

Structural properties of discs and bulges of early-type galaxies

Roelof S. de Jong,^{1*} Luc Simard,² Roger L. Davies,³ R. P. Saglia,⁴ David Burstein,⁵ Matthew Colless,⁶ Robert McMahan⁷ and Gary Wegner⁸

¹*Space Telescope Science Institute, 3700 San Martin Drive, Baltimore, MD 21218, USA*

²*Herzberg Institute of Astrophysics, National Research Council of Canada, Victoria, BC V9E 2E7, Canada*

³*University of Oxford, Astrophysics, Keble Road, Oxford OX1 3RH*

⁴*Max-Planck Institut für extraterrestrische Physik, Giessenbachstrasse D-85748, Garching, Germany*

⁵*Department of Physics and Astronomy, Arizona State University, Tempe, AZ 85287-1504, USA*

⁶*Anglo-Australian Observatory, PO Box 296, Epping, NSW 1710, Australia*

⁷*Department of Physics and Astronomy, University of North Carolina, CB#3255 Phillips Hall, Chapel Hill, NC 27599-3255, USA*

⁸*Department of Physics and Astronomy, Dartmouth College, Wilder laboratory, Hanover, NH 03755, USA*

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ABSTRACT

We have used the EFAR sample of galaxies to investigate the light distributions of early-type galaxies. We decompose the two-dimensional light distribution of the galaxies in a flattened spheroidal component with a Sérsic radial light profile and an inclined disc component with an exponential light profile. We show that if we assume that all galaxies can have a spheroidal *and* a disc component, then the brightest, bulge-dominated elliptical galaxies have a fairly broad distribution in the Sérsic profile shape parameter n_B , with a median of approximately 3.7 and with $\sigma \sim 0.9$. Other galaxies have smaller n_B values. This means that spheroids are in general less concentrated than the de Vaucouleurs $R^{1/4}$ -law profile, which has $n_B = 4$.

While the result of our light decomposition is robust, we cannot prove without kinematic information that these components are spheroids and discs, in the usual sense of pressure- and rotation-supported stellar systems. However, we show that the distribution of disc inclination angles is consistent with a random orientation if we take our selection effects into account. If we assume that the detected spheroids and discs are indeed separate components, we can draw the following conclusions: (1) the spheroid and disc scale sizes are correlated; (2) bulge-to-total luminosity ratios, bulge effective radii and bulge n_B values are all positively correlated; (3) the bivariate space density distribution of elliptical galaxies in the (luminosity, scale size)-plane is well described by a Schechter luminosity function in the luminosity dimension and a lognormal scale-size distribution at a given luminosity; (4) at the brightest luminosities, the scale size distribution of elliptical galaxies is similar to those of bright spiral galaxies, but extending to brighter magnitudes; at fainter luminosities the scale size distribution of elliptical galaxies peaks at distinctly smaller sizes than the size distribution of spiral galaxies; and (5) bulge components of early-type galaxies are typically a factor of 1.5–2.5 smaller than the discs of spiral galaxies with a slight luminosity dependence, while disc components of early-type galaxies are typically twice as large as the discs of spiral galaxies at all luminosities.

Key words: galaxies: elliptical and lenticular, cD – galaxies: fundamental parameters – galaxies: statistics – galaxies: structure.

1 INTRODUCTION

Studying the light distribution of galaxies in order to determine and compare their structural properties has a long tradition (e.g. de Vaucouleurs 1948; Fish 1964; Freeman 1970). Traditionally,

the light distributions of elliptical galaxies have been fitted by the $R^{1/4}$ -law profile of de Vaucouleurs (1948), while spiral galaxies have been fitted with exponential light profiles for their discs and often, due to their similarities to early-type galaxies, $R^{1/4}$ -law bulges (e.g. de Vaucouleurs 1959; Kormendy 1977). In the past decade we have seen some shift from this paradigm. A generalization of the $R^{1/4}$ and exponential light profile laws described first by Sérsic (1968), resulting in an $R^{1/n}$ law, has more often been fitted to early-type

*E-mail: dejong@stsci.edu

galaxies to account for observed deviations from the $R^{1/4}$ profiles (e.g. Caon, Capaccioli & D’Onofrio 1993; Graham & Colless 1997). Based on this development and based on theoretical predictions for some forms of bulge formation (e.g. Combes et al. 1990; Pfenniger & Norman 1990), more general forms of light profiles have been fitted to the bulge light of spiral galaxies as well (Andredakis, Peletier & Balcells 1995; de Jong 1996a; MacArthur, Courteau & Holtzman 2003). With the realization that many early-type galaxies have disc-like components (e.g. Bender et al. 1989; Rix & White 1990; Jørgensen & Franx 1994), bulge/disc decompositions of elliptical galaxies have been performed, but in general with $R^{1/4}$ bulges and exponential discs (e.g. Saglia et al. 1997b) except for a few notable exceptions on small samples of galaxies (D’Onofrio 2001; Balcells, Graham & Peletier 2004; Gutiérrez et al. 2004). Here we perform two-dimensional (2D) bulge/disc decompositions on the large EFAR sample of early-type galaxies (Wegner et al. 1996, hereafter Paper I) using an $R^{1/n}$ spheroid and exponential disc light profiles, determine relations between their structural parameters and their frequency of occurrence, and compare those to the properties of disc-dominated galaxies.

In an earlier paper in this series (Saglia et al. 1997a, Paper III hereafter) we presented the luminosity profiles for the 776 galaxies observed in the EFAR project to which we fitted seeing convolved $R^{1/4}$ -law bulge plus exponential disc models. In Paper III, 31 per cent of the sample proved to be spiral or barred spiral galaxies and 69 per cent of the sample were classified as early-type galaxies. Of those early types 18 per cent were well fitted by an $R^{1/4}$ alone, the great majority (70 per cent) were very much better fitted by the combination of an exponential disc plus an $R^{1/4}$ bulge than by a bulge alone. The remaining 12 per cent of the early-type galaxies were classified as cD, without reference to their profile type, as they were the brightest in their cluster and had half-total-light radii larger than 10 kpc.

In Saglia et al. (1997b, hereafter Paper IV) we showed that the combination of an $R^{1/4}$ -law spheroid plus an exponential disc is degenerate to an $R^{1/n}$ -law spheroid, unless one has data for a very large range in surface brightness and radius. Given that only 12 per cent of the EFAR galaxies were best fitted with an $R^{1/4}$ law alone, and that the remainder of the galaxies were better fitted with an $R^{1/4}$ bulge and an exponential disc, one might wonder about the validity of the $R^{1/4}$ law in general. To address this issue, we have refitted all EFAR galaxies again with an $R^{1/n}$ -law spheroid and an exponential disc. If early-type galaxies are dominated by $R^{1/n}$ spheroids, most fits should converge to largely spheroid-dominated systems with small discs. If the $R^{1/4}$ -law spheroids plus disc picture is correct, most fits should converge towards spheroids with a Sérsic parameter n close to 4. Obviously, the truth could be somewhere in between with spheroids with Sérsic parameter n covering a large range and with a large range in bulge-to-disc ratios.

We have to realize that even when we find that most early-type systems have light components that are well described by exponential light profiles, this in itself will never prove the existence of exponential discs. First of all, there is no reason why the exponential discs observed in most spiral galaxies should be present in a similar way in bulge-dominated systems, especially realizing that exponential discs in spiral galaxies are still poorly understood. Secondly, even if light is present in an exponential distribution, it might not be in a disc configuration flattened by rotation as discs of spiral galaxies are normally understood. One can probably make many stable configurations in bulge-like potentials that have small exponential-like stellar distributions on top of them, without the need for any

net rotation. The best confirmation of exponential light distributions in early-type galaxies being rotationally flattened discs as in spiral galaxies has to come from measurements of their kinematics.

Even when the kinematics is available, it still might be difficult to reach an unambiguous result. Scorza & Bender (1995) made dynamical models for a small sample of early-type galaxies based on their photometric bulge/disc decompositions and showed that the observed asymmetries in the long-slit stellar absorption-line profiles were consistent with the dynamical model predictions. However, with their data and modelling technique they could not prove this was a unique kinematic interpretation. Better kinematic constraints on disc components can be obtained with full 2D velocity profile measurements, such as those obtained with the SAURON project (Emsellem et al. 2004; Falcón-Barroso et al. 2004). However, so far these studies have been limited to the high surface brightness central regions of galaxies.

In recent years a number of studies have been performed where the galaxy samples were large enough and the selection bias understood well enough that bivariate space density distributions of structural parameters could be calculated. Cross et al. (2001) presented the bivariate (luminosity, surface brightness distribution) for the two-degree Field Galaxy Redshift Survey (2dFGRS) (Colless et al. 2001b). Blanton et al. (2003) showed many different bivariate space density distributions derived from the Sloan Digital Sky Survey (SDSS; York et al. 2000), and presented some distributions in bins of colour and concentration index to split their sample into different galaxy types. Kauffmann et al. (2003) derived mass-to-light ratios for the SDSS galaxies and presented stellar mass bivariate distributions. Shen et al. (2003) also investigated the bivariate distribution of structural parameters of different subsets of SDSS galaxies. However, even though calculated, they did not include the absolute space densities in their analysis. While based on a large sample to provide good statistics, in none of these studies were bulge/disc decompositions performed to derive space density distributions for the individual components. Such an analysis was performed on the discs of a sample of approximately 1000 Sb-Sdm galaxies by de Jong & Lacey (2000). However, no such analysis of early-type galaxies has appeared so far.

A fundamental question for galaxy formation is the origin of the Hubble sequence. Why do some galaxies have discs and others not? If in fact almost all galaxies have a disc component but one of varying mass, luminosity and size, the question of the origin of the Hubble sequence takes on a different perspective. In semi-analytic models of galaxy formation early-type galaxies go through phases where they have substantial discs which are disrupted and re-built during the cycle of hierarchical merging. Determining galaxy disc and bulge parameters and their bivariate distributions will help us to test and constrain such hierarchical semi-analytic models of galaxy formation.

In this paper we will use a 2D fitting algorithm to measure the bulge and disc properties in the EFAR sample of early-type galaxies. We estimate the frequency, luminosity and size of discs and bulges and explore the properties of the discs in the context of S0 and spiral galaxy discs.

In the next section we describe our method of 2D model fitting and parameter error estimation and compare our new results with the results of Paper III. In Section 3 we characterize the properties of the discs and their relation to the host bulge component. Our conclusions and the implications of this work for models of galaxy evolution are summarized in Section 4. We use a cosmology of $\Omega_m = 0.3$, $\Omega_\Lambda = 0.7$ and $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ throughout.

2 GIM2D: 2D GALAXY MODELLING

We decided to use full 2D modelling of the galaxy images, as this provides extra bulge/disc separation constraints due to the different flattening of bulges and discs, sometimes combined with the constraints resulting from the different position angles of the two components. Performing the modelling in 2D also allows proper weighting of data points reflecting uncertainties due to photon statistics, flat-fielding uncertainties, sky level uncertainties and model imperfections (i.e. real galaxies are not as smooth as our models).

To extract the best possible bulge and disc parameters and to estimate their errors we performed the following steps.

Many of the EFAR galaxies have several repeat observations, so for the 775 objects more than 2500 images are available. We only used the *R*-band images, because the *B*-band imaging of the EFAR sample is incomplete. We determined the seeing on all EFAR images by finding isolated star-like objects in each of them, measuring the FWHM of each object by fitting a Gaussian profile and taking the error weighted average of the eight objects in each image with the smallest errors after removing the objects with smallest and largest FWHM. At this stage we selected for each galaxy the image with the best seeing. All galaxies with too few good foreground stars (<3) to make a reliable seeing measurement were removed from the sample (23 galaxies).

SEXTRACTOR (Bertin & Arnouts 1996) was used to create a mask of all objects on the images, splitting the images of overlapping objects into separate components. This mask is used by the fitting routine to exclude all pixels with objects detected except for the pixels with the galaxy of interest and the sky pixels. A very low level SEXTRACTOR mask was also used to determine the sky level of the image, using a kappa-sigma clipped median of the pixels deemed free of objects, and this sky level was kept fixed during the fitting of the galaxy.

