



# CKM Fits : What the Data Say (Focused on B Physics)

S. T'Jampens

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# Beauty 2006

The 11th International Conference on B-physics  
at Hadron Machines  
September 25th-29th 2006 at Keble College,  
University of Oxford, UK

## CKM Fits: What the Data Say (focused on B-Physics)

Stéphane T'JAMPENS

LAPP (CNRS/IN2P3 & Université de Savoie)



# Outline

- ☀ CKM phase invariance and unitarity
- ☀ Statistical issues
- ☀ CKM metrology
  - 🖥 Inputs
    - 🖥 Tree decays:  $|V_{ub}|, |V_{cb}|$
    - 🖥 Loop decays:  $\Delta m_d, \Delta m_s, \epsilon_K$
    - 🖥 UT angles:  $\alpha, \beta, \gamma$
  - 🖥 The global CKM fit
- ☀ What about New Physics?
- ☀ Conclusion

■ Charm is interesting in several special areas, but I will concentrate on b's

# The Unitary Wolfenstein Parameterization

↪ The standard parameterization uses Euler angles and one CPV phase → unitary !

$$V = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}$$

Chau and Keung  
PRL **53**, 1802 (1984)  
[and PDG]

↪ Now, define

$$s_{12} \equiv \lambda$$

$$s_{23} \equiv A\lambda^2$$

$$s_{13}e^{-i\delta} \equiv A\lambda^3(\rho - i\eta)$$

Buras *et al.*,  
PRD **50**, 3433 (1994)

↪ And insert into  $V \rightarrow V$  is still unitary ! With this one finds (to all orders in  $\lambda$ ) :

$$\rho + i\eta = \frac{\sqrt{1 - A^2\lambda^4}(\bar{\rho} + i\bar{\eta})}{\sqrt{1 - \lambda^2} [1 - A^2\lambda^4(\bar{\rho} + i\bar{\eta})]}$$

where:  $\bar{\rho} + i\bar{\eta} = -\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*}$

Charles *et al.*  
EPJC **41**, 1 (2005)

$$\lambda^2 \equiv \frac{|V_{us}|^2}{|V_{ud}|^2 + |V_{us}|^2} \quad A^2\lambda^4 \equiv \frac{|V_{cb}|^2}{|V_{ud}|^2 + |V_{us}|^2}$$

Physically meaningful quantities are phase-convention invariant

→ Four unknowns [unitary-exact and phase-convention invariant]:

$$A, \lambda, \bar{\rho}, \bar{\eta}$$

# The CKM Matrix: Four Unknowns

## Measurement of Wolfenstein parameters:

- ✱  $\lambda$  from  $|V_{ud}|$  (nuclear transitions) and  $|V_{us}|$  (semileptonic  $K$  decays)  
→ combined precision: 0.5%
- ✱  $A$  from  $|V_{cb}|$  (inclusive and exclusive semileptonic  $B$  decays)  
→ combined precision: 2%
- ✱  $\bar{\rho}, \bar{\eta}$  from (mainly) CKM angle measurements:  
→ combined precision: 20% ( $\rho$ ), 7% ( $\eta$ )

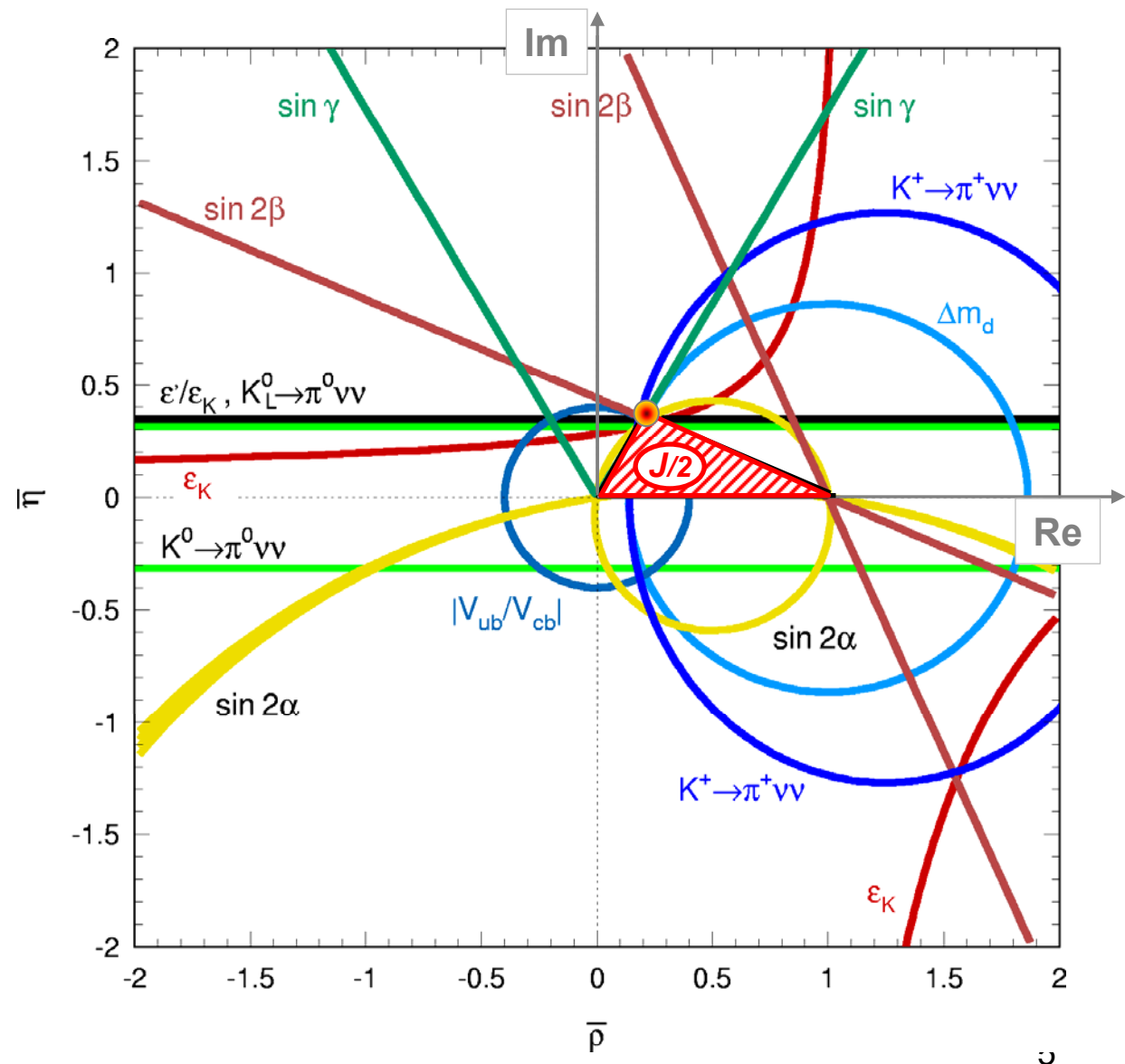
# Predictive Nature of KM Mechanism

All measurements  
must agree

Pre B-Factory:

Can the KM mechanism  
describe flavor dynamics of  
many constraints from vastly  
different scales?

This is what matters and not  
the measurement of the  
CKM phase's value *per se*



# The (rescaled) Unitarity Triangle: The $B_d$ System

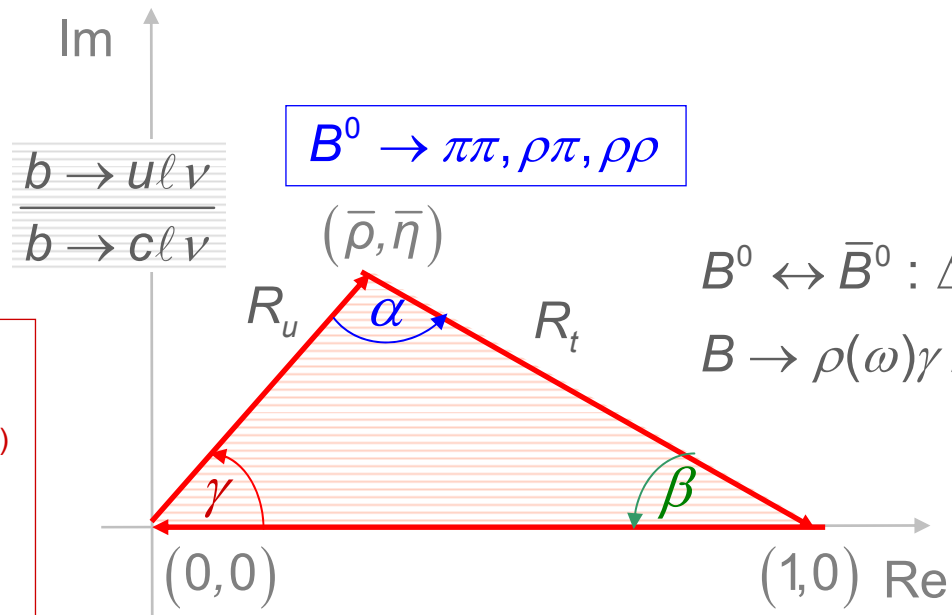
Convenient method to illustrate (dis-)agreement of observables with CKM predictions

$$\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*} + 1 + \frac{V_{td}V_{tb}^*}{V_{cd}V_{cb}^*} = 0$$

⏟  
phase invariant :  $\bar{\rho} + i\bar{\eta}$

“There is no such thing as  $\alpha/\phi_2$ ”  
 $[\alpha = \pi - (\beta + \gamma)]$

- $B \rightarrow D^{(*)}K^{(*)}$
- $B \rightarrow D_{K_S^0 \pi^+ \pi^-} K^{(*)}$
- $B^0 \rightarrow DK_S^0, \dots$
- $B^0 \rightarrow D^* \pi(\rho)$



$B^0 \leftrightarrow \bar{B}^0 : \Delta m_d$   
 $B \rightarrow \rho(\omega)\gamma \mid B \rightarrow K^{(*)}\gamma$

- $b \rightarrow c\bar{c}s$
- $B^0 \rightarrow J/\psi K_S^0, \dots$
- $B^0 \rightarrow \phi K_S^0, \dots$
- $b \rightarrow s\bar{s}s$

# The Unitarity Triangle: The $B_s$ System (hadron machines)

(sb) triangle (“ $B_s$  triangle”):

$$V_{us}V_{ub}^* + V_{cs}V_{cb}^* + V_{ts}V_{tb}^* = 0$$

$$O(\lambda^4) + O(\lambda^2) + O(\lambda^2) = 0$$

→ squashed triangle

$$\chi = \beta_s = \arg\left[-\frac{V_{cs}V_{cb}^*}{V_{ts}V_{tb}^*}\right]$$

Attention: sign

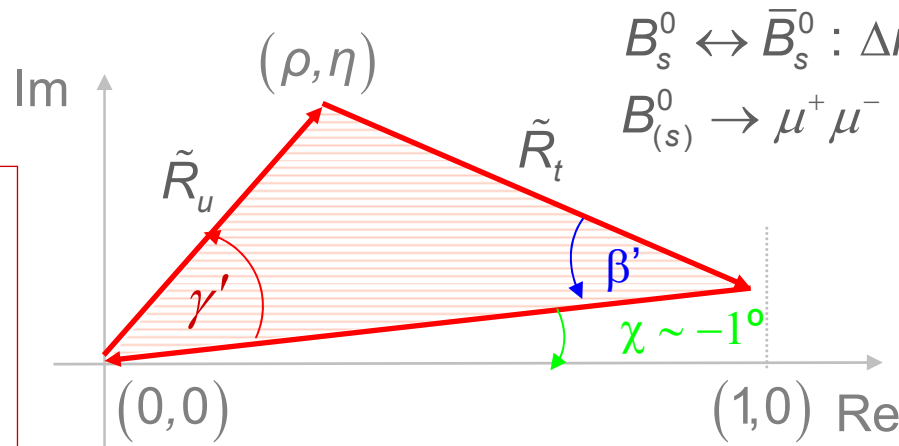
(ut) triangle:

$$V_{td}V_{ud}^* + V_{ts}V_{us}^* + V_{tb}V_{ub}^* = 0$$

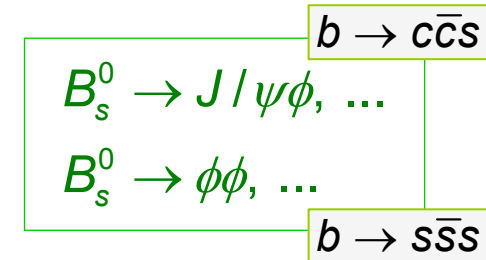
$$O(\lambda^3) + O(\lambda^3) + O(\lambda^3) = 0$$

→ non-squashed triangle

- $B_s^0 \rightarrow D_s K$
- $B_s^0 \rightarrow K^+ K^-$
- $B_s^0 \rightarrow J/\psi K_S^0$
- $B_s^0 \rightarrow D_s^+ D_s^-, \dots$



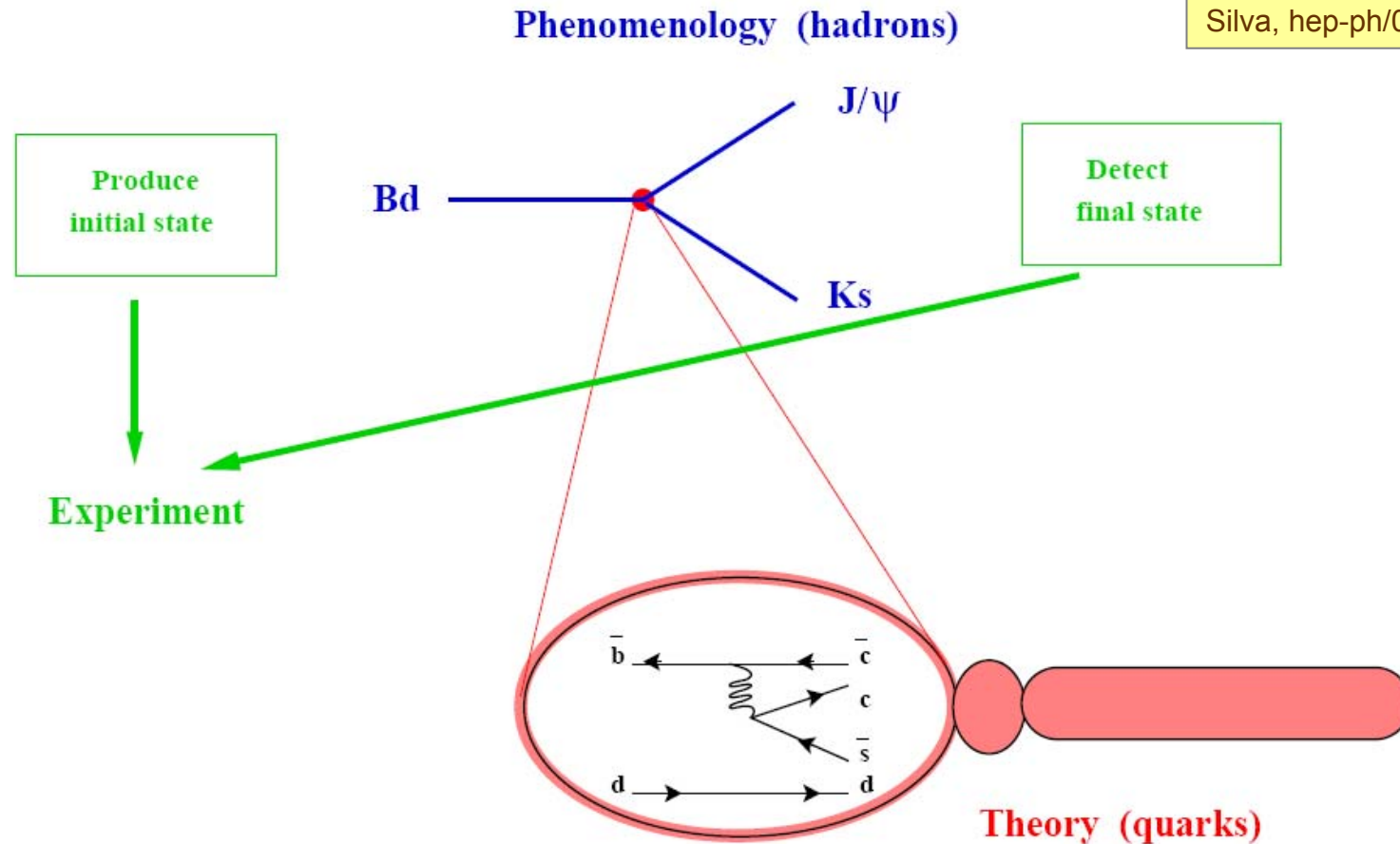
$B_s^0 \leftrightarrow \bar{B}_s^0 : \Delta m_s / \Delta m_d$   
 $B_{(s)}^0 \rightarrow \mu^+ \mu^-$  (BR for  $B_d^0 \sim 1 \times 10^{-10}$ )  
 (BR for  $B_s^0 \sim 4 \times 10^{-9}$ )





# Generic B physics experiment

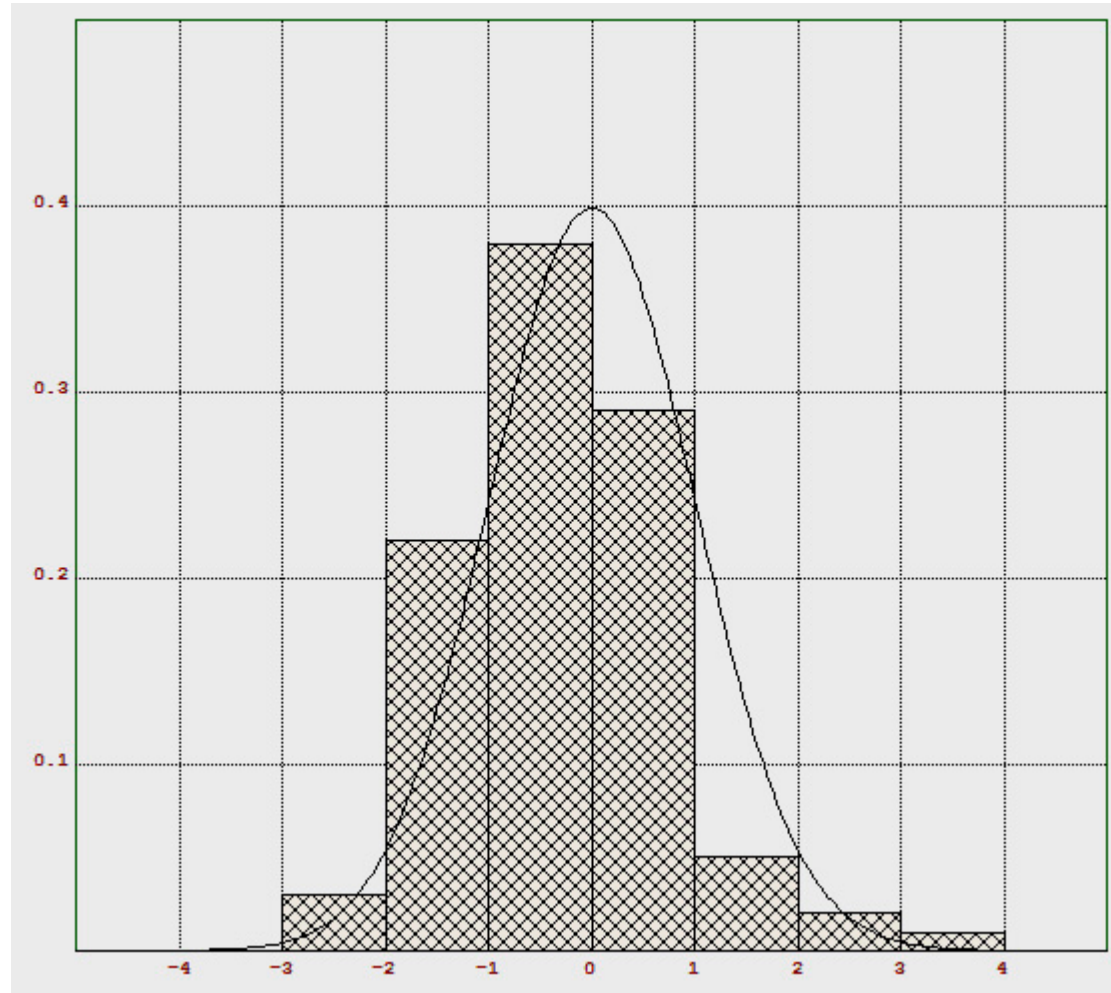
Silva, hep-ph/0410351



Probing short distance (quarks) but confined in hadrons (what we observe)

- QCD effects must be under control (various tools: HQET, SCET, QCDF, LQCD,...)
- “**Theoretical uncertainties**” have to be controlled **quantitatively** in order to test the Standard Model. There is however no systematic method to do that.

# Digression: Statistics



# Digression: Statistics

D.R. Cox, Principles of Statistical Inference, CUP (2006)

W.T. Eadie et al., Statistical Methods in Experimental Physics, NHP (1971)

www.phystat.org

Statistics tries answering a wide variety of questions → two main different! frameworks:

**Frequentist:** probability about the data (randomness of measurements), given the model

$$P(\text{data}|\text{model}) \quad [\text{only repeatable events} \\ (\text{Sampling Theory})]$$

*Hypothesis testing:* given a model, assess the **consistency** of the data with a particular parameter value → 1-CL curve (by varying the parameter value)

**Bayesian:** probability about the model (degree of belief), given the data

$$P(\text{model}|\text{data}) \propto \text{Likelihood}(\text{data}, \text{model}) \times \text{Prior}(\text{model})$$

$P(\text{data}|\text{model}) \neq P(\text{model}|\text{data})$ :  $P(\text{pregnant} | \text{female}) \sim 3\%$

model: Male or Female

data: pregnant or not pregnant

but

$P(\text{female} | \text{pregnant}) \gg \gg 3\%$

Lyons – CDF  
Stat Committee

Although the graphical displays appear similar: the meaning of the “Confidence level” is not the same. It is especially important to understand the difference in a time where one seeks 10 deviation of the SM.

