-Sturm-Villani curvature-dimension condition $C D(K, N)$ (see a brief discussion of $C D(K, N)$ in Remark 3.5). In his sequential paper [Oht07b], he also provides a necessary condition for the maximal diameter bound (which can be thought of as a weak version of Cheng's rigidity). For a summary of both results, we refer to [Oht14].

Theorem 4.4 (generalized Bonnet-Myers). Suppose that a metric measure space $(X, d, m)$ satisfies $C D(K, N)$ for some $K>0$ and $N \in(1, \infty)$. Then
(i) $\operatorname{diam}(X) \leq \pi \sqrt{(N-1) / K}$.
(ii) Each $x \in X$ has at most one point of distance $\pi \sqrt{(N-1) / K}$ from $x$.

Theorem 4.5 (maximal diameter). Suppose $(X, d, m)$ satisfies $C D(K, N)$ for some $K>0$ and $N \in(1, \infty)$, and $\operatorname{diam}(X)=\pi \sqrt{(N-1) / K}$. Let $x_{N}, x_{S} \in X$
be two points such that $d\left(x_{N}, x_{S}\right)=\operatorname{diam}(X)$. Then for all $z \in X$, we have $d\left(x_{N}, z\right)+d\left(z, x_{S}\right)=\operatorname{diam}(X)$. In particular, there exists a minimal geodesic from $x_{N}$ to $x_{S}$ passing through $z$.

If $X$ with the above maximal diameter is assumed to be non-branching (meaning there is no pairs of minimal geodesic $\gamma_{1}, \gamma_{2}:[0, \ell] \rightarrow X$ such that $\gamma_{1}(0)=\gamma_{2}(0)$ and $\gamma_{1}(t)=\gamma_{2}(t)$ for some $t>0$ but $\gamma_{1}(\ell)=\gamma_{2}(\ell)$ ), then the minimal geodesic from $x_{N}$ to $x_{S}$ passing through $z$ is unique. It implies further that $X$ is the spherical suspension of some topological measure space. This maximal diameter and spherical suspension result is not discussed here in details, but we would like to mention that an analogous result will appear later in the context of Ollivier Ricci curvature on graphs.

## Chapter 5

## Analytic consequences of Ricci curvature: Bakry-Émery theory

In this chapter, we first provide background knowledge about differential operators (gradient, divergence, Laplacian and Hessian) which is necessary in later discussions including Bochner's formula, Laplacian eigenvalues and Lichnerowicz's theorem, and Heat equation and heat kernel. These mentioned topics are building blocks for the Bakry-Émery Theory [BE84] which essentially gives rise to another Ricci curvature notion, known as Bakry-Émery curvature and related to an important concept about the gradient estimate of the heat flow.

### 5.1 Differential operators

Let $\left(M^{n},\langle\cdot, \cdot\rangle\right)$ be a smooth connected complete Riemannian manifold, and let $C^{\infty}(M)$ and $\mathfrak{X}(M)$ denote the spaces of all smooth real functions and smooth vector fields on $M$, respectively.

Definition 5.1 (Gradient, divergence and Laplacian).
The gradient operator grad : $C^{\infty}(M) \rightarrow \mathfrak{X}(M)$ is defined for any $f \in C^{\infty}(M)$ uniquely via the following inner product:

$$
\langle\operatorname{grad} f, X\rangle:=X f \quad \forall X \in \mathfrak{X}(M)
$$

The divergence operator div : $\mathfrak{X}(M) \rightarrow C^{\infty}(M)$ is defined for any $X \in \mathfrak{X}(M)$ by

$$
(\operatorname{div} X)(x):=\operatorname{tr}_{T_{x} M}\left(v \mapsto \nabla_{v} X\right) \quad \forall x \in M
$$

where the mapping is from the tangent space $T_{x} M$ onto itself and $\nabla$ denotes the Levi-Civita connection.

The Laplacian operator $\Delta: C^{\infty}(M) \rightarrow C^{\infty}(M)$ is defined to be the composition $\Delta:=\operatorname{div}(\mathrm{grad})$.

The product rule for gradient, divergence and Laplacian is given in the following proposition.
Proposition 5.2 (Product rule). Let $f, h \in C^{\infty}(M)$ and $X \in \mathfrak{X}(M)$. Then
(a) $\operatorname{grad}(f h)=f \operatorname{grad} h+h \operatorname{grad} f$
(b) $\operatorname{div}(f X)=\langle\operatorname{grad} f, X\rangle+f \operatorname{div} X$
(c) $\Delta(f h)=2\langle\operatorname{grad} f, \operatorname{grad} h\rangle+f \Delta h+h \Delta f$

Proof. The product rule for gradient and divergence is induced from the product rule of directional derivative and of Levi-Civita connection, respectively. For Laplacian, we have

$$
\begin{aligned}
\Delta(f h) & =\operatorname{div}(\operatorname{grad}(f h)) \\
& =\operatorname{div}(f \operatorname{grad} h)+\operatorname{div}(h \operatorname{grad} f) \\
& =\langle\operatorname{grad} f, \operatorname{grad} h\rangle+f \operatorname{div}(\operatorname{grad} h)+\langle\operatorname{grad} h, \operatorname{grad} f\rangle+h \operatorname{div}(\operatorname{grad} f) \\
& =2\langle\operatorname{grad} f, \operatorname{grad} h\rangle+f \Delta h+h \Delta f .
\end{aligned}
$$

Divergence is often referred to as the negative adjoint of gradient as it can be seen from the divergence theorem on a closed manifold $M$ that

$$
\begin{equation*}
0=\int_{M} \operatorname{div}(f X) d \mathrm{vol}=\int_{M}\langle\operatorname{grad} f, X\rangle d \mathrm{vol}+\int_{M} f \operatorname{div} X d \mathrm{vol} \tag{5.1}
\end{equation*}
$$

Definition 5.3 (Hessian). The Hessian tensor of $f \in C^{\infty}(M)$ is a bilinear form defined as

$$
\text { Hess } f(X, Y):=\left\langle\nabla_{X} \operatorname{grad} f, Y\right\rangle
$$

for any $X, Y \in \mathfrak{X}(M)$.
Recall the Riemannian property: $X\langle Y, Z\rangle=\left\langle\nabla_{X} Y, Z\right\rangle+\left\langle Y, \nabla_{X} Z\right\rangle$. Hessian can then be reformulated as

$$
\text { Hess } f(X, Y)=\left\langle\nabla_{X} \operatorname{grad} f, Y\right\rangle=X(Y f)-\left\langle\operatorname{grad} f, \nabla_{X} Y\right\rangle
$$

A fundamental property of the Hessian is the symmetry property.

Proposition 5.4. Hess $f(X, Y)=$ Hess $f(Y, X)$
Proof. Recall the Riemannian property: $X\langle Y, Z\rangle=\left\langle\nabla_{X} Y, Z\right\rangle+\left\langle Y, \nabla_{X} Z\right\rangle$. We derive

$$
\begin{aligned}
\operatorname{Hess}(f)(X, Y) & =\left\langle\nabla_{X} \operatorname{grad} f, Y\right\rangle \\
& =X\langle\operatorname{grad} f, Y\rangle-\left\langle\operatorname{grad} f, \nabla_{X} Y\right\rangle \\
& =X(Y f)-\left\langle\operatorname{grad} f, \nabla_{X} Y\right\rangle,
\end{aligned}
$$

and similarly, we have

$$
\operatorname{Hess}(f)(Y, X)=Y(X f)-\left\langle\operatorname{grad} f, \nabla_{Y} X\right\rangle,
$$

Since Hess $f(X, Y)=X(Y f)-\left\langle\operatorname{grad} f, \nabla_{X} Y\right\rangle$ and similarly Hess $f(Y, X)=Y(X f)-$ $\left\langle\operatorname{grad} f, \nabla_{Y} X\right\rangle$, we have

$$
\text { Hess } \begin{aligned}
f(X, Y)-\text { Hess } f(Y, X) & =X(Y f)-Y(X f)-\left\langle\operatorname{grad} f, \nabla_{X} Y-\nabla_{Y} X\right\rangle \\
& =[X, Y](f)-\langle\operatorname{grad} f,[X, Y]\rangle \\
& =[X, Y](f)-[X, Y](f) \\
& =0
\end{aligned}
$$

due to the torsion-freeness of $\nabla: \nabla_{X} Y-\nabla_{Y} X=[X, Y]$, the Lie bracket of vector fields.

The Hessian tensor can also be represented by a matrix $A=\left[a_{i j}\right]$ w.r.t. an arbitrary orthonormal frame $\left\{E_{i}\right\}_{i=1}^{n}$, that is,

$$
a_{i j}=\operatorname{Hess}(f)\left(E_{i}, E_{j}\right),
$$

and the norm $\|$ Hess $f \|$ is defined as in Hilbert-Schmidt norm:

$$
\| \text { Hess } f \|:=\sqrt{\operatorname{tr}\left(A A^{\top}\right)}=\sqrt{\sum_{i, j} a_{i j}^{2}}
$$

which is independent of the choice of orthonormal frame $E_{i}$ 's.
Proposition 5.5. The following two relations hold between Hessian and Laplacian.
(a) $\operatorname{tr}($ Hess $f)=\Delta f$
(b) $\|$ Hess $f \|^{2} \geq \frac{1}{n}(\Delta f)^{2}$ where $n$ is the dimension of $M$.

Proof. Part (a) follows directly from definitions:

$$
\operatorname{tr}(\operatorname{Hess} f)=\sum_{i=1}^{n} a_{i i}=\sum_{i=1}^{n}\left\langle\nabla_{E_{i}} \operatorname{grad} f, E_{i}\right\rangle=\operatorname{div}(\operatorname{grad} f)=\Delta f,
$$

and the part (b) follows by applying Cauchy-Schwarz's inequality:

$$
\| \text { Hess } f \|^{2}=\sum_{i, j} a_{i j}^{2} \geq \sum_{i=1}^{n} a_{i i}^{2} \geq \frac{1}{n}\left(\sum_{i=1}^{n} a_{i i}\right)^{2}=\frac{1}{n}(\Delta f)^{2} .
$$

### 5.2 Bochner's formula

Bochner's formula is an identity which connects differential operators to Ricci curvature. This formula is a fundamental motivation for the Bakry-Émery curvature notion. In this section, we give the statement of Bochner's formula and its proof, which can also be found in standard Riemannian geometry books, e.g. [GHL90, Proposition 4.15] and [Jos08, Theorem 4.5.1].

Theorem 5.6 (Bochner's formula). Let $\left(M^{n}, g\right)$ be a complete Riemannian manifold. Then for any function $f \in C^{\infty}(M)$, the equation

$$
\begin{equation*}
\frac{1}{2} \Delta|\operatorname{grad} f|^{2}=\|\operatorname{Hess} f\|^{2}+\langle\operatorname{grad} \Delta f, \operatorname{grad} f\rangle+\operatorname{Ric}(\operatorname{grad} f) \tag{5.2}
\end{equation*}
$$

holds pointwise on $M$.
Here we use the notation $\operatorname{Ric}(v)=\operatorname{Ric}(v, v)$.
Proof. Let $x \in M$ and $\left\{e_{i}\right\}_{i=1}^{n}$ an orthonormal basis of $T_{x} M$ consider a local orthonormal frame $\left\{E_{i}\right\}_{i=1}^{n}$ such that the basis $\left\{e_{i}\right\}_{i=1}^{n}$ of $T_{x} M$ defines a Riemannian normal coordinate system at $x$. More precisely, for neighborhood $U \subset M$ containing $x$ and $V \subset \mathbb{R}^{n}$ containing 0 , the coordinate map $\phi: U \rightarrow V$ has its inverse

$$
\phi^{-1}\left(x_{1}, \ldots, x_{n}\right)=\exp _{x}\left(\sum_{i} x_{i} e_{i}\right)
$$

which maps straight lines through origin to geodesics through $x$. For $y \in U,\left.\frac{\partial}{\partial x_{i}}\right|_{y}$ is an element of $T_{y} M$, and $\left.\frac{\partial}{\partial x_{i}}\right|_{x}=e_{i}$. In particular, when evaluating at $x$, we have $\left\langle e_{i}, e_{j}\right\rangle=\delta_{i j}$ and $\nabla_{e_{i}} E_{j}=0$ for all $i$ and $j$ (and hence $\nabla_{u} E_{j}=0$ for all $u \in T_{x} M$ by linearity).

By definition of Hessian and its relation to Laplacian, we have the following evaluation,

$$
\begin{align*}
\frac{1}{2} \Delta|\operatorname{grad} f|^{2} & =\frac{1}{2} \operatorname{tr}\left(\operatorname{Hess}\|\operatorname{grad} f\|^{2}\right)=\frac{1}{2} \sum_{i=1}^{n} \operatorname{Hess}\|\operatorname{grad} f\|^{2}\left(E_{i}, E_{i}\right) \\
& =\frac{1}{2} \sum_{i=1}^{n} E_{i}\left(E_{i}\|\operatorname{grad} f\|^{2}\right)-\left\langle\operatorname{grad}\|\operatorname{grad} f\|^{2}, \nabla_{E_{i}} E_{i}\right\rangle \tag{5.3}
\end{align*}
$$

where the last inner product vanishes at $x$ since $\nabla_{e_{i}} E_{i}=0$. Moreover, by the definition and symmetry of Hessian, we have

$$
\begin{align*}
\frac{1}{2} E_{i}\left(\|\operatorname{grad} f\|^{2}\right) & =\left\langle\nabla_{E_{i}} \operatorname{grad} f, \operatorname{grad} f\right\rangle=\operatorname{Hess} f\left(E_{i}, \operatorname{grad} f\right) \\
& =\operatorname{Hess} f\left(\operatorname{grad} f, E_{i}\right)=\left\langle\nabla_{\operatorname{grad} f} \operatorname{grad} f, E_{i}\right\rangle \tag{5.4}
\end{align*}
$$

We then substitute (5.4) into (5.3) and derive it at $x$ in terms of Riemann curvature tensor:

$$
\begin{align*}
\frac{1}{2} \Delta|\operatorname{grad} f|^{2}(x)= & \sum_{i=1}^{n} e_{i}\left\langle\nabla_{\operatorname{grad} f} \operatorname{grad} f, E_{i}\right\rangle \\
= & \sum_{i=1}^{n}\left\langle\nabla_{e_{i}} \nabla_{\operatorname{grad} f} \operatorname{grad} f, e_{i}\right\rangle+\langle\nabla_{\operatorname{grad} f} \operatorname{grad} f(x), \underbrace{\nabla_{e_{i}} E_{i}}_{=0}\rangle \\
= & \sum_{i=1}^{n}\left\langle R\left(e_{i}, \operatorname{grad} f(x)\right) \operatorname{grad} f(x), e_{i}\right\rangle+  \tag{5.5}\\
& \sum_{i=1}^{n}\left\langle\nabla_{\operatorname{grad} f(x)} \nabla_{E_{i}} \operatorname{grad} f, e_{i}\right\rangle+  \tag{5.6}\\
& \sum_{i=1}^{n}\left\langle\nabla_{\left[E_{i}, \operatorname{grad} f\right](x)} \operatorname{grad} f, e_{i}\right\rangle \tag{5.7}
\end{align*}
$$

The term in (5.5) is $\operatorname{Ric}(\operatorname{grad} f(x))$ by definition of Ricci curvature. We will finish the proof by showing that the terms in (5.6) and (5.7) are equal to $\langle\operatorname{grad} \Delta f(x), \operatorname{grad} f(x)\rangle$ and $\|$ Hess $f \|^{2}(x)$, respectively. The terms in (5.6) can be computed as

$$
\begin{aligned}
\sum_{i=1}^{n}\left\langle\nabla_{\operatorname{grad} f(x)} \nabla_{E_{i}} \operatorname{grad} f, e_{i}\right\rangle= & \sum_{i=1}^{n}\left(\operatorname{grad} f(x)\left\langle\nabla_{E_{i}} \operatorname{grad} f, E_{i}\right\rangle\right. \\
& -\langle\nabla_{e_{i}} \operatorname{grad} f, \underbrace{\nabla_{\operatorname{grad} f(x)} E_{i}}_{=0}\rangle) \\
& =\operatorname{grad} f(x)(\operatorname{tr} \operatorname{Hess} f)
\end{aligned}
$$

$$
=\operatorname{grad} f(x)(\Delta f)=\langle\operatorname{grad} \Delta f(x), \operatorname{grad} f(x)\rangle .
$$

Due to torsion-freeness $\left[E_{i}, \operatorname{grad} f\right](x)=\nabla_{E_{i}} \operatorname{grad} f(x)-\underbrace{\nabla_{\operatorname{grad} f(x)} E_{i}}_{=0}$, we can compute the term in line (5.7) as

$$
\begin{aligned}
\sum_{i=1}^{n}\left\langle\nabla_{\left[E_{i}, \operatorname{grad} f\right](x)} \operatorname{grad} f, e_{i}\right\rangle & =\sum_{i=1}^{n} \operatorname{Hess} f\left(\left[E_{i}, \operatorname{grad} f\right](x), e_{i}\right) \\
& =\sum_{i=1}^{n} \operatorname{Hess} f\left(e_{i}, \nabla_{e_{i}} \operatorname{grad} f\right) \\
& =\sum_{i=1}^{n}\left\langle\nabla_{e_{i}} \operatorname{grad} f, \nabla_{e_{i}} \operatorname{grad} f\right\rangle=\| \text { Hess } f \|^{2}(x) .
\end{aligned}
$$

### 5.3 Eigenvalues of Laplacian and Lichnerowicz spectral gap theorem

In this section, we give a brief overview of the Laplacian eigenvalues of manifolds and the Lichnerowicz theorem which relates the smallest eigenvalue to the lower bound of the Ricci curvature. For more details, we refer to Chavel's book [Cha84].
Suppose for simplicity that $\left(M^{n}, g\right)$ is a closed and connected Riemannian manifold. The eigenvalues of the Laplacian operator on $M$ are all the real numbers $\lambda$ such that there exists a nontrivial solution $f \in C^{2}(M)$ to the following eigenvalue problem

$$
\Delta f+\lambda f=0
$$

For each eigenvalue $\lambda$, a corresponding solution $f$ is called an eigenfunction. The eigenspace of $\lambda$ is the vector space of all eigenfunctions with respect to $\lambda$.

It is well-known that the set of these eigenvalues (with their multiplicities) consists of a sequence

$$
0=\lambda_{0}<\lambda_{1} \leq \lambda_{2} \leq \ldots, \text { where } \lambda_{n} \rightarrow \infty \text { as } n \rightarrow \infty,
$$

and each eigenspace has a finite dimension (which is the multiplicity of the corresponding $\lambda$ ). The trivial eigenvalue $\lambda_{0}$ corresponds to the constant function $f$. The smallest nonzero eigenvalue $\lambda_{1}$ is known as the spectral gap of the Laplacian.

Remark 5.7. Here we list some facts in case $M$ is not closed and connected.

1. If $M$ is not connected, the multiplicity of the trivial eigenvalue equals the number of connected components of $M$.
2. If $M$ is compact but not closed (i.e., $\partial M \neq \emptyset$ ), one may consider different eigenvalue problems. For example, the Neumann problem comes with the boundary condition $\hat{\mathrm{n}}(f)=0$ on $\partial M$ (where $\hat{\mathrm{n}}$ is the unit normal vector). On the other hand, the Dirichlet problem requires $f=0$ on $\partial M$, which means constant functions are not eigenfunctions.
3. If $M$ is non-compact, then the Laplacian may not have discrete spectrum. For example, the spectrum of the Euclidean space $\mathbb{R}^{n}$ is $[0, \infty)$, and the spectrum of the hyperbolic space $\mathbb{H}^{n}$ is $\left[\frac{(n-1)^{2}}{4}, \infty\right)$.

The Lichnerowicz spectral gap theorem [Lic58] asserts that the smallest nonzero eigenvalue $\lambda_{1}$ can be estimated from below in terms of the positive lower bound of Ricci curvature. The proof of the Lichnerowicz theorem is a straightforward implication of the Bochner's formula. In a general case where $M$ has nonempty boundary, we refer to a proof in, e.g. [GHL90, Theorem 4.70].

Theorem 5.8 (Lichnerowicz). Let $\left(M^{n}, g\right)$ be a closed Riemannian manifold. Suppose that $K:=\operatorname{Ric}(M)>0$. Then the smallest nonzero Laplacian eigenvalue satisfies

$$
\lambda_{1} \geq \frac{n}{n-1} K
$$

Proof. Consider an eigenfunction $f$ satisfying $\Delta f+\lambda f=0$. Without loss of generality, we assume that $\int_{M} f^{2}=1$. Bochner's formula (5.2) gives

$$
\begin{aligned}
\frac{1}{2} \Delta|\operatorname{grad} f|^{2} & =\|\operatorname{Hess} f\|^{2}+\langle\operatorname{grad}(-\lambda f), \operatorname{grad} f\rangle+\operatorname{Ric}(\operatorname{grad} f) \\
& =\|\operatorname{Hess} f\|^{2}-\lambda|\operatorname{grad} f|^{2}+\operatorname{Ric}(\operatorname{grad} f) \\
& \geq \frac{1}{n}(\Delta f)^{2}-\lambda|\operatorname{grad} f|^{2}+K|\operatorname{grad} f|^{2} \\
& =\frac{\lambda^{2}}{n} f^{2}+(K-\lambda)|\operatorname{grad} f|^{2} \\
& =\frac{\lambda^{2}}{n} f^{2}+(K-\lambda)\left(\frac{1}{2} \Delta f^{2}+\lambda f^{2}\right)
\end{aligned}
$$

where the inequality is due to $\|$ Hess $f \|^{2} \geq \frac{1}{n}(\Delta f)^{2}$ and the lower curvature bound assumption, and the last equality is the product rule : $\Delta f^{2}=2|\operatorname{grad} f|^{2}+2 f \Delta f$.

Integrating the above inequality over $M$ and using the fact from the divergence theorem that $\int_{M} \Delta|\operatorname{grad} f|^{2}=0=\int_{M} \Delta f^{2}$ by, we derive

$$
0 \geq \frac{\lambda^{2}}{n}+(K-\lambda) \lambda
$$

In particular, for $\lambda_{1}>0$, we conclude $\lambda_{1} \geq \frac{n}{n-1} K$ as desired.
Furthermore, like in Cheng's rigidity to the Bonnet-Myers diamter bound, there is also a rigidity theorem for the Lichnerowicz spectral bound, known as Obata's rigidity [Oba62] which confirms that the bound is sharp if and only if the manifold is isometric to a round sphere.

Theorem 5.9 (Obata's rigidity, [Oba62]). Let $\left(M^{n},\langle\cdot, \cdot\rangle\right)$ be a compact Riemannian manifold with positive lower Ricci curvature bound $K:=\operatorname{Ric}(M)>0$. Then $\lambda_{1}(M)=\frac{n}{n-1} K$ if and only if $M$ is isometric to the sphere $S_{r}^{n}$.

### 5.4 Curvature-dimension inequality and $\Gamma$-calculus

Due to the Bochner's formula and $\|\operatorname{Hess} f\|^{2} \geq \frac{1}{n}(\Delta f)^{2}$, the pointwise lower Ricci curvature bound $\operatorname{Ric}_{x} \geq K$ immediately implies the following curvaturedimension inequality in the sense of Bakry-Émery [BE84]:

$$
\begin{equation*}
\frac{1}{2} \Delta|\operatorname{grad} f|^{2}(x)-\langle\operatorname{grad} \Delta f(x), \operatorname{grad} f(x)\rangle \geq \frac{1}{n}(\Delta f(x))^{2}+K|\operatorname{grad} f(x)|^{2} \tag{5.8}
\end{equation*}
$$

Bakry and Émery also introduce the following $\Gamma$-calculus, consisting of the first iteration $\Gamma$ and the second iteration $\Gamma_{2}$, which helps to reformulate the curvaturedimension inequality.

Definition $5.10\left(\Gamma\right.$ and $\left.\Gamma_{2}\right)$. For $f, g \in C^{\infty}(M)$, define bilinear operators

$$
\begin{align*}
2 \Gamma(f, g) & :=\Delta(f g)-f \cdot \Delta g-\Delta f \cdot g  \tag{5.9}\\
2 \Gamma_{2}(f, g) & :=\Delta \Gamma(f, g)-\Gamma(f, \Delta g)-\Gamma(\Delta f, g) \tag{5.10}
\end{align*}
$$

and for shortened notation, we write $\Gamma f:=\Gamma(f, f)$ and $\Gamma_{2} f:=\Gamma_{2}(f, f)$.
By the product rule, one can realize that

$$
\begin{equation*}
\Gamma(f, g)=\langle\operatorname{grad} f, \operatorname{grad} g\rangle, \text { and } \Gamma f=|\operatorname{grad} f|^{2} . \tag{5.11}
\end{equation*}
$$

The inequality (5.8) can then be rewritten as $\Gamma_{2} f(x) \geq \frac{1}{n}(\Delta f(x))^{2}+K \Gamma f(x)$, which motivates the following definition of Bakry-Émery curvature notion.

Definition 5.11. Let $\left(M^{n}, g\right)$ be a Riemannian manifold and let $K \in \mathbb{R}$ and $N \in[0, \infty]$. A point $x \in M$ is said to satisfy $B E(K, N)$ if the following inequality holds true:

$$
\begin{equation*}
\Gamma_{2} f(x) \geq \frac{1}{N}(\Delta f(x))^{2}+K \Gamma f(x) \quad \forall f \in C^{\infty}(M) \tag{5.12}
\end{equation*}
$$

Here $K$ is an lower bound of the Ricci curvature and $N$ is an upper bound for the dimension of $M$. In the case that $N=\infty$, the $B E(K, \infty)$ inequality at $x$ is read as

$$
\begin{equation*}
\Gamma_{2} f(x) \geq K \Gamma f(x) \tag{5.13}
\end{equation*}
$$

In fact, the following result follows immediately from Bochner's formula.
Theorem 5.12. Let $\left(M^{n}, g\right)$ be a Riemannian manifold and $x \in M$ such that $\operatorname{Ric}_{x}(v) \geq K|v|^{2}$ for all $v \in T_{x} M$. Then $x$ satisfies $B E(K, N)$ for all $N \geq n$.

### 5.5 Heat kernel and heat semigroup operator on manifolds

In this section, we briefly overview the heat kernel $p_{t}$ and the heat semigroup operator $P_{t}$. They are main ingredients to another characterization of $B E(K, n)$ in terms of gradient estimate, which is the topic of the next section. Basic and useful properties of $p_{t}$ and $P_{t}$ will be provided here without proofs, and we refer to Grigoryan's book [Gri09] for details.
We consider the following Cauchy problem for the heat equation on a complete Riemannian manifold. Given a bounded continuous function $f \in C_{b}(M)$, the problem is to solve $u:[0, \infty) \times M \rightarrow \mathbb{R}$ (which is continuous, $C^{1}$ in $t \in(0, \infty)$ and $C^{2}$ in $x \in M$ ) such that

$$
\begin{cases}\frac{\partial}{\partial t} u & =\Delta u \quad \text { on }(0, \infty) \times M  \tag{5.14}\\ u(0, x) & =f(x)\end{cases}
$$

Definition 5.13 (Fundamental solution). A function $p:(0, \infty) \times M \times M \rightarrow \mathbb{R}$ is called a fundamental solution of the heat equation, if for every $f \in C_{b}(M)$, the function

$$
u(t, x)= \begin{cases}\int_{M} p(t, x, y) f(y) d \operatorname{vol}(y) & \text { for } t>0  \tag{5.15}\\ f(x) & \text { for } t=0\end{cases}
$$

is a solution to Cauchy problem (5.14) with initial data $f$.

The existence of solutions $u$ is proved via the following Dodziuk's construction [Dod83] (see also Chavel's book [Cha84]). First, we construct an exhaustion of $M$ by regular domains $\left\{D_{\ell}\right\}_{\ell=1}^{\infty}$ (relatively compact and open subsets of $M$ with smooth boundaries) so that $\bar{D}_{\ell} \subset D_{\ell+1}$ and $\bigcup_{\ell} D_{\ell}=M$. Then for each $\ell$, consider $p_{\ell}:(0, \infty) \times D_{\ell} \times D_{\ell} \rightarrow \mathbb{R}^{+}$to be the Dirichlet heat kernel, that is, $p_{\ell}$ is the unique fundamental solution of the following Dirichlet problem:

$$
\left\{\begin{array}{lll}
\frac{\partial}{\partial t} u & =\Delta u & \text { on }(0, \infty) \times D_{\ell}  \tag{5.16}\\
u(t, x) & =0 & \text { on }(0, \infty) \times \partial D_{\ell} \\
u(0, x) & =f(x) &
\end{array}\right.
$$

Two important properties of $p_{\ell}$ are that for all $t>0, \int_{D_{\ell}} p_{\ell}(t, x, y) d \operatorname{vol}(y) \leq 1$ for all $x \in D_{\ell}$ and $p_{\ell}(t, x, y) \leq p_{\ell+1}(t, x, y)$ for all $x, y \in D_{\ell}$ (which is a consequence of parabolic maximum principle); see [Dod83, Lemma 3.3]. It follows that the limit

$$
p(t, x, y):=\lim _{\ell \rightarrow \infty} p_{\ell}(t, x, y)
$$

exists, and the convergence is locally uniform due to Dini's theorem. This $p$ is indeed the smallest positive fundamental solution of the heat equation on $M$. We call this $p$ the heat kernel of $M$. Furthermore, the heat semigroup operator is defined via the kernel $p$ as follows.

Definition 5.14 (heat semigroup operator). The heat semigroup $\left\{P_{t}\right\}_{t \geq 0}$ is the family of operators $P_{t}: C_{b}(M) \rightarrow C_{b}(M)$ defined as $P_{0}=\mathrm{Id}$ and for $t>0$,

$$
P_{t} f(x):=\int_{M} p(t, x, y) f(y) d \operatorname{vol}(y),
$$

for all $f \in C_{b}(M)$ and $x \in M$.

It means the function $u(t, x):=P_{t} f(x)$ solves the Cauchy problem (5.14), or in short,

$$
\frac{\partial}{\partial t} P_{t}=\Delta P_{t} .
$$

Moreover, $P_{t} f(x)$ is $C^{\infty}$ in $(0, \infty) \times M$ with the limit $\lim _{t \rightarrow 0} P_{t} f(x)=f(x)$ locally uniformly in $x \in M$. It also satisfies

$$
\inf f \leq P_{t} f(x) \leq \sup f
$$

(see [Gri09, Theorem 7.16]). Basic properties of the heat kernel $p_{t}$ and the heat semigroup are summarized in the following two propositions.

Proposition 5.15. The heat kernel $p:(0, \infty) \times M \times M \rightarrow \mathbb{R}$ satisfies the following properties.

1. symmetry: $p(t, x, y)=p(t, y, x)$.
2. semigroup identity: $p(s+t, x, y)=\int_{M} p(s, x, z) p(t, z, y) d \operatorname{vol}(z)$.
3. density-like: $p(t, x, y)>0$ and

$$
\int_{M} p(t, x, y) d \operatorname{vol}(y) \leq 1
$$

Proposition 5.16. The heat semigroup operator $P_{t}: C_{b}(M) \rightarrow C_{b}(M)$ satisfies the following properties.

1. $P_{t} \circ P_{s}=P_{s+t}$.
2. $\Delta P_{t}=P_{t} \Delta$.
3. $f(x) \geq 0$ for all $x$ implies $P_{t} f(x) \geq 0$ for all $x$.
4. $\left\|P_{t}(f)\right\|_{\infty} \leq\|f\|_{\infty}$.

The uniqueness of a bounded solution to the Cauchy problem (5.14) requires a certain condition on the manifold $M$ called stochastic completeness, which roughly says that the heat is preserved within the system.

Definition 5.17 (Stochastic completeness). A Riemannian manifold $(M, g)$ is said to be stochastically complete if the heat kernel $p$ satisfies

$$
\int_{M} p(t, x, y) d \operatorname{vol}(y)=1
$$

for all $t>0$ and $x \in M$.
Theorem 5.18 ([Gri09, Theorem 8.18]). A Riemannian manifold $(M, g)$ is stochastically complete if and only if the Cauchy problem (5.14) has exactly one bounded solution $u$.

A criterion for the stochastic completeness by the volume test is provided in [Gri09, Theorem 11.8].

Theorem 5.19 ([Gri09]). Let ( $M, g$ ) be a complete connected Riemannian manifold. If for some $x_{0} \in M$,

$$
\begin{equation*}
\int_{0}^{\infty} \frac{r}{\log \operatorname{vol}\left(B_{r}\left(x_{0}\right)\right)} d r=\infty \tag{5.17}
\end{equation*}
$$

then $M$ is stochastically complete. Here $B_{r}\left(x_{0}\right)$ denotes the geodesic ball of radius $r$ centered at $x_{0}$.

The condition (5.17) holds in particular if $\operatorname{vol}\left(B_{r}(x)\right) \leq \exp \left(C r^{2}\right)$ for all $r$ large enough. For example, the Euclidean space $\mathbb{R}^{n}$ with vanishing curvature and the hyperbolic space $\mathbb{H}^{n}$ with constant negative curvature equal to -1 are stochastically complete. More generally, it is proved in [Yau76] that any complete manifold with Ricci curvature bounded from below is stochastically complete.
Remark 5.20. In [Gri09], the author considers the Cauchy problem (5.14) first for $L^{2}(M)$ before the case $C_{b}(M)$ is covered. In the case of a (geodesically) complete Riemannian manifold $(M, g)$, the Laplacian on functions with compact support is known to be essentially self-adjoint (see e.g., [Gaf54, Che73, Kar84]) and Spectral Theory can be employed to construct the self-adjoint heat semigroup operators $P_{t}: L^{2}(M) \rightarrow L^{2}(M)$, namely $P_{t}=e^{t \Delta}$. In the $L^{2}$-Cauchy problem, the solution is a unique: $u=P_{t} f$ with these uniquely determined operators $P_{t}: L^{2}(M) \rightarrow L^{2}(M)$, which are bounded, nonnegative, commute with $\Delta$ and map $L^{2}(M)$ to $C^{\infty}(M)$ for $t>0$. They are integral operators having smooth kernels which agree with the minimal heat kernel $p(t, x, y)$ from above.

### 5.6 Gradient estimate

An important result from Bakry and Émery [BE84] is the characterization of $B E(K, \infty)$ in terms of gradient estimate, which we present here in the following two variations.

Definition 5.21 (gradient estimates). A point $x \in M$ is said to satisfy the $\boldsymbol{g r a}$ dient estimate (or the $\Gamma$-gradient estimate, respectively) if for all $f \in C_{c}^{\infty}(M)$ and $t \in[0, \infty)$,

$$
\begin{equation*}
\left|\operatorname{grad} P_{t} f\right|(x) \leq \mathrm{e}^{-K t} P_{t}|\operatorname{grad} f|(x), \tag{GE}
\end{equation*}
$$

or

$$
\begin{equation*}
\Gamma\left(P_{t} f\right)(x) \leq \mathrm{e}^{-2 K t} P_{t}(\Gamma f)(x) \tag{GE-Г}
\end{equation*}
$$

respectively. Moreover, we say that $M$ satisfies gradient estimate (or $\Gamma$-gradient estimate) if such condition holds for all $x \in M$.

The condition (GE- $\Gamma$ ) can be considered as a weaker version of (GE) since

$$
\left(P_{t}|\operatorname{grad} f|\right)^{2} \leq P_{t}\left(|\operatorname{grad} f|^{2}\right)=P_{t}(\Gamma f)
$$

In fact, it is among the most famous results in the original work of Bakry and Émery that these gradient estimates are equivalent to the curvature-dimension condition $B E(K, \infty)$ on manifolds.

Theorem 5.22. Let $(M, g)$ be a complete Riemannian manifold. Let $x \in M$ and $K \in \mathbb{R}$. Then the following properties are equivalent.

1. $B E(K, \infty)$ at $x: \Gamma_{2} f(x) \geq K \Gamma f(x)$ for all $f$.
2. $\Gamma$-gradient estimate at $x: \Gamma\left(P_{t} f\right)(x) \leq \mathrm{e}^{-2 K t} P_{t}(\Gamma f)(x)$ for all $f$.
3. gradient estimate at $x: \sqrt{\Gamma\left(P_{t} f\right)}(x) \leq \mathrm{e}^{-K t} P_{t}(\sqrt{\Gamma f})(x)$ for all $f$.

The proof of this theorem, provided below, relies on the well-known semigroup interpolation argument. The equivalence (1) $\Leftrightarrow(2)$ will make an appearance again when we discuss the Bakry-Émery curvature on graphs and the Erbar-Maas entropic Ricci curvature. Although the equivalence (1) $\Leftrightarrow$ (3) looks similar, it indeed requires an additional diffusion property of $\Delta$ on manifolds (referred to [Bak97, pp. 4]) as we shall see in the proof.

Proof. (1) $\Rightarrow$ (2): Fix $t \in[0, \infty)$ and $f \in C_{c}^{\infty}(M)$. Define a function $F:[0, t] \rightarrow \mathbb{R}$ as

$$
F(s):=e^{-2 K s} P_{s}\left(\Gamma\left(P_{t-s} f\right)\right)(x)
$$

Every function here is evaluated at $x$ and we omit it from the writing. Moreover, let $g=P_{t-s} f$. The derivative of $F$ is computed as:

$$
\begin{aligned}
F^{\prime}(s) & =e^{-2 K s}\left(-2 K P_{s}(\Gamma g)+\left(\frac{\partial}{\partial s} P_{s}\right)(\Gamma g)+P_{s}\left(\frac{\partial}{\partial s} \Gamma g\right)\right) \\
& =e^{-2 K s} P_{s}\left(-2 K \Gamma P_{t-s} f+\Delta(\Gamma g)+\frac{\partial}{\partial s} \Gamma g\right),
\end{aligned}
$$

where in the second line we recall that $\frac{\partial}{\partial s} P_{s}=\Delta P_{s}=P_{s} \Delta$. Next, the term $\frac{\partial}{\partial s} \Gamma g$ is equal to

$$
\begin{align*}
\frac{\partial}{\partial s} \Gamma\left(P_{t-s} f, P_{t-s} f\right)=2 \Gamma\left(\frac{\partial}{\partial s} P_{t-s} f, P_{t-s} f\right) & =-2 \Gamma\left(\Delta P_{t-s} f, P_{t-s} f\right) \\
& =2 \Gamma_{2} g-\Delta(\Gamma g) \tag{5.18}
\end{align*}
$$

Substituting (5.18) into the calculation of $F^{\prime}(s)$ yields

$$
\begin{equation*}
F^{\prime}(s)=2 e^{-2 K s} P_{s}\left(\Gamma_{2} g-K \Gamma g\right) \tag{5.19}
\end{equation*}
$$

Under the assumption $B E(K, \infty)$ at $x$, we have $F^{\prime}(s) \geq 0$ for all $s \in[0, t]$. Thus $F(0) \leq F(t)$ which yields (GE-Г).
$(2) \Leftarrow(1)$ : Suppose $x$ satisfies (GE- $\Gamma$ ). Differentiating (GE- $\Gamma$ ) near $t=0$ gives

$$
\begin{aligned}
0 & \leq \lim _{t \rightarrow 0} \frac{1}{t}\left(e^{-2 K t} P_{t}(\Gamma f)-\Gamma\left(P_{t} f\right)\right)=\left.\frac{\partial}{\partial t}\right|_{t=0}\left(e^{-2 K t} P_{t}(\Gamma f)-\Gamma\left(P_{t} f\right)\right) \\
& =\Delta(\Gamma f)-2 K \Gamma f-\left.\frac{\partial}{\partial t}\right|_{t=0} \Gamma\left(P_{t} f, P_{t} f\right) \\
& \left.=\Delta(\Gamma f)-2 K \Gamma f+\left[2 \Gamma_{2}\left(P_{t} f\right)-\Delta\left(\Gamma\left(P_{t} f\right)\right)\right]\right]_{t=0}=2 \Gamma_{2} f-2 K \Gamma f .
\end{aligned}
$$

Thus $x$ satisfies $B E(K, \infty)$ as desired.
Similarly, the proof of $(1) \Rightarrow(3)$ starts by considering the function

$$
\tilde{F}(s):=e^{-K s} P_{s}\left(\sqrt{\Gamma P_{t-s} f}\right)(x) .
$$

The derivative $\tilde{F}(s)$ (with $g:=P_{t-s}$ ) gives

$$
\begin{aligned}
& \tilde{F}^{\prime}(s)=e^{-K s}\left(-K P_{s}(\sqrt{\Gamma g})+\left(\frac{\partial}{\partial s} P_{s}\right)(\sqrt{\Gamma g})+P_{s}\left(\frac{\partial}{\partial s} \sqrt{\Gamma g}\right)\right) \\
&=e^{-K s} P_{s}\left(-K \sqrt{\Gamma g}+\Delta \sqrt{\Gamma g}+\frac{1}{2 \sqrt{\Gamma g}} \cdot \frac{\partial}{\partial s} \Gamma g\right) \\
& \stackrel{(5.18)}{=} e^{-K s} P_{s}\left[\frac{1}{2 \sqrt{\Gamma g}}\left(2 \Gamma_{2} g-2 K \Gamma g+2 \sqrt{\Gamma g} \Delta \sqrt{\Gamma g}-\Delta(\Gamma g)\right)\right] \\
&=e^{-K s} P_{s}\left[\frac{1}{\sqrt{\Gamma g}}\left(\Gamma_{2} g-K \Gamma g-\Gamma(\sqrt{\Gamma g})\right)\right] .
\end{aligned}
$$

It was shown in [Bak97, Lemmas 2.4 and 2.5] that $B E(K, \infty)$ together with the diffusion property gives

$$
\Gamma_{2} g-K \Gamma g-\Gamma(\sqrt{\Gamma g}) \geq 0
$$

(This inequality is a self-improvement of $\Gamma_{2}-K \Gamma \geq 0$, and it is proved by employing the coordinate-system to formulate $\Delta, \Gamma, \Gamma_{2}$ and Hessian, which we do not cover here). Thus $\tilde{F}^{\prime}(s) \geq 0$, yielding (GE). Lastly, $(3) \Rightarrow(1)$ can be viewed as $(\mathrm{GE}) \Rightarrow(\mathrm{GE}-\Gamma) \Rightarrow B E(K, \infty)$.

Finally, we conclude the equivalent characterizations of the global lower Ricci curvature bound in spirit of Bakry-Émery.

Theorem 5.23. Let $\left(M^{n}, g\right)$ be a complete Riemannian manifold of dimension $n$ and let $K \in \mathbb{R}$. Then the following conditions on $M$ are equivalent:
(i) $\operatorname{Ric}(M) \geq K$,
(ii) $B E(K, n)$,
(iii) $B E(K, \infty)$,
(iv) gradient estimate: $\sqrt{\Gamma\left(P_{t} f\right)} \leq \mathrm{e}^{-K t} P_{t}(\sqrt{\Gamma f})$ for all $f$.
(v) $\Gamma$-gradient estimate: $\Gamma\left(P_{t} f\right) \leq \mathrm{e}^{-2 K t} P_{t}(\Gamma f)$ for all $f$.

Proof. The implication (i) $\Rightarrow$ (ii) $\Rightarrow$ (iii) follows from the Bochner's formula, the curvature-dimension inequality, and the work of $\Gamma$-calculus. The implication (iii) $\Rightarrow$ (iv) $\Rightarrow$ (v) follows from Theorem 5.22. The implication (v) $\Rightarrow$ (i) is due to [vRS05, Theorem 1.3].

## Chapter 6

## Bakry-Émery manifolds

A Bakry-Émery manifold, also known as weighted manifold, is a smooth Riemannian manifold $(M,\langle\rangle$,$) together with a reference measure d \nu(x)=e^{-V(x)} d \operatorname{vol}(x)$ for a fixed function $V \in C^{2}(M)$.

The divergence operator is given by $\operatorname{div}_{V}(X):=\operatorname{div} X-\langle\operatorname{grad} V, X\rangle$ and it is the negative adjoint of gradient (with respect to $\nu$ ), that is,

$$
\begin{aligned}
\int_{M}\langle\operatorname{grad} f, X\rangle d \nu & =\int_{M}\left\langle\operatorname{grad} f, e^{-V} X\right\rangle d \operatorname{vol}=-\int_{M} f \operatorname{div}\left(e^{-V} X\right) d \mathrm{vol} \\
& =-\int_{M} f\left(e^{-V} \operatorname{div} X+\left\langle\operatorname{grad} e^{-V}, X\right\rangle\right) d \mathrm{vol} \\
& =-\int_{M} f\left(e^{-V} \operatorname{div} X-e^{-V}\langle\operatorname{grad} V, X\rangle\right) d \operatorname{vol} \\
& =-\int_{M} f(\underbrace{\operatorname{div} X-\langle\operatorname{grad} V, X\rangle}_{=: \operatorname{div}_{V}(X)}) d \nu .
\end{aligned}
$$

The Laplacian $\Delta_{V}$ is then given by

$$
\Delta_{V} f:=\operatorname{div}_{V}(\operatorname{grad} f)=\Delta f-\langle\operatorname{grad} V, \operatorname{grad} f\rangle,
$$

and the generalized Ricci tensor (with $\infty$ dimension) is defined to be

$$
\operatorname{Ric}_{V, \infty}:=\operatorname{Ric}+\operatorname{Hess} V
$$

We provide two justifications of this generalized Ricci tensor. The first one is the weighted Bochner's formula (c.f. [OS14, Formula (1.2)] and [Vil09, Chapter 14]). The second one is the generalized result of the main theorem (Theorem 3.3) on weighted manifolds as presented in [LV09] and [Stu06].

Theorem 6.1 (weighted Bochner's formula). Let $\left(M^{n}, g\right)$ be a complete Riemannian manifold. Then

$$
\begin{equation*}
\frac{1}{2} \Delta_{V}|\operatorname{grad} f|^{2}=\|\operatorname{Hess} f\|^{2}+\left\langle\operatorname{grad} \Delta_{V} f, \operatorname{grad} f\right\rangle+\operatorname{Ric}_{V, \infty}(\operatorname{grad} f) \tag{6.1}
\end{equation*}
$$

holds pointwise for any function $f \in C^{\infty}(M)$.
Proof. Subtracting (6.1) from the original Bochner's formula (5.2), it is left to show that
$\left.\left.\frac{1}{2}\langle\operatorname{grad} V, \operatorname{grad}| \operatorname{grad} f\right|^{2}\right\rangle=\langle\operatorname{grad}\langle\operatorname{grad} V, \operatorname{grad} f\rangle, \operatorname{grad} f\rangle-\operatorname{Hess} V(\operatorname{grad} f, \operatorname{grad} f)$.
By the definition of the gradient and the Hessian and the Riemannian property, we have

$$
\begin{aligned}
\left.\frac{1}{2}\langle\operatorname{grad} V, \operatorname{grad}| \operatorname{grad} f\right|^{2} & =\frac{1}{2} \operatorname{grad} V\langle\operatorname{grad} f, \operatorname{grad} f\rangle \\
& =\left\langle\nabla_{\operatorname{grad} V} \operatorname{grad} f, \operatorname{grad} f\right\rangle \\
& =\operatorname{Hess} f(\operatorname{grad} V, \operatorname{grad} f)
\end{aligned}
$$

Similarly,

$$
\begin{aligned}
\langle\operatorname{grad}\langle\operatorname{grad} V, \operatorname{grad} f\rangle, \operatorname{grad} f\rangle & =\operatorname{grad} f\langle\operatorname{grad} V, \operatorname{grad} f\rangle \\
& =\left\langle\nabla_{\operatorname{grad} f} \operatorname{grad} V, \operatorname{grad} f\right\rangle+\left\langle\operatorname{grad} V, \nabla_{\operatorname{grad} f} \operatorname{grad} f\right\rangle \\
& =\operatorname{Hess} V(\operatorname{grad} f, \operatorname{grad} f)+\operatorname{Hess} f(\operatorname{grad} f, \operatorname{grad} V) .
\end{aligned}
$$

The desired identity then follows from the symmetry of the Hessian.
Theorem 6.2. For any smooth complete Riemannian manifold $M$ and any $K \in \mathbb{R}$, the following properties are equivalent:
(i) $\operatorname{Ric}_{V, \infty}(M) \geq K$, that is $\operatorname{Ric}_{x}(v)+\operatorname{Hess}_{x} V(v, v) \geq K|v|^{2}$ for all $x \in M$, $v \in T_{x} M$.
(ii) The relative entropy $\mathcal{E}_{\nu}(\cdot)$ with respect to $d \nu(x)=e^{-V(x)} d \operatorname{vol}(x)$ is displacement $K$-convex on $\mathscr{P}_{2}(M)$.

