# Rubidium and lead abundances in giant stars of the globular clusters M4 and M5 ${ }^{1}$ 

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#### Abstract

We present measurements of the neutron-capture elements Rb and Pb for bright giants in the globular clusters M4 and M5. The clusters are of similar metallicity $([\mathrm{Fe} / \mathrm{H}] \simeq-1.2)$ but M 4 is decidedly $s$-process enriched relative to $\mathrm{M} 5:[\mathrm{Ba} / \mathrm{Fe}]=+0.6$ for M4 but 0.0 for M5. The Rb and Pb abundances were derived by comparing synthetic spectra with high-resolution, high signal-to-noise ratio spectra obtained with MIKE on the Magellan telescope. Abundances of Y, Zr , La , and Eu were also obtained. In M 4 , the mean abundances from 12 giants are $[\mathrm{Rb} / \mathrm{Fe}]=0.39 \pm 0.02(\sigma=0.07),[\mathrm{Rb} / \mathrm{Zr}]=0.17 \pm 0.03(\sigma=0.08)$, and $[\mathrm{Pb} / \mathrm{Fe}]=0.30 \pm 0.02(\sigma=0.07)$. In M5, the mean abundances from two giants are $[\mathrm{Rb} / \mathrm{Fe}]=0.00 \pm 0.05(\sigma=0.06),[\mathrm{Rb} / \mathrm{Zr}]=0.08 \pm 0.08(\sigma=0.11)$, and $[\mathrm{Pb} / \mathrm{Fe}]=-0.35 \pm 0.02(\sigma=0.04)$. Within the measurement uncertainties, the


abundance ratios $[\mathrm{Rb} / \mathrm{Fe}],[\mathrm{Pb} / \mathrm{Fe}]$ and $[\mathrm{Rb} / \mathrm{X}]$ for $\mathrm{X}=\mathrm{Y}, \mathrm{Zr}$, La are constant from star-to-star in each cluster and none of these ratios are correlated with O or Na abundances. While M4 has a higher Rb abundance than M 5 , the ratios $[\mathrm{Rb} / \mathrm{X}]$ are similar in both clusters indicating that the nature of the $s$-products are very similar for each cluster but the gas from which M4's stars formed had a higher concentration of these products.

Subject headings: Galaxy: Abundances, Galaxy: Globular Clusters: Individual: Messier Number: M4, Galaxy: Globular Clusters: Individual: Messier Number: M5, Stars: Abundances

## 1. Introduction

Globular clusters continue to provide a source of fascination and frustration to both theorists and observers. Two notable accomplishments include the use of globular clusters to (a) check the age of the Universe (e.g., Gratton et al. 2003) and to (b) test and refine our understanding of stellar evolution (e.g., Renzini \& Fusi Pecci 1988). Despite these successes, globular clusters continue to present bewildering puzzles. The most persistent puzzle relates to chemical composition.

For many years, globular clusters have been known to exhibit star-to-star abundance variations for the light elements C, N, O, Na, Mg, and Al (e.g., see reviews by Smith 1987, Kraft 1994, and Gratton et al. 2004). While the amplitude of the star-to-star abundance dispersion can vary from cluster to cluster, the now familiar anticorrelations between C and $\mathrm{N}, \mathrm{O}$ and Na , and Mg and Al reveal that the abundance variations are likely produced during hydrogen burning at high temperatures via the $\mathrm{CNO}, \mathrm{Ne}-\mathrm{Na}$, and $\mathrm{Mg}-\mathrm{Al}$ cycles. (The O-Na and $\mathrm{Mg}-\mathrm{Al}$ anticorrelations are not seen in field stars.) However, the stars responsible for the nucleosynthesis and the nature of the pollution mechanism(s) remain poorly understood (see Lattanzio \& Tout 2006 for a recent summary).

One possible explanation for the observed abundance anomalies is internal mixing and nucleosynthesis (e.g., Sweigart \& Mengel 1979; Charbonnel 1995) within the present cluster members, the so-called evolutionary scenario. The systematic variation of the C and N (Suntzeff \& Smith 1991) and Li (Grundahl et al. 2002) abundances with luminosity along

[^0]the red giant branch demand an evolutionary component to the star-to-star abundance variations. Dredge-up of CN-cycled material accounts for the C and N variations. Development of a giant's convective envelope leading to mixing with highly Li-depleted gas accounts for the decline of the Li abundance with increasing luminosity. The proton-capture reactions causing the $\mathrm{O}, \mathrm{Na}, \mathrm{Mg}$, and Al variations demand much higher temperatures and much deeper mixing than those required for CN -cycling. Such mixing is not predicted by standard theoretical models of red giants and the discovery of the $\mathrm{O}, \mathrm{Na}, \mathrm{Mg}$, and Al anomalies in main sequence stars (e.g., Briley et al. 1996, Gratton et al. 2001) eliminates deep mixing as a viable explanation for the $\mathrm{O}-\mathrm{Al}$ variations. The interiors of main sequence stars are too cool to process Ne to Na or Mg to Al . Therefore, the cluster gas must have been inhomogeneous when the present stars were formed. This alternative explanation for the abundance anomalies is the so-called primordial scenario.

In the primordial scenario, intermediate-mass asymptotic giant branch stars (IM-AGBs) from the generation to which the observed stars belong have long been considered candidates for synthesizing the abundance variations (Cottrell \& Da Costa 1981). In IM-AGBs, the convective envelope can reach the top of the hydrogen-burning shell, a process called hotbottom burning. For sufficiently massive and metal-poor AGBs, the temperatures at the base of the convective envelope can exceed 100 million degrees thereby allowing the efficient operation of the CNO, Ne-Na, and Mg-Al cycles (e.g., Karakas \& Lattanzio 2003). That IM-AGBs do not alter the abundances of the alpha or iron-peak elements (as required by observations) adds to their qualitative appeal. However, quantitative tests reveal problems with the IM-AGB primordial scenario. Theoretical yields from IM-AGBs combined with a chemical evolution model (Fenner et al. 2004) suggest that O is not sufficiently depleted, Na is overproduced, Mg is produced rather than destroyed, the isotope ratios of Mg do not match the observations, and the sum of $\mathrm{C}+\mathrm{N}+\mathrm{O}$ increases substantially in contrast to the observations. Ventura \& D'Antona (2005a.b.c) find that many of the flaws noted above can be alleviated when IM-AGB yields are calculated using a revised treatment for convection and mass-loss. However, Ventura \& D'Antona note that problems persist, namely with the Mg isotope ratios, and warn that the predictive power of the current AGB models is limited. Recently, Prantzos \& Charbonnel (2006), Smith (2006), and Decressin et al. (2007) suggest that the winds from massive stars may be more promising candidates than IM-AGBs. There is no satisfactory explanation for the complex patterns for the light element abundances exhibited by every well studied Galactic globular cluster. Therefore, our present understanding of globular cluster chemical evolution and/or stellar nucleosynthesis is incomplete.

Determinations of the stellar abundances of the trans-iron or heavy elements offer clues to the history behind the chemical evolution of globular clusters. Here, we provide novel information - the Rb and Pb abundances - for giants in M4 and M5, a pair of clusters of
similar metallicity but with distinctly different levels of $s$-process products. The quintessential $r$-process element Eu has similar abundances in the two clusters and, indeed, across the collection of Galactic globular clusters. In sharp contrast, the $s$-process products are more evident in M4 than in M5 and other clusters of similar metallicity: $[\mathrm{Ba} / \mathrm{Fe}]$ is about +0.6 in M4 but 0.0 in M5. The questions - Are there differences in the Rb and Pb abundances between this pair of clusters? and Are the star-to-star variations in the abundances of light elements ( $\mathrm{O}, \mathrm{Na}, \mathrm{Mg}$, and Al ) reflected in variations among the abundances of Rb and Pb ? - seem likely to probe the origins of the $s$ - and $r$-process products for globular clusters.

Due to a critical branching point in the $s$-process path at ${ }^{85} \mathrm{Kr}$, the abundance of Rb relative to Sr , Y , or Zr can differ by a factor of 10 depending upon the neutron density at the $s$-process site. In the case of AGB stars, the neutron density in the He-shell is dependent on the stellar mass (e.g., see Tomkin \& Lambert 1983, Lambert et al. 1995, Busso et al. 1999, and Abia et al. 2001 for further details). Since the isotopes of Pb and Bi are the last stable nuclei on the $s$-process path, the $s$-process terminates at these elements and overabundances of Pb and Bi will arise if seed nuclei are shuffled by neutron captures down the entire $s$ process path. In particular, metal-poor AGB stars may produce large overabundances of Pb and Bi if the neutron supply per seed exceeds a critical value (e.g., see Goriely \& Siess 2001, Travaglio et al. 2001, and Busso et al. 2001 for further details). The suspicion is that the star-to-star abundance variations for light elements are due to contamination by IM-AGBs. Some contend that IM-AGBs also synthesize $s$-process nuclides and then one might expect to see star-to-star variations in the Rb and Pb abundances as well as correlations with light element abundances.

To further examine the possible role of IM-AGBs in the chemical evolution of globular clusters, Yong et al. (2006b) measured Rb and Pb in NGC 6752 and M13, the two clusters that exhibit the largest amplitude for Al variations. It was found that the abundance ratios $[\mathrm{Rb} / \mathrm{Zr}]$ and $[\mathrm{Pb} / \mathrm{Fe}]$ were constant from star-to-star within the measurement uncertainties. If IM-AGBs do synthesize Rb and Pb , then they may not be responsible for the abundance variations. On the other hand, if IM-AGBs are responsible for the abundance variations, they cannot synthesize Rb or Pb .

In this paper, we extend the measurements of Rb and Pb to the globular clusters M4 and M5. While these clusters are more metal-rich than NGC 6752 or M13, both M4 and M5 are known to exhibit large dispersions and correlations for the light element abundances [see pioneering studies on CN bimodality by Norris (1981) and Smith \& Norris (1983) as well as recent high-resolution spectroscopic studies by Ivans et al. (1999, 2001), Ramírez \& Cohen (2003), and references therein]. In particular, as noted above, M4 is remarkably, perhaps uniquely among globular clusters, enriched in $s$-process products.

## 2. Observations and data reduction

The targets included 12 stars in M4 previously studied by Ivans et al. (1999) and two stars in M5 previously studied by Ivans et al. (2001) and Ramírez \& Cohen (2003). The main focus was to observe a large number of stars in M4 and though we were restricted to the brightest giants, we note that the sample spans a considerable range of the known star-to-star elemental abundance variations. The smaller sample in M5 is due to the fact that those observations were conducted during other observing programs when primary targets were unavailable. In Table [1, we list the program stars. The stellar identifications for M4 and M5 are from Lee (1977) and Arp (1962) respectively.