The 2D modelling of the galaxy images was performed by the publicly available GIM2D fitting package (version 2.2.1), described in detail in Simard et al. (2002). The spheroidal component of each galaxy was fitted with a Sérsic law:

$$\Sigma_B(r) = \Sigma_e e^{-c[(r/R_{e,B})^{1/n_B} - 1]}, \quad (1)$$

where the effective radius ($R_{e,B}$) encloses half the total luminosity and Σ_e is the surface brightness at this radius. The parameter c is set to $1.9992n_B - 0.3271$ so that $R_{e,B}$ remains the half total light radius of the bulge. Better approximations of c were published after we performed our calculations (e.g. MacArthur et al. 2003), but these improvements are only relevant for n_B values much less than 1, which are rare in our sample. The bulge component was allowed a flattening with ellipticity $\epsilon = 1 - b/a$, where a and b are the semi-major and minor axis, respectively. The $R_{e,B}$ provided by GIM2D is the bulge half-total-light radius of the major axis, which we multiplied by $\sqrt{1 - \epsilon}$ to be consistent with our previous work, which uses the effective radius on a circular aperture.

The disc was modelled using an exponential disc:

$$\Sigma_D(r) = \Sigma_0 e^{-r/h}, \quad (2)$$

which was assumed to be infinitely thin, transparent and to have an inclination i . The two galaxy components had a common centre which was fitted as well, and an independent flux, here expressed in as a bulge-to-total light ratio (B/T). Each component also has an independent position angle. The summed galaxy components were convolved with a Gaussian with a FWHM equal to the determined seeing, using a fast Fourier transform (FFT) algorithm. GIM2D fits this large 11-dimensional parameter space for the two components using the Metropolis algorithm (Metropolis et al. 1953). The

Metropolis algorithm is very CPU intensive, but it tends to be more robust than methods based on gradient searches. The Metropolis algorithm is ideal for cases such as the current one where many fits have to be performed unsupervised. The CPU intensiveness was solved by using 16 fast Linux dual processor PCs in a cluster.

The fitted scale size parameters in pixels and intensities in counts were converted to arcsec and magnitudes using calibrations described in Paper III. To set the nomenclature, the galaxies are parametrized in the remainder of the paper as follows: for the total galaxy we have luminosity M_{tot} , effective (half-total-light) radius R_e and effective surface brightness within this radius (SB_e); for the bulge luminosity M_B , effective radius $R_{e,B}$ and effective surface brightness ($SB_{e,B}$); for the disc luminosity M_D , scalelength h and central surface brightness μ_0 .

After all galaxies had been fitted, galaxies with high- χ^2 residuals and/or suspicious model-subtracted residual images were identified. In many of these cases SEXTRACTOR had not detected objects near the nucleus of the target galaxy or performed an insufficient splitting of nearby objects, resulting in incorrect masks and causing poor fits. For most of these objects we were able to modify the masks by hand to give a satisfactory result, but a number of objects with bright foreground stars, many bright companions or the occasional galaxies with double nuclei were removed from the sample at this stage (20 galaxies).

We used the model-subtracted residual images to identify galaxies with spiral structure. Even though not necessarily spiral galaxies based on their B/T ratio, these galaxies are marked with special symbols in the diagrams. We will use the term ‘non-spiral galaxies’ in the remainder of this paper for those galaxies that had no clear spiral structure in the residual images (454 of the 558 galaxies used in the final selection). The brightest galaxy of each of the clusters with at least three members as identified in Paper VII of this series (Colless et al. 2001a) is also marked with a special symbol in the diagrams.

The GIM2D fitting package provides 68 per cent confidence limits on the fitted parameters, based on the topology of the parameter space that is built up during the fitting process. These errors only take the uncertainties in the fitted parameters into account, not uncertainties in fixed parameters such as the sky level and seeing. As errors in these parameters are often the main error in the determined bulge and disc parameters (de Jong 1996a), we repeated all fits with the maximum error expected in our determination of sky and seeing values.

We estimated the typical error in our sky estimates by dividing the image into four and measuring the sky in each of the quadrants separately using the same procedure as for the total image (using object masking and kappa-sigma clipping). The rms in the four sky averages provides an estimate in the sky level uncertainty on a galaxy size scale. This is probably a slight overestimate, as most galaxies are near the centre of the frame, so near the average sky of the total frame, while the sky rms of the quadrants will be more affected by sky level gradients across the frame. Most images used have small sky level uncertainties: 80 per cent of the galaxies have quadrant errors of less than 0.5 per cent, only 8 per cent of the galaxies have quadrant rms values larger than 1 per cent of the sky, mainly due to sky gradients. We repeated all GIM2D fits with 1 per cent sky added and subtracted, and the resulting fitting parameters should be seen as the maximum errors in the parameters that can be expected to be caused by sky uncertainty.

For each galaxy we repeated the fit with the seeing increased and decreased by 2 per cent to determine the errors in the fitted parameters due to seeing measurement uncertainties. Judging from

the rms in the seeing measurements of different stars in the different frames, less than 20 per cent of the galaxies have uncertainties in their seeing larger than 2 per cent. Again, these errors should be seen as the maximum error in the fitted parameters due to seeing measurement uncertainties.

Where appropriate, we show typical errorbars on our graphs, which are the median errors for all galaxies in the determined parameters, where the parameter error for an individual galaxy is the maximum error due to the fitting parameter space, sky uncertainty or seeing uncertainty. We chose the maximum error here, because it is hard to assign relative weights to the different error sources. To assign proper weights would require a full set of Monte Carlo simulations, which is prohibited given the large computation times involved.

We excluded all galaxies with no measured redshift from our analysis, mainly spiral galaxies which were excluded from the spectroscopy phase of the EFAR project as being not interesting for the main aim of the project. The final sample of galaxies used contains 558 galaxies.

2.1 Testing the 2D fitting routine

We tested the GIM2D fitting program by performing fits on 500 artificial galaxy images with characteristics similar to our real data. The images were 500×500 pixel² and the pixel size was 0.65 arcsec. The sky level was set at 20.5 mag arcsec⁻². Galaxy images were created in a Monte Carlo fashion with total magnitudes sampled linearly between 12 and 16 mag, with B/T ratios sampled linearly between 0.05 and 1, and having bulges with $18 \leq \langle SB_{e,B} \rangle \leq 22$ and $1 \leq n_B \leq 6.5$, and discs with $20 \leq \mu_0 \leq 23$. The CCD read noise and gain were set comparable to most of our observations and Poisson noise was added appropriately for the detected flux level. The galaxies were convolved with a Gaussian seeing with a FWHM of 1.8 arcsec, which is at the high end of the typical seeing of the used images (90 per cent with a seeing of less than 1.8 arcsec).

A comparison of the fitted GIM2D parameters with the input parameters is presented in the first five columns of Fig. 1. All parameters are reasonably well reproduced, but some systematic trends can be seen. As could be expected, the largest errors in the bulge

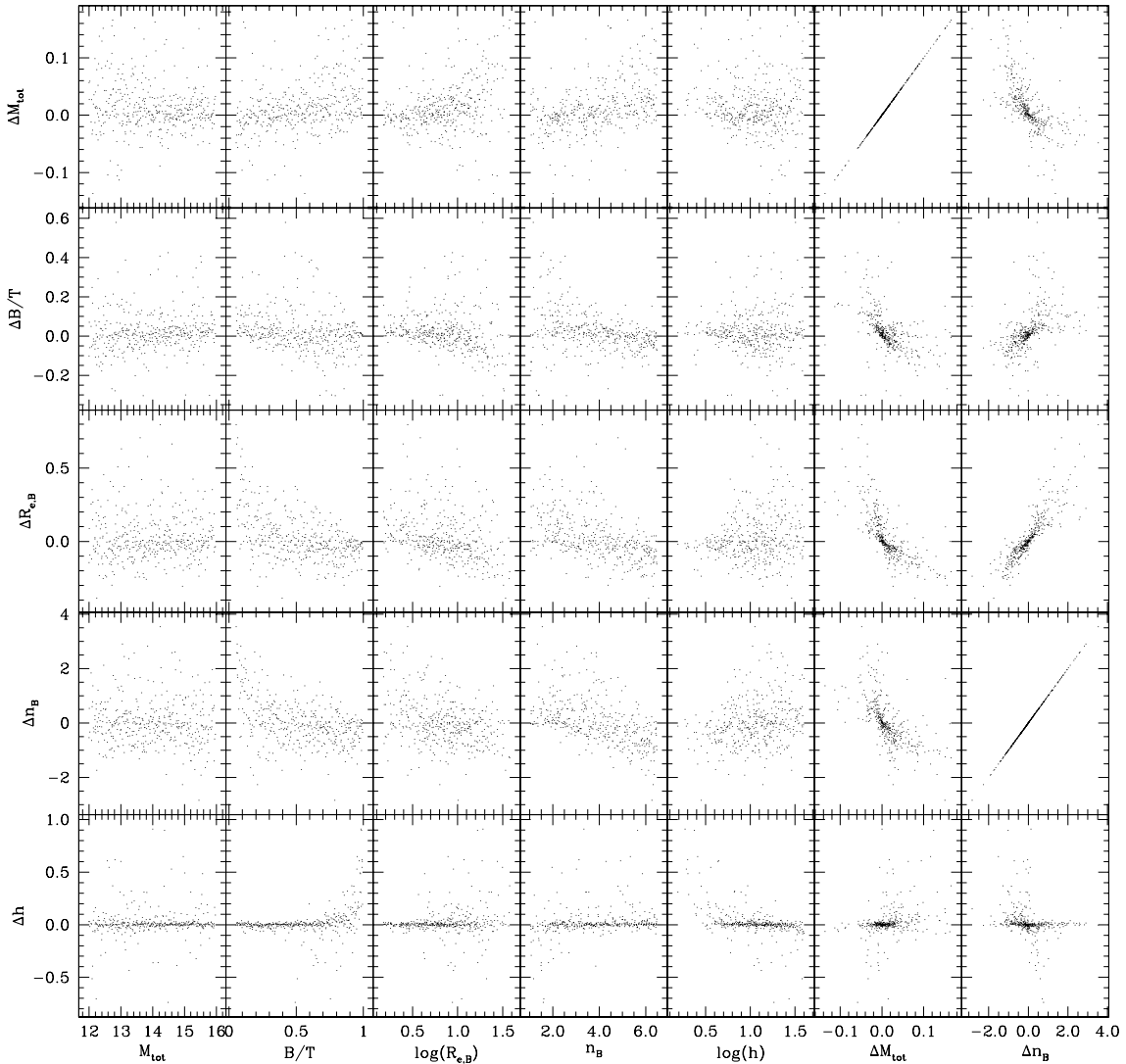


Figure 1. Results of GIM2D fits to the Monte Carlo simulated artificial galaxy images. The first five columns on the x-axis show the input parameters of the galaxy images, the y-axis and the last two columns on the x-axis show the differences between fitted and input parameters, defined as $\Delta M_{\text{tot}} = M_{\text{tot}}^{\text{fit}} - M_{\text{tot}}^{\text{in}}$, $\Delta B/T = B/T^{\text{fit}} - B/T^{\text{in}}$, $\Delta n_B = n_B^{\text{fit}} - n_B^{\text{in}}$, $\Delta R_{e,B} = \log(R_{e,B}^{\text{fit}}/R_{e,B}^{\text{in}})$ and $\Delta h = \log(h^{\text{fit}}/h^{\text{in}})$.

parameters ($R_{e,B}$, n_B) occur at small B/T values, while the largest disc errors can be found at large B/T values and small h values.