Proof. The proof here is similar and uses the same setup as the one given in Theorem 3.3. Again, we only prove that $(i) \Rightarrow$ (ii). Suppose that $\mu_{0}, \mu_{1} \ll \mathrm{vol}$. Then there exists a unique geodesic $\left(\mu_{t}\right)_{t \in[0,1]}$ in $\mathscr{P}_{2}(M)$ from $\mu_{0}$ to $\mu_{1}$ given by $\mu_{t}=\left(T_{t}\right)_{\#} \mu_{0}$. For each fixed $x$, we consider the curve $\gamma_{x}:[0,1] \rightarrow M$ given by $\gamma_{x}(t):=T_{t}(x)$. We also write $J_{x}(t):=D T_{t}(x)$.

For all $t, \mu_{t} \ll$ vol, so we can write

$$
d \mu_{t}(x)=\rho_{t}(x) d \operatorname{vol}(x)=\rho_{t}(x) e^{V(x)} d \nu(x)
$$

The relative entropy can be written as

$$
\begin{aligned}
\mathcal{E}_{\nu}\left(\mu_{t}\right) & =\int_{M} \log \left(\rho_{t}(x) e^{V(x)}\right) d \mu_{t}(x) \\
& =\mathcal{E}_{\mathrm{vol}}\left(\mu_{t}\right)+\int_{M} V(x) d \mu_{t}(x) \\
& =\mathcal{E}_{\mathrm{vol}}\left(\mu_{0}\right)-\int_{M} \log \left(\operatorname{det} J_{x}(t)\right) d \mu_{0}(x)+\int_{M} V(x) d \mu_{t}(x) \\
& =\mathcal{E}_{\mathrm{vol}}\left(\mu_{0}\right)+\int_{M}\left(-\log \left(\operatorname{det} J_{x}(t)\right)+V\left(\gamma_{x}(t)\right)\right) d \mu_{0}(x),
\end{aligned}
$$

where the third equality comes from (3.5) and the last equality is the change of variables formula for pushforward $\mu_{t}=\left(T_{t}\right)_{\#} \mu_{0}$ together with the identification $T_{t}(x)=\gamma_{x}(t)$. To compute the second derivative of $\mathcal{E}_{\nu}\left(\mu_{t}\right)$, we need the derivative of the two terms in the above integrand. First, recall from (3.6) and (3.7) that

$$
\begin{equation*}
-\frac{d^{2}}{d t^{2}} \log \left(\operatorname{det} J_{x}(t)\right) \geq-\operatorname{tr}\left(J_{x}^{\prime \prime}(t) J_{x}^{-1}(t)\right)=\operatorname{Ric}_{\gamma_{x}(t)}\left(\gamma_{x}^{\prime}(t), \gamma_{x}^{\prime}(t)\right) . \tag{6.2}
\end{equation*}
$$

For the second term $V\left(\gamma_{x}(t)\right)$, its second derivative can then be expressed in terms of Hessian as

$$
\begin{align*}
\frac{d^{2}}{d t^{2}} V\left(\gamma_{x}(t)\right) & =\frac{d}{d t}\left\langle\operatorname{grad} V\left(\gamma_{x}(t)\right), \gamma_{x}^{\prime}(t)\right\rangle \\
& =\left\langle\nabla_{\gamma_{x}^{\prime}(t)}(\operatorname{grad} V), \gamma_{x}^{\prime}(t)\right\rangle+\left\langle\operatorname{grad} V\left(\gamma_{x}(t)\right), \frac{D}{d t} \gamma_{x}^{\prime}(t)\right\rangle \\
& =\operatorname{Hess} V\left(\gamma_{x}^{\prime}(t), \gamma_{x}^{\prime}(t)\right), \tag{6.3}
\end{align*}
$$

where the latter inner-product vanishes because $\frac{D}{d t} \gamma_{x}^{\prime}(t)=0$ as $\gamma_{x}$ is a geodesic.
Now the assumption $\operatorname{Ric}_{V, \infty}(M) \geq K$ can be applied to equations (6.2) and (6.3), yielding

$$
\begin{aligned}
\frac{d^{2}}{d t^{2}}\left(-\log \left(\operatorname{det} J_{x}(t)\right)+V\left(\gamma_{x}(t)\right)\right) & \geq \operatorname{Ric}_{\gamma_{x}(t)}\left(\gamma_{x}^{\prime}(t), \gamma_{x}^{\prime}(t)\right)+\operatorname{Hess} V\left(\gamma_{x}^{\prime}(t), \gamma_{x}^{\prime}(t)\right) \\
& \geq K\left|\gamma_{x}^{\prime}(t)\right|^{2}=K d\left(x, T_{1}(x)\right)^{2}
\end{aligned}
$$

where the last equality $\left|\gamma_{x}^{\prime}(t)\right|=d\left(x, T_{1}(x)\right)$ is, again, due to the fact that $\gamma_{x}$ is a minimal constant speed geodesic whose length is $d\left(x, T_{1}(x)\right)$. Finally, we obtain

$$
\frac{d^{2}}{d t^{2}} \mathcal{E}_{\nu}\left(\mu_{t}\right)=\int_{M} \frac{d^{2}}{d t^{2}}\left(-\log \left(\operatorname{det} J_{x}(t)\right)+V\left(\gamma_{x}(t)\right)\right) d \mu_{0}(x)
$$

$$
\geq \int_{M} K d\left(x, T_{1}(x)\right)^{2} d \mu_{0}(x)=K W_{2}\left(\mu_{0}, \mu_{1}\right)^{2}
$$

since $T_{1}$ is optimal transport map from $\mu_{0}$ to $\mu_{1}$. and we can derive (by using the same trick as in the proof of Theorem 3.3) that

$$
\mathcal{E}_{\nu}\left(\mu_{t}\right) \leq(1-t) \mathcal{E}_{\nu}\left(\mu_{0}\right)+t \mathcal{E}_{\nu}\left(\mu_{1}\right)-\frac{1}{2} K t(1-t) W_{2}\left(\mu_{0}, \mu_{1}\right)^{2} .
$$

## Chapter 7

## Ollivier's coarse Ricci curvature

Another synthetic notion of Ricci curvature on metric measure spaces was proposed by Ollivier in [Oll09], which is called coarse Ricci curvature. While the curvature notion by Lott-Sturm-Villani is based on the displacement convexity of the entropy functional on the $L^{2}$-Wasserstein space, Ollivier's coarse Ricci is defined via the contraction property of $L^{1}$-Wasserstein metric.

The definition of Ollivier's is motivated from the following remarkable observation by von Renesse and Sturm (see [vRS05, Theorem 1.5 (xii)]. Let $\left(M^{n}, g\right)$ be a Riemannian manifold with lower curvature bound $\operatorname{Ric}(M) \geq K$, and let $m_{r, x}$ denote a uniform probability measure (with respect to the volume measure) on a geodesic ball of radius $r$ centered at $x \in M$. The $L^{1}$-Wasserstein distance $W_{1}\left(m_{r, x}, m_{r, y}\right)$ satisfies the asymptotic estimate:

$$
W_{1}\left(m_{r, x}, m_{r, y}\right) \leq\left(1-\frac{K}{2(n+2)} r^{2}+o\left(r^{2}\right)\right) \cdot d(x, y) .
$$

In particular, if $K>0$ then this Wasserstein distance is shorter than the distance between the centers. A refined statement by Ollivier [Oll09] of this observation will be discussed in the upcoming section.
For now, we shall provide the definition Ollivier's notion of coarse Ricci curvature for a metric measure space. Here we work in the setting that $(X, d)$ is a Polish space with $\sigma$-Borel algebra, and $\mathscr{P}_{1}(X)$ denotes the set of all Borel probability measures with finite 1-moments.

Definition 7.1 (random walk). A random walk $m$ on $X$ is a family of probability measures $m_{x}(\cdot) \in \mathscr{P}_{1}(X)$ for each $x \in X$ such that the measure $m_{x}$ depends measurably on $x \in X$, that is, for any Borel set $A \subset X$ and $c \in \mathbb{R}$, the set $\left\{x \in X: m_{x}(A)<c\right\}$ is Borel.

Definition 7.2 (coarse Ricci curvature). Let $(X, d, m)$ be a Polish space with a random walk $m$. Given two different points $x, y \in X$, the coarse Ricci curvature along $x y$ is defined as

$$
\begin{equation*}
\kappa(x, y):=1-\frac{W_{1}\left(m_{x}, m_{y}\right)}{d(x, y)} . \tag{7.1}
\end{equation*}
$$

### 7.1 Distance between two balls

Now we explain two important results from [Oll09] which give justification to the terminology of coarse Ricci curvature defined as in (7.1).

Proposition 7.3 ([Oll09, Proposition 6]). Let $\left(M^{n}, g\right)$ be a smooth connected and complete Riemannian manifold and let $x \in M$. Let $\delta, \varepsilon>0$ small enough, and let $v, w \in T_{x} M$ be unit tangent vectors at $x$. Let $y:=\exp _{x}(\delta v)$ and let $w^{\prime} \in T_{y} M$ be the tangent vector at $y$ obtained by the parallel transport of $w$ along the geodesic $\exp _{x}(t v)$. Moreover, let $x^{\prime}:=\exp _{x}(\varepsilon w)$ and $y^{\prime}:=\exp _{y}\left(\varepsilon w^{\prime}\right)$. Then

$$
\begin{equation*}
d\left(x^{\prime}, y^{\prime}\right) \leq \delta\left(1-\varepsilon^{2} K(v, w)+O\left(\varepsilon^{3}+\varepsilon^{2} \delta\right)\right) \tag{7.2}
\end{equation*}
$$

where $K(v, w)$ is the sectional curvature of the two-dimensional plane spanned by $\{v, w\}$.

Example 7.4 ([Oll09, Example 7]). Let $\left(M^{n}, g\right)$ be a smooth connected and complete Riemannian manifold. For some small $r>0$, let the Markov chain $m^{r}$ be defined as

$$
m_{x}^{r}(A):=\frac{\operatorname{vol}\left(A \cap B_{r}(x)\right)}{\operatorname{vol} B_{r}(x)} \quad \forall \text { Borel } A \subset M
$$

(that is $m_{x}^{r}$ represents the uniform probability measure on the geodesic ball $B_{r}(x)$ ). Let $x \in M$ and let $v \in T_{x} M$ be a unit tangent vector. Let $y$ be a point on the geodesic issuing from $v$, with $d(x, y)$ small enough. Then the $L^{1}$-Wasserstein distance between $m_{x}^{r}$ and $m_{y}^{r}$ satisfies the following estimate:

$$
\begin{equation*}
W_{1}\left(m_{x}^{r}, m_{y}^{r}\right)=d(x, y)\left(1-\frac{r^{2} \operatorname{Ric}(v, v)}{2(n+2)}+O\left(r^{3}+r^{2} d(x, y)\right)\right), \tag{7.3}
\end{equation*}
$$

or equivalently, the coarse Ricci curvature satisfies

$$
\kappa(x, y)=\frac{r^{2} \operatorname{Ric}(v, v)}{2(n+2)}+O\left(r^{3}+r^{2} d(x, y)\right) .
$$

The idea to prove the statement in Example 7.4 is to average (7.2) over all $w$ in the tangent space $T_{x} M$ and over $\varepsilon \in(0, r]$, and then derive the $\leq$ inequality of (7.3). The converse inequality is proved by employing the Kantorovich duality. We omit the proof of this example which can be found in [Oll09, Section 8], and we only provide Ollivier's proof of the above proposition.

Proof of Proposition 7.3. Let $c:[0, \delta] \rightarrow M$ be the geodesic $c(s):=\exp _{x}(s v)$, and let $v_{s}:=c^{\prime}(s) \in T_{c(s)} M$. For each $s$, let $w_{s} \in T_{c(s)} M$ obtained by the parallel transport of $w$ along $c$. Thus $\left\langle v_{s}, w_{s}\right\rangle$ is constant in $s \in[0, \delta]$. Moreover, $v_{0}=v$, $w_{0}=w$, and $w_{\delta}=w^{\prime}$. Consider $F:[0, \delta] \times[0, \varepsilon] \rightarrow M$ a geodesic variation defined by

$$
F(s, t):=c_{s}(t):=\exp _{c(s)}\left(t w_{s}\right)
$$

For each fixed $s_{0}$, let $J_{s_{0}}$ be the variational vector field associated to $F$ of the geodesic $c_{s_{0}}$, that is,

$$
\left.\frac{\partial}{\partial s}\right|_{s=s_{0}} F(s, t)=J_{s_{0}}(t)
$$

Hence $J_{s_{0}}$ is a Jacobi field along $c_{s_{0}}$ and satisfies the Jacobi equation

$$
J_{s_{0}}^{\prime \prime}(t)+R\left(c_{s_{0}}^{\prime}(t), J_{s_{0}}(t)\right) c_{s_{0}}^{\prime}(t)=0
$$

where $J_{s_{0}}^{\prime}(t)=\frac{D}{d t} J_{s_{0}}(t)=\nabla_{c_{s_{0}}^{\prime}(t)} J_{s_{0}}$ and $J_{s_{0}}^{\prime \prime}(t)=\frac{D^{2}}{d t^{2}} J_{s_{0}}(t)=\nabla_{c_{s_{0}}^{\prime}(t)}\left(\nabla_{c_{s_{0}}^{\prime}(t)} J_{s_{0}}\right)$.
Let $\gamma:[0, \delta] \rightarrow M$ be the curve $\gamma(s):=c_{s}(\varepsilon)=\exp _{c(s)}\left(\varepsilon w_{s}\right)$ from $\gamma(0)=x^{\prime}$ to $\gamma(\delta)=y^{\prime}$. We aim to compute the length of $\gamma$, which in turn gives an upper bound for the distance $d\left(x^{\prime}, y^{\prime}\right)$. We have

$$
\gamma^{\prime}\left(s_{0}\right)=\left.\frac{d}{d s}\right|_{s_{0}} F(s, \varepsilon)=J_{s_{0}}(\varepsilon)
$$

The Taylor's expansion of $f(t)=\left\|J_{s_{0}}(t)\right\|^{2}$ is given by

$$
\left\|\gamma^{\prime}\left(s_{0}\right)\right\|^{2}=f(\varepsilon)=f(0)+\varepsilon f^{\prime}(0)+\frac{\varepsilon^{2}}{2} f^{\prime \prime}(0)+O\left(\varepsilon^{3}\right)
$$

where $f(0), f^{\prime}(0)$, and $f^{\prime \prime}(0)$ are computed as

1. $f(0)=\left\|J_{s_{0}}(0)\right\|^{2}=\left\|c^{\prime}\left(s_{0}\right)\right\|^{2}=1$.
2. $f^{\prime}(0)=\left.\frac{d}{d t}\right|_{t=0}\left\langle J_{s_{0}}(t), J_{s_{0}}(t)\right\rangle=2\left\langle J_{s_{0}}^{\prime}(0), J_{s_{0}}(t)\right\rangle=0$ because

$$
J_{s_{0}}^{\prime}(0)=\left.\frac{D}{d t}\right|_{t=0} J_{s_{0}}(t)=\left.\left.\frac{D}{d t}\right|_{t=0} \frac{d}{d s}\right|_{s_{0}} F(s, t)
$$

$$
\begin{array}{ll}
=\left.\left.\frac{D}{d s}\right|_{s_{0}} \frac{d}{d t}\right|_{t=0} F(s, t) & (\text { symmetry lemma }) \\
=\left.\frac{D}{d s}\right|_{s_{0}} w_{s}=0 & \left(w_{s} \text { is parallel along } c\right) .
\end{array}
$$

3. $\frac{1}{2} f^{\prime \prime}(0)=\left\|J_{s_{0}}^{\prime}(0)\right\|^{2}+\left\langle J_{s_{0}}(0), J_{s_{0}}^{\prime \prime}(0)\right\rangle$

$$
\begin{aligned}
& =0-\left\langle J_{s_{0}}(0), R\left(c_{s_{0}}^{\prime}(0), J_{s_{0}}(0)\right) c_{s_{0}}^{\prime}(0)\right\rangle \quad \text { (Jacobi equation) } \\
& =-\left\langle v_{s_{0}}, R\left(w_{s_{0}}, v_{s_{0}}\right) w_{s_{0}}\right\rangle \\
& =-\langle R(w, v) w, v\rangle+O(\delta) \quad \forall s_{0} \in[0, \delta]
\end{aligned}
$$

where the last equality holds true by a linear approximation of a continuously differentiable function $A(s):=\left\langle v_{s}, R\left(w_{s}, v_{s}\right) w_{s}\right\rangle$ around $s=0$.

Thus

$$
\begin{aligned}
\left\|\gamma^{\prime}\left(s_{0}\right)\right\|^{2} & =f(\varepsilon)=f(0)+\varepsilon f^{\prime}(0)+\frac{\varepsilon^{2}}{2} f^{\prime \prime}(0)+O\left(\varepsilon^{3}\right) \\
& =1-\varepsilon^{2}\langle R(w, v) w, v\rangle+O\left(\varepsilon^{3}+\varepsilon^{2} \delta\right) \\
& \leq 1-\varepsilon^{2} K(v, w)+O\left(\varepsilon^{3}+\varepsilon^{2} \delta\right),
\end{aligned}
$$

for all $s_{0} \in[0, \delta]$. Finally, we obtain the desired upper bound of $d\left(x^{\prime}, y^{\prime}\right)$ via the length of $\gamma$ :

$$
d\left(x^{\prime}, y^{\prime}\right) \leq \int_{0}^{\delta}\left\|\gamma^{\prime}(s)\right\|^{2} d s \leq \delta\left(1-\varepsilon^{2} K(v, w)+O\left(\varepsilon^{3}+\varepsilon^{2} \delta\right)\right) .
$$

## 7.2 $W_{1}$-contraction property of random walks

Given a probability measure $\mu \in \mathscr{P}_{1}(X)$, the image of $\mu$ by a random walk $m$ is a measure $\mu * m$ defined via the following convolution with $m$ :

$$
\mu * m:=\int_{x \in X} d \mu(x) m_{x} .
$$

One can also see from the definition that $\delta_{x} * m=m_{x}$. Moreover, a measure obtaind after $n$-step random walk $m$ starting from $\mu$ is given by the iteration $\mu * m^{* n}:=$ $\left(\mu * m^{* n-1}\right) * m$. A stationary distribution (or an invariant distribution) is a probability measure $\pi$ such that $\pi * m=\pi$.

Next we present the so-called Bubley-Dyer theorem [BD97] about $W_{1}$ contraction property. Here the result is reformulated in the context of spaces with coarse Ricci curvature bound, [Oll09, Proposition 20].

Theorem 7.5. Let $(X, d, m)$ be a metric space with a random walk $m$. Let $\kappa \in \mathbb{R}$. Then we have $\kappa(x, y) \geq \kappa$ for all $x, y \in X$ if and only if for any two probability distributions $\mu, \nu$ one has

$$
W_{1}(\mu * m, \nu * m) \leq(1-\kappa) W_{1}(\mu, \nu) .
$$

Moreover, $\mu * m \in \mathscr{P}_{1}(X)$.
An important consequence of the above theorem is the exponential convergence of multi-step random walk.

Corollary 7.6. Assume $\kappa(x, y) \geq \kappa$ for all $x, y \in X$. For $\mu, \nu \in \mathscr{P}_{1}(X)$,

$$
\begin{equation*}
W_{1}\left(\mu * m^{* n}, \nu * m^{* n}\right) \leq(1-\kappa)^{n} W_{1}(\mu, \nu) \tag{7.4}
\end{equation*}
$$

Consequently, if $\kappa>0$, then the random walk has a unique stationary distribution $\mu_{0} \in \mathscr{P}_{1}(X)$. In fact, for all $\mu \in \mathscr{P}_{1}(X)$, the $n$-step random walk $\mu * m^{* n}$ converges in $W_{1}$ exponentially to $\mu_{0}$, i.e.,

$$
\begin{equation*}
W_{1}\left(\mu * m^{* n}, \mu_{0}\right) \leq(1-\kappa)^{n} W_{1}\left(\mu, \mu_{0}\right) \tag{7.5}
\end{equation*}
$$

We would like to remark, in comparison to the $W_{1}$-contraction property of random walks for Ollivier's coarse Ricci curvature, the following continous-time $W_{r^{-}}$ contraction of the heat semigroup for Ricci curvature on Riemannian manifolds; see [vRS05, Corollary 1.4] and also [DS08, Formulae (2.3) and (2.4)] for the proof which uses the fact that heat flow is the gradient flow of entropy.

Theorem 7.7. Let $(M, g)$ be a Riemannian manifold with $\operatorname{Ric}(M) \geq K$. Let $f d \mathrm{vol}, g d \mathrm{vol} \in \mathscr{P}(M)$ be two absolutely continuous measures with density function $f, g \in C^{\infty}(M)$. Then for any $r \in[1, \infty]$,

$$
\begin{equation*}
W_{r}\left(P_{t} f d \mathrm{vol}, P_{t} g d \mathrm{vol}\right) \leq e^{-K t} W_{r}(f d \mathrm{vol}, g d \mathrm{vol}) \tag{7.6}
\end{equation*}
$$

where $P_{t}$ denotes the heat semigroup operator.
This exponential convergence of random walks in case of positive coarse Ricci curvature plays a role of the spectral gap. The following proposition from [Oll09, Proposition 30] provides an analogous result to Lichnerowicz theorem (c.f. Theorem 5.8) in the sense of coarse Ricci curvature.

Definition 7.8 (averaging operator). An averaging operator $M$ is defined to be

$$
M f(x):=\int_{y \in X} f(y) d m_{x}(y)
$$

and the associated random walk Laplacian is $\Delta_{M}:=\mathrm{Id}-M$.
Proposition 7.9 (spectral gap for coarse Ricci curvature). Let $(X, d, m)$ be a metric space with random walk $m$, with invariant distribution $\pi$. Suppose that the coarse Ricci curvature of $X$ is at least $\kappa>0$ and that $\sigma<\infty$. Suppose that $\varphi$ is reversible, or that $X$ is finite. Then the first nonzero eigenvalue $\lambda_{1}$ of the Laplacian $\Delta_{M}$ satisfies

$$
\lambda_{1}\left(\Delta_{M}\right) \geq \kappa .
$$

## Chapter 8

## Original proof of Ricci curvature via displacement convexity of entropy

The goal of this chapter is to explain the original proofs by von Renesse and Sturm [vRS05] and Cordero-Erausquin, McCann and Schmuckenschläger [CEMS01] that the lower Ricci curvature bound can be characterized via the displacement convexity on the entropy functional.
To simplify matters, we assume here that all objects are smooth and we ignore conjugate points and cut loci. Technical details can be found in the original papers.

### 8.1 Differential of exponential maps, Hessian of squared distance and Jacobi fields

Let us start with a technical result about the differential of exponential maps and the Hessian of the squared distance function. This result and its corollary about a particular Jacobi field will help to clarify some proof ideas from [CEMS01] in the later section.

Theorem 8.1. Let $M$ be a complete manifold with no conjugate points and let $u \in T_{x} M$ be a unit vector. Let $x_{s}:=\exp _{x}(s u)$ and let $Z:[0,1] \rightarrow T M$ be a vector field along the geodesic $x_{s}$. We introduce the curve $z_{s}$ in $M$ given by

$$
z_{s}:=\exp _{x_{s}}(Z(s))
$$

and denote $z:=z_{0}=\exp _{x}(Z(0))$. Then

$$
\dot{z}_{s=0}=D \exp _{x}(Z(0))\left(\operatorname{Hess}_{x} \frac{d_{z}^{2}}{2}(u)+\left.\frac{D}{d s}\right|_{s=0} Z\right) .
$$

Here we recall that $d_{p}(\cdot):=d(p, \cdot)$ is the distance function from a point $p \in M$. Moreover, instead of viewing the $\operatorname{Hessian} \operatorname{Hess}_{x} f$ as symmetric bilinear form (c.f. Definition 5.3), we can also consider it as a self-adjoint linear operator $\operatorname{Hess}_{x} f$ : $T_{x} M \rightarrow T_{x} M$ defined via

$$
\left\langle\operatorname{Hess}_{x} f(v), w\right\rangle=\operatorname{Hess} f(v, w),
$$

for all $v, w \in T_{x} M$.

Proof. We consider $z_{s}=\exp (Z(s))$ for the exponential map $\exp : T M \rightarrow M$ on tangent bundle, whose derivative at $w \in T_{x} M$ is the map $D \exp (w): T_{w} T M \rightarrow$ $T_{w} M$. Let us recall the canonical identification $T_{w} T M \cong T_{x} M \times T_{x} M$ for any $w \in T_{x} M$, which is explained as follows.
Any $\xi \in T_{w} T M$ can be represented by $\xi=\dot{C}(0)=\left.\frac{d}{d s}\right|_{s=0} C(s)$ where $C:[0,1] \rightarrow$ $T M$ is a curve in $T M$ such that $C(0)=w$. Let $c:[0,1] \rightarrow M$ be the footpoint curve of $C$, that is, $C(s) \in T_{c(s)} M$ and $c(0)=x$. In other words, $C:[0,1] \rightarrow T M$ is a vector field along $c$. Then the identification is given by

$$
\begin{align*}
T_{w} T M \ni \xi=\left.\frac{d}{d s}\right|_{s=0} C(s) & =\left(\left.\frac{d}{d s}\right|_{s=0} c(s),\left.\frac{D}{d s}\right|_{s=0} C(s)\right) \\
& =\left(\dot{c}(0),\left.\frac{D}{d s}\right|_{s=0} C(s)\right) \in T_{x} M \times T_{x} M \tag{8.1}
\end{align*}
$$

where $\dot{c}(0) \in T_{x} M$ and $\left.\frac{D}{d s}\right|_{s=0} C(s) \in T_{x} M$ are called horizontal and vertical components of $\xi=\dot{C}(0) \in T_{w} T M$, respectively.
Now we differentiate the curve $s \mapsto z_{s}=\exp (Z(s))$ at $s=0$ and use the chain rule to obtain

$$
\begin{equation*}
\left.T_{z} M \ni \frac{d}{d s}\right|_{s=0} z_{s}=D \exp (Z(0))\left(\left.\frac{d}{d s}\right|_{s=0} Z(s)\right) \tag{8.2}
\end{equation*}
$$

The trick is to write $Z(s)=\tilde{Z}(s)+W(s)$ where $\tilde{Z}(s)=-\operatorname{grad} \frac{d_{z}^{2}}{2}\left(x_{s}\right)$, and note that the curve given by $s \mapsto \tilde{z}_{s}:=\exp \tilde{Z}(s)$ is indeed a point curve since

$$
\tilde{z}_{s}=\exp \tilde{Z}(s)=\exp _{x_{s}}\left(-\operatorname{grad} \frac{d_{z}^{2}}{2}\left(x_{s}\right)\right)=\exp _{x_{s}}\left(-d_{z}\left(x_{s}\right) \operatorname{grad} d_{z}\left(x_{s}\right)\right)=z
$$

It follows that

$$
\begin{equation*}
T_{z} M \ni 0_{z}=\left.\frac{d}{d s}\right|_{s=0} \tilde{z}_{s}=D \exp (\tilde{Z}(0))\left(\left.\frac{d}{d s}\right|_{s=0} \tilde{Z}(s)\right) . \tag{8.3}
\end{equation*}
$$

Moreover, the fact that $\tilde{z}_{0}=z_{0}$ implies $Z(0)=\tilde{Z}(0)$ since exponential maps are bijective. Both $\left.\frac{d}{d s}\right|_{s=0} Z(s)$ and $\left.\frac{d}{d s}\right|_{s=0} \tilde{Z}(s)$ have the same horizontal component, namely $\dot{x}_{s=0}$, so the subtraction (8.3) from (8.2) yields only the vertical component:

$$
\begin{aligned}
\dot{z}_{s=0} & =D \exp (Z(0))\left(\left.\frac{d}{d s}\right|_{s=0} Z(s)-\left.\frac{d}{d s}\right|_{s=0} \tilde{Z}(s)\right) \\
& =D \exp (Z(0))\left(\dot{x}_{s=0}-\dot{x}_{s=0},\left.\frac{D}{d s}\right|_{s=0}(Z(s)-\tilde{Z}(s))\right) \\
& =D \exp (Z(0))\left(0_{x},\left.\frac{D}{d s}\right|_{s=0}(Z(s)-\tilde{Z}(s))\right) \\
& =D \exp _{x}(Z(0))\left(\left.\frac{D}{d s}\right|_{s=0}(Z(s)-\tilde{Z}(s))\right) \\
& =D \exp _{x}(Z(0))\left(\left.\frac{D}{d s}\right|_{s=0} Z(s)+\operatorname{Hess}_{x} \frac{d_{z}^{2}}{2}(u)\right) .
\end{aligned}
$$

The above theorem has the following consequence.
Corollary 8.2. Let $u \in T_{x} M$ be a unit vector and $c(s)=\exp _{x}(s u)$. Let $v \in T_{x} M$ be arbitrary and let $\gamma:[0,1] \rightarrow M$ denote the geodesic $\gamma(x):=\exp _{x}(t v)$ with $y=\gamma(1)$. Consider a geodesic variation given by

$$
\gamma_{s}(t):=\exp _{c(s)}\left(-t \operatorname{grad} \frac{d_{y}^{2}}{2}(c(s))\right)
$$

where $d_{y}(\cdot):=d(y, \cdot)$ denotes the distance function from $y$. Then $\gamma_{s}$ is a geodesic variation of $\gamma$, that is, $\gamma_{0}(t)=\gamma(t)$. Moreover, define

$$
V(t):=\left.\frac{\partial}{\partial s}\right|_{s=0} \gamma_{s}(t)
$$

to be the corresponding variational vector field, which is a Jacobi field along $\gamma$. Then we have

$$
V(t)=D \exp _{x}(t v)\left(\operatorname{Hess}_{x} \frac{d_{\gamma(t)}^{2}}{2}(u)-t \operatorname{Hess}_{x} \frac{d_{y}^{2}}{2}(u)\right)
$$

Proof. First, it is straightforward to check that $\gamma_{0}(t)=\gamma(t)$ :

$$
\gamma_{0}(t)=\exp _{x}\left(-t \operatorname{grad} \frac{d_{y}^{2}}{2}(x)\right)=\exp _{x}\left(-t d_{y}(x) \operatorname{grad} d_{y}(x)\right)=\gamma(t)
$$

since $\operatorname{grad} d_{y}(x)=-\frac{\gamma^{\prime}(0)}{d_{y}(x)}$.
Next, we aim to apply Theorem 8.1, so we start with the same setup by letting $x_{s}:=\exp _{x}(s u)=c(s)$. For a fixed $t$, we choose a vector field $Z:[0,1] \rightarrow T M$ along the geodesic $c(s)$ to be

$$
Z(s):=-t \operatorname{grad} \frac{d_{y}^{2}}{2}(c(s)) .
$$

Moreover, let $z_{s}:=\exp _{x_{s}}(Z(s))=\gamma_{s}(t)$ and $z:=z_{0}=\gamma(t)$. We note that $c(0)=x$ and $c^{\prime}(0)=u$, so we derive

$$
\begin{equation*}
\left.\frac{D}{d s}\right|_{s=0} Z=-t \nabla_{c^{\prime}(0)} \operatorname{grad} \frac{d_{y}^{2}}{2}=-t \operatorname{Hess}_{x} \frac{d_{y}^{2}}{2}(u) . \tag{8.4}
\end{equation*}
$$

Applying Theorem 8.1, we obtain

$$
\begin{aligned}
V(t)=\left.\frac{\partial}{\partial s}\right|_{s=0} \gamma_{s}(t)=\dot{z}_{s=0} & =D \exp _{x}(Z(0))\left(\operatorname{Hess}_{x} \frac{d_{z}^{2}}{2}(u)+\left.\frac{D}{d s}\right|_{s=0} Z\right) \\
& \stackrel{(8.4)}{=} D \exp _{x}(Z(0))\left(\operatorname{Hess}_{x} \frac{d_{z}^{2}}{2}(u)-t \operatorname{Hess}_{x} \frac{d_{y}^{2}}{2}(u)\right) .
\end{aligned}
$$

### 8.2 Volume distortion coefficients and Jacobi fields

In this section, we survey two important results in [CEMS01]. The first one describes the volume distortion coefficient in terms of the differential of exponential maps and the Hessian of squared distance function. The second one then uses these coefficients to give an estimation of Jacobian determinant of the interpolating map.
Throughout this section, let $M$ be a smooth connected complete Riemannian manifold of dimension $n$.

Definition 8.3. For $x \in M, Y \subset M$ and $t \in[0,1]$, define the following locus of points

$$
Z_{t}(x, Y):=\{z \in M \mid \exists y \in Y, d(x, z)=t d(x, y) \text { and } d(z, y)=(1-t) d(x, y)\}
$$

Let $B_{r}(y) \subset M$ be the open ball of radius $r$ centered at $y \in M$. The volume distortion coefficient between $x$ and $y$ is defined as

$$
\begin{equation*}
v_{t}(x, y):=\lim _{r \rightarrow 0} \frac{\operatorname{vol}\left[Z_{t}\left(x, B_{r}(y)\right)\right]}{\operatorname{vol}\left[B_{t r}(y)\right]}>0 . \tag{8.5}
\end{equation*}
$$


(a) In negatively curved space, the coefficient is less than 1 .

(b) In positively curved space, the coefficient is greater than 1.

Figure 8.1: Distortion due to curvature effects: the volume distortion coefficient is the ratio of the volume in red to the volume of in blue.

In the case that $\operatorname{Ric}(M) \geq K, K \in \mathbb{R}$, Bishop's comparison theorem [BC64] gives the following estimate of the volume distortion coefficient :

$$
\begin{equation*}
v_{t}(x, y) \geq\left(\frac{\mathbf{s}_{K}(t d(x, y))}{\mathbf{s}_{K}(d(x, y))}\right)^{n-1} \tag{8.6}
\end{equation*}
$$

where $\mathbf{s}_{K}$ is defined by

$$
\mathbf{s}_{K}(r):= \begin{cases}\frac{\sin \left(\sqrt{\frac{K}{n-1}} \cdot r\right)}{\sqrt{\frac{K}{n-1}} \cdot r} & \text { if } \mathrm{K}>0, \\ 1 & \text { if } \mathrm{K}=0, \\ \frac{\sinh \left(\sqrt{\frac{-K}{n-1}} \cdot r\right)}{\sqrt{\frac{-K}{n-1}} \cdot r} & \text { if } \mathrm{K}<0 .\end{cases}
$$

Intuitively, the term $\mathbf{s}_{K}(r)^{n-1}$ is proportional to the area of the sphere of radius $r$ in the $n$-dimensional space form of constant sectional curvature $K /(n-1)$ (and consequently Ricci curvature $K$ ).
Next we present the first result about volume distortion coefficients in terms of the determinants of certain Jacobi matrices.
Definition 8.4. Let $v \in T_{x} M$ and let $\gamma:[0,1] \rightarrow M$ be the geodesic given by $\gamma(0)=x$ and $\gamma(t):=\exp _{x}(t v)$. A Jacobi matrix along $\gamma$ is a map $A(t): T_{x} M \rightarrow$ $T_{\gamma(t)} M$ such that $A(t) u$ is a Jacobi field along $\gamma$ for all $u \in T_{x} M$.
Remark 8.5. The Jacobi matrix $A(t)$ is uniquely determined by $A(0)$ and $A(1)$ if there are no conjugate points along $\gamma$. A well-known example of Jacobi matrices is $A(t):=t D \exp _{x}(t v)$ (see e.g. [GHL90]).


Figure 8.2: The diffeomorphic map $G_{t}: C_{r} \rightarrow Z_{t}\left(x, B_{r}(y)\right)$.
Lemma 8.6 (Lemma 2.1 of [CEMS01]). Fix $x, y \in M$ and let $\gamma(t):=\exp _{x}(t v)$ be the minimal geodesic joining $x=\gamma(0)$ and $y=\gamma(1)$. For $t \in[0,1]$, define $Y(t):=D \exp _{x}(t v)$ and $H(t):=\operatorname{Hess}_{x} \frac{d_{\gamma}^{2}(t)}{2}$. Then the volume distortion coefficient satisfies

$$
\begin{equation*}
v_{t}(x, y)=\frac{\operatorname{det} Y(t)}{\operatorname{det} Y(1)}=\operatorname{det} Y(t) Y(1)^{-1}>0 \tag{8.7}
\end{equation*}
$$

and for $t \neq 1$,

$$
\begin{equation*}
v_{1-t}(y, x)=\operatorname{det} \frac{Y(t)(H(t)-t H(1))}{1-t} \tag{8.8}
\end{equation*}
$$

Proof. For each $t \in[0,1]$, consider a map $G_{t}: T_{x} M \rightarrow M$ given by $G_{t}(u)=$ $\exp _{x}(t u)$. Its differential at $v$ is $D G_{t}(v)=t D \exp _{x}(t v)=t Y(t)$. For small $r>0$, we introduce the set $C_{r}:=\left\{u \in T_{x} M \mid \exp _{x}(u) \in B_{r}(y)\right\}$ and note that $G_{t}$ maps from $C_{r}$ diffeomorphically to $Z_{t}\left(x, B_{r}(y)\right)$, as illustrated in Figure 8.2.

The Jacobian transformation yields

$$
\operatorname{vol}\left[Z_{t}\left(x, B_{r}(y)\right)\right]=\int_{G_{t}\left(C_{r}\right)} 1=\int_{u \in C_{r}}\left|\operatorname{det} D G_{t}(u)\right|=\operatorname{vol}\left(C_{r}\right) \cdot\left|\operatorname{det} D G_{t}\left(u^{\prime}\right)\right|
$$

for some $u^{\prime} \in C_{r}$ by the mean value theorem. As $r \rightarrow 0, u^{\prime}$ converges $v$, and we obtain

$$
t^{n} \operatorname{det} Y(t)=\operatorname{det}(t Y(t))=\operatorname{det}\left(D G_{t}(v)\right)= \pm \lim _{r \rightarrow 0} \frac{\operatorname{vol}\left[Z_{t}\left(x, B_{r}(y)\right)\right]}{\operatorname{vol}\left[C_{r}\right]}
$$

where the sign is positive because $\operatorname{det} Y(t) \neq 0$ for all $t \in[0,1]$ and $Y(0)=\mathrm{Id}$. Since $Z_{1}\left(x, B_{r}(y)\right)=B_{r}(y)$, it follows that

$$
\begin{equation*}
\frac{\operatorname{det} Y(t)}{\operatorname{det} Y(1)}=\lim _{r \rightarrow 0} \frac{\operatorname{vol}\left[Z_{t}\left(x, B_{r}(y)\right)\right]}{t^{n} \operatorname{vol}\left[B_{r}(y)\right]}=\lim _{r \rightarrow 0} \frac{\operatorname{vol}\left[Z_{t}\left(x, B_{r}(y)\right)\right]}{\operatorname{vol}\left[B_{t r}(y)\right]}=v_{t}(x, y), \tag{8.9}
\end{equation*}
$$

which proves the first half of the theorem.
Let $P_{t}^{\gamma}: T_{x} M \rightarrow T_{\gamma(t)} M$ be the parallel transport along $\gamma$ (and note that $\operatorname{det} P_{t}^{\gamma}=$ 1). Consider the unique Jacobi matrix $A(t): T_{x} M \rightarrow T_{\gamma(t)} M$ along $\gamma$ such that $A(0)=0$ and $A(1)=P_{1}^{\gamma}$.
We note from Remark 8.5 that $t Y(t)=t D \exp _{x}(t v): T_{x} M \rightarrow T_{\gamma(t)} M$ is a Jacobi matrix along $\gamma$ and that $Y(1)^{-1} P_{1}^{\gamma} \in \operatorname{End}\left(T_{x} M\right)$. Hence, the following map $A_{\gamma}(t)$ : $T_{x} M \rightarrow T_{\gamma(t)} M$ defined as

$$
A_{\gamma}(t):=t Y(t) Y(1)^{-1} P_{1}^{\gamma}
$$

is a Jacobi matrix along $\gamma$ with $A_{\gamma}(0)=0$ and $A_{\gamma}(1)=P_{1}^{\gamma}$.
In particular, it follows from (8.9) that, for $t \neq 0$,

$$
\begin{equation*}
v_{t}(x, y)=\frac{\operatorname{det} Y(t)}{\operatorname{det} Y(1)}=\operatorname{det} \frac{A_{\gamma}(t)}{t} \tag{8.10}
\end{equation*}
$$

Next, consider the geodesic $\gamma^{-}:[0,1] \rightarrow M$ from $y$ to $x$ given by $\gamma^{-}(t):=\gamma(1-t)$. We obtain similarly to (8.10) that, for $t \neq 1$,

$$
v_{1-t}(y, x)=\operatorname{det} \frac{A_{\gamma^{-}}(1-t)}{1-t}
$$

where $A_{\gamma^{-}}$is a Jacobi matrix along $\gamma^{-}$with $A_{\gamma^{-}}(0)=0$ and $A_{\gamma^{-}}(1)=\left(P_{1}^{\gamma}\right)^{-1}$.
Furthermore, define another Jacobi matrix $B(t): T_{x} M \rightarrow T_{\gamma(t)} M$ along $\gamma$ as $B(t):=A_{\gamma^{-}}(1-t) P_{1}^{\gamma}$. We have $B(0)=\mathrm{Id}, B(1)=0$ and

$$
\begin{equation*}
v_{1-t}(y, x)=\operatorname{det} \frac{A_{\gamma^{-}}(1-t)}{1-t}=\operatorname{det} \frac{B(t)}{1-t} . \tag{8.11}
\end{equation*}
$$

On the other hand, we define $\tilde{B}(t):=Y(t)(H(t)-t H(1)): T_{x} M \rightarrow T_{\gamma(t)} M$ and recall that for all $u \in T_{x} M$,

$$
\tilde{B}(t)(u)=D \exp _{x}(t v)\left(\operatorname{Hess}_{x} \frac{d_{\gamma(t)^{2}}}{2}(u)-t \operatorname{Hess}_{x} \frac{d_{y^{2}}}{2}(u)\right)
$$

is a Jacobi field along $\gamma$ due to Corollary 8.2. Thus $\tilde{B}(t)$ is a Jacobi matrix along $\gamma$, and it has boundary values $\tilde{B}(0)=\operatorname{Id}$ and $\tilde{B}(1)=0$. Under the assumption of no conjugate points, we conclude that $B(t)$ coincides with $\tilde{B}(t)$, and (8.11) yields the desired inequality.

To prepare for upcoming results, we start with the setup and notation given as follows. Let $\phi: M \rightarrow \mathbb{R}$ be a smooth function. For $t \in[0,1]$, define

$$
\begin{align*}
F_{t}(x) & :=\exp _{x}(-t \operatorname{grad} \phi(x)),  \tag{8.12}\\
Y(t) & :=D \exp _{x}(-t \operatorname{grad} \phi(x)), \\
H(t) & :=\operatorname{Hess}_{x} \frac{d_{F_{t}(x)}^{2}}{2} \\
J_{t}(x) & :=\operatorname{det} D F_{t}(x) .
\end{align*}
$$

This function $F_{t}$ plays a role of the optimal transport map as in McCann's theorem (Theorem 1.24).

Lemma 8.7. For all $t \in[0,1], H(t)-t H(1)$ and $H(1)-\operatorname{Hess}_{x} \phi$ are positive semidefinite.

This lemma is described in more details in [CEMS01, Lemma 2.3 and Proposition 4.1(a)]

Proof. For any given $u \in T_{x} M$, let $x_{s}:=\exp _{x}(s u)$. By expressing the Hessian as the second derivative of the function evaluated along the representing curve, we have

$$
\langle(H(t)-t H(1)) u, u\rangle=\left.\frac{d^{2}}{d s^{2}}\right|_{s=0}\left(\frac{d_{F_{t}(x)}^{2}}{2}-t \frac{d_{F_{1}(x)}^{2}}{2}\right)\left(x_{s}\right) .
$$

To check that this term is nonnegative, it suffices to show that $\left(\frac{d_{F_{t}(x)}^{2}}{2}-t \frac{d_{F_{1}(x)}^{2}}{2}\right)(\cdot)$ is minimum at $x$, that is,

$$
d^{2}\left(F_{t}(x), m\right)-t d^{2}\left(F_{1}(x), m\right) \geq d^{2}\left(F_{t}(x), x\right)-t d^{2}\left(F_{1}(x), x\right) .
$$

holds true for all $m \in M$.
For shortened notations, we denote $\ell_{m}:=d\left(F_{1}(x), m\right)$ and $\ell:=d\left(x, F_{1}(x)\right)$, and we recall that $d\left(F_{t_{1}}(x), F_{t_{2}}(x)\right)=\left|t_{1}-t_{2}\right| \ell$ for all $t_{1}, t_{2}$ since $F .(x)$ is a constant speed geodesic. We apply the triangle inequality and derive

$$
\begin{aligned}
d^{2}\left(F_{t}(x), m\right)-t d^{2}\left(F_{1}(x), m\right) & \geq\left(d\left(F_{1}(x), m\right)-d\left(F_{t}(x), F_{1}(x)\right)\right)^{2}-t d^{2}\left(F_{1}(x), m\right) \\
& =(1-t)\left(\ell_{m}^{2}+(1-t) \ell^{2}-2 \ell_{m} \ell\right) \\
& =(1-t)\left(\left(\ell_{m}-\ell\right)^{2}-t^{2} \ell^{2}\right) \\
& \geq-t(1-t) \ell^{2}
\end{aligned}
$$

$$
=d^{2}\left(F_{t}(x), x\right)-t d^{2}\left(F_{1}(x), x\right),
$$

as desired. Similarly, to see that $H(1)-\operatorname{Hess}_{\phi}$ is positive semidefinite, we need to show that $\left(\frac{d_{F_{1}(x)}^{2}}{2}-\phi\right)(\cdot)$ is locally minimum at $x$, that is, to show

$$
\operatorname{grad} \frac{d_{F_{t}(x)}^{2}}{2}(x)-\operatorname{grad} \phi(x)=0
$$

This holds true as we recall $F_{1}(x)=\exp _{x}(-\operatorname{grad} \phi(x))$ and deduce that $\operatorname{grad} \frac{d_{F_{1}(x)}^{2}}{2}(x)=d_{F_{1}(x)}(x) \cdot \operatorname{grad} d_{F_{1}(x)}(x)=|\operatorname{grad} \phi(x)| \cdot \frac{\operatorname{grad} \phi(x)}{|\operatorname{grad} \phi(x)|}=\operatorname{grad} \phi(x)$.

Now we are ready for the second result about an estimate of the Jacobian determinant in terms of volume distortion coefficients.