The observations were performed with the Magellan Telescope using the Magellan Inamori Kyocera Echelle spectrograph (MIKE; Bernstein et al. 2003) on 2004 June 12-13, July 16 , and July 18. A $0.35^{\prime \prime}$ slit was used providing a resolving power of $\mathrm{R} \equiv \lambda / \Delta \lambda=55,000$ in the red and $\mathrm{R}=65,000$ in the blue per 4 pixel resolution element with wavelength coverage from $3800 \AA$ to $8500 \AA$. The IRAF 2 package of programs was used for most of the data reduction. In order to correct for the fact that the lines are severely tilted with respect to the orders and the tilt varies across the CCD, we used the mTOOLS 3 set of tasks written by Jack Baldwin to extract the spectral orders. In the one-dimensional wavelength-calibrated normalized spectra, the typical signal-to-noise ratio was 70 per pixel at $4050 \AA$ ( 140 per resolution element) and 400 per pixel at $7800 \AA$ ( 800 per resolution element).

## 3. Analysis

### 3.1. Stellar parameters and the iron abundance

The required stellar parameters for an abundance analysis are the effective temperature $\left(T_{\text {eff }}\right)$, the surface gravity $(\log g)$, and the microturbulent velocity $\left(\xi_{t}\right)$. Our analysis techniques closely follow Yong et al. (2006b) in which we determined these values adopting a traditional spectroscopic approach. Using routines in IRAF, we measured the equivalent widths (EWs) for a set of FeI and Feir lines. The set of Fe lines was identical to those used by Yong et al. (2006a) for which the $q f$-values were taken from Lambert et al. (1996), Biémont et al. (1991), Blackwell et al. (1995) and references therein. Model atmospheres

[^1]were taken from the Kurucz (1993) local thermodynamic equilibrium (LTE) stellar atmosphere grid and we interpolated within the grid when necessary. We used the LTE stellar line analysis program moog (Sneden 1973) to determine abundances for a given line. We adjusted $T_{\text {eff }}$ until there was no trend between the abundances from Fe ilines and the lower excitation potential. We adjusted $\log g$ until the abundances from Fei and Fe ir lines were in agreement. We adjusted $\xi_{t}$ until there was no trend between the abundances from FeI lines and EW. This process was iterated until all three parameters were simultaneously constrained. The final $[\mathrm{Fe} / \mathrm{H}]$ was the mean from all Fe lines assuming a solar abundance $\log \epsilon(\mathrm{Fe})=7.50$. The stellar parameters are given in Table 1 and we estimate the internal errors to be $T_{\text {eff }} \pm 50 \mathrm{~K}, \log g \pm 0.2 \mathrm{dex}$, and $\xi_{t} \pm 0.2 \mathrm{~km} \mathrm{~s}^{-1}$.

As an additional check on our surface gravities (and analysis techniques), we compared our $\log g$ values with the $Y^{2}$ isochrones (Demarque et al. 2004). We adopted the 13 Gyr isochrones with $[\alpha / \mathrm{Fe}]=+0.3$, our spectroscopic $T_{\text {eff }}$, and interpolated between the two closest metallicities $\mathrm{z}=0.004$ and $\mathrm{z}=0.001$. We found that our surface gravities were in fair agreement with the $Y^{2}$ isochrones, $\log g_{\text {spec }}-\log g_{\text {isochrone }}=-0.22 \pm 0.07(\sigma=0.26)$.

The stellar parameters are in good agreement with those obtained by Ivans et al. (1999, 2001) and Ramírez \& Cohen (2003). For the 12 M 4 giants, the mean differences (This Study - Ivans et al. (1999)) are $\Delta T_{\text {eff }}=71 \pm 17(\sigma=58 \mathrm{~K}), \Delta \log g=0.02 \pm 0.05(\sigma=0.17)$, $\Delta \xi_{t}=0.04 \pm 0.04\left(\sigma=0.15 \mathrm{~km} \mathrm{~s}^{-1}\right)$, and $\Delta[\mathrm{Fe} / \mathrm{H}]=-0.05 \pm 0.01(\sigma=0.05)$. For the 2 M5 giants, the mean differences (This Study - Ramírez \& Cohen (2003)) are $\Delta T_{\text {eff }}=20 \pm$ $5(\sigma=7 \mathrm{~K}), \Delta \log g=-0.15 \pm 0.15(\sigma=0.21), \Delta \xi_{t}=-0.16 \pm 0.03\left(\sigma=0.04 \mathrm{~km} \mathrm{~s}^{-1}\right)$, and $\Delta[\mathrm{Fe} / \mathrm{H}]=0.05 \pm 0.02(\sigma=0.03)$. For M5 IV-81, our $T_{\text {eff }}$ is 105 K higher, $\log g$ is 0.3 dex higher, $\xi_{t}$ is $0.02 \mathrm{~km} \mathrm{~s}^{-1}$ lower, and $[\mathrm{Fe} / \mathrm{H}]$ is 0.12 dex higher than the values determined by Ivans et al. (2001).

### 3.2. Rubidium abundances

The abundances for Rb were determined via spectrum synthesis of the $7800 \AA \mathrm{Rb}$ l line (see Figure 1). For all stars in our sample, this Rb line is blended with a Si I line as well as with weak CN lines and therefore an equivalent width analysis is not possible. Even in the coolest stars, the Rb line is only $10 \%$ deep relative to the continuum such that accurate abundances can be derived only from high-resolution, high signal-to-noise ratio spectra. The adopted wavelengths and relative strengths for the isotopic and hyperfine-structure components of Rb were identical to Lambert \& Luck (1976) and Tomkin \& Lambert (1999). We assumed a solar isotope ratio ${ }^{85} \mathrm{Rb} /{ }^{87} \mathrm{Rb}=3$ and the macroturbulent broadening was fixed by fitting the profile of the nearby $7798 \AA$ NiI line. Synthetic spectra were generated using Moog
and the Si and Rb abundances were varied to obtain the best fit to the observed spectrum. Our tests confirmed the finding by Lambert \& Luck (1976) that Rb isotope ratios cannot be measured from the $7800 \AA$ line due to the hyperfine structure and the small isotopic shift. Not surprisingly, our tests showed that the derived Rb abundances are insensitive to the assumed isotope ratio. In Yong et al. (2006b), we took the Kurucz et al. (1984) solar atlas and measured an abundance $\log \epsilon(\mathrm{Rb})=2.58$ using a model atmosphere with $T_{\text {eff }} / \log g / \xi_{t}=5770 / 4.44 / 0.85$. Our solar Rb abundance is in excellent agreement with the Grevesse \& Sauval (1998) value, $\log \epsilon(\mathrm{Rb})=2.60$. The weaker Rbi resonance line near $7947 \AA$ is detected but in a region affected by unidentified blends, atmospheric absorption, and fringing. While preliminary analyses suggest that the abundances derived from the $7947 \AA$ line are similar to those derived from the $7800 \AA$ line (see Yong et al. 2006 b for a more detailed comparison), we restrict our analyses to the $7800 \AA$ line.

In this study, we included CN molecular lines when fitting the $7800 \AA \mathrm{Rb}$ line. The CN lines were taken from Kurucz (2006). While the inclusion of these lines resulted in an improved fit to the local continuum, we note that the Rb abundances were unaffected. Even for the coolest stars in M4 for which the CN line strength was at a maximum, the Rb abundances increased by $\leq 0.02$ dex if the CN lines were omitted. For the more metal-poor cluster NGC 6752, we re-analyzed a subset of the Yong et al. (2006b) spectra and found that the Rb abundances were unchanged when CN lines were included in the analysis.

### 3.3. Lead abundances

The abundances for Pb were determined via spectrum synthesis of the $4058 \AA \mathrm{~Pb}$ l line (see Figure 2). High-resolution, high signal-to-noise ratio spectra are essential for measuring Pb abundances because the region near $4058 \AA$ is crowded with molecular lines of CH as well as $\mathrm{Mg}, \mathrm{Ti}, \mathrm{Mn}, \mathrm{Fe}$, and Co atomic lines. Since M4 and M5 are more metal-rich than NGC 6752 and M13, the $4058 \AA$ region is more crowded and the uncertainties in the derived Pb abundances are greater. However, we note that our syntheses provide a very good fit to the observed spectra. The macroturbulent broadening was fixed by fitting the profiles of nearby lines. The $g f$-value of the Pb i line was the same as that used by Aoki et al. (2002) as were the hyperfine-structure and isotopic components. We assumed a solar isotope ratio and our tests confirmed that the derived Pb abundances do not depend upon the assumed isotope ratio. For the solar Pb abundance, we adopted $\log \epsilon(\mathrm{Pb})=1.95$ from Grevesse \& Sauval (1998).

### 3.4. Additional elements

The abundances for $\mathrm{O}, \mathrm{Na}, \mathrm{Mg}, \mathrm{Al}, \mathrm{Y}, \mathrm{Zr}, \mathrm{La}$, and Eu were also measured in the program stars using the same lines as Yong et al. (2005). Y and Zr were selected since the ratios $[\mathrm{Rb} / \mathrm{Y}]$ and $[\mathrm{Rb} / \mathrm{Zr}]$ are sensitive to the neutron density at the site of the $s$-process. $\mathrm{O}, \mathrm{Na}, \mathrm{Mg}$, and Al were measured since they are known to vary from star-to-star in these and other globular clusters. Re-measuring these elements in M4 and M5 provides a check to see if our abundance determinations are on the same scale as other investigators. Similarly, the neutron-capture elements La and Eu were measured (hyperfine and/or isotopic splitting was included). When deriving O abundances, we generated synthetic spectra to account for possible blending from CN molecular lines. We assumed C and N abundances interpolated as a function of Na with the C and N abundances taken from Ivans et al. (1999, 2001) and Smith et al. (1997). A subset of abundance determinations were conducted using spectrum synthesis to account for possible blends. In Table 2 we present the measured elemental abundances for $\mathrm{O}, \mathrm{Na}, \mathrm{Mg}, \mathrm{Al}, \mathrm{Rb}, \mathrm{Y}, \mathrm{Zr}, \mathrm{La}, \mathrm{Eu}$, and Pb . The adopted solar abundances were $8.69,6.33,7.58,6.47,2.60,2.24,2.60,1.13,0.52$, and 1.95 respectively. These values are identical to those we used in previous papers and were taken from Grevesse \& Sauval (1998) for all elements except O (Allende Prieto et al. 2001), La (Lawler et al. 2001a), and Eu (Lawler et al. 2001b). Table 2 also contains the Zr abundances shifted onto the Smith et al. (2000) scale.

