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## Interphase Dynamics of photoanode in Micro-energy generation in Photovoltaics

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#### Abstract-

Dye-sensitized solar cells (DSCs)as photovoltaics have been subject of intensive research in the last two decades owing to their outstanding features of relative thermal stability, low cost, optical tunability, ability to operate under conditions of poor lighting, and environment friendliness as such, they offer a promising substitute for conventional monocrystalline silicon solar cells. Their poor output efficiency has been incessantly associated with grain boundaries, good quality films are associated with lower recombination decays. and it has been claimed that grain boundaries have a negligible influence on performance. We used microscopy interrelated with scanning electron microscopy to resolve spatially the recombination dynamics from organic films of Murraya Koenigii. Stoichiometric organic films, varied in different grains even for the same film. In regions where grain boundaries were dimmer, nonradiative decay occurred faster. Ultraviolet spectroscopy showed positive correlation consistent with photon absorption and regions of the electromagnetic spectrum. The photovoltaic efficiency of M. Koenigii DSCs based on TiO2 sensitized with ionic electrolytes was compared. The output efficiency revealed that the ionic potential of Br- was more than 200% value of Cl-, and a final end reaction of above 100% that of I- under similar conditions of atmosphere and experimental set-up.

Key words: Energy harvesting, Electron Occupancy, Affine correction

#### 1. Introduction

In DSSC, photoanode which essentially comprises of a metal oxide semiconducting base is one of most important components for enhancing sunlight absorption, electron excitation and light harvesting efficiency [1]. Photoanode material plays a deterministic role in charge carrier kinematics, dye loading, electron injection, and subsequent power conversion efficiency due to photocurrent and voltage [2]. In view of this, intense research focus is being centered on details of design, chemical composition, structure and morphology of the semiconductor photoanode [3]. Progressive investigations on photoanode type has evolved massive deployment of myriads of photoanode materials such as Nb<sub>2</sub>O<sub>5</sub>, SnO<sub>2</sub>, ZnO, TiO<sub>2</sub>, Al<sub>2</sub>O<sub>3</sub>, graphene, gold and so on [4].

A central theme for this diversity aims at high efficiency output thus, interrelationship of photoanode materials in conjunction with other DSC mechanisms have also become an object

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of great scientific work and scrutiny. In effect, morphology, structure and optical properties of metal oxide semiconducting materials remain pertinent subject for research enquiry [5-6].

Influence of ionic strength on thermodynamic environment and properties is also of topmost priority on performance of metal oxides of semiconductors. Inception of instability in colloids, mutual Brownian diffusion quotient of colloidal particles are factors which compete for special attention owing to competition in the choice of particles or ions to be adsorbed. In effect, there is a perceptible decline in particle potential, charge carrier volume and ionic strength between molecular layers. In general, increase in ionic concentration promotes efficient separation into colloidal phases, which accentuates the communal diffusivity, resulting in an appreciable increase in photovoltaic performance [7].

In order to enhance photoelectric output, semiconductors with wide bandgaps are considered. Foremost amongst such photoanode applications are TiO<sub>2</sub> and ZnO nanomaterials because, their wide band gap enables them create highly impressive interfacial, dynamic and structural properties with other DSC components both at room temperature and high sintering experimental conditions [8]. TiO<sub>2</sub> photoanode gained universal acceptability due to availability, low-cost, biocompatibility and non-toxicity. A great significance is attributed to structural isomerism of ZnO photoanode material, encompassing abundant nanostructure morphologies with high electron mobility in nano-particle film and in bulk [9]. In practice, ZnO has a similar conduction band edge and work function relative to TiO<sub>2</sub>, but has better charge carrier mobility than TiO2. This has made ZnO in some instances, become regarded as a more promising photoanode materials for DSCs [10]. Conversely, peculiar problems of ZnO instability in acidic medium, and development of aggregates with dyes on its surface depreciate its photoelectric performances [11]. A faster rate of electron transport rate is directly proportional to reduction in electron-hole recombination rate and enhanced conversion efficiency. Consequently, heterojunction and hybrid photoanodes are being used extensively to counteract the limitations of a monotype photoanode [12]. Some other semiconductor materials have gained influence for their higher photovoltaic ability, examples of such include, Zn2SnO4 which has a more positive conduction band edge position, high versatility and multi-functionality compared to TiO2 [13]. Heterogenous nanomaterials exhibit vast practical and technological applications in light emitting diodes, chemical and gas sensors and present opportunities to re-define optoelectronics, photocatalysis, memory resistors, pyroelectric and piezo-electronic spheres with their optical tunable features [14]. Thus, this study seeks to explore the effect of ionic variation on semiconductor oxide of TiO<sub>2</sub> nanomaterial.

