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▶ To cite this version:

Antoine Girard, Pierre Bastie, Jean-Louis Soubeyroux. Real time in situ hard X-ray texture evolution during the annealing of rolled CuNi tapes. Applied Physics A, Springer Verlag, 2007, 90 (1), pp.179-184. <10.1007/s00339-007-4254-8>. <hal-00264753>

HAL Id: hal-00264753 <https://hal.archives-ouvertes.fr/hal-00264753>

Submitted on 18 Mar 2008

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Real Time in situ hard X-ray texture evolution during the annealing of rolled CuNi tapes

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Abstract: A new real time in-situ scattering method to study texture evolution kinetic is exposed. The technique is adapted from an existing set up used at the Institut Laue-Langevin neutron facility to check monochromators. It uses a white hard X ray beam and works in transmission geometry. The 2D detector allows to follow recrystallization phenomena. A study of the annealing behavior of copper-nickel alloy rolled tapes, used as substrate for high temperature superconductor, is presented.

hard X-rays; diffraction; in-situ; annealing; recrystallization; CuNi 61.66.Dk / 61.10.Eq 61.18.-j / 81.05.Bx / 81.40.Ef

1. Introduction

In material science field it is often necessary to check crystallographic property evolution with temperature. Basic experiments could be done with ex-situ measurements, i.e. material properties are tested after a specific heat treatment. Two disadvantages can be noticed: first, the experiments are done at room temperature after cooling, which means that the samples should be quenched to reveal the crystallographic structure at high temperature. In some cases, these quenches are not possible. Secondly, to get fine evolution of the phase structure a lot of experiments are needed. An answer to this ex-situ problem is to develop real time in-situ measurements. Some common experiments are available to look for phase changes on polycrystalline material like X-ray diffraction used with a Debye-Scherrer geometry, a position-sensitive detector and an heated sample holder [1]. But this device can only reveals phase transformation of

polycrystalline materials and is not sensitive to texture changes. In the case of a monocrystal or a not random distributed orientated polycristal, crystallographic changes versus temperature are much more difficult to analyse. Calorimetric experiments can give some information about phase transformation but are few sensitive to texture evolutions due to the small energy involved. Moreover, this can not inform about the crystallographic orientation.

The purpose of this paper is to describe an original real time in-situ experiment able to show texture evolution versus temperature of a polycrystalline tape. This experiment is adapted from an existing set-up used at the Institut Laue-Langevin neutron facility of Grenoble by the monochromator group [2]. It uses hard X-ray photons which allow working in transmission through the all sample thickness, and give information of the whole sample structure [3]. To illustrate possibilities given by this new method, analysis were done on two copper-nickel alloys rolled tapes which are developed as substrate for high temperature coated superconductors. These substrates should present a biaxially textured surface to favor the epitaxial growth of upper layers, whose one of them is the superconductor ceramic (i.e. YBCO). Indeed this oxide should be well oriented to present good superconducting properties. Cu₅₅Ni₄₄Mn (known as Constantan) and $Cu₇₀Ni₃₀$ alloys were chosen for their cell parameters closed to the YBCO one and for their mechanical and magnetic properties [4].

2. Experimental specifications

A white X-ray beam issued from a 1.5×1.5 mm² focus size was used. The Philips MCN 421 X-ray generator with a tungsten anticathode may supply photons with an energy up to 420 kV. In this experiment the maximum voltage was 300 kV with an intensity of 3 mA. The spectrum has a continuous high energy background which corresponds to the bremsstrahlung (braking radiation) on which the two characteristic lines of the W tube superimpose. They are located at 59 keV (K_a=0.21 Å) and 67 keV (K_β=0.18 Å), see figure 1. The beam is collimated with a 3 mm diameter hole located at 3.2 m from the source, giving a 0.08° divergent beam.

Figure 1 Energy X-ray spectrum used for experiment.