For M4, the measured elemental abundances are in good agreement with Ivans et al. (1999). The mean differences (This Study - Ivans et al. (1999)) are $\Delta[\mathrm{O} / \mathrm{Fe}]=0.34 \pm 0.03$ $(\sigma=0.11), \Delta[\mathrm{Na} / \mathrm{Fe}]=0.14 \pm 0.03(\sigma=0.12), \Delta[\mathrm{Mg} / \mathrm{Fe}]=0.14 \pm 0.02(\sigma=0.06), \Delta[\mathrm{Al} / \mathrm{Fe}]$ $=0.05 \pm 0.03(\sigma=0.10), \Delta[\mathrm{La} / \mathrm{Fe}]=0.01 \pm 0.03(\sigma=0.10)$, and $\Delta[\mathrm{Eu} / \mathrm{Fe}]=0.03 \pm 0.02$ $(\sigma=0.07)$. The agreement for $\mathrm{Al}, \mathrm{La}$, and Eu is excellent. The largest discrepancy is for O and differences in the adopted solar abundances can account for 0.24 dex leaving a 0.10 dex residual. While both studies use the same O lines and $g f$ values, they employ spectrum synthesis which presumably involves a different set of atomic and molecular lines. A small offset also exists for Na and Mg which could be due to the different set of lines employed.

For M5, the elemental abundances are in fair agreement with Ramírez \& Cohen (2003). The mean differences (This Study - Ramírez \& Cohen (2003)) are $\Delta[\mathrm{O} / \mathrm{Fe}]=0.27 \pm 0.14$ $(\sigma=0.20), \Delta[\mathrm{Na} / \mathrm{Fe}]=-0.05 \pm 0.10(\sigma=0.14), \Delta[\mathrm{Mg} / \mathrm{Fe}]=0.11 \pm 0.06(\sigma=0.08)$, $\Delta[\mathrm{Zr} / \mathrm{Fe}]=-0.05 \pm 0.19(\sigma=0.27), \Delta[\mathrm{La} / \mathrm{Fe}]=-0.03 \pm 0.03(\sigma=0.04)$, and $\Delta[\mathrm{Eu} / \mathrm{Fe}]$ $=-0.06 \pm 0.09(\sigma=0.13)$. For O , the difference can be attributed to the adopted solar abundance. For the single star also studied by Ivans et al. (2001), the measured abundances are in good agreement $\Delta[\mathrm{X} / \mathrm{Fe}] \leq 0.07$ dex. The difference for La is 0.14 dex (possibly due to the different set of lines) and 0.58 dex for O . The adopted solar abundances can account
for 0.24 dex leaving a 0.34 dex residual. Dr I. Ivans kindly sent the synthetic and observed spectra near the $6300 \AA$ and $6363 \AA$ used to derive O abundances. The difference in the measured O abundances between the two studies is $\Delta \log \epsilon(\mathrm{O})=0.14$ dex. The difference in the measured Fe abundance is 0.12 dex. The remaining 0.08 dex residual may be attributed to differences in the spectrum synthesis analysis.

The adopted line lists and measured equivalent widths are presented in Table 3. In Table 4, the abundance dependences on the model parameters are given.

## 4. Discussion

### 4.1. The heavy-element canvas

The heavy elements are synthesized by neutron-captures in the $s$ - and the $r$-process. (We ignore here the nuclides, all of low abundance, referred to as $p$-nuclides.) In all but material dominated by $s$-processed products, Eu is a signature element for the $r$-process. There is evidence from analyses of field stars, especially those severely enriched in $r$-process products that the relative abundances for the $r$-process of nuclides from Ba to Eu are invariant with metallicity (e.g., see Cowan \& Sneden 2006). This invariance does not extend to the ratio of heavy products (i.e., Eu ) to light products (i.e., $\mathrm{Sr}, \mathrm{Y}$, and Zr ). In a solar mix of elements, Ba is taken as a measure of the $s$-process with but a slight contamination from the $r$-process (Burris et al. 2000). The [Ba/Eu] ratio of a star is widely taken to indicate the relative mix of $s$ - to $r$-processed material.

In the case of ratios such as $[\mathrm{Eu} / \mathrm{Fe}]$ and $[\mathrm{Ba} / \mathrm{Fe}]$ for globular cluster stars, three questions arise: Is there an intercluster spread in the ratios? Is there an intra-cluster spread, especially for clusters exhibiting a large spread among light element abundances? How do the cluster ratios compare with those found for the field stars, as a function of $[\mathrm{Fe} / \mathrm{H}]$ ?

### 4.1.1. The $[E u / F e]$ ratio

With but the single known exception of Ruprecht 106 , the $[\mathrm{Eu} / \mathrm{Fe}]$ ratio across the collection of examined Galactic globular clusters with $[\mathrm{Fe} / \mathrm{H}]$ from about -1.0 to -2.5 appears to be single-valued. Values of $[\mathrm{Eu} / \mathrm{Fe}]$ in the literature for well-sampled clusters with $[\mathrm{Fe} / \mathrm{H}]$ $<-1$ span a reported range of only 0.2 dex from about +0.35 to +0.55 with no obvious trend with $[\mathrm{Fe} / \mathrm{H}]$. [Ruprecht 106 is an outlier at 0.0 from analyses of two stars (Brown et al. 1997). M80 is probably also an outlier at +1.0 , but the analyses were based on less than ideal
spectra, $\mathrm{R}=18,000$ (Cavallo et al. 2004).] This range is gathered from examination of the extensive literature; references cited by Gratton et al. (2004) were read and supplemented by more recent papers. No attempt has been made by us to correct published values to a common standard. It would be premature to claim that the 0.2 dex spread in $[\mathrm{Eu} / \mathrm{Fe}]$ is real and thus not entirely attributable to cumulative errors of measurement. Our results $[\mathrm{Eu} / \mathrm{Fe}]$ $=+0.41$ for M 4 and +0.53 for M5 fall within this range and, as noted above, are consistent with previously published results.

Although our samples are small (see also Yong et al. 2006), the stars chosen for their contrasting light element abundances show no evidence of a dependence on these light element abundances. This decoupling of light element abundances from that of Eu is confirmed by previous studies of well-sampled clusters. Several authors have compared $[\mathrm{Eu} / \mathrm{Fe}]$ for globular clusters with results for Galactic field stars. The evidence is that the $[\mathrm{Eu} / \mathrm{Fe}]-$ $[\mathrm{Fe} / \mathrm{H}]$ trends (i.e., the apparent independence of $[\mathrm{Eu} / \mathrm{Fe}]$ on $[\mathrm{Fe} / \mathrm{H}]$ for $[\mathrm{Fe} / \mathrm{H}]<-0.7$ ) are the same for both samples (see, for example, Gratton et al. (2004), James et al. (2004), and Pritzl et al. (2005)).

In summary, globular cluster and field stars appear with remarkably few exceptions and to within measurement errors to have the same $[\mathrm{Eu} / \mathrm{Fe}]$ ratio. Our measurements of this ratio for M4 and M5 are consistent with previous measurements for these clusters and the common cluster-field star set of data.

### 4.1.2. The $[Y / F e],[Z r / F e],[B a / F e]$, and $[L a / F e]$ ratios

Among the few heavy elements in abundance analyses of globular cluster stars from which the $s$-process contribution may be assessed, Ba is the most widely reported. Other elements for which abundance data are available for several clusters include $\mathrm{Y}, \mathrm{Zr}$, and La.

For most globular clusters, the $[\mathrm{Ba} / \mathrm{Fe}]$ ratio is positive $([\mathrm{Ba} / \mathrm{Fe}] \simeq+0.2)$ and independent of $[\mathrm{Fe} / \mathrm{H}]$ over the interval -0.7 to about -1.5 . For the few investigated clusters with $[\mathrm{Fe} / \mathrm{H}] \leq-1.8,[\mathrm{Ba} / \mathrm{Fe}]$ declines to roughly 0.0 at $[\mathrm{Fe} / \mathrm{H}]$ of -2.0 to -2.5 . This run with $[\mathrm{Fe} / \mathrm{H}]$ for the clusters is matched by the field stars (James et al. 2004). The spread in $[\mathrm{Ba} / \mathrm{Fe}]$ at a given $[\mathrm{Fe} / \mathrm{H}]$ is about $\pm 0.2$ dex, a value consistent with the errors of measurement, with several outliers. The sole outlier with $[\mathrm{Ba} / \mathrm{Fe}]$ above the mean trend is M4 with an excess of about 0.4 dex in $[\mathrm{Ba} / \mathrm{Fe}]$. The outliers below the mean trend are NGC 3201, Ruprecht 106, and Palomar 12 (Brown et al. 1997; Gonzalez \& Wallerstein 1998) as well as NGC 5694 (Lee et al. 2006). Data for [La/Fe] is less extensive but suggests [La/Fe] $\simeq+0.1$ with M4 again as an outlier $([\mathrm{La} / \mathrm{Fe}]=+0.45)$ but with the Ba-outliers NGC 3201
and Pal 12 as conformers to the mean value. NGC 5694 is again an outlier with $[\mathrm{La} / \mathrm{Fe}]=$ -0.26 based on analysis of a single star (Lee et al. 2006). The [Y/Fe] ratio is close to zero (James et al. 2004). Our result for M5 at +0.2 from just two stars may be an (unlikely) outlier. At $[\mathrm{Y} / \mathrm{Fe}]=+0.7, \mathrm{M} 4$ is certainly an outlier.

The discussed data, as well as less extensive data on other heavy elements, show that the vast majority of the clusters are not distinguishable from field stars by their abundances of Ba and other heavy elements. For a given cluster, there is no compelling evidence for a star-to-star variation correlated with the variations seen among the light elements for many clusters. There are, however, examples of individual cluster stars with $s$-process enrichments, e.g., the CH stars in M2 and M55 (Smith \& Mateo 1990; Brilev et al. 1993), the possible CH star in M22 (McClure \& Norris 1977; Vanture \& Wallerstein 1992), a Y-rich star in M13 (Cohen \& Meléndez 2005), and a hint of Zr abundance variations in 47 Tuc (Wylie et al. 2006), M5 (Ramírez \& Cohen 2003), and NGC 6752 (Yong et al. 2005). While M15 exhibits variations of Ba and Eu , these variations are not correlated with the light element abundances and the ratio $[\mathrm{Ba} / \mathrm{Eu}]$ is constant from star-to-star (Sneden et al. 1997; Otsuki et al. 2006).