# Digression: Statistics (cont.)

Mayo – Error and the Growth of Experimental Knowledge, UCP(1996)

The Bayesian approach in physical science fails in the sense that nothing guarantees that *my* uncertainty assessment is any good for *you* - I'm just expressing an **opinion (degree of belief)**. To convince you that it's a good uncertainty assessment, I need to show that the **statistical model** I created makes good predictions in situations where we know what the **truth** is, and the process of **calibrating predictions** against reality is inherently frequentist."

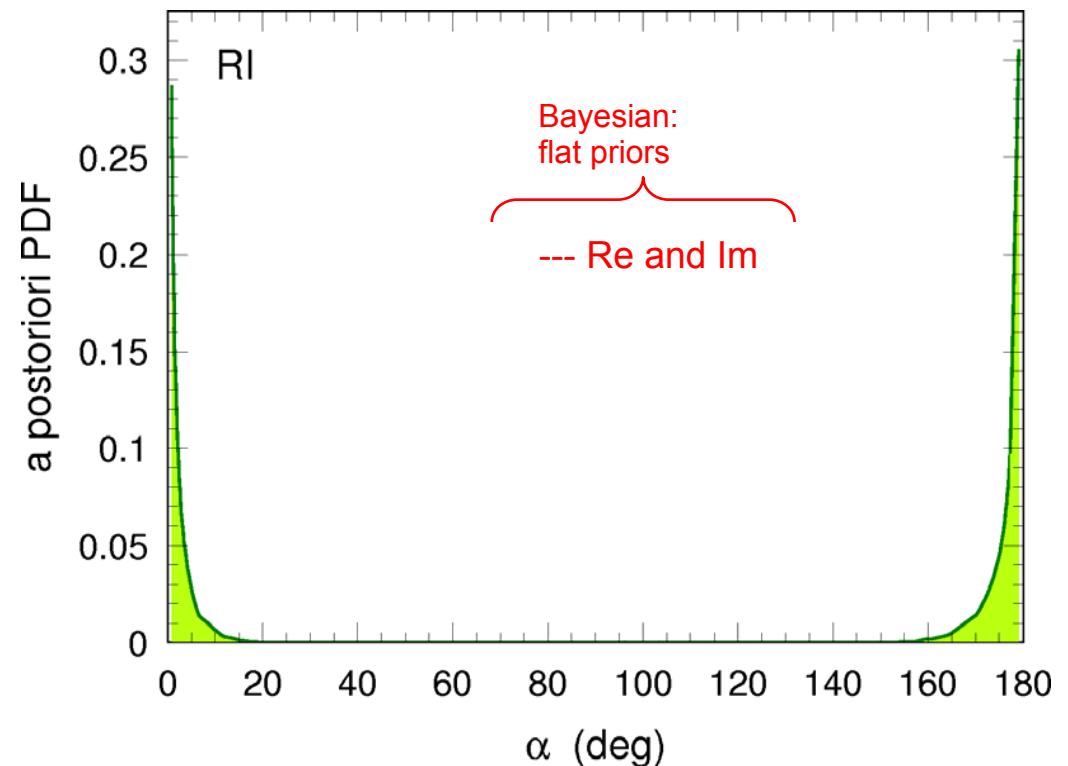
hep-ph/0607246: "Bayesian Statistics at Work: the Troublesome Extraction of the CKM Angle  $\alpha$ " (J. Charles *et al.*)

## How to read a Posterior PDF?

→ updated belief (after seeing the data) of the plausible values of the parameter

↪ it's a bet on a proposition to which there is no scientific answer

$B \rightarrow \rho\rho$  (w/o theoretical errors):



My talk is about “**What the Data say**”, thus I will stick to the frequentist approach



# Metrology: Inputs to the Global CKM Fit

## I) Direct Measurement: magnitude

$|V_{ud}|$  and  $|V_{us}|$  [not discussed here]

$|V_{ub}|$  and  $|V_{cb}|$

$B^+ \rightarrow \tau^+ \nu$

CPV in  $K^0$  mixing [not discussed here]

$B_d$  and  $B_s$  mixing

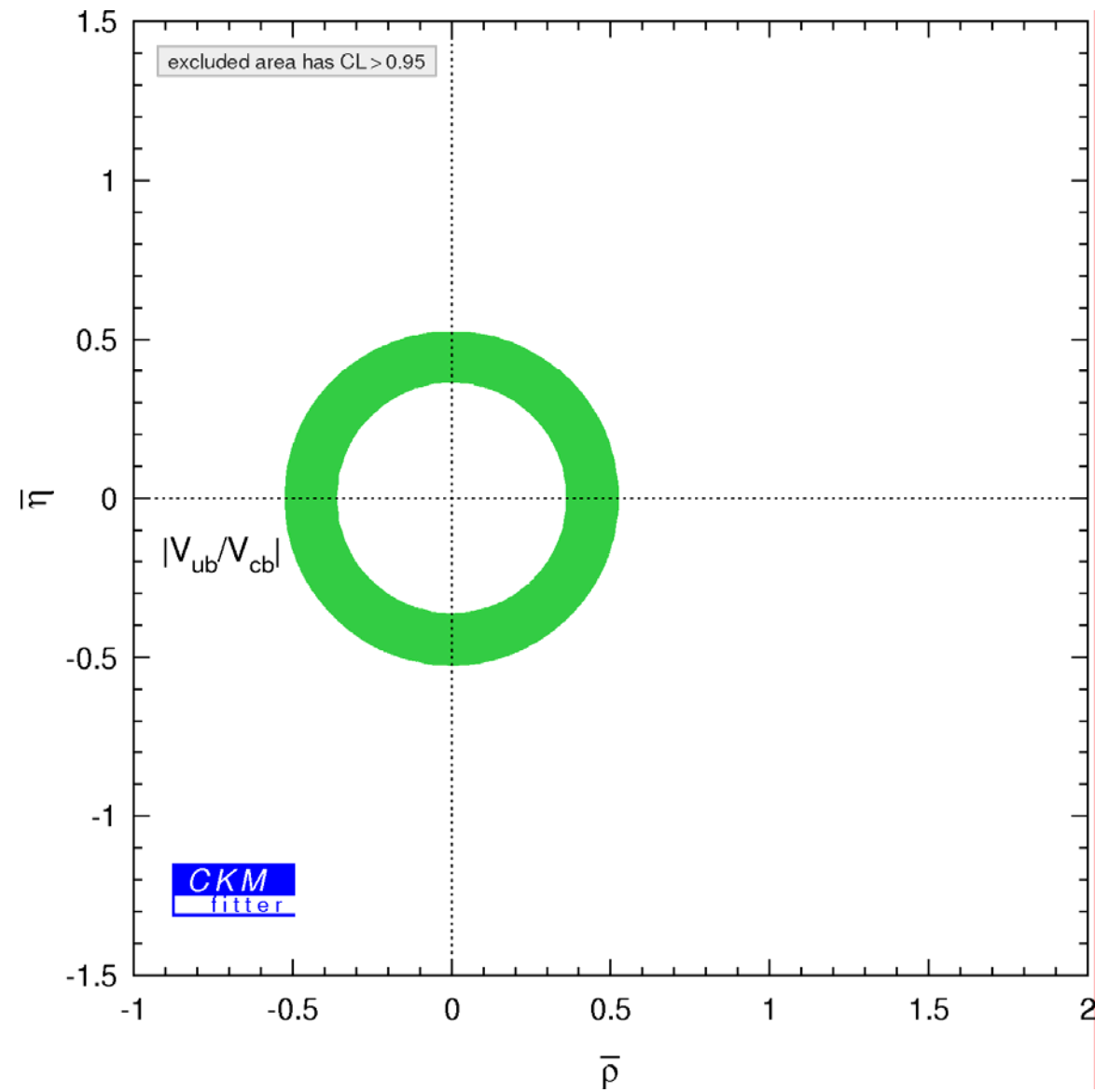
## II) Angle Measurements:

$\sin 2\beta$

$\alpha: (B \rightarrow \pi\pi, \rho\rho, \rho\pi)$

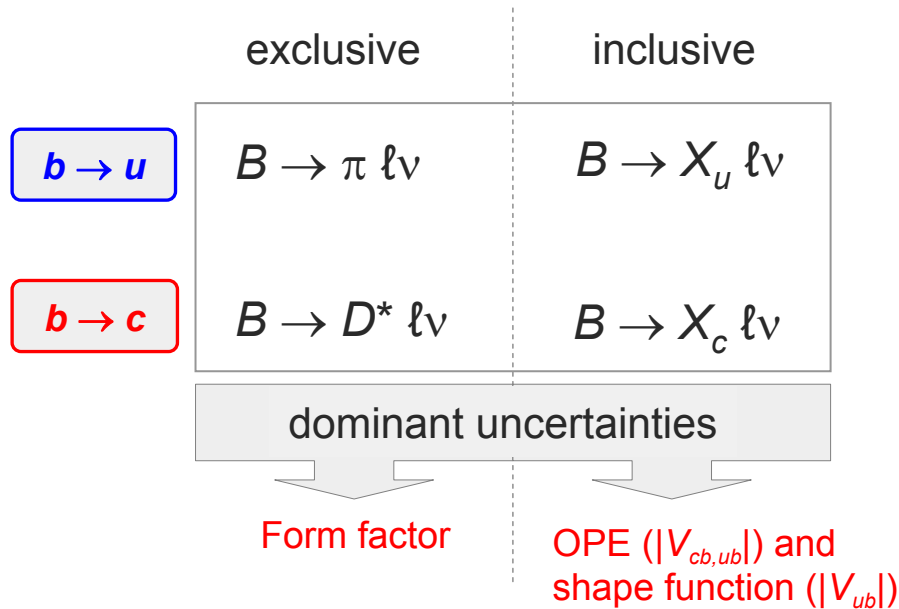
$\gamma: \text{ADS, GLW, Dalitz (GGSZ)}$

# $|V_{cb}|$ and $|V_{ub}|$



# $|V_{cb}| (\rightarrow A)$ and $|V_{ub}|$

For  $|V_{cb}|$  and  $|V_{ub}|$  exist **exclusive and inclusive** semileptonic approaches (complementary)



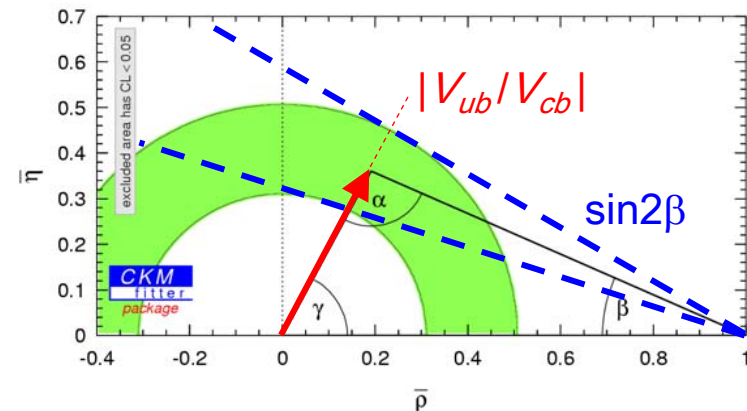
OPE parameters measured from data (spectra and moments of  $b \rightarrow s \gamma$  and  $b \rightarrow c \ell \nu$  distributions)

- $|V_{ub}| (\rightarrow \rho^2 + \eta^2)$  is crucial for the SM prediction of  $\sin(2\beta)$
- $|V_{cb}| (\rightarrow A)$  is important in the kaon system ( $\varepsilon_K, BR(K \rightarrow \pi \nu \nu), \dots$ )

Complication for charmless decays:

$$\frac{\Gamma(b \rightarrow u \ell \nu)}{\Gamma(b \rightarrow c \ell \nu)} \approx \frac{|V_{ub}|^2}{|V_{cb}|^2} \approx \frac{1}{50}$$

- need to apply kinematic cuts to suppress  $b \rightarrow c \ell \nu$  background
- measurements of partial branching fractions in restricted phase space regions
- theoretical uncertainties more difficult to evaluate

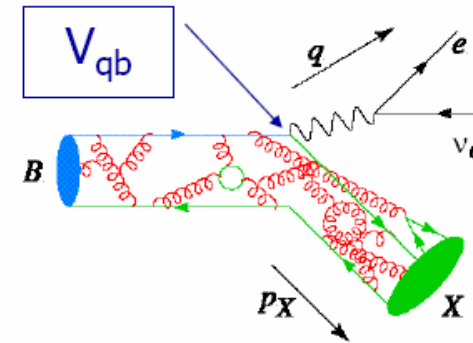


# $|V_{cb}|$ and $|V_{ub}|$

$|V_{cb}|$ : Precision measurement: 1.7% !

$|V_{cb}|_{\text{incl.}} [10^{-3}] = 41.70 \pm 0.70$  PDG06

$|V_{cb}|_{\text{excl.}} [10^{-3}] = 39.7 \pm 2.0$   
w/ FF=0.91±0.04 ICHEP06

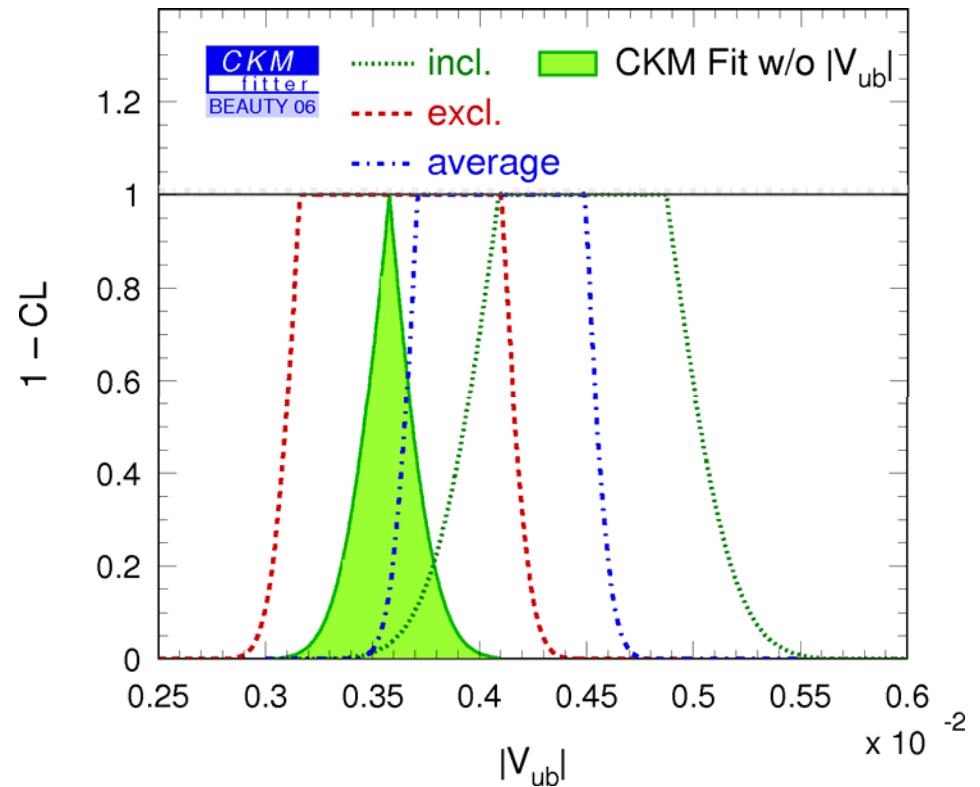


$|V_{ub}|$ :

- SF params. from  $b \rightarrow c l \nu$ , OPE from BLNP
- BR precision ~8%,  $|V_{ub}|$  excl. ~ 16%: theory dominated
- HFAG with our error budget

our average

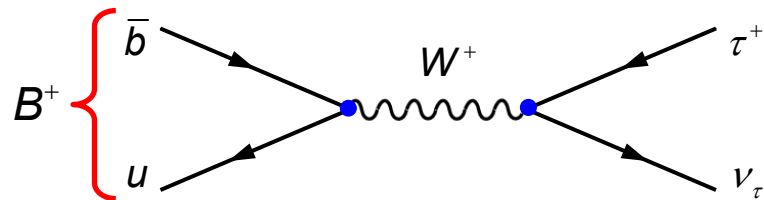
$|V_{ub}| [10^{-3}] = 4.10 \pm 0.09_{\text{exp}} \pm 0.39_{\text{theo}}$





$$B^+ \rightarrow \tau^+ \nu_\tau$$

- ☀ helicity-suppressed annihilation decay sensitive to  $f_B \times |V_{ub}|$
- ☀ Powerful together with  $\Delta m_d$  : removes  $f_B$  (Lattice QCD) dependence
- ☀ Sensitive to charged Higgs replacing the  $W$  propagator



$$\text{BR}(B^+ \rightarrow \tau^+ \nu) = \frac{G_F^2 m_B \tau_B}{8\pi} m_\tau^2 \left(1 - \frac{m_\tau^2}{m_B^2}\right)^2 f_B^2 |V_{ub}|^2$$

ICHEP06



$$\text{BR}[10^{-4}] = 0.88^{+0.68}_{-0.67} \text{ (stat)} \pm 0.11 \text{ (syst)}$$

(320m)

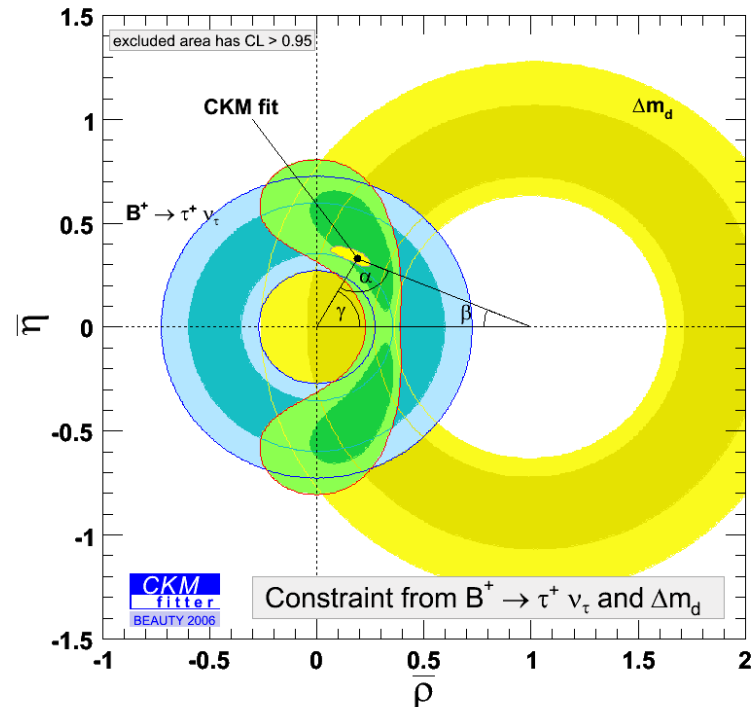


$$\text{BR}[10^{-4}] = 1.79^{+0.56}_{-0.49} \text{ (stat)}^{+0.39}_{-0.46} \text{ (syst)}$$