Obviously, for real galaxies we do not know the ‘input’ parameter values, but examining the errors in the fitted parameters as a function of the fitted output parameters, shows parameter ranges we should not trust. Therefore, here we exclude from the artificial data and later from the real data all fitted $R_{e,B}$ values for galaxies with $B/T < 0.2$ or $R_{e,B} < 1.0$ FWHM seeing. Likewise, we exclude all fitted h values for fits where $B/T > 0.95$ or $h < 1.5$ FWHM seeing. Applying these cuts results in rms errors between input and fitted parameters of the order of 5 per cent in M_{tot} , 10 per cent in B/T , 30 per cent in $R_{e,B}$, 20 per cent in n_B and 15 per cent in h .

It should be noted that some of the errors are strongly correlated, as shown in the last two columns of Fig. 1. Especially the bulge parameters $R_{e,B}$ and n_B are strongly correlated, with some weaker correlations between these bulge parameters and M_{tot} and B/T .

The n_B values show a systematic trend with the input n_B in Fig. 1. For a considerable fraction this is the result of the limits imposed on the fitted n_B values, which were constrained to $1 \leq n_B \leq 6.5$. This means that artificial galaxies with n_B values close to 1 will have fitted n_B values preferentially scattered to larger n_B fit values, and likewise for galaxies with input n_B values close to 6.5. A similar effect can be seen in the $\Delta B/T$ versus B/T plot. To minimize this effect on real galaxies we increased the limits to $1 \leq n_B \leq 8$.

At small n_B values a relatively large number of galaxies show ‘catastrophic’ errors. These are cases of mistaken identity; the bulge and disc shapes are so similar that they are swapped. Without prior assumptions concerning relative scale sizes or surface brightnesses this is unavoidable.

2.2 Comparing profile fitting and 2D fitting

In Paper III we derived bulge and disc parameters of the EFAR galaxies using one-dimensional (1D) luminosity profiles. Circularly averaged luminosity profiles were fitted by the combination of an $R^{1/4}$ and an exponential luminosity light distribution, convolved with seeing, which was also fitted. Multiple exposures of the same object were optimally combined and an optional sky-fitting procedure was used when necessary to correct for sky subtraction errors.

We will now compare the Paper III 1D parameters with the current GIM2D parameters to study the differences between 1D and 2D fitting and $R^{1/4}$ - and $R^{1/n}$ -law bulges. We removed the fit quality $Q = 3$ fits from the Paper III results, as those were considered to be of poor quality.

Fig. 2 shows the difference in integrated magnitude versus the effective radius difference [defined as $\Delta R_e = 2(R_e^{1D} - R_e^{\text{GIM2D}})/(R_e^{1D} + R_e^{\text{GIM2D}})$] when comparing the 1D and the GIM2D results. The GIM2D magnitudes are 0.06 mag fainter in the median and agree to within 0.17 mag rms. The effective radii agree to within 28 per cent rms, without taking the strong correlation into account. The correlated differences are for a considerable fraction due to the change in bulge parameter shape allowed by the $R^{1/n}$ law. Even though the flux in the inner region of the galaxies is well determined, the total flux (and hence half-light effective radius) is determined by extrapolating the model to infinity. This is demonstrated more clearly in the next figure.

Fig. 3 shows the old versus new bulge effective radius difference as a function of Sérsic bulge shape parameter n_B . When the 2D fitted n_B is getting smaller than the $R^{1/4}$ $n_B = 4$ value, light is extending less far and $R_{e,B}$ is getting smaller in the $R^{1/n}$ -law case than in the $R^{1/4}$ -law case. The opposite happens when $n_B > 4$. Taking this effect into account, the bulge parameters for the galaxies agree to

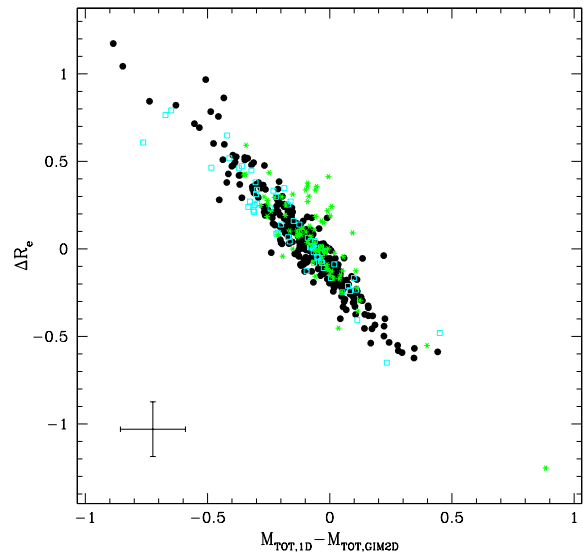


Figure 2. The integrated magnitude versus effective radius difference comparing Paper III with the current GIM2D determinations. The effective radius difference is defined as $\Delta R_e = 2(R_e^{1D} - R_e^{\text{GIM2D}})/(R_e^{1D} + R_e^{\text{GIM2D}})$. Open squares denote the brightest galaxies in EFAR clusters, star symbols denote galaxies with a clear spiral structure in their model-subtracted residual images. The errorbars indicate the median error for all galaxies in the parameters, where the parameter error for an individual galaxy is the maximum error due to 2D fit parameter space, sky subtraction error and seeing error.

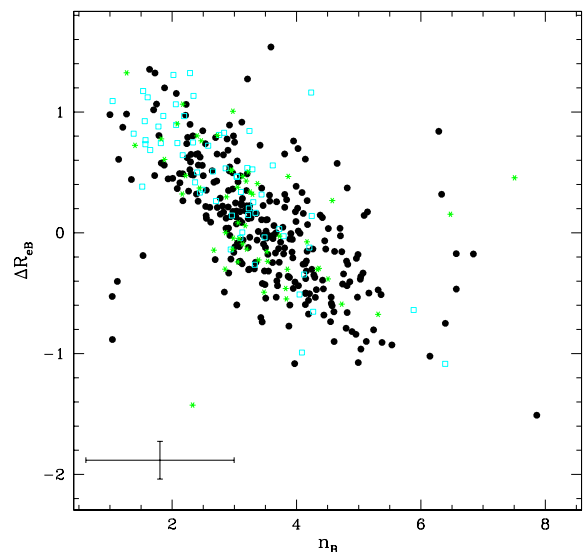


Figure 3. The bulge effective radius difference between 1D Paper III determinations and the current GIM2D determinations as a function of bulge shape n_B . The effective radius difference is defined as $\Delta R_{e,B} = 2(R_{e,B}^{1D} - R_{e,B}^{\text{GIM2D}})/(R_{e,B}^{1D} + R_{e,B}^{\text{GIM2D}})$. Symbols and errorbars are as in Fig. 2.

within 50 per cent rms of each other, i.e. a larger error than for the total galaxy parameters as might be expected.

Fig. 4 compares the old and new disc scalelengths. Some of the extreme outliers are due to the swapping interpretation of which component is a disc and which is a bulge. What used to be a small disc in a large $R^{1/4}$ -law bulge has become a small $R^{1/n}$ -law bulge with, in general, low n_B and sometimes vice versa. The disc scalelengths

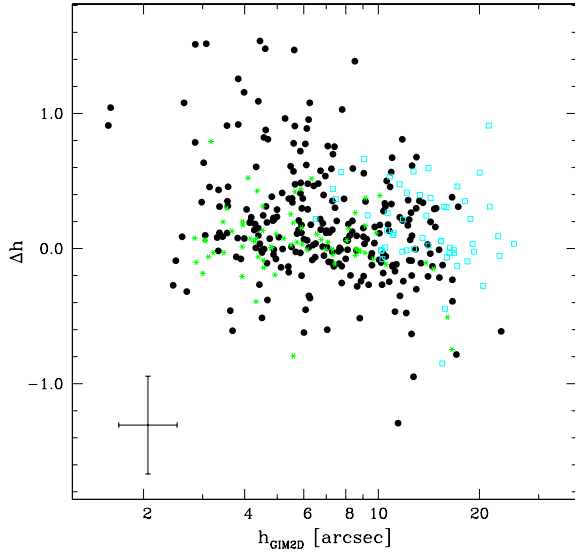


Figure 4. The disc scalelength difference between 1D Paper III determinations and the current GIM2D determinations as a function of bulge shape n_B . The effective radius difference is defined as $\Delta h = 2(h_i^{1D} - h^{GIM2D}) / (h_i^{1D} + h^{GIM2D})$, where $h_i^{1D} = h^{1D} / \sqrt{\cos(i^{GIM2D})}$. Symbols and errorbars are as in Fig. 2.

agree with each other to almost 35 per cent rms, which suggests that many early-type galaxies seem to have a light component that is robustly fitted by an exponential light profile, to some extent independent of the shape of the spheroid component. This does not prove that these are real discs: other radial disc light profiles may be more appropriate (see, e.g., Scorza & Bender 1995) and kinematics is required to confirm a disc configuration for an individual galaxy.

In this section we showed that while the 1D $R^{1/4}$ and the 2D $R^{1/n}$ fits result in systematically different structural parameters, these deviations are correlated with the Sérsic shape parameter. In determining uncertainties in the parameters we have to take these systematic chances into account. Therefore, under the assumption that the combination of an $R^{1/n}$ -law spheroid and exponential disc is a reasonably good description of the intrinsic light distribution of a galaxy and by giving some more weight to the artificial images test of Section 2.1, we conclude that bulge and disc parameters can be determined with a typical 1σ accuracy of approximately 30 per cent. However, one should bear in mind that some extreme outliers with much larger errors will be common in large samples. In the remainder of the paper we will use the GIM2D-determined parameters to investigate correlations between bulge and disc parameters and the space density of early-type galaxies as a function of these parameters. The fitted GIM2D parameters are available in electronic format from the Centre de Données astronomiques de Strasbourg (CDS) web site.

3 PROPERTIES OF DISCS AND BULGES

3.1 Sample selection and selection effects

The selection of the EFAR galaxy sample has been described in detail in Wegner et al. (1996), and we will repeat here the most important aspects relevant for the current study. All galaxies were selected with the intention of being cluster members, with the clusters lying in two selected regions in the sky which are rich in galaxy clusters. The target clusters had redshifts in the range from 6000 to

15 000 km s^{-1} . Photographic enlargements of POSS and SERC Sky Survey plates of the target clusters were inspected for galaxies with elliptical/S0 galaxy-like morphology and a minimum diameter of approximately 16 arcsec. A morphological selection was performed to err on the safe side, including many spiral galaxies which were identified as such later with subsequent CCD imaging. The sample is therefore intended to be complete in large, early-type galaxies (elliptical/S0), but is not complete in any other galaxy type. In most of the analysis presented here, we will only use the galaxies from the EFAR target clusters, excluding the galaxies from the comparison sample in Coma, Virgo and the field, as they suffer from completely different selection effects.

To maximize the number of early-type galaxies, galaxies were selected with rather round isophotes and the sample is therefore rather face-on. The sample has only eight galaxies with ellipticity $\epsilon > 0.4$ at the outer isophotes. The distribution of intrinsic shapes of early-type galaxies is poorly known, and we have therefore not attempted to correct the bulge parameters to face-on values using the bulge isophote shapes. However, the disc parameters shown in the diagrams are face-on values, assuming fully transparent discs.

While we cannot prove that our fitted disc light distributions are really discs, we can use the minor over major axis ratio (b/a) distribution to make a statistical test. For a sample of infinitely thin discs at random viewing angles the distribution of b/a is expected to be flat. We show this distribution for all non-spiral galaxies in Fig. 5 (line). The distribution is rather flat for $b/a > 0.6$, but steeply declines for smaller values of b/a . This is the result of our selection bias, with very few galaxies having $\epsilon > 0.4$ (i.e. $b/a < 0.6$). Assuming that discs are in general aligned with the spheroid they are sitting in, we expect the b/a distribution to be peaked toward face-on. This effect will be less for small discs in large bulges and indeed the distribution becomes relatively more spread out when looking at all non-spiral galaxies with $B/T > 0.7$ in Fig. 5 (solid histogram). While not conclusive, this test shows that our fitted disc components have orientation angles consistent with the random distribution expected.