Lemma 8.8 (Lemma 6.1 of [CEMS01]). Given the same setup as in (8.12), The Jacobian determinant $J_{t}(x):=\operatorname{det} D F_{t}(x)$ can be identified as

$$
\begin{equation*}
J_{t}(x)=\operatorname{det} Y(t)\left(H(t)-t \operatorname{Hess}_{x} \phi\right) . \tag{8.13}
\end{equation*}
$$

Moreover, it satisfies

$$
\begin{equation*}
J_{t}(x)^{\frac{1}{n}} \geq(1-t) v_{1-t}\left(F_{1}(x), x\right)^{\frac{1}{n}}+t v_{t}\left(x, F_{1}(x)\right)^{\frac{1}{n}} J_{1}(x)^{\frac{1}{n}} \tag{8.14}
\end{equation*}
$$

Proof of Lemma 8.8. First, we aim to apply Theorem 8.1 to verify the following identity:

$$
D F_{t}(x)=Y(t)\left(H(t)-t \operatorname{Hess}_{x} \phi\right)
$$

For any $u \in T_{x} M$, we follow most setup from Theorem 8.1: let $x_{s}:=\exp _{x}(s u)$ and choose $Z:[0,1] \rightarrow T M$ to be

$$
Z(s):=-t \operatorname{grad} \phi\left(x_{s}\right) \in T_{x_{s}} M
$$

and let $z_{s}:=\exp _{x_{s}}(Z(s))$ with $z:=z_{0}=\exp _{x}(-t \operatorname{grad} \phi(x))=F_{t}(x)$.
As $x_{0}=0$ and $\dot{x}_{s=0}=u$, we have

$$
D F_{t}(x)(u)=\left.\frac{d}{d s}\right|_{s=0} F_{t}\left(x_{s}\right)=\frac{d}{d s} \exp _{x_{s}} Z(s)=\dot{z}_{s=0}
$$

Applying Theorem 8.1 then yields

$$
D F_{t}(x)(u)=\dot{z}_{s=0}=D \exp _{x}(Z(0))\left(\operatorname{Hess}_{x} \frac{d_{z}^{2}}{2}(u)+\left.\frac{D}{d s}\right|_{s=0} Z\right)
$$

$$
\begin{aligned}
& =Y(t)\left(\operatorname{Hess}_{x} \frac{d_{F_{t}(x)}^{2}}{2}(u)-\left.t \frac{D}{d s}\right|_{s=0} \operatorname{grad} \phi\left(x_{s}\right)\right) \\
& =Y(t)\left(H(t)(u)-t \operatorname{Hess}_{x} \phi(u)\right),
\end{aligned}
$$

as desired, and (8.13) follows immediately.
To see (8.14), we write for $t \in[0,1)$,

$$
J_{t}(x)=\operatorname{det} Y(t) \operatorname{det}\left[(1-t) \frac{H(t)-t H(1)}{1-t}+t\left(H(1)-\operatorname{Hess}_{x} \phi\right)\right],
$$

and we recall from Lemma 8.7 that both matrices $H(t)-t H(1)$ and $H(1)-\operatorname{Hess}_{x} \phi$ are positive semidefinite.
Using the fact that $\operatorname{det}^{\frac{1}{n}}(\cdot)$ is a concave function over $n \times n$ positive semidefinite real symmetric matrices, and the fact that $\operatorname{det} Y(t)>0$, we have

$$
\begin{aligned}
J_{t}(x)^{\frac{1}{n}} & \geq \operatorname{det}^{\frac{1}{n}} Y(t) \cdot\left[(1-t) \operatorname{det}^{\frac{1}{n}} \frac{H(t)-t H(1)}{1-t}+t \operatorname{det}^{\frac{1}{n}}\left(H(1)-\operatorname{Hess}_{x} \phi\right)\right] \\
& =(1-t) \operatorname{det}^{\frac{1}{n}} \frac{Y(t)(H(t)-t H(1))}{1-t}+t \operatorname{det}^{\frac{1}{n}}\left(Y(t) Y^{-1}(1)\right) \cdot J_{1}(x)^{\frac{1}{n}} \\
& =(1-t) v_{1-t}\left(F_{1}(x), x\right)^{\frac{1}{n}}+t v_{t}\left(x, F_{1}(x)\right)^{\frac{1}{n}} J_{1}(x)^{\frac{1}{n}},
\end{aligned}
$$

where the last equality is due to (8.7) and (8.8).

### 8.3 Curvature equivalence on manifolds (revisited)

Now we are properly prepared to follow and understand the proof ideas of the main theorem (c.f. Theorem 3.3) as originally given in [CEMS01].

Theorem 8.9 (Theorem 1.1 of [vRS05]). For any smooth complete Riemannian manifold $M$ and any $K \in \mathbb{R}$, the following properties are equivalent:

1. $\operatorname{Ric}(M) \geq K$, that is $\operatorname{Ric}_{x}(v) \geq K|v|^{2}$ for all $x \in M, v \in T_{x} M$.
2. The entropy $\mathcal{E}_{\mathrm{vol}}(\cdot)$ is displacement $K$-convex on $\mathscr{P}^{2}(M)$.

Original Proof of (i) $\Rightarrow$ (ii). Assume that $\mu_{0}$ and $\mu_{1}$ are absolutely continuous with respect to volume measure. By McCann's theorem (Theorem 1.24), there exists a unique geodesic $\left(\mu_{t}\right)_{t \in[0,1]}$ in $\mathscr{P}^{2}(M)$ connecting $\mu_{0}$ and $\mu_{1}$, namely $\mu_{t}=$
$\left(F_{t}\right)_{\#} \mu_{0}$ where $F_{t}(x)=\exp _{x}(-t \operatorname{grad} \phi(x))$ with some function $\phi: M \rightarrow \mathbb{R}$. Here we recall that the pushforward $\mu_{t}=\left(F_{t}\right)_{\#} \mu_{0}$ means

$$
\int_{M} b(y) d \mu_{t}(y)=\int_{M} b\left(F_{t}(x)\right) d \mu_{0}(x)
$$

for all Borel functions $b: M \rightarrow \mathbb{R}$. Moreover, $\mu_{t}$ is absolutely continuous and can be written as $d \mu_{t}(x)=\rho_{t}(x) d \operatorname{vol}(x)$ with finite probability density $\rho_{t}(x)$. The Jacobian determinant $J_{t}(x):=\operatorname{det} D F_{t}(x)$ satisfies the Monge-Ampère equation (1.7):

$$
\rho_{0}(x)=\rho_{t}\left(F_{t}(x)\right) J_{t}(x)
$$

Applying the pushforward relation and the Monge-Ampère equation, we can derive the entropy term as

$$
\begin{aligned}
\mathcal{E}_{\mathrm{vol}}\left(\mu_{t}\right)=\int_{M} \log \left(\rho_{t}(y)\right) d \mu_{t}(y) & =\int_{M} \log \left(\rho_{t}\left(F_{t}(x)\right)\right) d \mu_{0}(x) \\
& =\int_{M} \log \left(\frac{\rho_{0}(x)}{J_{t}(x)}\right) d \mu_{0}(x) \\
& =\mathcal{E}_{\text {vol }}\left(\mu_{0}\right)-\int_{M} \log J_{t}(x) d \mu_{0}(x) .
\end{aligned}
$$

Next we derive a lower estimate for the convex combination of entropy terms, namely

$$
\mathcal{E}(t):=-\mathcal{E}_{\mathrm{vol}}\left(\mu_{t}\right)+(1-t) \mathcal{E}_{\mathrm{vol}}\left(\mu_{0}\right)+t \mathcal{E}_{\mathrm{vol}}\left(\mu_{1}\right)
$$

by applying Lemma 8.8 and Bishop's comparison theorem as follows.

$$
\begin{aligned}
& \mathcal{E}(t)= \int_{M} \log J_{t}(x) d \mu_{0}(x)-t \int_{M} \log J_{1}(x) d \mu_{0}(x) \\
& \stackrel{(8.14)}{\geq} n \int_{M} \log \left[(1-t) v_{1-t}\left(F_{1}(x), x\right)^{\frac{1}{n}}+t v_{t}\left(x, F_{1}(x)\right)^{\frac{1}{n}} J_{1}(x)^{\frac{1}{n}}\right] d \mu_{0}(x) \\
&-t \int_{M} \log J_{1}(x) d \mu_{0}(x) \\
& \geq \int_{M}\left[(1-t) \log v_{1-t}\left(F_{1}(x), x\right)+t \log \left(v_{t}\left(x, F_{1}(x)\right) J_{1}(x)\right)\right] d \mu_{0}(x)
\end{aligned}
$$

$$
\begin{aligned}
&-t \int_{M} \log J_{1}(x) d \mu_{0}(x) \quad \quad \text { (using concavity of } \log \text { ) } \\
&= \int_{M}\left[(1-t) \log v_{1-t}\left(F_{1}(x), x\right)+t \log v_{t}\left(x, F_{1}(x)\right)\right] d \mu_{0}(x) \\
& \stackrel{(8.6)}{\geq}(n-1) \int_{M}\left[(1-t) \log \mathbf{s}\left((1-t) d\left(x, F_{1}(x)\right)\right)+t \log \mathbf{s}\left(t d\left(x, F_{1}(x)\right)\right)\right. \\
&\left.-\log \mathbf{s}\left(d\left(x, F_{1}(x)\right)\right)\right] d \mu_{0}(x) .
\end{aligned}
$$

In order conclude that $\mathcal{E}_{\text {vol }}(\cdot)$ is displacement $K$-convex, we need to prove that

$$
\mathcal{E}(t) \geq \frac{K}{2} t(1-t) W_{2}^{2}\left(\mu_{0}, \mu_{1}\right)=\frac{K}{2} t(1-t) \int_{M} d^{2}\left(x, F_{1}(x)\right) d \mu_{0}(x) .
$$

Compared to the lower estimate of $\mathcal{E}(t)$ derived earlier, we are left to show in general that for all $r>0$,

$$
(1-t) \log \mathbf{s}((1-t) r)+t \log \mathbf{s}_{K}(t r)-\log \mathbf{s}_{K}(r) \geq \frac{t(1-t)}{2} \frac{K}{n-1} r^{2}
$$

which can be rewritten as

$$
(1-t) \lambda((1-t) r)+t \lambda(t r)-\lambda(r) \geq 0
$$

where $\lambda(r):=\log \mathbf{s}_{K}(r)+\frac{1}{6} \frac{K}{n-1} r^{2}$.
It suffices to show that $\lambda^{\prime}(r) \leq 0$ because the monotonicity of $\lambda$ would then imply $\lambda((1-t) r) \geq \lambda(r)$ and $\lambda(t r) \geq \lambda(r)$ and hence the desired inequality.
With the substitution $\sqrt{\frac{|K|}{n-1}} r=\mathbf{r}$, we obtain

$$
\lambda\left(\frac{n-1}{|K|} \mathbf{r}\right)= \begin{cases}\log \frac{\sin \mathbf{r}}{\mathbf{r}}+\frac{1}{6} \mathbf{r}^{2} \quad \text { if } K>0  \tag{8.15}\\ \log \frac{\sinh \mathbf{r}}{\mathbf{r}}-\frac{1}{6} \mathbf{r}^{2} \quad \text { if } K<0\end{cases}
$$

To verify that its derivative is nonpositive, we need to check that

$$
\begin{equation*}
\mathbf{r} \cos \mathbf{r}-\sin \mathbf{r}+\frac{1}{3} \mathbf{r}^{2} \sin \mathbf{r} \leq 0 \tag{8.16}
\end{equation*}
$$

in case $K>0$, and that

$$
\begin{equation*}
\mathbf{r} \cosh \mathbf{r}-\sinh \mathbf{r}+\frac{1}{3} \mathbf{r}^{2} \sinh \mathbf{r} \leq 0 \tag{8.17}
\end{equation*}
$$

For $K<0$, by differentiating of (8.17) and dividing by $\frac{\mathbf{r}}{3}$, we need $\mathbf{r} \cos \mathbf{r}-\sin \mathbf{r}+$ $\frac{1}{3} \mathbf{r}^{2} \sin \mathbf{r} \leq 0$. Again by differentiation, this follows from $-\mathbf{r} \sinh \mathbf{r} \leq 0$, which is also true for all $\mathbf{r} \leq 0$.

Similarly, in case $K>0$, to check (8.16) requires that $-\mathbf{r} \sin \mathbf{r} \leq 0$. This is true when $\mathbf{r} \in[0, \pi]$, and such restriction is allowed due to Bonnet-Myers theorem.

## Chapter 9

## Overview of curvature notions on discrete spaces

In this chapter, we discuss brief motivations which give rise to the three notions of Ricci curvature on discrete spaces, namely, Ollivier Ricci curvature, Bakry-Émery curvature, and Erbar-Maas entropic Ricci curvature. Each of the following three parts of this thesis is focusing one of these curvature notions.

### 9.1 Setting of discrete spaces

Let us start with the standard convention for three different settings of discrete spaces in which these curvature notions will be discussed.

First is a (combinatorial) graph, denoted by $G=(V, E)$, where $V$ is the set of vertices (or nodes), and $E$ is the set of undirected edges. Two vertices $x, y \in V$ are said to be adjacent (or neighbors) if there is an edge between $x$ and $y$, that is, $\{x, y\} \in E$, and we denote this edge or this adjacency by $x \sim y$ (or $y \sim x$ ). For $x \in V$, the vertex degree of $x$, denoted by $\operatorname{deg}(x)$, is the number of neighbors of $x$. All graphs $G$ are assumed to be simple (i.e., no loops nor multiple edges), connected and locally finite (i.e., every vertex has finite degree). The graph distance function $d: V \times V \rightarrow \mathbb{Z}_{\geq 0}$ is defined for a pair of vertices $x, y \in V$ to be the length (i.e., the number of edges) of a shortest path (also called a geodesic) between $x$ and $y$. The diameter of $G$ is given by $\operatorname{diam}(G):=\sup _{x, y \in V} d(x, y) \in \mathbb{N} \cup\{\infty\}$. For $r \in \mathbb{N}$, the $r$-sphere and the $r$-ball centered at a vertex $x$ consists of all vertices whose distance from $x$ is equal to, and less than or equal to $r$, respectively:

$$
S_{r}(x):=\{y \in V \mid d(x, y)=r\},
$$

$$
B_{r}(x):=\{y \in V \mid d(x, y) \leq r\} .
$$

In particular, $S_{1}(x)$ is the set of all neighbors of $x$. An interval $[x, y]$ is the set of all vertices lying on some geodesic(s) between $x$ and $y$, that is,

$$
\begin{equation*}
[x, y]=\{z \in V \mid d(x, z)+d(z, y)=d(x, y)\} . \tag{9.1}
\end{equation*}
$$

Second is a weighted graph, denoted by $G=(V, \mu, w)$, consisting of a vertex measure $\mu: V \rightarrow \mathbb{R}_{+}$and a symmetric edge-weight function $w: V \times V \rightarrow \mathbb{R}_{\geq 0}$, where we write $u_{x}=\mu(x)$ and $w_{x y}:=w(x, y)$. Two vertices $x, y$ are adjacent if and only if $w(x, y)>0$. Similarly to the case of graphs, all weighted graphs are also assumed to have no loops (that is, $w_{x x}=0$ for all $x \in V$ ) and to be connected and locally finite (that is, for each $x \in V$, there exists finitely many $y \in V$ such that $w_{x y}$ is nonzero). The definition of graph distance, geodesics, diameter, spheres and balls is carried from its underlying non-weighted graph. Furthermore, we define the transition rate from $x$ to $y$ as $p_{x y}:=\frac{w_{x y}}{\mu_{x}}$ and define the weighted vertex degree as $\operatorname{Deg}(x):=\frac{1}{\mu_{x}} \sum_{y \in V} w_{x y}=\sum_{y \in V} p_{x y}$.
Third is a discrete Markov chain $(X, Q, \pi)$, consists of a finite set of vertices $X$ and a Markov kernel $Q: X \times X \rightarrow \mathbb{R}_{\geq 0}$ satisfying $\sum_{y \in X} Q(x, y)=1$. We always assume $Q$ to be irreducible and reversible (see further discussion in Chapter 20), and consequently there exists a unique stationary probability measure (also known as a steady state) $\pi: X \rightarrow \mathbb{R}_{+}$satisfying $\sum_{x \in X} \pi(x)=1, \sum_{x \in X} Q(x, y) \pi(x)=$ $\pi(y)$ and the detailed balance equation $Q(x, y) \pi(x)=Q(y, x) \pi(y)$. The vertex $\pi$-degree of $x$ is defined as $\operatorname{Deg}_{\pi}(x):=\frac{1}{\pi(x)} \sum_{y \in S_{1}(x)} \pi(y)$.
Furthermore, a Markov chain $(X, Q, \pi)$ can be regarded as a weighted graph $G=$ $(X, w, \mu)$ by viewing $w_{x y}=Q(x, y) \pi(x)=w_{y x}$ as a symmetric edge-weight, and $\mu_{x}=\pi(x)$ as a vertex measure. The transition rate is then given by $p_{x y}:=\frac{w_{x y}}{\mu_{x}}=$ $Q(x, y)$, which is the transitional probability to go from $x$ to $y$. Here we allow $Q(x, x)$ to possibly be nonzero (which means the Markov chain can have selfloops), but whether $Q(x, x)=0$ is required or not has no effect on the associated Laplacian $\Delta f:=\sum_{y}(f(y)-f(x)) Q(x, y)$ and the entropic Ricci curvature.

### 9.2 Lin-Lu-Yau's modification of Ollivier Ricci curvature for graphs

The notion of coarse Ricci curvature $\kappa(x, y)=1-\frac{W_{1}\left(\mathfrak{m}_{x}, \mathfrak{m}_{y}\right)}{d(x, y)}$, which was introduced by Ollivier, have been studied in the context of combinatorial graphs $G=(V, E)$ by Lin, Lu and Yau [LLY11]. They consider, for each vertex $x \in V$, the probability
measure $\mathfrak{m}_{x}=\mu_{x}^{\alpha}$ to be obtained by a one-step lazy simple random walk from a vertex $x$ with the probability $\alpha \in(0,1]$ to stay idle. More precisely, $\mu_{x}^{\alpha}(x)=\alpha$ and $\mu_{x}^{\alpha}(v)=(1-\alpha) / \operatorname{deg}(x)$ if $v$ is a neighbor of $x$. Then the $\alpha$-Ollivier Ricci curvature between two different vertices $x, y \in V$ is defined as

$$
\kappa_{\alpha}(x, y):=1-\frac{W_{1}\left(\mu_{x}^{\alpha}, \mu_{y}^{\alpha}\right)}{d(x, y)} .
$$

This curvature notion is the central focus of Part II.

### 9.3 Bakry-Émery curvature on graphs

Recall from Section 5.4 the following Bakry-Émery curvature-dimension condition on a Riemannian manifold $\left(M^{n},\langle\cdot, \cdot\rangle\right)$. Given $K \in \mathbb{R}$ and $N \in[0, \infty]$, a point $x \in M$ satisfies $B E(K, N)$ if the following inequality holds for all $f \in C^{\infty}(M)$,

$$
\begin{equation*}
\Gamma_{2} f(x) \geq \frac{1}{N}(\Delta f(x))^{2}+K \Gamma f(x) \tag{9.2}
\end{equation*}
$$

where

$$
\begin{aligned}
2 \Gamma(f, g) & :=\Delta(f g)-f \Delta g-g \Delta f \\
2 \Gamma_{2}(f, g) & :=\Delta(\Gamma(f, g))-\Gamma(f, \Delta g)-\Gamma(g, \Delta f)
\end{aligned}
$$

and $\Gamma f:=\Gamma(f, f)$ and $\Gamma_{2} f:=\Gamma_{2}(f, f)$. The fact that the condition (9.2) and two iterations $\Gamma$ and $\Gamma_{2}$ are defined via $\Delta$ allows one to generalize the Bakry-Émery curvature-dimension condition $B E(K, N)$ for any spaces admitting the Laplace operator. Early works by Elworthy [Elw91], by Schmuckenschlaäger [Sch99], and by Lin and Yau [LY10] study this Bakry-Émery curvature-dimension condition on weighted graphs $G=(V, w, \mu)$ with the graph Laplacian $\Delta$ defined as

$$
\Delta f(x):=\frac{1}{\mu_{x}} \sum_{y \in V} w_{x y}
$$

In particular, a vertex $x \in V$ is said to satisfy $B E(K, N)$ if (9.2) holds for all functions $f: V \rightarrow \mathbb{R}$. Furthermore, the Bakry-Émery curvature at $x$, denoted as $\mathcal{K}(G, x ; N)$, is the maximum $K$ such that $B E(K, N)$ is satisfied at $x$. This curvature notion is the central focus of Part III.

### 9.4 Erbar-Maas' approach to adapt LSV curvature for discrete Markov chains

Recall the definition by Lott-Sturm-Villani (see Definition 3.4) which defines the lower Ricci curvature bound for a metric measure space ( $X, d, m$ ) by the displacement convexity of the entropy functional. This definition requires firstly that there exists a geodesic connecting any two given probability measures; in other words, the Wasserstein space $\mathscr{P}_{2}(X)$ needs to be a geodesic space, which is the case for when $X$ is geodesic space (see Theorem 1.20). On the other hand, when $X$ is a discrete space, there is no constant speed geodesic joining any two distinct probability measures.

Proposition 9.1. Let $(X, d)$ be a discrete metric space. Then there is no constant speed geodesic $\left(\mu_{t}\right)_{t \in[0,1]}$ in the Wasserstein space $\left(\mathscr{P}_{2}(X), W_{2}\right)$, except trivial geodesics where $\mu_{t}$ is constant on $t$.

Proof. Assume that $\left(\mu_{t}\right)_{t \in[0,1]}$ is a constant speed geodesic. Fix an arbitrary point $x \in X$. Since $X$ is discrete, there exist $\delta>0$ such that $\delta<d(x, y)$ for all $y \in X \backslash\{x\}$. For any transport plan $\pi \in \prod\left(\mu_{s}, \mu_{t}\right)$, its marginal constraints imply that $\pi(x, x) \leq \mu_{t}(x)$ and $\pi(x, x) \leq \mu_{s}(x)$, and that

$$
\begin{gathered}
\sum_{y \in X \backslash\{x\}} \pi(x, y)=\mu_{s}(x)-\pi(x, x) \geq \max \left\{\mu_{s}(x)-\mu_{t}(x), 0\right\} \\
\sum_{y \in X \backslash\{x\}} \pi(y, x)=\mu_{t}(x)-\pi(x, x) \geq \max \left\{\mu_{t}(x)-\mu_{s}(x), 0\right\}
\end{gathered}
$$

The quadratic cost of $\pi$ then satisfies

$$
\sum_{X \times X} d^{2}(u, v) \pi(u, v)>\delta^{2} \sum_{y \in X \backslash\{x\}}(\pi(x, y)+\pi(y, x)) \geq \delta^{2}\left|\mu_{s}(x)-\mu_{t}(x)\right| .
$$

The constant speed property of $\left(\mu_{t}\right)_{t \in[0,1]}$ then gives

$$
|s-t| W_{2}\left(\mu_{0}, \mu_{1}\right)=W_{2}\left(\mu_{s}, \mu_{t}\right) \geq \delta \sqrt{\left|\mu_{s}(x)-\mu_{t}(x)\right|}
$$

which implies $\left|\mu_{s}(x)-\mu_{t}(x)\right| \leq L|s-t|^{2}$ for all $s, t \in[0,1]$, where $L:=W_{2}\left(\mu_{0}, \mu_{1}\right)^{2} / \delta$ is a constant. By fixing $s$, one has the limit $\lim _{t \rightarrow s}\left|\frac{\mu_{t}(x)-\mu_{s}(x)}{t-s}\right| \leq \lim _{t \rightarrow s} L|s-t|=0$, which means $t \mapsto \mu_{t}(x)$ is differentiable at $t=s$ and its derivative is zero. It follows that $\mu_{t}(x)$ is constant on $t$ for every $x \in X$.

There were two approaches to circumvent the lack of constant speed geodesics in the Wasserstein space $\left(\mathscr{P}_{2}(X), W_{2}\right)$ for discrete spaces $X$. The first approach by Bonciocat and Sturm [BS09] introduces "approximated points" along "rough geodesics". The second approach by Erbar and Maas [EM12], modifies $W_{2}$-metric into a new transport metric, called $\mathcal{W}$ by applying the concept of Otto's calculus and Benamou-Brenier formula in such a way that $\left(\mathscr{P}_{2}(X), \mathcal{W}\right)$ is a geodesic space. This approach by Erbar and Maas gives rise the entropic Ricci curvature on discrete Markov chains, which is the central focus of Part IV.

## Part II

Ollivier Ricci curvature on graphs

## Chapter 10

## $W_{1}$ on graphs

In this chapter, we briefly recall basic concepts of optimal transport, including the Kantorovich cost-minimizing transportation problem and its duality. These concepts were already discussed in Section 1.1, but we would like to present them here in the context of $L^{1}$-Wasserstein distance $W_{1}$ on graphs. Readers are encouraged to consult the book by Peyré and Cuturi [PC19, Chapter 6].

### 10.1 Cost-minimizing problem and dual problem

As usual, let $G=(V, E)$ be a simple, connected and locally finite graph and let

$$
\mathscr{P}(V):=\left\{\mu: V \rightarrow[0,1]\left|\sum_{x \in V} \mu(x)=1,|\operatorname{supp}(\mu)|<\infty\right\}\right.
$$

be the set of finitely supported probability measures on $G$. For $\mu, \nu \in \mathscr{P}(V)$, the 1-Wasserstein distance $W_{1}$ is defined via the following Kantorovich optimal transportation problem:

$$
\begin{equation*}
W_{1}(\mu, \nu):=\inf _{\pi \in \Pi(\mu, \nu)}\left\{\operatorname{cost}(\pi):=\sum_{x, y \in V} d(x, y) \pi(x, y)\right\} \tag{10.1}
\end{equation*}
$$

where the infimum runs over the set of all transport plans from $\mu$ to $\nu$,

$$
\Pi(\mu, \nu):=\left\{\begin{array}{l|ll}
\pi: V \times V \rightarrow[0,1] & \begin{array}{ll}
\sum_{y \in V} \pi(x, y)=\mu(x) & \forall x \in V \\
\sum_{x \in V} \pi(x, y)=\nu(y) & \forall y \in V
\end{array} \tag{10.2}
\end{array}\right\}
$$

Note that the marginal constraints in (10.2) implies that $\operatorname{supp}(\pi) \subset \operatorname{supp}(\mu) \times$ $\operatorname{supp}(\nu)$, and consequently

$$
\sum_{y \in \operatorname{supp}(\nu)} \pi(x, y)=\mu(x) \quad \text { and } \quad \sum_{x \in \operatorname{supp}(\mu)} \pi(x, y)=\nu(y) .
$$

Any minimizer $\pi$ to the problem (10.1) is called an optimal transport plan from $\mu$ to $\nu$, and we denote by $\Pi_{o p t}(\mu, \nu)$ the set of all such optimal transport plans. On the other hand, the 1-Wasserstein distance can also be reformulated via the following Kantorovich dual problem:

$$
\begin{equation*}
W_{1}(\mu, \nu)=\sup _{\phi \in \operatorname{Lip}_{1}(V)} \sum_{x \in V} \phi(x)(\mu(x)-\nu(x)), \tag{10.3}
\end{equation*}
$$

where $\operatorname{Lip}_{1}(V):=\{\phi: V \rightarrow \mathbb{R} \mid \phi(x)-\phi(y) \leq d(x, y) \forall x, y\}$ is the set of all 1-Lipschitz functions on $V$. Any maximizer $\phi$ to the problem (10.3) is called an optimal Kantorovich potential.

The Existence of optimal plans and optimal Kantorovich potentials is asserted from a general optimal transport theory on metric spaces (see Theorem 1.10 and Remark 1.11). Alternatively, one can regard (10.1) as a standard finite-dimensional linear optimization problem with a variable $\pi=(\pi(x, y))_{\operatorname{supp}(\mu) \times \operatorname{supp}(\nu)} \geq 0$, whose minimizer always exists.

An important aspect of the Kantorovich duality is the relation between optimal transport plans and optimal Kantorocih potential, which can be described by the following complementary slackness theorem.

Theorem 10.1 (Complementary Slackness). Let $G=(V, E)$ be a graph. Suppose that $\pi$ is an optimal transport plan and $\phi$ is an optimal Kantorovich potential from $\mu$ to $\nu$ in $\mathscr{P}(V)$. Then

$$
\phi(x)-\phi(y)=d(x, y) \quad \forall(x, y) \in \operatorname{supp}(\pi) .
$$

### 10.2 Triangle inequality and transport geodesics

It holds in general that for every $r \in[1, \infty)$ and a metric measure space $(X, d)$, the Wasserstein distance $W_{r}$ defines a metric on the finite- $r$-moment probability measure space $\mathscr{P}_{r}(X)$ (see Proposition 1.13). Here we decide to present a proof in our particular case of $W_{1}$ on graphs. Minkowski's inequality for $L^{1}$ spaces is simply the triangle-inequality, and the gluing lemma shares the same concept as the concatenation of transport plans, which we will introduce as follows.

Proposition 10.2. Let $G=(V, E)$ be a graph. Then $W_{1}$ defines a metric on the space of probability measures $\mathscr{P}(V)$.

Proposition 10.3 (concatenation of transport plans). Let $G=(V, E)$ be a graph and $\mu, \nu, \rho \in \mathscr{P}(V)$ be three probability measures. For any transport plans $\pi_{1} \in$ $\Pi(\mu, \nu)$ and $\pi_{2} \in \Pi(\nu, \rho)$, the concatenation of plans $\pi_{1}$ and $\pi_{2}$, denoted as $\pi_{2} \circ \pi_{1}$, is defined to be the function $\pi_{3}: V \times V \rightarrow[0,1]$ such that

$$
\begin{equation*}
\pi_{3}(x, z):=\sum_{y \in \operatorname{supp}(\nu)} \frac{\pi_{1}(x, y) \pi_{2}(y, z)}{\nu(y)} \quad \forall x, z \in V \tag{10.4}
\end{equation*}
$$

Then $\pi_{3}$ defines a transport plan from $\mu$ to $\rho$.
Moreover, if $\pi_{1}, \pi_{2}$ are given by transport maps $T_{1}, T_{2}: V \rightarrow V$, then the concatenation $\pi_{3}=\pi_{2} \circ \pi_{1}$ is also given by a composite map $T_{3}=T_{2} \circ T_{1}$.

Proof of Proposition 10.3. We need to verify the marginal constraints for $\pi_{3} \in$ $\Pi(\mu, \rho)$. First, for all $x \in V$, we have

$$
\begin{aligned}
\sum_{z} \pi_{3}(x, z) & =\sum_{z} \sum_{y \in \operatorname{supp}(\nu)} \frac{\pi_{1}(x, y) \pi_{2}(y, z)}{\nu(y)} \\
& =\sum_{y \in \operatorname{supp}(\nu)} \frac{\pi_{1}(x, y)}{\nu(y)} \cdot \nu(y)=\sum_{y \in \operatorname{supp}(\nu)} \pi_{1}(x, y)=\mu(x),
\end{aligned}
$$

where we apply the marginals $\sum_{z} \pi_{2}(y, z)=\nu(y)$ and $\sum_{y} \pi_{1}(x, y)=\mu(x)$. Similarly, it can be shown that $\sum_{x} \pi_{3}(x, z)=\rho(z)$, and hence $\pi_{3} \in \Pi(\mu, \rho)$ as desired.

Moreover, if transport maps $T_{1}, T_{2}: V \rightarrow V$ induces the plans $\pi_{1}, \pi_{2}$, namely,

$$
\pi_{1}(x, y)=\left\{\begin{array}{ll}
\mu(x) & \text { if } y=T_{1}(x), \\
0 & \text { otherwise }
\end{array} ; \pi_{2}(y, z)= \begin{cases}\nu(y) & \text { if } z=T_{2}(y) \\
0 & \text { otherwise }\end{cases}\right.
$$

Then $\pi_{3}(x, z)$ defined via (10.4) is nonzero if and only if $\pi_{1}(x, y)>0$ and $\pi_{2}(y, z)>$ 0 , which occurs exactly when $y=T_{1}(x)$ and $z=T_{2}(y)$. In this case, $\sum_{z} \pi_{3}(x, z)=$ $\pi_{3}\left(x, T_{2}\left(T_{1}(x)\right)\right)=\mu(x)$. Therefore,

$$
\pi_{3}(x, z)= \begin{cases}\mu(x) & \text { if } z=T_{2} \circ T_{1}(x) \\ 0 & \text { otherwise }\end{cases}
$$

which means $\pi_{3}$ is induced from the composition $T_{2} \circ T_{1}: V \rightarrow V$.
Now we are ready to prove that $W_{1}$ satisfies the three metric axioms.

Proof of Proposition 10.2. The symmetry property $W_{1}(\mu, \nu)=W_{1}(\nu, \mu)$ is obvious, since for every transport plan $\pi \in \Pi(\mu, \nu)$, the reverse plan $\pi^{(-1)} \in \Pi(\nu, \mu)$ defined by $\pi^{(-1)}(x, y):=\pi(y, x)$ has the same cost as $\pi$. The identity property $W_{1}(\mu, \nu)=0 \Rightarrow \mu=\nu$ can also be proved easily by a contrapositive argument. We are left to check the triangle inequality, namely

$$
W_{1}(\mu, \nu)+W_{1}(\nu, \rho) \geq W_{1}(\mu, \rho) \quad \forall \mu, \nu, \rho \in \mathscr{P}(V) .
$$

Assume $\pi_{1} \in \Pi(\mu, \nu)$ and $\pi_{2} \in \Pi(\nu, \rho)$, and consider the concatenation $\pi_{3}:=$ $\pi_{2} \circ \pi_{1} \in \Pi(\mu, \rho)$. As one would expect, the cost of $\pi_{3}$ is no more than that of $\pi_{1}$ and $\pi_{2}$ combined:

$$
\begin{align*}
\operatorname{cost}\left(\pi_{3}\right) & =\sum_{x, z} d(x, z) \pi_{3}(x, z) \\
& =\sum_{x, z} \sum_{y \in \operatorname{supp}(\nu)} d(x, z) \cdot \frac{\pi_{1}(x, y) \pi_{2}(y, z)}{\nu(y)} \\
& \leq \sum_{x, z} \sum_{y \in \operatorname{supp}(\nu)}(d(x, y)+d(y, z)) \frac{\pi_{1}(x, y) \pi_{2}(y, z)}{\nu(y)}  \tag{10.5}\\
& =\sum_{x} \sum_{y \in \operatorname{supp}(\nu)} d(x, y) \pi_{1}(x, y)+\sum_{z} \sum_{y \in \operatorname{supp}(\nu)} d(y, z) \pi_{2}(y, z) \\
& =\sum_{x, y} d(x, y) \pi_{1}(x, y)+\sum_{y, z} d(y, z) \pi_{2}(y, z)=\operatorname{cost}\left(\pi_{1}\right)+\operatorname{cost}\left(\pi_{2}\right) .
\end{align*}
$$

By choosing $\pi_{1}$ and $\pi_{2}$ to be optimal transport plans, we conclude that

$$
\begin{equation*}
W_{1}(\mu, \rho) \leq \operatorname{cost}\left(\pi_{3}\right) \leq W_{1}(\mu, \nu)+W_{1}(\nu, \rho) . \tag{10.6}
\end{equation*}
$$

Remark 10.4. We may refer to the Wasserstein space ( $\left.\mathscr{P}(V), W_{1}\right)$ as the space of probability measures $\mathscr{P}(V)$ equipped with the Wasserstein metric $W_{1}$. Moreover, the underlying metric space $(V, d)$ is embedded isometrically into the Wasserstein space $\left(\mathscr{P}(V), W_{1}\right)$ by $x \mapsto \delta_{x}$ since $W_{1}\left(\delta_{x}, \delta_{y}\right)=d(x, y)$ (compared to Definition 1.14 and Remark 1.16).

Corollary 10.5. Given $\mu, \nu, \rho \in \mathscr{P}(V)$, suppose

$$
W_{1}(\mu, \rho) \leq W_{1}(\mu, \nu)+W_{1}(\nu, \rho) .
$$

Then for any optimal transport plans $\pi_{1} \in \prod_{\text {opt }}(\mu, \nu)$ and $\pi_{2} \in \prod_{o p t}(\nu, \rho)$, we have

1. the concatenation $\pi_{3}=\pi_{2} \circ \pi_{1}$ is optimal, and
2. $d(x, y)+d(y, z)=d(x, z)$ for all pairs $(x, y) \in \operatorname{supp}\left(\pi_{1}\right)$ and $(y, z) \in \operatorname{supp}\left(\pi_{2}\right)$.

Proof of Corollary 10.5. The fact that $W_{1}(\mu, \rho)=W_{1}(\mu, \nu)+W_{1}(\nu, \rho)$ implies that the inequalities (10.6) hold with equality, which means $\pi_{3}$ is optimal. Moreover, the triangle inequality (10.5) holds with equality, so we must have $d(x, z)=d(x, y)+$ $d(y, z)$ whenever $\pi_{1}(x, y)>0$ and $\pi_{2}(y, z)>0$.

Moreover, we imitate the convention of intervals, introduced as in (9.1), for Wasserstein space $\left(\mathscr{P}(V), W_{1}\right)$ and say that an interval

$$
[\mu, \rho]:=\left\{\nu \in \mathscr{P}(V) \mid W_{1}(\mu, \nu)+W_{1}(\nu, \rho)=W_{1}(\mu, \rho)\right\}
$$

consists of all measures $\nu$ lying on a $W_{1}$-geodesic from $\mu$ to $\rho$. Then one can reformulate the above corollary and extend the result inductively as follows.

Corollary 10.6. Suppose that $\mu_{1}, \mu_{2}, \ldots, \mu_{n} \in \mathscr{P}(V)$ lie orderly on a $W_{1}$-geodesic, that is,

$$
\sum_{i=1}^{n-1} W_{1}\left(\mu_{i}, \mu_{i+1}\right)=W_{1}\left(\mu_{1}, \mu_{n}\right) .
$$

For all $1 \leq i \leq n-1$, let $\pi_{i} \in \prod_{\text {opt }}\left(\mu_{i}, \mu_{i+1}\right)$ be an optimal transport plan. Then the concatenation of the supports of $\pi_{1}, \pi_{2}, \ldots, \pi_{n}$ forms a geodesic on $V$, that is,

$$
\sum_{i=1}^{n-1} d\left(x_{i}, x_{i+1}\right)=d\left(x_{1}, x_{n}\right)
$$

for all pairs $\left(x_{i}, x_{i+1}\right) \in \operatorname{supp}\left(\pi_{i}\right)$.
Definition 10.7 (transport geodesics). Suppose that $\mu_{1}, \mu_{2}, \ldots, \mu_{n} \in \mathscr{P}(V)$ lie orderly on a $W_{1}$-geodesic. Then any such sequence of vertices $x_{1}, x_{2}, \ldots, x_{n}$ (which may contain repetitions among consecutive vertices) from Corollary 10.6 is called a transport geodesic associated to the sequence $\mu_{1}, \mu_{2}, \ldots, \mu_{n}$.

## Chapter 11

## Ollivier Ricci curvature (ORC) with Lin-Lu-Yau modification

### 11.1 Relevant results about Ollivier Ricci curvature

$\mathrm{Lin}, \mathrm{Lu}$ and Yau introduce in [LLY11] a convention for Ollivier Ricci curvature on a combinatorial graph by considering a "small ball" centered at a vertex $x$ to be obtained after a one-step lazy simple random walk with idleness/laziness $\alpha$ from $x$.

Definition 11.1 ([LLY11]). Let $G=(V, E)$ be a locally finite graph. Let $\alpha \in$ $[0,1]$ be an idleness parameter. For any vertex $x \in V$, the probability measure $\mu_{x}^{\alpha}: V \rightarrow[0,1]$ is given by

$$
\mu_{x}^{\alpha}(v):= \begin{cases}\alpha & \text { if } v=x \\ \frac{1-\alpha}{\operatorname{deg}(x)} & \text { if } v \in S_{1}(x) \\ 0 & \text { otherwise }\end{cases}
$$

The $\alpha$-Ollivier Ricci curvature between two different vertices $x, y \in V$ is given by

$$
\kappa_{\alpha}(x, y):=1-\frac{W_{1}\left(\mu_{x}^{\alpha}, \mu_{y}^{\alpha}\right)}{d(x, y)}
$$

and the Lin-Lu-Yau curvature is defined to be

$$
\kappa(x, y):=\lim _{\alpha \rightarrow 1} \frac{\kappa_{\alpha}(x, y)}{1-\alpha} .
$$

The following lemma provides a useful fact that Ollivier Ricci curvature between two vertices with great distance apart can be bounded from below by the curvature between some pair of vertices which are closer.

Lemma 11.2. Suppose that $y \in[x, z]$. Then for all $\alpha \in[0,1]$, we have

$$
\kappa_{\alpha}(x, z) \geq \min \left\{\kappa_{\alpha}(x, y), \kappa_{\alpha}(y, z)\right\}
$$

Proof. Fix any $\alpha \in[0,1]$, and let $K_{\alpha}:=\min \left\{\kappa_{\alpha}(x, y), \kappa_{\alpha}(y, z)\right\}$. The triangle inequality $W_{1}\left(\mu_{x}^{\alpha}, \mu_{z}^{\alpha}\right) \leq W_{1}\left(\mu_{x}^{\alpha}, \mu_{y}^{\alpha}\right)+W_{1}\left(\mu_{y}^{\alpha}, \mu_{z}^{\alpha}\right)$ implies that

$$
\begin{aligned}
\left(1-\kappa_{\alpha}(x, z)\right) d(x, z) & \leq\left(1-\kappa_{\alpha}(x, y)\right) d(x, y)+\left(1-\kappa_{\alpha}(y, z)\right) d(y, z) \\
& \leq\left(1-K_{\alpha}\right)(d(x, y)+d(y, z)) \\
& =\left(1-K_{\alpha}\right) d(x, z)
\end{aligned}
$$

due to $y \in[x, z]$. We can then conclude that $\kappa_{\alpha}(x, z) \geq K_{\alpha}$ as desired.
An immediate consequence of the above lemma is that the infimum of curvature among all pairs of different vertices is equal to the infimum of curvature between any two adjacent vertices. Therefore, when we discuss the lower bound of curvature on graphs, it makes sense to consider Ollivier Ricci curvature restrictively on edges of graphs.

Corollary 11.3. For all $\alpha \in[0,1]$, we have

$$
\inf _{x \neq y} \kappa_{\alpha}(x, y)=\inf _{x \sim y} \kappa_{\alpha}(x, y) .
$$

Example 11.4 (hypercube). The hypercube $\mathcal{Q}^{n}$ can be viewed as the graph whose vertices are elements of $\{0,1\}^{n}$, and two vertices are adjacent if and only if their Hamming distance (i.e., the number of coordinates at which the corresponding values are different) is one. Consider a vertex $x=(0,0, \ldots, 0)=\overline{0}$ with $n$ neighbors represented by $\left\{\mathrm{e}_{i} \mid 1 \leq i \leq n\right\}$, where $\mathrm{e}_{i}$ is the standard unit vector with a 1 in the $i$-th coordinate and 0 's elsewhere. Let $y=\mathrm{e}_{1}$, and we aim to compute $\kappa_{\alpha}(x, y)=1-W_{1}\left(\mu_{x}^{\alpha}, \mu_{y}^{\alpha}\right)$, where the probability measures $\mu_{x}^{\alpha}, \mu_{y}^{\alpha}$ are

$$
\mu_{x}^{\alpha}(z)=\left\{\begin{array}{ll}
\alpha, & \text { if } z=\overline{0}, \\
\frac{1-\alpha}{n}, & \text { if } z=\mathrm{e}_{1}, \\
\frac{1-\alpha}{n}, & \text { if } z=\mathrm{e}_{i} \\
& \text { for } 2 \leq i \leq n, \\
0, & \text { otherwise },
\end{array} \quad \text { and } \mu_{y}^{\alpha}(z)= \begin{cases}\alpha, & \text { if } z=\mathrm{e}_{1}, \\
\frac{1-\alpha}{n}, & \text { if } z=\overline{0}, \\
\frac{1-\alpha}{n}, & \text { if } z=\mathrm{e}_{1}+\mathrm{e}_{i} \\
& \text { for } 2 \leq i \leq n, \\
0, & \text { otherwise }\end{cases}\right.
$$

Denote $m:=\min \left\{\alpha, \frac{1-\alpha}{n}\right\} \in \mathbb{R}_{\geq 0}$. One may be convinced that to transport optimally from $\mu_{x}^{\alpha}$ to $\mu_{y}^{\alpha}$, the masses can stay put at $\overline{0}$ and $\mathrm{e}_{1}$ for at most $m$ units each. The other masses of $1-2 m$ units, collectively, need to move with the distance one (between $\overline{0}$ and $\mathrm{e}_{1}$, and from $\mathrm{e}_{i}$ to $\mathrm{e}_{1}+\mathrm{e}_{i}$ for $i \geq 2$ ), so the total cost is $1-2 m$. Explicitly, the corresponding transport plan $\pi \in \Pi\left(\mu_{x}^{\alpha}, \mu_{y}^{\alpha}\right)$ is given by $\pi(\overline{0}, \overline{0})=\pi\left(\mathrm{e}_{1}, \mathrm{e}_{1}\right)=m, \pi\left(\overline{0}, \mathrm{e}_{1}\right)=\alpha-m, \pi\left(\mathrm{e}_{1}, \overline{0}\right)=\frac{1-\alpha}{n}-m$ and $\pi\left(\mathrm{e}_{i}, \mathrm{e}_{i}\right)=\frac{1-\alpha}{n}$ for all $2 \leq i \leq n$. This gives the total cost of $\pi$ to be

$$
\operatorname{cost}(\pi)=(\alpha-m)+\left(\frac{1-\alpha}{n}-m\right)+(n-1)\left(\frac{1-\alpha}{n}\right)=1-2 m
$$

which means $W_{1}\left(\mu_{x}^{\alpha}, \mu_{y}^{\alpha}\right) \leq 1-2 m$.
To verify that $W_{1}\left(\mu_{x}^{\alpha}, \mu_{y}^{\alpha}\right) \geq 1-2 m$, we need to employ the Kantorovich Duality. We construct a 1-Lipschitz function $\phi: V \rightarrow \mathbb{R}$ as follows:

$$
\left(\phi(\overline{0}), \phi\left(\mathrm{e}_{1}\right)\right)= \begin{cases}(0,1) & \text { if } \alpha \in\left[0, \frac{1}{n+1}\right] \\ (1,0) & \text { if } \alpha \in\left[\frac{1}{n+1}, 1\right]\end{cases}
$$

and $\phi\left(\mathrm{e}_{i}\right)=1$ and $\phi\left(\mathrm{e}_{1}+\mathrm{e}_{i}\right)=0$ for all $2 \leq i \leq n$. So far, $\phi$ satisfies the 1-Lipschitz condition among vertices in $\operatorname{supp}\left(\mu_{x}^{\alpha}\right) \cup \operatorname{supp}\left(\mu_{y}^{\alpha}\right)$, and it can be 1Lipschitz extended to all other vertices. Then we can deduce from the dual problem that

$$
\begin{aligned}
W_{1}\left(\mu_{x}^{\alpha}, \mu_{y}^{\alpha}\right) & \geq \sum_{z \in V} \phi(z)\left(\mu_{x}^{\alpha}(z)-\mu_{y}^{\alpha}(z)\right) \\
& =\frac{1-\alpha}{n}(n-1)+\left(\phi(\overline{0})-\phi\left(\mathrm{e}_{1}\right)\right)\left(\alpha-\frac{1-\alpha}{n}\right) \\
& = \begin{cases}1-2 \alpha & \text { if } \alpha \in\left[0, \frac{1}{n+1}\right], \\
1-\frac{2(1-\alpha)}{n} & \text { if } \alpha \in\left[\frac{1}{n+1}, 1\right]\end{cases}
\end{aligned}
$$

Therefore, $W_{1}\left(\mu_{x}^{\alpha}, \mu_{y}^{\alpha}\right)=1-2 m$. In conclusion, the $\alpha$-Ollivier Ricci curvature on an edge of the hypercube $\mathcal{Q}^{n}$ is given by

$$
\kappa_{\alpha}(x, y)=1-W_{1}\left(\mu_{x}^{\alpha}, \mu_{y}^{\alpha}\right)=2 m= \begin{cases}2 \alpha & \text { if } \alpha \in\left[0, \frac{1}{n+1}\right] \\ \frac{2(1-\alpha)}{n} & \text { if } \alpha \in\left[\frac{1}{n+1}, 1\right]\end{cases}
$$

and the Lin-Lu-Yau curvature is $\kappa(x, y)=\lim _{\alpha \rightarrow 1} \frac{\kappa_{\alpha}(x, y)}{1-\alpha}=\frac{2}{n}$.
Note that the non-lazy curvature $\kappa_{0}$ of a hypercube is zero. However, hypercubes (as the discrete counterpart of round spheres) are expected to have positive curvature, so the curvature with a nonzero idleness parameter $\alpha$ may serve as a better
representation of curvature. Common choices of $\alpha$ are $\frac{1}{2}$, or $\frac{1}{d+1}$ in case of a $d$ regular graph. For readers' interest, general curvature notions for hypercubes are also discussed by Ollivier and Villani in [OV12].
In fact, Bourne et al. $\left[\mathrm{BCL}^{+} 18\right.$, Theorem 1.1 and Remark 5.4] discovers that for any graph, the $\alpha$-Ollivier Ricci curvature on edges with large enough $\alpha$ is a linear function on $\alpha$, and hence it only differs from the Lin-Lu-Yau curvature by a multiplication factor. This result is extended further for the curvature defined on any pair of different vertices; see [CK19, Corollary 3.4].