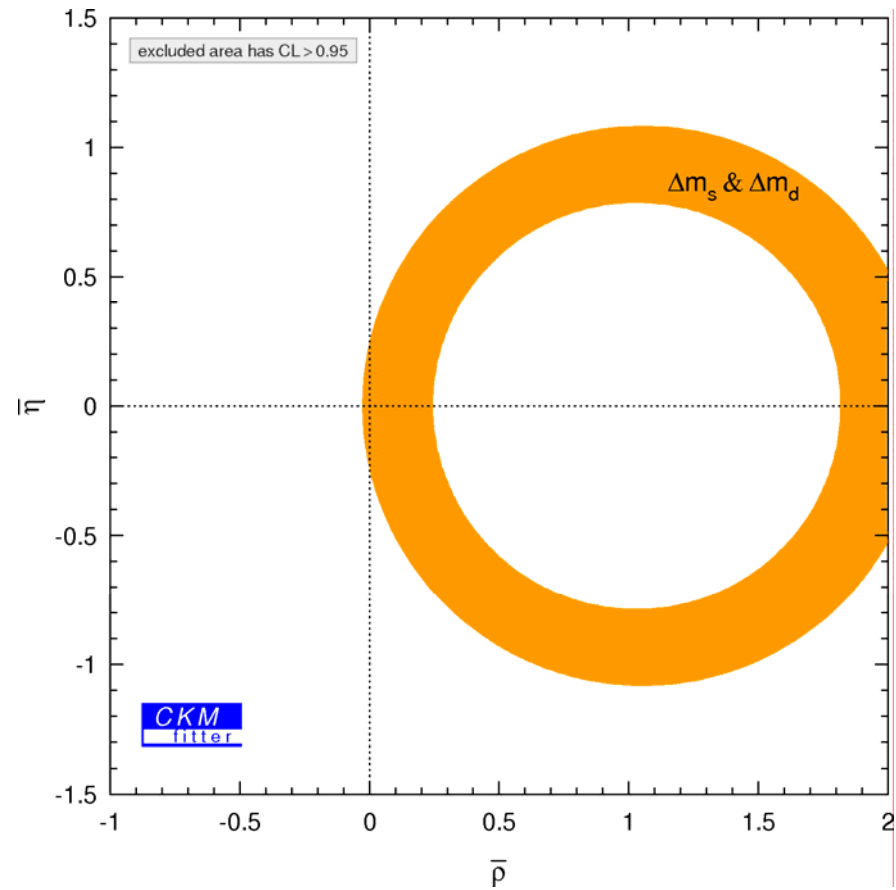
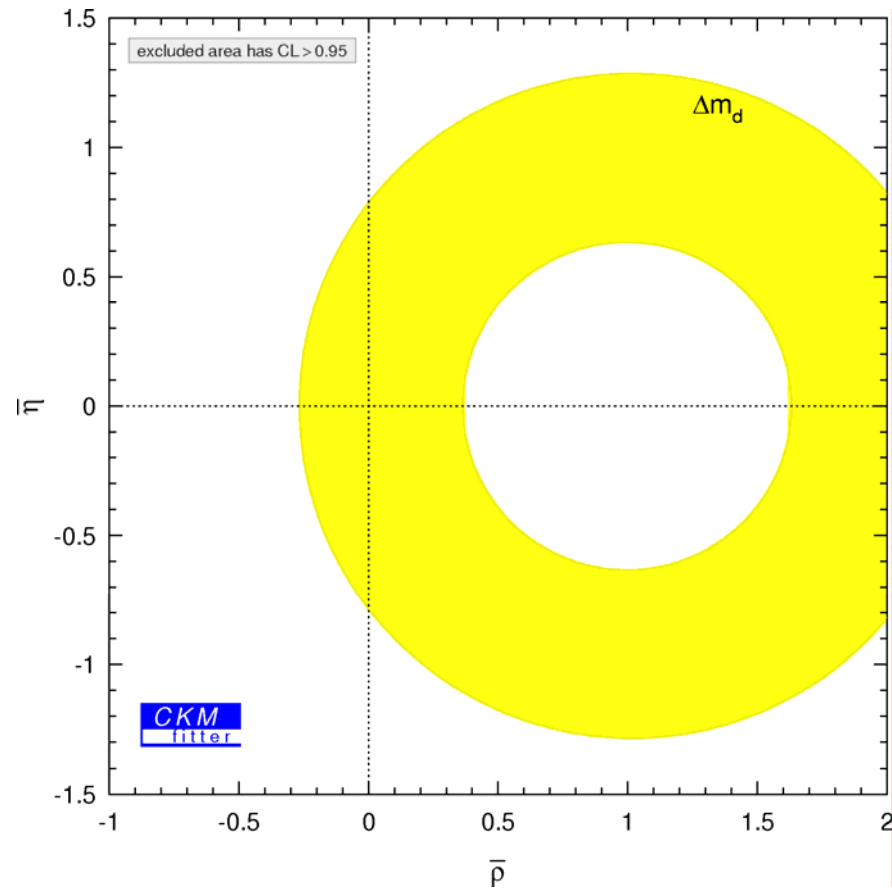
(447m)

- ☀ Prediction from global CKM fit :

$$\text{BF}(B^+ \rightarrow \tau^+ \nu_\tau) = (0.87^{+0.13}_{-0.20}) \times 10^{-4}$$



# $\Delta m_d$ and $\Delta m_s$



## $\Delta m_d$ and $\Delta m_s$ : constraints in the $(\bar{\rho}-\bar{\eta})$ plane

$$\Delta m_s = \frac{G_F^2}{6\pi^2} m_{B_s} m_W^2 \eta_B S_0(x_t) f_{B_s}^2 B_s |V_{ts} V_{tb}^*|^2$$

Very weak dependence  
on  $\bar{\rho}$  and  $\bar{\eta}$

The point is:

$$f_{B_s}^2 B_s = \frac{f_{B_s}^2 B_s}{f_{B_d}^2 B_d} f_{B_d}^2 B_d = \xi^2 f_{B_d}^2 B_d$$

$\xi$ : SU(3)-breaking corrections

Measurement of  $\Delta m_s$  **reduces the uncertainties** on  $f_{B_d}^2 B_d$  since  $\xi$  is better known

from Lattice QCD  $\sigma_{\text{rel}}(f_{B_d/s}^2 B_{d/s}) = 36\% \rightarrow \sigma_{\text{rel}}(\xi^2 = f_{B_s}^2 B_s / f_{B_d}^2 B_d) = 10\%$

→ Leads to improvement of the constraint from  $\Delta m_d$  measurement on  $|V_{td} V_{tb}^*|^2$

$$\Delta m_d = \frac{G_F^2}{6\pi^2} m_{B_d} m_W^2 \eta_B S_0(x_t) f_{B_d}^2 B_d |V_{td} V_{tb}^*|^2 \propto A^2 \lambda^6 [(1 - \bar{\rho})^2 + \bar{\eta}^2]$$

$\Delta m_s$



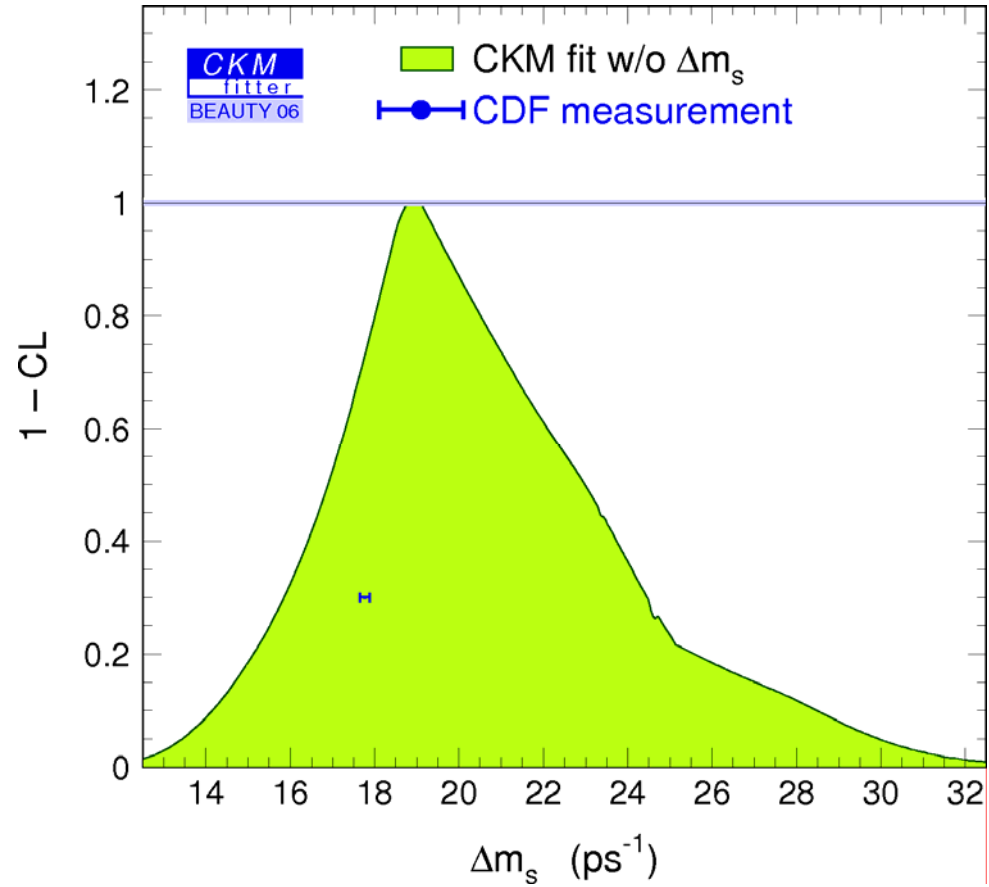
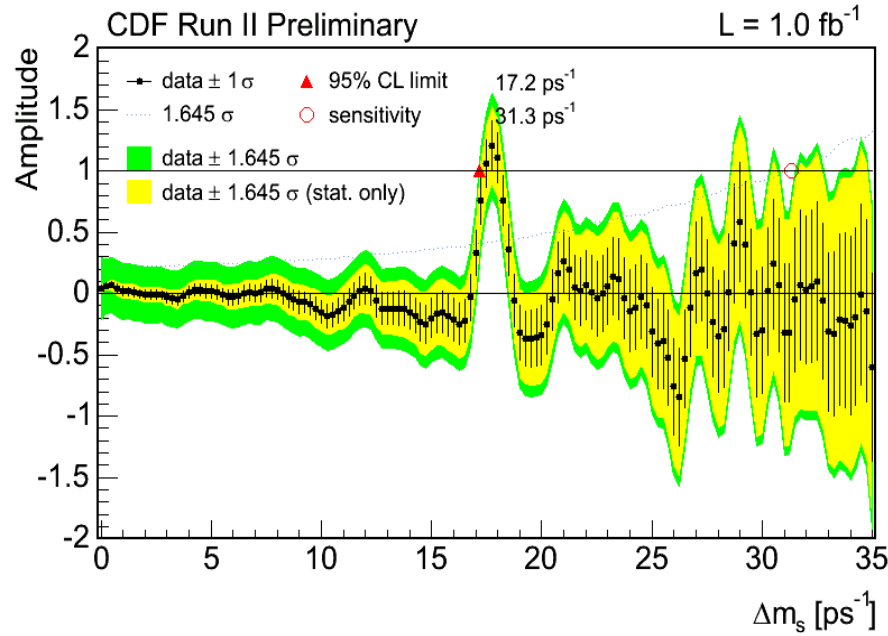
hep-ex/0603029

$17 < \Delta m_s < 21 \text{ ps}^{-1}$  @90 C.L.



hep-ex/0609040

$\Delta m_s : 17.77 \pm 0.10(\text{stat.}) \pm 0.07(\text{syst.}) \text{ ps}^{-1}$

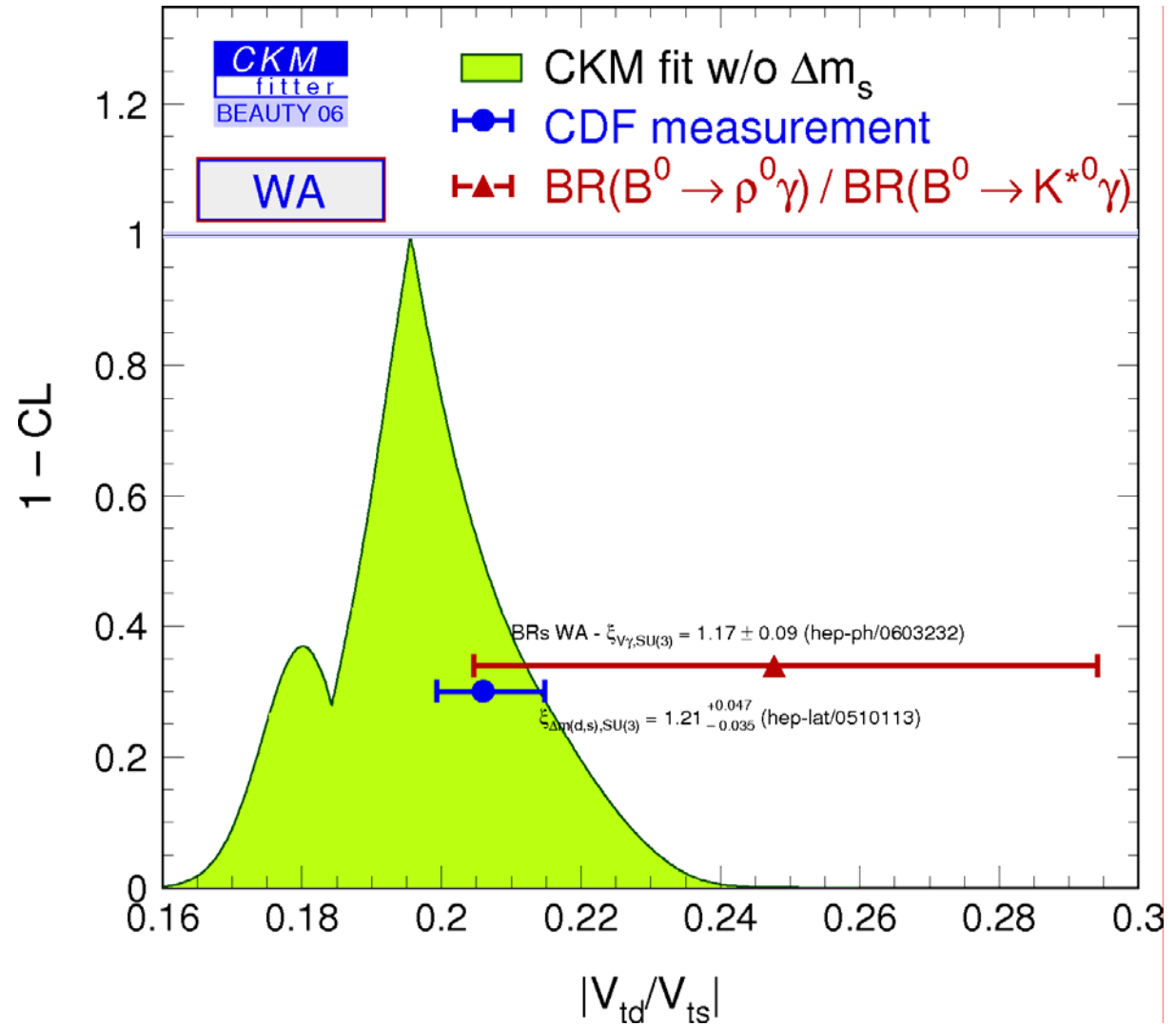


The signal has a significance of  $5.4\sigma$

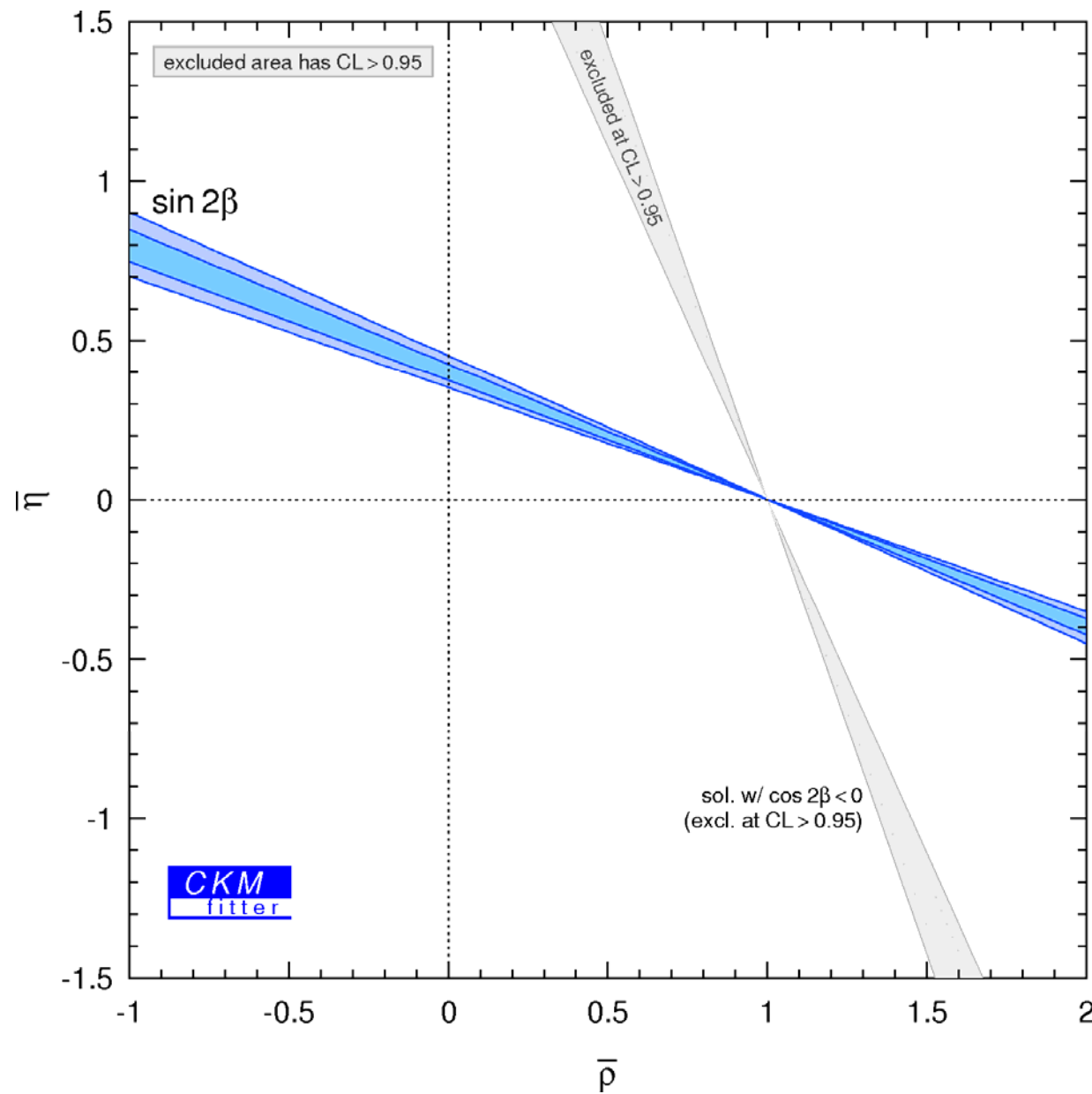
# Constraint on $|V_{td}/V_{ts}|$

$$\frac{\Delta m_d}{\Delta m_s} = \frac{m_{Bd}}{m_{Bs}} \xi_{\Delta m}^{-2} \frac{|V_{td}|^2}{|V_{ts}|^2}$$

→ First strong indication that  $B_s - \bar{B}_s$  mixing is probably SM-like.



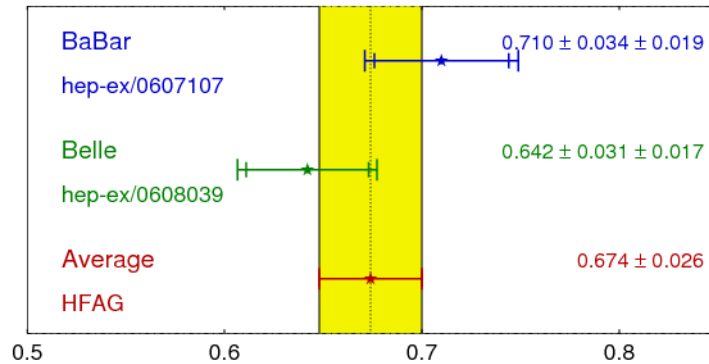
# angle $\beta$



# sin 2β

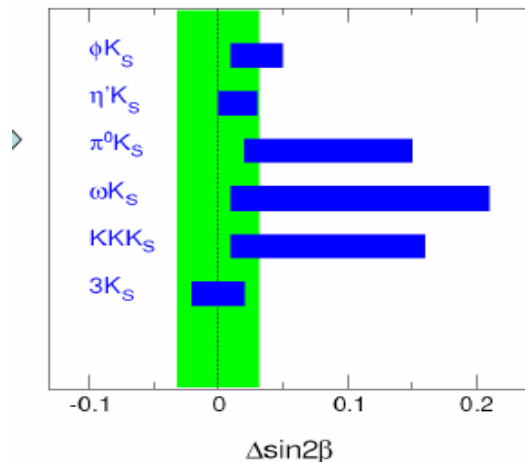
☀ “The” *raison d’être* of the *B* factories:

$$\sin(2\beta) \equiv \sin(2\phi_1) \quad \text{HFAG ICHEP 2006 PRELIMINARY}$$

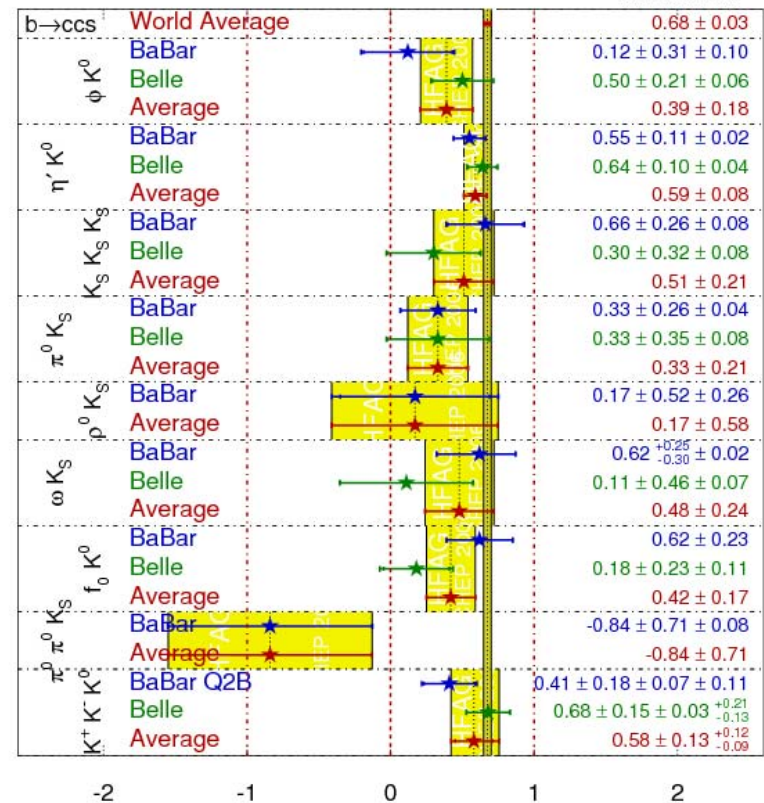


☀ Conflict with  $\sin 2\beta_{\text{eff}}$  from s-penguin modes ? (New Physics (NP)?)