The two-dimensional detector is composed of an X-ray image intensifier with 22 cm diameter (Thomson TH 9428) coupled with a cooled CCD camera with weak noise (Princeton 512×512 pixels). This detector can be moved laterally and vertically to be aligned with the direct beam. The important X-ray intensity, reinforced by the focusing effect and the very high detector sensitivity, allows a short acquisition time, typically 100 s for one measurement. A "beam stop" aligned on the detector center attenuates the transmitted rays to avoid the saturation of the detector but also allows to follow the direct beam. The tape was vertically applied on a sample-holder which was fixed at a stick equipped with thermocouples (figure 2). The furnace was used up to 1030°C, with a heating rate of 5°C/min. The outer part of the furnace was cooled by a water jacket, and two silicon monocrystal windows allow the X-rays beam path. The furnace was placed on a motorized goniometer allowing an adjustment of the angular sample positioning. Vertical and horizontal shifts were ensured manually by moving the lifting truck. A turbomolecular pump ensures a 10^{-4} to 10^{-5} mbar vacuum inside the furnace.

Figure 2 Sample-holder, with the sample.

Currently used to perform mosaïcity measurements on monocrystals, the line is originally built with a symmetrical ratio, i.e. the source-sample distance is equal to the sample-detector distance (usually 3.6 m). With this particular condition, if the crystal is oriented with an angle satisfying the Bragg conditions, all diffracted X-Rays are focused on a thin line [5-7]. In our case, measurements were done to follow the texture evolution in a polycristal sample. The main important fact is then to collect intensities from the two characteristic emission peaks. When the sample is too thick to enable the overlapping of the K_{α} and K_{β} lines, a local energy analyzer is used to separate intensities from each family plane (NaI scintillator for low energy resolution or cooled Ge diode for high resolution analysis are used). But in this study no energy detector was used as samples are 150 μm thick and allowed the transmission of these characteristic K_α and K_β lines. Considering the first interreticular distances of the alloy cells $(d_{111}=2.06 \text{ Å}$ and d_{200} =1.78 Å) and the K_α, K_β wavelengths, the 2θ angles are between 5^o and 14^o. With a 22 cm detector diameter, the sample-detector distance should be smaller than 90 cm. In this experiment a non-symmetrical transmission geometry was adopted with a 60 cm sample-detector distance. A picture of the assembly is shown figure 3 where all experimental devices are visible except the controlacquisition board.

Figure 3 View of the hard X-ray scattering experimental assembly.

When the beam and the detector are aligned, a "white experiment" (without sample) is carried out (figure 4). With this diagram, intrinsic intensities of the alignment are revealed, like: absorption of the beam-stop and its holder, a more or less homogeneous background and small spots (due to the diffraction of the monocrystal silicon windows of the furnace). This background has been substrate to all experimental diffraction patterns.

Figure 4 2D diffraction diagram without sample.

Finally, a radiographic image (i.e. using a broad beam with low intensity) is performed to ensure the good sample positioning in the X-ray beam. All the data computations were done with the LAMP software [8].

3. Experimental results and discussion

2D diffraction diagrams of $Cu_{70}Ni_{30}$ tape are recorded each 10°C during the annealing (5° C/min). Figure 5 presents four of them. a) is the diffraction pattern at the beginning: 30°C. That means that the tape shows the rolled structure after 98% deformation. b) and c) are respectively recorded at 610°C and 1000°C (maximum temperature of the experiment). d) is the diffraction diagram of the sample after cooling at room temperature (RT).

To better understand the diagram, figure 6 shows the theoretical diffraction circles calculated for a $Cu_{70}Ni_{30}$ random powder. Each circle corresponds to the diffraction of a specific wavelength with different interreticular distances. The smaller dashed circle is the theoretical position for the (111) diffraction plane by the smallest wavelength, i.e. the highest energy: 300 keV. Inside this circle no intensity from the sample can diffract. The solid circles are the scattering positions for the (111), (200) and (220) planes with K_{α} and K_{β} energies. The $(111)K_\alpha$, and the $(200)K_\beta$ circles are very close and it is not possible to distinguish their respective contribution. Plans that have $h^2 + k^2 + l^2$ higher than 8 can not diffract in the 2D detector with K_{α} and K_{β} energies. In addition to the intensities diffracted by the W characteristic energies, the radial streak intensities are due to the diffraction by wavelengths of the bremsstrahlung.