Theorem 11.5 ( $\alpha$-ORC and Lin-Lu-Yau). Let $G=(V, E)$ be a graph. Consider two different vertices $x, y \in V$ with vertex degrees $\operatorname{deg}(x)=d_{x}$ and $\operatorname{deg}(y)=d_{y}$.

1. If $x \sim y$, then $\kappa(x, y)=\frac{\kappa_{\alpha}(x, y)}{1-\alpha}$ for all $\alpha \in\left[\frac{1}{\max \left(d_{x}, d_{y}\right)+1}, 1\right)$. In particular, if $d_{x}=d_{y}=d$,

$$
\kappa(x, y)=2 \kappa_{\frac{1}{2}}(x, y)=\frac{d+1}{d} \kappa_{\frac{1}{d+1}}(x, y) .
$$

2. In general, for any $x \neq y$, we have $\kappa(x, y)=\frac{\kappa_{\alpha}(x, y)}{1-\alpha}$ for all $\alpha \in\left[\frac{1}{2}, 1\right)$. In particular,

$$
\begin{equation*}
\kappa(x, y)=2 \kappa_{\frac{1}{2}}(x, y) . \tag{11.1}
\end{equation*}
$$

From this relation between $\kappa_{\alpha}$ and $\kappa$, we see that Lemma 11.2 and Corollary 11.3 also hold with $\kappa$ instead of $\kappa_{\alpha}$.

Another important result is the computation of the curvature for the Cartesian product of two regular graphs. We refer [LLY11, Theorem 3.1] for a proof, and $\left[\mathrm{BCL}^{+} 18\right]$ for a generalized result on $\kappa_{\alpha}$, and [CK19, Theorem 6.2] for an extended result to nonadjacent vertices.

Theorem 11.6 (Cartesian product). Let $G=\left(V_{G}, E_{G}\right)$ be a $d_{G}$-regular graph and $H=\left(V_{H}, E_{H}\right)$ be a $d_{H}$-regular graph. Let $x_{1}, x_{2} \in V_{G}$ with $x_{1} \sim x_{2}$ and $y_{1}, y_{2} \in V_{H}$ with $y_{1} \sim y_{2}$. Then

$$
\begin{aligned}
& \kappa^{G \times H}\left(\left(x_{1}, y_{1}\right),\left(x_{2}, y_{1}\right)\right)=\frac{d_{G}}{d_{G}+d_{H}} \kappa^{G}\left(x_{1}, x_{2}\right), \\
& \kappa^{G \times H}\left(\left(x_{1}, y_{1}\right),\left(x_{1}, y_{2}\right)\right)=\frac{d_{H}}{d_{G}+d_{H}} \kappa^{H}\left(y_{1}, y_{2}\right) .
\end{aligned}
$$

Remark 11.7. By viewing the hypercube $\mathcal{Q}^{n}$ as the Cartesian products of $n$ copies of the complete graph $K_{2}$, one may apply Theorem 11.6 inductively and deduce that the curvature of $\mathcal{Q}^{n}$ is $\kappa(x, y)=\frac{2}{n}$ for all $x \sim y$.

### 11.2 Normalized Laplacian and Lichnerowicz theorem

The normalized graph Laplacian $\Delta$ is defined on a graph $G=(V, E)$ as

$$
\Delta f(x):=\frac{1}{\operatorname{deg}(x)} \sum_{z \in N(x)}(f(z)-f(x)) \quad \forall f \in \mathbb{R}^{V}, \forall x \in V .
$$

In case $G=(V, E)$ is a finite and connected graph, the Laplacian eigenvalues $\lambda \in \mathbb{R}$ solving $\Delta f+\lambda f=0$ are known to be nonnegative real numbers and can be arranged in an increasing order with multiplicities: $0=\lambda_{0}<\lambda_{1} \leq \ldots \leq \lambda_{n-1}$, where $n=|V|$. The trivial eigenvalue $\lambda_{0}=0$ corresponds to the constant eigenfunction, and the second smallest eigenvalue $\lambda_{1}$ is strictly positive as we assume the connectivity of the graph $G$. The next theorem is a discrete analogue to the Lichnerowicz spectral gap theorem on manifolds (c.f. Theorem 5.8 and Proposition 7.9). Here we follow a proof from [LLY11, Theorem 4.2].

Theorem 11.8 (Lichnerowicz for ORC). Let $G=(V, E)$ be a connected graph. Assume that $K:=\inf _{x \neq y} \kappa(x, y)>0$. Then the smallest nonzero eigenvalue satisfies

$$
\lambda_{1} \geq K
$$

Proof. First, we note that $G$ must be a finite graph due to the Bonnet-Myer diameter bound result (see Theorem 12.2). Fix $\alpha \in\left[\frac{1}{2}, 1\right.$ ), and recall from Theorem 11.5 that $\kappa(x, y)=\frac{\kappa_{\alpha}(x, y)}{1-\alpha}$ for all $x \neq y$. We define an average operator $M_{\alpha}: \mathbb{R}^{V} \rightarrow$ $\mathbb{R}^{V}$ as $M_{\alpha} f(x):=\sum_{z \in V} \mu_{x}^{\alpha}(z) f(z)$, which is simplified to

$$
M_{\alpha} f(x)=\alpha f(x)+\sum_{z \in S_{1}(x)} \frac{1-\alpha}{\operatorname{deg}(x)} f(z)=f(x)+(1-\alpha) \Delta f(x)
$$

Let $f_{1}$ be an eigenfunction satisfying $\Delta f_{1}+\lambda_{1} f_{1}=0$. Then we have

$$
\begin{equation*}
M_{\alpha} f_{1}(x)=\left(1-(1-\alpha) \lambda_{1}\right) f_{1}(x) . \tag{11.2}
\end{equation*}
$$

Consider the following Lipschitz constant of $f_{1}$, namely

$$
\ell:=\max _{x, y \in V} \frac{\left|f_{1}(x)-f_{1}(y)\right|}{d(x, y)}
$$

and suppose that the maximum is attained at $\left(x_{1}, y_{1}\right)$. Note that $\ell \neq 0$ since $f_{1}$ is not constant. We then have

$$
\begin{equation*}
\left|M_{\alpha}\left(\frac{f_{1}}{\ell}\right)\left(x_{1}\right)-M_{\alpha}\left(\frac{f_{1}}{\ell}\right)\left(y_{1}\right)\right|=\left|1-(1-\alpha) \lambda_{1}\right| d\left(x_{1}, y_{1}\right) . \tag{11.3}
\end{equation*}
$$

On the other hand, since $\frac{f_{1}}{\ell}$ is 1 -Lipschitz, the Kantorovich duality yields

$$
\begin{align*}
\left|M_{\alpha}\left(\frac{f_{1}}{\ell}\right)\left(x_{1}\right)-M_{\alpha}\left(\frac{f_{1}}{\ell}\right)\left(y_{1}\right)\right| & =\left|\sum_{z \in V} \frac{f(z)}{\ell}\left(\mu_{x_{1}}^{\alpha}(z)-\mu_{y_{1}}^{\alpha}(z)\right)\right| \\
& \leq W_{1}\left(\mu_{x_{1}}^{\alpha}, \mu_{y_{1}}^{\alpha}\right) \\
& =\left(1-\kappa_{\alpha}\left(x_{1}, y_{1}\right)\right) d(x, y) \tag{11.4}
\end{align*}
$$

Comparing (11.4) to (11.3) gives $\lambda_{1} \geq \frac{\kappa_{\alpha}\left(x_{1}, y_{1}\right)}{1-\alpha}=\kappa\left(x_{1}, y_{1}\right) \geq K$ as desired.
Remark 11.9. For a hypercube $\mathcal{Q}^{n}$, the eigenvalues of the adjacency matrix $A_{\mathcal{Q}^{n}}$ are ( $n-2 i$ ) with multipicities $\binom{n}{i}$ for all $0 \leq i \leq n$. The smallest nonzero Laplacian eigenvalue (corresponging to $n-2$ of $A_{\mathcal{Q}^{n}}$ ) is $\lambda_{1}=\frac{2}{n}$, and hence the Lichnerowicz spectral bound is sharp. In fact, the class of all Lichnerowicz sharp graphs are much larger and not yet fully classified. For readers' interest, some partial classification result is provided in [CKK ${ }^{+}$20, Section 6].

## Chapter 12

## Bonnet-Myers diameter bound and Maximal diameter theorem

In this chapter, we present a Bonnet-Myers type diameter bound theorem for Ollivier Ricci curvature. The diameter bound result appears in Ollivier's original work [Oll09, Proposition 23] as well as in the paper by Lin-Lu-Yau [LLY11, Theorem 4.1], and its proof features a simple application of the triangle inequality. In the second half of this chapter, we introduce the concept of the maximal diameter, which will prepare us for the discussion of the rigidity result in the next chapter.

All results in this chapter are stated in terms of the Lin-Lu-Yau curvature $\kappa$, but most proofs use arguments based on $\kappa_{\alpha}$ due to the relation $\kappa_{\alpha}=(1-\alpha) \kappa$ for $\alpha \in\left[\frac{1}{2}, 1\right]$.
Lemma 12.1. Let $G=(V, E)$ be a graph, and let $x, y \in V$ with $x \neq y$. Then

$$
\kappa(x, y) \leq \frac{2}{d(x, y)}
$$

Proof. We aim to prove that $\kappa_{\alpha}(x, y) \leq \frac{2(1-\alpha)}{d(x, y)}$ for any $\alpha$.
Let $\ell:=d(x, y)$ and denote a geodesic in $G$ from $x$ to $y$ by $x_{0} \sim x_{1} \sim \cdots \sim x_{\ell}$, where $x_{0}=x$ and $x_{\ell}=y$.

The triangle inequality for the metric $W_{1}$ gives

$$
\begin{align*}
\ell=W_{1}\left(\delta_{x}, \delta_{y}\right) & \leq W_{1}\left(\delta_{x_{0}}, \mu_{x_{0}}^{\alpha}\right)+W_{1}\left(\mu_{x_{0}}^{\alpha}, \mu_{x_{\ell}}^{\alpha}\right)+W_{1}\left(\mu_{x_{\ell}}^{\alpha}, \delta_{x_{\ell}}\right) \\
& =(1-\alpha)+\ell\left(1-\kappa_{\alpha}(x, y)\right)+(1-\alpha), \tag{12.1}
\end{align*}
$$

where $W_{1}\left(\delta_{v}, \mu_{v}^{\alpha}\right)=1-\alpha$ for any $v \in V$ since the total mass of $1-\alpha$ unit is transported from $v$ to its neighbors).

Theorem 12.2 (Bonnet-Myers). If a locally finite graph $G=(V, E)$ has a positive $\alpha$-Ollivier Ricci lower curvature bound for some $\alpha \in[0,1)$, namely $K_{\alpha}:=$ $\inf _{x \sim y} \kappa_{\alpha}(x, y)>0$, then the diameter of $G$ is bounded above by

$$
\begin{equation*}
\operatorname{diam}(G) \leq \frac{2(1-\alpha)}{K_{\alpha}} \tag{12.2}
\end{equation*}
$$

and hence $G$ is finite.
Similarly, if $G$ has a positive Lin-Lu-Yau lower curvature bound $K:=\inf _{x \sim y} \kappa(x, y)>$ 0 , then

$$
\begin{equation*}
\operatorname{diam}(G) \leq \frac{2}{K} \tag{12.3}
\end{equation*}
$$

The short version of the proof follows from Lemma 12.1 by choosing $x, y \in V$ such that $d(x, y)=\operatorname{diam}(G)$ and using Corollary 11.3. However for later purposes involving transport geodesics, we provide another proof which is direct and does not use the Lemma 12.1.

Proof. For arbitrary $x, y \in V$, let $L:=d(x, y)$ and denote a geodesic in $G$ from $x$ to $y$ by $x_{0} \sim x_{1} \sim \cdots \sim x_{L}$, where $x_{0}=x$ and $x_{L}=y$.
The triangle inequality gives

$$
\begin{align*}
L=W_{1}\left(\delta_{x}, \delta_{y}\right) & \leq W_{1}\left(\delta_{x_{0}}, \mu_{x_{0}}^{\alpha}\right)+\sum_{i=0}^{L-1} W_{1}\left(\mu_{x_{i}}^{\alpha}, \mu_{x_{i+1}}^{\alpha}\right)+W_{1}\left(\mu_{x_{L}}^{\alpha}, \delta_{x_{L}}\right) \\
& \leq(1-\alpha)+L\left(1-K_{\alpha}\right)+(1-\alpha), \tag{12.4}
\end{align*}
$$

where $W_{1}\left(\mu_{x_{i}}^{\alpha}, \mu_{x_{i+1}}^{\alpha}\right)=1-\kappa_{\alpha}\left(x_{i}, x_{i+1}\right) \leq 1-K_{\alpha}$ for all $i$. Rearranging (12.4) yields $L \leq \frac{2(1-\alpha)}{K_{\alpha}}$, which means the diameter of $G$ is no more than $\frac{2(1-\alpha)}{K_{\alpha}}$. A finite diameter bound for a locally finite graph $G$ then implies that $G$ is also finite.
The second statement of the theorem follows immediately due to the relation $\kappa=\frac{\kappa_{\alpha}}{1-\alpha}$ for $\alpha \in\left[\frac{1}{2}, 1\right)$.
Definition 12.3. A graph $G=(V, E)$ is said to be Bonnet-Myers sharp if (12.3) holds with equality, that is,

$$
\inf _{x \sim y} \kappa(x, y)=\frac{2}{\operatorname{diam}(G)}
$$

Equivalently, $G$ is Bonnet-Myers sharp iff (12.2) holds with equality for some (and for all) $\alpha \in\left[\frac{1}{2}, 1\right)$, that is,

$$
\inf _{x \sim y} \kappa_{\alpha}(x, y)=\frac{2(1-\alpha)}{\operatorname{diam}(G)}
$$

Moreover, a graph is called $(D, L)$-Bonnet-Myers sharp if it is $D$-regular and has the diameter $L$.

Hypercubes are examples of Bonnet-Myers sharp graphs. A complete classification of Bonnet-Myers sharp graphs is open, but we present in a full classification under the additional assumption of self-centeredness later in Chapter 13. Hypercubes were also examples of Lichnerowicz sharp graphs. In fact, regular Bonnet-Myers sharp graphs are Lichnerowicz sharp; see [CKK ${ }^{+}$20, Theorem 1.5].
An important characteristic of Bonnet-Myers sharp graphs is related to $W_{1}$-geodesics between the Dirac measures of the two opposite poles as explained in the following theorem.

Theorem 12.4. Let $G=(V, E)$ be a Bonnet-Myers sharp graph and let $x, y \in V$ with $d(x, y)=\operatorname{diam}(G)=: L$. Assume that a sequence of vertices $x=v_{0}, v_{1}, v_{2}, \ldots$, $v_{n-1}, v_{n}=y$ lie orderly on a geodesic in $G$. Then for any $\alpha \in\left[\frac{1}{2}, 1\right)$,
(a) the sequence of probability measures $\delta_{x}, \mu_{x}^{\alpha}, \mu_{v_{1}}^{\alpha}, \ldots, \mu_{v_{n-1}}^{\alpha}, \mu_{y}^{\alpha}, \delta_{y} \in \mathscr{P}(V)$ also lie orderly on a $W_{1}$-geodesic, and
(b) $\kappa_{\alpha}\left(v_{i}, v_{j}\right)=\frac{2(1-\alpha)}{L}$ and $\kappa\left(v_{i}, v_{j}\right)=\frac{2}{L}$ for all $v_{i} \neq v_{j}$.

Additionally, if $G$ is $D$-regular, then (a) and (b) also hold with $\alpha=\frac{1}{D+1}$.
Proof. Being Bonnet-Myers sharp means $\inf _{u \sim v} \kappa_{\alpha}(u, v)=\frac{2(1-\alpha)}{L}$ for all $\alpha \in\left[\frac{1}{2}, 1\right]$. In view of Corollary 11.3, it implies that for all $u, v \in V$ with $u \neq v$,

$$
\frac{2(1-\alpha)}{L} \leq \kappa_{\alpha}(u, v)=1-\frac{W_{1}\left(\mu_{u}^{\alpha}, \mu_{v}^{\alpha}\right)}{d(u, v)}
$$

that is, $W_{1}\left(\mu_{u}^{\alpha}, \mu_{v}^{\alpha}\right) \leq \frac{L-2(1-\alpha)}{L} \cdot d(u, v)$.
Recall also that $W_{1}\left(\delta_{x}, \mu_{x}^{\alpha}\right)=1-\alpha=W_{1}\left(\mu_{y}^{\alpha}, \delta_{y}\right)$. Therefore, by the triangle inequality on $\left(\mathscr{P}(V), W_{1}\right)$, we have

$$
\begin{align*}
L=W_{1}\left(\delta_{x}, \delta_{y}\right) & \stackrel{\Delta}{\leq} W_{1}\left(\delta_{x}, \mu_{x}^{\alpha}\right)+\sum_{i=0}^{n-1} W_{1}\left(\mu_{v_{i}}^{\alpha}, \mu_{v_{i+1}}^{\alpha}\right)+W_{1}\left(\mu_{y}^{\alpha}, \delta_{y}\right) \\
& \leq 2(1-\alpha)+\frac{L-2(1-\alpha)}{L} \sum_{i=0}^{n-1} d\left(v_{i}, v_{i+1}\right)=L \tag{12.5}
\end{align*}
$$

where the last equality comes from the assumption that $\sum_{i=0}^{n-1} d\left(v_{i}, v_{i+1}\right)=d(x, y)=$ $L$. The equality of (12.5) implies that both involved inequalities must hold with
equality. The first one means $\delta_{x}, \mu_{x}^{\alpha}, \mu_{v_{1}}^{\alpha}, \ldots, \mu_{v_{n-1}}^{\alpha}, \mu_{y}^{\alpha}, \delta_{y}$ lie orderly on a $W_{1^{-}}$ geodesic, which yields (a). The second one implies that $W_{1}\left(\mu_{v_{i}}^{\alpha}, \mu_{v_{j}}^{\alpha}\right)=\frac{L-2(1-\alpha)}{L}$. $d\left(v_{i}, v_{j}\right)$ for any $v_{i} \neq v_{j}$. This means $\kappa_{\alpha}\left(v_{i}, v_{j}\right)=\frac{2(1-\alpha)}{L}$ and $\kappa\left(v_{i}, v_{j}\right)=\frac{2}{L}$, which yields (b). Additionally, if $G$ is $D$-regular, then $\inf _{u \sim v} \kappa_{\alpha}(u, v)=\frac{2(1-\alpha)}{L}$ also holds for $\alpha=\frac{1}{D+1}$ due to Theorem 11.5. The statements (a) and (b) are therefore proved with the same arguments.

Our approach to understand more the property of Bonnet-Myers sharp graphs is to apply the concept of transport geodesics (Corollary 10.6 and Definition 10.7) to the $W_{1}$-geodesics obtained from Theorem 12.4. This approach, which led us to our classification result, will be discuss in more details in the next chapter. Let us finish this chapter by presenting the following simple application of the transport geodesic concept: Bonnet-Myers sharp graphs must satisfy the "spherical suspension" property in the sense that every vertex must lie on a geodesic from any pair of two opposite poles (compared to Theorem 4.5). This property has already been proved by Matsumoto in his Master thesis [Mat10] written in Japanese, and it was proved alternatively with a Laplacian method in [CKK ${ }^{+}$20, Theorem 5.5] in the special case of regular graphs. Similar result was proved in the setting of directed graphs in [OSY20, Theorem 1.2].

Theorem 12.5 (maximal diameter). Let $G=(V, E)$ be a Bonnet-Myers sharp graph, i.e., $\inf _{x \sim y} \kappa(x, y)=\frac{2}{\operatorname{diam}(G)}$. Then for any $x, y \in V$ with $d(x, y)=\operatorname{diam}(G)$, we have $[x, y]=V$.

Proof. By connectivity of $G$, it suffices to show that if a vertex $v$ lies in the interval $[x, y]$ then any neighbor of $v$ also lies in $[x, y]$.
Fix any $\alpha \in\left[\frac{1}{2}, 1\right)$. Suppose $v \in[x, y]$, that is, $x, v, y$ lie orderly on a geodesic. Theorem 12.4 implies in particular that $\delta_{x}, \mu_{v}^{\alpha}, \delta_{y}$ lie orderly on a $W_{1}$-geodesic. Consider optimal transport plans $\pi_{1} \in \Pi_{o p t}\left(\delta_{x}, \mu_{v}^{\alpha}\right)$ and $\pi_{2} \in \Pi_{o p t}\left(\mu_{v}^{\alpha}, \delta_{y}\right)$.
For any neighbor of $v$, namely $z\left(\right.$ so $\left.z \in \operatorname{supp}\left(\mu_{v}^{\alpha}\right)\right)$, we know that $\pi_{1}(x, z)>0$ and $\pi_{2}(z, y)>0$. Corollary 10.6 then implies that $x, z, y$ lie on a geodesic, i.e., $z \in[x, y]$ as desired.

## Chapter 13

## Rigidity of Bonnet-Myers sharp graphs

In this chapter, we aim to present the main result in $\left[\mathrm{CKK}^{+} 20\right]$ to find all possible Bonnet-Myers sharp graphs, i.e., satisfying $\operatorname{diam}(G)=\frac{2}{\inf _{x \sim y} \kappa(x, y)}$ (under the additional assumption of self-centeredness).
We restrict our interest to Bonnet-Myers sharp graphs which are regular graphs (i.e., every vertex has the same degree). By considering a $(D, L)$-Bonnet-Myers sharp graph $G=(V, E)$, where $D=\operatorname{deg}(G)$ and $L:=\operatorname{diam}(G)$, the most natural choice of the idleness parameter is $\alpha=\frac{1}{D+1}$, which gives the curvature on an edge $x \sim y$ as

$$
\kappa(x, y)=\frac{D+1}{D} \kappa_{\alpha}(x, y)=\frac{D+1}{D}\left(1-W_{1}\left(\mu_{x}, \mu_{y}\right)\right),
$$

where here and henceforth $\mu_{x}:=\mu_{x}^{\frac{1}{D+1}}$ is the probability measure uniformly distributed in the one-ball $B_{1}(x)$, that is,

$$
\mu_{x}(z)= \begin{cases}\frac{1}{D+1}, & \text { if } z \in B_{1}(x) \\ 0 & \text { otherwise }\end{cases}
$$

Remark 13.1. Due to the fact that all individual masses in $\mu_{x}$ and $\mu_{y}$ are equal to $\frac{1}{D+1}$, there always exists an optimal transport plan $\pi \in \Pi_{o p t}\left(\mu_{x}, \mu_{y}\right)$ without splitting masses. In other words, $\pi$ is induced by a bijective optimal transport map $T: B_{1}(x) \rightarrow B_{1}(y)$, where for each $v \in B_{1}(x)$, the image $T(v)$ is the corresponding vertex in $B_{1}(y)$ such that $\pi(v, T(v))=\frac{1}{D+1}$ (and $\pi(v, z)=0$ for all vertices $z$ other than $T(v)$ ). Moreover, we can assume without loss of generality that we do not need to move masses from origin that already lie in the destination; in other words, $T(v)=v$ for all $v \in B_{1}(x) \cap B_{1}(y)$ (see arguments in, e.g., $\left[\mathrm{BCL}{ }^{+} 18\right.$, Lemma 4.1]).

### 13.1 Transport geodesics of a regular Bonnet-Myers sharp graph

Definition 13.2. Consider a $(D, L)$-Bonnet-Myers sharp graph $G=(V, E)$ and assume a full-length geodesic of $G$, namely,

$$
g: x_{0} \sim x_{1} \sim x_{2} \sim \cdots \sim x_{L} .
$$

For every $1 \leq j \leq L$, consider an optimal transport map $T_{j}$ from $\mu_{x_{j-1}}$ to $\mu_{x_{j}}$ satisfying the following properties:
(P1) $T_{j}: B_{1}\left(x_{j-1}\right) \rightarrow B_{1}\left(x_{j}\right)$ is a bijection.
(P2) $T_{j}(v)=v$ if and only if $v \in B_{1}\left(x_{j-1}\right) \cap B_{1}\left(x_{j}\right)$.
Such a map $T_{j}$ is called a good optimal transport map. The existence of good maps $T_{j}$ is explained in Remark 13.1 (but these maps are not necessarily unique a priori).
Moreover, define maps $T^{j}: B_{1}\left(x_{0}\right) \rightarrow B_{1}\left(x_{j}\right)$ as the composition

$$
T^{j}:=T_{j} \circ \cdots \circ T_{1} \quad \forall 1 \leq j \leq L
$$

and $T^{0}$ is the identity map on $B_{1}\left(x_{0}\right)$ by convention. These maps $T^{j}$ are also bijections. For each $z \in B_{1}\left(x_{0}\right)$, we define $z(0):=z$ and for $1 \leq j \leq L$,

$$
z(j):=T^{j}(z):=T_{j} \circ \cdots \circ T_{1}(z) \in B_{1}\left(x_{j}\right) .
$$

Note in particular that $x_{0}(0)=x_{0}(1)=x_{0}$ by the condition (P2). One may define transport geodesics along the core geodesic $g$ as follows.

Proposition 13.3. Let $G=(V, E)$ be a $(D, L)$-Bonnet-Myers sharp graph with a full-length geodesic $g$ and maps $T_{j}$ and $T^{j}$ defined as in Definition 13.2. Then for every $z \in B_{1}\left(x_{0}\right)$, the vertices $x_{0}, z(0), z(1), z(2), \ldots, z(L), x_{L}$ lie orderly on a geodesic, that is,

$$
\begin{aligned}
L=d\left(x_{0}, x_{L}\right)= & d\left(x_{0}, z(0)\right)+d(z(0), z(1))+d(z(1), z(2))+\ldots+ \\
& d(z(L-1), z(L))+d\left(z(L), x_{L}\right) .
\end{aligned}
$$

This sequence of vertices $x_{0}, z(0), z(1), z(2), \ldots, z(L), x_{L}$ (which contains some repetitions among consecutive vertices) is called a transport geodesic along geodesic $g$.

### 13.2. TRANSPORT GEODESICS OF A SELF-CENTERED BONNET-MYERS SHARP GRAPH99

Proof of Proposition 13.3. We apply Theorem 12.4 in the case of $D$-regular BonnetMyers sharp graph and $\alpha=\frac{1}{D+1}$, and deduce that the sequence of probability measures $\delta_{x_{0}}, \mu_{x_{0}}, \mu_{x_{1}}, \ldots, \mu_{x_{L-1}}, \mu_{x_{L}}, \delta_{x_{L}} \in \mathscr{P}(V)$ lie orderly on a $W_{1}$-geodesic. Note also that along this sequence of probability measures, the mass of $\frac{1}{D+1}$ unit travels in the following chain of optimal transport plan/maps:

$$
x_{0} \mapsto z(0) \stackrel{T_{7}}{\mapsto} z(1) \stackrel{T_{2}}{\mapsto} z(2) \stackrel{T_{3}}{\mapsto} \ldots \stackrel{T_{L}}{\mapsto} z(L) \mapsto x_{L} .
$$

In view of Theorem 10.6 and Definition 10.7, the sequence of vertices $x_{0}, z(0)$, $z(1), z(2), \ldots, z(L), x_{L}$ is then a transport geodesic associated to the probability measures $\delta_{x_{0}}, \mu_{x_{0}}, \mu_{x_{1}}, \ldots, \mu_{x_{L-1}}, \mu_{x_{L}}, \delta_{x_{L}}$.

Corollary 13.4. For any $z \in B_{1}\left(x_{0}\right)$, we have $d\left(z, T_{1}(z)\right) \leq 1$.
Proof. Proposition 13.3 asserts in particular that $x_{0}, z(0), z(1)$ lie on a geodesic; in other words, $d\left(x_{0}, T_{1}(z)\right)=d(x, z)+d\left(z, T_{1}(z)\right)$ for all $z \in B_{1}\left(x_{0}\right)$. In the case of $z=x_{0}$, then $T_{1}\left(x_{0}\right)=x_{0}$, so $d\left(x_{0}, T_{1}\left(x_{0}\right)\right)=0$ as desired. Otherwise, if $z \in S_{1}\left(x_{0}\right)$, then $d\left(z, T_{1}(z)\right)=d\left(x_{0}, T_{1}(z)\right)-1$. Recall further that $T_{1}(z) \in B_{1}\left(x_{1}\right) \subset B_{2}\left(x_{0}\right)$, so $d\left(x_{0}, T_{1}(z)\right) \leq 2$. Therefore, $d\left(z, T_{1}(z)\right) \leq 1$ as desired.

### 13.2 Transport geodesics of a Self-centered BonnetMyers sharp graph

In this section, we assume that our $(D, L)$-Bonnet-Myers sharp graph $G=(V, E)$ has the extra condition that $G$ is self-centered: every vertex is a pole, i.e., for every $x \in V$, there exists $y \in V$ such that $d(x, y)=L$. Such a vertex $y$ is called an antipole of $x$.

As a consequence of $G$ being self-centered, the following proposition asserts that any vertex $v \in B_{1}\left(x_{j-1}\right)$ must be adjacent or equal to the image $T_{j}(v)$ for any $1 \leq j \leq L$.
For a shortened notion, we henceforth write $u \simeq v$ to represent that $u$ is adjacent or equal to $v$.

Proposition 13.5. Let $G=(V, E)$ be a self-centered $(D, L)$-Bonnet-Myers sharp with a full-length geodesic $g$ and maps $T_{j}: B_{1}\left(x_{j-1}\right) \rightarrow B_{1}\left(x_{j}\right)$ defined as in Definition 13.2. For any $1 \leq j \leq L$ and any $v \in B_{1}\left(x_{j-1}\right)$, we have $v \simeq T_{j}(v)$. In particular,
(P2) $v=T_{j}(v)$ if $v \in B_{1}\left(x_{j-1}\right) \cap B_{1}\left(x_{j}\right)$.
(P3) $v \sim T_{j}(v)$ if $v \in B_{1}\left(x_{j-1}\right) \backslash B_{1}\left(x_{j}\right)$.
Proof. Fix $1 \leq j \leq L$. By treating $x_{j-1}$ as a pole together with an antipole $y$ (such that $d\left(x_{j-1}, y\right)=L$ ), it is implied by the Maximal diameter theorem that $\left[x_{j-1}, y\right]=V$. Now since $x_{j} \in\left[x_{j-1}, y\right]$, one can construct a full-length geodesic from $x_{j-1}$ to $y$ passing through $x_{j}$, namely, $g^{\prime}: x_{j-1} \sim x_{j} \sim \ldots \sim y$.
Then Corollary 13.4 (with roles of $x_{0}, z, T_{1}$ replaced by $x_{j-1}, v, T_{j}$, respectively) asserts that $d\left(v, T_{j}(v)\right) \leq 1$, which means $v \simeq T_{j}(v)$. Since the property (P2) was assumed, the property (P3) follows immediately.

In words, $(\mathrm{P} 1)-(\mathrm{P} 3)$ tell us that the bijection $T_{j}: B_{1}\left(x_{j-1}\right) \rightarrow B_{1}\left(x_{j}\right)$ is based on triangles and a perfect matching: that is, $T_{j}$ fixes every vertex $z$ such that $x_{j-1} x_{j} z$ forms a triangle (including $x_{j-1}$ and $x_{j}$ ), and $T_{j}$ forms a perfect matching the other points in its domain and range (i.e., $B_{1}\left(x_{j-1}\right) \backslash B_{1}\left(x_{j}\right)$ and $\left.B_{1}\left(x_{j}\right) \backslash B_{1}\left(x_{j-1}\right)\right)$. This property was proved differently via Laplacian arguments in [CKK ${ }^{+}$20, Theorem 5.7 and Corollary 5.10].
The combination of Propositions 13.3 and 13.5 immediately yields the following result about the transport geodesics along $g$ of a self-centered Bonnet-Myers sharp graph.

Proposition 13.6. Let $G=(V, E)$ be a self-centered $(D, L)$-Bonnet-Myers sharp. Given a full-length geodesic $g$ and maps $T_{j}$ and $T^{j}$ (for $1 \leq j \leq L$ ) defined as above. Then every $z \in B_{1}\left(x_{0}\right)$ induces a transport geodesic along $g$, namely,

$$
x_{0} \simeq z(0) \simeq z(1) \simeq z(2) \simeq \cdots \simeq z(L) \simeq x_{L},
$$

and we denote this geodesic by $g_{z}$.
Remark 13.7. The definition of a transport geodesic $g_{z}$ depends on a full-length geodesic $g$ and sets of transport maps $\left\{T_{j}\right\}_{j=1}^{L}$. Each $T_{j}$ is a priori not uniquely defined, since the definition of $T_{j}$ is based on triangles and a perfect matching, the latter of which is not necessarily unique. We will see later (cf. Remark 13.13) that in fact the maps $\left\{T_{j}\right\}_{j=1}^{L}$ are already uniquely determined by $g$ in the case of self-centered Bonnet-Myers sharp graphs.

### 13.3 Antipoles of intervals in a self-centered BonnetMyers sharp graph

We still assume that our graph $G=(V, E)$ is a self-centered $(D, L)$-Bonnet-Myers sharp graph. Henceforth we will use the following notation related to intervals:

Given an interval $[x, y] \subset V$ in $G$ and a vertex $z \in[x, y]$, we call a vertex $\bar{z} \in[x, y]$ an antipole of $z$ w.r.t. $[x, y]$ if $d(x, y)=d(z, \bar{z})$. Note that antipoles were already introduced for graphs and this definition simply means that $z$ and $\bar{z}$ are antipoles of the induced subgraph of $[x, y]$. We now focus on identifying antipoles w.r.t. intervals via the method of transport geodesics.

Theorem 13.8. Let $G=(V, E)$ be a self-centered ( $D, L$ )-Bonnet-Myers sharp with a full-length geodesic $g: \quad x_{0} \sim x_{1} \sim \cdots \sim x_{L}$. Then for any $2 \leq k \leq L$, $x_{1}$ has a unique antipole w.r.t. the interval $\left[x_{0}, x_{k}\right]$, which we will then denote as $\operatorname{ant}_{\left[x_{0}, x_{k}\right]}\left(x_{1}\right)$. In fact, we show that

$$
\operatorname{ant}_{\left[x_{0}, x_{k}\right]}\left(x_{1}\right)=x_{0}(k)=T^{k}\left(x_{0}\right) \in B_{1}\left(x_{k}\right)
$$

for any fixed $\left\{T_{j}, T^{j}\right\}_{j=1}^{L}$ defined in Definition 13.2.
Proof. First, fix a set of transport maps $\left\{T_{j}, T^{j}\right\}_{j=1}^{L}$ associated to $g$. Suppose that there exists $z \in\left[x_{0}, x_{k}\right]$ which is an antipole of $x_{1}$ w.r.t. $\left[x_{0}, x_{k}\right]$, that is $z \in\left[x_{0}, x_{k}\right]$ and $d\left(x_{1}, z\right)=d\left(x_{0}, x_{k}\right)=k$. Since $x_{1} \sim x_{0}$ and $d\left(x_{1}, z\right)=k$, we have $d\left(x_{0}, z\right) \geq k-1$. Since $z \in\left[x_{0}, x_{k}\right]$ and $z \neq x_{k}$, we must have $d\left(x_{0}, z\right)=k-1$ and $d\left(z, x_{k}\right)=1$. Since $z \in B_{1}\left(x_{k}\right)$, there is a unique $a \in B_{1}\left(x_{0}\right)$ such that $a(k)=z$, that is $a=\left(T^{k}\right)^{-1}(z)$ (because $T^{k}$ is a bijective map). By Proposition 13.6, we obtain a geodesic
as a part of the geodesic $g_{a}$. Therefore it satisfies

$$
\begin{equation*}
k-1=d\left(x_{0}, z\right)=d\left(x_{0}, a(0)\right)+d(a(0), a(1))+d(a(1), z) \tag{13.1}
\end{equation*}
$$

On the other hand, since $d\left(x_{1}, z\right)=k$ and $a(1) \in B_{1}\left(x_{1}\right)$, the triangle inequality gives $d(a(1), z) \geq k-1$. Equation (13.1) then implies $x_{0}=a(0)=a(1)$, which means $a=x_{0}$. Thus $z=a(k)=x_{0}(k)$. So far we have shown that, for every $2 \leq k \leq L, x_{0}(k)$ is the only candidate for an antipole of $x_{1}$ w.r.t. $\left[x_{0}, x_{k}\right]$. It remains to show that $x_{0}(k)$ is in fact the antipole of $x_{1}$ w.r.t $\left[x_{0}, x_{k}\right]$.
In particular, when $k=L$, the antipole of $x_{1}$ w.r.t. $\left[x_{0}, x_{L}\right]=V$ exists by the assumption that $G$ is self-centered. Denote this antipole by $\bar{x}_{1}$ (which obviously differs from $x_{L}$ ). By the previous argument, $x_{0}(L)$ must be $\bar{x}_{1}$.
Consider the transport geodesic $g_{x_{0}}$ and recall that $x_{0}=x_{0}(0)=x_{0}(1)$ and $x_{0}(L)=$ $\bar{x}_{1} \neq x_{L}$ :

$$
\begin{array}{r}
g_{x_{0}}: \quad x_{0}=x_{0}(0)=x_{0}(1) \simeq x_{0}(2) \simeq \cdots \simeq x_{0}(L) \sim x_{L} . \\
\bar{x}_{1}
\end{array}
$$

Computing the length of $g_{x_{0}}$ yields the inequality

$$
L=d\left(x_{0}, x_{L}\right)=\sum_{j=1}^{L-1} d\left(x_{0}(j), x_{0}(j+1)\right)+d\left(x_{0}(L), x_{L}\right) \leq(L-1)+1=L
$$

which holds with equality. It means $x_{0}(j) \sim x_{0}(j+1)$ for all $1 \leq j \leq L-1$.
By removing $x_{L}$ from the geodesic $g_{x_{0}}$, we obtain a new geodesic $x_{0} \sim x_{0}(2) \sim$ $\cdots \sim x_{0}(L)$ of length $L-1$. Next, by extending this geodesic by $x_{1} \sim x_{0}$, we obtain a new sequence of vertices, namely,
Next we consider a new sequence of vertices obtained by removing $x_{L}$ and extending $x_{1} \sim x_{0}$ to the left of the geodesic $g_{x_{0}}$, namely,

$$
g^{\prime}: \quad x_{1} \sim \underbrace{x_{0} \sim x_{0}(2) \sim \cdots \sim x_{0}(L)}_{\text {geodesic of length } L-1} .
$$

Together with the assumption that $d\left(x_{1}, x_{0}(L)\right)=L$, we know $g^{\prime}$ is also a geodesic. Consequently, we can read off from the geodesic $g^{\prime}$ that for every $k \in\{2, \ldots, L\}$, we have

1. $d\left(x_{1}, x_{0}(k)\right)=k$, and
2. $x_{0}(k) \in\left[x_{0}, x_{k}\right]$, because $k=d\left(x_{0}, x_{k}\right) \leq d\left(x_{0}, x_{0}(k)\right)+d\left(x_{0}(k), x_{k}\right) \leq(k-$ 1) +1 .

Therefore, $x_{0}(k)$ is the unique antipole of $x_{1}$ w.r.t. $\left[x_{0}, x_{k}\right]$ as desired.
Let us first discuss an immediate consequence of Theorem 13.8. Note that the theorem implies that there is a well-defined antipole map

$$
\operatorname{ant}_{[x, y]}:[x, y] \cap B_{1}(x) \rightarrow[x, y] \cap B_{1}(y)
$$

The existence and uniqueness of antipoles for neighbours of $x$ w.r.t. $[x, y]$ implies the following result.

Corollary 13.9. Let $G=(V, E)$ be a self-centered Bonnet-Myers sharp graph, $x, y \in V$ be two different vertices. Then the antipole map

$$
\operatorname{ant}_{[x, y]}:[x, y] \cap B_{1}(x) \rightarrow[x, y] \cap B_{1}(y)
$$

is bijective and, consequently,

$$
\left|[x, y] \cap B_{1}(x)\right|=\left|[x, y] \cap B_{1}(y)\right| .
$$

Remark 13.10. Let $x, y \in V$ be two different vertices and $x^{\prime} \in[x, y] \cap B_{1}(x)$ with its antipole $y^{\prime}=\operatorname{ant}_{[x, y]}\left(x^{\prime}\right)$. Observe that then $x, y \in\left[x^{\prime}, y^{\prime}\right]$ and $y=\operatorname{ant}_{\left[x^{\prime}, y^{\prime}\right]}(x)$.

Another immediate consequence of Theorem 13.8 is the following corollary about the intervals of length two.

Corollary 13.11. Let $G=(V, E)$ be a self-centered Bonnet-Myers sharp graph. Then for any $x, y \in V$ such that $d(x, y)=2$, the induced subgraph of $[x, y]$ is a cocktail party graph, that is, for any $z \in[x, y]$, there exists a unique $z^{\prime} \in[x, y]$ such that $d\left(z, z^{\prime}\right)=2$.

Proof. Let $x, y \in V$ with $d(x, y)=2$ and let $z \in[x, y]$. If $z=x$ (or resp. $z=y$ ), it follows immediately from the definition of intervals that such vertex $z^{\prime} \in[x, y]$ with $d\left(z, z^{\prime}\right)=2$ must be $z^{\prime}=y$ (or resp. $z^{\prime}=x$ ). Now we consider the case of $z \notin\{x, y\}$. Let $x=x_{0}$ has an antipole $x_{L}$, and let $z=x_{1}$ and $y=x_{2}$. Now we have $d\left(x_{0}, x_{L}\right)=L, d\left(x_{0}, x_{1}\right)=d\left(x_{1}, x_{2}\right)=1$. By the maximal diameter theorem, $x_{2} \in\left[x_{0}, x_{L}\right]$, that is $d\left(x_{2}, x_{L}\right)=L-2$. Thus $x_{0}, x_{1}, x_{2}, x_{L}$ lie orderly on some geodesic $g$ of length $L$, namely,

$$
\begin{array}{clll}
g: & x_{0} \sim & x_{1} \sim & x_{2} \sim \cdots \sim x_{L} \\
& \| & \| & \| \\
& x & z & y
\end{array}
$$

Applying Theorem 13.8 with $k=2$, we conclude that there is a unique $z^{\prime} \in[x, y]$ satisfying $d\left(z, z^{\prime}\right)=2$ as desired.

Remark 13.12. The fact that all intervals of length two are cocktail party graphs allows us to naturally introduce a switching map, defined as follows. Consider a pair $x, y \in V$ with $d(x, y)=2$. We define $N_{x y}:=S_{1}(x) \cap S_{1}(y)$ to be the set of all common neighbors of $x$ and $y$. Then the switching map $\sigma_{x y}: N_{x y} \rightarrow N_{x y}$ is defined by $\sigma_{x y}(z):=\operatorname{ant}_{[x, y]}(z)$ and satisfies $\sigma_{x y}^{2}=\operatorname{Id}_{N_{x y}}$.
Remark 13.13. Recall from Remark 13.7 that for general Bonnet-Myers sharp graphs the perfect matchings defining the maps $T_{j}$ are not necessarily unique. However, under the additional condition of self-centeredness, the fact that all intervals of length two are cocktail party graphs implies the uniqueness of these perfect matchings and the associated transport maps $T_{j}$. Therefore, the definition of a transport geodesic $g_{z}$ depends only on the geodesic $g$.
In particular, the transport geodesic $g_{x_{0}}$ containing all antipoles of $x_{1}$ w.r.t. increasing intervals $\left[x_{0}, x_{k}\right]$ (see Theorem 13.8) can be also understood as been generated via the following recursive process of switching maps, as illustrated in Figure 13.1:

$$
x_{0}(2)=\sigma_{x_{0} x_{2}}\left(x_{1}\right),
$$

$$
\begin{aligned}
x_{0}(3) & =\sigma_{x_{0}(2) x_{3}}\left(x_{2}\right), \\
\vdots & \\
x_{0}(k) & =\sigma_{x_{0}(k-1) x_{k}}\left(x_{k-1}\right), \\
\vdots & \\
x_{0}(L) & =\sigma_{x_{0}(L-1) x_{L}}\left(x_{L-1}\right) .
\end{aligned}
$$



Figure 13.1: Transport geodesic $g_{x_{0}}$ (along $g$ ) shown in bold

### 13.4 Self-centered Bonnet-Myers sharp graphs are strongly spherical (strategy)

In this section and the next section, we aim to prove a rigidity result that all selfcentered Bonnet-Myers sharp graphs are strongly spherical. Strongly spherical graphs are a special class of graphs which generalize the hypercubes via the combinatorial property about intervals. A complete classification of strongly spherical graphs is given by Koolen, Moulton and Stevanović [KMS04]. This leads to the main result in $\left[\mathrm{CKK}^{+} 20\right]$ about the classification of all self-centered Bonnet-Myers sharp graphs (see Theorem 13.26).

Let us first give the definitions of relevant combinatorial properties of graphs. For a finite connected graph $G=(V, E)$,

- $G$ is self-centered if for every $x \in V$ there exists $\bar{x} \in V$ such that $d(x, \bar{x})=$ $\operatorname{diam}(G)$. The vertex $\bar{x}$ is then called an antipole of $x$.
- $G$ is antipodal if for every $x \in V$ there exists $\bar{x} \in V$ such that $[x, \bar{x}]=V$. The vertex $\bar{x}$ is then called an antipode of $x$.
- $G$ is strongly spherical if $G$ is antipodal, and the induced subgraph of every interval of $G$ is antipodal.

Remark 13.14. It is important to notice the distinction between the notions "antipole" and "antipode". Here are basic facts about antipodes for any graph $G=$ $(V, E)$ :

- Antipodes are also antipoles: Let $\bar{x}$ be an antipode of $x$ in $G$, that is, $[x, \bar{x}]=$ $V$. We choose arbitrary $y, z \in V$ such that $d(y, z)=\operatorname{diam}(G)$. Then we have by $y, z \in[x, \bar{x}]$ and the triangle inequality

$$
\begin{aligned}
& \operatorname{diam}(G) \geq d(x, \bar{x})=\frac{1}{2}(d(x, y)+d(y, \bar{x}))+\frac{1}{2}(d(x, z)+d(z, \bar{x})) \\
& =\frac{1}{2}(d(y, x)+d(x, z))+\frac{1}{2}(d(y, \bar{x})+d(\bar{x}, z)) \geq d(y, z)=\operatorname{diam}(G) .
\end{aligned}
$$

- Antipodes are necessarily unique: Assume $\bar{x}_{1}$ and $\bar{x}_{2}$ are antipodes of $x$. Then $\bar{x}_{2}$ lies on a geodesic from $x$ to $\bar{x}_{1}$. Since $d\left(x, \bar{x}_{1}\right)=d\left(x, \bar{x}_{2}\right)$, this implies $\bar{x}_{1}=\bar{x}_{2}$.

Remark 13.15. The maximal diameter theorem (Theorem 12.5) asserts that every antipole of a Bonnet-Myers sharp graph $G$ must be an antipode. With an additional assumption that $G$ is self-centered, it implies that $G$ is antipodal. We want to show that self-centered BM-sharps are not only antipodal but also "strongly spherical", which mean all intervals are also antipodal.