### 4.2. Rubidium

In M4 the mean Rb abundance is $[\mathrm{Rb} / \mathrm{Fe}]=0.39 \pm 0.02(\sigma=0.07)$ and in M 5 the mean abundance is $[\mathrm{Rb} / \mathrm{Fe}]=0.00 \pm 0.05(\sigma=0.06)$. As expected, M 4 has a higher abundance of Rb than M5. While $50 \%$ of the solar Rb abundance can be attributed to the $r$-process (Burris et al. 2000), M4 has a slightly lower [Eu/Fe] ratio than does M5 and therefore, it would be difficult to attribute any Rb excess in M4 relative to M5 as being due to a larger $r$ process contribution. Neither cluster shows any evidence for a dispersion in Rb abundances. That is, the scatter in the Rb abundances within M4 and M5 is small and can be attributed entirely to the measurement uncertainties. The Rb abundances are not correlated with O or Na (the two elements that exhibit large star-to-star abundance variations).

The Rb abundances of M4 and M5 are compared with the limited data in the literature for other clusters and field stars. In Figure 3, we compare the $[\mathrm{Rb} / \mathrm{Fe}]$ abundance ratios with field dwarfs and giants (Gratton \& Sneden 1994; Tomkin \& Lambert 1999) as well as the globular clusters NGC 3201 (Gonzalez \& Wallerstein 1998), $\omega$ Cen (Smith et al. 2000), and NGC 6752 (Yong et al. 2006b). M5 has [Rb/Fe] ratios that are essentially identical to $\omega$ Cen giants at the same metallicity. [We note that $\omega$ Cen is now regarded as the nucleus of an accreted dwarf spheroidal galaxy (Smith et al. 2000).] These Rb abundances are comparable to the lowest values seen in field stars at the same metallicity. Given that M4 is s-process enriched relative to other globular clusters and field stars, it is puzzling that at the metallicity
of M4, the majority of field stars share the same $[\mathrm{Rb} / \mathrm{Fe}]$ ratio. With the addition of M4 and M 5 , the $[\mathrm{Rb} / \mathrm{Fe}]$ ratios in NGC 6752 appear rather low compared to other globular clusters and field stars. While the sample sizes remain small and the analyses have been performed by different investigators, the globular clusters cover a wider range of $[\mathrm{Rb} / \mathrm{Fe}]$ values than do the field stars in the same metallicity regime.

The relative abundances of $\mathrm{Rb}, \mathrm{Y}$, and Zr are very similar for M 4 and M5. In Figure 4, the upper panel shows $[\mathrm{Rb} / \mathrm{Y}]$ versus $[\mathrm{Y} / \mathrm{Fe}]$ for mean values for M 4 , M 5 , and NGC 6752. The dotted line corresponds to $[\mathrm{Rb} / \mathrm{Y}]=-0.25$, the mean value for the three clusters. Published results for NGC 3201 and $\omega$ Cen do not fall on the line but straddle it. This may indicate zero-point differences between our and other analyses or a real cluster-to-cluster difference among the $\mathrm{Rb}, \mathrm{Y}$, and Zr abundances. Considering the individual measurements, we find $[\mathrm{Rb} / \mathrm{Y}]=-0.29 \pm 0.02(\sigma=0.06)$ for M 4 and $[\mathrm{Rb} / \mathrm{Y}]=-0.17 \pm 0.10(\sigma=0.13)$ for M5. For Zr , we find $[\mathrm{Rb} / \mathrm{Zr}]=0.17 \pm 0.03(\sigma=0.08)$ for M 4 and $[\mathrm{Rb} / \mathrm{Zr}]=0.08 \pm 0.08(\sigma=0.11)$ for M5. (These Zr abundances have been shifted by +0.3 dex onto the Smith et al. 2000 scale, see Yong et al. 2006b for details.) Neither the $[\mathrm{Rb} / \mathrm{Y}]$ nor the $[\mathrm{Rb} / \mathrm{Zr}]$ ratio exhibits a star-to-star dispersion and neither is correlated with O or Na abundances. In Figure 55 we compare the $[\mathrm{Rb} / \mathrm{Zr}]$ abundance ratio with field dwarfs and giants as well as globular cluster giants. Compared to the $\omega$ Cen giants at a similar metallicity, M4 and M5 have $[\mathrm{Rb} / \mathrm{Zr}]$ ratios 0.4-0.5 dex higher (which may be due to systematic offsets). For M4, the difference arises due to $\omega$ Cen having less Rb and more Zr. For M5, the difference results from $\omega$ Cen having similar Rb and considerably more Zr .

The similarity in heavy element abundance ratios involving Rb extends to [ $\mathrm{Rb} / \mathrm{La}]$ - see middle panel of Figure 4. The indication from our analyses is that the relative abundances of Rb to Y, Zr, and La are the same for M5 and NGC 6752 as well as M4, the cluster clearly enriched in $s$-process elements. This demonstrates that the source of the M4 enrichment is identical to that providing the $s$-process elements for the natal material for M5 and NGC 6752. The abundance of Rb relative to the other elements is slightly subsolar, i.e., $[\mathrm{Rb} / \mathrm{Y}]$ $=-0.25$. There is a well known difference between the spectroscopically determined Rb abundance and that obtained from carbonaceous meteorites: $\log \epsilon(\mathrm{Rb})=2.60$ versus 2.33, respectively (Asplund et al. 2005). If the meteoritic value is adopted $[\mathrm{Rb} / \mathrm{Y}]$ and similar ratios are increased by 0.27 dex and rendered indistinguishable from the solar ratios. Then, $s$-process donors to the natal clouds of the globular clusters and of local stars including the Sun must have been similar.

### 4.3. Lead

In M4 the mean Pb abundance is $[\mathrm{Pb} / \mathrm{Fe}]=0.30 \pm 0.02(\sigma=0.07)$ and in M 5 the mean abundance is $[\mathrm{Pb} / \mathrm{Fe}]=-0.35 \pm 0.02(\sigma=0.04)$. We note that $20 \%$ of the solar Pb abundance is attributable to the $r$-process (Burris et al. 2000) but the $[\mathrm{Eu} / \mathrm{Fe}]$ ratio is comparable for M4 and M5. Again, since M4 has a lower $[\mathrm{Eu} / \mathrm{Fe}]$ ratio than M5, the excess Pb in M4 relative to M5 cannot be due to an increased $r$-process contribution for M4. The excess likely arises from more complete processing down the $s$-process chain to its termination at Pb . The difference in $[\mathrm{Pb} / \mathrm{Fe}]$ between M 4 and M 5 is 0.65 dex. For the other neutron-capture elements synthesized via the $s$-process, the difference in $[\mathrm{X} / \mathrm{Fe}]$ between M 4 and M5 is typically 0.4 dex. Regarding the scatter in Pb abundances in M4 and M5, we again find that the small dispersion can be entirely explained by the measurement uncertainties. The Pb abundances in M4 and M5 are not correlated with O or Na .

In Figure 6, we compare the $[\mathrm{Pb} / \mathrm{Fe}]$ abundance ratios with field stars (Sneden et al. 1998; Travaglio et al. 2001) as well as the globular clusters M13 and NGC 6752 (Yong et al. 2006b). Data on Pb in the field stars is obviously very limited. The one certain data point is for HD 23439A, a star enriched in s-process products (Tomkin \& Lambert 1999), with $[\mathrm{Pb} / \mathrm{Fe}]=+0.6$. The other three data points not indicated as upper limits in Figure 6 are marked with the customary ' $\because$ ' by (Travaglio et al. 2001) and inspection of their published spectra shows that an alternative designation as an upper limit is possibly more appropriate. Our Pb abundance determination for HD 141531 (Yong et al. 2006b) is consistent with the sub-solar $[\mathrm{Pb} / \mathrm{Fe}]$ ratios for NGC 6752 , M5, and M13. Although more Pb abundances are needed for field stars, we suppose that the Pb abundances of the globular clusters NGC 6752, M5, and M13 and normal field stars are in fair agreement, as found for $\mathrm{Rb}, \mathrm{Y}, \mathrm{Zr}, \mathrm{Ba}$, and other heavy elements.

In the lower panel of Figure 4, we plot $[\mathrm{Rb} / \mathrm{Pb}]$ versus $[\mathrm{Pb} / \mathrm{Fe}]$ for $\mathrm{M} 4, \mathrm{M} 5$, and NGC 6752. Interestingly, the $[\mathrm{Rb} / \mathrm{Pb}]$ ratios for M 4 and NGC 6752 appear similar which again suggests that the source of the $s$-process enrichment in M4 may be identical to the source of the $s$-process elements for NGC 6752. However, in this Figure, M5 now appears to be the outlier.

Given the small comparison sample, clearly it would be of great interest to expand the measurements of Pb to larger samples of field stars. The synthetic spectra in this study and in Yong et al. (2006b) indicate that measurements of Pb in cool giants in the range $-2.0 \leq$ $[\mathrm{Fe} / \mathrm{H}] \leq-1.0$ are feasible.

### 4.4. Consequences for the primordial scenario

Formation of the Galactic and extragalactic globular clusters remains a subject with many unanswered questions (Brodie \& Strader 2006). From the point of view of the quantitative spectroscopist interested in the composition of stars in the Galactic globular clusters, two questions seem preeminent: (i) Why do the cluster and field stars follow the same $[\mathrm{X} / \mathrm{Fe}]$ versus $[\mathrm{Fe} / \mathrm{H}]$ relations? (ii) What causes the light element abundance variations first observed among cluster giant stars but now found also in subgiant and main sequence stars?

The simplest answer to question (i) would appear to be that the clusters are formed during episodic mergers in the hierarchical formation of the Galaxy, as discussed by Bekki et al. (2002) and Beasley et al. (2002). If the gas content of the protocluster is dominated by that from the Galaxy and not the infalling system, the composition of the cluster stars will be very similar to that of the Galactic stars formed just preceding the merger. Gas remaining after star formation in the cluster will be ejected by supernovae from the massive stars. Star formation will be restricted to a single generation of stars, i.e., stars within a given cluster will have the identical metallicity. Indeed the complete composition of all stars will be the same but for changes resulting from stellar evolution and the agents responsible for the star-to-star variations among light element abundances. These agents comprise what is widely referred to as 'the primordial scenario'.