some of recent QCDF estimates  
 $\sin 2\beta_{\text{eff}}^f - \sin 2\beta$



$$\sin(2\beta^{\text{eff}}) \equiv \sin(2\phi_1^{\text{eff}}) \quad \text{HFAG ICHEP 2006 PRELIMINARY}$$

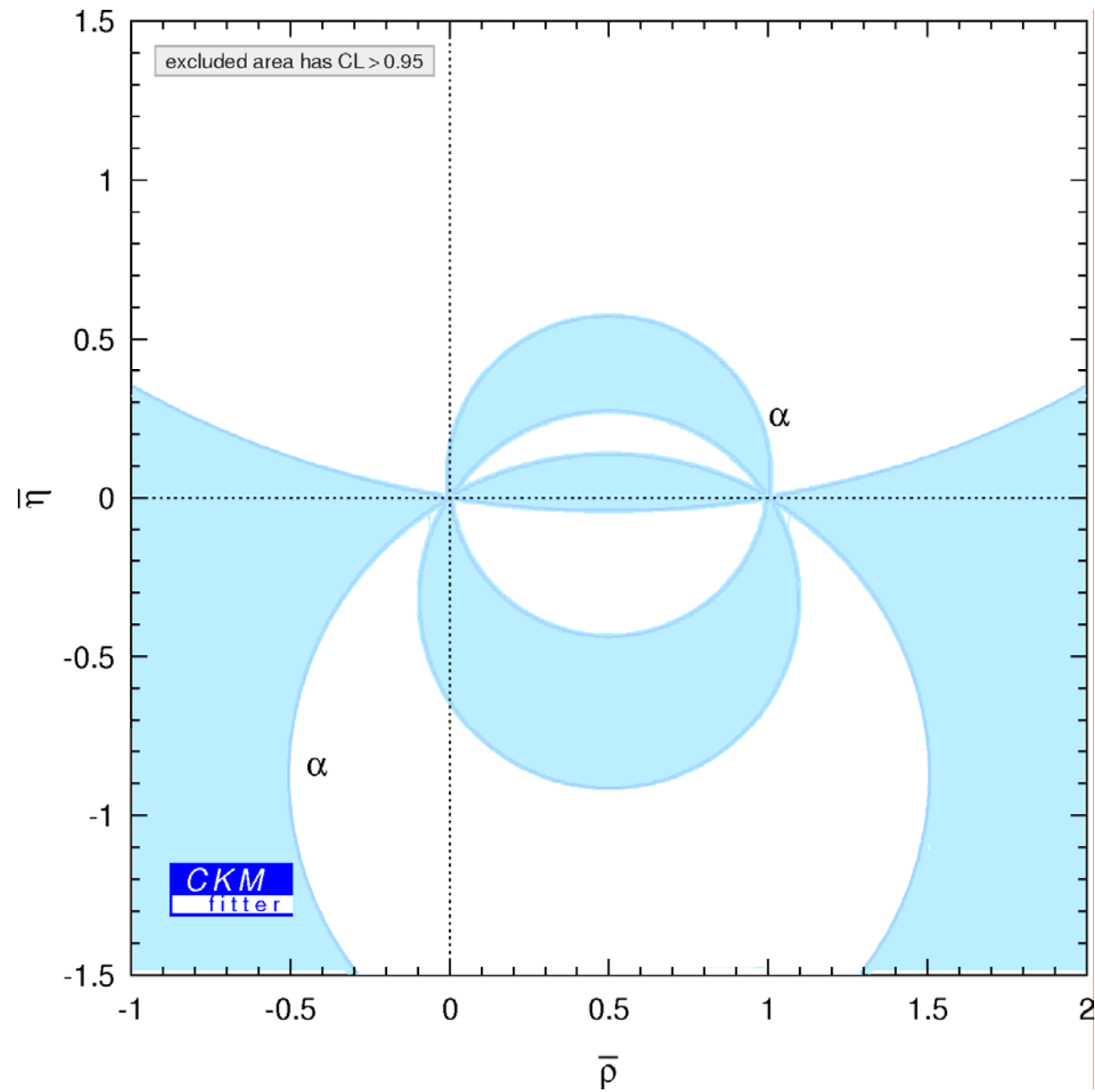


NP can contribute differently among the various s-penguin modes (Naïve average:  $0.52 \pm 0.05$ ).

NB: a disagreement would falsify the SM. The interference NP/SM amplitudes introduces hadronic uncertainties

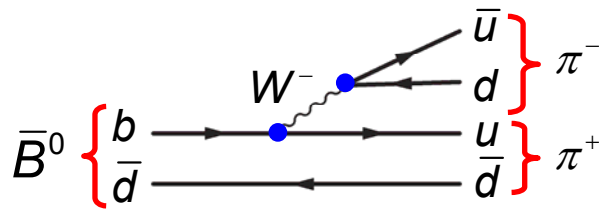
→ Cannot determine the NP parameters cleanly

# angle $\alpha$





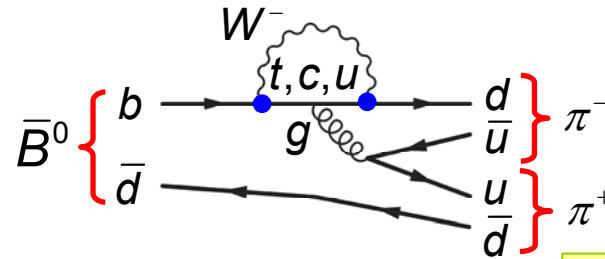
# angle $\alpha$



Tree : dominant

$$\propto V_{ub} V_{ud}^*$$

$$\propto \lambda^3$$



Penguin : competitive ?

$$\propto V_{tb} V_{td}^*$$

$$\propto \lambda^3$$

☀ Time-dependent CP observable :

$$A_{h^+h^-}(t) = S_{h^+h^-} \sin(\Delta m_d t) - C_{h^+h^-} \cos(\Delta m_d t)$$

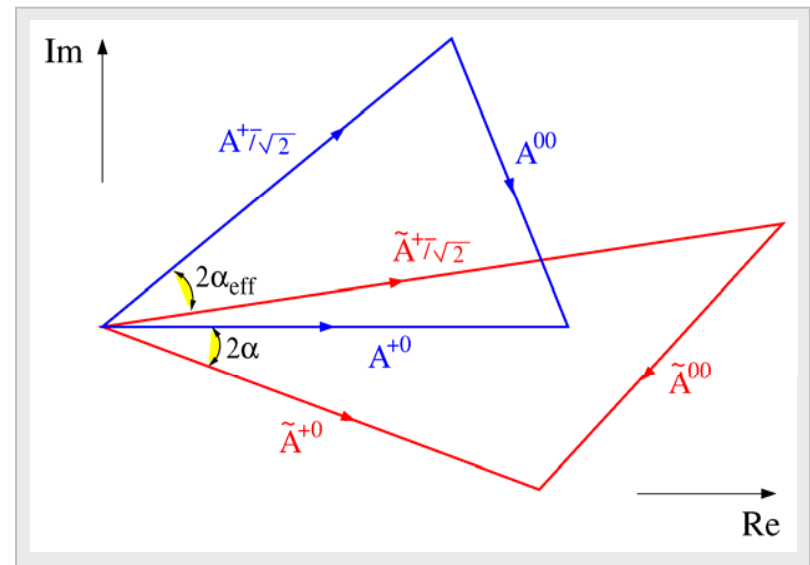
$$= \sqrt{1 - C_{h^+h^-}^2} \sin(2\alpha_{\text{eff}}) \cdot \sin(\Delta m_d t) - C_{h^+h^-} \cos(\Delta m_d t)$$

Time-dependent CP analysis of  $B^0 \rightarrow \pi^+\pi^-$  alone determines  $\alpha_{\text{eff}}$  : but, we need  $\alpha$  !



**Isospin analysis**

( $\alpha$  can be resolved up to an 8-fold ambiguity within  $[0, \pi]$ )

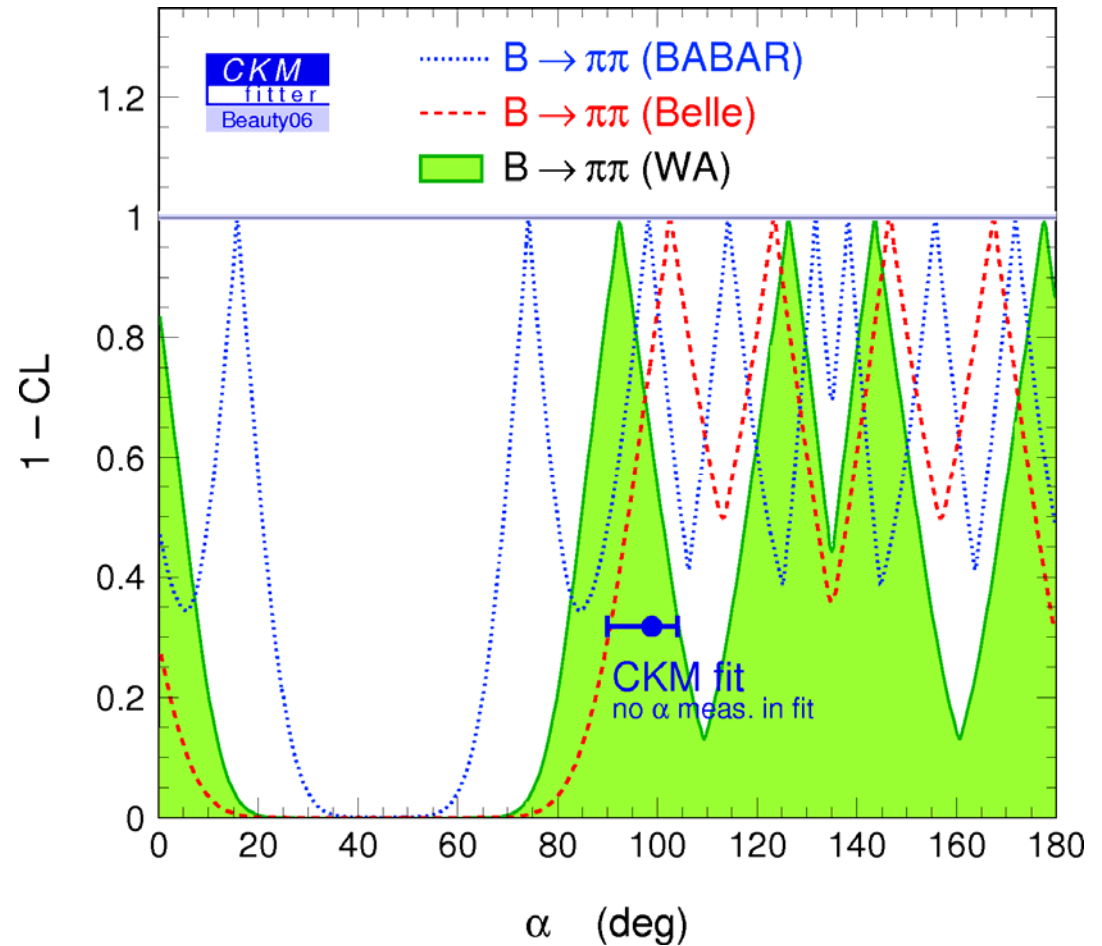
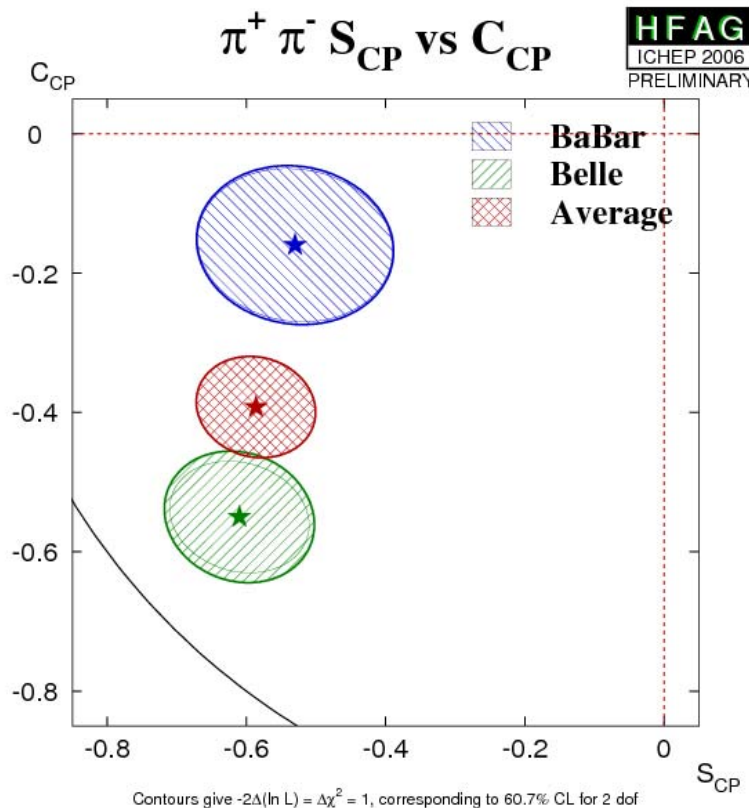


# Isospin Analysis: $B \rightarrow \pi\pi$

	BABAR (347m)	Belle (532m)	Average
$S_{\pi\pi}$	$-0.53 \pm 0.14 \pm 0.02$	$-0.61 \pm 0.10 \pm 0.04$	$-0.58 \pm 0.09$
$C_{\pi\pi}$	$-0.16 \pm 0.11 \pm 0.03$	$-0.55 \pm 0.08 \pm 0.05$	$-0.39 \pm 0.07$

“agreement”:  $2.6\sigma$

BABAR & Belle



# Isospin Analysis: $B \rightarrow \rho\rho$

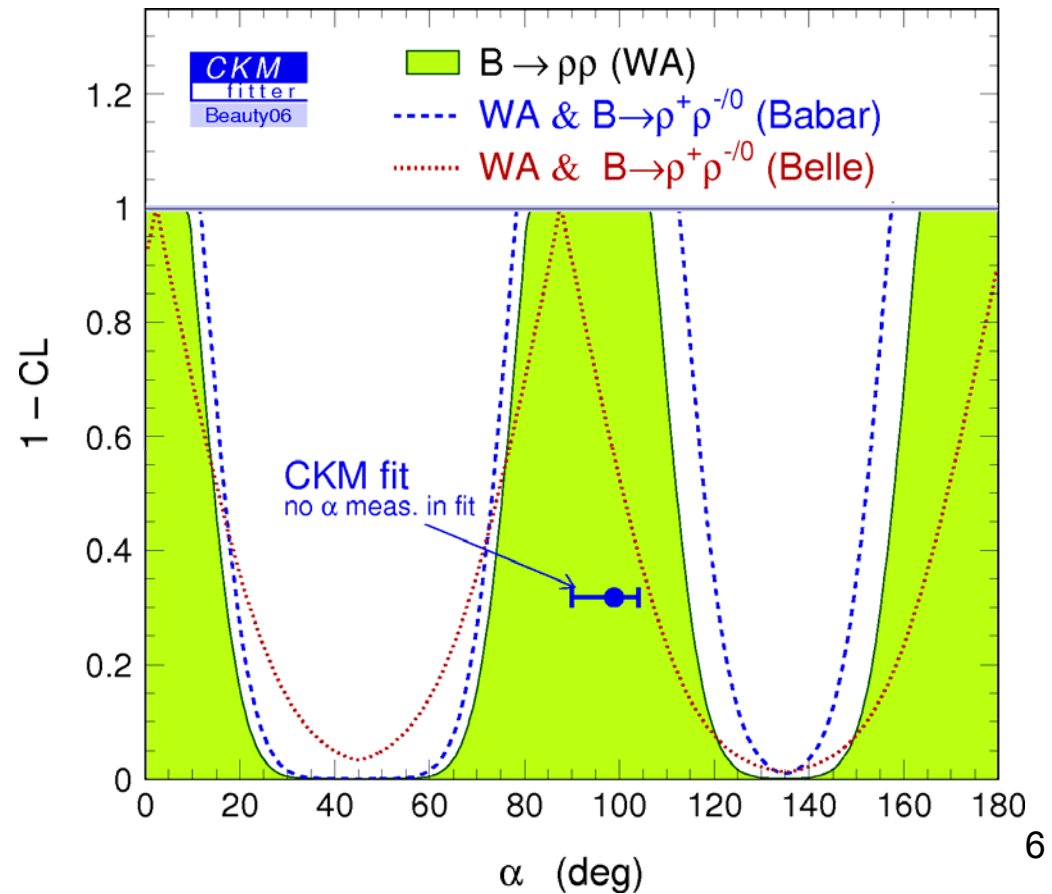
	BABAR (347m)	Belle (275m)	Average
$S_{\rho\rho}$	$-0.19 \pm 0.21^{+0.05}_{-0.07}$	$0.08 \pm 0.41 \pm 0.09$	$-0.13 \pm 0.19$
$C_{\rho\rho}$	$-0.07 \pm 0.15 \pm 0.06$	$0.0 \pm 0.3 \pm 0.09$	$-0.06 \pm 0.14$

BABAR & Belle

BABAR (347m)	
$f_L^{00}$	$0.86^{+0.11}_{-0.13} \pm 0.06$
$BR^{00}$	$(1.2 \pm 0.4 \pm 0.3) \times 10^{-6}$

☀ Isospin analysis :

$$\alpha = [ 94 \pm 21 ]^\circ$$



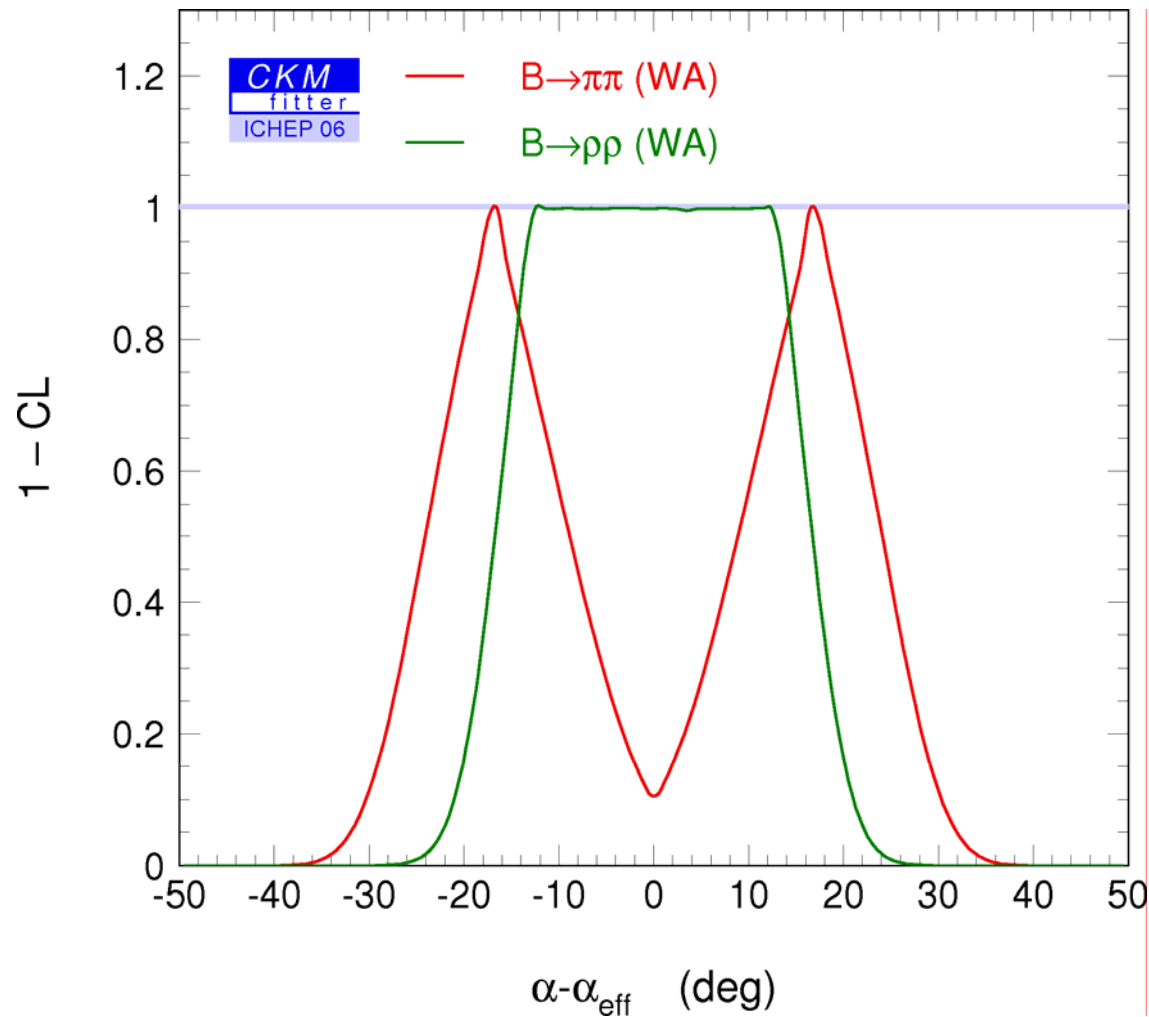
# Isospin Analysis: angle $\alpha_{\text{eff}}$ [ $B \rightarrow \pi\pi/\rho\rho$ ]

☀ Isospin analysis  $B \rightarrow \pi\pi$ :