The following four steps provides an outline of our proof that a self-centered Bonnet-Myers sharp graph $G=(V, E)$ is strongly spherical, i.e., not only that $G$ itself is antipodal but every interval $[x, y]$ of $G$ is also antipodal.

Step 1: Let $x^{\prime} \in[x, y] \cap S_{1}(x)$ with antipole $y^{\prime}=\operatorname{ant}_{[x, y]}\left(x^{\prime}\right)$. We prove for every $z \in[x, y] \cap B_{1}(x)$ that $z \in\left[x^{\prime}, y^{\prime}\right]$ (see Theorem 13.16).

Step 2: Let $x^{\prime} \in[x, y] \cap S_{1}(x)$ with antipole $y^{\prime}=\operatorname{ant}_{[x, y]}\left(x^{\prime}\right)$. We prove for every $z \in[x, y]$ that $z \in\left[x^{\prime}, y^{\prime}\right]$ (see Theorem 13.20).

Step 3: Let $x^{\prime} \in[x, y] \cap S_{1}(x)$ with antipole $y^{\prime}=\operatorname{ant}_{[x, y]}\left(x^{\prime}\right)$. We prove that $[x, y]=\left[x^{\prime}, y^{\prime}\right]$ (see Corollary 13.21).

Step 4: Let $x^{\prime} \in[x, y]$. We prove that there exists $y^{\prime} \in[x, y]$ such that $[x, y]=$ $\left[x^{\prime}, y^{\prime}\right]$ (see Theorem 13.22).

### 13.5 Self-centered Bonnet-Myers sharp graphs are strongly spherical (proof)

Let us now prove the statements in Steps 1-4 in this logical order. Recall that the existence of antipoles of vertices in $[x, y] \cap B_{1}(x)$ w.r.t. $[x, y]$ is guaranteed by Corollary 13.9.

Theorem 13.16. Let $G=(V, E)$ be self-centered Bonnet-Myers sharp. Let $x, y \in$ $V$ be two different vertices, and consider any $x^{\prime} \in[x, y] \cap S_{1}(x)$ with its antipole $y^{\prime}=\operatorname{ant}_{[x, y]}\left(x^{\prime}\right)$. Then every $z \in[x, y] \cap B_{1}(y)$ satisfies $z \in\left[x^{\prime}, y^{\prime}\right]$.

We start with the set-up and introduce particular sets $A, A_{1}, A_{2}, Z, Z_{1}, Z_{2}$ and a function $F$ which will be important for the proof of the above theorem.
Let $k=d(x, y)$. We re-label the vertices as $x=x_{0}$ and $y=x_{k}$ and $x^{\prime}=x_{1}$ and $y^{\prime}=\bar{x}_{1}$, as illustrated in Figure 13.2. Keep in mind that $\bar{x}_{1}=\operatorname{ant}_{\left[x_{0}, x_{k}\right]}\left(x_{1}\right)$ and $x_{0} \sim x_{1}$ and $x_{k} \sim \bar{x}_{1}$.


Figure 13.2: The interval $[x, y]$ with the re-labelled vertices in bold, and $z \in$ $\left[x_{0}, x_{k}\right] \cap B_{1}\left(x_{k}\right)$.

We define the following sets

$$
\begin{array}{ll}
A:=\left[x_{0}, x_{k}\right] \cap B_{1}\left(x_{0}\right), & Z:=\left[x_{0}, x_{k}\right] \cap B_{1}\left(x_{k}\right), \\
A_{1}:=A \cap S_{1}\left(x_{1}\right) \backslash\left\{x_{0}\right\}, & Z_{1}:=Z \cap S_{1}\left(\bar{x}_{1}\right) \backslash\left\{x_{k}\right\}, \\
A_{2}:=A \cap S_{2}\left(x_{1}\right), & Z_{2}:=Z \cap S_{2}\left(\bar{x}_{1}\right) .
\end{array}
$$

Note that the sets $A$ and $Z$ can be partitioned into

$$
A=\left\{x_{0}, x_{1}\right\} \sqcup A_{1} \sqcup A_{2} \quad \text { and } \quad Z=\left\{x_{k}, \bar{x}_{1}\right\} \sqcup Z_{1} \sqcup Z_{2} .
$$

Now fix an arbitrary full-length geodesic $g$ from $x_{0}$ to $x_{L}$ (the antipole of $x_{0}$ ) which passes through $x_{1}$ and $x_{k}$ (this can be done since $x_{1} \in\left[x_{0}, x_{k}\right]$ and $x_{k} \in\left[x_{0}, x_{L}\right]$ due to the maximal diameter theorem: Theorem 12.5), namely

$$
g: \quad x_{0} \sim x_{1} \sim x_{2} \sim \cdots \sim x_{k} \sim x_{k+1} \sim \cdots \sim x_{L} .
$$

Consider the transport map $T^{k}: B_{1}\left(x_{0}\right) \rightarrow B_{1}\left(x_{k}\right)$ introduced in Subsection 13.2. Recall that $T^{k}$ is bijective. Then define a function $F: Z \rightarrow A$ to be $F(z):=$
$\left(T^{k}\right)^{-1}(z)$ for all $z \in Z \subset B_{1}\left(x_{k}\right)$. Lemma 13.17 below guarantees that $F(Z) \subseteq A$, hence $F$ is well-defined.

In order to conclude Theorem 13.16, we need to prove that $\forall z \in Z: z \in\left[x_{1}, \bar{x}_{1}\right]$, which is divided into Lemma 13.18 (dealing with the case $z \in Z_{2} \sqcup\left\{x_{k}, \overline{x_{1}}\right\}$ ) and Lemma 13.19 (dealing with the case $z \in Z_{1}$ ).

Lemma 13.17. $F(Z) \subseteq A$ and $F: Z \rightarrow A$ is bijective.
Lemma 13.18. $F\left(Z_{2} \sqcup\left\{x_{k}, \bar{x}_{1}\right\}\right)=A_{2} \sqcup\left\{x_{0}, x_{1}\right\}$ and $\forall z \in Z_{2} \sqcup\left\{x_{k}, \bar{x}_{1}\right\}: z \in$ $\left[x_{1}, \bar{x}_{1}\right]$.

Lemma 13.19. $F\left(Z_{1}\right)=A_{1}$ and $\forall z \in Z_{1}: \quad z \in\left[x_{1}, \bar{x}_{1}\right]$.
Now we will prove the above three lemmas in order, and then conclude Theorem 13.16.

Proof of Lemma 13.17. First we show that $F(z) \in A$ for all $z \in Z$. Let $a:=$ $F(z) \in B_{1}\left(x_{0}\right)$, that is $a=a(0)$ and $z=a(k)$. By Proposition 13.6, we know that

$$
x_{0} \simeq a(0) \simeq a(1) \simeq \ldots \simeq a(k)
$$

is a geodesic. Moreover, since $a(k)=z \in\left[x_{0}, x_{k}\right]$, this geodesic can be extended to another geodesic $\gamma$, namely

$$
\begin{gather*}
\gamma: \quad x_{0} \simeq a(0) \simeq a(1) \simeq \ldots \simeq a(k) \simeq x_{k} .  \tag{13.2}\\
\| \\
a
\end{gather*}
$$

Therefore, $a$ must lie in the interval $\left[x_{0}, x_{k}\right]$, which means $a \in A$ and we have $F(Z) \subseteq A$.
Next, note that the function $F: Z \rightarrow A$, which is a restriction of $\left(T^{k}\right)^{-1}$, must be injective (because $T^{k}$ is bijective). Note also that $|A|=|Z|$ because of Corollary 13.9. Therefore, $F$ must be bijective.

Proof of Lemma 13.18. A main feature of the following proof is to show $A_{2} \sqcup$ $\left\{x_{0}, x_{1}\right\} \subseteq F\left(Z_{2} \sqcup\left\{x_{k}, \bar{x}_{1}\right\}\right)$. For that reason we start with an element $a \in A_{2} \sqcup$ $\left\{x_{0}, x_{1}\right\}$. Then there exists a uniqe $z \in Z$ with $F(z)=a$. Consequently, $z=a(k)$ and $z \in\left[x_{0}, x_{k}\right]$. Consider the following two cases.
$\underline{\text { Case } a=x_{0}}$ : From Theorem 13.8, we have $a=x_{0}=\left(T^{k}\right)^{-1}\left(\bar{x}_{1}\right)=F\left(\bar{x}_{1}\right)$, so $z=\bar{x}_{1}$ and $z \in\left[x_{1}, \bar{x}_{1}\right]$.

Case $a \in A_{2} \sqcup\left\{x_{1}\right\}$ : As in the proof of Lemma 13.17, we have the following geodesic $\gamma$ of length $k$ (referred to the one in (13.2)):

$$
\gamma: \quad x_{0} \sim a(0) \simeq a(1) \simeq \ldots \simeq a(k) \simeq x_{k} .
$$

From this geodesic $\gamma$ and an observation that

$$
d\left(x_{0}, a(1)\right)= \begin{cases}1, & \text { if } a=x_{1}\left(\text { and therefore also } a(1)=x_{1}\right), \\ 2, & \text { if } a \in A_{2}(\text { and therefore } a(1) \neq a(0)),\end{cases}
$$

we conclude

$$
d\left(a(1), x_{k}\right)= \begin{cases}k-1, & \text { if } a=x_{1},  \tag{13.3}\\ k-2, & \text { if } a \in A_{2}\end{cases}
$$

Now we extend the geodesic

$$
a(1) \simeq \ldots \simeq a(k) \simeq x_{k}
$$

to

$$
\begin{gathered}
x_{1} \simeq a(1) \simeq \ldots \simeq a(k) \simeq x_{k} \sim \bar{x}_{1}, \\
\| \\
z
\end{gathered}
$$

which is, again a geodesic because of (13.3) (and recall that $d\left(x_{1}, \bar{x}_{1}\right)=k$ ). We can then read off from the above geodesic that $z=x_{k}$ or $z \in S_{2}\left(\bar{x}_{1}\right) \cap\left[x_{0}, x_{k}\right]=Z_{2}$ and $z \in\left[x_{1}, \bar{x}_{1}\right]$.
We conclude from both cases that $A_{2} \sqcup\left\{x_{0}, x_{1}\right\} \subseteq F\left(Z_{2} \sqcup\left\{x_{k}, \bar{x}_{1}\right\}\right)$. Since $F$ is bijective, it follows that $\left|A_{2}\right| \leq\left|Z_{2}\right|$. By switching the roles between $x_{0}$ and $x_{k}$ and between the antipoles $x_{1}$ and $\bar{x}_{1}$ w.r.t. $\left[x_{0}, x_{k}\right]$, we obtain the opposite inequality $\left|Z_{2}\right| \leq\left|A_{2}\right|$. Therefore, we have $\left|Z_{2}\right|=\left|A_{2}\right|$, and thus $A_{2} \sqcup\left\{x_{0}, x_{1}\right\}=$ $F\left(Z_{2} \sqcup\left\{x_{k}, \bar{x}_{1}\right\}\right)$, as desired.
Consequently, if we consider any $z \in Z_{2} \sqcup\left\{x_{k}, \bar{x}_{1}\right\}$, then $a \in A_{2} \sqcup\left\{x_{0}, x_{1}\right\}$ falls into one of the above cases, in which we have shown $z \in\left[x_{1}, \bar{x}_{1}\right]$.

Proof of Lemma 13.19. Since $F: Z \rightarrow A$ is bijective and $F\left(Z_{2} \sqcup\left\{x_{k}, \bar{x}_{1}\right\}\right)=$ $A_{2} \sqcup\left\{x_{0}, x_{1}\right\}$, we conclude $F\left(Z_{1}\right)=A_{1}$.
Moreover, consider $z \in Z_{1}$. It follows that $z \sim \bar{x}_{1}$ and $d\left(x_{1}, z\right) \leq k-1$, because $z \neq \bar{x}_{1}=\operatorname{ant}_{\left[x_{0}, x_{k}\right]}\left(x_{1}\right)$. Therefore

$$
d\left(x_{1}, z\right)+d\left(z, \bar{x}_{1}\right)=d\left(x_{1}, z\right)+1 \leq(k-1)+1=k,
$$

which means $z \in\left[x_{1}, \bar{x}_{1}\right]$.

Proof of Theorem 13.16. Recalling the original set-up and notation, we only need to show that $z \in\left[x_{1}, \bar{x}_{1}\right]$. This follows immediately from Lemma 13.18 and Lemma 13.19.

The next theorem generalizes Theorem 13.16 by removing the restriction $z \in$ $B_{1}(y)$.

Theorem 13.20. Let $G=(V, E)$ be self-centered Bonnet-Myers sharp. Let $x, y \in$ $V$ be two different vertices, and consider any $x^{\prime} \in[x, y] \cap S_{1}(x)$ with its antipole $y^{\prime}=\operatorname{ant}_{[x, y]}\left(x^{\prime}\right)$. Then every $z \in[x, y]$ satisfies $z \in\left[x^{\prime}, y^{\prime}\right]$.

Proof of Theorem 13.20. Let $d_{1}=d(x, y)$ and $d_{2}=d(z, y)$ (note that $0 \leq d_{2} \leq d_{1}$ ). We will prove the statement of the theorem by induction on $d_{1}$ and $d_{2}$.

Base step: For any value of $d_{1}$, the cases $d_{2}=0,1$ are both covered by Theorem 13.16.

Inductive step: Assume that the statement is true for $d_{1}=k-1$ and all $d_{2}$, and assume that the statement is true for $d_{1}=k$ and $d_{2}=j-1$ for some $2 \leq$ $j \leq k-1$. Now consider $d(x, y)=k$ and $z \in[x, y] \cap S_{j}(y)$. Choose an arbitrary $z_{1} \in[z, y] \cap S_{1}(y)$. Hence $x, z, z_{1}, y$ lies in a geodesic, see Figure 13.3. In particular, $z \in\left[x, z_{1}\right]$.


Figure 13.3: The interval $[x, y]$ and $z \in[x, y]$ with $d(z, y)=j \geq 2$ and $z_{1} \in$ $[z, y] \cup S_{1}(y)$.

Now consider the following three cases whether $d\left(z_{1}, y^{\prime}\right)$ is 0,1 , or 2 .
Case $z_{1}=y^{\prime}:$ It follows immediately that $z \in\left[x, z_{1}\right]=\left[x, y^{\prime}\right] \subseteq\left[x^{\prime}, y^{\prime}\right]$ where the last inclusion is due to $x \in\left[x^{\prime}, y^{\prime}\right]$.

Case $z_{1} \sim y^{\prime}:$ Since $z_{1} \in[x, y]$, by Theorem 13.8 there is a unique $a_{1}=\operatorname{ant}_{[x, y]}\left(z_{1}\right) \in$ $\overline{[x, y]}$. Since $a_{1} \in[x, y] \cap B_{1}(x)$ and $z_{1} \in[x, y] \cap B_{1}(y)$, by Theorem 13.16, $z_{1}, a_{1} \in$ $\left[x^{\prime}, y^{\prime}\right]$.

The fact that $a_{1}, z_{1} \in\left[x^{\prime}, y^{\prime}\right]$ and that $d\left(a_{1}, z_{1}\right)=d(x, y)=d\left(x^{\prime}, y^{\prime}\right)$ altogether implies that $a_{1}$ must be the unique antipole ant ${ }_{\left[x^{\prime}, y^{\prime}\right]}\left(z_{1}\right)$ by Corollary 13.9 since $z_{1} \sim y^{\prime}$. This is illustrated in Figure 13.4. By Remark 13.10, it implies that $y^{\prime}=\operatorname{ant}_{\left[a_{1}, z_{1}\right]}\left(x^{\prime}\right)$.
Observe also that $z \in\left[x, z_{1}\right] \subset\left[a_{1}, z_{1}\right]$ with $d\left(z, z_{1}\right)=j-1$.
We are now in a position to apply the induction hypothesis for the interval $\left[a_{1}, z_{1}\right]$ (instead of $[x, y]$ ) and $z \in\left[a_{1}, z_{1}\right]$ with $d\left(z, z_{1}\right)=j-1$. Note that $d\left(a_{1}, z_{1}\right)=k$. Note also that $x^{\prime} \in\left[a_{1}, z_{1}\right] \cap S_{1}\left(a_{1}\right)$ and $y^{\prime}=\operatorname{ant}_{\left[a_{1}, z_{1}\right]}\left(x^{\prime}\right)$. Then the induction hypothesis implies $z \in\left[x^{\prime}, y^{\prime}\right]$, finishing this case.


Figure 13.4: Picture for Case $z_{1} \sim y^{\prime}$. The bold cycle represents the fact that $a_{1}$ and $z_{1}$ are antipoles w.r.t. not only $[x, y]$ but also $\left[x^{\prime}, y^{\prime}\right]$.

Case $d\left(z_{1}, y^{\prime}\right)=2$ : Since $z_{1} \in[x, y] \cap B_{1}(y)$, by Theorem 13.16, we have $z_{1} \in\left[x^{\prime}, y^{\prime}\right]$. The condition $d\left(z_{1}, y^{\prime}\right)=2$ then implies that $d\left(x^{\prime}, z_{1}\right)=d\left(x^{\prime}, y^{\prime}\right)-2=k-2$. It follows that

$$
d\left(x, x^{\prime}\right)+d\left(x^{\prime}, z_{1}\right)+d\left(z_{1}, y\right)=1+(k-2)+1=k=d(x, y)
$$

which means that $x^{\prime}$ and $z_{1}$ lie on a geodesic from $x$ to $y$. Let us denote this geodesic by $g^{*}$ :

$$
g^{*}: \quad x \sim x^{\prime} \sim \cdots \sim z_{1} \sim y .
$$

In particular, $x^{\prime} \in\left[x, z_{1}\right]$. Then $y^{\prime \prime}:=\operatorname{ant}_{\left[x, z_{1}\right]}\left(x^{\prime}\right)$ exists by Corollary 13.9. The situation is illustrated in Figure 13.5.
Next we apply the induction hypothesis for the interval $\left[x, z_{1}\right]$ (instead of $[x, y]$ ) and $z \in\left[x, z_{1}\right]$ with $d\left(z, z_{1}\right)=j-1$. Note that $d\left(x, z_{1}\right)=k-1$. Note also that $x^{\prime} \in\left[x, z_{1}\right] \cap S_{1}(x)$ and $y^{\prime \prime}=\operatorname{ant}_{\left[x, z_{1}\right]}\left(x^{\prime}\right)$. Then the induction hypothesis implies $z \in\left[x^{\prime}, y^{\prime \prime}\right]$.
So far we have that $d\left(x^{\prime}, z\right)+d\left(z, y^{\prime \prime}\right)=d\left(x^{\prime}, y^{\prime \prime}\right)=d\left(x, z_{1}\right)=k-1$. It remains to show that $d\left(y^{\prime \prime}, y^{\prime}\right)=1$ which would imply

$$
k=d\left(x^{\prime}, y^{\prime}\right) \leq d\left(x^{\prime}, z\right)+d\left(z, y^{\prime \prime}\right)+d\left(y^{\prime \prime}, y^{\prime}\right)=(k-1)+1=k,
$$



Figure 13.5: Picture for Case $d\left(z_{1}, y^{\prime}\right)=2$. The bold cycle represents $y^{\prime \prime}=$ $\operatorname{ant}_{\left[x, z_{1}\right]}\left(x^{\prime}\right)$.
that is $z \in\left[x^{\prime}, y^{\prime}\right]$, as desired.
To prove $d\left(y^{\prime \prime}, y^{\prime}\right)=1$ we use transport geodesic techniques. Therefore, we relabel the vertices of the geodesic $g^{*}$ and extend $g^{*}$ to a full-length geodesic $g$ in $G$ starting from $x=x_{0}$ as follows:
and consider the transport geodesic along $g$ starting at $x_{0}$.
Theorem 13.8 guarantees that $x_{0}(m)=\operatorname{ant}_{\left[x_{0}, x_{m}\right]}\left(x_{1}\right)$ for all $2 \leq m \leq L$. In particular, we have $y^{\prime \prime}=\operatorname{ant}_{\left[x_{0}, x_{k-1}\right]}\left(x_{1}\right)=x_{0}(k-1)$ and $y^{\prime}=\operatorname{ant}_{\left[x_{0}, x_{k}\right]}\left(x_{1}\right)=x_{0}(k)$. Therefore, $y^{\prime \prime}=x_{0}(k-1)$ and $y^{\prime}=x_{0}(k)$ must be adjacent vertices (as illustrated in Figure 13.6), thus completing the proof.


Figure 13.6: Transport geodesic $g_{x_{0}}$ and the antipoles of $x_{1}$ w.r.t. increasing intervals $\left[x_{0}, x_{m}\right]$.

An immediate but important consequence of the above theorem is the following corollary.

Corollary 13.21. Let $G=(V, E)$ be self-centered Bonnet-Myers sharp. Let $x, y \in$ $V$ be two different vertices, and consider any $x^{\prime} \in[x, y] \cap S_{1}(x)$ with its antipole $y^{\prime}=\operatorname{ant}_{[x, y]}\left(x^{\prime}\right)$. Then $\left[x^{\prime}, y^{\prime}\right]=[x, y]$.

Proof of Corollary 13.21. Theorem 13.20 can be rephrased as $[x, y] \subseteq\left[x^{\prime}, y^{\prime}\right]$. Since $y=\operatorname{ant}_{\left[x^{\prime}, y^{\prime}\right]}(x)$ by Remark 13.10, we can interchange the roles of $x, y$ and $x^{\prime}, y^{\prime}$ to obtain the opposite inclusion $\left[x^{\prime}, y^{\prime}\right] \subseteq[x, y]$. Therefore $\left[x^{\prime}, y^{\prime}\right]=[x, y]$, as desired.

Now, we are ready to conclude the ultimate result of this section by applying Corollary 13.21 inductively.

Theorem 13.22. Let $G=(V, E)$ be self-centered Bonnet-Myers sharp. Then for two different vertices $x, y \in V$, the induced subgraph of the interval $[x, y]$ is antipodal. Therefore $G$ is strongly spherical.

Proof of Theorem 13.22. Let $x^{\prime} \in[x, y]$. We will prove by induction on the distance $d\left(x, x^{\prime}\right)$ that there exists a vertex $y^{\prime} \in[x, y]$ such that $[x, y]=\left[x^{\prime}, y^{\prime}\right]$.
Base step: The case $d\left(x, x^{\prime}\right)=0$ is trivial and $d\left(x, x^{\prime}\right)=1$ is covered by Corollary 13.21.

Inductive step: We assume the statement of the Theorem is true for all $d\left(x, x^{\prime}\right) \leq$ $m-1$ with $2 \leq m \leq \operatorname{diam}(G)$. Let $x^{\prime} \in[x, y]$ with $d\left(x, x^{\prime}\right)=m$. We choose a vertex $x_{1} \in[x, y] \cap S_{1}(x)$ such that $x_{1}, x^{\prime}$ lie on a geodesic from $x$ to $y$. By the induction hypothesis, there exists $y_{1} \in[x, y]$ such that

$$
[x, y]=\left[x_{1}, y_{1}\right] .
$$

Since $d\left(x_{1}, x^{\prime}\right)=m-1$, the induction hypothesis again implies the existence of $y^{\prime} \in\left[x_{1}, y_{1}\right]=[x, y]$ such that

$$
\left[x_{1}, y_{1}\right]=\left[x^{\prime}, y^{\prime}\right]=[x, y] .
$$

This finishes the proof.

### 13.6 Classification of Self-centered Bonnet-Myers sharp graphs

Since we have shown that all regular self-centered Bonnet-Myers sharp graphs must be strongly spherical (Theorem 13.22), the task to find all such graphs is reduced to checking through the list of all the strongly spherical graphs, which is classified by Koolen, Moulton and Stevanović [KMS04] as follows.

Theorem 13.23 ([KMS04]). Strongly spherical graphs are precisely all the Cartesian products of the following graphs:

1. hypercubes $Q^{n}, n \geq 1$;
2. cocktail party graphs $C P(n), n \geq 3$;
3. the Johnson graphs $J(2 n, n), n \geq 3$;
4. even-dimensional demi-cubes $Q_{(2)}^{2 n}, n \geq 3$;
5. the Gosset graph.

For readers' convenience, a brief description of these graphs is given below.

1. The hypercube $Q^{n}$ has $2^{n}$ vertices indexed by $\{0,1\}^{n}$ and edges between them if Hamming distance equals one.
2. The cocktail party graph $C P(n)$ is obtained by removal of a perfect matching from the complete graph $K_{2 n}$;
3. the Johnson graph $J(2 n, n)$ has vertices corresponding to $n$-element subsets of $\{1,2, \ldots, 2 n\}$ and edges between them if they overlap in $n-1$ elements.
4. The even-dimensional demi-cube $Q_{(2)}^{2 n}$ is one of the two isomorphic connected components of the vertex set $\{0,1\}^{2 n}$ and edges between them if Hamming distance equals two.
5. The Gosset graph with 56 vertices: the vertices are in one-to-one correspondence with the edges $\{i, j\}$ and $\{i, j\}^{\prime}$ of two disjoint copies of $K_{8}$, respectively, and $\{i, j\} \sim\{k, l\}$ if $|\{i, j\} \cap\{k, l\}|=1$ and $\{i, j\} \sim\{k, l\}^{\prime}$ if $\{i, j\} \cap\{k, l\}=\emptyset$.

These graphs (with $n$ prescribed as in Theorem 13.23) are non-isomorphic to one another, and they are irreducible (with the exception of the hypercubes $Q^{n}, n \geq 2$ ). Moreover, they are vertex- and edges-transitive and self-centered. Their specific information including the curvature on an edge (see $\left[\mathrm{CKK}^{+} 20\right]$ for calculation details) is listed in the following table:
One can see that $\kappa(x, y)=\frac{2}{L}$ in all examples above, confirming that they are all ( $D, L$ )-Bonnet-Myers sharp. Interestingly, unlike in the case of strongly regular graphs, Bonnet-Myers sharpness is not necessarily preserved under taking any Cartesian products. In order for the product to be Bonnet-Myers sharp, it requires an extra condition that each component has the same ratio of degree and diameter, as explained in the following theorem.

| $G$ | $\|V\|$ | $D=\operatorname{deg}(G)$ | $L=\operatorname{diam}(G)$ | $\kappa(x, y)$ |
| :--- | :--- | :--- | :--- | :---: |
| $Q^{n}$ | $2^{n}$ | $n$ | n | $\frac{2}{n}$ |
| $C P(n)$ | $2 n$ | $2 n-2$ | 2 | 1 |
| $J(2 n, n)$ | $\binom{2 n}{n}$ | $n^{2}$ | n | $\frac{2}{n}$ |
| $Q_{(2)}^{2 n}$ | $2^{2 n-1}$ | $2 n^{2}-n$ | n | $\frac{2}{n}$ |
| Gosset | 56 | 27 | 3 | $\frac{2}{3}$ |

Table 13.1: Examples of $(D, L)$-Bonnet-Myers sharp graphs

Theorem $13.24\left(\left[\mathrm{CKK}^{+} 20\right]\right)$. Let $\left\{G_{i}=\left(V_{i}, E_{i}\right)\right\}_{i=1}^{N}$ be a family of regular graphs where $G_{i}$ has valency $D_{i}$ for each $i$. Let $L_{i}$ be the diameter of $G_{i}$. Let $G=G_{1} \times$ $\cdots \times G_{N}$. The following are equivalent:

1. $G$ is Bonnet-Myers sharp.
2. Each $G_{i}$ is Bonnet-Myers sharp and $\frac{D_{1}}{L_{1}}=\cdots=\frac{D_{N}}{L_{N}}$.

Proof. We may assume without loss of generality that $N=2$ (and the case for $N>2$ can be argued inductively). Recall from Theorem 11.6 that the curvature of the Cartesian product is given by

$$
\begin{align*}
\inf _{\substack{u, v \in V(G) \\
u \sim v}} \kappa^{G}(u, v) & =\min \left\{\inf _{\substack{x_{1}, x_{2} \in V_{1} \\
x_{1} \sim x_{1} \\
y \in V_{2}}} \kappa^{G}\left(\left(x_{1}, y\right),\left(x_{2}, y\right)\right), \inf _{\substack{y_{1}, y_{2} \in V_{2} \\
y_{1} \sim y_{2} \\
x \in V_{1}}} \kappa^{G}\left(\left(x, y_{1}\right),\left(x, y_{2}\right)\right)\right\} \\
& =\min \left\{\inf _{\substack{x_{1}, x_{2} \in V_{1} \\
x_{1} \sim x_{2}}} \frac{D_{1}}{D_{1}+D_{2}} \kappa^{G_{1}}\left(x_{1}, x_{2}\right), \inf _{\substack{y_{1}, y_{2} \in V_{2} \\
y_{1} \sim y_{2}}} \frac{D_{2}}{D_{1}+D_{2}} \kappa^{G_{2}}\left(y_{1}, y_{2}\right)\right\} . \tag{13.4}
\end{align*}
$$

First we prove that (i) implies (ii). Since $G$ is Bonnet-Myers sharp, togther with the Bonnet-Myers diameter bound on $G_{1}$, we have

$$
\frac{2}{L_{1}+L_{2}}=\inf _{\substack{u, v \in V(G) \\ u \sim v}} \kappa^{G}(u, v) \leq \inf _{\substack{x_{1}, x_{2} \in V_{1} \\ x_{1} \sim v_{2}}} \frac{D_{1}}{D_{1}+D_{2}} \kappa^{G_{1}}\left(x_{1}, x_{2}\right) \leq \frac{D_{1}}{D_{1}+D_{2}} \cdot \frac{2}{L_{1}},
$$

which is equivalent to $\frac{D_{2}}{L_{2}} \leq \frac{D_{1}}{L_{1}}$.
On the other hand, the Bonnet-Myers diameter bound on $G_{2}$ gives

$$
\frac{2}{L_{1}+L_{2}}=\inf _{\substack{u, v \in V(G) \\ u \sim v}} \kappa^{G}(u, v) \leq \inf _{\substack{y_{1}, y_{2} \in V_{2} \\ y_{1} \not y_{2}}} \frac{D_{2}}{D_{1}+D_{2}} \kappa^{G_{2}}\left(y_{1}, y_{2}\right) \leq \frac{D_{2}}{D_{1}+D_{2}} \cdot \frac{2}{L_{2}}
$$

which is equivalent to $\frac{D_{1}}{L_{1}} \leq \frac{D_{2}}{L_{2}}$.
Therefore, we can conclude that $\frac{D_{1}}{L_{1}}=\frac{D_{2}}{L_{2}}$, and all the inequalities above are sharp, that is $G_{1}$ and $G_{2}$ are Bonnet-Myers sharp as well.

To prove (ii) implies (i): we simply plug into (13.4)

$$
\inf _{\substack{x_{1}, x_{2} \in V_{1} \\ x_{1} \sim x_{2}}} \kappa^{G_{1}}\left(x_{1}, x_{2}\right)=\frac{2}{L_{1}} \quad \text { and } \quad \inf _{\substack{y_{1}, y_{2} \in V_{2} \\ y_{1} y_{2}}} \kappa^{G_{2}}\left(y_{1}, y_{2}\right)=\frac{2}{L_{2}}
$$

and use the assumption that $\frac{D_{1}}{L_{1}}=\frac{D_{2}}{L_{2}}$. As a result, we obtain $\inf _{\substack{u, v \in V(G) \\ u \sim v}} \kappa^{G}(u, v)=$ $\frac{2}{L_{1}+L_{2}}$.
Remark 13.25. In contrast to the necessary and sufficient condition for BonnetMyers sharpness in Theorem 13.24, much less is required for the Cartesian product $G=G_{1} \times G_{2}$ to be Lichnerowicz sharp. In fact, $G$ is Lichnerowicz sharp already if $G_{1}$ is Lichenerowicz sharp and $G_{2}$ is an arbitrary graph with its curvature lower bound large enough, as explained in the following argument.
Let $\lambda_{1}^{G_{1}}, \lambda_{1}^{G_{2}}, \lambda_{1}^{G}$ be the smallest positive eigenvalues of the Laplacians on $G_{1}, G_{2}$, $G$. We have

$$
\begin{align*}
\inf _{\substack{u, v \in V(G) \\
u \sim v}} \kappa^{G}(u, v) & =\min \left\{\inf _{\substack{x_{1}, x_{2} \in V_{1} \\
x_{1} \sim x_{2}}} \frac{D_{1}}{D_{1}+D_{2}} \kappa^{G_{1}}\left(x_{1}, x_{2}\right), \inf _{\substack{y_{1}, y_{2} \in V_{2} \\
y_{1} \sim y_{2}}} \frac{D_{2}}{D_{1}+D_{2}} \kappa^{G_{2}}\left(y_{1}, y_{2}\right)\right\} \\
& \leq \min \left\{\frac{D_{1}}{D_{1}+D_{2}} \lambda_{1}^{G_{1}}, \frac{D_{2}}{D_{1}+D_{2}} \lambda_{2}^{G_{2}}\right\}=\lambda_{1}^{G} \tag{13.5}
\end{align*}
$$

where the inequality comes from Lichnerowicz' Theorem on each graph $G_{i}$ : $\inf \kappa^{G_{i}} \leq$ $\lambda_{1}^{G_{i}}$. In order to obtain the equality in (13.5), a sufficient condition is inf $\kappa^{G_{1}}=\lambda_{1}^{G_{1}}$ (i.e. $G_{1}$ is Lichnerowicz sharp) and $\frac{D_{1}}{D_{2}} \inf \kappa^{G_{1}} \leq \inf \kappa^{G_{2}}$.

Finally, we can combine Theorems 13.22, 13.23 and 13.24 to conclude the complete classification of regular self-centered Bonnet-Myers sharp graphs.

Theorem 13.26 (main result, $\left[\mathrm{CKK}^{+} 20\right]$ ). Regular self-centered Bonnet-Myers sharp graphs are precisely the Cartesian products $G=G_{1} \times G_{2} \times \cdots \times G_{N}$, where each factor $G_{i}$ is any of the following graphs:

1. hypercubes $Q^{n}, n \geq 1$;
2. cocktail party graphs $C P(n), n \geq 3$;
3. the Johnson graphs $J(2 n, n), n \geq 3$;
4. even-dimensional demi-cubes $Q_{(2)}^{2 n}, n \geq 3$;
5. the Gosset graph,
and satisfies the condition $\frac{D_{1}}{L_{1}}=\frac{D_{2}}{L_{2}}=\cdots=\frac{D_{N}}{L_{N}}$.

## Chapter 14

## Outlook for Ollivier Ricci curvature

Let us conclude Part II with some interesting questions, especially about Cheng's and Obata's rigidity for Ollivier Ricci curvature.

1. While the main classification result of Bonnet-Myers sharp graphs (Theorem 13.26) assumes the graphs in consideration to be regular and self-centered, it is a natural question to ask if we could drop either or both of these assumptions. From an optimistic viewpoint, we conjecture that all regular Bonnet-Myers sharp graphs must be self-centered. This conjecture holds in the extremal cases where the graph diameter is equal to two (in this case, the graph has to be a Cocktail party graph) or to the vertex degree (in this case, the graph has to be a hypercube); for the proof see [CKK $\left.{ }^{+} 20\right]$. Moreover, in [Kam20a] we confirm this conjecture for all such graphs of diameter three.
2. It is shown in $\left[\mathrm{CKK}^{+} 20\right]$ that Bonnet-Myers sharp graphs are Lichnerowicz sharp, and the class of all Lichnerowicz sharp graphs is much larger. The paper also presents some partial results about the classification of Lichnerowicz sharp graphs which are distance-regular with a certain value of the second largest adjacency eigenvalue. It would be interesting to find further classification results of Lichnerowicz sharp graphs.
3. It was brought to our attention by Paul Terwilliger that the classification found for the Bonnet-Myers sharp graphs agrees with the class of graphs related hypermetric spaces mentioned in [TD87]. This agreement is likely not to be coincidental, and there could potentially be some connection between these two classes.

## Part III

## Bakry-Émery curvature on graphs

## Chapter 15

## Discrete Bakry-Émery theory

The application of the Bakry and Émery theory to define the Ricci curvature notion for discrete spaces was initiated and promoted by, e.g., Elworthy [Elw91], Schmuckenschläger [Sch99] and Lin-Yau [LY10]. In this chapter, we revisit the $\Gamma$-calculus and the definition Bakry-Émery curvature which were previously given in Chapter 5 but are now discussed in the context of weighted graphs.

### 15.1 Setting of weighted graphs and Laplacians

Let $G=(V, w, \mu)$ be a weighted graph consisting of a countable vertex set $V$, a vertex measure $\mu: V \rightarrow \mathbb{R}^{+}$, and an edge-weight function $w: V \times V \rightarrow \mathbb{R}_{\geq 0}$ which is a symmetric function with $w_{x x}=0$ for all $x \in V$. Two vertices $x, y \in V$ are adjacent (denoted as $x \sim y$ ) if and only if $w_{x y}>0$. The graph $G$ is assumed to be connected, and locally finite, that is, each vertex has only finitely many neighbors. For $r \in \mathbb{N}$, the sphere (resp. ball) of radius $r$ centered at $x \in V$, denoted by $S_{r}(x)$ (resp. $B_{r}(x)$ ), is the set of all vertices whose minimum number of edges from $x$ is equal to (resp. less than or equal to) $r$. In particular, $S_{1}(x)$ contains all neighbors of $x$.
Furthermore, let $p_{x y}:=\frac{w_{x y}}{\mu_{x}}$ be the transition rate from $x$ to $y$. Let $\operatorname{Deg}(x):=$ $\frac{1}{\mu_{x}} \sum_{y \in V} w_{x y}=\sum_{y \in V} p_{x y}$ be the vertex degree of $x$, and mostly we assume maximal degree to be finite, i.e., $\operatorname{Deg}_{\max }:=\sup _{x \in V} \operatorname{Deg}(x)<\infty$. In the special case of $\operatorname{Deg}(x)=1$ for all $x \in V$, the terms $p_{x y}$ can be understood as transition probabilities of a reversible Markov chain (which is the topic of discussion in Chapter 20). Another special situation is a non-weighted (or combinatorial) graph $G=(V, E)$ where $E$ is the set of edges (without loops and multiple edges), that is, $\mu \equiv 1$ and $w_{x y}=1$ iff $x$ is adjacent to $y$, and $w_{x y}=0$ otherwise. In this
case, $\operatorname{Deg}(x)=\sum_{y \in V} w_{x y}$ equals the number of neighbors of $x$, which is the usual definition of the vertex degree.
The Laplacian $\Delta: \mathbb{R}^{V} \rightarrow \mathbb{R}^{V}$ (where $\mathbb{R}^{V}$ is the vector space of all functions $f: V \rightarrow \mathbb{R}$ ) is a linear operator on $V$ given by

$$
\begin{aligned}
\Delta f(x):=\frac{1}{\mu_{x}} \sum_{y \in V} w_{x y}(f(y)-f(x)) & =\sum_{y \in V} p_{x y}(f(y)-f(x)) \\
& =\sum_{y \in S_{1}(x)} p_{x y}(f(y)-f(x)),
\end{aligned}
$$

where the summation can be restricted to only those $y \in S_{1}(x)$ since by definition $x \sim y$ if and only if $w_{x y}>0$. The Laplacian associated to non-weighted graphs is also known as the non-normalized Laplacian.

## 15.2 $\Gamma$-calculus on graphs

The Laplacian $\Delta$ gives rise to the symmetric bilinear forms $\Gamma$ and $\Gamma_{2}$, namely,

$$
\begin{align*}
2 \Gamma(f, g) & :=\Delta(f g)-f \Delta g-g \Delta f  \tag{15.1}\\
2 \Gamma_{2}(f, g) & :=\Delta(\Gamma(f, g))-\Gamma(f, \Delta g)-\Gamma(g, \Delta f)
\end{align*}
$$

for all $f, g \in \mathbb{R}^{V}$. Moreover, we write $\Gamma f:=\Gamma(f, f)$ and $\Gamma_{2} f:=\Gamma_{2}(f, f)$.
In words, $\Gamma$ represents how much the Laplacian is far away from the usual chain rule for derivative $D(f g)=f D g+g D f$, and $\Gamma_{2}$ is an iteration from $\Gamma$ where the product of two functions $f g$ are replaced by $\Gamma(f, g)$.
Some basic properties of $\Delta$ and $\Gamma$ on graphs are given in the following proposition.
Proposition 15.1. Let $G=(V, w, \mu)$ be a weighted graph. For all $f, g \in \mathbb{R}^{V}$ and $x \in V$, we have
(L1) $2 \Gamma(f, g)(x)=\sum_{y \in S_{1}(x)} p_{x y}(f(y)-f(x))(g(y)-g(x))$, and particularly, $2 \Gamma f(x)=$ $\sum_{y \in S_{1}(x)} p_{x y}(f(y)-f(x))^{2} ;$
(L2) $(\Delta f(x))^{2} \leq 2 \operatorname{Deg}(x) \Gamma f(x)$;
(L3) If $G$ is finite, then $\langle\Delta f, g\rangle=\langle f, \Delta g\rangle$, where the norm is denoted by $\langle f, g\rangle:=\sum_{x \in V} \mu_{x} f(x) g(x)$. In particular, $\langle\Delta f, 1\rangle=$ $\sum_{x \in V} \mu_{x} \Delta f(x)=0$.

Remark 15.2. The properties (L1) and (L3) are discrete analogues of properties in Riemannian manifolds. With an additional notion of discrete gradient given by $\nabla f(x, y):=f(y)-f(x)$, the property (L1) can be written as $2 \Gamma(f, g)=\langle\nabla f, \nabla g\rangle$ and $2 \Gamma f=|\nabla f|^{2}$, which is a discrete analogue to (5.11) (but notice the factor 2 discrepancy). The property (L3) means $\Delta$ is self-adjoint on $L^{2}$-norm, and its consequence $\sum_{x \in V} \mu_{x} \Delta f(x)=0$ is a discrete version of the divergence theorem for closed manifolds. The property (L2), on the other hand, is a specific property in the discrete setting of graphs and does not have a counterpart in Riemannian manifolds (as one should not expect a bound like $|\Delta f| \leq C|\operatorname{grad} f|^{2}$ since the Laplacian is a second-order differential operator, while gradient is only a firstorder differential operator).

Proof of Proposition 15.1. (L1) A straightforward calculation gives

$$
\begin{aligned}
2 \Gamma(f, g)(x)= & \Delta(f g)(x)-f(x) \cdot \Delta g(x)-\Delta f(x) \cdot g(x) \\
= & \sum_{y \in S_{1}(x)} p_{x y}[f(y) g(y)-f(x) g(x) \\
& \quad-f(x)(g(y)-g(x))-g(x)(f(y)-f(x))] \\
= & \sum_{y \in S_{1}(x)} p_{x y}(f(y)-f(x))(g(y)-g(x)),
\end{aligned}
$$

where the identity $2 \Gamma f(x)=\sum_{y \in S_{1}(x)} p_{x y}(f(y)-f(x))^{2}$ follows immediately.
(L2) We apply Cauchy-Schwarz inequality to (L1) and obtain

$$
\begin{aligned}
(\Delta f(x))^{2} & =\left(\sum_{y \in S_{1}(x)} p_{x y}(f(y)-f(x))\right)^{2} \\
& \leq\left(\sum_{y \in S_{1}(x)} p_{x y}(f(y)-f(x))^{2}\right) \cdot \sum_{y \in S_{1}(x)} p_{x y} \stackrel{(L 1)}{=} 2 \operatorname{Deg}(x) \Gamma f(x) .
\end{aligned}
$$

(L3) Suppose $G$ is finite. Then

$$
\begin{aligned}
\sum_{x \in V} \mu_{x} \Delta f(x) g(x) & =\sum_{x} \sum_{y} w_{x y}(f(y)-f(x)) g(x) \\
& =\sum_{x} \sum_{y} w_{x y} f(y) g(x)-\sum_{x} \sum_{y} w_{x y} f(x) g(x) \\
& =\sum_{y} \sum_{x} w_{x y} f(x) g(y)-\sum_{x} \sum_{y} w_{x y} f(x) g(x)
\end{aligned}
$$

$$
=\sum_{x \in V} \mu_{x} f(x) \Delta g(x) .
$$

### 15.3 Bakry-Émery curvature in local matrix form

The previous section has prepared us with the discrete version of bilinear operators $\Gamma$ and $\Gamma_{2}$. Now we are ready to present in the context of weighted graphs, the definition of Bakry-Émery curvature, which uses $\Gamma$ and $\Gamma_{2}$ as the key ingredients (cf. Section 5.4).
Definition 15.3 (Bakry-Émery curvature). Let $G=(V, w, \mu)$ be a locally finite weighted graph. Let $K \in \mathbb{R}$ and $N \in(0, \infty]$. We say that a vertex $x \in V$ satisfies $B E(K, N)$ if the following Bakry-Émery's curvature-dimension inequality holds for all $f \in \mathbb{R}^{V}$ :

$$
\begin{equation*}
\Gamma_{2} f(x) \geq K \Gamma f(x)+\frac{1}{N}(\Delta f(x))^{2} \tag{15.2}
\end{equation*}
$$

The Bakry-Émery curvature at a vertex $x$, denoted by $\mathcal{K}(G, x ; N)$, is then defined to be the largest $K$ such that $x$ satisfies $B E(K, N)$. Moreover, we say that $G$ satisfies $B E(K, N)$ if all $x \in V$ satisfy $B E(K, N)$.

Most of the times, saying that "a graph $G$ satisfies $B E(K, N)$ " is almost the same as $K=\inf _{x \in V} \mathcal{K}(G, x ; N)$, except that in the former statement $K$ does not need to be maximum possible one. We may use both of them interchangeably to express the global lower bound of Bakry-Émery curvature.
Here, the parameter $N$ acts as an upper bound for "dimension" of the weighted graph, and $K$ acts as a "Ricci lower curvature bound" at $x$. There is no unified notions of dimension and Ricci curvature for graphs, so this convention is only to be compared with Theorem 5.12 for manifolds.
Remark 15.4. Both $\Delta f(x)$ and $\Gamma(f, g)(x)$ only involves the information within $B_{1}(x)$ (i.e., neighbors of $x$ including $x$ itself). The iterated $\Gamma_{2}(f, g)$ involves the information within $B_{2}(x)$ (i.e., all vertices with distance $\leq 2$ from $\left.x\right)$. Therefore, the test functions for (15.2) can be restricted to those $f: B_{2}(x) \rightarrow \mathbb{R}$.
Remark 15.5. If $x$ satisfies $B E(K, N)$ then $x$ also satisfies $B E\left(K^{\prime}, N^{\prime}\right)$ for all $K^{\prime} \leq K$ and all $N^{\prime} \geq N$ (including $N^{\prime}=\infty$ ) because $\Gamma f(x) \geq 0$ for all $f$. Equivalently, the curvature at $x$ is given by

$$
\mathcal{K}(G, x ; N)=\inf _{f} \frac{\Gamma_{2} f(x)-\frac{1}{N}(\Delta f(x))^{2}}{\Gamma f(x)}
$$

where the infimum is taken over all functions $f \in \mathbb{R}^{V}$ which are non-constant on $B_{1}(x)$ and therefore $\Gamma f(x)>0$.