#### 2. Experimental Methods

In standard procedure, 50 g of  $\text{TiO}_2$  Degussa of commercial variety, was dissolved in 20 ml of 1.0M conc. trioxonitrate (V) acid and then 3 drops of ethanol were slowly added to the colloidal solution with continuous stirring, to constitute the  $\text{TiO}_2$  photoanode. Afterward, the entire reaction mixture was transferred into air tight reagent bottles where they were consequently applied by the doctor blade technique on indium doped tin oxide (ITO) conducting glass. The photoanode was affixed onto the ITO through high temperature sintering at  $450^{\circ}\text{C}$  for 1.5 h and allowed to subsequently cool down to room temperature. Dye

solution comprising 1ml: 100 of 1.0M conc. ethanol was sprayed onto the photoanode surface. The counter electrode was fabricated by masking an ITO with epithelial layers of soot over a naked Bunsen flame and then allowed to assume constant temperature. A pair comprising of photoanode and counter electrode were fastened together with binder clips, two drops of different ionic electrolytes constituted by standard practice; 1 g potassium iodide, potassium bromide and potassium chloride dissolved in 100 ml of deionized water was introduced inbetween the ITO. This experimental set-up was connected in parallel to variable load and a digital multimeter to obtain readings under standard conditions of 1 atmosphere [15].

#### 3. Results and discussion

Ultraviolet spectroscopy: The vibrionic transitions that occur in M.Koenigii molecules being excited by photons of energy is responsible for electronic transitions illustrated by Figure 1. Under standard conditions of 1 atmosphere, at room temperature, the molecule begins from the ground state where energy v=0, and consequent of absorption of quantized energy, transits to an excited energy state. This new state may cause a shift in the equilibrium of M.Koenigii nucleus, re-configuring its electron configuration. In Figure 1, this shift is depicted as different strata,  $E_0$  is the ground,  $E_1$  first excited state and  $E_2$  the highest excitation state. The distinct energy levels are representing different cartesian coordinates, characterized by internuclear separation [16]. A probability of a particular molecule ending up in a definite energy level is proportional to square of the vertical displacement of vibrational wavefunctions of  $E_1$  and  $E_2$  [17]. According to quantum principle,  $E_2$  excited state molecules can rapidly decrease to  $E_1$  vibrational level and decay to  $E_0$ , electronic ground state by photon emission. The latter is known as a relaxation process causing the release of phonons, this is a distinctive property of the lattice vibration [18].

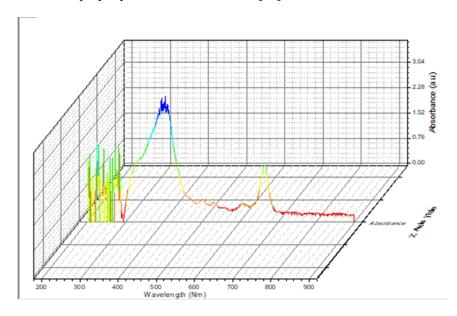


Figure 1: Molecular excitation of M. Koenigii coordinates with color shades

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Scanning Electron Microscopy (SEM): The performance on nanoscale was investigated using scanning electron microscopy with Gwydion software. In Figure 2 (a), electron seats are represented by the red dots, carrier transport is largely by hopping from one seat to the next. A large percentage of the electron seats are present in the mid-section, carrier transport will be higher in this region. At the fringes, the blue zone dominates, region where electron tunneling activities are promoted. Examination of result from surface probe of M.Koenigii is illustrated in Figure 2 (b) within  $\phi$  (x):  $\pm$  54.8 which corroborates Figure 2 (a) spectroscopy. Further spectroscopic investigation along the grains of M.Koenigii is revealed in Figure 3 (a), the etch observed transversely efficiently separates the mid-section of the micrograph from the boundaries. In order to obtain the coordinate geometry as shown on Table 1 of the grain boundaries, the affine correction was plotted as shown in Figure 3 (b). Figure 3 (a) shows the scalar proportions while Figure 3 (b) presents the vectoral dimensions in M.Koenigii nanocomposite.

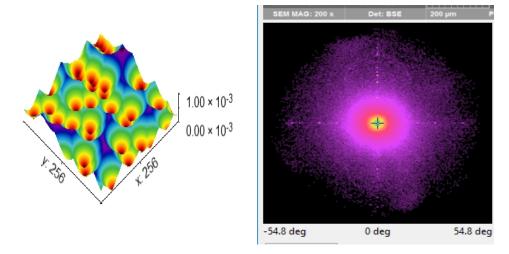


Figure 2 (a): *M.Koenigii* Electron occupancy and (b) Facet Analysis at 3.000 <sup>0</sup> tolerance level and facet plane size 3*px* 