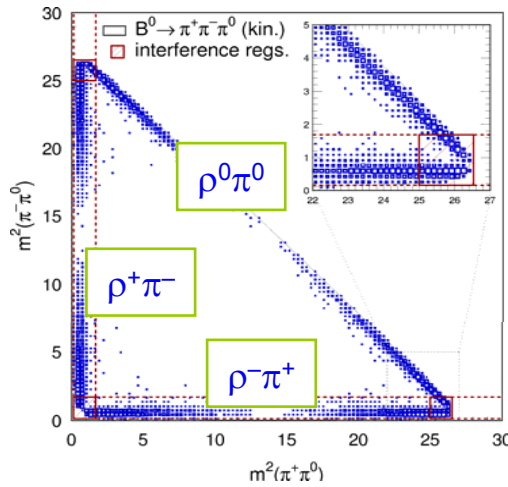
$$|\alpha - \alpha_{\text{eff}}| < 32.1^\circ \text{ (95\% CL)}$$

☀ Isospin analysis  $B \rightarrow \rho\rho$ :

$$|\alpha - \alpha_{\text{eff}}| < 22.4^\circ \text{ (95\% CL)}$$



# The $B \rightarrow \rho\pi$ System



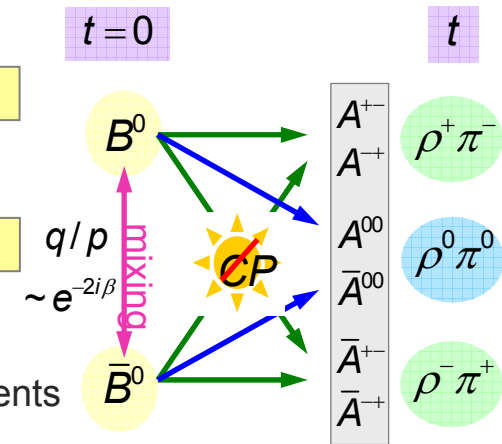
Dominant mode  $\rho^+\pi^-$  is not a  $CP$  eigenstate

Aleksan *et al.*, NP B361, 141 (1991)

Amplitude interference in Dalitz plot

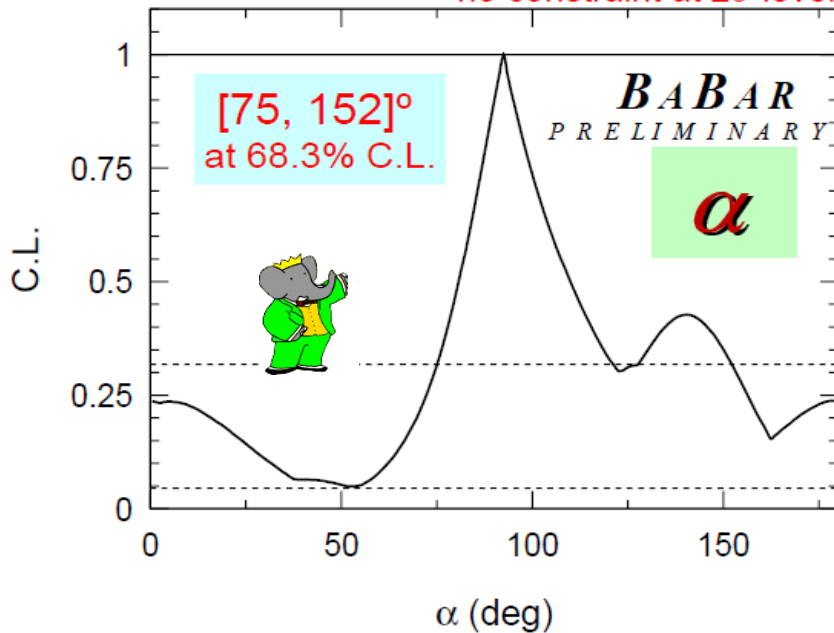
Snyder-Quinn, PRD 48, 2139 (1993)

- simultaneous fit of  $\alpha$  and strong phases
- Measure 26 (27) bilinear Form Factor coefficients
- correlated  $\chi^2$  fit to determine physics quantities



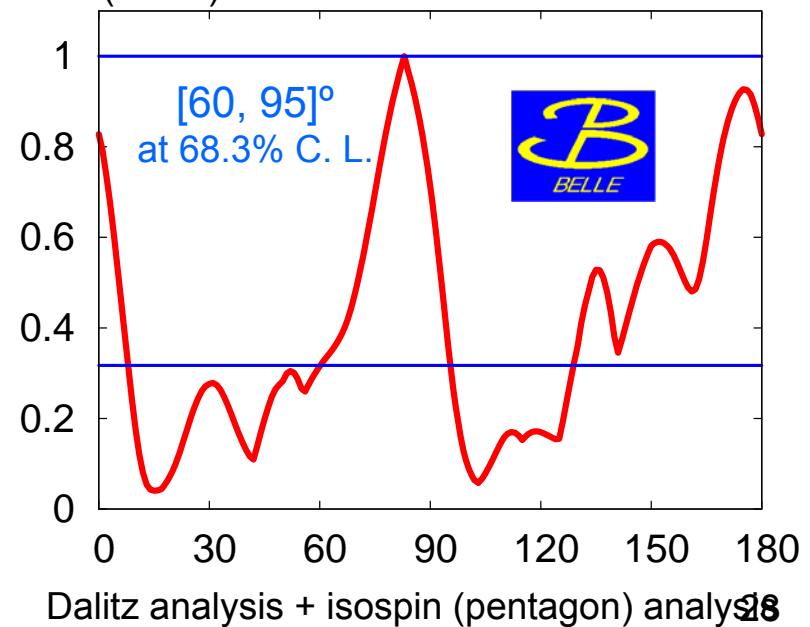
(347m)

no constraint at  $2\sigma$  level



(449m)

no constraint at  $2\sigma$  level



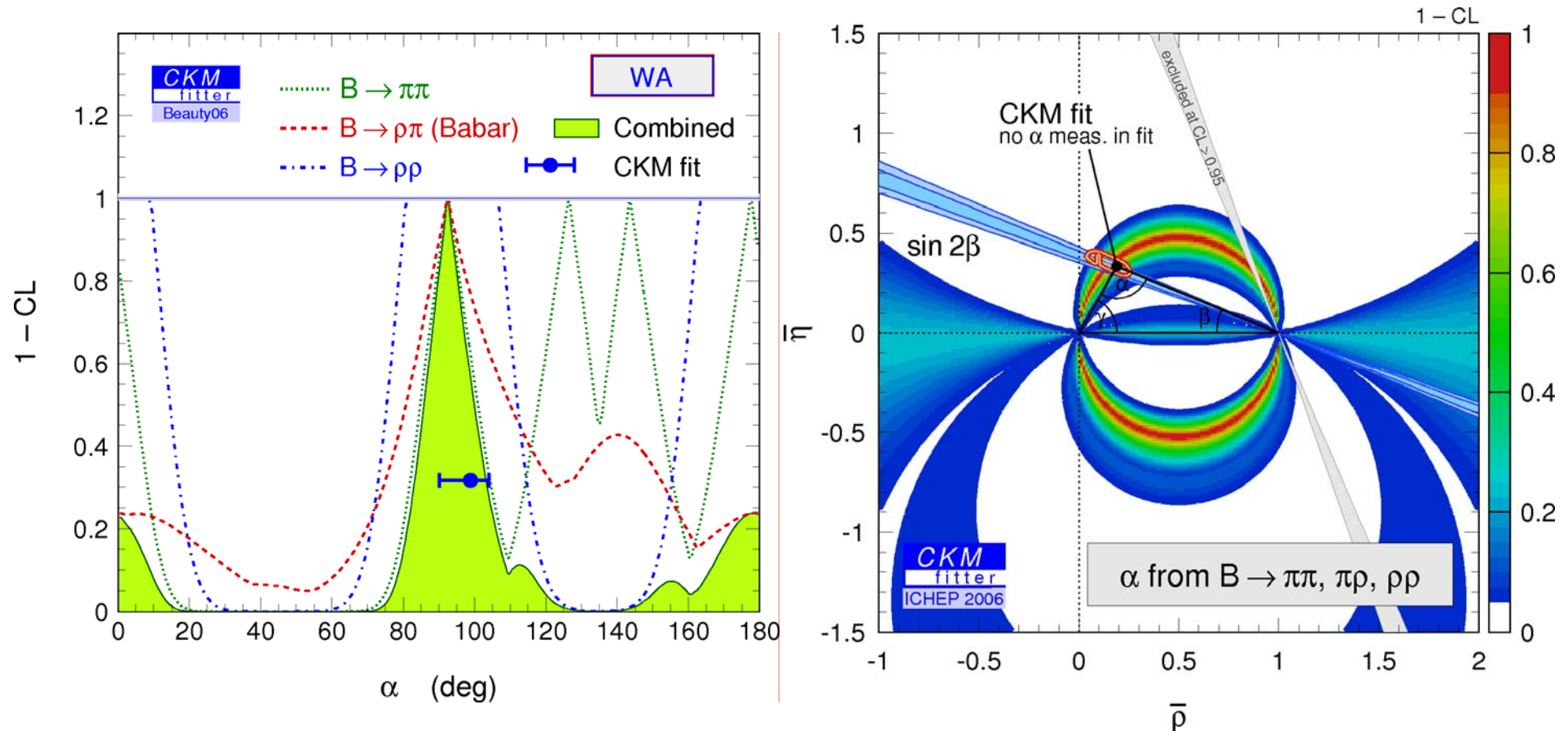
Lipkin *et al.*, PRD 44, 1454 (1991)

# Isospin Analysis: angle $\alpha$ [ $B \rightarrow \pi\pi$ | $\rho\pi$ | $\rho\rho$ ]

$$\alpha_{\text{B-Factories}} = [ 93^{+11}_{-9} ]^\circ$$



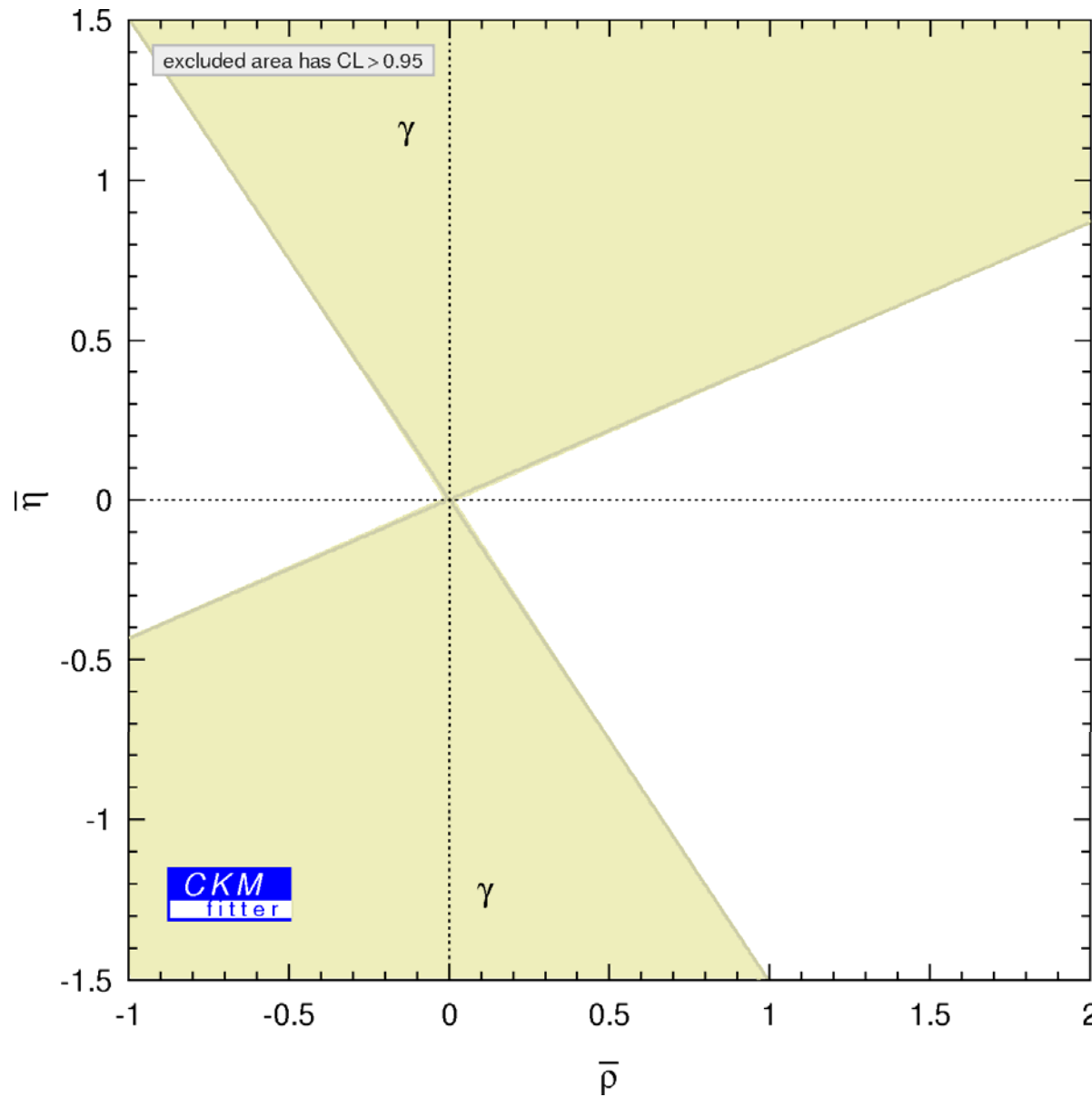
$$\alpha_{\text{Global Fit}} = [ 100^{+5}_{-7} ]^\circ$$



$B \rightarrow \rho\rho$ : at very large statistics, systematics and model-dependence will become an issue

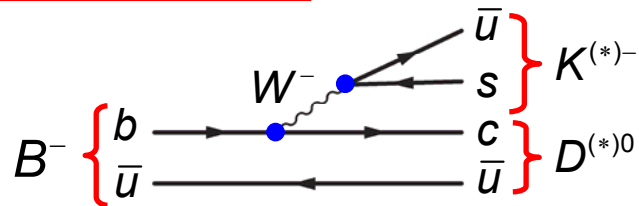
$B \rightarrow \rho\pi$  Dalitz analysis: model-dependence is an issue !

# angle $\gamma$



# angle $\gamma$ [ next UT input that is not theory limited ]

$$b \rightarrow c\bar{u}s, u\bar{c}s$$

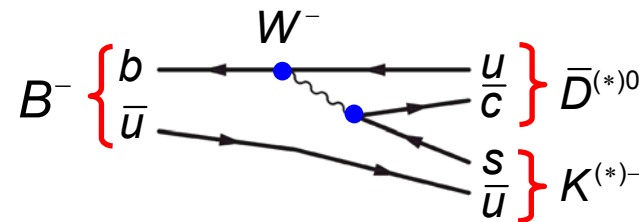


Tree: dominant

$$\propto V_{cb} V_{us}^*$$

$$\propto \lambda^3$$

Tree: color-suppressed



$$\propto V_{ub} V_{cs}^*$$

$$\propto \lambda^3 \sqrt{\rho^2 + \eta^2}$$

No Penguins 😊

relative CKM phase :  $\gamma$

Several variants:

- GLW :  $D^0$  decays into  $CP$  eigenstate
- ADS :  $D^0$  decays to  $K^-\pi^+$  (favored) and  $K^+\pi^-$  (suppressed)
- GGSZ :  $D^0$  decays to  $K_S\pi^+\pi^-$  (interference in Dalitz plot)

Gronau-London, PL B253, 483 (1991);  
Gronau-Wyler, PL B265, 172 (1991)

Atwood-Dunietz-Soni, PRL 78, 3257 (1997)

Giri *et al*, PRD 68, 054018 (2003)

➡ All methods fit simultaneously:  $\gamma$ ,  $r_B$  and  $\delta$  (different  $r_B$  and  $\delta$ )

$r_B$   
 $r_B^*$  } how small ?

$\sigma_\gamma$  depends significantly on the value of  $r_B$



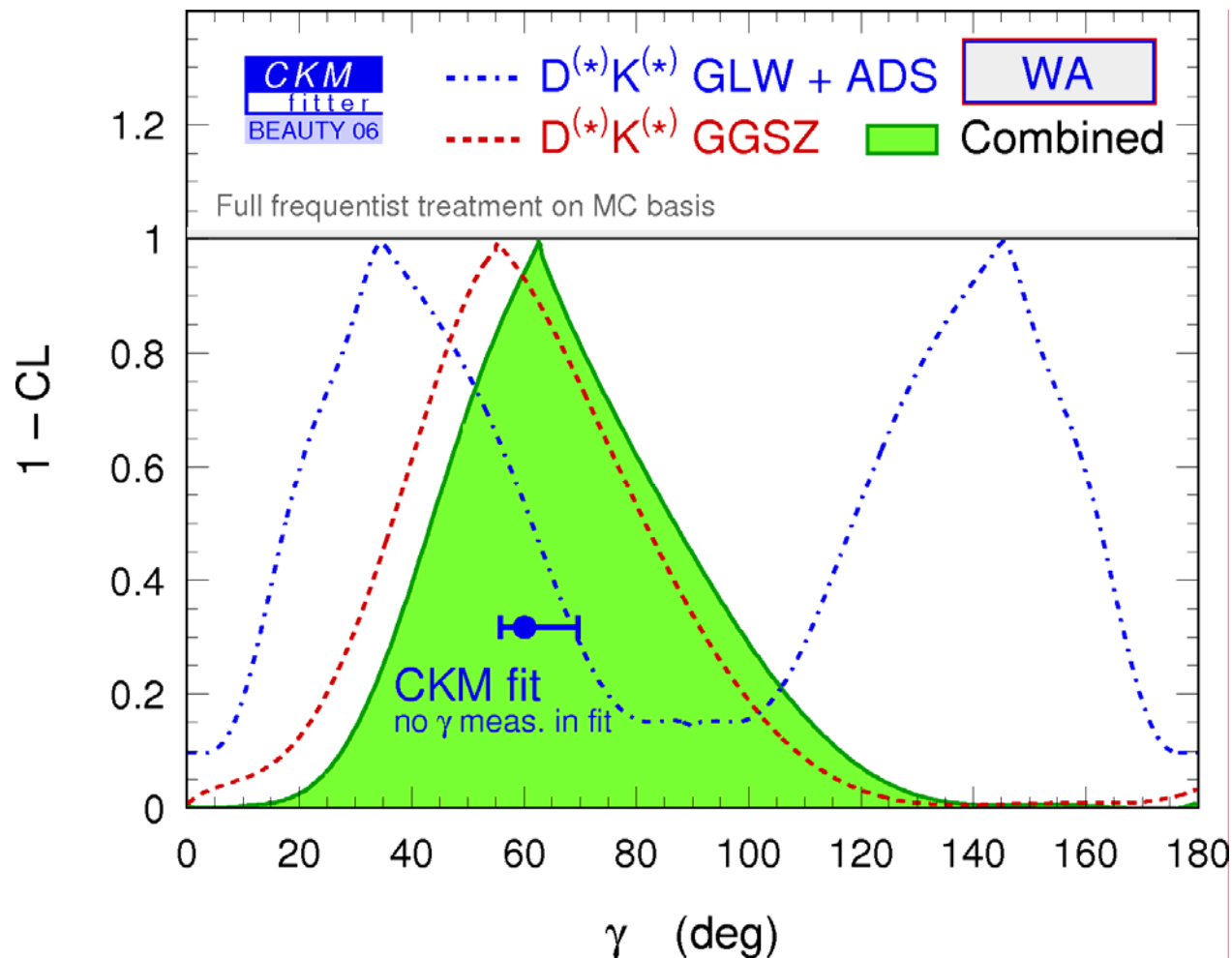
# Constraint on $\gamma$

$$r_B(DK) = 0.10^{+0.03}_{-0.04}$$

$$r_B(D^*K) = 0.10^{+0.04}_{-0.06}$$

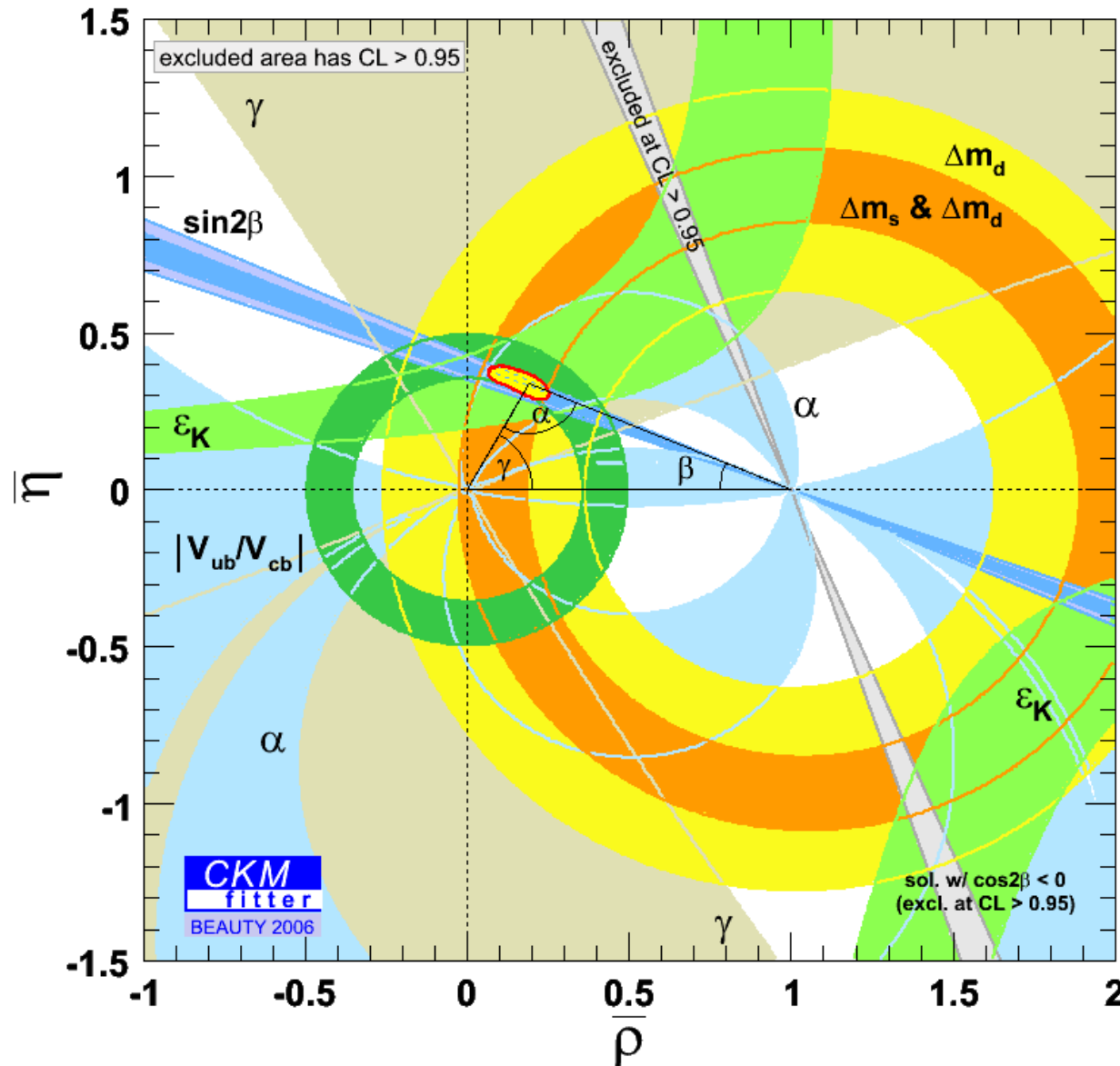
$$r_B(DK^*) = 0.11^{+0.09}_{-0.11}$$

$$\gamma_{B\text{-Factories}} = [60^{+38}_{-24}]^\circ \quad \Rightarrow \quad \gamma_{\text{Global Fit}} = [59^{+9}_{-4}]^\circ$$