To fully quantify the selection limits of the sample as a function of the structural parameters of the galaxies, we need not only to

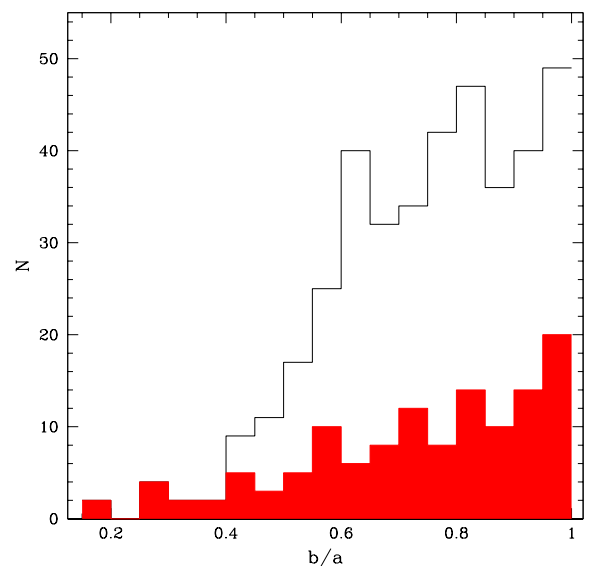


Figure 5. The distribution of minor over major axis ratios (b/a) for all non-spiral galaxies (line histogram) and non-spiral galaxies with $B/T > 0.7$ (solid histogram).

know the angular diameter limits, but also the surface brightness at the selection diameter. The average surface brightness at the selection diameter was $22.05 R\text{-mag arcsec}^{-2}$, with an rms of $0.5 R\text{-mag arcsec}^{-2}$, as derived from the 1D profiles of Paper III. This value is used, in combination with the typical $\log(D_w) > 1.2$ diameter limit (Paper I, see also equation 3) and the minimum 6000 km s^{-1} distance limit, to calculate the selection limits indicated in our diagrams where possible.

3.2 Galactic extinction and cosmological corrections

The observed parameters were corrected for Galactic extinction according to the precepts of Burstein (2003), which averages the estimates of Schlegel, Finkbeiner & Davis (1998) and Burstein & Heiles (1982) with appropriate offsets. No attempt was made to correct the parameters for internal extinction. We used the cluster redshifts as listed in Paper VII to calculate distances to individual galaxies. Cosmological corrections were made using $\Omega_m = 0.3$, $\Omega_\Lambda = 0.7$ and $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, and K-corrections were made as described in Paper III. All of these corrections are typically small for these galaxies and making different reasonable assumptions will not affect our results at all.

3.3 Structural parameter correlations

We will now investigate correlations in structural parameters for the bulge and disc. Obviously, many combinations of bulge and disc parameters can be made, and here we will show only those deemed most interesting.

In Fig. 6 we show the bulge effective radius versus the bulge effective surface brightness for our sample. There seems to be a reasonably tight correlation between these two parameters for the non-spiral galaxies, but we have to be wary of selection effects. The solid line in this diagram shows the selection cut-off for galaxies with an $R^{1/4}$ profile dominating at the selection radius, calculated using the average surface brightness at the selection radius. The

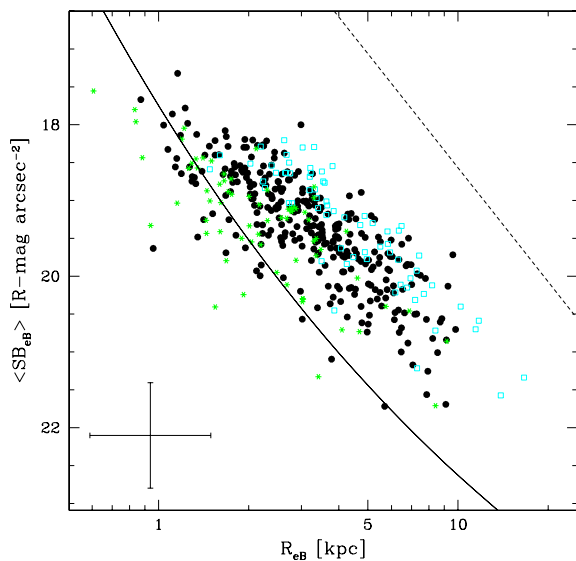


Figure 6. The bulge effective radius versus the bulge effective surface brightness. The solid line shows the selection limit for pure $R^{1/4}$ profile galaxies at 6000 km s^{-1} . The dashed line indicates $R^{1/4}$ profile galaxies of -25 absolute R -mag. Bulges with equal luminosities will lie on lines parallel to this line. Symbols and errorbars are as in Fig. 2.

rather sharp cut-off in the high surface brightness, long scale size region of the diagram is therefore real; there are no selection effects acting in this part of the diagram. This upper limit is related to the zone of exclusion in κ -space discussed in detail by Burstein et al. (1997).

At the low surface brightness end of the diagram we have to be more careful: most of the galaxy distribution seems to lie clearly separated to the right of the selection limit line, but the selection line is for a 6000 km s^{-1} cluster, and moving to the right for the higher-redshift clusters. The ‘visibility volume’ (Disney & Phillipps 1983) for these lower surface brightness galaxies is definitely much smaller than for high surface brightness galaxies, certainly taking further into account the measurement uncertainty in the selection diameter (Paper I). We will make full ‘visibility volume’ corrections in Section 3.4, to calculate the volume density of early-type galaxies as a function of structural parameters.

There are quite a number of galaxies, the bulges of which do not follow the main trend seen for the most non-spiral galaxies in Fig. 6. These bulges often lie to the left of the selection line, and these galaxies are included in the sample because their discs make a significant contribution at the selection diameter or because their profile does not follow an $R^{1/4}$ law. Many of these bulges, especially in some of the spiral galaxies, may have the more exponential-like bulges typical of late-type spiral galaxies (Andredakis et al. 1995; Courteau, de Jong & Broeils 1996; Graham 2001) instead of the more $R^{1/4}$ -like bulges found in early-type galaxies.

In Fig. 7 we show the disc scalelength versus disc central surface brightness for all fitted galaxies. The dotted line shows the $cz = 6000 \text{ km s}^{-1}$ selection limit for galaxies which have an exponential light distribution dominating at the selection radius (i.e. spiral galaxies). Elliptical galaxies will not adhere to this selection limit, as they were selected on the basis of their $R^{1/n}$ light distribution. Some interesting trends can be seen in this diagram, but we have to be wary of more hidden selection effects. At the very short scalelengths we are close to our scalelength cut-off, as 1 kpc is of the order of 2.5 arcsec for $cz = 6000 \text{ km s}^{-1}$ and even less for higher redshifts. The observed limit at the low surface brightness, small scalelength

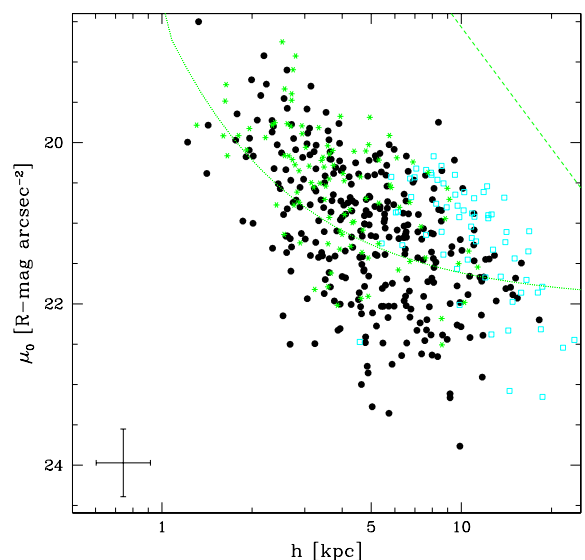


Figure 7. The disc scalelength versus disc central surface brightness. The dotted line shows the selection limit for exponential discs at $cz = 6000 \text{ km s}^{-1}$. The dashed line indicates exponential discs of -25 absolute R -mag. Symbols and errorbars are as in Fig. 2.

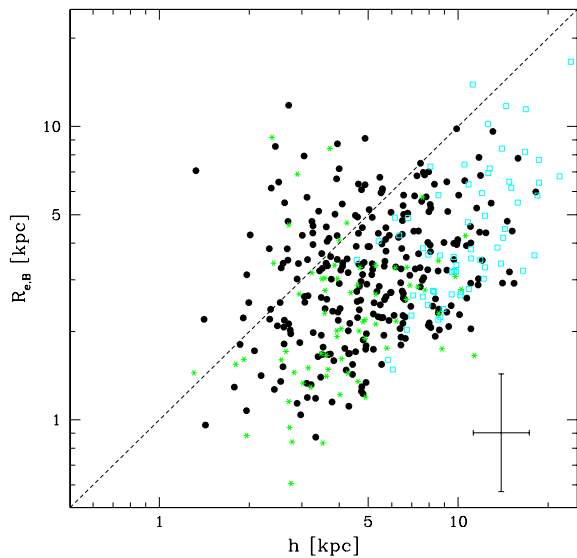


Figure 8. The disc scalelength versus bulge effective radius. To guide the eye, the line of equality is shown as a dashed line. Symbols and errorbars are as in Fig. 2.

end is more subtle and probably arises from the minimum contrast needed to see a low surface brightness disc on top of a bright bulge distribution. Using the relation between h and $R_{e,B}$ shown in Fig. 8 discussed in the next paragraph and the relation between $R_{e,B}$ and $\langle SB_{e,B} \rangle$ shown in Fig. 6, we find that we probe discs with central surface brightnesses at most approximately 2 mag fainter than the $\langle SB_{e,B} \rangle$ of the bulge. It is not unreasonable to suspect that the low surface brightness cut-off is caused by a lack of disc/bulge contrast. The only cut-off observed in this diagram that is definitely real is at large scalelength, high surface brightness, as no selection effects operate in this part of the diagram. This cut-off is similar to that observed for spiral galaxies (de Jong 1996b).

In Fig. 8 we show the correlation between bulge and disc scale sizes. No obvious selection effects operate in this diagram, except for a spiral galaxy having a large disc scalelength and a small bulge would probably not have entered the sample. This is not the case for elliptical galaxies and therefore we have to make sure that other effects are not causing this correlation, before we can accept it as real. Our Monte Carlo simulations showed that our GIM2D fitting routine is capable of recovering models of the full parameter space presented here, so it is unlikely that it is the result of a fitting routine artefact. Indeed, we do see a small number of galaxies with $R_{e,B}/h > 1$ in Fig. 8, and so the routine is capable of covering this parameter space in real data. Still, given the limits on the parameter space for a given galaxy (not larger than the image size, not smaller than the seeing), it is not completely surprising that the bulge and disc scale parameters are similar in size.

The $R_{e,B}/h > 1$ galaxies can be divided into three groups. A number of these galaxies are clearly disturbed by dust-lanes or are poorly fitted due to nearby bright galaxies or stars. We may have a number of cases of mistaken identity, where the bulge and the disc are swapped. This is especially the case when n_B is low and the bulge has an almost exponential light profile. Swapping the bulge and disc parameters on these galaxies would bring them in line with the general trend. Finally, there are a very few galaxies which seem to have a genuine small inner light component for their given bulge size, but these should not be mistaken for the nuclear discs found in *Hubble Space Telescope* (*HST*) images of early-type

galaxies, as these are a factor of 5–20 smaller again (Scorza & van den Bosch 1998). In a number of cases they might be identified with the inner break radius observed in *HST* images of early-type galaxies (Faber et al. 1997; Trujillo et al. 2004), as shown by D’Onofrio (2001).

In Fig. 9 we show in three sets of diagrams the bulge, disc, and combined bulge and disc parameters as a function of B/T ratio and Sérsic n_B values. Before we can draw any conclusions on any apparent correlation, we have to be aware of the many selection effects playing in these diagrams. Some are explicit, such as the limits we placed on some parameters based on the B/T ratio to reduce fitting errors. Others are more hidden, such as the exclusion of bright discs, as their more visible spiral structure would have precluded them from entering the sample in the first place.