Alternatively, it is easier to view the involved operators in (15.2) by their matrix representations. The linear operator $\Delta(\cdot)(x)$ and the bilinear forms $\Gamma(\cdot, \cdot)(x), \Gamma_{2}(\cdot, \cdot)(x)$ can be represented by a local vector $\Delta(x)$ and local matrices $\Gamma(x), \Gamma_{2}(x)$ as

$$
\begin{aligned}
\Delta f(x) & =\Delta(x)^{\top} \vec{f}, \\
\Gamma(f, g)(x) & =\vec{f}^{\top} \Gamma(x) \vec{g}, \\
\Gamma_{2}(f, g)(x) & =\vec{f}^{\top} \Gamma_{2}(x) \vec{g} .
\end{aligned}
$$

These vector and matrices are "local" in the sense that its entries are only indexed by vertices near $x$ instead the whole graph. More precisely, in the first two equations above, $\vec{f}$ and $\vec{g}$ are vector representations indexed by vertices in $B_{1}(x)$ as $\vec{f}=\left(\begin{array}{llll}f(x) & f\left(y_{1}\right) & \cdots & f\left(y_{m}\right)\end{array}\right)^{\top}$, where $S_{1}(x)=\left\{y_{1}, \ldots, y_{m}\right\}$. In the third equation, the vectors are indexed by vertices in $B_{2}(x)$ as

$$
\vec{f}=\left(\begin{array}{lllllll}
f(x) & f\left(y_{1}\right) & \cdots & f\left(y_{m}\right) & f\left(z_{1}\right) & \cdots & f\left(z_{n}\right)
\end{array}\right)^{\top},
$$

where $S_{2}(x)=\left\{z_{1}, \ldots, z_{n}\right\}$.
The Bakry-Émery curvature $\mathcal{K}(G, x ; N)$ can then be reformulated as the solution to the following semidefinite programming:

$$
\begin{align*}
& \operatorname{maximize} K  \tag{P}\\
& \text { subject to } \Gamma_{2}(x)-\frac{1}{N} \Delta(x) \Delta(x)^{\top}-K \Gamma(x) \succeq 0 .
\end{align*}
$$

Here $M \succeq 0$ (resp. $M \succ 0$ ) means $M$ is positive semidefinite (resp. positive definite). Note also that $\Gamma_{2}(x)-\frac{1}{N} \Delta(x) \Delta(x)^{\top}-K \Gamma(x)$ is a symmetric metric due to the symmetric bilinear forms $\Gamma$ and $\Gamma_{2}$ at $x$. The above computing method has been studied by Schmuckenschläger [Sch99], and later on in e.g. [CLP20].

## Chapter 16

## Curvature as a minimal eigenvalue problem

The goal of this chapter is to present one of the main results from [CKLP21] which gives a reformulation of Bakry-Émery curvature $\mathcal{K}(G, x ; N)$ (previously given by the semidifinite problem $(P)$ ) as a smallest eigenvalue problem by employing the Schur complement of a square block matrix $M_{22}$ in $M=\left(\begin{array}{ll}M_{11} & M_{12} \\ M_{21} & M_{22}\end{array}\right)$, namely $M / M_{22}:=M_{11}-M_{12} M_{22}^{-1} M_{21}$, applied to the matrix

$$
\Gamma_{2}(x)_{\hat{1}}=\left(\begin{array}{ll}
\Gamma_{2}(x)_{S_{1}, S_{1}} & \Gamma_{2}(x)_{S_{1}, S_{2}} \\
\Gamma_{2}(x)_{S_{2}, S_{1}} & \Gamma_{2}(x)_{S_{2}, S_{2}}
\end{array}\right) .
$$

Here, the matrix $\Gamma_{2}(x)_{\hat{1}}$ refers to the principle submatrix of $\Gamma_{2}(x)$ obtained by removing its first row and column corresponding to the vertex $x$. The matrix $\Gamma_{2}(x)_{S_{i}, S_{j}}$ refers to the submatrix of $\Gamma_{2}(x)$ whose rows and columns are indexed by the vertices of the combinatorial spheres $S_{i}(x)$ and $S_{j}(x)$.
We use the notation $Q(x):=\Gamma_{2}(x)_{\hat{1}} / \Gamma_{2}(x)_{S_{2}, S_{2}}$ for simplicity, and define

$$
\begin{align*}
A_{\infty}(x) & :=2 \operatorname{diag}\left(\mathbf{v}_{0}(x)\right)^{-1} Q(x) \operatorname{diag}\left(\mathbf{v}_{0}(x)\right)^{-1} \\
A_{N}(x) & :=A_{\infty}(x)-\frac{2}{N} \mathbf{v}_{0}(x) \mathbf{v}_{0}(x)^{\top} \tag{16.1}
\end{align*}
$$

where $\mathbf{v}_{0}(x):=\left(\sqrt{p_{x y_{1}}} \sqrt{p_{x y_{2}}} \ldots \sqrt{p_{x y_{m}}}\right)^{\top}$ with $S_{1}(x)=\left\{y_{1}, y_{2}, \ldots, y_{m}\right\}$ labelling the neighbours of $x$. Note that the matrices $Q(x), A_{\infty}(x), A_{N}(x)$ are all symmetric matrices, and that $A_{N}(x)$ is a rank one perturbation of $A_{\infty}(x)$. The formulation of Bakry-Émery curvature can be summarized as follows.

Theorem 16.1. Let $G=(V, w, \mu)$ be a weighted graph. For $x \in V$ and $N \in(0, \infty]$, the Bakry-Émery curvature $\mathcal{K}(G, x ; N)$ is the smallest eigenvalue of the symmetrix matrix $A_{N}(x)$ defined in (16.1), that is,

$$
\mathcal{K}(G, x ; N)=\lambda_{\min }\left(A_{N}(x)\right) .
$$

### 16.1 Curvature matrix $A_{\infty}$

This section is dedicated to the proof from [CKLP21] of Theorem 16.1 that the Bakry-Émery curvature at a vertex $x$, namely

$$
\mathcal{K}(G, x ; N)=\underset{K}{\arg \max }\left\{\Gamma_{2}(x)-\frac{1}{N} \Delta(x) \Delta(x)^{\top}-K \Gamma(x) \succeq 0\right\}
$$

is equal to the smallest eigenvalue of the so-called curvature matrix $A_{\infty}(x)$.
First, let us recall that $\mathcal{K}(G, x ; N)$ is a local concept and uniquely determined by the structure of the two-ball $B_{2}(x)$. In particular, the symmetric matrix $\Gamma_{2}(x)$ is of size $\left|B_{2}(x)\right|$, and the symmetric matrices $\Delta(x) \Delta(x)^{\top}$ and $\Gamma(x)$ are of sizes $\left|B_{1}(x)\right|$ (and trivially extended by zeros to matrices of sizes $\left|B_{2}(x)\right|$ ).
Schmuckenschläger [Sch99] observes that the size of these matrices can be reduced by one: since $\Gamma_{2}(f), \Gamma(f), \Delta f$ all vanish for constant functions $f$, the curvaturedimension inequality $B E(K, N)$ remains valid after shifting $f$ by an additive constant. It is therefore sufficient to verify (15.2) for all functions $f: V \rightarrow \mathbb{R}$ with $f(x)=0$. This observation allows us remove from these matrices the row and column corresponding to the vertex $x$, and we are able to reformulate the problem $(P)$ as
maximize $K$
$\left(P^{\prime}\right)$
subject to $M_{K, N}(x):=\left(\Gamma_{2}(x)-\frac{1}{N} \Delta(x) \Delta(x)^{\top}-K \Gamma(x)\right)_{S_{1} \cup S_{2}, S_{1} \cup S_{2}} \succeq 0$.
Next we recall the concept of the Schur complement, which allows us to further reduce the size of the involved symmetric matrices in $\left(P^{\prime}\right)$.
Lemma 16.2. Consider a real symmetric matrix $M=\left(\begin{array}{ll}M_{11} & M_{12} \\ M_{21} & M_{22}\end{array}\right)$, where $M_{11}$ and $M_{22}$ are square submatrices, and assume that $M_{22} \succ 0$. The Schur complement $M / M_{22}$ is defined as

$$
\begin{equation*}
M / M_{22}:=M_{11}-M_{12} M_{22}^{-1} M_{21} . \tag{16.2}
\end{equation*}
$$

Then $M / M_{22} \succeq 0$ if and only if $M \succeq 0$.

The proof of this lemma can be found in, e.g., $\left[\mathrm{G}^{+} 10\right.$, Proposition 2.1] or [CLP20, Proposition 5.13]. We aim to apply this lemma for the symmetric matrix $M_{K, N}(x)$ given in $\left(P^{\prime}\right)$. Since $\Delta(x)$ and $\Gamma(x)$ have zero entries in the $S_{2}(x)$-structure, it means the matrix $M_{K, N}(x)$ has the following block structure:

$$
M_{K, N}(x)=\left(\begin{array}{cc}
\Gamma_{2}(x)_{S_{1}, S_{1}}-\frac{1}{N} \Delta(x)_{S_{1}} \Delta(x)_{S_{1}}^{\top}-K \Gamma(x)_{S_{1}, S_{1}} & \Gamma_{2}(x)_{S_{1}, S_{2}} \\
\Gamma_{2}(x)_{S_{2}, S_{1}} & \Gamma_{2}(x)_{S_{2}, S_{1}}
\end{array}\right) .
$$

By folding $M_{K, N}(x)$ into the upper left block, the Schur complement is given by

$$
\begin{aligned}
& M_{K, N}(x) / \Gamma_{2}(x)_{S_{2}, S_{2}} \\
& =\Gamma_{2}(x)_{S_{1}, S_{1}}-\frac{1}{N} \Delta(x)_{S_{1}} \Delta(x)_{S_{1}}^{\top}-K \Gamma(x)_{S_{1}, S_{1}}-\Gamma_{2}(x)_{S_{1}, S_{2}} \Gamma_{2}(x)_{S_{2}, S_{2}}^{-1} \Gamma_{2}(x)_{S_{2}, S_{1}} \\
& =Q(x)-\frac{1}{N} \Delta(x)_{S_{1}} \Delta(x)_{S_{1}}^{\top}-K \Gamma(x)_{S_{1}, S_{1}}
\end{aligned}
$$

where $Q(x):=\left(\begin{array}{ll}\Gamma_{2}(x)_{S_{1}, S_{1}} & \Gamma_{2}(x)_{S_{1}, S_{2}} \\ \Gamma_{2}(x)_{S_{2}, S_{1}} & \Gamma_{2}(x)_{S_{2}, S_{2}}\end{array}\right) / \Gamma_{2}(x)_{S_{2}, S_{2}}$.
The importance of $\Gamma_{2}(x)_{S_{1}, S_{1}}$ for a lower curvature bound was already mentioned in Schmuckenschläger [Sch99, pp. 194-195] where he used the notation $A_{I I}$. It is then implied by Lemma 16.2 that

$$
\begin{equation*}
\mathcal{K}(G, x ; N)=\underset{K}{\arg \max }\left\{Q(x)-\frac{1}{N} \Delta(x)_{S_{1}} \Delta(x)_{S_{1}}^{\top}-K \Gamma(x)_{S_{1}, S_{1}} \succeq 0\right\} \tag{16.3}
\end{equation*}
$$

One can see from Proposition 15.1(L1) that $\Gamma(x)_{S_{1}, S_{1}}=\frac{1}{2} \operatorname{diag}\left(\Delta(x)_{S_{1}}\right)$ and $\Delta(x)_{S_{1}}=$ $\left(p_{x y_{1}} p_{x y_{2}} \ldots p_{x y_{m}}\right)^{\top}$, where $S_{1}(x)=\left\{y_{1}, y_{2}, \ldots, y_{m}\right\}$.
Denote the vector $\mathbf{v}_{0}:=\mathbf{v}_{0}(x)=\left(\sqrt{p_{x y_{1}}} \sqrt{p_{x y_{2}}} \ldots \sqrt{p_{x y_{m}}}\right)^{\top}$. The maximizer $K$ in (16.3) does not change under the multiplication by $\operatorname{diag}\left(\mathbf{v}_{0}\right)^{-1} \succ 0$ both from left and right sides, that is,

$$
\begin{equation*}
\mathcal{K}_{G, x}(N)=\underset{K}{\arg \max }\left\{\operatorname{diag}\left(\mathbf{v}_{0}\right)^{-1} Q(x) \operatorname{diag}\left(\mathbf{v}_{0}\right)^{-1}-\frac{1}{N} \mathbf{v}_{0} \mathbf{v}_{0}^{\top}-\frac{K}{2} \mathrm{Id} \succeq 0\right\} \tag{16.4}
\end{equation*}
$$

We finally deduce that

$$
\begin{aligned}
\mathcal{K}(G, x ; N) & =\lambda_{\min }\left(2 \operatorname{diag}\left(\mathbf{v}_{0}\right)^{-1} Q(x) \operatorname{diag}\left(\mathbf{v}_{0}\right)^{-1}-\frac{2}{N} \mathbf{v}_{0} \mathbf{v}_{0}^{\top}\right) \\
& =\lambda_{\min }\left(A_{\infty}(x)-\frac{2}{N} \mathbf{v}_{0} \mathbf{v}_{0}^{\top}\right)=\lambda_{\min }\left(A_{N}(x)\right)
\end{aligned}
$$

as desired.
Results in later sections and chapters are sometimes presented in terms of BakryÉmery curvature at dimension $N=\infty$. For convenience, we use a shortened notation for the curvature as defined below.

Definition 16.3 (curvature at $N=\infty)$. When a weighted graph $G=(V, w, \mu)$ is fixed and unambiguous, the Bakry-Émery curvature at $x \in V$ is given by

$$
\begin{aligned}
\mathcal{K}_{\infty}(x):=\mathcal{K}(G, x ; \infty) & =\max \left\{K \in \mathbb{R} \mid \Gamma_{2} f(x) \geq K \Gamma f(x) \quad \forall f \in \mathbb{R}^{V}\right\} \\
& =\lambda_{\min }\left(A_{\infty}(x)\right)
\end{aligned}
$$

Remark 16.4. Both the Bakry-Émery curvature $\mathcal{K}_{\infty}(x)$ and the curvature matrix $A_{\infty}(x)$ are completely determined by the weighted structure of the incomplete two-ball around $x$, namely $B_{2}^{\text {inc }}(x)$, which is obtained from the induced subgraph of $B_{2}(x)$ by removing all edges connecting vertices within $S_{2}(x)$. The explicit expression of $A_{\infty}(x)$ is provided in [CKLP21, Appendix A].

### 16.2 Curvature of Cartesian products and graph modifications

One important application of the curvature matrix $A_{\infty}$ is the curvature computation of the Cartesian products. Let us first provide the definition of (weighted) Cartesian product for weighted graphs.

Definition 16.5 (weighted Cartesian product). Given two weighted graphs $G, G^{\prime}$ and two fixed positive numbers $\alpha, \beta \in \mathbb{R}^{+}$, the weighted Cartesian product $G \times_{\alpha, \beta}$ $G^{\prime}$ is defined with the following weight function and vertex measure: for $x, y \in G$ and $x^{\prime}, y^{\prime} \in G^{\prime}$,

$$
\begin{aligned}
w_{\left(x, x^{\prime}\right)\left(y, x^{\prime}\right)} & :=\alpha w_{x y} \mu_{x^{\prime}}, \\
w_{\left(x, x^{\prime}\right)\left(x, y^{\prime}\right)} & :=\beta w_{x^{\prime} y^{\prime}} \mu_{x}, \\
\mu_{\left(x, x^{\prime}\right)} & :=\mu_{x} \mu_{x^{\prime}} .
\end{aligned}
$$

The parameters $\alpha$ and $\beta$ serve two purposes.

1. In the case of non-weighted graphs $G$ and $G^{\prime}$ (i.e., $\mu \equiv 1$ and $w \in\{0,1\}$ ), the choice of $\alpha=\beta=1$ gives the usual Cartesian product graph $G \times G^{\prime}$.
2. In the case of $G$ and $G^{\prime}$ representing Markov chains (i.e., when $\sum_{y} w_{x y}=\mu_{x}$ and $\left.\sum_{y^{\prime}} w_{x^{\prime} y^{\prime}}=\mu_{x^{\prime}}\right)$, the choice of $\alpha+\beta=1$ gives the weighted product $G \times_{\alpha, \beta} G$ which represents the random walk with probability $\alpha$ and $\beta$ following horizontal and vertical edges, respectively.

It was shown in [CKLP21] that the curvature matrix of the Cartesian product of two graphs is simply the direct sum of the curvature matrices of each graph.

The proof of this result, which we omit here, relies on the explicit formula for the curvature matrix.

Theorem 16.6. The curvature matrix of the product $G \times_{\alpha, \beta} G^{\prime}$ is the weighted direct sum of the curvature matrices $G$ and $G^{\prime}$ :

$$
A_{\infty}^{G \times_{\alpha, \beta} G^{\prime}}\left(\left(x, x^{\prime}\right)\right)=\alpha A_{\infty}^{G}(x) \oplus \beta A_{\infty}^{G^{\prime}}\left(x^{\prime}\right)
$$

As an immediate consequence, the curvature of the product is the minimum of the curvature in each factor.

Corollary 16.7. The curvature of the weighted product $H:=G \times_{\alpha, \beta} G^{\prime}$ is

$$
\mathcal{K}_{\infty}^{H}\left(\left(x, x^{\prime}\right)\right)=\min \left\{\alpha \mathcal{K}_{\infty}^{G}(x), \beta \mathcal{K}_{\infty}^{G^{\prime}}\left(x^{\prime}\right)\right\} .
$$

Example 16.8 (hypercubes). Consider the weighted graph $K_{2}$ consisting of two points $x$ and $y$ with the Laplacian given by

$$
\Delta f(x)=p_{x y}(f(y)-f(x)) \text { and } \Delta f(y)=p_{y x}(f(x)-f(y)),
$$

where $p_{x y}=\frac{w_{x y}}{\mu_{x}}$ and $p_{y x}=\frac{w_{y x}}{\mu_{y}}=\frac{w_{x y}}{\mu_{y}}$. Explicit expressions of $\Gamma$ and $\Gamma_{2}$ can be derived as

$$
\begin{aligned}
& 2 \Gamma(f, g)(x)=p_{x y}(f(y)-f(x))(g(y)-g(x)), \\
& 2 \Gamma(f, g)(y)=p_{y x}(f(y)-f(x))(g(y)-g(x)),
\end{aligned}
$$

and

$$
\begin{aligned}
2 \Gamma_{2}(f, g)(x)= & \Delta \Gamma(f, g)(x)-\Gamma(f, \Delta g)(x)-\Gamma(\Delta f, g)(x) \\
= & p_{x y}(\Gamma(f, g)(y)-\Gamma(f, g)(x))-\frac{1}{2} p_{x y}(f(y)-f(x))(\Delta g(y)-\Delta g(x)) \\
& -\frac{1}{2} p_{x y}(\Delta f(y)-\Delta f(x))(g(y)-g(x)) \\
= & \left(\frac{1}{2} p_{x y}\left(p_{y x}-p_{x y}\right)+2 \cdot \frac{1}{2} p_{x y}\left(p_{y x}+p_{x y}\right)\right)(f(y)-f(x))(g(y)-g(x)) \\
= & \left(\frac{3}{2} p_{x y} p_{y x}+\frac{1}{2} p_{x y}^{2}\right)(f(y)-f(x))(g(y)-g(x)) .
\end{aligned}
$$

In particular, $\Gamma_{2} f(x)=\left(\frac{3}{2} p_{y x}+\frac{1}{2} p_{x y}\right) \Gamma f(x)$ and thus the curvature is $\mathcal{K}_{\infty}(x)=$ $\frac{3}{2} p_{y x}+\frac{1}{2} p_{x y}$.
Alternatively, the matrix forms $\Gamma(x)$ and $\Gamma_{2}(x)$ are given by

$$
2 \Gamma(x)=p_{x y}\left(\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right), \quad 2 \Gamma_{2}(x)=\left(\frac{3}{2} p_{x y} p_{y x}+\frac{1}{2} p_{x y}^{2}\right)\left(\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right) .
$$

The Schur complement term is simply $Q(x)=\Gamma_{2}(x)_{S_{1}, S_{1}}$ since $S_{2}(x)$ is empty. The curvature matrix $A_{\infty}(x)$ (of size $\left|S_{1}(x)\right|=1$ ) then equals

$$
\begin{aligned}
A_{\infty}(x) & =2 \operatorname{diag}\left(\mathbf{v}_{0}\right)^{-1} Q(x) \operatorname{diag}\left(\mathbf{v}_{0}\right)^{-1} \\
& =2\left(\sqrt{p_{x y}}\right)^{-1}\left(\frac{3}{4} p_{x y} p_{y x}+\frac{1}{4} p_{x y}^{2}\right)\left(\sqrt{p_{x y}}\right)^{-1}=\left(\frac{3}{2} p_{y x}+\frac{1}{2} p_{x y}\right),
\end{aligned}
$$

and of course, $\mathcal{K}_{\infty}(x)=\lambda_{\text {min }}\left(A_{\infty}\right)=\frac{3}{2} p_{y x}+\frac{1}{2} p_{x y}$.
In the particular case of the non-weighted $K_{2}$ (so $p_{x y}=p_{y x}=1$ ), its curvature matrix is $A_{\infty}^{K_{2}}(x)=(2)$. By viewing the $n$-dimensional hypercube $Q^{n}$ as the Cartesian product $\left(K_{2}\right)^{n}$, its curvature matrix and Bakry-Émery curvature are given by

$$
A_{\infty}^{Q^{n}}(x)=\left(\begin{array}{ccc}
2 & & \\
& \ddots & \\
& & 2
\end{array}\right)_{n \times n} \quad, \text { and } \mathcal{K}_{\infty}^{Q^{n}}(x)=2 .
$$

We would also like to present without proof another application of the relation between $\mathcal{K}_{\infty}(x)$ and $A_{\infty}(x)$ to derive a certain graph modification result in [CKLP21].

Theorem 16.9. Let $G=(V, w, \mu)$ be a weighted graph and fix a vertex $x \in V$. Assume that $p^{-}(y)=p_{y x}$ is independent of $y \in S_{1}(x) .{ }^{1} \quad$ Consider a modified weighted graph $\widetilde{G}$ obtained from $G$ by one of the following operations:
(O1) Increase the edge-weight between a fixed pair $y, y^{\prime} \in S_{1}(x)$ with $y \neq y^{\prime}$ by $\tilde{w}_{y y^{\prime}}=w_{y y^{\prime}}+C_{1}$ for any constant $C_{1}>0$.
(O2) Delete a vertex $z_{0} \in S_{2}(x)$ and remove all of its incident edges, i.e., $\tilde{w}_{y z_{0}}=0$ for all $y \in S_{1}(x)$. Increase the edge-weight between all pairs $y, y^{\prime} \in S_{1}(x)$ with $y \neq y^{\prime}$ by

$$
\begin{equation*}
\tilde{w}_{y y^{\prime}}=w_{y y^{\prime}}+C_{2} w_{y z_{0}} w_{z_{0} y^{\prime}} \tag{16.5}
\end{equation*}
$$

with any constant $C_{2} \geq \frac{p^{-}(y)}{\mu_{x} p_{x z 0}^{(2)}}$.
Then $\mathcal{K}(\widetilde{G}, x ; N) \geq \mathcal{K}(G, x ; N)$ for all $N \in(0, \infty]$.

[^0]
## Chapter 17

## Discrete Laplacian and heat semigroup operator

In this chapter, we present results which go in parallel with Sections 5.3 and 5.5 about the Laplacian, the spectral gap and the heat semigroup operator.

As usual, let $G=(V, w, \mu)$ be a connected and locally finite weighted graph with its graph Laplacian $\Delta$ given by

$$
\Delta f(x)=\sum_{y \in S_{1}(x)} p_{x y}(f(y)-f(x))
$$

for all $f \in \mathbb{R}^{V}$ and all $x \in V$. The eigenvalues-eigenvectors of $\Delta$ are the solutions $(\lambda, f) \in \mathbb{R} \times \mathbb{R}^{V}$ to the equation $\Delta f+\lambda f=0$ on $G$. The eigenspace of each $\lambda$ is the vector space consisting of all eigenfunctions with respect to $\lambda$. If $G$ is finite, its Laplacian eigenvalues are discrete: $0=\lambda_{0}<\lambda_{1} \leq \lambda_{2} \leq \ldots \leq \lambda_{|V|-1}$ including their multiplicities. If $G$ is infinite, the Laplacian is self-adjoint operator (under suitable boundary conditions), and its spectrum is contained in $[0, \infty)$. For example, if $G$ is a non-weighted $d$-regular infinite tree, the spectrum of the non-normalized Laplacian is given by

$$
\sigma(-\Delta)=[d-2 \sqrt{d-1}, d+2 \sqrt{d-1}]
$$

Since the $d$-regular tree is the universal covering of any connected $d$-regular graph, this interval agrees precisely with the interval required for the nontrivial eigenvalues (excluding 0 , and $2 d$ in the bipartite case) of $d$-regular Ramanujan graphs.

Apart from the trivial $\lambda_{0}=0$ which corresponds to the constant eigenfunction $f$, the smallest nonzero eigenvalue $\lambda_{1}$ is the Laplacian spectral gap (and it is closely related to difference between the two largest eigenvalues of the graph's adjacency matrix).

### 17.1 Lichnerowicz spectral gap theorem

In the smooth setting of Riemannian manifolds, we already obtain Lichnerowicz's bound on the smallest nonzero eigenvalue by taking an integral of Bochner's formula and applying the divergence theorem (see Section 5.3). Here we present an analogous result on weighted graphs (see also, e.g. [BCLL17, Theorem 2.1]).

Theorem 17.1 (Lichnerowicz for $B E(K, N)$ ). Fix $N \in[0, \infty)$ and let $G=$ $(V, w, \mu)$ be a weighted graph with $\operatorname{Deg}_{\max }<\infty$ and $K:=\mathcal{K}(G, x ; N)>0$. Then the smallest nonzero Laplacian eigenvalue satisfies

$$
\lambda_{1} \geq \frac{N}{N-1} K
$$

Similarly, for $N=\infty$, we have $\lambda_{1} \geq K$.

Proof. First, we remark that $G$ is finite due to the Bonnet-Myers diameter bound (see Corollary 18.2), since $G$ satisfies $B E(K, N)$ (and hence $B E(K, \infty)$ ) with $K>0$.

Suppose $f$ is an eigenfunction with respect to $\lambda_{1}$, so $\Delta f=-\lambda f$. Using $B E(K, N)$ condition and the divergence theorem $\langle\Delta g, 1\rangle=1$ for all $g$ (see Proposition 15.1(L3)), we deduce that

$$
\begin{aligned}
0 & \leq\left\langle\Gamma_{2} f-K \Gamma f-\frac{1}{N}(\Delta f)^{2}, 1\right\rangle \\
& \leq\left\langle\frac{1}{2} \Delta \Gamma f-\Gamma(f, \Delta f)-K \Gamma f-\frac{\lambda^{2}}{N} f^{2}, 1\right\rangle \\
& =\left\langle\lambda \Gamma f-K \Gamma f-\frac{\lambda^{2}}{N} f^{2}, 1\right\rangle \\
& =\left\langle(\lambda-K)\left(\frac{1}{2} \Delta\left(f^{2}\right)-f \Delta f\right)-\frac{\lambda^{2}}{N} f^{2}, 1\right\rangle \\
& =\left\langle(\lambda-K)\left(\lambda f^{2}\right)-\frac{\lambda^{2}}{N} f^{2}, 1\right\rangle \\
& =\left(\frac{N-1}{N} \lambda^{2}-K \lambda\right)\left\langle f^{2}, 1\right\rangle .
\end{aligned}
$$

Since $\lambda_{1}>0$, we conclude that $\lambda_{1} \geq \frac{N}{N-1} K$ as desired.

### 17.2 Heat semigroup and gradient estimate

In this section, we introduce and discuss properties of the heat semigroup operators $P_{t}$ and the heat kernel $p_{t}$ on weighted graphs. They are very important in the analysis of Bakry-Émery curvature, especially the semigroup characterization via gradient estimate. The definitions and properties of the heat semigroup operators and the heat kernel on graphs are in a similar fashion to those on manifolds, which we discussed in Section 5.5. We refer to the arXiv version of Wojciechowski's PhD thesis [Woj08] for technical details in the case of graphs, including e.g. Dodzuik's construction of the heat kernel and the stochastic completeness. At the same time, we remark that our convention for the graph Laplacian $\Delta$ is weighted and has the opposite sign to the one in [Woj08], which is the combinatorial Laplacian $\Delta f=\sum_{y \in S_{1}(x)}(f(x)-f(y))$.
Let $G=(V, w, \mu)$ be a locally finite and connected weighted graph. We denote by $C_{b}(V):=\left\{f: V \rightarrow \mathbb{R} \mid\|f\|_{\infty}<\infty\right\}$ the space of all bounded functions on $V$.

Definition 17.2. The heat kernel $p:[0, \infty) \times M \times M \rightarrow[0,1]$ (written as $\left.p_{t}(x, y)=p(t, x, y)\right)$ is the smallest nonnegative fundamental solution of the heat equation, that is, for every $f \in C_{b}(V)$, the function

$$
\begin{equation*}
u(t, x)=\sum_{y \in V} \mu_{y} p_{t}(x, y) f(y) \tag{17.1}
\end{equation*}
$$

is a solution to the Cauchy problem

$$
\begin{cases}\frac{\partial}{\partial t} u & =\Delta u  \tag{17.2}\\ u(0, \cdot) & =f\end{cases}
$$

Equivalently, the heat kernel is the smallest nonnegative function $p$ which solves the heat equation

$$
\frac{\partial}{\partial t} p_{t}(x, y)=\Delta p_{t}(x, y)
$$

where the Laplacian $\Delta$ is taken with respect to either $x$ or $y$, and $p$ satisfies the initial condition

$$
p_{0}(x, y)=\delta_{x}(y) .
$$

Furthermore, the heat semigroup operator $P_{t}: C_{b}(V) \rightarrow C_{b}(V), t \geq 0$ is given by

$$
P_{t} f(x):=\sum_{y \in V} \mu_{y} p_{t}(x, y) f(y),
$$

for all $f \in C_{b}(V)$ and $x \in M$.

By definition, it means the function $u(t, x)=P_{t} f(x)$ solves the Cauchy problem (17.2), or in short, we write

$$
\frac{\partial}{\partial t} P_{t}=\Delta P_{t} .
$$

Fundamental and very useful properties of $p_{t}$ and $P_{t}$ are summarized in the following proposition.

Proposition 17.3. $P_{t}$ and $p_{t}$ satisfy the following properties:
(H1) $p_{t}(x, y)=p_{t}(y, x)$.
(H2) $p_{t}(x, y)>0$ for all $x, y$ and $t>0$.
(H3) $\sum_{y \in V} \mu_{y} p_{t}(x, y) \leq 1$.
(H4) $p_{s+t}(x, y)=\sum_{z \in V} \mu_{z} p_{s}(x, z) p_{t}(z, y)$.
(H5) If $f(x) \geq 0$ for all $x$, then $P_{t} f(x) \geq 0$ for all $x$ and $t \geq 0$. Additionally, if $f$ is not identically zero, then $P_{t} f(x)>0$ all $x$ and $t>0$.
(H6) $\left\|P_{t} f\right\|_{\infty} \leq\|f\|_{\infty}$ for all $f \in C_{b}(V)$.
(H7) $P_{t} \circ P_{s}=P_{s+t}$ for all $s, t \geq 0$.
(H8) $\Delta P_{t}=P_{t} \Delta$.

Proof. Properties (H1)-(H4) are due to [Woj08, Theorems 2.1.2 and 2.1.5]. (H5), (H6) and (H7) follow from (H2), (H3) and (H4), respectively, via
$P_{t} f(x)=\sum_{y \in V} \mu_{y} p_{t}(x, y) f(y)$. Lastly, (H8) follows from the proof in [Woj08, Theorem 2.2.1].

The next theorem gives a characterization of the $B E(K, \infty)$ in terms of ( $\Gamma$-type) gradient estimate of the heat semigroup $P_{t}$. This result should be compared to Theorem 5.23, and its proof also employs the same technique of the semigroup interpolation. We refer to a paper by Lin and Liu [LL15] for details about this characterization which also includes $B E(K, N)$ for finite $N$.

Theorem 17.4 (Semigroup characterization of $B E(K, \infty)$ ). Let $G=(V, w, \mu)$ be a weighted graph. A vertex $x$ satisfies $B E(K, \infty)$ if and only if the $\Gamma$-gradient estimate:

$$
\begin{equation*}
\Gamma\left(P_{t} f\right)(x) \leq e^{-2 K t} P_{t}(\Gamma f)(x) \tag{17.3}
\end{equation*}
$$

holds true for all $f \in C_{b}(V)$ and $t \geq 0$.

Proof. $(\Rightarrow)$ : Fix $t \geq 0$ and let $F:[0, t] \rightarrow \mathbb{R}$ be a smooth function given by

$$
F(s):=e^{-2 K s} P_{s}\left(\Gamma P_{t-s} f\right)(x)
$$

We must show that $F(0) \leq F(t)$, so and suffices to prove $F^{\prime}(s) \geq 0$ on $[0, t]$. Differentiation $F$ and using the product rule and the chain rule, we obtain

$$
F^{\prime}(s)=e^{-2 K s}\left[-2 K P_{s}\left(\Gamma P_{t-s} f\right)(x)+\left(\frac{\partial}{\partial s} P_{s}\right)\left(\Gamma P_{t-s} f\right)(x)+P_{s}\left(\frac{\partial}{\partial s} \Gamma\left(P_{t-s} f\right)\right)(x)\right]
$$

With the relation $\frac{\partial}{\partial s} P_{s}=\Delta P_{s}=P_{s} \Delta$ being replaced into the second summand, we obtain $F^{\prime}(s)=e^{-2 K s} P_{s} H f(x)$ where $H$ denotes the operator

$$
H f:=-2 K \Gamma P_{t-s} f+\Delta\left(\Gamma P_{t-s} f\right)+\frac{\partial}{\partial s} \Gamma\left(P_{t-s} f\right)
$$

Moreover, the last summand of $H$ is equal to

$$
\begin{aligned}
\frac{\partial}{\partial s} \Gamma\left(P_{t-s} f\right) & =\frac{\partial}{\partial s} \Gamma\left(P_{t-s} f, P_{t-s} f\right) \\
& =2 \Gamma\left(\frac{\partial}{\partial s}\left(P_{t-s} f\right), P_{t-s} f\right) \\
& =-2 \Gamma\left(\Delta P_{t-s} f, P_{t-s} f\right) \\
& =2 \Gamma_{2}\left(P_{t-s} f\right)-\Delta\left(\Gamma P_{t-s} f\right)
\end{aligned}
$$

Therefore $H f(x)=-2 K \Gamma\left(P_{t-s} f\right)(x)+2 \Gamma_{2}\left(P_{t-s} f\right)(x) \geq 0$, since $x$ satisfies $B E(K, \infty)$. Proposition 17.3(H5) then implies that $P_{s} H f \geq 0$, which gives $F^{\prime}(s) \geq 0$ as desired.
$(\Leftarrow)$ : We differentiate (17.3) at $t=0$ and obtain

$$
\begin{aligned}
0 & \leq \lim _{t \rightarrow 0} \frac{1}{t}\left(e^{-2 K t} P_{t}(\Gamma f)-\Gamma\left(P_{t} f\right)\right) \\
& =\left.\frac{\partial}{\partial t}\right|_{t=0}\left(e^{-2 K t} P_{t}(\Gamma f)-\Gamma\left(P_{t} f\right)\right) \\
& =\Delta(\Gamma f)-2 K \Gamma f-\left.\frac{\partial}{\partial t}\right|_{t=0} \Gamma\left(P_{t} f, P_{t} f\right) \\
& =\Delta(\Gamma f)-2 K \Gamma f+\left[2 \Gamma_{2}\left(P_{t} f\right)-\Delta\left(\Gamma\left(P_{t} f\right)\right)\right]_{t=0} \\
& =2 \Gamma_{2} f-2 K \Gamma f
\end{aligned}
$$

and therefore $x$ satisfies $B E(K, \infty)$
In [LMP17, Theorem 3.4], the authors provide equivalent conditions where the $\Gamma$-gradient estimate (17.3) is sharp. This result, which is stated and proved below, will be helpful when we discuss the rigidity of diameter bound in an upcoming chapter.

Theorem 17.5. Let $G=(V, w, \mu)$ be a weighted graph with $\mathrm{Deg}_{\max }<\infty$ that satisfies $B E(K, \infty)$ with $K>0$. Let $f \in C_{b}(V)$. Then the following are equivalent.
(i) There exists some $x_{0} \in V$ such that $\Gamma P_{t} f\left(x_{0}\right)=e^{-2 K t} P_{t} \Gamma f\left(x_{0}\right)$ for all $t \geq 0$.
(ii) $\Gamma P_{t} f=e^{-2 K t} P_{t} \Gamma f$ for all $t \geq 0$.
(iii) $\Gamma_{2} f=K \Gamma f$.
(iv) $f=\varphi+c$ for a constant $c \in \mathbb{R}$ and an eigenfunction $\varphi$ such that $\Delta \varphi+K \varphi=$ 0 .

Moreover, if any of the above holds, then
(a) $\Gamma f$ is constant.

Proof. Firstly, $G$ is finite due to the Bonnet-Myers diameter bound (Corollary 18.2). We will prove this equivalence in the following order:

$$
(i i) \Rightarrow(i) \Rightarrow(i i i) \Rightarrow(i v) \Rightarrow(i i)
$$

$(i i) \Rightarrow(i)$ is trivial.
$(i) \Rightarrow(i i i)$ : Let $F_{x}(s):=e^{-2 K s} P_{s}\left(\Gamma P_{t-s} f\right)(x)$ for $s \in[0, t]$, and recall that $F_{x}^{\prime}(s)=$ $2 e^{-2 K s} P_{s}\left(\Gamma_{2}\left(P_{t-s} f\right)-K \Gamma\left(P_{t-s} f\right)\right)(x)$. We know from (i) that

$$
0=F_{x_{0}}(t)-F_{x_{0}}(0)=\int_{0}^{t} 2 e^{-2 K s} P_{s}\left(\Gamma_{2}\left(P_{t-s} f\right)-K \Gamma\left(P_{t-s} f\right)\right)\left(x_{0}\right) d s
$$

Due to $B E(K, \infty)$, the above integrand is nonnegative, and hence it must be zero for all $s$. Particularly at $s=t$, we have $P_{t}\left(\Gamma_{2} f-K \Gamma f\right)\left(x_{0}\right)=0$ which holds for all $t>0$. Proposition 17.3(H5) then asserts that $\Gamma_{2} f-K \Gamma f$ is identically zero (otherwise $P_{t}\left(\Gamma_{2} f-K \Gamma f\right)\left(x_{0}\right)>0$ ), so (iii) is true.

Next, suppose for the sake of contradiction that (ii) is false, that is, $\Gamma P_{t} f<$ $e^{-2 K t} P_{t} \Gamma f<0$ at some $x_{1} \in V$ and $t>0$. Then

$$
0<F_{x_{1}}(t)-F_{x_{1}}(0)=\int_{0}^{t} 2 e^{-2 K s} P_{s}\left(\Gamma_{2}\left(P_{t-s} f\right)-K \Gamma\left(P_{t-s} f\right)\right)\left(x_{1}\right) d s
$$

which implies $\Gamma_{2}\left(P_{t-s} f\right)-K \Gamma\left(P_{t-s} f\right)$ is not identically zero. Proposition 17.3(H5) asserts that $P_{s}\left(\Gamma_{2} P_{t-s}-K \Gamma P_{t-s}\right)>0$ for all $x \in V$ and all $0<s<t$, which in turn gives $\Gamma P_{t} f-e^{-2 K t} P_{t} \Gamma f<0$ for all $x \in V$, contradicting to (1).
$(i i i) \Rightarrow(i v)$ : Using Proposition 15.1(L3), we derive

$$
\begin{aligned}
\left\langle\Gamma_{2} f, 1\right\rangle=\langle-\Gamma(f, \Delta f), 1\rangle & =\left\langle\frac{1}{2} f \Delta(\Delta f)+\frac{1}{2}(\Delta f)^{2}, 1\right\rangle \\
& =\frac{1}{2}\langle\Delta(\Delta f), f\rangle+\frac{1}{2}\langle\Delta f, \Delta f\rangle \\
& =\langle\Delta f, \Delta f\rangle .
\end{aligned}
$$

On the other hand, $\Gamma_{2} f=K \Gamma f$ gives

$$
\left\langle\Gamma_{2} f, 1\right\rangle=K\langle\Gamma f, 1\rangle=-K\langle f \Delta f, 1\rangle=-K\langle f, \Delta f\rangle,
$$

and therefore we have $-K\langle f, \Delta f\rangle=\langle\Delta f, \Delta f\rangle$.
Next, we spectrally decompose $f$ into $f=\sum_{i=0}^{|V|-1} \alpha_{1} f_{i}$, where $\alpha_{i} \in \mathbb{R}$ and $f_{i}$ is an eigenfunction corresponding to the Laplacian eigenvalue $\lambda_{i}$ (listed as $0=\lambda_{0}<$ $\lambda_{1} \leq \lambda_{2} \leq \ldots \leq \lambda_{|V|-1}$ ), and $\left\{f_{i}\right\}_{i=0}^{|V|-1}$ forms an orthonormal basis of $\mathbb{R}^{V}$, that is, $\left\langle f_{i}, f_{j}\right\rangle=\delta_{i j}$. This decomposition gives

$$
K \sum_{i} \lambda_{i} \alpha_{i}^{2}=-K\langle f, \Delta f\rangle=\langle\Delta f, \Delta f\rangle=\sum_{i} \lambda_{i}^{2} \alpha_{i}^{2},
$$

which means $\sum_{i} \lambda_{i}\left(\lambda_{i}-K\right) \alpha_{i}^{2}=0$. Recall that $\lambda_{0}=0$ and from Lichnerowicz theorem that $\lambda_{i} \geq K$ for all $i \geq 1$. It implies that $\alpha_{i}=0$ for all $i$ such that $\lambda_{i} \notin\{0, K\}$, and hence we can write $f=c+\varphi$ with some constant $c \in \mathbb{R}$ and $\Delta \varphi+K \varphi=0$.
Next, we prove $(i v) \Rightarrow(a)$, which will help proving $(i v) \Rightarrow(i i)$ later. We know from $f=\varphi+c$ that $\Gamma f=\Gamma \varphi$. The condition $B E(K, N)$ applied for $\varphi$ gives

$$
K \Gamma \varphi \leq \Gamma_{2} \varphi=\frac{1}{2} \Delta\left((\Gamma \varphi)-\Gamma(\varphi, \Delta \varphi)=\frac{1}{2} \Delta(\Gamma \varphi)+K \Gamma \varphi,\right.
$$

and thus $\Delta(\Gamma \varphi) \geq 0$. Due to the divergence theorem $\langle\Delta(\Gamma \varphi), 1\rangle=0$, it must be that $\Delta(\Gamma \varphi)=0$, and hence $\Gamma \varphi$ is constant as desired.
$(i v) \Rightarrow(i i):$ Fix any $x \in V$ and let $H(t):=P_{t} \varphi(x)-e^{-K t} \varphi(x)$. Note that $H(0)=0$ and $H^{\prime}(t)=\Delta P_{t} \varphi(x)+K e^{-K t} \varphi(x)$.
By applying $P_{t}=e^{t \Delta}=\sum_{n=0}^{\infty} \frac{t^{n} \Delta^{n}}{n!}$ to $\varphi$, we have

$$
P_{t} \varphi=\sum_{n=0}^{\infty} \frac{t^{n}(-K)^{n} \varphi}{n!}=e^{-K t} \varphi
$$

Using $f=\varphi+c$, we finally derive

$$
\Gamma P_{t} f=\Gamma P_{t} \varphi=\Gamma e^{-K t} \varphi=e^{-2 K t} \Gamma \varphi=e^{-2 K t} \Gamma f \stackrel{(a)}{=} e^{-2 K t} P_{t} \Gamma f
$$

## Chapter 18

## Bonnet-Myers diameter bound and rigidity of hypercube

In this chapter, we survey for a Bonnet-Myers diameter bound result in the context of Bakry-Émery curvature, and its correspoding rigidity result from the two papers by Liu, Münch and Peyerimhoff [LMP18, LMP17].

### 18.1 Diameter bound for $B E(K, \infty)$

For the sake of simplicity, the results in this survey are presented only in the case of $N=\infty$. The arguments for curvature at a finite dimension $N$ are carried out in a more complicated fashion and can be found in the original paper [LMP18].

Theorem 18.1 ([LMP18, Theorem 2.1]). Let $G=(V, w, \mu)$ be a weighted graph with the positive lower curvature bound $K:=\inf _{x \in V} \mathcal{K}_{\infty}(x)>0$ and $\operatorname{Deg}_{\max }<\infty$. Then for all $x_{0}, y_{0} \in V$, we have

$$
\begin{equation*}
d\left(x_{0}, y_{0}\right) \leq \frac{\sqrt{2 \operatorname{Deg}\left(x_{0}\right)}+\sqrt{2 \operatorname{Deg}\left(y_{0}\right)}}{K} \tag{18.1}
\end{equation*}
$$

Proof. Let $L:=d\left(x_{0}, y_{0}\right)$ and define a function $f(x):=\max \left\{L-d\left(x_{0}, x\right), 0\right\}$, which is bounded and 1-Lipschitz with $f\left(x_{0}\right)=L$ and $f\left(y_{0}\right)=0$. Moreover,

$$
\Gamma f(x)=\frac{1}{2} \sum_{y} p_{x y}(f(y)-f(x))^{2} \leq \frac{1}{2} \sum_{y} p_{x y}=\frac{\operatorname{Deg}(x)}{2}
$$

which means $\|\Gamma f\|_{\infty} \leq \frac{\text { Deg }_{\text {max }}}{2}$.

The triangle inequality gives, for all $t \geq 0$,

$$
\begin{align*}
L & =\left|f\left(x_{0}\right)-f\left(y_{0}\right)\right| \\
& \triangle\left|f\left(x_{0}\right)-P_{t} f\left(x_{0}\right)\right|+\left|P_{t} f\left(x_{0}\right)-P_{t} f\left(y_{0}\right)\right|+\left|P_{t} f\left(y_{0}\right)-f\left(y_{0}\right)\right| . \tag{18.2}
\end{align*}
$$

For any $x \in V$, the term $\left|f\left(x_{0}\right)-P_{t} f\left(x_{0}\right)\right|$ is bounded from above by

$$
\begin{equation*}
\left|f(x)-P_{t} f(x)\right| \leq \int_{0}^{t}\left|\frac{\partial}{\partial s} P_{s} f(x)\right| d s=\int_{0}^{t}\left|\Delta P_{s} f(x)\right| d s \tag{18.3}
\end{equation*}
$$

due to the fundamental theorem of calculus. Moreover,

$$
\begin{align*}
\left(\Delta P_{s} f(x)\right)^{2} \leq 2 \operatorname{Deg}(x) \Gamma\left(P_{s} f\right)(x) & \leq 2 \operatorname{Deg}(x) e^{-2 K s} P_{s}(\Gamma f)(x) \\
& \leq 2 \operatorname{Deg}(x) e^{-2 K s}\|\Gamma f\|_{\infty} \leq \operatorname{Deg}(x)^{2} e^{-2 K s}, \tag{18.4}
\end{align*}
$$

due to Proposition 15.1(L2), Theorem 17.4, and Proposition 17.3(H6).
Combining (18.3) and (18.4) yields for all $x \in V$,

$$
\left|f(x)-P_{t} f(x)\right| \leq \operatorname{Deg}(x) \int_{0}^{t} e^{-K s} d s=\frac{\operatorname{Deg}(x)}{K}\left(1-e^{-K t}\right) \leq \frac{\operatorname{Deg}(x)}{K}
$$

Thus $\left|f\left(x_{0}\right)-P_{t} f\left(x_{0}\right)\right| \leq \frac{\operatorname{Deg}\left(x_{0}\right)}{K}$ and $\left|P_{t} f\left(y_{0}\right)-f\left(y_{0}\right)\right| \leq \frac{\operatorname{Deg}\left(y_{0}\right)}{K}$. We also know from (18.4) that $\Gamma P_{t} f \rightarrow 0$ as $t \rightarrow \infty$, which implies that $\left|P_{t} f(y)-P_{t} f(x)\right| \rightarrow 0$ for all $x \sim y$. Due to connectivity of $G$, the triangle inequality can be applied to deduce that $\left|P_{t} f(y)-P_{t} f(x)\right| \rightarrow 0$ also for all $x, y \in V$. Finally, we can conclude from (18.2) that $L \leq \frac{\operatorname{Deg}\left(x_{0}\right)+\operatorname{Deg}\left(y_{0}\right)}{K}$ as desired.

As an immediate sequence of the above theorem, we obtain the following Bonnet-Myers-type diameter bound result.