For the majority of the Galactic globular clusters subjected to abundance analysis, the compositions do closely resemble field stars of the cluster's metallicity. This is as expected on the merger hypothesis. M4 is an apparent exception with its overabundance of heavy elements (Y, Zr, Ba, La etc, but not Eu) matched by very few field stars. Some of the field stars with the cluster's degree of overabundance of $s$-process products are binaries in which the products were transferred to the present star when its companion was an AGB star. There are very few known examples where mass transfer across a binary is not a viable explanation. The case of HD 23439A and B (Tomkin \& Lambert 1999) would appear to be such a case. HD 23439A shows no sign of radial velocity variation during 14 years of monitoring (Latham et al. 2002). On the other hand, HD 23439B is a binary but its period and orbital eccentricity do not suggest mass transfer (Latham et al. 2002).

Other clusters with few counterparts among field stars include Rup 106 and Pal 12. In their case, the Sagittarius dwarf spheroidal galaxy now undergoing destruction by the Galaxy has giants with the peculiar composition of Rup 106 and Pal 12 (Sbordone et al. 2007) indicating that they belong to the dwarf spheroidal galaxy. A conjecture is that the cluster M4 was captured from a stellar system in which chemical evolution had led to $s$ process enrichment; a change in the IMF might have been immediately responsible. However, consideration of space velocities reveals that M4's orbit is restricted to the inner disk and
bulge (apocentric radius, $R_{a}=5.9 \mathrm{kpc}$ ) whereas M5's orbit may be more consistent with a capture event, $R_{a}=35.4 \mathrm{kpc}$ (Dinescu et al. 1999). Given M4's orbit, a comparison of the $s$-process elements in M4 and comparable metallicity stars of the inner disk and bulge would be of great interest.

In the primordial scenario, the star-to-star abundance variations apart from those attributable to dredge-up in giants arise from the accretion of gas (or contamination of natal clouds) to varying degrees by cluster stars. The accreted gas is assumed to be mass lost from the more massive and now dead stars of the cluster. Proposed candidates include IM-AGBs who experience extensive nucleosynthesis and lose their envelopes at low velocity so that the ejecta reside within the cluster's potential (Cottrell \& Da Costa 1981). Recently, a role for massive stars in polluting the cluster's environment has been proposed (Decressin et al. 2007; Smith 2006; Prantzos \& Charbonnel 2006).

Both M4 and M5 exhibit star-to-star abundance variations for the light elements, as does every well studied Galactic globular cluster. The variations involve $\mathrm{O}, \mathrm{Na}, \mathrm{Mg}$, and Al , as well as C and N for which there is a component related to dredge-up by giants. For the majority of the analyzed clusters (M4 is an obvious exception), the elemental abundances identified with the unpolluted cluster stars match those of field stars of the same metallicity. Two clusters have been studied for variations among the isotopic ratios of Mg but the normal stars in M13 and NGC 6752 (i.e., cluster stars with high O, high Mg, low Na, low Al and whose $\mathrm{O}-\mathrm{Al}$ abundances match field stars at the same metallicity) had Mg isotope ratios that exceeded measurements in comparable metallicity field stars (Yong et al. 2003a, b, 2006a).

In this study, we extended the elemental abundance measurements of Rb and Pb to the globular clusters M4 and M5. All four clusters that we have studied (M4, M5, M13, and NGC 6752) show light element variations, but constant Rb and Pb abundances within each cluster. Thus, the accreted gas providing the light element variations cannot contain significant amounts (or significant deficiencies) of Rb and Pb . The accreted gas must have the same proportions of Rb and Pb as the ambient material.

The source of the gas may be IM-AGBs that at least qualitatively account for the range of abundance variations. A quantitative test is impossible at present because the predicted yields from metal-poor IM-AGBs may be model dependent (Busso et al. 2001; Ventura \& D'Antona 2005a.b.c. As far as s-products are concerned which, as we have emphasised, do not show abundance variations, IM-AGBs remain a viable candidate because Lattanzio et al. (2004) and Straniero (2006, private communication) suggest that IM-AGBs will not produce any $s$-process elements since the mass of the He shell and the duration of the thermal pulse decreases with increasing AGB mass.

Massive stars achieve limited synthesis of $s$-process products but are held to be the sites of the weak $s$-process (and the $r$-process) contributing nuclides from the iron group up to about Rb , Y and Zr (Raiteri et al. 1991). Predicted yields are such that massive stars could serve to provide the light element abundance variations without leading to predicted variations for $\mathrm{Rb}, \mathrm{Y}$, and Zr .

In short, the lack of abundance variations for $\mathrm{Rb}, \mathrm{Y}, \mathrm{Zr}$ and other s-process elements may not exclude either IM-AGBs or massive stars as sources of the accreted pollutants.

## 5. Concluding remarks

In this paper we present measurements of the neutron-capture elements Rb and Pb in the globular clusters M4 and M5. While both clusters exhibit star-to-star abundance variations for the light elements, we find that the abundances of Rb and Pb are constant. None of the abundance ratios $[\mathrm{Rb} / \mathrm{Fe}],[\mathrm{Rb} / \mathrm{Zr}]$, and $[\mathrm{Pb} / \mathrm{Fe}]$ are correlated with O or Na abundances. In the primordial scenario, the abundance variations for the light elements are attributed to different levels of accretion of ejecta from IM-AGBs or massive stars. The fact that the heavy elements including Rb and Pb do not show abundance variations implies that the accreted material has the same composition as the ambient material for the heavy elements (i.e., the accreted material cannot be highly underabundant or overabundant in these elements). That the ratios $[\mathrm{Rb} / \mathrm{X}]$ for $\mathrm{X}=\mathrm{Y}, \mathrm{Zr}, \mathrm{La}$ are similar for M 4 and M 5 suggests that the source of the $s$-process elements are similar and that M4 had a greater concentration of these products.

There remains a need to pursue additional observational tests of the primordial scenario. In particular, present data on the Rb and Pb abundances in field and cluster stars are sparse. The indication that the Mg isotopic ratios of unpolluted or normal cluster stars differ from those of field stars of the same metallicity deserves closer scrutiny by, in particular, extending the measurement of these isotopic ratios to additional clusters.

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Fig. 1.- Observed spectra (filled circles) for M4 L3413 (upper panel), M4 L4611 (middle panel), and M5 IV-81 (lower panel) near the $7800 \AA$ Rb feature. Synthetic spectra with different Rb abundances are shown (the solid line represents the best fit).


Fig. 2.- Observed spectra (filled circles) for M4 L1411 (upper panel), M4 L3624 (middle panel), and M5 IV-82 (lower panel) near the $4058 \AA$ Pb feature. Synthetic spectra with different Pb abundances are shown (the solid line represents the best fit).


Fig. 3.- $[\mathrm{Rb} / \mathrm{Fe}]$ vs. $[\mathrm{Fe} / \mathrm{H}]$ for M4 (filled circles) and M5 (open circles). Green triangles show NGC 6752 (Yong et al. 2006b), red squares represent $\omega$ Cen (Smith et al. 2000), blue asterisks are NGC 3201 (Gonzalez \& Wallerstein 1998), and magenta plus signs and upper limits represent field stars from Tomkin \& Lambert (1999) and Gratton \& Sneden (1994). A representative error bar is shown. The abundances have been shifted onto the Smith et al. (2000) scale.


Fig. 4. $-[\mathrm{Rb} / \mathrm{Y}]$ vs. $[\mathrm{Y} / \mathrm{Fe}]$ (upper), $[\mathrm{Rb} / \mathrm{La}]$ vs. $[\mathrm{La} / \mathrm{Fe}]$ (middle), and $[\mathrm{Rb} / \mathrm{Pb}]$ vs. $[\mathrm{Pb} / \mathrm{Fe}]$ (lower) for M4, M5, and NGC 6752. In each panel, the dotted line represents the mean value.


Fig. 5.- $[\mathrm{Rb} / \mathrm{Zr}]$ vs. $[\mathrm{Fe} / \mathrm{H}]$. The symbols are the same as in Figure 3. The abundances have been shifted onto the Smith et al. (2000) scale. A representative error bar is shown.


Fig. 6. $-[\mathrm{Pb} / \mathrm{Fe}]$ vs. $[\mathrm{Fe} / \mathrm{H}]$ for M 4 (filled circles) and M 5 (open circles). Green triangles show NGC 6752, red squares represent M13, and blue asterisks are HD 141531 (Yong et al. 2006 b ). The magenta plus signs and upper limits represent field stars from Sneden et al. (1998) and Travaglio et al. (2001). A representative error bar is shown.

Table 1. Program stars, exposure times and stellar parameters.

| Star | Exposure Time (min) | $\begin{array}{r} \mathrm{S} / \mathrm{N}^{\mathrm{a}} \\ 4050 \AA \end{array}$ | $\begin{array}{r} \mathrm{S} / \mathrm{N}^{\mathrm{a}} \\ 7800 \AA \end{array}$ | $\begin{gathered} T_{\mathrm{eff}} \\ \mathrm{~K} \end{gathered}$ | $\log g$ | $\begin{gathered} \xi_{t} \\ \mathrm{~km} \mathrm{~s}^{-1} \end{gathered}$ | [Fe/H] |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| M4 L1411 | 40 | 61 | 397 | 4025 | 0.80 | 1.75 | $-1.23$ |
| M4 L1501 | 151 | 72 | 353 | 4175 | 1.00 | 1.55 | -1.29 |
| M4 L1514 | 35 | 59 | 432 | 3950 | 0.30 | 1.85 | -1.22 |
| M4 L2307 | 80 | 82 | 453 | 4125 | 0.95 | 1.65 | -1.19 |
| M4 L2406 | 30 | 72 | 395 | 4150 | 0.15 | 2.20 | -1.30 |
| M4 L2617 | 140 | 71 | 377 | 4275 | 1.25 | 1.65 | -1.20 |
| M4 L3209 | 35 | 68 | 407 | 4075 | 0.75 | 1.95 | -1.25 |
| M4 L3413 | 135 | 88 | 440 | 4225 | 1.10 | 1.75 | -1.23 |
| M4 L3624 | 80 | 39 | 285 | 4225 | 1.05 | 1.60 | -1.29 |
| M4 L4511 | 80 | 75 | 376 | 4150 | 1.05 | 1.70 | -1.22 |
| M4 L4611 | 60 | 29 | 393 | 3925 | 0.15 | 1.45 | -1.09 |
| M4 L4613 | 45 | 42 | 420 | 3900 | 0.20 | 1.70 | $-1.25$ |
| M5 IV-81 | 77 | 25 | 210 | 4050 | 0.30 | 1.90 | -1.28 |
| M5 IV-82 | 190 | 79 | 218 | 4400 | 1.20 | 1.75 | -1.33 |

${ }^{\mathrm{a}} \mathrm{S} / \mathrm{N}$ values are per pixel.