Concentrating first on the B/T ratios, we see that more bulge-dominated systems have larger bulge scale sizes, but lower effective surface intensities, which combines in a rather weak B/T trend with total bulge luminosity. On the other hand, looking at the disc parameters in the middle set of diagrams, we see a weak trend of disc scale sizes with B/T ratio, but at each B/T ratio there is a wide range in disc scale sizes. More bulge-dominated systems have discs of lower surface intensities. The combined trends of disc scale size and surface brightness result in a strong correlation between disc luminosity and B/T ratio for our sample dominated by bright, early-type galaxies.

Combining these results for this sample, the change in B/T ratio seems to be mainly driven by changes in disc luminosity, not bulge luminosity. This is opposite to what has been observed for disc-dominated systems as found by, for example, Trujillo et al. (2002) and Balcells et al. (2004), who showed that most of the change in B/T ratio in these systems is due to bulge luminosity variation. We have to be cautious though. Our selection criteria resulted in a very limited range in total galaxy luminosities. We are only sampling the upper decade in the luminosity function, the 5–10 brightest galaxies in a cluster. It could be that lower-luminosity elliptical galaxies reveal often equal-sized discs, changing the observed trends.

Finally, we see that the weak opposite trends in the bulge and disc scale sizes with B/T ratio combine in a strong correlation between $h/R_{e,B}$ and B/T . We should be cautious not to see too much in this relationship. The expected relation between $h/R_{e,B}$ and B/T under the assumption that there is a fixed bulge-to-disc surface brightness difference and no significant bulge profile shape change is indicated by the dashed line. The bulk of the galaxies indeed have a limited range in disc-to-bulge surface brightness difference as shown in Fig. 9. It is just hard to see a faint component on top of a bright component, unless their light distributions are completely different. The variation from $n_B = 1$ to 5 amounts to only a factor of 2 in bulge luminosity change at a fixed bulge-to-disc surface brightness ratio (see Graham 2001, his fig. 4) and this change is systematic with $h/R_{e,B}$ as shown in Fig. 9. All in all, the relation between B/T and $h/R_{e,B}$ with the observed scatter is expected, given the observed scatter in disc-to-bulge surface brightness contrast and the limited and systematic effect in n_B .

We note that D’Onofrio (2001) finds a different slope in the B/T versus $h/R_{e,B}$ relation. Many of his discs were fitted to the small inner components (either discs or cores) of the galaxies in his nearby sample. We are not able to resolve these components and, as D’Onofrio rightly indicates, it is unlikely that these inner components represent similar structures such as the more extended components fitted in discy galaxies. We do fit the more extended components and indeed find that they have a different relation to the main galaxy.

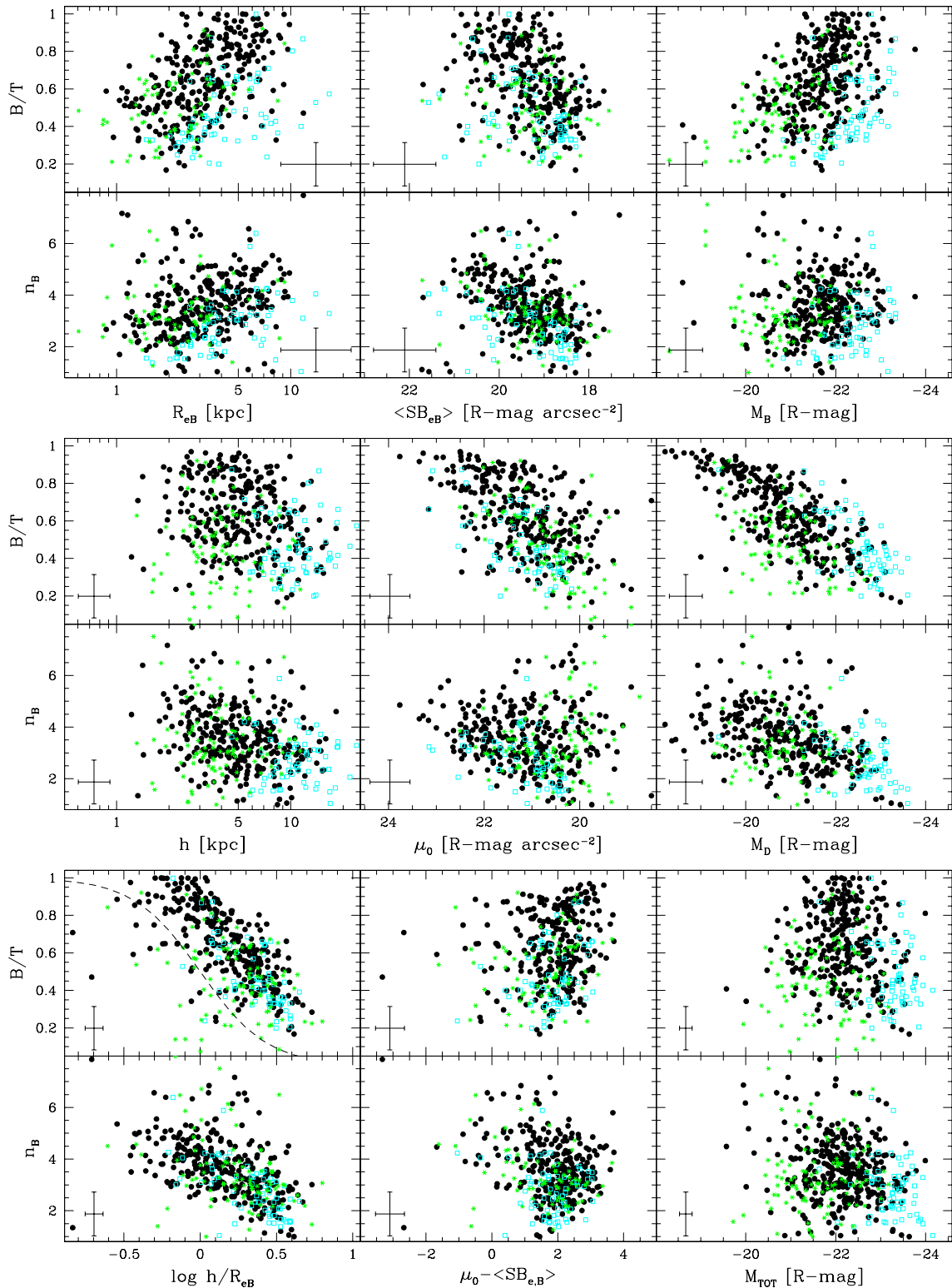


Figure 9. Bulge-to-total-light ratio and Sérsic profile shape parameter n_B as a function of the bulge parameters (top), disc parameters (middle) and bulge–disc combined parameters (bottom). Symbols and errorbars are as in Fig. 2. The dashed line in the $h/R_{e,B}$ versus B/T diagram indicates the expected relation between these parameters under the assumption that there is a fixed bulge-to-disc surface brightness difference with arbitrary zero-point.

If we now turn our attention to the relationships between n_B and the bulge and disc parameters, we note that n_B has a behaviour very similar to the B/T ratio with respect to these bulge and disc parameters. There are some more outliers in the n_B relations, most of which are for n_B values larger than 5.5.

The very similar behaviour of n_B and the B/T ratio in Fig. 9 suggests that these two parameters must be correlated, which is indeed the case as shown in Fig. 10. As mentioned before, all other things being equal a change in n_B from 1 to 5 only results in a factor of 2 change in the bulge-to-disc luminosity, so the change

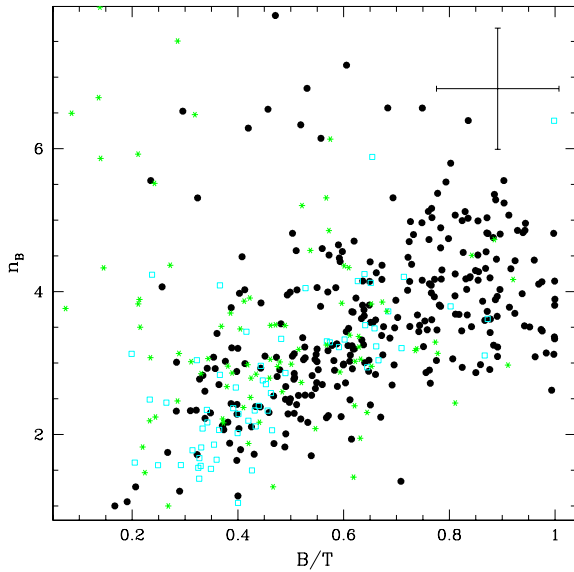


Figure 10. The Sérsic n_B bulge shape parameter versus B/T ratio. Symbols and errorbars are as in Fig. 2.

in n_B does not drive the variation in the B/T ratio on its own. For this set of early-type galaxies, a large fraction of the B/T change seems to come from the change in the bulge effective radius, and it is the correlation between n_B and $R_{e,B}$ that seems to drive the trend between n_B and the B/T ratio. Still, cause and effect are hard to disentangle based on this photometric data set alone. We note that these trends cannot solely be the result of the correlated errors noted in Fig. 1, as the slope of the errors is incorrect and the size of the effect is too small.

In Fig. 11 we show the observed distributions of the n_B parameter for different subsets of our sample. Galaxies that were fitted with n_B values larger than 6 were outliers in previous diagrams and are considered to be of dubious quality. They are excluded from further analysis here. The distributions are quite narrowly peaked, with 1σ dispersions ranging from approximately 1.05 for the total

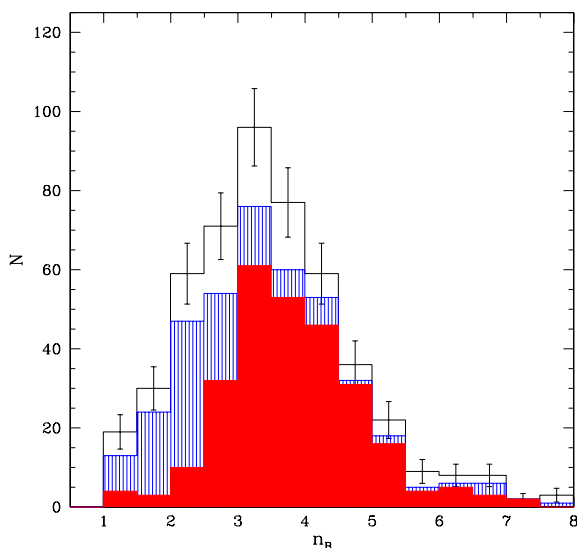


Figure 11. The distribution of Sérsic profile shape parameter n_B for our sample. The top solid line shows the total sample, the hashed histogram shows all galaxies without clear spiral structures and the solid histogram shows all non-spiral galaxies with $B/T > 0.5$.

sample to 0.88 for bulge-dominated non-spiral galaxies. This width is approximately twice the mean of the individual maximal error estimates of the n_B values, and is therefore significant. The median n_B value shifts from 3.24 for the total sample, to 3.31 for the non-spiral sample, and to 3.66 for the non-spiral, bulge-dominated $B/T > 0.5$ sample. For galaxies with $B/T > 0.7$ the median is 3.99. The real n_B values might be even somewhat lower. Balcells et al. (2003) showed that for a sample of early-type disc galaxies the n_B values decreased compared with ground-based data when a nuclear component revealed in their *HST* data was included as a separate component in their fits.

This result complicates bulge–disc decompositions of higher-redshift galaxies with typically a lower signal-to-noise ratio and resolution. With such lower-quality data a Sérsic-law bulge with an exponential disc does not reliably converge and one is forced to use a fixed n_B value. Given that in this morphologically preselected and bright galaxy sample there is a broad distribution of n_B independent of luminosity (see Fig. 9) peaking at values below $n_B = 4$, and given that spiral galaxies tend to have bulges with n_B values in the 0.5–3 range (e.g. Balcells et al. 2003; MacArthur et al. 2003), it may be more appropriate to use an n_B in the 1.5–3 range rather than the more customary $n_B = 4$ in fixed n_B fits. A similar conclusion was reached for early-type disc galaxies by Balcells et al. (2003), but is now extended to bright, early-type cluster galaxies here.