The use of transmission geometry and high energy radiation allowed to be sensitive to the diffraction planes which are almost parallel to the normal of the tape surface. These intensities correspond to the ones observed in the pole figure periphery. To simplify the correlation between intensities observed in figure 5 with a given orientation, theoretical pole figures of the tape are presented in figure 7. The three upper figures are the (111), (200) and (220) pole figures of a rolling tape. The represented intensities are those of a fcc metal with a "pure metal" type deformation structure. These structures are composed with 8 different orientations which are the S = $\{123\}$ <63 4>, the C = $\{112\}$ <11 1>, the B = $\{110\}$ <1 12> and their symmetric orientations. The three lower figures are the (111), (200) and

(220) pole figures of the cubic orientation {100}<001>. This structure is the annealed texture of a tape that acquires a C+S+B structure after rolling.

Figure 5 2D diffraction pattern of a $Cu_{70}Ni_{30}$ rolled tape a) not annealed, b) during the annealing process at 610°C, c) at 1000°C and d) after cooling at room temperature.

Figure 6 Theoretical diffraction circle positions for a $Cu_{70}Ni_{30}$ random powder for different energies and different planes. The measured pattern is the same that the figure 6b. On the figure 5a it is possible to observe all the intensities from the rolled tape diffraction pole figures periphery. For example on the transverse direction (TD) the intensities correspond to the B orientation diffraction of the (111) plane and are well located on the $(111)K_a$ and $(111)K_b$ circles. Similar observations can be made at 15° on each side of the rolling direction (RD) also on the $(111)K_a$ and $(111)K_B$ circles and correspond to the B and S orientation diffractions of the (111) plane. The weak intensities on the RD on the same circles are due to the C orientation diffractions of the (111) plane. On the $(200)K_a$ and $(200)K_\beta$ circles at 35° on each side of the RD, intensities are due to the B orientation diffractions of the (200) plane. Figure 5 are then well representative of a rolling deformation texture.

On the figure 5c, intensities from the cubic orientation are visible on the RD et TD if considering the (200)K_a and (200)K_β circles and on the bisecting lines if considering the (220)K_a and (220)K_β circles. No intensity from the K_α and K_β energy on the $(111)K_\alpha$ and $(111)K_\beta$ circles can be seen, as no intensity is observed in the (111) pole figure periphery. Intensities are less important in the half lower part because the tape is probably bent slightly with the temperature as the sample is only fixed to one side with screws.

Figure 5b is recorded during annealing at 610°C. Both textures, rolled and annealed, can be seen.

Figure 5d is the diffraction pattern recorded at RT after cooling down. Intensities are in the same place than these observed on figure 5c but they are less spread with a higher maximum level. The difference is probably not due to structure changes but should be the consequence of a better structure of the diffusion pattern at low temperature through the Debye-Waller coefficient [9].

Figure 7 (111), (200), (220) pole figures with theoretical intensity localizations which can be seen on a fcc material with a "pure metal" type after a heavy rolling (up), then annealed (bottom). Marks represent specific orientations (and their equivalences) = $S = \{123\} \le 63-4$, = C = ${112}{\leq11\text{-}1$ = B = {110} ${<}1\text{-}12$ >, = cubic = {100} ${<}001$ >

To allow a better visualization of the transition, intensities from a small horizontal stripe which goes through the centre for each recorded pattern, are plotted versus temperature (figure 8a). Figure 8b is the same but considering the vertical stripes. On these graphs, it is possible to follow the intensity evolutions of the $(111)K_{\beta}$, $(111)K_{\alpha}$, $(200)K_{\beta}$ and $(200)K_{\alpha}$ lines. As in one hand the $(111)K_{\beta}$ line has too small intensities and in the other hand $(111)K_\alpha$ and $(200)K_\beta$ lines are to close to separate their respective contribution, the better choice is to follow the $(200)K_a$ evolution. Contributions of each 4 lines are summed and plotted with temperature (figure 9, blue line). The same works have been done with a Constantan tape which has a higher nickel content: Cu₅₅Ni₄₄Mn (red line).

With these two curves, the recrystallisation kinetics can be followed. In the insert the beginning transformation temperatures may be evaluated. The intersection of the two tangents points out the temperatures: 582° C and 595° C for Cu₇₀Ni₃₀ and Constantan respectively. These values fit well with other measurements performed elsewhere with other diffraction techniques [10]. The small difference between these temperatures is probably due to higher nickel content of the Constantan which increases the thermic phenomena to higher temperature level.