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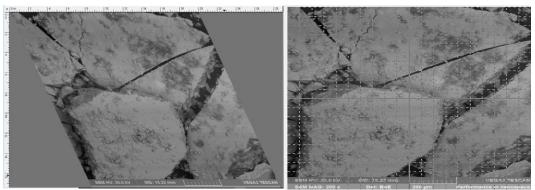
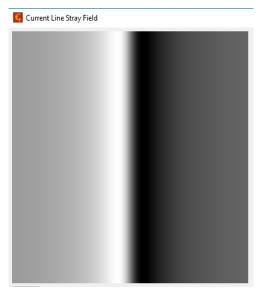


Figure 3 (a): Co-Planar spectroscopy of SEM and (b) Affine Correction of M.Koenigii

Table 1: Coordinates of lattice vectors of *M.Koenigii* plotted with Gwydion

Data Process	Lattice Vec						
Affine Correction	x (µm)	y (µm)	$\Phi$ (0)	Len	gth(	(μm)	
	a <sub>1</sub> : 5.000	-0.000	-0.00	5.00	00		
	a <sub>2</sub> : 0.000	5.39	90.00	5.39	91		
Lattice type: $a_1(1)$ , $\phi$ (2.035°)	$n  \Phi^{(0)}$	$\mathbf{O}(0)$	x(µm)	y(µ	$y(\mu m)   z(\mu m)$		
	1 36.47	169.12	-0.584	0.1	0.112 0.804		
Estimated Affine Correction	x (m)	y (m)	Length (	(m)	Angle (0)		
	a <sub>1: 1.0000</sub>	-	1.0000		-0.00		
		0.0000					
	a2: -0.4481	0.8940	1.0000		11	6.62	

Magnetic Flux Microscopy in *M.Koenigii*: The current line stray field investigated a graphical simulation of magnetic lines of force distribution in *M.Koenigii* particle with four domains acting under the influence of a field pulsation. These lines correlate with sphere of magnetism which consists of parallel components as depicted by the hue of shades to the long and short sides of the particle and are denoted by full and dashed lines, respectively. In Figure 4(a), a pictorial form of algorithms for stray field calculation using Fourier transform is illustrated. It is used to modify the field from one-layer height to another in multiple modular times to simulate a planar coil shown as two lines with opposite directions of the current. In order to optimize the charge carrier flow, a Hough space of *M.Koenigii* current field line is grouped into uniform sub-sections. The effect is a 2-D lattice defined by parameters along the x and y



cartesian coordinate system as shown in Figure 4(b). The result is an average vote for each electric field current line, resulting in the cluster observed which coincides with the pattern of electric flux of *M.Koenigii*.

Figure 4:

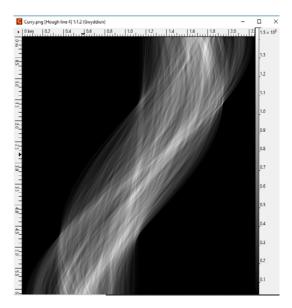
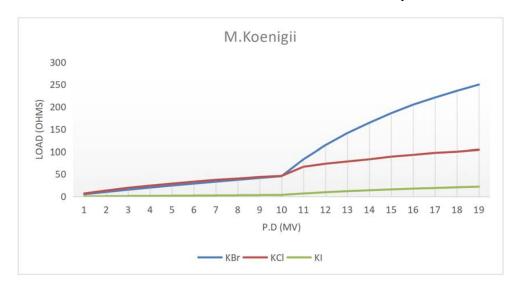


Figure 5: (a): Current line stray field and (b) Hough line of *M.Koenigii* Photoelectric Evaluation of *M.Koenigii* DSCs: The effect of sensitizing *M.Koenigii* with three different electrolytic solutions comprising Br<sup>-</sup>, Cl<sup>-</sup> and I<sup>-</sup> is illustrated in Figure 5. Under the same operating conditions, the kinematics of the reaction based on similar TiO<sub>2</sub> photoanode nanomaterial, is decidedly disparate. Br<sup>-</sup> shows a direct proportional relationship between increase in load and the ionic potential, this is a promising result for suitable application in photosensors. The activation energy required by I<sup>-</sup> is very large thus, the reaction kinematics does not favor the transport of I<sup>-</sup>. It is notable that the reaction of Cl<sup>-</sup> and Br<sup>-</sup> was the same at the initial stage then, Br<sup>-</sup> was preferred in the redox reaction whereas Cl<sup>-</sup> maintained almost the same level of ionic potential.



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Figure 6: Photoelectric response of *M.Koenigii* DSCs to different ions

#### 4. Conclusions

The use of grain boundary exposed the margin of active reaction in *M.Koenigii* DSCs, Br complements molecular kinematics in *M.Koenigii*. A preliminary check provided by the current stray line launches a novel pathway to boost the subsequent output performance of Br and every other DSC. Kinematics and Ionization potential of reactions are crucial determinants of charge transport in DSCs, future researches would reveal more insights into these dual phenomena. Also, co-sensitization of the ionic electrolyte sensitizers could be used in future applications on TiO<sub>2</sub> photoanode and hybrid photoanodes for subsequent application in high efficiency optoelectrical devices.

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