3.4 Space density distribution functions

Probably more important than correlations between structural parameters is the true space density of galaxies as a function of structural parameters, i.e. the distribution of structural parameters corrected for selection effects.

When calculating space densities of galaxies from the EFAR sample there are several points to note. The sample is not complete for spiral galaxies, so we cannot say anything concerning the space densities of those galaxies. The regions selected by EFAR were particularly rich in clusters, and therefore are not representative for the average space density of galaxies in the local Universe. Assuming that galaxies have similar properties in different clusters, we can still use the relative space density of structural parameters, but for absolute space densities we will have to renormalize the EFAR regions to the universal cluster density.

Calculating the space density of galaxies requires a good understanding of the selection function. The selection of the EFAR sample is described in detail in Paper I. The selection of EFAR galaxies was a two-step process, where first clusters were selected, believed to be in the right redshift range and to be rich enough to obtain accurate distances. In the next step galaxies were selected in each cluster according to a minimum angular diameter criterion.

Here we will use a simple V_{\max} correction method (Schmidt 1968) to calculate the relative space density of E/S0 galaxies as a function of their structural parameters. The V_{\max} correction method assumes that the average space density of an object with certain parameters is inversely proportional to the maximum ‘visibility’ volume V_{\max} of that object, given the selection criteria of the sample (for more detailed descriptions see Felten 1976; Disney & Phillipps 1983; de Jong & Lacey 2000). The selection of galaxies was a two-step process and therefore calculating the maximum visibility volume of a galaxy must reproduce this two-step process.

First, the V_{\max} of each individual galaxy was calculated, by integrating the total volume occupied between the minimum and maximum selection redshifts of $z_{\min} = 6000/c$ and $z_{\max} = 15\,000/c$, where c is the speed of light. While calculating the volume integral,

we have to take into account the diameter selection function determined in Paper I for each cluster as applied to each galaxy individually. The selection function describes the selection diameter (D_W) dependent probability of sample inclusion, parametrized by

$$P[\log(D_W)] = 0.5 \left\{ 1 + \operatorname{erf} \left[\log(D_W/D_W^0) / \delta_W \right] \right\}, \quad (3)$$

where D_W^0 is the mid-point and δ_W the width of the cut-off in the selection function. Furthermore, we took into account that the contribution to the volume has to drop to zero when R_e becomes less than 1.5 arcsec had the galaxy been at that redshift, because we excluded such galaxies from our analysis. The integral to calculate the total visibility volume is then

$$V_{\max} = 4\pi\Omega_f \int_{z_{\min}}^{z_{\max}} P \left\{ \log \left[D_W \frac{D_M(z_{\text{Cl}})}{D_M(z)} \right] \right\} \frac{cD_M^2(z)S(z)}{(1+z)H(z)} dz \quad (4)$$

with

$$S(z) = \begin{cases} 1 & \text{if } R_e \frac{D_M(z_{\text{Cl}})}{D_M(z)} \geq 1.5 \text{ arcsec,} \\ 0 & \text{otherwise,} \end{cases} \quad (5)$$

where z_{Cl} is the cluster redshift of the galaxy, $\Omega_f = 0.139$ is the fraction of the sky covered by our survey, $D_M(z)$ is the comoving distance and $H(z)$ is the Hubble constant at redshift z as defined in Hogg (1999).

In the second step, we have to take into account that objects were only selected if they belonged to a cluster, with at least three early-type galaxies obeying the diameter selection criterion. Therefore, the V_{\max} of a cluster is identical to the V_{\max} of the galaxy ranked third in diameter belonging to the cluster. The final V_{\max} of each galaxy is the minimum of its individual and cluster V_{\max} . This last step only modifies the V_{\max} of the top three diameter galaxies in each cluster, and only if the diameter limit of the cluster is reached before the third ranked galaxy is redshifted to $15\,000 \text{ km s}^{-1}$.

As mentioned before, the EFAR regions were chosen to be particularly rich in clusters at the target redshift range. We used the NASA/IPAC Extragalactic Data base (NED) to estimate the size of the overdensity. NED lists 140 original, non-supplement Abell clusters (Abell 1958; Abell, Corwin & Olowin 1989) for $|b| > 30^\circ$ ($\Omega_f = 0.5$) with $6000 < v < 15\,000 \text{ km s}^{-1}$ versus 87 in the EFAR regions ($\Omega_f = 0.139$). This translates to an overdensity of approximately 1.75. However, we observed only 32 of the 87 clusters identified by NED, and while not every NED cluster should be in our sample (some cluster redshifts are based on only one galaxy redshift), it is clear that the overdensity is probably offset by incompleteness. Our sample contains another 41 clusters selected by other means, which are not included in the Abell catalogue. We have no means to access the completeness of these, in general, smaller clusters. Given these uncertainties, we have decided not to correct for overdensity and incompleteness, and assume that the uncertainty in the absolute space density zero-point is a factor of 2. We do, however, correct for our ellipticity incompleteness, assuming that elliptical galaxies with $\epsilon < 0.4$ constitute 65 per cent of the total E+S0 population (Jørgensen & Franx 1994).

We will show the factor of 2 uncertainty in absolute zero-point of the space density distributions in the diagrams presented in the remainder of this section. This uncertainty will not affect the conclusions drawn from these diagrams. We will also show that our space density distribution and normalization is remarkably consistent with that of the SDSS, giving confidence that our normalization is indeed correct to within a factor of 2.

In Fig. 12 we show the combined bivariate space density of all non-spiral galaxies as a function of absolute magnitude and effective

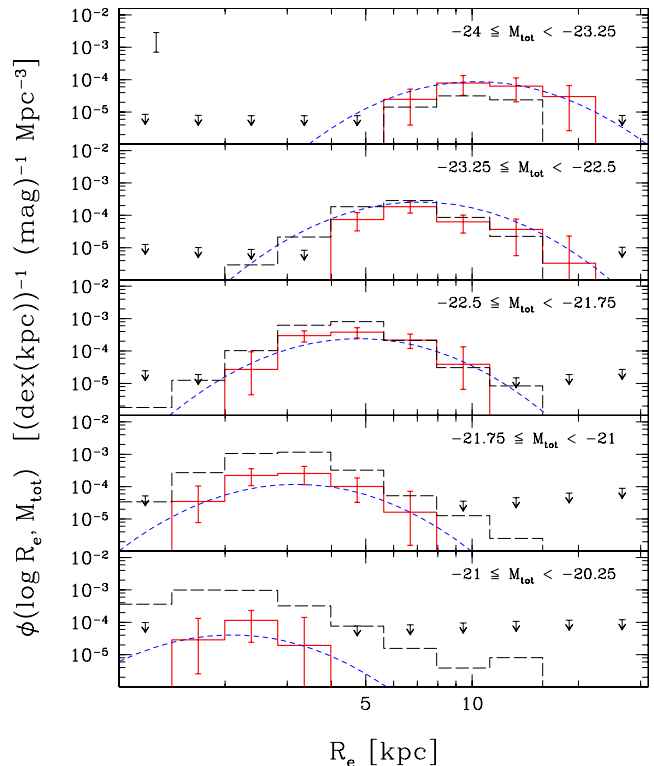


Figure 12. The space density of all non-spiral galaxies as a function of total magnitude and the effective radius of the whole galaxy is shown as the solid line histogram. Errorbars on the histogram show the 95 per cent confidence limits due to uncertainties in the selection diameter and distance, and due to Poisson statistics. The upper limit symbols indicate the 95 per cent confidence limits based on the non-detection of such galaxies and the visibility volume of $R^{1/4}$ profiles shape galaxies for our selection criteria. The errorbar in the top left-hand corner of the top panel shows the factor of 2 uncertainty in the absolute space density normalization of our distributions. The dashed curve shows the bivariate distribution function of the form presented in de Jong & Lacey (2000) fitted to the data. The dashed histogram is the space density distribution derived by Blanton et al. (2003) from SDSS data for galaxies with $n > 3$ and scaled as described in the text.

tive radius. The 95 per cent confidence limit error estimates on the number density distribution and the 95 per cent confidence upper limits in the regions with no galaxies in the sample were derived in a similar fashion to that described in de Jong & Lacey (2000), assuming a perfect $R^{1/4}$ -law luminosity profile for all galaxies. These error estimates and upper limits include the effects of distance and diameter uncertainties and Poisson statistics.

Fig. 12 also shows the fitted bivariate distribution function of the form described by de Jong & Lacey (2000). This function takes the shape of a Schechter (1976) luminosity function in the luminosity dimension (with the usual parameters ϕ_* , M_* and α), and a log-normal distribution in the scale size dimension, with a width σ_λ , a median scale size for galaxy with $M = M_*$ of R_{e*} , and a shift in median scale size with luminosity parametrized by β :

$$\begin{aligned} \phi[M, \log(R_e)] dM d \log R_e &= 0.4 \ln(10) \frac{\ln(10)}{\sqrt{2\pi}\sigma_\lambda} \\ &\times \phi_* 10^{-0.4(\alpha+1)(M-M_*)} \exp[-10^{-0.4(M-M_*)}] dM \\ &\times \exp \left\{ -\frac{1}{2} \left[\frac{\log(R_e/R_{e*}) - 0.4\beta(M-M_*)}{\sigma_\lambda / \ln(10)} \right]^2 \right\} d \log R_e. \quad (6) \end{aligned}$$

Table 1. Bivariate distribution function parameters. Errors indicate 95 per cent confidence limits.

Fit	ϕ_* ($\times 10^{-4} \text{ Mpc}^{-3}$)	α	M_* (R -mag)	R_{e*} (kpc)	σ_λ	β
Total galaxy	0.99 ± 0.19	1.06 ± 0.26	-21.81 ± 0.18	3.9 ± 0.5	0.34 ± 0.04	-0.62 ± 0.08
Bulge	0.83 ± 0.31	1.30 ± 0.50	-21.15 ± 0.30	2.3 ± 0.4	0.39 ± 0.04	-0.55 ± 0.10
Disc	0.91 ± 0.26	-0.17 ± 0.24	-21.98 ± 0.35	12.6 ± 2.2	0.42 ± 0.04	-0.40 ± 0.07

The second line in this equation represents the Schechter luminosity function in magnitudes, while the third line represents the lognormal scale size distribution.

The function was fitted using a maximum-likelihood technique and the resulting parameters are listed in Table 1, with the errors indicating the 95 per cent confidence limits. These confidence limits were determined by means of Monte Carlo bootstrap resampling of the galaxies (Press et al. 1993), where we also varied the size of the bins in both directions. The fitted function is a good description of the observed distribution: the goodness-of-fit parameter Q (Press et al. 1993) was larger than 0.1 in more than 28 per cent of the bootstrap resampled realizations and $Q > 0.001$ in more than 87 per cent of the realizations. A function similar in shape was fitted by Chołoniewski (1985) to a sample of 233 elliptical galaxies and our parameters do agree to within the quoted uncertainties.

In Fig. 12 we compare our bivariate distribution with the findings of Blanton et al. (2003), based on the SDSS data (see their fig. 13). They fitted Sérsic profiles to the total galaxy luminosity profiles, without an added exponential disc component and hence without performing a full bulge/disc decomposition. We show their $^{0.1}i$ distribution for all galaxies with a Sérsic $n > 3$. We scaled their distribution to our Hubble constant and using a $^{0.1}i$ to R conversion of 0.2 mag [from Blanton et al. 2003 $^{0.1}(r - i) \sim 0.4$, $^{0.1}r = ^{0.0}r - 2.5 \log(1.1)$, and $R \sim r - 0.3$ from Bell & Vandenberg 1987].