Corollary 18.2 (Bonnet-Myers for $B E(K, \infty))$. Let $G=(V, w, \mu)$ be a weighted graph with the positive lower curvature bound $K:=\inf _{x \in V} \mathcal{K}_{\infty}(x)>0$ and $\operatorname{Deg}_{\max }<$ $\infty$. Then $G$ is finite and

$$
\begin{equation*}
\operatorname{diam}(G) \leq \frac{2 \mathrm{Deg}_{\max }}{K} \tag{18.5}
\end{equation*}
$$

We remark from Example 16.8 that in the non-weighted case, the curvature of $n$-dimensional hypercube $\mathcal{Q}^{n}$ is $\mathcal{K}_{\infty}^{\mathcal{Q}^{n}}=2$, and $\operatorname{Deg}_{\max }=\operatorname{diam}\left(\mathcal{Q}^{n}\right)=n$. Thus, the diameter bound (18.5) is sharp for any hypercube. In fact, we will discuss in the next section that hypercubes are the only nonweighted graphs for which (18.5) is sharp.

Example 18.3 (antitrees). Now we would like to make a comparison to Remark 4.3 about a paraboloid being a noncompact space with positive curvature everywhere but the infimum of the curvature is zero (and thus the classical BonnetMyers theorem does not apply). The family of graphs that has the same analogy are known as antitrees, whose name and general definition come from a paper by Keller, Lenz and Wojciechowski [KLW13]. For an infinite sequence of positive integers $\left(a_{k}\right)_{k \in \mathbb{N}}$, an antitree with respect to $\left(a_{k}\right)$ is defined to be a simple and connected graph with the vertex set

$$
V=\bigsqcup_{k} V_{k}, \quad\left|V_{k}\right|=a_{k},
$$

and satisfies the following two properties.
(AT1) Any root vertex $x \in V_{1}$ is connected to all vertices in $V_{2}$ and no vertex in $V_{k}$, $k \geq 3$.
(AT2) Any vertex $x \in V_{k}, k \geq 2$, is connected to all vertices in $V_{k-1}$ and $V_{k+1}$ and no vertex in $V_{l},|l-k| \geq 2$.

In [CLMP20], the authors provide explicit formulae to calculate both the Ollivier Ricci curvature and (non-weighted) Bakry-Émery curvature for antitrees which satisfy additional properties that $\left(a_{k}\right)$ is an increasing sequence, and any $x \in V_{k}$ is connected to all other vertices in $V_{k}$. Such a unique antitree is denoted by $\mathcal{A} \mathcal{T}\left(\left(a_{k}\right)\right)$. In the particular case that $a_{k}=1+(k-1) t$ for a fixed $t \in \mathbb{N}$, the infinite antitree $\mathcal{A} \mathcal{T}\left(\left(a_{k}\right)\right)$ has positive curvature everywhere and the infimum of the curvature is zero (in both contexts of Bakry-Émery curvature and Ollivier Ricci curvature); see [CLMP20, Theorems 2.1.4 and 2.1.6]. Thus these antitrees serve to be an example of infinite graphs with positive curvature everywhere but a Bonnet-Myers-type diameter bound theorem is not applicable.

### 18.2 Rigidity of hypercubes

Continuing from the Bonnet-Myers-type diameter bound result in the previous section, we now present the main result in [LMP17] about the rigidity in the special case of non-weighted graphs; see Cheng's rigidity on manifolds (Theorem 4.2) for comparison. Readers are welcome to consult the original paper for more technical details when dealing with weighted graphs.
As usual, our non-weighted graph $G=(V, E)$ here is connected and locally finite. Moreover, we assume that $G$ has finite maximal vertex degree, $\operatorname{Deg}_{\max }<\infty$.

Theorem 18.4. Let $G=(V, E)$ be a graph with $K:=\inf _{x \in V} \mathcal{K}_{\infty}(x)>0$. Then $G$ satisfies the sharp Bonnet-Myers diameter bound, i.e.,

$$
\operatorname{diam}(G)=\frac{2 \mathrm{Deg}_{\max }}{K}
$$

if and only if $G$ is a hypercube.
We summarize the authors' ideas to prove the forward implication of the above theorem into the following three steps (see Theorems 18.5, 18.7 and 18.8). Afterwards, we present only the proof of Theorem 18.5 to demonstrate how to utilize the sharpness of inequalities involved in the derivation of the distance bound (18.1).

Theorem 18.5. Let $G=(V, E)$ be a graph with $K:=\inf _{x \in V} \mathcal{K}_{\infty}(x)>0$, and let $D:=\operatorname{Deg}_{\text {max }}$. Suppose $G$ satisfies the sharp Bonnet-Myers diameter bound, $\operatorname{diam}(G)=\frac{2 D}{K}$. Let $x_{0} \in V$ be a pole, that is, there exists $y_{0} \in V$ with $d\left(x_{0}, y_{0}\right)=$ $\operatorname{diam}(G)$. Then $K=2$ and $G$ has the hypercube shell structure $\operatorname{HSS}\left(D, 1, x_{0}\right)$, which means
(i) $G$ is $D$-regular and bipartite, and
(ii) the in-degree $d_{x_{0}}^{-}(x):=\left|\left\{v \in S_{1}(x) \mid d\left(x_{0}, v\right)=d\left(x_{0}, x\right)-1\right\}\right|$ is equal to $d\left(x_{0}, x\right)$.

The authors of [LMP17] prove further that a graph $G=(V, E)$ with $K:=$ $\inf _{x \in V} \mathcal{K}_{\infty}(x)=2$ must satisfy two combinatorial properties called the small twosphere property (SSP) and the non-clustering property ( $N C P$ ).

Definition 18.6. Let $G=(V, E)$ be a $D$-regular graph. We say a vertex $x \in V$ satisfy the small two-sphere property (SSP) if $\left|S_{2}(x)\right| \leq\binom{ D}{2}$. We say $x$ satisfies the non-clustering property $(N C P)$ if the following implication holds true: if $d_{x}^{-}(z)=2$ for all $z \in S_{2}(x)$, then for all different $y_{1}, y_{2} \in S_{1}(x)$ there is at most one $z \in S_{2}(x)$ that $y_{1} \sim z \sim y_{2}$. Moreover, $G$ satisfies (SSP) (or (NCP)) if such property is satisfied for all $x \in V$.

Theorem 18.7. Let $G=(V, E)$ be a $D$-regular bipartite graph. Suppose that a vertex $x \in V$ satisfies $\mathcal{K}_{\infty}(x) \geq 2$. Then $x$ satisfies both (SSP) and (NCP). In particular, if $K:=\inf _{x \in V} \mathcal{K}_{\infty}(x)=2$, then $G$ satisfies (SSP) and (NCP).

The authors then conclude that (SSP), (NCP) and $\operatorname{HSS}\left(D, 1, x_{0}\right)$ for some $x_{0} \in V$ altogether imply that $G$ must be a $D$-dimensional hypercube.

Theorem 18.8. Let $G=(V, E)$ be a graph with $\operatorname{HSS}\left(D, 1, x_{0}\right)$ at some $x_{0} \in V$. Suppose that $G$ satisfies both (SSP) and (NCP). Then $G$ is isomorphic to the $D$-dimensional hypercube $Q^{D}$.

Proof of Theorem 18.5. We remark $G$ is finite due to the Bonnet-Myers diameter bound theorem. Let $f_{0} \in \mathbb{R}^{V}$ be given by $f_{0}(x):=d\left(x_{0}, x\right)$ for all $x \in V$. We follow closely the arguments in the proof of Theorem 18.1 that

$$
\begin{align*}
\operatorname{diam}(G)=f_{0}\left(y_{0}\right)-f_{0}\left(x_{0}\right) & \leq \int_{0}^{\infty}\left|\Delta P_{t} f_{0}\left(x_{0}\right)\right|+\left|\Delta P_{t} f_{0}\left(y_{0}\right)\right| d t \\
& \leq \sqrt{2 D} \int_{0}^{\infty} \sqrt{\Gamma P_{t} f_{0}\left(x_{0}\right)}+\sqrt{\Gamma P_{t} f_{0}\left(y_{0}\right)} d t  \tag{18.6}\\
& \leq \sqrt{2 D} \int_{0}^{\infty} e^{-K t}\left(\sqrt{P_{t} \Gamma f_{0}\left(x_{0}\right)}+\sqrt{P_{t} \Gamma f_{0}\left(y_{0}\right)}\right) d t  \tag{18.7}\\
& \leq \sqrt{2 D} \int_{0}^{\infty} 2 e^{-K t} \sqrt{\left\|\Gamma f_{0}\right\|_{\infty}} d t \\
& \leq \sqrt{2 D} \int_{0}^{\infty} 2 e^{-K t} \sqrt{\frac{D}{2}} d t=\frac{2 D}{K} . \tag{18.8}
\end{align*}
$$

The assumption of sharpness $\operatorname{diam}(G)=\frac{2 D}{K}$ implies that all the involved inequalities must hold with equality. The equality of (18.6) implies $\operatorname{Deg}\left(x_{0}\right)=D$, and the equality of (18.7) implies that $\Gamma P_{t} f_{0}\left(x_{0}\right)=e^{-2 K t} P_{t} \Gamma f_{0}\left(x_{0}\right)$ for all $t \geq 0$. Theorem 17.5 asserts that $\Gamma f_{0}$ is constant, and particularly for all $x \in V$ we have

$$
2 \Gamma f_{0}(x)=2 \Gamma f_{0}\left(x_{0}\right)=\sum_{y \in S_{1}\left(x_{0}\right)}\left(d\left(x_{0}, y\right)-d\left(x_{0}, x_{0}\right)\right)^{2}=D
$$

It follows that

$$
D=2 \Gamma f_{0}(x)=\sum_{y \in S_{1}(x)}\left(d\left(x_{0}, y\right)-d\left(x_{0}, x\right)\right)^{2} \leq \sum_{y \in S_{1}(x)} 1=\operatorname{Deg}(x) \leq D
$$

The sharpness of the above inequalities implies that $\operatorname{Deg}(x)=D$ for all $x$ (i.e., $G$ is $D$-regular) and that $\left|d\left(x_{0}, y\right)-d\left(x_{0}, x\right)\right|=1$ for all $y \sim x$. This means there is no spherical edge ( $x \sim y$ such that $d\left(x_{0}, y\right)=d\left(x_{0}, x\right)$ ), so $G$ is indeed a bipartite graph (between vertices in $\bigcup_{i \text { even }} S_{i}\left(x_{0}\right)$ and $\bigcup_{i \text { odd }} S_{i}\left(x_{0}\right)$ ). Observe the Laplacian $\Delta f_{0}$ in terms of degrees:

$$
\begin{equation*}
\Delta f_{0}(x)=\sum_{y \in S_{1}(x)}\left(d\left(x_{0}, y\right)-d\left(x_{0}, x\right)\right)=d_{x_{0}}^{+}(x)-d_{x_{0}}^{-}(x)=D-2 d_{x_{0}}^{-}(x), \tag{18.9}
\end{equation*}
$$

where the last equality comes from the fact that $d_{x_{0}}^{+}(x)+d_{x_{0}}^{-}(x)=D$ since $G$ is $D$-regular and $d_{x_{0}}^{0}(x)=0$ (i.e., there is no spherical edge).
Recall again from Theorem 17.5 that $f_{0}=\varphi+c$ where $\Delta \varphi+K \varphi=0$, so

$$
\Delta f_{0}=\Delta \varphi=-K \varphi=-K f_{0}+c K
$$

To find the values of $c$ and $K$, we evaluate at $x_{0}$ and at its neighbor $x_{1}$. At $x_{0}$, we have $f_{0}\left(x_{0}\right)=0$ and $\Delta f_{0}\left(x_{0}\right)=D$, so the substitution into the above equation yields $c K=D$, which means $\Delta f_{0}=-K f_{0}+D$. At any $x_{1} \in S_{1}\left(x_{0}\right)$, we have $f_{0}\left(x_{1}\right)=1$ and from (18.9) that $\Delta f_{0}\left(x_{1}\right)=D-2$, so the substitution yields $K=2$. Therefore, $\Delta f_{0}=-2 f_{0}+D$, and by comparing this equation to (18.9) we can conclude that $d_{x_{0}}^{-}(x)=f_{0}(x)=d\left(x_{0}, x\right)$ as desired.

## Chapter 19

## Outlook for Bakry-Émery curvature

In this final chapter of Part III, we collect some interesting questions about BakryÉmery curvature on non-weighted graphs and its connection to Ollivier Ricci curvature.

1. The Bonnet-Myers-type diameter bounds with respect to Ollivier Ricci curvature and to Bakry-Émery curvature are similar in comparison:

$$
\begin{align*}
\operatorname{diam}(G) & \leq \frac{2}{\inf _{x, y} \kappa(x, y)}  \tag{19.1}\\
\operatorname{diam}(G) & \leq \frac{2}{\inf _{x} \mathcal{K}_{\infty}(x)} \tag{19.2}
\end{align*}
$$

However, we have learned that their corresponding rigidity results are different: (19.2) is only sharp for hypercubes whereas (19.1) is sharp for a larger class of graphs, including all self-centered Bonnet-Myers sharp graphs as listed in Theorem 13.26.

In [CKK ${ }^{+}$20, Section 10.3], we pose an interesting conjecture that for all regular graphs $G=(V, E)$, the Bakry-Emery curvature satisfies the following inequality

$$
\begin{equation*}
\inf _{x \in V} \mathcal{K}_{\infty}(x) \leq \frac{1}{\operatorname{deg}(G)}+\frac{1}{\operatorname{diam}(G)}, \tag{19.3}
\end{equation*}
$$

which is an improvement of (19.2) in the case of $\operatorname{deg}(G)>\operatorname{diam}(G)$. In fact, it is shown that the new inequality (19.3) is sharp for all self-centered Bonnet-Myers sharp graphs.
2. A broader question about the relation between Ollivier Ricci curvature and Bakry-Émery curvature is to understand graphs which have positive (or non-
negative) curvature in both contexts. Prominent examples are abelian Cayley graphs. Some partial results in this direction can be found in a paper by Ralli [Ral17] and our joint paper [CKK $\left.{ }^{+} 21\right]$. In the latter paper, we also draw connections of Ollivier Ricci and Bakry-Émery to various versions of Ricci-flatness, a curvature notion which was originally introduced by Chung and Yau in [CY96].

## Part IV

Erbar-Maas entropic Ricci curvature on Markov chains

## Chapter 20

## Setting of discrete Markov chains

### 20.1 Irreducible and reversible Markov chains and graph structure

We start with a finite Markov chain $(X, Q)$, where $X$ is a finite set, and $Q$ : $X \times X \rightarrow[0,1]$ is a Markov kernel, where $Q(x, y)$ serves as the transition probability from $x$ to $y$ and it satisfies $\sum_{y \in X} Q(x, y)=1$ for all $x \in V$. The kernel $Q$ can be represented by a square matrix of size $|X|$ with nonnegative entries whose sum in each row equals one. The kernel $Q$ induces a directed graph on $X$ by assigning an directed edge $x \rightarrow y$ whenever $Q(x, y)>0$.
Throughout this thesis, the Markov kernel $Q$ is always assumed to be irreducible and reversible, which we briefly discuss for readers' convenience.

Definition 20.1 (irreducibility). Let $A$ be an $n \times n$ square matrix with nonnegative entries. The matrix $A$ is irreducible if, for any $k, l \leq n$, there exists $m \in \mathbb{N}$ such that $A_{k, l}^{m}>0$.

The notion of irreducibility translates naturally for a Markov kernel. A Markov chain is irreducible means its induced directed graph is strongly connected, i.e., it has a directed path from any vertex to any vertex.

Definition 20.2 (aperiodicity). Let $A$ be an irreducible $n \times n$ matrix $A$. Each index $1 \leq k \leq n$ has a period $p \in \mathbb{N}$ to be the greatest common divisor of all $m \in \mathbb{N}$ such that $A_{k, k}^{m}>0$. The irreducibility asserts that all indexes share the same period $p$ (and hence $p$ is called the period of $A$ ). The matrix $A$ is called aperiodic if $p=1$.

Theorem 20.3 (Perron-Frobenius). Let $A$ be an irreducible square matrix with all entries nonnegative. Then the spectral radius $\sigma(A)>0$ is a real (left- and right-) eigenvalue $r>0$ with multiplicity of one. The (left-) eigenvector corresponding to $r$ has all entries positive. If $A$ is additionally aperiodic, then all other eigenvalues $\lambda$ has modulus strictly less than $r$.

In the case of aperiodicity, the infimum $\inf \{\sigma(A)-|\lambda|: \lambda$ is eigenvalue $\}$ is strictly positive and called the spectral gap. In the particular case of a Markov kernel $Q$, the left-spectral radius equals one due to $\sum_{y \in X} Q(x, y)=1$. Consequently, there exists a unique stationary probability measure $\pi$ on $X$ satisfying $\pi(x)>0$ for all $x, \sum_{x \in X} \pi(x)=1$ and $\sum_{x \in X} Q(x, y) \pi(x)=\pi(y)$ (or written compactly as $\pi^{\top} Q=\pi^{\top}$ ). The spectral gap in fact describes the rate of convergence (i.e., the mixing rate) from any initial probability distribution to the limiting distribution, which is $\pi$.

Definition 20.4 (reversibility). An irreducible Markov kernel $Q: X \times X \rightarrow \mathbb{R}_{\geq 0}$ (with the unique stationary probability measure $\pi$ ) is reversible if it satisfies the following detailed balance equation

$$
\begin{equation*}
Q(x, y) \pi(x)=Q(y, x) \pi(y) \quad \forall x, y \in X \tag{20.1}
\end{equation*}
$$

Henceforth, the kernel $Q$, assumed to be irreducible and reversible, always comes with a unique stationary probability measure $\pi$, which is treated as a reference measure.

One may define the symmetric function $w: X \times X \rightarrow \mathbb{R}_{\geq 0}$ by $w_{x y}:=Q(x, y) \pi(x)=$ $w_{y x}$, which can be regarded as the symmetric weight function for the induced weighted graph $G=(X, w, \pi)$ (compared to the setting of weighted graphs in Section 15.1). This induced weighted graph is no longer a directed graph, since there exists an edge $x \sim y$ if and only if any of the following equivalences holds true: $Q(x, y)>0 \Leftrightarrow w_{x y}>0 \Leftrightarrow Q(y, x)>0$. Moreover, we define the vertex $\pi$-degree to be $\operatorname{Deg}_{\pi}(x):=\frac{1}{\pi(x)} \sum_{y \in V} \pi(y)$. Note that this degree is not the same as the vertex degree for weighted graphs: $\operatorname{Deg}(x):=\frac{1}{\pi(x)} \sum_{y \in S_{1}(x)} w_{x y}=$ $\sum_{y \in S_{1}(x)} Q(x, y)=1$ for all $x$.
A special case is a simple random walk (without laziness) on a finite $d$-regular connected graph $G=(V, E)$ which is given by $X=V, w_{x y}=\frac{1}{d}$ iff $x \sim y$, and $\pi(x)=\frac{1}{|V|}$ for $x \in V$. In this case, we have $\operatorname{Deg}_{\pi}(x)=d$, the usual vertex degree.

### 20.2 Probability space and discrete differential operators

The space of our interest $\mathscr{P}(X) \subset \mathbb{R}^{X}$ is the space of all probability densities (with respect to the reference measure $\pi$ ), which is given by

$$
\mathscr{P}(X):=\left\{\rho: X \rightarrow \mathbb{R}_{\geq 0} \mid \sum_{x \in X} \rho(x) \pi(x)=1\right\} .
$$

Furthermore, a function $V: X \times X \rightarrow \mathbb{R}$ is called a discrete vector field if it satisfies

$$
V(x, y)= \begin{cases}-V(y, x) & \text { for } x \sim y \\ 0 & \text { for } x \nsim y\end{cases}
$$

We denote the space of all discrete vector fields by $\mathfrak{X}(X) \subset \mathbb{R}^{X \times X}$.
Next we give the explicit formulae of the operators discrete gradient $\nabla: \mathbb{R}^{X} \rightarrow$ $\mathfrak{X}(X)$, discrete divergence div $: \mathfrak{X}(X) \rightarrow \mathbb{R}^{X}$, and discrete Laplacian $\Delta:=$ div $\circ \nabla$ as well as the first $\Gamma$-iteration, namely $2 \Gamma(f, g):=\Delta(f g)-f \Delta g-g \Delta f$.

Definition 20.5. For all $f, g \in \mathbb{R}^{X}$ and $V \in \mathfrak{X}(X)$,

$$
\begin{aligned}
\nabla f(x, y) & := \begin{cases}f(y)-f(x) & \text { for } x \sim y, \\
0 & \text { for } x \nsim y,\end{cases} \\
\operatorname{div} V(x) & :=\frac{1}{2} \sum_{y \in X}(V(x, y)-V(y, x)) Q(x, y)=\sum_{y \in X} V(x, y) Q(x, y), \\
\Delta f(x) & :=\sum_{y \in X}(f(y)-f(x)) Q(x, y), \\
2 \Gamma(f, g)(x) & :=\sum_{y \in X}(f(y)-f(x))(g(y)-g(x)) Q(x, y) .
\end{aligned}
$$

The second $\Gamma$-iteration is defined by $2 \Gamma_{2}(f, g):=\Delta \Gamma(f, g)-\Gamma(f, \Delta g)-\Gamma(g, \Delta f)$. For shortened notations, we also write $\Gamma f(x):=\Gamma(f, f)(x)$ and $\Gamma_{2} f(x):=\Gamma_{2}(f, f)(x)$.

All these operators in the setting of Markov chains are the same ones as in the setting of weighted graphs where we discuss the Bakry-Émery curvature. Furthermore, inner-products for functions and for vector fields are defined as follows.

Definition 20.6. For $f_{i} \in \mathbb{R}^{X}$ and $V_{i} \in \mathfrak{X}(X)$,

$$
\begin{aligned}
\left\langle f_{1}, f_{2}\right\rangle & :=\sum_{x \in X} f_{1}(x) f_{2}(x) \\
\left\langle f_{1}, f_{2}\right\rangle_{\pi} & :=\sum_{x \in X} f_{1}(x) f_{2}(x) \pi(x) \\
\left\langle V_{1}, V_{2}\right\rangle_{\pi} & :=\frac{1}{2} \sum_{x, y \in X} V_{1}(x, y) V_{2}(x, y) Q(x, y) \pi(x) .
\end{aligned}
$$

The following proposition asserts that, with respect to the inner products $\langle\cdot, \cdot\rangle_{\pi}$, the divergence is the negative adjoint of the gradient (compared to (5.1) in the smooth setting of manifolds).

Proposition 20.7. For all $f \in \mathbb{R}^{X}$ and $V \in \mathfrak{X}(X)$, we have $\langle\nabla f, V\rangle_{\pi}=-\langle f, \operatorname{div} V\rangle_{\pi}$.
Proof. A straightforward calculation gives

$$
\begin{aligned}
\langle f, \operatorname{div} V\rangle_{\pi} & =\sum_{x \in X} f(x) \pi(x)\left(\frac{1}{2} \sum_{y \in X}(V(x, y)-V(y, x)) Q(x, y)\right) \\
& =\frac{1}{2} \sum_{x, y \in X} f(x) V(x, y) Q(x, y) \pi(x)-\frac{1}{2} \sum_{x, y \in X} f(x) V(y, x) Q(x, y) \pi(x) \\
& =\frac{1}{2} \sum_{x, y \in X} f(x) V(x, y) Q(x, y) \pi(x)-\frac{1}{2} \sum_{y, x \in X} f(y) V(x, y) Q(x, y) \pi(y) \\
& =\frac{1}{2} \sum_{x, y \in X}(f(x)-f(y)) V(x, y) Q(x, y) \pi(x)=-\langle\nabla f, V\rangle_{\pi} .
\end{aligned}
$$

where the third inequality is due to the interchange of $x$ and $y$, together with the detailed balance equation $Q(y, x) \pi(y)=Q(x, y) \pi(x)$.

While it is natural to multiply a function $\rho$ to a vector field $V$ on manifolds, such multiplication is not allowed in the discrete setting since a discrete vector field $V$ is defined on the product space $X \times X$. In order to make them compatible, we transform $\rho \in \mathbb{R}_{\geq 0}^{X}$ into $\hat{\rho} \in \mathbb{R}_{\geq 0}^{X \times X}$ via the logarithmic mean, explained as follows

Definition 20.8. [Maa11, EM12] The function $\theta: \mathbb{R}_{\geq 0} \times \mathbb{R}_{\geq 0} \rightarrow \mathbb{R}_{\geq 0}$ is assumed to be the logarithmic mean, that is,

$$
\theta(s, t)=\int_{0}^{1} s^{1-\alpha} t^{a} d \alpha=\frac{t-s}{\log t-\log s},
$$

where the convention in extremal cases is $\theta(0, t)=\theta(t, 0)=0$ and $\theta(t, t)=t$.
For $\rho \in \mathscr{P}(X)$, the average $\hat{\rho}: X \times X \rightarrow \mathbb{R}_{\geq 0}$ is defined as

$$
\hat{\rho}(x, y):=\theta(\rho(x), \rho(y)) .
$$

Another inner-product for vector fields is defined with respect to $\rho \in \mathscr{P}(X)$ as

$$
\begin{equation*}
\left\langle V_{1}, V_{2}\right\rangle_{\rho}:=\left\langle\hat{\rho} V_{1}, V_{2}\right\rangle_{\pi}=\frac{1}{2} \sum_{x, y \in X} V_{1}(x, y) V_{2}(x, y) \hat{\rho}(x, y) Q(x, y) \pi(x) . \tag{20.2}
\end{equation*}
$$

We remark that the symmetric property $\hat{\rho}(x, y)=\hat{\rho}(y, x)$ asserts that the multiplication $\hat{\rho} V$ is skew-symmetric and remains to be a vector field.

In [Maa11, EM12], the authors allow $\theta$ to be any suitable mean which satisfies certain properties: $\theta$ is continuous on $\mathbb{R}_{\geq 0} \times \mathbb{R}_{\geq 0}$ and $C^{\infty}$ on $(0, \infty) \times(0, \infty)$; $\theta(s, t)=\theta(t, s) ; \theta(s, t)>0$ for $s, t>0 ; \theta(r, t) \geq \theta(s, t)$ for $r \geq s ; \theta(\lambda s, \lambda t)=$ $\lambda \theta(s, t)$ and $\theta(1,1)=1$.

The importance of the choice of $\theta$ being logarithmic mean (which is a crucial ingredient for Erbar-Maas Ricci curvature notion later) is to re-establish the chain rule " $\operatorname{grad}(\log \rho)=\frac{1}{\rho} \operatorname{grad} \rho$ " that appears in the continuous setting:

$$
\begin{equation*}
\hat{\rho}(x, y)=\frac{\rho(y)-\rho(x)}{\log \rho(y)-\log \rho(x)}=\frac{\nabla \rho(x, y)}{\nabla \log \rho(x, y)} . \tag{20.3}
\end{equation*}
$$

Another interesting choice of $\theta$ is the arithmetic mean, which relates Erbar-Maas curvature to Bakry-Émery curvature (see Proposition 22.6).

## Chapter 21

## Otto's calculus for discrete spaces

In this chapter, we revisit the arguments of Otto's calculus (see Chapter 2) about the Riemannian structure of the probability space $\mathscr{P}(X)$. However, since now $X$ is a finite set, arguments certainly become rigorous for a finite-dimensional manifold $\mathscr{P}(X)$ (in comparison to Otto's formal calculus for an infinite-dimensional $\mathscr{P}(M)$ ).

### 21.1 Identification of tangent spaces of $\mathscr{P}(X)$

First, we regard as a manifold $\mathcal{M}$, the subspace of strictly positive probability densities

$$
\mathcal{M}:=\mathscr{P}_{*}(X):=\left\{\rho: X \rightarrow \mathbb{R}^{+} \mid \sum_{x \in X} \rho(x) \pi(x)=1\right\} \subset \mathscr{P}(X) .
$$

Note that $\mathcal{M}$ is an open and convex subspace of the affine hyperplane $\mathcal{V}:=$ $\left\{\rho \in \mathbb{R}^{X} \mid \sum_{x \in X} \rho(x) \pi(x)=1\right\}$, so it is indeed a smooth manifold of dimension $|X|-1$.

Its tangent space at $\rho \in \mathcal{M}$ can be naturally realized as

$$
T_{\rho} \mathcal{M}=\left\{(\rho, s) \text { where } s \in \mathbb{R}^{X} \mid \sum_{x \in X} s(x) \pi(x)=0\right\}
$$

where the curve in $\mathcal{M}$ which passes through the point $\rho$ in the direction $s$ (in the small neighborhood of $\rho$ ) is given by $\rho(t)=\rho_{t}=\rho+t s \in \mathcal{M}$ for small enough $|t|$.

The main goal of this chapter is to follow the ideas from Otto's calculus and identify this tangent space $T_{\rho} \mathcal{M}$ with the following space

$$
\operatorname{Tan}_{\rho} \mathcal{M}:=\left\{(\rho, \nabla \psi) \text { where } \psi \in \mathbb{R}^{X}\right\}
$$

through an identification $\operatorname{map} \mathcal{T}(\rho): \operatorname{Tan}_{\rho} \mathcal{M} \xrightarrow{\sim} T_{\rho} \mathcal{M}$ given by

$$
(\rho, \nabla \psi) \longmapsto(\rho, s=-\operatorname{div}(\hat{\rho} \nabla \psi)) .
$$

Theorem 21.1. The map $\mathcal{T}(\rho): \operatorname{Tan}_{\rho} \mathcal{M} \rightarrow T_{\rho} \mathcal{M}$ is well-defined, and it is a linear isomorphism.

Proof. Here the footpoint $\rho \in \mathcal{M}$ is fixed and omitted from calculation, that is, we only write

$$
\operatorname{Tan}_{\rho} \mathcal{M}=\left\{\nabla \psi: \psi \in \mathbb{R}^{X}\right\}, \text { and } T_{\rho} \mathcal{M}=\left\{s \in \mathbb{R}^{X}: \sum_{x \in X} s(x) \pi(x)=0\right\}
$$

First we check that the image $-\operatorname{div}(\hat{\rho} \nabla \psi)$ lies in $\operatorname{Tan}_{\rho} \mathcal{M}$. This can be seen as a consequence of the divergence theorem

$$
-\sum_{x \in X} \operatorname{div}(\hat{\rho} \nabla \psi) \pi(x)=-\langle 1, \operatorname{div}(\hat{\rho} \nabla \psi)\rangle_{\pi}=\langle\nabla 1, \hat{\rho} \nabla \psi\rangle_{\pi}=0
$$

Now consider the following commutative diagram

where the map $\psi \mapsto \nabla \psi$ is a surjection on $\operatorname{Tan}_{\rho} \mathcal{M}$ by construction, and $T_{\rho} \mathcal{M} \subset$ $\mathbb{R}^{X}$ is a canonical embedding. The map $\tilde{A}(\rho): \mathbb{R}^{X} \rightarrow \mathbb{R}^{X}$ is a linear operator given by $\tilde{A}(\rho) \psi:=-\operatorname{div}(\hat{\rho} \nabla \psi)$. It satisfies

$$
\begin{aligned}
-\operatorname{div}(\hat{\rho} \nabla \psi)(x) & =-\sum_{y \in X} \hat{\rho}(x, y)(\psi(y)-\psi(x)) Q(x, y) \\
& =\left(\sum_{y \neq x} \hat{\rho}(x, y) Q(x, y)\right) \psi(x)-\sum_{y \neq x} \hat{\rho}(x, y) Q(x, y) \psi(y)
\end{aligned}
$$

so its matrix representation is

$$
\tilde{A}(\rho)_{x, y}= \begin{cases}\sum_{z \neq x} \hat{\rho}(x, z) Q(x, z) & \text { if } x=y \\ -\hat{\rho}(x, y) Q(x, y) \psi(y) & \text { if } x \neq y\end{cases}
$$

Note that the matrix $\tilde{A}(\rho)$ has zero sum in each row, which corresponds to the fact that constant functions $\phi$ are in its kernel. We claim that the kernel of $\tilde{A}(\rho)$ has only one dimension and contains only the constant functions. Consider a vector $a=\left(a_{y}\right) \in \mathbb{R}^{X}$ in the kernel of $\tilde{A}(\rho)$ and denote $a_{\max }:=\max _{y \in X}\left(a_{y}\right)$. Pick any $x_{0} \in X$ such that $a_{x_{0}}=a_{\max }$. Since $-\tilde{A}(\rho)_{x_{0}, y}$ is nonnegative for all $y \neq x_{0}$ and vanishes for all $y \not \chi_{Q} x_{0}$, the relation $\sum_{y \in X} \tilde{A}(\rho)_{x_{0}, y} a_{y}=0$ implies that

$$
\tilde{A}(\rho)_{x_{0}, x_{0}} a_{x_{0}}=\sum_{\substack{y \neq x_{0} \\ y \sim Q x_{0}}}(-\tilde{A}(\rho))_{x_{0}, y} a_{y} \leq a_{x_{0}} \sum_{\substack{y \neq x_{0} \\ y \sim Q x_{0}}}(-\tilde{A}(\rho))_{x_{0}, y}=\tilde{A}(\rho)_{x_{0}, x_{0}} a_{x_{0}},
$$

where the last equality is due to the zero sum on $x_{0}$-row. Since this inequality holds with equality, we can infer that $a_{y}=a_{x_{0}}=a_{\max }$ for all $y \sim_{Q} x_{0}$. Then we may pick any $x_{1} \in\left\{y: y \sim_{Q} x_{0}\right\}$ (so that $a_{x_{1}}=a_{\max }$ ) and argue similarly that $a_{y}=a_{x_{1}}=a_{\max }$ for all $y \sim_{Q} x_{1}$. Since $Q$ is irreducible and hence the induced graph is connected, we may repeat this process until we can conclude that $a_{v}=a_{\max }$ for all $v \in X$, that is, the vector $a \in \operatorname{ker} \tilde{A}(\rho)$ must be a constant vector as claimed.
To see that $\mathcal{T}(\rho)=-\operatorname{div}(\hat{\rho} \cdot)$ is injective on $\operatorname{Tan}_{\rho} M$, assume that $-\operatorname{div}(\hat{\rho} \nabla \psi)=$ $0=\tilde{A}(\rho) \psi$. Then $\psi \in \operatorname{ker} \tilde{A}(\rho)$ must be constant, and hence $\nabla \psi=0$. Lastly, we observe that dimensions of $\operatorname{Tan}_{\rho} \mathcal{M}$ and of $T_{\rho} \mathcal{M}$ coincide:

$$
\operatorname{dim} \operatorname{Tan}_{\rho} \mathcal{M}=\operatorname{dim} \mathbb{R}^{X}-\operatorname{dim} \operatorname{ker} \nabla=|X|-1=\operatorname{dim} T_{\rho} \mathcal{M}
$$

since $T_{\rho} \mathcal{M}$ is defined via a nontrivial linear condition. Therefore, $\mathcal{T}(\rho): \operatorname{Tan}_{\rho} \mathcal{M} \rightarrow$ $T_{\rho} \mathcal{M}$ is indeed an isomorphism.

So far we have identified $s \in T_{\rho} \mathcal{M}$ with $\nabla \psi \in \operatorname{Tan}_{\rho} \mathcal{M}$ through the relation $s=-\operatorname{div}(\hat{\rho} \nabla \psi)$. As a consequence, for a smooth curve $\gamma:[a, b] \rightarrow \mathcal{M}$, the tangent vector field $\gamma^{\prime}$ along the curve $\gamma$, which is normally described by $\gamma^{\prime}(t)=$ $\gamma_{t}^{\prime}=\left(\rho_{t}, s_{t}\right) \in T_{\rho_{t}} \mathcal{M}$, can be identified with $\left(\rho_{t}, \nabla \psi_{t}\right) \in \operatorname{Tan}_{\rho_{t}} \mathcal{M}$ where $\frac{d}{d t} \rho_{t}=$ $s_{t}=-\operatorname{div}\left(\hat{\rho}_{t} \nabla \psi_{t}\right)$. In other words, a smooth curve in $\mathcal{M}$ can be represented by $\left(\rho_{t}, \psi_{t}\right)_{t \in[a, b]}$ which satisfies the following discrete continuity equation:

$$
\begin{equation*}
\frac{d}{d t} \rho_{t}+\operatorname{div}\left(\hat{\rho}_{t} \nabla \psi_{t}\right)=0 \tag{21.1}
\end{equation*}
$$

### 21.2 Induced metric from metric tensor

The smooth manifold $\mathcal{M}=\mathscr{P}_{*}(X):=\left\{\rho: X \rightarrow \mathbb{R}^{+} \mid \sum_{x \in X} \rho(x) \pi(x)=1\right\}$ can be equipped by the metric tensor $g_{\rho}: T_{\rho} \mathcal{M} \times T_{\rho} \mathcal{M} \rightarrow \mathbb{R}$ given by

$$
g_{\rho}\left(s_{1}, s_{2}\right):=\left\langle\nabla \psi_{1}, \nabla \psi_{2}\right\rangle_{\rho}
$$

$$
=:\left\langle A(\rho) \psi_{1}, \psi_{2}\right\rangle
$$

where $s_{i}=-\operatorname{div}\left(\hat{\rho} \nabla \psi_{i}\right)$ for $i \in\{1,2\}$, and the linear operator $A$ is defined uniquely via Riesz representation theorem for the symmetric bilinear form $\left(\psi_{1}, \psi_{2}\right) \mapsto$ $\left\langle\nabla \psi_{1}, \nabla \psi_{2}\right\rangle_{\rho}$ as follows.

Definition 21.2 (operator $A(\rho)$ ). For a fixed $\rho \in \mathscr{P}_{*}(X)$, a linear operator $A(\rho)$ : $\mathbb{R}^{X} \rightarrow \mathbb{R}^{X}$ is uniquely defined by

$$
\left\langle A(\rho) \psi_{1}, \psi_{2}\right\rangle:=\left\langle\nabla \psi_{1}, \nabla \psi_{2}\right\rangle_{\rho},
$$

for all $\psi_{1}, \psi_{2} \in \mathbb{R}^{X}$. Equivalently, $A(\rho)=-\pi \operatorname{div}(\hat{\rho} \nabla \cdot)$, which can be seen via the following relation:

$$
\left\langle\nabla \psi_{1}, \nabla \psi_{2}\right\rangle_{\rho}=\left\langle\hat{\rho} \nabla \psi_{1}, \nabla \psi_{2}\right\rangle_{\pi}=\left\langle-\operatorname{div}\left(\hat{\rho} \nabla \psi_{1}\right), \psi_{2}\right\rangle_{\pi}=\left\langle-\pi \operatorname{div}\left(\hat{\rho} \nabla \psi_{1}\right), \psi_{2}\right\rangle .
$$

With this metric tensor $g$, the length of a smooth curve $\gamma:[a, b] \rightarrow \mathcal{M}$ with the tangent vector field $\gamma^{\prime}(t)=\gamma_{t}^{\prime}=\left(\rho_{t}, s_{t}\right)$ is given by

$$
\ell(\gamma):=\int_{a}^{b} g_{\rho_{t}}\left(s_{t}, s_{t}\right)^{\frac{1}{2}} d t=\int_{a}^{b}\left\langle\nabla \psi_{t}, \nabla \psi_{t}\right\rangle_{\rho}^{\frac{1}{2}} d t,
$$

and its energy is given by

$$
E(\gamma):=\int_{a}^{b} g_{\rho_{t}}\left(s_{t}, s_{t}\right) d t=\int_{a}^{b}\left\langle\nabla \psi_{t}, \nabla \psi_{t}\right\rangle_{\rho_{t}} d t
$$

In view of Proposition 2.5 , the metric $\mathcal{W}$ on $\mathscr{P}_{*}(X)$ induced by $g$ can be characterized by the infimum of length of curves or by the infimum of energy. More precisely, the distance between any given $\rho_{0}, \rho_{1} \in \mathcal{M}=\mathscr{P}_{*}(X)$ can be written as

$$
\mathcal{W}\left(\rho_{0}, \rho_{1}\right)=\inf _{\gamma} \ell(\gamma)=\inf _{\gamma} E(\gamma)^{1 / 2},
$$

which the infimum is taken over all curves $\gamma:[0,1] \rightarrow \mathcal{M}$ whose endpoints are $\rho_{0}$ and $\rho_{1}$. This metric $\mathcal{W}$ was introduced and studied in [Maa11, EM12], and it can be formulated as follows.

Definition 21.3 (Discrete transport metric). The Riemannian metric tensor $g$ at $\rho \in \mathcal{M}=\mathscr{P}_{*}(X)$ is defined as

$$
\begin{equation*}
g_{\rho}\left(s_{1}, s_{2}\right):=\left\langle\nabla \psi_{1}, \nabla \psi_{2}\right\rangle_{\rho}, \tag{21.2}
\end{equation*}
$$

where $s_{i}=-\operatorname{div}\left(\hat{\rho} \nabla \psi_{i}\right)$ for $i \in\{1,2\}$. The metric $\mathcal{W}$ induced by $g$ is given, for all pairs $\rho_{0}, \rho_{1} \in \mathscr{P}_{*}(X)$, as

$$
\mathcal{W}\left(\rho_{0}, \rho_{1}\right):=\inf \left\{\int_{0}^{1}\left\|\nabla \psi_{t}\right\|_{\rho_{t}}^{2} d t \mid\left(\rho_{t}, \psi_{t}\right) \in C E\left(\rho_{0}, \rho_{1}\right)\right\}^{\frac{1}{2}}
$$

where $C E\left(\rho_{0}, \rho_{1}\right)$ consists of all smooth curves $\left(\rho_{t}\right)_{t \in[0,1]}$ in $\mathscr{P}_{*}(X)$ together with $\left(\psi_{t}\right)_{t \in[0,1]}$ in $\mathbb{R}^{X}$ which satisfies the discrete continuity equation (21.1): $\frac{d}{d t} \rho_{t}+$ $\operatorname{div}\left(\hat{\rho}_{t} \nabla \psi_{t}\right)=0$ with $\rho_{t=0}=\rho_{0}$ and $\rho_{t=1}=\rho_{1}$.

In [Maa11, Theorem 3.31], Maas formulates the following geodesic equation, which refers to the concept of co-geodesic flows. Here we discuss an alternative idea to prove this formula by applying the first variation formula to a constant speed geodesic from $\rho_{0}$ to $\rho_{1}$ which minimizes the energy: $\mathcal{W}\left(\rho_{0}, \rho_{1}\right)=\inf _{\gamma} E(\gamma)$.

Theorem 21.4. Given any $\rho_{0}, \rho_{1} \in \mathscr{P}_{*}(X)$, a curve $\left(\rho_{t}, \psi_{t}\right)_{t \in[0,1]} \in C E\left(\rho_{0}, \rho_{1}\right)$ is a geodesic from $\rho_{0}$ to $\rho_{1}$ if it satisfies (in addition to the discrete continuity equation) the following discrete geodesic equation:

$$
\begin{equation*}
\frac{d}{d t} \psi_{t}(x)+\frac{1}{2} \sum_{y \in X}\left(\psi_{t}(y)-\psi_{t}(x)\right)^{2} \partial_{1} \theta\left(\rho_{t}(x), \rho_{t}(y)\right) Q(x, y)=0 \quad \forall x \in X \tag{21.3}
\end{equation*}
$$

where $\partial_{1} \theta$ is the partial derivative of $\theta$ with respect to the first coordinate.
Partial proof. Here we will only prove that the function

$$
\Phi_{t}(x):=\frac{d}{d t} \psi_{t}(x)+\frac{1}{2} \sum_{y \in X}\left(\psi_{t}(y)-\psi_{t}(x)\right)^{2} \partial_{1} \theta\left(\rho_{t}(x), \rho_{t}(y)\right) Q(x, y)
$$

is a constant function on $x \in X$.
Consider a proper variation of a curve $\left(\rho_{t}\right)_{t \in[0,1]}$, namely $\rho_{t}^{s}:=\rho_{t}+s \mathfrak{s}_{t} \in \mathscr{P}_{*}(X)$ for $s \in(-\varepsilon, \varepsilon)$ (for some fixed $\varepsilon>0$ ) and $t \in[0,1]$ with $\mathfrak{s}_{t} \in T_{\rho_{t}} \mathscr{P}_{*}(X)$. Here, being a proper variation means that $\rho_{0}^{s}=\rho_{0}$ and $\rho_{1}^{s}=\rho_{1}$ for all $s$, that is, $\mathfrak{s}_{0}=\mathfrak{s}_{1}=0$. Moreover, for each $s$, we represent the curve $\left(\rho_{t}^{s}\right)_{t \in[0,1]}$ by $\left(\rho_{t}^{s}, \psi_{t}^{s}\right)$ which satisfies the continuity equation $\frac{d}{d t} \rho_{t}^{s}+\operatorname{div}\left(\hat{\rho_{t}^{s}} \nabla \psi_{t}^{s}\right)=0$, that is,

$$
\begin{equation*}
\frac{d}{d t} \rho_{t}+s \frac{d}{d t} \mathfrak{s}_{t}=-\operatorname{div}\left(\hat{\rho_{t}^{s}} \nabla \psi_{t}^{s}\right) \tag{21.4}
\end{equation*}
$$

At $s=0$, the continuity equation is simply $\frac{d}{d t} \rho_{t}=-\operatorname{div}\left(\hat{\rho}_{t} \nabla \psi_{t}\right)$, where the convention $\psi_{t}:=\psi_{t}^{0}$ is used here. Taking derivative of (21.4) at $s=0$ yields

$$
\begin{equation*}
\frac{d}{d t} \mathfrak{s}_{t}=-\operatorname{div}\left(\hat{\rho}_{t} \nabla\left(\partial_{s} \psi_{t}^{s}\right)\right)-\operatorname{div}\left(\left(\partial_{s} \hat{\rho}_{t}^{s}\right) \nabla \psi_{t}\right) . \tag{21.5}
\end{equation*}
$$

(Here and henceforth in the proof, we write in short $\partial_{s}:=\left.\frac{\partial}{\partial s}\right|_{s=0}$.)
For each $s$, the energy of the curve $\left(\rho_{t}^{s}\right)_{t \in[0,1]}$ is given by

$$
E\left(\rho_{.}^{s}\right)=\int_{0}^{1}\left\langle\nabla \psi_{t}^{s}, \nabla \psi_{t}^{s}\right\rangle_{\rho_{t}^{s}} d t=\int_{0}^{1}\left\langle-\operatorname{div}\left(\hat{\rho_{t}^{s}} \nabla \psi_{t}^{s}\right), \psi_{t}^{s}\right\rangle_{\pi} d t
$$

$$
=\int_{0}^{1}\left\langle\frac{d}{d t} \rho_{t}+s \frac{d}{d t} \mathfrak{s}_{t}, \psi_{t}^{s}\right\rangle_{\pi} d t
$$

Its derivative at $s=0$ can be computed as

$$
\begin{align*}
\left.\frac{d}{d s}\right|_{s=0} E\left(\rho_{.}^{s}\right) & =\int_{0}^{1}\left\langle\frac{d}{d t} \mathfrak{s}_{t}, \psi_{t}\right\rangle_{\pi}+\left\langle\frac{d}{d t} \rho_{t}, \partial_{s} \psi_{t}^{s}\right\rangle_{\pi} d t \\
& =\int_{0}^{1}\left\langle\frac{d}{d t} \mathfrak{s}_{t}, \psi_{t}\right\rangle_{\pi}+\left\langle-\operatorname{div}\left(\hat{\rho}_{t} \nabla \psi_{t}\right), \partial_{s} \psi_{t}^{s}\right\rangle_{\pi} d t \tag{21.6}
\end{align*}
$$

The inner product $\left\langle-\operatorname{div}\left(\hat{\rho}_{t} \nabla \psi_{t}\right), \partial_{s} \psi_{t}^{s}\right\rangle_{\pi}$ can be computed using (21.5) as follows:

$$
\begin{align*}
\left\langle-\operatorname{div}\left(\hat{\rho}_{t} \nabla \psi_{t}\right), \partial_{s} \psi_{t}^{s}\right\rangle_{\pi} & =\left\langle\hat{\rho}_{t} \nabla \psi_{t}, \nabla\left(\partial_{s} \psi_{t}^{s}\right)\right\rangle_{\pi}=\left\langle\nabla \psi_{t}, \hat{\rho}_{t} \nabla\left(\partial_{s} \psi_{t}^{s}\right)\right\rangle_{\pi} \\
& =\left\langle\psi_{t},-\operatorname{div}\left(\hat{\rho}_{t} \nabla\left(\partial_{s} \psi_{t}^{s}\right)\right)\right\rangle_{\pi} \\
& =\left\langle\psi_{t}, \frac{d}{d t} \mathfrak{s}_{t}+\operatorname{div}\left(\left(\partial_{s} \hat{\rho}_{t}^{s}\right) \nabla \psi_{t}\right)\right\rangle_{\pi} \\
& =\left\langle\psi_{t}, \frac{d}{d t} \mathfrak{s}_{t}\right\rangle_{\pi}-\left\langle\nabla \psi_{t},\left(\partial_{s} \hat{\rho}_{t}^{s}\right) \nabla \psi_{t}\right\rangle_{\pi} . \tag{21.7}
\end{align*}
$$

Next, we plug (21.7) into (21.6) and use integral by parts to obtain

$$
\begin{align*}
\left.\frac{d}{d s}\right|_{s=0} E\left(\rho_{.}^{s}\right) & =\int_{0}^{1} 2\left\langle\frac{d}{d t} \mathfrak{s}_{t}, \psi_{t}\right\rangle_{\pi}-\left\langle\nabla \psi_{t},\left(\partial_{s} \hat{\rho}_{t}^{s}\right) \nabla \psi_{t}\right\rangle_{\pi} d t \\
& =2\left\langle\mathfrak{s}_{1}, \psi_{1}\right\rangle_{\pi}-2\left\langle\mathfrak{s}_{0}, \psi_{0}\right\rangle_{\pi}-\int_{0}^{1} 2\left\langle\mathfrak{s}_{t}, \frac{d}{d t} \psi_{t}\right\rangle_{\pi}+\left\langle\nabla \psi_{t},\left(\partial_{s} \hat{\rho}_{t}^{s}\right) \nabla \psi_{t}\right\rangle_{\pi} d t \\
& =-\int_{0}^{1} 2\left\langle\mathfrak{s}_{t}, \frac{d}{d t} \psi_{t}\right\rangle_{\pi}+\left\langle\nabla \psi_{t},\left(\partial_{s} \hat{\rho}_{t}^{s}\right) \nabla \psi_{t}\right\rangle_{\pi} d t \tag{21.8}
\end{align*}
$$

where $\left\langle\mathfrak{s}_{0}, \psi_{0}\right\rangle_{\pi}=\left\langle\mathfrak{s}_{1}, \psi_{1}\right\rangle_{\pi}=0$ since $\mathfrak{s}_{0}=\mathfrak{s}_{1}=0$.
We recall that $\partial_{s} \rho_{t}^{s}=\partial_{s}\left(\rho_{t}+s \mathfrak{s}_{t}\right)=\mathfrak{s}_{t}$, and hence

$$
\partial_{s} \hat{\rho}_{t}^{s}(x, y)=\partial_{s} \theta\left(\rho_{t}^{s}(x), \rho_{t}^{s}(y)\right)=\partial_{1} \theta\left(\rho_{t}(x), \rho_{t}(y)\right) \mathfrak{s}_{t}(x)+\partial_{2} \theta\left(\rho_{t}(x), \rho_{t}(y)\right) \mathfrak{s}_{t}(y)
$$

Then

$$
\begin{align*}
& \left\langle\nabla \psi_{t},\left(\partial_{s} \hat{\rho}_{t}^{s}\right) \nabla \psi_{t}\right\rangle_{\pi} \\
& =\frac{1}{2} \sum_{x, y \in X}\left(\psi_{t}(y)-\psi_{t}(x)\right)^{2}\left(\partial_{1} \theta\left(\rho_{t}(x), \rho_{t}(y)\right) \mathfrak{s}_{t}(x)+\partial_{2} \theta\left(\rho_{t}(x), \rho_{t}(y)\right) \mathfrak{s}_{t}(y)\right) Q(x, y) \pi(x) \\
& =\sum_{x, y \in X}\left(\psi_{t}(y)-\psi_{t}(x)\right)^{2} \partial_{1} \theta\left(\rho_{t}(x), \rho_{t}(y)\right) \mathfrak{s}_{t}(x) Q(x, y) \pi(x) . \tag{21.9}
\end{align*}
$$

Plugging (21.9) into (21.8) then gives

$$
\begin{aligned}
& \left.\frac{d}{d s}\right|_{s=0} E\left(\rho_{.}^{s}\right) \\
& =-\int_{0}^{1} \sum_{x \in X} \mathfrak{s}_{t}(x) \pi(x)\left(\frac{d}{d t} \psi_{t}(x)+\sum_{y \in X}\left(\psi_{t}(y)-\psi_{t}(x)\right)^{2} \partial_{1} \theta\left(\rho_{t}(x), \rho_{t}(y)\right) Q(x, y)\right) d t \\
& =-\int_{0}^{1}\left\langle\mathfrak{s}_{t}, \Phi_{t}\right\rangle_{\pi} d t .
\end{aligned}
$$

The original curve $\left(\rho_{t}\right)_{t \in[0,1]}$ being a geodesic means that it minimizes the energy among its proper variation $\rho_{t}^{s}=\rho_{t}+s \mathfrak{s}_{t}$, that is, $0=\left.\frac{d}{d s}\right|_{s=0} E\left(\rho_{.}^{s}\right)=-\int_{0}^{1}\left\langle\mathfrak{s}_{t}, \Phi_{t}\right\rangle_{\pi} d t$. Since $\mathfrak{s}_{t}$ is chosen arbitrarily, we can conclude that $\Phi_{t}(x)=c_{t}$, a constant depending on $t$.