Table 2. Elemental abundances.

| Star | [ $\mathrm{O} / \mathrm{Fe}$ ] | [ $\mathrm{Na} / \mathrm{Fe}$ ] | [ $\mathrm{Mg} / \mathrm{Fe}$ ] | [ $\mathrm{Al} / \mathrm{Fe}$ ] | [ $\mathrm{Rb} / \mathrm{Fe}]$ | [Y/Fe] | $[\mathrm{Zr} / \mathrm{Fe}]^{\mathrm{a}}$ | $[\mathrm{Zr} / \mathrm{Fe}]^{\mathrm{b}}$ | [La/Fe] | [ $\mathrm{Eu} / \mathrm{Fe}$ ] | [ $\mathrm{Pb} / \mathrm{Fe}$ ] |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| M4 L1411 | 0.46 | 0.55 | 0.55 | 0.80 | 0.45 | 0.74 | 0.62 | 0.32 | 0.58 | 0.49 | 0.33 |
| M4 L1501 | 0.57 | 0.51 | 0.59 | 0.79 | 0.37 | 0.69 | 0.43 | 0.13 | 0.48 | 0.47 | 0.29 |
| M4 L1514 | 0.69 | 0.10 | 0.50 | 0.63 | 0.32 | 0.71 | 0.62 | 0.32 | 0.28 | 0.25 | 0.17 |
| M4 L2307 | 0.47 | 0.50 | 0.54 | 0.76 | 0.44 | 0.79 | 0.53 | 0.23 | 0.49 | 0.46 | 0.34 |
| M4 L2406 | 0.51 | 0.46 | 0.58 | 0.66 | 0.25 | . . | . . | ... | 0.28 | 0.18 | 0.20 |
| M4 L2617 | 0.44 | 0.48 | 0.52 | 0.76 | 0.39 | 0.63 | 0.51 | 0.21 | 0.51 | 0.44 | 0.25 |
| M4 L3209 | 0.67 | 0.42 | 0.60 | 0.75 | 0.49 | 0.67 | 0.65 | 0.35 | 0.53 | 0.44 | 0.35 |
| M4 L3413 | 0.76 | 0.06 | 0.57 | 0.67 | 0.38 | 0.62 | 0.51 | 0.21 | 0.49 | 0.44 | 0.33 |
| M4 L3624 | 0.78 | 0.18 | 0.64 | 0.66 | 0.40 | 0.68 | 0.49 | 0.19 | 0.49 | 0.41 | 0.29 |
| M4 L4511 | 0.49 | 0.50 | 0.60 | 0.76 | 0.37 | 0.70 | 0.42 | 0.12 | 0.51 | 0.48 | 0.27 |
| M4 L4611 | 0.25 | 0.67 | 0.47 | 0.75 | 0.39 | . . | . . | ... | 0.51 | 0.29 | 0.34 |
| M4 L4613 | 0.59 | 0.67 | 0.62 | 0.86 | 0.47 | . | . | $\ldots$ | 0.56 | 0.39 | 0.45 |
| M5 IV-81 | 0.50 | 0.15 | 0.35 | 0.33 | -0.04 | 0.22 | 0.25 | -0.05 | 0.17 | 0.46 | -0.32 |
| M5 IV-82 | 0.69 | -0.22 | 0.40 | 0.06 | 0.05 | 0.12 | 0.19 | -0.11 | 0.22 | 0.59 | -0.37 |

${ }^{\text {a }} \mathrm{Zr}$ abundances when using the $g f$ values and solar abundance adopted in Yong et al. (2005).
${ }^{\mathrm{b}} \mathrm{Zr}$ abundances when shifted onto the Smith et al. (2000) scale - see Yong et al. (2006b) for details.

Table 3. Line lists and equivalent widths for program stars.

| $\AA$ | Species | $\chi(\mathrm{eV})$ | $\log g f$ | L1411 | L1501 | L1514 | L2307 | L2406 | L2617 | L3209 | L3413 | L3624 | L4511 | L4611 | L4613 | IV-81 | IV-82 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 6300.30 | 8.0 | 0.00 | -9.75 | Abundances derived using synthetic spectra |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 6363.78 | 8.0 | 0.02 | -10.25 | Abundances derived using synthetic spectra |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 4982.83 | 11.0 | 2.10 | -0.91 | $\ldots$ | 93.4 | . | 122.5 | ... | 92.2 | ... |  | 69.5 | 104.7 | $\ldots$ | $\ldots$ | 74.4 | 29.7 |
| 5682.65 | 11.0 | 2.10 | -0.71 |  | 117.9 | 119.4 | ... |  | 121.1 | ... | 96.1 | 96.8 | ... | ... |  | 102.8 | 51.0 |
| 5688.22 | 11.0 | 2.10 | -0.40 |  |  |  |  |  |  |  | 114.7 | 114.7 |  |  |  |  | 69.6 |
| 6154.23 | 11.0 | 2.10 | -1.56 | 69.7 | 52.4 | 43.4 | 60.6 | 44.2 | 50.2 | 53.1 | 26.4 | 30.2 | 57.0 | 88.0 | 79.7 | 36.6 | 8.9 |
| 6160.75 | 11.0 | 2.10 | -1.26 | 95.2 | 75.0 | 66.5 | 85.5 | 69.7 | 75.0 | 79.0 | 47.0 | 49.6 | 82.4 | 112.8 | 104.4 | 57.4 | 19.7 |
| 5711.09 | 12.0 | 4.35 | -1.73 |  | ... | ... | $\ldots$ |  | , |  |  |  | . | $\ldots$ | ... |  | 101.8 |
| 6318.71 | 12.0 | 5.11 | -1.97 | Abundances derived using synthetic spectra |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 6319.24 | 12.0 | 5.11 | -2.20 | Abundances derived using synthetic spectra |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 5557.07 | 13.0 | 3.14 | -1.95 | 32.9 | 21.5 | 69.3 | 25.8 | 15.7 | 21.1 | 28.0 | 19.2 | 17.8 | 23.7 | 104.6 | 105.4 | 15.2 | 7.5 |
| 6696.02 | 13.0 | 3.14 | -1.57 | Abundances derived using synthetic spectra |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 6698.67 | 13.0 | 3.14 | -1.89 | Abundances derived using synthetic spectra |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 4602.00 | 26.0 | 1.61 | -3.15 |  | 121.3 |  |  |  |  |  |  |  |  |  |  |  | 109.4 |
| 4802.88 | 26.0 | 3.69 | -1.53 | 72.3 | 65.5 | 74.4 | 73.2 | 75.1 | 70.7 | 63.9 | 71.2 | $\ldots$ | $\ldots$ | 69.1 | 68.9 | 73.6 | 59.6 |
| 4930.31 | 26.0 | 3.96 | -1.26 | 79.9 | 69.4 | . | 78.0 | 75.3 | 72.2 | 86.2 | 73.1 | $\ldots$ | $\ldots$ | $\ldots$ | ... | 78.6 | 66.4 |
| 5054.64 | 26.0 | 3.64 | -1.94 | ... | 49.1 | $\ldots$ | ... | 53.8 | $\ldots$ | $\ldots$ | $\ldots$ | 52.5 | $\ldots$ | $\ldots$ | $\ldots$ | 62.5 | 37.3 |
| 5223.19 | 26.0 | 3.63 | $-1.80$ | 45.7 | 37.3 | $\ldots$ | 42.5 | \% | 36.5 | 49.5 | 39.4 | $\ldots$ | $\cdots$ | $\ldots$ | $\ldots$ | 42.8 |  |
| 5242.49 | 26.0 | 3.63 | -0.98 | 107.6 | 99.2 | 111.4 | . | $\ldots$ | . | 107.2 | 103.7 | 84.6 | 104.7 | 108.5 | 105.7 | 110.0 | 90.8 |