Surprisingly, the distributions agree well at the bright end, even though our absolute space density zero-point calibration is somewhat uncertain (see the errorbar in the top panel) due to the cluster overdensity in the EFAR region and our limited knowledge of cluster incompleteness. The cut-off at $R_e > 20$ kpc in the Blanton et al. (2003) distribution is not real, but due to the cut-off in their figure. At fainter luminosities our distribution is lower than that of Blanton et al. (2003). One explanation for this might be that we were stricter in excluding spiral galaxies from our distribution, while the Blanton et al. distribution includes all galaxies with $n > 3$, which may include spiral galaxies where the combined bulge and disc luminosity profile mimics that of an elliptical galaxy. On the other hand, incompleteness may affect our distribution at the faint end. Our sample was drawn from rich clusters, while the Blanton et al. (2003) SDSS distributions were calculated for all of the sky. This will have little effect on the bright end of the distribution, because the brightest elliptical galaxies are only found in clusters, but will affect the faint end. Notwithstanding, the size distributions at a given luminosity are very similar.

In recent years it has become clear that there are two types of elliptical galaxies (Kormendy & Bender 1996; Faber et al. 1997). On the one hand we have the bright elliptical galaxies with boxy isophotes, with a clear central core, dynamically mainly pressure supported, while on the other hand we have the lower-luminosity elliptical galaxies, often with more discy isophotes, cuspy cores and for a larger fraction supported and flattened by rotation. The boxy galaxies tend to be central dominant galaxies of clusters (cD galaxies) or at the centre of one of the subclumps of a cluster (resulting from cluster merging?). The transition from one type into the other

occurs at approximately -21.2 V -mag (Faber et al. 1997), i.e. approximately -20.5 R -mag, with an overlap of both types occurring for approximately 1.5 mag.

The bright, boxy-type elliptical galaxies are probably the result of the merging of many subclump units in the hierarchical galaxy formation scenario. The origin of the lower-luminosity elliptical galaxies is less clear. They could be the result of the merging of two similar sized gaseous galaxies (Barnes 1992) or the result of ‘failed’ spiral galaxies due to the lack of angular momentum in the protogalaxy (e.g. Dalcanton, Spergel & Summers 1997; de Jong & Lacey 2000). These two options might actually be somewhat related, as merging of two subunits is of course the result of a lack of orbital angular momentum. For S0 galaxies even more formation scenarios have been proposed: ram-pressure stripping of gas (Gunn & Gott 1972; Kent 1981; Quilis, Moore & Bower 2000), galaxy harassment via high-speed impulsive encounters (Moore et al. 1996), cluster tidal forces (Byrd & Valtonen 1990) which distort galaxies as they come close to the centre, interaction/merging of galaxies (Icke 1985; Bekki 1998), and removal and consumption of the gas due to the cluster environment (Larson, Tinsley & Caldwell 1980; Bekki, Couch & Shioya 2002). In all of these scenarios the origin of non-boxy elliptical galaxies and S0 galaxies lies in spiral galaxies or protospiral clumps. Therefore, it seems natural to compare the size–luminosity distribution of our elliptical sample with that of spiral galaxies.

In Fig. 13 we again show the space density distribution of EFAR galaxies as a function of size and luminosity, but this time we compare the distribution with that of Blanton et al. (2003) SDSS galaxies with Sérsic profile fits with $n < 1.5$, i.e. galaxies with exponential-like profiles, predominantly spiral galaxies. The Blanton et al. distributions were scaled for Hubble constant and passband in the same way as in the previous diagram for $n > 3$ galaxies. The disc effective radii were multiplied by 0.75 to take into account that for an $R^{1/4}$ -law galaxy the projected surface brightness effective radius is approximately 0.75 times the spherical effective radius. In this way, we make a rough comparison between the stellar mass distributions, ignoring the effect of other baryonic components (most notably the H I in the disc galaxies, which could be converted to stars during merging) and the systematic change in M/L with luminosity, especially for disc galaxies. When scaled for Hubble constant and passband, the Sc-Sd spiral galaxy distribution of de Jong & Lacey (2000) is very similar to that of Blanton et al. (2003), particularly at the bright end.

Fig. 13 immediately shows that to form the brightest elliptical galaxies one will need to merge several current-day bright spiral galaxies, because there is a substantial space density of bright elliptical galaxies that are at least 0.75 mag brighter than the brightest spiral galaxies. In the process the effective radius of the resulting galaxy is increased by a factor of a few compared with the originating galaxy, presumably because the kinematic energy of the galaxies orbiting each other is converted to heat the random kinematic energy of the stars in the merger product. Barnes (1992) finds for dissipationless equal-mass mergers a typical increase of effective radius of

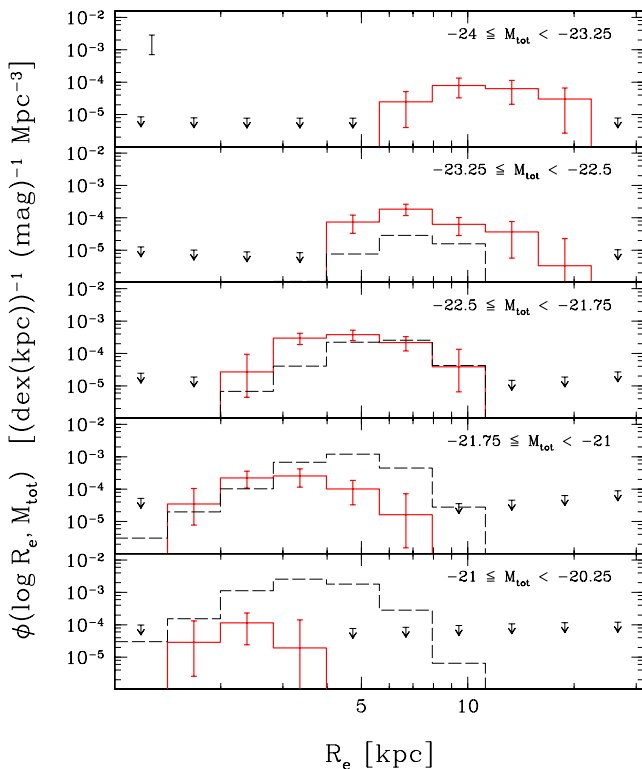


Figure 13. The space density of EFAR galaxies as a function of magnitude and effective radius of the total galaxy (solid histogram). Errorbars and upper limits are the same as in Fig. 12. The dashed histogram is the space density distribution derived by Blanton et al. (2003) from SDSS data for galaxies with $n < 1.5$ and scaled as described in the text.

approximately 1.4, similar to what one expects assuming conservation of energy and the virial theorem (Cole et al. 2000; Shen et al. 2003). We have stepped our luminosity bins in Fig. 13 by 0.75 mag, i.e. approximately a factor of 2 in luminosity (and roughly stellar mass), so each equal-mass dissipationless merger of two disc galaxies would result in a merger remnant one bin up with approximately 1.4 larger size. Merging of at least a few and probably many more of present-day spiral galaxies would be needed to form the brightest elliptical galaxies. Because merging of spiral galaxies will involve dissipation and therefore the scale size may be less, it may be more reasonable to think that the brightest elliptical galaxies formed from the merging of several smaller E/S0 galaxies.

A similar conclusion was reached by Shen et al. (2003). They showed that by having a series of mergers of small elliptical galaxies they could create the observed size–luminosity relation and its dispersion, when making reasonable assumptions concerning the transfer of orbital angular energy. However, they did not compare space densities, and furthermore, this process leaves the question of the origin of small elliptical galaxies unanswered.

The distribution of the lower-luminosity elliptical galaxies ($M_{\text{tot}} > -22.5$ R-mag) in Fig. 13 is more surprising. Even though in number density not completely unlike the spiral galaxies (remember the absolute zero-point for the elliptical distribution is somewhat uncertain), the effective radius distribution of the elliptical galaxies peaks at smaller radii than the distribution of disc galaxies. As argued before, dissipationless merging results in larger scale sizes, typically by a factor of approximately 1.4, and therefore the dissipationless merging of typical current-day spiral galaxies cannot result in these current-day lower-luminosity elliptical galaxies.

What other creation scenarios can we invoke to explain the lower-luminosity elliptical galaxies? First of all, the merging of two spiral galaxies is not going to be dissipationless. Simulations that include gas show that during the merger process gas will quickly stream to the centre of the merging galaxies, presumably creating a starburst and dragging along some of the existing stellar population, creating a more concentrated stellar remnant. Whether this process is enough to offset the effect of the added orbital kinematic energy is questionable, given that current bright spiral galaxies have gas mass fractions of 5–20 per cent. At higher redshifts the gas fractions will have been higher and therefore this process may have been more efficient then. In the hierarchical galaxy formation picture, galaxies at higher redshifts have smaller scale sizes, so the small-scale sizes of low-luminosity elliptical galaxies in clusters may be the result of high-redshift mergers which have subsequently not been able to grow any more in the cluster environment.

Another option to explain the small-scale sizes of lower-luminosity elliptical galaxies might come from the failed disc scenario proposed by several authors for different reasons (e.g. Dalcanton et al. 1997; McGaugh & de Blok 1998; Mo, Mao & White 1998; de Jong & Lacey 2000). de Jong & Lacey (2000) found that the scale size distribution of spiral galaxies used here was narrower than what one would expect in a hierarchical galaxy formation scenario where disc galaxy sizes are determined by their spin acquired from tidal torques with neighbours. They proposed (similar to Mo et al. 1998) that the lowest angular momentum protogalaxies were never able to form a proper disc and immediately collapsed to spheroids. The problem here is that we cannot compare the elliptical galaxy distribution drawn from a cluster sample with the distribution of spiral galaxies drawn mainly from the field.

Can spiral galaxies falling into a cluster result in an elliptical scale size distribution as observed here? Simple stripping of gas by ram-pressure or tidal interactions and subsequent fading of the stellar population is clearly not enough, elliptical galaxies are just too small compared with spiral galaxies. Subsequent stripping of the stars in the outer parts of the infalling disc galaxy by galaxy harassment (Moore et al. 1996) could result in smaller effective radii of the remnants. Still a lot of stripping would have to occur, as stripping half the stellar mass exclusively from the outer part of the galaxy would shift the disc distribution one bin down while dividing the size by two. A considerable fraction of this process would have to occur at quite high redshifts ($z > 1.5$) as only then were bright disc galaxies substantially smaller than current-day disc galaxies (de Jong & Lacey 1999; Ferguson et al. 2004; Ravindranath et al. 2004).

As is so often the case, a combination of several of the above-listed options may be the final answer. However, these relative elliptical and spiral galaxy scale size distributions suggest that a substantial fraction of the lower-luminosity elliptical galaxies were created at quite high redshift, not only in stars but also in structural parameters, when disc galaxies were still smaller and more gas-rich.

So far, we have only looked at the luminosities and sizes for the total systems, but in Figs 14 and 15 we show the luminosity–size distributions for the bulges and the discs, respectively. The luminosity and size parameters are much better determined for the total system than for its components, but none the less it will be instructive to compare the distributions of the components with that of spiral galaxies.

Comparison of the elliptical bulge parameters with the SDSS disc parameters in Fig. 14 shows that the distributions are somewhat similar, but that the bulges are approximately a factor of from ~ 1.5 to 2.5 smaller than the spiral discs from the brightest to the faintest

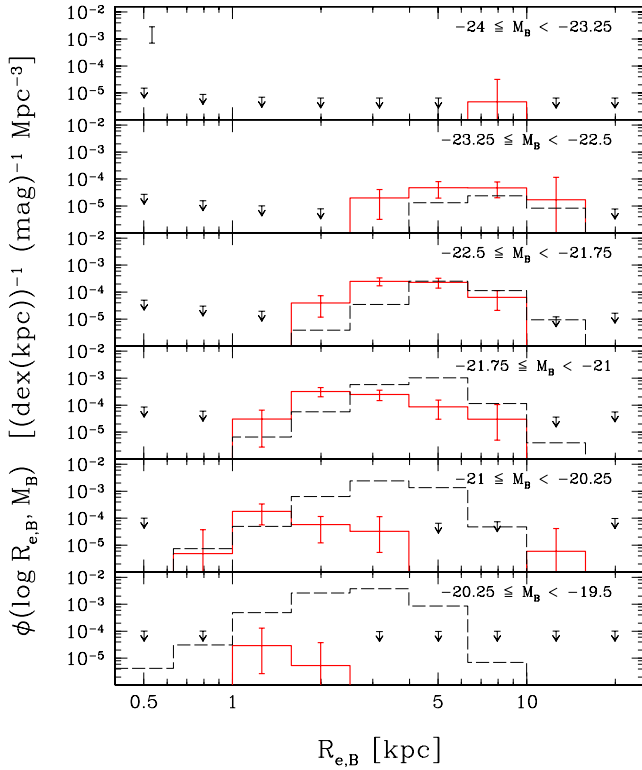


Figure 14. The space density of EFAR elliptical galaxies as a function of bulge magnitude and bulge scalelength (solid histogram). The errorbars and upper limits are the same as in Fig. 12. The dashed histogram is the space density distribution derived by Blanton et al. (2003) from SDSS data for galaxies with $n < 1.5$, scaled as described in the text.

luminosities. Here the scaling of the SDSS disc parameters is the same as in Fig. 13, including the factor of 0.75 decrease in disc scale sizes to account for the difference of a spherical bulge projected to a 2D effective radius versus a 2D disc effective radius.