# Putting it all together

the global CKM fit



Inputs:

$$\left| \frac{V_{ub}}{V_{cb}} \right|$$

$$\Delta m_d$$

$$\Delta m_s$$

$$B \rightarrow \tau \nu$$

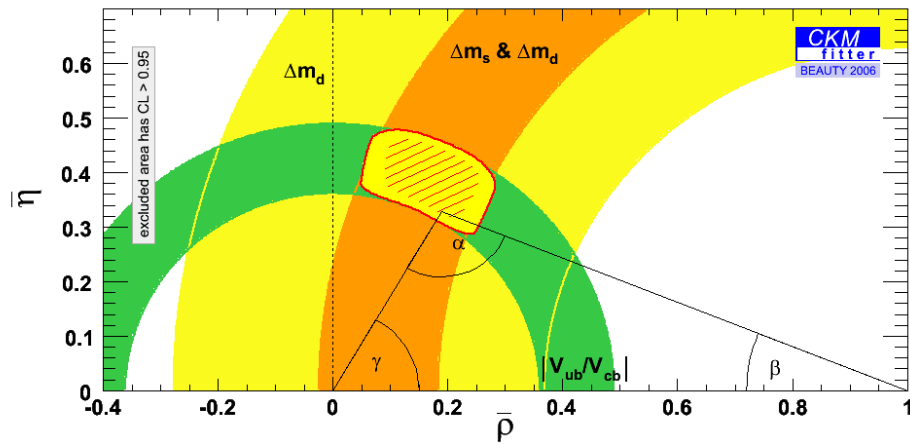
$$|\epsilon_K|$$

$$\sin 2\beta$$

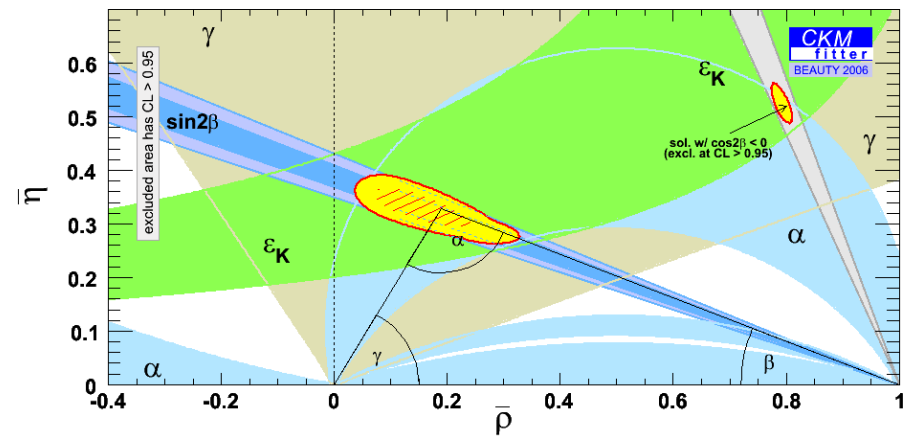
$$\alpha$$

$$\gamma$$

# The global CKM fit: Testing the CKM Paradigm

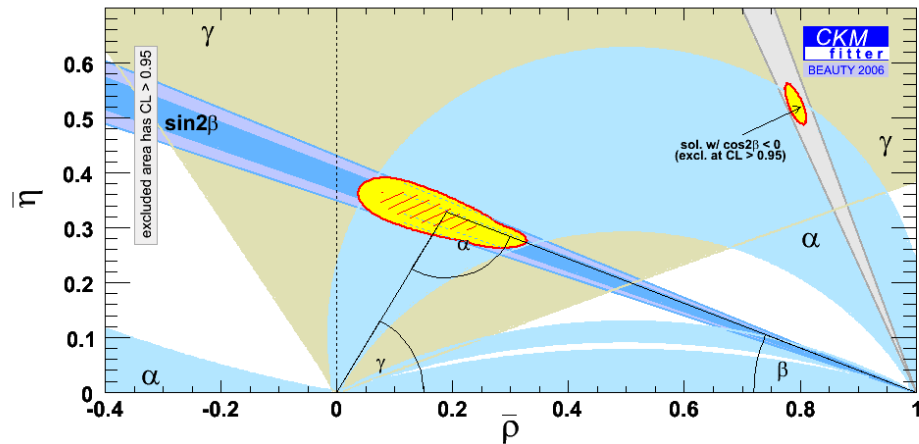


CP Conserving

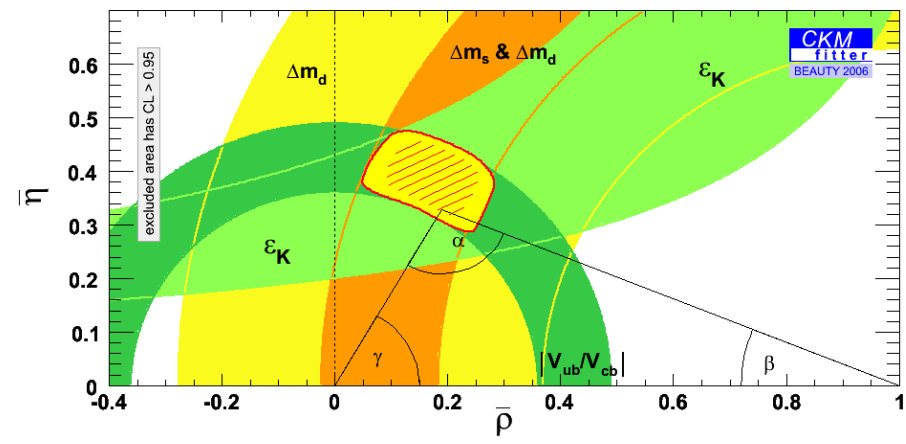


CP Violating

*CP-insensitive observables imply CP violation !*

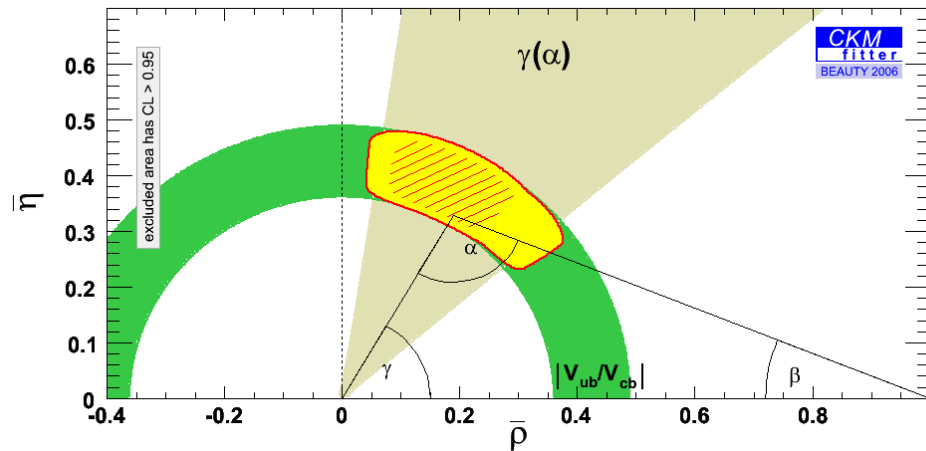


Angles (no theory)

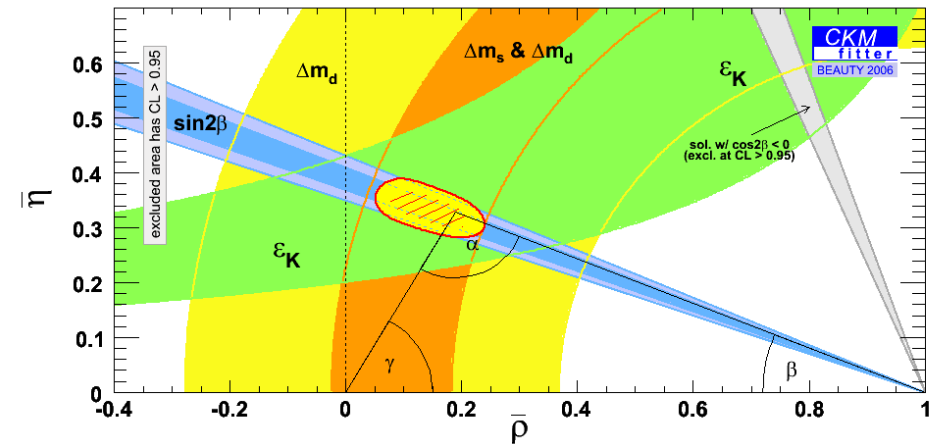


No angles (with theory)

# The global CKM fit: Testing the CKM Paradigm (cont.)



Tree (NP-Free) “Reference UT”



Loop

[No NP in  $\Delta I=3/2$   $b \rightarrow d$  EW penguin amplitude  
Use  $\alpha$  with  $\beta$  (charmonium) to cancel NP amplitude]

CKM mechanism: dominant source of CP violation

The global fit is not the whole story: several  $\Delta F=1$  rare decays are not yet measured

→ Sensitive to NP

# The global CKM fit: selected predictions

## Wolfenstein parameters:

$$A = 0.806_{-0.014}^{+0.014} \quad \lambda = 0.2272_{-0.0010}^{+0.0010} \quad \bar{\rho} = 0.195_{-0.055}^{+0.022} \quad \bar{\eta} = 0.326_{-0.015}^{+0.027}$$

## Jarlskog invariant:

$$J = (2.91_{-0.14}^{+0.25}) \times 10^{-5}$$

## UT Angles:

$$\alpha = (99.0_{-9.4}^{+4.0})^\circ \quad \beta = (22.03_{-0.62}^{+0.72})^\circ \quad \gamma = (59.0_{-3.7}^{+9.2})^\circ \quad \Sigma_{meas.} = (175_{-27}^{+40})^\circ$$

## UT sides:

$$R_u = 0.380_{-0.009}^{+0.011} \quad R_t = 0.868_{-0.025}^{+0.060}$$

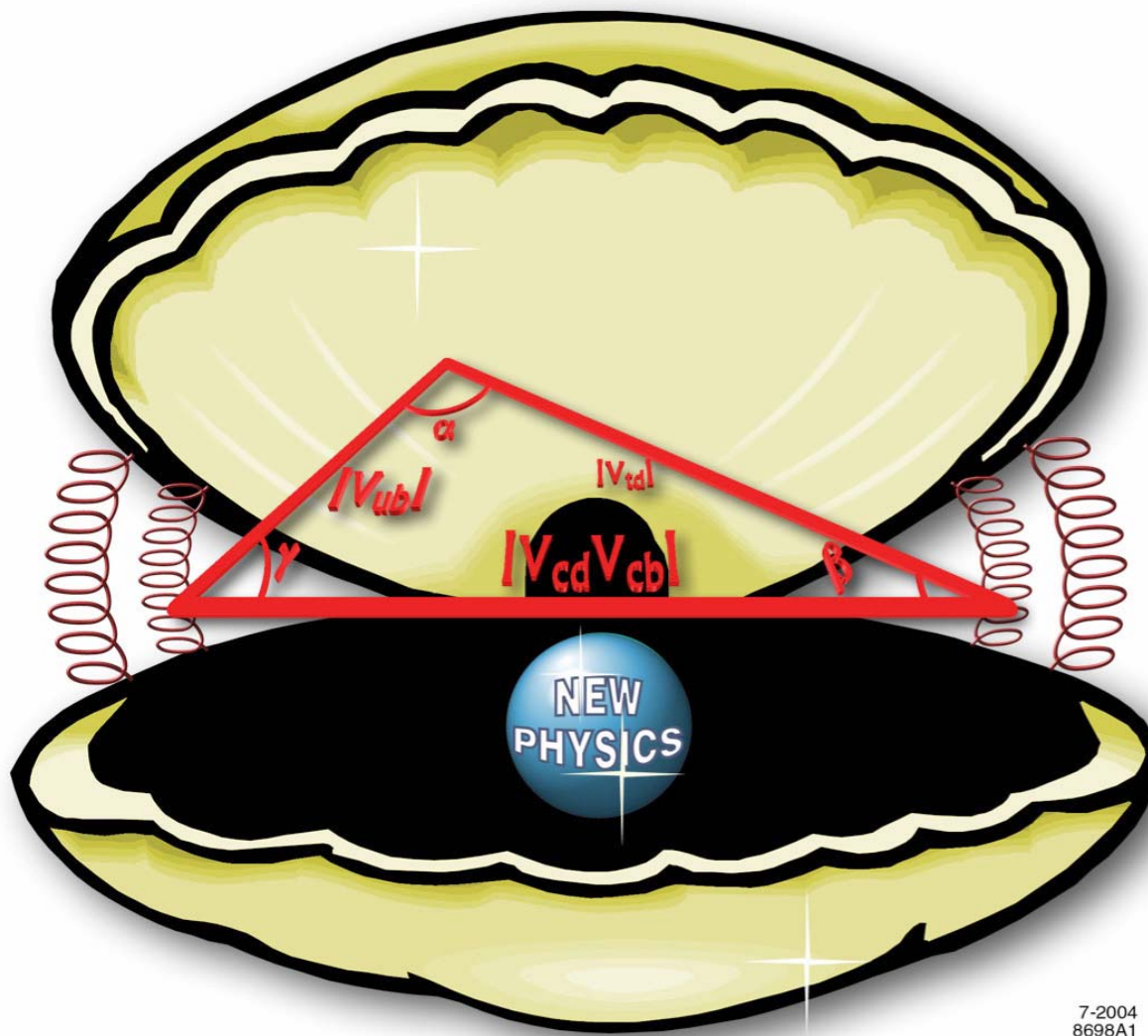
## B-B mixing:

$$\Delta m_s = (18.9_{-2.8}^{+5.7}) \text{ ps}^{-1} \quad (\text{CKM Fit}) \quad \Delta m_s : 17.77 \pm 0.1 (\text{stat.}) \pm 0.07 (\text{syst.}) \text{ ps}^{-1} \\ (\text{direct, CDF})$$

## $B \rightarrow \tau \nu$

$$\text{BF}(B^+ \rightarrow \tau^+ \nu_\tau) = (0.87_{-0.20}^{+0.13}) \times 10^{-4} \quad (\text{CKM Fit}) \quad (1.45_{-0.43}^{+0.46}) \times 10^{-4} \quad (\text{direct, WA})^{36}$$

# New Physics?

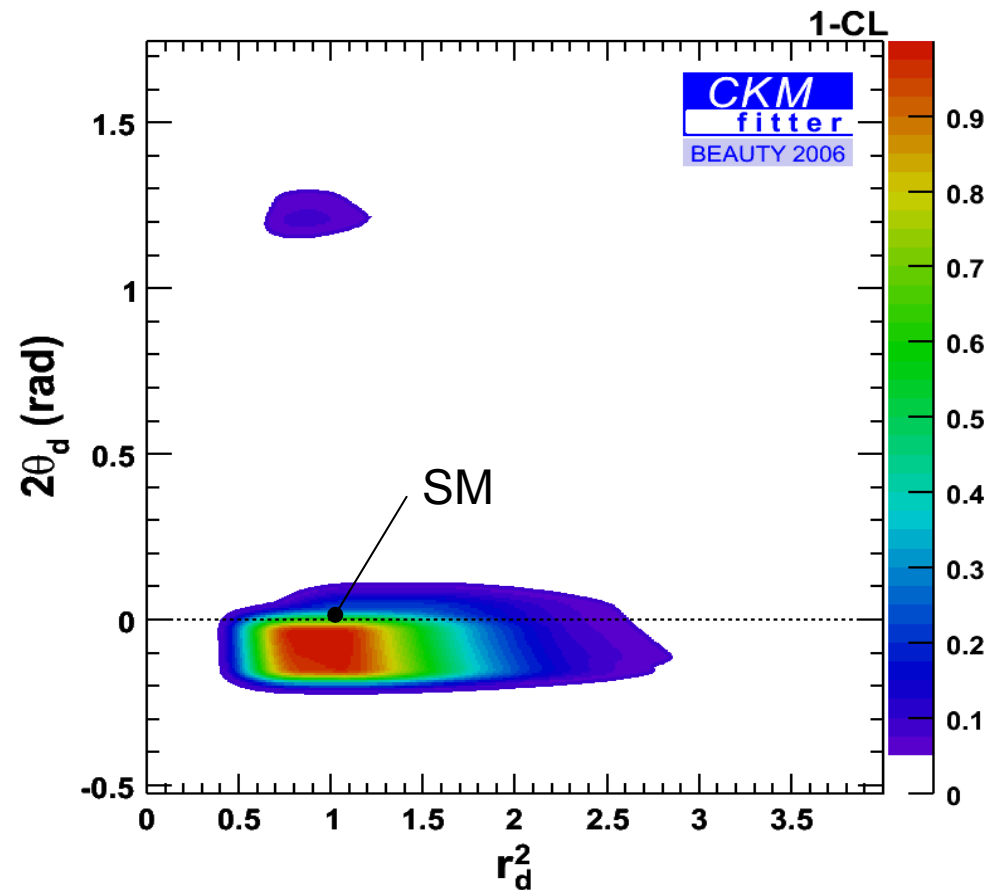


7-2004  
8698A1

# New Physics in $B_d$ - $\bar{B}_d$ Mixing?

$$r_d^2 \exp(2i\theta_d) = \frac{\langle B^0 | H_{eff}^{full} | \bar{B}^0 \rangle}{\langle B^0 | H_{eff}^{SM} | \bar{B}^0 \rangle}$$

No significant modification of the  $B$ - $\bar{B}$  mixing amplitude



# NP Parameterization in $B_s$ system

$$\frac{\langle B_s^0 | H_{eff}^{SM+NP} | \bar{B}_s^0 \rangle}{\langle B_s^0 | H_{eff}^{SM} | \bar{B}_s^0 \rangle} = r_s^2 e^{i2\theta_s} = 1 + h_s e^{i2\sigma_s}$$

Grossman, PL **B380**, 99 (1996)  
Dunietz, Fleischer, Nierste, PRD **63**, 114015 (2001)

Hypothesis: NP in loop processes only (negligible for tree processes)

$$\text{Mass difference: } \Delta m_s = (\Delta m_s)^{SM} r_s^2$$

$$\text{Width difference: } \Delta \Gamma_s^{CP} = (\Delta \Gamma_s)^{SM} \cos^2(2\chi - 2\theta_s)$$

Semileptonic asymmetry:

$$A_{SL}^s = -\text{Re}(\Gamma_{12}/M_{12})^{SM} \sin(2\theta_s)/r_s^2$$

$$S_{\psi\phi} = \sin(2\chi - 2\theta_s)$$

NP wrt to SM:

- reduces  $\Delta \Gamma_s$
- enhances  $\Delta m_s$

UT of  $B_d$  system: non-degenerated

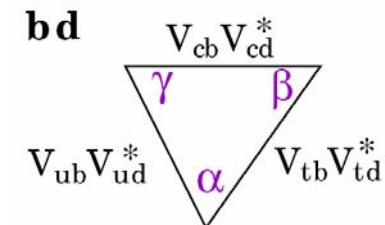
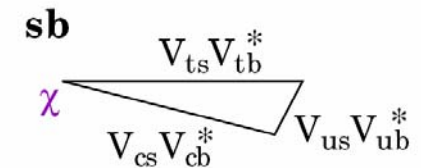
→  $(h_d, \sigma_d)$  strongly correlated to the determination of  $(\rho, \eta)$