Figure 8 TD (a) and RD (b) cross section of each recorded scattering pattern plot versus annealing temperature.

The curves show that the transformation is very abrupt, as the half of the maximum intensity is reached in 50°C which is equivalent to 10 min. The interpretation of the transformation around the maximum (between 900°C and 950°C) needs much more experiments to better understand the contributions to the phenomena. Indeed, in one hand biaxial texture breaks down after a critical temperature which leads the intensity to decrease, and in the other hand intensity normally decreases even for a stable crystal due to Debye-Waller coefficient.

Figure 9 Sum of the 4 contributions of the (200) K α peak is plotted with annealing temperature for both Constantan and $Cu_{70}Ni_{30}$ tapes.

Some quantitative informations about the sharpness of the annealed texture can also be extracted from the data. The out-of-plane disorientation, normally given by the rocking curve may be directly checked by measuring the full width half maximum (FWHM) of the $(200)K_\alpha$ intensity distribution (perpendicularly to the radial direction). Out-of-plan disorientation in the rolling direction (RD) and in the transverse direction (TD) can be separately measured by checking respectively the intensity distribution in the RD or TD. Figure 10 shows the evolution of the intensity distribution corresponding to a rocking curve in the RD.

Figure 10 (200) K_a intensity of the rolling direction recorded during the annealing of a Constantan tape. Arrows represent the FWHM of respective curves at different temperatures.

Despite an increasing of the intensity, FWHM does not present any significant variation between 700°C and 1000C°. The out-of plan disorientation (RD) is about 7° which corresponds more or less to figures find in a normal rocking curve measurement at RT on an annealed tape by X-ray pole figures.

4. Conclusions

This original real time in-situ hard X-ray diffraction method first developed for monocrystal measurements provides also interests for non-random polycristal studies. Indeed with this transmission geometry alignment, texture evolution could be followed in real time during heating treatments. Qualitative, as well as quantitative, results are available.

First experiments done on the annealing behaviour of copper-nickel alloys rolled tapes, have allowed to validate the method, the setting and the new dedicated furnace. Interesting biaxially textured recrystallisation kinetics data have been recorded. The beginning of the transformation has been found at 580-595°C depending of the nickel content. Fast evolution is then noticed before that the intensity slowly decreases after a maximum value. The interpretation of this last phenomena could be explained either by the untexturation of the material or the intrinsic decrease of the intensity with high temperature (Debye-Waller coefficient). Quantitative out-of-plan disorientation may also be checked with this method.

This real time in-situ, non destructive technique able to check texture evolution versus temperature in bulk material by using relatively cheap device is surely promising for future developments.

Acknowledgements: The authors wish to thank Mr. A. Daramsy and ILL staff for their technical assistance. This work has been carried out in the frame of the project: "Superflex" with the support of the "Région Rhône-Alpes".

References

- 1. P. Bénard, J.P. Auffrédic,. D. Louër: Materials Science Forum 228-231, 325 (1996)
- 2. B. Hamelin: Annual Report ILL 1998, 98 (1998)
- 3. T. Chauveau, P. Gerber, P. Bastie, B. Hamelin, J. Tarasiuk, B. Bacroix: Journal de Physique IV 12, 107 (2002)

4. A. Girard, C.E. Bruzek, J.L. Jorda, L. Ortega, J.L. Soubeyroux: Journal of Physics: Conference Series (EUCAS 2005) 43, 341 (2006)

5. B. Hamelin, P. Bastie: Journal de Physique IV 8, 3 (1998)

6. B. Hamelin, P. Bastie: Proceeding of SPIE 2002 4786, 29 (2002)

7. B. Hamelin, P. Bastie, D. Richard, A. Elaazzouzi: Journal de Physique IV 118, 437 (2004)

8. ILL-Computing-dept LAMP the Large Array Manipulation Program:

http://www.ill.fr/data_treat/lamp/front.html

9. E.F. Kaelble: Handbook of X-rays (Mc Graw-Hill Book Company New-York 1967)

10. J.L. Soubeyroux, C.E. Bruzek, A. Girard, J.L. Jorda: IEEE Transactions on Applied Superconductivity, ASC2004 15, 2687 (2005)