| 5267.28 | 26.0 | 4.37 | -1.66 | 33.5 | . | . | $\ldots$ | $\ldots$ | $\ldots$ | ... | . | . | . | . | . | . |  |
| 5288.53 | 26.0 | 3.69 | $-1.53$ | ... | $\cdots$ | $\ldots$ | $\ldots$ | 87.1 | $\ldots$ | ... | $\cdots$ | $\ldots$ | . | $\ldots$ | ... | . |  |
| $5321.11$ | 26.0 | 4.43 | -1.11 | 55.7 | $\ldots$ | $\ldots$ | 48.4 | 42.6 | 38.9 | $\ldots$ | $\ldots$ | 39.8 | $\cdots$ | $\ldots$ | . | $\cdots$ | 27.8 |
| $5326.14$ | 26.0 | 3.57 | -2.13 | 54.8 | 48.4 | 61.9 | $\ldots$ | 53.8 | 49.3 | $\ldots$ | $\ldots$ | 50.4 | $\ldots$ | $\ldots$ | 54.3 | $\ldots$ | . |
| 5364.86 | 26.0 | 4.44 | 0.21 | 114.5 | ... | ... | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\cdots$ | 105.3 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ |
| 5365.40 | 26.0 | 3.57 | -1.04 | 102.1 | 89.5 | 108.7 | 98.7 | $\ldots$ | . | $\ldots$ | $\ldots$ | 90.3 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ |
| 5367.48 | 26.0 | 4.41 | 0.43 | 116.2 | ... | $\ldots$ | $\ldots$ | ... | .. | $\ldots$ | $\cdots$ | 110.1 | 107.7 | 119.8 | 110.6 | $\ldots$ |  |
| 5379.57 | 26.0 | 3.69 | $-1.53$ | 81.6 | 68.7 | 87.3 | 77.7 | 83.7 | 75.3 | 82.2 | 75.5 | 71.6 | 76.8 | 83.2 | 83.2 | 84.2 | 63.0 |
| 5412.80 | 26.0 | 4.43 | $-1.78$ | . | ... | 87. | , | ... | . | , | . |  | ... |  | . | 27.6 | . |
| 5491.84 | 26.0 | 4.18 | -2.25 | $\ldots$ | 13.4 | $\ldots$ | 18.6 | 15.9 | 13.8 | $\ldots$ | $\ldots$ | 10.8 | 14.4 | 28.9 | $\ldots$ | ... |  |
| 5525.54 | 26.0 | 4.23 | -1.15 | $\ldots$ | 54.7 | 66.8 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 58.6 | 55.2 | 58.8 | ... | $\ldots$ | $\ldots$ | 41.1 |
| 5618.63 | 26.0 | 4.21 | -1.29 | $\ldots$ | 48.2 | $\ldots$ | $\ldots$ | 56.7 | 48.0 | $\ldots$ | $\ldots$ | ... | 54.7 | $\ldots$ | $\ldots$ | 55.2 | 40.4 |
| 5661.35 | 26.0 | 4.28 | $-1.82$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 24.9 | $\cdots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 15.0 |
| 5705.47 | 26.0 | 4.30 | -1.42 | 43.2 | 35.6 | 44.9 | 41.4 | 38.6 | 38.9 | 43.5 | 38.7 | $\ldots$ | 39.4 | 44.3 | 42.6 | 37.8 | 27.9 |
| 5775.08 | 26.0 | 4.22 | -1.31 | 67.0 | 58.1 | 68.8 | 65.5 | 69.5 | 62.5 | 69.0 | 65.8 | $\cdots$ | 63.8 | 62.2 | 65.4 | $\cdots$ | 52.6 |
| 5778.45 | 26.0 | 2.59 | -3.48 | ... | ... | $\ldots$ | ... | ... | .. | $\ldots$ | . | $\ldots$ | $\ldots$ | $\ldots$ | ... | 67.8 | 38.4 |
| 5855.09 | 26.0 | 4.60 | -1.55 | 17.1 | 13.3 | . | 18.3 | 13.8 | 14.7 | 20.1 | 16.7 | 14.9 | ... | 25.4 | 24.0 | 18.5 |  |
| 5909.97 | 26.0 | 3.21 | -2.64 | 79.7 | 62.5 | 85.1 | 73.7 | 71.6 | 63.9 | 80.3 | 68.1 | 60.6 | 69.3 | . | 85.0 | 77.7 | 43.9 |
| 5916.25 | 26.0 | 2.45 | -2.99 | ... | ... | ... | $\ldots$ | $\ldots$ | . | ... | $\cdots$ | 94.3 | ... | . | ... | $\ldots$ | 80.8 |
| 6012.21 | 26.0 | 2.22 | -4.07 | 65.7 | 52.6 | 73.2 | 60.6 | 59.4 | 50.6 | 66.7 | 54.7 | 50.8 | 56.5 | 70.7 | 68.2 | 65.4 | 34.4 |
| 6027.05 | 26.0 | 4.07 | -1.11 | 77.5 | . | 84.0 | 78.6 | 84.7 | 72.7 | 79.4 | 74.6 | 73.0 | 76.5 | 81.7 | 76.9 | 78.8 | 60.9 |
| 6120.24 | 26.0 | 0.91 | $-5.97$ | 66.9 | 43.3 | 82.9 | 56.6 | 55.5 | $\ldots$ | 66.7 | 46.0 | 41.0 | 48.5 | 87.9 | 84.1 | 72.1 | 25.2 |
| 6151.62 | 26.0 | 2.17 | -3.30 | 116.7 | 95.8 |  |  | ... |  | ... | 97.4 | 94.7 |  | . | . | . | 83.0 |
| 6165.36 | 26.0 | 4.14 | -1.49 | 52.3 | 44.3 | $\ldots$ | 50.4 | 49.1 | 48.5 | 53.7 | 46.4 | 43.5 | 50.4 | $\ldots$ | $\ldots$ | $\ldots$ | 33.4 |
| 6180.20 | 26.0 | 2.73 | -2.64 | . | 89.5 | . | . | . | 90.2 | . | . | 88.0 | . | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ |
| 6200.31 | 26.0 | 2.61 | -2.44 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 113.5 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ |  |
| 6229.23 | 26.0 | 2.84 | -2.85 | 79.2 | 65.6 | $\ldots$ | 73.8 | 77.4 | 62.2 | $\ldots$ | 68.0 | 63.3 | 70.4 | $\ldots$ | $\ldots$ | 81.8 | 51.6 |
| 6232.64 | 26.0 | 3.65 | -1.28 | 111.3 | ... | ... | ... | ... | $\ldots$ | 109.5 | 100.4 | 97.6 | 104.2 | $\ldots$ | $\ldots$ | 109.3 | 83.8 |
| 6270.22 | 26.0 | 2.86 | -2.51 | 96.0 | 85.0 | 98.0 | 91.7 | 103.5 | 87.1 | 94.2 | 87.5 | 83.8 | 90.4 | 97.0 | 93.6 | 100.1 | 70.1 |
| 6271.28 | 26.0 | 3.33 | $-2.76$ | 46.2 | ... | 51.7 | ... | 39.5 | 36.9 | 47.4 | 39.0 | 36.0 | 39.2 | 53.4 | 50.5 | 43.5 | 25.4 |
| 6336.82 | 26.0 | 3.68 | -0.92 | . | $\ldots$ | $\ldots$ | $\ldots$ |  | $\ldots$ | ... | $\ldots$ | 111.9 | $\ldots$ | . | $\ldots$ |  |  |
| 6353.84 | 26.0 | 0.91 | $-6.48$ | 34.1 | $\ldots$ | 43.5 | $\ldots$ | 28.2 | $\ldots$ | 34.4 | 22.1 | 18.9 | $\ldots$ | 49.8 | 45.2 | 35.7 | $\ldots$ |