UT of  $B_s$  system: highly degenerated

→  $(h_s, \sigma_s)$  almost independent of  $(\rho, \eta)$

$B_s$  mixing phase very small in SM:  $\chi = -1.02 \pm 0.06$  (deg)

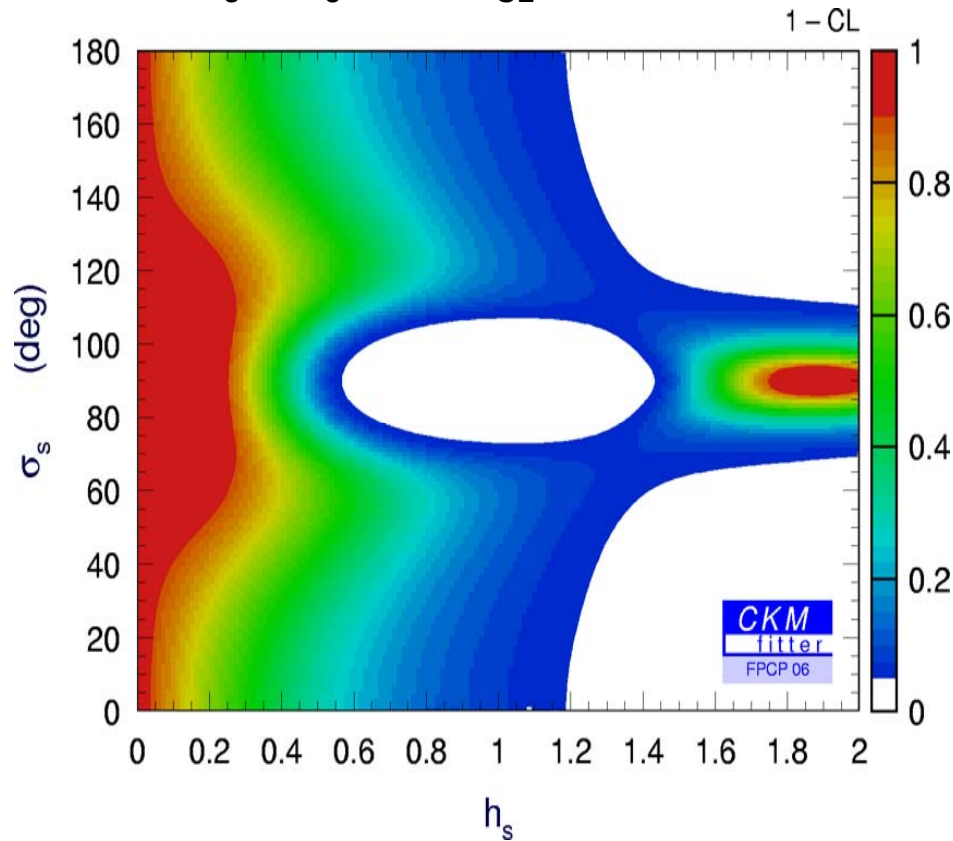
→ **Bs mixing: very sensitive probe to NP**



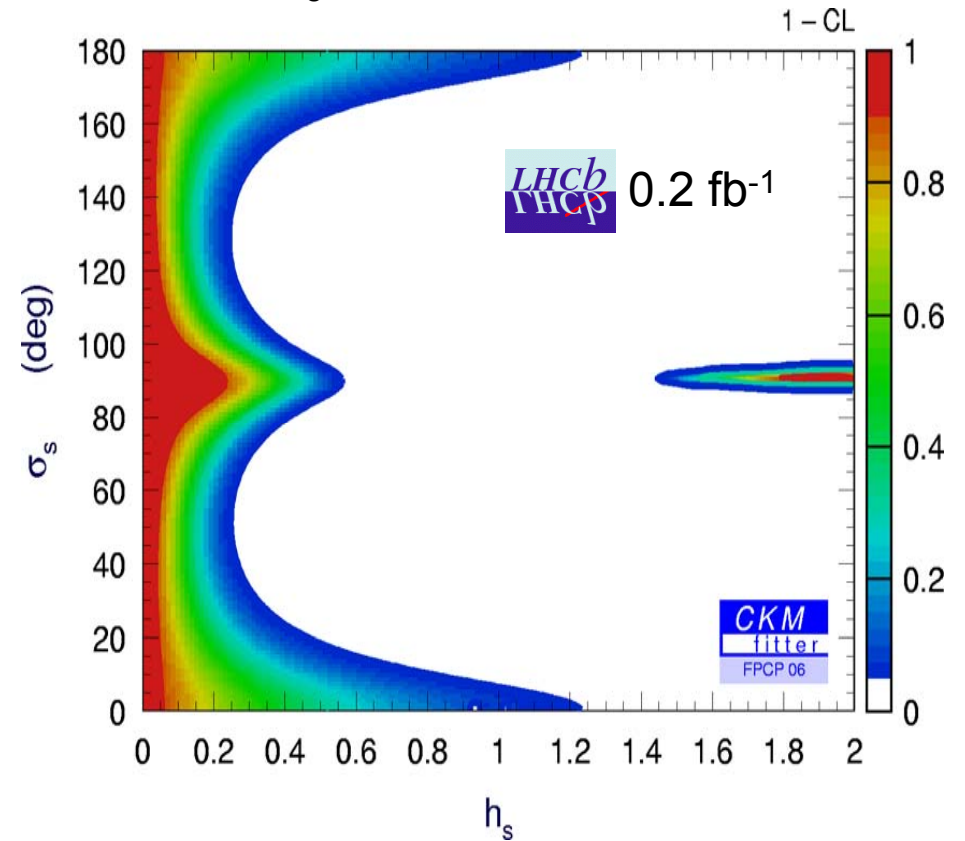


# NP in $B_s$ System

$\Delta m_s, \Delta \Gamma_s$  and  $A_{SL}^s$



$\sigma(\Delta m_s) = 0.035, \sigma(\sin(2\chi)) = 0.1$



First constraint for NP in the  $B_s$  sector  
 Still plenty of room for NP  
 Large theoretical uncertainties: LQCD

$$h_s \lesssim 3 \quad (h_d \lesssim 0.3, h_K \lesssim 0.6)$$

# B<sub>s</sub>-mixing phase



ICHEP06 – Conf note 5144

(Preliminary)

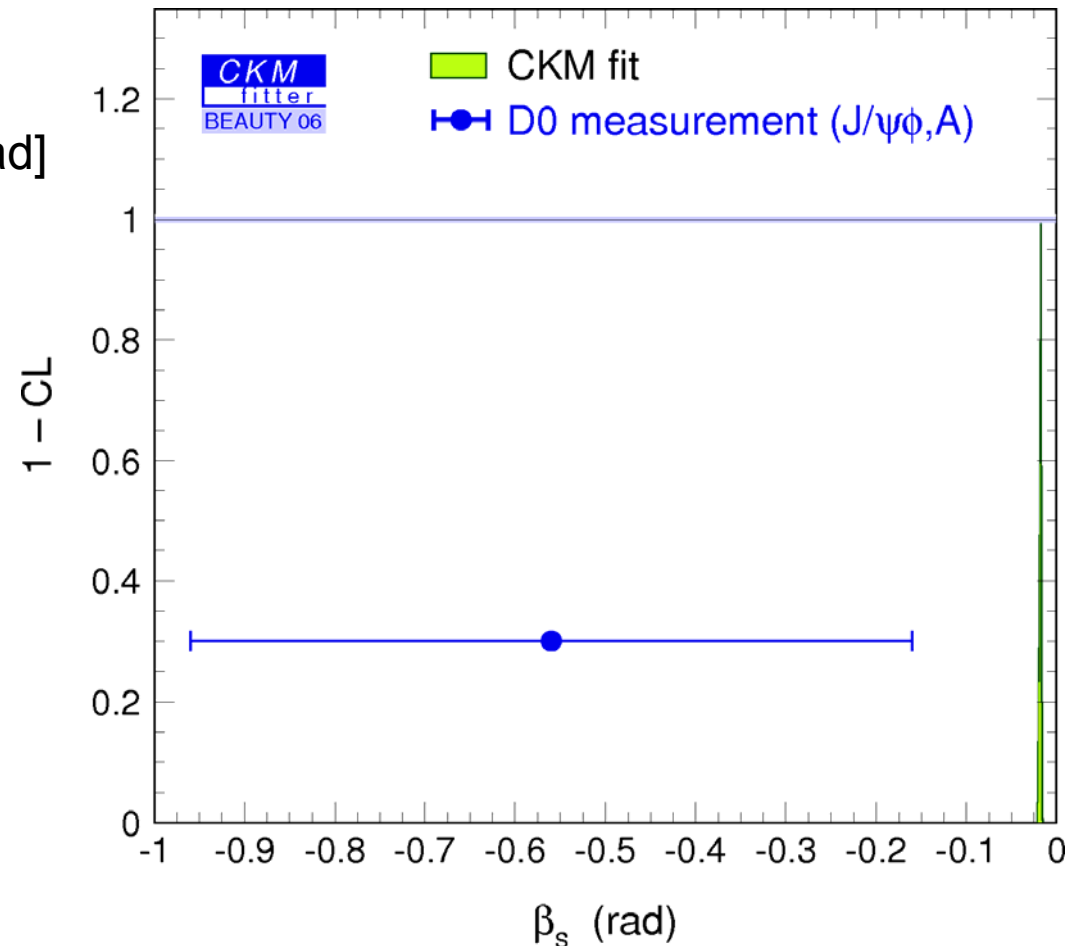
$$\beta_s = (-0.56^{+0.44}_{-0.41}) \text{ (stat+syst) [rad]}$$

Time-dependent angular  
distribution of untagged decays  
B<sub>s</sub> → J/ψφ + charge asymmetry

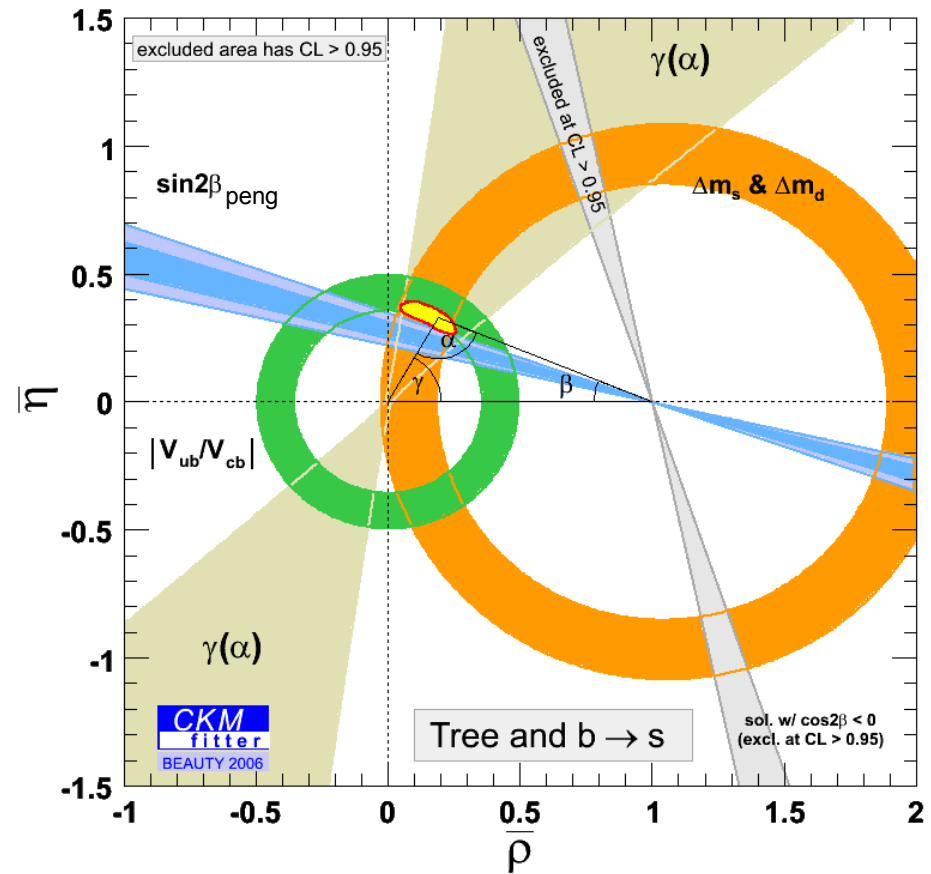
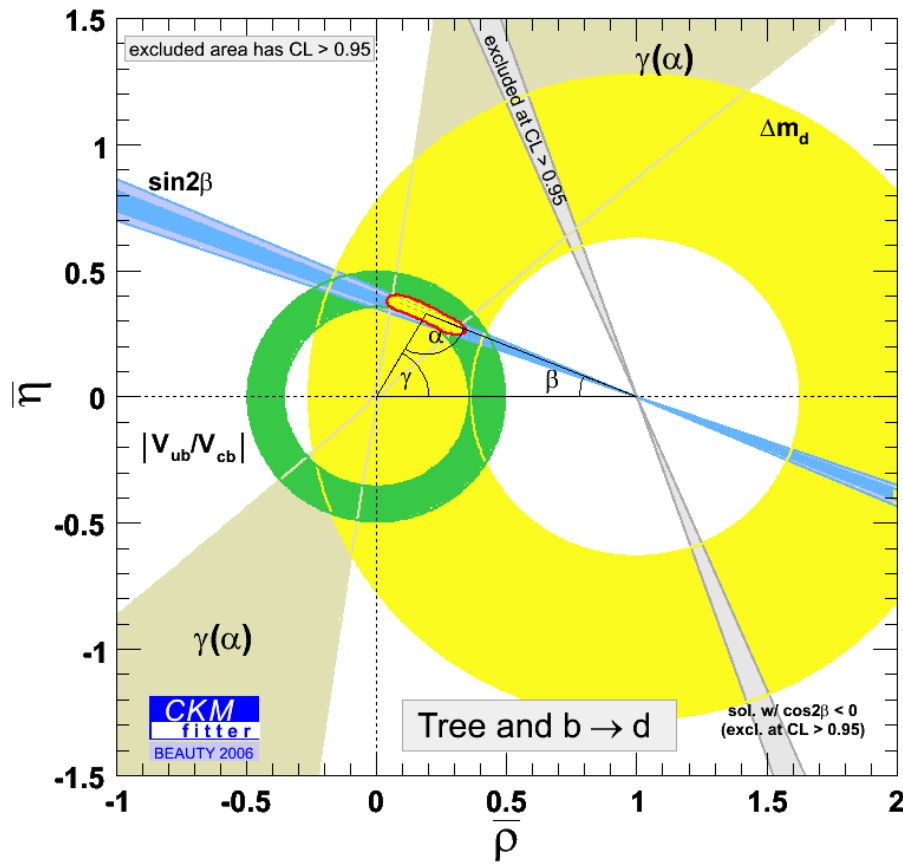
✱ Prediction from global CKM fit :

$$\beta_s = (-0.0175^{+0.0015}_{-0.0008}) \text{ [rad]}$$

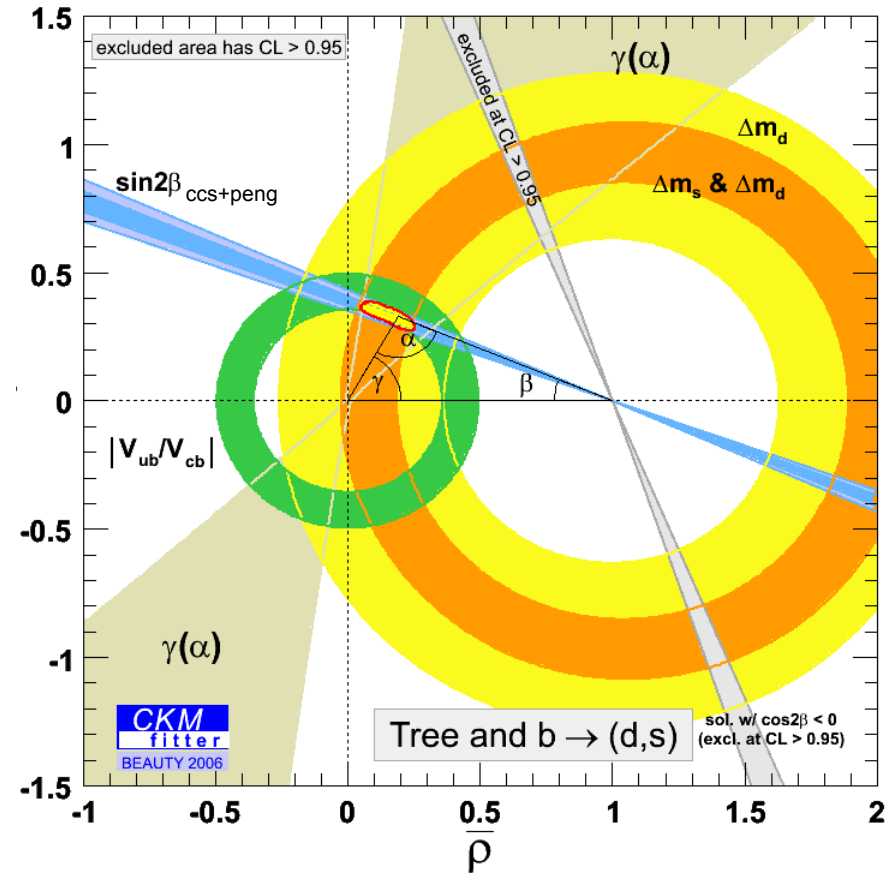
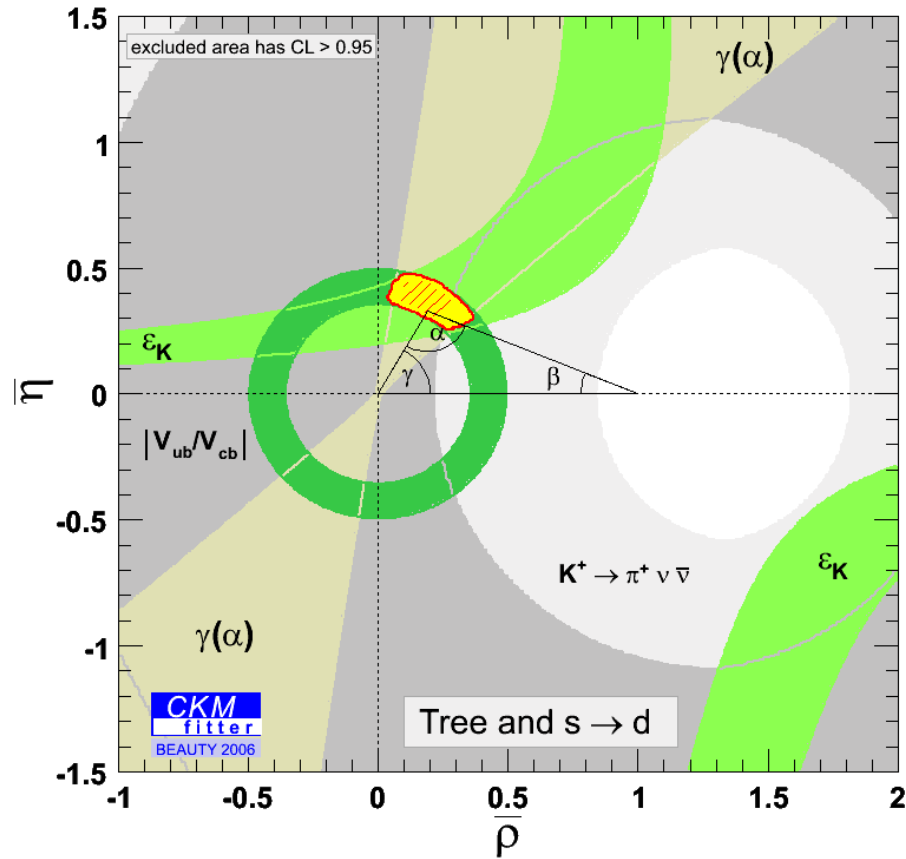
- ➔ Precision prediction
- ➔ Sensitive test to NP



# NP in $b \rightarrow s$ transitions?



# NP related solely to the third generations?



# Conclusion

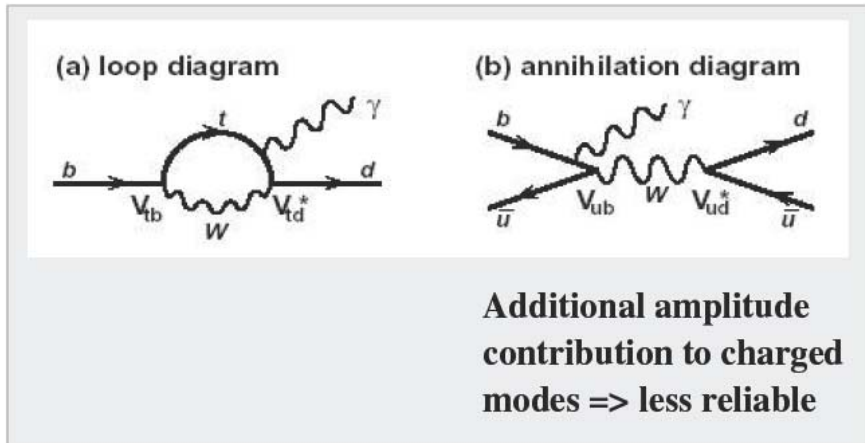
- CKM mechanism: success in describing flavor dynamics of many constraints from vastly different scales.
- Improvement of Lattice QCD is very desirable [Charm/tau factory will help]
- $B_s$ : an independent chapter in Nature's book on fundamental dynamics
  - there is no reason why NP should have the same flavor structure as in the SM
  - $B_s$  transitions can be harnessed as powerful probes for NP ( $\chi$ : "NP model killer")
- With the increase of statistics, lots of assumptions will be needed to be reconsidered [e.g., extraction of  $\alpha$  from  $B \rightarrow 3\pi, 4\pi$ , etc.,  $P_{EW}$ , ...]
- Before claiming NP discovery, be sure that everything is "under control" (assumptions, theoretical uncertainties, etc.)
  - null tests of the SM
- There are still plenty of measurements yet to be done