On the other hand, when we compare (in Fig. 15) the SDSS $n < 1.5$ disc parameter distribution with the distribution for the component identified as discs in our EFAR sample, we again see somewhat similar, parallel distributions, but this time the EFAR disc-like component is approximately a factor of 2 larger than the SDSS discs. We have to be careful, because as argued before in Section 3.3 when discussing Fig. 7, discs with low surface brightness and/or small scalelength might be hard to detect when superposed on a high surface brightness bulge and therefore this distribution may be incomplete. Still, it is clear that there are larger disc components in the EFAR non-spiral galaxies than seen in the SDSS spiral discs. Upper limits are somewhat hard to draw on this diagram, as the selection was dominated by the bulge light distribution, not by the disc light.

Without kinematic information proving that the B/D decompositions we made for our sample have an underlying physical meaning it is somewhat premature to interpret these distributions in detail, but we can make some speculations. Could the spheroidal components be the result of early, dissipationless mergers of roughly equal-sized galaxies from the time when galaxies were still much smaller ($z > 1-2$)? The disc-like component can clearly not be formed from normal disc galaxies, as their scale sizes are just too large for their luminosity. The disc-like component tends to dominate slightly in the most luminous galaxies (see Fig. 9). Are these the stars from smaller, infalling galaxies that maintain a larger component from

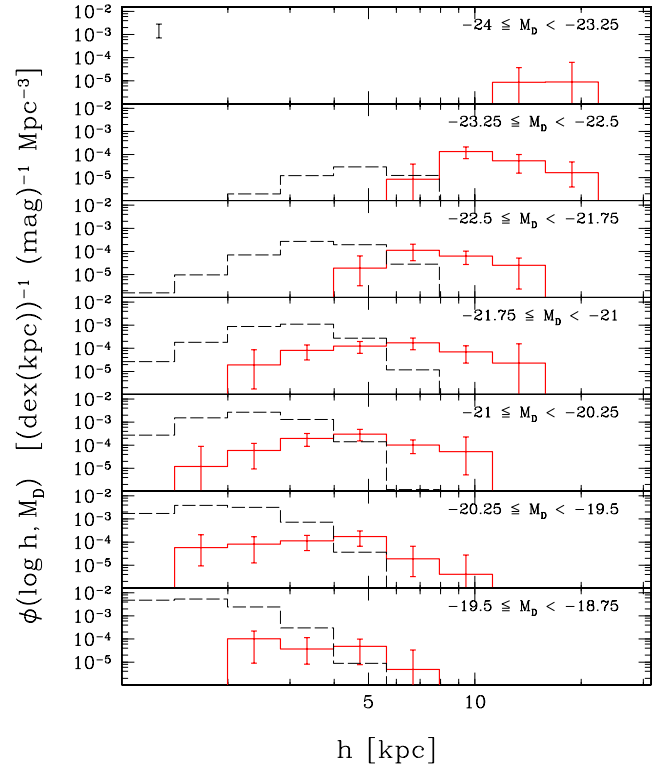


Figure 15. The space density of EFAR galaxies as a function of disc magnitude and disc scalelength. Errorbars are the same as in Fig. 12. The dashed histogram is the space density distribution derived by Blanton et al. (2003) from SDSS data for galaxies with $n < 1.5$, scaled as described in the text.

their initial orbital energy? Do they form disc-like components because they predominantly fall into the cluster on the central galaxies along the filamentary structures of the cosmic web?

While attractive from the galaxy structure point of view, this picture leaves some other properties of elliptical galaxies unexplained. Elliptical galaxies in clusters obey a colour–magnitude relation up to redshifts of ~ 0.9 (e.g. Stanford, Eisenhardt & Dickinson 1998), which means that this relation is driven by a mass–metallicity relation, with bigger galaxies being more metal-rich. This therefore means that the protogalaxies that formed the big spheroids in the above picture would have to have a higher metallicity before they merged than the protogalaxies of smaller galaxies. Big elliptical galaxies also have higher α -element over Fe abundance ratios than smaller galaxies, meaning that their star formation time-scales must have been shorter (e.g. Mehlert et al. 2003). In the hierarchical galaxy formation picture this could be accomplished by having the protogalaxies that will merge into big galaxies form early on top of an overdensity of what will later become a cluster of galaxies at redshifts much larger than 3. In a Λ -cold dark matter universe there would be a lot of early merging, causing vigorous and short time-scale star formation, leading to non-Solar rates of α -element overabundances. Because these protogalaxies would already live in the deeper potential well of what will later become a large galaxy in a cluster of galaxies, they would be better at holding on to their enriched gas and therefore be able to create the mass–metallicity relation. However, in the above picture galaxies falling in along the filaments would have lower (over)abundances and would potentially be somewhat younger, and hence the disc components should have a lower metallicity and age. While de Jong & Davies (1997) show that discy elliptical galaxies indeed have line indices

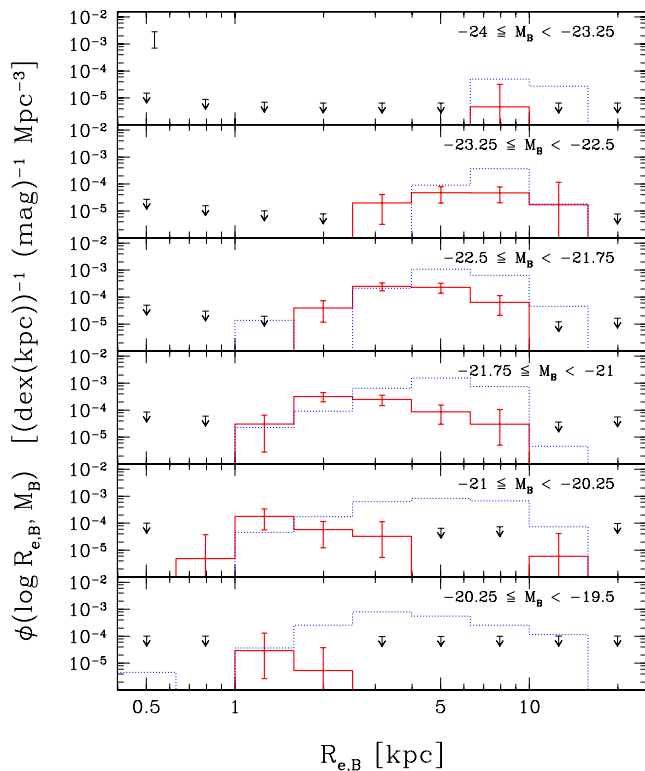


Figure 16. The space density as a function of bulge magnitude and bulge effective radius for the EFAR galaxies (solid histogram) and for the semi-analytic models of Cole et al. (2000) (dotted histogram). Errorbars and upper limits are the same as in Fig. 12.

indicating younger ages and/or lower metallicities, spatially resolved line indices mapping as performed by, for instance, the SAURON team (Falcón-Barroso et al. 2004) will be needed to fully access this model.

Fig. 16 shows the comparison of the (bulge magnitude, bulge effective radius)-bivariate space density of the EFAR elliptical galaxies with the models of Cole et al. (2000). The models are overpredicting the effective radii of early-type bulges by approximately a factor of 2, and even more so at fainter magnitudes. This is the same problem as identified in Fig. 13; the effective radii of low-luminosity early-type galaxies are smaller than those of spiral galaxies of similar luminosity. Cole et al. assumed that all early-type galaxies were created by merging in their models, and as their models do a reasonable job in modelling the spiral galaxy scale size distribution (de Jong & Lacey 2000), the models will overpredict the early-type galaxy scale size in the same way as was discussed for Fig. 13.

4 SUMMARY AND CONCLUSIONS

We have performed 2D bulge/disc decompositions on a sample of 558 early-type galaxies from the EFAR sample using the GIM2D package. In contrast with most earlier work, we have used a Sérsic luminosity profile for the spheroidal component, while using the common exponential light profiles for the disc-like component. We showed with extensive testing on model galaxies that the total galaxy parameters and disc parameters can be recovered at the 10–15 per cent level, but that the bulge parameters are only recovered at the 30 per cent level, mainly due to degeneracies resulting in correlated errors between n_B and R_e . Our galaxy selection criteria are well

known and we have calculated bivariate space densities of elliptical galaxies as a function of galaxy parameters.

Without kinematic or population information it is impossible to prove that our bulge and disc parameter are physically meaningful, but under the assumption that the spheroid and disc parameters represent real galaxy components we derive the following conclusions.

(i) The scale sizes of the bulge and disc components are correlated. This may be partly due to the exclusion of spiral galaxies from the sample, excluding objects with small bulges and large discs. However, it should be noted that spiral galaxies show a similar trend (Courteau et al. 1996; de Jong 1996b; MacArthur et al. 2003).

(ii) The bulge-to-disc B/T ratios, the Sérsic n_B bulge shape parameters, and the bulge effective radii show positive correlations with each other. We have no model that tells us which is cause and which is effect.

(iii) The median n_B value is 3.24 for all galaxies in our sample, 3.66 for all non-spiral, bulge-dominated galaxies in our sample and 4 for non-spiral galaxies with $B/T > 0.7$. Given that most spiral galaxies have bulge n_B values much lower than these values, this means that the standard $R^{1/4}$ de Vaucouleurs (i.e. $n_B = 4$) fitted to spheroids may not be the most appropriate choice when performing bulge/disc decompositions on a random set of galaxies. In cases where one does not have the signal-to-noise ratio or the spatial resolution to perform a full Sérsic bulge plus exponential disc fit, it will be more correct to use a fixed value of $n_B \sim 3.5$ for the brightest galaxies. Lower values should be used for less luminous galaxies.

(iv) The (luminosity, scale size) bivariate space density distribution of bright elliptical galaxies is well described by the analytic parametrization presented by de Jong & Lacey (2000).

(v) Comparing the total galaxy (luminosity, scale size) bivariate distributions of EFAR non-spiral galaxies with SDSS disc-like galaxies shows that the brightest early-type galaxies could, in principle, be formed by merging a few large, current-day spiral galaxies. Low-luminosity early-type galaxies have much smaller scale sizes than spiral galaxies of the same luminosity, and hence they can surely not be created from simple merging of two current-day spiral galaxies.

(vi) When comparing the EFAR bulge and disc (luminosity, scale size) bivariate distributions with those of the SDSS disc-like galaxies, we showed that the bulges of early-type galaxies are typically a factor of 1.5–2.5 smaller and the discs a factor of 2 larger than current-day disc galaxies at a given luminosity.

We speculate that the spheroidal components are the result of merging of similar sized small protogalaxies galaxies at high redshifts, while the disc-like component may be the result of smaller galaxies falling in later along the filaments of the cosmic web. While attractive from the galaxy structure point of view, this model requires careful tuning of the formation process of the protogalaxies at high redshift in order to reproduce the mass–metallicity relation and the α -element overabundance observed at lower redshifts.

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