Table 3-Continued

| A | Species | $\chi(\mathrm{eV})$ | $\log g f$ | L1411 | L1501 | L1514 | L2307 | L2406 | L2617 | L3209 | L3413 | L3624 | L4511 | L4611 | L4613 | IV-81 | IV-82 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 6355.03 | 26.0 | 2.84 | $-2.40$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 111.2 | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ | 89.5 |
| 6408.02 | 26.0 | 3.68 | -1.07 | ... | $\ldots$ | ... | ... | $\ldots$ | $\ldots$ | $\ldots$ | ... | 102.6 | $\ldots$ | ... | $\ldots$ |  | 91.2 |
| 6518.36 | 26.0 | 2.83 | $-2.50$ | $\ldots$ | 88.3 | $\ldots$ |  | . | $\ldots$ | $\ldots$ |  | 87.7 |  | ... | ... | 107.8 |  |
| 6574.23 | 26.0 | 0.99 | -5.00 | $\ldots$ | 103.3 | $\ldots$ | 119.5 | $\ldots$ | $\ldots$ | $\ldots$ | 109.7 | 100.4 | 110.3 | ... | $\ldots$ | ... | 82.6 |
| 6575.02 | 26.0 | 2.59 | -2.73 | $\ldots$ | ... | $\ldots$ | ... | $\ldots$ | $\ldots$ | ... | 110.0 | 103.0 | 113.5 | $\ldots$ | $\ldots$ | ... | 87.1 |
| 6581.21 | 26.0 | 1.48 | -4.71 | ... | $\ldots$ | $\ldots$ | $\ldots$ | 99.2 | $\ldots$ | $\ldots$ | ... |  | . | $\ldots$ | $\ldots$ |  | 40.0 |
| 6609.11 | 26.0 | 2.56 | -2.69 |  |  |  |  |  |  | $\ldots$ |  | 103.4 | ... | $\ldots$ |  |  | 94.2 |
| 6625.02 | 26.0 | 1.01 | -5.37 |  | $\ldots$ | $\ldots$ | ... | 114.5 |  | $\ldots$ | $\ldots$ | ... |  | $\cdots$ |  |  | 61.1 |
| 6648.08 | 26.0 | 1.01 | -5.92 | 70.7 | 54.8 | 79.9 | 65.5 | 66.4 | 52.3 | 71.4 | 47.9 | 42.7 | 60.1 | 86.6 | 81.3 | 70.8 | 25.5 |
| 6699.16 | 26.0 | 4.59 | -2.17 |  | 5.1 | 6.4 | . | . | 5.6 | 5.1 | ... | 4.7 | 4.9 | $\ldots$ | 6.0 | 5.5 | . |
| 6739.52 | 26.0 | 1.56 | -4.82 | 72.8 | 53.2 | . | $\ldots$ | $\ldots$ | 51.5 | 74.2 | 57.1 | 50.7 | 58.0 | $\ldots$ |  | 77.2 | $\cdots$ |
| 6793.25 | 26.0 | 4.07 | -2.39 |  | . | . |  |  |  |  | $\ldots$ | $\ldots$ | $\ldots$ | $\ldots$ |  | 17.0 |  |
| 6810.26 | 26.0 | 4.60 | -1.00 | 46.5 | 42.5 | 54.2 | 47.3 | 42.9 | 41.4 | 49.5 | 42.3 | 37.6 | 43.2 | 56.8 | 56.8 | 41.5 | 26.4 |
| 6837.02 | 26.0 | 4.59 | $-1.76$ |  | . | 16.3 | ... | ... | $\ldots$ | ... | 15.8 |  | ... | ... | ... | 15.6 | 8.9 |
| 6971.93 | 26.0 | 3.02 | -3.39 | 33.7 |  | 39.6 | $\ldots$ |  |  | 35.1 | 26.6 | 25.5 | $\ldots$ | ... | ... | 32.0 | ... |
| 7112.17 | 26.0 | 2.99 | -3.04 | 62.6 | 54.2 | . | 61.9 | 61.7 | 55.1 | 63.0 | 52.0 | 50.5 | $\ldots$ | $\ldots$ | $\ldots$ | 62.7 | 34.9 |
| 7223.66 | 26.0 | 3.02 | $-2.27$ | 110.8 | 96.1 | $\ldots$ |  | $\ldots$ | ... | 109.4 | 96.4 | 93.7 | $\ldots$ | $\ldots$ | $\ldots$ | 110.8 | 79.9 |
| 7401.69 | 26.0 | 4.18 | -1.66 | 45.0 | $\ldots$ | 46.3 | 44.8 | $\ldots$ | ... | 45.3 | $\ldots$ | $\ldots$ | $\ldots$ | .. |  | 45.4 | . |
| 7710.36 | 26.0 | 4.22 | -1.13 | 77.5 | 62.9 | 82.0 | 75.5 | 78.4 | 65.7 | 77.7 | 66.6 | 59.6 | 70.9 | 77.2 |  |  | ... |
| 7723.20 | 26.0 | 2.28 | $-3.62$ | 97.6 | 79.4 | ... | 90.0 | 103.4 | 78.8 | 97.6 | 85.8 | 81.3 | 88.1 | 110.6 | 111.2 | 111.0 | 69.2 |
| 7912.86 | 26.0 | 0.86 | -4.85 |  |  | $\ldots$ |  | ... |  |  |  |  |  |  |  |  | 111.4 |
| 7941.09 | 26.0 | 3.27 | $-2.33$ | 67.7 | 57.4 | 77.0 | 63.5 | $\ldots$ | 58.1 | 67.7 | 61.3 | 56.2 | 61.7 | 71.7 | 69.4 | 71.8 | 47.0 |
| 8075.15 | 26.0 | 0.91 | -5.09 | . | 118.6 | $\ldots$ | $\ldots$ | $\ldots$ | ... | $\ldots$ | ... | 114.7 | $\ldots$ | . |  | .. | 91.2 |
| 8204.10 | 26.0 | 0.91 | -6.05 | 77.4 | 46.4 | 94.8 | 62.4 | . | 43.8 | 77.4 | 50.4 | 41.4 | 55.4 | 94.5 | 93.1 | 77.2 | 28.4 |
| 4620.52 | 26.1 | 2.83 | -3.08 | 40.0 | 45.1 | 46.2 | 44.4 | 60.6 | ... | $\ldots$ | ... | ... | 41.4 | $\ldots$ | ... | ... | 54.0 |
| 4993.36 | 26.1 | 2.81 | -3.49 | 44.6 | 36.6 | 48.8 | 40.4 | ... | $\ldots$ | ... | ... | ... | $\ldots$ | 53.5 |  | 55.7 | 34.0 |
| 5234.62 | 26.1 | 3.22 | -2.15 | ... | 74.5 | 78.8 | 74.8 |  | 79.9 | $\ldots$ | $\ldots$ |  | 72.8 | ... | 66.6 | ... | 84.4 |
| 5325.55 | 26.1 | 3.22 | -3.22 | ... | ... | 34.2 | 37.0 | 60.5 | 39.9 | $\ldots$ | ... | 33.9 | 35.0 | $\ldots$ | 26.4 | ... | . |
| 5414.07 | 26.1 | 3.22 | $-3.75$ | $\ldots$ | . | $\ldots$ | $\ldots$ | 29.3 | ... | $\ldots$ | $\ldots$ | 18.0 | 17.1 | $\ldots$ |  | 20.0 | $\ldots$ |
| 5425.26 | 26.1 | 3.20 | $-3.37$ | 29.6 | 31.5 | 32.3 | 28.8 | 47.9 | 31.0 | 30.1 | 31.8 | 24.5 | 30.1 | 29.2 | 27.1 | 34.7 | 32.5 |
| 5991.38 | 26.1 | 3.15 | -3.56 | ... | ... | $\ldots$ | $\ldots$ | ... | $\cdots$ | $\ldots$ | $\ldots$ | ... | $\ldots$ | ... | . | $\cdots$ | 25.1 |
| 6084.11 | 26.1 | 3.20 | -3.81 | 12.8 | . | 15.8 | 14.3 | 24.4 | 15.5 | 12.0 | 13.8 | 12.8 | 13.8 | $\ldots$ | 9.9 | 20.3 | 17.9 |
| 6149.26 | 26.1 | 3.89 | -2.72 |  | 23.9 |  | 24.1 | 37.7 | 23.9 |  |  | 16.9 |  |  |  | 25.0 | 25.7 |
| 6247.56 | 26.1 | 3.89 | -2.33 | 29.0 | 33.5 | 33.4 | 31.8 | 56.7 | 35.2 | 29.0 | 34.7 | 33.4 | 33.3 | 35.0 | 31.6 | 36.5 | 41.8 |
| 6369.46 | 26.1 | 2.89 | -4.25 | 9.8 | 12.6 | 17.3 | 11.8 | 19.4 | 13.0 | 12.6 | 16.0 | 14.5 | 13.2 | 19.4 |  | 17.1 | 17.7 |
| 6416.92 | 26.1 | 3.89 | -2.74 | 23.5 | 26.2 | 29.6 | 26.5 | 37.0 | 24.2 | 24.6 | 25.7 | 25.6 | 26.3 | 31.2 | 29.2 | 29.6 | 27.3 |
| 6432.68 | 26.1 | 2.89 | -3.71 | 29.8 | 33.9 | 35.0 | 33.2 | 50.8 | 33.4 | 33.4 | 34.6 | 33.7 | 34.4 | 31.7 | 29.6 | 39.6 | 39.5 |
| 6456.38 | 26.1 | 3.90 | -2.08 | 36.6 | 41.9 | 41.7 | 40.6 | 74.4 | 43.1 | 36.4 | 42.3 | 41.9 | 42.0 | ... | 35.5 | 46.5 | 51.4 |
| 6516.08 | 26.1 | 2.89 | -3.45 | ... | ... | 50.2 | 49.3 | 68.3 | ... | 47.8 | $\ldots$ | . ${ }^{\text {a }}$ | ... | 55.6 | 49.7 | 57.2 | 51.4 |
| 7224.49 | 26.1 | 3.89 | -3.24 | $\ldots$ | $\ldots$ | ... |  | ... | $\ldots$ | $\ldots$ | 9.2 | 9.0 | $\ldots$ | ... | ... | 15.1 | 9.8 |
| 7479.69 | 26.1 | 3.89 | $-3.59$ | $\ldots$ | $\ldots$ | 5.8 | 8.7 | $\ldots$ | $\ldots$ | 7.4 | 6.0 | 6.6 | $\ldots$ | $\ldots$ |  | 7.9 |  |
| 7515.83 | 26.1 | 3.90 | -3.43 |  |  |  |  |  |  |  |  |  |  | $\ldots$ |  |  | 5.5 |
| 7711.72 | 26.1 | 3.90 | -2.54 | 22.3 | 26.3 | 25.5 | 24.7 | 42.1 | 29.1 | 23.3 | 26.2 | 28.2 | 28.1 | 25.6 | 25.5 | 29.0 | 31.6 |
| 5200.42 | 39.1 | 0.99 | $-0.57$ | $\ldots$ | 105.0 | $\ldots$ | 121.5 | $\ldots$ | 105.0 | $\ldots$ | 111.9 | 102.1 | 116.0 |  |  | 112.8 | 72.5 |
| 5509.91 | 39.1 | 0.99 | -1.01 | 110.0 | 97.0 | 116.0 | 109.0 | 114.0 | 96.0 | 113.0 | 100.0 | 99.0 | 101.0 | 120.0 | 114.0 | 90.0 | 68.0 |
| 5544.61 | 39.1 | 1.74 | -1.08 | 39.0 | 28.0 | 52.0 | 32.0 | 31.0 | 29.0 | 38.0 | 28.0 | 31.0 | 32.0 | 72.0 | 71.0 | 20.0 | 12.0 |
| 6127.44 | 40.0 | 0.15 | -1.06 | 107.0 | 58.0 | 126.0 | 80.0 | 47.0 | 55.0 | 105.0 | 63.0 | 56.0 | 68.0 | 145.0 | 143.0 | 76.0 | 15.0 |
| 6134.55 | 40.0 | 0.00 | -1.28 | 106.0 | 53.0 | 126.0 | 78.0 | 44.0 | 49.0 | 104.0 | 58.0 | 49.0 | 63.0 | 146.0 | 143.0 | 76.0 | 14.0 |
| 6143.20 | 40.0 | 0.07 | -1.10 | 109.0 | 57.0 | 128.0 | 84.0 | 49.0 | 55.0 | 108.0 | 63.0 | 55.0 | 69.0 | 149.0 | 147.0 | 78.0 | 14.0 |
| 5303.53 | 57.1 | 0.32 | -1.35 |  |  |  |  |  | andances | derived | ing synt | etic spec |  |  |  |  |  |
| 6390.49 | 57.1 | 0.30 | $-1.41$ |  |  |  |  |  | undances | derived | ing synt | tic spec |  |  |  |  |  |

## Table 3-Continued

| $\AA$ | Species | $\chi(\mathrm{eV})$ | $\log g f$ | L1411 | L1501 | L1514 | L2307 | L2406 | L2617 | L3209 | L3413 | L3624 | L4511 | L4611 | L4613 | IV-81 | IV-82 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 6645.13 | 63.1 | 1.37 | 0.20 |  |  |  |  |  | ndances | derived | ing synt | tic spec |  |  |  |  |  |

Table 4. Abundance dependences on model parameters for M4 L2307.

| Species | $T_{\text {eff }}+50$ | $\log g+0.2$ | $\xi_{t}+0.2$ |
| :--- | ---: | ---: | ---: |
| $[\mathrm{Fe} / \mathrm{H}]$ | -0.01 | 0.04 | -0.03 |
| $[\mathrm{O} / \mathrm{Fe}]$ | 0.01 | 0.02 | 0.02 |
| $[\mathrm{Na} / \mathrm{Fe}]$ | 0.06 | -0.07 | 0.01 |
| $[\mathrm{Mg} / \mathrm{Fe}]$ | 0.04 | -0.05 | 0.02 |
| $[\mathrm{Al} / \mathrm{Fe}]$ | 0.05 | -0.06 | 0.02 |
| $[\mathrm{Rb} / \mathrm{Fe}]$ | 0.08 | -0.05 | 0.03 |
| $[\mathrm{Y} / \mathrm{Fe}]$ | 0.00 | 0.03 | -0.02 |
| $[\mathrm{Zr} / \mathrm{Fe}]$ | 0.13 | -0.03 | 0.02 |
| $[\mathrm{La} / \mathrm{Fe}]$ | 0.02 | 0.02 | 0.01 |
| $[\mathrm{Eu} / \mathrm{Fe}]$ | -0.01 | 0.04 | 0.01 |
| $[\mathrm{~Pb} / \mathrm{Fe}]$ | 0.10 | -0.09 | -0.02 |


[^0]:    ${ }^{1}$ Based on observations made with the Magellan Clay Telescope at Las Campanas Observatory.

[^1]:    ${ }^{2}$ IRAF (Image Reduction and Analysis Facility) is distributed by the National Optical Astronomy Observatory, which is operated by the Association of Universities for Research in Astronomy, Inc., under contract with the National Science Foundation.
    ${ }^{3}$ http://www.lco.cl/lco/magellan/instruments/MIKE/reductions/mtools.html