**BACKUP SLIDES**

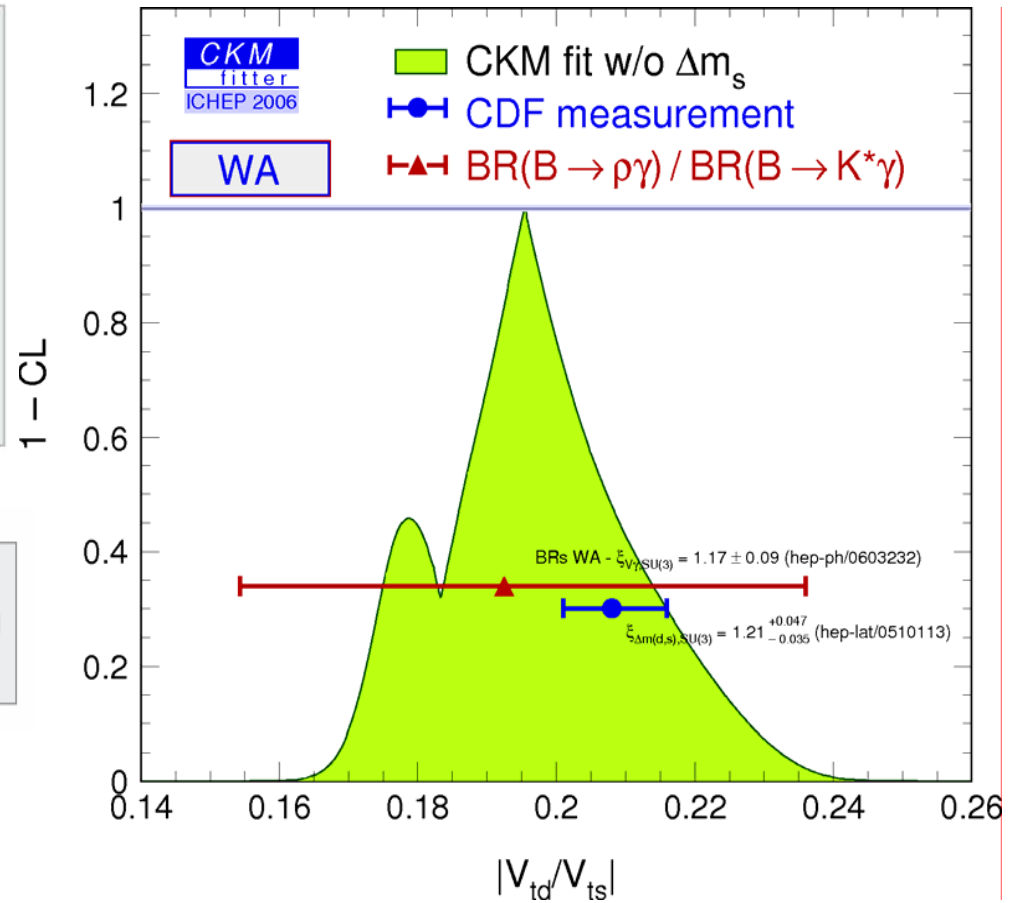
# Radiative Penguin Decays: $BR(B \rightarrow \rho \gamma) / BR(B \rightarrow K^* \gamma)$

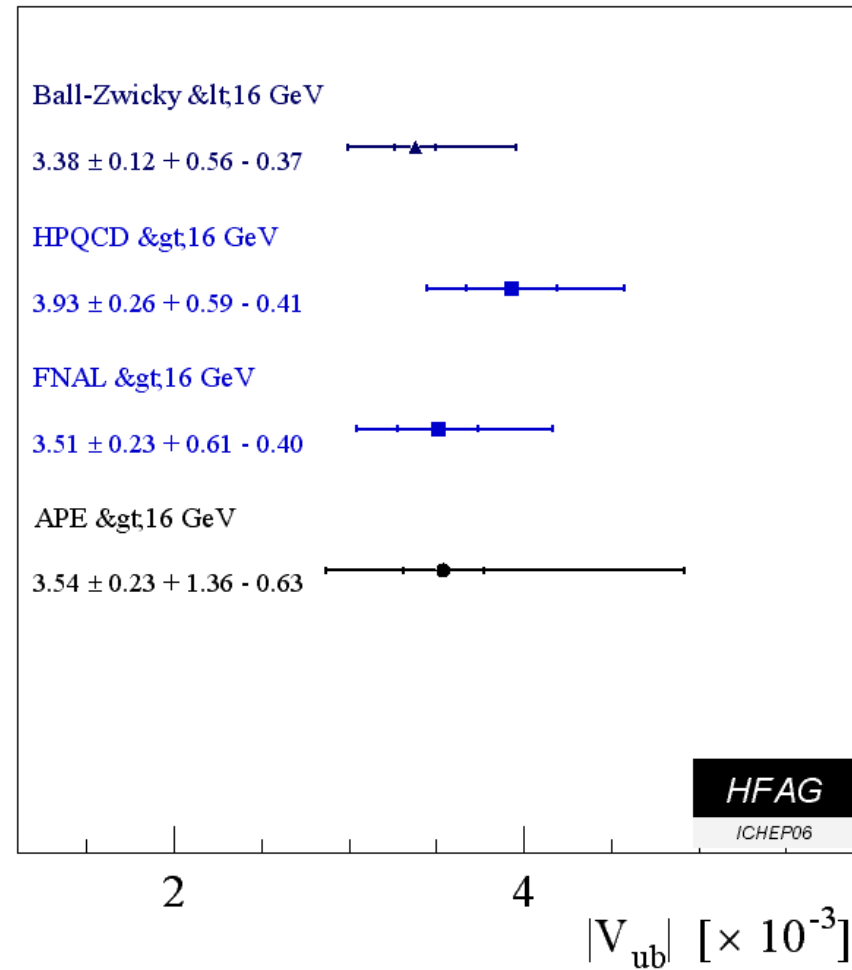
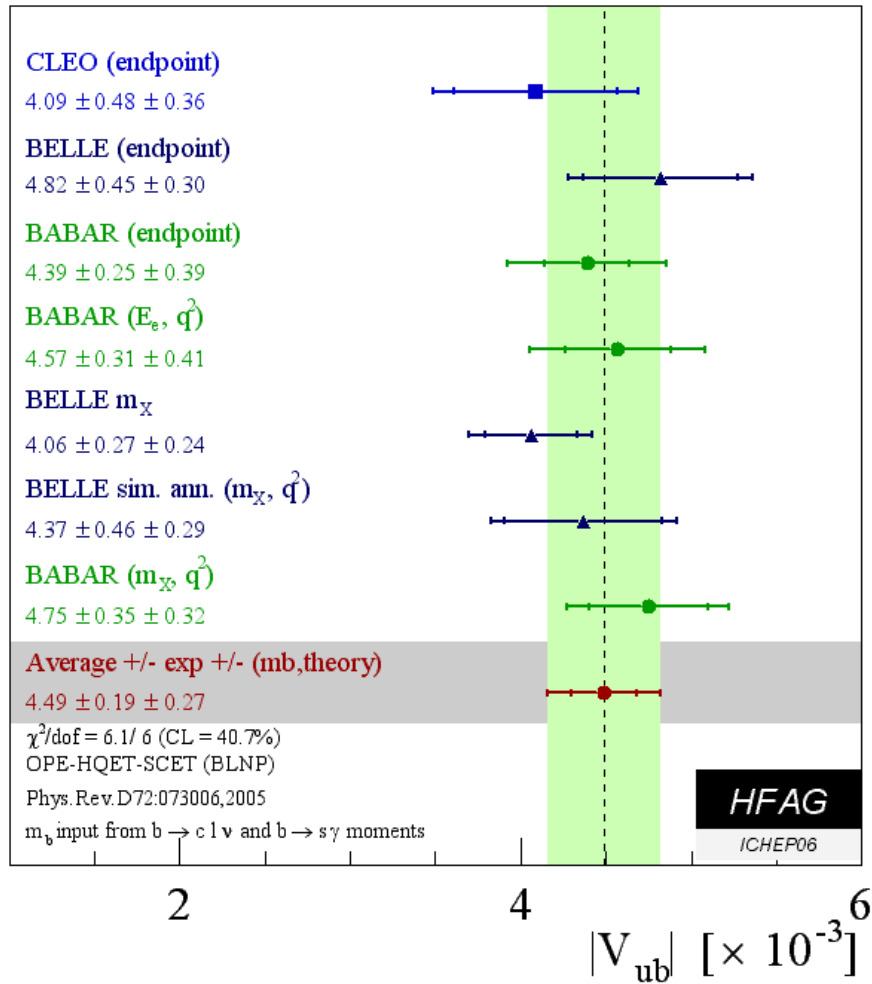
$B \rightarrow \rho \gamma (\propto |V_{td}|^2)$  &  $B \rightarrow K^* \gamma (\propto |V_{ts}|^2)$  sensitive to **New Physics**



$$\frac{BF(B^0 \rightarrow \rho^0 \gamma)}{BF(B^0 \rightarrow K^{*0} \gamma)} = 1.023 \frac{1}{2} \left| \frac{V_{td}}{V_{ts}} \right|^2 \xi^{-2} (1 + \Delta)$$

	BABAR (347m)	Belle (386m)
$\rho^0 \gamma$	$0.77^{+0.21}_{-0.19} \pm 0.07$	$1.25^{+0.37}_{-0.33} \pm 0.07$
$\rho^+ \gamma$	$1.06^{+0.35}_{-0.31} \pm 0.09$	$0.55^{+0.42}_{-0.36} \pm 0.09$







## FLAVOR STRUCTURE

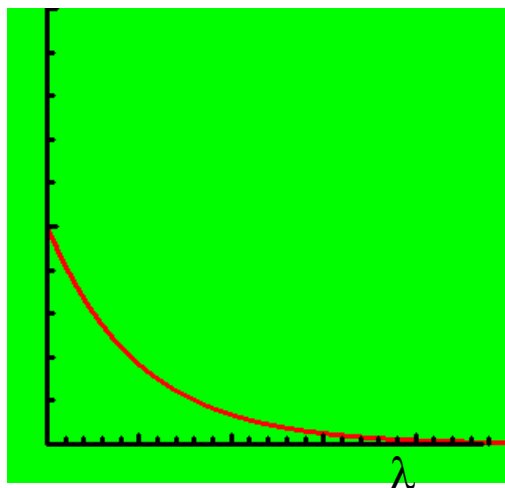
ELECTROWEAK  
STRUCTURE

	$b \rightarrow s$	$b \rightarrow d$	$s \rightarrow d$
$\Delta F=2$ box	$\Delta M_{B_s}$ $A_{CP}(B_s \rightarrow \psi\phi)$	$\Delta M_{B_d}$ $A_{CP}(B_d \rightarrow \psi K)$	$\Delta M_K, \epsilon_K$
$\Delta F=1$ 4-quark box	$B_d \rightarrow \phi K, B_d \rightarrow K\pi, \dots$	$B_d \rightarrow \pi\pi, B_d \rightarrow \rho\pi, \dots$	$\epsilon'/\epsilon, K \rightarrow 3\pi, \dots$
gluon penguin	$B_d \rightarrow X_s \gamma, B_d \rightarrow \phi K,$ $B_d \rightarrow K\pi, \dots$	$B_d \rightarrow X_d \gamma, B_d \rightarrow \pi\pi, \dots$	$\epsilon'/\epsilon, K_L \rightarrow \pi^0 \ell^+ \ell^-, \dots$
$\gamma$ penguin	$B_d \rightarrow X_s \ell^+ \ell^-, B_d \rightarrow X_s \gamma$ $B_d \rightarrow \phi K, B_d \rightarrow K\pi, \dots$	$B_d \rightarrow X_d \ell^+ \ell^-, B_d \rightarrow X_d \gamma$ $B_d \rightarrow \pi\pi, \dots$	$\epsilon'/\epsilon, K_L \rightarrow \pi^0 \ell^+ \ell^-, \dots$
$Z^0$ penguin	$B_d \rightarrow X_s \ell^+ \ell^-, B_s \rightarrow \mu\mu$ $B_d \rightarrow \phi K, B_d \rightarrow K\pi, \dots$	$B_d \rightarrow X_d \ell^+ \ell^-, B_d \rightarrow \mu\mu$ $B_d \rightarrow \pi\pi, \dots$	$\epsilon'/\epsilon, K_L \rightarrow \pi^0 \ell^+ \ell^-,$ $K \rightarrow \pi \nu \nu, K \rightarrow \mu\mu, \dots$
$H^0$ penguin	$B_s \rightarrow \mu\mu$	$B_d \rightarrow \mu\mu$	$K_{L,S} \rightarrow \mu\mu$

# Bayes at work

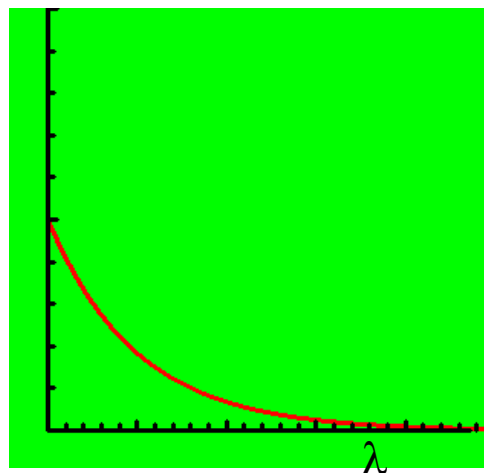
Zero events seen

$$P(n; \lambda) = e^{-\lambda} \lambda^n / n!$$



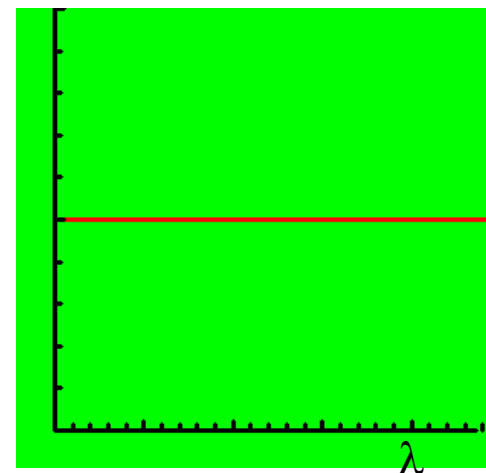
Posterior  $P(\lambda)$

=



$P(0 \text{ events} | \lambda)$   
(Likelihood)

x



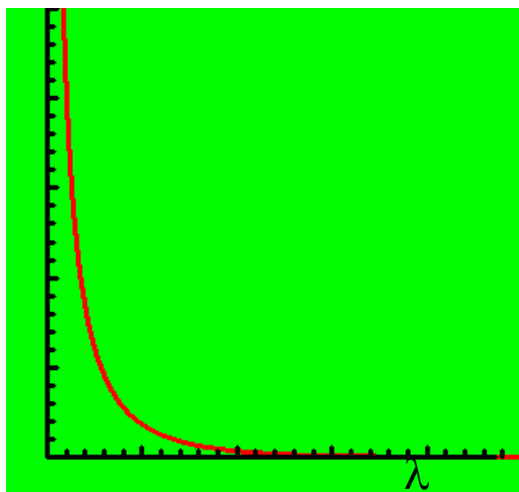
Prior: uniform

$$\int_0^3 P(\lambda) d\lambda = 0.95$$

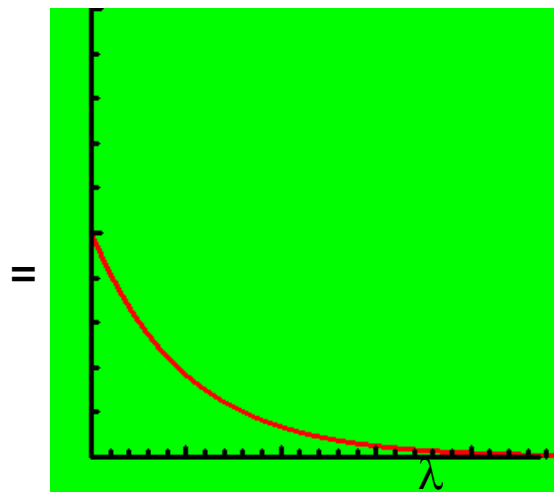
Same as Frequentist limit -  
Happy coincidence

# Bayes at work again

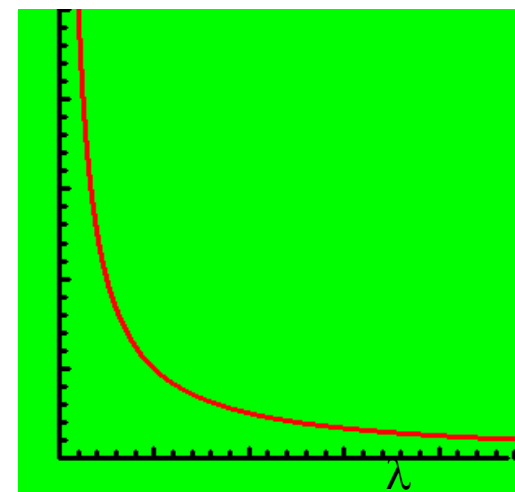
Is that uniform prior really credible?



Posterior  $P(\lambda)$



$P(0 \text{ events}|\lambda)$



Prior: uniform in  $\ln \lambda$

Upper limit totally different!

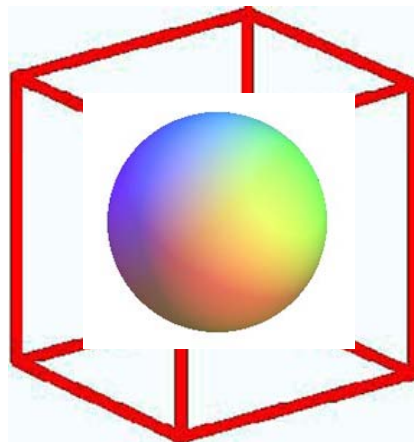
$$\int_0^3 P(\lambda) d\lambda \gg 0.95$$

# Bayes: the bad news

- The prior affects the posterior. It is your choice. That makes the measurement subjective. This is BAD. (We're physicists, dammit!)
- A Uniform Prior does not get you out of this.
- Beware snake-oil merchants in the physics community who will sell you Bayesian statistics (new – cool – easy – intuitive) and don't bother about robustness.

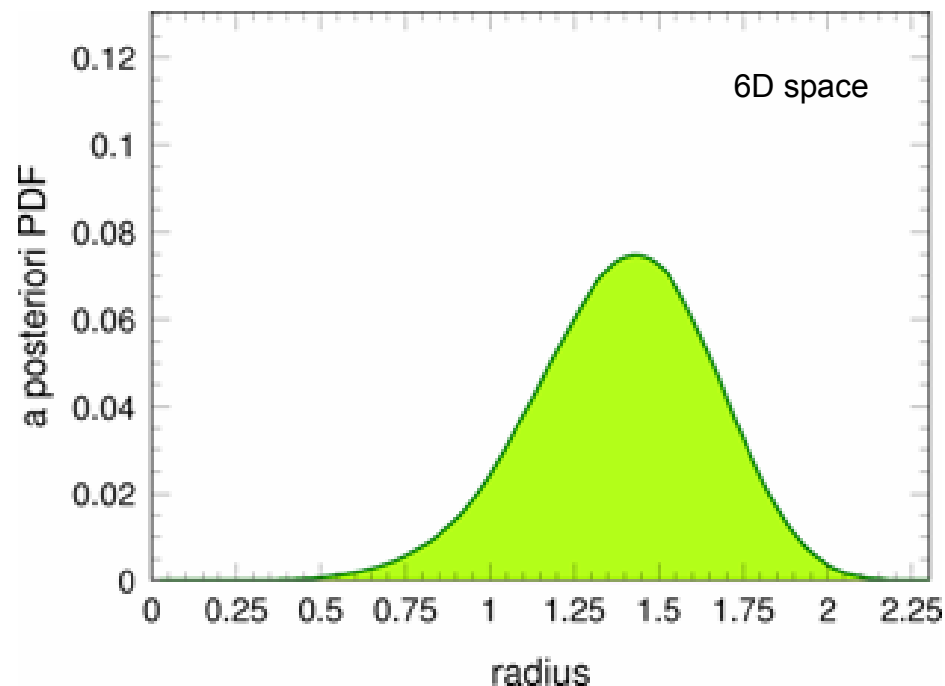
## Digression: Statistics(cont.)

Hypersphere:



One knows nothing about the individual Cartesian coordinates  $x, y, z, \dots$

What do we know about the **radius**  
 $r = \sqrt{(x^2 + y^2 + \dots)}$  ?



One has achieved the remarkable feat of learning something about the radius of the hypersphere, whereas one knew nothing about the Cartesian coordinates and without making any experiment.