The metric $\mathcal{W}$ in Definition 21.3 can be extended to a metric on $\mathscr{P}(X)$, and particularly $\mathcal{W}\left(\rho_{0}, \rho_{1}\right)<\infty$ for all $\rho_{0}, \rho_{1} \in \mathscr{P}(X)$; see [Maa11, Theorems 3.8 and 3.12]. Moreover, the metric space $(\mathscr{P}(X), \mathcal{W})$ is complete (see [Maa11, Theorems 3.22]). The mentioned metric and completeness results in the the original paper are stated with more technical details to deal with more general choices of $\theta$. In case of logarithmic mean however, things are simplified as $C_{\infty}:=\inf _{0}^{1} \frac{1}{\sqrt{\theta(1-r, 1+r)}} d r<\infty$ and $\mathscr{P}_{\alpha}(X):=\{\rho \mid \mathcal{W}(\rho, \alpha)<\alpha\}=\mathscr{P}(X)$.

### 21.3 Discrete JKO theorem

The entropy functional (relative to a reference measure $\pi$ ) is defined for a probability density $\rho \in \mathscr{P}(X)$ by

$$
\begin{equation*}
\operatorname{Ent}(\rho):=\sum_{x \in X} \rho(x) \log \rho(x) \pi(x)=\langle\rho, \log \rho\rangle_{\pi} . \tag{21.10}
\end{equation*}
$$

Maas states and proves in [Maa11, Theorem 4.7 and Corollary 4.8] the discrete analogue of theorem by Jordan-Kinderlehrer-Otto (JKO theorem; see details in Sections 2.3 and 2.4) which asserts that heat flow is gradient flow of the entropy.

Theorem 21.5 (Discrete JKO, [Maa11]). Let $\theta$ be the logarithmic mean defined by $\theta(a, b):=\int_{0}^{1} a^{1-t} b^{t} d t=\frac{a-b}{\log a-\log b}$. An integral curve $t \mapsto \rho_{t}$ of gradient flow of the entropy functional $\operatorname{Ent}(\cdot)$ (in the Riemannian structure of $(\mathscr{P}(X), \mathcal{W}))$ is the solution to the heat equation $\frac{\partial}{\partial t} \rho_{t}=\Delta \rho_{t}$.

Proof. Let $\rho_{t}$ be the integral curve of gradient flow of $\operatorname{Ent}(\cdot)$. For a fixed $t \in \mathbb{R}$ and $s \in T_{\rho_{t}} \mathcal{M}$, we first compute $D \operatorname{Ent}\left(\rho_{t}\right)(s)$ through a curve $c:(-\varepsilon, \varepsilon) \rightarrow \mathscr{P}(X)$ given by $c(\tau)=\rho_{t}+\tau s$ (so that $c(0)=\rho_{t}$ and $\left.c^{\prime}(0)=s\right)$. We have

$$
\begin{aligned}
D \operatorname{Ent}\left(\rho_{t}\right)(s)=\left.\frac{d}{d \tau}\right|_{\tau=0} \operatorname{Ent}(c(\tau)) & =\left.\frac{d}{d \tau}\right|_{\tau=0}\langle c(\tau), \log c(\tau)\rangle_{\pi} \\
& =\left[\left\langle 1+\log c(\tau), c^{\prime}(\tau)\right\rangle_{\pi}\right]_{\tau=0}=\left\langle 1+\log \rho_{t}, s\right\rangle_{\pi}
\end{aligned}
$$

With the identification $s=-\operatorname{div}\left(\hat{\rho}_{t} \nabla \psi\right)$, we have

$$
\begin{aligned}
g_{\rho_{t}}\left(\frac{d}{d t} \rho_{t}, s\right) & =-D \operatorname{Ent}_{\mathrm{vol}}\left(\rho_{t}\right)(s)=-\left\langle 1+\log \rho_{t}, s\right\rangle_{\pi} \\
& =\left\langle 1+\log \rho_{t}, \operatorname{div}\left(\hat{\rho}_{t} \nabla \psi\right)\right\rangle_{\pi} \\
& =-\left\langle\nabla \log \rho_{t}, \hat{\rho}_{t} \nabla \psi\right\rangle_{\pi} \\
& =-\left\langle\nabla \log \rho_{t}, \nabla \psi\right\rangle_{\rho_{t}}=-g_{\rho_{t}}(\mathfrak{s}, s),
\end{aligned}
$$

where $\mathfrak{s}=\operatorname{div}\left(\hat{\rho}_{t} \nabla \log \rho_{t}\right)$ via the defined metric tensor (21.2).
Since the above equality holds for all $s$, we can conclude that

$$
\frac{d}{d t} \rho_{t}=-\mathfrak{s}=\operatorname{div}\left(\hat{\rho}_{t} \nabla \log \rho_{t}\right)=\operatorname{div}\left(\nabla \rho_{t}\right)=\Delta \rho_{t}
$$

Here $\hat{\rho}=\frac{\nabla \rho}{\nabla \log \rho}$ from (20.3) due to the choice of $\theta$ being the logarithmic mean.

## Chapter 22

## Erbar-Maas entropic Ricci curvature

In [EM12], Erbar and Maas introduce for finite Markov chains the following definition of lower Ricci curvature bound in the sense of Lott-Sturm-Villani curvature notion (cf. Definition 3.4).

Definition 22.1 (Entropic Ricci curvature, [EM12]). An irreducible and reversible finite Markov chain $(X, Q, \pi)$ is said to have the entropic Ricci curvature bound below by $K \in \mathbb{R}$, or written as $\operatorname{Ric}(Q) \geq K$, if the relative entropy $\operatorname{Ent}(\cdot)$ is displacement $K$-convex, that is, every constant speed geodesic $\left(\rho_{t}\right)_{t \in[0,1]}$ in $(\mathscr{P}(X), \mathcal{W})$ satisfies the following $K$-convexity inequality:

$$
\begin{equation*}
\operatorname{Ent}\left(\rho_{t}\right) \leq(1-t) \operatorname{Ent}\left(\rho_{0}\right)+t \operatorname{Ent}\left(\rho_{1}\right)-K \frac{t(1-t)}{2} \mathcal{W}\left(\rho_{0}, \rho_{1}\right)^{2} \tag{22.1}
\end{equation*}
$$

for all $t \in[0,1]$.
It is remarked in [EM12] that it suffices to verify (22.1) for only those constant speed geodesics $\left(\rho_{t}\right)_{t \in[0,1]}$ in $\mathscr{P}_{*}(X)$. Since $\operatorname{Ent}(\cdot)$ is smooth on $\mathscr{P}_{*}(X)$, this is equivalent to the infinitesimal $K$-convexity: $\frac{d^{2}}{d t^{2}} \operatorname{Ent}\left(\rho_{t}\right) \geq K \mathcal{W}\left(\rho_{0}, \rho_{1}\right)^{2}=$ $K\left\|\nabla \psi_{t}\right\|_{\rho_{t}}^{2}$, where the rightmost equality is due to the fact that $\rho_{t}$ has constant speed. Hence we may reformulate the above definition of the entropic Ricci curvature as follows.

Definition 22.2. Let $(X, Q, \pi)$ be an irreducible and reversible finite Markov chain. A point $\rho \in \mathscr{P}_{*}(X)$ is said to satisfy $\operatorname{Ric}_{Q}(\rho) \geq K$ if the inequality

$$
\begin{equation*}
\left.\frac{d^{2}}{d t^{2}}\right|_{t=0} \operatorname{Ent}\left(\rho_{t}\right) \geq K\|\nabla \psi\|_{\rho}^{2} \tag{22.2}
\end{equation*}
$$

holds for all $\psi \in \mathbb{R}^{X}$ and all $\left(\rho_{t}, \psi_{t}\right)_{t \in(-\varepsilon, \varepsilon)}$ satisfying (21.1) and (21.3) with $\left(\rho_{0}, \psi_{0}\right)=(\rho, \psi)$.

Note that the entropic Ricci curvature given in Definition 22.2 is of local-type (i.e., it is defined at every point in $\mathscr{P}_{*}(X)$ ), in contrast to the curvature defined globally as in Definition 22.1. In fact, a Markov chain $(X, Q, \pi)$ satisfies $\operatorname{Ric}(Q) \geq K$ if and only if $\operatorname{Ric}_{Q}(\rho) \geq K$ holds for all $\rho \in \mathscr{P}_{*}(X)$.

### 22.1 Equivalent reformulations of the entropic Ricci curvature via Bochner's formula and heat semigroup

This section features an important result from Erbar and Maas [EM12] about the characterization of the entropic Ricci curvature (originally defined via displacement convexity of the Entropy) in terms of the curvature-dimension inequality derived from Bochner's formula (see Theorem 22.5 below). This inequality resembles the condition $B E(K, \infty)$ for Bakry-Émery curvature, and consequently it has another equivalent interpretation in terms of gradient estimate.

Let us first mention a key component introduced in [EM12], namely $\mathcal{B}(\rho, \psi)$.
Definition 22.3. Let $\hat{\rho}$ and $\hat{\Delta} \rho$ be defined as

$$
\begin{aligned}
\hat{\rho}(x, y) & :=\theta(\rho(x), \rho(y)) \\
\hat{\Delta} \rho & :=\partial_{1} \theta(\rho(x), \rho(y)) \Delta \rho(x)+\partial_{2} \theta(\rho(x), \rho(y)) \Delta \rho(y)
\end{aligned}
$$

where $\partial_{1} \theta$ and $\partial_{2} \theta$ is the partial derivatives with respect to the first and second coordinate of $\theta$, respectively. For $\rho \in \mathscr{P}_{*}(X)$ and $\psi \in \mathbb{R}^{X}$, the term $\mathcal{B}(\rho, \psi) \in \mathbb{R}$ is defined to be

$$
\begin{equation*}
\mathcal{B}(\rho, \psi):=\frac{1}{2}\langle\hat{\Delta} \rho \cdot \nabla \psi, \nabla \psi\rangle_{\pi}-\langle\hat{\rho} \cdot \nabla \psi, \nabla \Delta \psi\rangle_{\pi} . \tag{22.3}
\end{equation*}
$$

If compared to the Bochner's formula (5.2), $\mathcal{B}(\rho, \psi)$ vaguely represents the terms $\frac{1}{2} \Delta|\operatorname{grad} f|^{2}-\langle\operatorname{grad} \Delta f, \operatorname{grad} f\rangle$. In fact, it can be regarded as the following Hessian term (see also [EM12, Proposition 4.3])

Proposition 22.4 ([EM12]). For all $\rho \in \mathscr{P}_{*}(X)$ and $\psi \in \mathbb{R}^{X}$, the following identity

$$
\begin{equation*}
\mathcal{B}(\rho, \psi)=\operatorname{Hess}_{\mathcal{W}} \operatorname{Ent}(s, s) . \tag{22.4}
\end{equation*}
$$

holds where $s:=-\operatorname{div}(\hat{\rho} \nabla \psi) \in T_{\rho} \mathscr{P}_{*}(X)$.
Proof. Given any $\rho \in \mathscr{P}_{*}(X)$ and $\psi \in \mathbb{R}^{X}$, let $s=-\operatorname{div}(\hat{\rho} \nabla \psi) \in T_{\rho} \mathscr{P}_{*}(X)$. Let $\left(\rho_{t}\right)_{t \in(-\varepsilon, \varepsilon)}$ be a geodesic in $\left(\mathscr{P}_{*}(X), \mathcal{W}\right)$ such that $\rho_{0}=\rho$ and $\dot{\rho}_{0}=s$, and recall
the fact that the second derivative in time of a functional along a geodesic can be expressed in terms of Hessian, that is,

$$
\begin{aligned}
\frac{d^{2}}{d t^{2}} \operatorname{Ent}\left(\rho_{t}\right) & =\frac{d}{d t} g_{\rho_{t}}\left(\operatorname{grad}_{\mathcal{W}} \operatorname{Ent}\left(\rho_{t}\right), \dot{\rho}_{t}\right) \\
& =g_{\rho_{t}}\left(\nabla_{\dot{\rho}_{t}}\left(\operatorname{grad}_{\mathcal{W}} \operatorname{Ent}\right), \dot{\rho}_{t}\right)+g_{\rho_{t}}(\operatorname{grad}_{\mathcal{W}} \operatorname{Ent}\left(\rho_{t}\right), \underbrace{\frac{D}{d t} \dot{\rho}_{t}}) \\
& =0 \text { geodesic } \\
& =\operatorname{Hess}_{\mathcal{W}} \operatorname{Ent}\left(\dot{\rho}_{t}, \dot{\rho}_{t}\right)
\end{aligned}
$$

Hence it suffices to prove that $\mathcal{B}\left(\rho_{t}, \psi_{t}\right)=\frac{d^{2}}{d t^{2}} \operatorname{Ent}\left(\rho_{t}\right)$. First, take derivative in time of $\operatorname{Ent}\left(\rho_{t}\right)=\left\langle\rho_{t}, \log \rho_{t}\right\rangle_{\pi}$ and use the continuity equation (21.1) together with the fact that $\hat{\rho}_{t}=\frac{\nabla \rho_{t}}{\nabla \log \rho_{t}}$ from (20.3):

$$
\begin{aligned}
\frac{d}{d t} \operatorname{Ent}\left(\rho_{t}\right) & =\left\langle 1+\log \rho_{t}, \frac{d}{d t} \rho_{t}\right\rangle_{\pi} \\
& =-\left\langle 1+\log \rho_{t}, \operatorname{div}\left(\hat{\rho}_{t} \nabla \psi_{t}\right)\right\rangle_{\pi} \\
& =\left\langle\nabla \log \rho_{t}, \hat{\rho}_{t} \nabla \psi_{t}\right\rangle_{\pi}=\left\langle\nabla \rho_{t}, \nabla \psi_{t}\right\rangle_{\pi}
\end{aligned}
$$

The second derivative of $\operatorname{Ent}\left(\rho_{t}\right)$ is given by

$$
\begin{aligned}
\frac{d^{2}}{d t^{2}} \operatorname{Ent}\left(\rho_{t}\right) & =\left\langle\nabla\left(\frac{d}{d t} \rho_{t}\right), \nabla \psi_{t}\right\rangle_{\pi}+\left\langle\nabla \rho_{t}, \nabla\left(\frac{d}{d t} \psi_{t}\right)\right\rangle_{\pi} \\
& =-\left\langle\frac{d}{d t} \rho_{t}, \Delta \psi_{t}\right\rangle_{\pi}-\left\langle\Delta \rho_{t}, \frac{d}{d t} \psi_{t}\right\rangle_{\pi}
\end{aligned}
$$

where the continuity equation (21.1) gives

$$
\left\langle\frac{d}{d t} \rho_{t}, \Delta \psi_{t}\right\rangle_{\pi}=-\left\langle\operatorname{div}\left(\hat{\rho}_{t} \nabla \psi_{t}\right), \Delta \psi_{t}\right\rangle_{\pi}=\left\langle\hat{\rho}_{t} \nabla \psi_{t}, \nabla \Delta \psi_{t}\right\rangle_{\pi}
$$

and the geodesic equation (21.3) gives

$$
\begin{aligned}
\left\langle\Delta \rho_{t}, \frac{d}{d t} \psi_{t}\right\rangle_{\pi}= & -\frac{1}{2} \sum_{x, y \in X} \Delta \rho_{t}(x)\left(\psi_{t}(y)-\psi_{t}(x)\right)^{2} \partial_{1} \theta\left(\rho_{t}(x), \rho_{t}(y)\right) Q(x, y) \pi(x) \\
= & -\frac{1}{4} \sum_{x, y \in X}\left(\psi_{t}(y)-\psi_{t}(x)\right)^{2}\left(\partial_{1} \theta\left(\rho_{t}(x), \rho_{t}(y)\right) \Delta \rho_{t}(x)+\right. \\
& \left.\partial_{2} \theta\left(\rho_{t}(x), \rho_{t}(y)\right) \Delta \rho_{t}(y)\right) Q(x, y) \pi(x) \\
= & -\frac{1}{2}\left\langle\hat{\Delta} \rho \cdot \nabla \psi_{t}, \nabla \psi_{t}\right\rangle_{\pi} .
\end{aligned}
$$

Hence $\frac{d^{2}}{d t^{2}} \operatorname{Ent}\left(\rho_{t}\right)=-\left\langle\hat{\rho}_{t} \nabla \psi_{t}, \nabla \Delta \psi_{t}\right\rangle_{\pi}+\frac{1}{2}\left\langle\hat{\Delta} \rho \cdot \nabla \psi_{t}, \nabla \psi_{t}\right\rangle_{\pi}=\mathcal{B}\left(\rho_{t}, \psi_{t}\right)$ as desired.

Now we are ready to state and prove the equivalent reformulations of the entropic Ricci curvature. This result is due to [EM12] and Erbar and Fathi [EF18].
Theorem 22.5. Given an irreducible and reversible finite Markov chain $(X, Q, \pi)$, the following statements are equivalent.

1. $\operatorname{Ric}(Q) \geq K$, i.e., (22.1) holds.
2. All $\rho \in \mathscr{P}_{*}(X)$ satisfies $\operatorname{Ric}_{Q}(\rho) \geq K$, i.e., (22.2) holds.
3. The following Bochner's inequality holds for all $\rho \in \mathscr{P}_{*}(X)$ and $\psi \in \mathbb{R}^{X}$ :

$$
\begin{equation*}
\mathcal{B}(\rho, \psi) \geq K\|\nabla \psi\|_{\rho}^{2}, \tag{22.5}
\end{equation*}
$$

where $\mathcal{B}(\rho, \psi)$ is defined in (22.3) with $\theta$ being the logarithmic mean.
4. Global gradient estimate holds for for all $\rho \in \mathscr{P}(X)$ and $\psi \in \mathbb{R}^{X}$ and all $t \geq 0$,

$$
\left\|\nabla P_{t} \psi\right\|_{\rho}^{2} \leq e^{-2 \kappa t}\|\nabla \psi\|_{P_{t} \rho}^{2},
$$

or more explicitly

$$
\begin{align*}
& \frac{1}{2} \sum_{u, v}\left(P_{t} \psi(v)-P_{t} \psi(u)\right)^{2} \hat{\rho}(u, v) Q(u, v) \pi(u) \\
\leq & e^{-2 \kappa t} \frac{1}{2} \sum_{u, v}(\psi(v)-\psi(u))^{2} \widehat{P_{t} \rho}(u, v) Q(u, v) \pi(u) \tag{22.6}
\end{align*}
$$

where $P_{t}$ denotes the heat semigroup operator.
Proof. The equivalence (1) $\Leftrightarrow(2)$ is discussed in the beginning of this chapter. The equivalence $(2) \Leftrightarrow(3)$ is a consequence of Proposition 22.4 applied to (22.2). The equivalence $(3) \Leftrightarrow(4)$ is due to [EF18, Theorem 3.1], following a similar argument in Theorem 5.22 by differentiating $F(s):=e^{-2 K s}\left\|\nabla P_{t-s} \psi\right\|_{P_{s} \rho}^{2}$.

The above equivalence $(2) \Leftrightarrow(3)$ requires the fact that $\theta$ is the logarithmic mean. We can generalize this curvature notion by defining it directly from Bochner's inequality. More precisely, we define the entropic Ricci curvature at a fixed $\rho \in$ $\mathscr{P}(X)$ as

$$
\operatorname{Ric}_{Q, \theta}(\rho):=\sup \left\{K \in \mathbb{R} \mid(22.5) \text { holds for all } \psi \in \mathbb{R}^{X}\right\}
$$

with a given choice of the mean $\theta$. Another interesting choice of $\theta$ is the arithmetic mean (AM), where the resulting curvature notion is closely related to the BakryÉmery curvature. This observation is mentioned in [Maa17], and we clarify it in the following proposition.

Proposition 22.6. Let $(X, Q, \pi)$ be a finite irreducible and reversible Markov chain. For any $x \in X$, consider the probability density $\delta_{x}:=\frac{1}{\pi(x)} \mathbf{1}_{x} \in \mathscr{P}(X)$. Then $\operatorname{Ric}_{Q, \mathrm{AM}}\left(\delta_{x}\right) \geq K$ if and only if $x$ satisfies $B E(K, \infty)$.

Proof. The choice of $\theta$ being the arithmetic mean implies that $\Delta \hat{\rho}(u, v)=\frac{1}{2} \rho(u)+$ $\frac{1}{2} \rho(v)$. With the choice of $\rho=\delta_{x}$, we have nontrivial values of $\Delta \hat{\rho}$ and $\hat{\Delta} \rho$ as follows:

$$
\begin{aligned}
\hat{\rho}(x, y) & =\hat{\rho}(y, x)=\frac{1}{2 \pi(x)} \\
\Delta \hat{\rho}(x, y) & =\Delta \hat{\rho}(y, x)=-\frac{1}{2 \pi(x)}+\frac{Q(y, x)}{2 \pi(x)} \\
\Delta \hat{\rho}(y, z) & =\Delta \hat{\rho}(z, y)=\frac{Q(y, x)}{2 \pi(x)}=\frac{Q(x, y)}{2 \pi(y)} .
\end{aligned}
$$

for all $y \in S_{1}(x)$ and $z \in S_{2}(x)$. We must show that the inequality $\mathcal{B}(\rho, \psi) \geq$ $K\|\nabla \psi\|_{\rho}^{2}$ is simplified into $\Gamma_{2} \psi(x) \geq K \Gamma \psi(x)$. First, the term $\langle\hat{\Delta} \rho \cdot \nabla \psi, \nabla \psi\rangle_{\pi}$ is given by

$$
\begin{align*}
\langle\hat{\Delta} \rho \cdot \nabla \psi, \nabla \psi\rangle_{\pi}= & \frac{1}{2} \sum_{u, v}(\psi(v)-\psi(u))^{2} Q(u, v) \pi(u) \hat{\Delta} \rho(u, v) \\
= & \sum_{y \in S_{1}(x)}(\psi(y)-\psi(x))^{2} Q(x, y)\left(-\frac{1}{2}+\frac{Q(y, x)}{2}\right)+ \\
& \sum_{y \in S_{1}(x)} \sum_{z \in S_{1}(y) \backslash\{x\}}(\psi(z)-\psi(y))^{2} Q(y, z) \frac{Q(x, y)}{2} \\
= & \sum_{y \in S_{1}(x)}(\Gamma \psi(y)-\Gamma \psi(x)) Q(x, y)=\Delta(\Gamma \psi)(x) . \tag{22.7}
\end{align*}
$$

The second term $\langle\hat{\rho} \cdot \nabla \psi, \nabla \Delta \psi\rangle_{\pi}$ is given by

$$
\begin{align*}
\langle\hat{\rho} \cdot \nabla \psi, \nabla \Delta \psi\rangle_{\pi} & =\frac{1}{2} \sum_{u, v}(\psi(v)-\psi(u))(\Delta \psi(v)-\Delta \psi(u)) Q(u, v) \pi(u) \hat{\rho}(u, v) \\
& =\sum_{y \in S_{1}(x)}(\psi(y)-\psi(x))(\Delta \psi(v)-\Delta \psi(u)) \frac{Q(x, y)}{2} \\
& =\Gamma(\psi, \Delta \psi)(x) \tag{22.8}
\end{align*}
$$

and similarly, $\|\nabla \psi\|_{\rho}^{2}=\langle\hat{\rho} \cdot \nabla \psi, \nabla \psi\rangle_{\pi}=\Gamma \psi(x)$. Combining (22.7) and (22.8) yields $\mathcal{B}(\rho, \psi)=\frac{1}{2} \Delta(\Gamma \psi)(x)-\Gamma(\psi, \Delta \psi)(x)=\Gamma_{2} \psi(x)$ as desired.

### 22.2 Products of Markov chains and hypercube example

One of the main results in [EM12] which follows from the characterization of entropic curvature by the Bochner's inequalty $\mathcal{B}(\rho, \psi) \geq K\|\nabla \psi\|_{\rho}^{2}$ is the curvature estimate for the Cartesian product of two Markov chains.
The following definition of the Cartesian products of Markov chains is almost perfectly aligned with The definition of Cartesian products is previously discussed in the setting of weighted graphs (see Definition 16.5). The products of Markov chains are very similar (with additional exception that they are allowed to be lazy, i.e. $Q(x, x)>0)$. Products here are defined for $n \geq 2$ components.

Definition 22.7. Given $n$ irreducible and reversible finite Markov chains ( $X_{i}, Q_{i}, \pi_{i}$ ), $1 \leq i \leq n$ and a fixed $\alpha=\left(\alpha_{1}, \alpha_{2}, \ldots, \alpha_{n}\right) \in\left(\mathbb{R}_{\geq 0}\right)^{n}$ with $\sum_{i=1}^{n} \alpha_{i}=1$. Their Cartesian product $\left(X, Q_{\alpha}, \pi\right)$ is defined on the product space $X=\prod_{i} X_{i}$ with the Markov kernel given for $\mathbf{x}=\left(x_{1}, \ldots, x_{n}\right)$ and $\mathbf{y}=\left(y_{1}, \ldots, y_{n}\right)$ in $X$ as

$$
Q_{\alpha}(\mathbf{x}, \mathbf{y}):= \begin{cases}\sum_{i=1}^{n} \alpha_{i} Q_{i}\left(x_{i}, x_{i}\right) & \text { if } \mathbf{x}=\mathbf{y}, \\ \alpha_{i} Q_{i}\left(x_{i}, y_{i}\right) & \text { if } x_{i} \neq y_{i} \text { and } x_{j}=y_{j} \forall j \neq i, \\ 0 & \text { otherwise } .\end{cases}
$$

Moreover, $\pi(\mathbf{x}):=\prod_{i} \pi_{i}\left(x_{i}\right)$ is indeed the unique stationary distribution for this kernel.
Theorem 22.8. [EM12, Theorem 6.2] Assume that $\operatorname{Ric}\left(Q_{i}\right) \geq K_{i}$ for $1 \leq i \leq n$. Then

$$
\operatorname{Ric}\left(Q_{\alpha}\right) \geq \min _{i}\left\{\alpha_{i} K_{i}\right\} .
$$

The proof in the original paper applies the concept of mapping representations for the terms $\mathcal{B}(\rho, \psi)$ and $\|\nabla \psi\|_{\rho}^{2}$, which is not covered here in this survey. However, we would like to mention the resemblance between this result and the Cartesian product result for the Bakry-Émery curvature given in Corollary 16.7.

Now we revisit the classical example of the hypercubes and compute the lower curvature bound for (a simple random walk on) the hypercube.
Example 22.9 (hypercubes). We start with the Markov chain of the two-point space $K_{2}=\{x, y\}$ with $Q(x, y)=Q(y, x)=1$ (so it represent the (nonlazy) simple random walk of $K_{2}$ ) and $\pi(x)=\pi(y)=\frac{1}{2}$. Note that $\Delta f(x)=f(y)-$ $f(x)=-\Delta f(y)$ for all $f \in \mathbb{R}^{\{x, y\}}$. To derive $\mathcal{B}(\rho, \psi)$ and $\|\nabla \psi\|_{\rho}^{2}$, we first compute $\|\nabla \psi\|_{\rho}^{2}=\frac{1}{2}(\psi(y)-\psi(x))^{2} \hat{\rho}(x, y)$, and

$$
\langle\hat{\Delta} \rho \cdot \nabla \psi, \nabla \psi\rangle_{\pi}=\frac{1}{2}(\psi(y)-\psi(x))^{2} \hat{\Delta} \rho(x, y)
$$

$$
\begin{aligned}
\langle\hat{\rho} \cdot \nabla \psi, \nabla \Delta \psi\rangle_{\pi} & =\frac{1}{2}(\psi(y)-\psi(x))(\Delta \psi(y)-\Delta \psi(x)) \hat{\rho}(x, y) \\
& =-(\psi(y)-\psi(x))^{2} \hat{\rho}(x, y) .
\end{aligned}
$$

Thus $\mathcal{B}(\rho, \psi)=\left(\frac{1}{2} \frac{\hat{\Delta} \rho(x, y)}{\hat{\rho}(x, y)}+2\right)\|\nabla \psi\|_{\rho}^{2}$. And hence the optimal bound for the curvature is

$$
K=2+\inf _{\rho} \frac{1}{2} \frac{\hat{\Delta} \rho(x, y)}{\hat{\rho}(x, y)} .
$$

For shortened notations, we write $\rho(x)=a, \rho(y)=b$ and $\log b-\log a=c$. Thus

$$
\begin{aligned}
\hat{\Delta} \rho(x, y) & =\partial_{1} \theta(\rho(x), \rho(y))(\rho(y)-\rho(x))+\partial_{2} \theta(\rho(x), \rho(y))(\rho(x)-\rho(y)) \\
& =\frac{1}{c^{2}}\left(-c+\frac{b-a}{a}\right)(b-a)+\frac{1}{c^{2}}\left(c-\frac{b-a}{b}\right)(a-b) \\
& =\frac{b-a}{c}\left(-2+\frac{(b-a)}{c}\left(\frac{1}{a}+\frac{1}{b}\right)\right) \geq 0,
\end{aligned}
$$

where the last inequality is due to the logarithmic mean being greater or equal to the harmonic mean: $\frac{b-a}{\log b-\log a} \geq \frac{2}{\frac{1}{a}+\frac{1}{b}}$. Moreover, choosing $\rho(x)=\rho(y)$ gives $\hat{\Delta} \rho(x, y)=0$. Therefore, the entropic Ricci curvature for the simple random walk on $K_{2}$ is given by the sharp estimate $\operatorname{Ric}\left(K_{2}, s r w\right) \geq 2$.

The simple random walk on the $n$-dimensional hypercube $\mathcal{Q}^{n}$ is the Cartesian product of $n$ copies of the simple random walk on $K_{2}$ with weights $\alpha_{i} \equiv \frac{1}{n}$. It then follows immediately that $\operatorname{Ric}\left(\mathcal{Q}^{n}, s r w\right) \geq \frac{2}{n}$. In fact, it is further discussed in [EM12, Corollary 7.10] that $\frac{2}{n}$ is the sharp curvature bound.

## Chapter 23

## Diameter bound for Markov chain with positive curvature

In this chapter, we present a generalization of the main result in [Kam20b] that for a Markov chain $(X, Q, \pi)$ with positive entropic Ricci curvature $\operatorname{Ric}(Q) \geq K>0$, the diameter of the underlying graph is bounded from above by

$$
\operatorname{diam}(X) \leq \frac{2}{K} \sqrt{\frac{D \log D}{D-1}}
$$

where $D$ is the maximal vertex $\pi$-degree. The proof is divided into two steps. The first step is to follow the ideas by Erbar and Fathi [EF18] to derive a similar local gradient estimate as follows.

Theorem 23.1 (Local gradient estimate). If $\operatorname{Ric}(Q) \geq K$, then for all $x \in X$ and $\psi \in \mathbb{R}^{X}$ and all $\varepsilon \geq 0$ we have

$$
\begin{equation*}
\Gamma\left(P_{t} \psi\right)(x) \pi(x) \leq \frac{e^{-2 K t}}{2 \theta(1, \varepsilon)}\left[P_{t}(\Gamma \psi)(x) \pi(x)+\varepsilon \sum_{y \in N(x)} P_{t}(\Gamma \psi)(y) \pi(y)\right] \tag{23.1}
\end{equation*}
$$

where $2 \Gamma f(x):=\sum_{y}(f(y)-f(x))^{2} Q(x, y)$.
Proof. The proof follows ideas from [EF18, Corollary 3.4]. For convenience, we recall the global gradient estimate (22.6) here:

$$
\begin{array}{r}
\frac{1}{2} \sum_{u, v}\left(P_{t} \psi(v)-P_{t} \psi(u)\right)^{2} \hat{\rho}(u, v) Q(u, v) \pi(u) \\
\leq e^{-2 K t} \frac{1}{2} \sum_{u, v}(\psi(v)-\psi(u))^{2} \widehat{P_{t} \rho}(u, v) Q(u, v) \pi(u) .
\end{array}
$$

We localize (22.6) by choosing $\rho=\mathbf{1}_{x}+\varepsilon \sum_{y \in S_{1}(x)} \mathbf{1}_{y}$ for a fixed $x \in V$ and a parameter $\varepsilon \in[0, \infty)$. Note that we do not require $\rho$ to be a probability density (i.e., $\sum_{x} \rho(x) \pi(x)=1$ ) since (22.6) is homogeneous in $\rho$.

For our particular choice of $\rho$, we have $\hat{\rho}(x, y)=\hat{\rho}(y, x)=\theta(1, \varepsilon)$ for all $y \in S_{1}(x)$. The left-hand-side of (22.6) is bounded from below by

$$
\begin{equation*}
\text { L.H.S. } \geq \sum_{y \in S_{1}(x)}\left(P_{t} \psi(y)-P_{t} \psi(x)\right)^{2} \theta(1, \varepsilon) Q(x, y) \pi(x)=2 \theta(1, \varepsilon) \Gamma\left(P_{t} \psi\right)(x) \pi(x) . \tag{23.2}
\end{equation*}
$$

On the other hand, the right-hand-side of (22.6) is bounded from above by

$$
\begin{align*}
\text { R.H.S. } & \leq e^{-2 K t} \frac{1}{2} \sum_{u, v}(\psi(v)-\psi(u))^{2} \frac{P_{t} \rho(u)+P_{t} \rho(v)}{2} Q(u, v) \pi(u) \\
& =e^{-2 K t} \frac{1}{2} \sum_{u, v}(\psi(v)-\psi(u))^{2} P_{t} \rho(u) Q(u, v) \pi(u) \tag{23.3}
\end{align*}
$$

due to $\theta(s, t) \leq(s+t) / 2$ and the symmetry from interchanging $u$ and $v$.
We now apply the heat kernel $p_{t}(\cdot, \cdot)$ given by $P_{t} g(u)=\sum_{z} p_{t}(u, z) g(z) \pi(z)$ for every function $g$. With our choice of $\rho$, we obtain

$$
P_{t} \rho(u)=p_{t}(u, x) \rho(x) \pi(x)+\varepsilon \sum_{y \in S_{1}(x)} p_{t}(u, y) \rho(y) \pi(y),
$$

which we substitute into (23.3) and use the symmetry of heat kernel: $p_{t}(u, v)=$ $p_{t}(v, u)$ to derive

$$
\begin{align*}
\text { R.H.S. } \leq & \frac{e^{-2 K t}}{2}\left[\pi(x) \sum_{u} p_{t}(u, x) \pi(u) \sum_{v}(\psi(v)-\psi(u))^{2} Q(u, v)+\right. \\
& \left.\varepsilon \sum_{y \in S_{1}(x)} \pi(y) \sum_{u} p_{t}(u, y) \pi(u)\left(\sum_{v}(\psi(v)-\psi(u))^{2} Q(u, v)\right)\right] \\
= & e^{-2 K t}\left[\pi(x) P_{t}(\Gamma \psi)(x)+\varepsilon \sum_{y \in S_{1}(x)} \pi(y) P_{t}(\Gamma \psi)(y)\right] . \tag{23.4}
\end{align*}
$$

The desired inequality then follows from combining (23.2) and (23.4).
We have the following corollary as an immediate consequence of the above theorem.
Corollary 23.2. If $\operatorname{Ric}(Q) \geq K$, then for all $x \in X$ and $\psi \in \mathbb{R}^{X}$, we have

$$
\begin{equation*}
\Gamma\left(P_{t} \psi\right)(x) \leq c \cdot e^{-2 K t}\left\|P_{t}(\Gamma \psi)\right\|_{\infty} \tag{23.5}
\end{equation*}
$$

where $c:=\frac{D \log D}{D-1}$ and $D$ is the maximal vertex $\pi$-degree,

$$
D:=\max _{x \in X} \operatorname{Deg}_{\pi}(x)=\max _{x \in X} \frac{1}{\pi(x)} \sum_{y \in S_{1}(x)} \pi(y)
$$

Proof. It is implied by Theorem 23.1 that $\Gamma\left(P_{t} \psi\right)(x) \leq c_{\varepsilon, x} \cdot e^{-2 K t}\left\|P_{t}(\Gamma \psi)\right\|_{\infty}$, where the constant term $c_{\varepsilon, x}$ (depending on $\varepsilon$ and $x$ ) is given by

$$
c_{\varepsilon, x}:=\frac{1}{2 \theta(1, \varepsilon)}\left(1+\varepsilon \sum_{y \in S_{1}(x)} \frac{\pi(y)}{\pi(x)}\right)=\frac{1+\varepsilon \operatorname{Deg}_{\pi}(x)}{2 \theta(1, \varepsilon)} \leq \frac{1+\varepsilon D}{2 \theta(1, \varepsilon)} .
$$

In particular when choosing $\varepsilon=\frac{1}{D}$, we have $c_{\varepsilon, x} \leq \frac{D \log D}{D-1}$.
The second step is to prove a Bonnet-Myers-type diameter bound on the underlying graph of ( $X, Q, \pi$ ) using the same technique from [LMP18] (see Theorem 18.1 for comparison).

Theorem 23.3 (Diameter bound). Let $(X, Q, \pi)$ be an irreducible and reversible finite Markov chain with strictly positive entropic Ricci curvature $\operatorname{Ric}(Q) \geq K>0$. Then the diameter of the induced weighted graph $X$ is bounded from above by

$$
\begin{equation*}
\operatorname{diam}(X) \leq \frac{2}{K} \sqrt{\frac{D \log D}{D-1}} \tag{23.6}
\end{equation*}
$$

where $D$ is the maximal vertex $\pi$-degree,

$$
D:=\max _{x \in X} \operatorname{Deg}_{\pi}(x)=\max _{x \in X} \frac{1}{\pi(x)} \sum_{y \in S_{1}(x)} \pi(y)
$$

Remark 23.4. We previously discuss that the entropic Ricci curvature of the hypercube $\mathcal{Q}^{n}$ is sharply estimated by $\frac{2}{n}$. Therefore, in view of the hypercube $\mathcal{Q}^{n}$, the bound (23.6) is not optimal:

$$
n=\operatorname{diam}\left(\mathcal{Q}^{n}\right) \leq n \sqrt{\frac{n}{n-1} \log n}
$$

Proof of Theorem 23.3. Consider a function $f \in \mathbb{R}^{X}$ given by $f(x):=d\left(x, x_{0}\right)$ for an arbitrary reference point $x_{0} \in X$. Since $f$ is a 1-Lipschitz function, it follows that $2 \Gamma f(x)=\sum_{y \in S_{1}(x)} Q(x, y)(f(y)-f(x))^{2} \leq 1$ for all $x \in X$, i.e., $\|2 \Gamma f\|_{\infty} \leq 1$, which then implies $\left\|2 P_{t}(\Gamma f)\right\|_{\infty} \leq\|2 \Gamma f\|_{\infty} \leq 1$.

Moreover, the Cauchy-Schwartz inequality and the inequality (23.5) yield

$$
\begin{align*}
\left|\Delta P_{t} f(x)\right|^{2} & =\left(\sum_{y \in S_{1}(x)} Q(x, y)\left(P_{t} f(y)-P_{t} f(x)\right)\right)^{2} \\
& \leq\left(\sum_{y \in S_{1}(x)} Q(x, y)\left(P_{t} f(y)-P_{t} f(x)\right)^{2}\right)\left(\sum_{y \in S_{1}(x)} Q(x, y)\right) \\
& =2 \Gamma\left(P_{t} f\right)(x) \leq e^{-2 K t} c\left\|2 P_{t}(\Gamma f)\right\|_{\infty} \leq c e^{-2 K t} \tag{23.7}
\end{align*}
$$

From the fundamental theorem of calculus and the definition of $P_{t}$, we then obtain

$$
\left|f(x)-P_{T} f(x)\right| \leq \int_{0}^{T}\left|\frac{\partial}{\partial t} P_{t} f(x)\right| d t=\int_{0}^{T}\left|\Delta P_{t} f(x)\right| d t \leq \int_{0}^{T} \sqrt{c} e^{-K t} d t \leq \frac{\sqrt{c}}{K}
$$

which holds true for all $T>0$.
Moreover, (23.7) implies that $\left|P_{t} f(y)-P_{t} f(x)\right| \rightarrow 0$ as $t \rightarrow \infty$ for all edges $y \sim x$, and the same result is extended for all pairs of $x, y$ which are not necessarily neighbors.

Passing to the limit $T \rightarrow \infty$, we conclude from triangle inequality that

$$
\begin{aligned}
d\left(x, x_{0}\right) & =\left|f(x)-f\left(x_{0}\right)\right| \\
& \leq\left|f(x)-P_{T} f(x)\right|+\left|f\left(x_{0}\right)-P_{T} f\left(x_{0}\right)\right|+\left|P_{T} f(x)-P_{T} f\left(x_{0}\right)\right| \\
& \leq \frac{2 \sqrt{c}}{K}
\end{aligned}
$$

which implies the desired diameter bound.

## Chapter 24

## Outlook for Erbar-Maas entropic Ricci curvature

Let us finish this Part IV with further questions about the entropic Ricci curvature on discrete Markov chains.

1. Curvature in general is a local concept. Ollivier Ricci curvature is defined on edges of graphs, while Bakry-Émery is defined on vertices. The entropic Ricci curvature, on the other hand, is defined to be the global lower bound of Ricci curvature for whole spaces. Previously, we mentioned about the possibility to define the entropic Ricci curvature "pointwise on $\mathscr{P}(X)$ ", namely $\operatorname{Ric}_{Q}(\rho)$ for each $\rho \in \mathscr{P}(X)$. However, it is unlikely that this local property can be pushed further to "pointwise or locally on $X$ " $\left(\operatorname{since}^{\operatorname{Ric}}{ }_{Q}(\mu)=+\infty\right.$ for any $\mu$ which does not have full support in the case of the logarithmic mean). However, it would be interesting to investigate whether there are measures $\rho_{0} \in \mathscr{P}_{*}(X)$ with

$$
\operatorname{Ric}_{Q}\left(\rho_{0}\right)=\inf _{\rho \in \mathscr{P}(X)} \operatorname{Ric}_{Q}(\rho),
$$

and if they exist, to provide a characterization of them.
2. To derive the Bonnet-Myers type diameter bound in Theorem 23.3, our approach applies the gradient estimate only to those $\rho$ with local structure, namely $\rho=1_{x}+\varepsilon \sum_{y \in S_{1}(x)} 1_{y}$. This could be a main reason that this diameter bound is likely not optimal. It would be interesting to investigate whether utilizing non-local measures could lead to a better diameter bound, which possibly becomes sharp in the case of a simple random walk on hypercubes.

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[^0]:    ${ }^{1}$ A vertex $x$ satisfying this property is called $\boldsymbol{S}_{1}$-in regular.

